

Double Beta Decay Experiments



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Outline



part 1. Double Beta Decay

introduction
DBD and neutrino mass
NME
decay signature
detector choice
experimental sensitivity
background sources

part 2. DBD experiments:

Ge experiments
NEMO3
Cuoricino

part 3. DBD experiments: the new generation

on-going: GERDA - CUORE - EXO ...
proposed ...
conclusions

Part1.



Double Beta Decay and Majorana neutrino

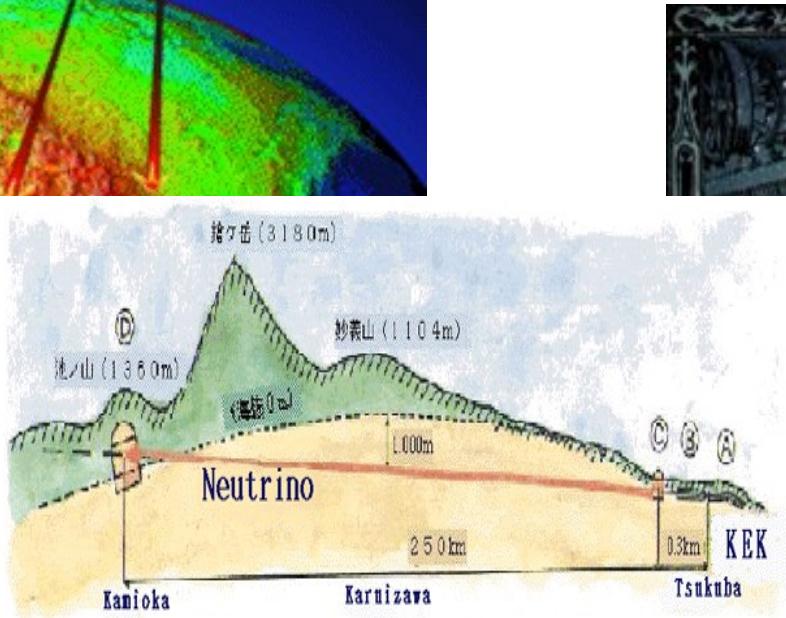
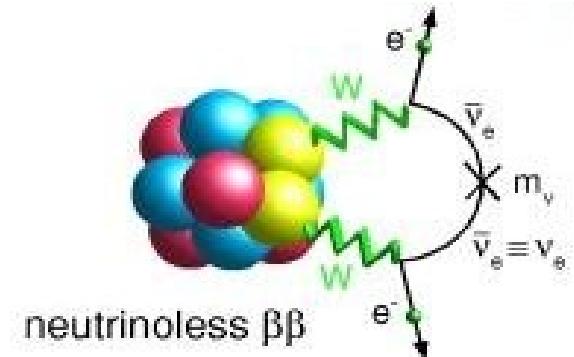
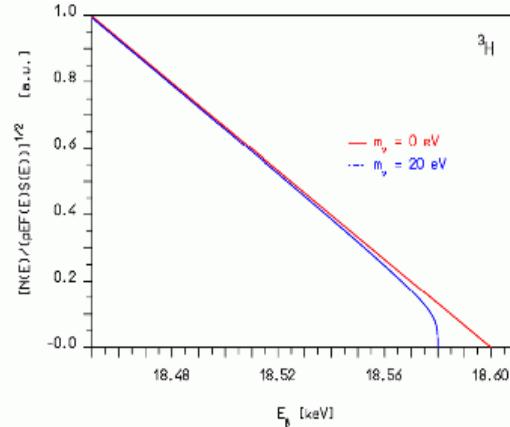
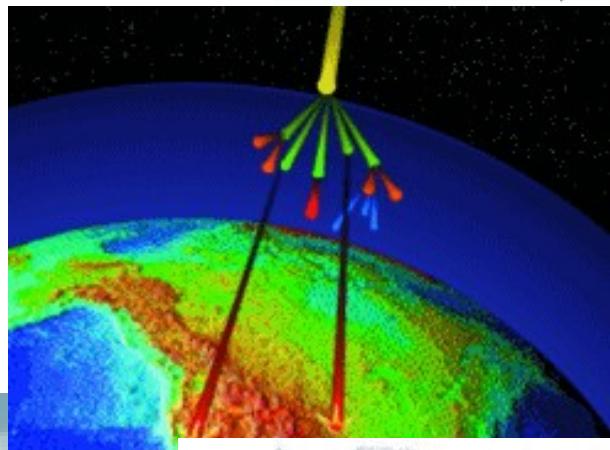
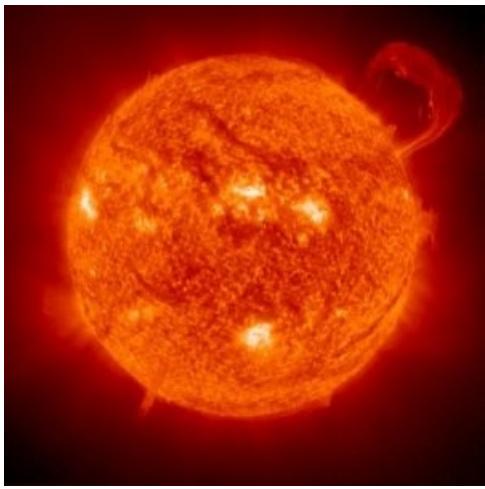
<< There are several categories of scientists in the world;
those of second or third rank do their best but never get
very far. Then there is the first rank, those who make
important discoveries, fundamental to scientific progress.
But then there are the geniuses, like Galilei and Newton.
Majorana was one of these.

Majorana had what no-one in the world had. Unfortunately he
lacked what other people generally have: common sense >>

(Enrico Fermi)



Introduction: the “neutrino industry”



the “neutrino industry”: proves mixing and masses

ν oscillations are a reality, we have distinguishable eigenvalues for:

flavor $\nu_e \nu_\mu \nu_\tau$

mass $\nu_1 \nu_2 \nu_3$

neutrino matrix is not diagonal:

$$\begin{pmatrix} \nu_e \\ \nu_\mu \\ \nu_\tau \end{pmatrix} = \begin{pmatrix} U_{e1} & U_{e2} & U_{e3} \\ U_{\mu 1} & U_{\mu 2} & U_{\mu 3} \\ U_{\tau 1} & U_{\tau 2} & U_{\tau 3} \end{pmatrix} \begin{pmatrix} \nu_1 \\ \nu_2 \\ \nu_3 \end{pmatrix}$$

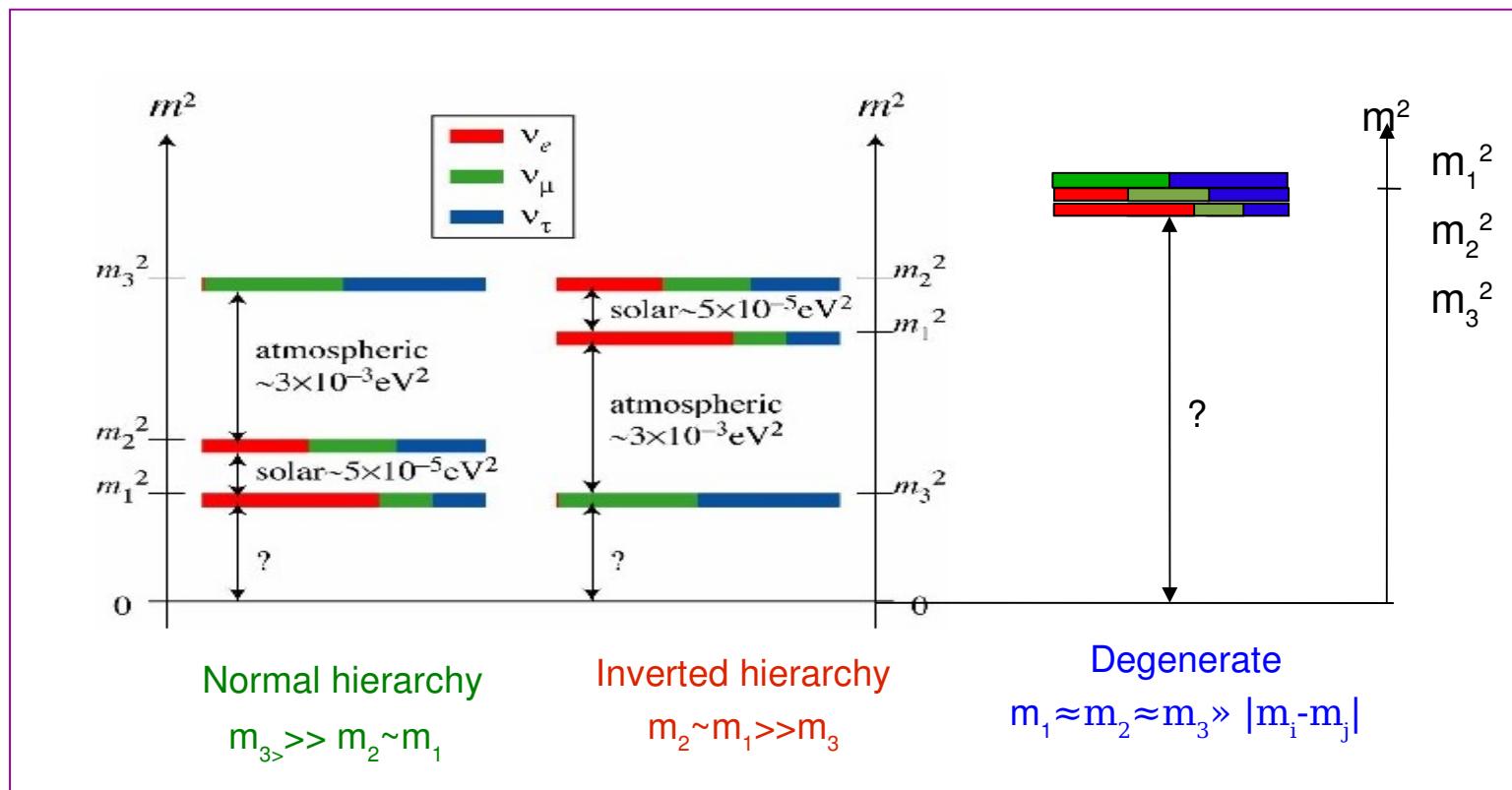
we can experimentally measure the values of the matrix elements

<i>Atmospheric + K2K</i>	<i>Reactors (CHOOZ)</i> <i>Accelerators (JPARC)</i>	<i>Solar + Reactors</i>
$U = \begin{pmatrix} 1 & 0 & 0 \\ 0 & \cos\theta_{23} & \sin\theta_{23} \\ 0 & -\sin\theta_{23} & \cos\theta_{23} \end{pmatrix} \times \begin{pmatrix} \cos\theta_{13} & 0 & e^{-i\delta_{CP}} \sin\theta_{13} \\ 0 & 1 & 0 \\ -e^{i\delta_{CP}} \sin\theta_{13} & 0 & \cos\theta_{13} \end{pmatrix} \times \begin{pmatrix} \cos\theta_{12} & \sin\theta_{12} & 0 \\ -\sin\theta_{12} & \cos\theta_{12} & 0 \\ 0 & 0 & 1 \end{pmatrix} \times \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\alpha/2} & 0 \\ 0 & 0 & e^{i\alpha/2+i\beta} \end{pmatrix}$		
$\theta_{23} \sim 45^\circ$	$\theta_{13} < \sim 13^\circ \quad \delta_{CP} = CP \text{ violation}$	$\theta_{12} \sim 30^\circ \quad \alpha, \beta : Majorana phase$

the “neutrino industry”: proves mixing and masses

... and measure Δm^2

$$\begin{array}{ccc} \nu_\mu & \rightarrow & \nu_\tau \\ \nu_e & \rightarrow & \nu_\mu \end{array}$$
$$\Delta m^2_{23} \sim 2 \cdot 10^{-3} \text{ eV}^2$$
$$\Delta m^2_{12} \sim 5 \cdot 10^{-5} \text{ eV}^2$$



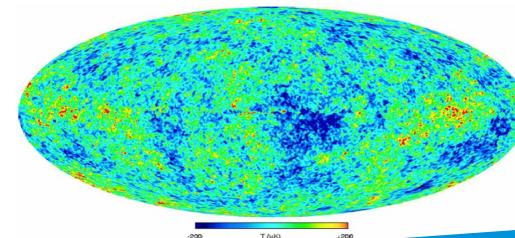
"how to measure the mass?"

Cosmology

$$\Sigma = m_1 + m_2 + m_3$$



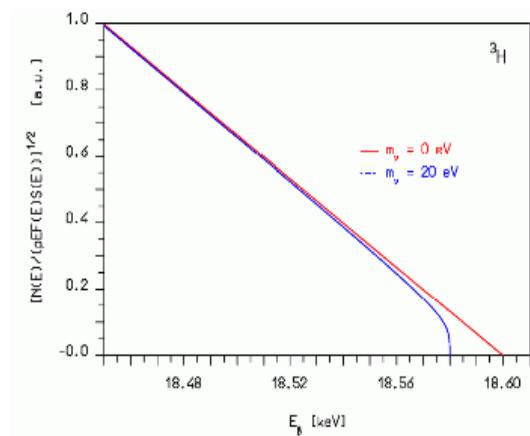
simple sum
pure kinematic effect



Beta decay

$$m_{\nu e} = (\sum |U_{ei}|^2 m_i^2)^{1/2}$$

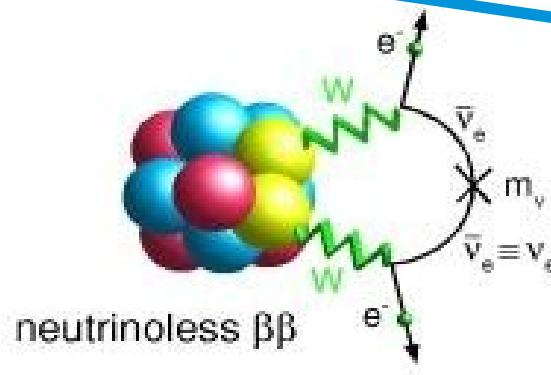
incoherent sum
real neutrino



Double Beta Decay

$$|\langle m_\nu \rangle| = |\sum U_{ei}^2 m_i|$$

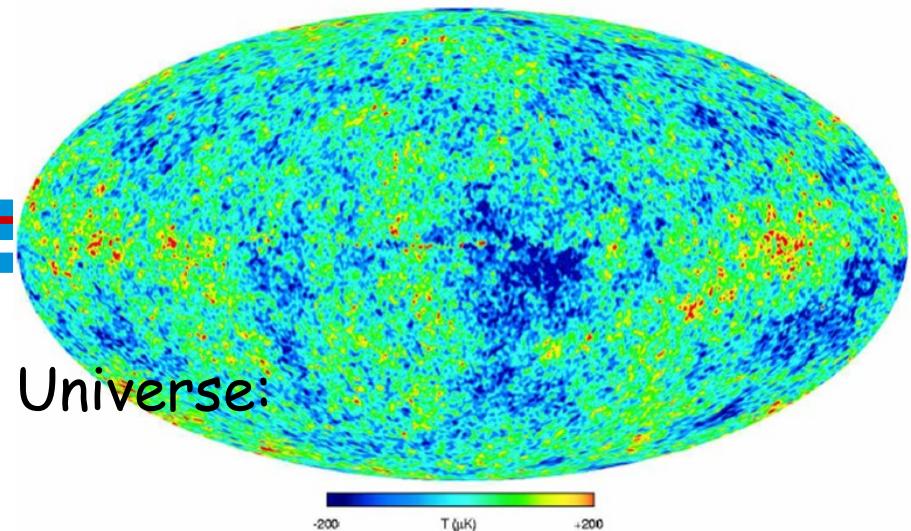
coherent sum
virtual neutrino
Majorana phases



Astrophysical bounds:

ν contribute to the energy density of our Universe:

$$\Omega_\nu h^2 \sim \Sigma / 94 \text{ (eV)}$$



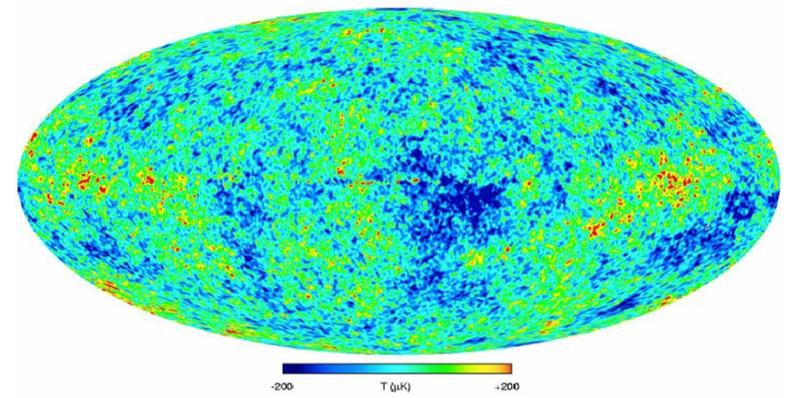
they influence various astronomical observables:

- CMB (Cosmic Microwave Background) power spectrum
- LSS (Large Scale Structures) matter distribution reconstructed by redshift survey of galaxies
- Ly- α (Lyman α) distant quasars light absorbed at Ly- α frequency by intervening matter
- ...

BUT

the upper bounds on Σ depend on data set included, priors and statistical treatments

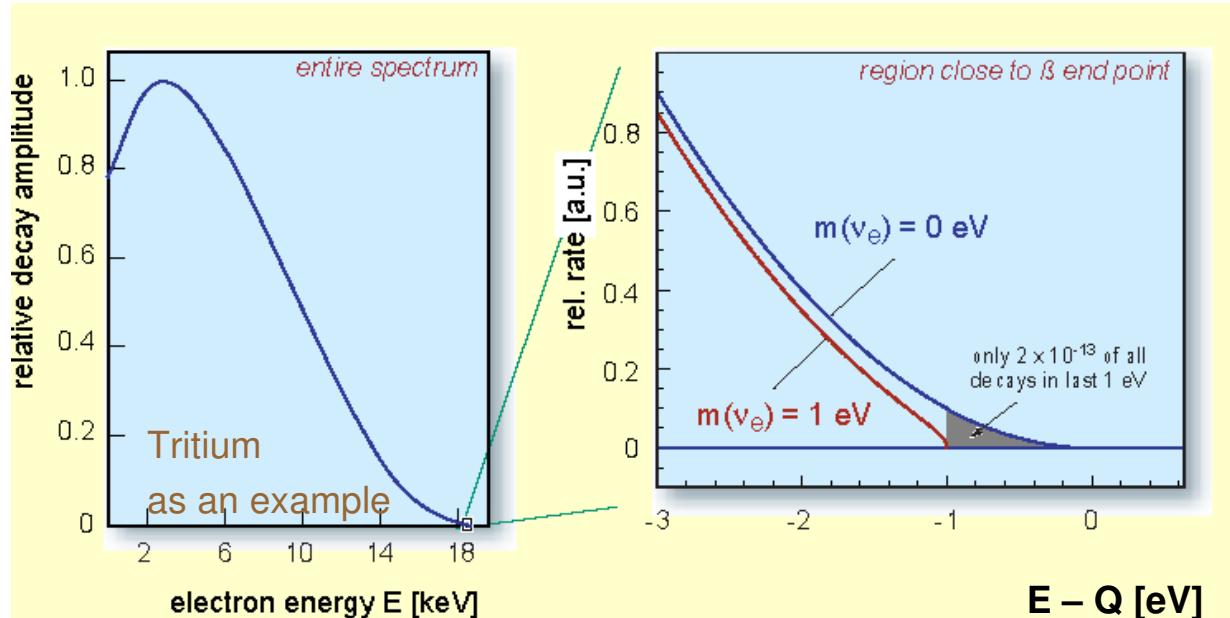
Astrophysical bounds:



- (1) $\Sigma < 1.3$ eV (WMAP5)
- (2) $\Sigma < 1.19$ eV (WMAP5+ACBAR+VSA+CBI+BOOMERANG)
- (3) $\Sigma < 0.75$ eV (CMB+HST+SN-Ia)
- (4) $\Sigma < 0.60$ eV (CMB+HST+SN-Ia+BAO)
- (5) $\Sigma < 0.19$ eV (CMB+HST+SN-Ia+BAO+Ly)

[Fogli et al., arXiv:0805.2517]

Direct measurements: deformation of Kurie plot



Tiny effect due to finite neutrino mass:

$$\sim (m_\beta c^2/Q)^3$$

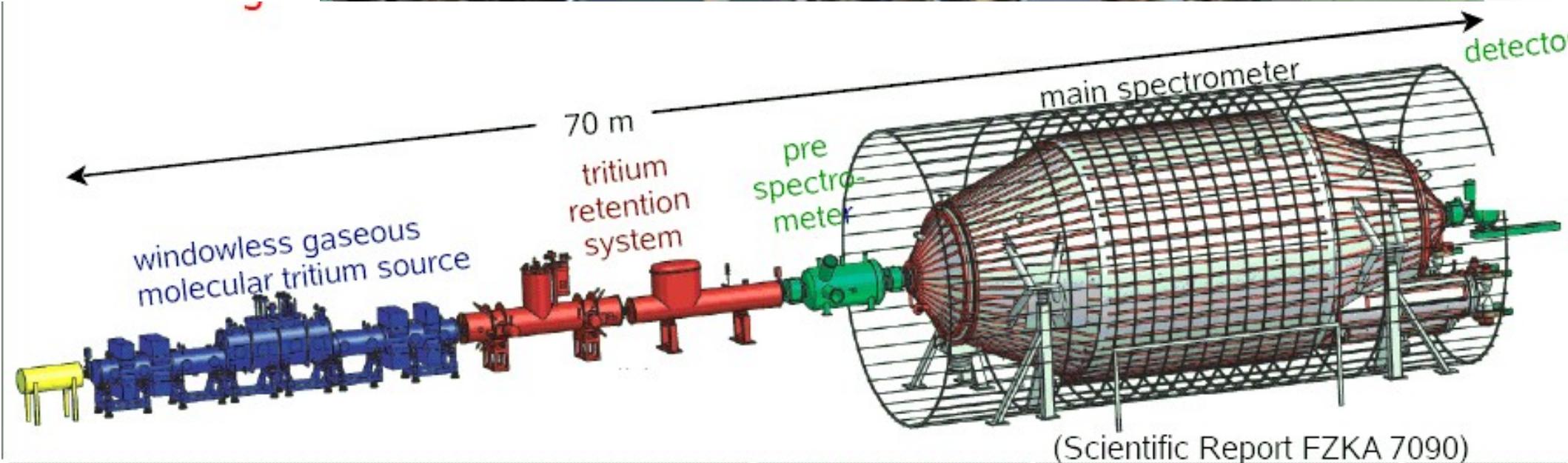
Mass sensitivity:

$$m_\beta \sim \{Q^3 \Delta E / A T\}^{1/4}$$

- | | |
|---------------|---|
| Present limit | = 2.3 eV (Mainz, Troitzk spectrometers) |
| Near future | = 0.2 eV (Katrean spectrometer) |
| Far future | = beyond 0.1 eV with bolometers ??? |

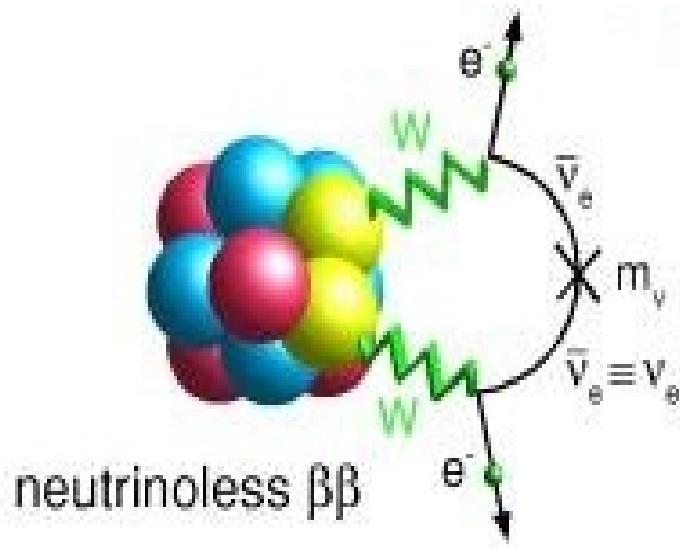
Bolometers use ^{187}Re ($Q=2.47 \text{ keV } \tau=43 \text{ Gyr}$)
to go beyond 0.2 eV
they require technological improvement
accurate calculation of ^{187}Re beta spectrum

KATRIN



(Scientific Report FZKA 7090)

Direct measurement: neutrinoless double beta decay



Phase space factor

Nuclear matrix element

$$T_{1/2}^{-1} = F(Q_{\beta\beta}^5, Z) |M|^2 \langle m_\nu \rangle^2$$

Effective Majorana mass

$0\nu\beta\beta$ allow to study ν properties and measure ν mass
BUT

ν mass can be extracted from $T_{1/2}$ ONLY after $|M|^2$

mass sensitivities compared: TODAY

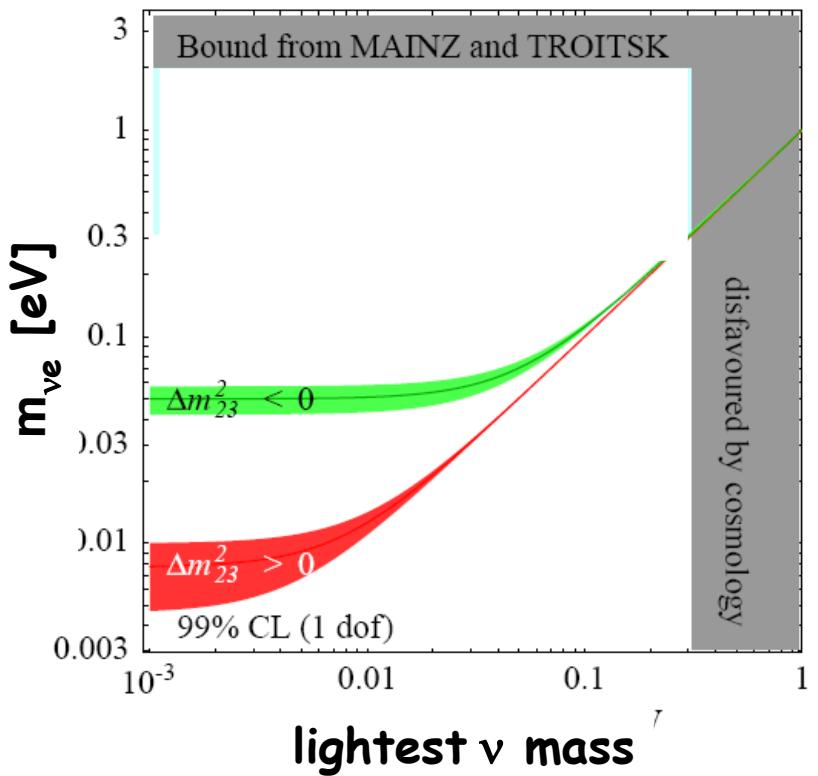
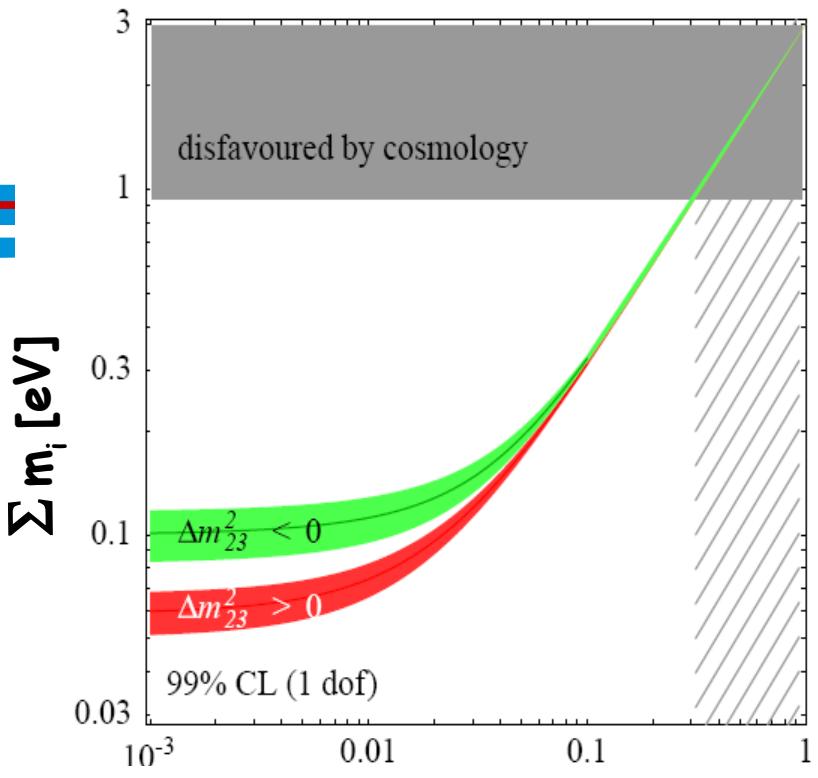
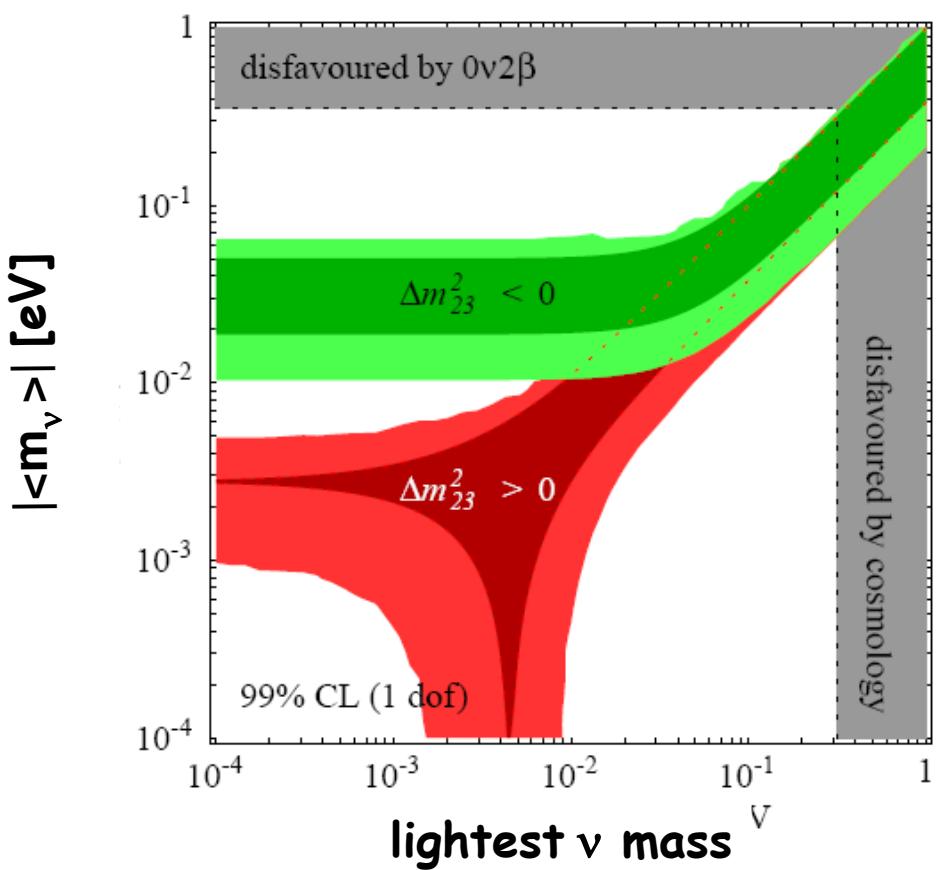
Cosmology

$$\sum m_i = m_1 + m_2 + m_3 \leq 1 \text{ eV}$$

Beta decay

$$m_{ve} = (\sum |U_{ei}|^2 m_i^2)^{1/2} < 2.3 \text{ eV}$$

Double Beta Decay $|\langle m_\nu \rangle| = |\sum U_{ei}^2 m_i| \leq 2.3 \text{ eV}$



mass sensitivities compared: TOMORROW

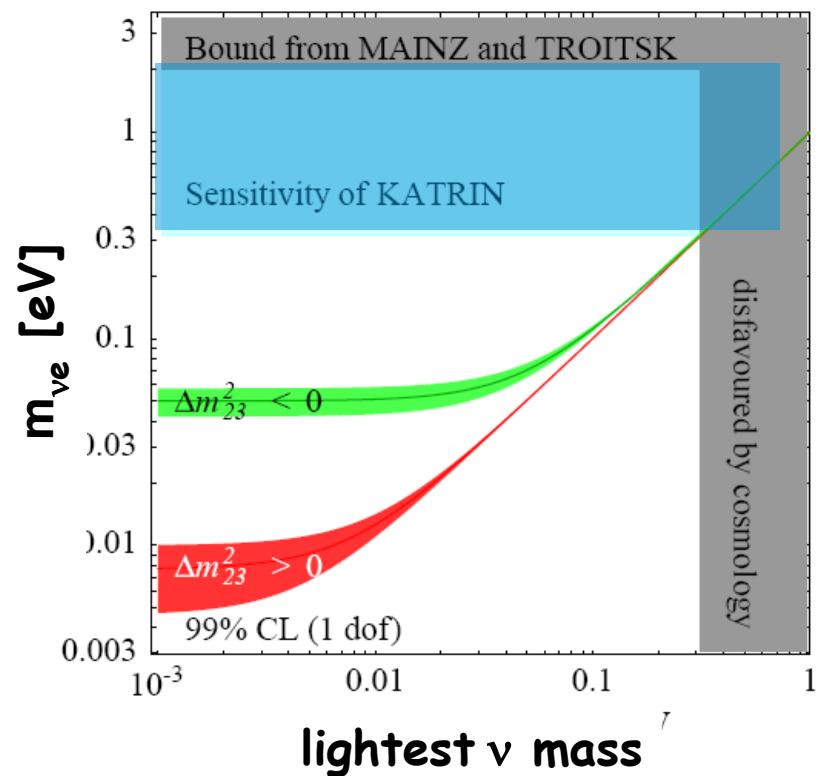
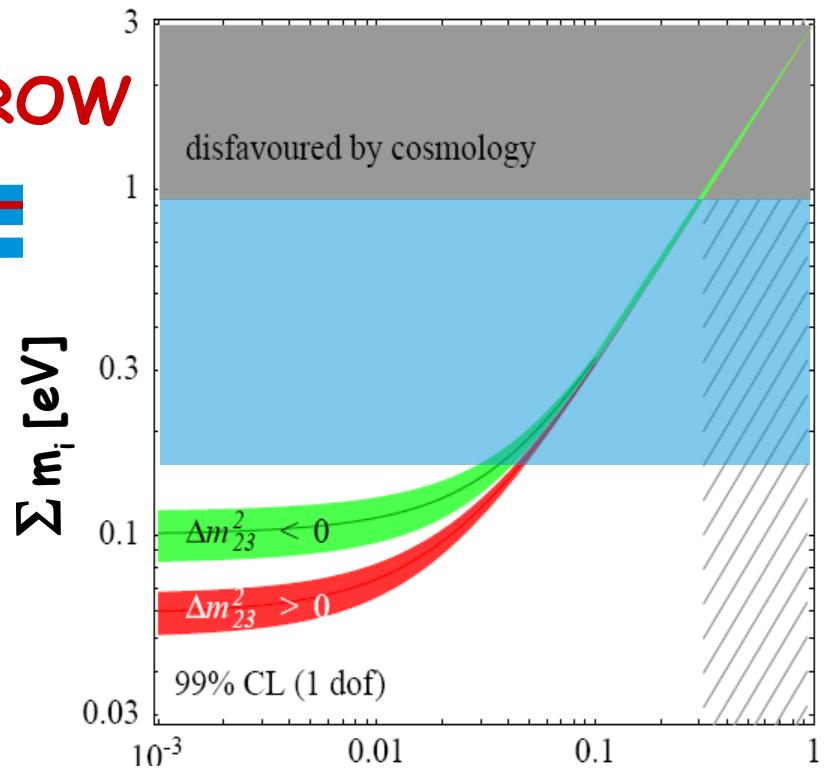
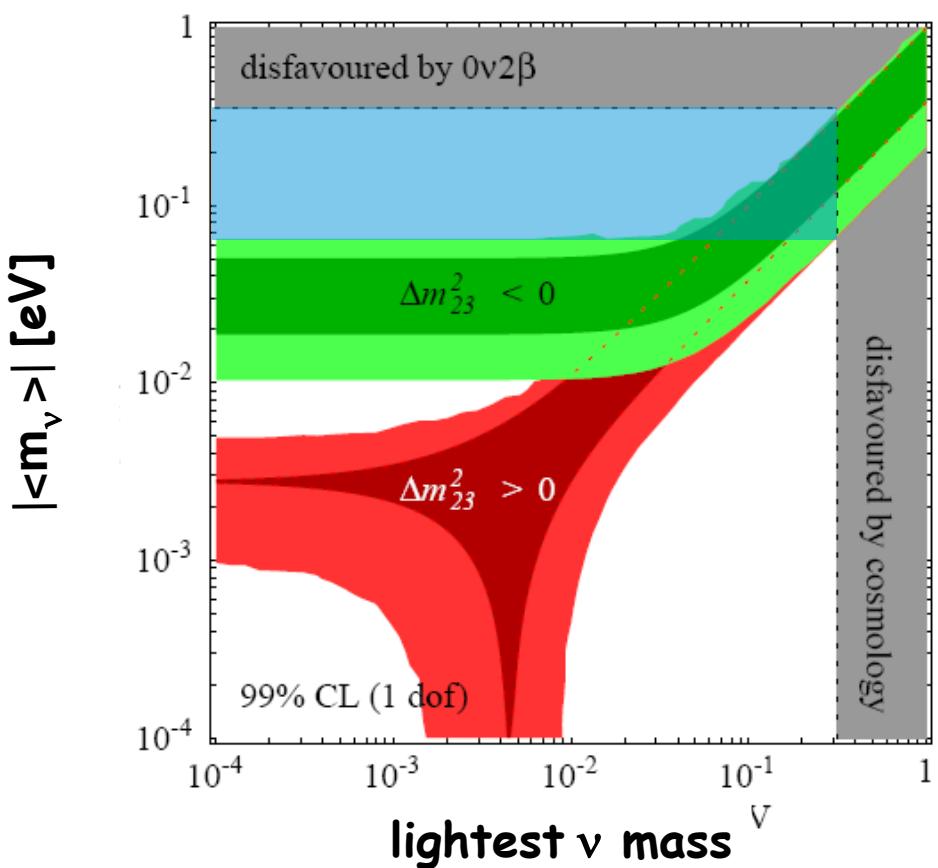
Cosmology

$$\sum m_i = m_1 + m_2 + m_3 \leq 0.1 \text{ eV}$$

Beta decay

$$m_{ve} = (\sum |U_{ei}|^2 m_i^2)^{1/2} < 0.3 \text{ eV}$$

Double Beta Decay $|\langle m_\nu \rangle| = |\sum U_{ei}^2 m_i| \leq 0.05 \text{ eV}$



Tools for the investigation of the mass scale

Tools	Present sensitivity	Future sensitivity (a few year scale)
Cosmology (CMB + LSS)	0.7 - 1 eV	0.1 eV
Neutrinoless Double Beta Decay	0.5 eV	0.05 eV
Single Beta Decay	2.2 eV	0.2 eV

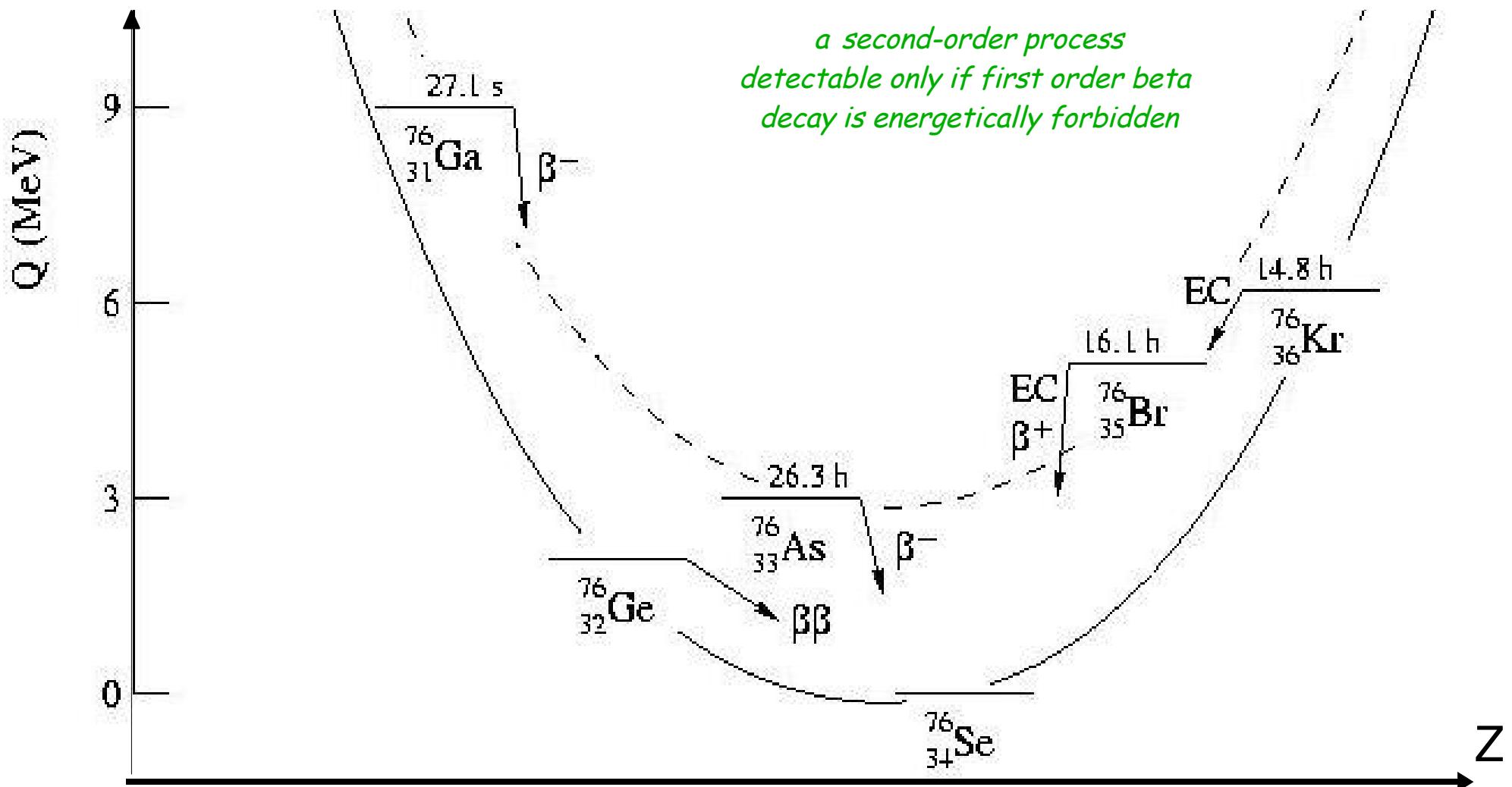
Model dependent (Orange double-headed arrow)

Laboratory measurements (Blue double-headed arrow)

Direct determination (Green double-headed arrow)

DBD and neutrino mass

even(A)-even(Z) nuclei whose single beta decay is energetically forbidden can however "double beta decay" on a lower mass isomer

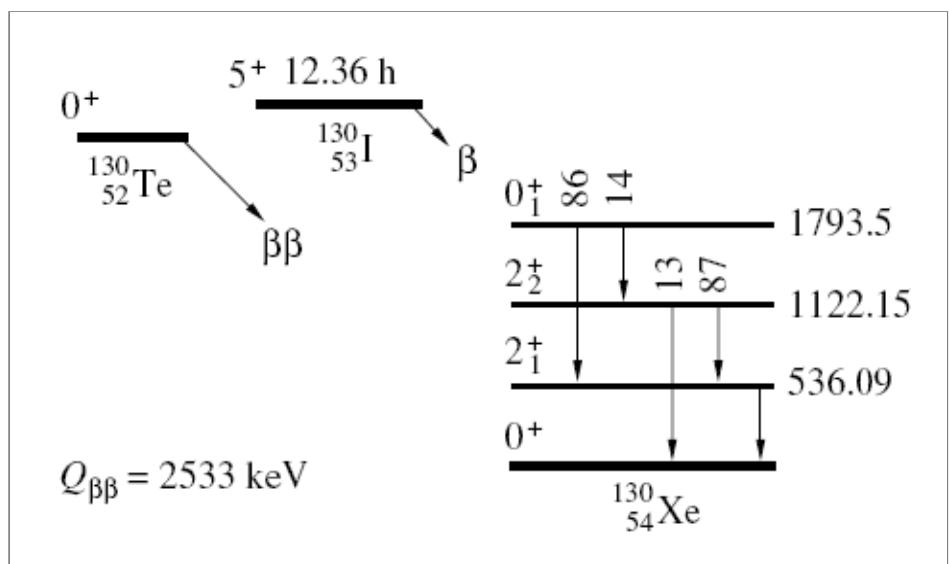


candidate nuclei with $Q > 2$ MeV

Candidate	Q (MeV)	Abund. (%)
$^{48}\text{Ca} \rightarrow ^{48}\text{Ti}$	4.271	0.187
$^{76}\text{Ge} \rightarrow ^{76}\text{Se}$	2.040	7.8
$^{82}\text{Se} \rightarrow ^{82}\text{Kr}$	2.995	9.2
$^{96}\text{Zr} \rightarrow ^{96}\text{Mo}$	3.350	2.8
$^{100}\text{Mo} \rightarrow ^{100}\text{Ru}$	3.034	9.6
$^{110}\text{Pd} \rightarrow ^{110}\text{Cd}$	2.013	11.8
$^{116}\text{Cd} \rightarrow ^{116}\text{Sn}$	2.802	7.5
$^{124}\text{Sn} \rightarrow ^{124}\text{Te}$	2.228	5.64
$^{130}\text{Te} \rightarrow ^{130}\text{Xe}$	2.533	34.5
$^{136}\text{Xe} \rightarrow ^{136}\text{Ba}$	2.479	8.9
$^{150}\text{Nd} \rightarrow ^{150}\text{Sm}$	3.367	5.6

all these nuclei decay by $\beta^-\beta^-$ emissions
but also candidates decaying
via $\beta^+\beta^+$ and EC EC do exist

the decay can be either on the 0^+ ground state
or on the excited state/s



decay processes

there is one SM allowed process:

$2\nu\beta\beta$



two electrons sum energy $E_1+E_2 < Q$

and various energetically allowed but not standard 0ν processes:

$0\nu\beta\beta$ through $\langle m \rangle$ or RH currents or



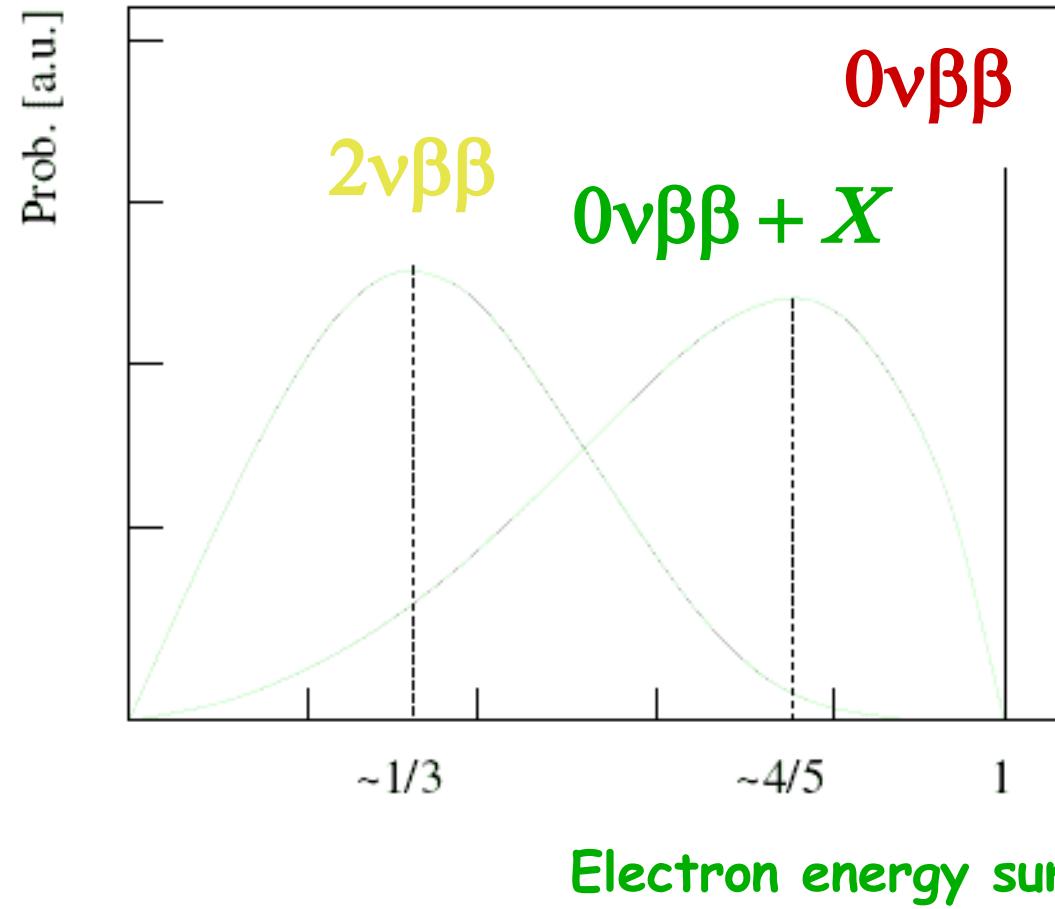
two electrons sum energy $E_1+E_2 = Q$

$0\nu\beta\beta +$ Majoron



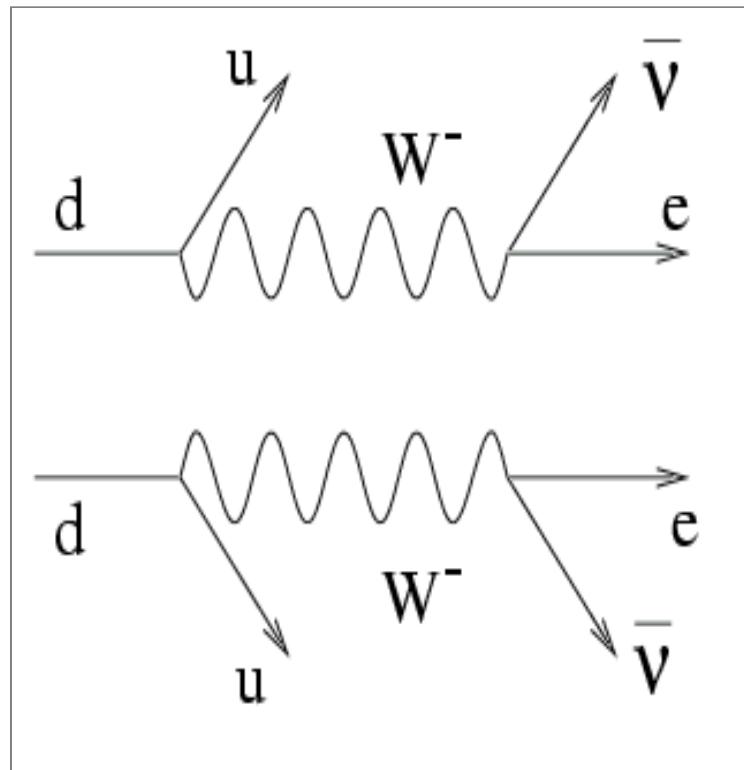
two electrons sum energy $E_1+E_2 < Q$

...



two neutrino decay

$$(A, Z) \rightarrow (A, Z+2) + 2 e^- + 2 \bar{\nu}$$



this process is studied to get info concerning the NME

Isotope	$Q\beta\beta$ (MeV)	Isotopic abundance (%)	Half life $T_{1/2}, 2\nu$ (y)
^{48}Ca	4.27	0	$\sim 4.0 \cdot 10^{19}$
^{76}Ge	2.04	7.8	$\sim 1.4 \cdot 10^{21}$
^{82}Se	3	9.2	$\sim 0.9 \cdot 10^{20}$
^{96}Zr	3.35	2.8	$\sim 2.1 \cdot 10^{19}$
^{100}Mo	3.03	9.6	$\sim 8.0 \cdot 10^{18}$
^{116}Cd	2.8	7.5	$\sim 3.3 \cdot 10^{19}$
^{128}Te	0.87	31.7	$\sim 2.5 \cdot 10^{24}$
^{130}Te	2.53	34.5	$\sim 0.9 \cdot 10^{21}$
^{136}Xe	2.48	8.9	not observed yet
^{150}Nd	3.37	5.6	$\sim 7.0 \cdot 10^{18}$

neutrinoless decay: $(A, Z) \rightarrow (A, Z+2) + 2e^-$

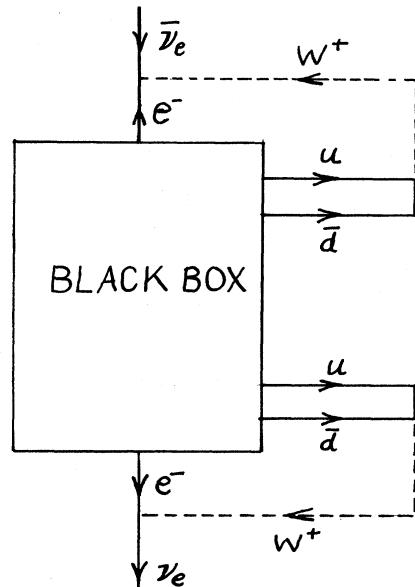


FIG. 2. Diagram showing how any neutrinoless double- β decay process induces a $\bar{\nu}_e$ -to- ν_e transition, that is, an effective Majorana mass term.

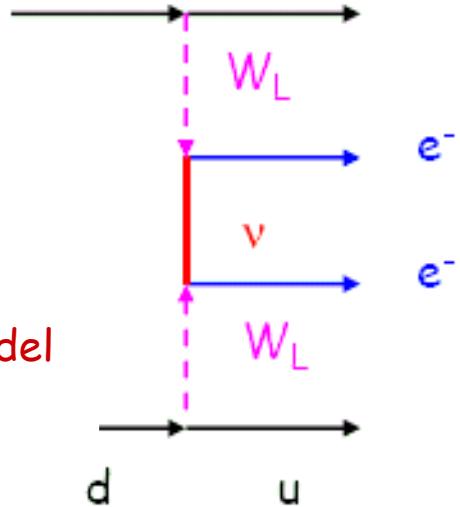
Whatever processes cause $0\nu\beta\beta$,
its observation
would imply the existence of
a Majorana mass term"

Schechter and Valle,
Phys. Rev. D, Vol. 25 N.11, 1982

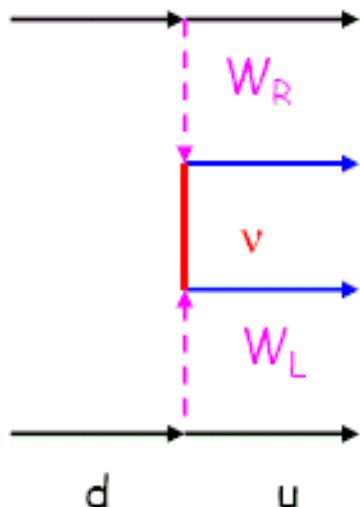
Hence observing the $0\nu\beta\beta$ decay guarantees that ν are massive Majorana particles.

neutrinoless decay: $(A, Z) \rightarrow (A, Z+2) + 2e^-$

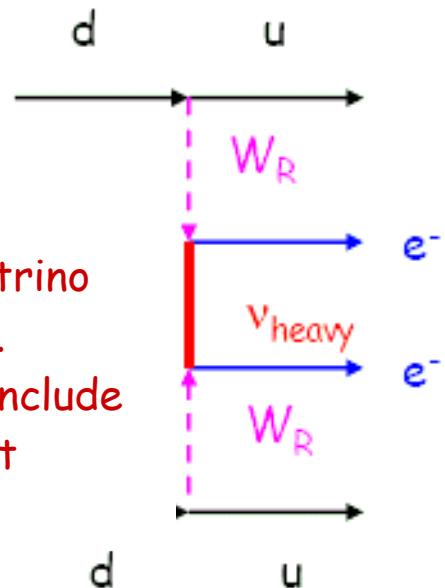
(I)
Light Majorana neutrino,
only Standard Model
weak interactions



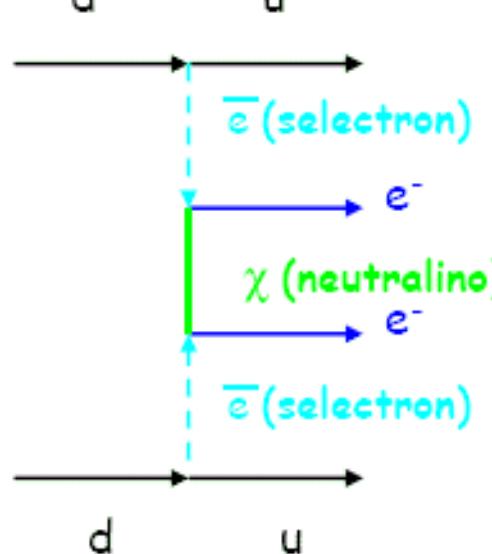
(II)
Light or heavy Majorana neutrino. Model extended to include right-handed WR.
Mixing extended between the left and right-handed neutrinos.



(III)
Heavy Majorana neutrino interacting with WR.
Model extended to include right-handed current interactions.



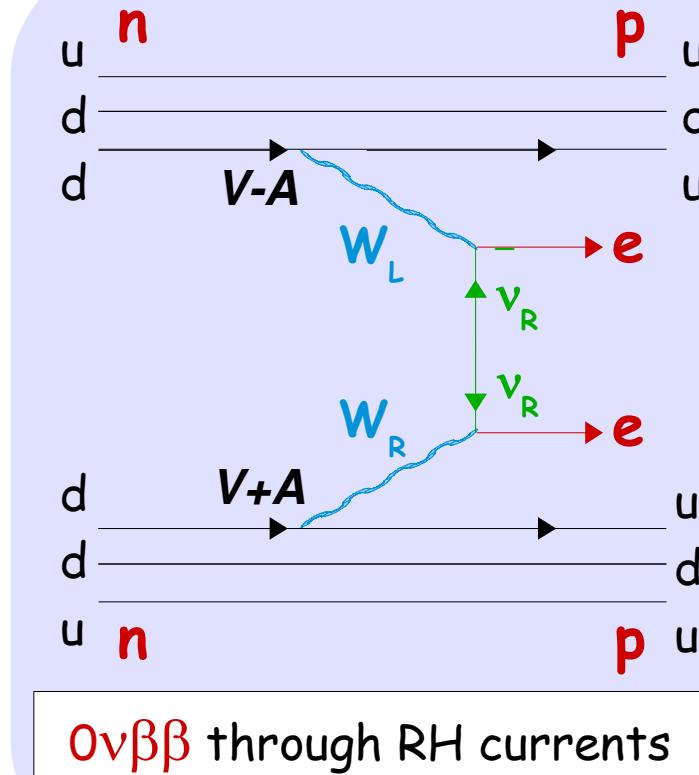
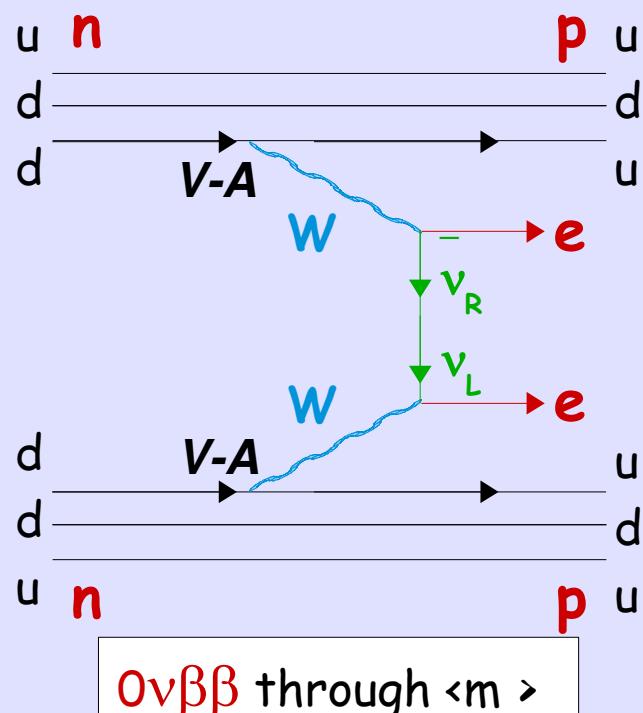
(IV)
Supersymmetry with R-parity violation. Many new particles invoked. Light Majorana neutrinos exist also.



$$(A, Z) \rightarrow (A, Z+2) + 2e^-$$

the observation of
this decay proves:

- Lepton number violation
- Majorana character of neutrinos
- measurement of the Majorana mass of neutrino
- infer the absolute mass scale and hierarchy
- test existence of RH currents

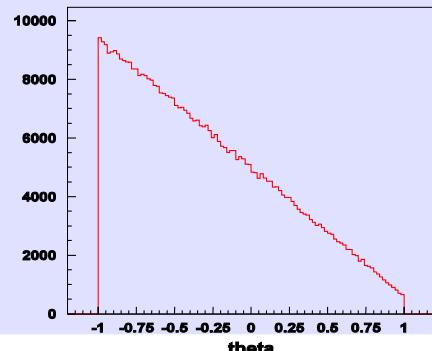
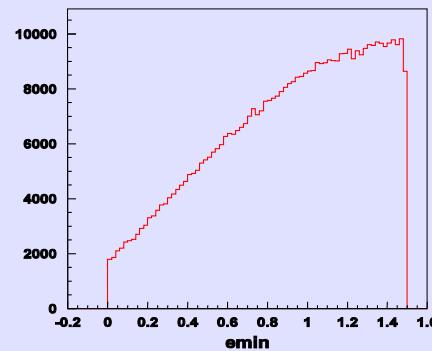




Light neutrino exchange

$0\nu\beta\beta$ through $\langle m_\nu \rangle$

$$T_{1/2}^{-1} = F(Q_{\beta\beta}, Z) |M|^2 \langle m_\nu \rangle^2$$



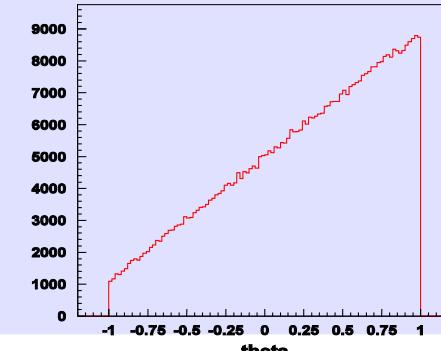
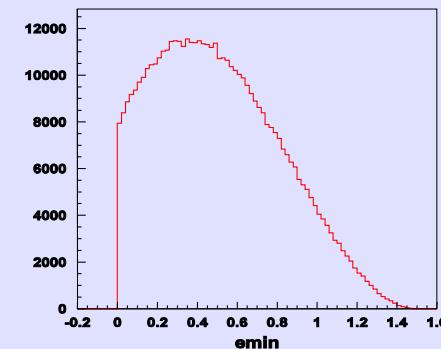
Minimum electron energy

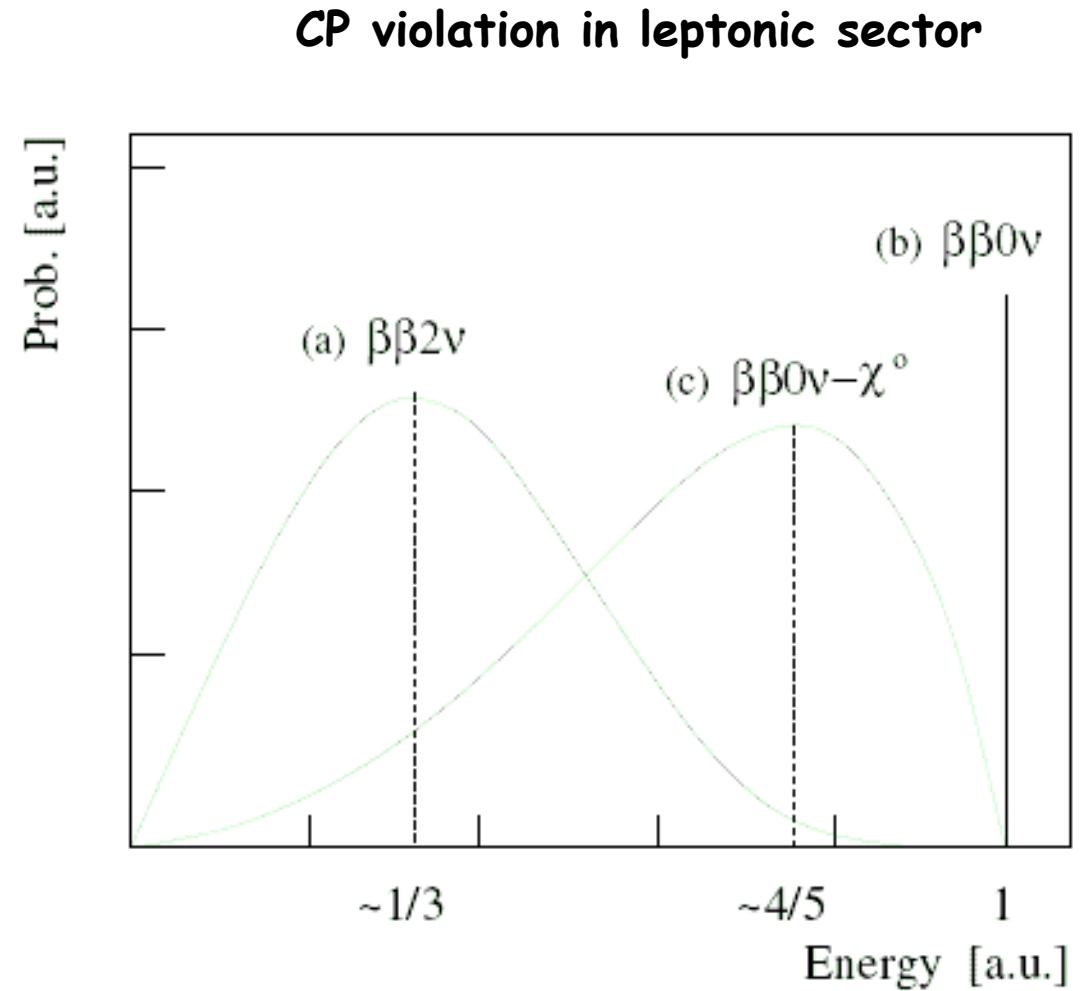
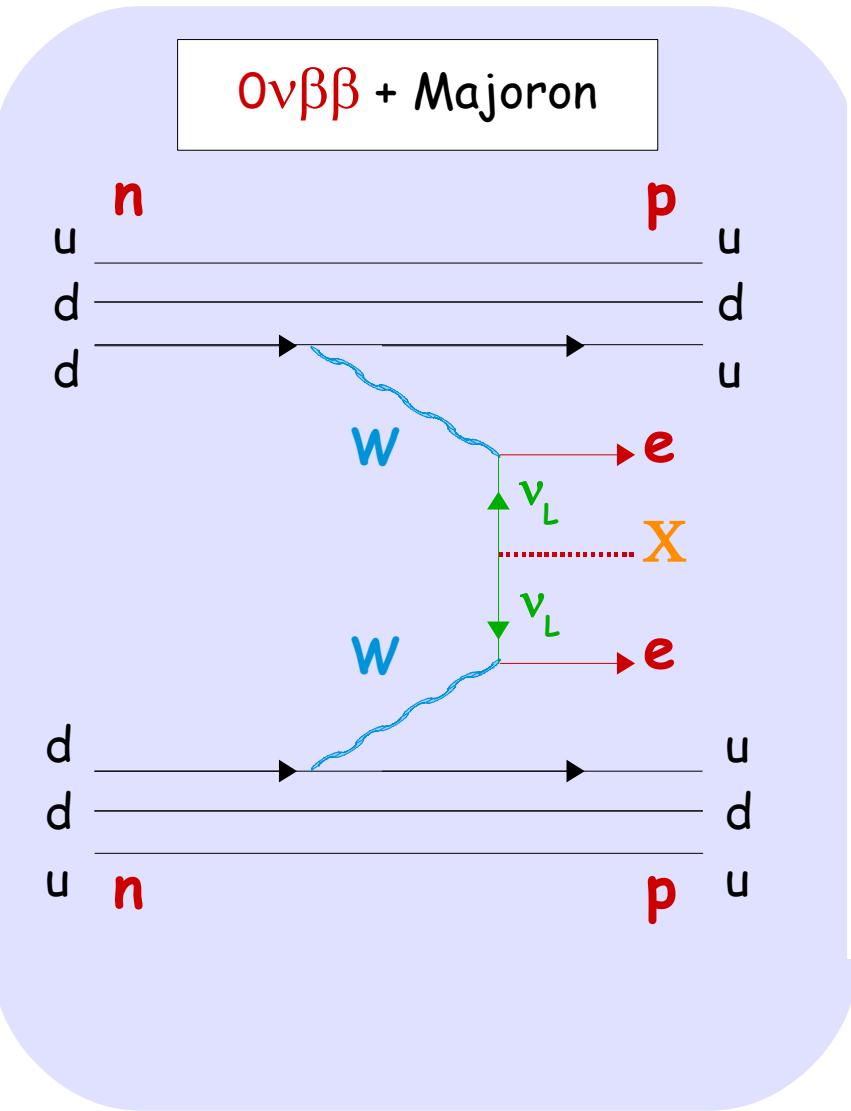
Angular distribution between the 2 electrons

V+A current

$0\nu\beta\beta$ through RH currents

$$T_{1/2}^{-1} = C_1 \langle m_\nu \rangle^2 + f(\langle m_\nu \rangle, \lambda, \eta)$$





M. Pavan, Sept. 2008, Beijing "DBD Experiments"

$0\nu\beta\beta$ test lepton number conservation

A partial list of processes where the lepton number would be violated:

Ov $\beta\beta$ decay:	$(Z, A) \rightarrow (Z \pm 2, A) + 2e^{(\pm)}$	$T_{1/2} > \sim 10^{25}$ y
Muon conversion:	$\mu^- + (Z, A) \rightarrow e^+ + (Z-2, A)$	$BR < 10^{-12}$
Anomalous kaon decays:	$K^+ \rightarrow \pi^- \mu^+ \mu^+$	$BR < 10^{-9}$
Flux of ν_e from the Sun:		$BR < 10^{-4}$
Flux of ν_e from a nuclear reactor:		$BR < ?$

Observing any of these processes would mean that the lepton number is not conserved, and that neutrinos are massive Majorana particles.

It turns out that the study of the $0\nu\beta\beta$ decay is by far the most sensitive test of the total lepton number conservation.

$0\nu\beta\beta$ and neutrino mass (assume diagram I dominates)

$0\nu\beta\beta$ life-time is written as:

$$T_{1/2}^{-1} = F(Q_{\beta\beta}, Z) |M|^2 \langle m_\nu \rangle^2$$

$F(Q_{\beta\beta}, Z)$ = phase space
decay cinematic

$|M|^2$ = nuclear matrix element
theoretically evaluated (shell model, QRPA models ...)
different results according to the nuclear model used

$$|\sum U_{ei}^2 m_i|^2 = \langle m_\nu \rangle^2 = \frac{1}{T_{1/2} F(Q_{\beta\beta}, Z) |M|^2}$$

- $T_{1/2}^{-1}$ is experimentally measured → neutrino is Majorana particle
- $\langle m_\nu \rangle^2$ is extracted from the experimental value → mass absolute scale

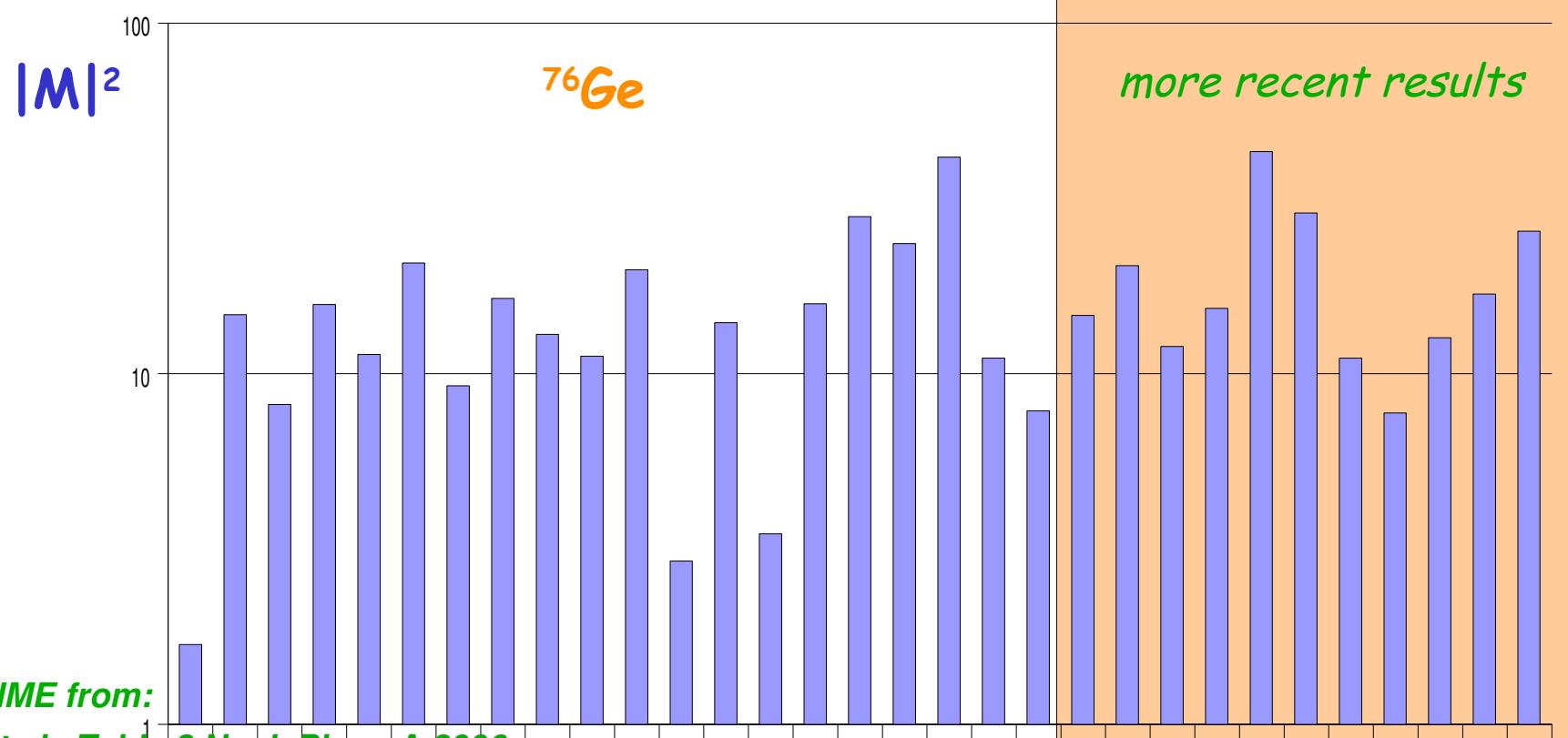
however this ($\langle m_\nu \rangle^2$) result is model (NME) dependent !!!

Nuclear Matrix Element

Nuclear matrix elements are calculated using two models:

QRPA (RQRPA, SQRPA,), Shell model ...

with sometimes (particularly in the past) quite different results



QRPA NME from:

Rodin et al. Table 3 Nucl. Phys. A 2006

+ erratum nucl-th:0706.4304v1

selection of QRPA result

Nuclear Matrix Element



suggestion from Bahcall et al.

use the nuclear matrix range as an uncertainty: « Democratic approach »

BUT

does not take into account the improvements of the Models

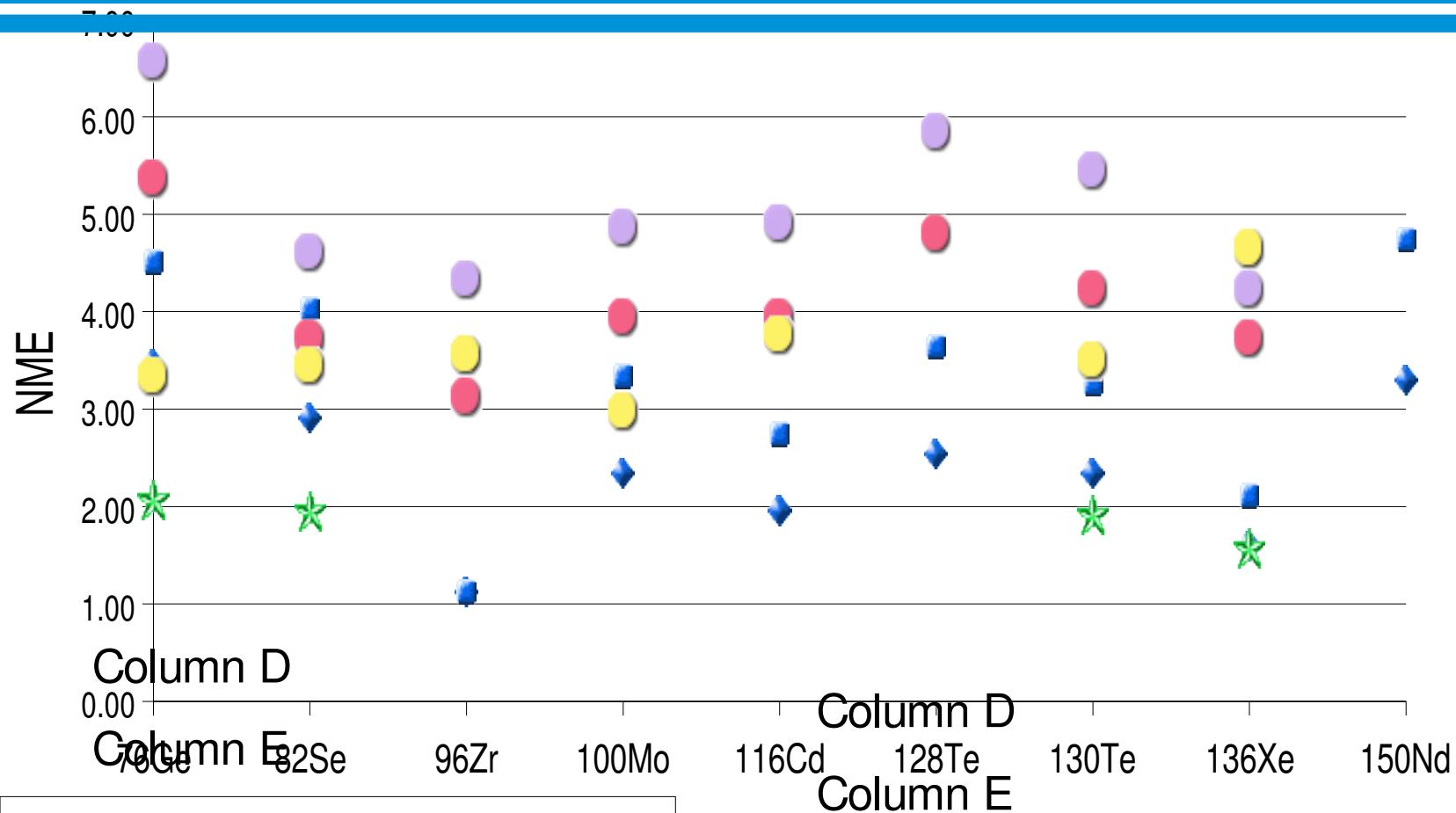
does not help in the choice of the best candidate for an experiment

TODAY

exchanges between groups to understand discrepancies and evaluate errors
use of β and $2\nu\beta\beta$ decay data to fix parameters in QRPA, new results from
Shell Model

new results are more similar than in the past !!!

$0\nu\beta\beta$: Nuclear Matrix Evaluation



QRPA upper and lower values in Rodin Faessler Simkovic Vogel, nucl-th:0706.4304v1

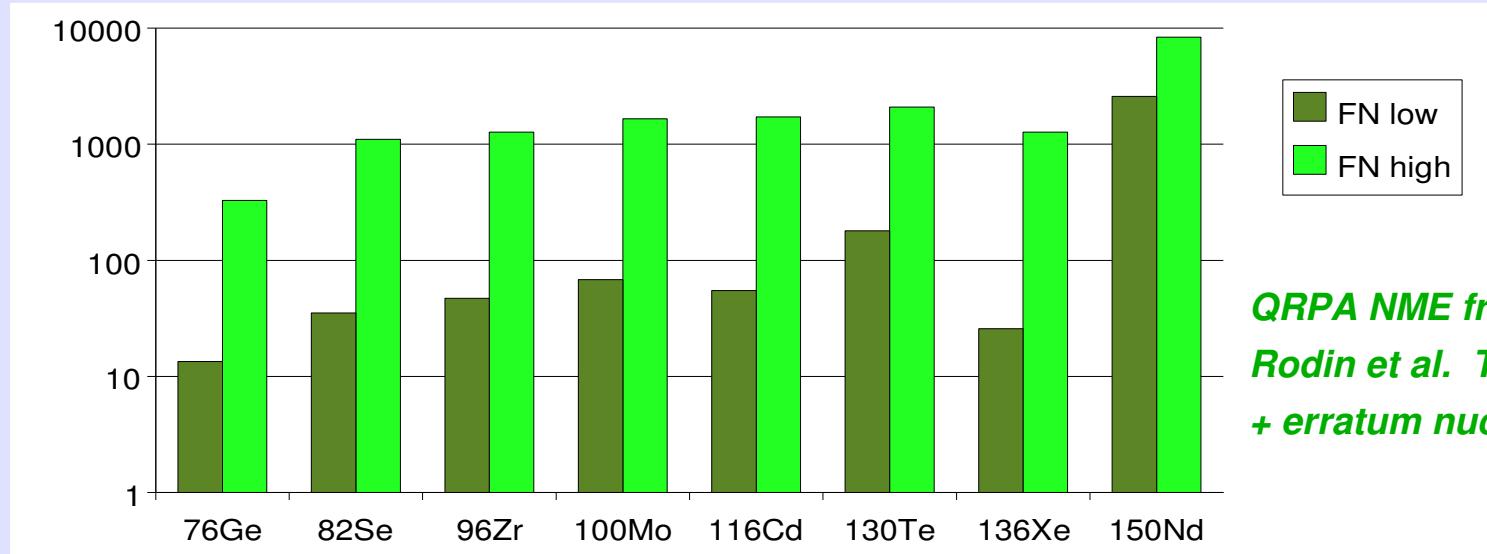
Nuclear Shell model
Caurier et al. nucl-th/0801.3760

QRPA upper and lower values in Kortelainen Suhonen 2007, Phys. Rev C

QRPA Civitarese Shuonen
nucl-th/0208005

extracting m_ν from $T_{1/2}^{0\nu\beta\beta}$

$$\langle m_\nu \rangle^2 = \frac{1}{T_{1/2} F(Q_{\beta\beta}, Z) |M|^2} \quad \rightarrow F_N \text{ nuclear factor of merit}$$



QRPA NME from:
Rodin et al. Table 3 Nucl. Phys. A 2006
+ erratum nucl-th:0706.4304v1

→ m_ν range: $\left(\frac{1}{T_{1/2}^{0\nu\beta\beta} F_{N\text{high}}} ; \frac{1}{T_{1/2}^{0\nu\beta\beta} F_{N\text{low}}} \right)$

but which selection of NME values should be used? my solution for these slides was
to quote the more recent results of the 2 QRPA "schools" and of the SM

Decay signature



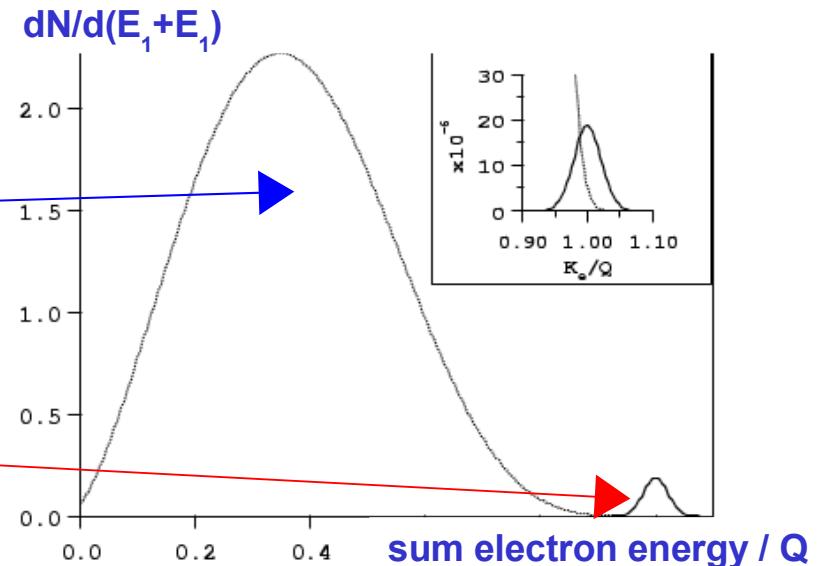
search for the decay looking for the 2 e-



signature = the sum energy spectrum of the 2 e- is:

$2\nu 2\beta =$
a continuum with end point at Q-value

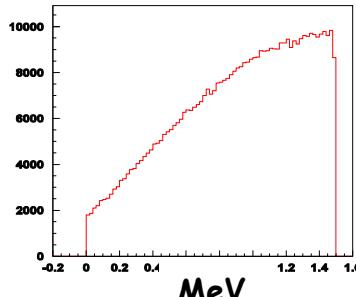
$0\nu 2\beta =$
a peak at Q-value, enlarged by det. res.



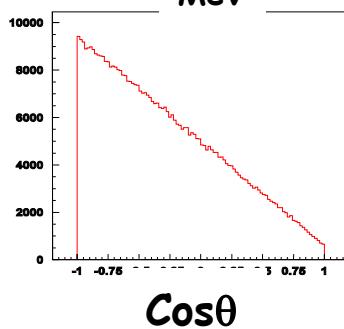
the ratio between the two half-life depends on the isotope !!

other observables

Light neutrino exchange

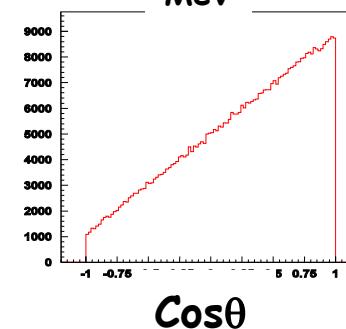
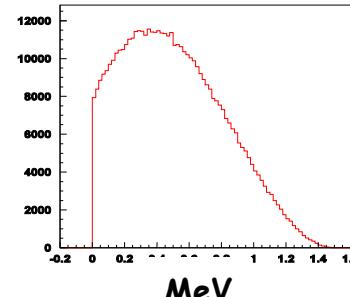


Minimum electron
energy



Angular distribution
between the 2 electrons

V+A current



Available only in
tracking experiments

further signatures:

the signal is due to electrons = charge -1

short range particle "single site"

the daughter nucleus

Novel techniques
under development

what does happen if you find the signal?

in most experiments the signature of the decay will be rather poor, it is unlikely - but possible - that background event can mimic all the observables of Ov $\beta\beta$

it is commonly accepted by the " $\beta\beta$ community" that the discovery of Ov $\beta\beta$ would require:

the decay shows up

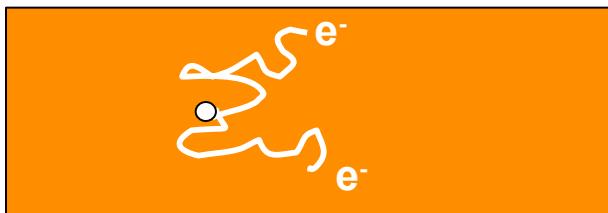
in more than one experiment
for more than one isotope

the comparison of experiments studying the same isotope is trivial
(so it is possible to clearly state that a certain result is confirmed or not,
however this does not prove that the observed events are due to Ov $\beta\beta$ and
not to some background source)

in the case of different isotopes the comparison goes through the NME
(only observing the decay in different isotopes we can be reasonably
confident that Ov $\beta\beta$ decay have been discovered)

Detector choice

*two electrons each with a continuous spectrum
and a monochromatic sum energy*



Source ≡ Detector
(calorimetric technique)

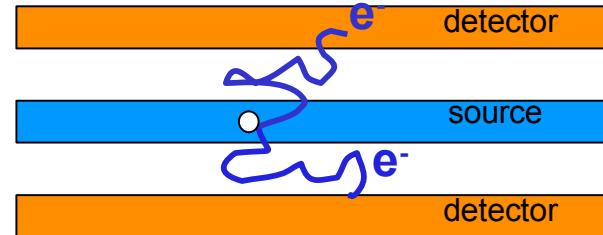
*semiconducting devices Ge, CdTe
scintillators*

solid CaF_2 , CdWO_4

liquid/gas Xe

bolometers

TeO_2 , CdWO_4 , CaF_2 , MoWO_4 ...

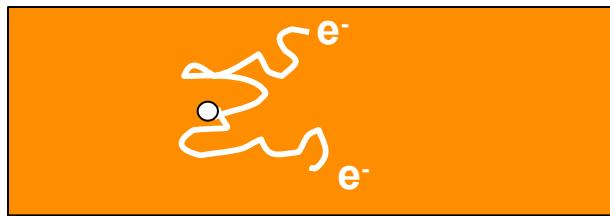


Source ≠ Detector

*very thin sheets containing the source
ionization detectors for tracking and
energy measurement*

almost any candidate

source = detector



Source = Detector
(calorimetric technique)



constraints on detector materials



very large masses are possible
demonstrated: up to ~ 50 kg
proposed: up to ~ 1000 kg



with proper choice of the detector,
very high energy resolution

Ge-diodes

bolometers

BUT

no event topology
no or limited event identification

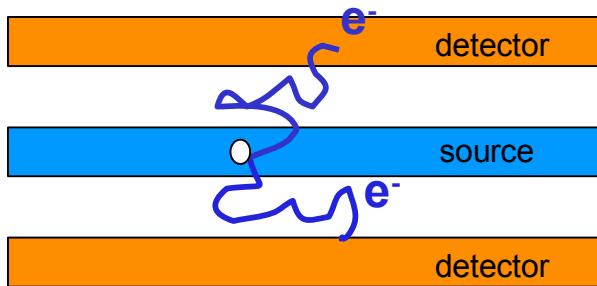


in gaseous/liquid Xenon detector,
indication of event topology
and event identification

BUT

poor energy resolution

source \neq detector



Source \neq Detector

- ☺ neat reconstruction of **event topology**
(allow to study the decay mechanism)
- ☹ it is **difficult** to get large source mass
- ☺ **several candidates** can be studied
with the same detector
- ☹ generally poor **energy resolution**
($2\nu\beta\beta$ could be the dominant background)

experimental sensitivity $S_{0\nu}$ [years]

half-life corresponding to the minimum detectable number of events above background at a given C.L

$$T_{1/2}^{0\text{nbb}} = \frac{\ln 2 \cdot N_{0\text{nbb}}}{rate_{0\text{nbb}}}$$

b= background counting rate in the ROI (Region Of interest) [counts/keV/kg/y]

FWHM ~ size of the ROI [keV]

M = mass of the $\beta\beta$ isotopes in the source [kg]

counts in the ROI = 0

90% upper limit on the rate

$$rate_{0\text{nbb}} = \frac{N_{Poisson}(90\%)}{time}$$

$$S_{0n} \sim M \text{ time}$$

counts in the ROI \neq 0

statistical fluctuation of bkg in the ROI

$$rate_{0\text{nbb}} = \frac{\sqrt{b \cdot time \cdot FWHM}}{time}$$

$$S_{0n} \sim \sqrt{\frac{M \text{ time}}{FWHM \cdot b}}$$

fighting for sensitivity

$$S_{0n} = \ln 2 \cdot N \cdot i.a. \cdot \sqrt{\frac{time}{(bgk \cdot FWHM)}} \cdot eff.$$

i.a. = isotopic abundance of the $\beta\beta$ candidate

N = number of nuclei

time = experiment live time

eff. = % of Onbb decay detected

FWHM = measures detector capability to distinguish
a $0\nu\beta\beta$ event from a bkg event

bkg = background of events capable of mimic a $0\nu\beta\beta$ decay

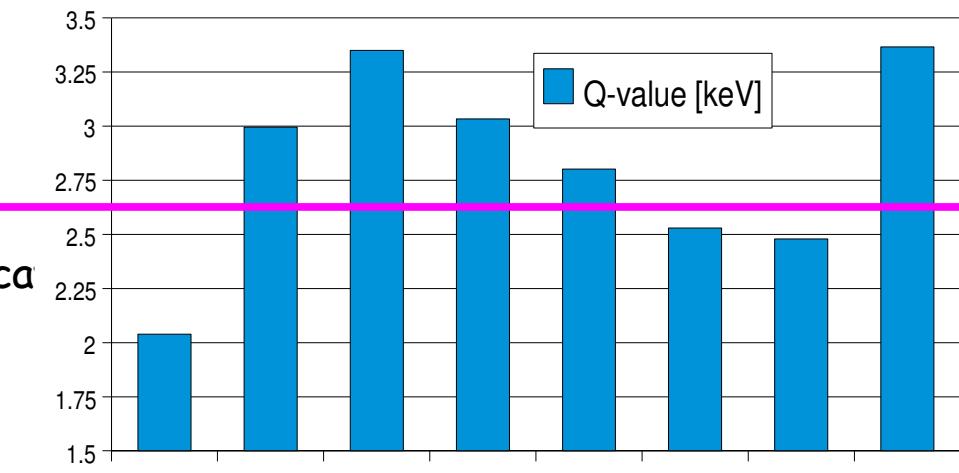
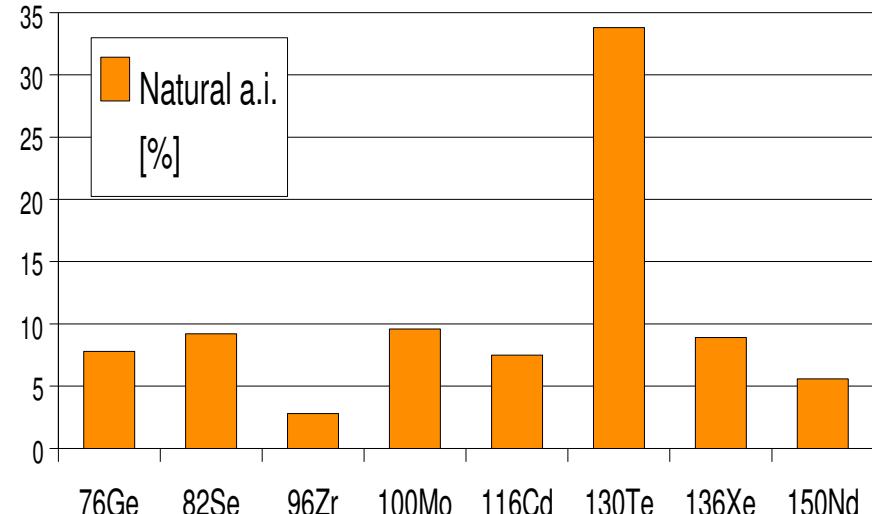
Q-value is a rather good indicator for bkg,

"the higher the Q value, the lower the bkg"

most natural γ radioactivity is negligible above 2.6 MeV

location underground and deep to shield from cosmic rays

Nature or Money
Detector choice
Location + Shielding
Material selection



The enemy: background sources able to mimic a $O\nu\beta\beta$ event

from the Universe: cosmic rays

→ deep underground laboratories

here only muons survive muon energy and flux depends on the lab depth

μ direct interaction in the apparatus

→ veto

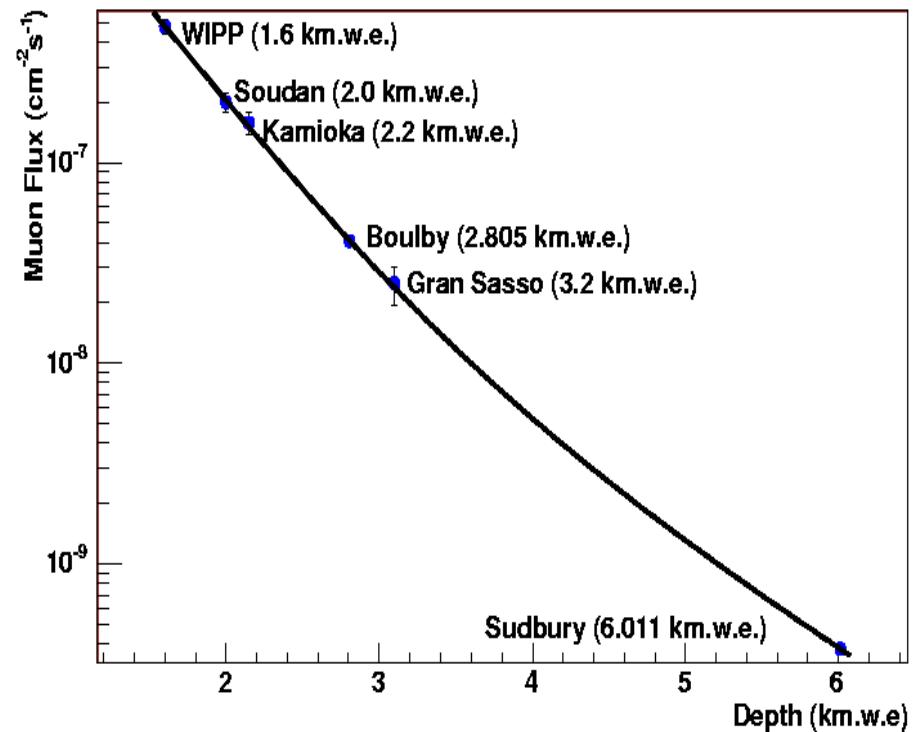
surround the exp. apparatus with a μ detector

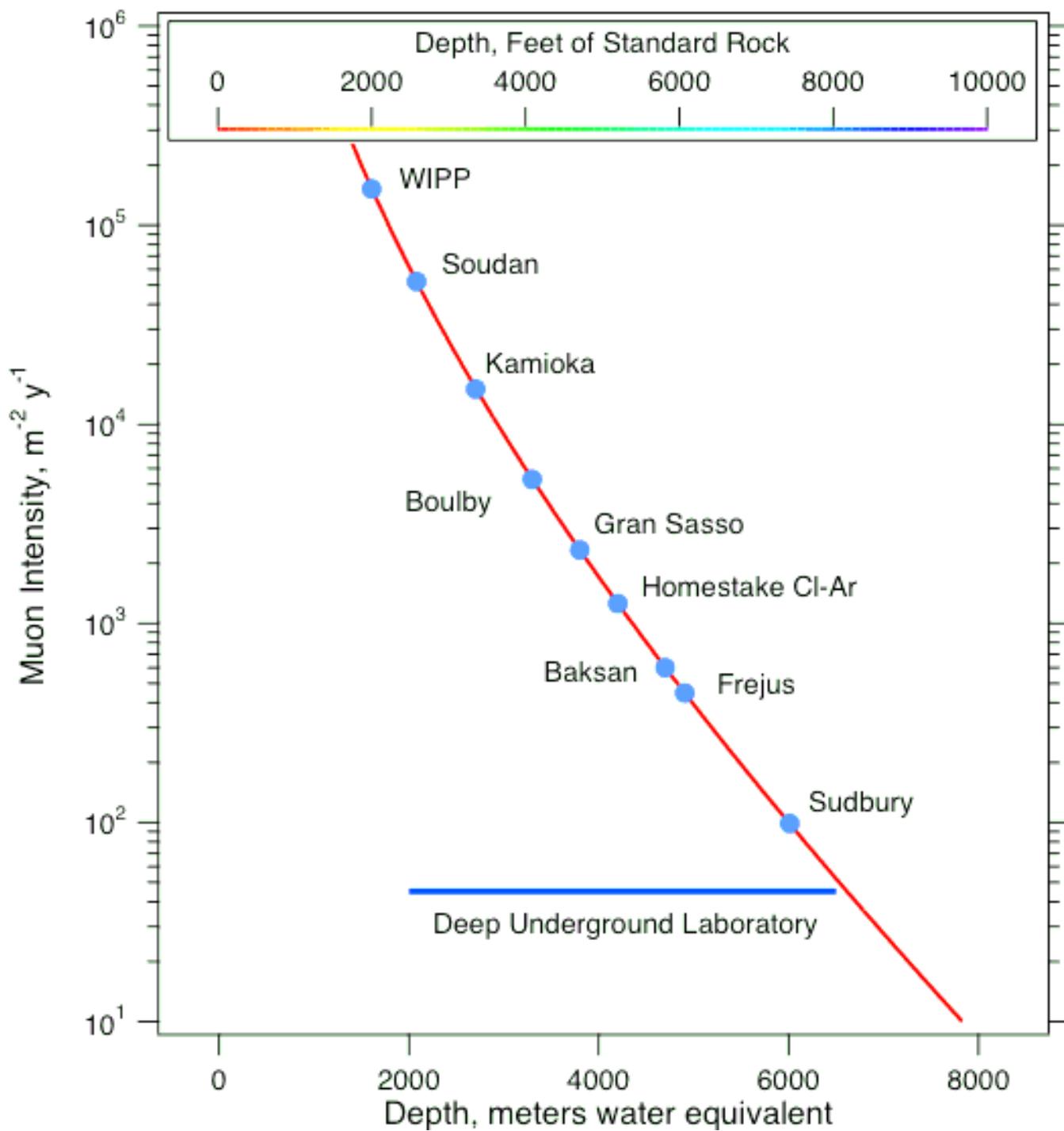
NOTE: does not solve the problem of metastable nuclei produced by μ interaction in the materials

μ interaction in the rock, relevant secondaries are gammas
→ shields neutrons

these have an harder spectrum than the one due to radioactivity (but extremely low fluxes)

if veto and shields are not enough the only solution is a deeper location





Radioactivity

natural

long living isotopes (U, Th, K ...) α , β , γ and n



cosmogenically activated isotopes β and γ

n are produced by (α, n) reactions and spontaneous fission, they have $E < 10 \text{ MeV}$

anthropogenic (nuclear explosions, nuclear plants , industry, medical ...) β and γ

radioactivity of the rock:

γ and n



shield the apparatus

labs located in mountains with low radioactivity rock (salt mines)

radioactivity of the apparatus:

γ , n and also α and β (when emitted near to the active volume)

material selection



avoid cosmic ray activation (underground storage of materials)

develop active rejection technique

(pulse shape rejection, double reading, coincidences ...)

fighting background

a typical experimental apparatus is

deep underground

shielded against gamma and neutrons

built with specially selected low activity materials

tools useful for the study of background components

MonteCarlo simulation

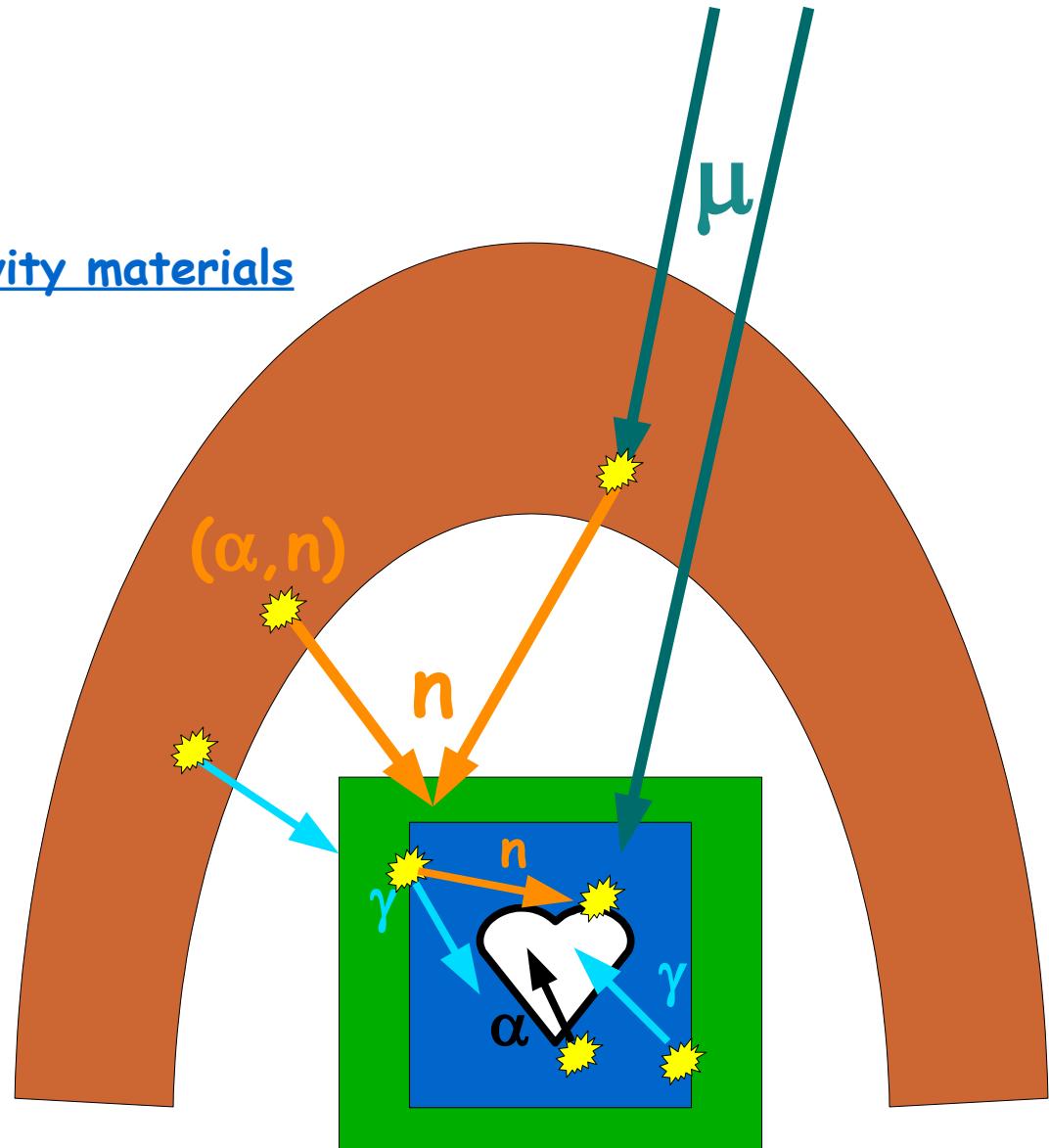
--> require several exp. inputs

--> validation

μ and n flux measurements
(or simulation)

HPGe, ICPMS, n activation ...

--> material selection



Part2.

Present and Past Experiments

“prehistory”

- 1948 - first counter experiment (Geiger counters,
 $T_{1/2}(0\nu) > 3 \cdot 10^{15} \text{ y}$)
- 1950 - first evidence for $2\beta 2\nu$ decay of ^{130}Te in
first geochemical experiment: $T_{1/2} \sim 1.4 \cdot 10^{21} \text{ y}$
- 1950-1965 - a few tens experiments with
sensitivity $\sim 10^{16}-10^{19} \text{ y}$

“history”

- $T_{1/2}(^{76}\text{Ge}) > 5 \cdot 10^{21} \text{ y}$; Ge(Li) detector, 1973 (E. Fiorini et al.)
- $T_{1/2}(^{48}\text{Ca}) > 2 \cdot 10^{21} \text{ y}$; streamer chamber + magnetic field + plastic scint., 1970 (C. Wu et al.)
- $T_{1/2}(^{82}\text{Se}) > 3.1 \cdot 10^{21} \text{ y}$; streamer chamber + magnetic field + plastic scint., 1975 (C. Wu et al.)
- Geochemical experiments with ^{130}Te , ^{128}Te , ^{82}Se (“positive” results)

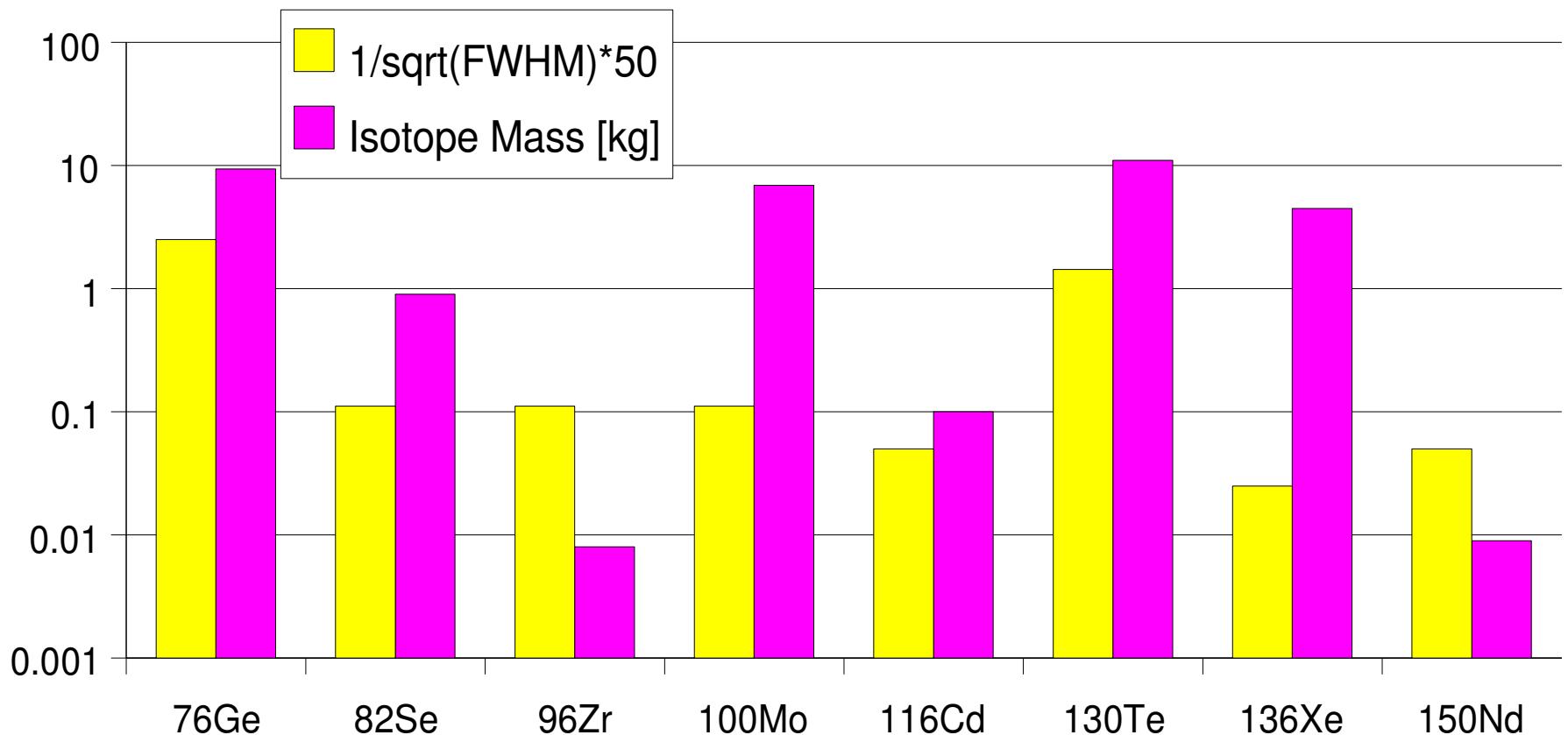
yesterday

Emitter	Experiment	$T_{1/2}^{0\nu} >$	C.L. %
^{48}Ca	ELEGANT VI	$1.4 \times 10^{22} \text{ y}$	90
^{76}Ge	MPIH/KIAE	$1.9 \times 10^{25} \text{ y}$	90
	IGEX	$1.6 \times 10^{25} \text{ y}$	90
^{82}Se	UCI	$2.7 \times 10^{22} \text{ y}$	68
	NEMO 3	$4.7 \times 10^{22} \text{ y}$	90
^{90}Zr	NEMO 2	$1.3 \times 10^{21} \text{ y}$	90
^{100}Mo	LBL/MHC/UNM	$2.2 \times 10^{22} \text{ y}$	68
	UCI	$2.6 \times 10^{21} \text{ y}$	90
	Osaka	$5.5 \times 10^{22} \text{ y}$	90
	NEMO 3	$6 \times 10^{22} \text{ y}$	90
^{116}Cd	Kiev	$1.7 \times 10^{23} \text{ y}$	90
	Osaka	$2.9 \times 10^{21} \text{ y}$	90
	NEMO 3	$1.6 \times 10^{22} \text{ y}$	90
^{130}Te	Milano	$2.1 \times 10^{23} \text{ y}$	90
	CUORICINO	$1 \times 10^{24} \text{ y}$	90
^{136}Xe	Caltech/UN/PSI	$4.4 \times 10^{23} \text{ y}$	90
^{136}Xe	Rome	$1.2 \times 10^{24} \text{ y}$	90
^{150}Nd	UCI	$1.2 \times 10^{21} \text{ y}$	90
	NEMO 3	$1.4 \times 10^{21} \text{ y}$	90

Table 1.4: Limits on Neutrinoless Decay Modes

a selection of present/past experiments:

$$S_{0n} = \ln 2 \cdot N \cdot ai \cdot \sqrt{\frac{time}{(bgk \cdot FWHM)}} \cdot eff.$$



27 November 1967

A SEARCH FOR ANTONIUS-CONSERVATION IN DOUBLE BETA DECAY
WITH A GERMANIUM DETECTOR.

E. FRUNZI and A. PELLEA
Istituto di Fisica dell'Università e INFN, Milano, Italy

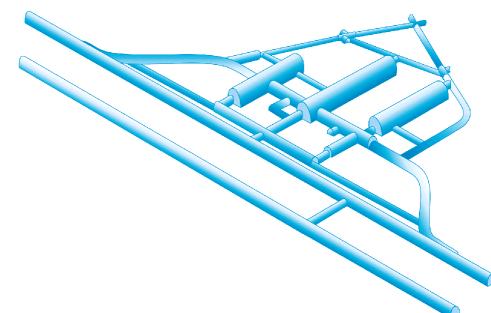
C. BERTOLINI, R. CAPPELLANI and G. FRATELLI
Istituto di Fisica, Roma, Italy

Resolution Ge(Li) mass 90 g.
Running time = 112 h
Resolution 4.7 keV @ 1.32 MeV

Why $\beta\beta$ searches of ⁷⁶Ge have so long history?

- Germanium has high intrinsic purity (required also by other applications)
- Ge diodes detector is a well established technology
- source = detector high efficiency!
- ⁷⁶Ge isotopic abundance = 7.44%, but enrichment of ⁷⁶Ge possible at centrifuge up to >80%.
- Ge density = 5.3 g cm⁻³ compact setup
- low Atomic Weight (1 kg of ⁷⁶Ge = 13.1 Moles = 7.9×10^{24} nuclei)

^{76}Ge : Heidelberg-Moscow experiment



source=detector experiment

- *tech. suggested in '60 by E. Fiorini, dominating this field since now*
- *exp. started in 1990*

detector = 5 Ge diodes

source = 10.9 kg diodes

enriched in ^{76}Ge (i.a. 86%)

Q - value = 2039 keV

Location = Lab. Naz. del Gran Sasso - Italy

Countries = Germania + Russia



The experiment is closed but can be considered the precursor (together with IGEX) of the next generation expt. GERDA and Majorana

^{76}Ge : Heidelberg-Moscow experiment

statistics=

53.9 kg γ , with PSA = 35.5 kg γ

performances =

4.2 keV FWHM resolution at DBD Q-value

background in $0\nu\beta\beta$ region =

0.19 +/- 0.01 c/keV/kg/ γ

0.06 +/- 0.01 c/keV/kg/ γ with PSA

Pulse Shape Analysis, used to identify and reject multi-site events (gamma background)

the entire collaboration (14 authors):

Klapdor-Kleingrothaus et al.

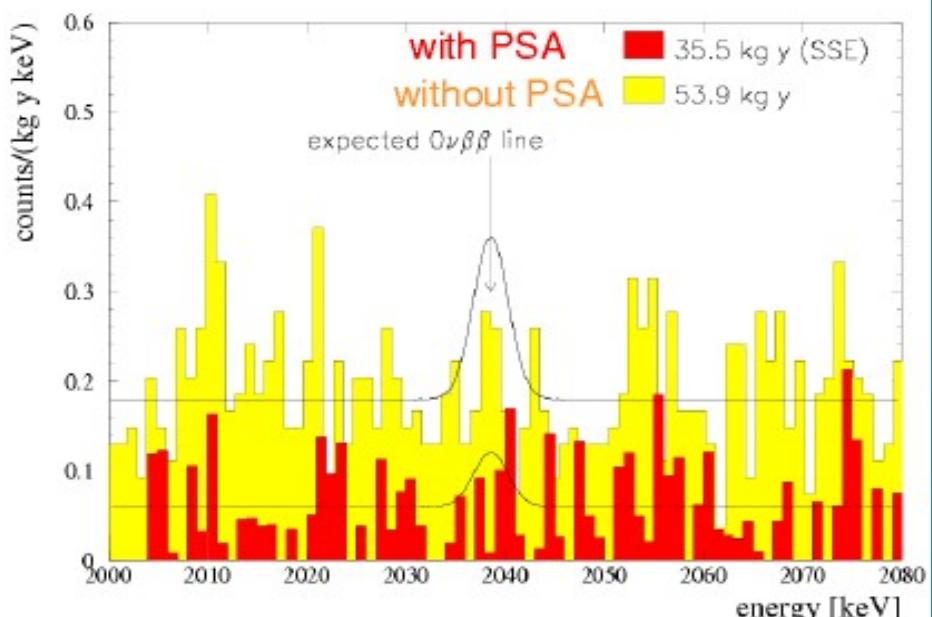
Eur. Phys. J. 12 (2001) 147

$\tau_{1/2}^{0n} > 1.9 \cdot 10^{25} \text{ y}$ at 90% C.L.

$\langle m_\nu \rangle < 0.35 \text{ eV}$ *

* using nuclear calculations of Staudt et al.

Europ. Lett. 13 (1990) 31



^{76}Ge : Heidelberg-Moscow experiment a controversial result

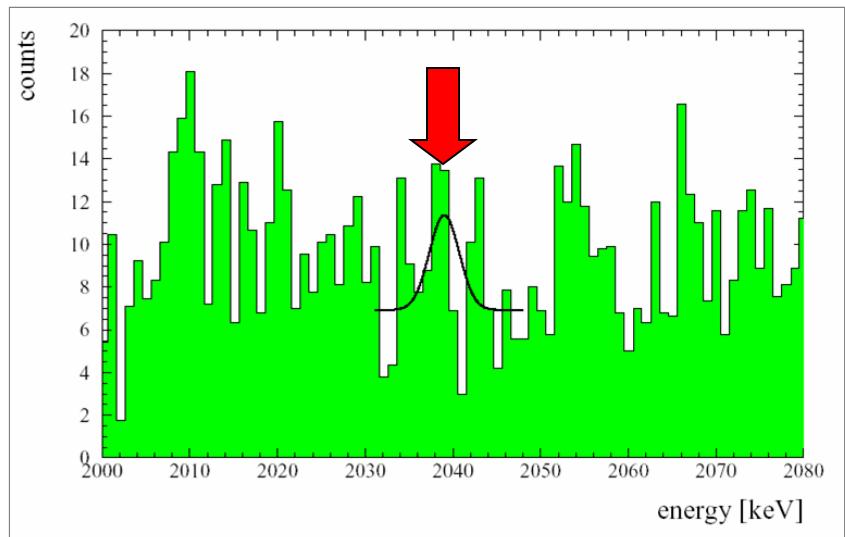
2001: a re-analysis by part of the collaboration (4 authors):

"Evidence for Neutrinoless DBD"

Klapdor-K. et al. *Mod. Phys. Lett. A* 16 (2001) 2049

$$\tau_{1/2}^{0n} = (0.8 - 18.3) 10^{25} \text{ y} @ 95\% \text{ C.L.}$$

→ $\langle m \rangle_\nu = 0.11 - 0.56 \text{ eV}$

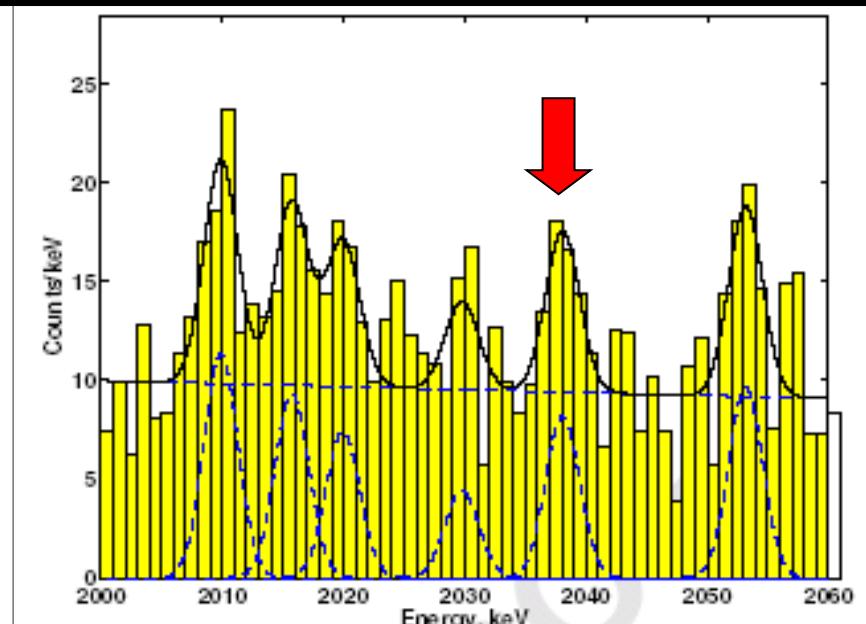


2004: data with higher statistics and with better quality show an increase of the statistical significance of the "peak":

Klapdor-K et al. *Phys. Lett. B* 586 (2004) 198-212

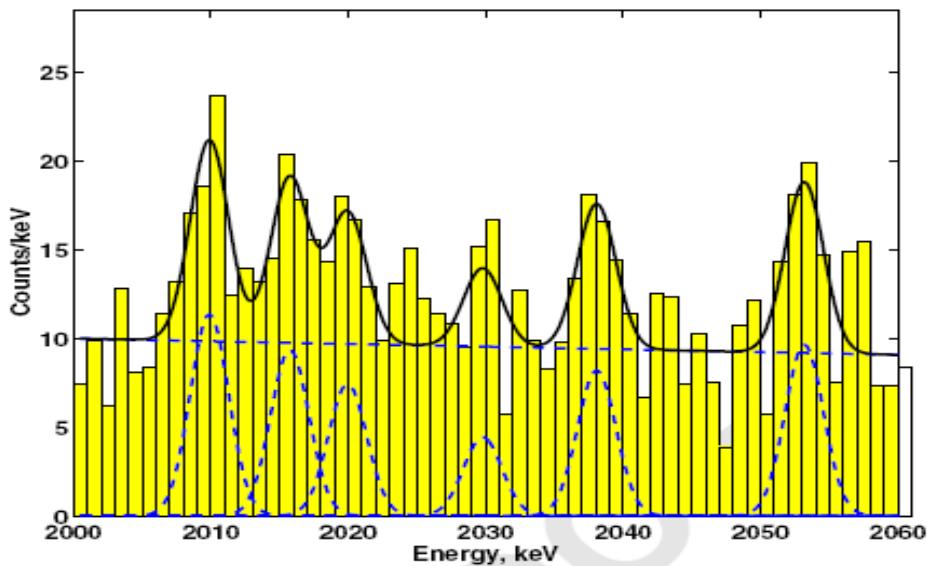
$$\tau_{1/2}^{0n} = (0.69 - 4.18) 10^{25} \text{ y } (3\sigma)$$

→ $\langle m \rangle_\nu = 0.24 - 0.58 \text{ eV}$



^{76}Ge : Heidelberg-Moscow experiment a controversial result

skepticism in DBD community



$\langle m \rangle [\text{eV}]$

0.2-0.6

0.1-0.4

0.3-0.7

0.4-1.1

comments and reanalysis:

Aalseth CE et al., Mod. Phys. Lett. A 17 (2002) 1475

Feruglio F et al., Nucl. Phys. B 637 (2002) 345

Zdezenko Yu G et al., Phys. Lett. B546(2002)206

replies:

Klapdor-Kleingrothaus HV hep-ph/0205228

H.L. Harney, hep-ph/0205293

- One peak not explained
- Looking at a larger range, many structures resemble the DBD "peak" and need to be explained
- The statistical significance depends on the flat component of the background

QRPA Rodin Faessler Simkovic Vogel, nucl-th:0706.4304v1

QRPA Kortelainen Suhonen 2007, Phys. Rev C

QRPA Civitarese Shuonen , nucl-th/0208005

Nuclear Shell model, Caurier et al. nucl-th/0801.3760

^{76}Ge : IGEX

Collaboration: (ITEP, INR, U. South Carolina;
PNNL, U. of Zaragoza, Yerevan)

Location: Canfranc UL (Spain)

Detectors: two enriched Ge diodes

Pulse shape discrimination

$$117 \text{ mol} \cdot \gamma = 8.9 \text{ kg} \cdot \gamma$$

Bkg at 2 Mev

~ 0.1 (with PSA)

~ 0.2 (without PSA) [cts/(keV. γ)]

$T_{1/2} > 1.57 \cdot 10^{25} \text{ y}$ with PSA

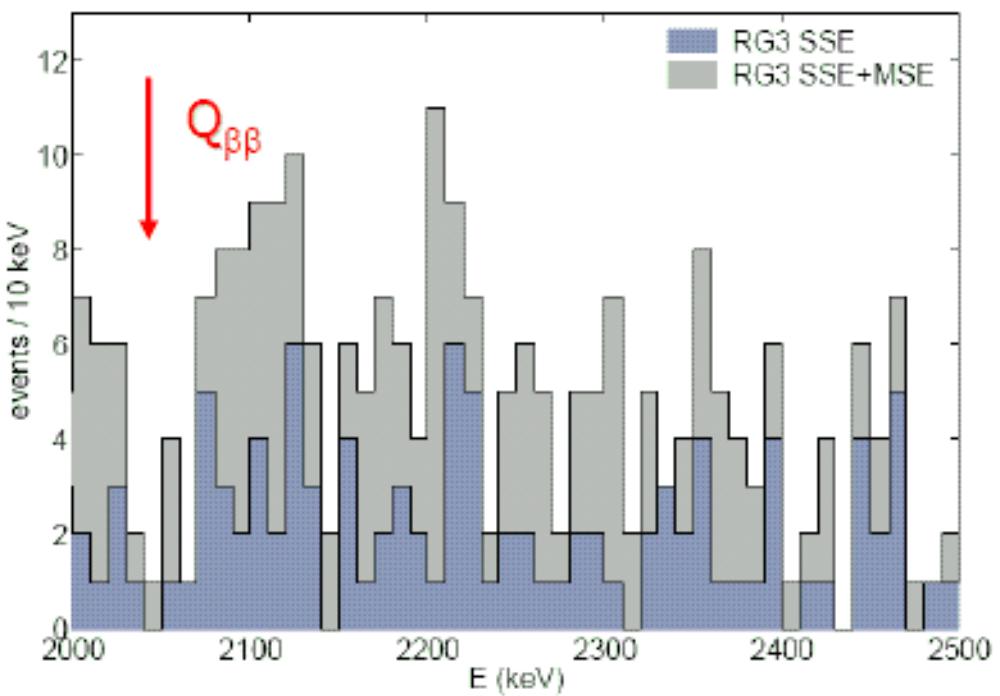
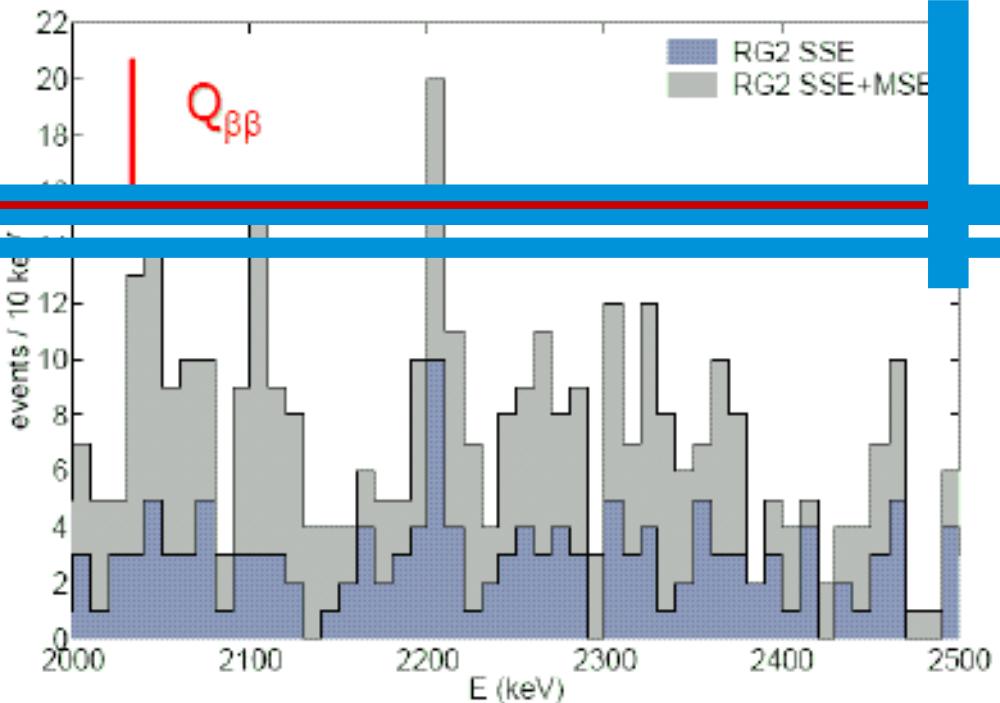


Figure 10: Background spectra before and after the PSD based on the counting of the number of lobes for detectors RG2 (top) and RG3 (bottom).

What if not Ge?

There are several reasons to study more than one isotope. The "Ge way" was the easiest and natural one: this detector is based on a very well known technology, with several different application justifying large efforts in its technical improvement (in purity, size, resolution).

How to choose an isotope/detector/experiment?

These are the ingredients:

NME

isotopic abundance

total mass

Q value

energy resolution

detector technology

I tell you the unusual story that gave birth to bolometers

"bolometer tale"

Once upon a time ... searching for the best DBD candidate a physicist discovered how Nature had been kind with ^{130}Te :

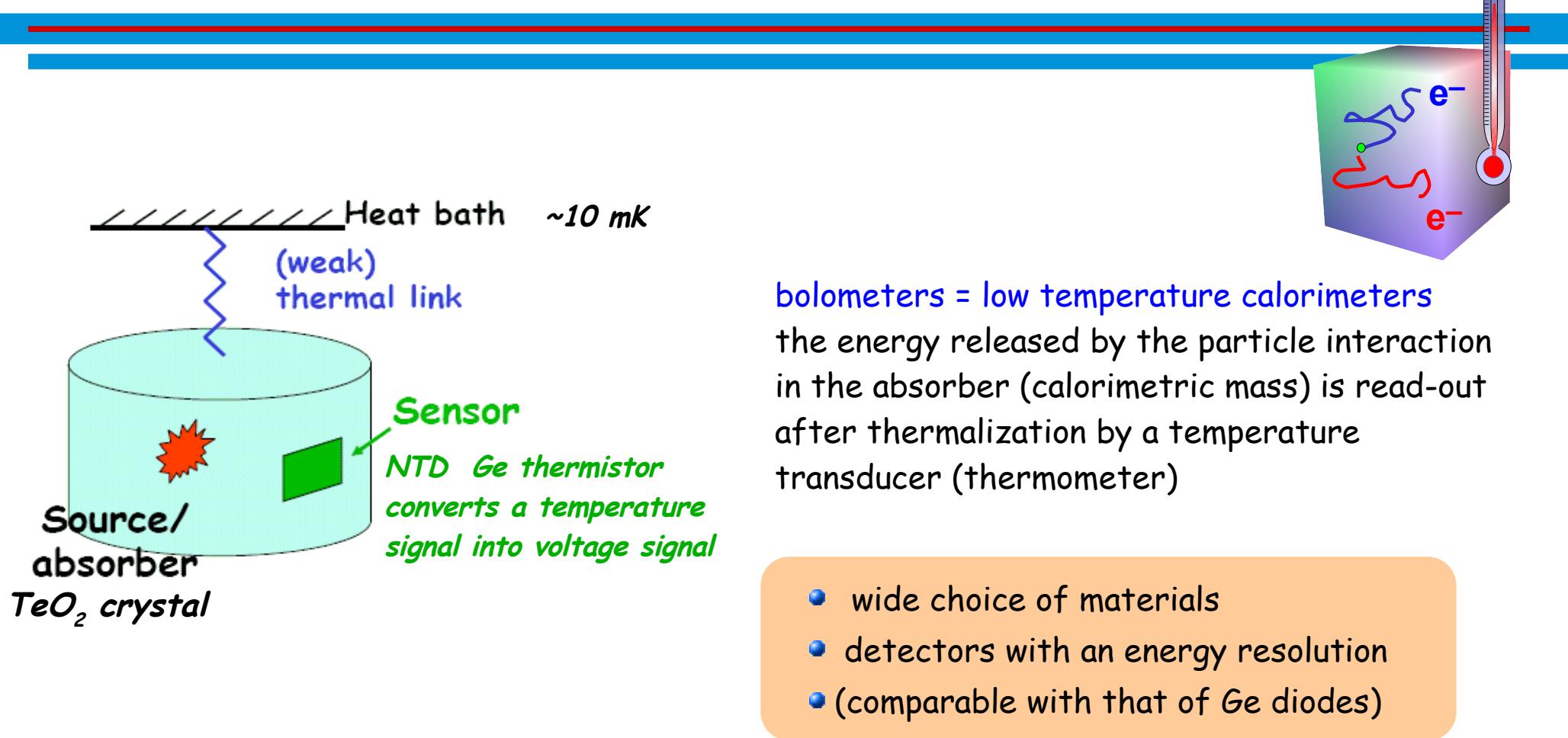
- a good F_N
- a rather high Q-value
- a high i.a. (quite relevant for the empty pockets of the experimentalist).

At that time the best DBD experiments used ^{76}Ge diodes, ^{130}Te was/is in principle better ^{76}Ge , the only (not negligible) point was that at that time no solid state detector - made of Te - with performances competitive with Ge diodes did exist.

The solution: **JUST "INVENT" IT (this looked like a DREAM)**

that is build a bolometric detector (at that time - '80 - maximum size was ~ grams) and reach energy resolution and masses as good as HPGe

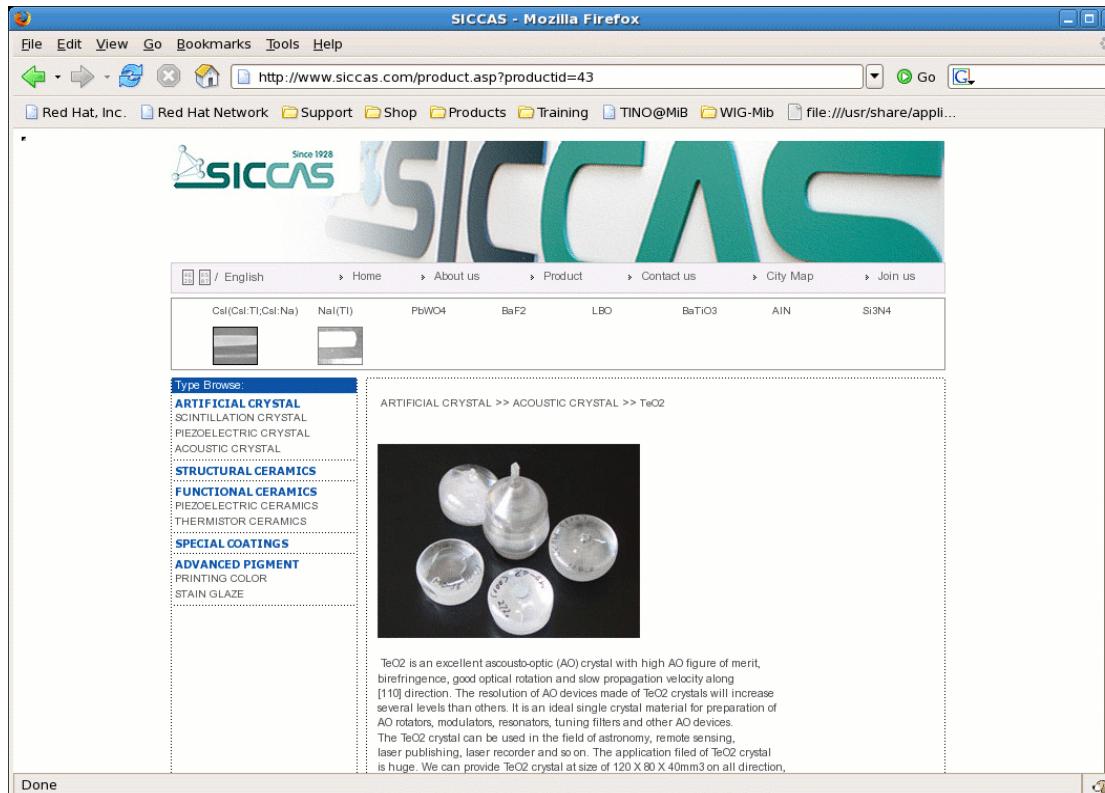
bolometers



- tech. suggested in 1985 by E. Fiorini and T.O. Niinikoski
- the first Te detector worked in 1989
- first bolometric DBD experiment in 1997
- predecessor of CUORICINO : 20 crystal array (MiDBD)
- other applications: Dark Matter detection (CDMS, Edelweiss, CRESST)

"bolometer tale"

Our physicist it well known to be a lucky man and he happened to be told during a meeting that there was a crystal producer in China that was growing beautiful TeO_2 crystals.



TeO₂ has good thermal, mechanical, radioactive properties

- easy to grow large size crystals (up to 6x6x6 cm³ = 1.2 kg)
- good intrinsic radiopurity (better than 10⁻¹³ g/g in U and Th)
- extremely good bolometric performances
(FWHM at the ¹³⁰Te DBD transition energy ~ 7 keV)

extensively tested by the Milano group during the last 10 years

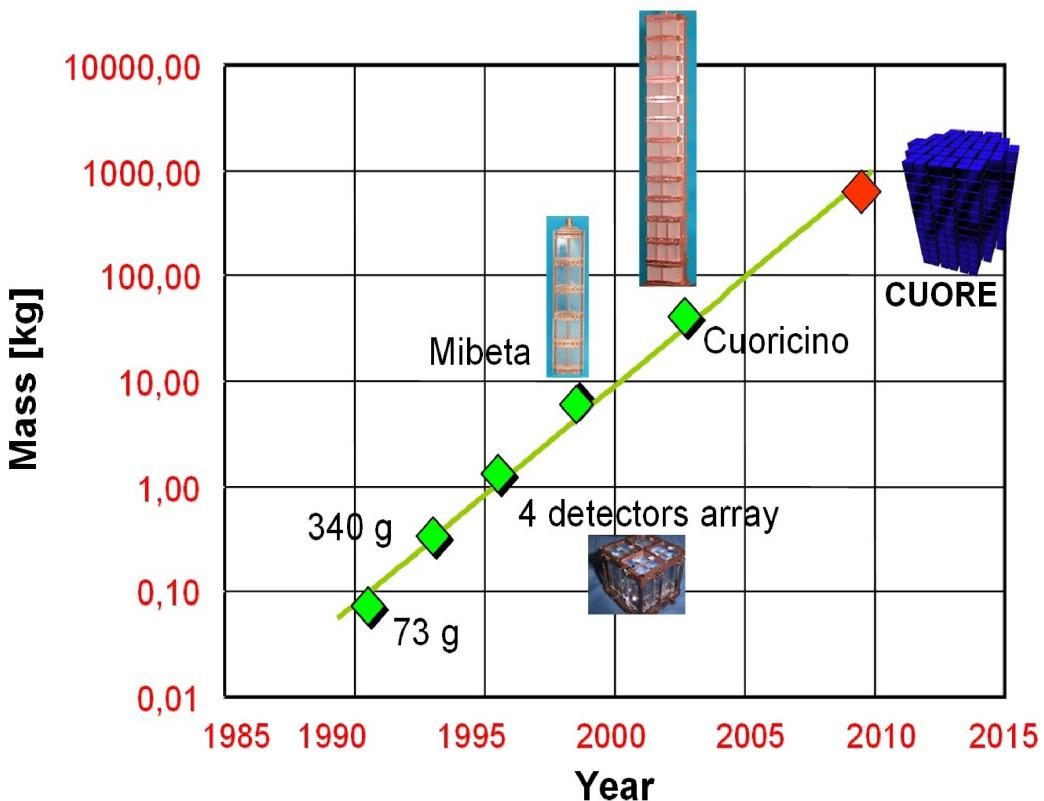
standard production by SICCAS (Shanghai China)

"bolometer tale"

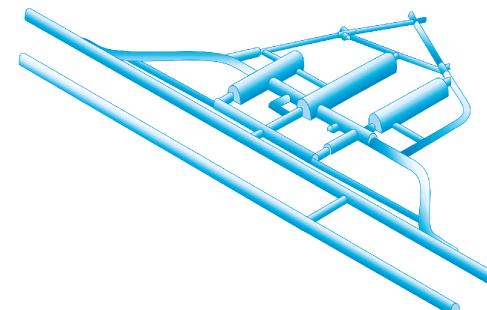
... AND THE REALITY OVERCAME THE DREAM ...

- the detector did work very well
- in the years detectors growing in size without showing deterioration of performances (resolution)
- the bulk radio-purity of these crystals have always be extremely high
- our physicist is now the spokesperson of a big collaboration (CUORE)

The Moore's law of TeO₂ bolometers



^{130}Te : Cuoricino



Lab. Naz. del Gran Sasso - Italy

source=detector experiment

- tech. suggested in 1985 by E. Fiorini and T.O. Niinikoski
- first experiment in 1997

source ~ 40 kg natural TeO_2

(^{130}Te i.a. 33.8 % => ~ 11 kg ^{130}Te)



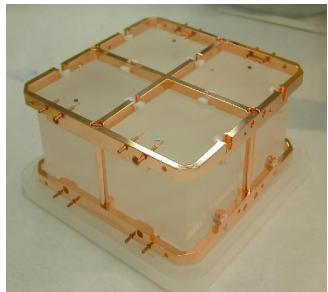
The experiment was recently closed (June 2008).

A 1 ton experiment (CUORE) is under construction

Cuoricino detector

40.7 kg TeO₂ ~ 11 kg ¹³⁰Te

44 TeO₂ crystals 5x5x5 cm³
+ 18 TeO₂ crystals 3x3x6 cm³



11 modules, 4 detector each,
crystal dimension 5x5x5 cm³
crystal mass 790 g

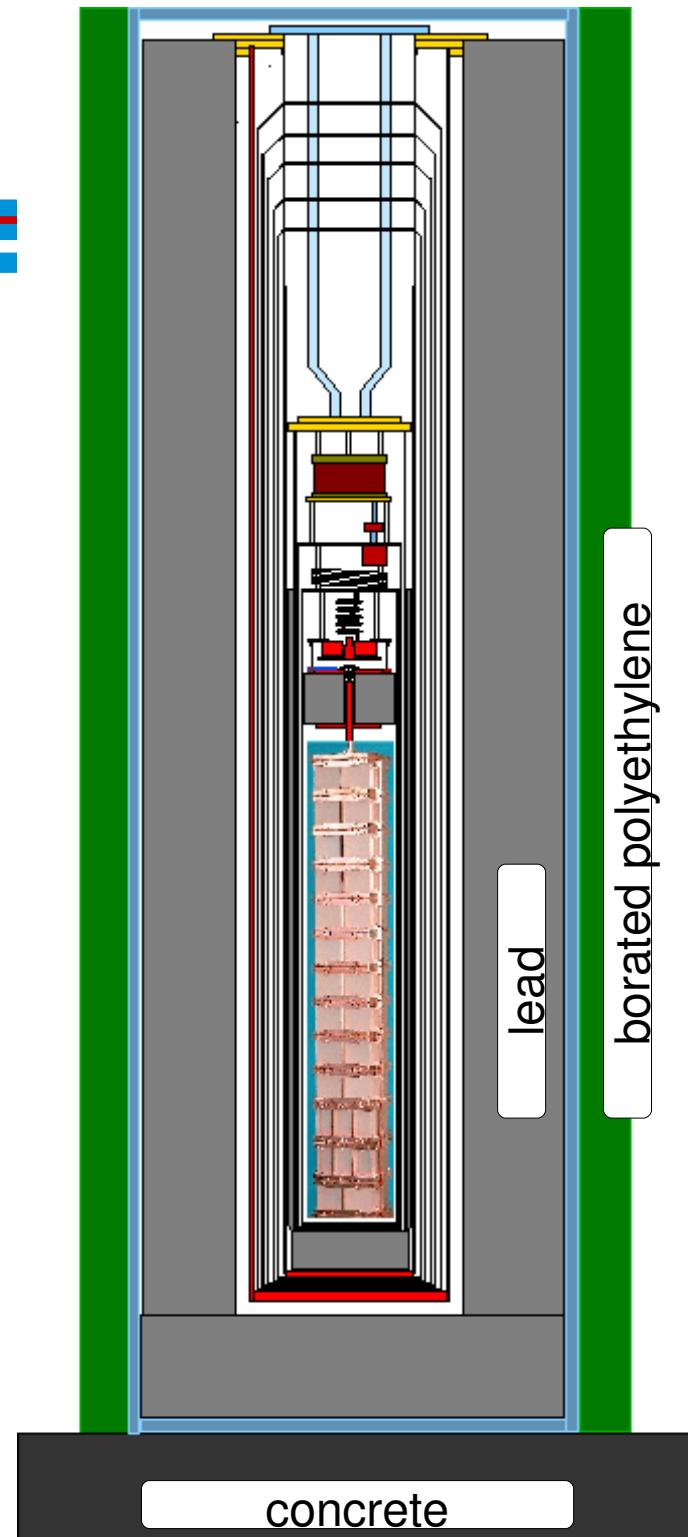
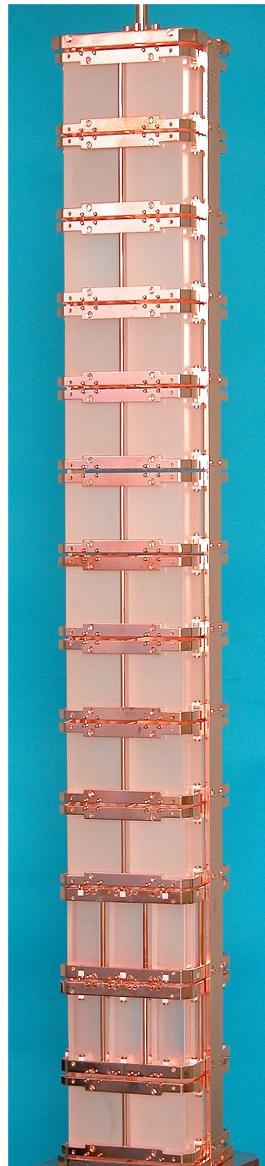
$$4 \times 11 \times 0.79 = 34.76 \text{ kg of TeO}_2$$



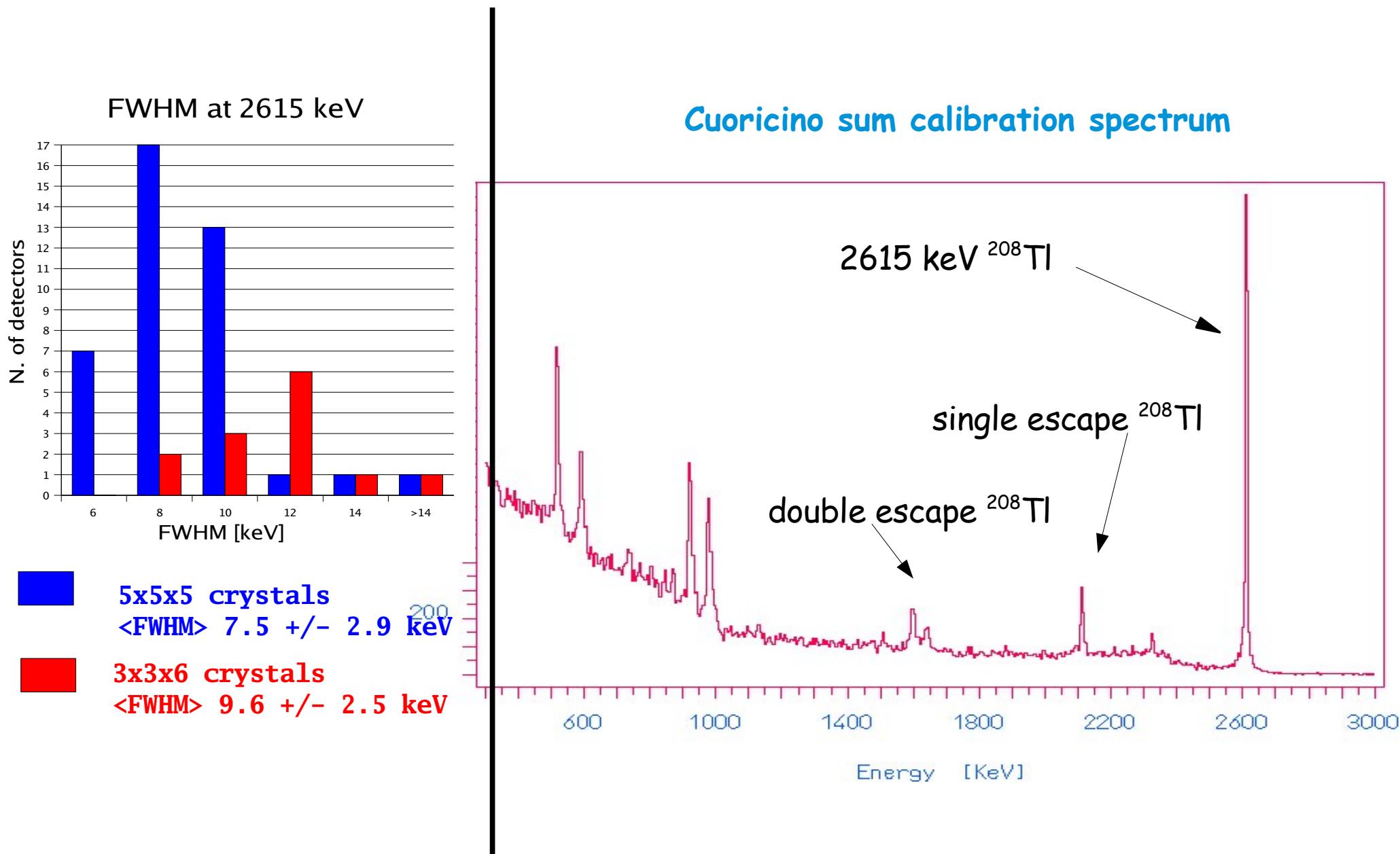
2 modules, 9 detector each,
crystal dimension 3x3x6 cm³
crystal mass 330 g

$$9 \times 2 \times 0.33 = 5.94 \text{ kg of TeO}_2$$

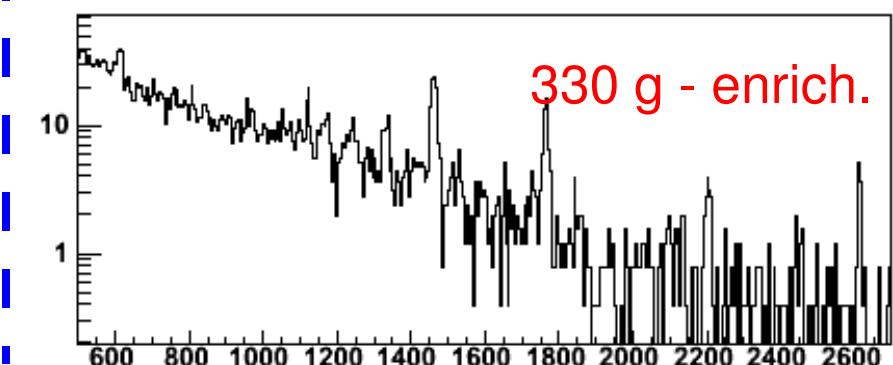
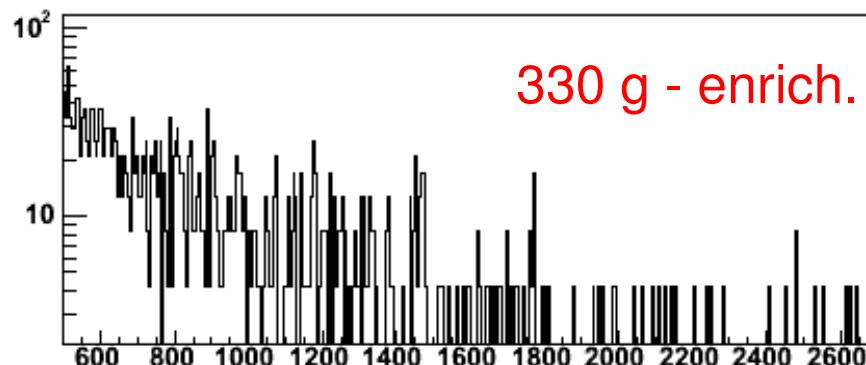
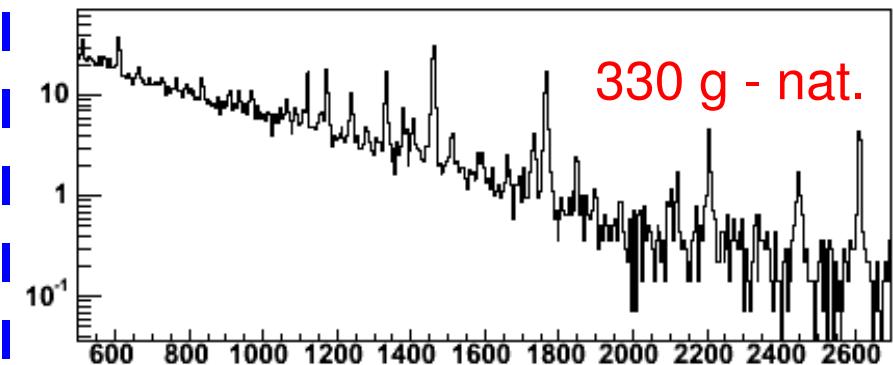
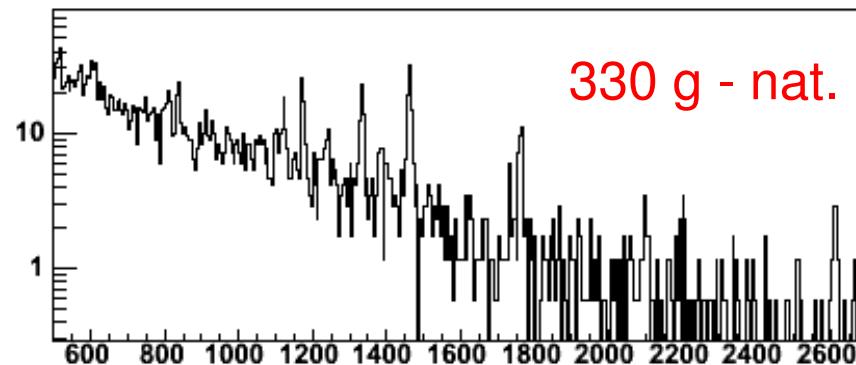
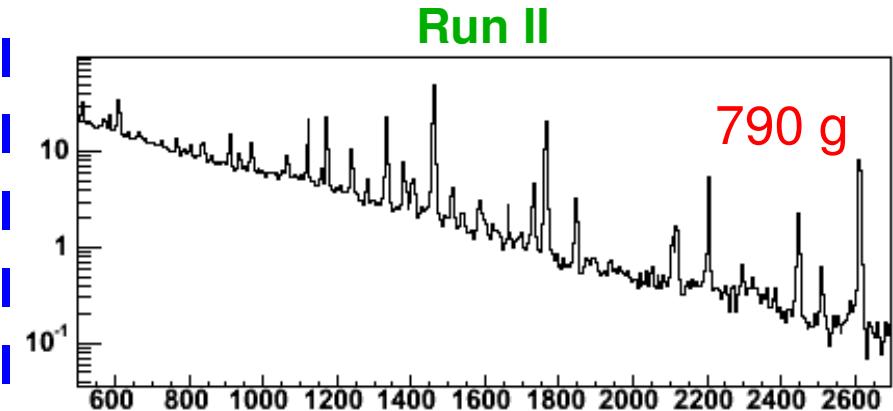
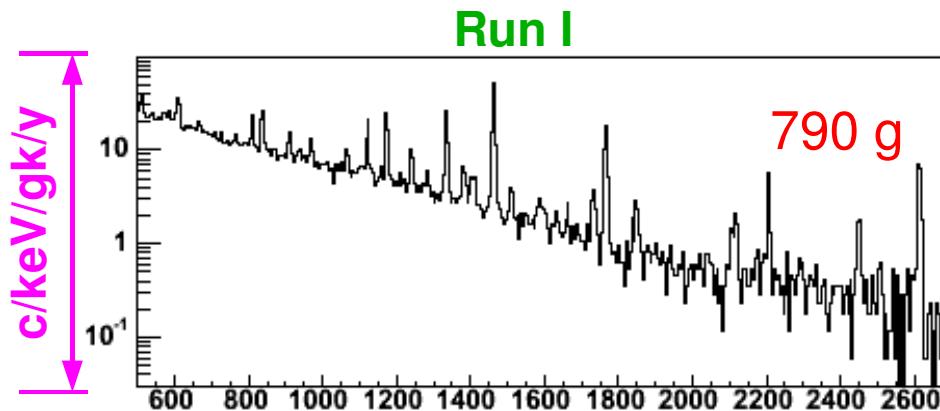
two small crystals enriched in ¹³⁰Te
two small crystals enriched in ¹²⁸Te



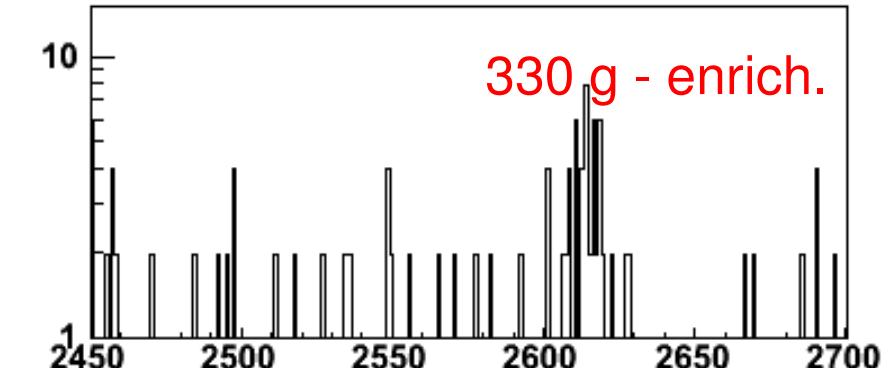
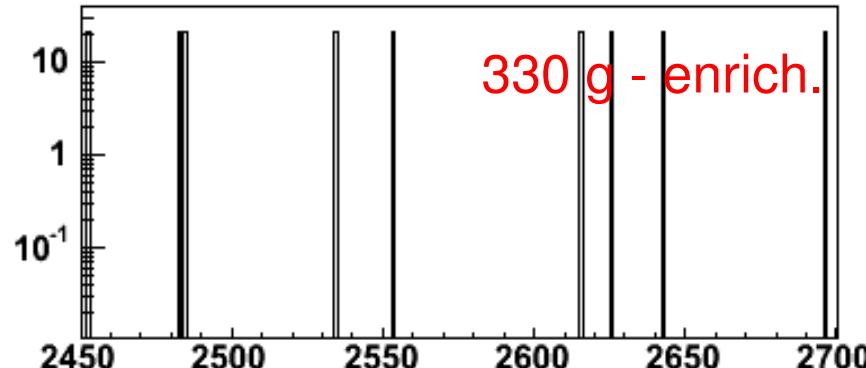
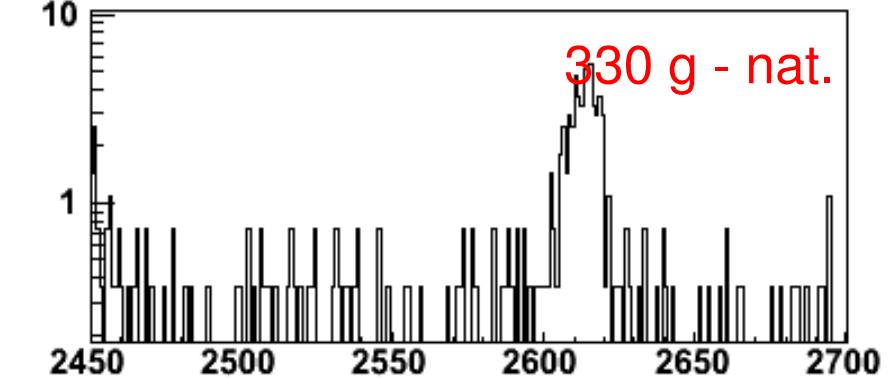
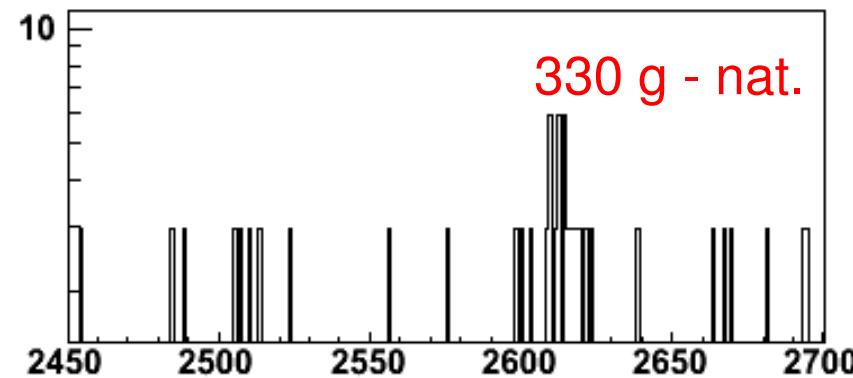
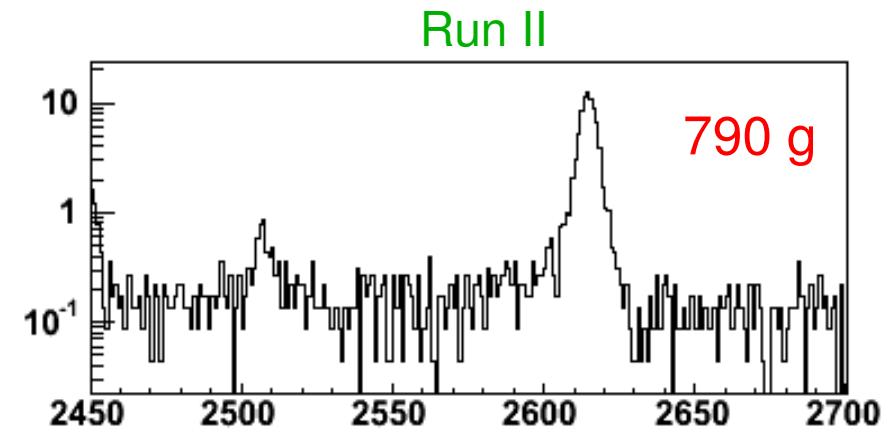
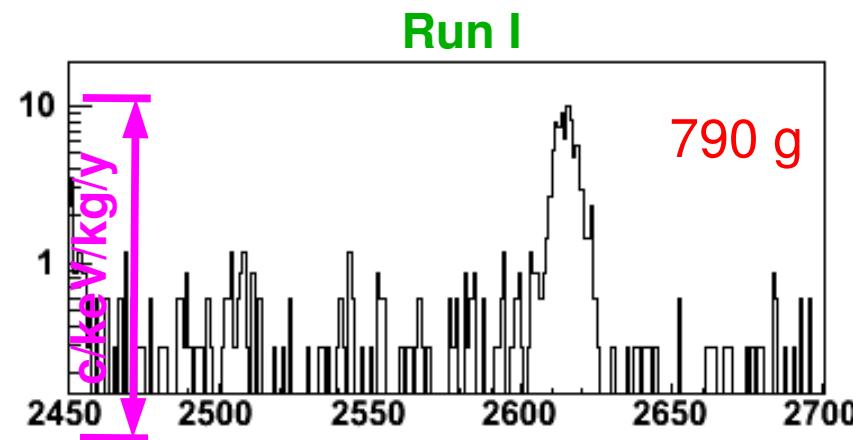
an array behaving as a single detector:



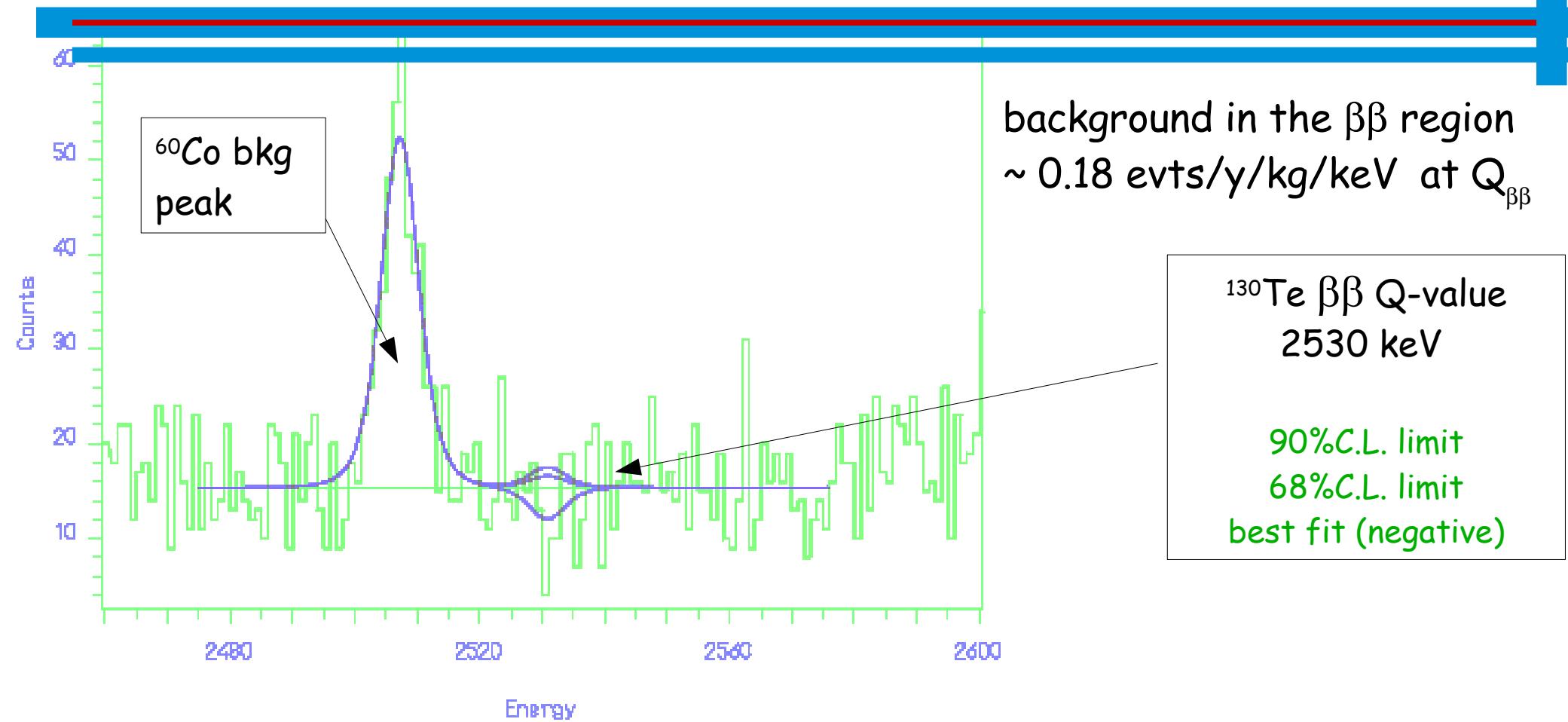
Cuoricino: sum spectra in the γ region



Cuoricino: sum spectra in the $\beta\beta$ region



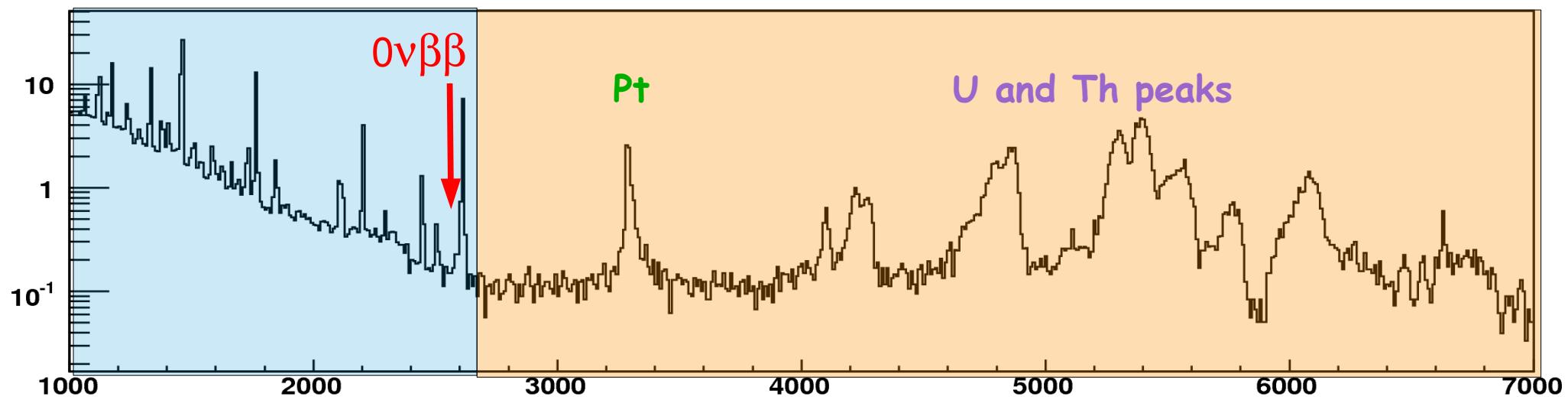
Cuoricino Ov $\beta\beta$ result:



$$\tau_{1/2} > 3.1 \cdot 10^{24} \text{ [y]} @ 90\% \text{ C.L.}$$
$$\langle m_n \rangle < (0.2 - 0.98) \text{ [eV]}$$

to be published
previous result Phys. Rev. C 78 (2008) 035502

Cuoricino background



γ region

peaks are due to

- ^{238}U , ^{232}Th contaminations in the cryostat
- ^{60}Co cosmogenic activation of Cu
- ^{40}K ? no info on its location

continuum accounted for by

- degraded γ

α region (internal or surface contaminations)

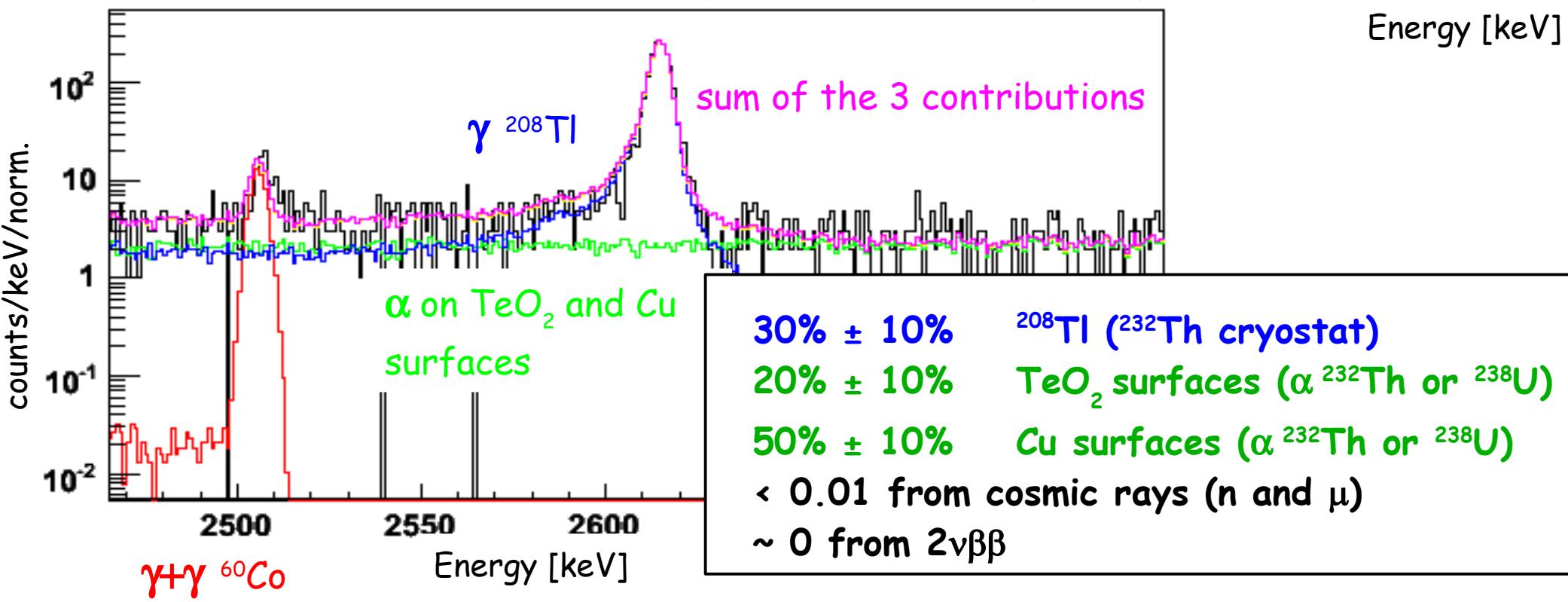
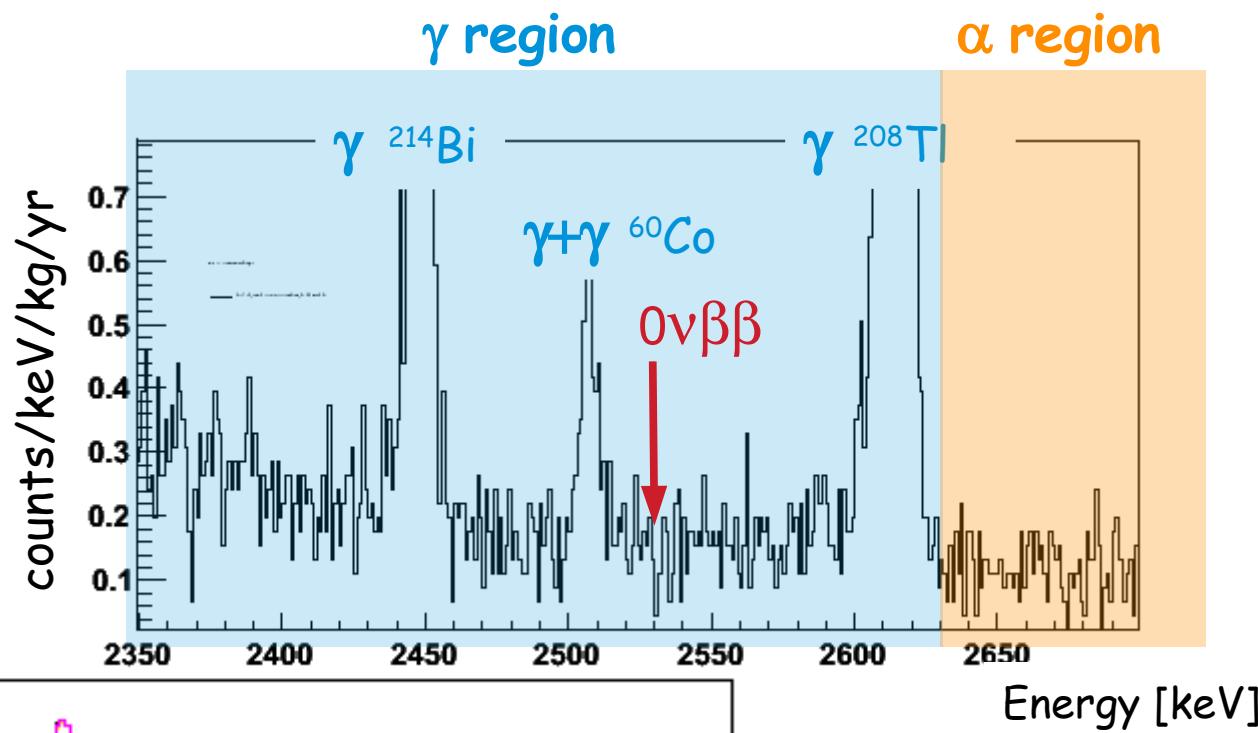
peaks are due to

- ^{238}U , ^{232}Th surf. contam of TeO_2 crystals
- ^{190}Pt bulk contam. of TeO_2
- minor contribution from ^{238}U , ^{232}Th bulk. contam of TeO_2 crystals

continuum accounted for by

- degraded α (surf. contam. of Cu?)

Cuoricino background in the $\beta\beta$ region



Cuoricino result vs. HM claim of evidence:

HM claim of evidence:

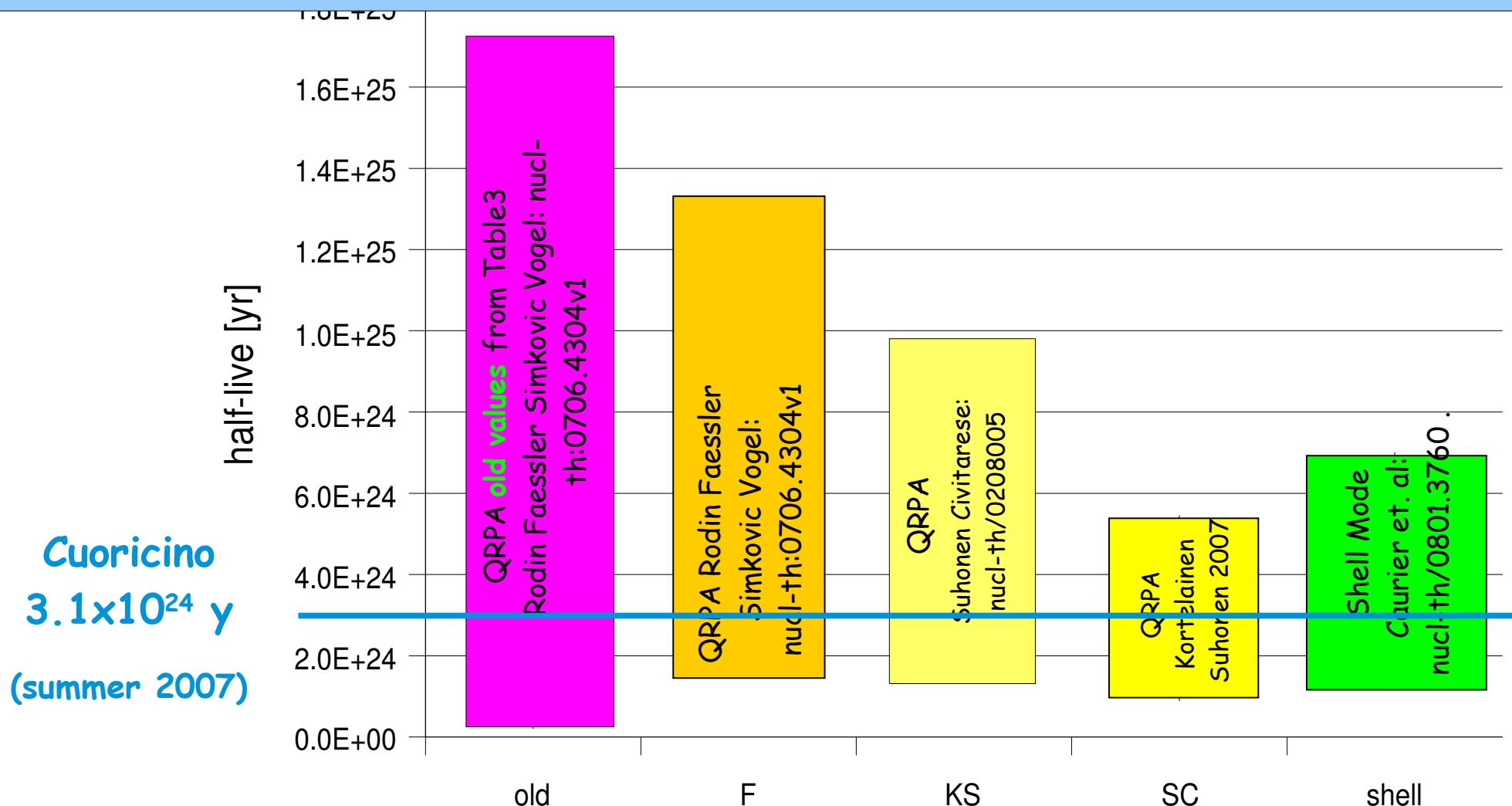
$$T_{1/2}^{0\nu} (y) = (0.69-4.18) \times 10^{25} \text{ (3}\sigma\text{ range)}$$



^{130}Te predicted

$T_{1/2}^{0\nu} (y)$ range

$$T_{1/2}^{-1} = F(Q_{\beta\beta}, Z) |M|^2 \langle m_\nu \rangle^2$$

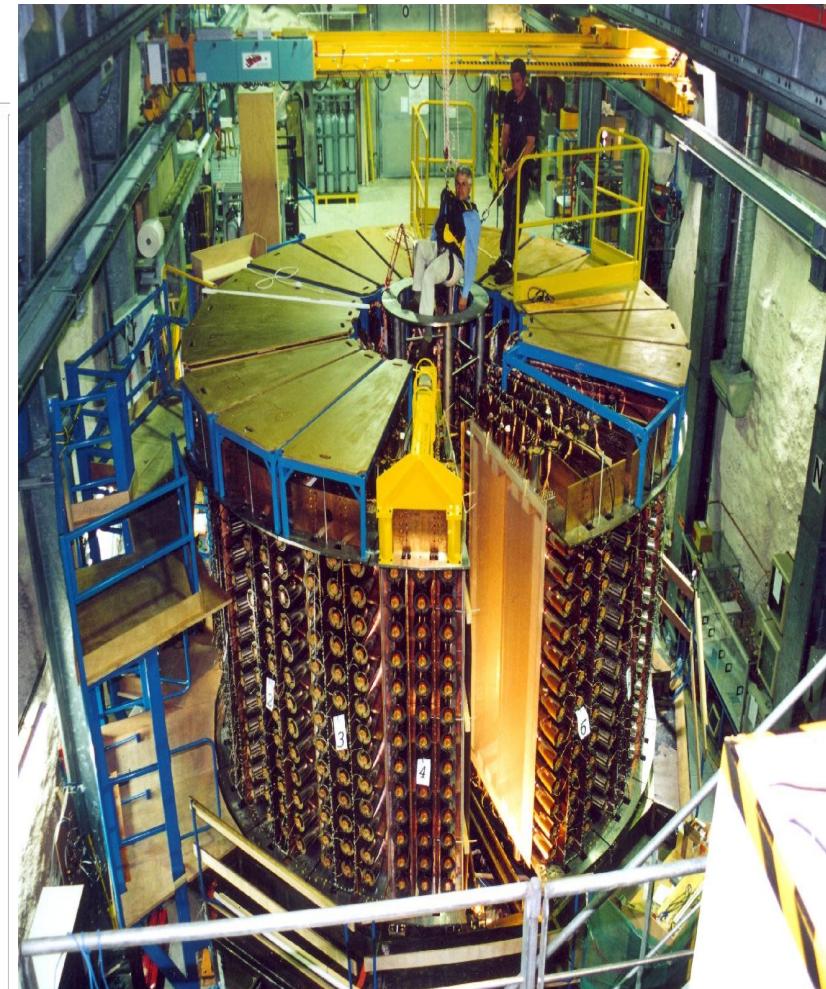


NEMO3: Neutrino ettore Majorana Observatory



**LS Modane, FR
Tunnel Frejus**

4800 m.w.e



During installation AUGUST 2001

F. Piquemal (CENBG)

NEMO3 detector

tracking detector He + 4% ethil alcohol + Ar
6180 vertical drift cells operated in Geiger mode

each cell has an octagonal shape and is 2.7 m long, 3D tracking is accomplished with the arrival time of the signals on the anode wires and the plasma propagation times to the ends of the drift cells

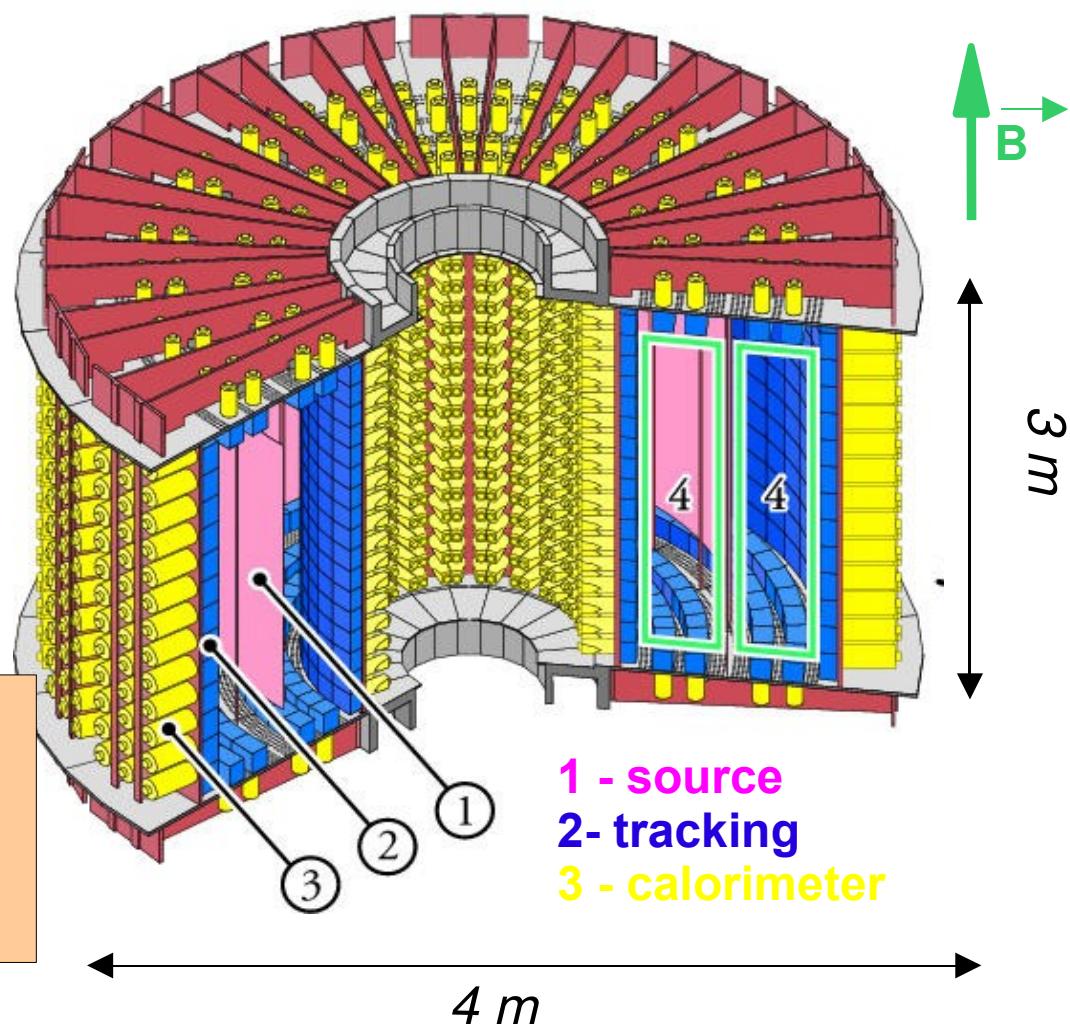
calorimeter

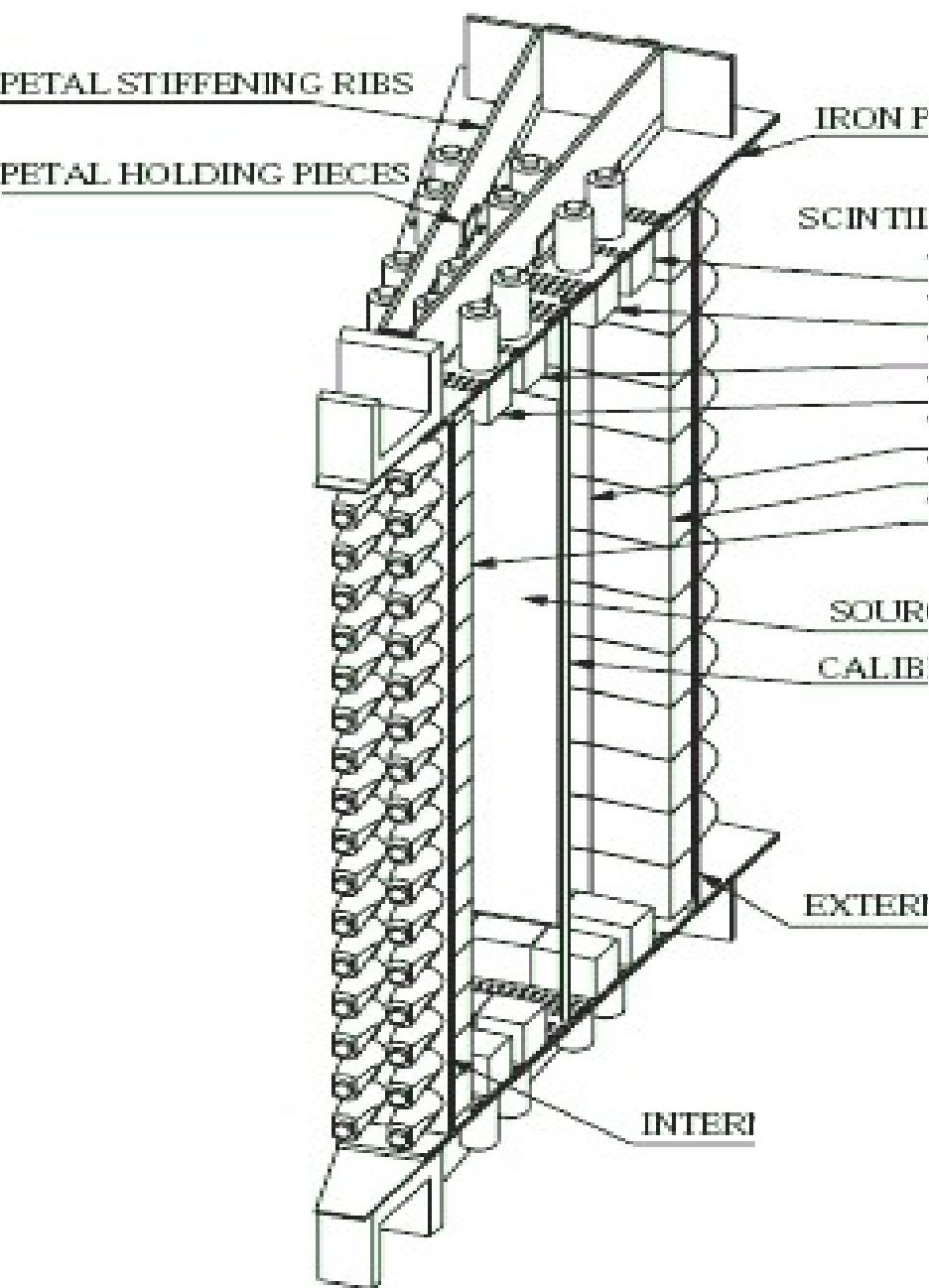
1940 plastic scintillators + PMTs
for energy and TOF measurement

magnetic field $B=25\text{ G}$
for charge identification

provides:

- track reconstruction
- energy resolution $\sigma_E/E \sim 3\%$ at 3 MeV
- particle identification

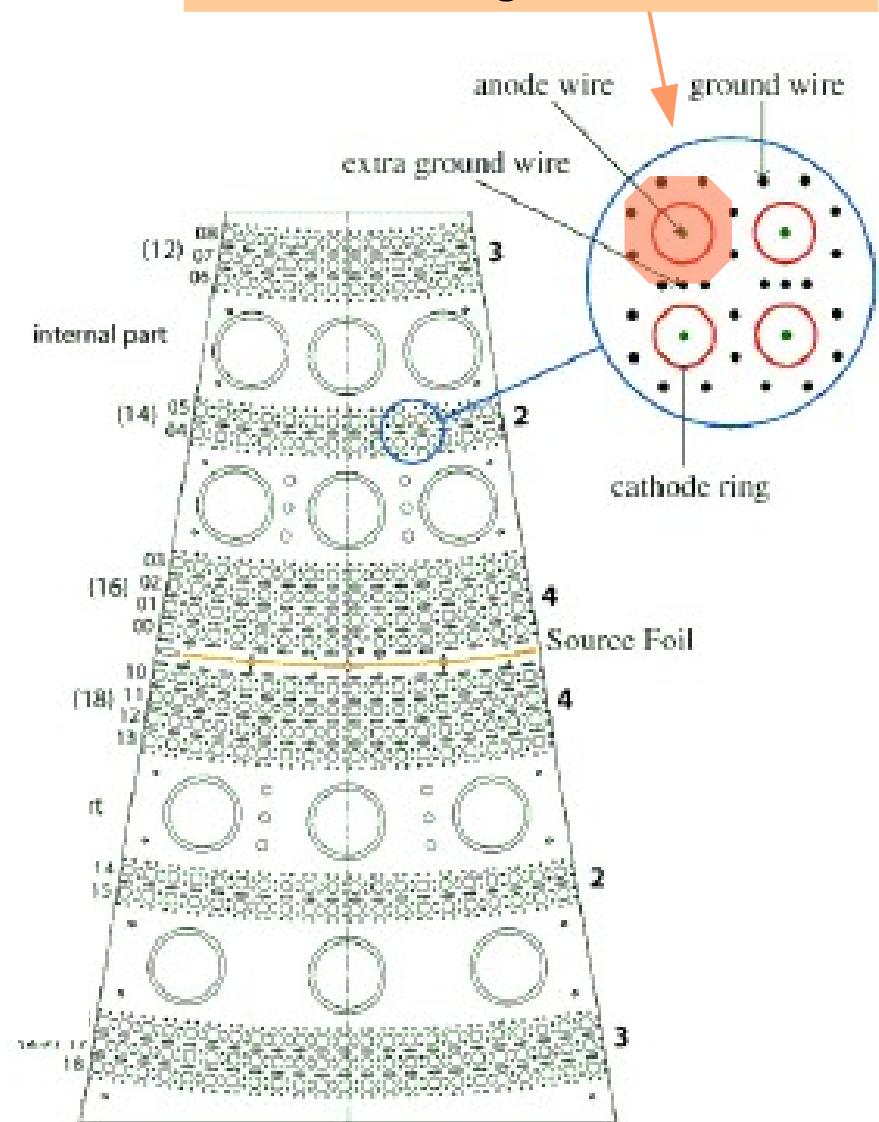




NEMO3 detector: a sector

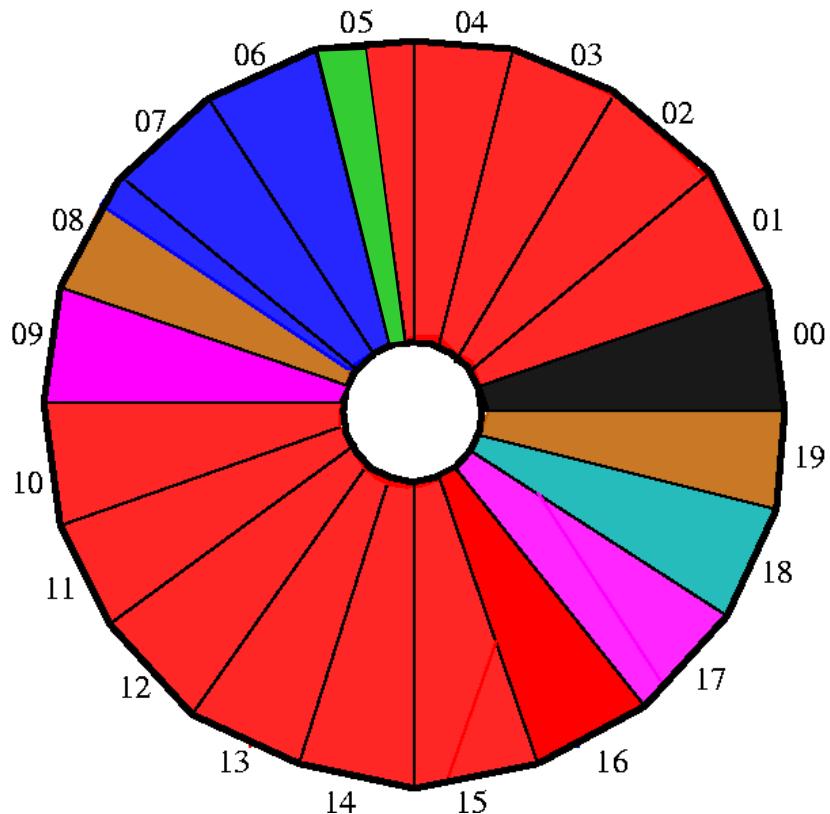
top view of a drift cell=

1 anode wire + 9 ground wires
a cathode ring



NEMO3: the sources

sources are foils (30-60 mg/cm²) of different material forming a cylindrical patchwork inserted in the middle of the detector



¹⁰⁰Mo **6.914 kg**

$Q_{\beta\beta} = 3034 \text{ keV}$

⁸²Se **0.932 kg**

$Q_{\beta\beta} = 2995 \text{ keV}$

$\beta\beta 0\nu$ search

$\beta\beta 2\nu$ measurement

¹¹⁶Cd **405 g**

$Q_{\alpha\alpha} = 2805 \text{ keV}$

⁹⁶Zr **9.4 g**

$Q_{\beta\beta} = 3350 \text{ keV}$

¹⁵⁰Nd **37.0 g**

$Q_{\beta\beta} = 3367$

⁴⁸Ca **7.0 g**

$Q_{\beta\beta} = 4272 \text{ keV}$

¹³⁰Te **454 g**

$Q_{\beta\beta} = 2529 \text{ keV}$

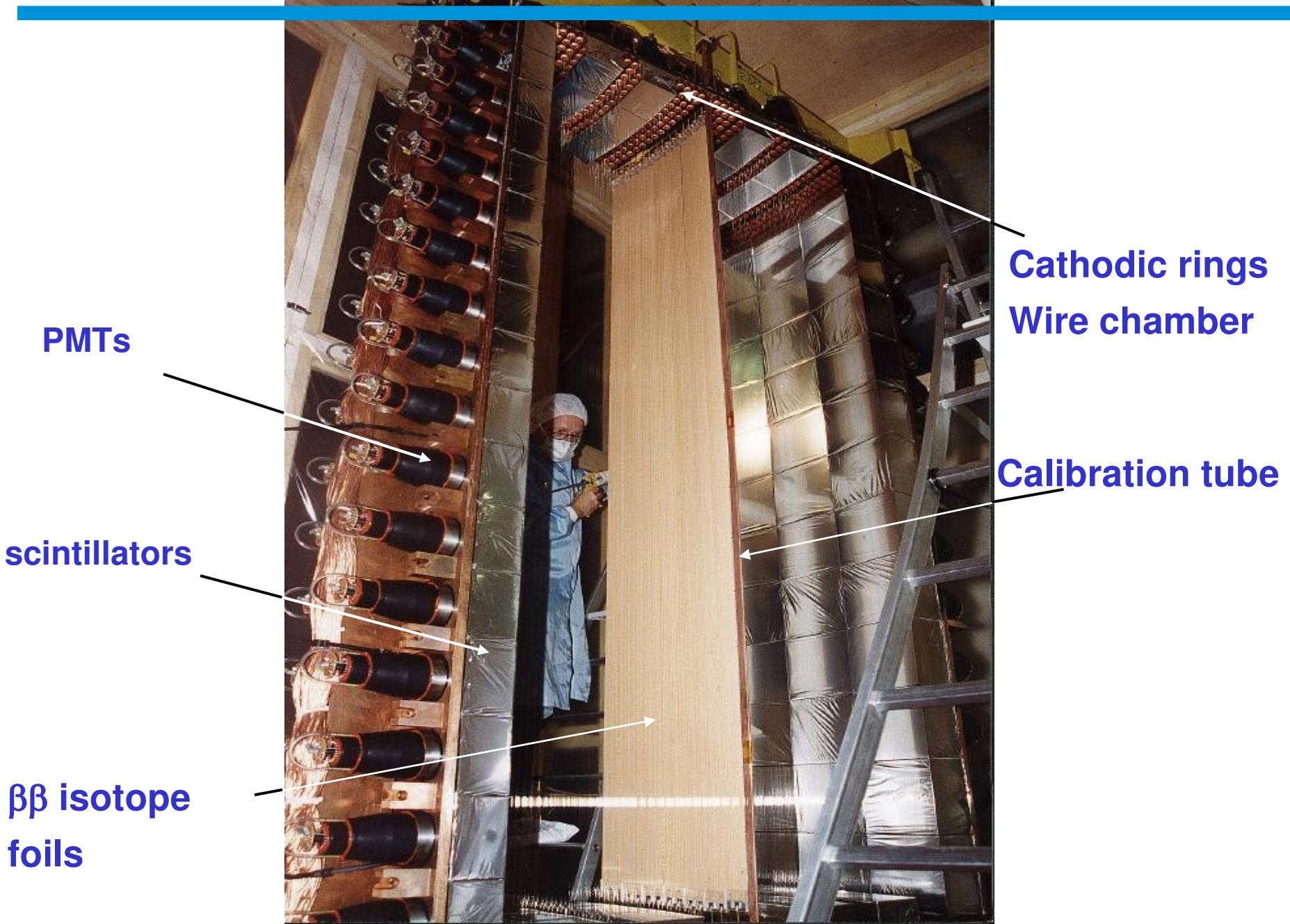
^{nat}Te **491 g**

Cu **621 g**

**External bkg
measurement**

(All enriched isotopes produced in Russia)

NEMO3: detector



NEMO3: shields



top/bottom wood shield
40 cm thick

moderates/absorbs neutrons

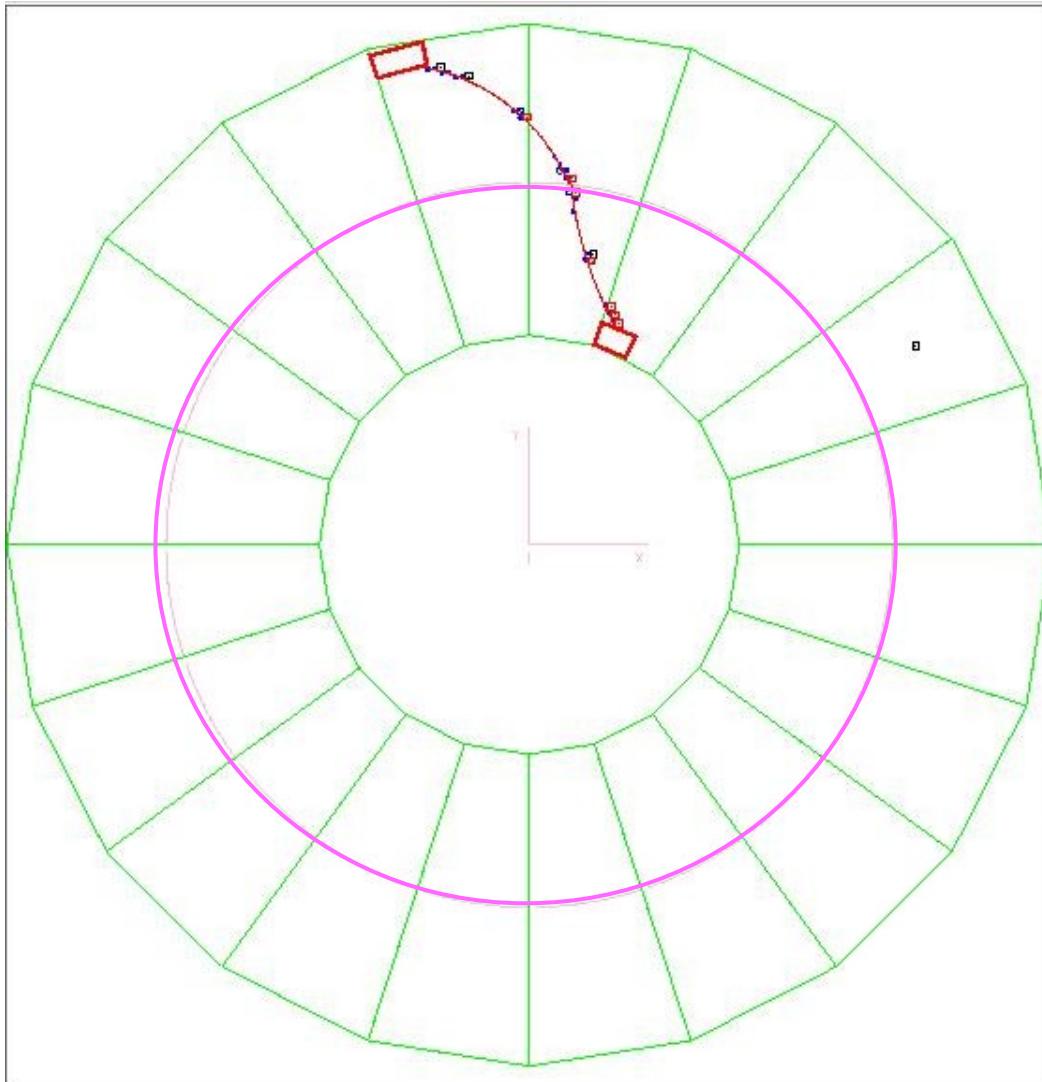
iron shield
18 cm thick

shields from
environmental γ 's

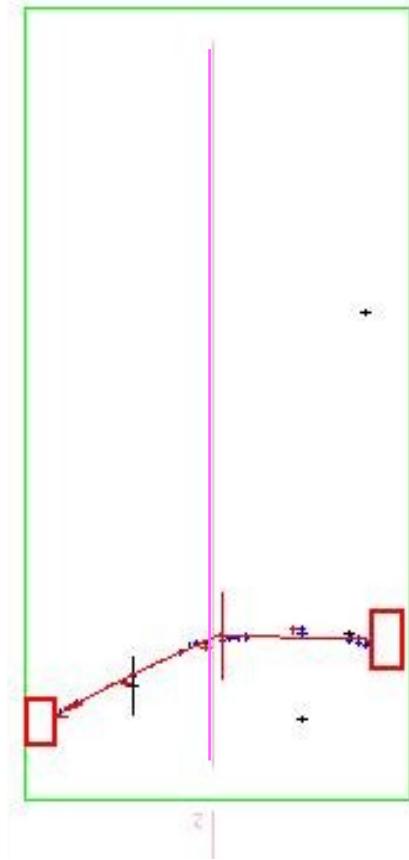
water tank
lateral, 30 cm thick

moderates/absorbs neutrons
since 2004 added boron

NEMO3: a $\beta\beta$ event

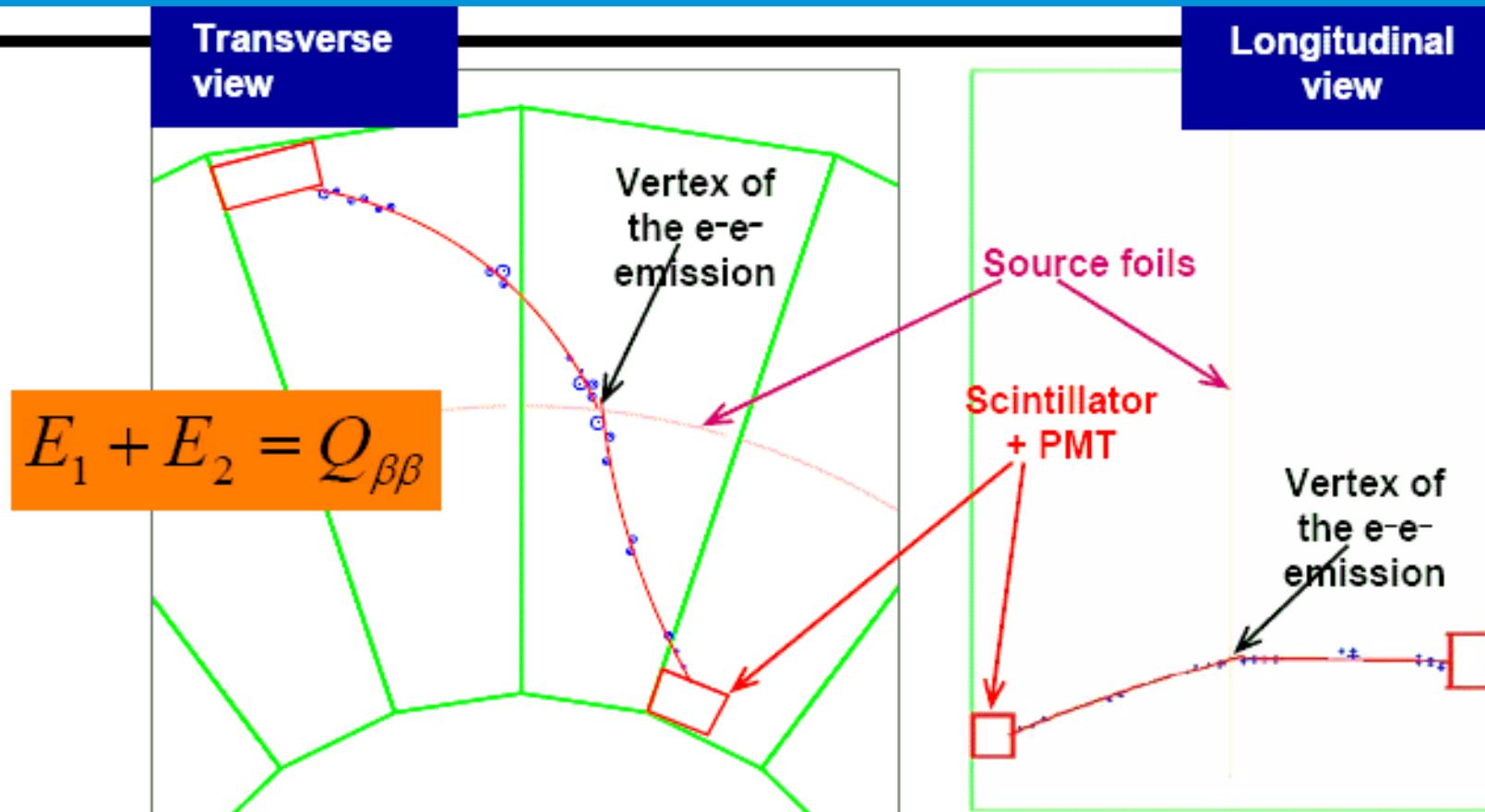


Top view



Side view

NEMO3: a $\beta\beta$ event

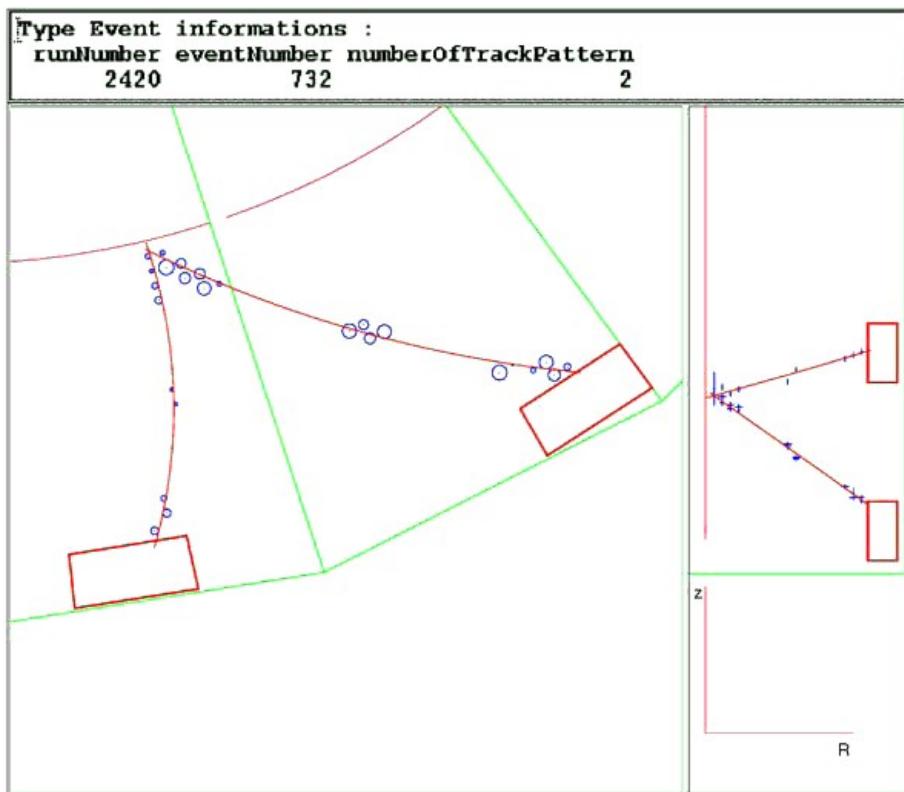


Observables of the final state

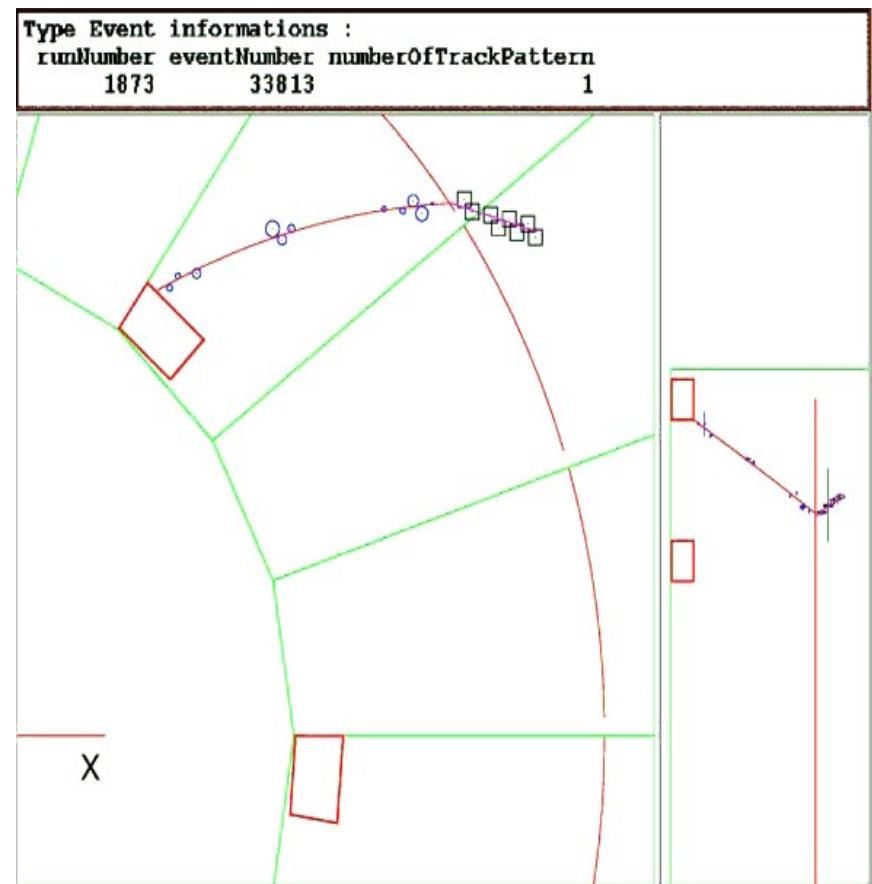
- Trajectories of the 2 electrons
- Energies of the 2 electrons
- Time of flight
- Curvature of the tracks in a B-field (+ or -).

NEMO3: background

most background sources produce events quite different from those due to a $\beta\beta$ decay



$e^+ e^-$ background event



Decay of some internal contamination producing a single electron event coming from a molybdenum source foil (sector 02) followed by a delayed alpha particle, which is the short straight track represented by open squares. Note the presence of one gamma-ray

NEMO3: background events that can mimic a $0\nu\beta\beta$ decay

^{100}Mo $0\nu\beta\beta$ decay - number of events in the $2.8 < E_1 + E_2 < 3.2 \text{ MeV}$ window

External Background ^{208}TI (PMTs)

$\sim 10^{-3} \beta\beta 0\nu$ -like events $\text{year}^{-1} \text{kg}^{-1}$

External Neutrons and High Energy gamma

$\sim 0.02 \beta\beta 0\nu$ -like events $\text{year}^{-1} \text{kg}^{-1}$

^{208}TI impurities inside the foils

$\sim 0.06 \beta\beta 0\nu$ -like events $\text{year}^{-1} \text{kg}^{-1}$

^{214}Bi impurities inside the foils

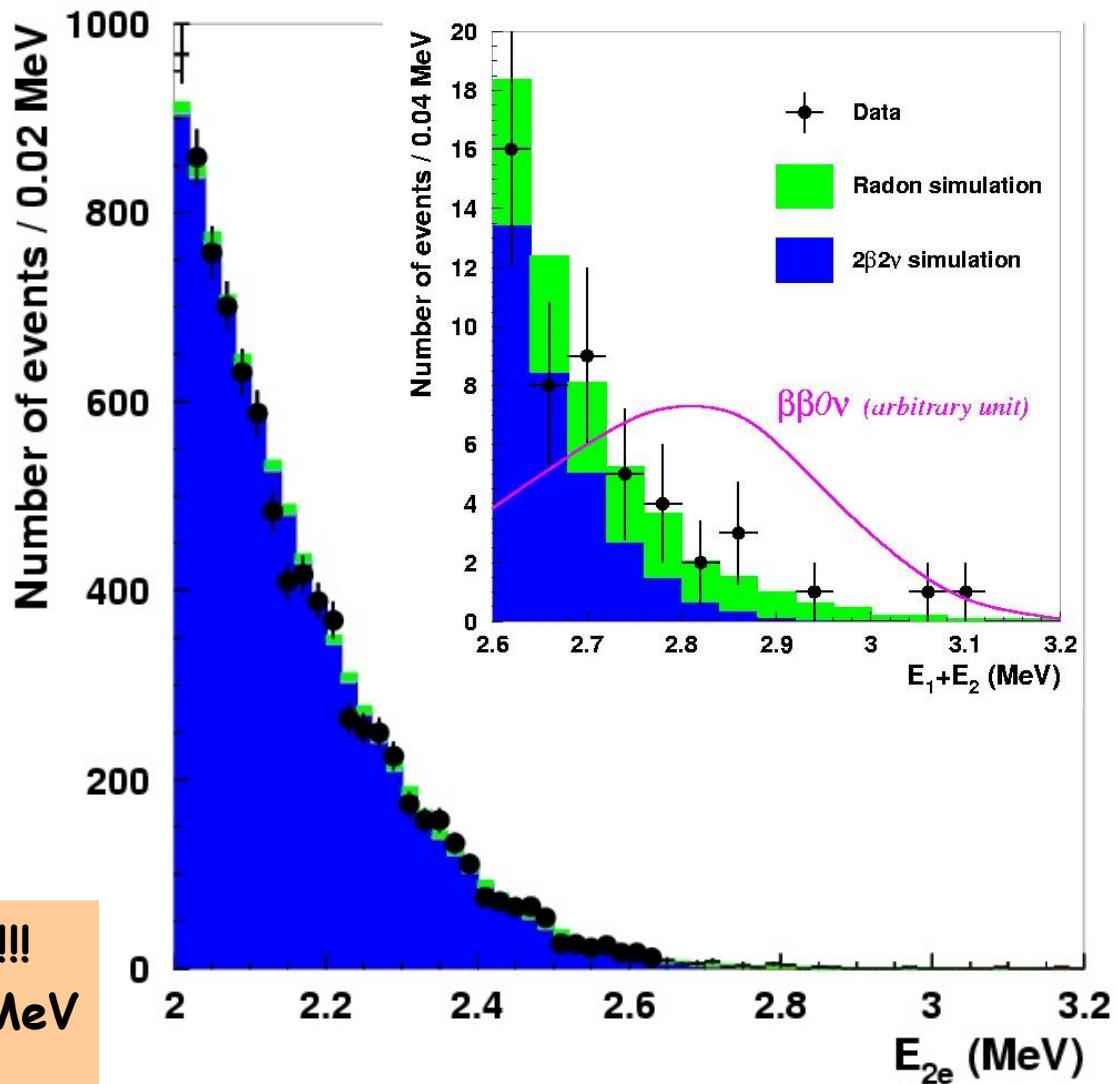
$< 0.1 \beta\beta 0\nu$ -like events $\text{year}^{-1} \text{kg}^{-1}$

Dominated by radon in first period, Radon suppress

^{100}Mo $\beta\beta 2\nu$ decay $T_{1/2} = 7.14 \cdot 10^{18} \text{ y}$

$\sim 0.3 \beta\beta 0\nu$ -like events $\text{year}^{-1} \text{kg}^{-1}$

the ultimate bkg comes from $2\nu\beta\beta$!!!
(energy resolution FWHM $\sim 8\%$ at 3 MeV
= 240 keV !!!)



NEMO3 $0\nu\beta\beta$ results (V-A):

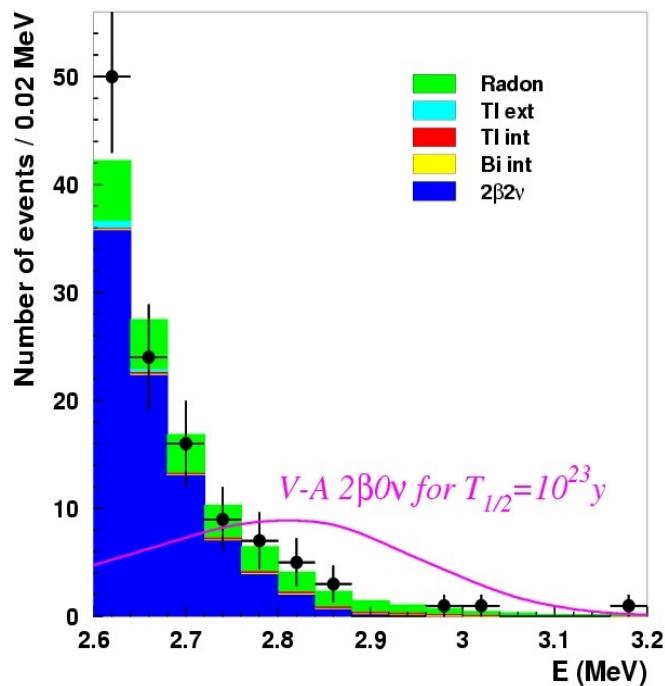
^{100}Mo

^{82}Se

$$T_{1/2}(\beta\beta 0\nu) > 5.8 \cdot 10^{23} \text{ (90 \% C.L.)}$$

$$T_{1/2}(\beta\beta 0\nu) > 2.1 \cdot 10^{23} \text{ (90 \% C.L.)}$$

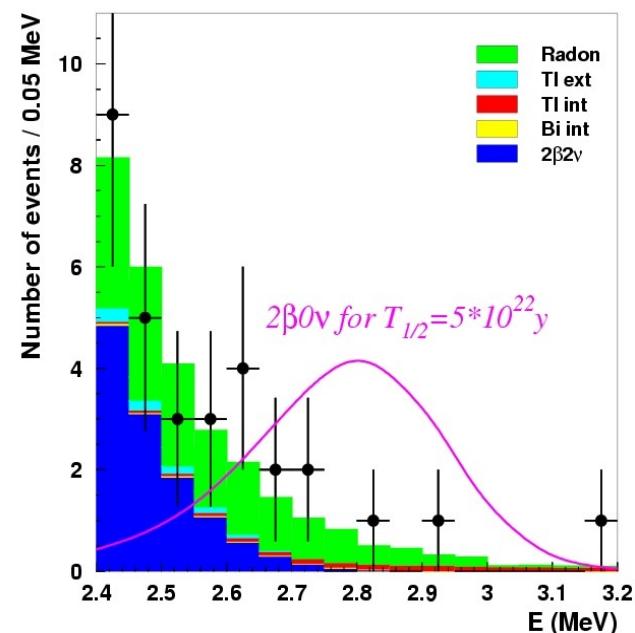
(evts $\text{y}^{-1} \text{ kg}^{-1}$ [2.8-3.2] MeV)



Phase I + II
693 days

1

$$m_\nu < (0.5 - 2.5) \text{ eV}$$



$$m_\nu < (1 - 5.9) \text{ eV}$$

NEMO3 $0\nu\beta\beta$ results ($V+A$) and Majoron:

old values to be updated

Limit on Majoron

^{100}Mo : $T_{1/2}(\beta\beta 0\nu M) > 1.8 \ 10^{22} \text{ y}$

$g_M < (5.3 - 8.5) \ 10^{-5}$ (best limit)

Simkovic (1999), Stoica (1999)

^{82}Se : $T_{1/2}(\beta\beta 0\nu M) > 1.5 \ 10^{22} \text{ y}$

$g_M < (0.7 - 1.6) \ 10^{-4}$

Simkovic (1999), Stoica (2001)

Limit on $V+A$

^{100}Mo : $T_{1/2}(\beta\beta 0\nu V+A) > 2.3 \ 10^{23} \text{ y}$

$\lambda < (1.5 - 2.0) \ 10^{-6}$

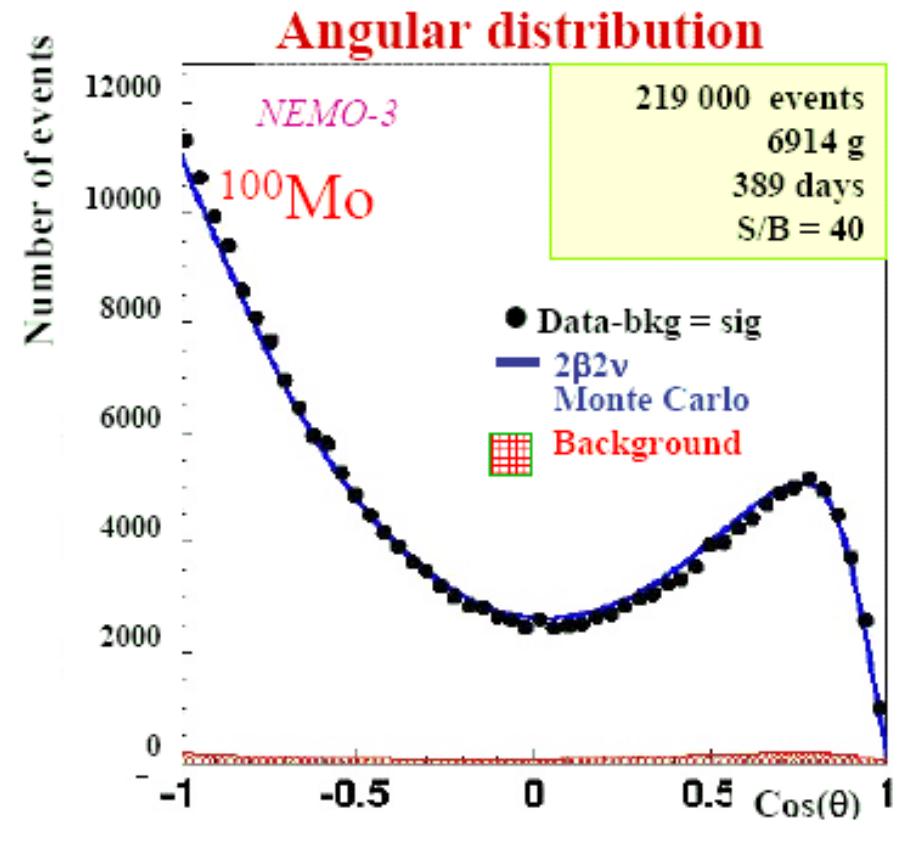
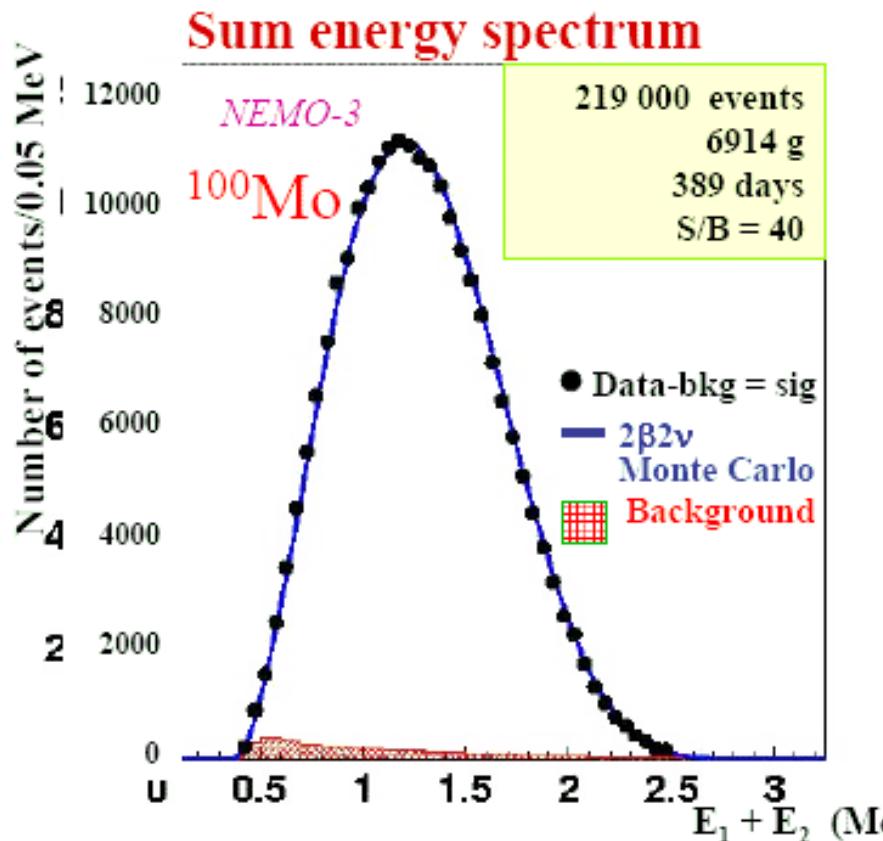
Tomoda (1991), Suhonen (1994)

^{82}Se : $T_{1/2}(\beta\beta 0\nu V+A) > 1.0 \ 10^{23} \text{ y}$

$\lambda < 3.2 \ 10^{-6}$

Tomoda (1991)

NEMO3 2νββ results 100Mo 2νββ results



$$T_{1/2}(2\nu\beta\beta) = 7.11 \pm 0.02 \text{ (stat)} \pm 0.54 \text{ (syst)} \times 10^{18} \text{ years}$$

Phase 1: Feb. 2003 – Dec. 2004
“High Radon”

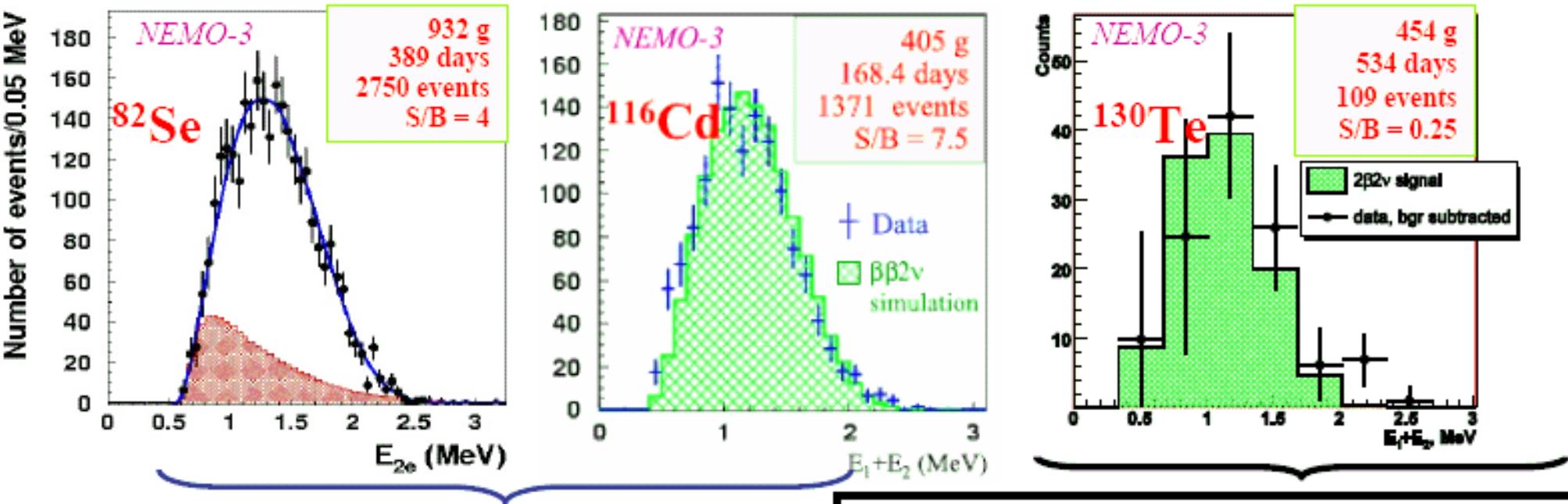
Phys. Rev. Lett. 95 182302 (2005)

28 May 2008

NEMO-3 Neutrino08

Now have in excess of
0.5M events and will
update later this year.

NEMO3 $2\nu\beta\beta$ results



Results for Phase I data. Additional statistics are being analysed and to be published soon.

^{82}Se :

$$T_{1/2} = [9.6 \pm 0.3 \text{ (stat)} \pm 1.0 \text{ (syst)}] \times 10^{19} \text{ y}$$

^{116}Cd :

$$T_{1/2} = [2.8 \pm 0.1 \text{ (stat)} \pm 0.3 \text{ (syst)}] \times 10^{19} \text{ y}$$

Preliminary:
Result for Phase 1 and 2 data.

^{130}Te :

$$T_{1/2} = [7.6 \pm 1.5 \text{ (stat)} \pm 0.8 \text{ (syst)}] \times 10^{20} \text{ y}$$

$2\nu\beta\beta$ is important: 1) Experimental input to NME calculation
2) Ultimate background for $0\nu\beta\beta$

NEMO3 result vs. HM claim of evidence:

HM claim of evidence:

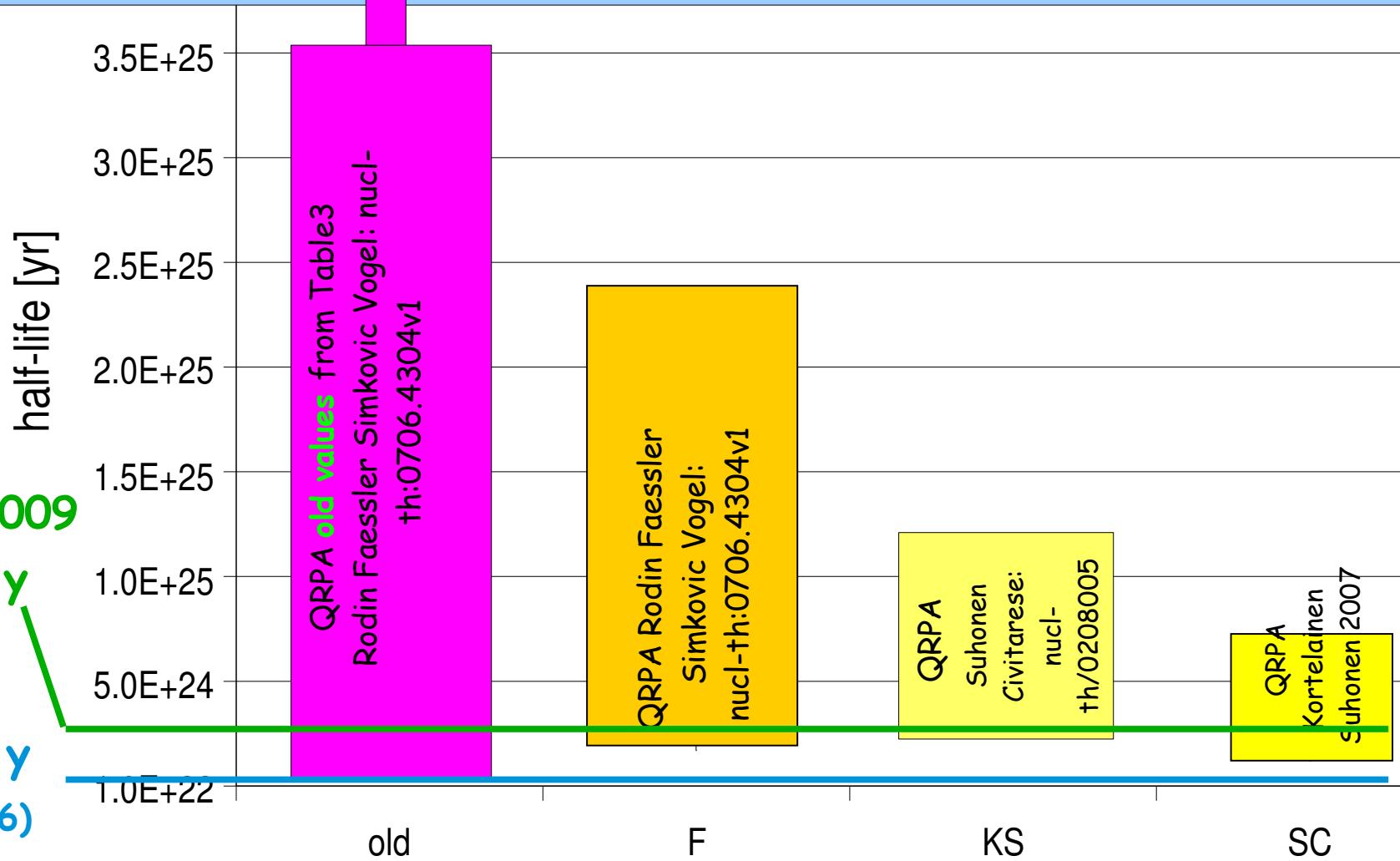
$$T_{1/2}^{0\nu} (y) = (0.69-4.18) \times 10^{25} \text{ (3}\sigma\text{ range)}$$



^{100}Mo predicted

$T_{1/2}^{0\nu} (y)$ range

$$T_{1/2}^{-1} = F(Q_{\beta\beta}, Z) |M|^2 \langle m_\nu \rangle^2$$

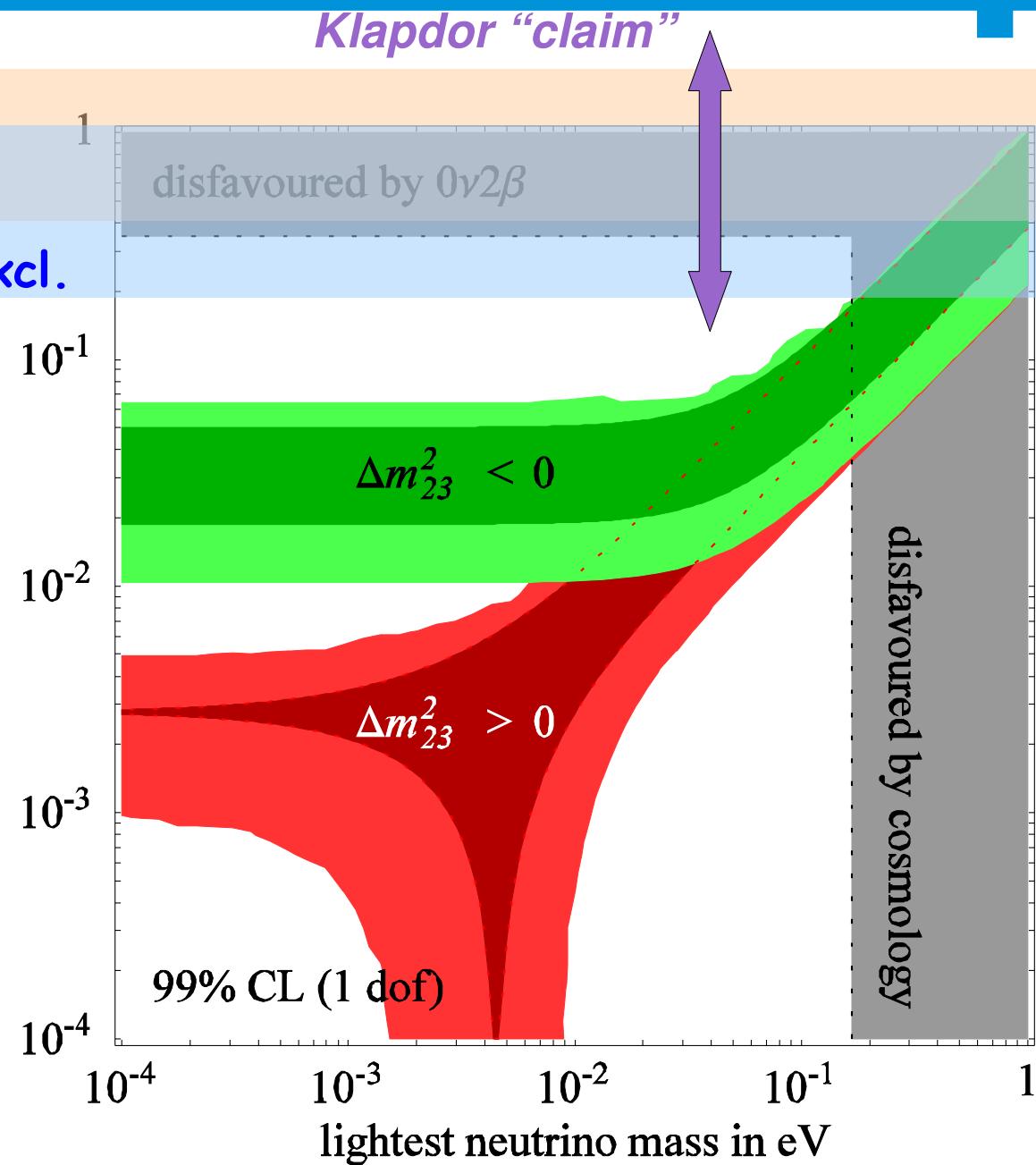


Past/Present experiments

Isotope	τ	90% limit [y]	Experiment	Low [eV]	High [eV]
^{76}Ge	1.9E+25		HM	1.0	0.2
^{82}Se	1.2E+23		NEMO3	5.9	1.1
^{100}Mo	5.8E+23		NEMO3	2.6	0.5
^{130}Te	3.0E+24		Cuoricino '07	0.7	0.2

QRPA NME from:

Rodin et al. Table 3 Nucl. Phys. A 2006
+ erratum nucl-th:0706.4304v1



Bibliography

several slides adapted from talks at:

Neutrino 2006

<http://neutrinosantafe06.com/>

Neutrino 2008

<http://www2.phys.canterbury.ac.nz/~jaa53/>

Varenna school 2008

http://www.sif.it/SIF/en/portal/activities/fermi_school/mmviii_en

NOW 2008

<http://www.ba.infn.it/~now2008/>