

Tensor force and Deltas for the structure of light nuclei

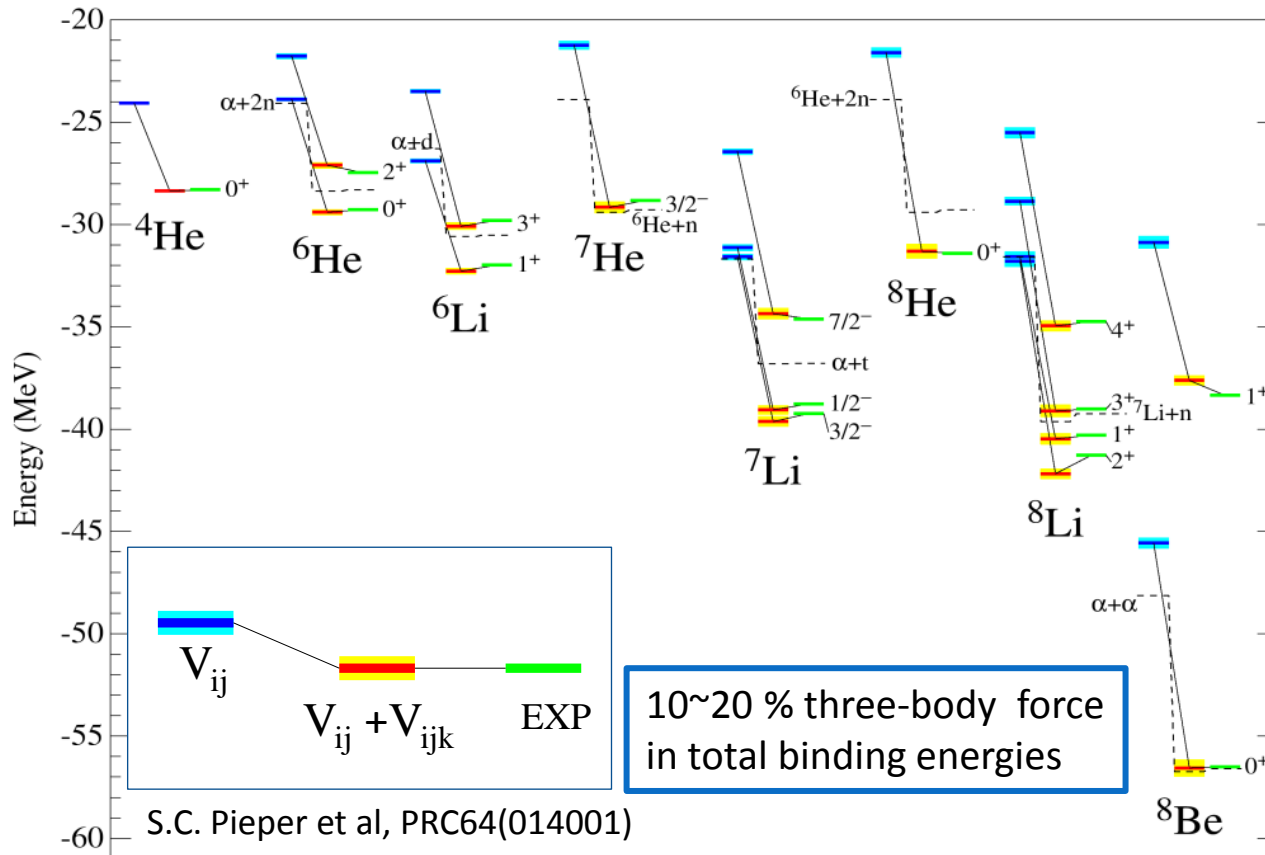
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in collaboration with Takayuki Myo, Hiroshi Toki

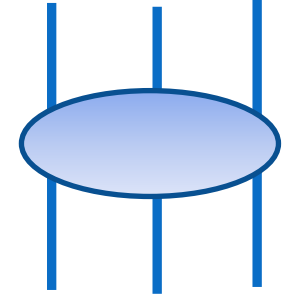
RCNP, Osaka University

Osaka Institute of Technology University

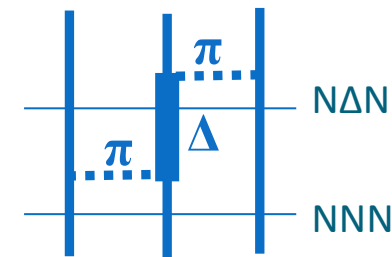
Three-body force and $\Delta(1232)$



Phenomenological three-body force



Fujita-Miyazawa three-body force model



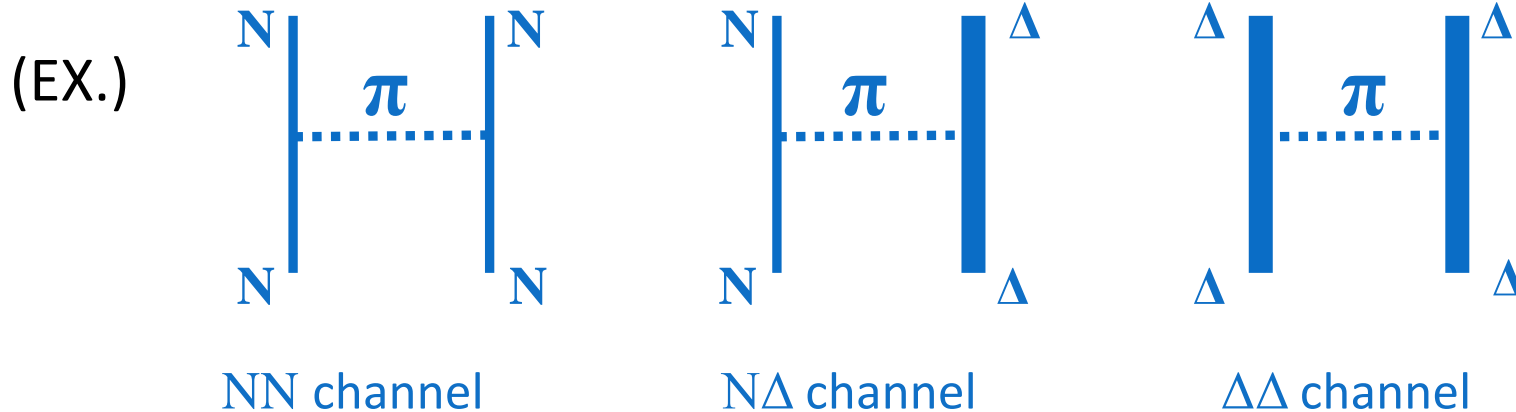
Delta can be generated by π exchange and play important roles in 3-body force.

Delta (1232) Spin $S=3/2$, Isospin $T=3/2$

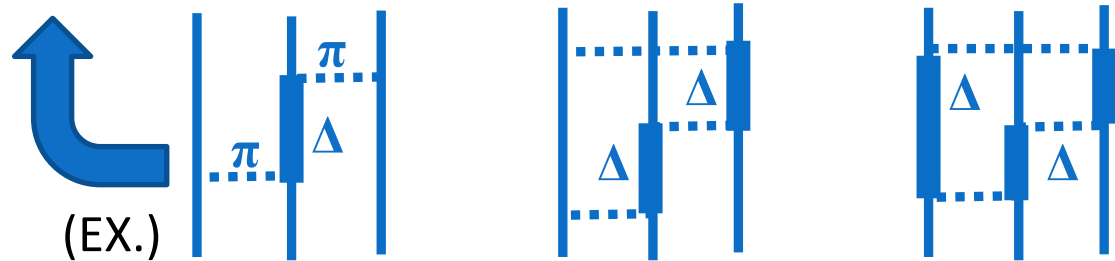
We introduce Δ in two-body force as the origin of three-body force.

Treatment of the delta state

We add the delta degrees of freedom in two-body interaction.



Many-body force processes can be treated by the two-body correlations with delta.



Two-body NN interaction with delta degrees of freedom \rightarrow AV28 potential

We study the three-body system with delta degrees of freedom.

Effect of three-body force ? Tensor force from delta ?

AV28 potential

R.B. Wiringa et al, PRC 29, 1207(1984) AV14 & AV28

AV14 → bare NN interaction, 14 operators term

AV28 → bare two-body interaction with Δ degree of freedom

NN, N Δ , $\Delta\Delta$ channel, 28 operators term

Structure of Δ part potential

N Δ , $\Delta\Delta$ without transition

$$V = V_S + V_I + V_\pi$$

(Ex.)



N $\leftrightarrow\Delta$ transition

$$V = V_\pi$$

(Ex.)



$V_S + V_I$ The short- and intermediate range phenomenological part : same NN potential

V_π Pion exchange part : Yukawa & Tensor type function

Coupling constant $(f^2_{\pi N\Delta}/4\pi) = 4(f^2_{\pi NN}/4\pi), (f^2_{\pi\Delta\Delta}/4\pi) = 0.04(f^2_{\pi NN}/4\pi)$

Brown and Weise, PRC(22)279 (1975)

Reproduce the deuteron properties and NN scattering phase-shifts

AV28 potential

R.B. Wiringa et al, PRC 29, 1207(1984) AV14 & AV28

AV14 → bare NN interaction, 14 operators term

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NN, N Δ , $\Delta\Delta$ channel, 28 operators term

Structure of Δ part potential

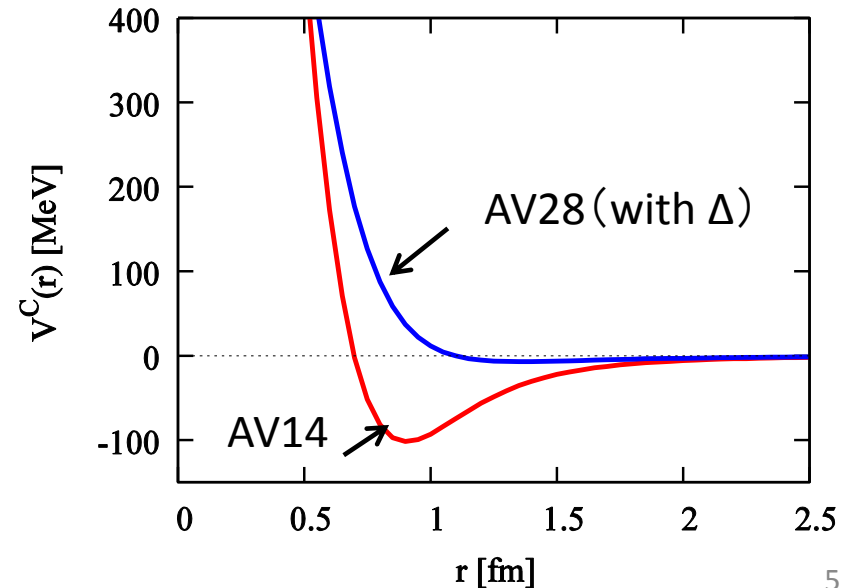
N Δ , $\Delta\Delta$ without transition

$$V = V_S + V_I + V_\pi$$

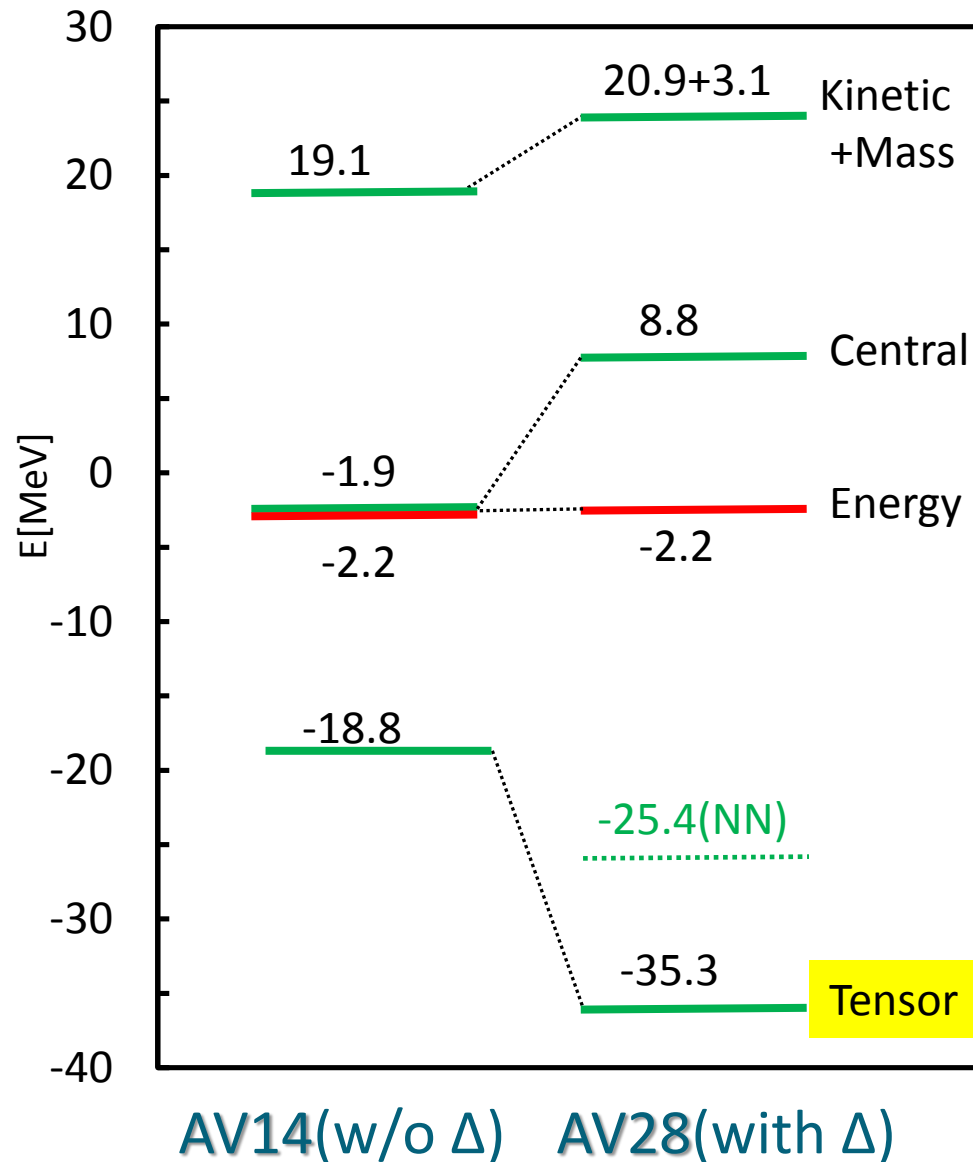
N $\leftrightarrow\Delta$ transition

$$V = V_\pi$$

Difference of the central force ($S=1, T=0$ (3E))
between AV14 & AV28 →



Effect of Δ in deuteron



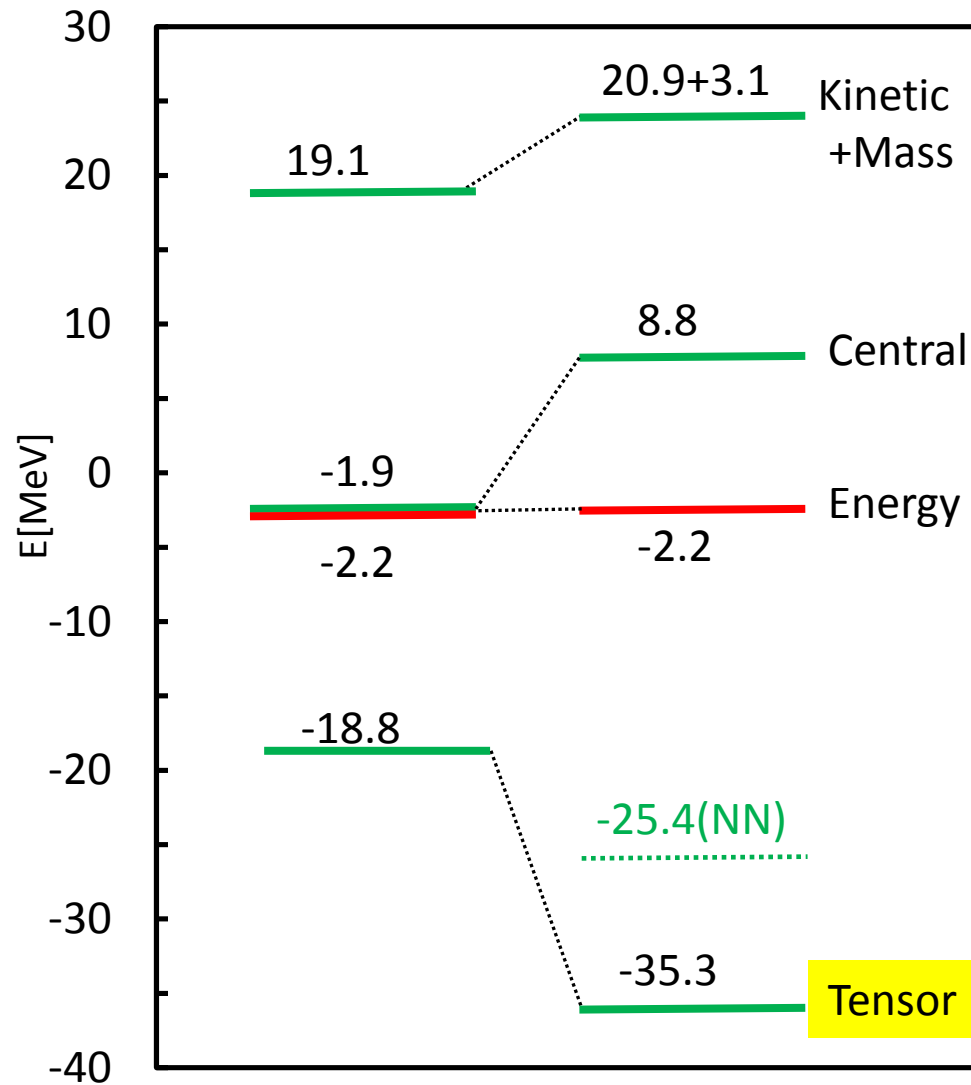
Deuteron 1^+	AV14	AV28
$L \cdot S$	0.4	0.8
L^2	3.1	3.6
$(L \cdot S)^2$	- 4.0	- 4.1
$P_{NN} [^3S_1] \%$	93.9	93.3
$P_{NN} [^3D_1]$	6.1	6.2
$P_{\Delta\Delta} [^3S_1]$		0.04
$P_{\Delta\Delta} [^3D_1]$		0.02
$P_{\Delta\Delta} [^7D_1]$		0.42
$P_{\Delta\Delta} [^7G_1]$		0.04

Wave function

$$\Psi_{NN} = |^3S_1\rangle + |^3D_1\rangle$$

$$\Psi_{\Delta\Delta} = |^3S_1\rangle + |^3D_1\rangle + |^7D_1\rangle + |^7G_1\rangle$$

Effect of Δ in deuteron



AV14(w/o Δ) AV28(with Δ)

Deuteron 1^+	AV14	AV28
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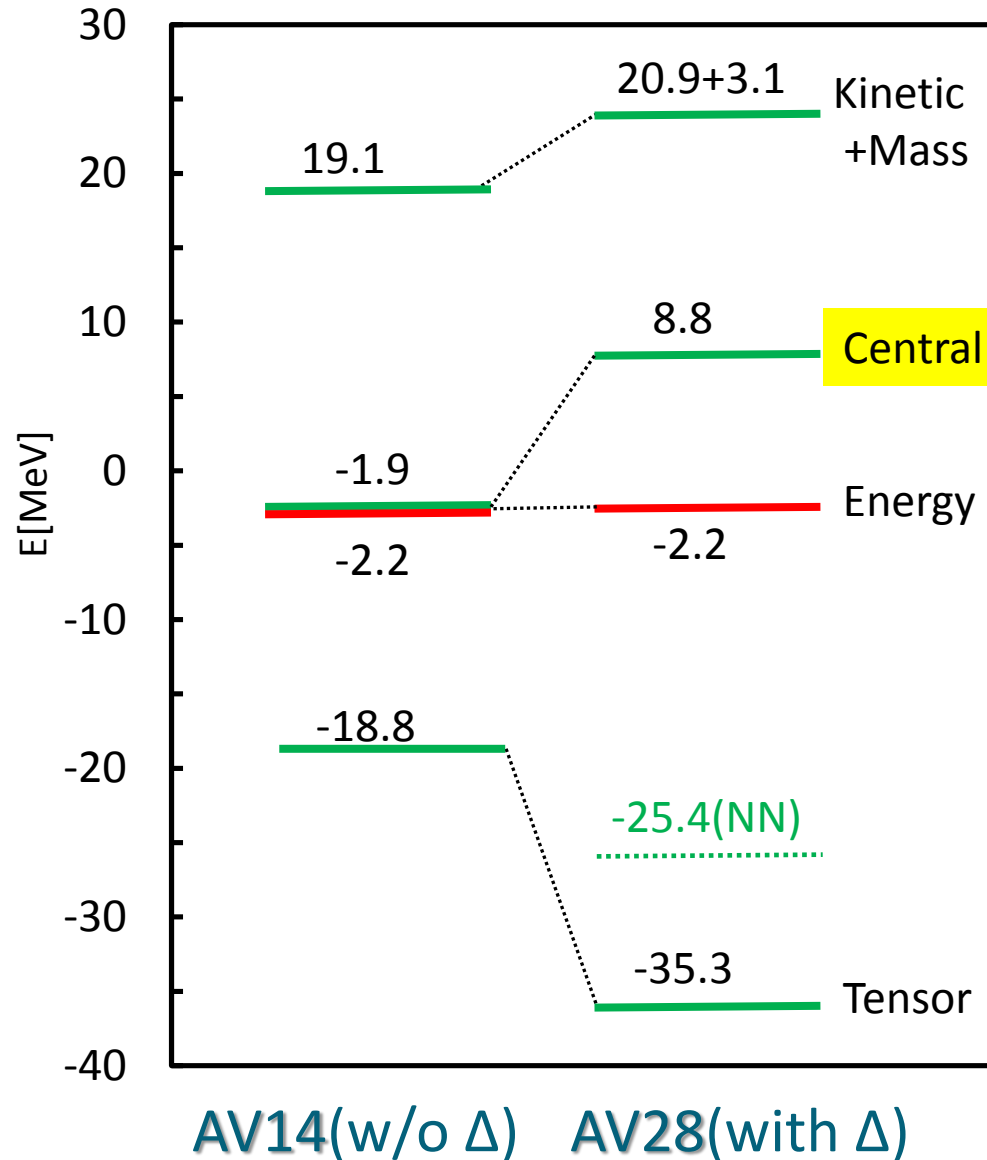
Wave function

$$\Psi_{NN} = |^3S_1\rangle + |^3D_1\rangle$$

$$\Psi_{\Delta\Delta} = |^3S_1\rangle + |^3D_1\rangle + |^7D_1\rangle + |^7G_1\rangle$$

Tensor couplings between 3S_1 (NN) and 7D_1 ($\Delta\Delta$) provide larger attractions.

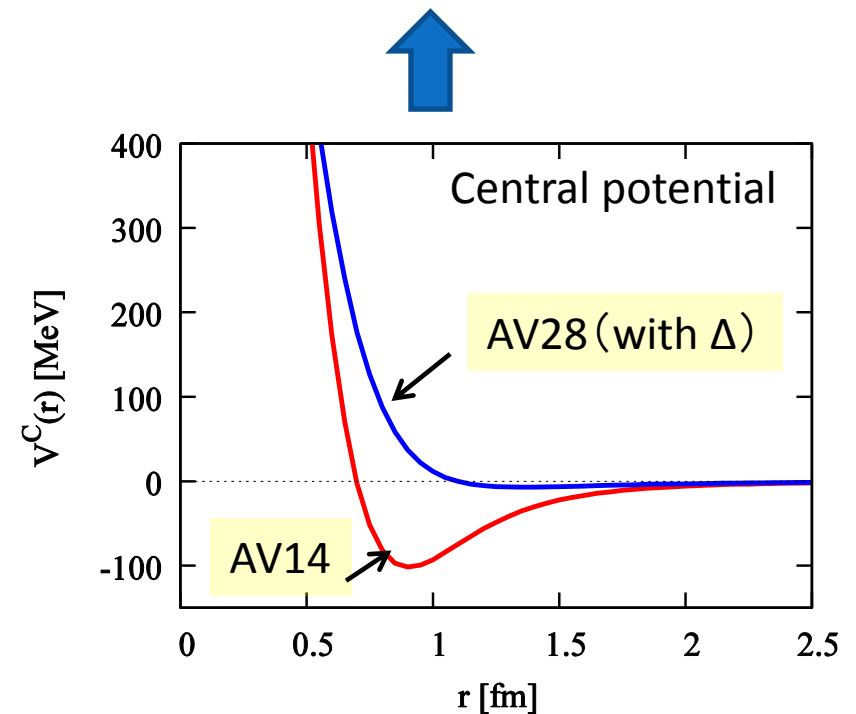
Effect of Δ in deuteron



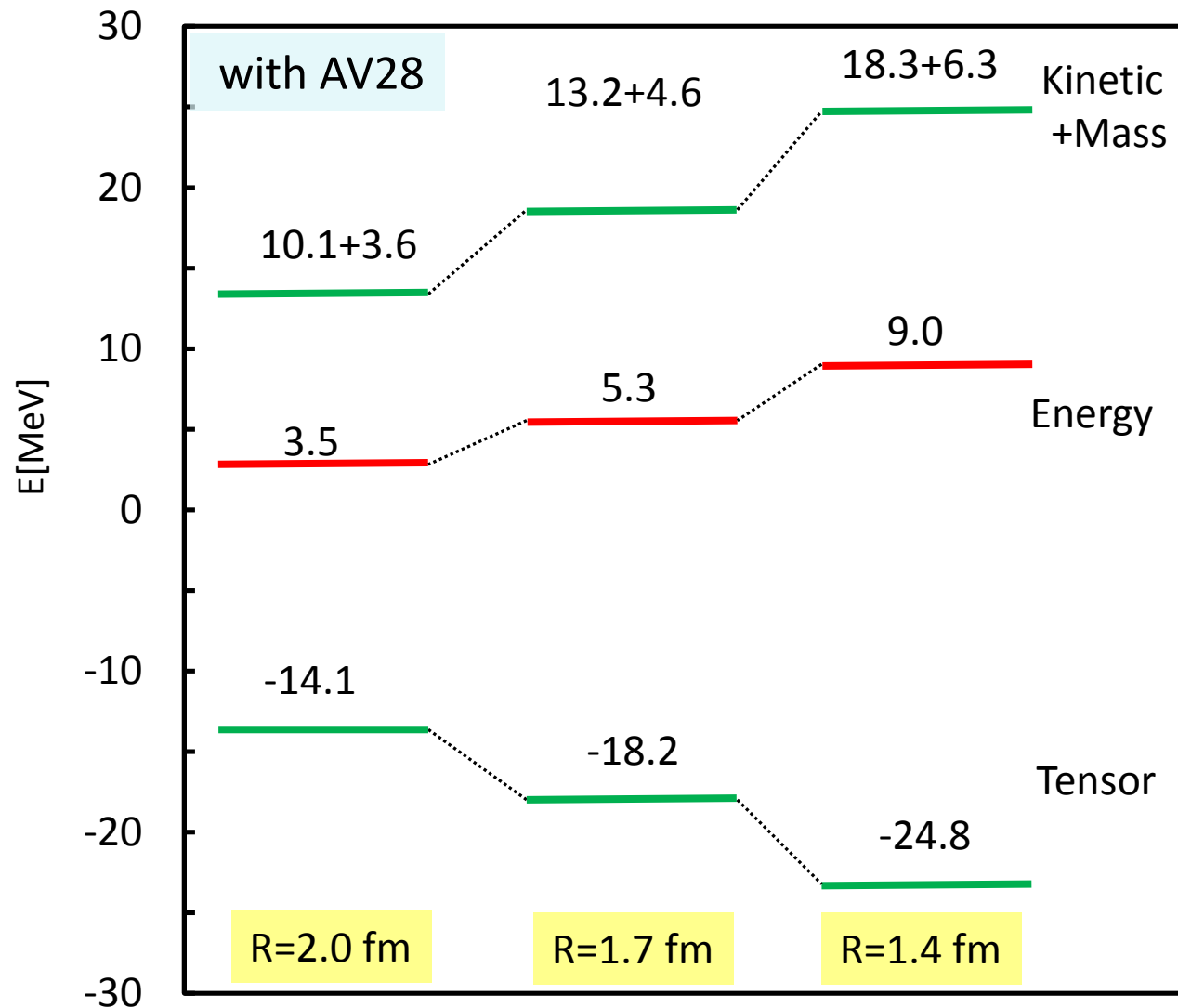
Central force components

AV14 = -1.9 MeV (attractive)

AV28 = 8.8 MeV (repulsive)



Effect of Δ in 1E channel ($T=1$)



Wave function

$$\Psi_{NN} = |{}^1S_0\rangle$$

$$\Psi_{\Delta\Delta} = |{}^1S_0\rangle + |{}^5D_0\rangle$$

$$\Psi_{N\Delta} = |{}^5D_0\rangle$$

Hamiltonian with radius constraint

$$\tilde{H} = H + \lambda \hat{R}^2$$

$$\begin{aligned} &\langle S(NN) | V_T | D(N\Delta) \rangle \\ &\langle S(NN) | V_T | D(\Delta\Delta) \rangle \end{aligned}$$

Previous work for triton including Δ

Two groups have studied ${}^3\text{H}$ including delta degree of freedom by Faddeev calculation.

1. Hannover Group (A. Deltuva et al PRC68,024005,(2003))

$\text{NN} + \text{N}\Delta$ model space with Bonn potential $\rightarrow (E_B = -8.2 \text{ MeV})$

$\Delta\Delta$ states are not taken into account !!

2. Los Alamos Group (A. Picklesminer et al [1]~[5])

$\text{NN} + \text{N}\Delta + \Delta\Delta$ model space with AV28 potential $\rightarrow (E_B = -7.3 \text{ MeV}, E_B(\text{exp}) = -8.4 \text{ MeV})$

Model space restriction? They truncated $L=4$ (G-wave) contributions in $\Delta\Delta$ state.

Deuteron wave function

$$\Psi_{NN} = |{}^3S_1\rangle + |{}^3D_1\rangle$$

$$\Psi_{\Delta\Delta} = |{}^3S_1\rangle + |{}^3D_1\rangle + |{}^7D_1\rangle + |{}^7G_1\rangle \leftarrow 0.04\%$$

$$\langle \Psi_S | S_{12} | \Psi_G \rangle = 0$$

Transitions between D-wave(NN) and G-wave ($\Delta\Delta$) are not negligible !!

$$\langle \Psi_D | S_{12} | \Psi_G \rangle \neq 0$$

NO G-wave $E_B = -1.6 \text{ MeV}$, Tensor = -28 MeV

Full $E_B = -2.2 \text{ MeV}$, Tensor = -35 MeV

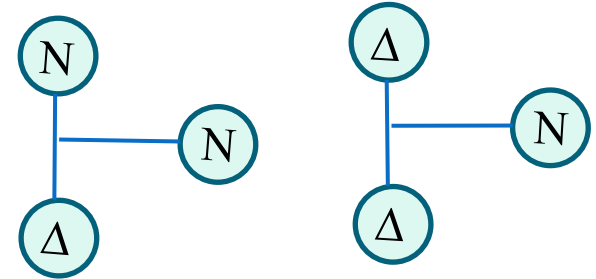
Triton including Δ degrees of freedom

^3H wave function

$$\Psi = \Psi_{NNN} + \Psi_{N\Delta N} + \Psi_{\Delta\Delta N}$$

$$\Psi_{N\Delta N} = \Psi_{N\Delta_2 N} + \Psi_{N\Delta_3 N} + \Psi_{N\Delta_1 N}$$

$$\Psi_{\Delta\Delta N} = \Psi_{\Delta_1\Delta_2 N} + \Psi_{\Delta_2\Delta_3 N} + \Psi_{\Delta_3\Delta_1 N}$$



We take into account all nucleon can excite the delta state.

Valiational calculation with Stocathtic Variational Method(SVM)

Few Body System 42(2008)33 Y.Suzuki, et al

Most important configurations for tensor correlation (TOSM concept)

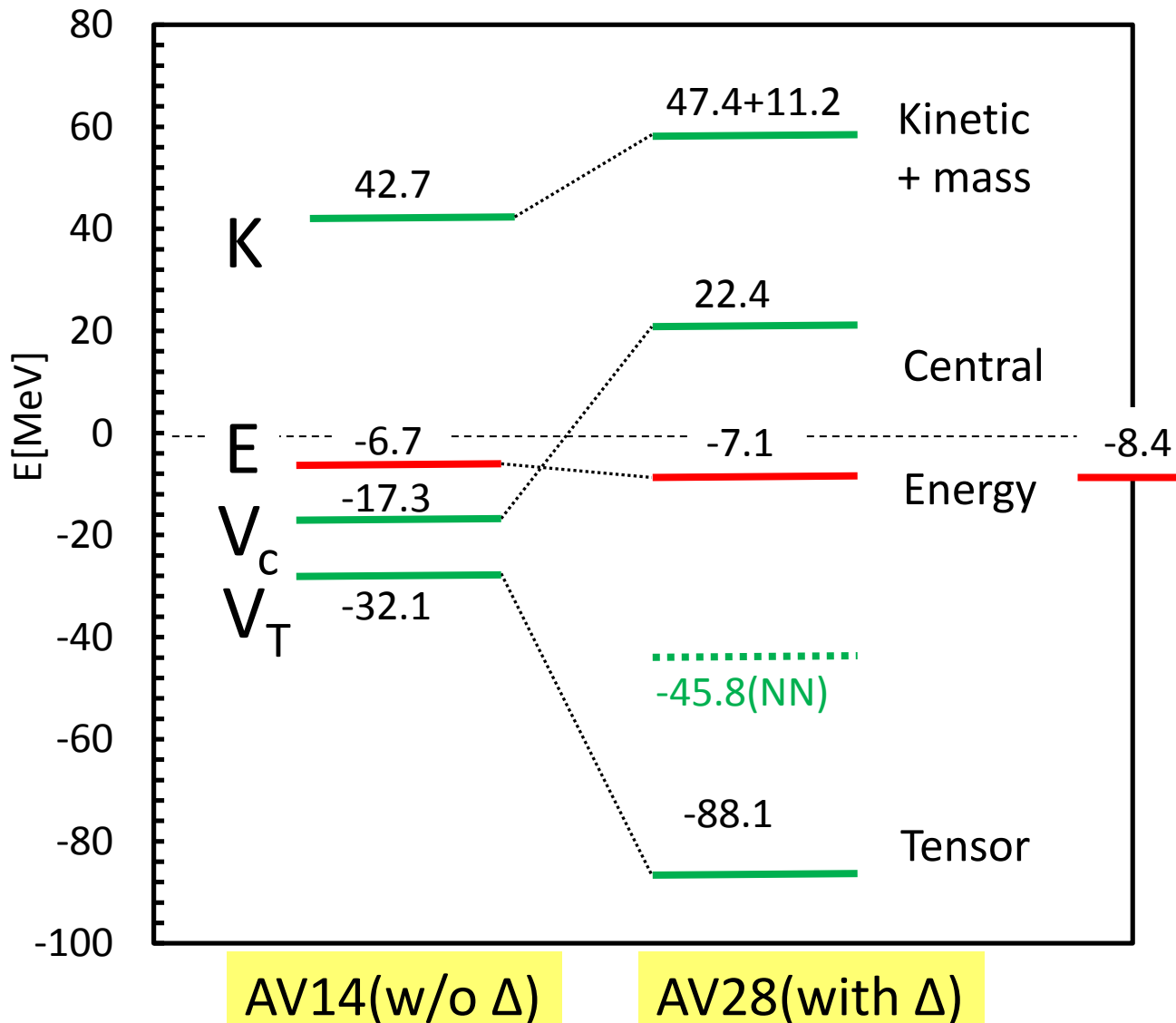
$$\left\{ \begin{array}{l} \Psi_{NNN} = \Psi_S + \Psi_D \\ \Psi_{N\Delta N} = \Psi_S + \Psi_D \\ \Psi_{\Delta\Delta N} = \Psi_S + \Psi_D + \Psi_G \end{array} \right.$$

Myo, Toki, Ikeda PTP 121 (2009) 511

K.H., Myo, Toki, Ikeda PTP 127 (2011) 1109

Result ^3H with AV14 & AV28

^3H $J^\pi = 1/2^+$

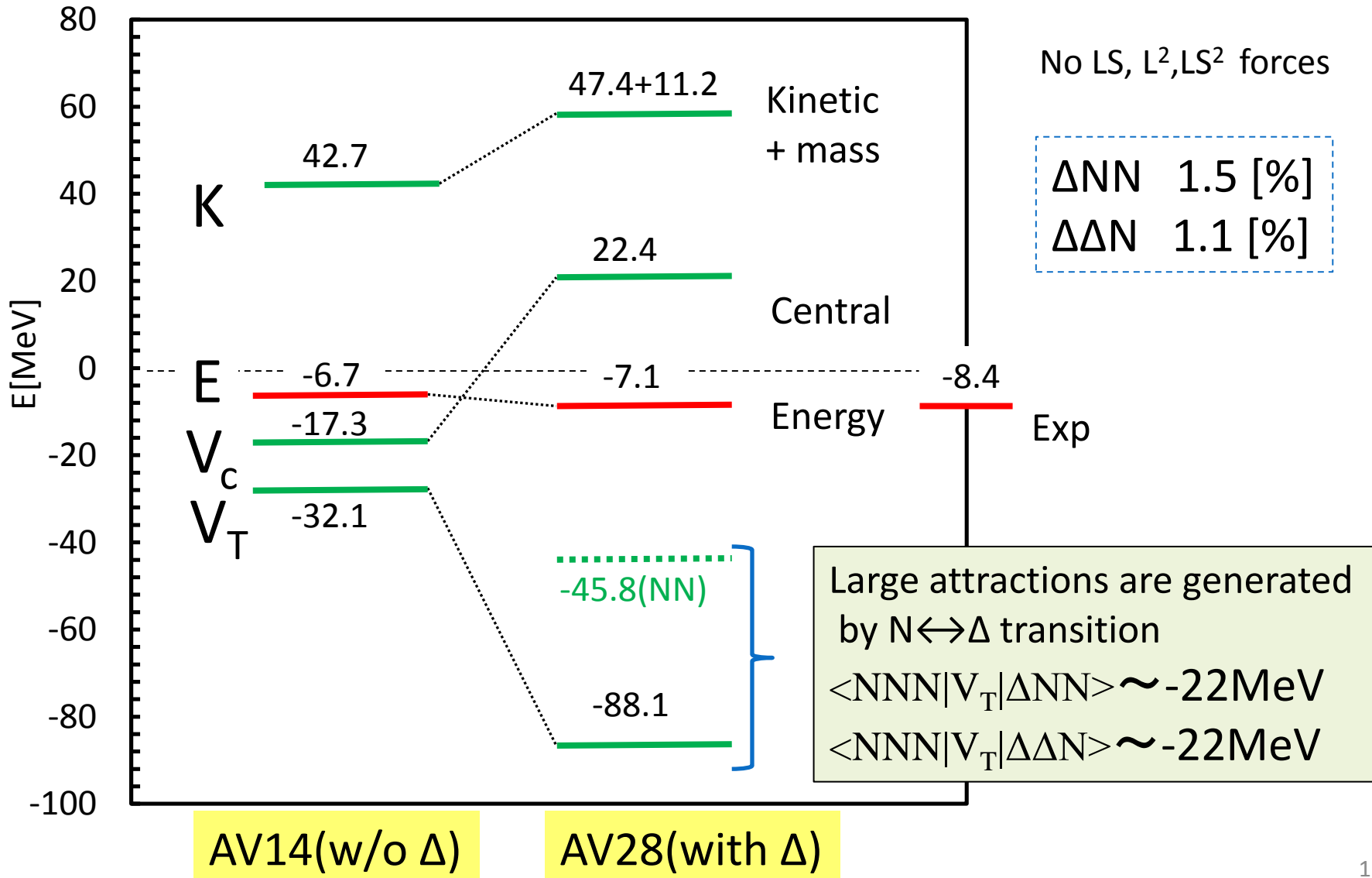


No LS, L^2 , LS^2 forces

ΔNN 1.5 [%]
 $\Delta\Delta N$ 1.1 [%]

Result ^3H with AV14 & AV28

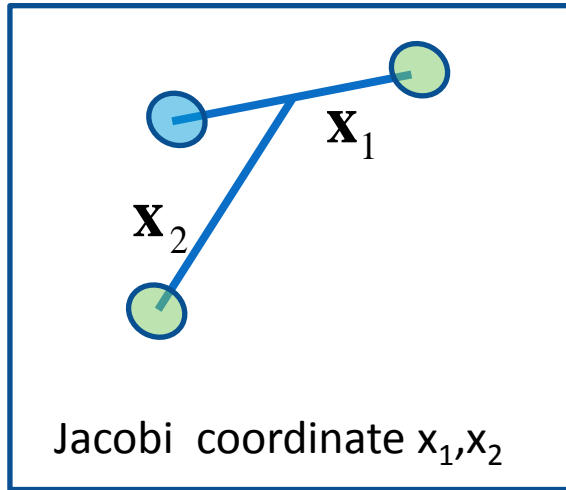
^3H $J^\pi = 1/2^+$



Summary & Outlook

1. As the origin of many-body force, we introduced the delta degrees of freedom in two-body interaction .
2. We used the two-body interaction AV28 including delta.
3. The transition between nucleon and delta provide large tensor correlations in deuteron and triton.
4. In the triton result, the tensor component extremely large and the changes of hamiltonian components are similar to the deuteron results.
5. We will develop the treatment of the delta state in the frameworks of TOFM • TOSM • EBHF.

Tensor Optimized Few-body Model for s-shell



TOFM w.f. $|\Psi\rangle = |\Psi\rangle_S + |\Psi\rangle_D$

$$\Psi = \sum_i C_i \mathcal{A}[\psi_L^{space} \chi_S^{spin}]_J \chi_T^{isospin}$$

For ${}^4\text{He}$ Total $J=0$, S-wave ($L=0, S=0$), D-wave ($L=2, S=2$)

Basis function Correlated gaussian basis with the global vector

$$\psi_L^{space} = \exp\left(-\frac{1}{2} \tilde{\mathbf{x}} \mathbf{A} \mathbf{x}\right) |\tilde{u} \mathbf{x}|^L Y_{LM}(\tilde{u} \mathbf{x})$$

Few Body System 42(2008)33 Y.Suzuki, et al

Correlated gaussian $\tilde{\mathbf{x}} \mathbf{A} \mathbf{x} = \sum_{i,j}^{N-1} A_{ij} \mathbf{x}_i \cdot \mathbf{x}_j$ global vector $\tilde{u} \mathbf{x} = \sum_{i=1}^{N-1} u_i \mathbf{x}_i$

Variational calculation with the Stochastic Variational Method (SVM)

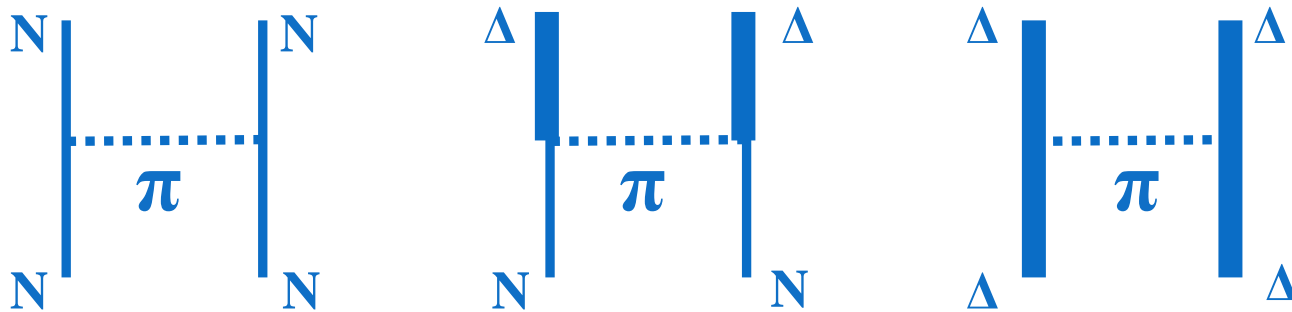
*All parameter of basis function are decided to obtain the energy minimum.

*The amplitude are decided by diagonalization.

Effect of Δ in deuteron

We would like to discuss the effect of delta in deuteron.

Deuteron $J^\pi = 1^+$, $T=0$ (No $N\Delta$ channel)



Wave function $\Psi = \Psi_{NN} + \Psi_{\Delta\Delta}$

$$\Psi_{NN} = |^3S_1\rangle + |^3D_1\rangle$$

$$\Psi_{\Delta\Delta} = |^3S_1\rangle + |^3D_1\rangle + |^7D_1\rangle + |^7G_1\rangle$$

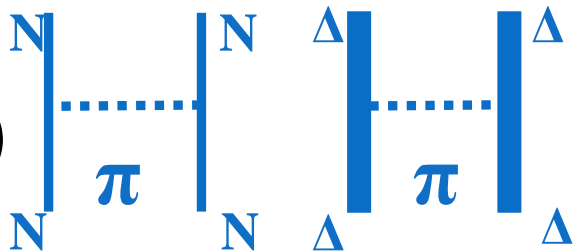
Hamiltonian

$$H = (m_1 + m_2 - 2m_N) + \frac{\hbar^2 (m_1 + m_2)}{2m_1 m_2} \nabla_{12}^2 + V_{12}$$

Effect of Δ in two-body system

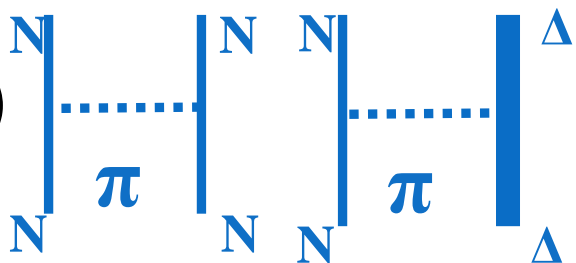
Effects of large tensor force by the transition between N and Δ

3E channel (T=0)



$$\langle \Psi_{NN} | V_T | \Psi_{\Delta\Delta} \rangle \sim \underline{-10MeV}$$

1E channel (T=1)



$$\langle \Psi_{NN} | V_T | \Psi_{N\Delta} \rangle \sim \underline{-8MeV}$$

$N\Delta$ states (T=1 channel) also generate large tensor correlations.
 → These transitions play important roles in three-body system.

Effect of three-body force by Δ

We analyze the effect of three-body force by the Δ excitation.

We separate all the Δ excitation processes into 2-body & 3-body force contributions.

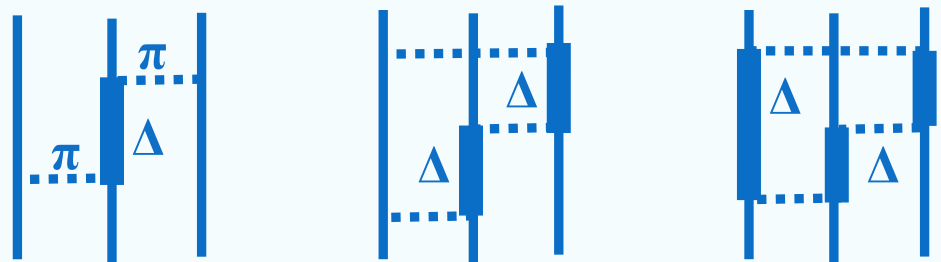
$$\langle \Psi | \sum_{i,j} V_{ij} | \Psi \rangle = \langle 2body \rangle + \langle 3body \rangle$$

2-body processes



The last one nucleon does not take part in the interaction.

3-body processes



All nucleon contributes to the Δ excitations.

Feshbach projection operator method

Projection operator P and Q

Normalization $P + Q = 1$

Wave function with delta

$$\Psi = C\Psi_N + D\Psi_\Delta$$

$$\begin{aligned} P\Psi &= C\Psi_N && \text{Nucleon states} \\ Q\Psi &= D\Psi_\Delta && \text{Delta states} \end{aligned}$$

Schrodinger equation $H\Psi = E\Psi$

$$PHP\Psi + PVQ\Psi = EP\Psi$$

$$QVP\Psi + QHQ\Psi = EQ\Psi$$



$$H^P C + G D = E C$$

$$H^Q D + G^\dagger C = E D$$

We eliminate the D using the second equation and obtain,

$$(H^P - G(H^Q - E)^{-1}G^\dagger - E) C = 0$$

We can obtain the effective potential.

$$V_{eff} = V - G(H^Q - E)^{-1}G^\dagger$$

Separation 2 and 3 body force with PQ method

Projection operator P and Q

$$P\Psi = C\Psi_N$$

$$Q\Psi = D\Psi_\Delta$$

$$(H^P - G(H^Q - E)^{-1}G^\dagger - E)C = 0$$

$$G = \langle P|V|Q\rangle \quad G^\dagger = \langle Q|V|P\rangle$$

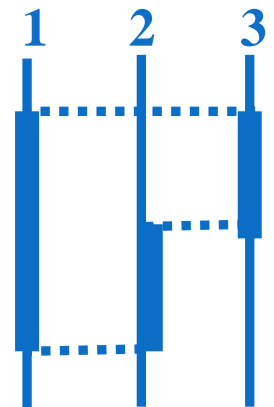
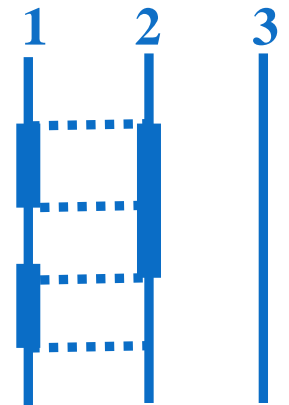
Two-body forces are constructed by 1-2 pair interaction

$$(H^P - G_{12}(H_{12}^Q - E)^{-1}G_{12}^\dagger - E)C = 0$$

$$G_{12} = \langle P|V_{12}|Q\rangle \quad G_{12}^\dagger = \langle Q|V_{12}|P\rangle \quad H_{12}^Q = \langle Q|H_{12}|Q\rangle$$

Three-body forces are constructed by 1-2 and other pairs interaction

$$(H^P - G_{12}(H_{23}^Q - E)^{-1}G_{13}^\dagger - E)C = 0$$



ダブルΔに効くテンソル力



質量差
約600MeV

$$\Psi_{NN} = |^3S_1\rangle + |^3D_1\rangle$$
$$\Psi_{\Delta\Delta} = |^3S_1\rangle + |^3D_1\rangle + |^7D_1\rangle + |^7G_1\rangle$$

NNの 3S_1 からダブルΔの 7D_1 に働く
テンソル力が大きい

テンソル力の行列要素から評価すると、

$$\langle ^7D_{\Delta\Delta} | v_{\pi} S_{12} | ^3S_{NN} \rangle \sim (f_{\pi N\Delta}^2 / 4\pi) M_S \sim 4.0$$

$$\langle ^3D_{\Delta\Delta} | v_{\pi} S_{12} | ^3S_{NN} \rangle \sim (f_{\pi N\Delta}^2 / 4\pi) M_S \sim 0.7$$

$$\langle ^3D_{\Delta\Delta} | v_{\pi} S_{12} | ^3S_{\Delta\Delta} \rangle \sim (f_{\pi\Delta\Delta}^2 / 4\pi) M_S \sim 0.1$$

結果 - ^1O チャンネル ---

J=1	Radius 1.70 fm
Energy [MeV]	25.6
Kinetic	23.25(NN=23.16)
Central	2.76(NN=2.79)
Tensor	-0.92
L·S	-----
L ²	-----
(L·S) ²	-----
P _{NN} [1P_1] %	99.93
P _{ΔΔ} [1P_1]	0.005
P _{ΔΔ} [5P_1]	0.032
P _{ΔΔ} [5F_1]	0.029

^1O チャンネル L=odd, S=even, T=0

$$\Psi_{NN} = |^1P_1\rangle$$

$$\Psi_{\Delta\Delta} = |^1P_1\rangle + |^5P_1\rangle + |^5F_1\rangle$$

ダブルΔとのテンソル力は、
ほとんど働かない。
3E(重陽子)で経験したように、
スピンは $[3/2 \times 3/2]=3$ の状態に
組むのがすき。2に組む状態は、
好まれない。

結果 - ^3O チャンネル ---

^3O チャンネル L=odd, S=odd, T=1

J=0	Radius 1.70 fm	N Δ なし
Energy [MeV]	20.1	20.4
Kinetic	23.18(NN=22.99)	23.19 (NN=23.07)
Central	1.27(NN=1.31)	1.34 (NN=1.32)
Tensor	-6.87 (NN=-5.41)	-6.26 (NN=-5.31)
L·S	1.12	1.06
L ²	-0.07	-0.07
(L·S) ²	0.82	0.77
P _{NN} [$^3\text{P}_0$] %	99.85	99.94
P _{$\Delta\Delta$} [$^3\text{P}_0$]	0.0004	0.0009
P _{$\Delta\Delta$} [$^7\text{F}_0$]	0.06	0.05
P _{NΔ} [$^3\text{P}_0$]	0.08	-----

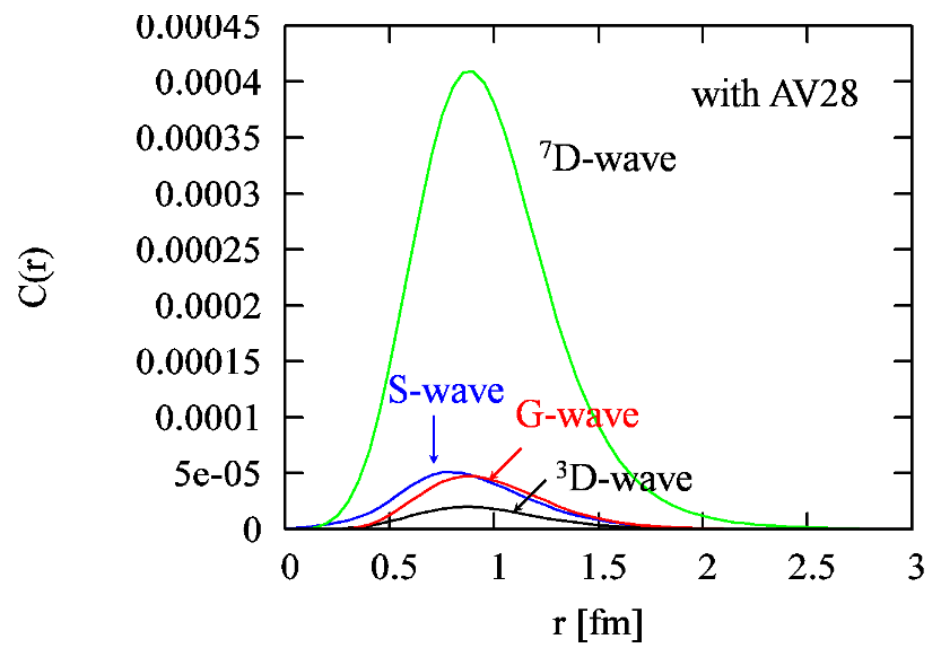
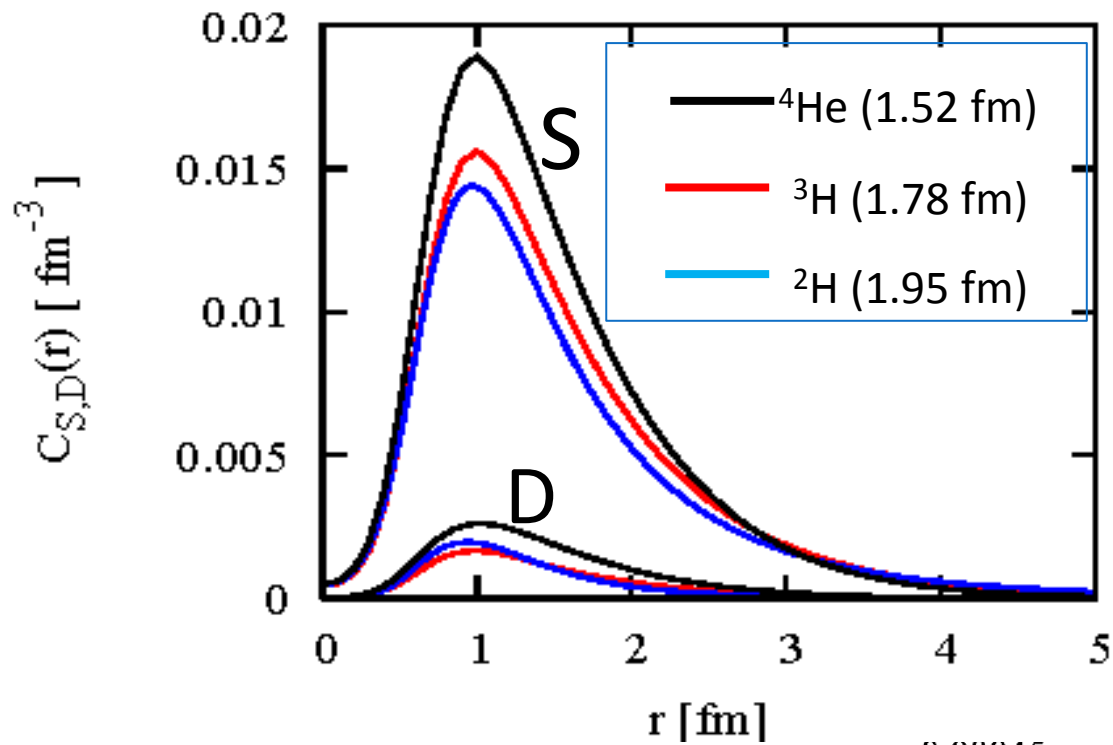
$$\Psi_{NN} = |^3P_0\rangle$$

$$\Psi_{\Delta\Delta} = |^3P_0\rangle + |^7F_0\rangle$$

$$\Psi_{N\Delta} = |^3P_0\rangle$$

NNの $^3\text{P}_0$ からダブル Δ の $^7\text{F}_0$ に働くテンソル力が大きい

J=1,2の状態には、
テンソル力が効くような
チャンネルはなかった



実験にみえる三体力の効果

青線・2体力 赤線・2体力+3体力

