

A Roadmap to Reconstruct Higgs Potentials @ colliders



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CEPC-SPPC Symposium
IHEP/CAS, April 8-9, 2016, Beijing

Outline

- Introduction
- Higgs Potentials of a few Models
- Higgs Potentials at future colliders
- Prospects

Based on:

Q.Li, Z. Li, QY, X.R. Zhao, PRD92(2015)1,014015, arXiv:1503.07611

C.Y. Chen, QY, X.R. Zhao, Z.J. Zhao, Y.M. Zhong, PRD93 (2016)1,
013007, arXiv:1510.04013

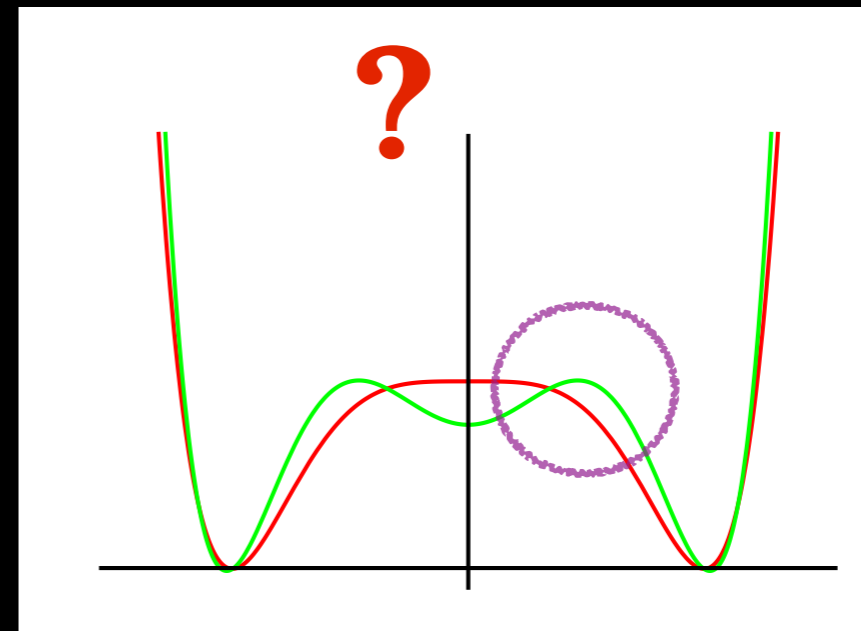
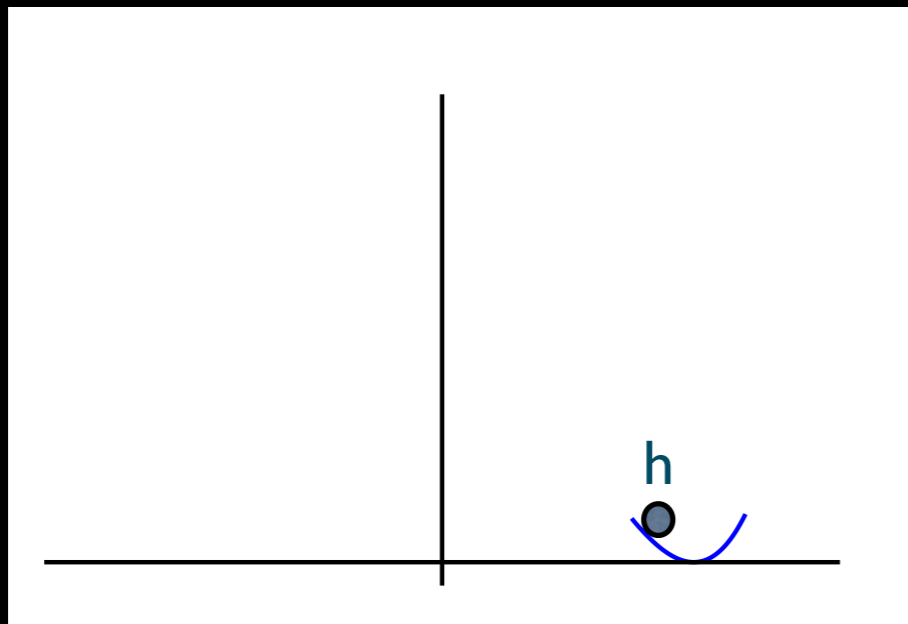
X.R. Zhao, Q. Li, Z. Li, QY, in preparation

Introduction

Why Higgs Potential?

- 1) It determines the EWSB mechanisms, mass generation of Higgs bosons, Higgs self-couplings
- 2) It relate to the EW phase transition in the early universe and might leave a finger print in the primordial gravitational wave
- 3) It could relate to the EW baryogenesis scenario, spontaneous CP breaking, spontaneous parity breaking
- 4) It directly connects with new physics (new particles, dark matter,)

Can it make difference?



A big question: the shape of Higgs potential is crucial for the strong first order electroweak phase transition, which is needed in the EW Baryogenesis scenarios.

CEPC Pre-CDR report

M. Trodden, Rev.Mod.Phys.71(1999)1463

D.E. Morrissey, M.J. Musolf, NJP14(2012)125003

N.Arkani-Hamed, T.Han, M.Mangano, L.T.Wang, 1511.06495

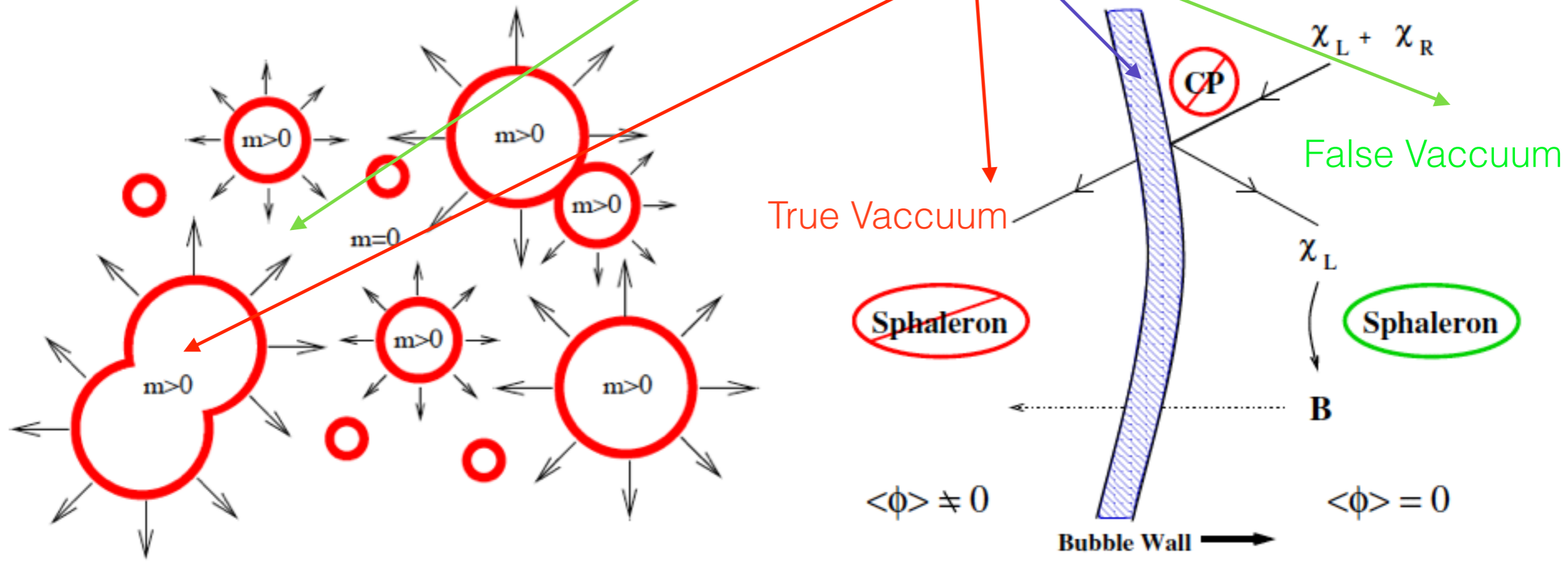
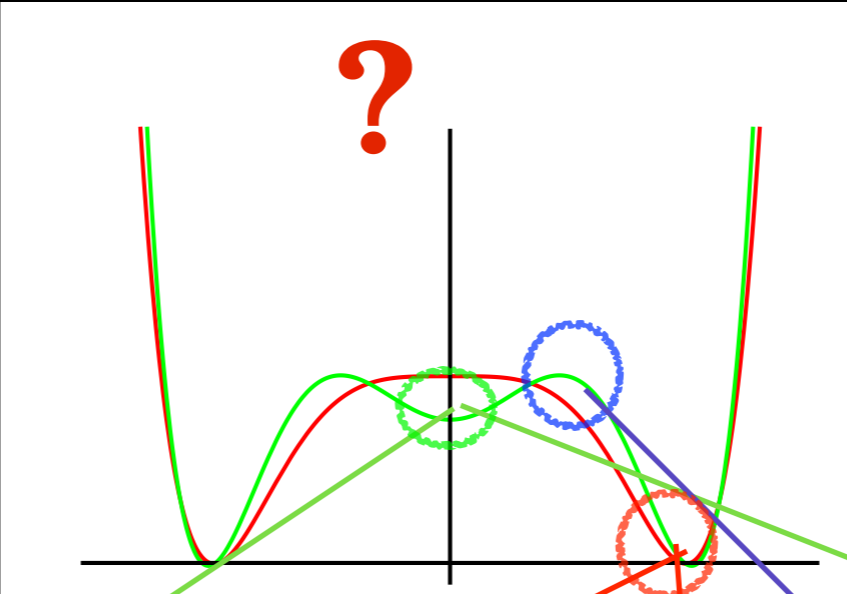


Fig. 11. Bubble nucleation during a first-order EWPT.

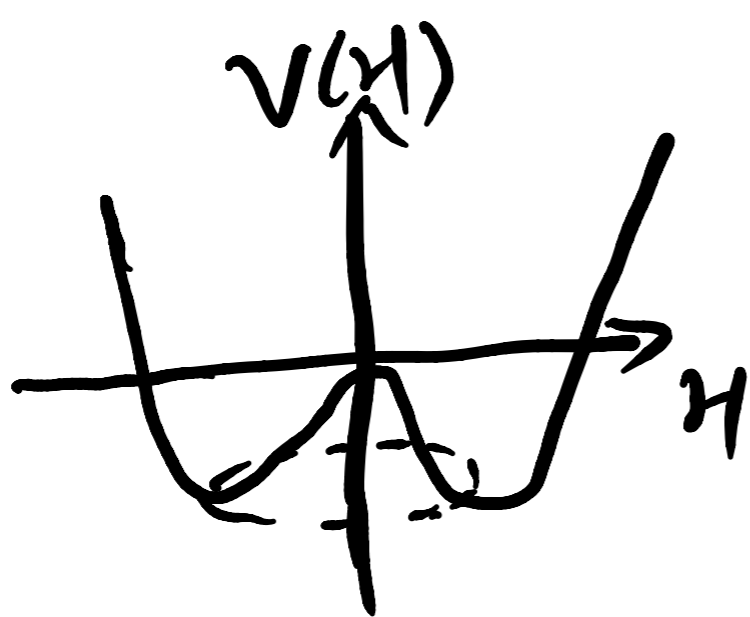
The bubble wall needs the little bump in potential shape

J.M. Cline, hep-ph/0609145
 D.E. Morrissey, M.J. Musolf, NJP14(2012)125003

2. Higgs Potentials of a few models

① Higgs Potential in the SM

$$V(H) = -m_H^2 H^\dagger H + \frac{\lambda}{4} (H^\dagger H)^2$$



$$\langle H \rangle = \begin{pmatrix} 0 \\ v/\sqrt{2} \end{pmatrix}$$

$$H = \begin{pmatrix} G^- \\ \frac{v+h+iG^0}{\sqrt{2}} \end{pmatrix}$$

$$V(h) = \frac{1}{2} m_h^2 h^2 + \frac{\lambda}{4} v h^3 + \frac{\lambda}{16} h^4$$

$$m_h^2 = \frac{\lambda}{4} v^2$$

$$*) (D_\mu H)^\dagger (D^\mu H) \longrightarrow \frac{1}{2} \partial_\mu h \partial^\mu h$$

$$+ (v+h)^2 \left[\frac{g^2}{4} w^+ w^- + \frac{g_z^2}{8} z z \right]$$

$$*) \mathcal{L}_{\text{Yukawa}} = Y_{t\bar{t}h} \bar{t} t h + Y_{b\bar{b}h} \bar{b} b h \\ + Y_{c\bar{c}h} \bar{c} c h + \dots$$

Effective theories:

$$\mathcal{L}_Y^{\text{eff}} = Y_t \bar{t} (a + b i \gamma_5) t h +$$

$$g_{tthh} \frac{1}{v} \bar{t} t h^2 + \dots$$

$$\begin{aligned} \mathcal{L}_V^{\text{eff}} = & d_1 v h \left[\frac{g^2}{2} w^\dagger w + \frac{g_Z^2}{4} z \cdot z \right] \\ & + d_2 h^2 \left[\frac{g^2}{4} w^\dagger w + \frac{g_Z^2}{8} z \cdot z \right] \\ & + \frac{h}{v} \left(g_{hgg} G_{\mu\nu}^2 + g_{hZZ} Z_{\mu\nu}^2 + g_{hrr} A_{\mu\nu}^2 \right) \\ & + \hat{g}_{hgg} \hat{G}_{\mu\nu}^2 + \hat{g}_{hZZ} \hat{Z}_{\mu\nu}^2 + \hat{g}_{hrr} \hat{A}_{\mu\nu}^2 \end{aligned}$$

$$\begin{aligned}
& + g_{\text{non}} W_{\mu\nu} W^{\mu\nu} + \tilde{g}_{\text{non}} W_{\mu\nu} \tilde{W}^{\mu\nu} \\
& + g_{\text{HZA}} Z^{\mu\nu} A_{\mu\nu} + \tilde{g}_{\text{HZA}} Z^{\mu\nu} \tilde{A}^{\mu\nu} \\
& + \frac{h^2}{v^2} (g_{\text{HGG}} G_{\mu\nu}^2 + \dots)
\end{aligned}$$

\mathcal{L}_{eff}

$$= -\lambda_3 \lambda_{\text{SU}}^3 h^3 - \lambda_4 \lambda_{\text{SU}}^4 h^4$$

$$+ \frac{\lambda}{v} \partial h \partial h h + \frac{\lambda_5}{v} h^5 + \dots$$

$$\Delta\mathcal{L}_S = -\frac{1}{2}m_S^2 S^2 - \frac{1}{4}\lambda_S S^4 - \frac{1}{4}\lambda_{hSS} H^\dagger H S^2 ,$$

$$\Delta\mathcal{L}_V = \frac{1}{2}m_V^2 V_\mu V^\mu + \frac{1}{4}\lambda_V (V_\mu V^\mu)^2 + \frac{1}{4}\lambda_{hVV} H^\dagger H V_\mu V^\mu ,$$

$$\Delta\mathcal{L}_f = -\frac{1}{2}m_f \bar{f} f - \frac{1}{4} \frac{\lambda_{hff}}{\Lambda} H^\dagger H \bar{f} f .$$

J. Baglio, A. Djouadi, J. Quevillon, 1511.07853

② Higgs potential in the Singlet + SM

$$V = V(H) + V(S) + V_X(S, H) + V_{ph}$$

$$V(H) = -m_H^2 H^\dagger H + \frac{\lambda}{4} (H^\dagger H)^2$$

$$V(S) = -m_S^2 S^\dagger S + \frac{\lambda_S}{4} (S^\dagger S)^2$$

$$V_X(S, H) = \lambda_X (S^\dagger S) (H^\dagger H)$$

$$V_{ph} = 2(\mu^2 + \lambda_1 S^\dagger S + \lambda_2 H^\dagger H) \cdot \text{Re} S^2 \\ + 2\lambda_3 \text{Re}(S^4)$$

$$\mathcal{L} = \mathcal{M} \bar{D}_2 D_2 + (f_j S + f'_j S^*) \bar{D}_2 d_k^j + \text{h.c.}$$

$$\langle H \rangle = \begin{pmatrix} 0 \\ \frac{v}{\sqrt{2}} \end{pmatrix}$$

$$\langle S \rangle = \frac{v_s}{\sqrt{2}} e^{i\alpha} \uparrow$$

$$V_{ph} = a \cos 2\alpha + b \cos 4\alpha$$

Spontaneous CP violation!

$$H = \begin{pmatrix} G^- \\ \frac{v+h+iG^0}{\sqrt{2}} \end{pmatrix}$$

$$S = \left[\frac{(v_R + s_R) + (v_I + s_I)}{\sqrt{2}} \right]$$

$$\mathcal{L}_2 = (h \quad s_R \quad s_I) \mathcal{M} \begin{pmatrix} h \\ s_R \\ s_I \end{pmatrix}$$

$$\mathcal{M} = \begin{pmatrix} \frac{1}{2}v^2 & (2\lambda_2 + \lambda_k) v v_R & (-2\lambda_2 + \lambda_k) v v_I \\ \cdot & \frac{1}{2}(\lambda_1 + \lambda_3 + \lambda_5) v^2 & \frac{1}{2}(-2\lambda_3 + \lambda_5) v v_R \\ \cdot & \cdot & \frac{1}{2}(-\lambda_1 + \lambda_3 + \lambda_5) v^2 \end{pmatrix}$$

Triple/quartic Higgs Coupling in the model.

$$\mathcal{L} = C_{klj} S_k S_l S_j + C_{kljmn} S_k S_l S_j S_m$$

$$C_{klj} = \frac{T_{1e} m_e^2 T_{lk} T_{lj}}{2v} + \frac{T_{2e} m_e^2 T_{2k} T_{2j}}{2v_R} + \frac{T_{3e} m_e^2 T_{3k} T_{3j}}{2v_I}$$

$$C_{kljmn} = \frac{1}{8} \sum_i \left[\frac{m_i^2}{v^2} T_{ii} T_{ii} T_{lk} T_{lc} T_{lj} T_{lm} + \frac{m_i^2}{v_R^2} T_{2i} T_{2i} T_{2k} T_{2l} \right.$$

$$T_{2j} T_{2m} + \frac{m_i^2}{v_I^2} T_{3i} T_{3i} T_{3k} T_{3l} T_{3j} T_{3m} +$$

$$\frac{m_i^2}{v v_R} T_{ii} T_{2i} (T_{lk} T_{lc} T_{2j} T_{2m} + T_{lj} T_{lm} T_{2k} T_{2l}) +$$

$$\frac{m_i^2}{v v_I} T_{ii} T_{3i} (T_{lk} T_{ll} T_{3j} T_{3m} + T_{lj} T_{lm} T_{3k} T_{3l}) +$$

$$\frac{m_i^2}{v v_R} T_{2i} T_{3i} (T_{2k} T_{2l} T_{3j} T_{3m} + T_{2j} T_{2m} T_{3k} T_{3l}) \left. \right]$$

$$\begin{array}{ccc}
 \left(\begin{array}{l} m_H^2, m_S^2, \mu^2 \\ \lambda, \lambda_5, \lambda_x \\ \lambda_1, \lambda_2, \lambda_3 \end{array} \right) & \longrightarrow & \left(\begin{array}{l} m_1^2, m_2^2, m_3^2 \\ \underline{\alpha, \beta, \gamma} \\ \nu, \nu_R, \nu_I \end{array} \right) \longrightarrow \underline{I} \\
 9 & & 7+2
 \end{array}$$

A simple roadmap for this model:

- 1) To discover extra neutral Higgs bosons
- 2) To measure three mixing angles
- 3) 2 must be extracted from T/Q coupling measurement

③ Higgs potential of the 2HDM

$$\begin{aligned} V = & m_{11}^2 \phi_1^\dagger \phi_1 + m_{22}^2 \phi_2^\dagger \phi_2 - (m_{12}^2 \phi_1^\dagger \phi_2 + \text{h.c.}) \\ & + \frac{\lambda_1}{2} (\phi_1^\dagger \phi_1)^2 + \frac{\lambda_2}{2} (\phi_2^\dagger \phi_2)^2 + \lambda_3 (\phi_1^\dagger \phi_1) (\phi_2^\dagger \phi_2) \\ & + \lambda_4 (\phi_1^\dagger \phi_2) (\phi_2^\dagger \phi_1) \\ & + \left[\frac{\lambda_5}{2} (\phi_1^\dagger \phi_2) (\phi_1^\dagger \phi_2) + \lambda_6 (\phi_1^\dagger \phi_2) (\phi_1^\dagger \phi_1) \right. \\ & \left. + \lambda_7 (\phi_1^\dagger \phi_2) (\phi_2^\dagger \phi_2) + \text{h.c.} \right] \end{aligned}$$

Higgs Basis

Triple/Quartic Couplings in 2HDM

$$\begin{aligned} \mathcal{L} &= v H^\dagger H^- C_i S_i \\ &+ C_{ijk} S_i S_j S_k \\ &+ \frac{\lambda_2}{2} H^\dagger H^- H^\dagger H^- \\ &+ C_{kl} H^\dagger H^- S_k S_l \\ &+ C_{kl, jmn} S_k S_l S_j S_m \end{aligned}$$

$$G_j = T_{1j} \lambda_3 + T_{2j} \operatorname{Re} \lambda_7 - T_{3j} \operatorname{Im} \lambda_7$$

$$C_{jkl} = \frac{v}{2} G_j (\delta_{kl} - T_{lk} T_{il})$$

$$+ \frac{1}{v} (m_4^2 - \frac{m_j^2}{2}) T_{1j} T_{lk} T_{il}$$

$$+ \frac{1}{v} (m_k^2 - m_4^2) T_{1j} \delta_{kl}$$

$$C_{kl} = \frac{\lambda_2}{2} \delta_{kl} - \frac{\lambda_2 t \lambda_3}{2} T_{lk} T_{il} + C_e T_{lk}$$

$$C_{kljm} = \frac{m_i^2 T_{ii} T_{ii}}{8v^2} T_{lk} T_{il} T_{lj} T_{im} + \left[\frac{\lambda_2}{8} (\delta_{kl} - T_{lk} T_{il}) \right.$$

$$\left. - \frac{\lambda_3}{4} T_{lk} T_{il} + \frac{C_e}{2} T_{il} + \frac{m_j^2 - m_4^2}{2v^2} T_{lk} T_{il} \right] \cdot ($$

$$\delta_{jm} - T_{lj} T_{im})$$

$$\left(\begin{array}{l} m_{11}^2, m_{22}^2, \underline{m_{12}^2}, \\ \lambda_1, \lambda_2, \lambda_3, \lambda_4, \\ \underline{\lambda_5, \lambda_6, \lambda_7} \end{array} \right) \rightarrow \left(\begin{array}{l} m_4^2, m_1^2, m_2^2, m_3^2 \\ \nu, \alpha, \beta, \gamma \\ \lambda_2, \lambda_3, \text{Re}\lambda_4, \\ \text{Im}\lambda_7 \end{array} \right)$$

$$14 - 3$$

$U(2)$ \uparrow

11
3 must be extracted
from T/Q coupling
measurements

A simple roadmap for 2HDM:

- 1) To discover 3 extra neutral HB, 2 charged HB
- 2) To measure three mixing angles
- 3) 3 must be extracted from T/Q couplings

④ n -Higgs Doublet model.

$$\begin{aligned} \mathcal{L} = & m_{ii}^2 \phi_i^\dagger \phi_i + [m_{ij}^2 \phi_i^\dagger \phi_j + \text{h.c.}] \\ & + a_{ij} (\phi_i^\dagger \phi_i) (\phi_j^\dagger \phi_j) \\ & + b_{ij} (\phi_i^\dagger \phi_j) (\phi_j^\dagger \phi_i) \\ & + [C'_{ijk} \phi_i^\dagger \phi_j \phi_k^\dagger \phi_k + C_{ij}^2 \phi_i^\dagger \phi_j \phi_i^\dagger \phi_j \\ & + d_{ijkl} \phi_i^\dagger \phi_j \phi_k^\dagger \phi_l + \text{h.c.}] \end{aligned}$$

$$\frac{n^4}{2} + \frac{3n^2}{2} - 2n + 1$$

No. of CP phases: $\frac{1}{4}n^2(n^2-1) - (n-1)$

No. of charged Higgs bosons:

$n-1$, T matrix $\in SO(n-1)$

No. of Neutral Higgs bosons:

$2n-1$, T matrix $\in SO(2n-1)$

No. of parameters must be probed via triple/quartic couplings:

$$N_{t/q} = \frac{n^4}{2} - n^2 - \frac{n}{2} + 1$$

$$n=3, N_{t/q} = 31$$

$$n=4, N_{t/q} = 111$$

⑤ Higgs potential of the left-right model

$$V = V(\Delta_L, \Delta_R) + V(\phi_1, \phi_2) + V(\Delta, \phi)$$

$$V(\Delta_L, \Delta_R) = -\mu^2 \text{Tr}(\Delta_L^\dagger \Delta_L + \Delta_R^\dagger \Delta_R) \\ + \beta_1 [\text{Tr}(\Delta_L^\dagger \Delta_L)^2 + \text{Tr}(\Delta_R^\dagger \Delta_R)^2]$$

$$+ \beta_2 [\text{Tr}(\Delta_L^\dagger \Delta_L \Delta_L^\dagger \Delta_L) + \text{Tr}(\Delta_R^\dagger \Delta_R \Delta_R^\dagger \Delta_R)]$$

$$+ \beta_3 [\text{Tr}(\Delta_L^\dagger \Delta_L) \text{Tr}(\Delta_R^\dagger \Delta_R)]$$

$$+ \beta_4 [\text{Tr}(\Delta_L^\dagger \Delta) \text{Tr}(\Delta_L \Delta_L) + \\ \text{Tr}(\Delta_R^\dagger \Delta_R) \text{Tr}(\Delta_R \Delta_R)]$$

$$\Delta_{LR} = \begin{pmatrix} \delta^+/\sqrt{2} & \delta^{++} \\ \delta^0 & -\delta^+/\sqrt{2} \end{pmatrix}_{LR}$$

5 free parameters

$$\begin{aligned}
V(\phi_1, \phi_2) = & -\mu_{ij}^2 \text{Tr}(\phi_i^\dagger \phi_j) \\
& + \lambda_{ijkl} \text{Tr}(\phi_i^\dagger \phi_j) \text{Tr}(\phi_k^\dagger \phi_l) \\
& + \lambda'_{ijkl} \text{Tr}(\phi_i^\dagger \phi_j \phi_k^\dagger \phi_l) \quad (14)
\end{aligned}$$

$$\phi_1 = \begin{pmatrix} \phi_1^0 & \phi_1^+ \\ \phi_1^- & \phi_1^0 \end{pmatrix} \quad \phi_2 = \tau_2 \phi_2^* \tau_2$$

$$\begin{aligned}
V(\Delta, \phi) = & \alpha_{ij} \text{Tr}(\phi_i^\dagger \phi_j) \text{Tr}(\Delta_L^\dagger \Delta_L + \Delta_R^\dagger \Delta_R) \\
& + \beta_{ij} \text{Tr}(\phi_i^\dagger \phi_j \Delta_L^\dagger \Delta_L + \phi_i \phi_j^\dagger \Delta_R^\dagger \Delta_R) \\
& + (\gamma_{ij} \text{Tr}(\Delta_L^\dagger \phi_i \Delta_R \phi_j^\dagger) + \text{h.c.}) \quad (12)
\end{aligned}$$

$$\langle \Delta_{LR} \rangle = \begin{pmatrix} 0 & 0 \\ v_{LR} & 0 \end{pmatrix}$$

$$\langle \phi \rangle = \begin{pmatrix} k & 0 \\ 0 & k' \end{pmatrix} e^{i\alpha}$$

$$v_L = \left(\frac{\gamma_{12}}{2(\rho_1 + \rho_2) - \rho_3} \right) \frac{k^2}{v_R}$$

Rich Higgs plane:

$$2: \quad \delta_1^{++}, \delta_2^{++} \leftarrow \delta_L^{++}, \delta_R^{++}$$

$$2: \quad \frac{(\delta_1^+ \delta_2^+)}{(G_1^+ G_2^+)} \leftarrow \delta_L^+, \delta_R^+, \phi_1^+, \phi_2^+$$

$$6: \quad \delta_1^0 \dots \delta_6^0 \leftarrow \delta_L^0, \delta_R^0, \phi_1^0, \phi_2^0$$
$$(G_1^0, G_2^0)$$

No. of physical parameters

$$(2+1) + (2+1) + (6 + \frac{6 \times 5}{2}) = 27$$

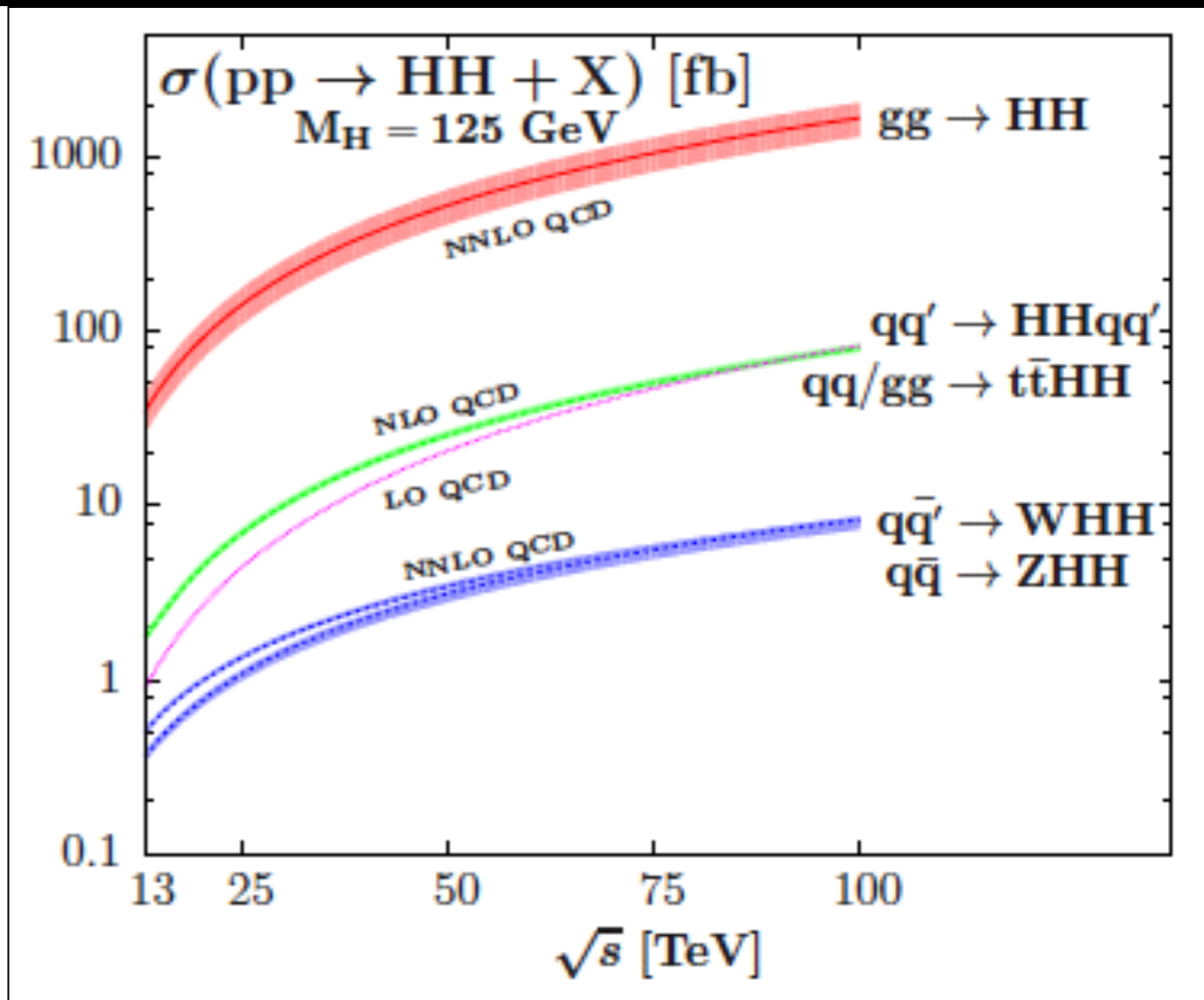
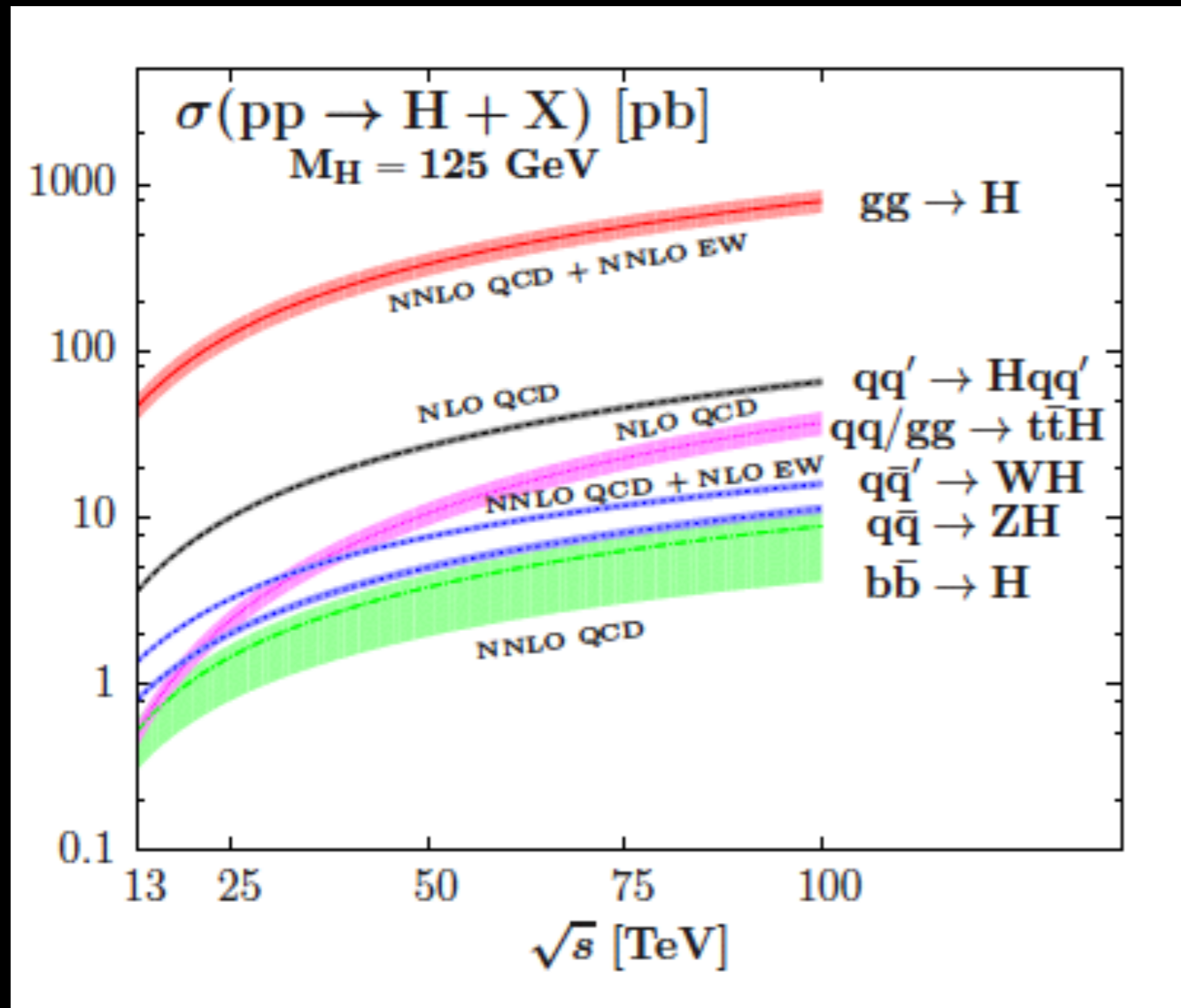
No. of independent parameters in T/Q couplings

$$31 - 27 = 4$$

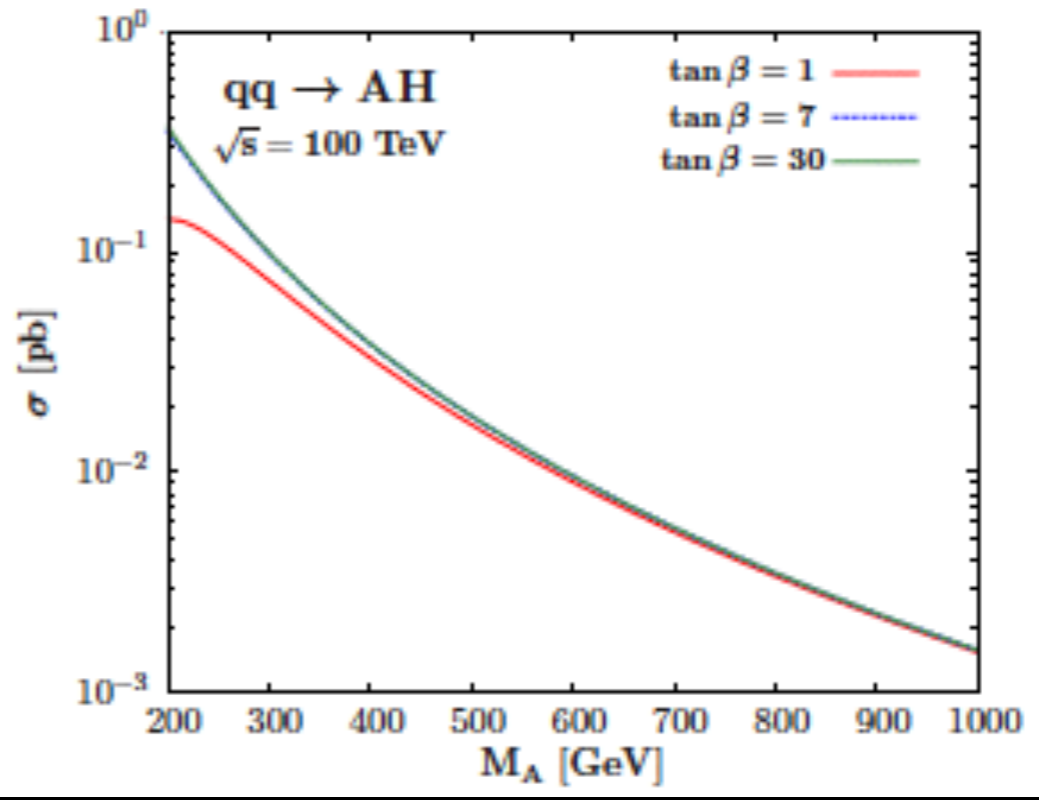
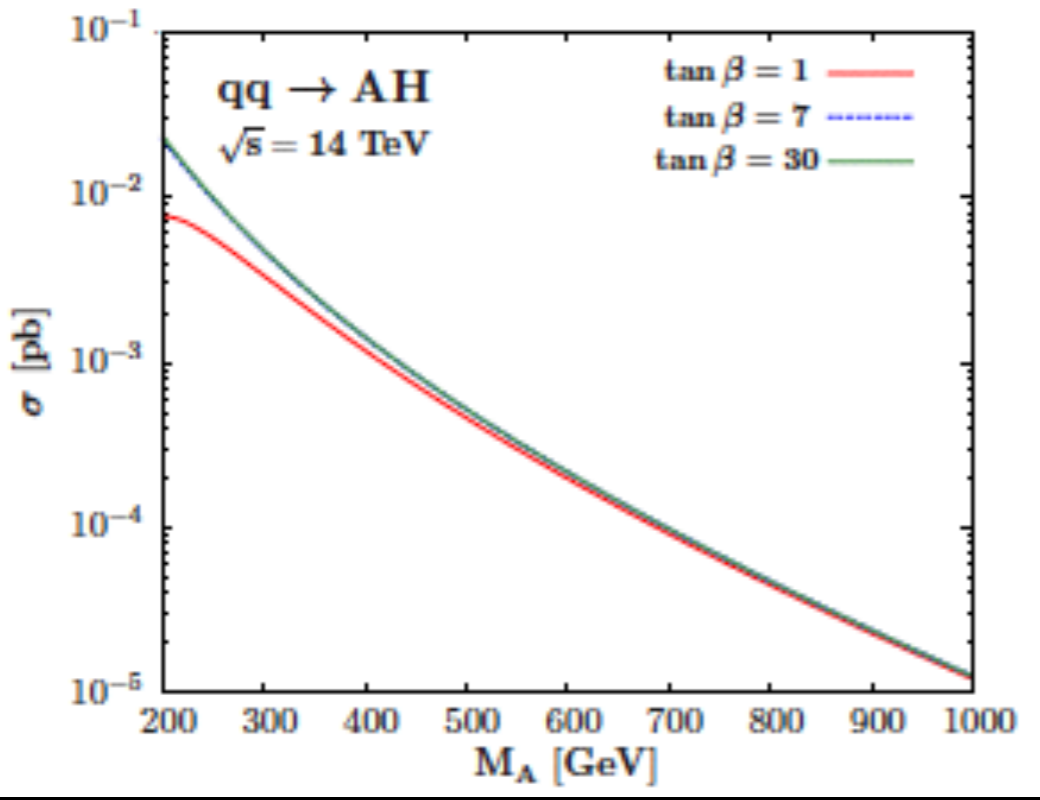
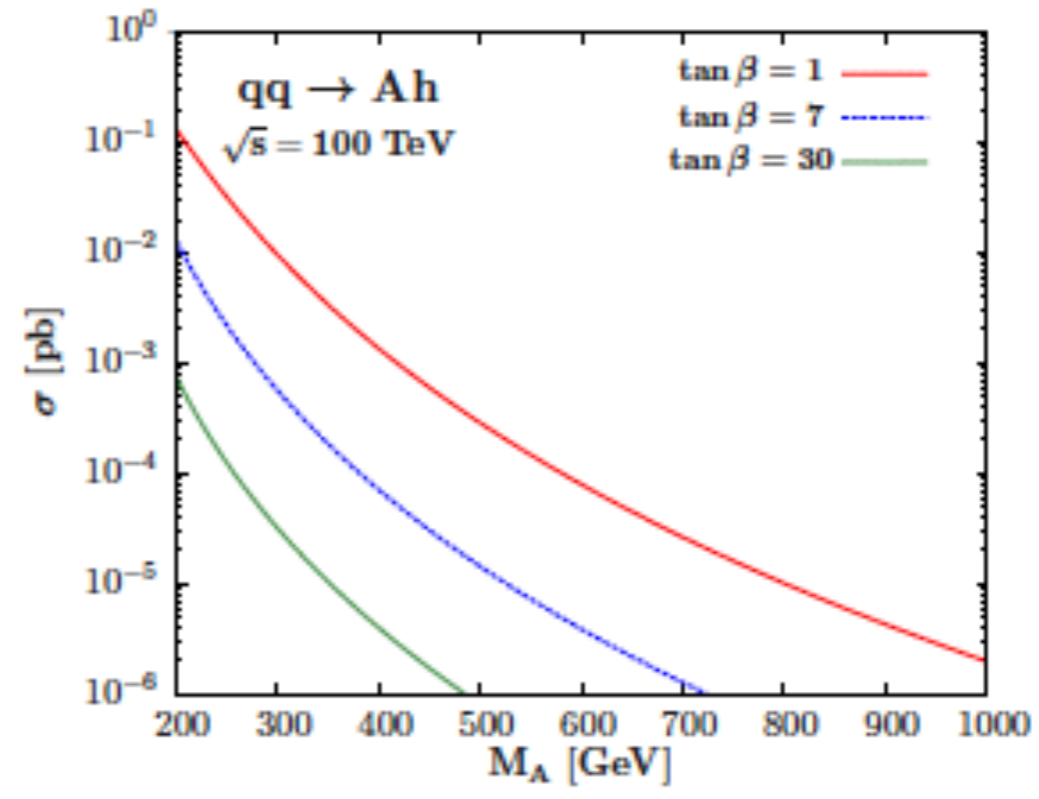
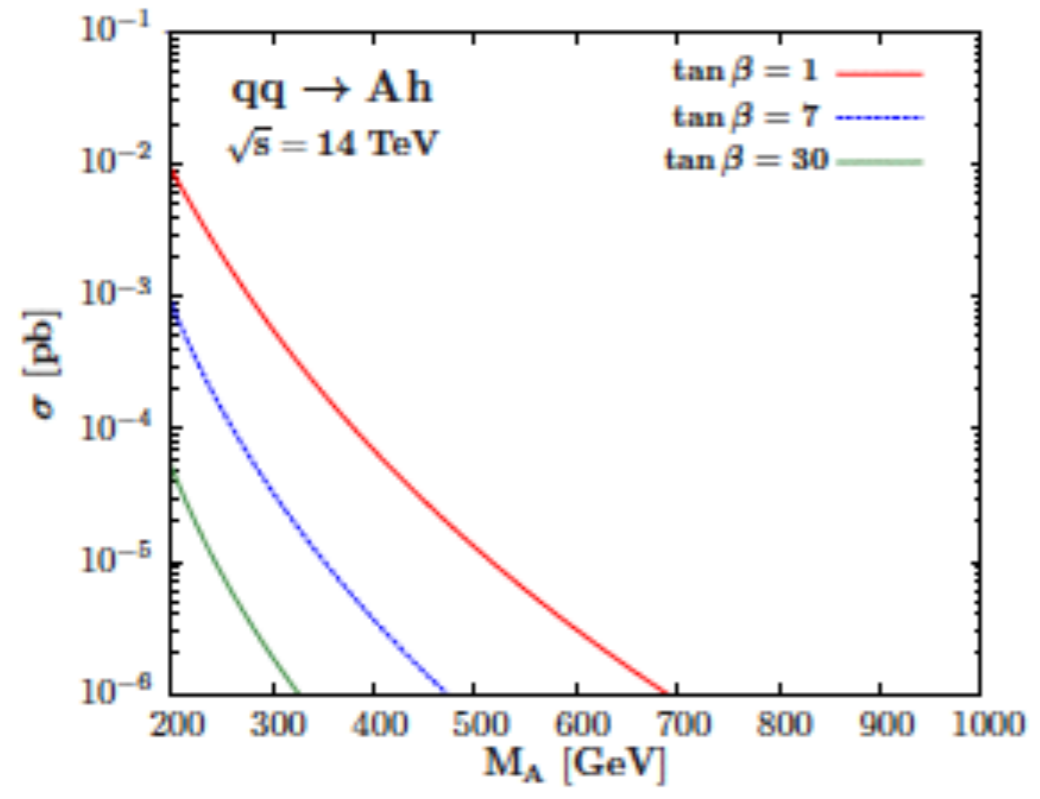
(a naive count)

3. Higgs Potentials@Colliders

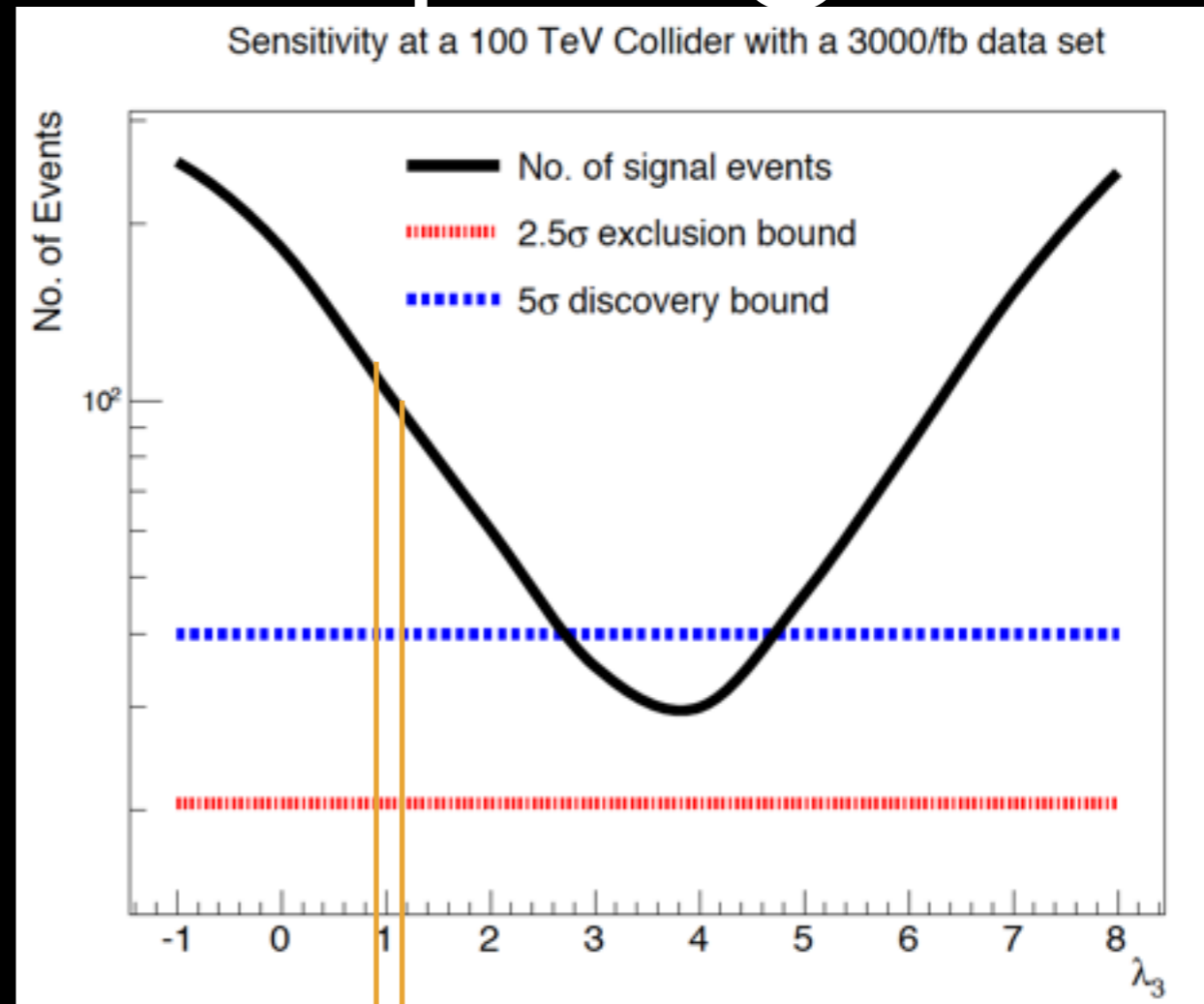
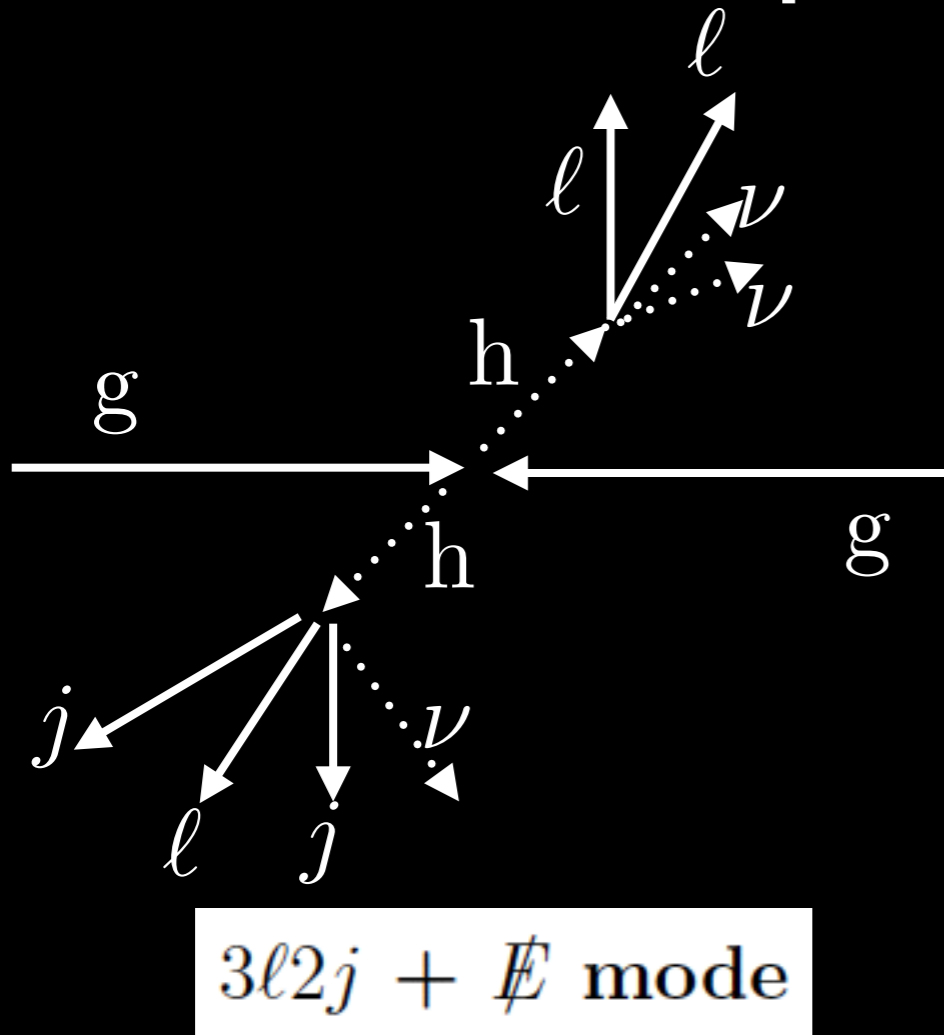
- The production and decay of new Higgs bosons
- The determinations of mixing angles
- The measurement of T/Q Couplings
- Global fit to the Higgs potential



J. Baglio, A. Djouadi, J. Quevillon, 1511.07853



Triple Coupling

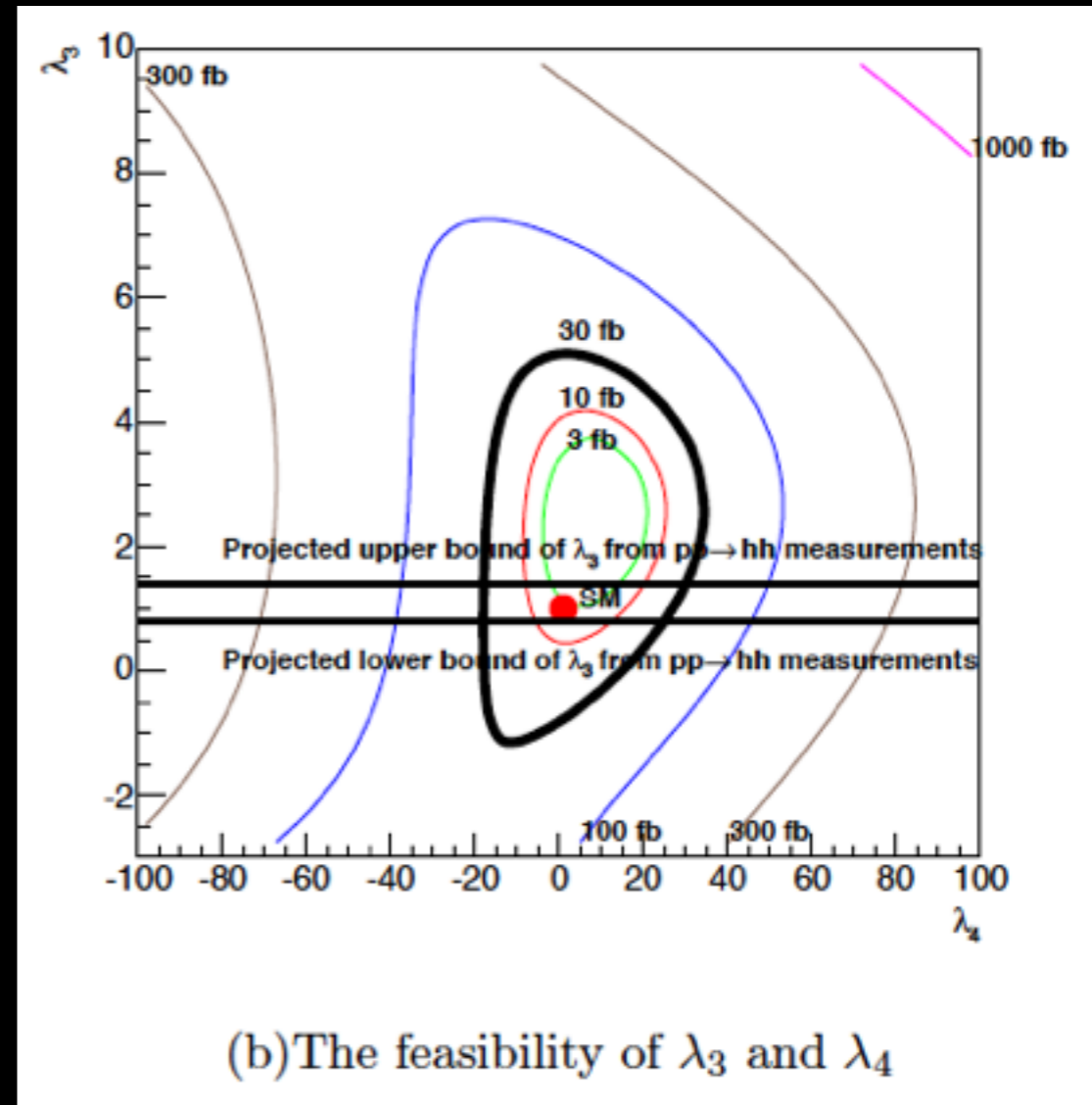
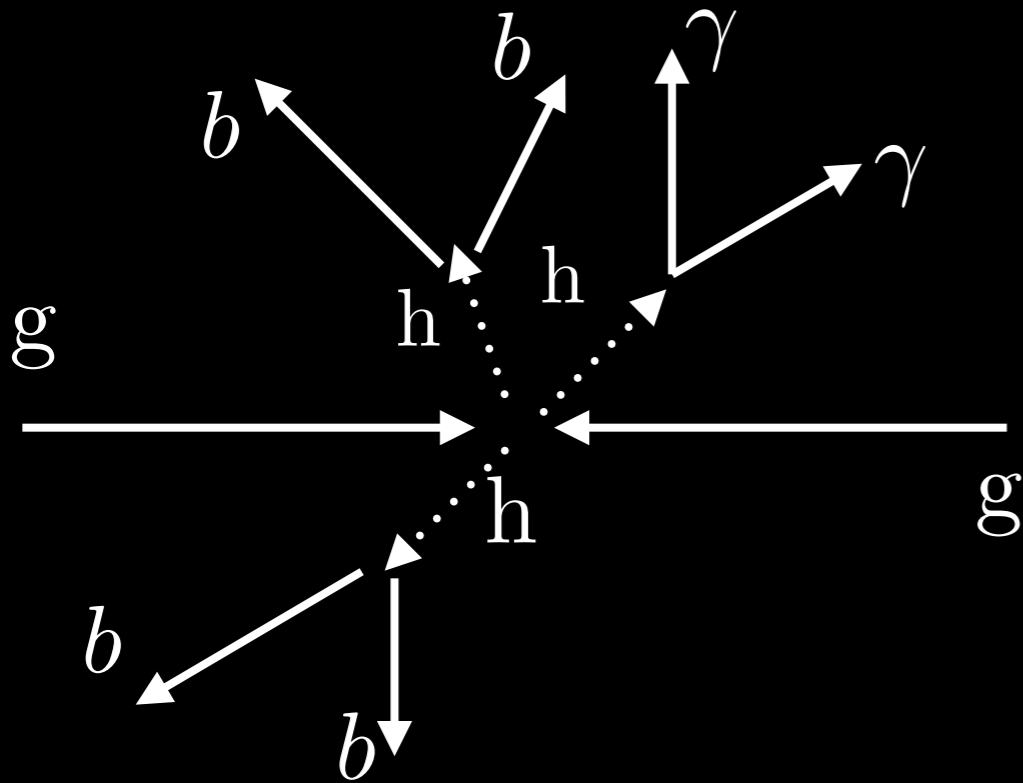


To overcome the b mistag and photo mistag issues,
 $gg \rightarrow hh \rightarrow WW^*WW^* \rightarrow 3l + 2j + MET$ is proposed.

By using this mode, SPPC can determine
 λ_3 to the window [0.9,1.2]

Q. Li, Z. Li, QY, X.R. Zhao, PRD92(2015)1,014015, arXiv:1503.07611

Quartic Coupling



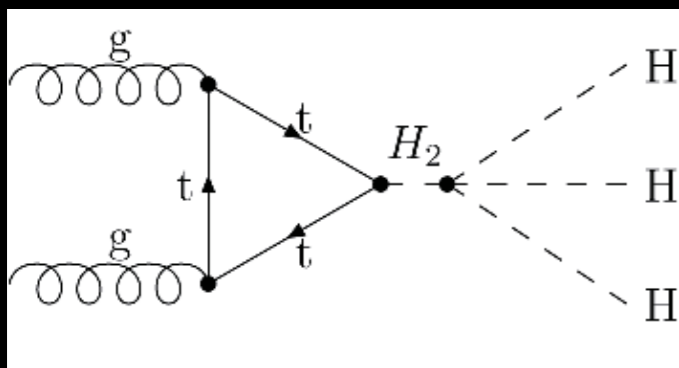
A 100 TeV collider can set a loose bound to λ_4
via $gg \rightarrow hhh \rightarrow 4b2\gamma$ mode.

Better modes are under searching.

Heavy Higgs in single+SM

	B1	B2	B3
m_{H_2} (GeV)	460	500	490
θ	0.354	0.354	0.354
a_2	3.29	3.48	3.43
b_3 (GeV)	-706	-612	-637
b_4	8.38	8.38	8.38

	B1	B2	B3
$\Gamma_{\text{tot}}(H_2)$ (GeV)	5.6	7.5	7.0
$BR(H_2 \rightarrow W^+W^-)$	0.57	0.56	0.57
$BR(H_2 \rightarrow ZZ)$	0.27	0.27	0.27
$BR(H_2 \rightarrow t\bar{t})$	0.15	0.16	0.16
$BR(H_2 \rightarrow b\bar{b})$	3.4×10^{-4}	2.8×10^{-4}	2.9×10^{-4}
$BR(H_2 \rightarrow HH)$	5.3×10^{-7}	8.8×10^{-7}	1.5×10^{-7}
$BR(H_2 \rightarrow HHH)$	1.0×10^{-3}	1.4×10^{-3}	1.3×10^{-3}
$\sigma(gg \rightarrow H_2)$ @ 14 TeV (fb)	3.2×10^2	2.3×10^2	2.5×10^2
$\sigma(gg \rightarrow HHH)$ @ 14 TeV (fb)	0.70	0.69	0.71
$\sigma(gg \rightarrow H_2)$ @ 100 TeV (fb)	1.4×10^4	1.1×10^4	1.2×10^4
$\sigma(gg \rightarrow HHH)$ @ 100 TeV (fb)	37	38	39

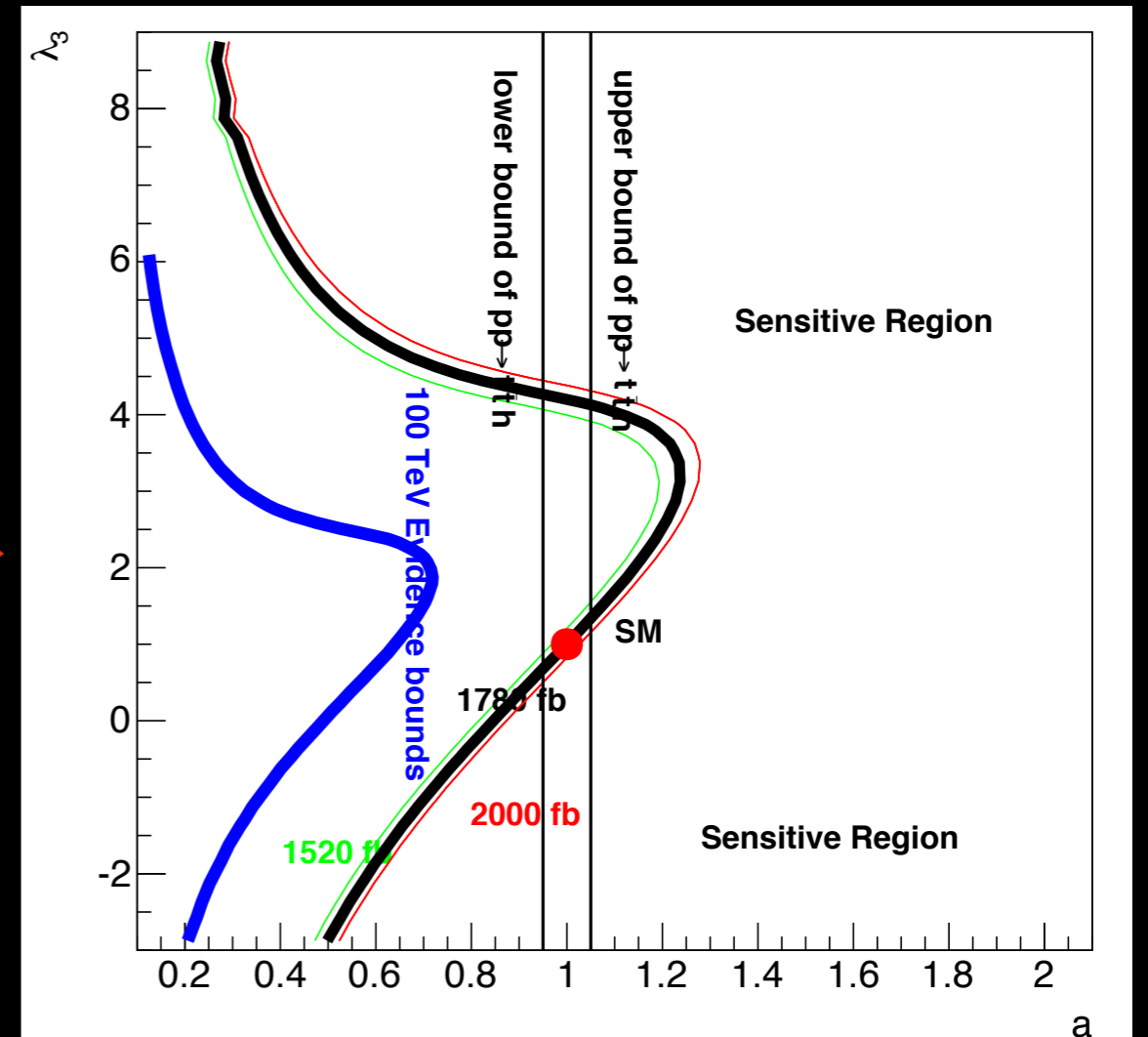
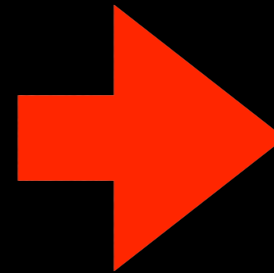
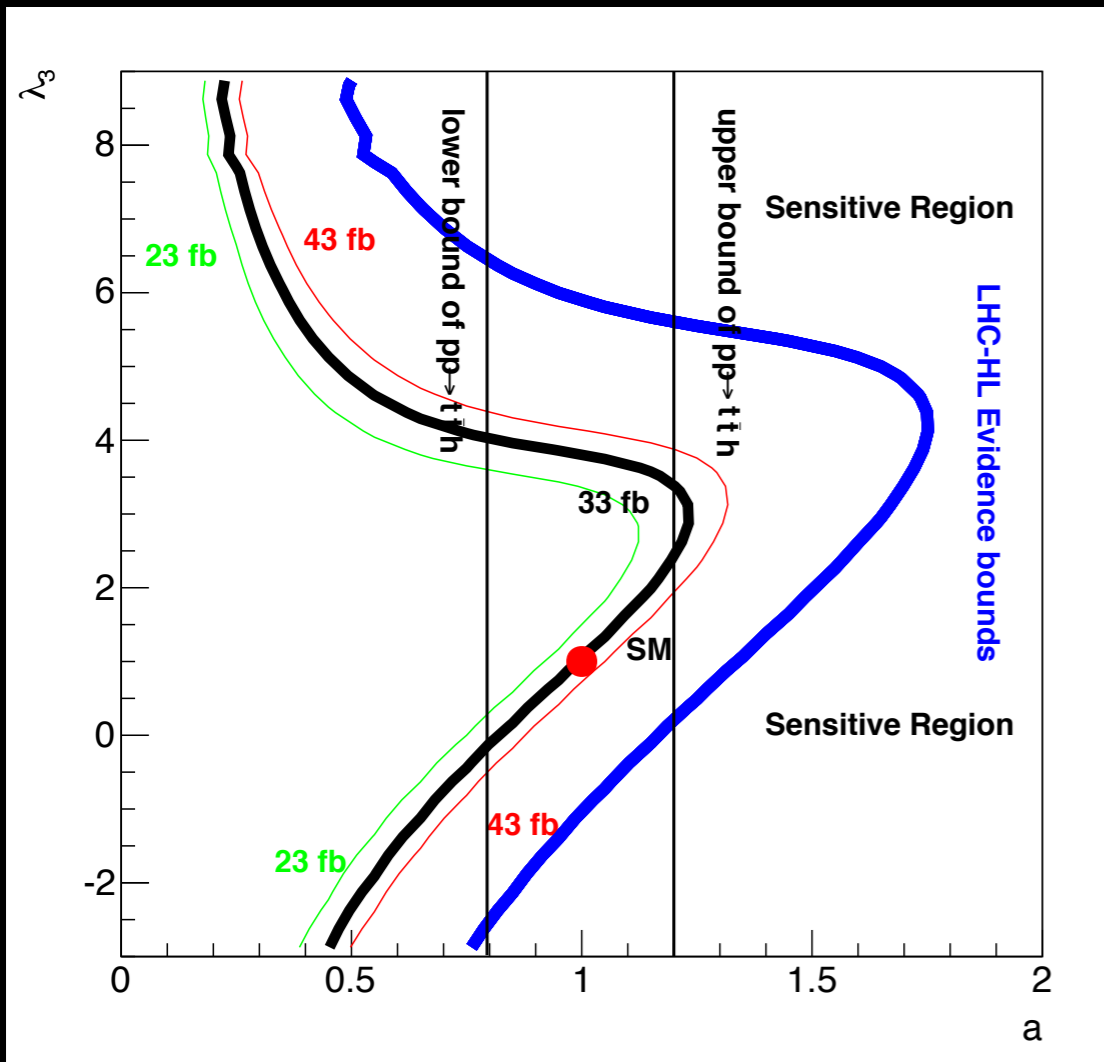


Production rate
is enhanced!

Interplay between $t\bar{t}h$ and hh

LHC-HL run

100 TeV SPPC

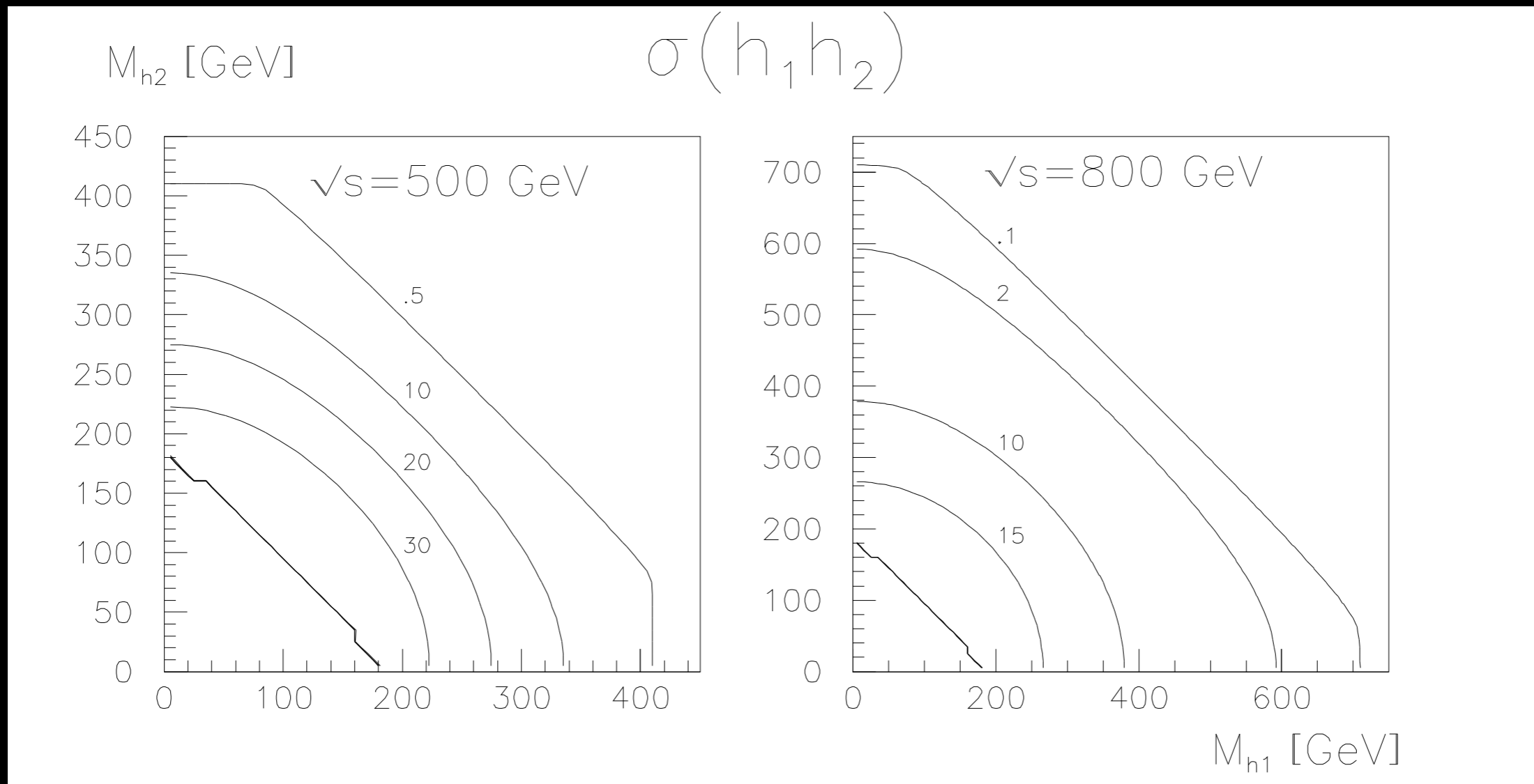


$$\Delta\mathcal{L} = Y_t (a \bar{t}t + i b \bar{t}\gamma_5 t) h + \lambda_3 \lambda_{SM} v h h h + \dots$$

J.F. Gunion, B. Grzadkowski, X.G. He, PRL77(1996)5172

Q. Li, Z. Li, QY, X.R. Zhao, in preparation

$$e^+ e^- \rightarrow h_i h_j$$



$$\frac{g_2}{2} Z_\mu \epsilon_{ijk} T_{i2} S_j \partial^\mu S_k$$

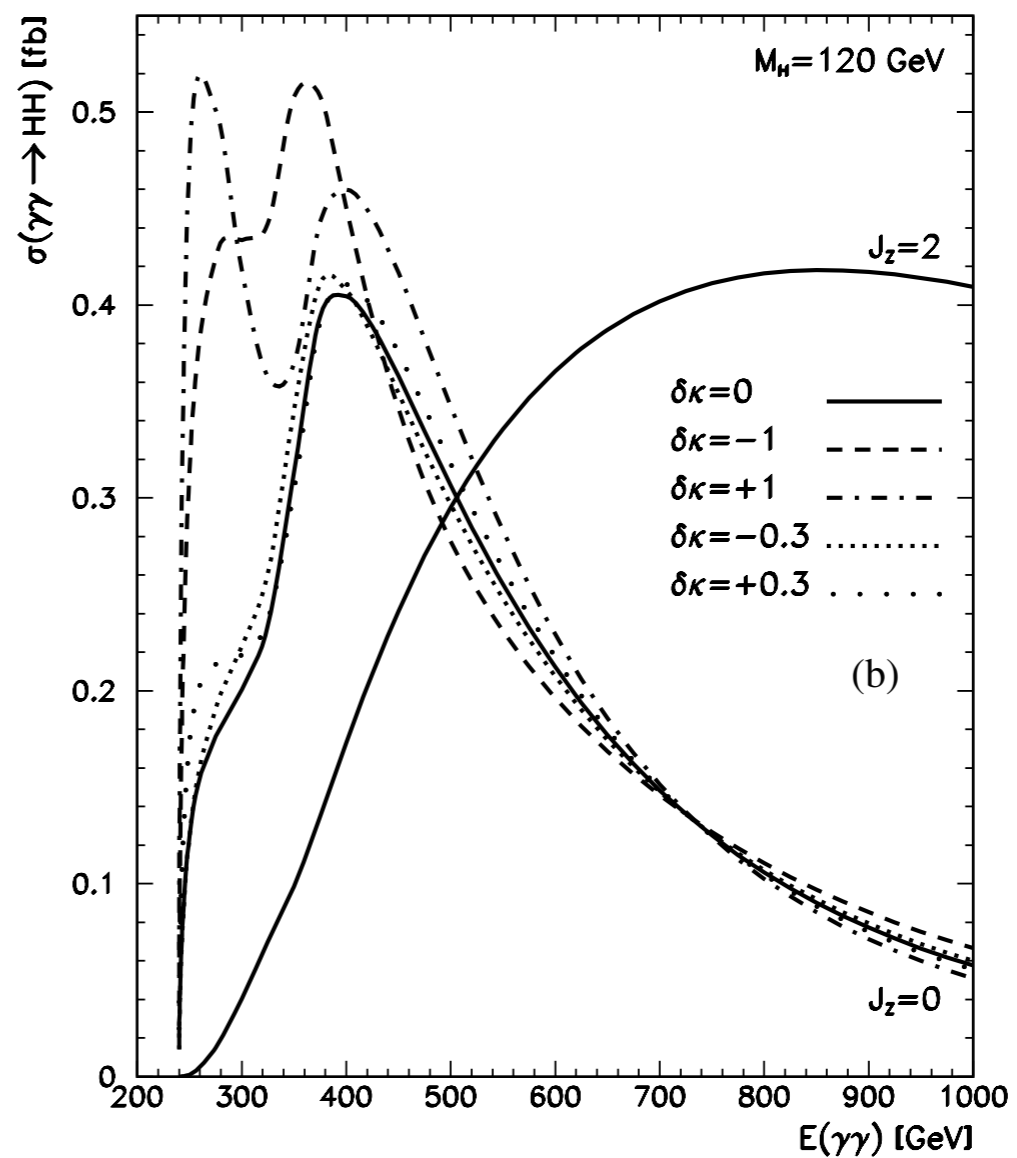
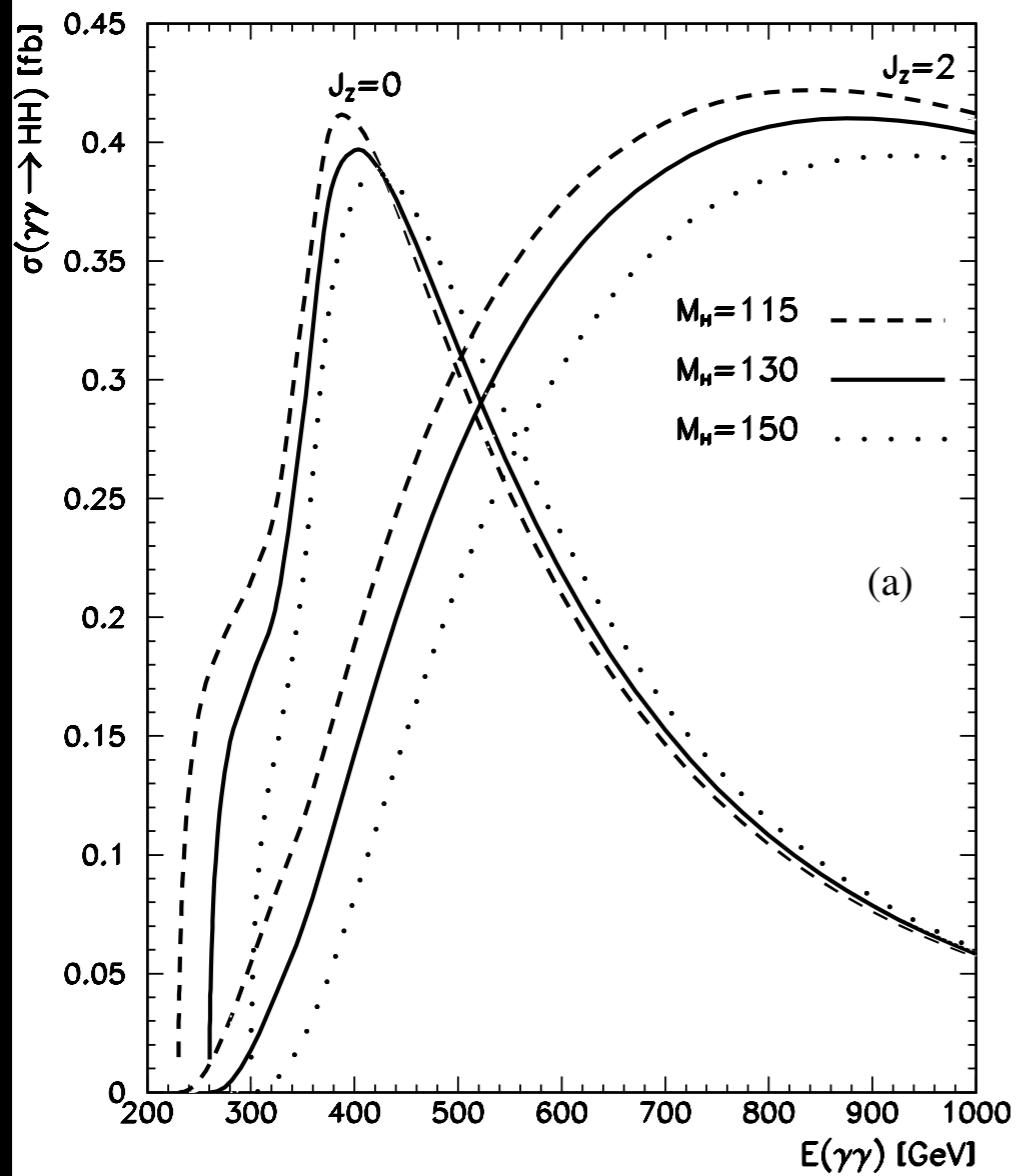
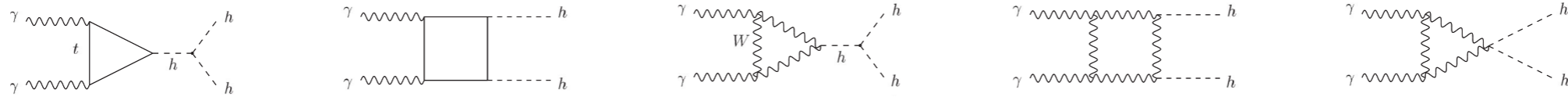
$$J_1 = -v^2 \text{Im}(\lambda_6^* \lambda_5^2)$$

$$= (m_1^2 - m_2^2)(m_2^2 - m_3^2)(m_3^2 - m_1^2) T_{11} T_{12} T_{13}$$

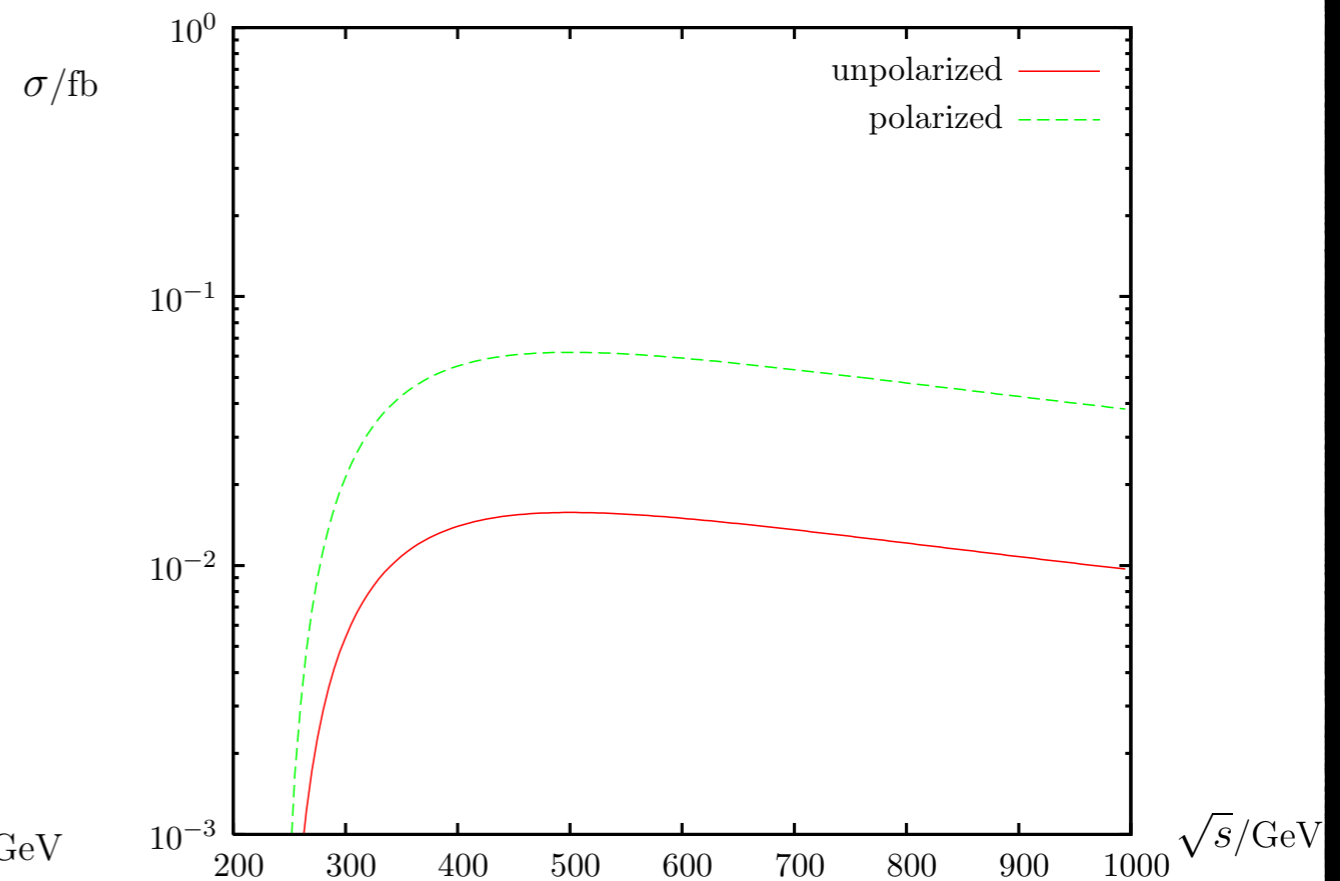
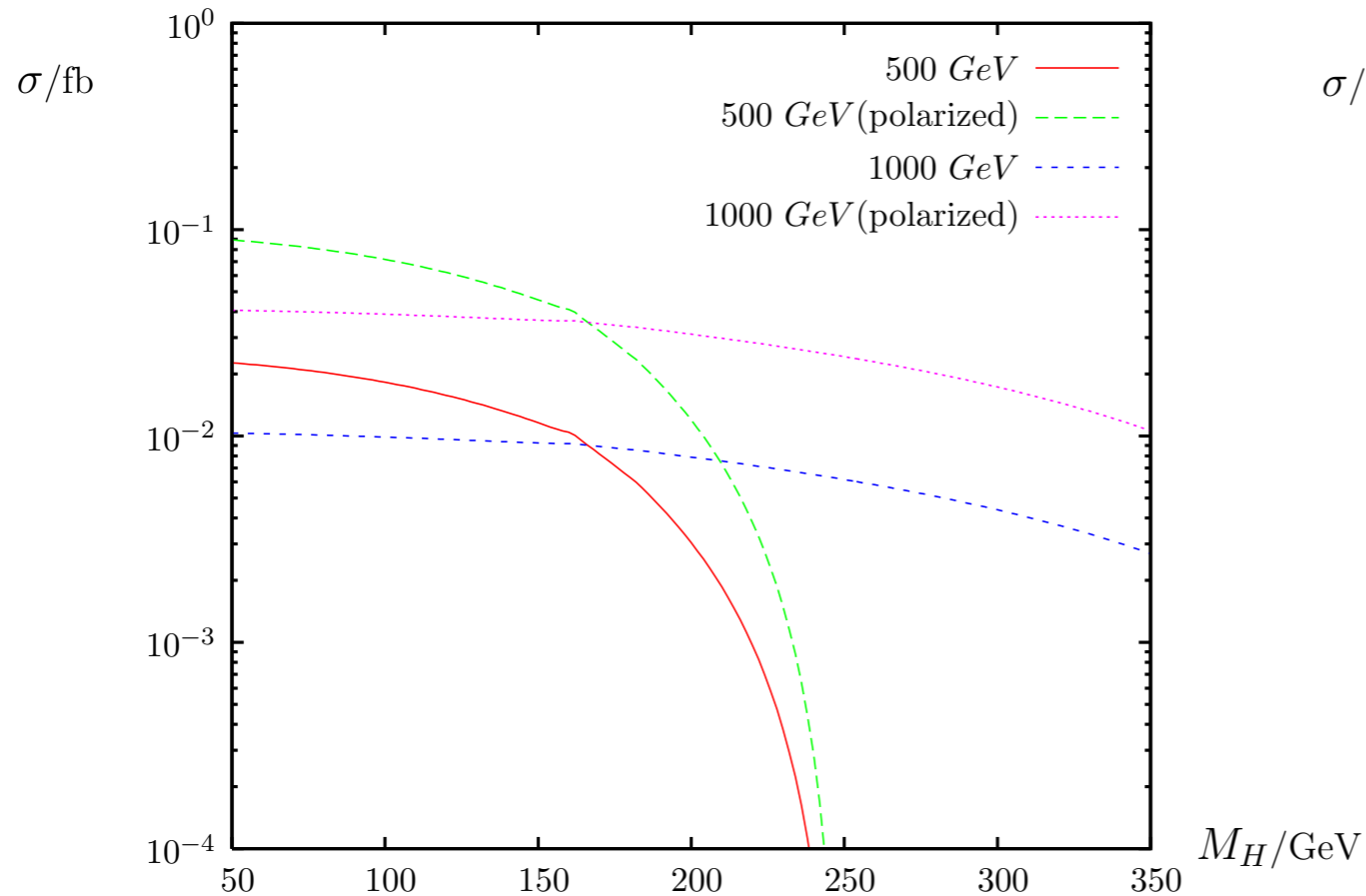
A. Mendez, A. Pomarol, PLB272(1991)313

B. Grzadkowski, G.F. Gunion, J. Kalinowski, PLB480(2000)287

$$\gamma\gamma \rightarrow hh$$

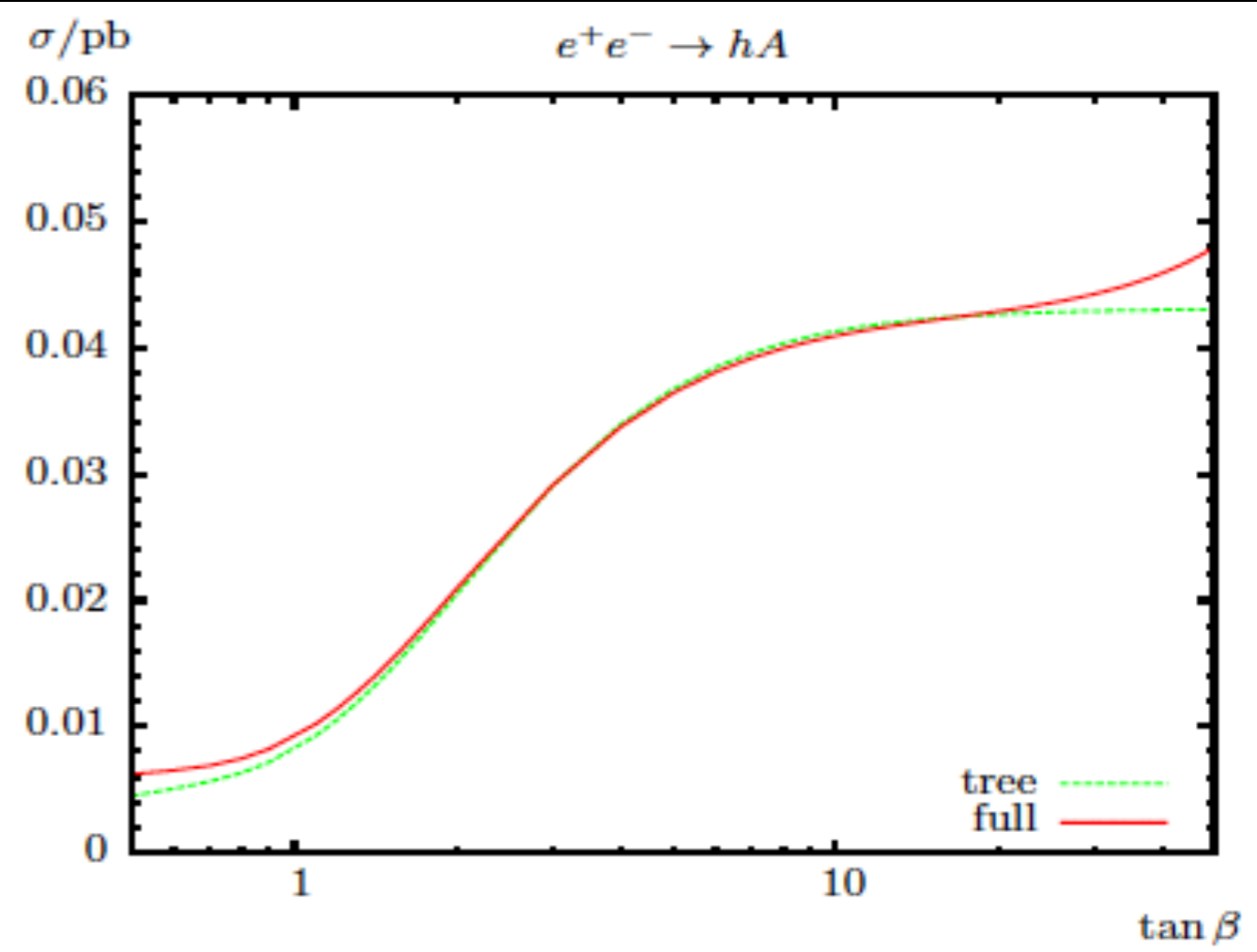
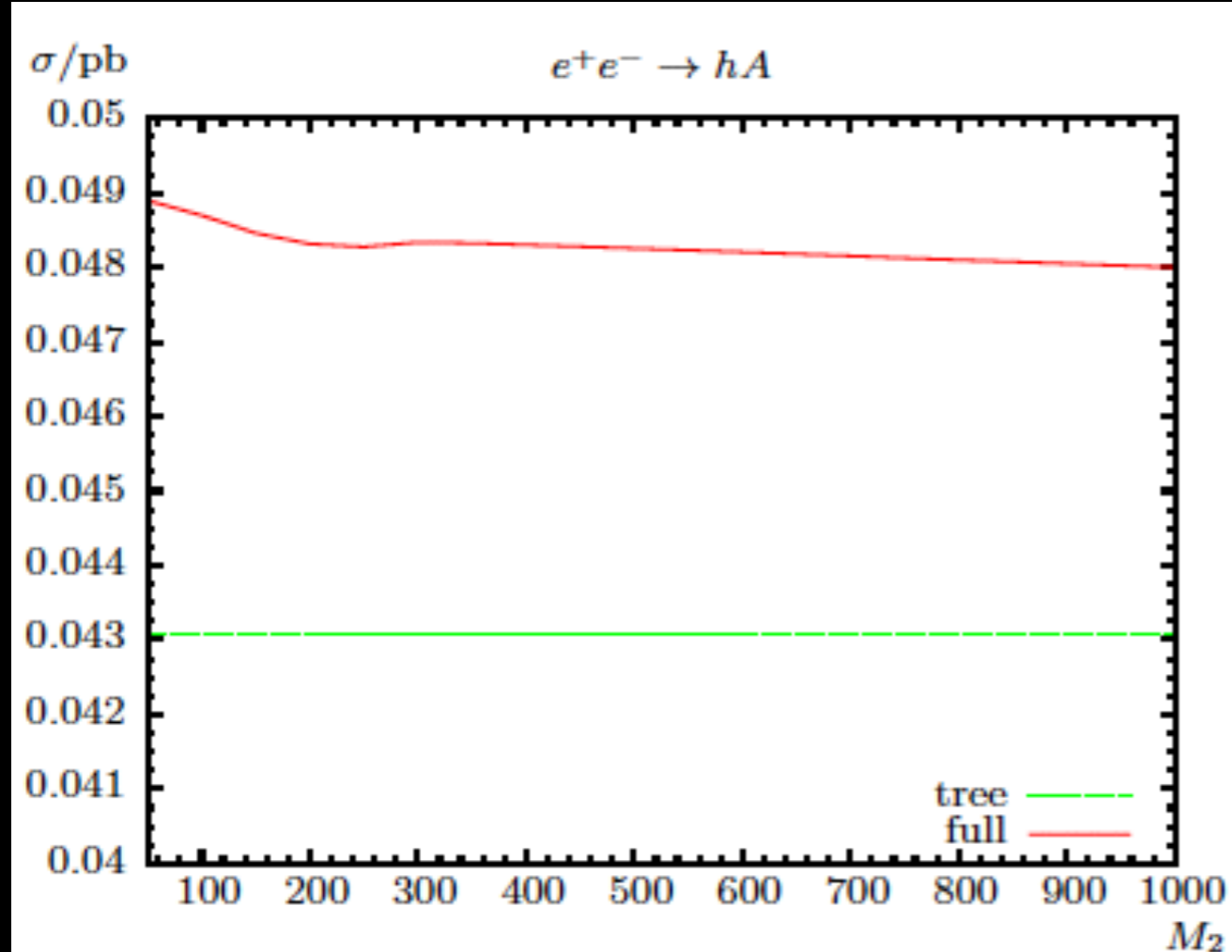


$$e^+e^- \rightarrow hh$$



J.J.Lopez-Villarejo and J.A.M. Vermaseren, PLB675,356

S. Heinemeyer and C. Schappacher, 1511.06002



S. Heinemeyer and C. Schappacher, 1511.06002

Higgs Decay :

	Order	Scale uncertainty	ref.	Comment
$b\bar{b}, c\bar{c}$	N^4LO, QCD	0.1%	0511063	massless
$\tau\bar{\tau}$	NLO, EW	1-2%	Z.Phys. C53,507	
gg	N^3LO, QCD	3%	0604194	Large m_t
$\gamma\gamma$	$N^2LO QCD +$ $NLO EW$	< 1%	1212.6233	
$Z\gamma$	$NLO QCD$	< 5%	PLB262 350	
$WW, ZZ \rightarrow H$	$NLO QCD$ $NLO EW$	$\sim 0.5\%$	0708.4123	interference included

Higgs Production @ HC			
	Order	Scale uncertainty	ref.
GF	N^3LO QCD	4%	1603.0800
VBF	N^2LO QCD	2%	1003.4451
HV	NNLO	2%	0307206
HH	N^2LO	10-15%	0211352

Prospects

- High Precision predictions for Higgs Physics
- Discovery new Higgs Bosons
- Multi-Higgs boson Final States, Discovery and Precision measurement
- CP phases determination of Higgs potentials
- Higgs portal and dark matter