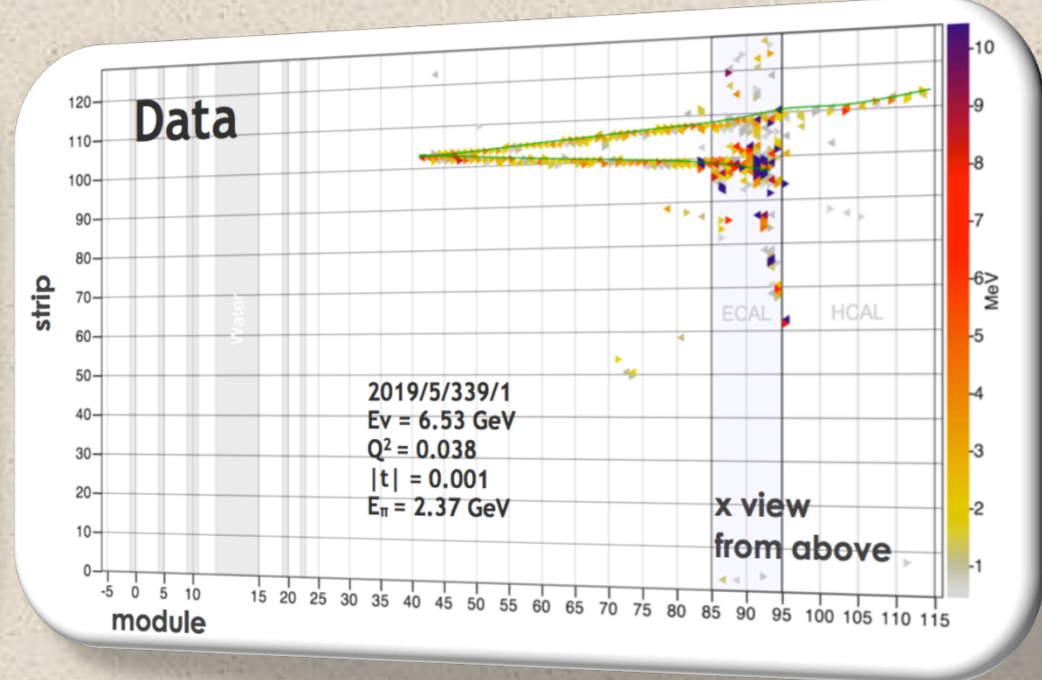
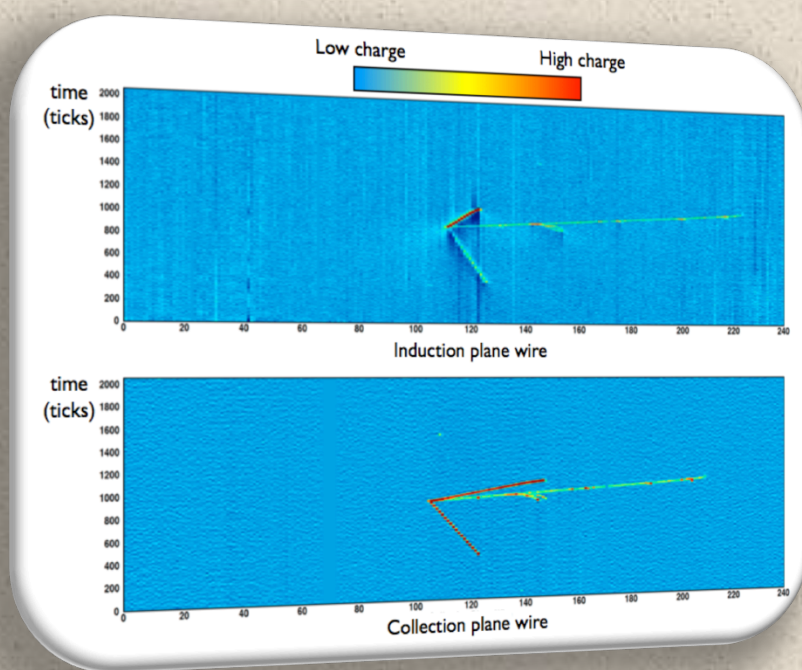


Experimental Overview of Neutrino-Nucleus Scattering



ν_μ CC candidate in ArgoNeuT

ν_μ CC candidate in MINERvA

NuFact 2013: International Workshop on Neutrino Factories, Super Beams and Beta Beams
23 August 2013, IHEP, Beijing, China

David Schmitz, University of Chicago

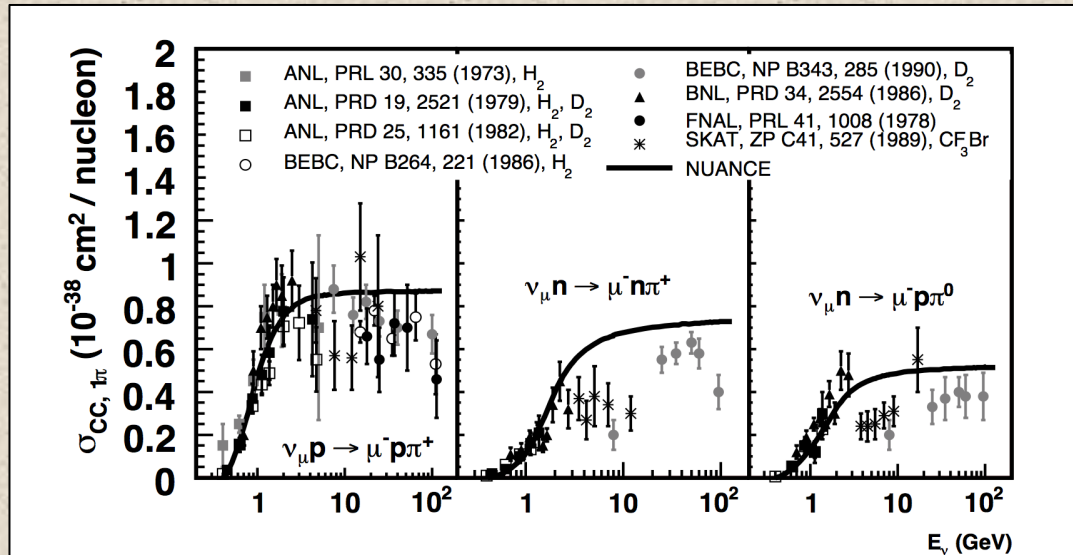
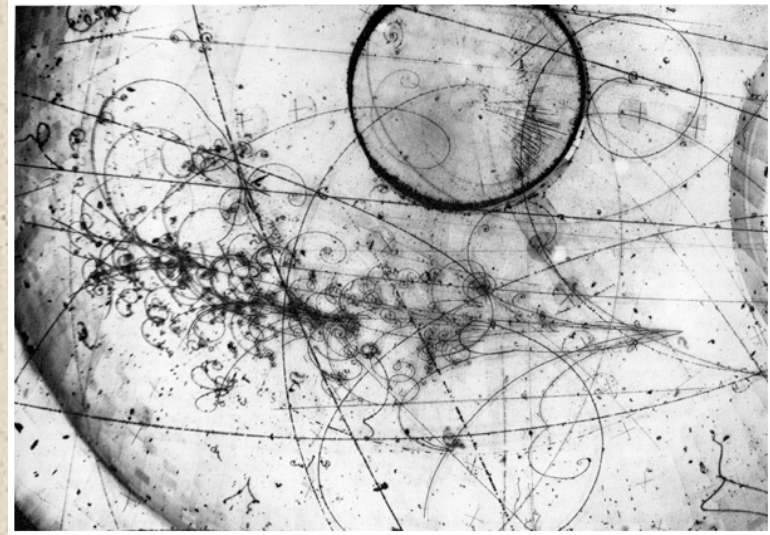
Outline

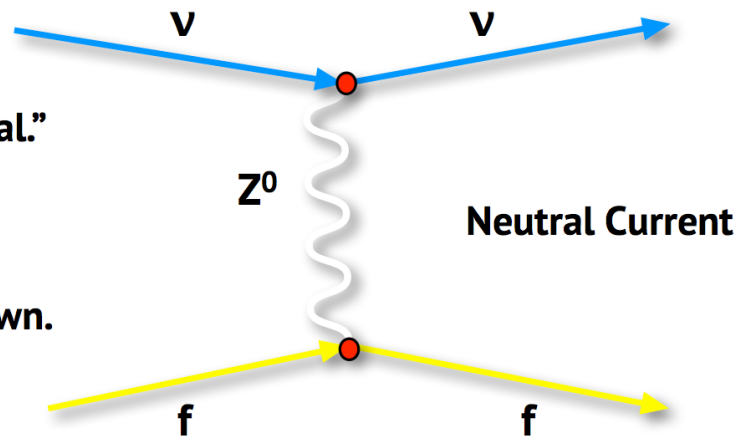
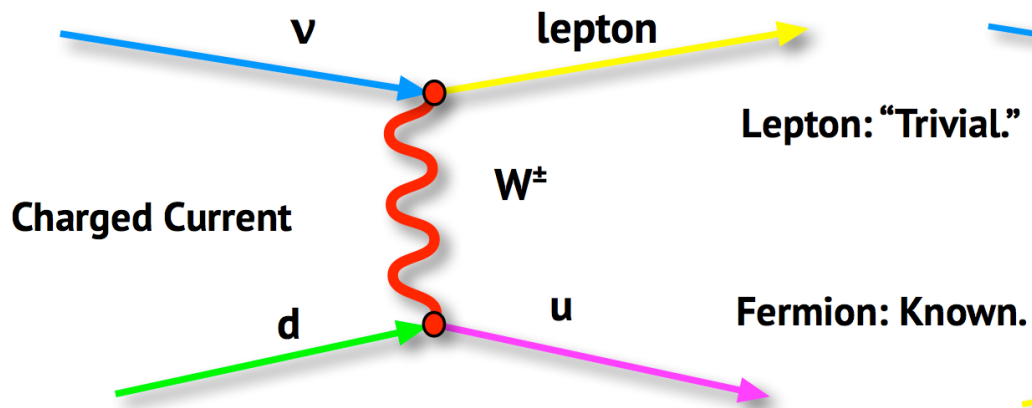
- Why do we need to measure neutrino-nucleus scattering?
- Recent measurements from:
 - MiniBooNE
 - MINERvA
 - ArgoNeuT
 - T2K
- Future possibilities

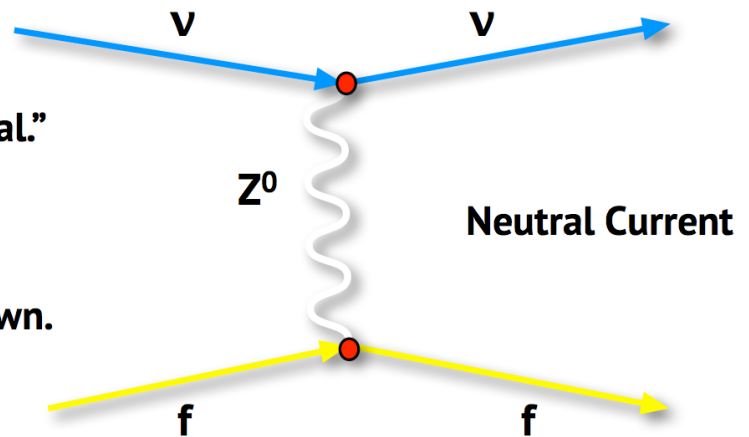
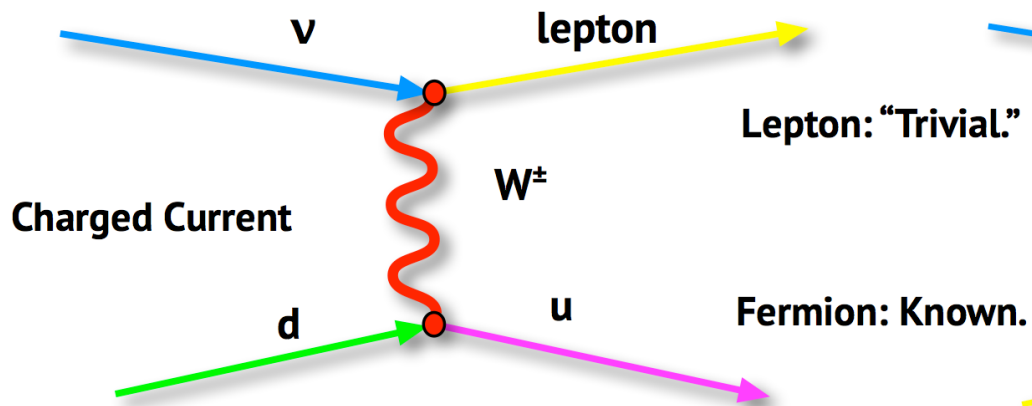
First, a Question

Physicists have been scattering neutrinos off nuclei and studying what happens for decades

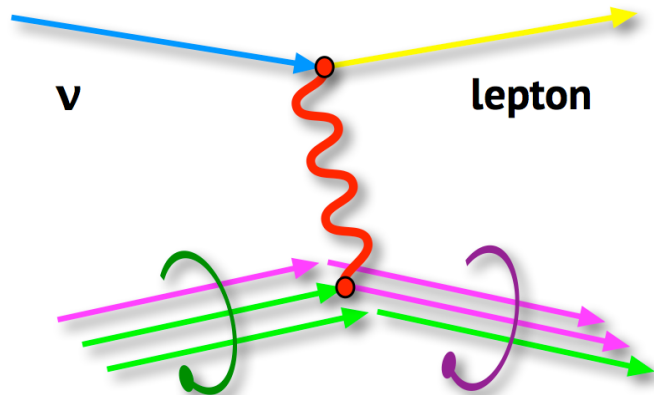
Why are we still doing this!?!

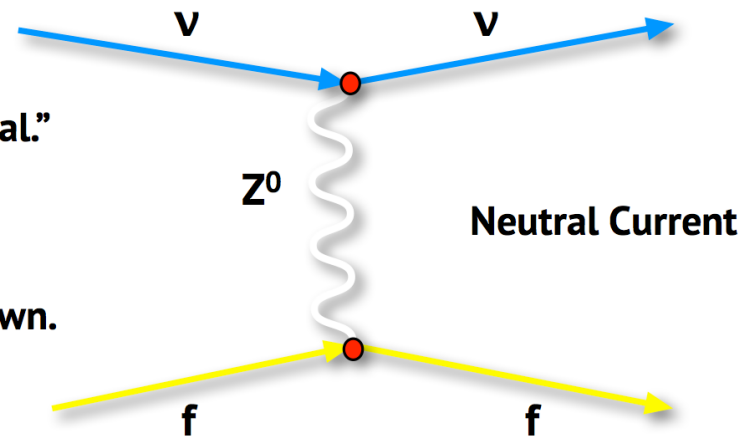
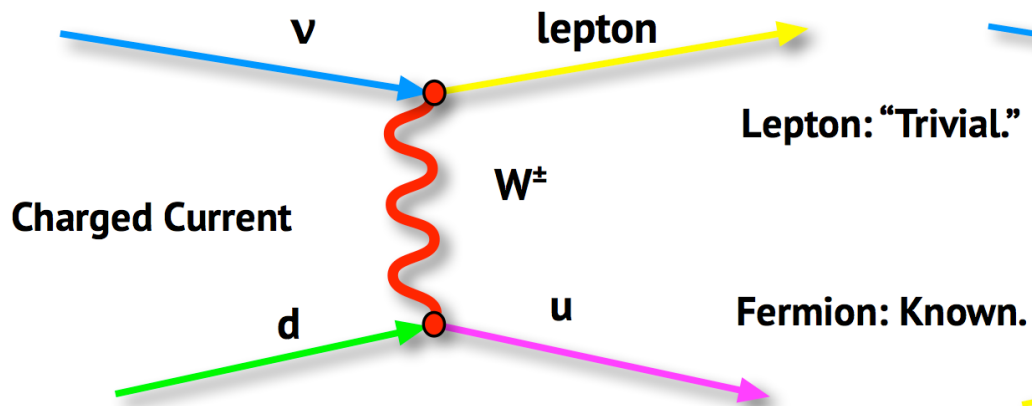




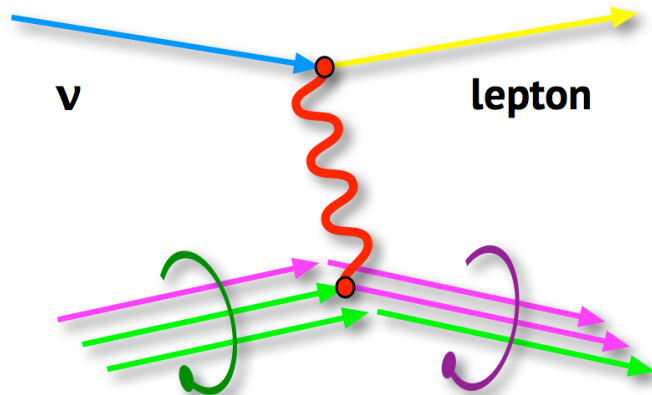


Nucleon: Parameterize
w/ Form Factors.

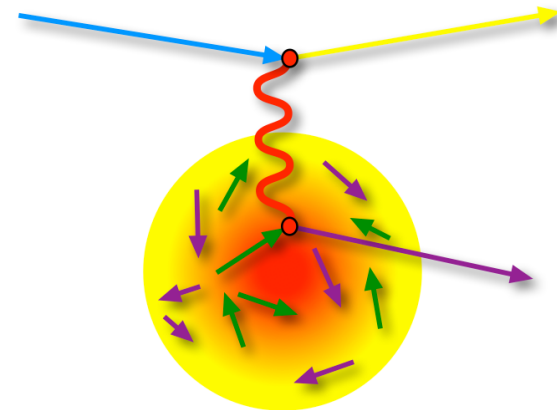




**Nucleon: Parameterize
w/ Form Factors.**



Nucleus: Hard!
Very complex nuclear physics.
But this is where we want σ .

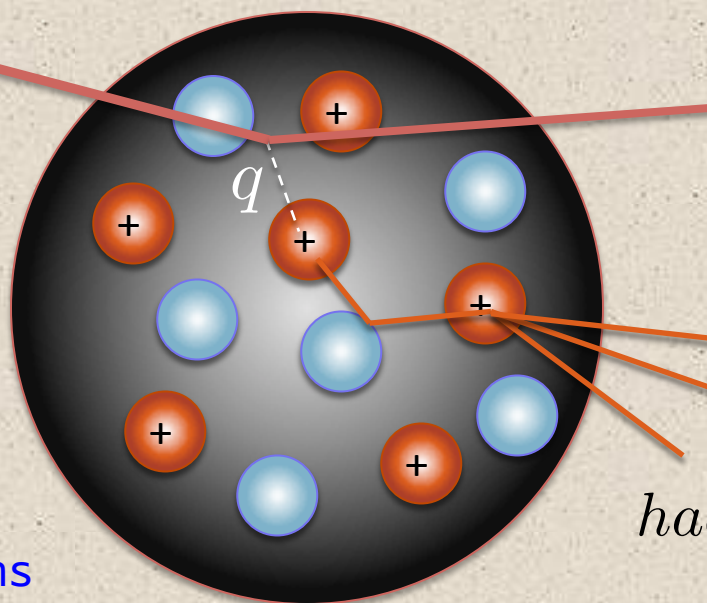


Physics of GeV ν -N Interactions

- 1 Cross section models for all exclusive ν -nucleon interaction channels (elastics, resonance productions, DIS ...)

ν

lepton



- 2 Models of nucleons within the nucleus (Relativist Fermi Gas, spectral functions, nucleon correlations, ...)

- 3 Final state interaction models which alter the hadronic final state (rescattering, absorption, charge exchange, ...)

hadrons

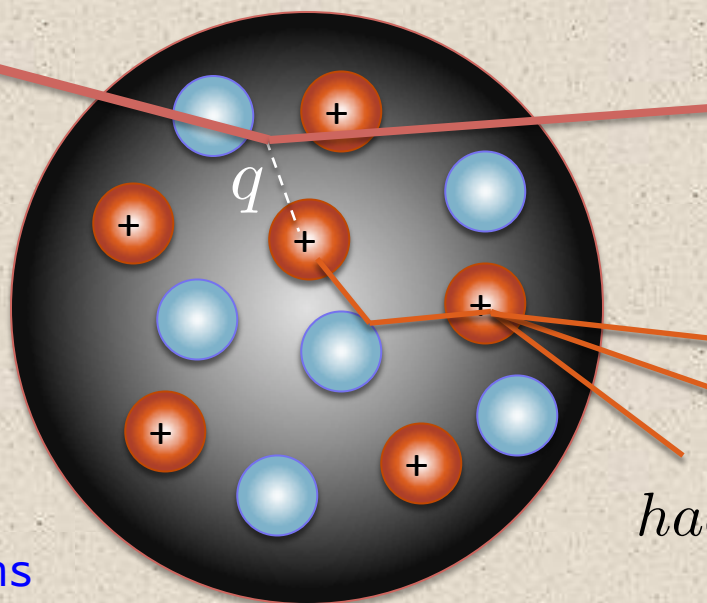
Physics of GeV ν -N Interactions

- 1 Cross section models for all exclusive ν -nucleon interaction channels (elastics, resonance productions, DIS ...)

ν

lepton

E_{ν}^{true}



$E_{visible}$

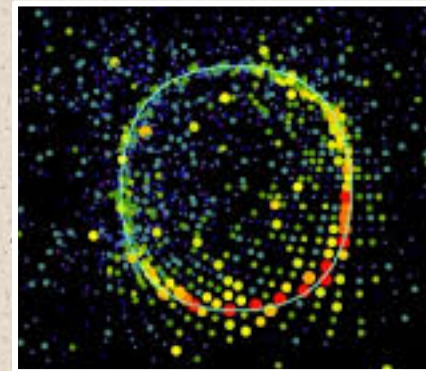
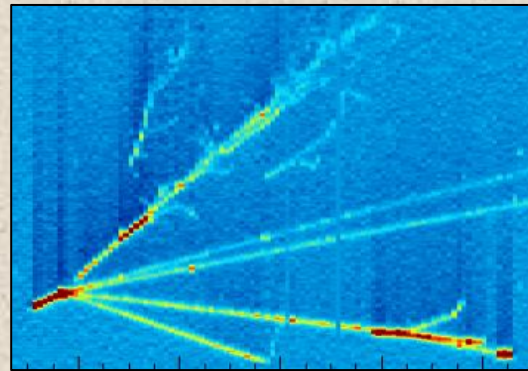
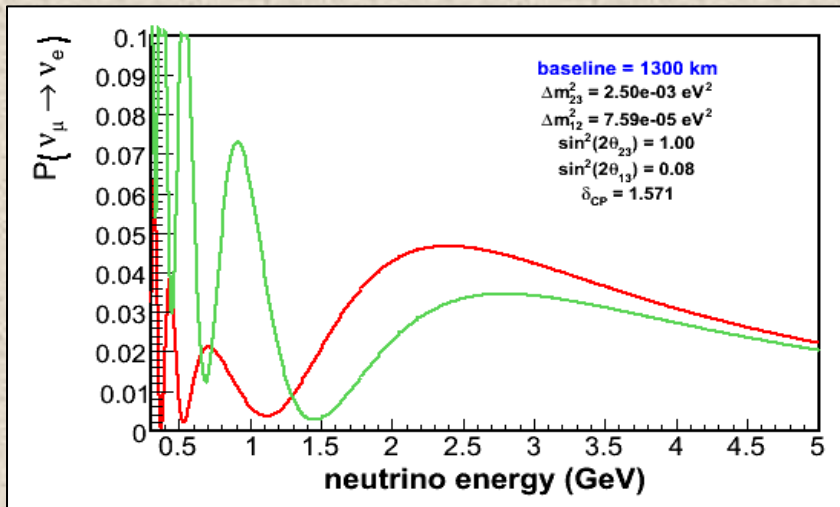
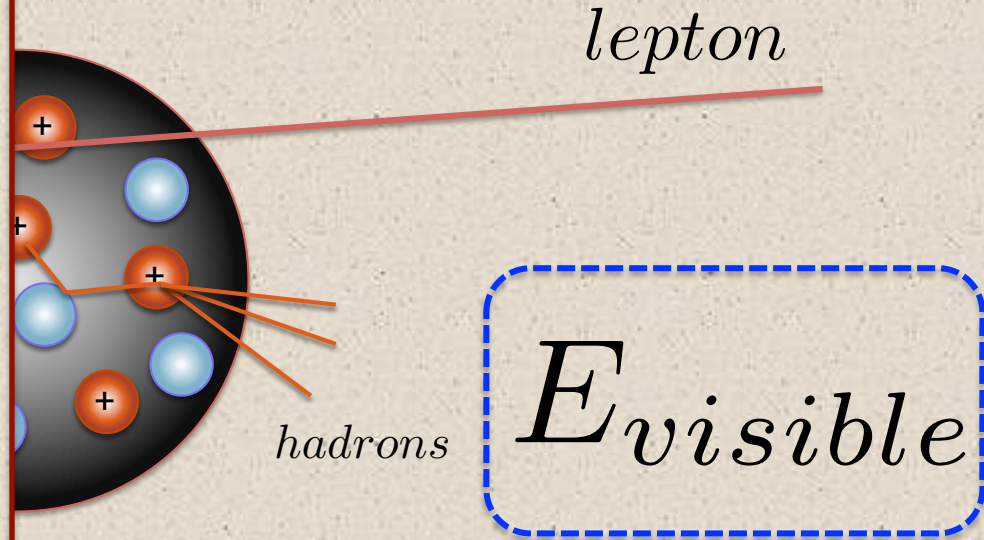
hadrons

- 2 Models of nucleons within the nucleus (Relativist Fermi Gas, spectral functions, nucleon correlations, ...)

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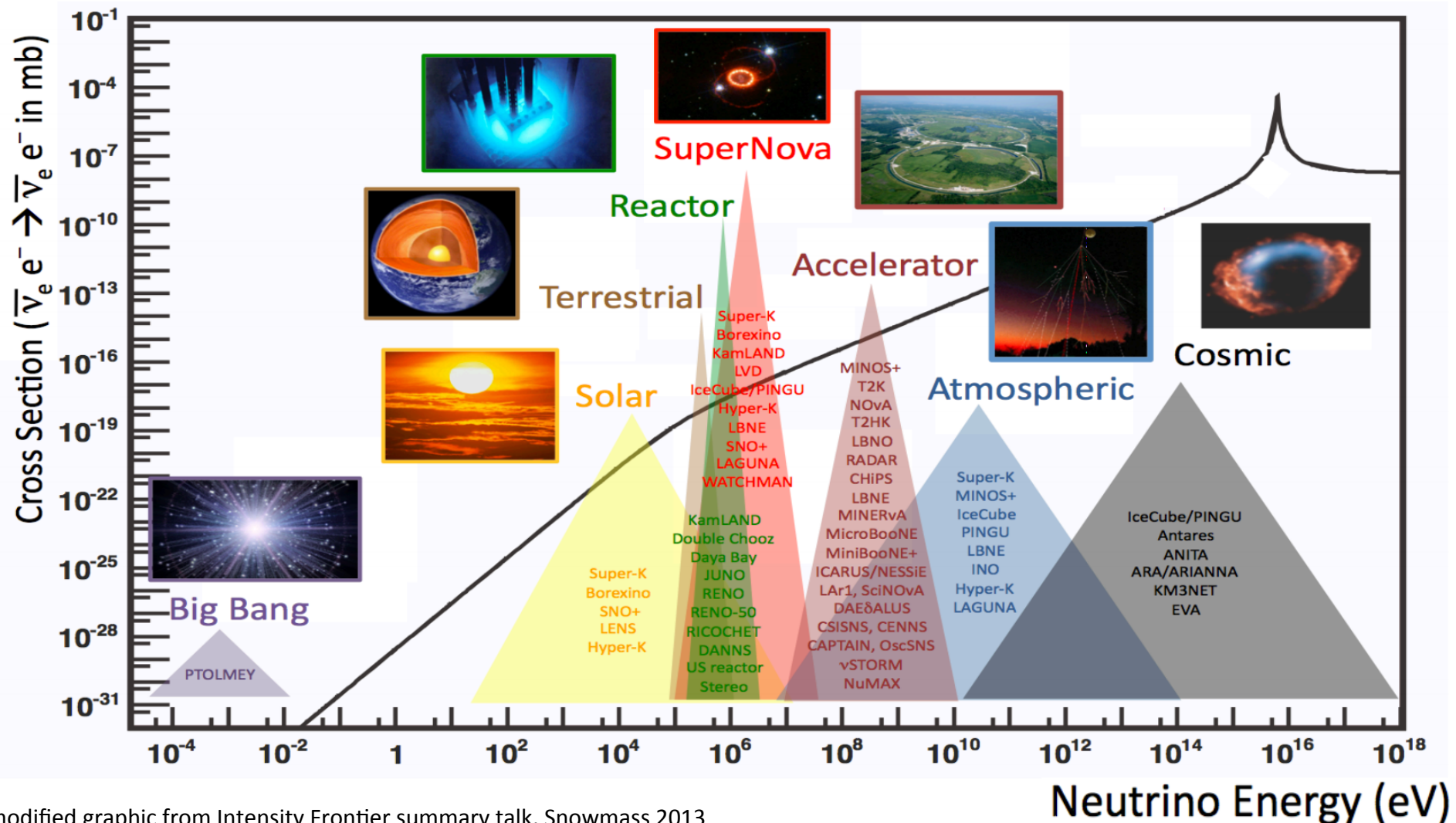
Physics of GeV ν -N Interactions

$$E_{\nu}^{true}$$



Neutrino Sources

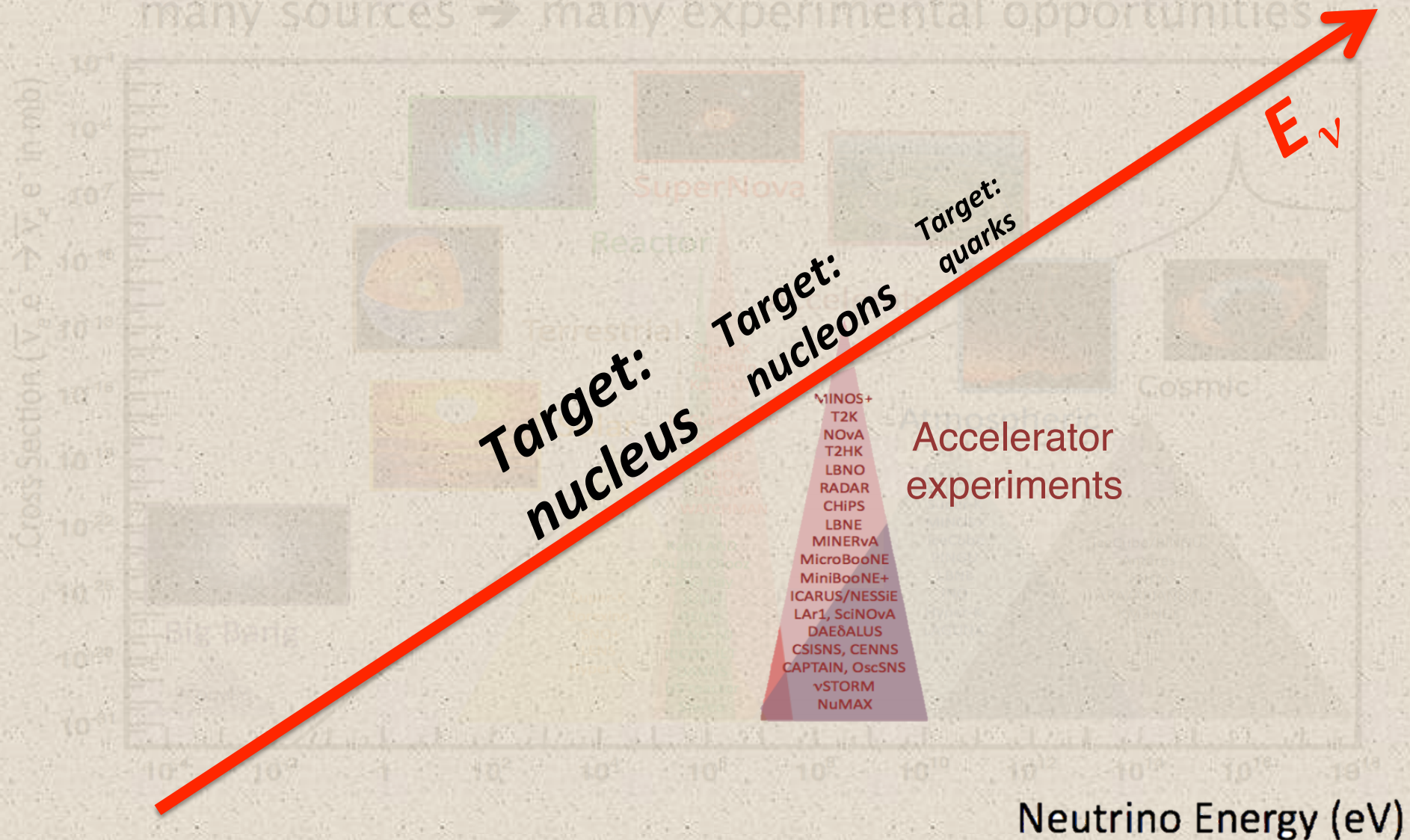
many sources → many experimental opportunities



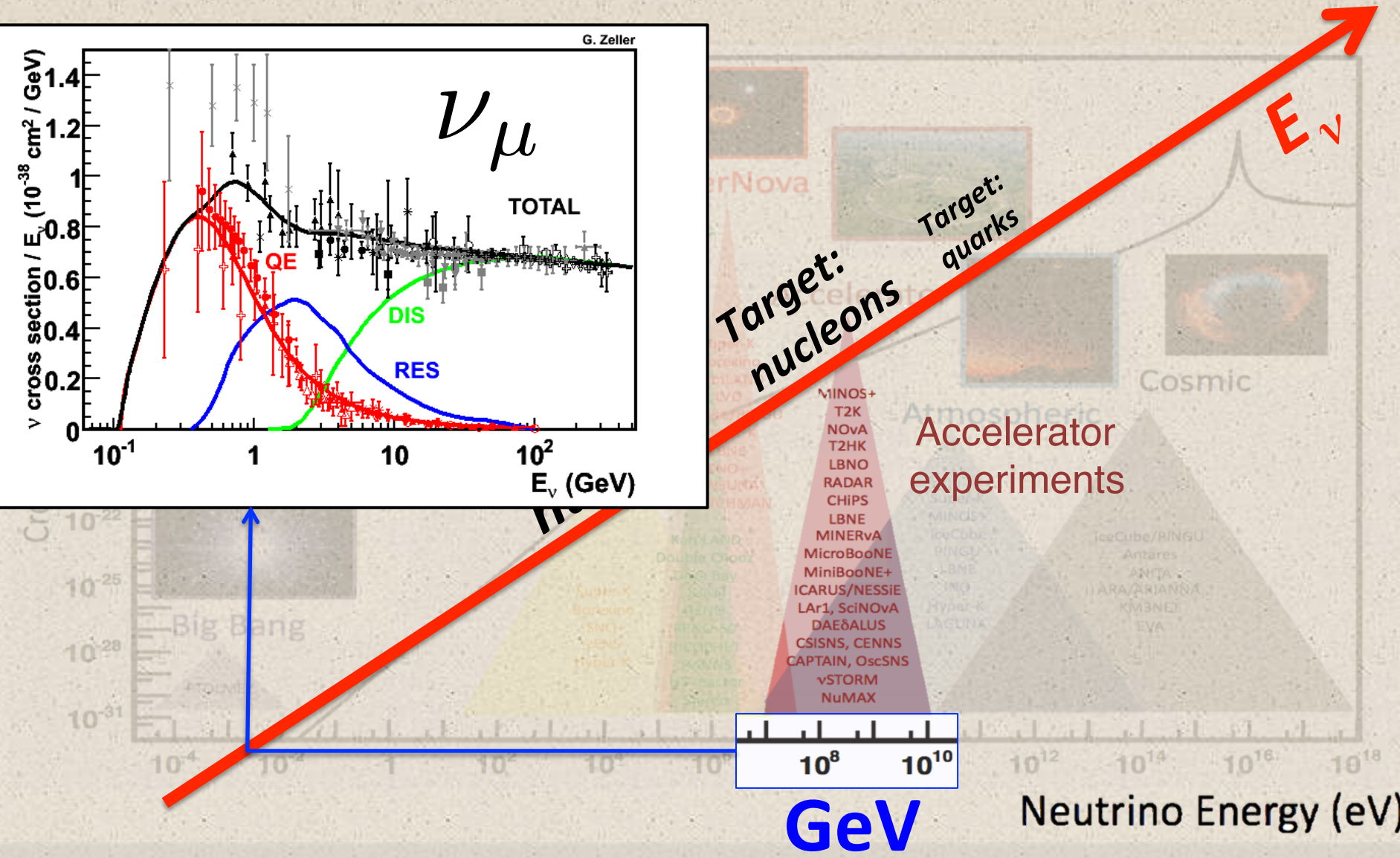
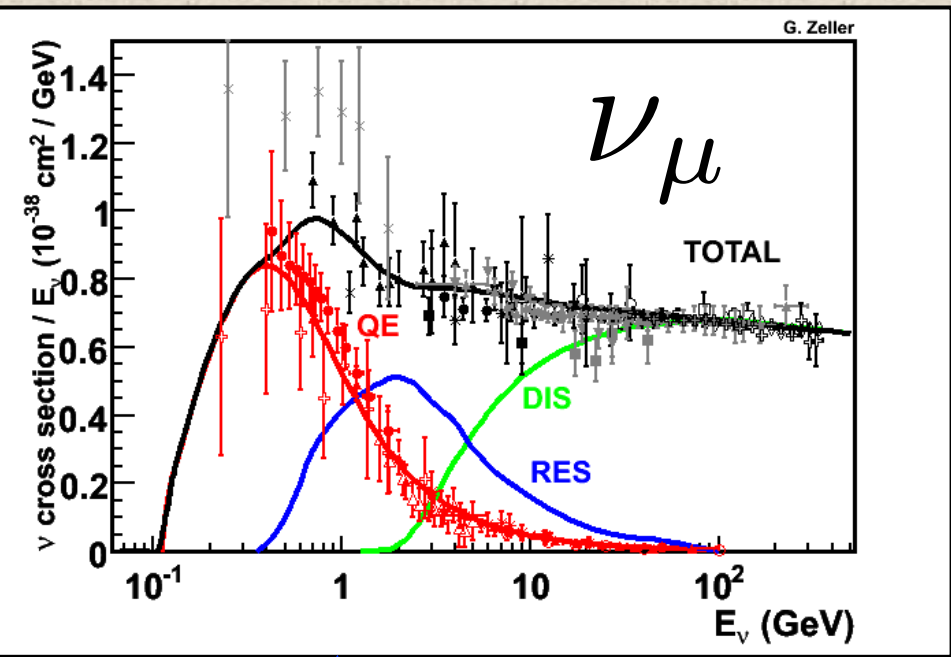
modified graphic from Intensity Frontier summary talk, Snowmass 2013
original from J.A. Formaggio and G.P. Zeller, Rev. Mod. Phys. 84, 1307-1341, 2012

Neutrino Targets

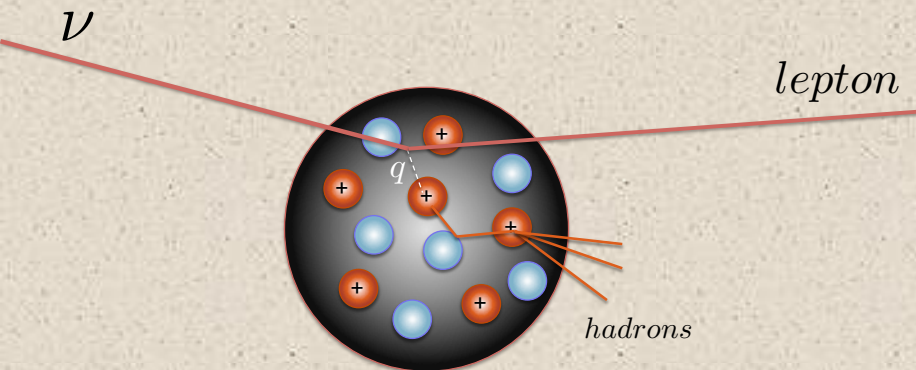
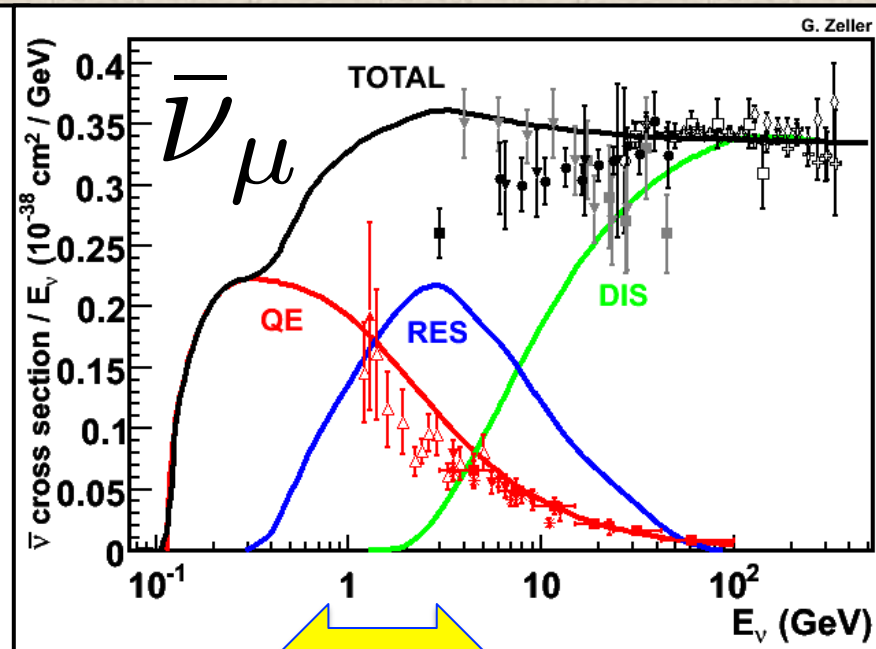
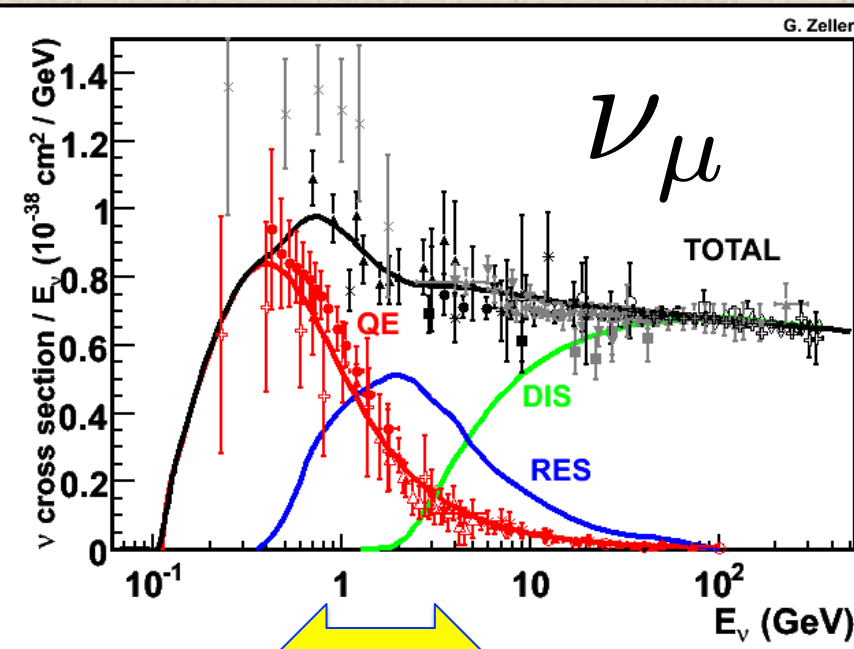
many sources \rightarrow many experimental opportunities



- QE** – charged-current quasi-elastic: $\nu n \rightarrow \ell^- p$ $\bar{\nu} p \rightarrow \ell^+ n$
- RES** – excite a nucleon resonance, decays to pions & nucleons
- DIS** – scatter off quarks instead of nucleon, high multiplicity



- QE** – charged-current quasi-elastic: $\nu n \rightarrow \ell^- p$ $\bar{\nu} p \rightarrow \ell^+ n$
- RES** – excite a nucleon resonance, decays to pions & nucleons
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Accelerator-based experiments
carried out in a sweet spot
of maximally
complicated nuclear effects
in neutrino interactions!

For more quantitative look at
impact, see Pilar Coloma's
talk later in this session

In Summary: Nuclear Physics Meets Neutrino Physics



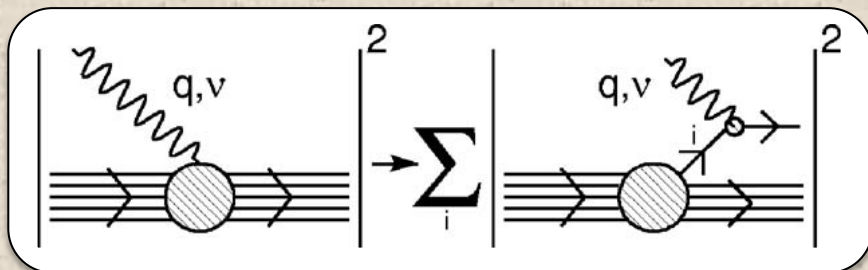
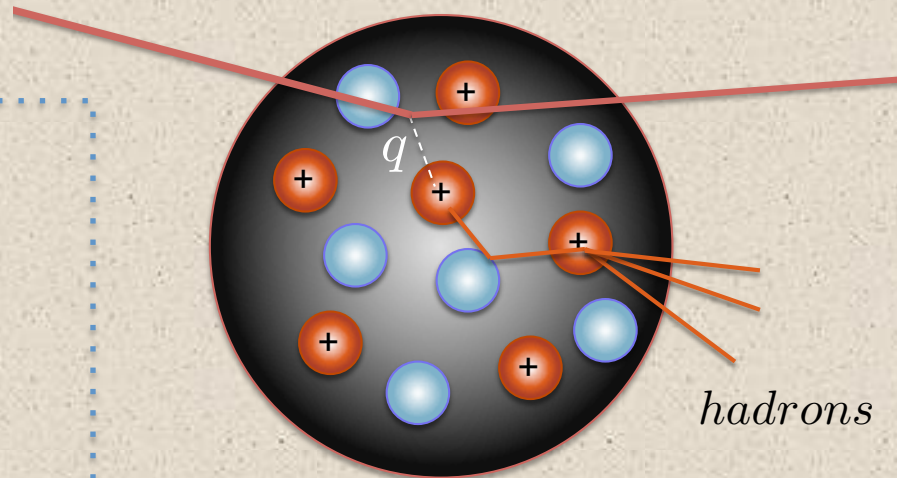
In Summary: Nuclear Physics Meets Neutrino Physics

The results I will highlight today are steps toward understanding how the nuclear environment alters the models of neutrino scattering we have been using for decades.

CC inclusive, CC quasi-elastic, studies of final state nucleons

Relativistic Fermi Gas For Nucleus

- Nuclear model in neutrino scattering event generators same since 1970s
- In the *Impulse Approximation*, scatter off independent single nucleons incoherently summed over all nucleons in the nucleus



NEUTRINO REACTIONS ON NUCLEAR TARGETS.[‡]

R. A. SMITH^{‡‡} and E. J. MONIZ^{‡‡‡}

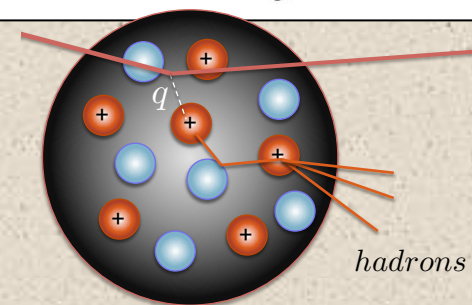
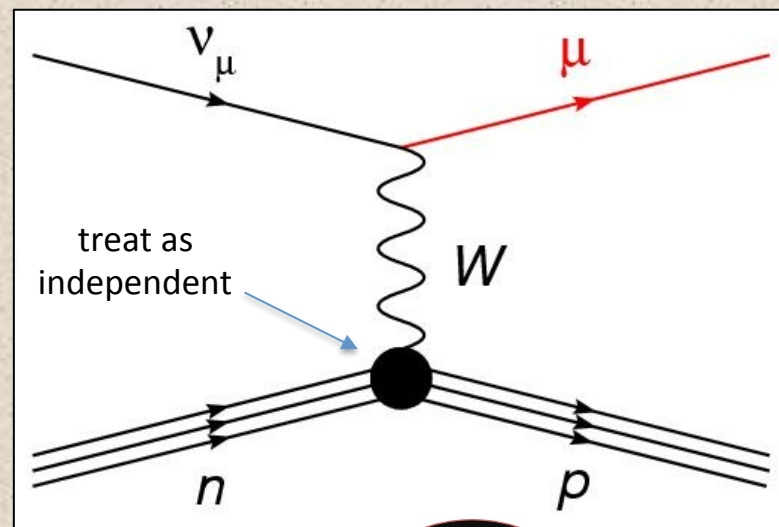
*Institute of Theoretical Physics, Department of Physics,
Stanford University, Stanford, California 94305*

with energies lowered from free particle energies by an average nuclear potential. Although this choice of wave functions corresponds to a nucleus without detailed structure, the dominant features of the nuclear cross section are consistently represented.

Smith, R. A., and E. J. Moniz, 1972, Nucl. Phys. B43, 605.

Relativistic Fermi Gas For Nucleus

- For quasi-elastic scattering, if we further assume the *nucleon is at rest*, we can determine E_ν and Q^2 from lepton kinematics only (“2-body interaction”)
 - Technique used by many oscillation experiments, particularly when blind to the hadronic final state



neutrino energy

$$E_\nu^{QE} = \frac{2(M_n - E_B) E_\ell - \left[(M_n - E_B)^2 + m_\ell^2 - M_p^2 \right]}{2[M_n - E_B - E_\ell + p_\ell \cos(\theta_\ell)]}$$

$$Q_{QE}^2 = -m_\ell^2 + 2E_\nu^{QE} \left(E_\ell - \sqrt{E_\ell^2 - m_\ell^2} \cos(\theta_\ell) \right)$$

4-momentum transferred

M_n = neutron mass
 M_p = proton mass
 E_B = separation energy
 m_ℓ = lepton mass
 E_ℓ, θ_ℓ = lepton energy and angle

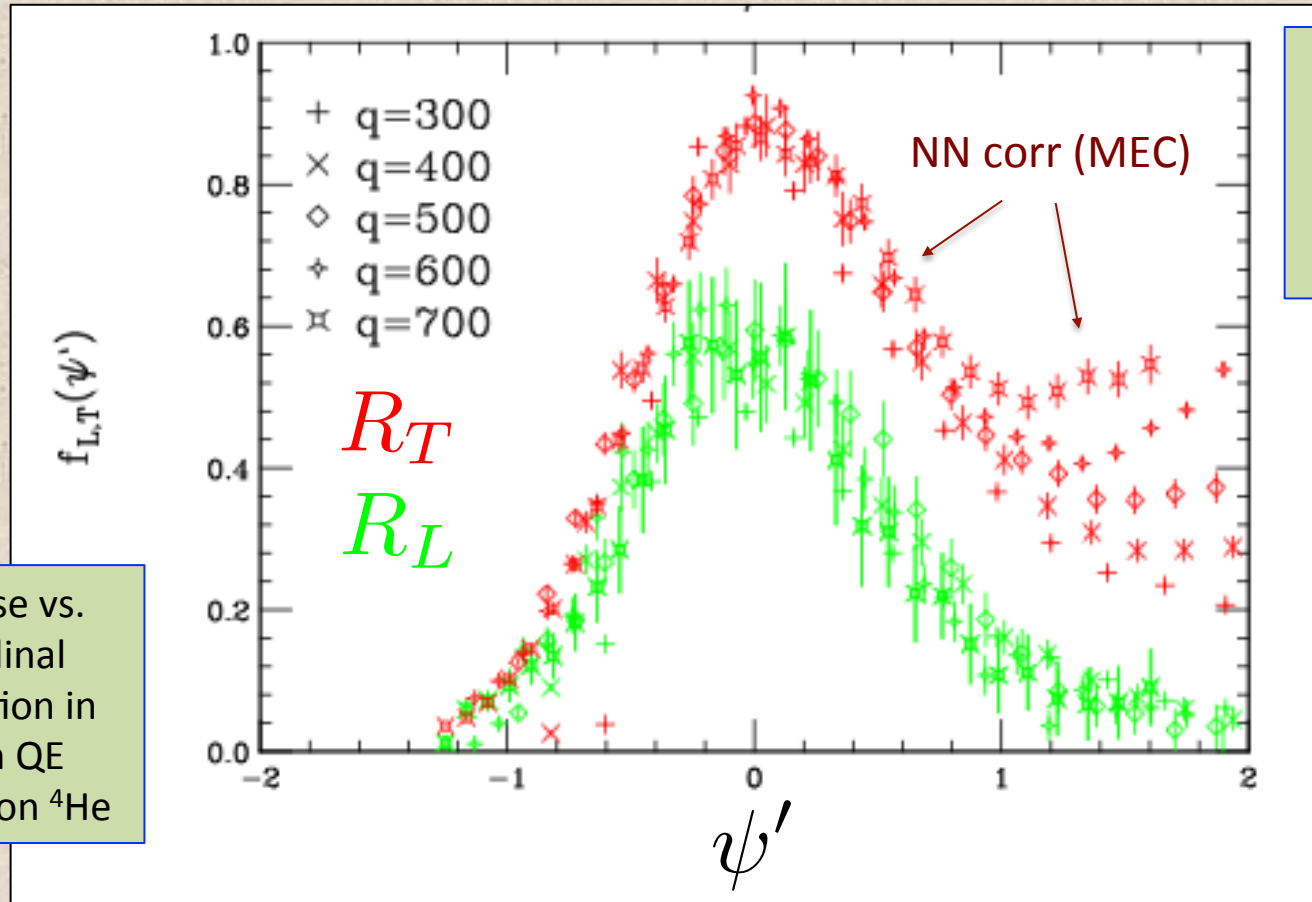
But Is It Good Enough?

- Both theory and experimental hints tell us that it is not
- Nucleon-nucleon interactions and 2-body currents are ignored
 - RPA effects ('nucleus as a whole') expected to suppress rate at low Q^2
 - NN correlations such as *meson exchange currents (MEC)* which may enhance part of the cross section significantly
 - Correlated nucleon(s) may be ejected when partner is scattered by a probe
 - This change on the hadronic side of the interaction impacts the kinematics and spoils your neutrino energy estimation
 - For $A \geq 12$ *short range correlations (SRC)* affect ~20% of nucleons
 - These can lead to nucleon momentum well above the Fermi sea cutoff
 - Again, multi-nucleon emission and impacted energy reconstruction are consequences

See Jan Sobczyk's theory summary coming up next

Transverse and Longitudinal Strength

- Impact of correlations seen in *electron scattering* data



$R_L = R_T$
for RFG of
independent
nucleons!

J. Carlson, et al., PRC 65, 024002 (2002)

Transverse vs.
longitudinal
cross-section in
electron QE
scattering on ^4He

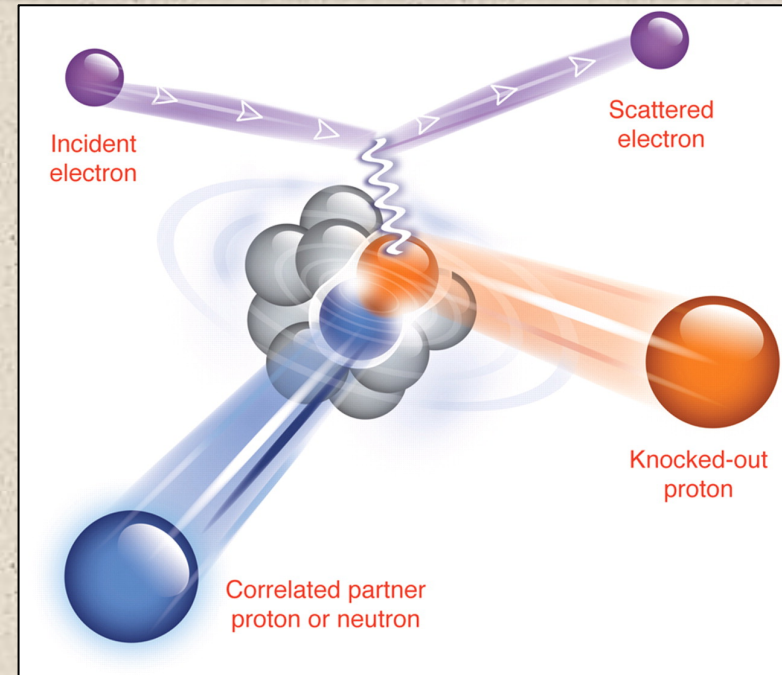
$$\psi' = \frac{\sqrt{\omega^2 + 2M\omega} - q}{k_F}$$

ω = energy transfer to nucleus

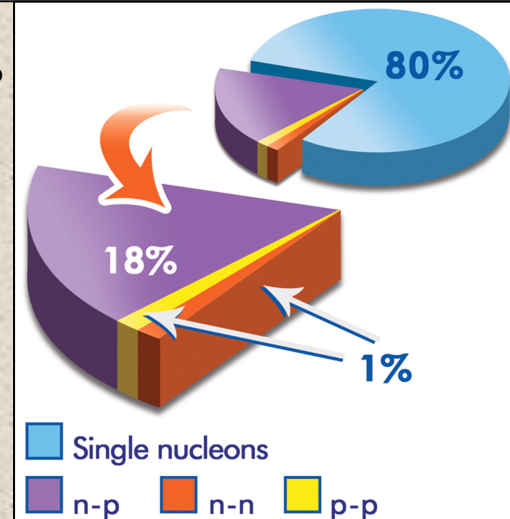
q = 3-momentum transfer to nucleus

Short Range NN Correlations

- Electron scattering
 - Measurements on ^{12}C indicate 20% correlated nucleons with mostly **np** pairs in the initial state
- Neutrino scattering
 - Implies final states in neutrino scattering of **nn** in **antineutrino** and **pp** in **neutrino** CC scattering
 - For other forms of correlation, final state depends on model
 - Of course, strongly coupled to Final State Interactions when interpreting neutrino scattering data

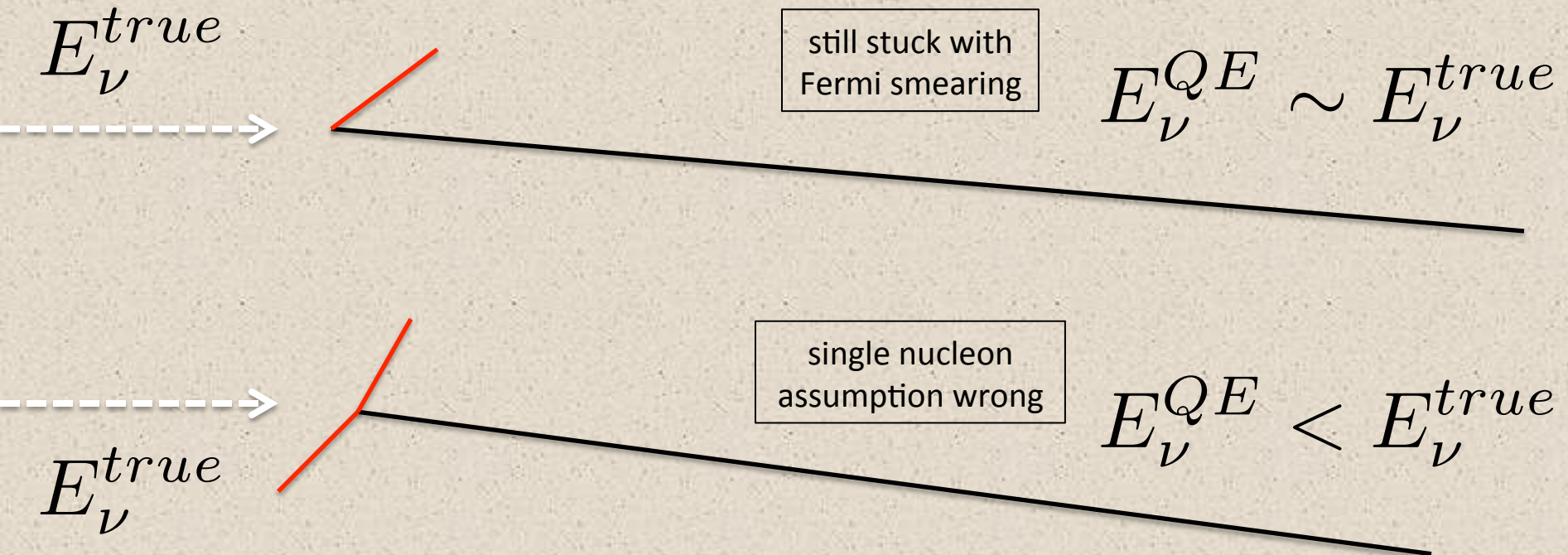


R. Subedi et al.,
Science 320, 1476
(2008)



Kinematic

$$E_{\nu}^{QE} = \frac{2(M_n - E_B) E_{\ell} - \left[(M_n - E_B)^2 + m_{\ell}^2 - M_p^2 \right]}{2[M_n - E_B - E_{\ell} + p_{\ell} \cos(\theta_{\ell})]}$$



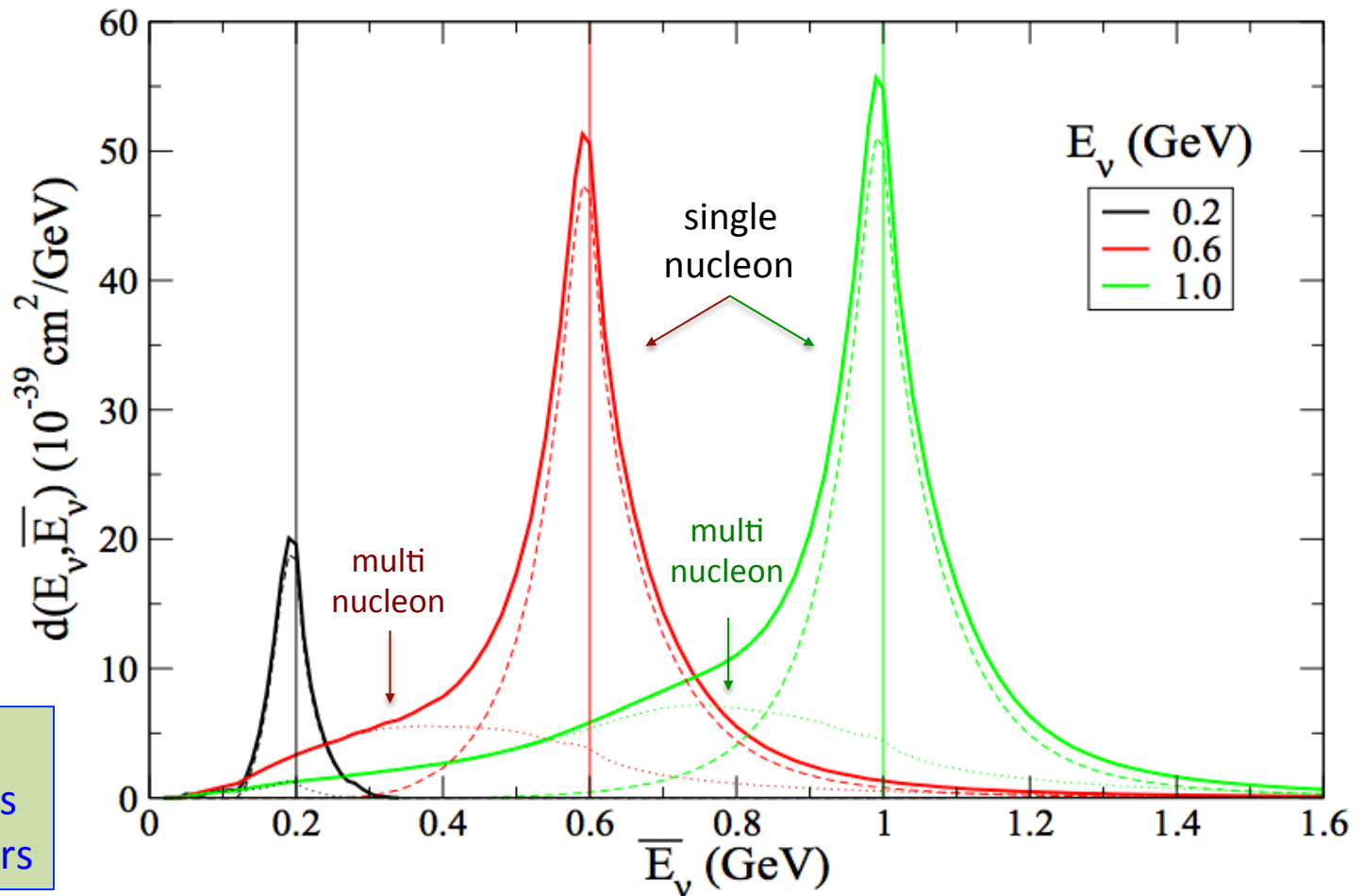
Calorimetric

$$E_{\mu} + E_{visible}^{had} \leq E_{\nu}^{true}$$

invisible
component

Impact For Neutrino Experiments

Kinematic



Recall, up to
20-30% of events
off correlated pairs

Martini et al. arXiv:1211.1523

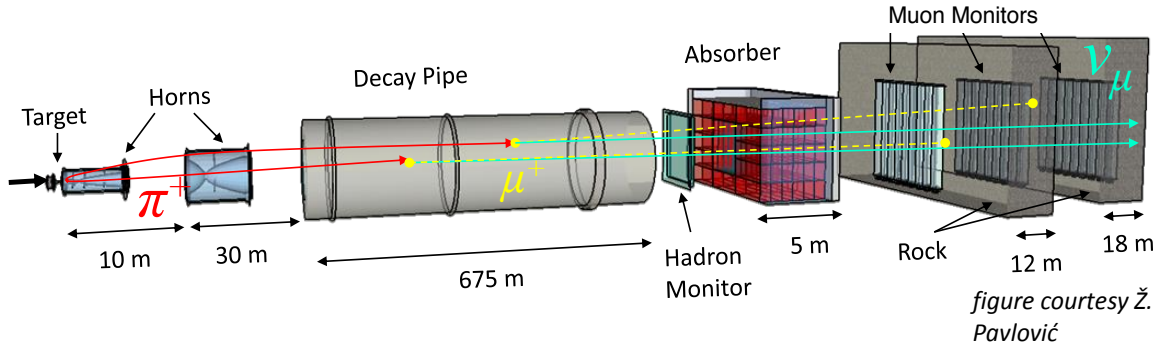
Also: J. Sobczyk arXiv:1201.3673,
Lalakulich et al. arXiv:1208.3678,
Nieves et al. arXiv:1204.5404

Elements of a ν -N Program

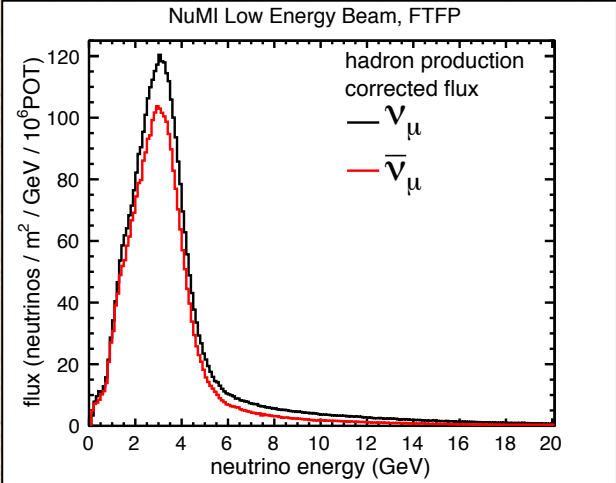
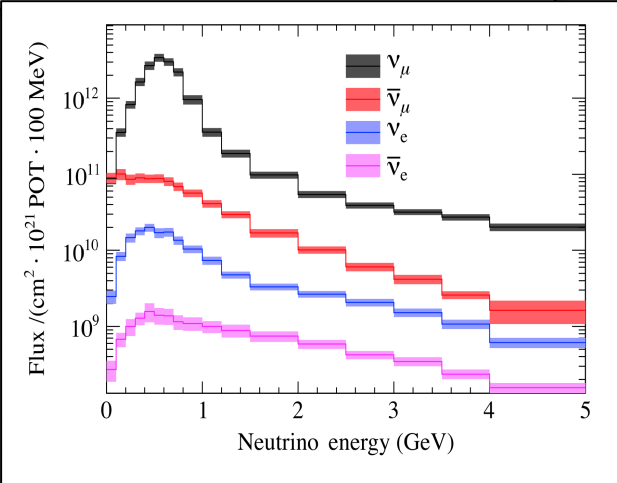
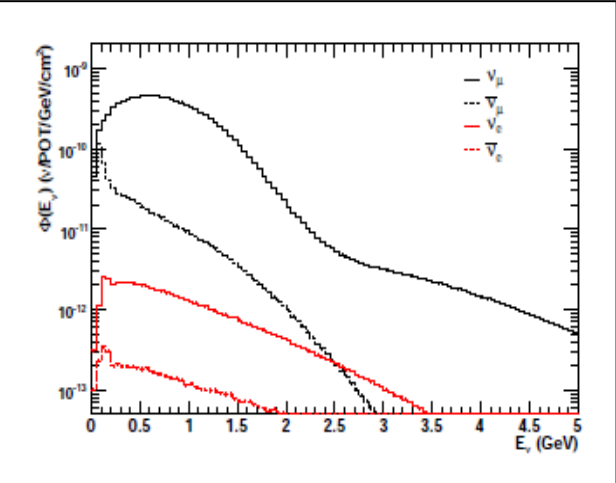
1. Span of neutrino energies (~ 100 MeV to 10 GeV)
 - With minimized flux uncertainties (spectrum and normalization)
2. Range of nuclear targets
3. High resolution detectors
 - Good resolution of leptonic and hadronic sides of the final state
4. Differential cross sections \rightarrow statistics
 - Required to untangle underlying physics and validate models
5. Close collaboration with theoretical community
 - Much of this physics is at the cross roads of particle and nuclear
 - Improvement of event generators is key to utilizing in osc. experiments

New Results in 2013

- **MiniBooNE:** antineutrino charged-current quasi-elastic differential cross sections on carbon
PRD 88, 032001 (2013)
- **MINERvA:** neutrino and antineutrino charged-current quasi-elastic $d\sigma/dQ^2$ on carbon
PRL 111, 022501 (2013)
PRL 111, 022502 (2013)
- **ArgoNeuT:** detailed studies of proton multiplicities in charged-current interactions on argon
PRL 108, 161802 (2012)
+ preliminary proton data
- **T2K ND280:** neutrino charged-current inclusive differential cross sections on carbon
PRD 87, 092003 (2013)



Flux predictions a tough problem for all super beam facilities.
Heroic progress, really.



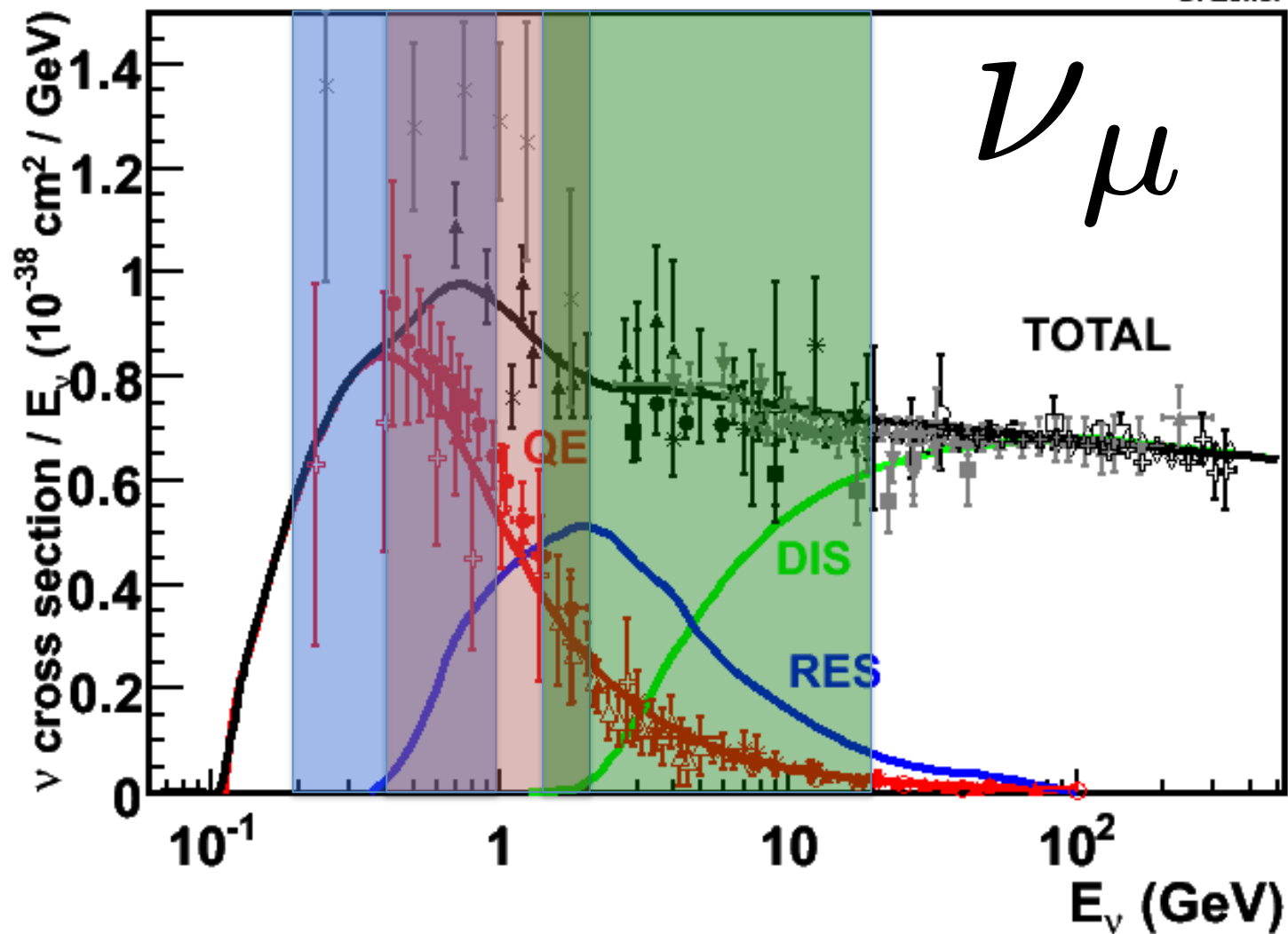
BooNE	T2K	NuMI
8 GeV p+beryllium	30 GeV p+carbon	120 GeV p+carbon
HARP π prod.	NA61 $\pi/K/p$ prod.	NA49 158 GeV p prod.
2 nd interactions < 10%	2 nd interactions ~ 35%	2 nd interaction > 40%
		configurable beamline
≥ 9% uncertainty	10-15% uncertainty	10-15% uncertainty

T2K

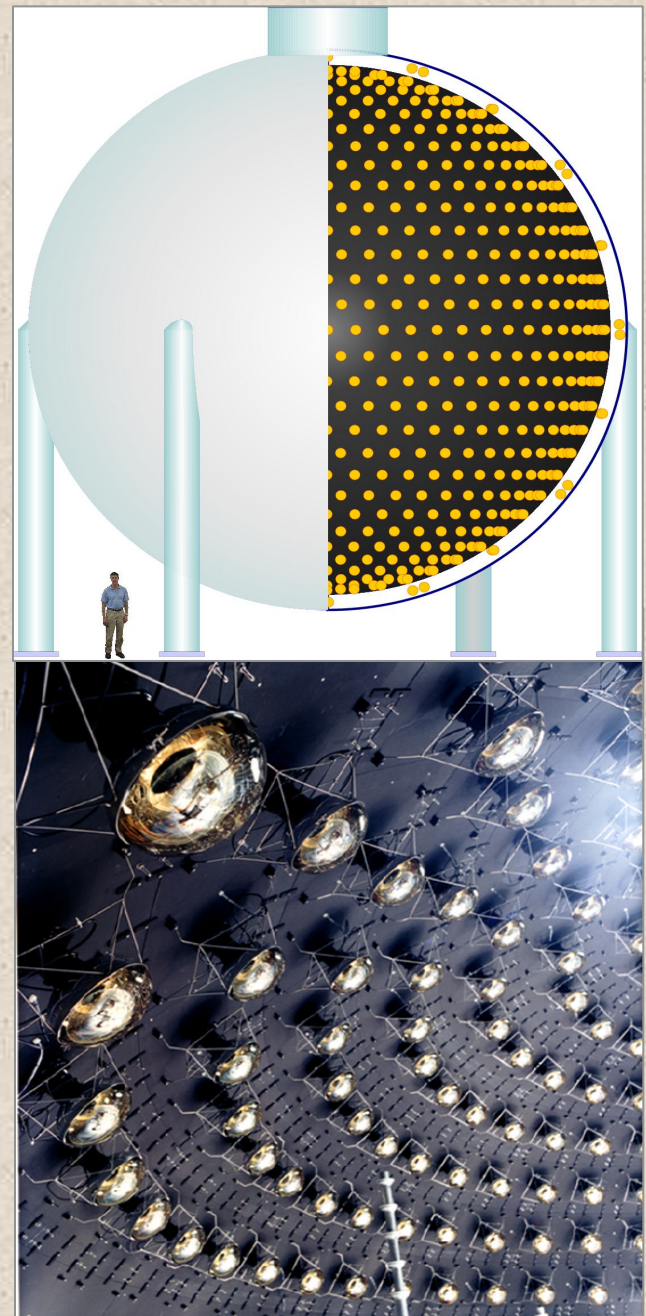
BooNE

NuMI

G. Zeller



- CH₂ target
- 4 π detector, complete angular coverage
- Good lepton reconstruction & pion rejection
- Essentially blind to details of the nucleon final state in CC events



- MiniBooNE has been a prolific producer of neutrino-carbon cross section measurements in the ≤ 1 GeV region
 - CCQE MA – PRL 100, 032301 (2008)
 - NC π^0 – PL B664, 41 (2008)
 - CC π^+ /QE – PRL 103, 081801 (2009)
 - NC π^0 – PRD 81, 013005 (2010)
 - QE – PRD 81, 092005 (2010)
 - NC elastic – PRD 82, 092005 (2010)
 - CC π^0 – PRD 83, 052009 (2011)
 - CC π^+ – PRD 83, 052007 (2011)
 - **Antineutrino QE – PRD 88, 032001 (2013)**
 - CC inclusive and antineutrino NC elastic coming soon

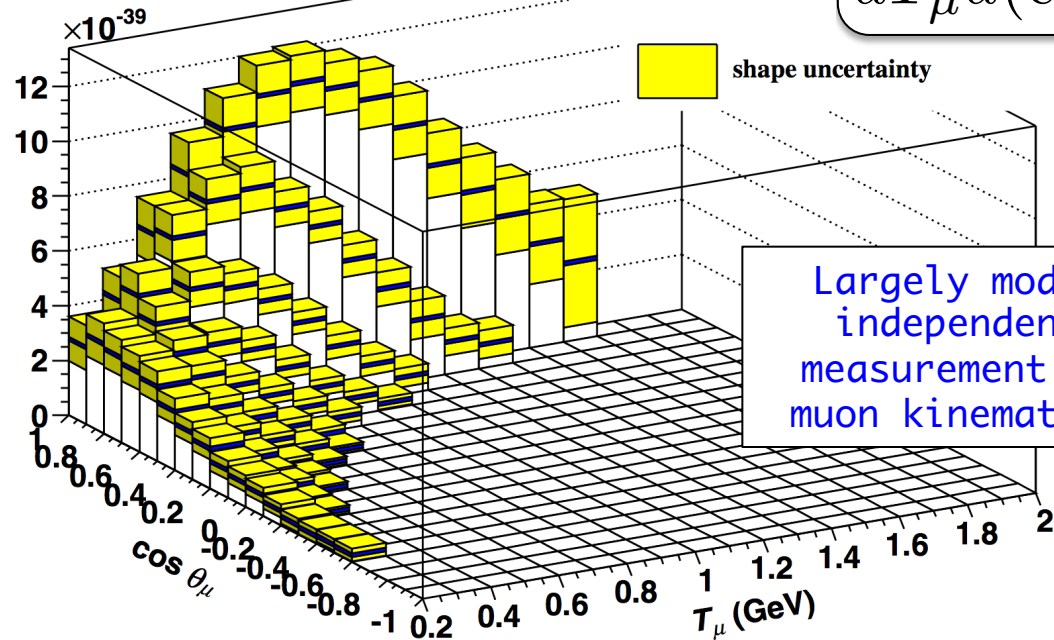
Several examples of
first measurements of
differential cross
sections

See Zarko Pavlovic's talk
from WG2 on Tuesday

First Measurement of the Muon Antineutrino Double-Differential Charged-Current Quasielastic Cross Section

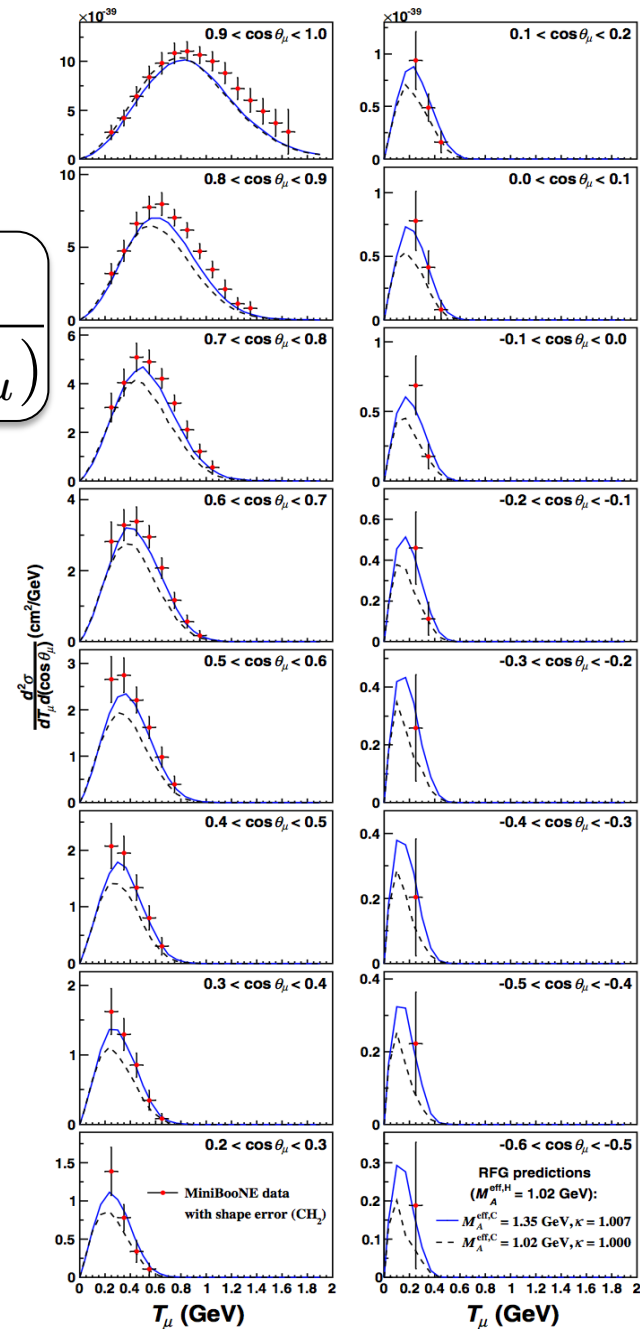
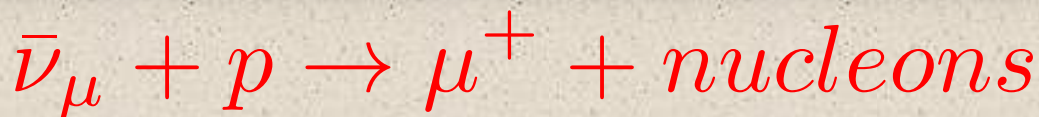
PRD 88, 032001 (2013)

$$\frac{d^2\sigma}{dT_\mu d(\cos\theta_\mu)} \text{ (cm}^2/\text{GeV)}$$



Largely model independent measurement of muon kinematics

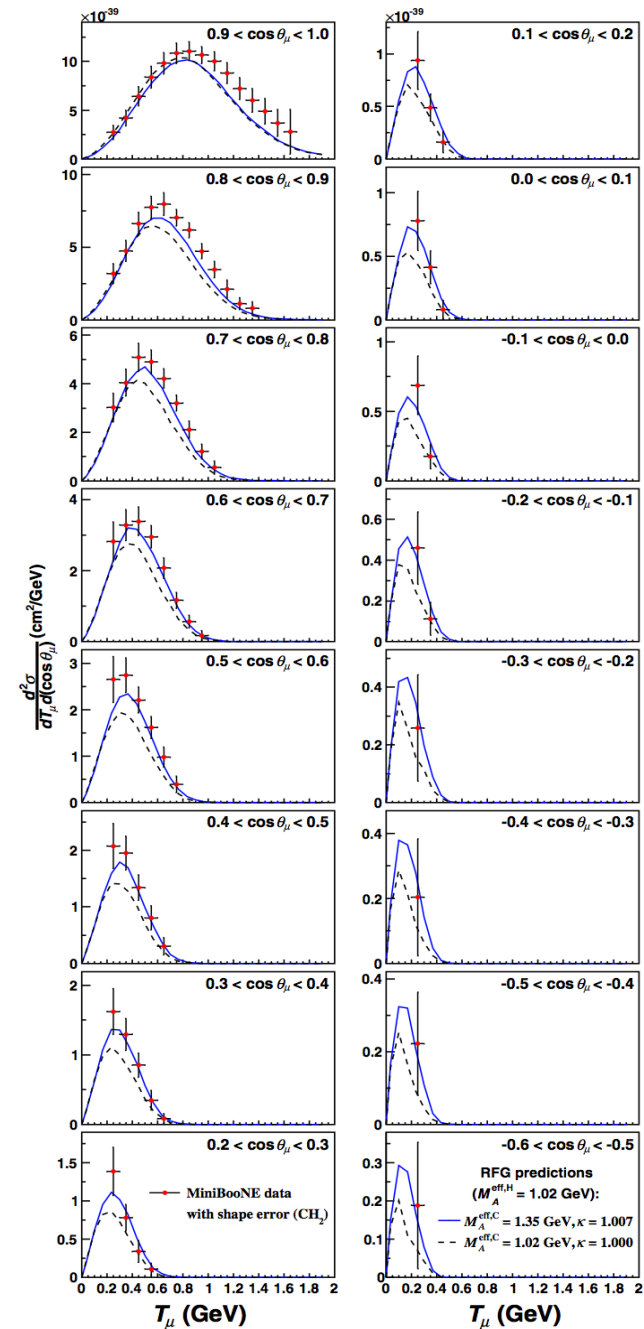
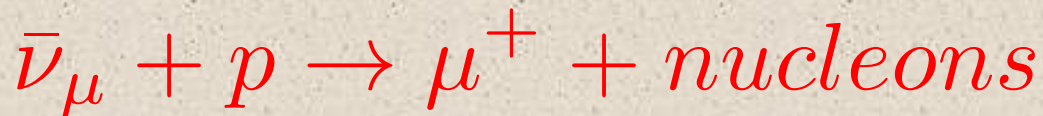
$$\frac{d^2\sigma}{dT_\mu d(\cos\theta_\mu)}$$



First Measurement of the Muon Antineutrino Double-Differential Charged-Current Quasielastic Cross Section

PRD 88, 032001 (2013)

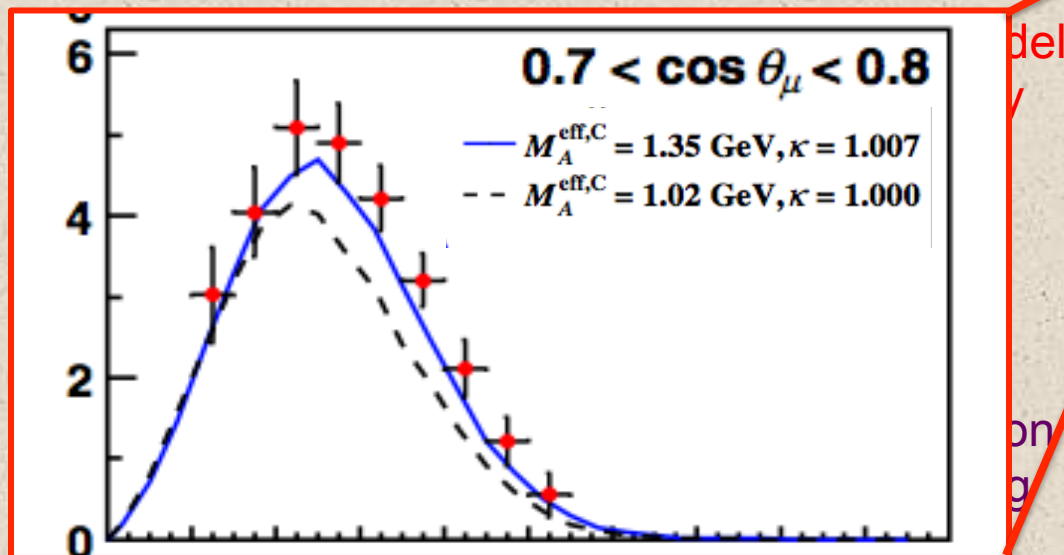
- Strong conclusions about the RFG model:
 - “It is clear in Fig. 9 [right] that the RFG model assuming $M_A \sim 1$ GeV does not adequately describe these data in shape or in normalization.”
 - “Consistent with other recent CCQE measurements on nuclear material, a significant enhancement in the normalization that grows with decreasing muon scattering angle is observed compared to the expectation with $M_A = 1.0$ GeV.”



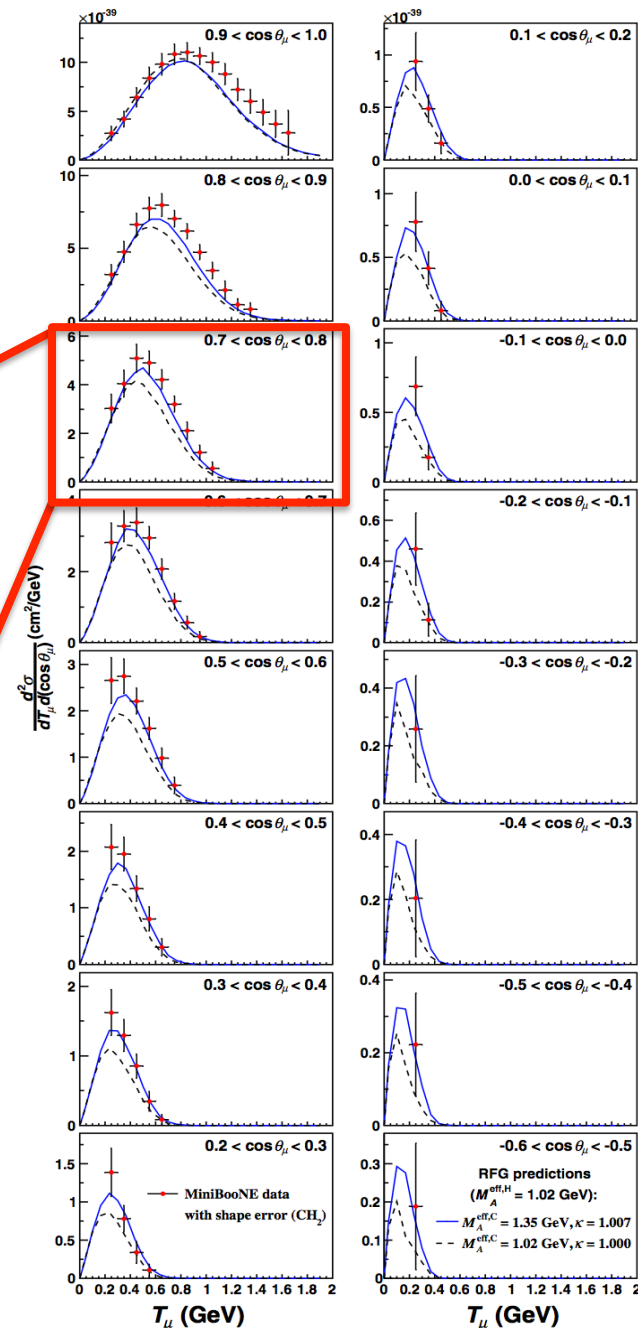
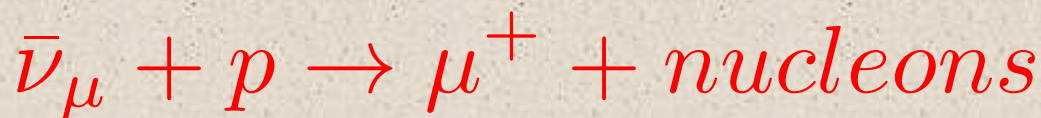
First Measurement of the Muon Antineutrino Double-Differential Charged-Current Quasielastic Cross Section

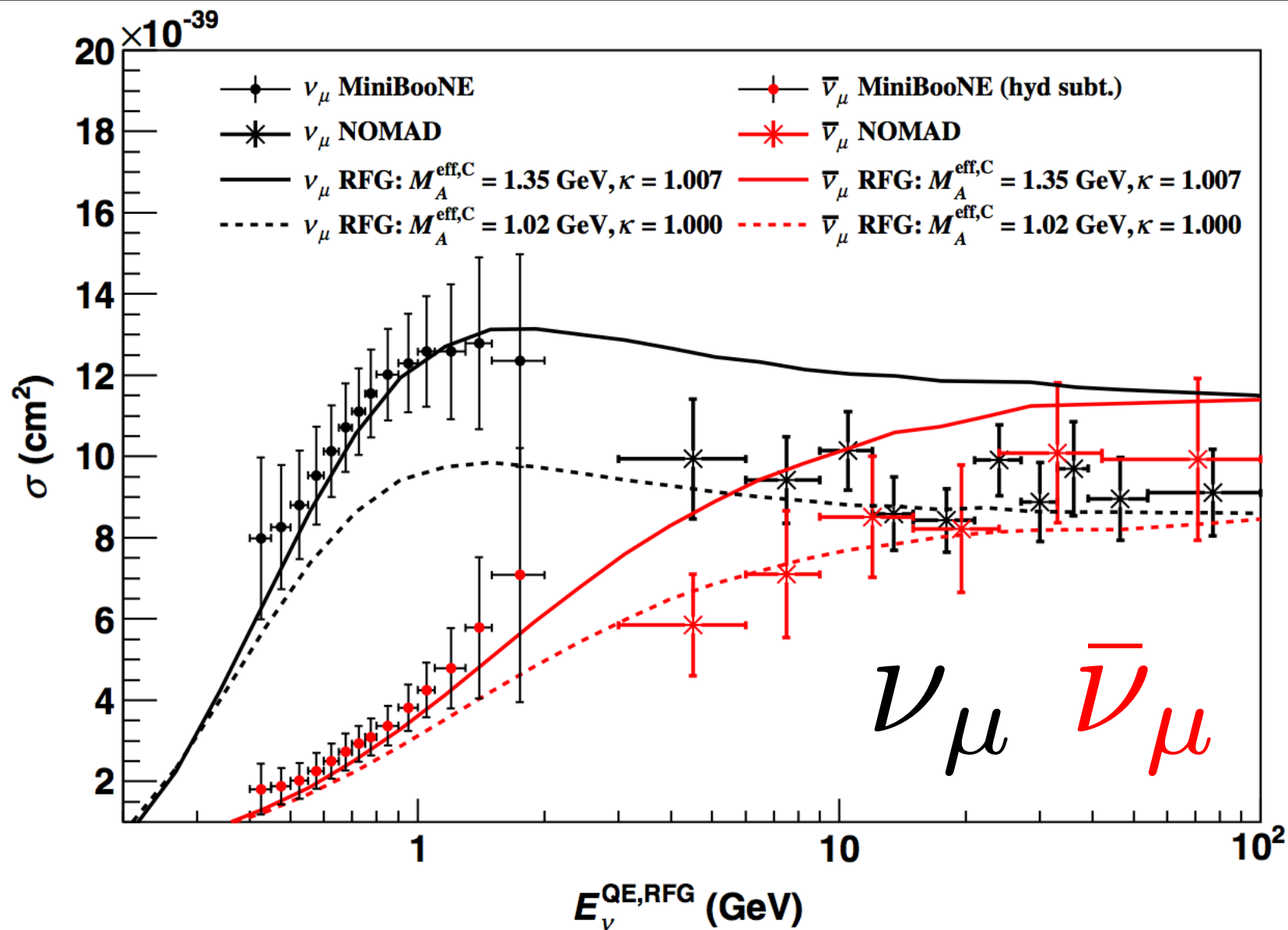
PRD 88, 032001 (2013)

- Strong conclusions about the RFG model:



expectation with $M_A = 1.0$ GeV.”



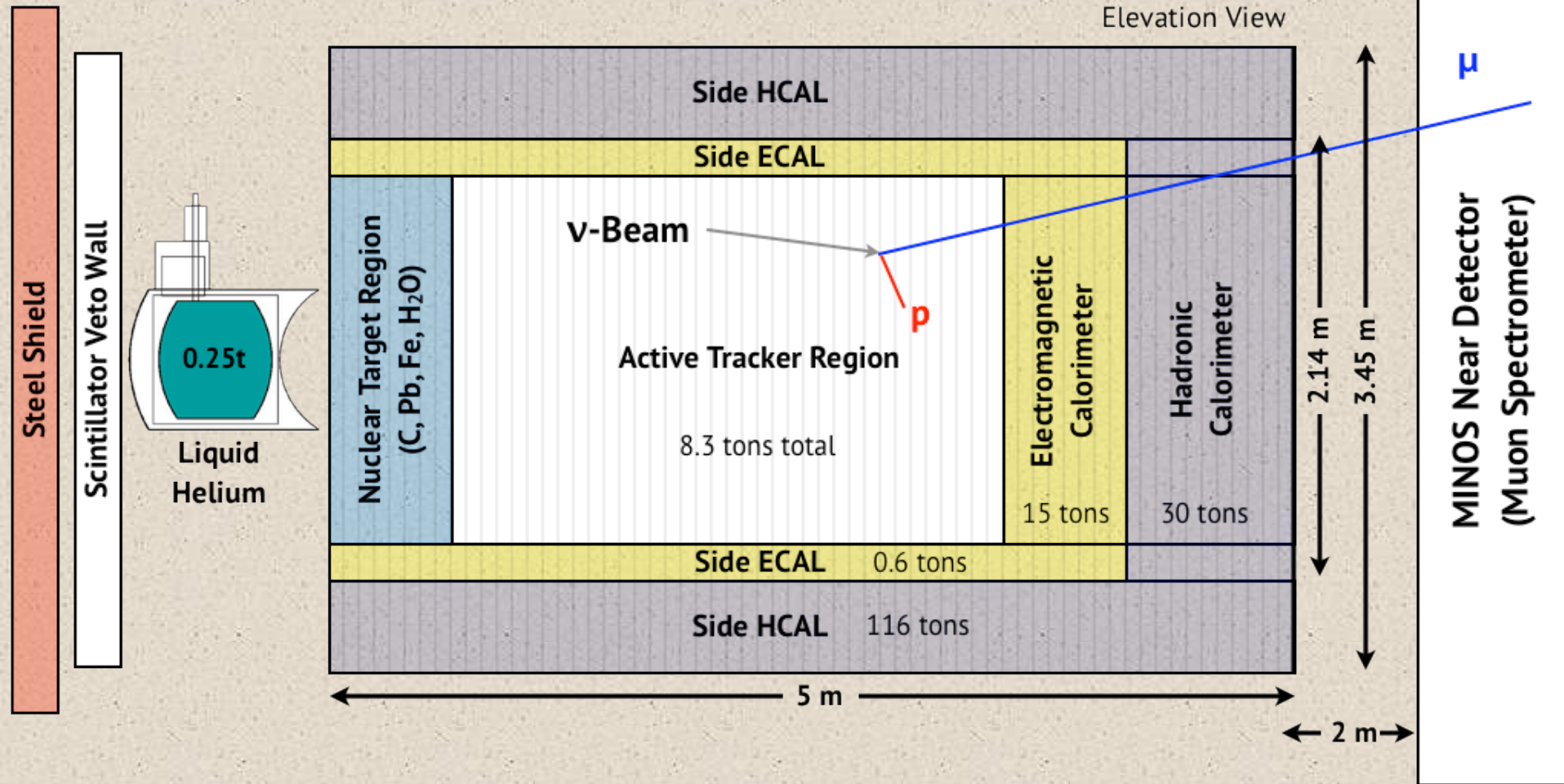




MINERvA

See WG2 talks from
C. Marshall, A. Higuera,
A. Bravar

MINOS ND serves as
muon spectrometer

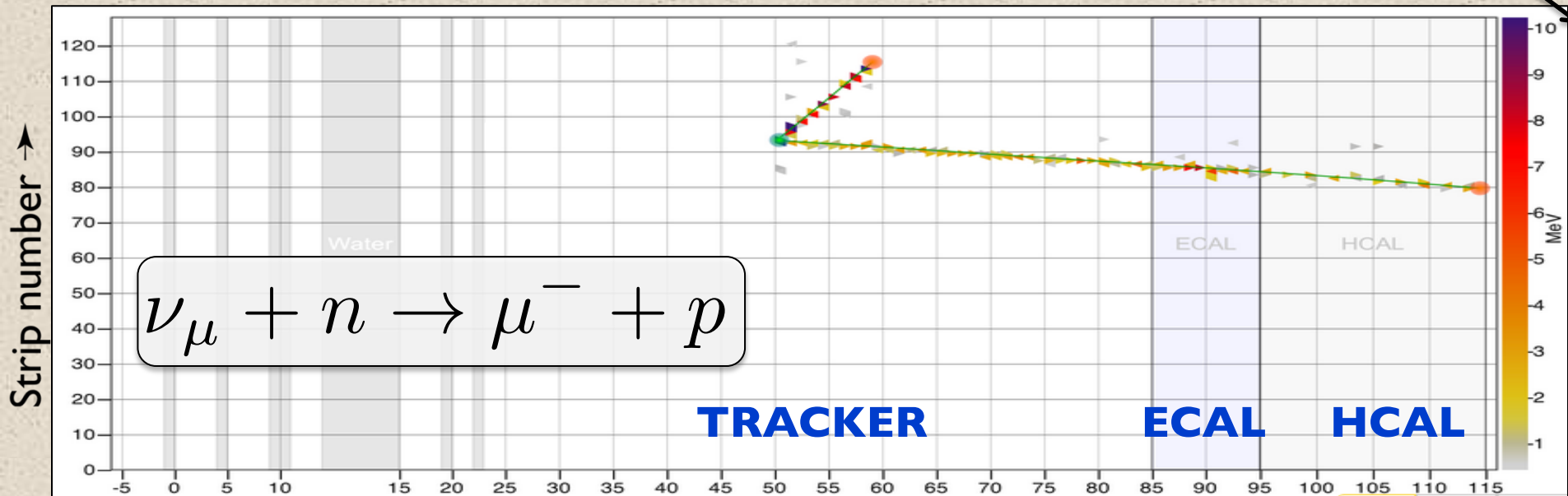
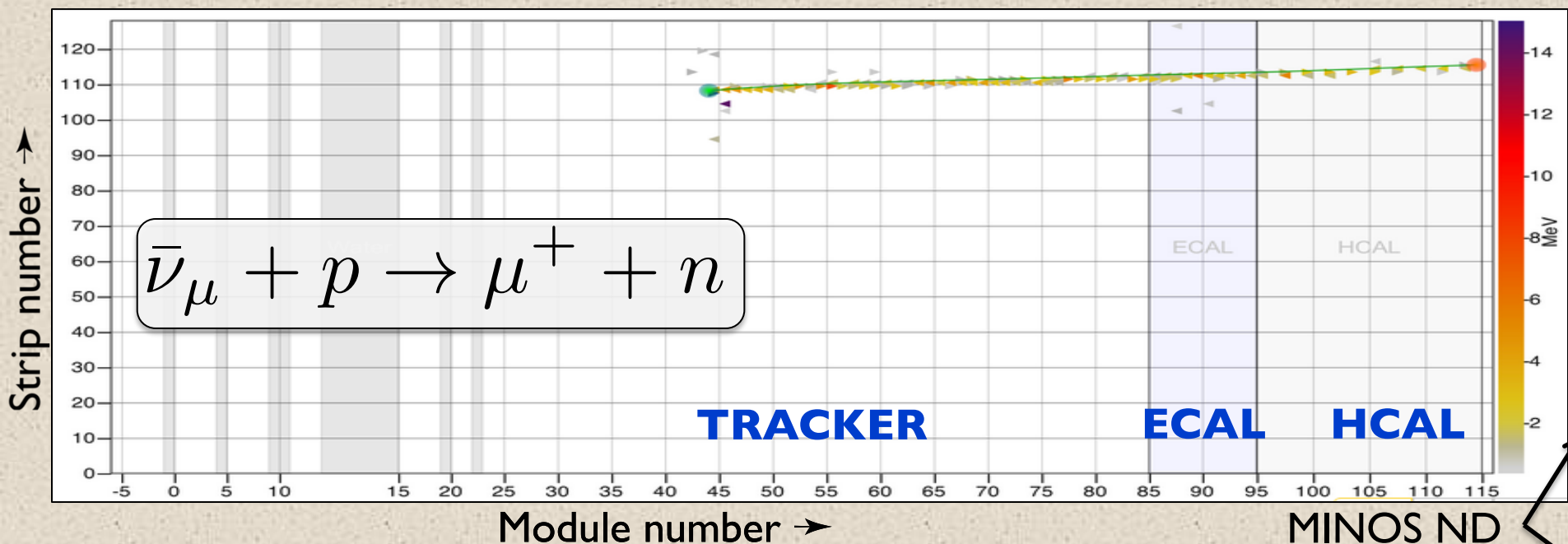


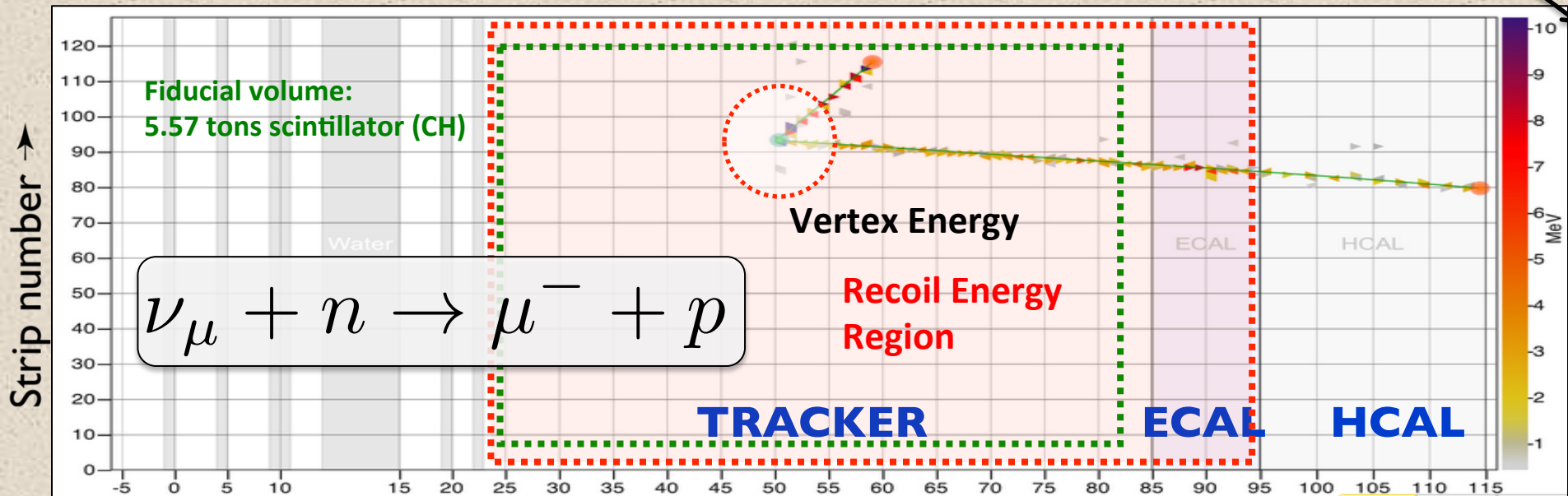
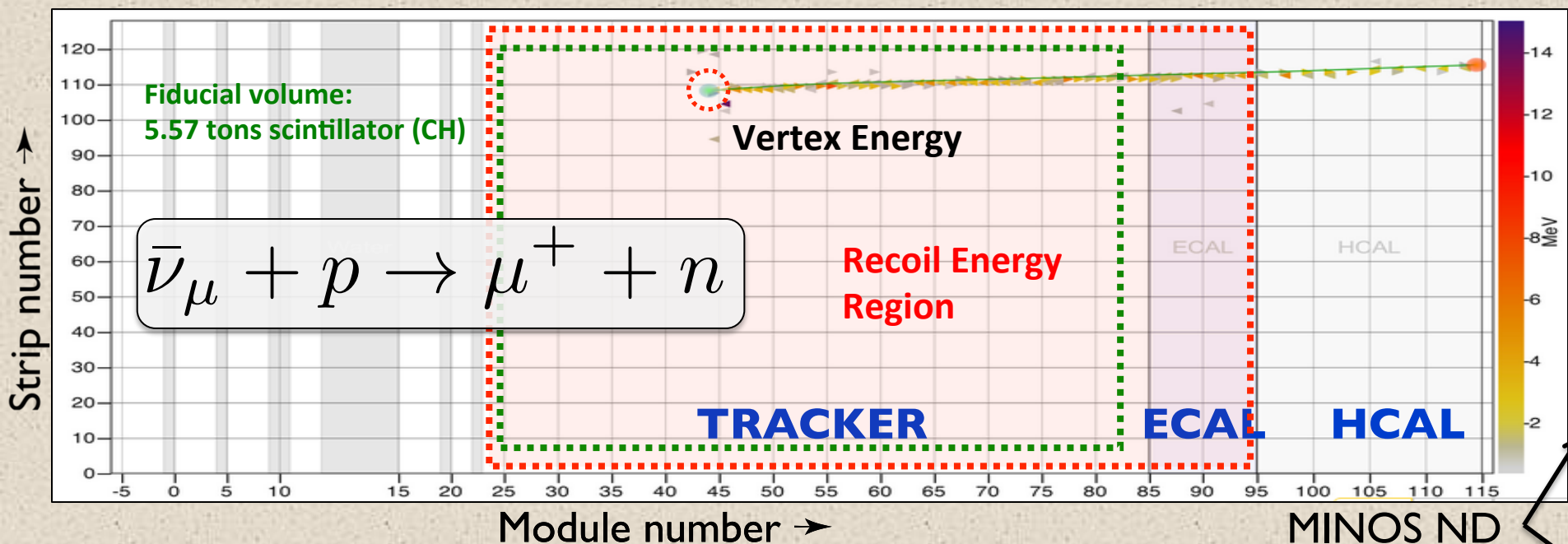
MINERvA detector comprised of 120 modules of varying composition stacked along the beam direction. Finely segmented tracking (~32k channels) with nuclear targets (C, CH, Fe, Pb, He, H₂O)

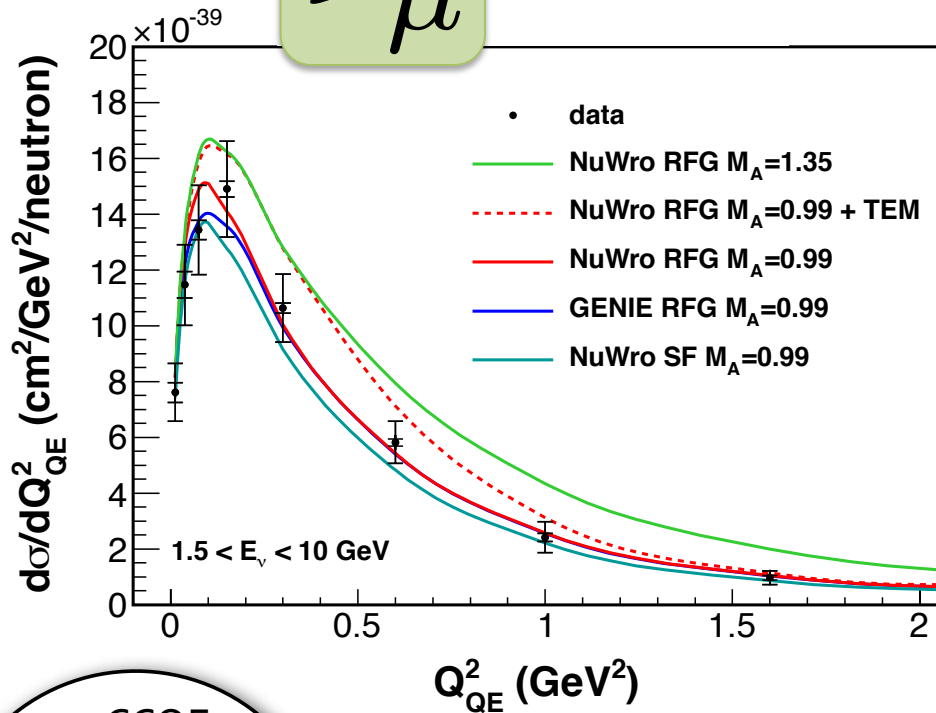
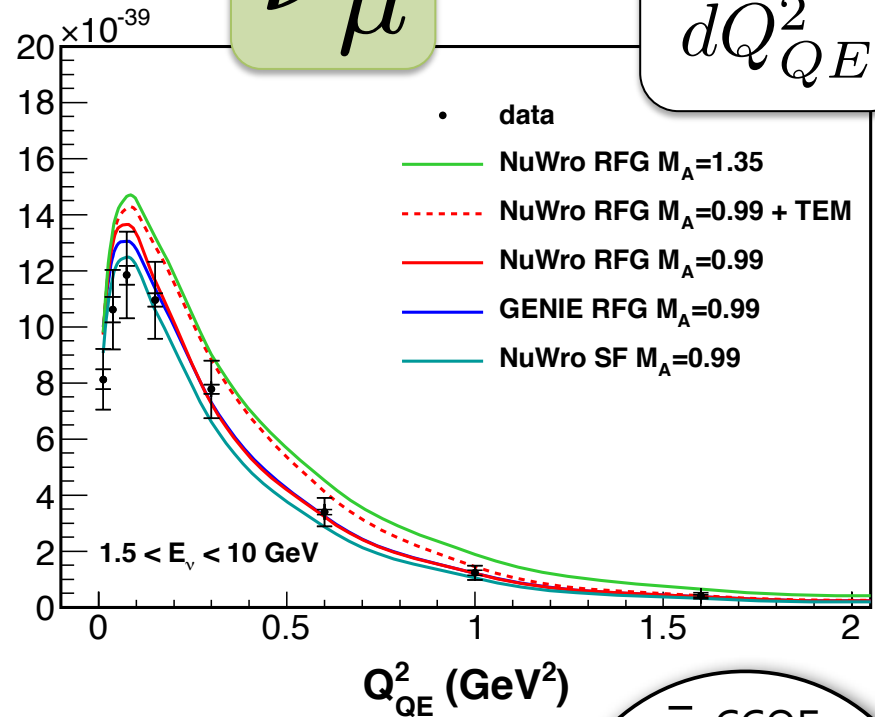
$\bar{\nu}$ Beam \longrightarrow

charged-current quasi-elastic scattering

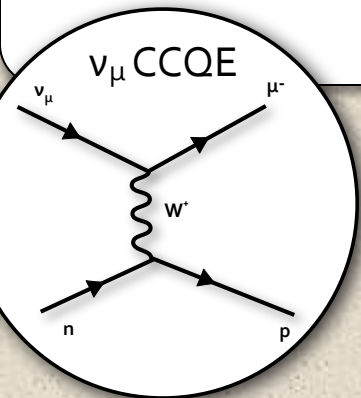
MeV



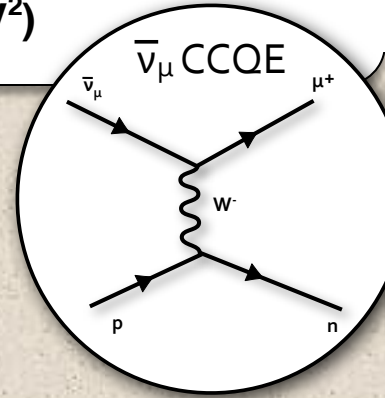


ν_μ  $\bar{\nu}_\mu$ 

$$\frac{d\sigma}{dQ_{QE}^2}$$



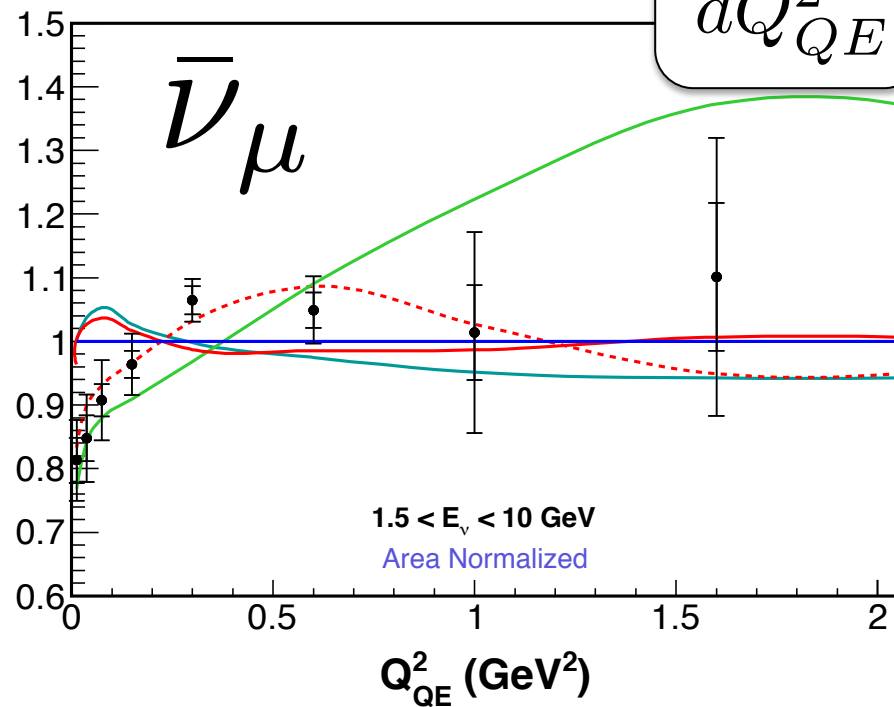
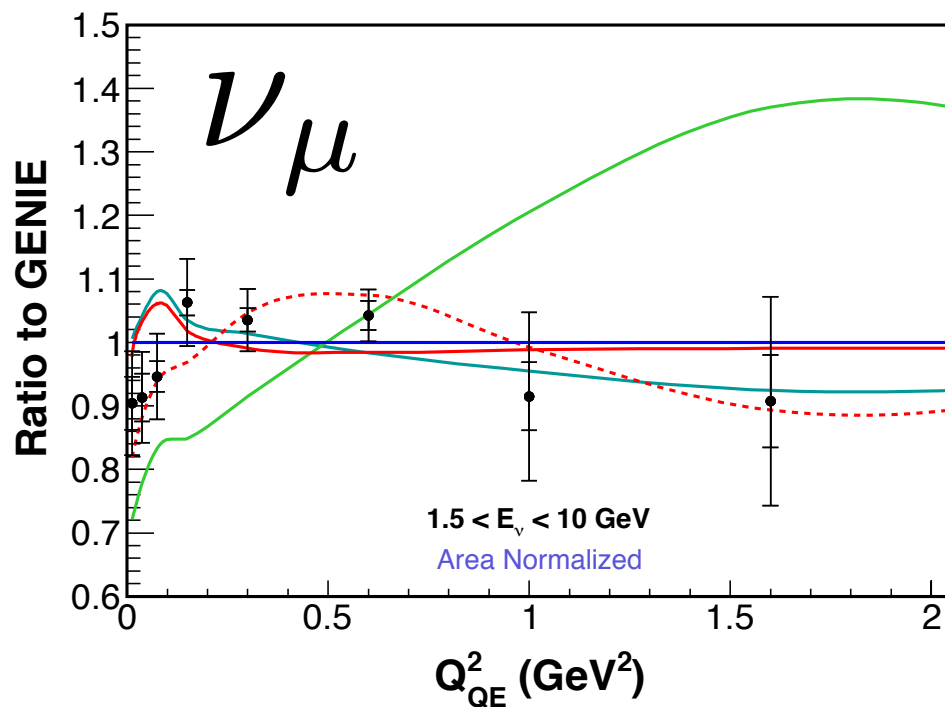
	Events	Efficiency	Purity
Neutrinos	29,620	47%	49%
Antineutrinos	16,467	54%	77%



MINERvA

Emphasize Shape: Normalize each prediction and the measured cross section to GENIE prediction and form ratio

$$\frac{d\sigma}{dQ_{QE}^2}$$



- GENIE** — independent nucleons in a mean field ($M_A = 0.99$ GeV)
- $M_A = 1.35$ GeV** — best fit to MiniBooNE data
- SF** — improved nucleon momentum-energy relation
- TEM** - - - empirical model based on electron scattering data



MINERvA

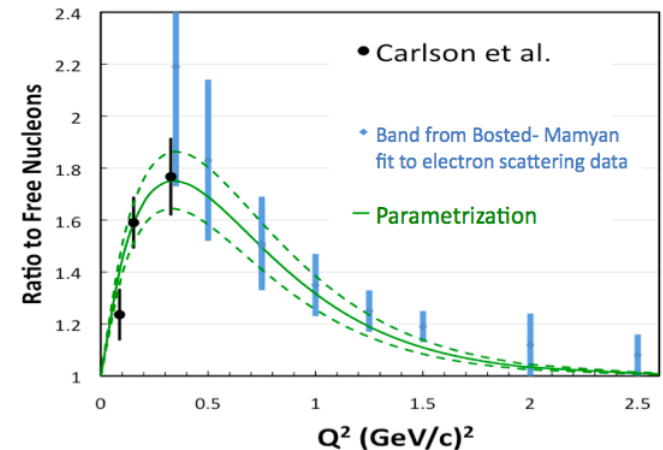
Transverse Enhancement Model

Bodek, Budd, Christy, Eur.
Phys. J. C 71:1726 (2011),
arXiv:1106.0340

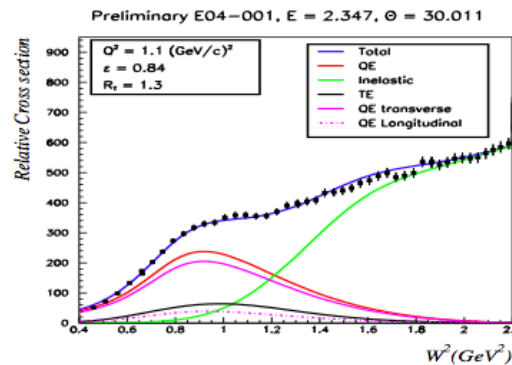
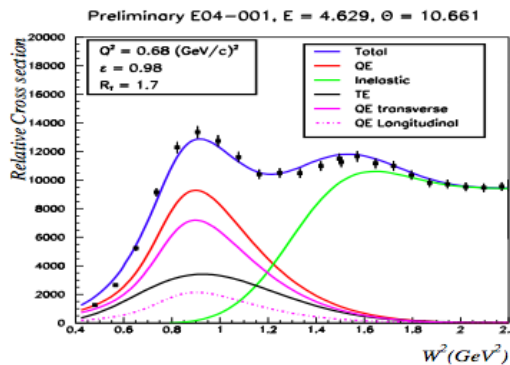
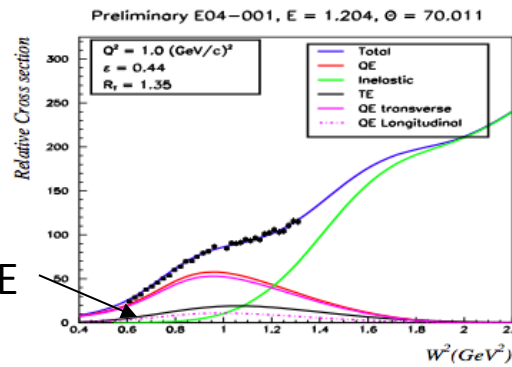
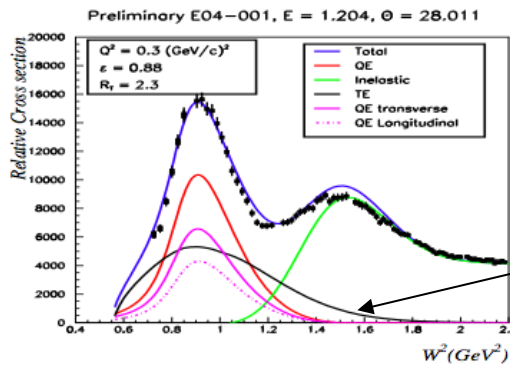
Fits in Q^2 bins

$$R_T = \frac{QE_{transverse} + TE}{QE_{transverse}}$$

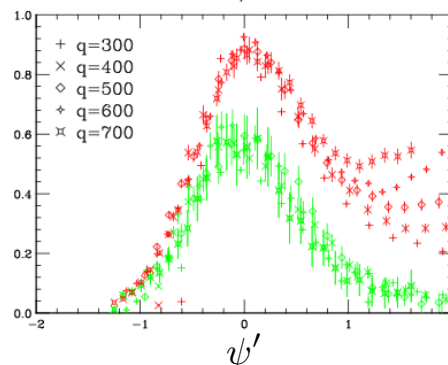
Transverse Enhancement Carbon 12



Applied as modifications of the
**magnetic form factors for
bound nucleons**



An attempt to
parameterize
this feature
we saw in
electron
scattering



J. Carlson, et al., PRC 65, 024002 (2002)

$$G_M^p(Q^2)$$

$$G_M^n(Q^2)$$

Ratio to GENIE

ν_μ

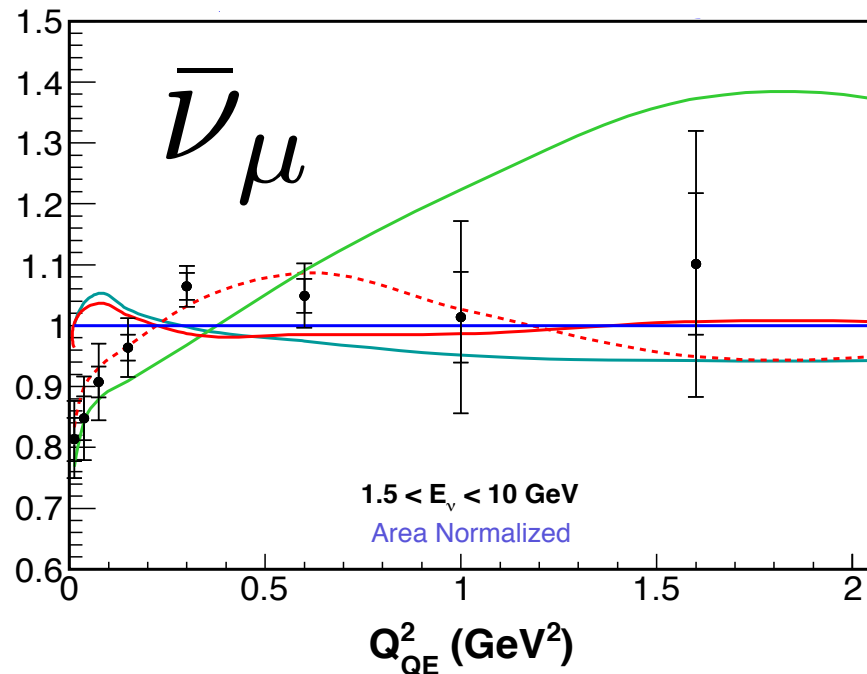
$1.5 < E_\nu < 10 \text{ GeV}$

Area Normalized

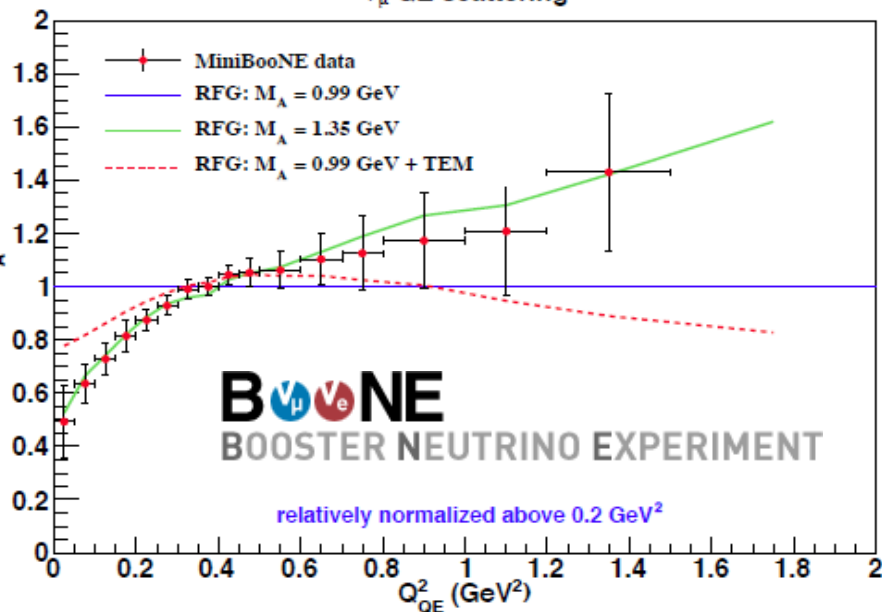
$Q_{QE}^2 \text{ (GeV}^2\text{)}$



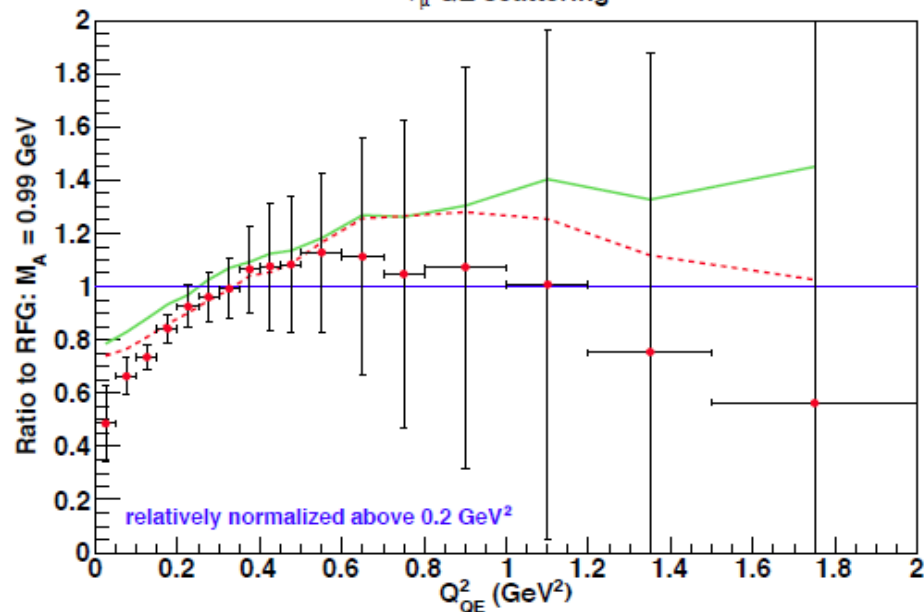
MINERvA



ν_μ QE scattering

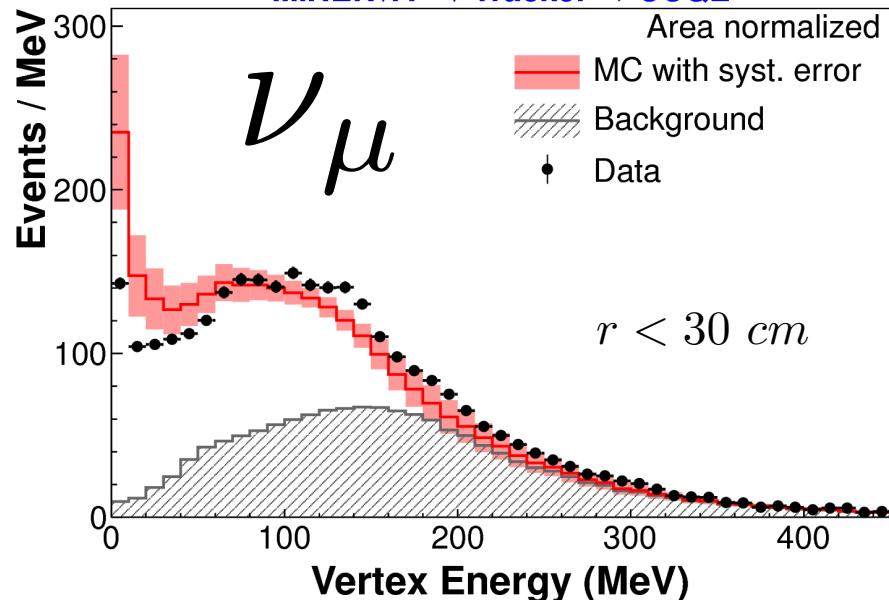
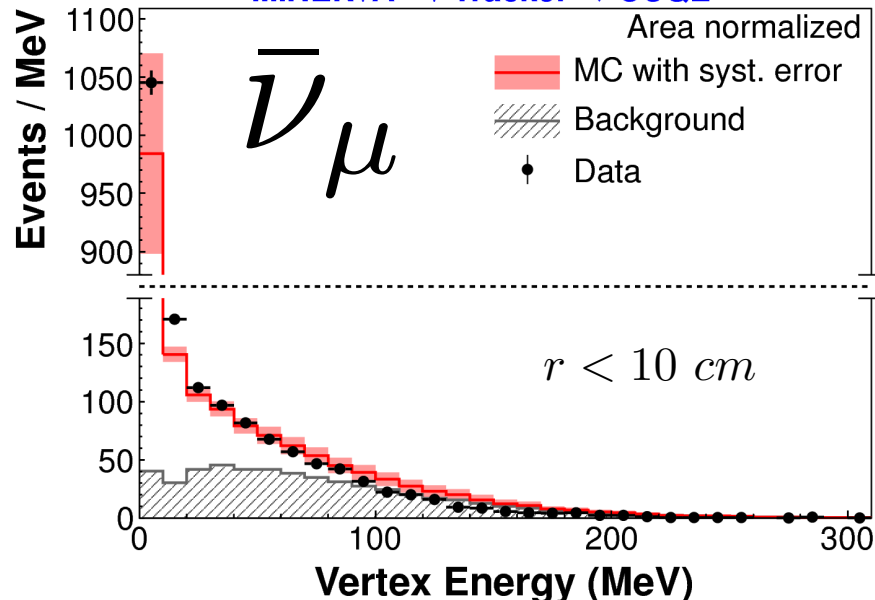


$\bar{\nu}_\mu$ QE scattering





Vertex Energy

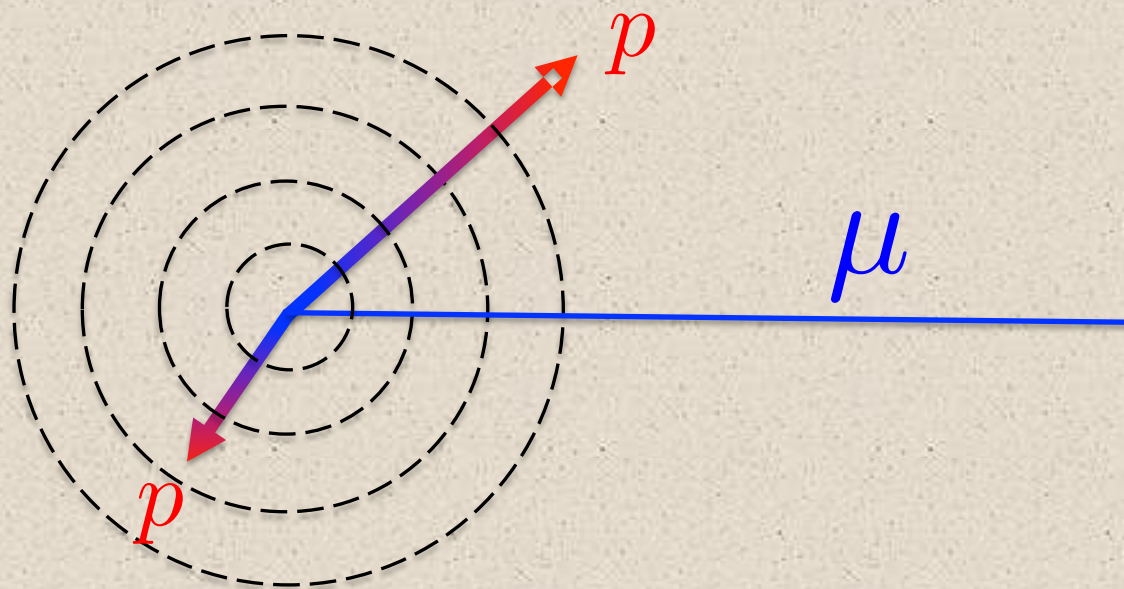
MINERvA • ν Tracker \rightarrow CCQEMINERvA • $\bar{\nu}$ Tracker \rightarrow CCQE

- A harder spectrum of vertex energy is observed in neutrinos
- All systematics considered, including energy scale errors on charged hadrons and FSI model uncertainties
- At this point, we make the *working assumption* that the additional vertex energy per event in data is *due to protons*



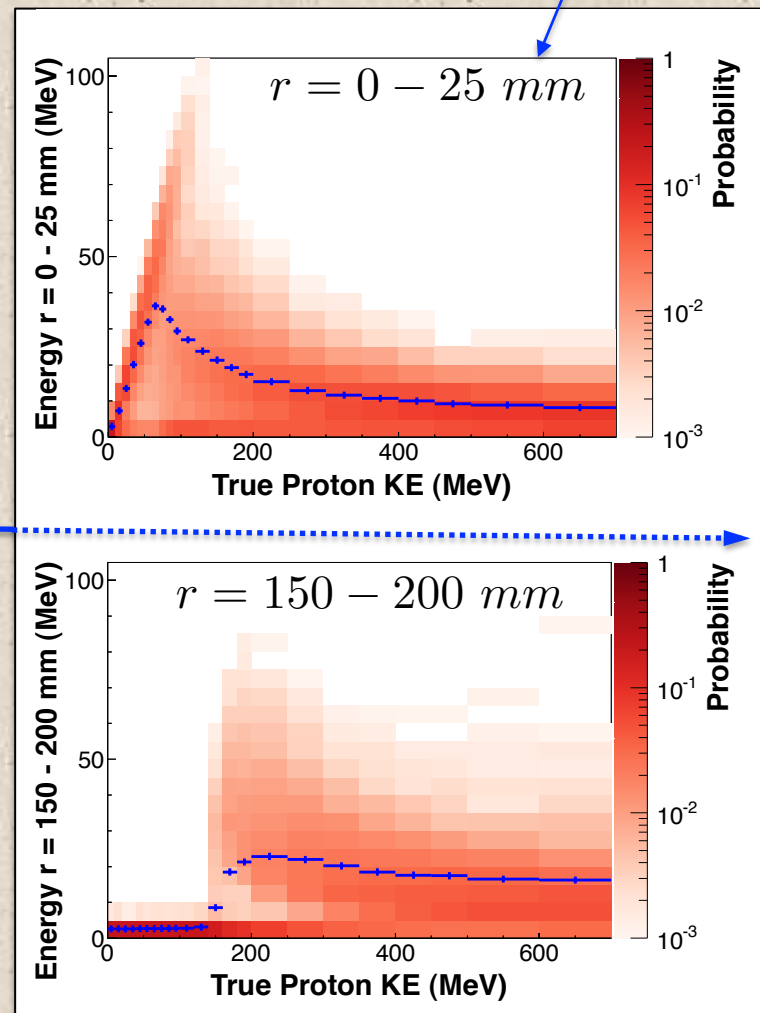
Vertex Energy

- Examine annular rings around the reconstructed vertex
 - Out to 10 cm for antineutrino (~ 120 MeV proton)
 - Out to 30 cm for neutrino (~ 225 MeV proton)



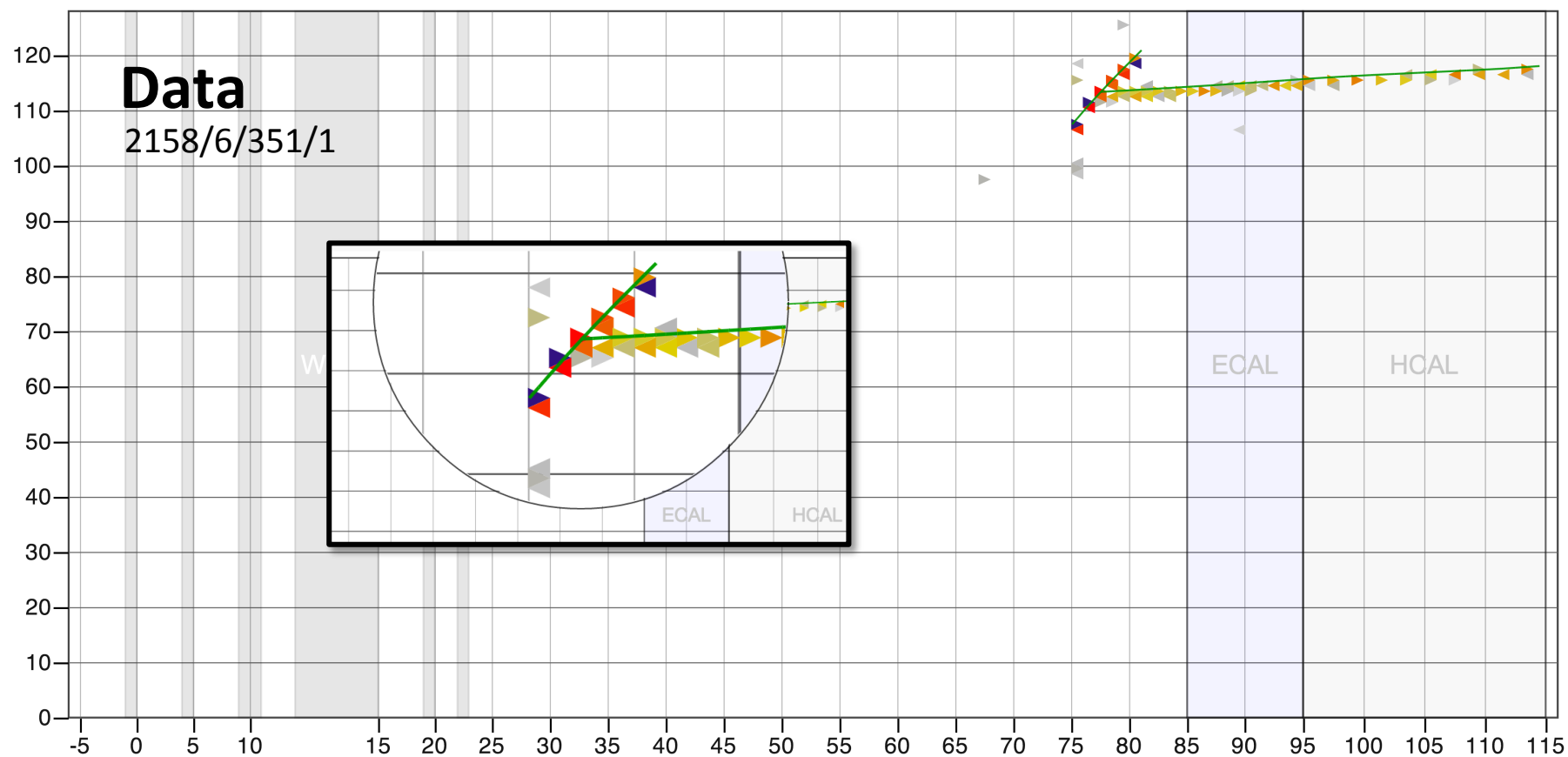
Note: to add visible energy to an inner annulus you must **add a charged hadron**, not just increase energy of an existing one

E_{vis} in that
annulus vs.
true $\text{KE}_{\text{proton}}$

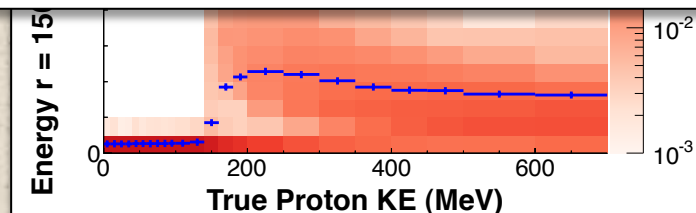




Vertex Energy

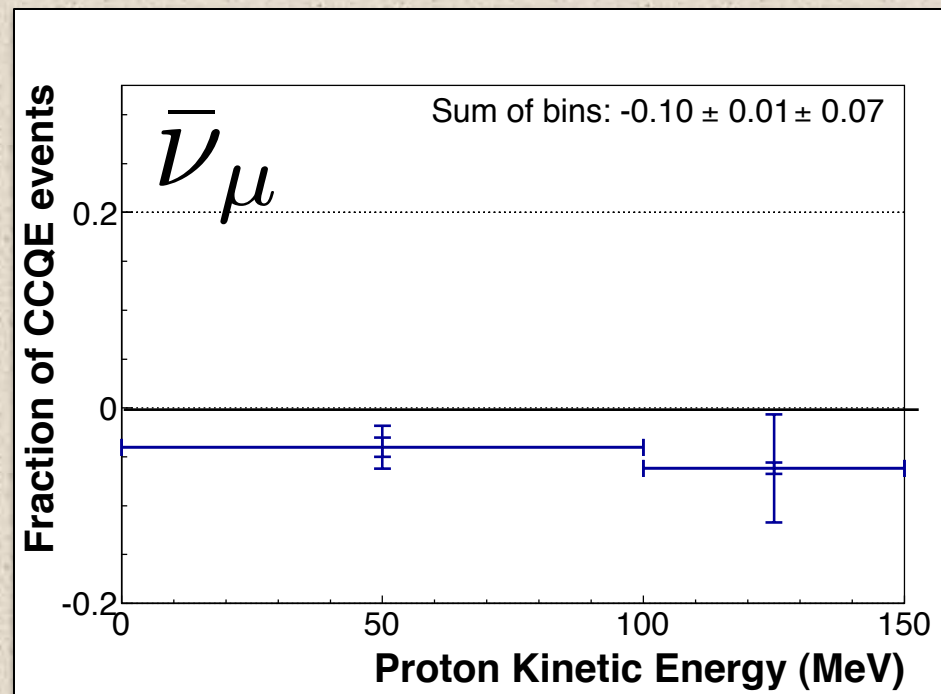
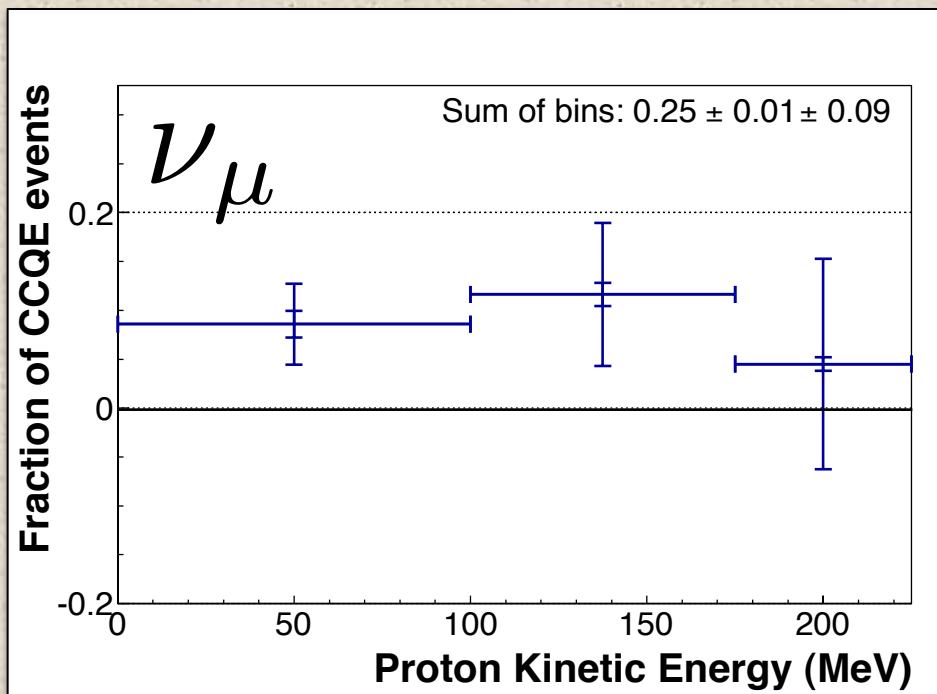


Note: to add visible energy to an inner annulus you must **add a charged hadron**, not just increase energy of an existing one



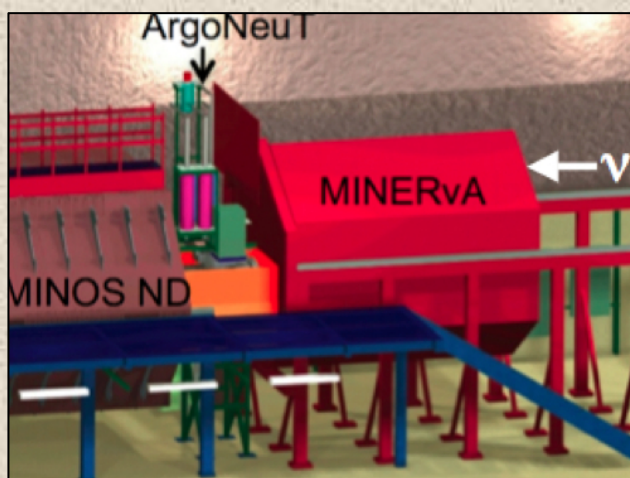


Vertex Energy



Find that adding an additional low-energy proton (KE < 225 MeV) to **$(25 \pm 9)\%$ of QE events** improves agreements with data

No such addition required for antineutrinos. Slight reduction if anything. **$(-10 \pm 7)\%$ of QE events**

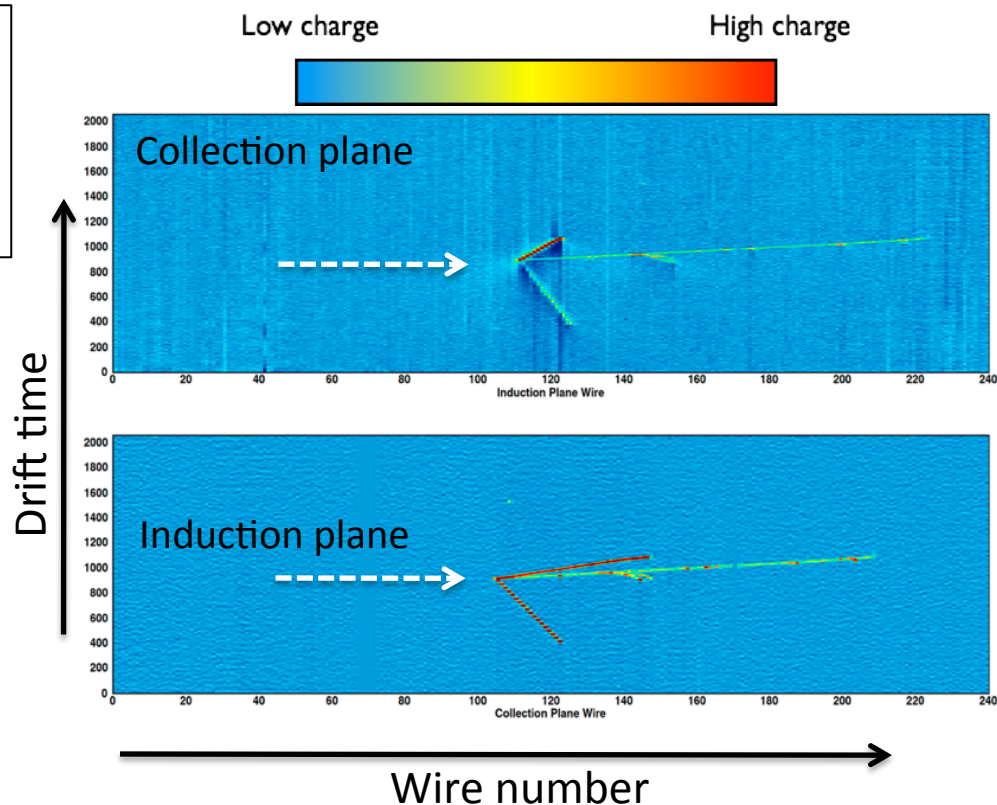
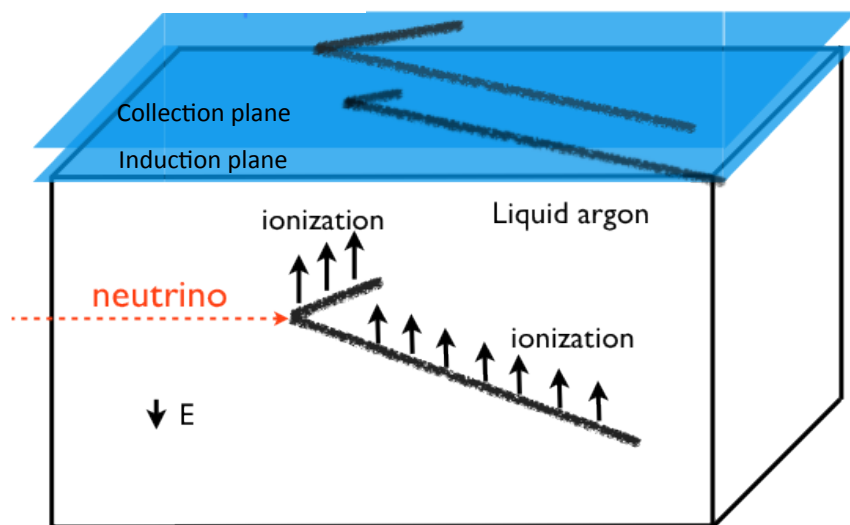


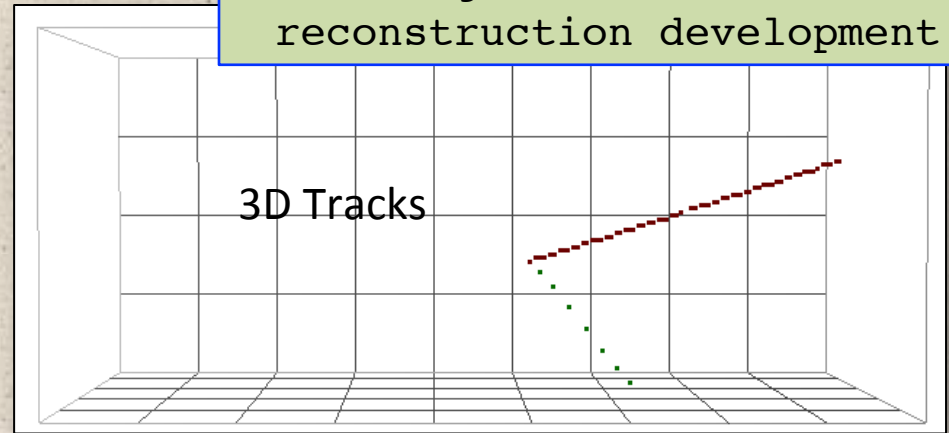
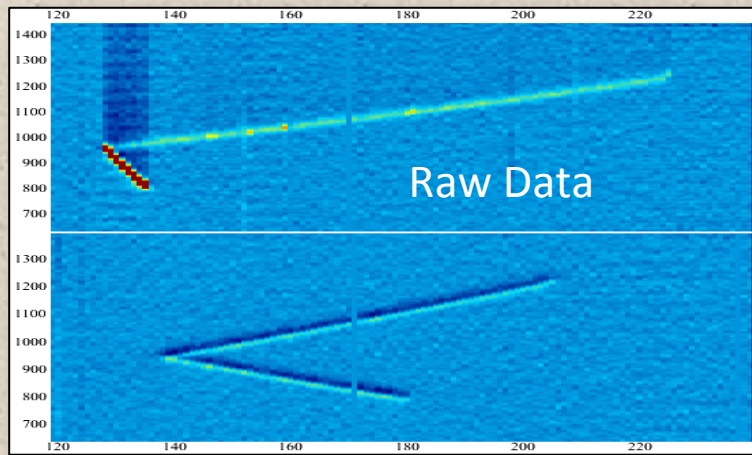
- Liquid argon time projection chambers offer a great opportunity for neutrino physics, including detailed study of neutrino-nucleus scattering
- ArgoNeuT detector exposed to NuMI beam
 - 0.085e20 POT neutrino mode
 - 1.2e20 POT antineutrino mode

See Tingjun Yang's talk from WG2 on Tuesday

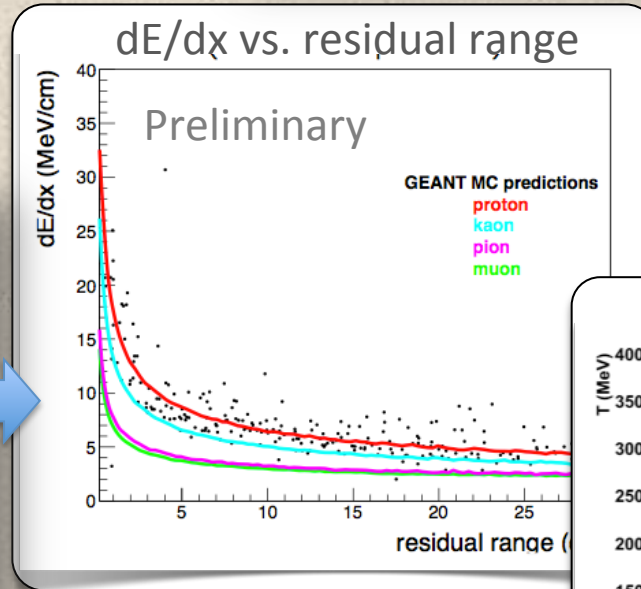
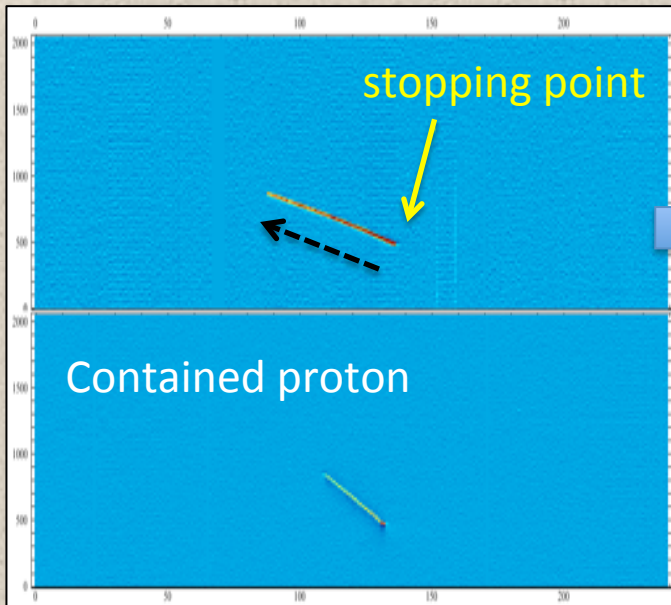


TPC Volume: 175 L
Wire Pitch: 4 mm
Max Drift: 0.5 m (330 μ s)
Electric Field: 500 V/m

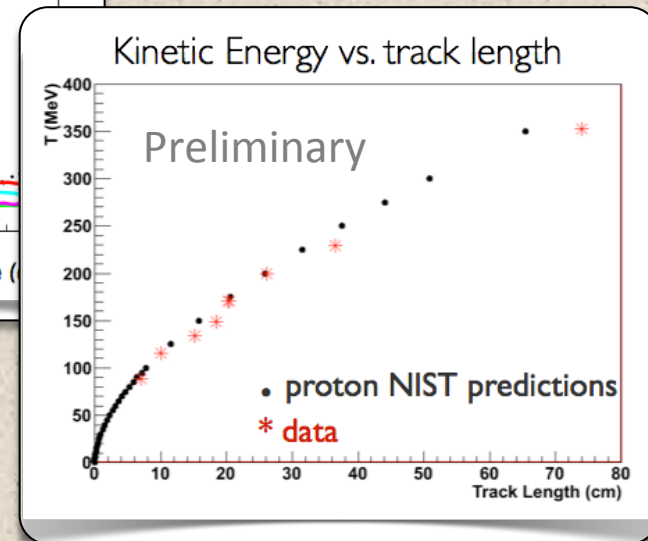


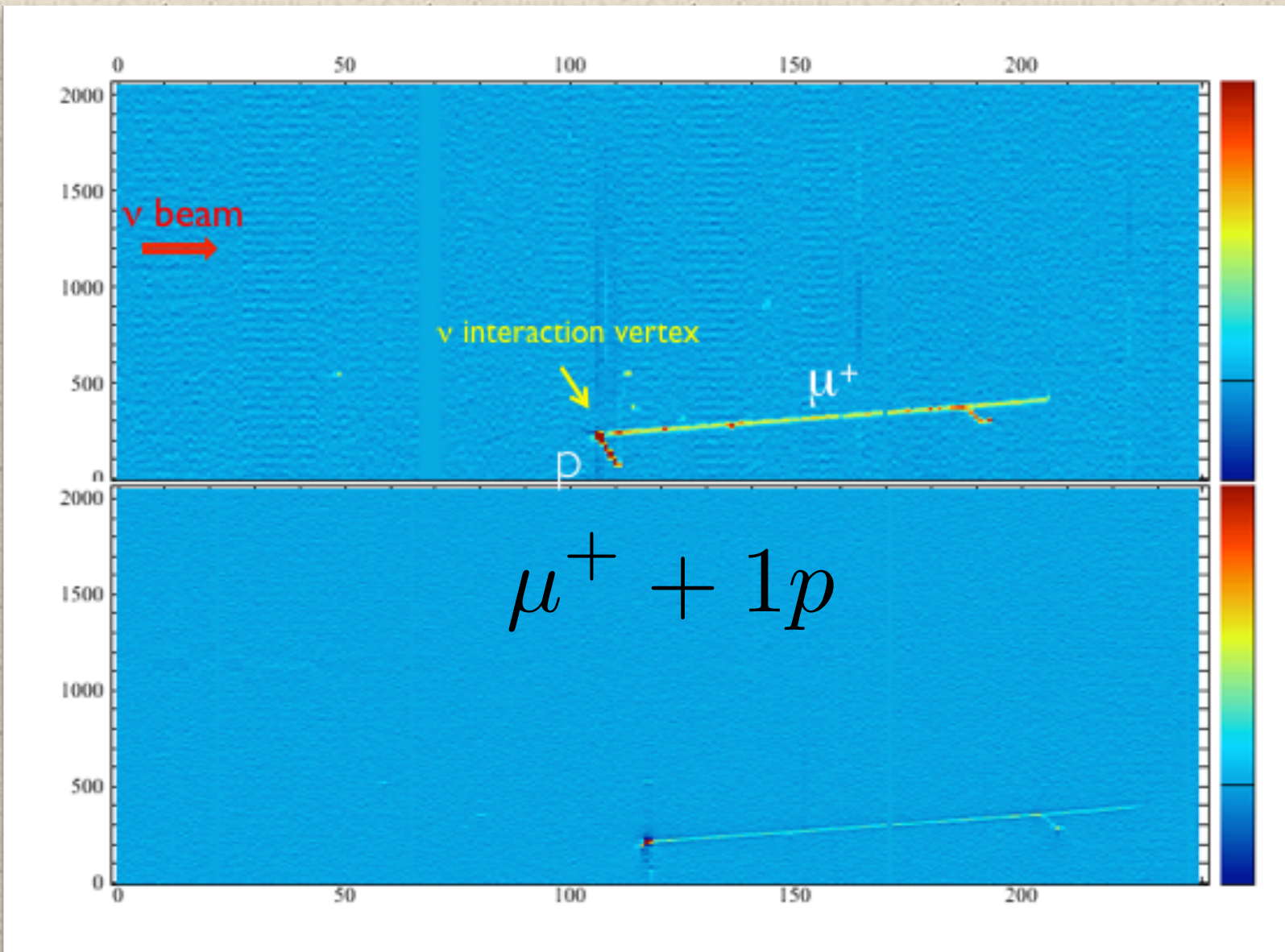


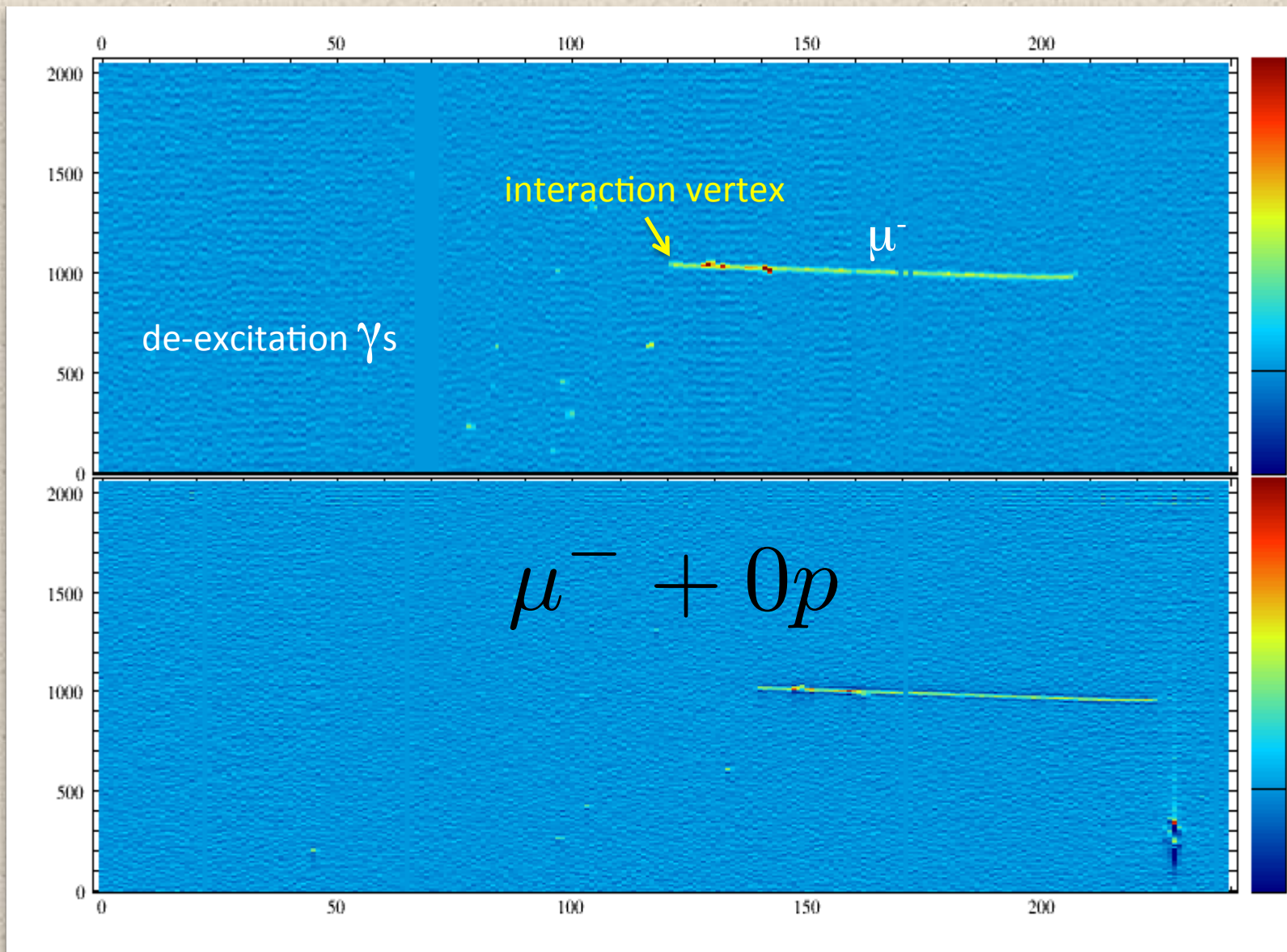
Providing test bed for event reconstruction development

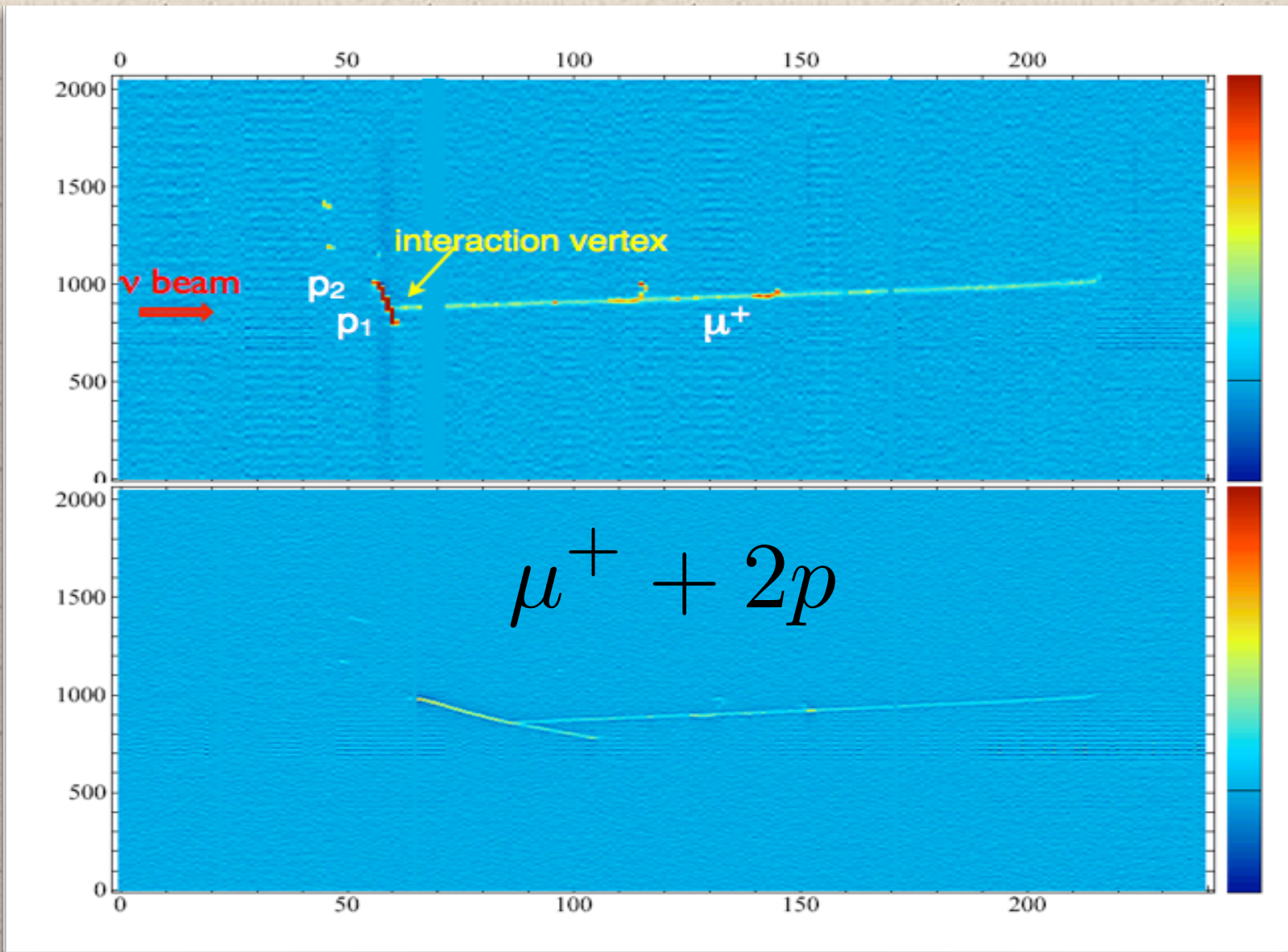


proton threshold:
 $T_p > 21 \text{ MeV}$

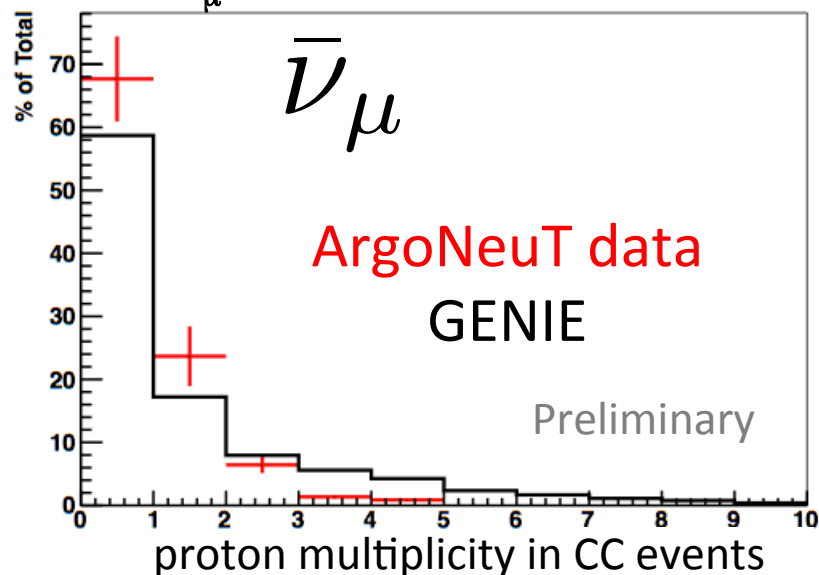




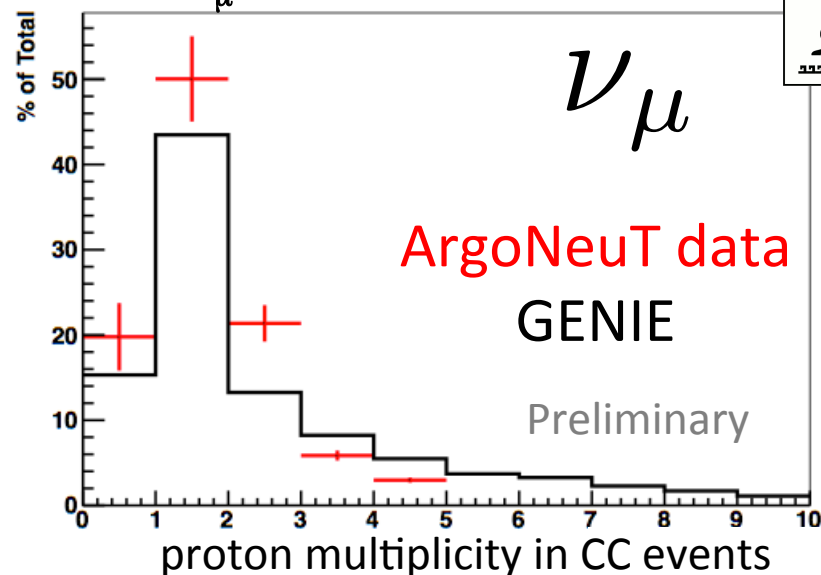




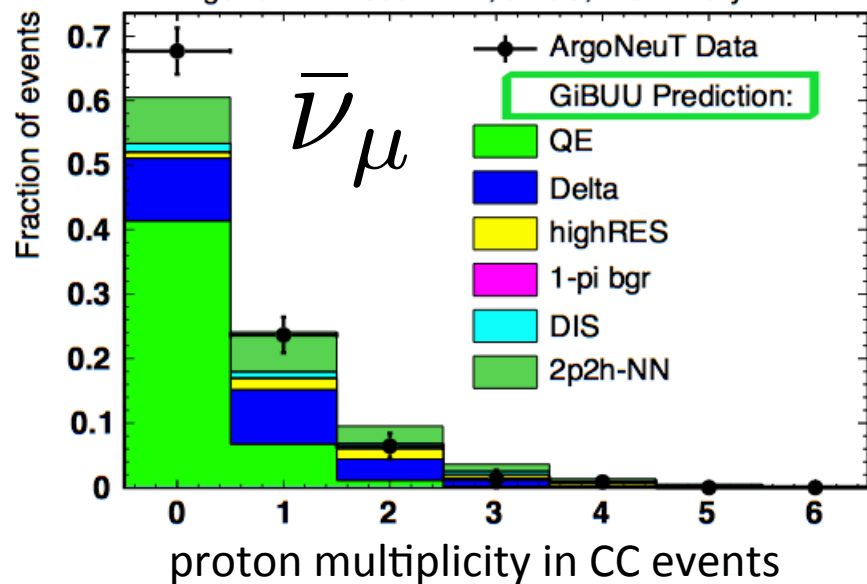
$\bar{\nu}_\mu$ – antineutrino mode run



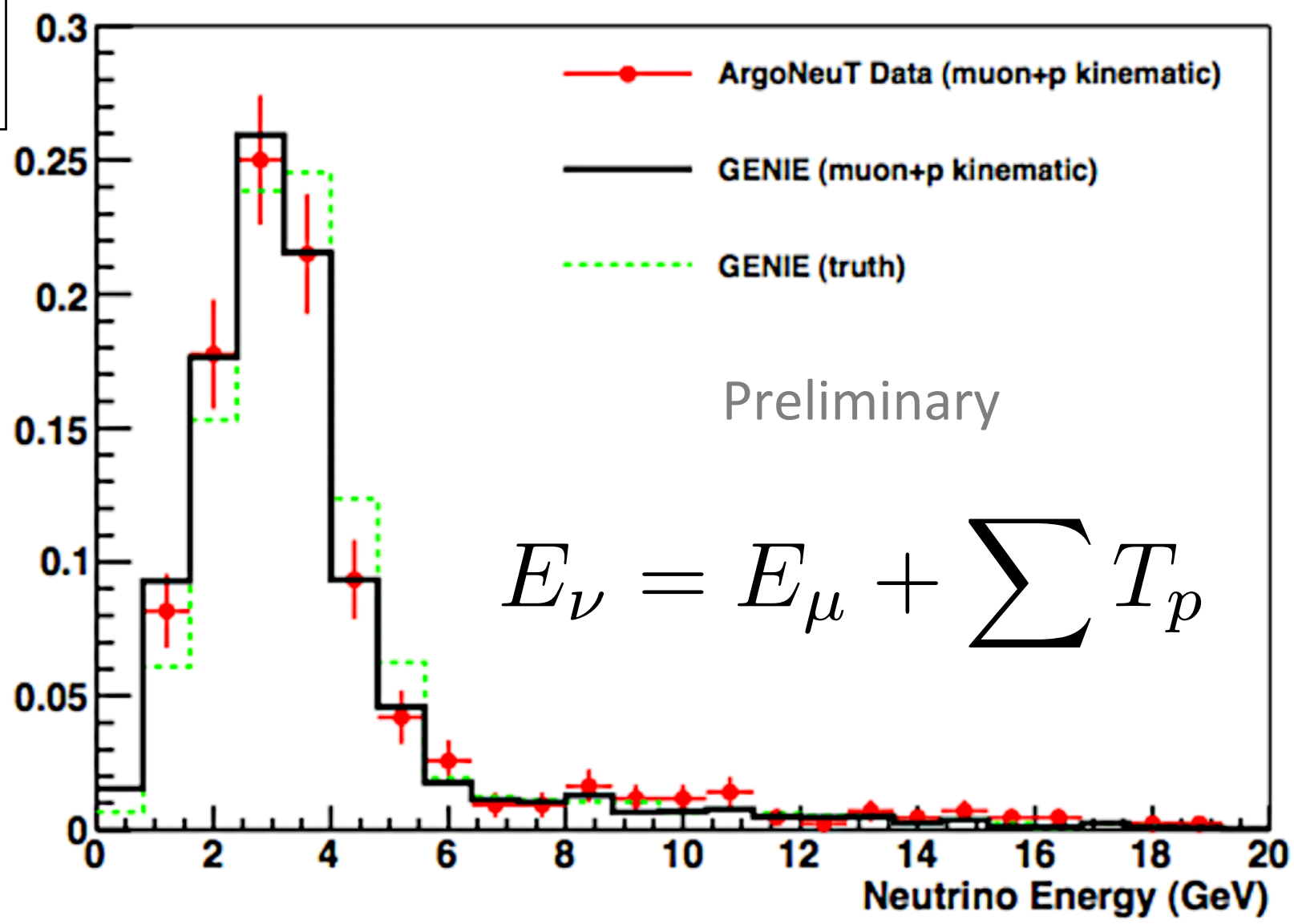
ν_μ – antineutrino mode run



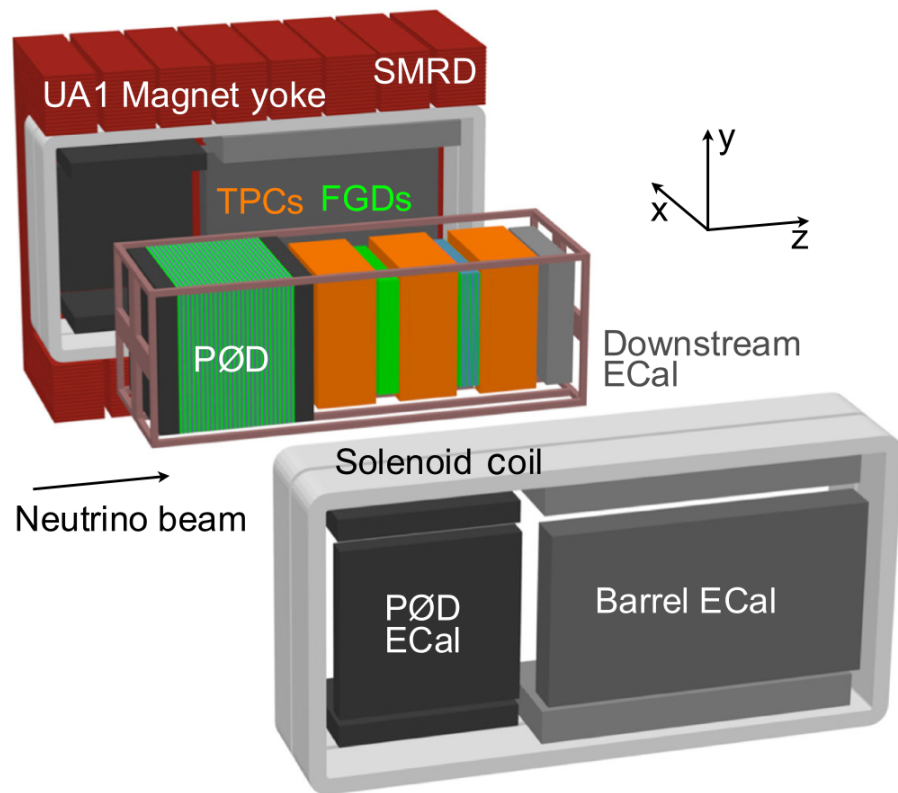
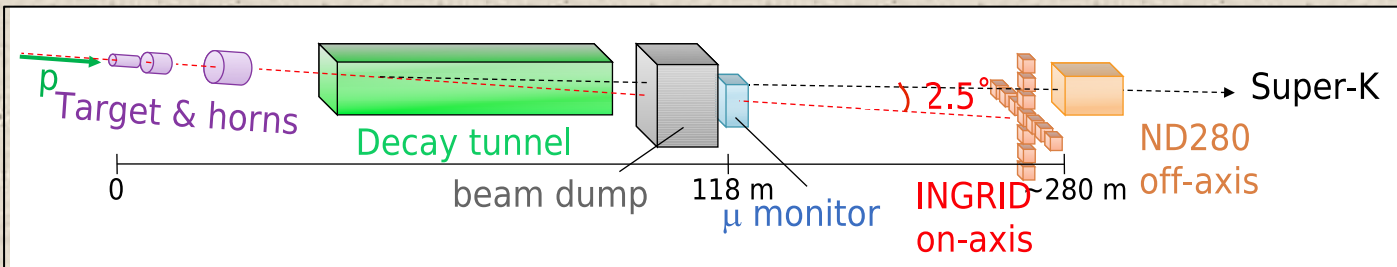
ArgoNeuT $\bar{\nu}$ -mode $\bar{\nu}$ -flux, 0π -CC, Preliminary



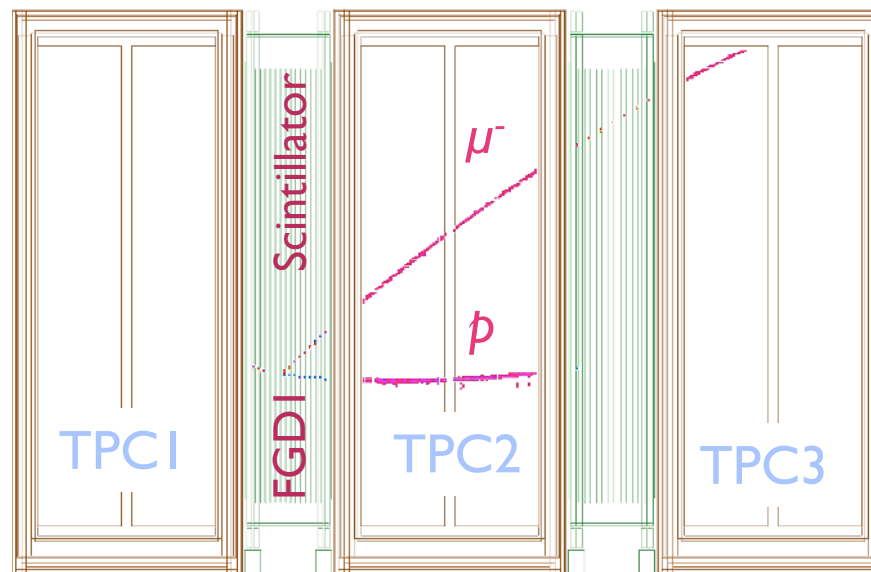
Low threshold allows
model independent
reconstruction of
complete final state
for detailed testing
of models



T2K



Run #: 4200 Evt #: 24083 Time: Sun 2010-03-21 22:33:25 JST

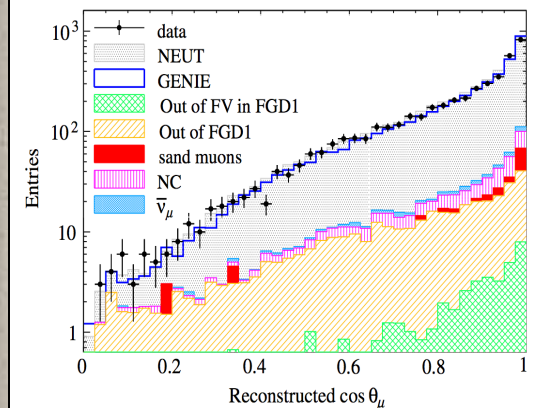
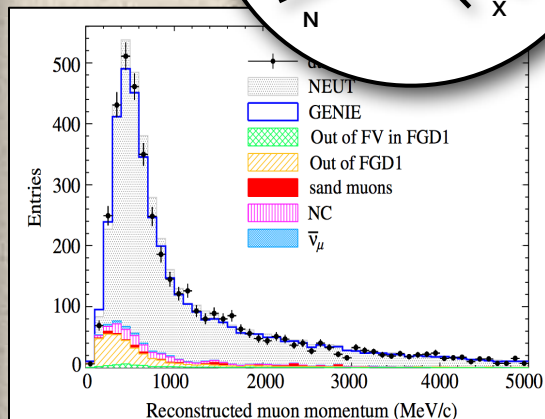
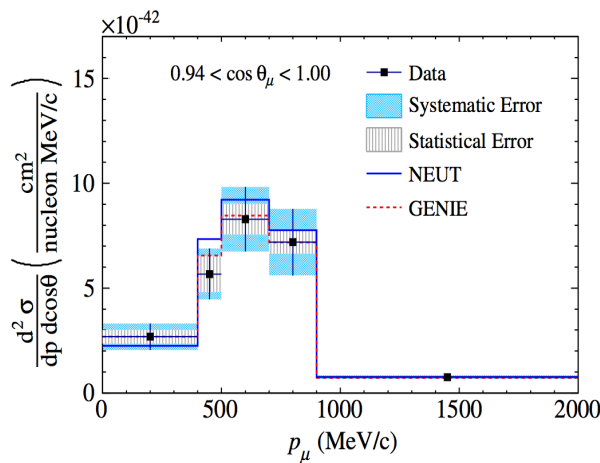
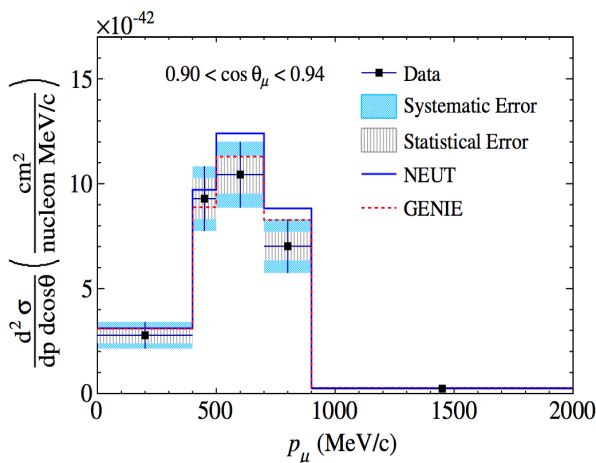
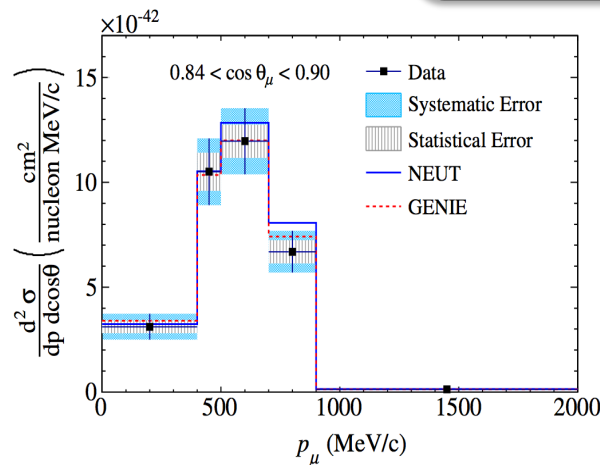
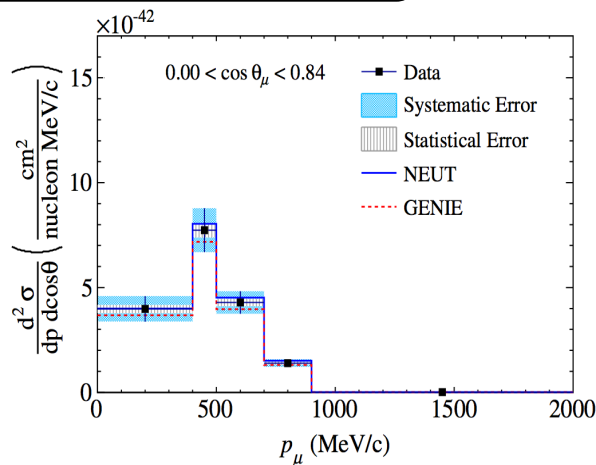
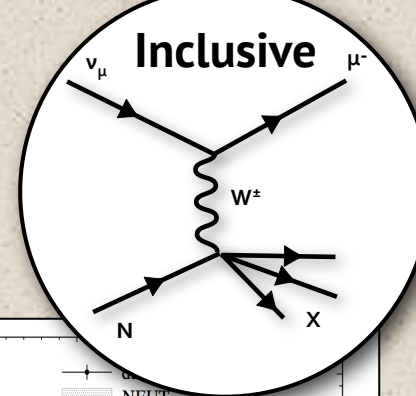


CCQE event candidate in the tracker region of the near detector. Muon reconstructed angle 40° and reconstructed momentum: 566 MeV/c

T2K

Differential cross sections very valuable!

$$\frac{d^2\sigma}{dT_\mu d(\cos\theta_\mu)}$$



CCQE and **NC elastic** measurements shown for first time here this week!

See WG2 talks from
D. Hadley (CC),
D. Ruterbories (NC)

Much More to Come

- MiniBooNE
 - CC inclusive cross sections
 - Antineutrino NC elastic
- MINERvA
 - More CCQE results
 - CC inclusive cross sections, comparisons between nuclear targets
 - Pion production processes (charged, neutral, coherent, resonant)
 - Electron scattering, electron neutrino CC interactions
- ArgoNeuT
 - Analysis of higher statistics antineutrino mode data (nu and nubar)
 - NC π^0 , de-excitation g, pion production, electron events
- T2K
 - CCQE cross sections
 - NC channels

Future Possibilities

- The current generation of experiments will tell us a lot, but flux uncertainties in 5-10% range probably best we can do
- In some cases, analysis techniques used to reduce sensitivity to flux uncertainties

1. New sources would go a long way

- nuSTORM $\delta\Phi(E) < 1\%$
- Narrow band beams

2. Fine-grained, low-density detectors

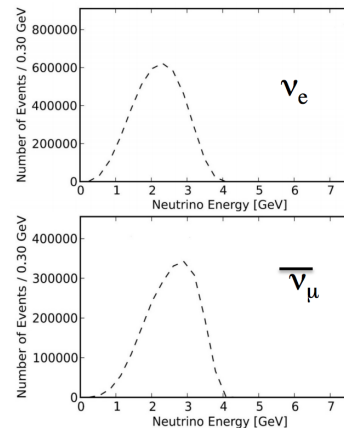
- The resolution and thresholds of LAr are fantastic, but we only measure on one pretty massive, non-isoscalar nucleus
- Coupling trackable hydrogen/deuterium detector (bubble chamber) with well known ν_μ/ν_e fluxes represents the ideal complex for untangling cross sections and nuclear effects

The nuSTORM Neutrino Beam

$$\mu^+ \rightarrow \bar{\nu}_\mu + \nu_e + e^+$$

$$\mu^- \rightarrow \nu_\mu + \bar{\nu}_e + e^-$$

- ◆ nuSTORM will provide a **very well-known** ($\delta\Phi(E) \leq 1\%$) beam of ν and $\bar{\nu}$.
- ◆ nuSTORM will provide a **high-intensity source of ν_e events!**



3.8 GeV μ^+ stored, 175m straight, flux at 50m

Jorge G. Morfin - Fermilab

μ^+		μ^-	
Channel	N_{evts}	Channel	N_{evts}
$\bar{\nu}_\mu$ NC	844,793	$\bar{\nu}_e$ NC	709,576
ν_e NC	1,387,698	ν_μ NC	1,584,003
$\bar{\nu}_\mu$ CC	2,145,632	$\bar{\nu}_e$ CC	1,784,099
ν_e CC	3,960,421	ν_μ CC	4,626,480

event rates per 1E21 POT -
100 tons at 50m

See WG2 talk from
J. Morfin on Wednesday

3

Motivations Summary

- Precision matters (P. Huber's talk from Monday)
- Systematics are key (M. Mezzetto's talk from Monday)
 - Uncertainty in interaction and nuclear models often largest systematic on oscillation parameters measured through appearance. Are we even accounting for all the things we don't know in these models?
- Accelerator-based neutrino experiments performed in a very tricky region for neutrino-nucleus interactions
 - $E_\nu \sim 1$ GeV off nuclear targets
 - Need to study both signal and background channels
 - Relationship between visible energy and the true neutrino energy has direct impact on extraction of oscillation parameters from data
- • Be prepared...

Conclusions

- Busy times in neutrino interaction physics
 - Significant progress in past year from new experiments MINERvA, ArgoNeuT and T2K
 - Past year's results focused on inclusive samples and understanding nuclear effects in neutrino scattering (especially QE)
- Guaranteed to be continued progress by NuFact 2014. See you there!

End