# Dynamical Holographic QCD Model for Chiral Symmetry Breaking and Confinement

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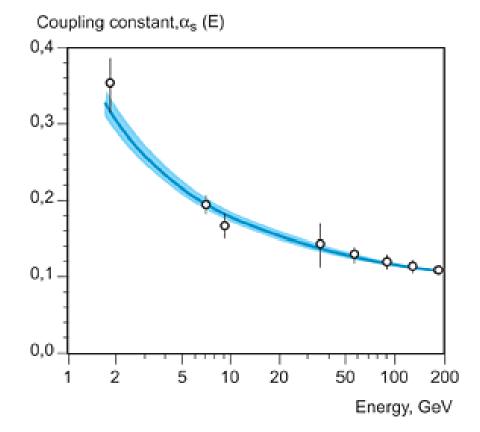
Beihang, Oct, 28, 2013

## Outlines

- I.A brief introduction to holographic QCD
- II.Graviton-Dilaton system and pure gluon system
- III. Accommodate chiral symmetry breaking and linear confinement in a dynamical holographic QCD model
- IV.Conclusion

### I.A brief introduction to holographic QCD

### Why holographic QCD?



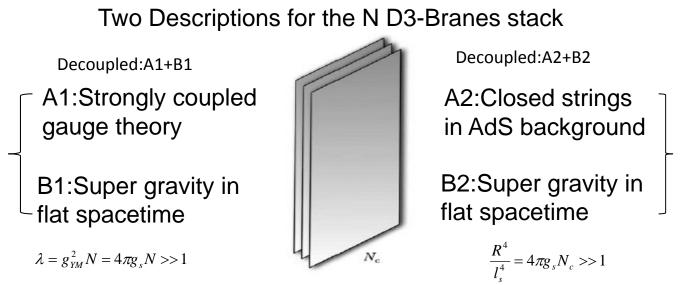
Perturbation calculation fails in low energy region

Non-perturbative methods: Lattice simulation, DSE,FRG, Effective field theory...

Holographic QCD method: analytic understanding, easier calculation

### AdS/CFT Correspondence

Juan Maldacena, Adv. Theor. Math. Phys. 2:231-252, 1998

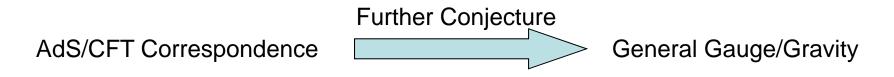


Conjecture: Duality of the two theories

Equivalence of the partition function:

$$< e^{\int \phi_0^i O_i} >_{CFT} = Z_{string} [\phi^i(x, z \mid_{boundary}) = \phi_0^i(x)] \sim e^{-S_{Gravity}}$$

### **Towards Realistic Models**



**Task:** find the proper non-conformal gravity background for QCD:

#### **Top-Down Approach:**

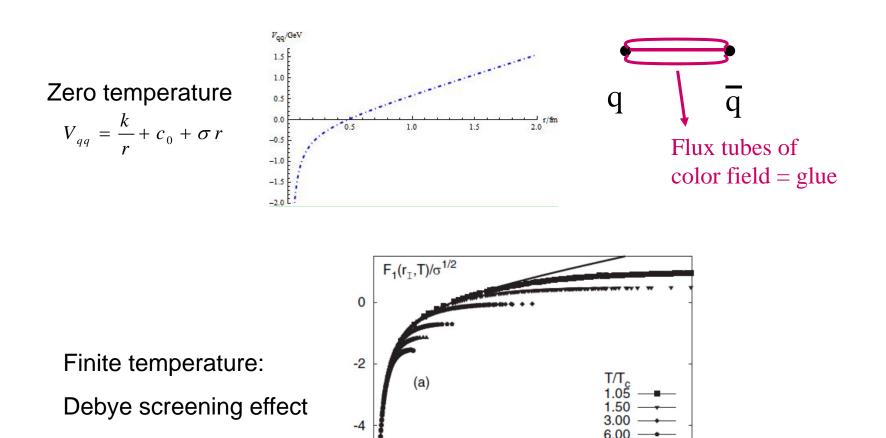
derive the background from fundamental string theory

#### **Bottom-Up Approach:**

find the background in a phenomenological way provide clues to the Top-down

# II.Graviton-Dilaton system and pure gluon system

# Confinement in QCD vacuum and deconfinement at finite temperature



0.2

0

0.6

12.0

0.8

r<sub>⊤</sub> [fm]

0.4

### **Graviton-Dilaton System**

$$S_{GD} = \frac{1}{2\kappa^2} \int dx^5 \sqrt{-g} e^{-2\phi} (R + 4\partial_m \phi \partial^m \phi + V_s(\phi))$$

Zero temperature:

deform the background metric: the warp factor

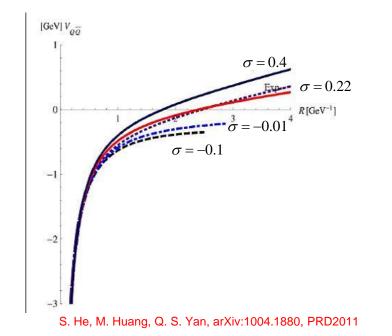
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$$dS^{2} = g_{mn}dx^{m}dx^{n} = \frac{e^{2A_{s}}}{z^{2}}(-dt^{2} + dz^{2} + dx_{i}dx^{i})$$

Andreev-Zakharov

Quadratic warp factor

$$A_s = \frac{\sigma z^2}{2}$$



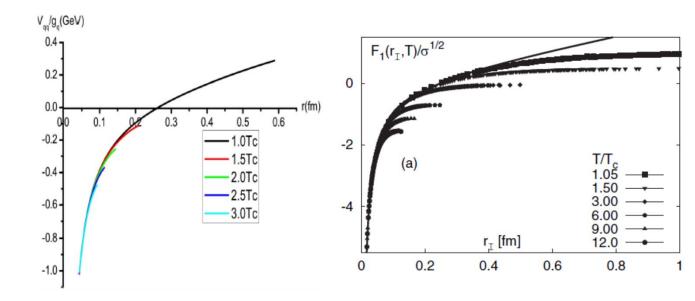
Study the finite T behavior: Black hole solution

$$dS^{2} = g_{mn}dx^{m}dx^{n} = \frac{e^{2A_{s}}}{z^{2}}(-f(z)dt^{2} + \frac{1}{f(z)}dz^{2} + dx_{i}dx^{i})$$
$$T = \frac{|f'(z_{h})|}{4\pi}.$$

$$A_s = \frac{\sigma z^2}{2}$$

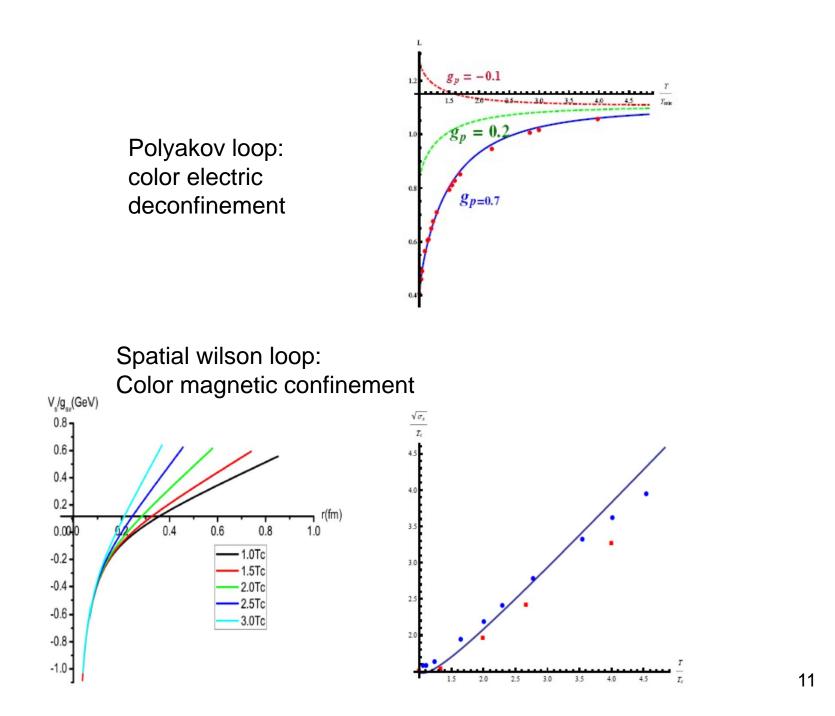
Dilaton profile and potential can be self-consistently solved from the equation of motion

Finite temperature behavior in the gravity background

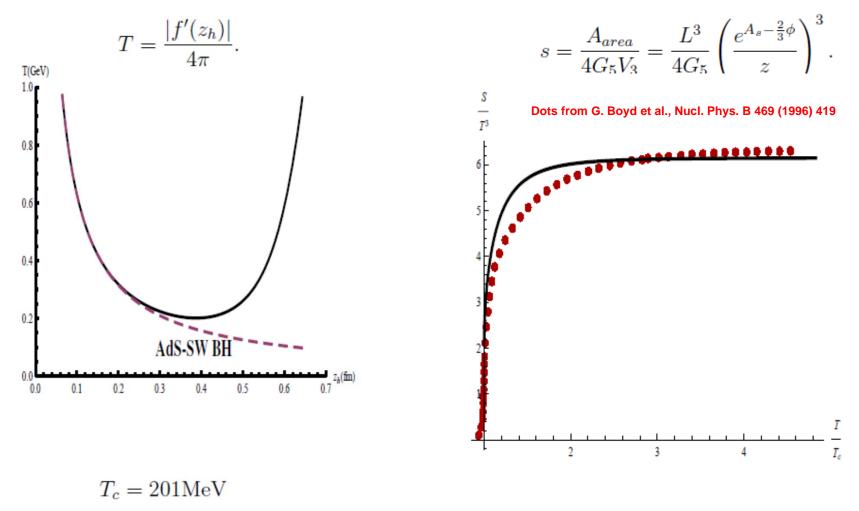


D.N. Li, S. He, M.Huang, Q. S. Yan, arXiv:1103.5389, JHEP2011

O.Kaczmarek, F.Karsch, and F.Zantow, Phys.Rev.D.70.074505



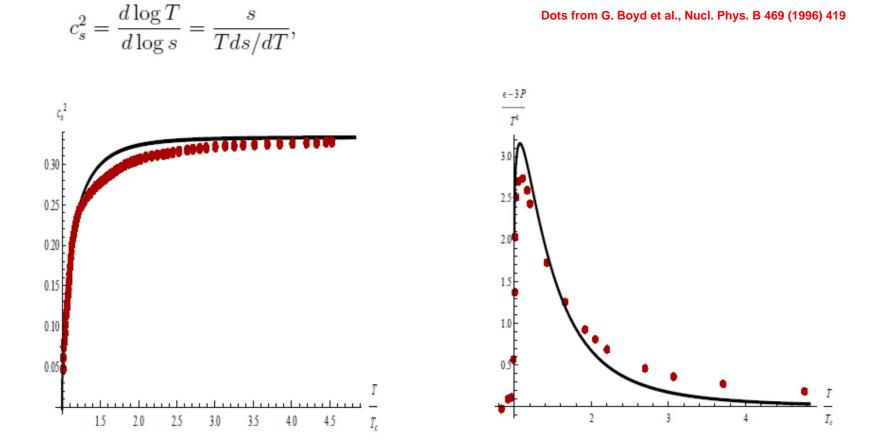
### EOS V.S. Quenched Lattice Results



D.N. Li, S. He, M.Huang, Q. S. Yan, arXiv:1103.5389, JHEP2011

12

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D.N. Li, S. He, M.Huang, Q. S. Yan, arXiv:1103.5389, JHEP2011

13

Dots from G. Boyd et al., Nucl. Phys. B 469 (1996) 419

- At zero temperature and finite temperature, such a graviton-dilaton coupled system can describe the heavy quark potential behavior very well.
- The equation of states fits the quenched lattice data well (above Tc), which indicate that such a system can describe the pure gluon dynamics.

III. Accommodate chiral symmetry breaking and linear confinement in a dynamic holography QCD model

# Chiral symmetry breaking and confinement from meson spectra

#### Chiral symmetry breaking:

Chiral condensate spontaneously breaks the chiral symmetry, and provide quarks dynamical masses.

Mass split in chiral partner

### 

#### **Confinement:**

Linear spectra in highly excited states

Data from C.Amsler et.al, Phys.Lett.B667(2008)1

### **KKSS** model

Karch, Katz, Son, Stephonov, Phys. Rev. D74:015005, 2006

$$S_{kkss} = \int dx^{5} \sqrt{-g} e^{-\phi} ||DX|^{2} + 3|X|^{2} - \frac{1}{4g_{5}^{2}} (F_{L}^{2} + F_{R}^{2})| \phi = \mu^{2} Z^{2} \qquad m_{\rho,n}^{2} = 4\mu^{2}n$$
  
Chiral symmetry is broken by  
the expectation value of X  

$$X = \frac{\chi}{2} I_{2x2}, \quad \chi = m_{q} z + \sigma z^{3} + o(z^{4})$$

$$V_{\rho} = \frac{A_{s}^{"} - \phi^{"}}{2} + \frac{(A_{s}^{'} - \phi^{'})^{2}}{4}$$

$$V_{a_{1}} = \frac{A_{s}^{"} - \phi^{"}}{2} + \frac{(A_{s}^{'} - \phi^{'})^{2}}{4} + g_{5}^{2} \chi^{2} e^{2A_{s}}$$
Chiral Symmetry  
breaking
$$I^{T}$$

$$Drive a linear part
$$I^{T}$$$$

A successful holographic QCD model

- should incoperate the linear potential and linear spectra (two aspects of confinement) simultaneously
- should realize chiral symmetry breaking and confinement simultaneously

# Accommodate the two aspects at zero temperature

$$S = S_{GD} + \frac{N_f}{N_c} S_{KKSS}$$

$$N_f = 2, N_c = 3$$

Zero temperature

Dilaton profile

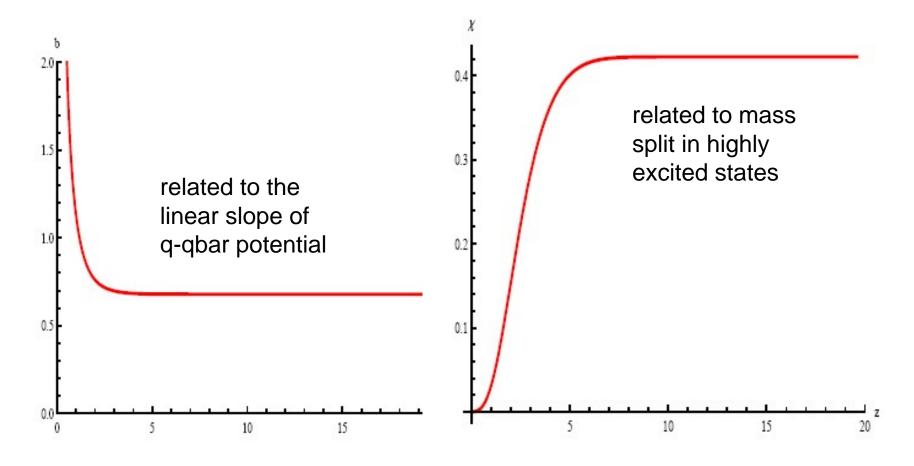
$$dS^{2} = e^{2A_{s}(z)}(-dt^{2} + dz^{2} + dx_{i}dx^{i})$$

$$\phi = \mu^2 z^2$$

UV behavior of  $\chi$ 

$$\chi = m_q z + \sigma z^3 + o(z^4)$$

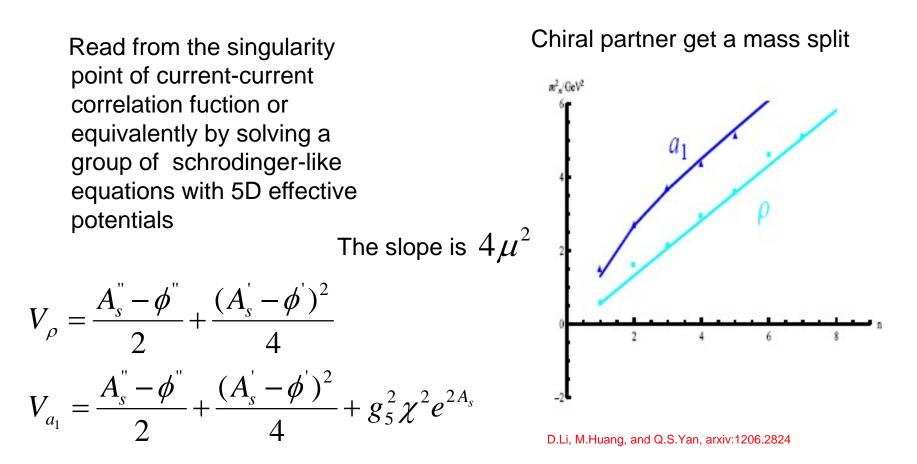
## The Background Solution



D.Li, M.Huang, and Q.S.Yan, arxiv:1206.2824

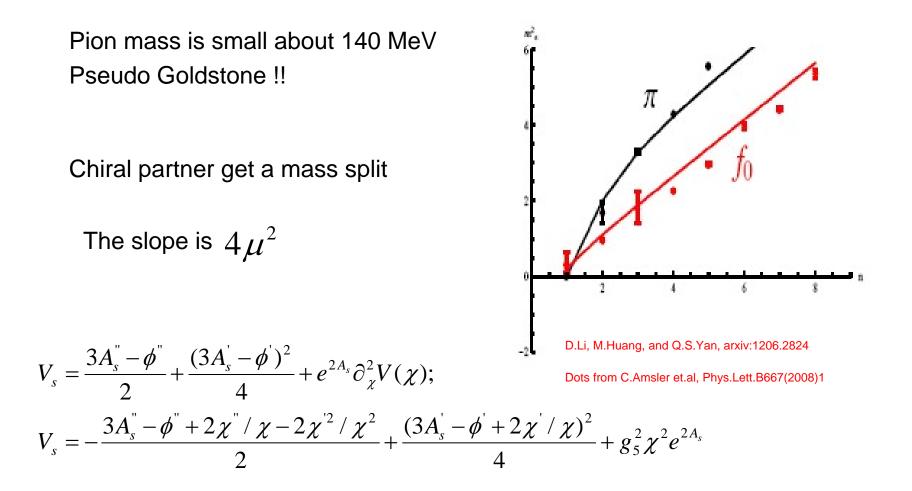
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### Rho-a masses



Dots from C.Amsler et.al, Phys.Lett.B667(2008)1

## pi-f0 masses



22

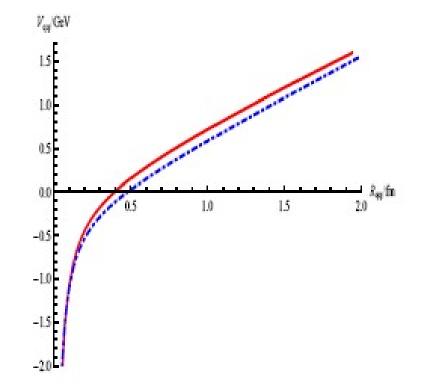
# The quark antiquark potential

The linear potential is produced.

The slope of the linear part is:  $\sigma_s \approx 4\mu^2$ 

$$\phi = \mu^2 z^2$$
$$m_{\phi}^2 L^2 = -4 \int_{\Delta_{\phi}} \Delta_{\phi} = 2$$

Some dimension-2 operatror related to gluon dynamics ? Need further study



D.Li, M.Huang, and Q.S.Yan, arxiv:1206.2824

## **IV.Conclusion**

- The graviton-dilaton system describes the pure gluon dynamics above Tc quite well.
- Graviton-dilaton-scalar system can cover the gluon and chiral dynamics well. In this system, we accommodate chiral symmetry breaking, linear spectra and linear potential selfconsistently for the first time in holographic QCD frame. The slope of the linear potential and the linear spectra are the same and related to dilaton, which is consistent with the consideration that the two phenomena have the same origin---the gluon dynamics.
- Studying the confinement-deconfinement phase transition and chiral phase transition simultaneously is in progress.

# Thanks for your attention !