

Constraining Natural SUSY at the HL-LHC, ILC and CEPC

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References:

- ▶ T. Li, S. Raza and K. Wang, Phys. Rev. D **93**, no. 5, 055040 (2016) [arXiv:1601.00178 [hep-ph]].

Outline

Introduction

The MSSM with GmSUGRA

Numerical Results

Conclusion

Outline

Introduction

The MSSM with GmSUGRA

Numerical Results

Conclusion

The Supersymmetric Standard Models:

- ▶ Solving the gauge hierarchy problem
- ▶ Gauge coupling unification
- ▶ Radiatively electroweak symmetry breaking
- ▶ Natural dark matter candidates
- ▶ Electroweak baryogenesis
- ▶ Electroweak precision: R parity

Supersymmetry

- ▶ The most promising new physics beyond the Standard Model.
- ▶ Gauge coupling unification in the SSMs strongly suggests the Grand Unified Theories (GUTs)
 $SU(5)$, and $SO(10)$, etc.
- ▶ The SUSY GUTs can be constructed from superstring theory

Heterotic $E_8 \times E_8$ string theory, D-brane models, free-fermionic string models, F-theory models, etc.

Supersymmetry is a bridge between the low energy phenomenology and high-energy fundamental physics.

Higgs boson mass in the MSSM:

- ▶ The SM-like Higgs boson mass is around 126 GeV.
- ▶ The tree-level Higgs boson mass is smaller than M_Z .
- ▶ The Higgs boson mass is enhanced by the top quarks/squarks loop corrections.
- ▶ The maximal stop mixing is needed to relax the fine-tuning.

The LHC Supersymmetry Search Constraints:

- ▶ The gluino and squark mass low bounds are around 1.7 TeV in the CMSSM/mSUGRA
- ▶ The gluino mass low bound is around 1.3 TeV.
- ▶ The stop/sbottom mass low bounds are around 600 GeV.

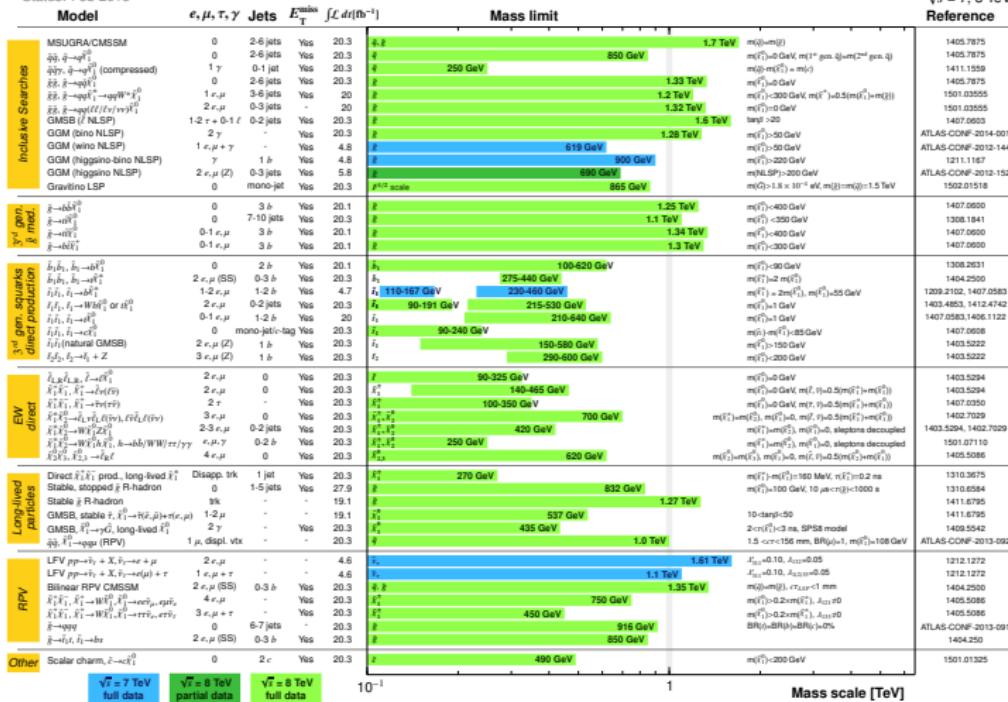
The SSMs are fine-tuned!

Introduction
The MSSM with GmSUGRA
Numerical Results
Conclusion

ATLAS SUSY Searches* - 95% CL Lower Limits

Status: Feb 2015

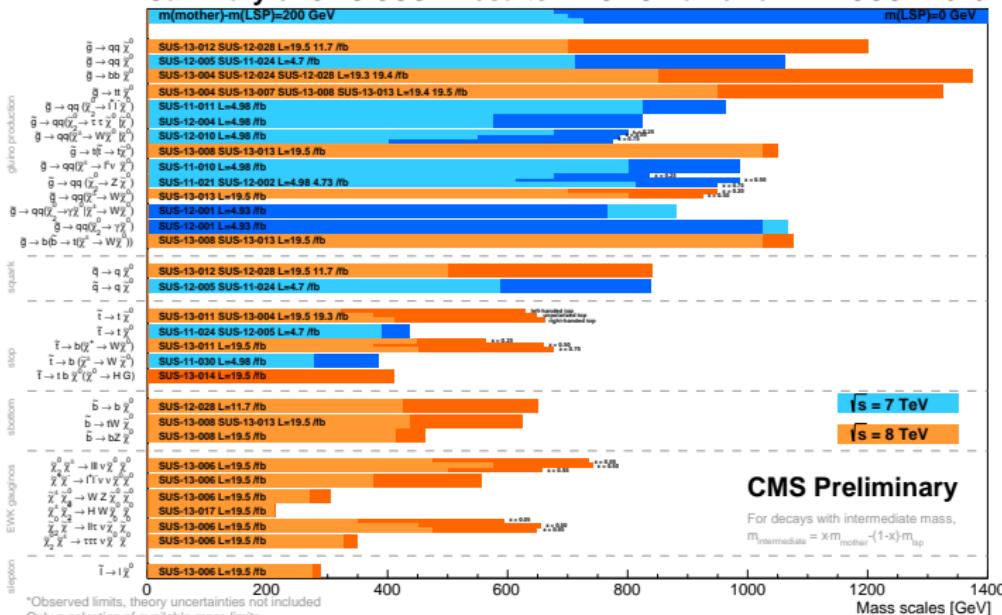
ATLAS Preliminary
 $\sqrt{s} = 7, 8 \text{ TeV}$
Reference



*Only a selection of the available mass limits on new states or phenomena is shown. All limits quoted are observed minus 1 σ theoretical signal cross section uncertainty.

Summary of CMS SUSY Results* in SMS framework

SUSY 2013



*Observed limits, theory uncertainties not included
 Only a selection of available mass limits
 Probe "up to" the quoted mass limit

Fine-Tuning Definition

- ▶ The minimization condition for electroweak symmetry breaking

$$\frac{M_Z^2}{2} = \frac{m_{H_d}^2 + \Sigma_d^d - (m_{H_u}^2 + \Sigma_u^u) \tan^2 \beta}{\tan^2 \beta - 1} - \mu^2 .$$

- ▶ The fine-tuning measure ¹

$$\Delta_{\text{EW}} \equiv \text{Max}\left\{\frac{2C_i}{M_Z^2}\right\} .$$

We consider $\Delta_{\text{EW}} < 100$.

¹H. Baer, V. Barger, P. Huang, D. Mickelson, A. Mustafayev and X. Tata, Phys. Rev. D **87**, no. 11, 115028 (2013) [arXiv:1212.2655 [hep-ph]].

Outline

Introduction

The MSSM with GmSUGRA

Numerical Results

Conclusion

Grand Unified Theories with Gravity Mediated Supersymmetry Breaking (mSUGRA/CMSSM):²

- The supersymmetry breaking soft terms: $M_{1/2}$, M_0 , A_0 , $\tan\beta$, $\text{sgn}(\mu)$.

$$\frac{1}{\alpha_3} = \frac{1}{\alpha_2} = \frac{1}{\alpha_1},$$

$$\frac{M_3}{\alpha_3} = \frac{M_2}{\alpha_2} = \frac{M_1}{\alpha_1}.$$

² A. H. Chamseddine, R. L. Arnowitt and P. Nath, Phys. Rev. Lett. **49**, 970 (1982); H. P. Nilles, Phys. Lett. B **115**, 193 (1982); L. E. Ibanez, Phys. Lett. B **118**, 73 (1982); R. Barbieri, S. Ferrara and C. A. Savoy, Phys. Lett. B **119**, 343 (1982); H. P. Nilles, M. Srednicki and D. Wyler, Phys. Lett. B **120**, 346 (1983); J. R. Ellis, D. V. Nanopoulos and K. Tamvakis, Phys. Lett. B **121**, 123 (1983); J. R. Ellis, J. S. Hagelin, D. V. Nanopoulos and K. Tamvakis, Phys. Lett. B **125**, 275 (1983); L. J. Hall, J. D. Lykken and S. Weinberg, Phys. Rev. D **27**, 2359 (1983).

Modification of Gauge Coupling Unification ³

- ▶ The SM gauge couplings need not be unified at the GUT Scale due to the high-dimensional operators

$$\mathcal{L} \supset \frac{c}{M_*} \text{Tr}(\Phi F_{\mu\nu} F^{\mu\nu}) .$$

$$\delta(1/g_3^2) : \delta(1/g_2^2) : \delta(1/g_1^2) = 2 : -3 : -1$$

- ▶ M_* can be the reduced Planck scale, string scale, or compactification scale.

³C. T. Hill, Phys. Lett. B **135**, 47 (1984); Q. Shafi and C. Wetterich, Phys. Rev. Lett. **52**, 875 (1984).
J. R. Ellis, C. Kounnas and D. V. Nanopoulos, Nucl. Phys. B **247**, 373 (1984); J. R. Ellis, K. Enqvist,
D. V. Nanopoulos and K. Tamvakis, Phys. Lett. B **155**, 381 (1985); M. Drees, Phys. Lett. B **158**, 409 (1985).

- ▶ In the GUTs with large number of fields, the renormalization effects significantly decrease the scale at which quantum gravity (or fundamental scale) becomes strong ⁴

$$M_* = \frac{M_{\text{Pl}}}{\Delta}, \quad \Delta = \sqrt{1 + \frac{N}{12\pi}}, \quad N = N_0 + N_{1/2} - 4N_1.$$

- ▶ This can be realized in the $SU(5)$ Models with $U(1)_Y$ flux ⁵ and $SO(10)$ Models with $U(1)_{B-L}$ flux ⁶.

⁴ X. Calmet, S. D. H. Hsu and D. Reeb, Phys. Rev. Lett. **101**, 171802 (2008).

⁵ R. Donagi and M. Wijnholt, Adv. Theor. Math. Phys. **15**, no. 6, 1523 (2011) [arXiv:0808.2223 [hep-th]];

R. Blumenhagen, Phys. Rev. Lett. **102**, 071601 (2009).

⁶ T. Li, Phys. Rev. D **81**, 065018 (2010) [arXiv:0905.4563 [hep-th]].

The Generalized mSUGRA (GmSUGRA):⁷

- ▶ Gauge coupling relation

$$\frac{1}{\alpha_2} - \frac{1}{\alpha_3} = k \left(\frac{1}{\alpha_1} - \frac{1}{\alpha_3} \right),$$

- ▶ Gaugino mass relation

$$\frac{M_2}{\alpha_2} - \frac{M_3}{\alpha_3} = k \left(\frac{M_1}{\alpha_1} - \frac{M_3}{\alpha_3} \right),$$

- ▶ Index k might be obtained from the LHC and ILC. We choose $k = 5/3$, and consider M_1 and M_2 as free parameters.

⁷ TL and D. V. Nanopoulos, Phys. Lett. B **692**, 121 (2010) [arXiv:1002.4183 [hep-ph]]; C. Balazs, TL, D. V. Nanopoulos and F. Wang, JHEP **1009**, 003 (2010) [arXiv:1006.5559 [hep-ph]].

The GmSUGRA

- ▶ Choosing slepton masses as input parameters, we can parametrize the squark masses as follows

$$\begin{aligned} m_{\tilde{Q}_i}^2 &= \frac{5}{6}(m_0^U)^2 + \frac{1}{6}m_{\tilde{E}_i^c}^2 , \\ m_{\tilde{U}_i^c}^2 &= \frac{5}{3}(m_0^U)^2 - \frac{2}{3}m_{\tilde{E}_i^c}^2 , \\ m_{\tilde{D}_i^c}^2 &= \frac{5}{3}(m_0^U)^2 - \frac{2}{3}m_{\tilde{L}_i}^2 . \end{aligned}$$

- ▶ The Higgs soft masses $m_{\tilde{H}_u}$ and $m_{\tilde{H}_d}$, and the trilinear soft terms A_U , A_D and A_E .

The Higgs couplings

- ▶ The deviations from the SM Higgs couplings

$$k_i \equiv g_{hii}^{\text{SUSY}} / g_{hii}^{\text{SM}} ,$$

where $i = W, Z, b, \tau, t, g, \gamma$.

- ▶ The Higgs couplings to gauge bosons

$$k_V \equiv \frac{g_{hVV}^{\text{SUSY}}}{g_{hVV}^{\text{SM}}} \sim 1 - \mathcal{O} \left(\frac{m_Z^4}{m_A^4} \frac{2}{\tan^2 \beta} \right) , \text{ for } V = W, Z.$$

$$k_g \approx 1 + \frac{m_t^2}{4} \left[\frac{1}{m_{\tilde{t}_1}^2} + \frac{1}{m_{\tilde{t}_2}^2} - \frac{X_t^2}{m_{\tilde{t}_1}^2 m_{\tilde{t}_2}^2} \right] ,$$

where $X_t = |A_t - \mu / \tan \beta|$ is the stop mixing parameter.

The Higgs couplings

- ▶ The Higgs couplings to the SM fermions

$$k_{b,\tau} \sim 1 + \mathcal{O}\left(2 \frac{m_Z^2}{m_A^2}\right) ,$$
$$k_t \sim 1 - \mathcal{O}\left(\frac{m_Z^2}{m_A^2} \frac{2}{\tan^2 \beta}\right) .$$

Thus, $\delta k_{b,\tau}$ and δk_t are predicted to have opposite sign.

The anomalous magnetic moment of the muon

- The discrepancy of $a_\mu = (g - 2)_\mu / 2$

$$\Delta a_\mu \equiv a_\mu(\text{exp}) - a_\mu(\text{SM}) = (28.7 \pm 8.0) \times 10^{-10} .$$

- The neutralino-smuon loop contribution

$$\Delta a_\mu^{\tilde{\chi}^0 \tilde{\mu}} \simeq \frac{1}{192\pi^2} \frac{m_\mu^2}{M_{SUSY}^2} \left(\text{sgn}(\mu M_1) g_1^2 - \text{sgn}(\mu M_2) g_2^2 \right) \tan \beta .$$

The anomalous magnetic moment of the muon

- ▶ The chargino-sneutrino loop contribution

$$\Delta a_\mu^{\tilde{\chi}^\pm \tilde{\nu}} \simeq \text{sgn}(\mu M_2) \frac{1}{32\pi^2} \frac{m_\mu^2}{M_{SUSY}^2} g_2^2 \tan \beta .$$

- ▶ $\Delta a_\mu \sim 10^{-9} \left(\frac{100 \text{ GeV}}{M_{SUSY}} \right)^2 \tan \beta$ for $\text{sgn}(\mu M_2) > 0$.
- ▶ The 2σ bound on Δa_μ can be achieved for $\tan \beta = 10$ if four relevant sparticles are lighter than 600 – 700 GeV. While for smaller $\tan \beta$ (~ 3), the lighter sparticles ($\lesssim 500$ GeV) are needed.

Outline

Introduction

The MSSM with GmSUGRA

Numerical Results

Conclusion

The parameter space for scan

$0 \text{ GeV} \leq m_0^U \leq 9000 \text{ GeV},$
 $100 \text{ GeV} \leq M_1 \leq 2000 \text{ GeV},$
 $100 \text{ GeV} \leq M_2 \leq 2100 \text{ GeV},$
 $100 \text{ GeV} \leq m_{\tilde{L}} \leq 1200 \text{ GeV},$
 $100 \text{ GeV} \leq m_{\tilde{E}^c} \leq 1200 \text{ GeV},$
 $100 \text{ GeV} \leq \mu \leq 1500 \text{ GeV},$
 $-16000 \text{ GeV} \leq A_U = A_D \leq 18000 \text{ GeV},$
 $-6000 \text{ GeV} \leq A_E \leq 6000 \text{ GeV},$
 $0 \text{ GeV} \leq m_A \leq 9500 \text{ GeV}, \quad 2 \leq \tan\beta \leq 60,$
 $\mu > 0, \quad m_t = 173.3 \text{ GeV}, \quad m_b^{\overline{\text{DR}}}(M_Z) = 2.83 \text{ GeV}.$

Basic constraints (I)

- The Radiative Electroweak Symmetry Breaking (REWSB).
- One of the neutralinos is the LSP.
- The Higgs boson mass

$$123 \text{ GeV} \leq m_h \leq 127 \text{ GeV} .$$

- The fine-tuning measure $\Delta_{\text{EW}} \leq 100$.

Basic constraints (I)

► The sparticle mass bounds

$$\begin{aligned} m_{\tilde{t}_1}, m_{\tilde{b}_1} &\geq 100 \text{ GeV}, \\ m_{\tilde{\tau}_1} &\geq 105 \text{ GeV}, \\ m_{\tilde{\chi}_1^\pm} &\geq 103 \text{ GeV}. \end{aligned}$$

$$\begin{aligned} 1.7 \text{ TeV} &\leq m_{\tilde{g}} \text{ (for } m_{\tilde{g}} \sim m_{\tilde{q}}) \quad , \\ 1.3 \text{ TeV} &\leq m_{\tilde{g}} \text{ (for } m_{\tilde{g}} \ll m_{\tilde{q}}) \quad , \\ 300 \text{ GeV} &\leq m_A \quad . \end{aligned}$$

Basic constraints (I)

► B-physics

$$1.6 \times 10^{-9} \leq \text{BR}(B_s \rightarrow \mu^+ \mu^-) \leq 4.2 \times 10^{-9} \quad (2\sigma) \quad ,$$

$$2.99 \times 10^{-4} \leq \text{BR}(b \rightarrow s\gamma) \leq 3.87 \times 10^{-4} \quad (2\sigma) \quad ,$$

$$0.70 \times 10^{-4} \leq \text{BR}(B_u \rightarrow \tau\nu_\tau) \leq 1.5 \times 10^{-4} \quad (2\sigma) \quad .$$

Muon anomalous magnetic moment constraint (II)

$$4.7 \times 10^{-10} \leq \Delta a_\mu \leq 52.7 \times 10^{-10} \text{ (} 3\sigma \text{)} .$$

Higgs coupling constraints (III)

Collider	HL-LHC	ILC	CEPC (2 IP)	FCC-ee (4 IP)
\sqrt{s} (GeV)	14000	500	240	240
\mathcal{L} (fb^{-1})	3000	500	5000	10000
(e^-, e^+)	-	(-0.8, +0.3)	(0, 0)	(0, 0)
k_g	9.1	2.3	1.5	1.1
k_W	5.1	1.2	1.2	0.85
k_Z	4.4	1.0	0.26	0.16
k_γ	4.9	8.4	4.7	1.7
k_b	12	1.7	1.3	0.88
k_τ	9.7	2.4	1.4	0.94
k_t	11	14	-	-

Table: Summary for the precisions of Higgs boson coupling measurements in percentage at different colliders.

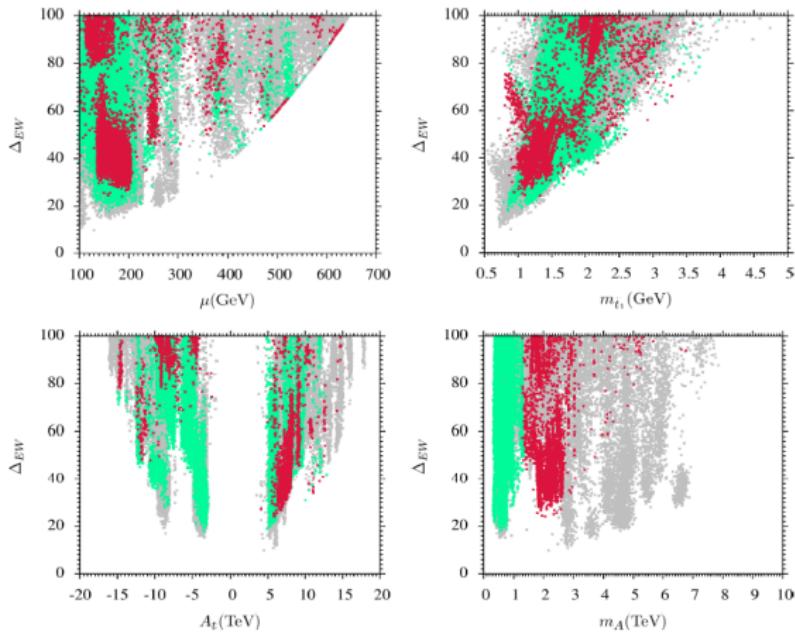


Figure: Δ_{EW} vs. μ , $m_{\tilde{t}_1}$, A_t and m_A : Grey points satisfy the basic constraints (I). Green points are a subset of grey points and satisfy the muon magnetic moment constraint (II). Red points are a subset of green points and satisfy the Higgs coupling constraint (III) from CEPC (2 IP) only.

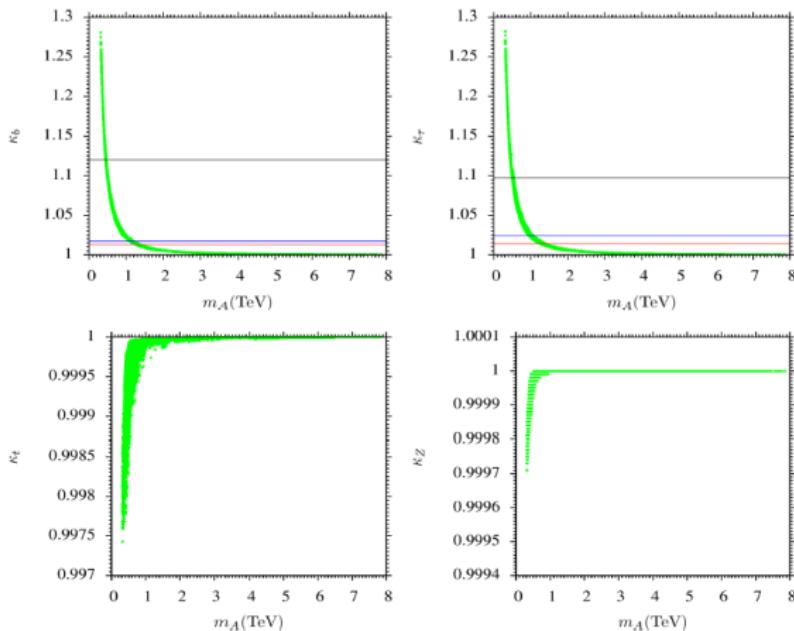


Figure: k_b , k_T , k_t and k_Z vs. m_A : Green points satisfy the basic constraints (I). The horizontal lines (if shown) in these plots label the precisions at different colliders: HL-LHC (black), ILC (blue) and CEPC (2 IP) (red) respectively.

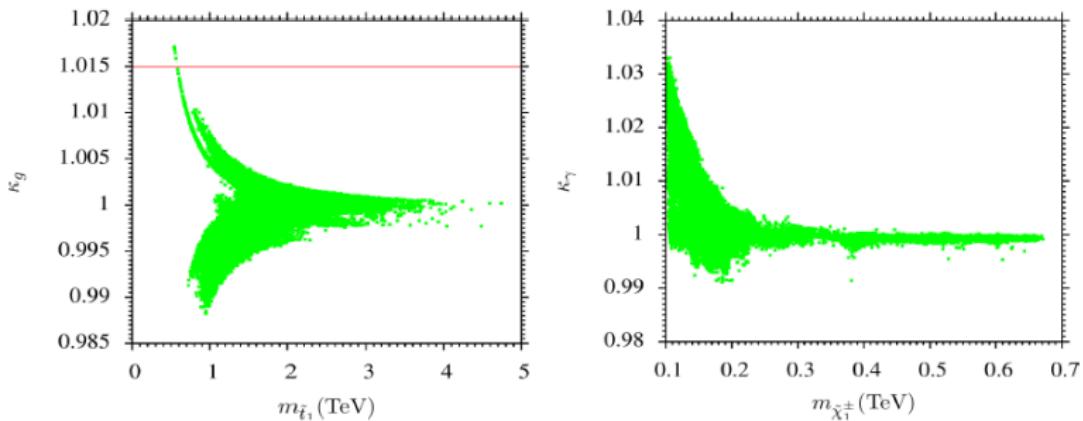


Figure: k_g vs. $m_{\tilde{t}_1}$ and k_γ vs. $m_{\tilde{\chi}_1^\pm}$.

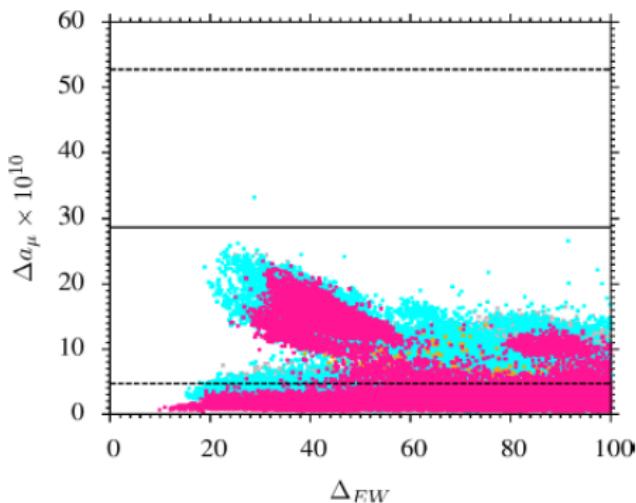


Figure: Δa_μ vs. Δ_{FW} : Grey points satisfy the basic constraints (I). Cyan, orange and red points are subsets of grey points and satisfy the Higgs coupling constraint (III) from HL-LHC, ILC, and CEPC (2 IP), respectively. The solid horizontal line displays the central value of Δa_μ from the muon magnetic moment experiment, while the dash lines show the 3σ values.

Introduction
The MSSM with GmSUGRA
Numerical Results
Conclusion

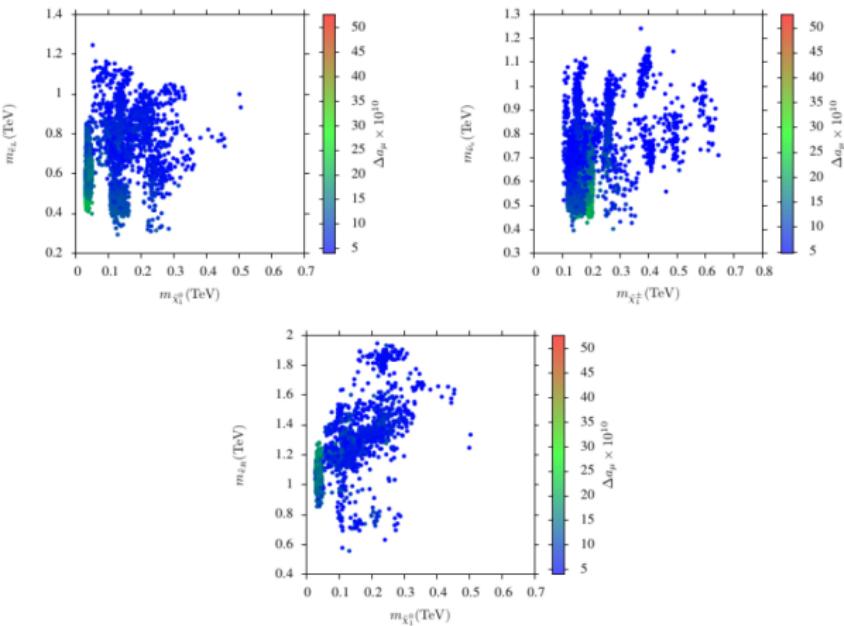


Figure: $m_{\tilde{e}_{R,L}}$ vs. $m_{\tilde{\chi}_1^0}$ and $m_{\tilde{\nu}_e}$ vs. $m_{\tilde{\chi}_1^\pm}$: Points satisfy both the basic constraints (I), the muon magnetic moment constraint (II) and the Higgs coupling constraint (III) from the CEPC (2 IP) only. Color represents the Δa_μ values.

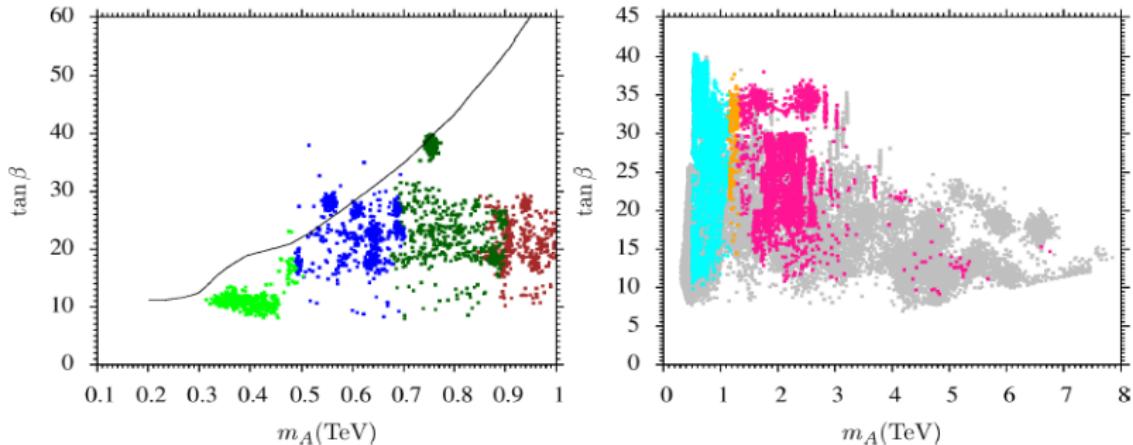
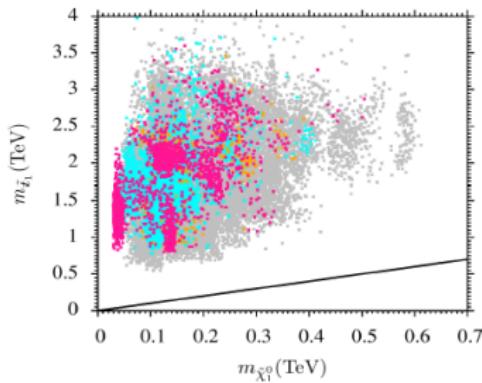
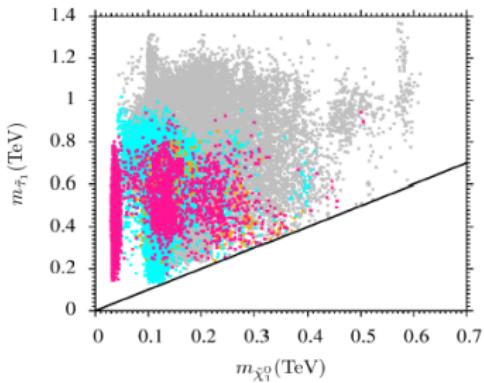
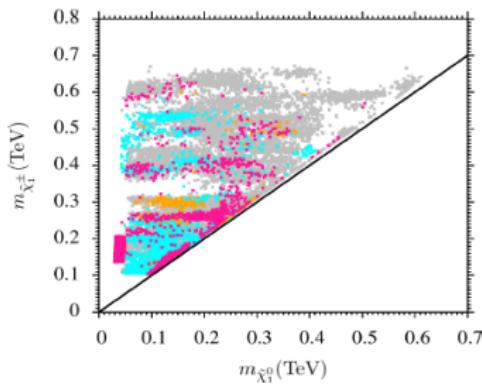
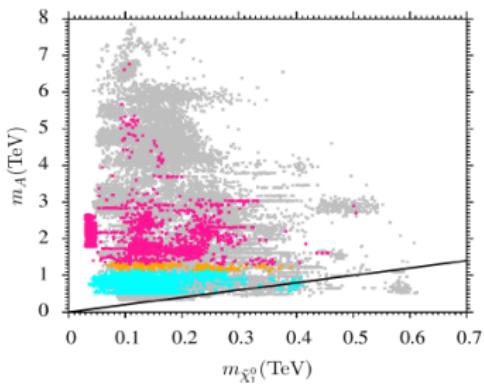


Figure: $\tan \beta$ vs. m_A : In the left plot, all points satisfy the basic constraints (I) and $m_A \leq 2\mu$; green, blue, dark green and brown points represent $k_b \geq 1.1$, $1.05 \leq k_b \leq 1.1$, $1.03 \leq k_b \leq 1.05$ and $k_b \leq 1.03$ respectively; black curve shows the $1 - \sigma$ bound from the ATLAS direct heavy Higgs search. In the right plot, grey points satisfy the basic constraints (I); cyan, orange and red points are subsets of grey points, and satisfy both the muon magnetic moment constraint (II) and the Higgs coupling constraint (III) from HL-LHC, ILC, and CEPC (2 IP), respectively.



Introduction
The MSSM with GmSUGRA
Numerical Results
Conclusion

	Point 1	Point 2	Point 3	Point 4
$m_0 v$	1037	2582	5057	5327
$m_{\tilde{Q}}$	1003.3	2366.4	4631	4879.2
$m_{\tilde{D}^c}$	1162.1	3306.7	6487	6830.6
$m_{\tilde{B}^c}$	1311.7	3326.3	6474.9	6841.3
$m_{\tilde{L}}$	328.1	264.8	1023	858.2
$m_{\tilde{\nu}^c}$	814.1	514.8	900.9	978.6
M_1	857.1	350.9	602.7	686.5
M_2	656	993	634.3	828.8
M_3	1158.8	-612.25	555.3	473.05
$A_t = A_b$	-3292	5390	-9095	-9684
A_f	-691.6	1192	-1570	-2937
$\tan \beta$	11.6	33.8	19.7	21.4
μ	183	172.5	377.5	168.2
m_A	338.9	754.9	603	846
Δ_{HS}	1883	3151	10873	11903.67
Δ_{EW}	27	23	87	87
Δa_μ	4.731×10^{-10}	16.63×10^{-10}	7.683×10^{-10}	12.161×10^{-10}
m_h	123	123	125	125
m_H	342	751	607	856
m_{H^\pm}	348	750	608	850
κ_b, κ_t	1.22541, 0.99817	1.04205, 0.99996	1.06459, 0.99983	1.03296, 0.99992
$\kappa_\tau, \kappa_W = \kappa_Z$	1.22756, 0.99981	1.04787, 0.99999	1.06644, 0.99999	1.03454, 0.99999
κ_g, κ_γ	0.99083, 1.00369	1.00473, 0.99763	0.99954, 1.00008	1.00038, 0.99714
$m_{\tilde{\chi}_{1,2}^0}$	171, 191	131, 183	260, 376	161, 178
$m_{\tilde{\chi}_{3,4}^0}$	368, 541	195, 834	390, 567	312, 717
$m_{\tilde{\chi}_{1,2}^\pm}$	185, 534	179, 823	378, 555	175, 702
$m_{\tilde{g}}$	2548	1539	1498	1317
$m_{\tilde{u}_{L,R}}$	2419, 2558	2705, 3476	4721, 6513	4942, 6832
$m_{\tilde{u}_{1,2}}$	1036, 1798	1151, 1762	2311, 3401	2326, 3543
$m_{\tilde{d}_{L,R}}$	2421, 2532	2706, 3547	4722, 6595	4943, 6938
$m_{\tilde{b}_{1,2}}$	1771, 2468	1212, 3089	2380, 6340	2401, 6630
$m_{\tilde{e}_{1,2}}$	726	521	698	507
$m_{\tilde{e}_3}$	718	343	645	308
$m_{\tilde{\tau}_{1,2}}$	737, 525	514, 787	651, 1409	398, 1487
$m_{\tilde{\tau}_{1,2}}$	518, 728	353, 568	627, 1338	185, 1333
$\Omega_{CDM} h^2$	0.0014	0.0846	0.1017	0.0099

Outline

Introduction

The MSSM with GmSUGRA

Numerical Results

Conclusion

The Natural Supersymmetry from GmSUGRA at the CEPC

- ▶ $\Delta_{EW} \sim 30$, consistence with all constraints, and having supersymmetric contributions to the muon anomalous magnetic moment within 1σ can be achieved.
- ▶ The precision of k_b and k_T measurements at CEPC can bound m_A to be above 1.2 TeV and 1.1 TeV respectively.
- ▶ The combination of the Higgs coupling measurement and muon anomalous magnetic moment measurement constrain \tilde{e}_R mass to be in the range from 0.6 TeV to 2 TeV. The range of both \tilde{e}_L and $\tilde{\nu}_e$ masses is $0.4 \text{ TeV} \sim 1.2 \text{ TeV}$. And the $\tilde{\chi}_1^0$ mass needs to be small (mostly $\leq 400 \text{ GeV}$).
- ▶ The comparison of bounds in the $\tan \beta - m_A$ plane shows that the Higgs coupling measurement is complementary to the direct collider searches for heavy Higgs.

Thank You Very Much
for Your Attention!