

Progress on Hyperon Resonances

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Puze Gao, Jun Shi, B. S. Zou, Phys. Rev. C86 (2012) 025201

Jun Shi, B. S. Zou, ArXiv: 1411.0486 [hep-ph]

C. S. An, B. Metsch, B. S. Zou, Phys. Rev. C87 (2013) 065207

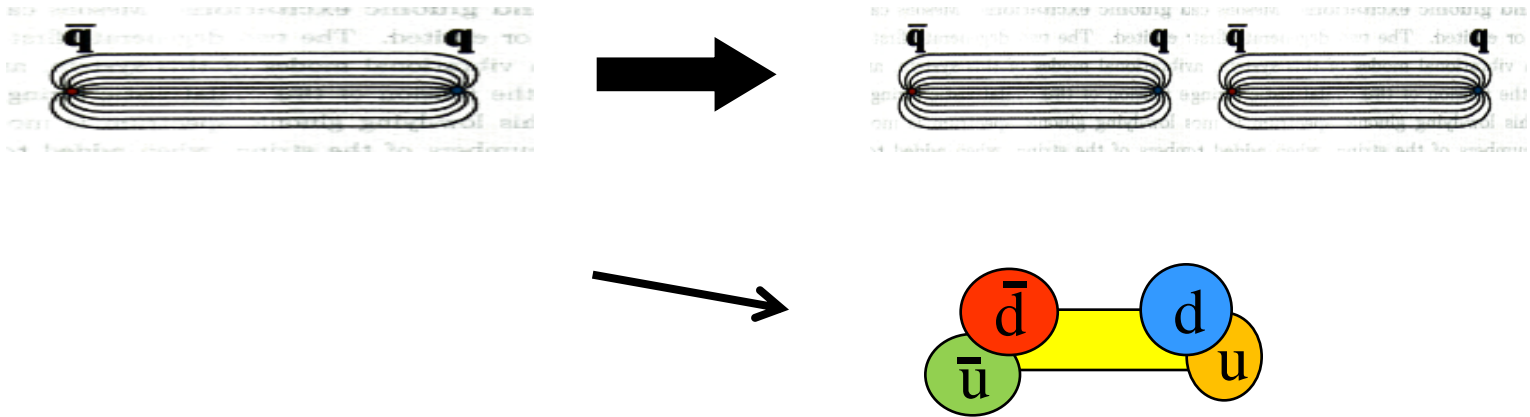
C. S. An, B. S. Zou, Phys. Rev. C89 (2014) 055209

Outline :

- 1. Why hyperon resonances ?**
- 2. New results on Σ^* & Λ^* from CB data**
- 3. Prediction on Ω^* from NJL unquenched quark model**
- 4. Conclusions and Prospects**

1. Why hyperon resonances ?

Unquenched dynamics: **gluons** \rightarrow $\bar{q}q$
crucial for quark confinement & hadron structure



quenched or unquenched quark models give very different predictions of hyperon spectrum

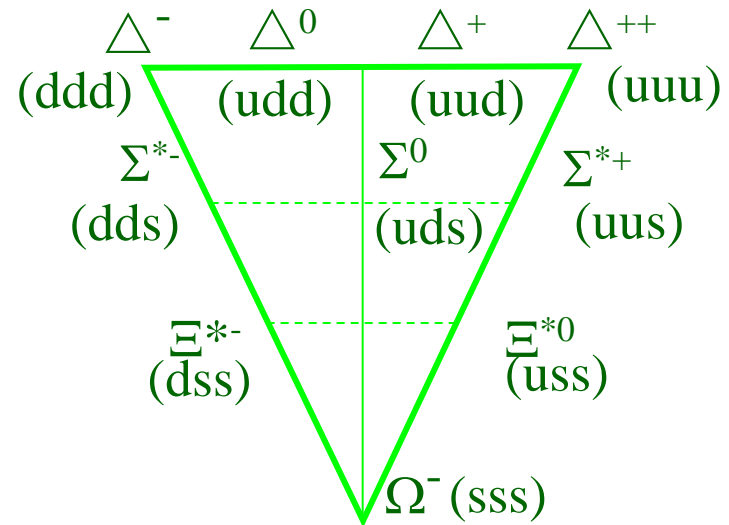
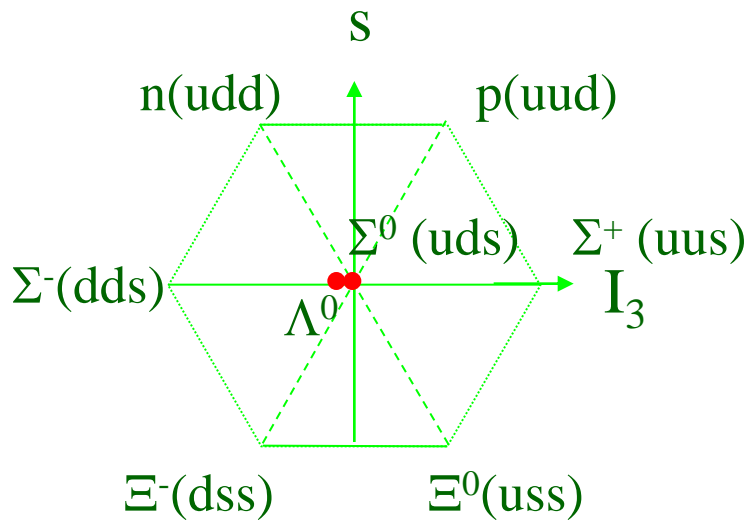
Hadron spectroscopy with strangeness

SU(3) 3q-quark model for baryons

1/2 +

spin-parity

3/2 +



Successful for spatial ground states !

Prediction $m_{\Omega^-} \cong 1670 \text{ MeV}$

experiment $m_{\Omega^-} \cong 1672.45 \pm 0.29 \text{ MeV}$

1/2⁻ baryon nonet with strangeness

- Mass pattern : quenched or unquenched ?

$$\text{uds (L=1) } 1/2^- \sim \Lambda^*(1670) \sim [\text{us}][\text{ds}] \bar{s}$$

$$\text{uud (L=1) } 1/2^- \sim \text{N}^*(1535) \sim [\text{ud}][\text{us}] \bar{s}$$

$$\text{uds (L=1) } 1/2^- \sim \Lambda^*(1405) \sim [\text{ud}][\text{su}] \bar{u}$$

$$\text{uus (L=1) } 1/2^- \sim \Sigma^*(1390) \sim [\text{us}][\text{ud}] \bar{d}$$

Zou et al, NPA835 (2010) 199 ; CLAS, PRC87(2013)035206

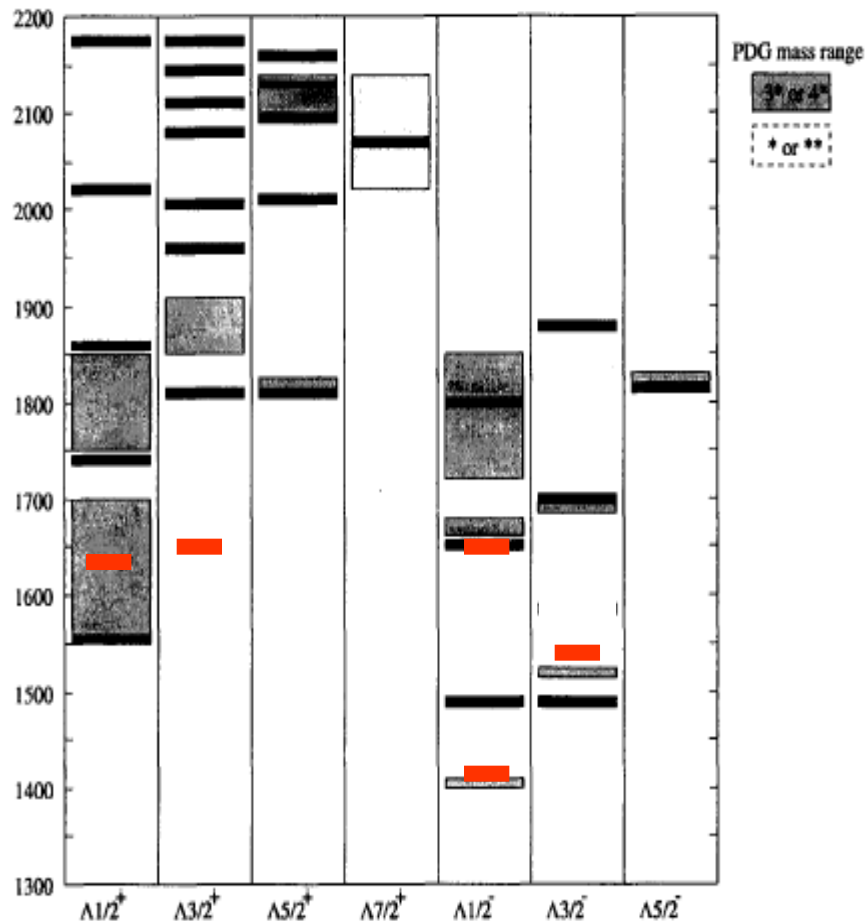
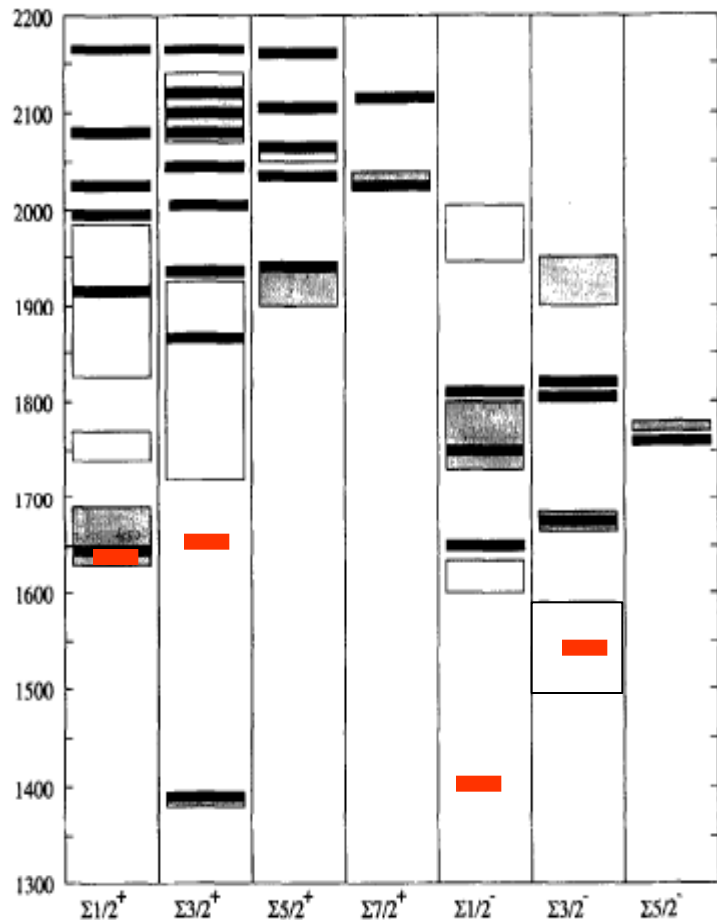
- Strange decays of N*(1535) and $\Lambda^*(1670)$:

N*(1535) large couplings $g_{\text{N}^*\text{N}\eta}$, $g_{\text{N}^*\text{K}\Lambda}$, $g_{\text{N}^*\text{N}\eta'}$, $g_{\text{N}^*\text{N}\phi}$

$\Lambda^*(1670)$ large coupling $g_{\Lambda^*\Lambda\eta}$

Distinctive

Predictions by quenched - & unquenched - quark models



Quenched quark model: Capstick-Roberts, Prog.Part.Nucl.Phys. 45 (2000) S241-S331

Unquenched model: Helminen-Riska, Nucl. Phys. A 699 (2002) 624

A.Zhang, S.L.Zhu et al., HEPNP 29 (2005) 250

Alternative pictures :

Hadronic molecules

$$N^*(1440) \sim N\sigma$$

$$N^*(1535) \sim K\Sigma-K\Lambda$$

$$\Lambda^*(1405) \sim KN-\Sigma\pi$$

Penta-quark states

$$N^*(1440) \sim [ud][ud] \bar{q}$$

$$N^*(1535) \sim [ud][us] \bar{s}$$

$$\Lambda^*(1405) \sim [ud][sq] \bar{q}$$

**Kaiser, Weise, Oset, Ramos, Oller,
Meissner, Hyodo, Jido, Hosaka, Oh, ...**

Distinguishable model predictions for Σ^* of $3/2^-$ and $1/2^+$

qqq

$\bar{q}q^6$ or $\bar{K}\pi N-\pi\pi Y$

$3/2^-$ $\Sigma^*(1650)$

$\Sigma^*(1570)$

Gal 2011

$1/2^+$ $\Sigma^*(1720)$

$\Sigma^*(1630)$ & $\Sigma^*(1656)$ Oset 2008

Experiment knowledge on hyperon states still very poor !

Ω^* in PDG:

- **** $\Omega(1672) 3/2^+$,
- *** $\Omega(2250)$
- ** $\Omega(2380), \Omega(2470)$

Ξ^* in PDG:

- **** $\Xi(1320) 1/2^+, \Xi(1530) 3/2^+$
- *** $\Xi(1690), \Xi(1820) 3/2^-, \Xi(1950), \Xi(2030)$
- ** $\Xi(2250), \Xi(2370)$
- * $\Xi(1620), \Xi(2120), \Xi(2500)$

Σ^* in PDG2012

**** $\Sigma(1189)1/2^+$ $\Sigma^*(1385)3/2^+$ $\Sigma^*(1670)3/2^-$
 $\Sigma^*(1775)5/2^-$ $\Sigma^*(1915)5/2^+$ $\Sigma^*(2030)7/2^+$

*** $\Sigma^*(1660)1/2^+$ $\Sigma^*(1750)1/2^-$ $\Sigma^*(1940)3/2^-$
 $\Sigma^*(2250)??$

** $\Sigma^*(1620)1/2^-$ $\Sigma^*(1690)??$ $\Sigma^*(1880)1/2^+$
 $\Sigma^*(2080)3/2^+$ $\Sigma^*(2455)??$ $\Sigma^*(2620)??$

* $\Sigma^*(1480)??$ $\Sigma^*(1560)??$ $\Sigma^*(1580)3/2^-$
 $\Sigma^*(1770)1/2^+$ $\Sigma^*(1840)3/2^+$ $\Sigma^*(2000)3/2^-$
 $\Sigma^*(2070)5/2^+$ $\Sigma^*(2100)7/2^-$ $\Sigma^*(3000)??$
 $\Sigma^*(3170)??$

All from old experiments of 1970-1985 !!

No established $1/2^- \Sigma^*$, Ξ^* , Ω^* !

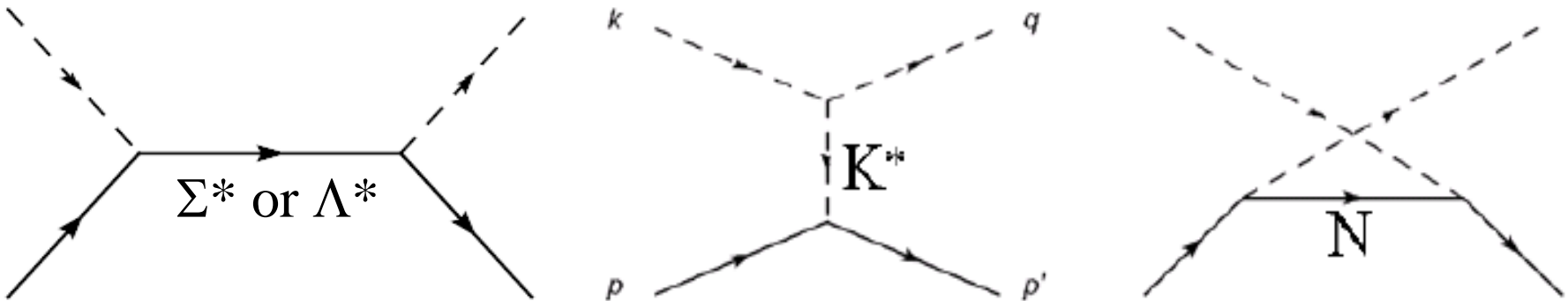
2. New results on Σ^* & Λ^* from CB data

Crystal Ball: Prakhov et al., **PRC 80**(2009) 025204

$$K^- + p \rightarrow \pi^0 + \Lambda \quad \& \quad K^- + p \rightarrow \pi^0 + \Sigma^0$$

$$p_K = 514\text{--}750 \text{ MeV}, \quad \sqrt{s} = 1569\text{--}1676 \text{ MeV}$$

The high precision new data can give valuable information on Σ^* & Λ^*



$\Sigma^*(1620)1/2^- \rightarrow$ supporting evidence for quenched qqq models ?

Problem : evidence for its existence is very shaky !

Among 4 references listed in PDG for it:

One without PWA for J^P

Two based on multi-channel analysis gave contradicted BRs

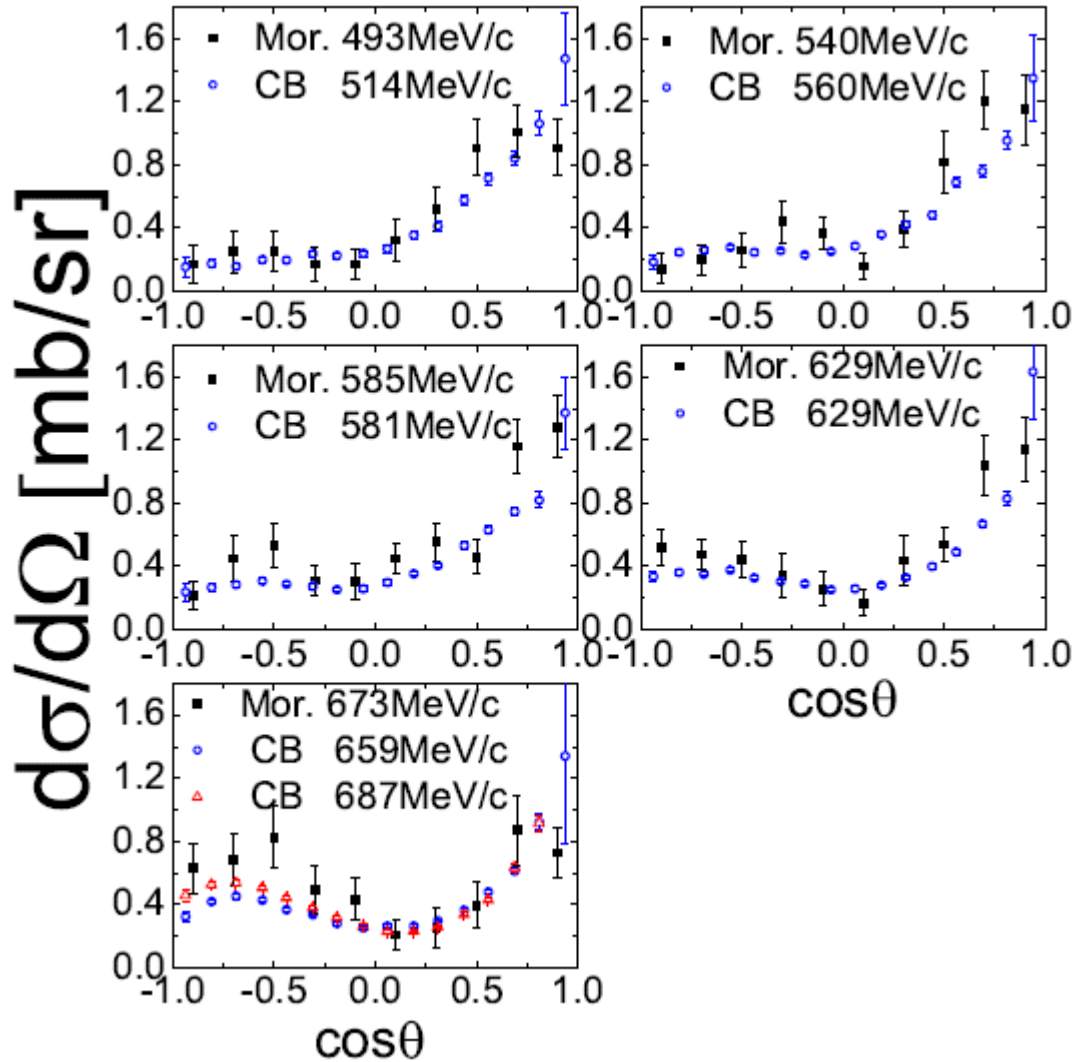
Other later multi-channel analyses claim to $\Sigma^*(1660)1/2^+$

**The 4-th gave two comparable solutions with and without it
by fitting $K^- n \rightarrow \pi^- \Lambda$ data**

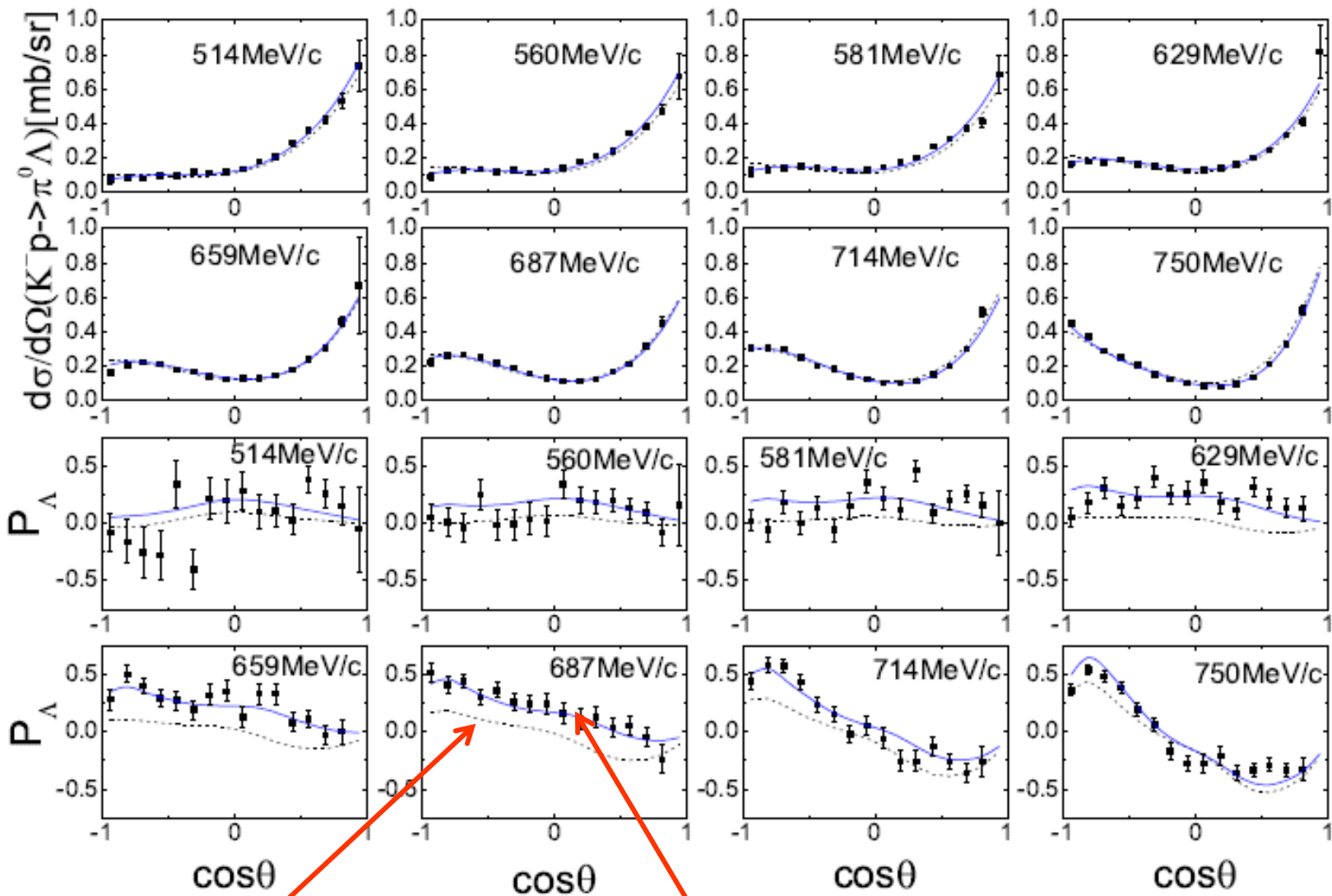
W.A. Morris et al., PRD17, 55 (1978)

**Is the new CB data compatible with the old $K^- n \rightarrow \pi^- \Lambda$ data
analyzed by W.A. Morris et al., claiming possible $\Sigma^*(1620)1/2^-$?**

new CB data on $K^-p \rightarrow \pi^0\Lambda$ vs old $K^-n \rightarrow \pi^-\Lambda$ data



new CB data on $K^-p \rightarrow \pi^0 \Lambda$



with basic ingredients

adding $\Sigma(1635) 1/2^+$

Adding *** $\Sigma(1660)1/2^+$, $\chi^2=572$ for 348 data points with 18 tunable parameters. $\Delta\chi^2=1008$!

	mass(MeV)(PDG estimate)	Γ_{tot} (MeV)(PDG estimate)	$\sqrt{\Gamma_{\pi\Lambda}\Gamma_{\bar{K}N}}/\Gamma_{tot}$ (PDG range)
$\Sigma(1670)\frac{3}{2}^-$	$1673 \pm 1(1665, 1685)$	$52_{-2}^{+5}(40, 80)$	$0.081_{-0.004}^{+0.002}(0.018, 0.17)$
$\Sigma(1660)\frac{1}{2}^+$	$1633 \pm 3(1630, 1690)$	$121_{-7}^{+4}(40, 200)$	$-0.064_{-0.003}^{+0.005}(-0.065, 0.24)$

Replace $\Sigma(1660)1/2^+$ by $\Sigma(1/2^-)$: $\chi^2 = 899$ with $M < 1360$ MeV
 Not $\Sigma(1620) 1/2^-$

by $\Sigma(3/2^+)$: $\chi^2 = 943$

by $\Sigma(3/2^-)$: $\chi^2 = 1392$

Add both $\Sigma(1660)1/2^+$ & $\Sigma(1/2^-)$: $\Delta\chi^2 = 24$ with $M = 1432$ MeV
 $\Gamma > 1000$ MeV

No $\Sigma(1620) 1/2^-$!!

CB Λ Polarization data is crucial for discriminating $\Sigma(1620)1/2^-$ from $\Sigma(1635) 1/2^+$.

Add $\Sigma(3/2^+)$: $\Delta\chi^2 = 85$ with $M \sim 1840$ MeV PDG $\Sigma(1840)3/2^+$?

Add $\Sigma(3/2^-)$: $\Delta\chi^2 = 37$ with $M \sim 1542$ MeV PDG-GAL $\Sigma(1580) 3/2^-$?

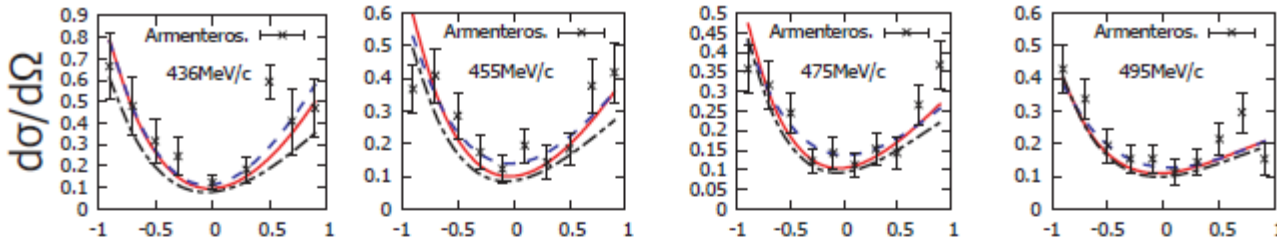
Add $\Sigma(1/2^+)$: $\Delta\chi^2 = 31$ with $M \sim 1610$ MeV two $\Sigma(1/2^+) \sim 1633$ MeV ?
as claimed by Oset et al.

For details : Puze Gao, Jun Shi, B. S. Zou, PRC86 (2012) 025201

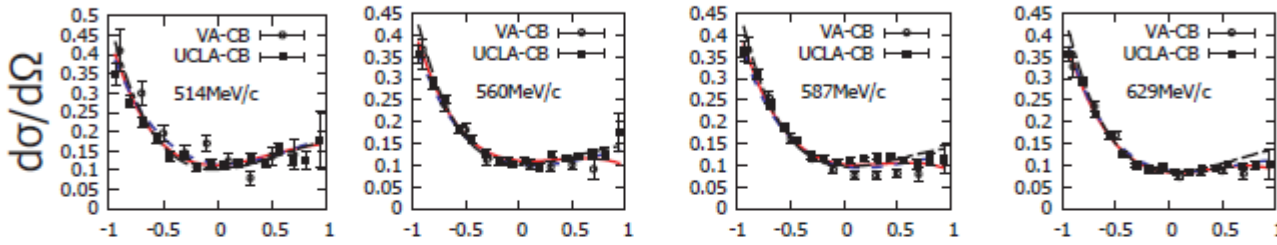
PDG2014 downgrades $\Sigma(1620)1/2^-$ from ** to *

Fits to new CB data on $K^-p \rightarrow \pi^0 \Sigma^0$ J. Shi, B. S. Zou, ArXiv: 1411.0486

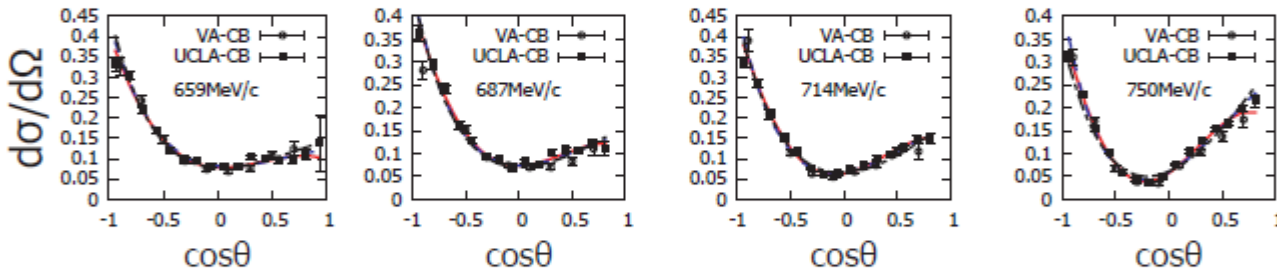
$\Lambda^*(1670)1/2^-$ **** + $\Sigma^*(1600)1/2^+$ *** $\rightarrow \chi^2 = 763$ for 236 data points



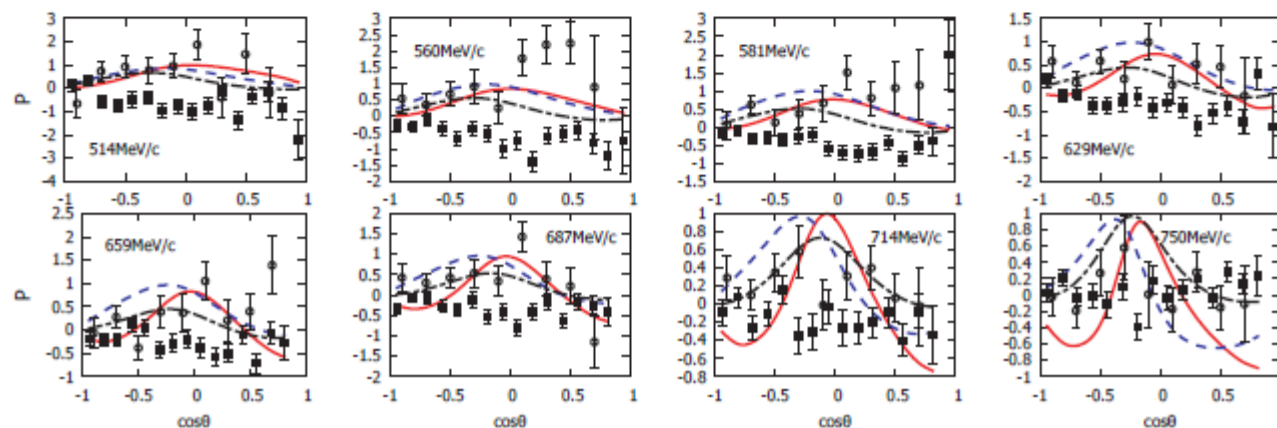
$\Lambda^*(1670)1/2^-$ & $\Lambda^*(1600)1/2^-$
 + $\Lambda^*(1690)3/2^-$ ****
 $\rightarrow \chi^2 = 540$



$\Lambda^*(1670)1/2^-$ & $\Lambda^*(1600)1/2^-$
 + $\Lambda^*(1680)3/2^+$ (new)
 $\rightarrow \chi^2 = 419$



$\Lambda^*(1680)3/2^+$ replaces
 $\Lambda^*(1690)3/2^-$ ****



Strong support for unquenched quark model!

3. Ω^* from NJL unquenched quark model

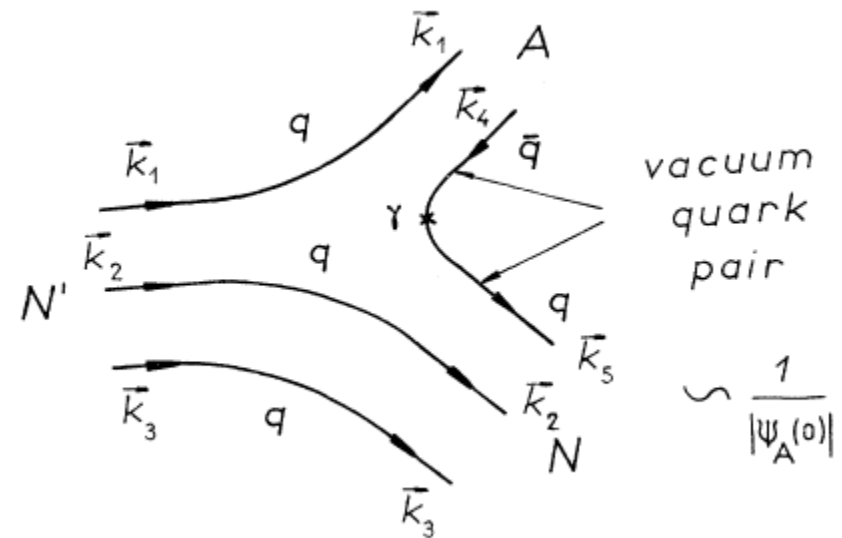
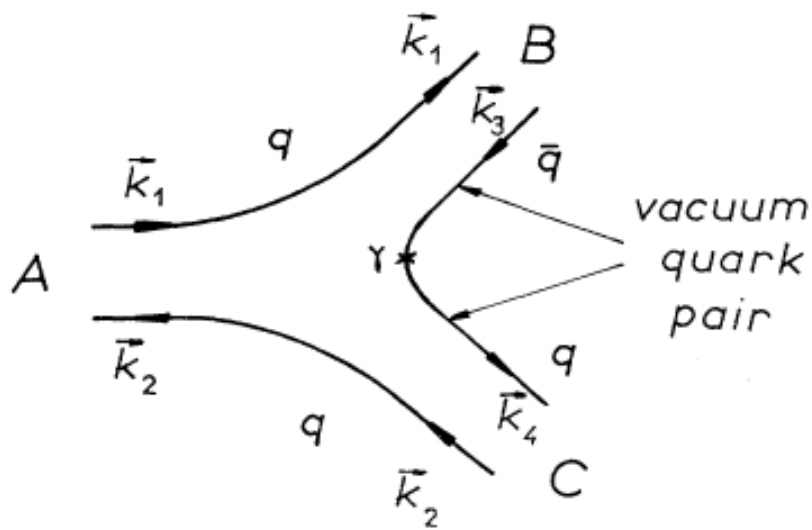
Mechanisms for $\bar{q}q$ pair production

1) Perturbative 3S_1

failed for 1^- and 1^+ decays

2) Non-perturbative 3P_0

quite successful & popular



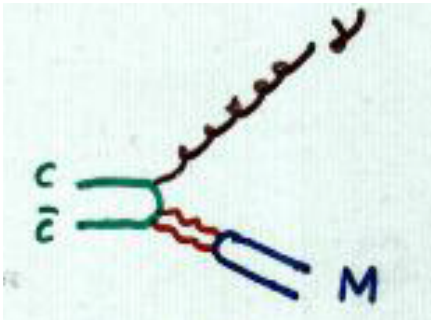
A. Le Yaouanc et al., **Phys.Rev. D8 (1973) 2223**

B. Aubert et al. (BABAR Collaboration), PRD 78 (2008) 112002:

$$\Gamma(Y(4S) \rightarrow \eta Y(1S)) / \Gamma(Y(4S) \rightarrow \pi^+ \pi^- Y(1S)) = 2.41 \pm 0.40_{\text{stat}} \pm 0.12_{\text{syst}}$$

M. Ablikim et al. (BESIII Collaboration), PRD 86 (2012) 071101

$$\Gamma(\psi(4040) \rightarrow \eta J/\psi) / \Gamma(\psi(4040) \rightarrow \pi^+ \pi^- J/\psi) > 2$$



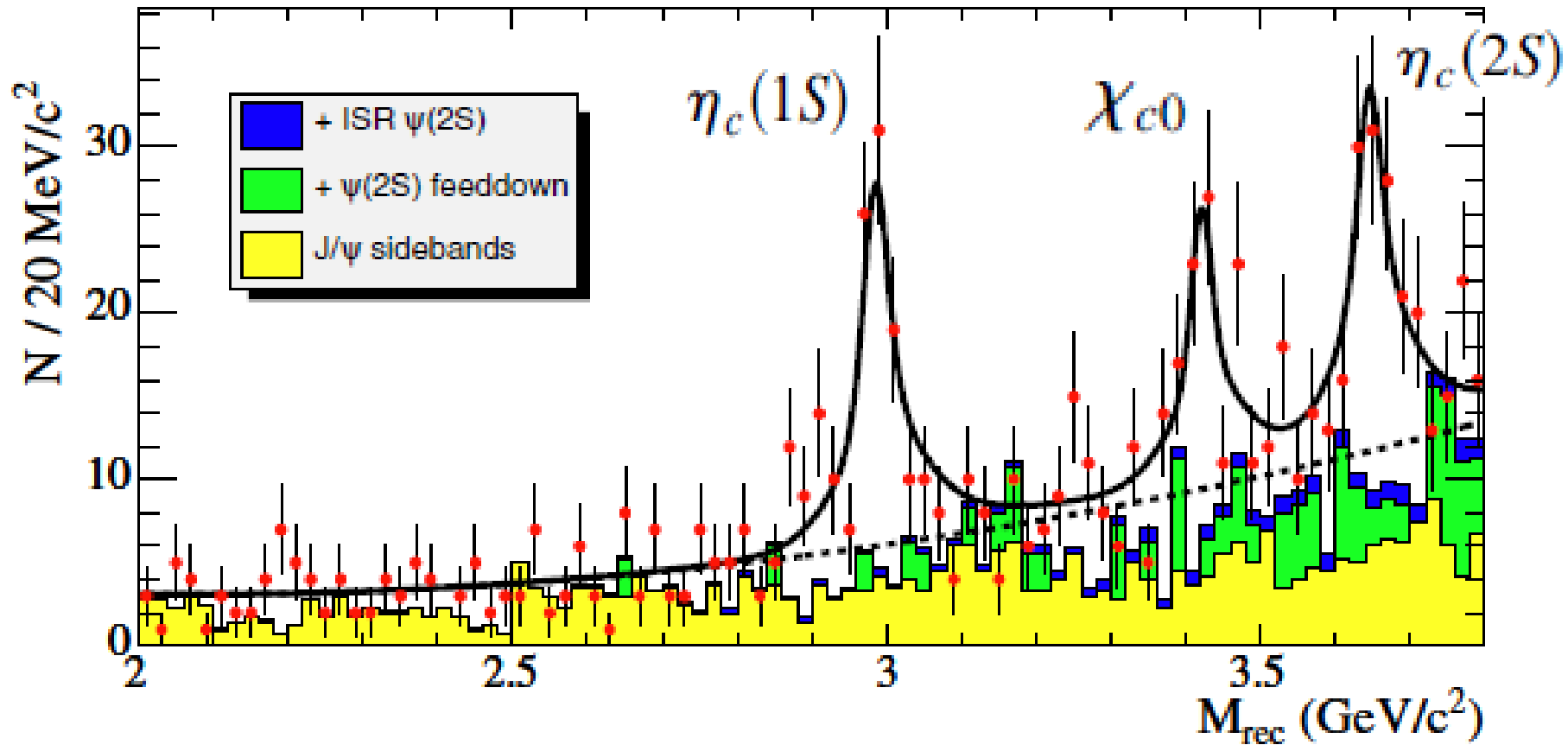
PDG:

$$\Gamma(J/\psi \rightarrow \gamma \eta) > \Gamma(J/\psi \rightarrow \gamma \sigma)$$

Glucos more favor to produce $\bar{q}q(^1S_0)$ sometimes !

$$e^+ e^- \rightarrow J/\psi c\bar{c}$$

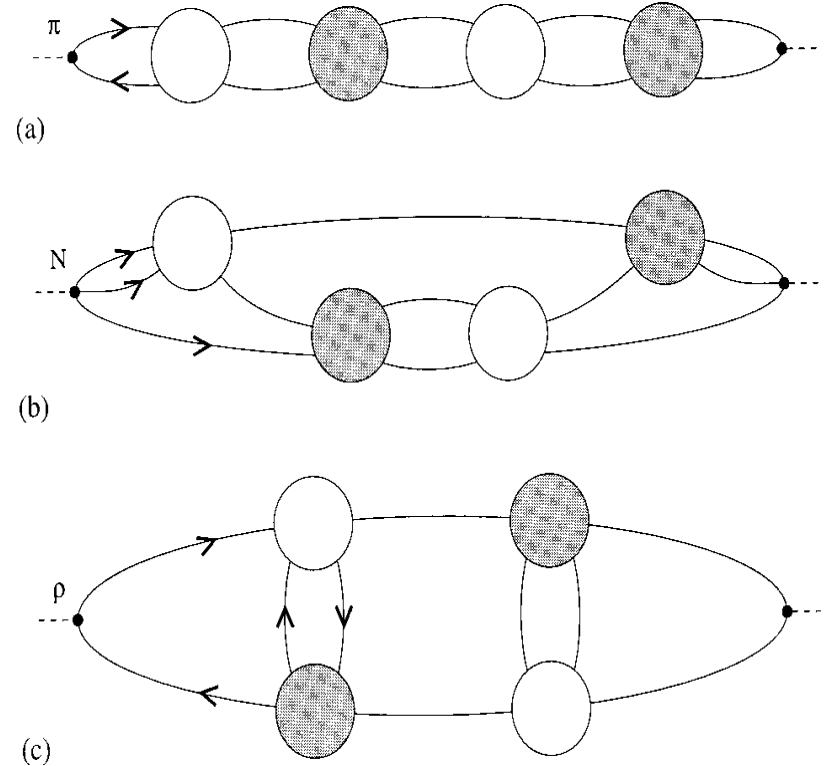
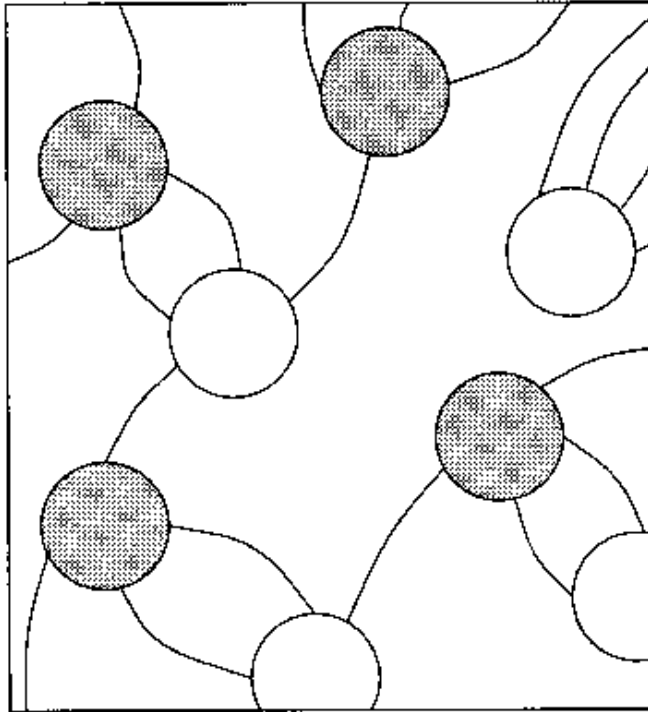
only 0^- and 0^+ $c\bar{c}$ observed !



BaBar Collaboration, **Phys.Rev. D72 (2005) 031101**

Phenomenology of instantons in QCD

T. Schafer, E. Shuryak



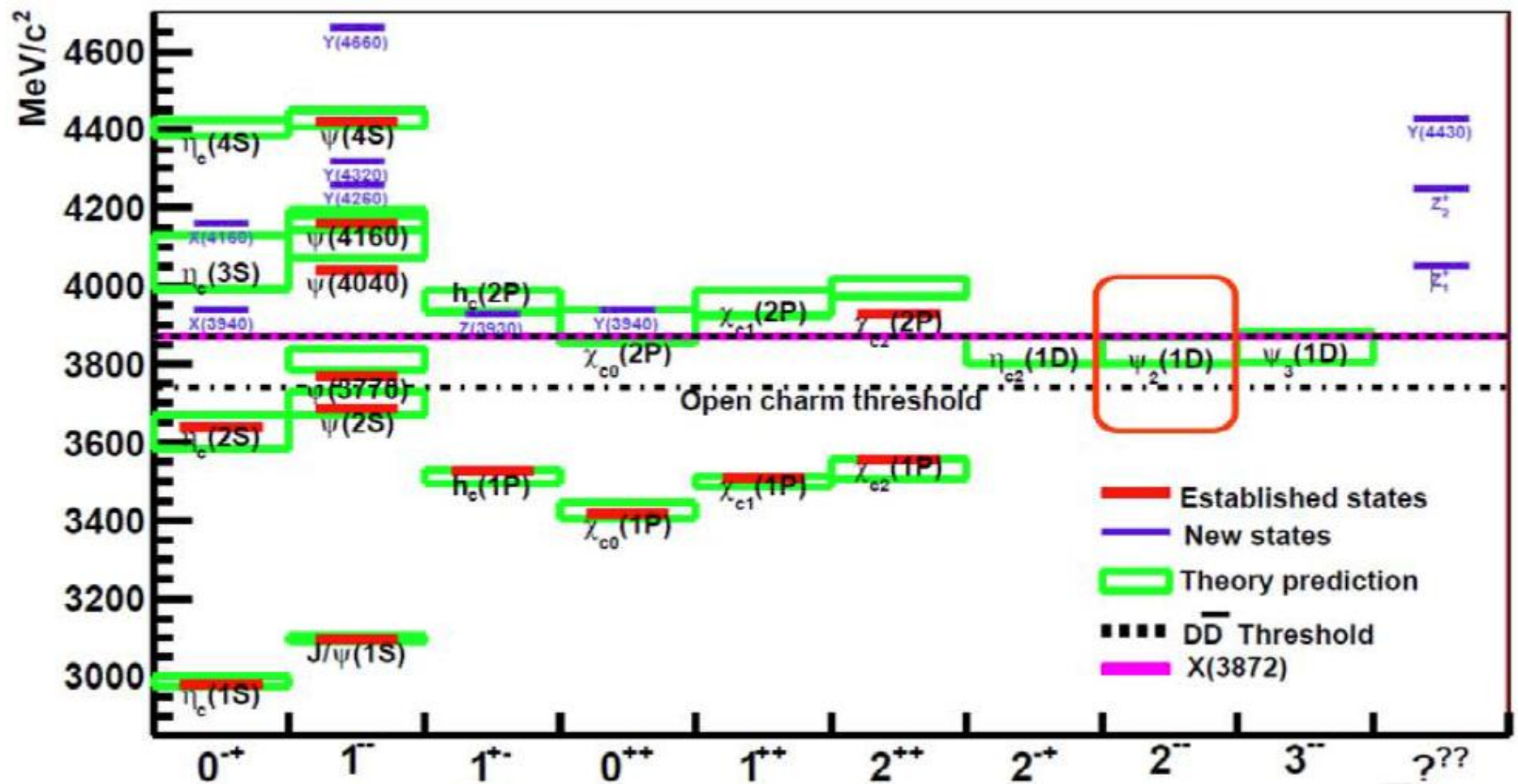
NJL model : **both 0^- & 0^+ important !**

$$\mathcal{L} = i\bar{\psi}\not{\partial}\psi + \frac{\lambda}{4} [(\bar{\psi}\psi)(\bar{\psi}\psi) - (\bar{\psi}\gamma^5\psi)(\bar{\psi}\gamma^5\psi)] = i\bar{\psi}_L\not{\partial}\psi_L + i\bar{\psi}_R\not{\partial}\psi_R + \lambda(\bar{\psi}_L\psi_R)(\bar{\psi}_R\psi_L).$$

Best playgrounds for unquenched quark models:

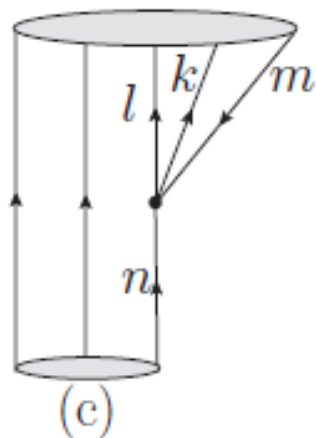
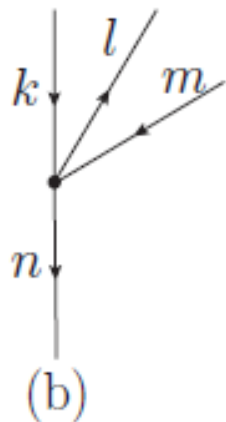
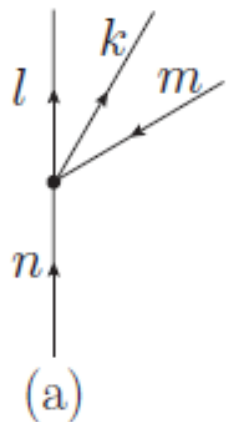
for baryon $sss \rightarrow sss \bar{q}q$

for meson $\bar{c}c \rightarrow \bar{c}c \bar{q}q$



for baryon $sss \rightarrow sss \bar{q}q$

$$H = \begin{pmatrix} H_3 & V_{\Omega_3 \leftrightarrow \Omega_5} \\ V_{\Omega_3 \leftrightarrow \Omega_5} & H_5 \end{pmatrix}$$



$$H_N = H_o + H_{hyp} + \sum_{i=1}^N m_i$$

$$H_o = \sum_{i=1}^N \frac{\vec{p}_i^2}{2m_i} + \sum_{i < j}^N V_{conf}(r_{ij})$$

$$H_{qq}^{NJL} = \sum_{i < j}^N \sum_{a=0}^8 \hat{g}_{ij} \lambda_i^a \lambda_j^a \left[1 + \frac{1}{4m_i m_j} \hat{\sigma}_i \cdot (\vec{p}_i' - \vec{p}_i) \hat{\sigma}_j \cdot (\vec{p}_j' - \vec{p}_j) \right]$$

from $\mathcal{L}_{NJL} = \frac{1}{2} g_s \sum_{a=0}^8 [(\bar{q} \lambda^a q)^2 + (\bar{q} i \lambda^a \gamma_5 q)^2]$

Predictions for the lowest Ω^* by various models:

$\Omega^*(\mathbf{x}/2^-)$ as sss ($L=1$) : ~ 2020 MeV

Chao, Isgur, Karl, PRD38(1981)155

$\Omega^*(1/2^-)$ as $\bar{K}\Xi$ bound state: ~ 1805 MeV

W.L.Wang, F.Huang, Z.Y.Zhang, F.Liu, JPG35 (2008) 085003

$\Omega^*(\mathbf{x}/2^-)$ as $\bar{u}uss$ ($L=0$) : ~ 1820 MeV

Yuan-An-Wei-Zou-Xu, PRC87(2013)025205

$\Omega^*(3/2^-)$ as $sss - \bar{u}uss$ mixture : ~ 1780 MeV
by instanton/NJL interaction

An-Metsch-Zou, PRC87(2013) 065207; An-Zou, PRC89 (2014) 055209

Very important to find the lowest Ω^* ($1/2^-$ or $3/2^-$)

$$\psi(2S) \rightarrow \bar{\Omega}\Omega \quad \text{BR} = (5 \pm 2) \times 10^{-5}$$

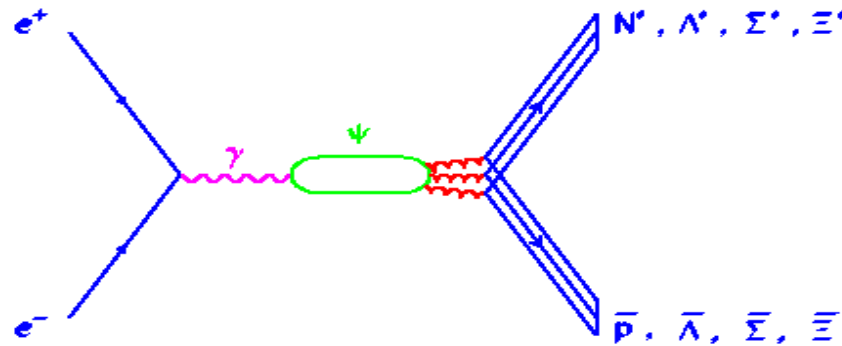
M. Ablikim et al. (BESII Coll.), CPC36(2012)1040

$$\psi(2S) \rightarrow \bar{\Omega}\Omega^* \quad \text{with } \Omega^* \rightarrow \gamma \Omega$$

\rightarrow excitation mechanism for sss states

$\bar{c}c$ decays -- a important new source for baryons

$$\psi \rightarrow \bar{B}BM \Rightarrow N^*, \Lambda^*, \Sigma^*, \Xi^*, \Omega^*$$



an ideal isospin and low spin filter from $\bar{c}c$ annihilation

No contamination from t/u-channel scattering as in πN and γN

high statistics extension to ψ', χ_{cJ}, η_c

3/7 new N^* from PDG92 to PDG14 are from BESII & BESIII

4. Conclusions and Prospects

- **New hyperons support unquenched quark picture**
 - new $\Sigma^*(1380)1/2^-$ replaces $\Sigma^*(1620)1/2^-$ **
 - new $\Lambda^*(1680)3/2^+$ replaces $\Lambda^*(1690)3/2^-$ ****
 - new $\Lambda^*(1670)3/2^-$ with width of 1.5 MeV [us]{ds} \bar{s}
 $\rightarrow \Lambda\eta$ Liu&Xie, PRC86(2012)055202
- **Both 1S_0 and 3P_0 are important for non-perturbative qq pair production from gluon field**
- **Distinguishable prediction for hyperon spectroscopy is yelling for experimental confirmation :**
Very important to find the lowest Ω^* ($1/2^-$ or $3/2^-$) at BES3 or super τ -charm

Many more interesting channels at super τ -charm :

$$\bar{\Omega} \Xi \bar{K}, \bar{\Xi} \Xi \pi, \bar{\Lambda} \Lambda \gamma, \bar{\Sigma} \Lambda \gamma, \bar{\Sigma} \Sigma \gamma, \bar{\Xi} \Xi \gamma, \dots$$

with $\Omega \rightarrow \Lambda K, \Xi \rightarrow \Lambda \pi$

S.Dulat, J.J.Wu, B.S.Zou, PRD83 (2011) 094032

“Proposal and theoretical formalism for studying baryon radiative decays from $J/\psi \rightarrow \bar{B}B^* + \bar{B}^*B \rightarrow \bar{B}B\gamma$ ”.

JLAB : $N^*, \Delta^* \rightarrow \gamma N$

Super τ -c: $\Lambda^* \rightarrow \gamma \Lambda, \gamma \Sigma ; \Sigma^* \rightarrow \gamma \Lambda, \gamma \Sigma ; \Xi^* \rightarrow \gamma \Xi; \Omega^* \rightarrow \gamma \Omega !$

Thanks !