

V_{ub} and weak annihilation in inclusive semileptonic D/D_s decays

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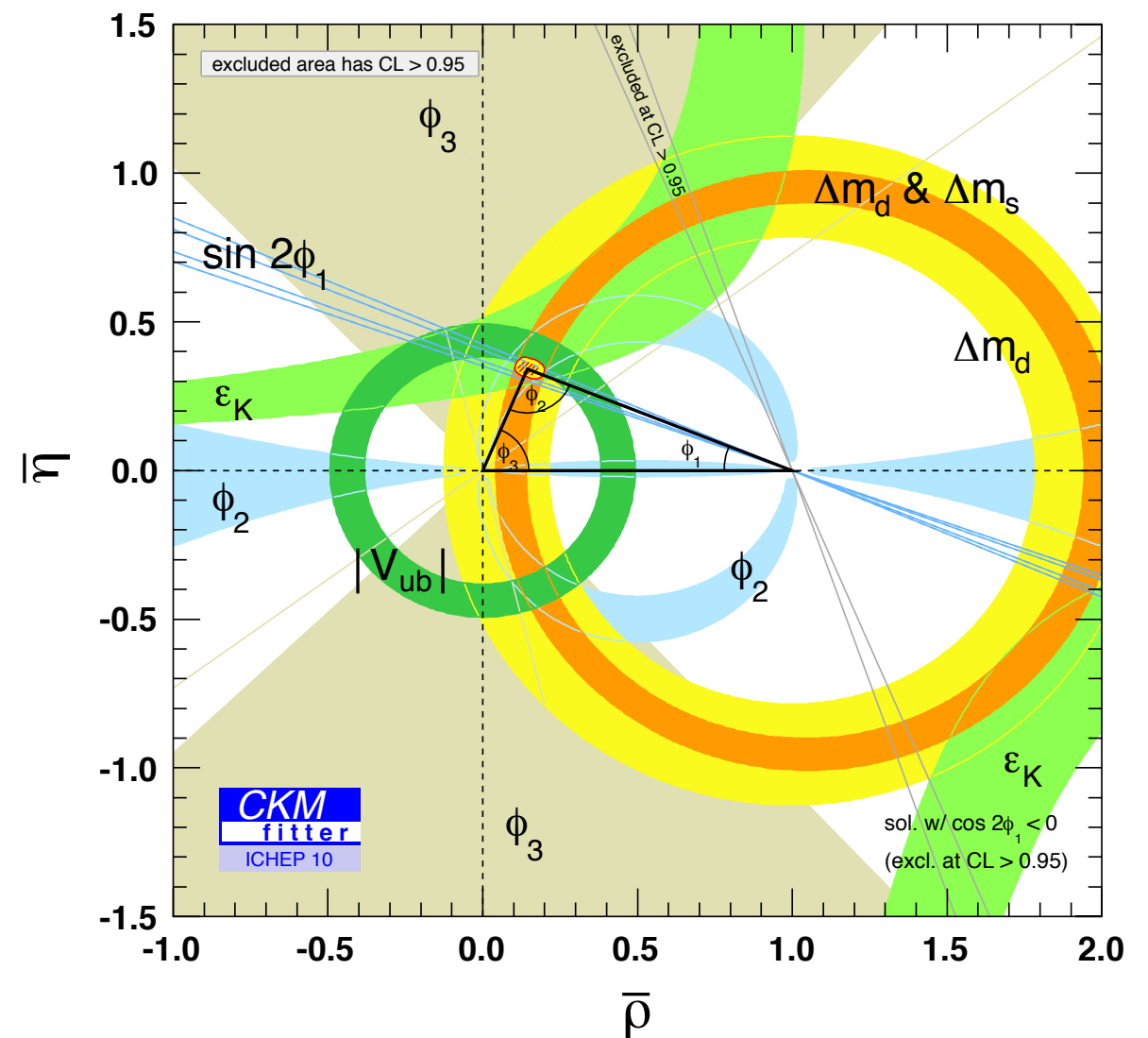
Motivation: CKM Unitarity Analysis

- UTA within the SM

$$\epsilon_K, \Delta m_d, \left| \frac{\Delta m_s}{\Delta m_d} \right|, \left| \frac{V_{ub}}{V_{cb}} \right|$$

- relying on theoretical calculations of hadronic matrix elements

CKMFitter @
ICHEP 2010

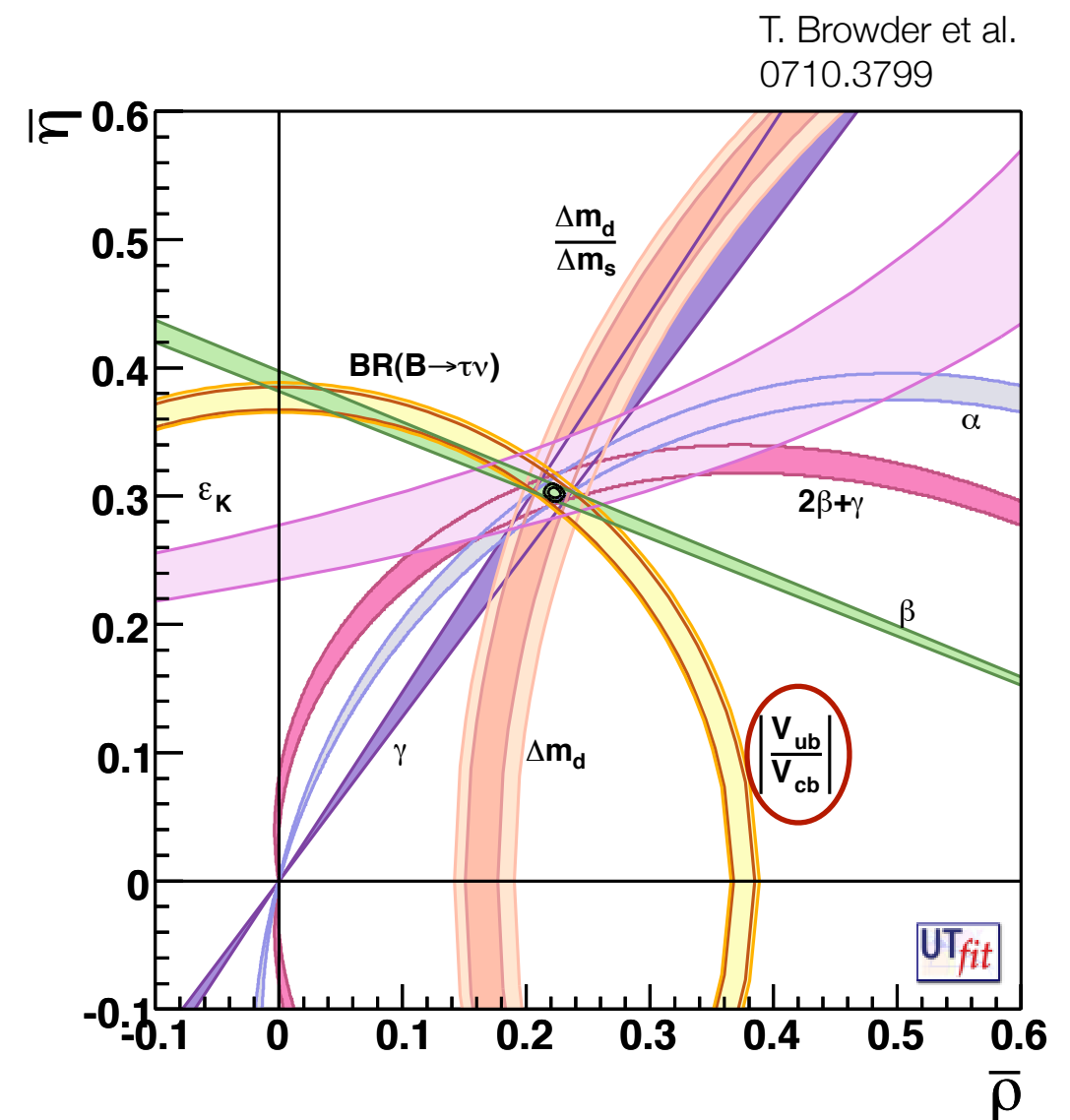


Motivation: CKM Unitarity Analysis

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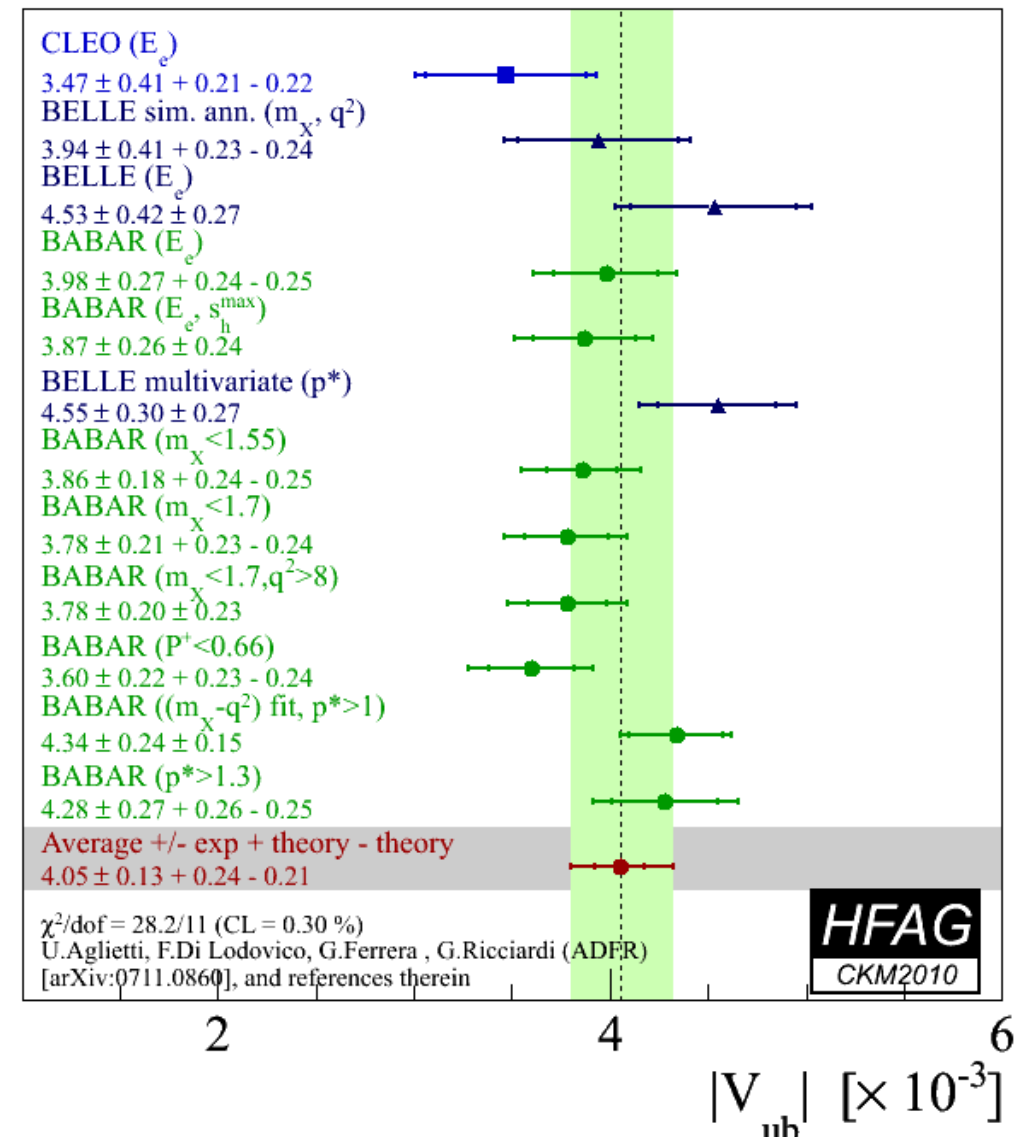
- relying on theoretical calculations of hadronic matrix elements
- Projected Super Flavour Factory sensitivity
 - V_{ub} (exclusive): 3-5%
 - V_{ub} (inclusive): 2-6%



Status of $B \rightarrow X_u \ell \nu$

- Inclusive determination of V_{ub} using OPE and HQE
 - Expansion in α_s and $1/m_b$
- Present precision around 6-7%
 - however 15% tension with UTA
 - dominant source of theoretical uncertainty due to shape-function modeling (kinematical cuts)
- A fully inclusive analysis would carry a tiny 2-3% theoretical error

Antonelli et al.
0907.5386



Status of $B \rightarrow X_u \ell \nu$

- At $1/m_b^3$ leading spectator effects due to dimension 6 four quark operators (WA contributions)
 - $16\pi^2$ phase space enhanced compared to LO & NLO contributions*
 - Affect both the total rate and spectra (expected to populate the q^2 / lepton energy endpoint region)
 - Cannot be extracted from inclusive $B \rightarrow X_c \ell \nu$ analysis
 - Nor completely from comparing B^+ and B^0 decay modes
- Difficult to study non-perturbatively

Bigi & Uraltsev
hep-ph/9310285

Dikeman & Uraltsev
hep-ph/9703437

Bigi et al.
hep-ph/9706520

Not present at dim=7*
[Dassinger et al.
hep-ph/0611168]

Uraltsev
hep-ph/9905520

Voloshin
hep-ph/0106040

D. Becirevic
hep-ph/0110124

D. Becirevic et al.
0804.1750

Existing estimates spread between 3-10%

Inclusive Semileptonic Charm Decays

- Recently determined experimentally

$$\mathcal{B}(D^+ \rightarrow Xe\nu) = (16.13 \pm 0.20 \pm 0.33)\%$$

$$\mathcal{B}(D^0 \rightarrow Xe\nu) = (6.46 \pm 0.17 \pm 0.13)\%$$

N. E. Adam et al.
[CLEO]
hep-ex/0604044

- Similar results for muons

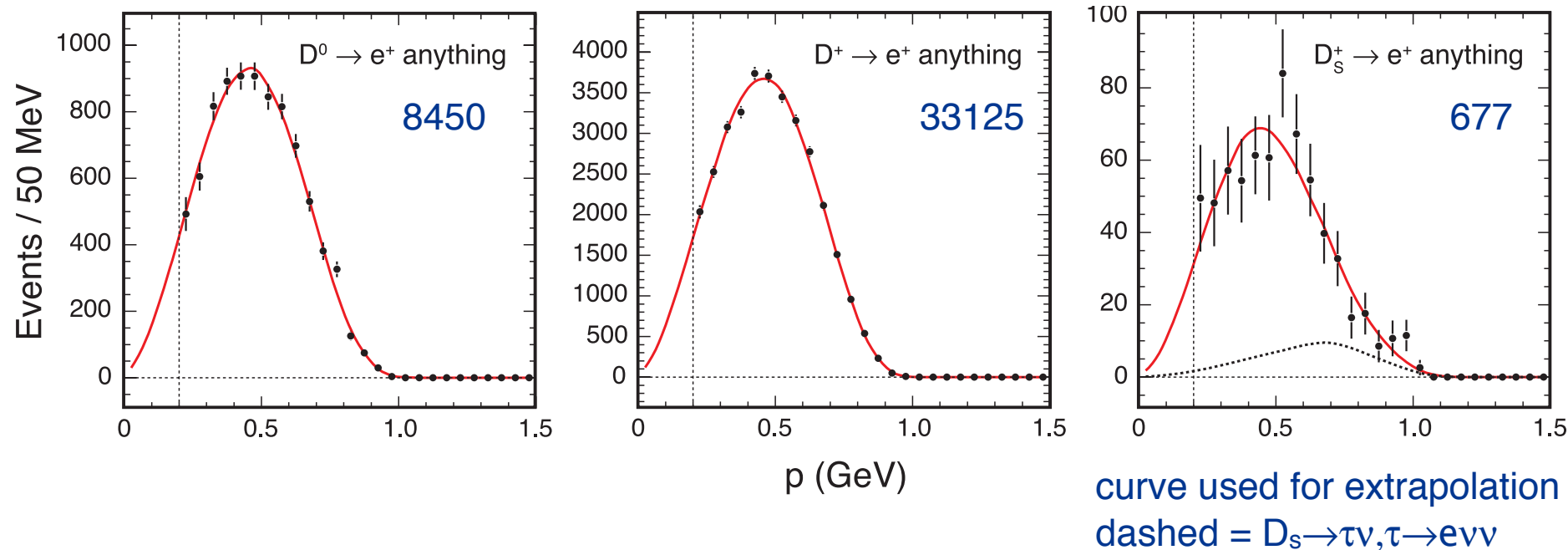
M. Ablikim et al.
[BES]
arXiv:0804.1454

- Very recently results also for D_s decays

$$\mathcal{B}(D_s \rightarrow Xe\nu) = (6.52 \pm 0.39 \pm 0.15)\%$$

Asner et al.
[CLEO]
0912.4232

- Including spectra



Inclusive Semileptonic Charm Decays

- Ratio of D_s and D^0 rates shows significant [17(6)%] deviation from unity

Asner et al.
[CLEO]
0912.4232

$$\begin{aligned}\Gamma(D^+ \rightarrow Xe^+\nu)/\Gamma(D^0 \rightarrow Xe^+\nu) &= 0.985(28), \\ \Gamma(D_s^+ \rightarrow Xe^+\nu)/\Gamma(D^0 \rightarrow Xe^+\nu) &= 0.828(57)\end{aligned}$$

- Signs of WA in D_s decays?
- How to disentangle from possible SU(3) violation?

SU(3) violation in Charm (Two examples)

- Hyperfine mass splitting $\Delta_{D_q}^{hf} = 3(m_{D_q^*}^2 - m_{D_q}^2)/4$

$$\Delta_{D^+}^{hf} = 0.409(1) \text{ GeV}^2, \quad \Delta_{D^0}^{hf} = 0.413(1) \text{ GeV}^2, \quad \Delta_{D_s}^{hf} = 0.440(2) \text{ GeV}^2$$

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- SU(3) violation at 10%

- Decay constants

- Lattice estimates: $f_{D_s} = 260(10) \text{ MeV}, \quad f_D = 217(10) \text{ MeV}$

Bazavov et al.
[Fermilab & MILC]
0912.5221

$$f_{D_s} = 260(10) \text{ MeV} \quad f_D = 217(10) \text{ MeV}$$

- SU(3) violation at 20%

Inclusive Semileptonic Charm Decays in OPE

- Treating charm quark mass as heavy, one can attempt an expansion in $\alpha_s(m_c)$, Λ/m_c
 - Need to estimate local operator matrix elements between hadronic states
 - First appear at $1/m_c^2$ \Leftarrow sources of SU(3) violation
 - Heavy quark symmetry relates these estimates between the charm and beauty sectors
 - Quantitative translation (renormalization) not straight-forward
- Alternative approach involves an educated sum over known exclusive modes

I. I. Bigi & N. G. Uraltsev,
Phys. Lett. B 280 (1992)

OPE for the rate & leptonic moments

- Rate & leptonic energy moments in HQE & OPE

- $x=2E/m_c, \quad r=(m_s/m_c)^2$

$$\Gamma^{(n)} \equiv \int_0^{(1-r)} \frac{d\Gamma}{dx} x^n dx = \frac{G_F^2 m_c^5}{192\pi^3} |V_{cs}|^2 \left[f_0^{(n)}(r) + \frac{\alpha_s}{\pi} f_1^{(n)}(r) + \frac{\alpha_s^2}{\pi^2} f_2^{(n)}(r) + \frac{\mu_\pi^2}{m_c^2} f_\pi^{(n)}(r) + \frac{\mu_G^2}{m_c^2} f_G^{(n)}(r) + \frac{\rho_{LS}^3}{m_c^3} f_{LS}^{(n)}(r) + \frac{\rho_D^3}{m_c^3} f_D^{(n)}(r) + \frac{32\pi^2}{m_c^3} B_{WA}^{(n)s} \right],$$

A. Pak & A. Czarnecki
0803.0960,

K. Melnikov
0803.0951

V. Aquila et al.
hep-ph/0503083

Czarnecki & Jezabek
hep-ph/9402326

- α_s corrections known up to α_s^2 for the total rate ($\alpha_s^2 \beta_0$ for the higher moments)

- $1/m_c$ corrections known up to $1/m_c^4$ (all present analyses use $1/m_c^3$)

Gremm and Kapustin
hep-ph/9603448

Dassinger et al.
hep-ph/0611168

- Cabibbo suppressed modes contribute to the total rate at the level of 5%, but their effect is highly suppressed in the normalized moments

WA in OPE

- WA contributions to the rate can be related to matrix elements of dim=6 four quark operators

$$\langle H_{Q\bar{q}} | O_{V-A}^{q'} | H_{Q\bar{q}} \rangle \equiv \langle H_{Q\bar{q}} | \bar{Q} \gamma_\mu (1 - \gamma_5) q' \bar{q}' \gamma^\mu (1 - \gamma_5) Q | H_{Q\bar{q}} \rangle$$

$$\langle H_{Q\bar{q}} | O_{S-P}^{q'} | H_{Q\bar{q}} \rangle \equiv \langle H_{Q\bar{q}} | \bar{Q} (1 - \gamma_5) q' \bar{q}' (1 - \gamma_5) Q | H_{Q\bar{q}} \rangle$$

- In the SU(3) limit one distinguishes between isosinglet/triplet contributions - **only the later can be estimated from the rate differences of B⁺ and B⁰**
 - Conventionally one parametrizes deviations from VSA: bag parameters
- $$\langle D | O_{V-A} | D \rangle = f_D^2 m_D^2 B_1, \quad \langle D | O_{S-P} | D \rangle = f_D^2 m_D^2 B_2$$
- Renormalization scale dependent, mix with the Darwin contributions at LO

$$\delta\Gamma \sim \left[C_{WA} B_{WA}(\mu_{WA}) - \left(8 \ln \frac{m_c^2}{\mu_{WA}^2} - \frac{77}{6} \right) \frac{\rho_D^3}{m_c^3} + \mathcal{O}(\alpha_s) \right]$$

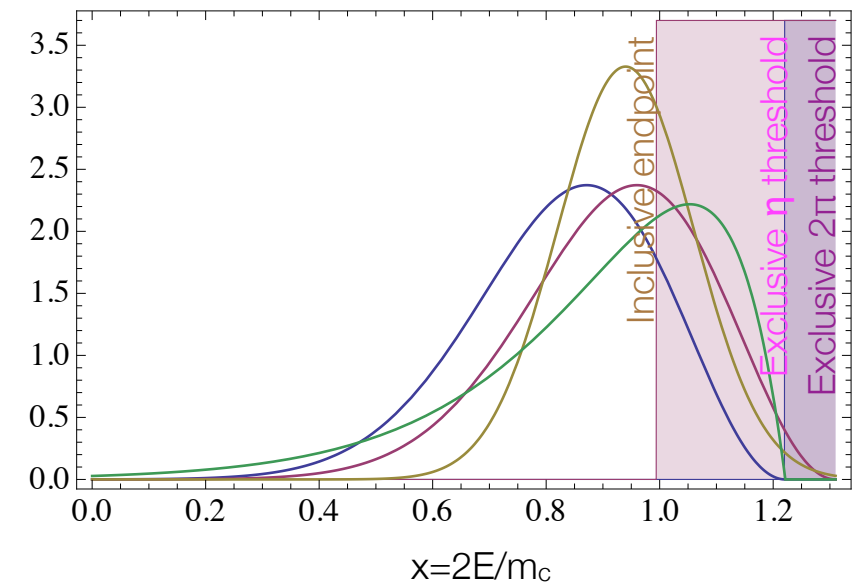
P. Gambino et al.
hep-ph/0505091,
0707.2493

- can be used to estimate WA contributions to the rate

I. I. Bigi et al.
0911.3322

Modeling WA in leptonic moments

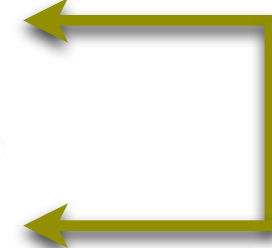
- WA contributions to the weak current correlators vanish in the OPE - need to model
- Expected to populate the spectrum endpoint
Bigi & Uraltsev
hep-ph/9310285
- Develop a perturbative tail & non-perturbative smearing
A. K. Leibovich et al.
hep-ph/0205148]
- Possible phase-space suppression by hadronic thresholds
 - Can be studied directly using exclusive channels
($D_s \rightarrow \omega \ell \nu$)
Gronau & Rosner
0902.1363



The WA interpretation of rate differences

- Without resorting to quantitative OPE predictions, one can estimate WA from rate differences

$$\begin{aligned}\Gamma_{WA}(D^0) &\propto \cos^2 \theta_c B_{WA}^s(D^0) + \sin^2 \theta_c B_{WA}^d(D^0), \\ \Gamma_{WA}(D^+) &\propto \cos^2 \theta_c B_{WA}^s(D^+) + \sin^2 \theta_c B_{WA}^d(D^+), \\ \Gamma_{WA}(D_s) &\propto \cos^2 \theta_c B_{WA}^s(D_s) + \sin^2 \theta_c B_{WA}^d(D_s),\end{aligned}$$



Bigi et al.
0911.3322

- By equating the difference between D_s and D^0 rates with the isotriplet component of WA
 - assumes SU(3) violating effects are sub-leading
- Isosinglet component unconstrained

Confronting OPE convergence in charm

- In order to constrain WA fully, need to explicitly compute semileptonic rates and/or distribution moments - compare with exp.

Ligeti et al.
1003.1351

- Perturbative corrections known in the pole scheme

J.F.K.
0909.2755

Gambino & J.F.K.
1004.0114

$$\begin{aligned}
 \Gamma &= \Gamma_0 \left[1 - 0.72 \alpha_s - 0.29 \alpha_s^2 \beta_0 - 0.60 \mu_G^2 - 0.20 \mu_\pi^2 + 0.42 \rho_D^3 + 0.38 \rho_{LS} + 80 \bar{B}_{WA}^{(0)} \right], \\
 \langle E \rangle &= \langle E \rangle_0 \left[1 - 0.03 \alpha_s - 0.03 \alpha_s^2 \beta_0 - 0.07 \mu_G^2 + 0.20 \mu_\pi^2 + 1.4 \rho_D^3 + 0.29 \rho_{LS} + 135 \bar{B}_{WA}^{(1)} \right], \\
 \langle E^2 \rangle &= \langle E^2 \rangle_0 \left[1 - 0.07 \alpha_s - 0.05 \alpha_s^2 \beta_0 - 0.14 \mu_G^2 + 0.52 \mu_\pi^2 + 3.5 \rho_D^3 + 0.66 \rho_{LS} + 204 \bar{B}_{WA}^{(2)} \right], \\
 \sigma_E^2 &= (\sigma_E^2)_0 \left[1 - 0.09 \alpha_s - 0.05 \alpha_s^2 \beta_0 - 0.14 \mu_G^2 + 1.7 \mu_\pi^2 + 9.4 \rho_D^3 + 1.4 \rho_{LS} + 641 \bar{B}_{WA}^{(\sigma)} \right],
 \end{aligned}$$

- Renormalon (Λ/m_c) ambiguity of pole mass
 - all moments affected (n-th scales as m_c^n)

c.f. Antonelli et al.
0907.5386

- Better to use a short distance - threshold mass definition

Convergence of perturbative corrections

- Marginal in the pole scheme ($\alpha_s(m_c) \approx 0.35$)

Ligeti et al.
1003.1351

$$\frac{\Gamma}{\Gamma_0[m_c^{\text{pole}}]} = 1 - 0.269 \epsilon - 0.360 \epsilon_{\text{BLM}}^2 + 0.069 \epsilon^2 + \dots, \quad (\epsilon[=1] \text{ - pert. order counting parameter})$$

- Improves in short distance m_c schemes

$$\frac{\Gamma}{\Gamma_0[m_c^{1S}]} = 1 - 0.133 \epsilon - 0.006 \epsilon_{\text{BLM}}^2 - 0.017 \epsilon^2.$$

- One can try to soften the strong dependence on the charm quark mass using information from inclusive B decays

$(\Delta = m_b - m_c)$

$$\frac{\Gamma}{\Gamma_0[m_b^{1S} - \Delta]} = 1 - 0.075 \epsilon - 0.013 \epsilon_{\text{BLM}}^2 - 0.021 \epsilon^2, \quad (\Delta = m_b - m_c)$$

Convergence of perturbative corrections

- In schemes with explicit IR cut-off, one needs to choose proper (low) IR scale (0.5-0.8 GeV)

Gambino & J.F.K
1004.0114

- Need to translate OPE parameters as well (from global B fits)

using HFAG
winter '09 update

- Perturbative and OPE corrections translated to kinetic scheme

$$\begin{aligned}\Gamma_{kin} &= 1.2(3)10^{-13}\text{GeV} \left\{ 1 + 0.23\alpha_s + 0.18\alpha_s^2\beta_0 - 0.79\mu_G^2 - 0.26\mu_\pi^2 + 1.45\rho_D^3 + 0.56\rho_{LS}^3 + 120\bar{B}_{WA}^{(0)} \right\} \\ \langle E_\ell \rangle_{kin} &= 0.415(21)\text{GeV} \left\{ 1 + 0.03\alpha_s + 0.02\alpha_s^2\beta_0 - 0.09\mu_G^2 + 0.26\mu_\pi^2 + 2.7\rho_D^3 + 0.44\rho_{LS}^3 + 203\bar{B}_{WA}^{(1)} \right\} , \\ \langle E_\ell^2 \rangle_{kin} &= 0.192(20)\text{GeV}^2 \left\{ 1 + 0.001\alpha_s + 0.02\alpha_s^2\beta_0 - 0.18\mu_G^2 + 0.68\mu_\pi^2 + 6.6\rho_D^3 + 0.99\rho_{LS}^3 + 307\bar{B}_{WA}^{(2)} \right\} \\ \sigma_{E,kin}^2 &= 0.019(2)\text{GeV}^2 \left\{ 1 - 0.53\alpha_s - 0.17\alpha_s^2\beta_0 - 0.18\mu_G^2 + 2.2\mu_\pi^2 + 17\rho_D^3 + 2.1\rho_{LS}^3 + 961\bar{B}_{WA}^{(\sigma)} \right\} ,\end{aligned}$$

- Rate uncertainty dominated by m_c & μ_G
- Higher leptonic moments by ρ_D

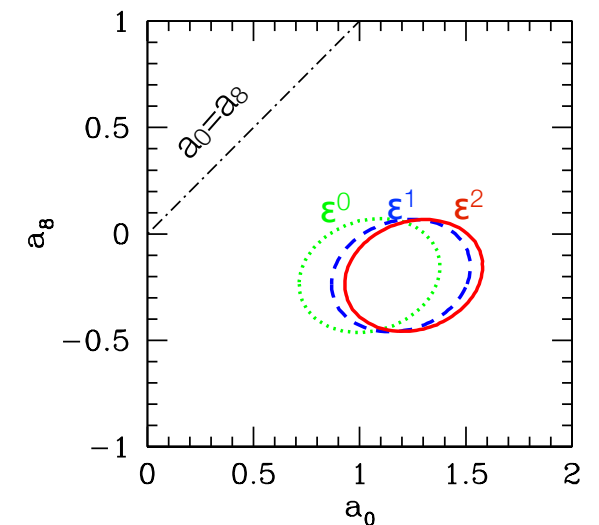
Extraction of WA contributions

- Comparing theoretical expressions with experimental rates (in 1S scheme)
 - using OPE parameters and masses as extracted from global B decay fits
 - neglecting possible SU(3) violations
- Indication of a non-zero isosinglet WA contribution

Ligeti et al.
1003.1351

$$a_0 = 1.25 \pm 0.15,$$

$$a_8 = -0.20 \pm 0.12,$$



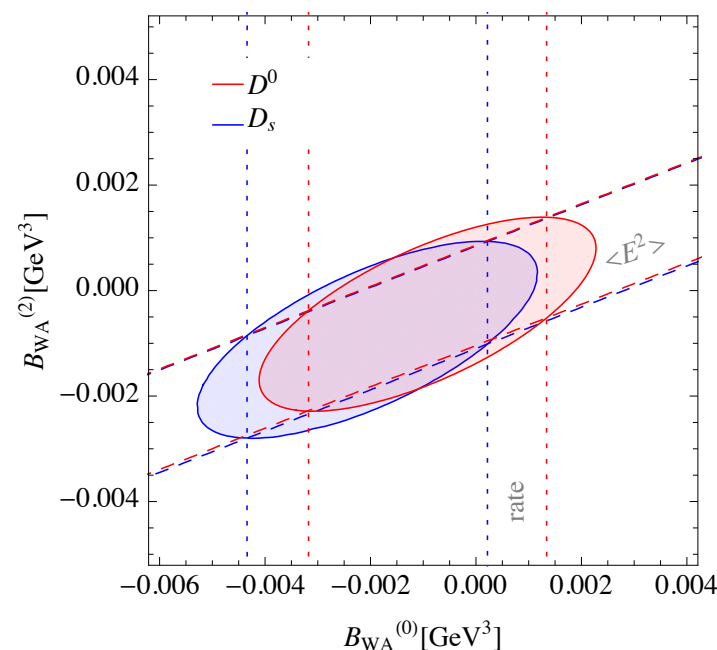
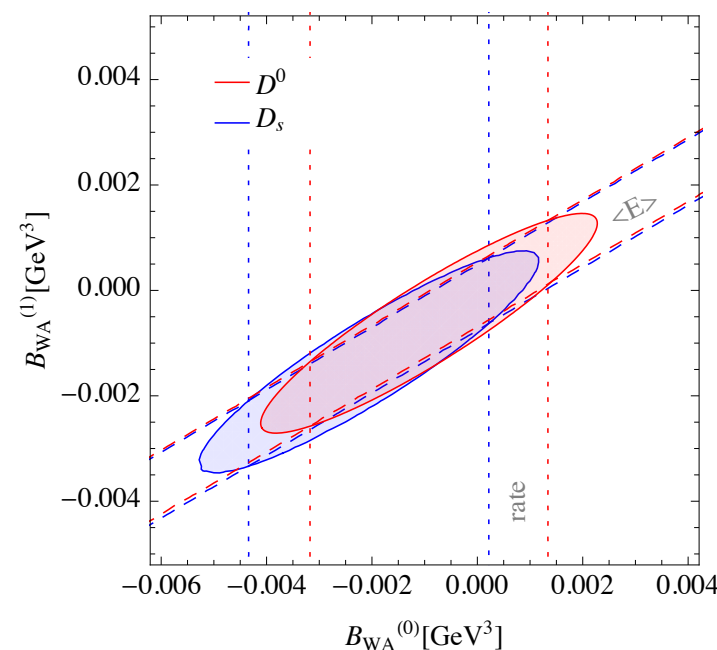
$$a_{0,8} = \frac{m_c^2 m_D f_D^2}{m_c^5} 16\pi^2 (B_2^{s,ns} - B_1^{s,ns}),$$

- Translates into O(1-2%) effect in $B \rightarrow X_u l \nu$ rate

Extraction of WA contributions

- Including information on the leptonic energy moments
- Different dependence of moments on the OPE parameters allows to possibly disentangle SU(3) violating effects from WA contributions
- Introduces dependence due to the modeling of the WA shape in the spectra
- Correlated WA determination from the rate and the moments

Gambino & J.F.K
1004.0114



Extraction of WA contributions

- Including information on the leptonic energy moments
- Allowing for O(20%) SU(3) violation in OPE parameters
- Largest uncertainty due to ρ_D - linear (scale dependent) combination of ρ_D and WA contributions determined precisely
- For $\mu_{WA} \approx 1 \text{ GeV}$ no clear indication of non-zero WA contributions

Gambino & J.F.K
1004.0114

$$B_{WA}^s = -0.0003(25) \text{ GeV}^3$$

- Translates into O(2%) uncertainty in $B \rightarrow X_u \ell \nu$ decay rate

$$B_{WA}^s = -0.0003(25) \text{ GeV}^3$$

Conclusions

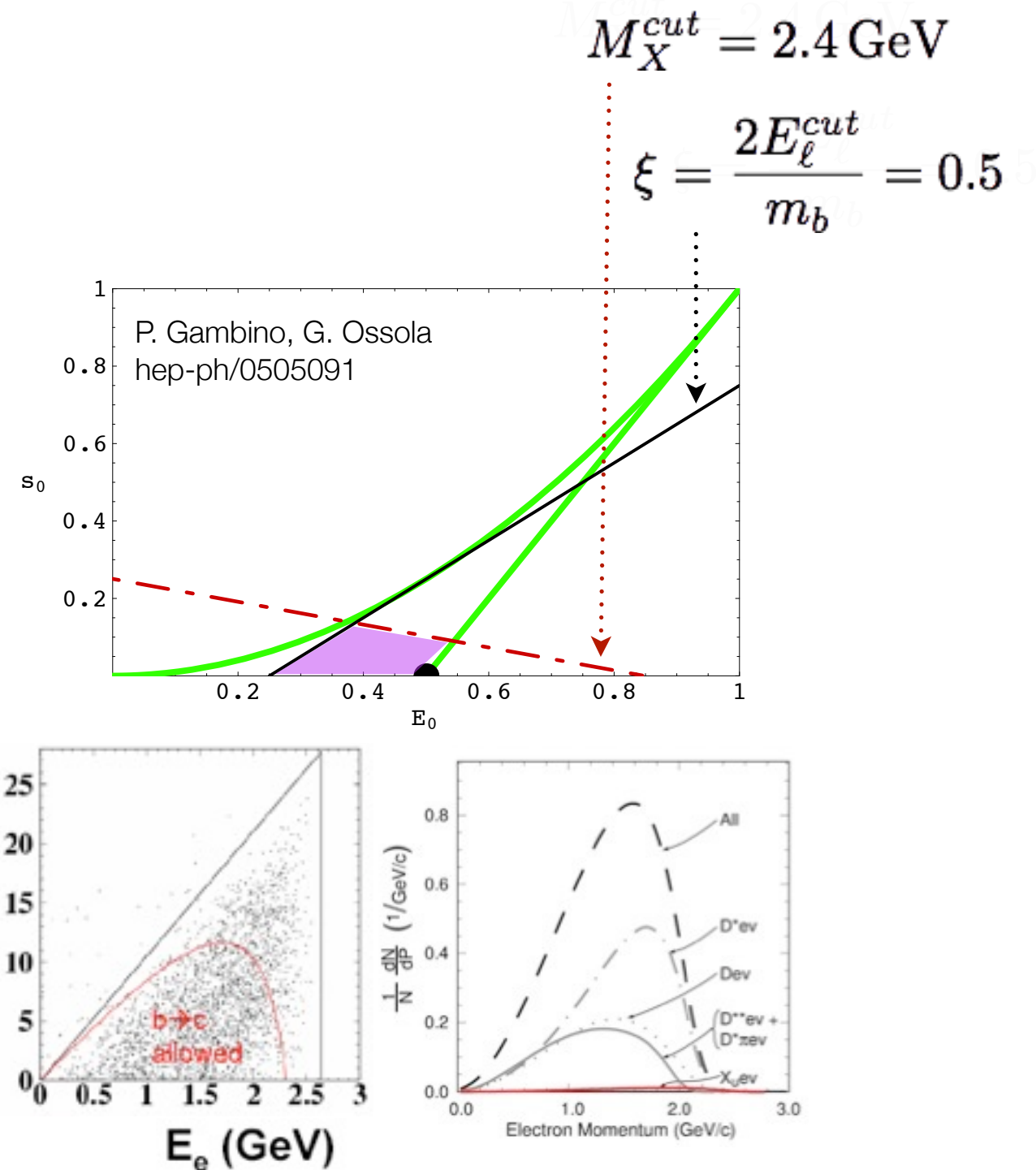
- Inclusive semileptonic charm decays can be used as a laboratory to test the OPE techniques used in the extraction of $|V_{ub}|$ and $|V_{cb}|$ from inclusive B decays
 - perturbative convergence seems to be surprisingly good
- Use several observables to over-constrain the OPE parameter uncertainties and test OPE convergence
- Indications that WA related uncertainties in inclusive $|V_{ub}|$ extraction smaller than previously expected [O(1%)]
- More tests possible in the future with additional experimental inputs (**experimentally determined leptonic energy and hadronic invariant mass moments**) from Cleo and BESIII

Backup Slides

Status of $B \rightarrow X_u \ell \nu$

- Experimental cuts on the leptonic energy and hadronic invariant mass to suppress dominant charm final state contributions
- Introduce theoretical sensitivity to effects beyond the OPE
- Modeled by s.c. shape-functions
- A fully inclusive analysis would carry a tiny **2-3%** theoretical error

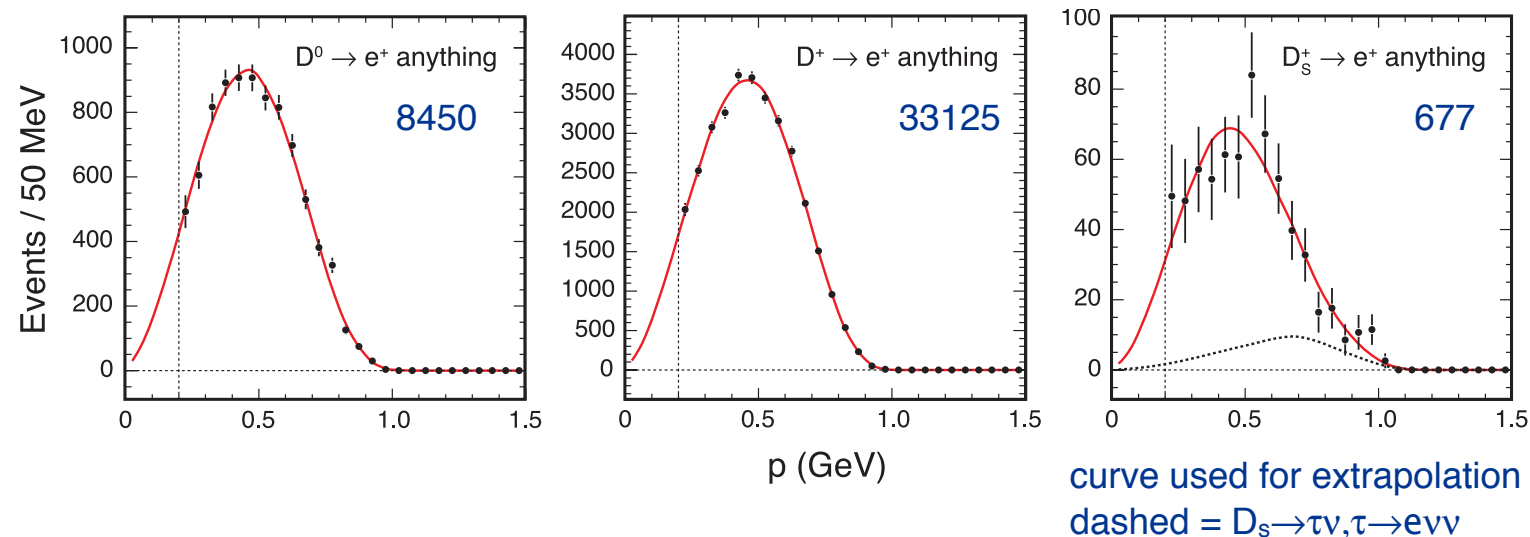
Antonelli et al.
0907.5386



Playing the experimentalist

- One would want to compare completely inclusive leptonic energy moments in the rest-frame of the decaying hadron
- This is not what Cleo presently provide:
 - do not compute the leptonic energy moments
 - spectra given in the lab frame
 - involve a lower $E_e=0.2$ GeV cut
 - do subtract the $D_s \rightarrow \tau \nu$ leptonic background

Asner et al.
[CLEO]
0912.4232



Playing the experimentalist

- One would want to compare completely inclusive leptonic energy moments in the rest-frame of the decaying hadron
- We try to compensate: Gambino & J.F.K
1004.0114
 - extrapolate the spectra down to $E_e=0$ using inclusive model shapes
 - compute the leptonic energy moments from extrapolated spectra (in the lab frame)
 - boost the moments to the D frame by directional averaging
$$\langle E'_e \rangle = \gamma \langle E_e \rangle \quad \langle E'^2_e \rangle = \gamma^2 (1 + \beta^2/3) \langle E_e^2 \rangle$$
 - D's produced in pairs at $E_{CM}=3774\text{MeV}$
 - D_s 's produced associated with D_s^* 's and through their decays

OPE and heavy quark expansion

- Optical theorem: $\Gamma(H_{Q\bar{q}}) = \frac{1}{2m_H} \langle H_{Q\bar{q}} | \mathcal{T} | H_{Q\bar{q}} \rangle$

$$\mathcal{T} = \text{Im } i \int d^4x T\{\mathcal{H}_{eff}(x)\mathcal{H}_{eff}(0)\}$$

- (Global) quark-hadron duality, HQE & OPE

- Equations of motion

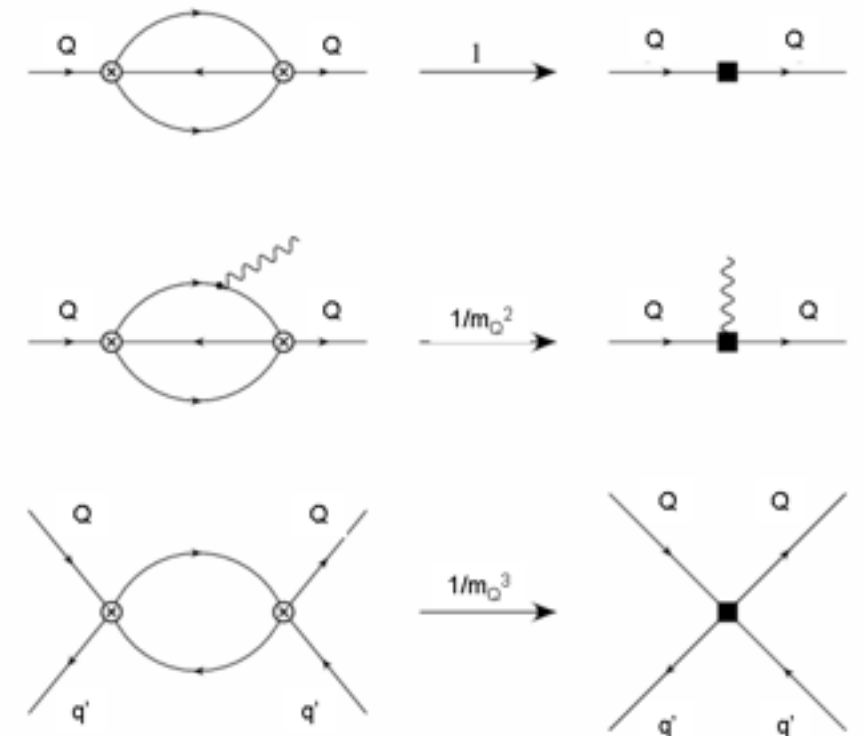
$$\bar{c}c = \bar{c}\not{p}c + \frac{1}{2m_c^2} \left(\bar{c}(iD_\perp)^2 c + \bar{c} \frac{g_s}{2} \sigma \cdot G c \right) + \mathcal{O}(1/m_c^3)$$

- HQE parameters: $\mu_\pi^2 = -\frac{1}{2m_D} \langle D | \bar{c}(iD_\perp)^2 c | D \rangle$
 $\mu_G^2 = \frac{1}{2m_D} \langle D | \bar{c} \frac{g_s}{2} \sigma \cdot B c | D \rangle$

Bigi et al.
[hep-ph/9207214]

Manohar and Wise,
[hep-ph/9308246]

...



- Only applicable for the total rate

OPE and heavy quark expansion

- Analogously define current correlator whose imaginary part gives the hadronic tensor contributing to inclusive semileptonic spectra
- Again use HQE & OPE
- Requires local quark-hadron duality to hold
- Can be softened by instead computing spectral moments
- Any spectral cuts will reintroduce sensitivity to contributions beyond OPE

Bigi et al.
[hep-ph/9207214]

Manohar and Wise,
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