Pions and tensor force in heavy quark hadrons

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Contents

1. Introduction

From Yukawa to Nambu Problems in hadron physics

2. Heavy Quark Hadrons

Heavy *exotic* baryons and mesons

as hadronic *Molecules*

- **Z**_b: Recently observed at Belle
- 3. Summary

1. Introduction

What binds the protons and nucleon? 1934 at Osaka Univ

• Finite (long) range

$$\Delta \phi = 0 \implies (\Delta + m_{\pi}^2)\phi = 0$$
$$m_{\pi} \sim \frac{1}{\langle r \rangle} \sim 200 \text{ MeV} \quad \text{This is light}$$



H. Yukawa

• Pseudoscalar: $J^{P} = 0^{T}$



Spontaneous breaking of chiral symmetry

(1) Light bare quarks=> Massive constituent quarks

(2) Appearance of the *massless pion*







Problems in hadron physics

 $R(s) = \sigma(e^+e^- \rightarrow hadrons, s) / \sigma(e^+e^- \rightarrow \mu^+\mu^-, s).$



Problems in hadron physics

 $R(s) = \sigma(e^+e^- \rightarrow hadrons, s) / \sigma(e^+e^- \rightarrow \mu^+\mu^-, s).$



Threshold region

2011 Oct. 21

2. Heavy Quark Hadrons

u	С	t
d	S	b

Light ~ q Heavy ~ Q

π: ud^{bar}, etcJP = 0-ρ:JP = 1-

P(0-) and **P*(1-)**

- $\overline{\mathbf{D}}$: C^{bar} u, etc JP = 0- $\overline{\mathbf{D}}^*$: JP = 1-
 - **B**: B^{bar} u, etc JP = 0-**B**^{*}: JP = 1-





Pions appear where light quarks are

π-Yukawa vertex

πDD, πBB vertex is forbidden
BUT πDD*, πD*D* (πBB*, πB*B*) are allowed



$$\mathcal{L}_{vHH} = -i\beta \operatorname{Tr} \left[H_b v^{\mu}(\rho_{\mu})_{ba} \bar{H}_a \right] + i\lambda \operatorname{Tr} \left[H_b \sigma^{\mu\nu} F_{\mu\nu}(\rho)_{ba} \bar{H}_a \right]$$
$$H_a \sim q_a \bar{Q} = \frac{1+\psi}{2} \left[P_{a\mu}^* \gamma^{\mu} - P_a \gamma_5 \right] ,$$
$$\bar{H}_a \sim \gamma_0 (q_a \bar{Q})^{\dagger} \gamma_0 = \gamma_0 H_a^{\dagger} \gamma_0$$
$$A_{ab}^{\mu} = \frac{1}{2} (\xi^{\dagger} \partial^{\mu} \xi - \xi \partial^{\mu} \xi^{\dagger}) \simeq \frac{\imath}{f_{\pi}} \partial^{\mu} \hat{\pi}_{ab}$$
10



Yamaguchi, Ohkoda,Yasui, Hosaka Phys.Rev.D84:014032,2011. e-Print: 1105.0734 [hep-ph]

$$m_{K^*} - m_K \sim 400 \ MeV$$

 $m_{D^*} - m_D \sim 140 \ MeV$
 $m_{B^*} - m_B \sim 35 \ MeV$

Heavy quark symmetry Degeneracy SD mixing

	$\bar{D}N(\pi ho\omega)$	$BN(\pi\rho\omega)$
E_B [MeV]	2.14	23.04
$\langle r^2 \rangle^{1 \over 2} ~[{ m fm}]$	3.2	1.2





Exotic P^(*)N molecules

- Pion exchange potential dominates
- Tensor force causes SD mixing for heavy Q systems => Strong attraction
- These features are for m_Q ≥ m_C

Z_b resonance at **Belle** KEK





Exotic hadrons (mesons)

	State	Mass (MeV)	Width (MeV)	Decay	Production
	Ys(2175)	2175±8	58±26	ff ₀	ISR
	X(3872)	3871.84±0.33	<0.95	J/ypp, J/yg	B decay
	X(3872)	3872.8 +0.7/-0.6	3.9 +2.8/-1.8	D*0D0	B decay
	Z(3940)	3929±5	29±10	DD	gg
	X(3940)	3942±9	37±17	DD*	Double-charm
	Y(3940)	3942±17	87±34	J/yw	B decay
	Y(4008)	4008 +82/-49	226 +97/-80	Ј/урр	ISR
	Z(4051)+	4051 +24/-43	82 +51/-28	pc _{c1}	B decay
	X(4160)	4156±29	139 +113/-65	D*D*	Double-charm
	Z(4248)+	4248 +185/-45	177 +320/-72	pc _{c1}	B decay
	Y(4260)	4264±12	83±22	J/ypp	ISR
	Y(4350)	4361±13	74±18	y'pp	ISR
\rightarrow	Z(4430)+	4433±5	45 +35/-18	y'p	B decay
	Y(4660)	4664±12	48±15	y'pp	ISR
	Y _b (10890)	10889.6±2.3	54.7 +8.9/-7.6	ppƳ(nS)	e⁺e⁻ annihilation
	Y(3915)	3915±4	17±10	J/yw	gg
	X(4350)	4350 +4.7/-5.1	13 +18/-14	J/yf	gg
Twin	h _b (1P)	9898.3±1.5		MM(pp)	$\Upsilon(5S)/Y_{b}$ decay
	h _b (2P)	10259.3 +1.6/-1.2		MM(pp)	Ƴ(5S) /Y _b decay
	Z _b (10610)	10608.4±2.0	15.6±2.5	$(\Upsilon(nS) \text{ or } h_b)p$	$\Upsilon(5S)/Y_{b}$ decay
2011 Nov.	Z _b (10650)	10653.2±1.5	14.4±3.2	$(\Upsilon(nS) \text{ or } h_b)_P$	$\Upsilon(5S)/Y_{b}$ decay



Invariant mass of πY(nS)



What are interesting?

Heavy quark limit -> Heavy quark spin is conserved



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Heavy quark limit -> Heavy quark spin is conserved



 $h_h \pi$

$B^{(*)}B^{(*)}$ classification (I = 1)

J^{PC}	components		exoticness	
		I = 0	I = 1	
0++	${ m B}ar{ m B}({}^1S_0),{ m B}^*ar{ m B}^*({}^1S_0),{ m B}^*ar{ m B}^*({}^5D_0)$	$\chi_{ m b0}$	\checkmark	
1+-	$\frac{1}{\sqrt{2}} \left(\mathbf{B}\bar{\mathbf{B}}^* - \mathbf{B}^*\bar{\mathbf{B}} \right) ({}^3S_1), \ \frac{1}{\sqrt{2}} \left(\mathbf{B}\bar{\mathbf{B}}^* - \mathbf{B}^*\bar{\mathbf{B}} \right) ({}^3D_1), \ \mathbf{B}^*\bar{\mathbf{B}}^* ({}^3S_1), \ \mathbf{B}^*\bar{\mathbf{B}}^* ({}^3D_1)$	h_b	\mathbf{Z}_{b}	
1++	$\frac{1}{\sqrt{2}} \left(\mathbf{B}\bar{\mathbf{B}}^* + \mathbf{B}^*\bar{\mathbf{B}} \right) ({}^3S_1), \ \frac{1}{\sqrt{2}} \left(\mathbf{B}\bar{\mathbf{B}}^* + \mathbf{B}^*\bar{\mathbf{B}} \right) ({}^3D_1), \ \mathbf{B}^*\bar{\mathbf{B}}^* ({}^5D_1)$	$\chi_{ m b1}$	\checkmark	
2^{++}	$B\bar{B}(^{1}D_{2}), \ \frac{1}{\sqrt{2}}\left(B\bar{B}^{*}+B^{*}\bar{B}\right)(^{3}D_{2}), \ B^{*}\bar{B}^{*}(^{1}D_{2}), \ B^{*}\bar{B}^{*}(^{5}S_{2}), \ B^{*}\bar{B}^{*}(^{5}D_{2}), \ B^{*}\bar{B}^{*}(^{5}G_{2})$	$\chi_{ m b2}$		

$B^{(*)}B^{(*)}$ classification (I = 1)





Four states are found

J^{PC}	components		exoticness	
		I = 0	I = 1	
0++	$B\bar{B}(^{1}S_{0}) B^{*}\bar{B}^{*}(^{1}S_{0}), B^{*}\bar{B}^{*}(^{5}D_{0})$	$\chi_{ m b0}$	\checkmark	
1+-	$\frac{1}{\sqrt{2}} \left(\mathrm{B}\bar{\mathrm{B}}^* - \mathrm{B}^*\bar{\mathrm{B}} \right) ({}^3S_1), \ \frac{1}{\sqrt{2}} \left(\mathrm{B}\bar{\mathrm{B}}^* - \mathrm{B}^*\bar{\mathrm{B}} \right) ({}^3D_1), \ \mathrm{B}^*\bar{\mathrm{B}}^* ({}^3S_1), \ \mathrm{B}^*\bar{\mathrm{B}}^* ({}^3D_1) \right)$	h _b	Z_b	
1++	$\sum_{\sqrt{2}} \left(\mathbf{B}\bar{\mathbf{B}}^* + \mathbf{B}^*\bar{\mathbf{B}} \right) \left({}^3S_1 \right), \frac{1}{\sqrt{2}} \left(\mathbf{B}\bar{\mathbf{B}}^* + \mathbf{B}^*\bar{\mathbf{B}} \right) \left({}^3D_1 \right), \mathbf{B}^*\bar{\mathbf{B}}^* \left({}^5D_1 \right)$	$\chi_{ m b1}$	\checkmark	
2^{++}	$B\bar{B}(^{1}D_{2}), \frac{1}{\sqrt{2}} (B\bar{B}^{*} + B^{*}\bar{B}) (^{3}D_{2}), B^{*}\bar{B}^{*}(^{1}D_{2}), B^{*}\bar{B}^{*}(^{5}S_{2}) B^{*}\bar{B}^{*}(^{5}D_{2}), B^{*}\bar{B}^{*}(^{5}G_{2})$	$\chi_{ m b2}$	\checkmark	



More for I = 0 but with exotic J^{PC}

Ohkoda, Yamaguchi, Yasui and Hosaka, in preparation



Summary

- In the threshold region, **light quarks** appear and the **pion** starts to play.
- Exotic baryons are very likely to exist in the heavy quark region, m_Q ≥ m_c, as a hadronic molecule.
- The Z_b may well be a BB molecule bound by the pion