# BESIII上A<sup>+</sup>强子衰变的研究

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# **BEPCII & BESIII**







First HEP collider in China (1988) c.m.s energy: 2 ~ 5 GeV Max luminosity: 1×10<sup>33</sup>cm<sup>-2</sup>s<sup>-1</sup> Non-perturbative  $\tau - charm$  region  $\tau^{\pm} \ D/D_s \ \Lambda_c^{\pm}...$ 

 $J/\psi$ : 2.97 fb<sup>-1</sup>(10B)  $\psi$ (3686): 4.07 fb<sup>-1</sup>(2.7B)  $\psi$ (3770): 20 fb<sup>-1</sup> 4.6~4.95GeV: 6.4 fb<sup>-1</sup>

# $\Lambda_c^+$ : The lightest charmed baryon

- Most of the charmed baryons will eventually decay to Λ<sup>+</sup><sub>c</sub>.
- The Λ<sup>+</sup><sub>c</sub> is one of important tagging hadrons in c-quark counting in the productions at high energy experiment.
- Λ<sup>+</sup><sub>c</sub> may reveal more information of strong- and weak-interactions in charm region, complementary to D/D<sub>s</sub>.



### New data samples in 2020 and 2021

#### **Two major changes in BEPCII machine:**

- max beam energy: 2.30→2.35(2020)→ 2.48 GeV(2021)
- **top-up injection:** data taking efficiency increased by 20~30%



#### Available data for charmed baryons

- ✓ 0.567 fb<sup>-1</sup> at 4.6 GeV (35 days in 2014)
- ✓ 3.9 fb<sup>-1</sup> scan at 4.61, 4.63, 4.64, 4.66, 4.68, 4.7 GeV (186 days in 2020)
- ✓ 1.93 fb<sup>-1</sup> scan at 4.74, 4.75, 4.78, 4.84, 4.92, 4.95 GeV (99 days in 2021)
- 8x  $\Lambda_c$  data that those at 4.6GeV.(~0.77M  $\Lambda_c^+\overline{\Lambda}_c^-$ )
- accessible to  $\Sigma_c / \Xi_c / \Lambda_c^*$  prod. & decays

### Production measurement near threshold

•  $e^+e^- \rightarrow \Lambda_c^+ \overline{\Lambda_c^-}$  cross section are measured at twelve energy points from 4.612-4.951 GeV.

$$\sigma_{\pm} = \frac{N_{\rm ST}^{\pm}}{\varepsilon_{\rm ST}^{\pm} f_{\rm ISR} f_{\rm VP} \mathcal{L}_{\rm int} N_{\rm DT}} \sum_{n=1}^{9} \left( \frac{N_{\rm ST}^{\pm,n} \varepsilon_{\rm DT}^{n}}{\varepsilon_{\rm ST}^{\pm,n}} \right)$$

- Indicate no enhancement around Y(4630) resonance.
   => Conflict with Belle.
- $|G_E/G_M|$  ratio are derived by fitting to angular distribution.
- The oscillations on  $|G_E/G_M|$  ratio is significantly observed with higher frequency than of the proton.



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# Studies on the $\Lambda_c^+$ hadronic measurements at BESIII using data 20/21

- Two-body decays
  - $\begin{array}{c} \Box \quad \Lambda_c^+ \rightarrow n\pi^+ \checkmark \\ \Box \quad \Lambda_c^+ \rightarrow \Sigma^0 K^+, \Sigma^+ K_S^0 \\ \Box \quad \Lambda_c^+ \rightarrow p\pi^0 \\ \Box \quad \Lambda_c^+ \rightarrow p\eta, p\omega \\ \Box \quad \Lambda_c^+ \rightarrow \Lambda K^+ \\ \Box \quad \Lambda_c^+ \rightarrow p\eta' \\ \Box \quad \Lambda_c^+ \rightarrow \Xi^0 K^+ \checkmark \end{array}$
- $\begin{array}{l} & \text{Multi-body decays} \\ & \square \ \Lambda_c^+ \rightarrow nK_S^0\pi^+\pi^0 \\ & \square \ \Lambda_c^+ \rightarrow nK_S^0\pi^+, nK_S^0K^+ \\ & \square \ \overline{\Lambda_c^-} \rightarrow \overline{n}X \\ & \square \ \Lambda_c^+ \rightarrow \Lambda K^+\pi^0, \Lambda K^+\pi^+\pi^- \\ & \square \ \Lambda_c^+ \rightarrow \Xi^0K^+\pi^0 \\ & \square \ \Lambda_c^+ \rightarrow \Sigma^-K^+\pi^+ \\ & \square \ \Lambda_c^+ \rightarrow \Sigma^+K^+K^-, \Sigma^+\phi, \Sigma^+K^+\pi^-(\pi^0) \\ & \square \ \Lambda_c^+ \rightarrow n\pi^+\pi^0, n\pi^+\pi^-\pi^+, nK^-\pi^+\pi^+ \\ & \square \ \Lambda_c^+ \rightarrow \Lambda\pi^+\pi^0 \end{array}$

- : PRL 128.142001 (2022).
- : PRD 106.052003 (2022).
- : arXiv 2311.06883.
- : JHEP 11.137 (2023).
- : PRD 106.L111101 (2022).
- : PRD 106.072002 (2022).
- : PRL 132.031801 (2024).
- : PRD 109.053005 (2024). : PRD 109.072010 (2024). : PRD 108.L031101 (2023). : PRD 109.032003 (2024). : PRD 109.052001 (2024). : PRD 109.L071103 (2024). : JHEP 09.125 (2023). : CPC 47.023001 (2023).
- : JHEP 12.033 (2022).

## First observation of $\Lambda_c^+ \rightarrow n\pi^+$

#### PRL 128.142001 (2022)



• First singly Cabibbo-suppressed  $\Lambda_c^+$  decay involved neutron was observed (7.3 $\sigma$ ).

• Absolute BF is measured to be  $\mathcal{B}(\Lambda_c^+ \to n\pi^+) = (6.6 \pm 1.2_{stat.} \pm 0.4_{syst.}) \times 10^{-4}$ . =>Consistent with SU(3) flavor symmetry prediction. [PLB790,225 (2019)]

=>Twice larger than the dynamical calculation based on Pole model and CA. [PRD97,074028 (2018)]

•  $\mathcal{B}(\Lambda_c^+ \to \Lambda \pi^+) = (1.31 \pm 0.08_{stat.} \pm 0.05_{syst.}) \times 10^{-2}$ . => Consistent with previous BESIII results

• 
$$\mathcal{B}(\Lambda_c^+ \to \Sigma^0 \pi^+) = (1.22 \pm 0.08_{stat.} \pm 0.07_{syst.}) \times 10^{-2}$$
. => Consistent with previous BESIII results

•  $R = \frac{\mathcal{B}(\Lambda_c^+ \to n\pi^+)}{\mathcal{B}(\Lambda_c^+ \to p\pi^0)} > 7.2 @90\% C.L. (\mathcal{B}(\Lambda_c^+ \to p\pi^0) < 8.0 \times 10^{-5} @90\% C.L.$  From Belle) =>Disagrees with SU(3) flavor symmetry and dynamical calculation (2.0-4.7) while in consistent with SU(3) plus topological-diagram approach (9.6).

### Decay Asymmetry of $\Lambda_c^+ \to \Xi^0 K^+$

- $\Lambda_c^+ \to \Xi^0 K^+$  is pure W-exchange process which have significant contributions in charmed baryon decay.  $\Lambda_c^+$
- Nonfactorizable W-exchange diagram cannot be calculated using theoretical approaches.
- Long-standing puzzle on how large the S-wave amplitude.

PRL 132.031801	(2024)
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FIG. 1. Feynman diagrams for  $\Lambda_c^+ \to \Xi^0 K^+$ .

Theory or experiment	$\mathcal{B}(\Lambda_c^+ \to \Xi^0 K^+) \; (\times 10^{-3})$	$lpha_{\Xi^0K^+}$	$ A  \; (\times 10^{-2} G_F \; \mathrm{GeV}^2)$	$ B  \; (\times 10^{-2} G_F \; \mathrm{GeV}^2)$	$\delta_p - \delta_s \text{ (rad)}$
Körner (1992), CCQM [7]	2.6	0			
Xu (1992), Pole [8]	1.0	0	0	7.94	
Źencaykowski (1994), Pole [9]	3.6	0			
Ivanov (1998), CCQM [10]	3.1	0			
Sharma (1999), CA [11]	1.3	0			
Geng (2019), SU(3) [12]	$5.7\pm0.9$	$0.94\substack{+0.06\\-0.11}$	$2.7\pm0.6$	$16.1\pm2.6$	
Zou (2020), CA [6]	7.1	0.90	4.48	12.10	
Zhong (2022), $SU(3)^a$ [13]	$3.8^{+0.4}_{-0.5}$	$0.91\substack{+0.03\\-0.04}$	$3.2\pm0.2$	$8.7^{+0.6}_{-0.8}$	
Zhong (2022), $SU(3)^b$ [13]	$5.0^{+0.6}_{-0.9}$	$0.99\pm0.01$	$3.3^{+0.5}_{-0.7}$	$12.3^{+1.2}_{-1.8}$	
BESIII (2018) [14]	$5.90 \pm 0.86 \pm 0.39$		• • •		
PDG fit (2022) [2]	$5.5\pm0.7$		••••		

• Experimental measurement of decay asymmetry is crucial and urgent.

### Decay Asymmetry of $\Lambda_c^+ \to \Xi^0 K^+$

$\alpha_{BP}$ =	$=\frac{2\operatorname{Re}(s^*p)}{ s ^2+ p ^2}\qquad\beta_{BP}=$	$= \frac{2 \text{Im}(s^* p)}{ s ^2 +  p ^2}$	$\gamma_{BP} = \frac{ s ^2 -  p ^2}{ s ^2 +  p ^2}$
	$\beta_{BP} = \chi$	$\sqrt{1-\alpha_{BP}^2}\sin\Delta_{BP}$	$\gamma_{BP} = \sqrt{1 - \alpha_{BP}^2 \cos\Delta_{BP}}$
Level	Decay	Helicity angle	Helicity amplitude
0	$e^+e^- \to \Lambda_c^+(\lambda_1) \bar{\Lambda}_c^-(\lambda_2)$	$( heta_0)$	$\mathcal{A}_{\lambda_1,\lambda_2}$
1	$\Lambda_c^+ \to \Xi^0(\lambda_3)  K^+$	$_{( heta_1,\phi_1)}$	$\mathcal{B}_{\lambda_3}$
2	$\Xi^0  o \Lambda(\lambda_4)  \pi^0$	$( heta_2,\!\phi_2)$	$\mathcal{C}_{\lambda_4}$
3	$\Lambda  o p(\lambda_5) \pi^-$	$_{( heta_3,\phi_3)}$	$\mathcal{D}_{\lambda_5}$

 $\Lambda \to p(\lambda_5) \pi^ (\theta_3,\phi_3)$ 

- $d\cos\theta_0 \ d\cos\theta_1 \ d\cos\theta_2 \ d\cos\theta_3 \ d\phi_1 \ d\phi_2 \ d\phi_3$
- $\propto 1 + \alpha_0 \cos^2 \theta_0$

- +  $(1 + \alpha_0 \cos^2 \theta_0) \alpha_{\equiv^0 K^+} \alpha_{\Lambda \pi^0} \cos \theta_2$
- +  $(1 + \alpha_0 \cos^2 \theta_0) \alpha_{\Xi^0 K^+} \alpha_{p\pi^-} \cos \theta_2 \cos \theta_3$
- +  $(1 + \alpha_0 \cos^2 \theta_0) \alpha_{\Lambda \pi^0} \alpha_{n\pi^-} \cos \theta_3$
- $-(1+\alpha_0\cos^2\theta_0) \alpha_{\Xi^0K^+} \sqrt{1-\alpha_{\Lambda\pi^0}^2} \alpha_{p\pi^-}\sin\theta_2\sin\theta_3\cos(\Delta_{\Lambda\pi^0}+\phi_3)$
- $+\sqrt{1-\alpha_0^2}\sin\Delta_0\sin\theta_0\cos\theta_0\alpha_{\Xi^0K^+}\sin\theta_1\sin\phi_1$
- +  $\sqrt{1 \alpha_0^2 \sin \Delta_0 \sin \theta_0 \cos \theta_0 \alpha_{\Lambda \pi^0} \sin \theta_1 \sin \phi_1 \cos \theta_2}$
- $-\alpha_0^2 \sin \Delta_0 \sin \theta_0 \cos \theta_0 \alpha_{\Xi^0 K^+} \alpha_{\Lambda \pi^0} \alpha_{p \pi^-} \sin \theta_1 \sin \phi_1 \cos \theta_3$
- $+\sqrt{1-\alpha_0^2}\sin\Delta_0\sin\theta_0\cos\theta_0\alpha_{p\pi}-\sin\theta_1\sin\phi_1\cos\theta_2\cos\theta_3$
- $-\sqrt{1-\alpha_0^2}\sin\Delta_0\sin\theta_0\cos\theta_0\sqrt{1-\alpha_{\Lambda\pi^0}^2}\alpha_{p\pi^-}\sin\theta_1\sin\phi_1\sin\theta_2\sin\theta_3\cos(\Delta_{\Lambda\pi^0}+\phi_3)$  $+\sqrt{1-\alpha_0^2}\sin\Delta_0\sin\theta_0\cos\theta_0\sqrt{1-\alpha_{\pm 0\,K^+}^2}\alpha_{\Lambda\pi^0}\cos\phi_1\sin\theta_2\sin(\Delta_{\pm 0\,K^+}+\phi_2)$
- $+\sqrt{1-\alpha_0^2}\sin\Delta_0\sin\theta_0\cos\theta_0\sqrt{1-\alpha_{\pm 0\,K^+}^2}\alpha_{\Lambda\pi^0}\cos\theta_1\sin\phi_1\sin\theta_2\cos(\Delta_{\pm 0\,K^+}+\phi_2)$
- $+\sqrt{1-\alpha_0^2}\sin\Delta_0\sin\theta_0\cos\theta_0\sqrt{1-\alpha_{\Xi^0K^+}^2}\alpha_{p\pi^-}\cos\theta_1\sin\phi_1\sin\theta_2\cos(\Delta_{\Xi^0K^+}+\phi_2)\cos\theta_3$
- $+\sqrt{1-\alpha_0^2}\sin\Delta_0\sin\theta_0\cos\theta_0\sqrt{1-\alpha_{\Xi^0K^+}^2}\alpha_{p\pi^-}\cos\phi_1\sin\theta_2\sin(\Delta_{\Xi^0K^+}+\phi_2)\cos\theta_3$
- $-\sqrt{1-\alpha_0^2}\,\sin\Delta_0\sin\theta_0\cos\theta_0\sqrt{1-\alpha_{\Xi^0\,K^+}^2}\sqrt{1-\alpha_{\Lambda\pi^0}^2}\,\alpha_{p\pi^-}\cos\theta_1\sin\phi_1\sin(\Delta_{\Xi^0\,K^+}+\phi_2)\sin\theta_3\sin(\Delta_{\Lambda\pi^0}+\phi_3)$
- $+\sqrt{1-\alpha_0^2}\sin\Delta_0\sin\theta_0\cos\theta_0\sqrt{1-\alpha_{\Xi^0\,K^+}^2}\sqrt{1-\alpha_{\Lambda\pi^0}^2}\alpha_{p\pi^-}\cos\theta_1\sin\phi_1\cos\theta_2\cos(\Delta_{\Xi^0\,K^+}+\phi_2)\sin\theta_3\cos(\Delta_{\Lambda\pi^0}+\phi_3)$  $+\sqrt{1-\alpha_0^2}\,\sin\Delta_0\sin\theta_0\cos\theta_0\sqrt{1-\alpha_{\Xi^0\,K^+}^2}\sqrt{1-\alpha_{\Lambda\pi^0}^2}\,\alpha_{p\pi^-}\cos\phi_1\cos(\Delta_{\Xi^0\,K^+}+\phi_2)\sin\theta_3\sin(\Delta_{\Lambda\pi^0}+\phi_3)$
- $+\sqrt{1-\alpha_0^2}\sin\Delta_0\sin\theta_0\cos\theta_0\sqrt{1-\alpha_{\Xi^0K^+}^2}\sqrt{1-\alpha_{\Lambda\pi^0}^2}\alpha_{p\pi^-}\cos\phi_1\cos\theta_2\sin(\Delta_{\Xi^0K^+}+\phi_2)\sin\theta_3\cos(\Delta_{\Lambda\pi^0}+\phi_3)$



The joint angular distribution for  $\Lambda_c^+ \to \Xi^0 K^+$  is derived based on helicity amplitude.

#### 2024/5/11

#### BESIII上A<sup>+</sup>强子衰变的研究

## Decay Asymmetry of $\Lambda_c^+ \to \Xi^0 K^+$

#### PRL 132.031801 (2024)



- From the fit, we obtain  $\alpha_{\Xi^0K^+} = 0.01 \pm 0.16_{stat.} \pm 0.03_{syst.},$  $\Delta_{\Xi^0K^+} = 3.84 \pm 0.90_{stat.} \pm 0.17_{syst.}.$
- α<sub>Ξ<sup>0</sup>K<sup>+</sup></sub> is in good agreement with zero.
   => strong identification for theoretical predictions.

$$\begin{split} \Gamma_{\Xi^{0}K^{+}} &= \frac{\mathcal{B}(\Lambda_{c}^{+} \to \Xi^{0}K^{+})}{\tau_{\Lambda_{c}^{+}}} = \frac{|\vec{p}_{c}|}{8\pi} \Big[ \frac{(m_{\Lambda_{c}^{+}} + m_{\Xi^{0}})^{2} - m_{K^{+}}^{2}}{m_{\Lambda_{c}^{+}}^{2}} |A|^{2} + \frac{(m_{\Lambda_{c}^{+}} - m_{\Xi^{0}})^{2} - m_{K^{+}}^{2}}{m_{\Lambda_{c}^{+}}^{2}} |B|^{2} \Big] \\ \alpha_{\Xi^{0}K^{+}} &= \frac{2\kappa |A| |B| \cos(\delta_{p} - \delta_{s})}{|A|^{2} + \kappa^{2} |B|^{2}} \\ \Delta_{\Xi^{0}K^{+}} &= \arctan \frac{2\kappa |A| |B| \sin(\delta_{p} - \delta_{s})}{|A|^{2} - \kappa^{2} |B|^{2}} \end{split}$$

- Especially, cos(δ<sub>p</sub> δ<sub>s</sub>) is measured to close to zero. => Not considered in previous literature.
- Fills the long-standing puzzle on how to model  $\alpha_{\Xi^0K^+}$  and  $\mathcal{B}(\Lambda_c^+ \to \Xi^0K^+)$  simultaneously.
- After considered the phase shift, some calculations
  - Geng (2023), SU(3)  $-0.15 \pm 0.14$  : PRD 109.L071302 (2024).
    - Zhong (2024), TDA  $-0.16 \pm 0.13$  : arXiv 2401.15926.
  - ✓ Zhong (2024), IRA
- $-0.19 \pm 0.12$  : arXiv 2401.15926.

# BF measurement of $\Lambda_c^+ \to \Xi^0 K^+ \pi^0$

#### PRD 109.052001 (2024)



✓  $\mathcal{B}(\Lambda_c^+ \to \Xi(1530)^0 K^+) \Rightarrow$  Consistent with previous BESIII results.

✓  $\mathcal{B}(\Lambda_c^+ \to \Xi^0 K^+ \pi^0)$  => Lower than prediction based on SU(3) symmetry.

### Observation of SCS decay $\Lambda_c^+ \rightarrow \Sigma^- K^+ \pi^+$

#### PRD 109.L071103 (2024)



- Singly Cabibbo-suppressed decay  $\Lambda_c^+ \to \Sigma^- K^+ \pi^+$  was observed for the first time (5.4 $\sigma$ ).
- Absolute BF is measured to be  $\mathcal{B}(\Lambda_c^+ \to \Sigma^- K^+ \pi^+) = (3.8 \pm 1.2_{stat.} \pm 0.2_{syst.}) \times 10^{-4}$ . =>Consistent with SU(3) flavor symmetry prediction  $(3.3 \pm 2.3) \times 10^{-4}$ . [PRD99, 073003 (2019)]
- $\mathcal{B}(\Lambda_c^+ \to \Xi^- K^+ \pi^+) = (7.74 \pm 0.76_{stat.} \pm 0.54_{syst.}) \times 10^{-3}$ . => Consistent with PDG Fit
- $\mathcal{B}(\Lambda_c^+ \to \Xi(1530)^0 K^+) = (5.03 \pm 0.77_{stat.} \pm 0.20_{syst.}) \times 10^{-3}$ . => Consistent with previous BESIII results [PLB783, 200-206 (2018)]

• 
$$\frac{\mathcal{B}(\Lambda_c^+ \to \Sigma^- K^+ \pi^+)}{\mathcal{B}(\Lambda_c^+ \to \Sigma^- \pi^+ \pi^+)} = (2.03 \pm 0.73)\% \simeq (0.4 \pm 0.1)s_c^2 \ (s_c^2 \equiv sin\theta_c = 0.2248)$$
$$=> \text{Close to } \frac{\mathcal{B}(\Xi_c^0 \to \Xi^- K^+)}{\mathcal{B}(\Xi_c^0 \to \Xi^- \pi^+)} \text{ and deviates significantly from } 1.0s_c^2.$$

=>Ratio of isospin partner modes  $\frac{\mathcal{B}(\Lambda_c^+ \to \Sigma^+ K^+ \pi^-)}{\mathcal{B}(\Lambda_c^+ \to \Sigma^+ \pi^+ \pi^-)}$  consistent with  $1.0s_c^2$ .

#### Observation of multi-body SCS/CF decay involved neutron

#### CPC 47.023001 (2023)



• Absolute BF is measured to be

$$\begin{aligned} \mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^0) &= (0.64 \pm 0.09 \pm 0.02)\%, \\ \mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^-\pi^+) &= (0.45 \pm 0.07 \pm 0.03)\%, \\ \mathcal{B}(\Lambda_c^+ \to nK^-\pi^+\pi^+) &= (1.90 \pm 0.08 \pm 0.09)\%. \end{aligned}$$

- $\frac{\mathcal{B}(\Lambda_c^+ \to p\pi^-\pi^+)}{\mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^0)} = 0.72 \pm 0.11$  Provides key SU(3) symmetry constraint
- $\frac{\mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^0)}{\mathcal{B}(\Lambda_c^+ \to n\pi^+)} = 9.2 \pm 2.4$  Rich intermediate resonance states
- $\frac{\mathcal{B}(\Lambda_c^+ \to n\pi^+\pi^-\pi^+)}{\mathcal{B}(\Lambda_c^+ \to nK^-\pi^+\pi^+)} = 0.24 \pm 0.04$  Consistent with  $1.0s_c^2$

PWA for  $\Lambda_c^+ \to \Lambda \pi^+ \pi^0$ 

#### JHEP 12.033 (2022)



- About 10K events survived which purity is larger than 80%.
- PWA based on helicity amplitude is performed.
- Interference mostly exist  $\Lambda \rho(770)$  and  $\Sigma(1385)^{0/+}\pi^{+/0}$ .

# PWA for $\Lambda_c^+ \to \Lambda \pi^+ \pi^0$

$\frac{1}{2}^{+}(\Lambda_{c}^{+}) \to \frac{3}{2}^{+}(\Sigma(1385)^{+}) + 0^{-}(\pi^{0})$			$\frac{1}{2}^{+}(\Lambda_{c}^{+})$ -	$\rightarrow \frac{3}{2}^+ (\Sigma(1385))$	$)^{0}) + 0^{-}(\pi^{+})$	
Amplitude	Magnitude	Phase $\phi$ (rad)	Amplitude	Magnitude	Phase $\phi$ (rad)	
$g_{1,\frac{3}{2}}^{\Sigma(1385)^+}$	1.0  (fixed)	0.0 (fixed)	$g_{1,\frac{3}{2}}^{\Sigma(1385)^0}$	1.0 (fixed)	0.0 (fixed)	
$g_{2,\frac{3}{2}}^{\Sigma(1385)^+}$	$1.29\pm0.25$	$2.82\pm0.18$	$g_{2,\frac{3}{2}}^{\Sigma(1385)^0}$	$1.70\pm0.38$	$2.70\pm0.22$	$\alpha_{I}$
$\frac{1}{2}^+(\Lambda_c^+$	$T^{-}) \rightarrow \frac{3}{2}^{-}(\Sigma(1670)^{+})$	$) + 0^{-}(\pi^{0})$	$\frac{1}{2}^{+}(\Lambda_{c}^{+})$ -	$\rightarrow \frac{3}{2}^{-}(\Sigma(1670))$	$)^{0}) + 0^{-}(\pi^{+})$	
Amplitude	Magnitude	Phase $\phi$ (rad)	Amplitude	Magnitude	Phase $\phi$ (rad)	
$g_{1,\frac{3}{2}}^{\Sigma(1670)^+}$	1.0 (fixed)	0.0 (fixed)	$g_{1,\frac{3}{2}}^{\Sigma(1670)^0}$	1.0 (fixed)	0.0 (fixed)	
$g_{2,\frac{3}{2}}^{\Sigma(1670)^+}$	$1.39\pm0.42$	$0.85\pm0.26$	$g_{2,\frac{3}{2}}^{\Sigma(1670)^0}$	$0.74 \pm 0.18$	$0.29 \pm 0.24$	
$\frac{1}{2}^{+}(\Lambda_{c}^{+}) \rightarrow \frac{1}{2}^{-}(\Sigma(1750)^{+}) + 0^{-}(\pi^{0})$			$\frac{1}{2}^+(\Lambda_c^+)$ -	$\rightarrow \frac{1}{2}^{-}(\Sigma(1750))$	$)^{0}) + 0^{-}(\pi^{+})$	a
Amplitude	Magnitude	Phase $\phi$ (rad)	Amplitude	Magnitude	Phase $\phi$ (rad)	
$g_{0,\frac{1}{2}}^{\Sigma(1750)^+}$	1.0 (fixed)	0.0 (fixed)	$g_{0,\frac{1}{2}}^{\Sigma(1750)^0}$	1.0 (fixed)	0.0 (fixed)	
$g_{1,\frac{1}{2}}^{\Sigma(1750)^+}$	$0.45\pm0.10$	$-2.28\pm0.22$	$g_{1,\frac{1}{2}}^{\Sigma(1750)^0}$	$0.38\pm0.10$	$-2.03\pm0.20$	
$\frac{1}{2}^{+}(\Lambda_{c}^{+}) \rightarrow \frac{1}{2}^{+}(\Lambda) + 1^{-}(\rho(770)^{+})$			$\frac{1}{2}^+(\Lambda_c^+)$	$) \rightarrow \frac{1}{2}^+(\Lambda) +$	$1^{-}(NR_{1^{-}})$	
Amplitude	Magnitude	Phase $\phi$ (rad)	Amplitude	Magnitude	Phase $\phi$ (rad)	
$g^{\rho}_{0,\frac{1}{2}}$	1.0 (fixed)	0.0 (fixed)	$g^{N\!R}_{0,\frac{1}{2}}$	1.0 (fixed)	0.0 (fixed)	
$g^{\rho}_{1,\frac{1}{2}}$	$0.48\pm0.12$	$-1.69\pm0.12$	$g_{1,\frac{1}{2}}^{N\!R}$	$0.94\pm0.12$	$-0.49\pm0.16$	
$g^{\rho}_{1,\frac{3}{2}}$	$0.90\pm0.10$	$0.48\pm0.13$	$g_{1,\frac{3}{2}}^{N\!R}$	$0.21\pm0.09$	$-2.84\pm0.53$	
$g^{\rho}_{2,\frac{3}{2}}$	$0.55\pm0.08$	$-0.04\pm0.18$	$g^{N\!R}_{2,\frac{3}{2}}$	$0.33\pm0.14$	$-1.92\pm0.30$	
1/2	$^{+}(\Lambda) \to \frac{1}{2}^{+}(p) + 0^{-}$	$(\pi^{-})$				
Amplitude	Magnitude	Phase $\phi$ (rad)				
$g^{\Lambda}_{0,\frac{1}{2}}$	1.0 (fixed)	0.0 (fixed)				
$g_{1,\frac{1}{2}}^{\Lambda}$	0.435376 (fixed)	0.0 (fixed)				

#### JHEP 12.033 (2022)

$$\begin{split} {}_{\Lambda\rho(770)^{+}} &= \frac{|H^{\rho}_{\frac{1}{2},1}|^{2} - |H^{\rho}_{-\frac{1}{2},-1}|^{2} + |H^{\rho}_{\frac{1}{2},0}|^{2} - |H^{\rho}_{-\frac{1}{2},0}|^{2}}{|H^{\rho}_{\frac{1}{2},1}|^{2} + |H^{\rho}_{-\frac{1}{2},-1}|^{2} + |H^{\rho}_{\frac{1}{2},0}|^{2} + |H^{\rho}_{-\frac{1}{2},0}|^{2}} \\ &= \frac{\sqrt{\frac{1}{9}} \cdot 2 \cdot \Re\left(g^{\rho}_{0,\frac{1}{2}} \cdot \bar{g}^{\rho}_{1,\frac{1}{2}} - g^{\rho}_{1,\frac{3}{2}} \cdot \bar{g}^{\rho}_{2,\frac{3}{2}}\right) - \sqrt{\frac{8}{9}} \cdot 2 \cdot \Re\left(g^{\rho}_{0,\frac{1}{2}} \cdot \bar{g}^{\rho}_{1,\frac{3}{2}} + g^{\rho}_{1,\frac{1}{2}} \cdot \bar{g}^{\rho}_{2,\frac{3}{2}}\right)}{|g^{\rho}_{0,\frac{1}{2}}|^{2} + |g^{\rho}_{1,\frac{1}{2}}|^{2} + |g^{\rho}_{1,\frac{3}{2}}|^{2} + |g^{\rho}_{2,\frac{3}{2}}|^{2}} \end{split}$$

$$\alpha_{\Sigma(1385)\pi} = \frac{|H_{0,\frac{1}{2}}^{\Sigma(1385)}|^2 - |H_{0,-\frac{1}{2}}^{\Sigma(1385)}|^2}{|H_{0,\frac{1}{2}}^{\Sigma(1385)}|^2 + |H_{0,-\frac{1}{2}}^{\Sigma(1385)}|^2} = \frac{2\Re\left(g_{1,\frac{3}{2}}^{\Sigma(1385)} \cdot \bar{g}_{2,\frac{3}{2}}^{\Sigma(1385)}\right)}{|g_{1,\frac{3}{2}}^{\Sigma(1385)}|^2 + |g_{2,\frac{3}{2}}^{\Sigma(1385)}|^2}.$$

 Decay asymmetry parameters can be obtained by the fit results of the partial wave amplitudes.

# PWA for $\Lambda_c^+ \to \Lambda \pi^+ \pi^0$

**Table 9.** The comparison among this work, various theoretical calculations and PDG results. Here,the uncertainties of this work are the combined uncertainties. "—" means unavailable.

	Theoretical c	This work	PDG	
$10^2 \times \mathcal{B}(\Lambda_c^+ \to \Lambda \rho(770)^+)$	$4.81 \pm 0.58$ [13]	$4.0 \ [14, \ 15]$	$4.06\pm0.52$	< 6
$10^3 \times \mathcal{B}(\Lambda_c^+ \to \Sigma(1385)^+ \pi^0)$	$2.8 \pm 0.4$ [16]	$2.2 \pm 0.4$ [17]	$5.86 \pm 0.80$	
$10^3 \times \mathcal{B}(\Lambda_c^+ \to \Sigma(1385)^0 \pi^+)$	$2.8 \pm 0.4$ [16]	$2.2 \pm 0.4$ [17]	$6.47\pm0.96$	
$lpha_{\Lambda ho(770)^+}$	$-0.27 \pm 0.04$ [13]	-0.32 [14, 15]	$-0.763 \pm 0.070$	
$lpha_{\Sigma(1385)^+\pi^0}$	$-0.91^{+0.4}_{-0.1}$	$-0.917 \pm 0.089$		
$lpha_{\Sigma(1385)^0\pi^+}$	$-0.91^{+0.4}_{-0.1}$	$-0.79\pm0.11$		

- No theoretical models is able to explain both BFs and decay asymmetries simultaneously.
- Fruitful results are extracted which provide crucial input to extend the understanding of dynamics of charmed baryon hadronic decays.

### Some new results



0.54



FIG. 3. Combined simultaneous fit results to the distributions of  $M_{\rm BC}$  for (a)  $\Lambda_c^+ \rightarrow \Lambda K^+ \pi^0$  and (b)  $\Lambda_c^+ \rightarrow \Lambda K^+ \pi^- \pi^-$  at 13 energy points.

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FIG. 4. Combined simultaneous fit results to the distributions of  $M_{\rm BC}$  for (a)  $\Lambda_c^+ \to \Lambda \pi^+ \pi^0$  and (b)  $\Lambda_c^+ \to \Lambda \pi^+ \pi^- \pi^-$  at 13 energy points.

✓  $\mathcal{B}(\Lambda_c^+ \to nK_S^0 \pi^+ \pi^0) = (0.85 \pm 0.13 \pm 0.03)\%$ ✓ Differ from theoretical prediction based on isospin by 4.4*σ*.

$$\checkmark \frac{\mathcal{B}(\Lambda_{c}^{+} \to \Lambda K^{+} \pi^{0})}{\mathcal{B}(\Lambda_{c}^{+} \to \Lambda \pi^{+} \pi^{0})} = (2.09 \pm 0.39 \pm 0.07) \times 10^{-2}$$

$$\checkmark \frac{\mathcal{B}(\Lambda_{c}^{+} \to \Lambda K^{+} \pi^{+} \pi^{-})}{\mathcal{B}(\Lambda_{c}^{+} \to \Lambda \pi^{+} \pi^{+} \pi^{-})} = (1.13 \pm 0.41 \pm 0.06) \times 10^{-2}$$

$$\checkmark \mathcal{B}(\Lambda_{c}^{+} \to \Lambda K^{+} \pi^{0}) = (1.49 \pm 0.27 \pm 0.05 \pm 0.08_{\text{ref.}}) \times 10^{-3}$$

$$= > \text{Lower than prediction based on SU(3)}$$

$$\checkmark \mathcal{B}(\Lambda_{c}^{+} \to \Lambda K^{+} \pi^{+} \pi^{-}) = (4.13 \pm 1.48 \pm 0.20 \pm 0.33_{\text{ref.}}) \times 10^{-4}$$

$$= > \text{Consistent with BaBar experiment}$$

## Summary

- > BESIII have collected the largest data samples with 6.4fb<sup>-1</sup> integrated luminosity from 4.60 to 4.95 GeV near the  $\Lambda_c^+ \overline{\Lambda}_c^-$  production threshold.
- > Many singly Cabibbo-suppressed  $\Lambda_c^+$  decay involved neutron were observed for the first time.
- ➤ The polarization of pure W-exchange process Λ<sup>+</sup><sub>c</sub> → Ξ<sup>0</sup>K<sup>+</sup> was measured for the first time, which fills the long-standing puzzle on how to model α<sub>Ξ<sup>0</sup>K<sup>+</sup></sub> and B(Λ<sup>+</sup><sub>c</sub> → Ξ<sup>0</sup>K<sup>+</sup>) simultaneously.
- → The process  $\Lambda_c^+ \to \Lambda K^+ \pi^0$  was observed and BF of  $\Lambda_c^+ \to \Xi^0 K^+ \pi^0$  was updated, which are both lower than prediction based on SU(3) symmetry.
- ➤ The SCS decay  $\Lambda_c^+ \to \Sigma^- K^+ \pi^+$  was observed for the first time and consistent with SU(3) flavor symmetry prediction  $(3.3 \pm 2.3) \times 10^{-4}$ .
- ➤ The polarization of two-body intermediate channels in the three-body decay  $\Lambda_c^+ \to \Lambda \pi^+ \pi^0$  was measured, and more similar analyses are ongoing.
  ✓  $\Lambda_c^+ \to pK^-\pi^+$ ✓  $\Lambda_c^+ \to pK_s^0\pi^0$ ✓  $\Lambda_c^+ \to \Lambda_c^0\pi^+\eta$ ✓  $\Lambda_c^+ \to \Sigma^-\pi^+\pi^+$ ✓  $\Lambda_c^+ \to pK_s^0\pi^+$ ✓  $\Lambda_c^+ \to \Sigma^+\pi^+\pi^-$ ✓ ...
- → More polarization information about  $\Lambda_c^+ \rightarrow pK_s^0/\Lambda^0 \pi^+/\Sigma^0 \pi^+/\Sigma^+ \pi^0$  will be released soon.

# Thanks!