

ILC Physics

Keisuke Fujii

KEK

arXiv: 1506.05992 (ILC Physics Case)

arXiv: 1506.07830 (ILC Run Scenarios)

arXiv: 1306.6352 (ILC TDR: Physics)

EPJC (2015) 75:371 (LC Physics)

International *Linear Collider* (ILC)

Energy Frontier
in e^+e^- collisions

$E_{cm} = 500\text{GeV}$ (total length: 31km)
Expandable to 1TeV (~50km)

of DRFS Klystron: 7280 total
of Cryomodules : 1680 total
of Cavities : 14560 total

Tunnel Layout Plan for a Japanese Mountain Site

Ultra-low emittance

normalized emittance = 37nm

Nano-beam collisions

High gradient

world highest gradient as with
super-conducting cavities
= 31.5 MV/m
beam current = 5.8 mA

Cryomodules housing
Super Cond. Cavities

Technologies at hand!

Damping Ring

Beam Delivery System

Detectors

ILD

ILD

SiD

High resolution and high
granularity detectors

International Effort for the ILC project

2004 technology choice (SCRF)

basic specifications

2005 **GDE** started

2007 Reference Design Report

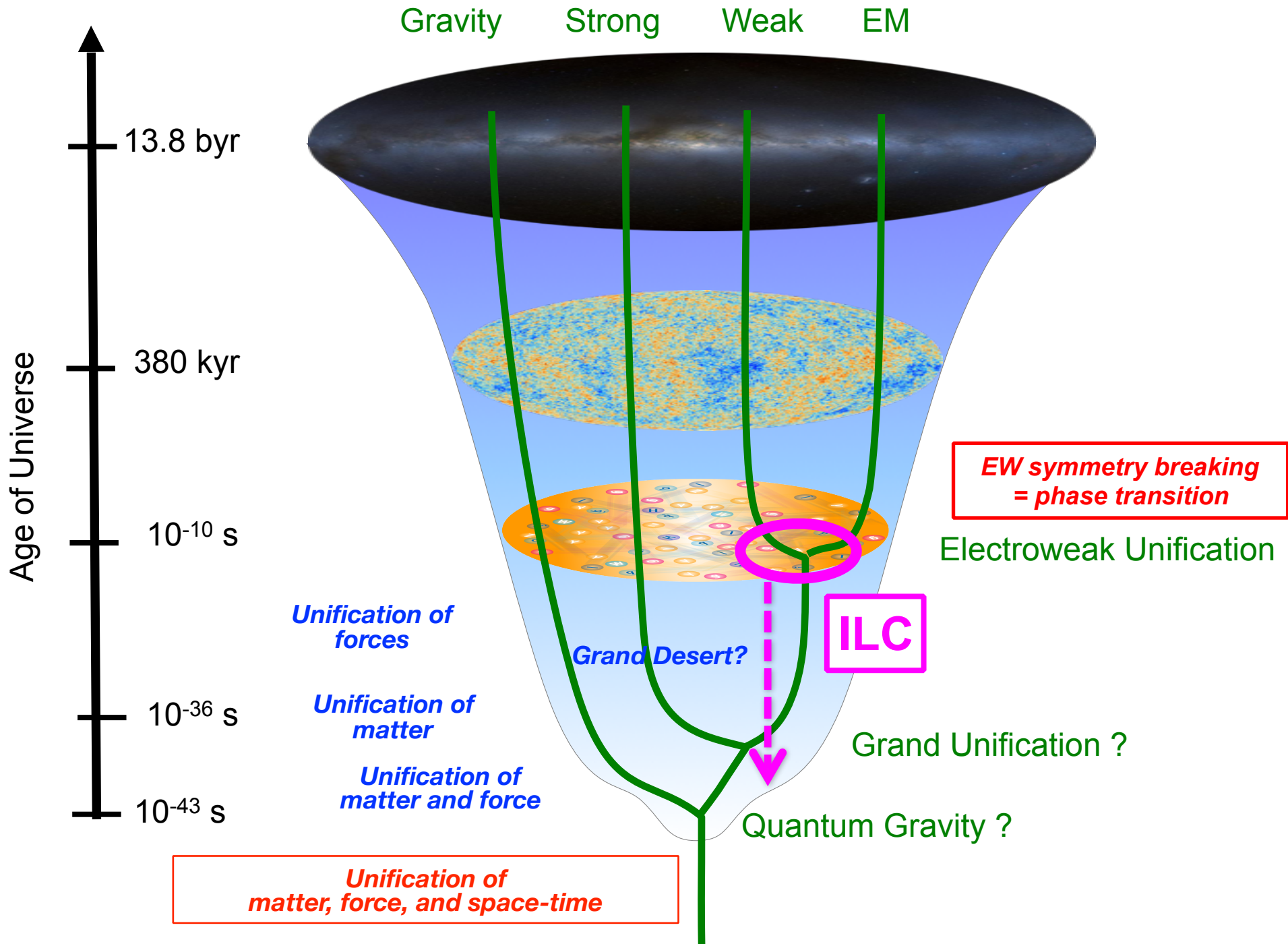
2012 Technical Design Report

2013 GDE completed its mission
by publication of the TDR

→ Linear Collider Collab. (**LCC**) ²

Physics at ILC

Towards ultimate unification

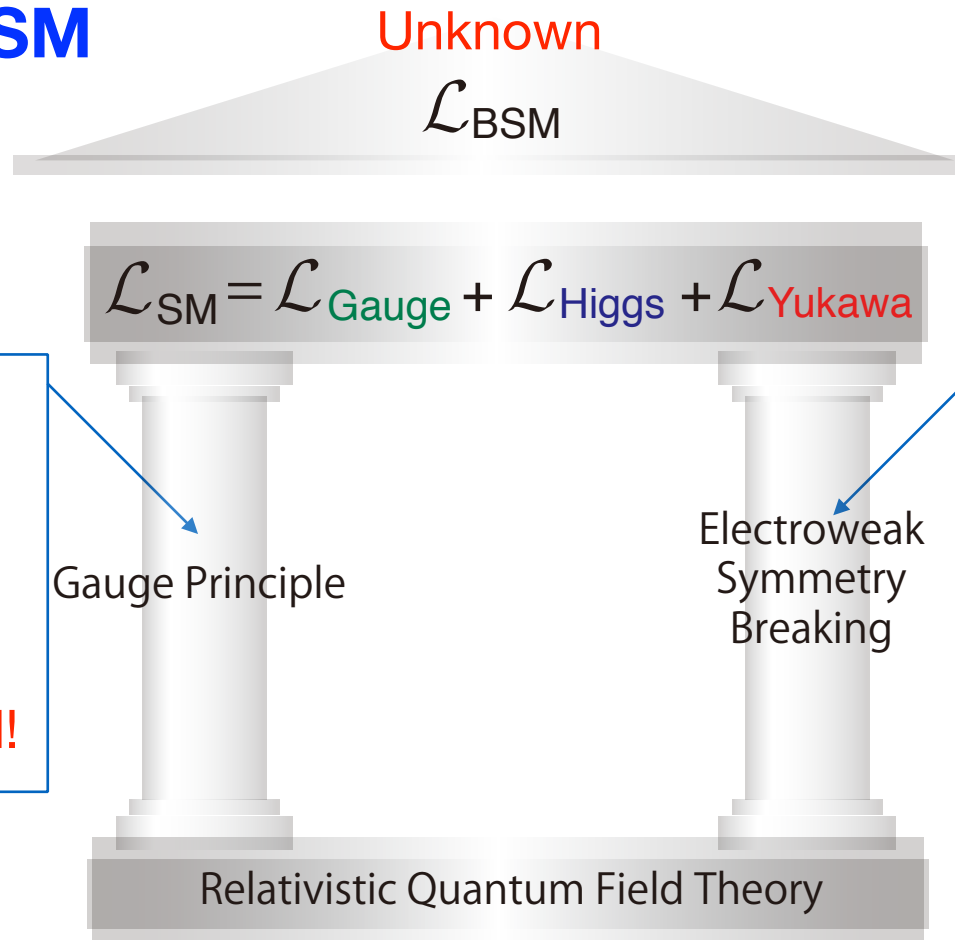


**Why is the EW scale
so important ?**

Why is the EW scale so important?

Mystery of something in the vacuum

2 Pillars of SM



Success of SM
= success of
gauge theory
(left pillar)

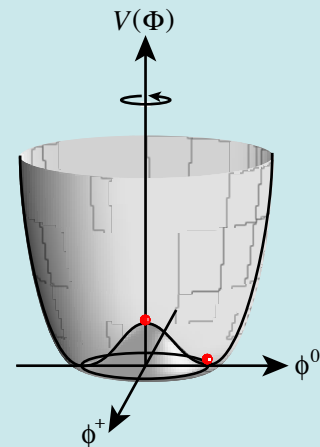
Precisely tested!

Vacuum filled with weak
charge (evidence: H125)

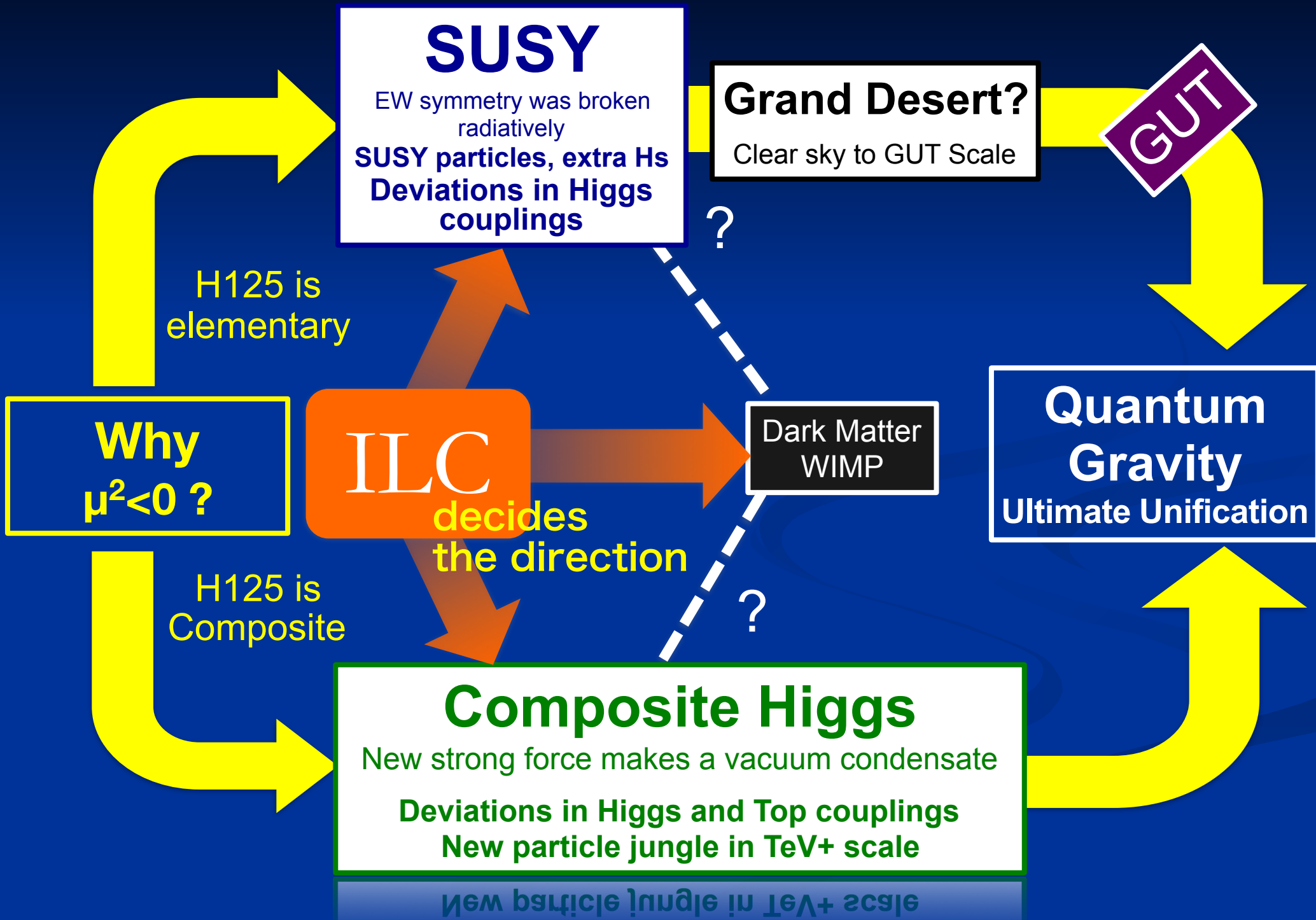
The nature of the
Higgs field - its
multiplet structure &
dynamics behind it -
is all unknown!

The SM does not explain **why the Higgs field developed a vacuum expectation value** (*Why $\mu^2 < 0$?*)! The answer forks depending on whether **H125 is elementary or composite!**

$$V(\phi) = \mu^2 |\phi|^2 + \lambda |\phi|^4$$



Big Branching Point at the EW Scale

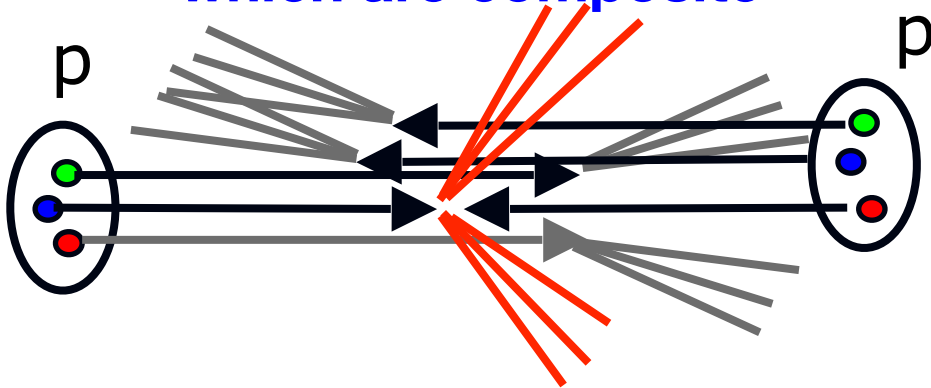


**The 3 major probes
for BSM at ILC:**

***Higgs, Top, and*
search for
*New Particles***

3 Powerful Tools

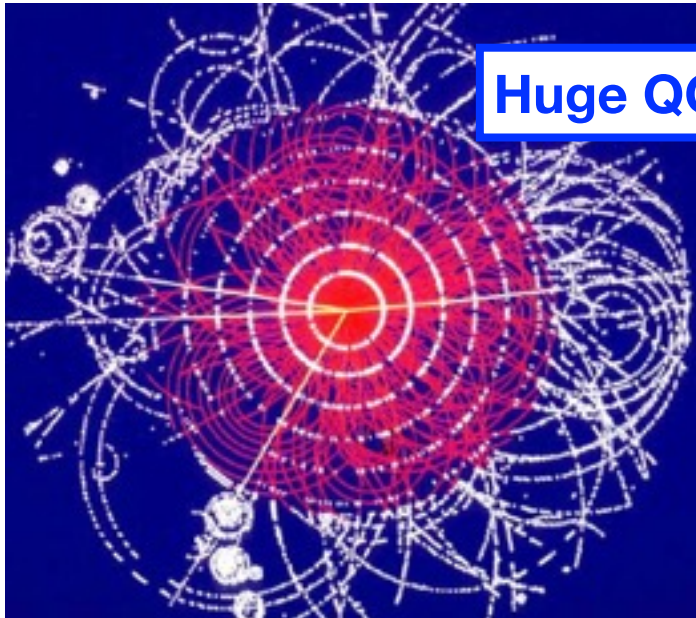
**LHC: Collision of protons
which are composite**



E_{cm} 7-14 TeV

Pileup

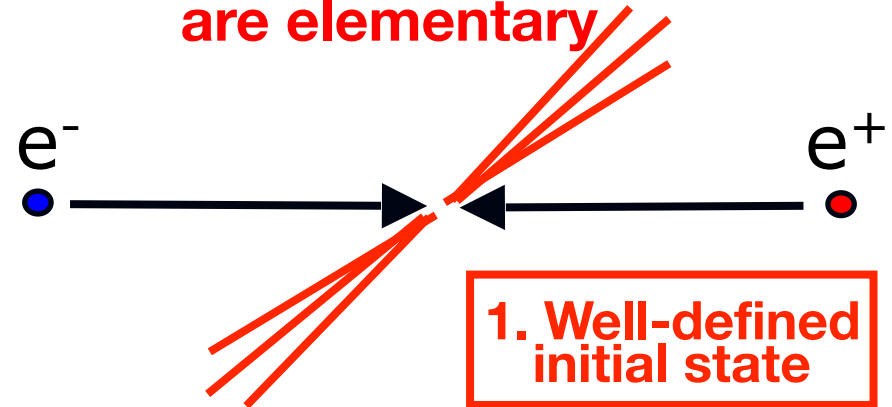
Initial state not very well defined



Huge QCD BG

proton is composite \Rightarrow events are complicated but
maximum reachable energy is high!

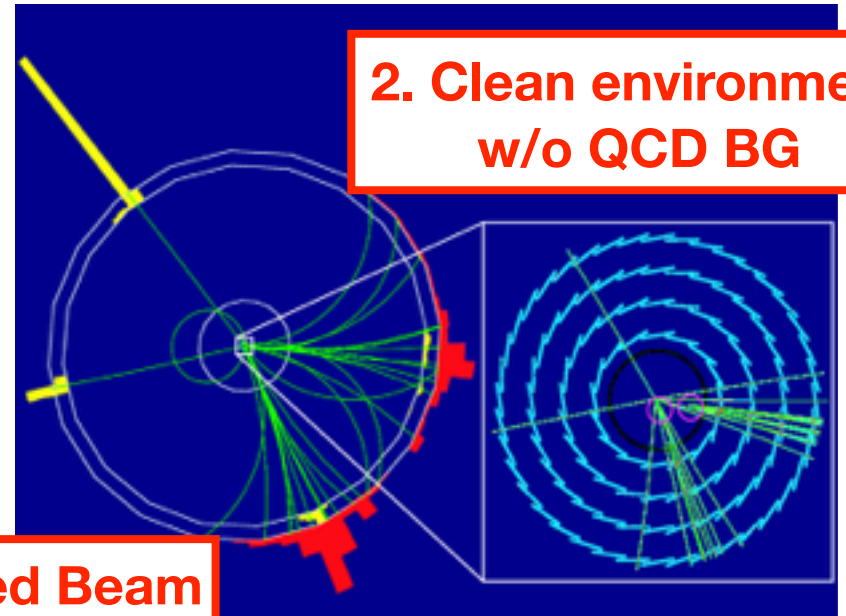
**ILC: Collision of e^+e^- which
are elementary**



**1. Well-defined
initial state**

E_{cm} 0.25-1 TeV

Lab. frame = CM frame



**2. Clean environment
w/o QCD BG**

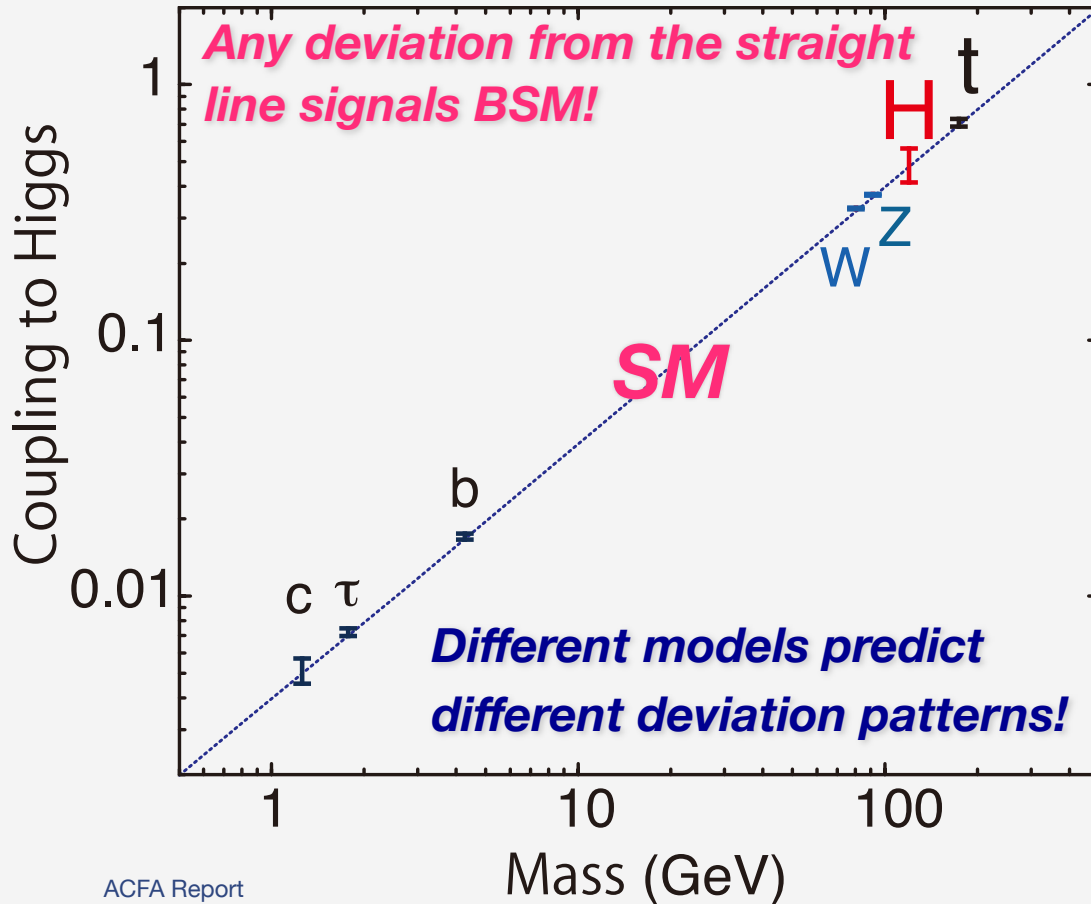
3. Polarized Beam

clean and and able to detect everything produced!

Higgs

Deviation in Higgs Couplings

Mass-coupling relation



The size of the deviation depends on the new physics scale (Λ)!

Decoupling Theorem:
 $\Lambda \uparrow \rightarrow SM$

example 1: **Minimal SUSY**

(MSSM : $\tan\beta=5$, radiative correction factor ≈ 1)

$$\frac{g_{hbb}}{g_{h_{SM}bb}} = \frac{g_{h\tau\tau}}{g_{h_{SM}\tau\tau}} \simeq 1 + 1.7\% \left(\frac{1 \text{ TeV}}{m_A} \right)^2$$

heavy Higgs mass

example 2: **Minimal Composite Higgs Model**

$$\frac{g_{hVV}}{g_{h_{SM}VV}} \simeq 1 - 8.3\% \left(\frac{1 \text{ TeV}}{f} \right)^2$$

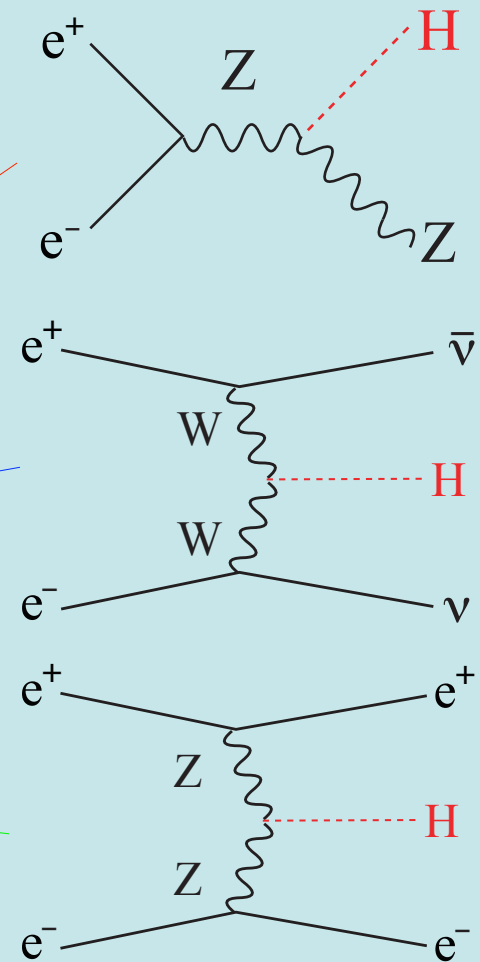
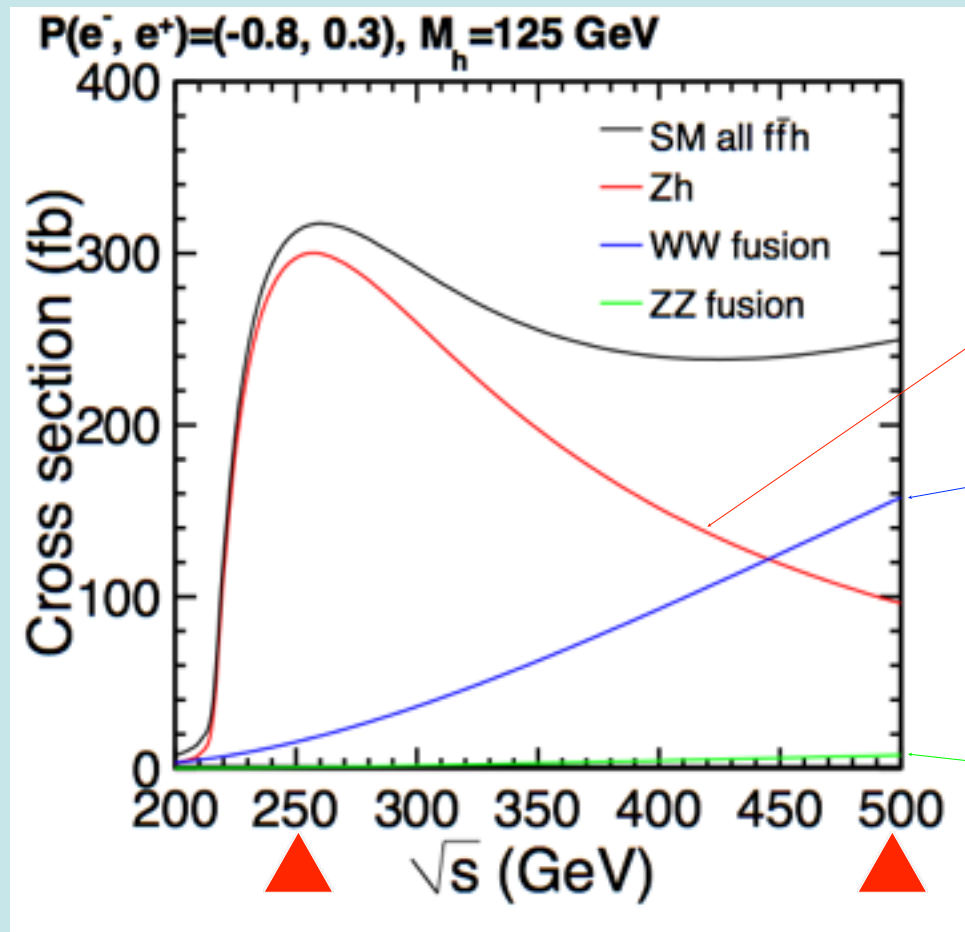
composite scale

New physics at 1 TeV \rightarrow deviation is at most $\sim 10\%$
 We need a %-level precision \rightarrow LHC is not enough \rightarrow **ILC**

Main Production Processes

Single Higgs Production

Production cross section



ZH dominates at 250 GeV
(~80k ev: 250 fb⁻¹)

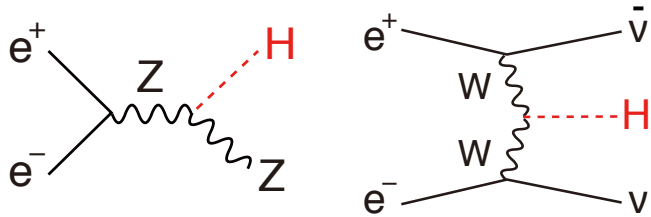
vvH takes over at 500 GeV
(~125k ev: 500 fb⁻¹)

200k w/ TDR baseline, eventually >1M Higgs events!

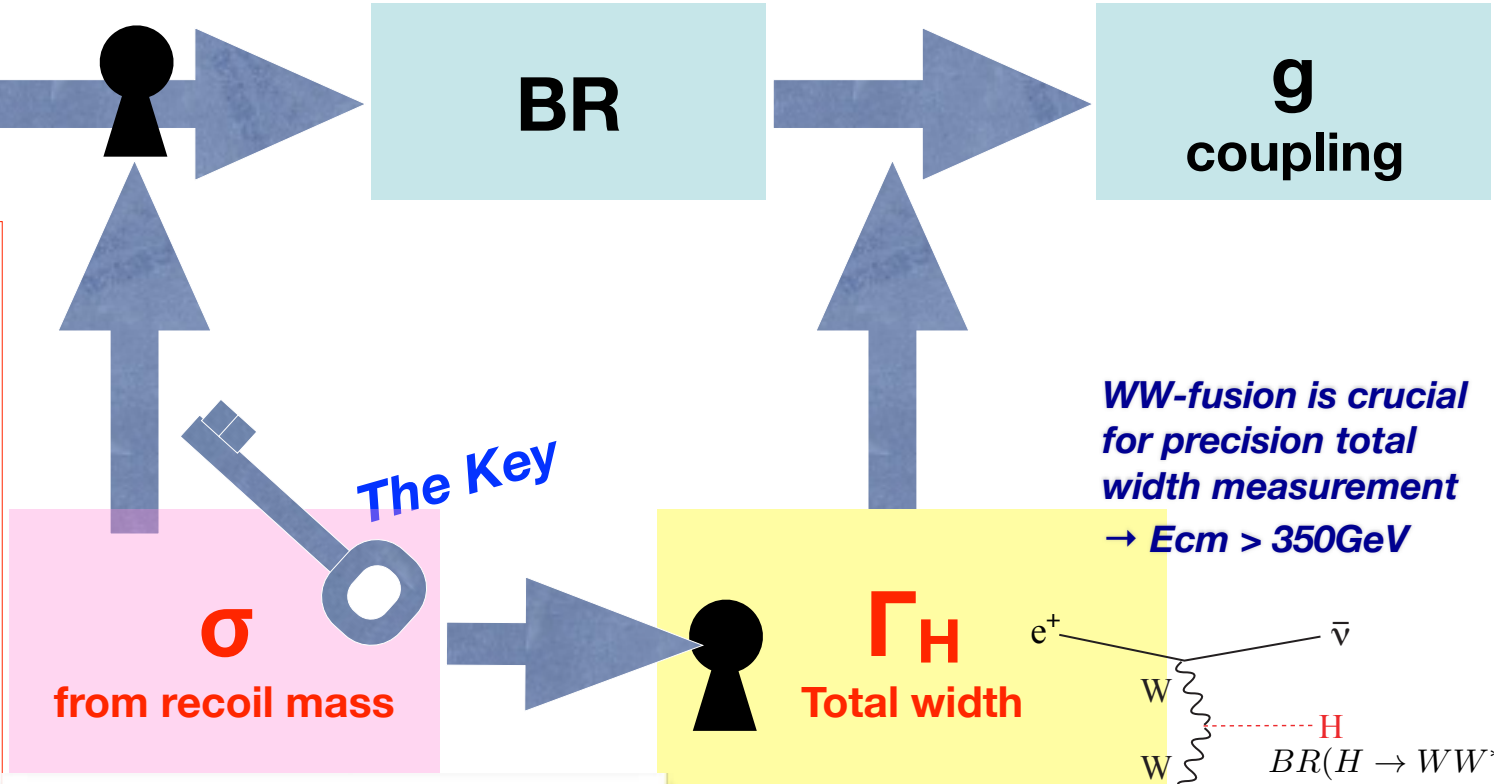
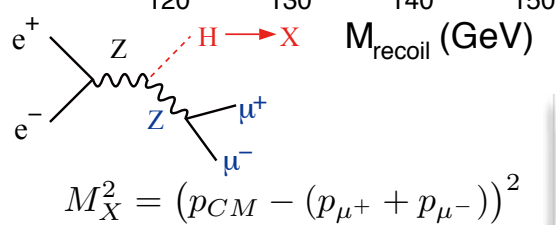
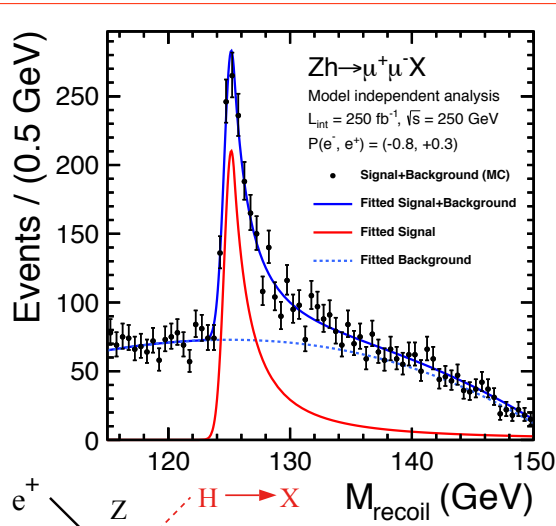
Key Point

At LHC all the measurements are $\sigma \times BR$ measurements.

At ILC all but **the σ measurement using recoil mass technique** is $\sigma \times BR$ measurements.



$$g_{HAA}^2 \propto \Gamma(H \rightarrow AA) = \Gamma_H \cdot BR(H \rightarrow AA)$$



Can detect even if Higgs decays invisibly!

Higgs Couplings

Model-independent coupling fit, impossible at LHC

H20 Scenario

arXiv: 1506.05992
arXiv: 1506.07830

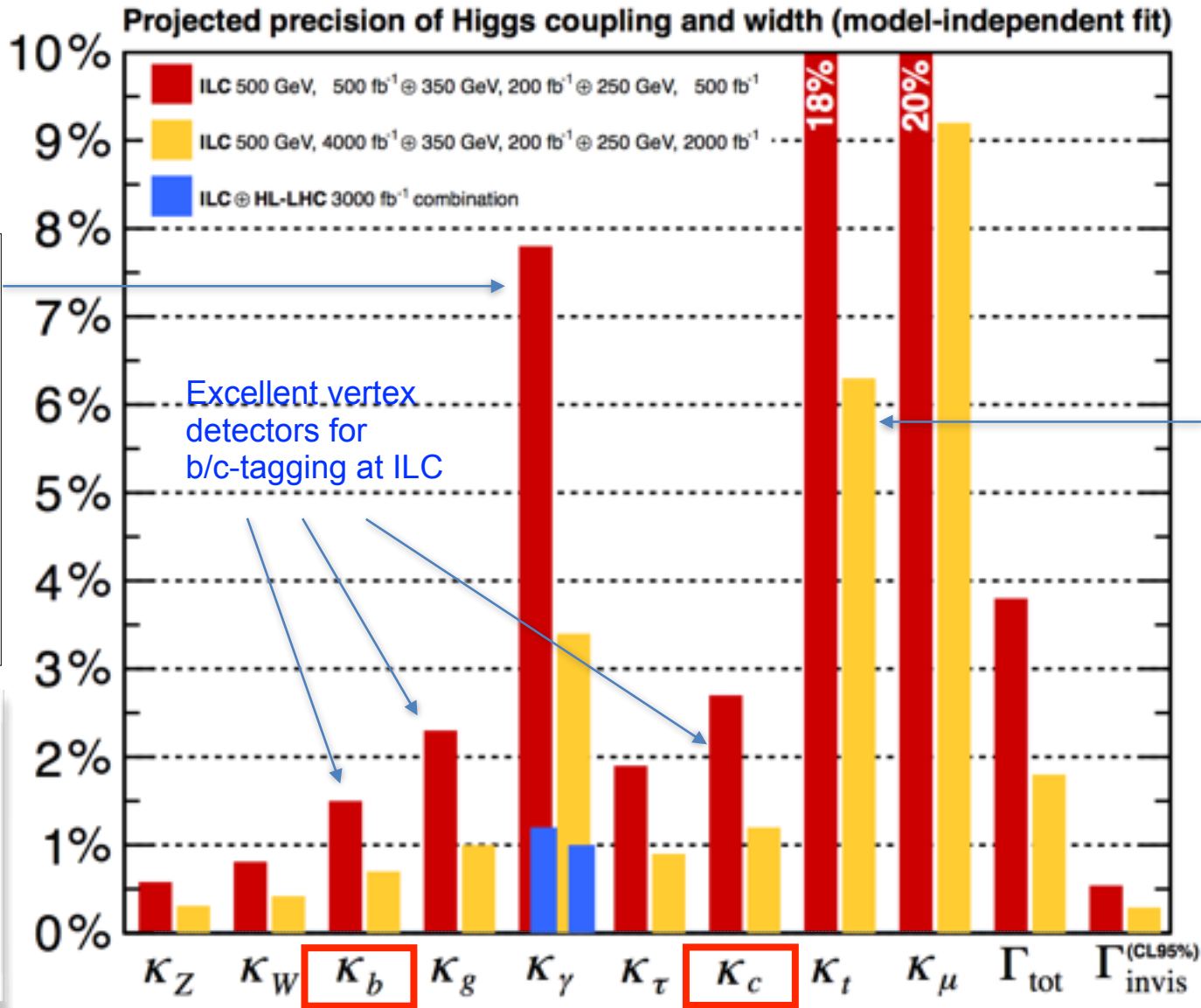
Better hyy with LHC/ILC synergy

LHC can precisely measure

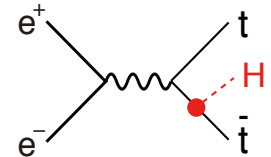
$$BR(h \rightarrow \gamma\gamma) / BR(h \rightarrow ZZ^*) = (K_\gamma / K_Z)^2$$

ILC can precisely measure K_Z

All of major Higgs decay modes accessible at ILC with 250-500GeV!



Top Yukawa improves by going to 550 GeV



Near threshold → a factor of 4 enhancement of σ_{th} by going from 500GeV to 550 GeV

→ 3%

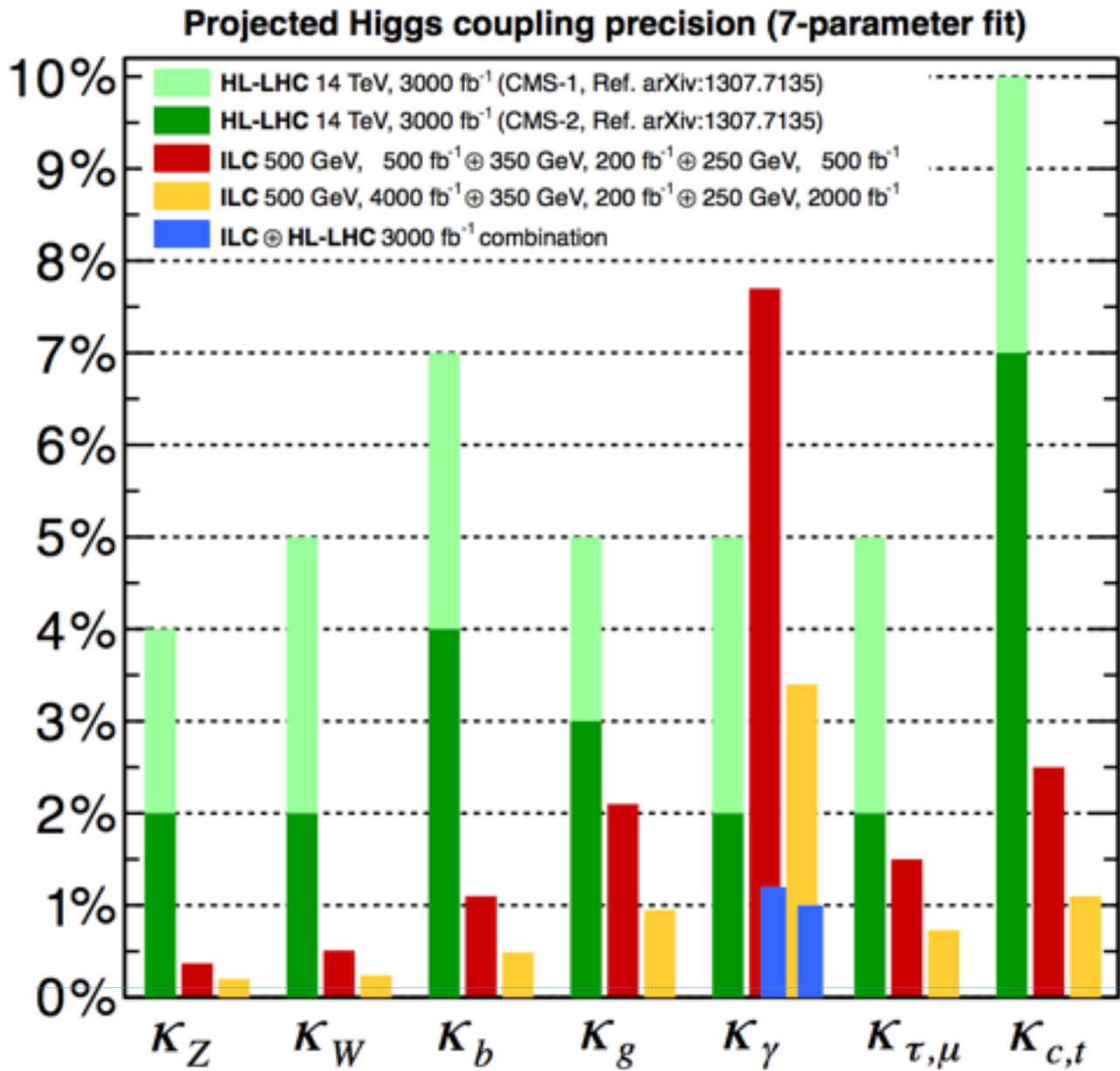
500 GeV already excellent except for K_t , K_μ , and K_γ

~1% or better for most couplings!

Model-dependent coupling fit (LHC-style 7-parameter fit)

H20 Scenario
 arXiv: 1506.05992
 arXiv: 1506.07830

$\Sigma_{SM} BR = 1$

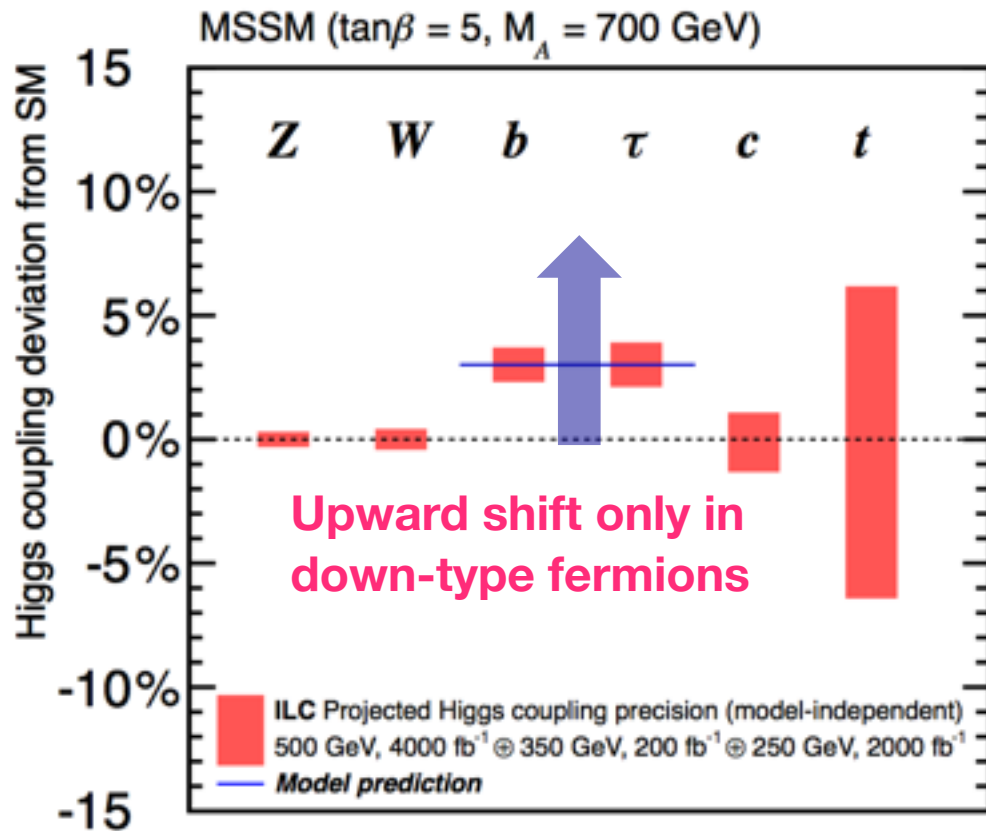


Possible to achieve precision far exceeding LHC!

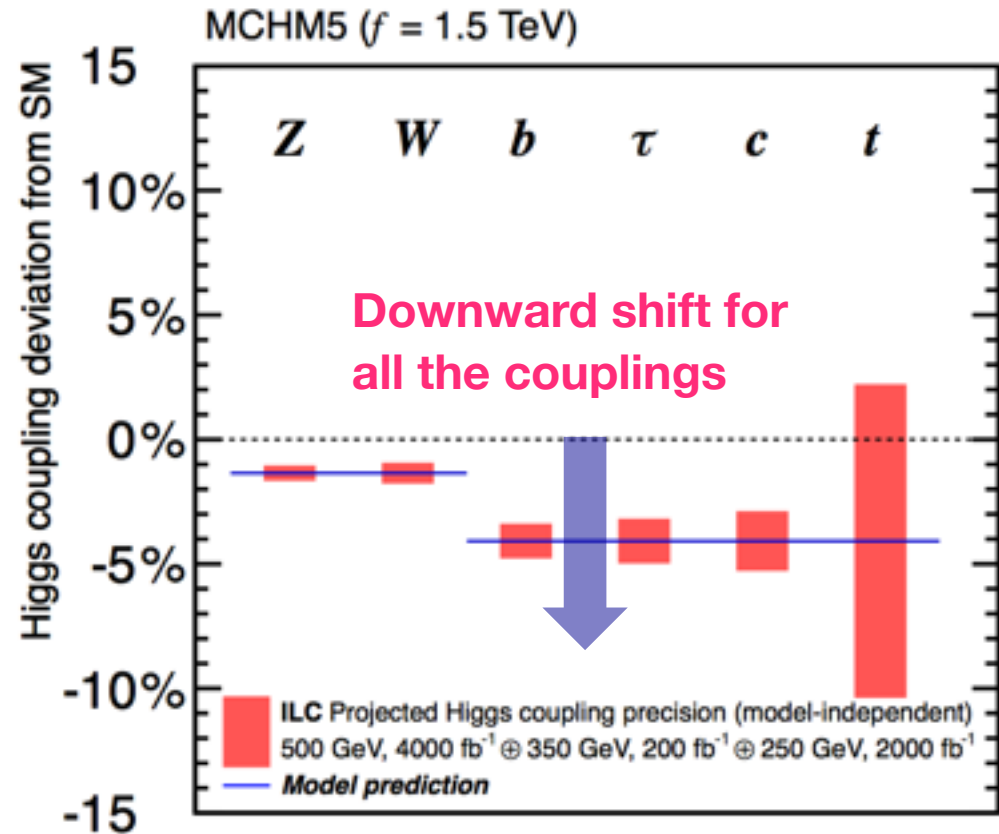
Fingerprinting

Elementary v.s. Composite?

Supersymmetry (MSSM)



Composite Higgs (MCHM5)



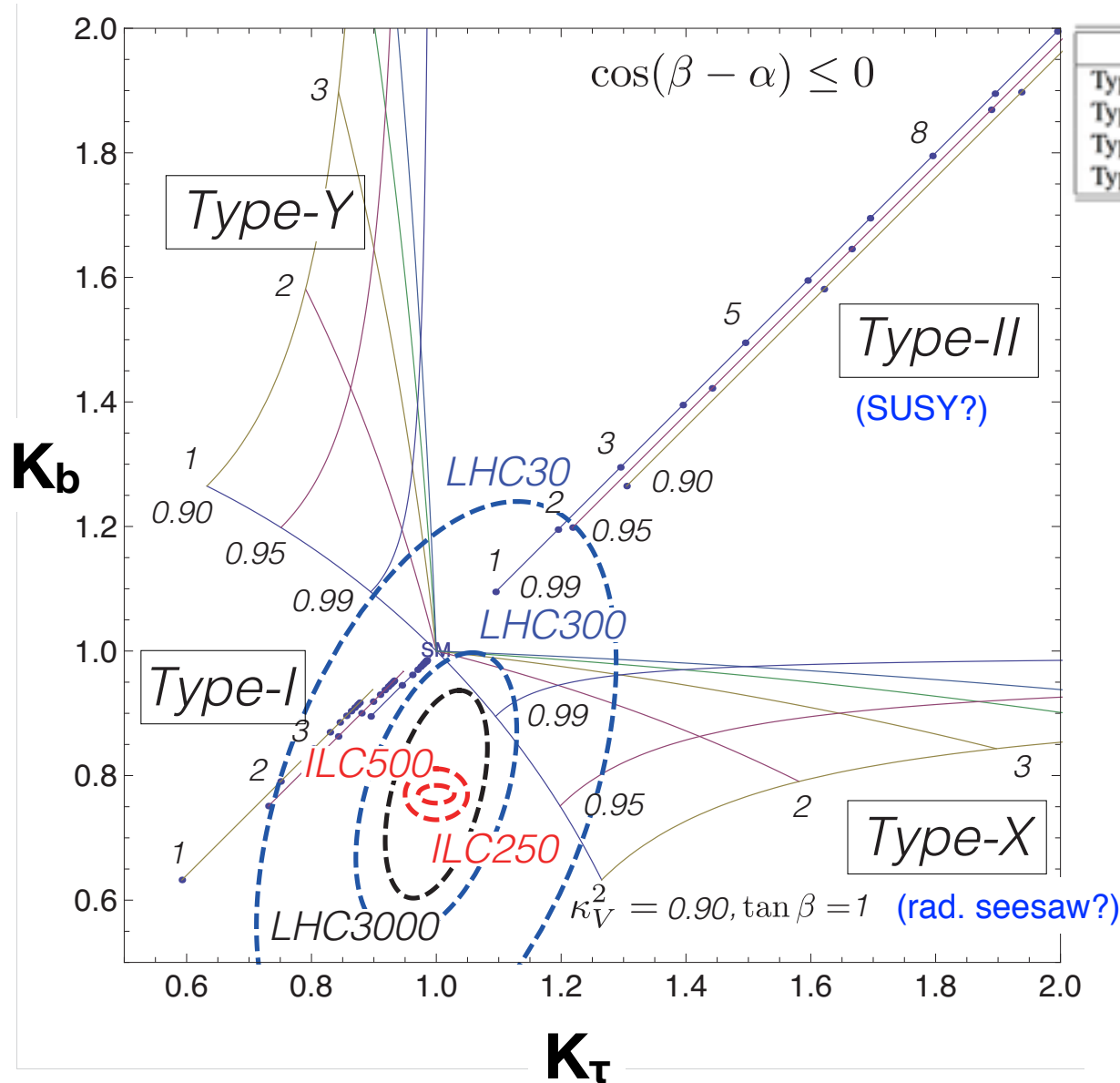
ILC 250+500 LumiUP

Complementary to direct searches at LHC: Depending on parameters, ILC's sensitivity far exceeds that of LHC!

Fingerprinting

2HDM

Multiplet Structure



	Φ_1	Φ_2	u_R	d_R	ℓ_R	Q_L, L_L
Type I	+	-	-	-	-	+
Type II (SUSY)	+	-	-	+	+	+
Type X (Lepton-specific)	+	-	-	-	+	+
Type Y (Flipped)	+	-	-	+	-	+

4 Possible Z_2 Charge Assignments that forbids tree-level Higgs-induced FCNC

$$\kappa_V^2 = \sin(\beta - \alpha)^2 = 1 \Leftrightarrow \text{SM}$$

Given a deviation of the Higgs to Z coupling: $\Delta \kappa_V^2 = 1 - \kappa_V^2 = 0.01$ we will be able to **discriminate the 4 models!**

Model-dependent
7-parameter fit
ILC: Baseline lumi.

ILC TDR

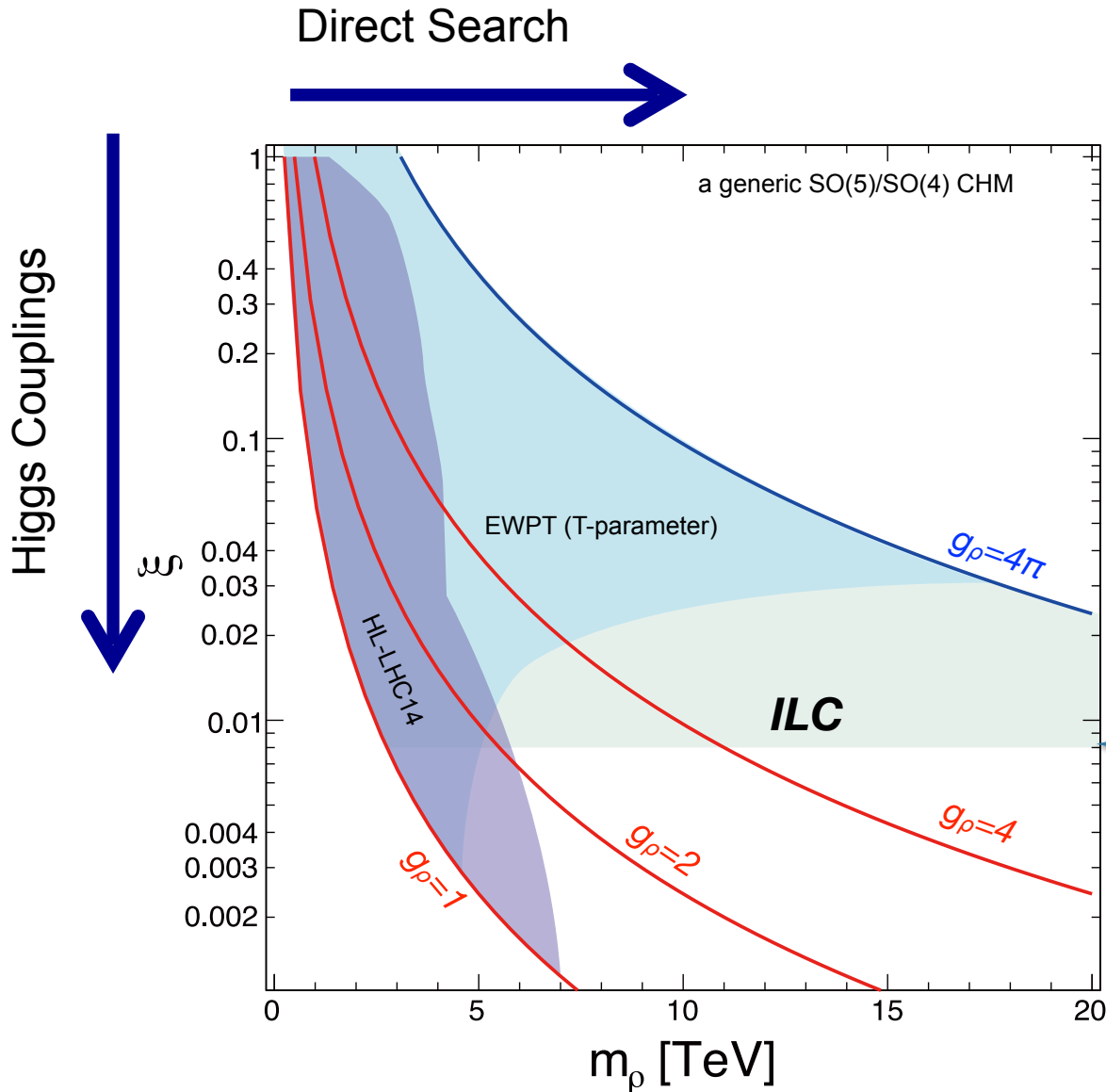
Snowmass ILC Higgs White Paper (arXiv: 1310.0763)

Kanemura et al (arXiv: 1406.3294)

Composite Higgs: Reach

Complementary approaches to probe composite Higgs models

- Direct search for heavy resonances at the LHC
 - Indirect search via Higgs couplings at the ILC
- Comparison depends on the coupling strength (g_*)



Based on Contino, et al, JHEP 1402 (2014) 006
 Torre, Thamm, Wulzer 2014
 Grojean @ LCWS 2014

$$\xi = \frac{g_\rho^2}{m_\rho^2} v^2 = \frac{v^2}{f^2}$$

$$\frac{g_{hVV}}{g_{h_{SM}VV}} = \sqrt{1 - \xi}$$

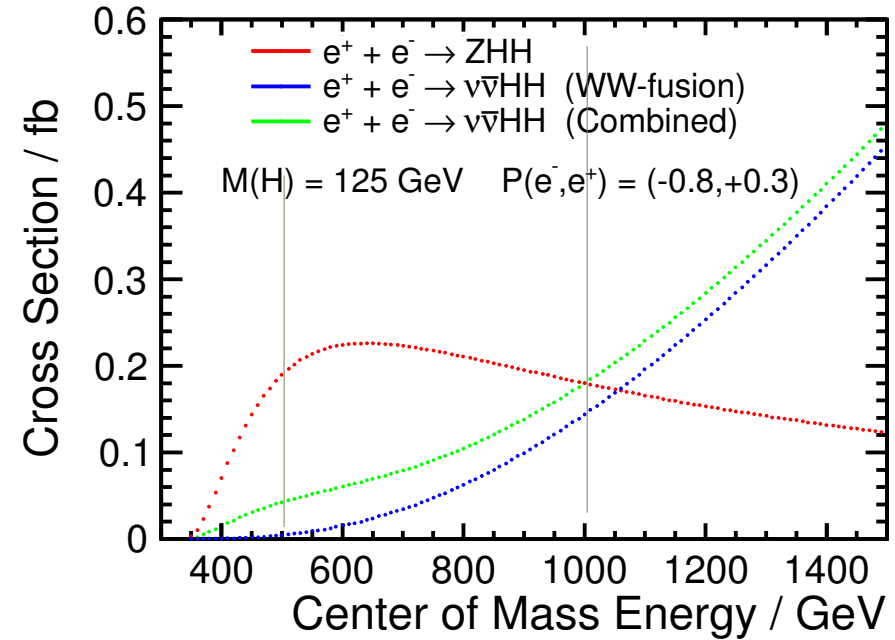
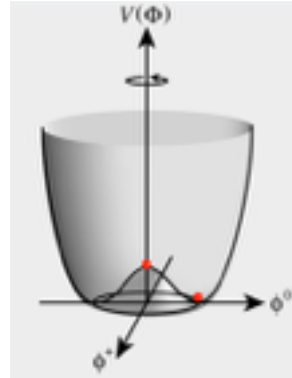
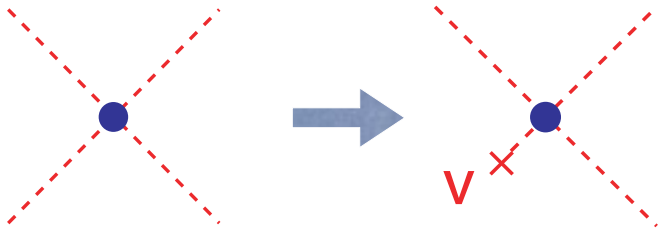
ILC (250+500 LumiUP)

$$\Delta \frac{g_{hVV}}{g_{hVV}} = 0.4\%$$

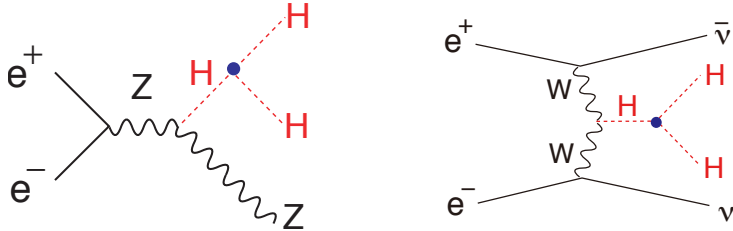
Higgs Self-coupling

Higgs Self-Coupling

The **Higgs 3-point self-coupling** is at the heart of EWSB!



There are **two ways to measure it** at ILC



arXiv:1310.0763	ILC500	ILC500-up	ILC1000	ILC1000-up
\sqrt{s} (GeV)	500	500	500/1000	500/1000
$\int \mathcal{L} dt$ (fb ⁻¹)	500	1600 [‡]	500+1000	1600+2500 [‡]
$P(e^-, e^+)$	(-0.8, 0.3)	(-0.8, 0.3)	(-0.8, 0.3/0.2)	(-0.8, 0.3/0.2)
$\sigma(ZHH)$	42.7%		42.7%	23.7%
$\sigma(\nu\bar{\nu}HH)$	-	-	26.3%	16.7%
λ	83%	46%	21%	13%

27% (H20)

Challenging even at ILC because of

- Small cross section
- **Presence of irreducible BG diagrams**

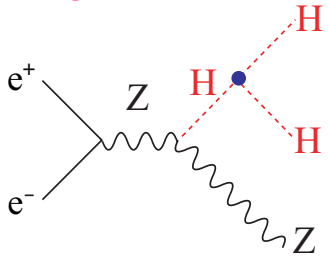
CLIC (arXiv: 1307.5288)

1.4 TeV (1.5 ab ⁻¹)	+3 TeV (2 ab ⁻¹)
21%	10%

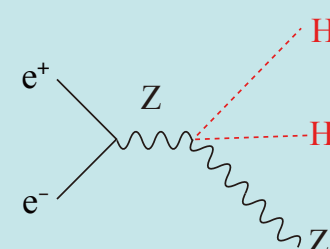
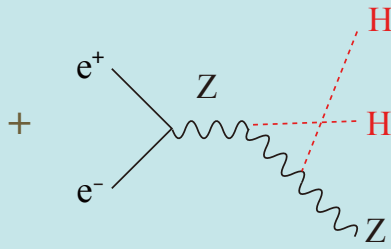
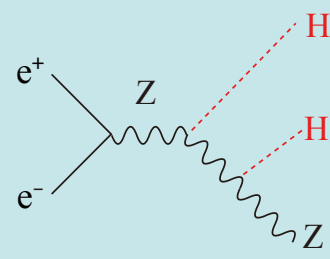
Ongoing analysis improvements **towards O(10)% measurement**

The Problem : BG diagrams dilute self-coupling contribution

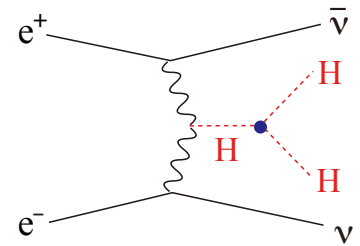
Signal diagram



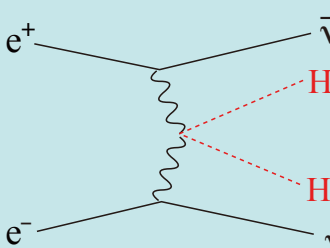
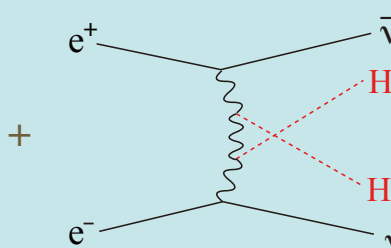
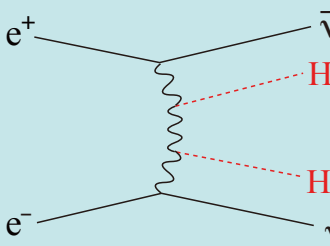
Irreducible BG diagrams



Signal diagram



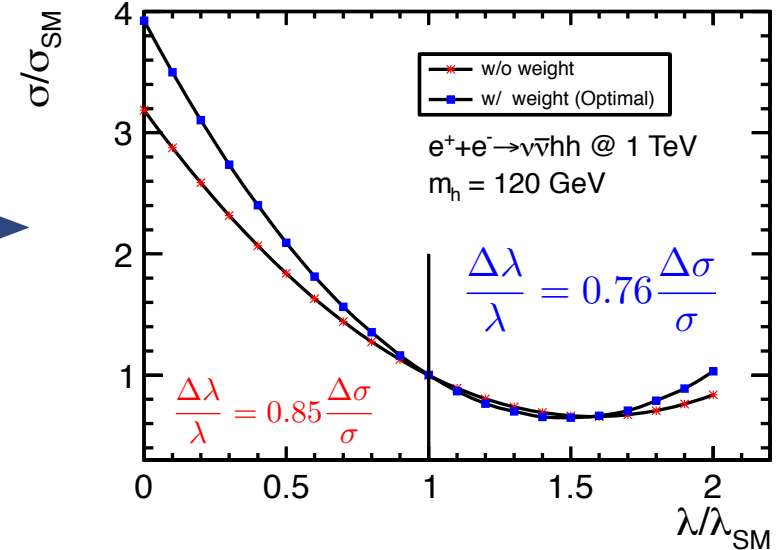
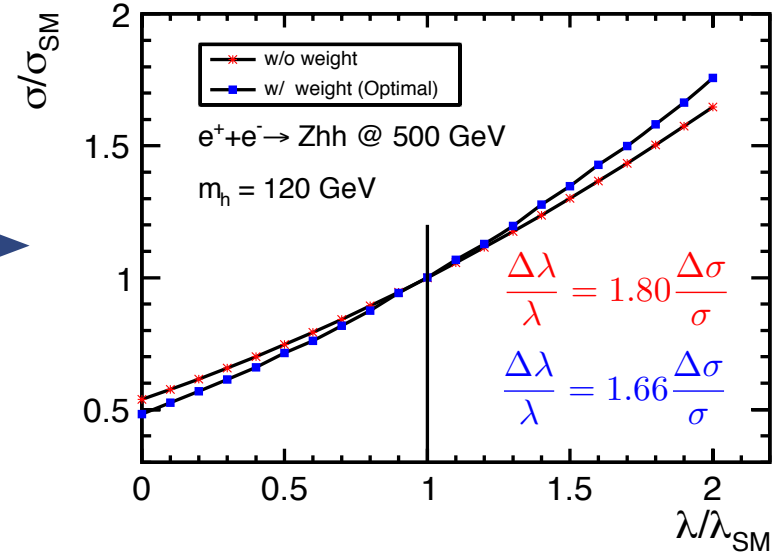
Irreducible BG diagrams



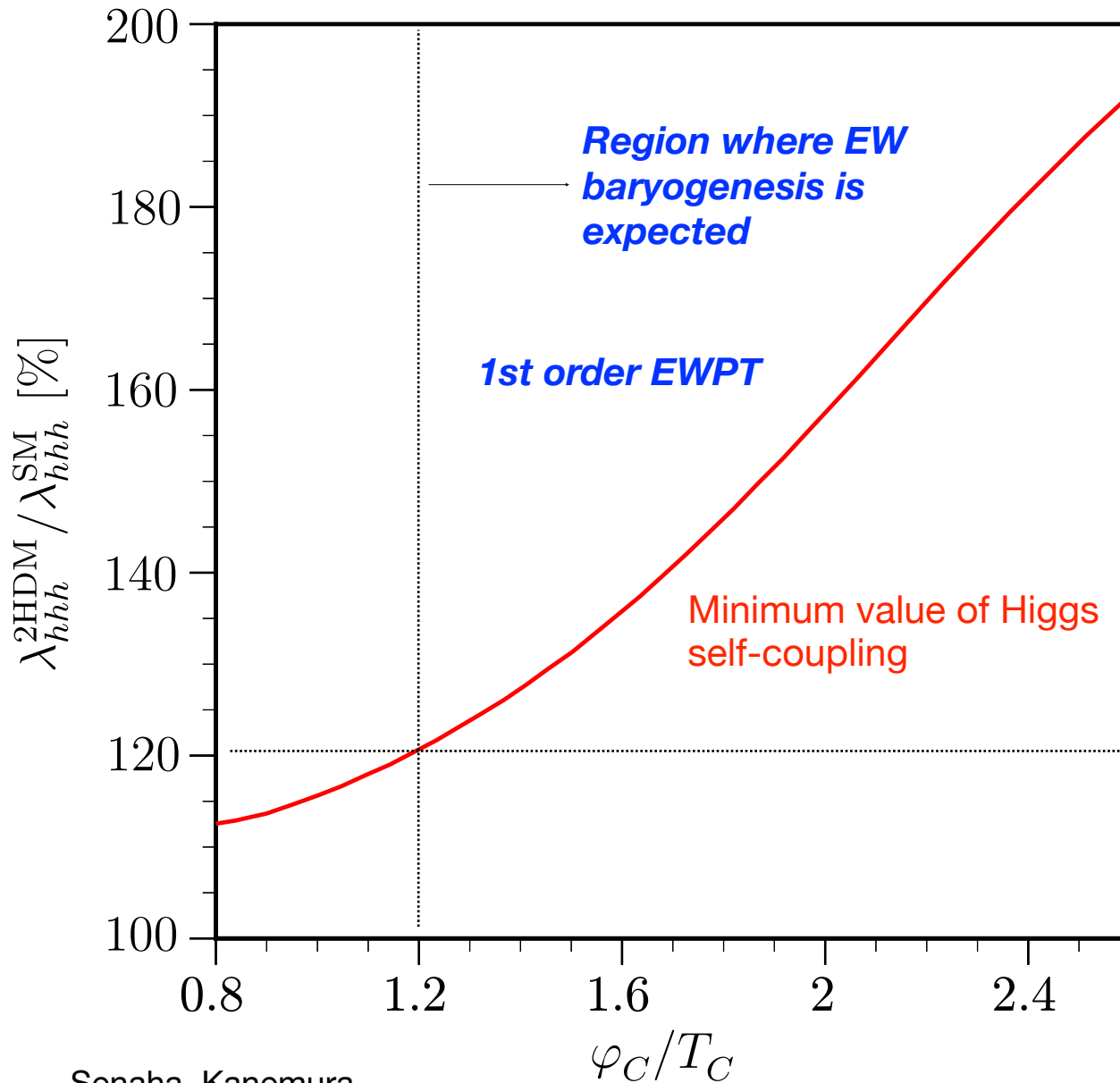
$$\sigma = \lambda^2 S + \lambda I + B$$

$$\frac{\Delta\lambda}{\lambda} = F \cdot \frac{\Delta\sigma}{\sigma}$$

$F=0.5$ if no BG diagrams



Electroweak Baryogenesis



Example:

Electroweak baryogenesis in a **Two Higgs Doublet Model**

Large deviations in Higgs self-coupling

→ **1st order EW phase transition**

→ **Out of equilibrium**

+ **CPV in Higgs sector**

→ **EW baryogenesis possible**

Constructive interference between signal and BG diagrams:

→ **if +100% deviation, then 14% precision expected on λ at 500GeV.**

ILC can address the idea of **baryogenesis occurring at the electroweak scale.**

Top

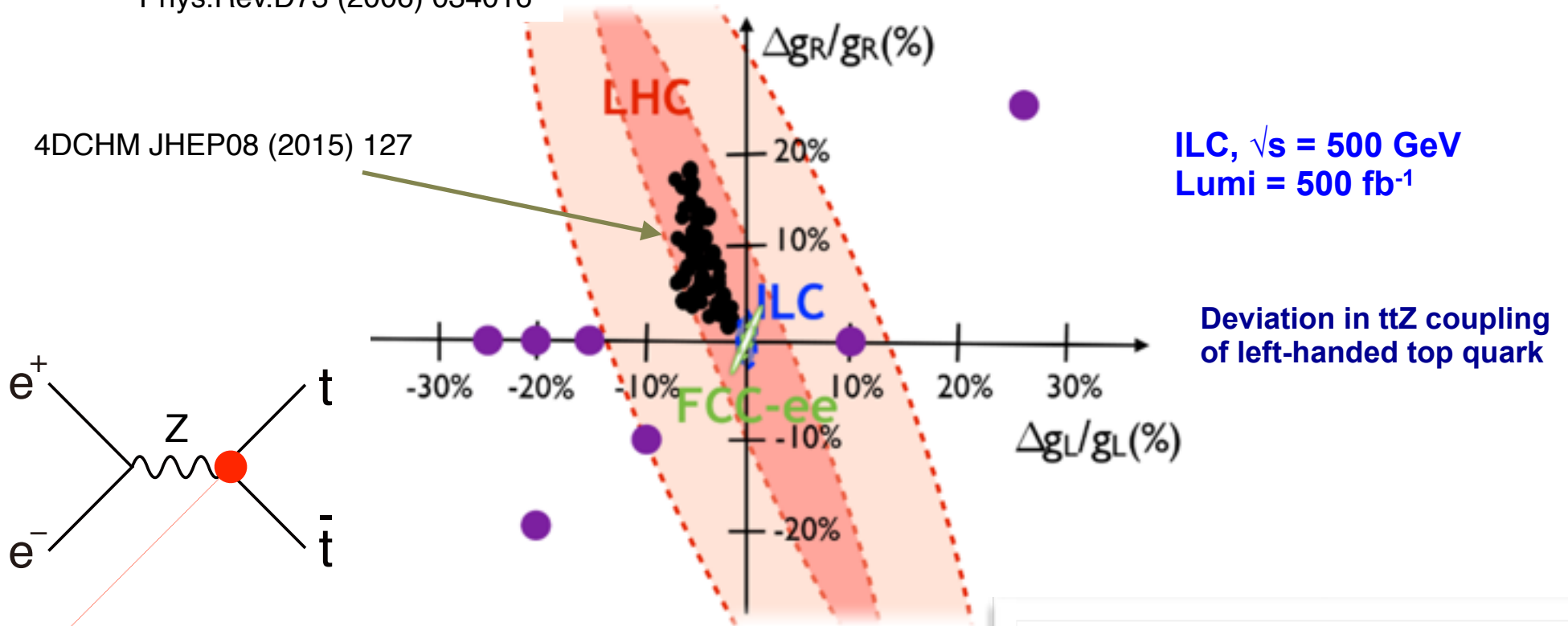
Search for Anomalous tZZ Couplings

- Top: Heaviest in SM → Must couple strongly to EWSB sector (source of $\mu^2 < 0$)!
- **Specific deviation pattern** expected in **ttZ form factors** depending on new physics.
 - **Beam polarization essential** to separate L- and R-couplings (Strength of ILC)

Deviation in ttZ coupling of right-handed top quark

Phys.Rev.D73 (2006) 034016

4DCHM JHEP08 (2015) 127



ILC, $\sqrt{s} = 500$ GeV
Lumi = 500 fb⁻¹

Deviation in ttZ coupling of left-handed top quark

$$\Gamma_{\mu}^{ttX}(k^2, q, \bar{q}) = ie \left\{ \gamma_{\mu} \left(\tilde{F}_{1V}^X(k^2) + \gamma_5 \tilde{F}_{1A}^X(k^2) \right) + \frac{(q - \bar{q})_{\mu}}{2m_t} \left(\tilde{F}_{2V}^X(k^2) + \gamma_5 \tilde{F}_{2A}^X(k^2) \right) \right\}$$

Purple points: deviations expected for various new physics models (new physics scale ~ 1 TeV) compiled in arXiv:1505.06020

ILC is sensitive to M_{KK} up to ~ 25 TeV for typical RS scenarios (even up to ~ 80 TeV in extreme cases)!

**What if no deviation from
the SM would be seen?**

Clarify the Range of Validity of SM

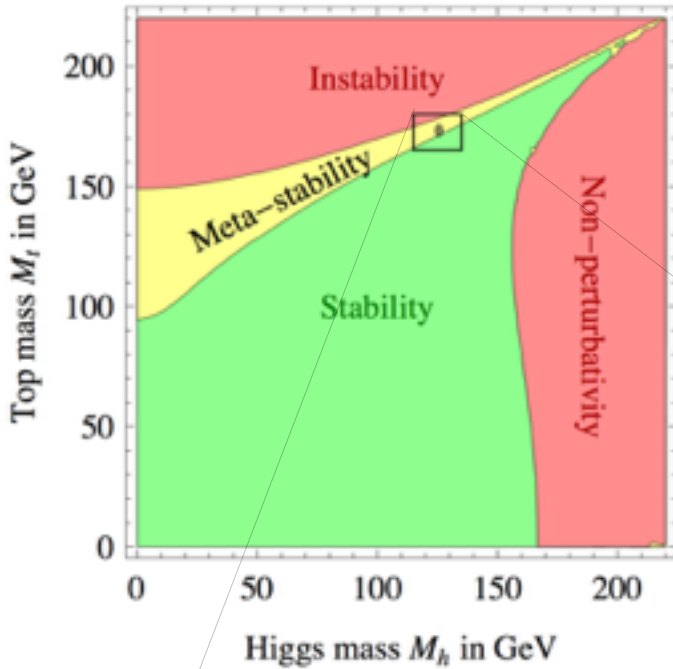
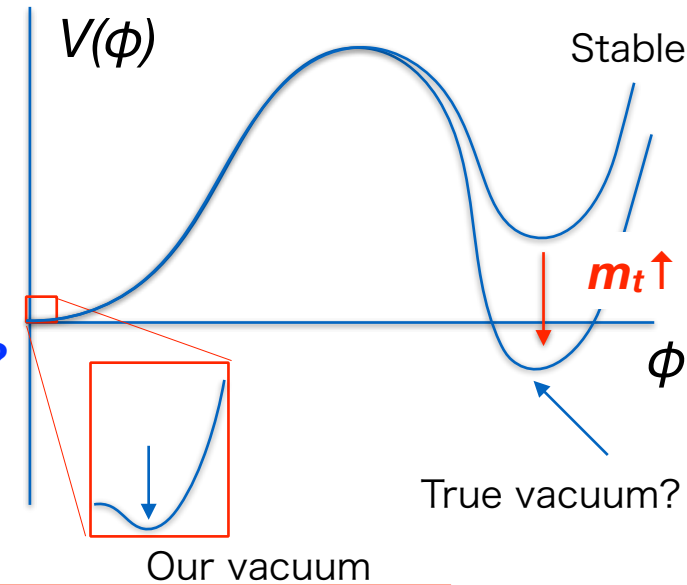
Stability of SM Vacuum

Top Yukawa coupling drives the 4-point Higgs coupling (λ) to negative!

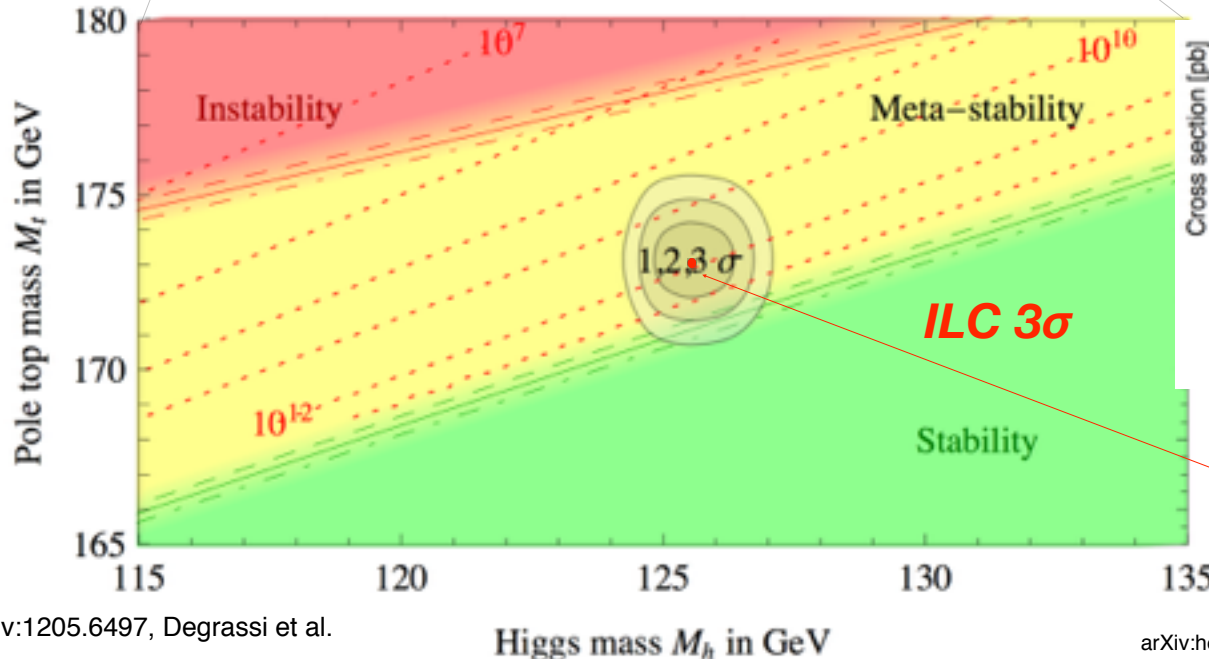
The current values of m_t and m_h :
Subtle point of meta-stability!

Does λ go to negative below Λ_P ?
 or $\lambda(\Lambda_P) = 0$?

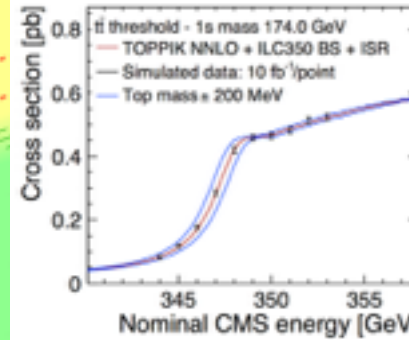
To answer this, we need
precision m_t measurement!



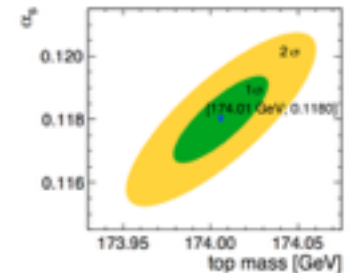
At LHC, theory error limits the precision to ~500 MeV.



TTbar Threshold Scan @ILC



Theoretically very clean measurement of m_t



$\Delta m_t(\overline{MS}) \lesssim 50 \text{ MeV}$
 $\Delta m_H = 30 \text{ MeV}$
ILC pinpoints the vacuum location

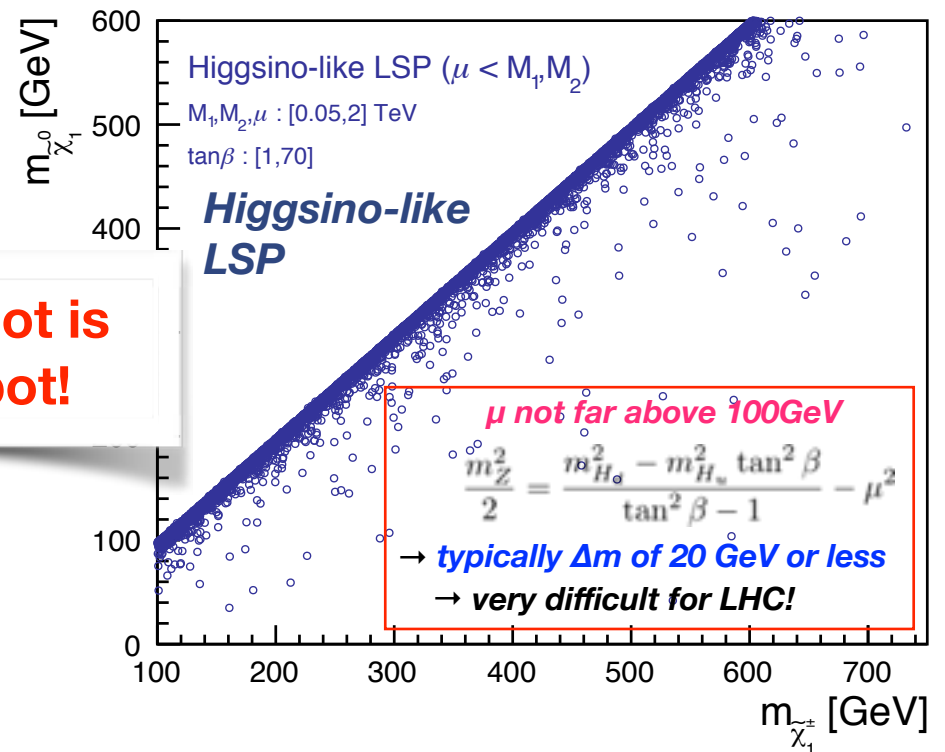
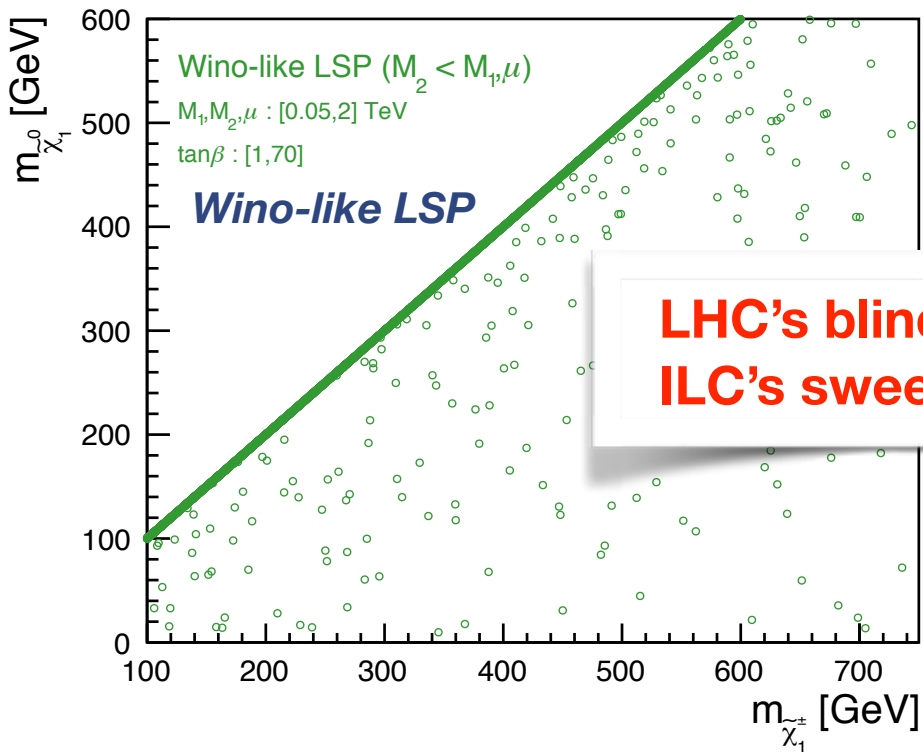
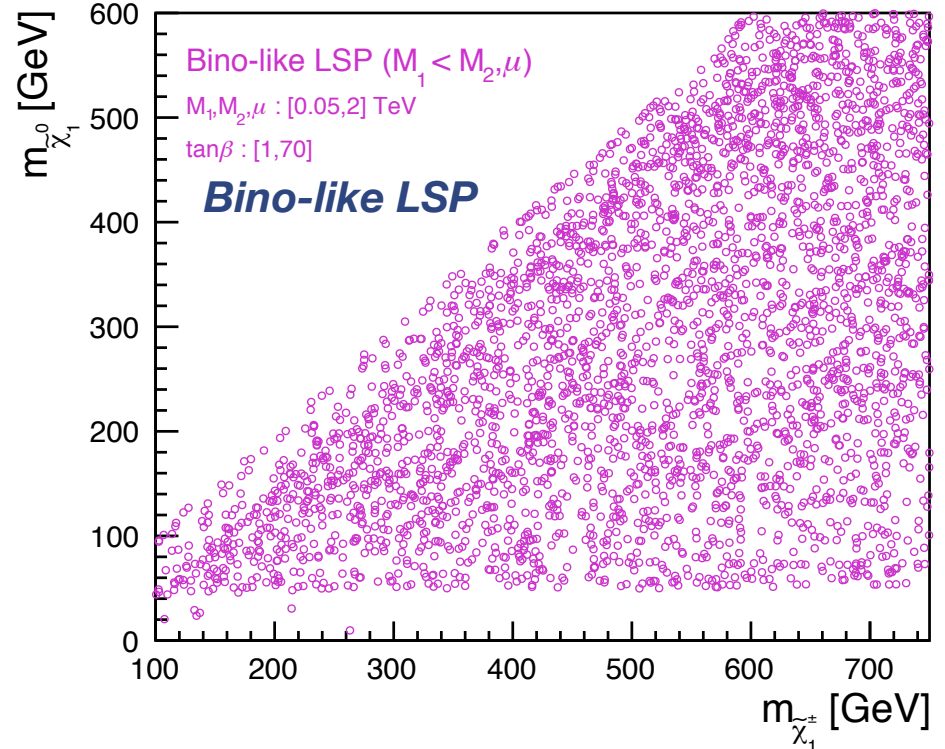
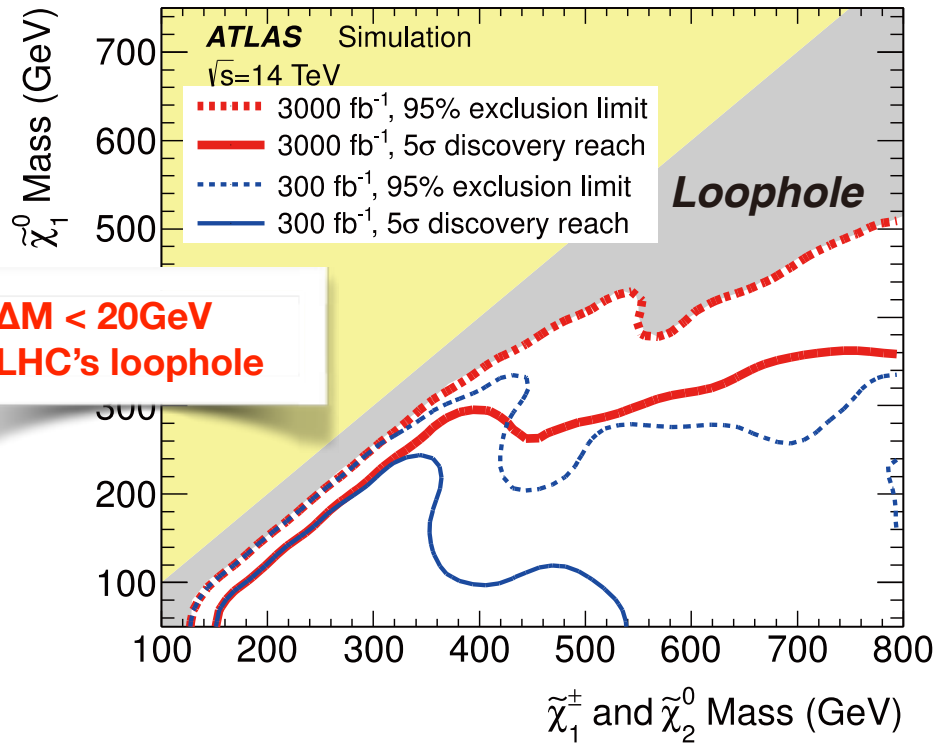
Direct Searches
for
New Particles

ILC, too, is an energy frontier machine!

*It will enter **uncharted waters of e^+e^- collisions***

Thanks to well-defined initial states,
clean environment w/o QCD BG, and polarized beams
ILC can cover blind spots of LHC

Chargino Search



LHC's blind spot is ILC's sweet spot!

Higgsinos in Natural SUSY ($\Delta M < \text{a few GeV}$)

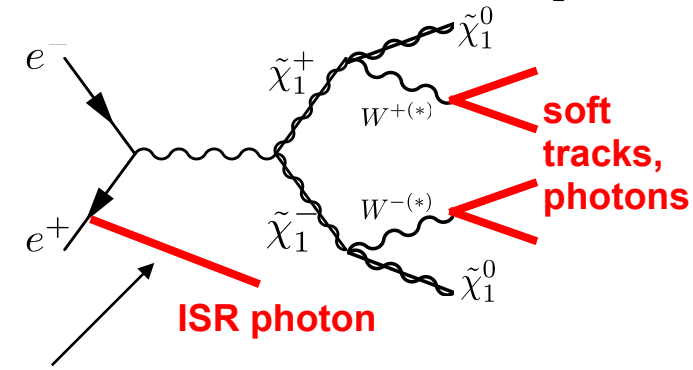
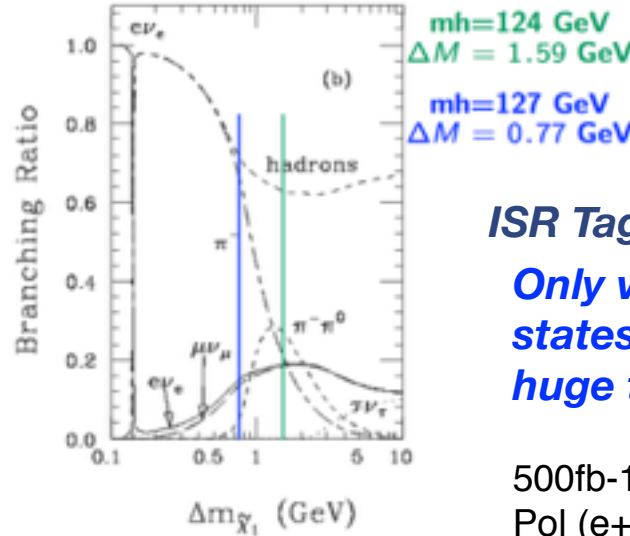
ISR Tagging

$$e^+e^- \rightarrow \tilde{\chi}_1^+ \tilde{\chi}_1^- \gamma$$

$$e^+e^- \rightarrow \tilde{\chi}_2^0 \tilde{\chi}_1^0 \gamma$$

ILC as a Higgsino Factory

Ref: C.-H. Chen et al. hep-ph:9512230

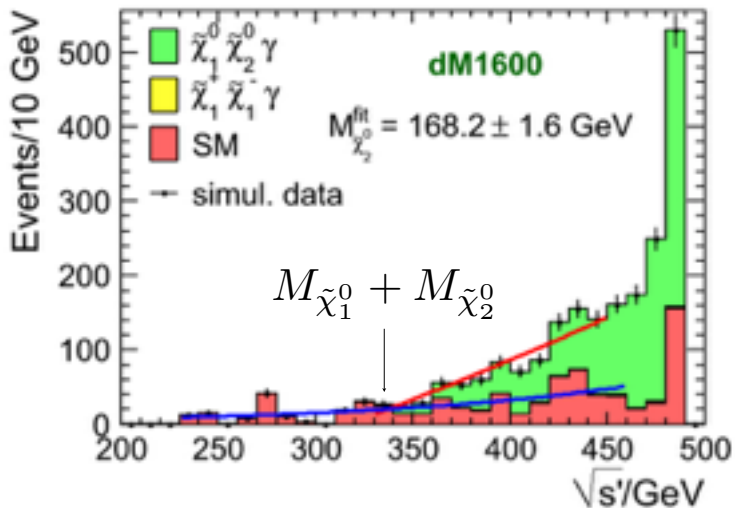
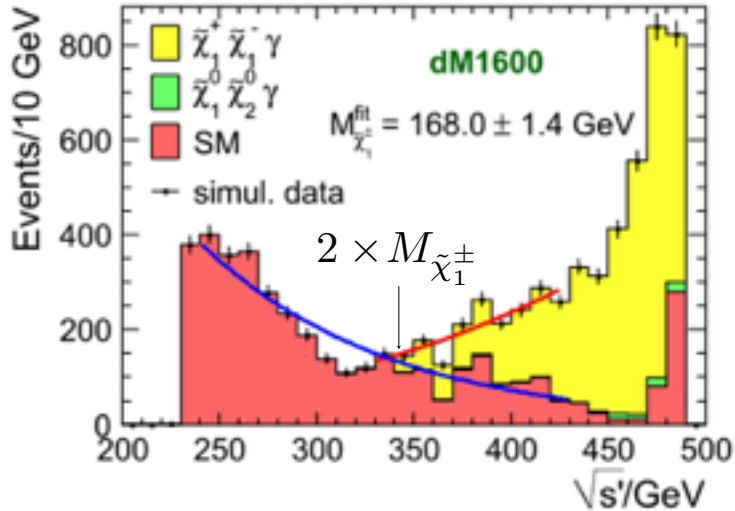


ISR Tagging

Only very soft particles in the final states → Require a hard ISR to kill huge two-photon BG!

500fb-1 @ Ecm=500GeV

Pol (e+,e-) = (+0.3,-0.8) and (-0.3,+0.8)



EPJC (2013) 73:2660

dm1600

Mass Spectrum	
Particle	Mass (GeV)
h	124
$\tilde{\chi}_1^0$	164.17
$\tilde{\chi}_1^\pm$	165.77
$\tilde{\chi}_2^0$	166.87
H^\pm 's	$\sim 10^3$
$\tilde{\chi}^\pm$'s	$\sim 2 - 3 \times 10^3$

$\Delta M(\tilde{\chi}_1^\pm, \tilde{\chi}_1^0) = 1.59 \text{ GeV}$

$$\delta(\sigma \times BR) \simeq 3\%$$

$$\delta M_{\tilde{\chi}_1^\pm}(M_{\tilde{\chi}_1^0}) \simeq 2.1(3.7) \text{ GeV}$$

$$\delta \Delta M(\tilde{\chi}_1^\pm, \tilde{\chi}_1^0) \simeq 70 \text{ MeV}$$

dm770

Mass Spectrum	
Particle	Mass (GeV)
h	127
$\tilde{\chi}_1^0$	166.59
$\tilde{\chi}_1^\pm$	167.36
$\tilde{\chi}_2^0$	167.63
H^\pm 's	$\sim 10^3$
$\tilde{\chi}^\pm$'s	$\sim 2 - 3 \times 10^3$

$\Delta M(\tilde{\chi}_1^\pm, \tilde{\chi}_1^0) = 0.77 \text{ GeV}$

$$\delta(\sigma \times BR) \simeq 1.5\%$$

$$\delta M_{\tilde{\chi}_1^\pm}(M_{\tilde{\chi}_1^0}) \simeq 1.5(1.6) \text{ GeV}$$

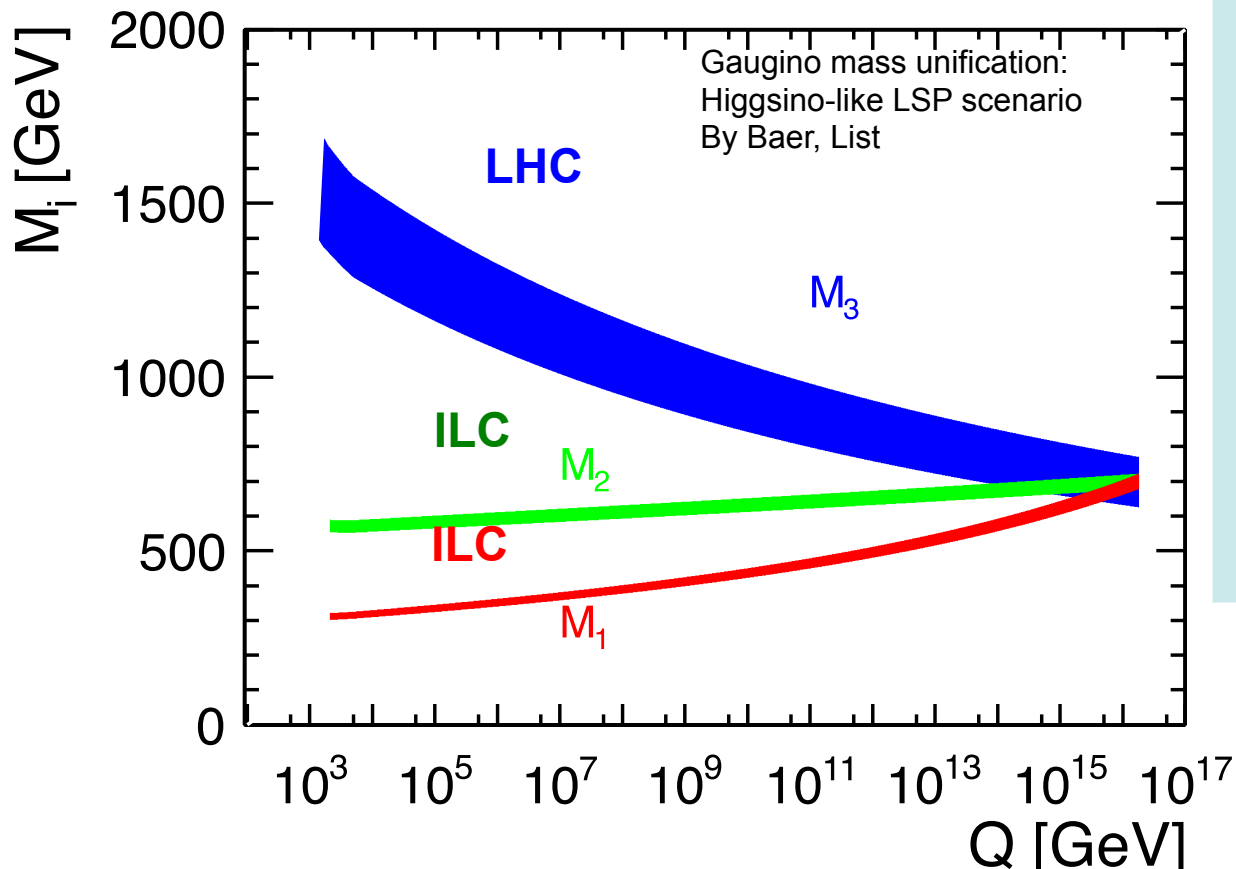
$$\delta \Delta M(\tilde{\chi}_1^\pm, \tilde{\chi}_1^0) \simeq 20 \text{ MeV}$$

GUT Scale Physics

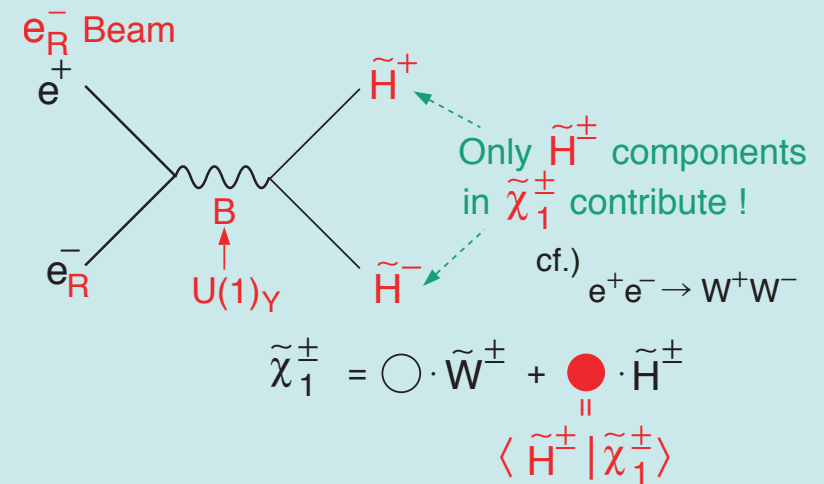
If we are lucky and the gluino is in LHC's mass reach and the lighter chargino and the neutralinos are in ILC's mass reach, *we will be able to test the gaugino mass unification!*

LHC: gluino discovery
 → mass determination

ILC: Higgsino-like EWkino discovery
 → M_1, M_2 via mixing between Higgsino and Bino/Wino



Chargino decomposition



Beam polarization is essential to decompose the EWkinos to bino, wino, and higgsino and extract M_1 and M_2 .

WIMP Dark Matter Search @ ILC

Weakly Interacting Massive Particle

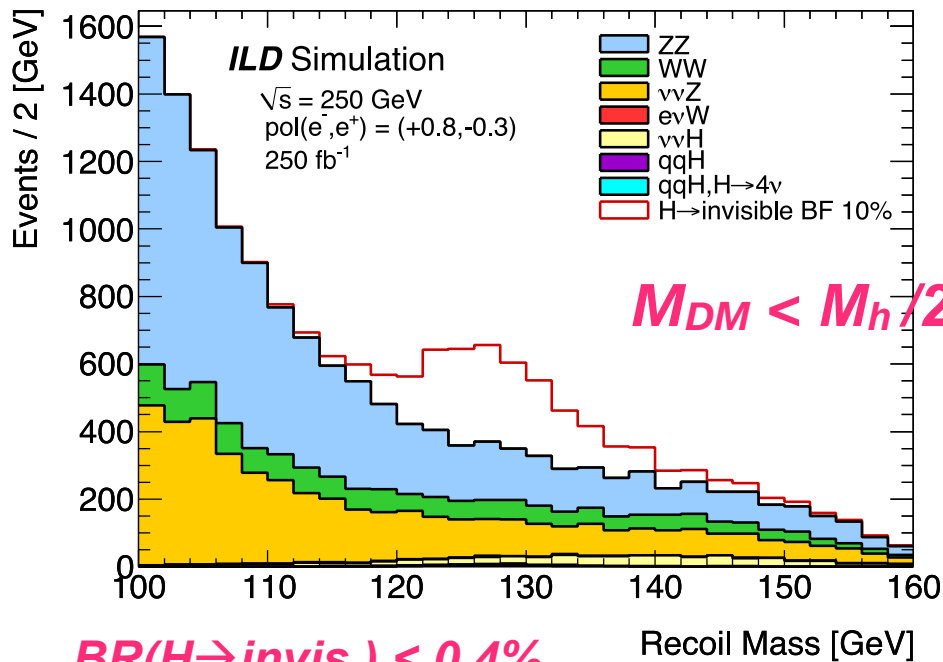
Decay of a new particle to Dark Matter (DM)

DM has a charged partner in many new physics models.

SUSY: The Lightest SUSY Particle (LSP) = DM → Its partner decays to a DM.

- Events with missing Pt (example: light chargino: see the previous page)

Higgs Invisible Decay

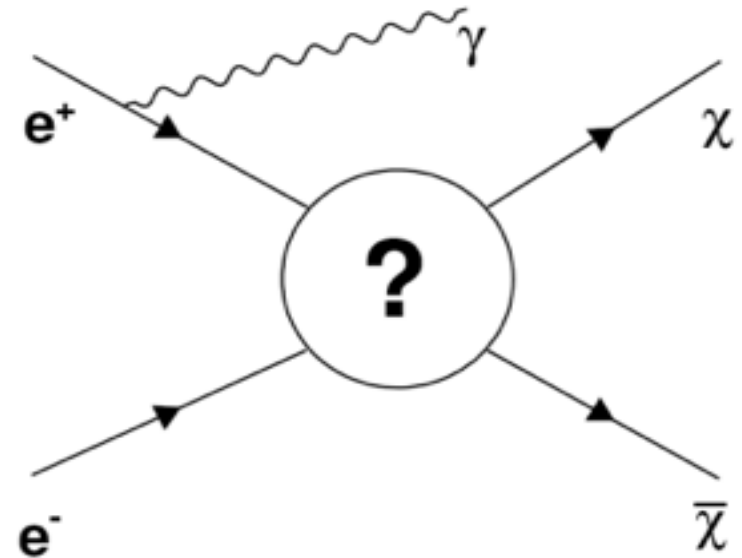


$BR(H \rightarrow \text{invis.}) < 0.4\%$

at 250 GeV, 1150 fb⁻¹ (<0.3% at 95%CL: H20)

Possible to access BR_{inv} to 0.4%!

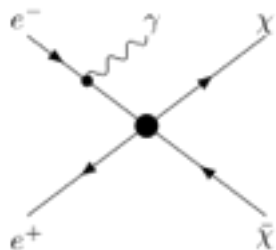
Mono-photon Search



→ $M_{DM} \text{ reach} \sim E_{cm}/2$

Possible to access DM to $\sim E_{cm}/2$!

DM: Effective Operator Approach



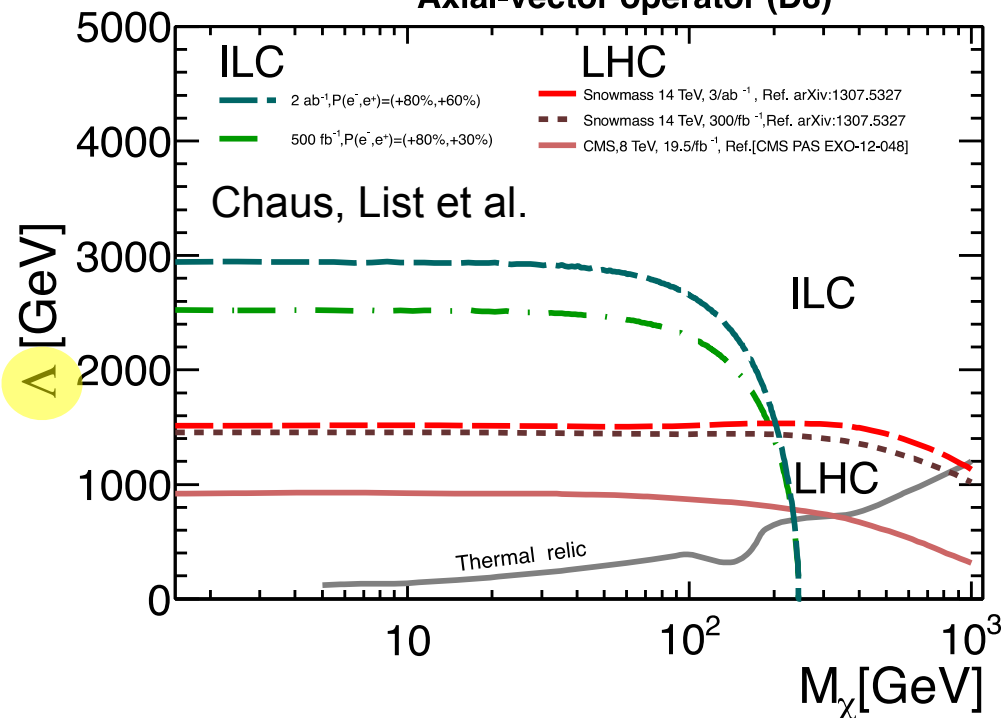
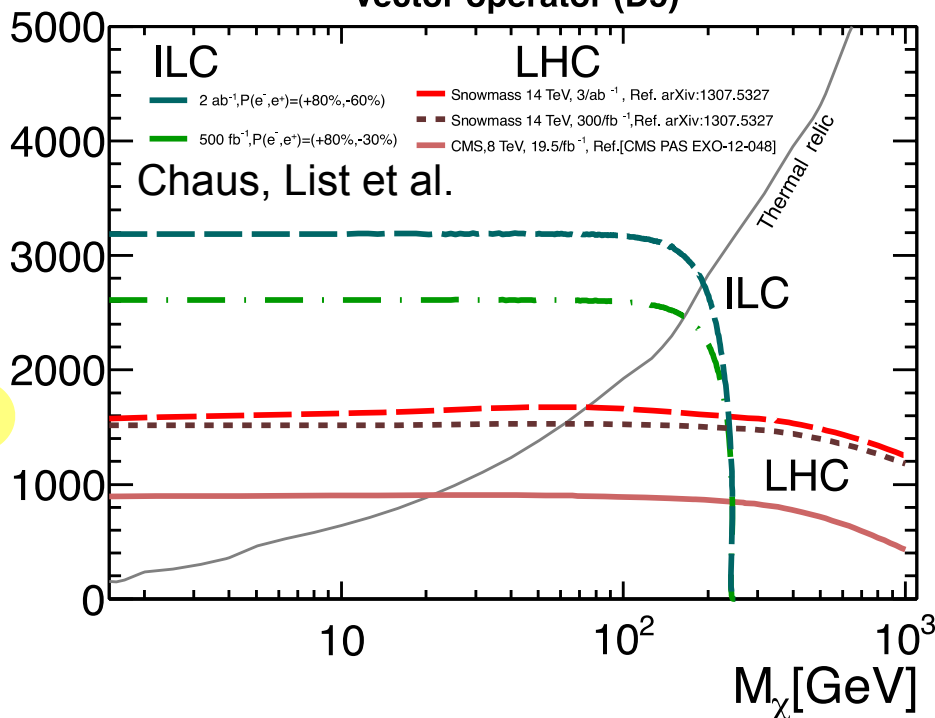
$$\mathcal{L}_{\text{int}} = \frac{1}{\Lambda^2} \mathcal{O}_i$$

$$\mathcal{O}_V = (\bar{\chi} \gamma_\mu \chi) (\bar{l} \gamma^\mu l)$$

Vector operator (D5)

$$\mathcal{O}_A = (\bar{\chi} \gamma_\mu \gamma_5 \chi) (\bar{l} \gamma^\mu \gamma^5 l)$$

Axial-vector operator (D8)



LHC sensitivity: Mediator mass up to $\Lambda \sim 1.5$ TeV for **large DM mass**

ILC sensitivity: Mediator mass up to $\Lambda \sim 3$ TeV for **DM mass up to $\sim \sqrt{s}/2$**



LHC-ILC synergy!

Summary

- The primary goal for the next decades is ***to uncover the secret of the EW symmetry breaking***. The discovery of H(125) completed the SM particle spectrum and taught us how the EW symmetry was broken. However, it does not tell us why it was broken. ***Why $\mu^2 < 0$?*** To address this question we need to go beyond the SM.
- There is a big branching point concerning the question: ***Is H(125) elementary or composite?*** There are ***two powerful probes*** in hand: ***H(125) itself and the top quark***. Different models predict different deviation patterns in Higgs and top couplings. ***ILC will measure these couplings with unprecedented precision***.
- This will open up ***a window to BSM*** and ***fingerprint BSM models***, otherwise will ***set the energy scale for the E-frontier machine that will follow LHC and ILC***.
- ***Cubic self-coupling measurement*** will decide whether the EWSB was strong 1st order phase transition or not. If it was, it will provide us the possibility of understanding ***baryogenesis at the EW scale***.
- ***The ILC is an ideal machine to answer these questions*** (regardless of BSM scenarios) and we can do this ***model-independently***.
- It is also very important to stress that ***ILC, too, is an energy frontier machine***. It will ***access the energy region never explored with any lepton collider***. It is not a tiny corner of the parameter space that will be left after LHC. ***There is a wide and interesting region for ILC to explore (eg. Natural SUSY)***.
- Once a new particle is found at ILC, we can precisely determine its properties, making full use of ***polarized beams***. In the case of natural radiative SUSY scenario, we might even probe GUT scale physics using RGE.
- ***In this way, ILC will pave the way to BSM physics***.

Backup

I strongly believe that **ILC is worth building regardless of what LHC is going to discover.**

But the MEXT ILC Advisory Panel recommended **to closely monitor, analyze, and examine the development of LHC experiments.**

X750

Our (=LCC Physics WG) Attitudes towards X750

1. It's too early to get excited,
2. but if it is real, it is ***a good example of case 3*** in the ICFA letter to MEXT's ILC Advisory Panel:
case 3: LHC discovers relatively heavy new particles (which cannot be directly produced at the 500 GeV ILC)
3. Since the MEXT Panel recommended to ***closely monitor, analyze, and examine the development of LHC experiments***, this is ***a good opportunity to do exercise for case 3***. → motivation for this note
4. In LCC's letter to the panel, it is stated that "***While performing precision studies of the Higgs boson and the top quark, we will prepare for the energy upgrade of the ILC taking advantage of energy expandability enabled by its linear shape.***"
5. *The note is intended to show*
 1. ***The 500 GeV ILC has a lot to say about X750 through precision measurements plus possible discovery of NPs associated with X750.***
 2. ***Possible energy upgrade with PLC option will open up even greater opportunities to uncover the new physics operating behind X750 together with LHC.***

Implications of the 750 GeV $\gamma\gamma$ Resonance as a Case Study for the International Linear Collider

LCC PHYSICS WORKING GROUP

KEISUKE FUJII¹, CHRISTOPHE GROJEAN^{2,3} MICHAEL E. PESKIN⁴(CONVENERS);
TIM BARKLOW⁴, YUANNING GAO⁵, SHINYA KANEMURA⁶, HYUNGDO KIM⁷,
JENNY LIST², MIHOKO NOJIRI^{1,8}, MAXIM PERELSTEIN⁹, ROMAN PÖSCHL¹⁰,
JÜRGEN REUTER², FRANK SIMON¹¹, TOMOHIKO TANABE¹², JAEHOON YU¹³,
JAMES D. WELLS¹⁴; ADAM FALKOWSKI¹⁵, SHIGEKI MATSUMOTO⁸,
TAKEO MOROI¹⁶, FRANCOIS RICHARD¹⁰, JUNPING TIAN¹², MARCEL VOS¹⁷,
HIROSHI YOKOYA¹⁸; HITOSHI MURAYAMA^{8,19,20}, HITOSHI YAMAMOTO²¹

ABSTRACT

If the $\gamma\gamma$ resonance at 750 GeV suggested by 2015 LHC data turns out to be a real effect, what are the implications for the physics case and upgrade path of the International Linear Collider? Whether or not the resonance is confirmed, this question provides an interesting case study testing the robustness of the ILC physics case. In this note, we address this question with two points: (1) Almost all models proposed for the new 750 GeV particle require additional new particles with electroweak couplings. The key elements of the 500 GeV ILC physics program—precision measurements of the Higgs boson, the top quark, and 4-fermion interactions—will powerfully discriminate among these models. This information will be important in conjunction with new LHC data, or alone, if the new particles accompanying the 750 GeV resonance are beyond the mass reach of the LHC. (2) Over a longer term, the energy upgrade of the ILC to 1 TeV already discussed in the ILC TDR will enable experiments in $\gamma\gamma$ and e^+e^- collisions to directly produce and study the 750 GeV particle from these unique initial states.

In this note X750 is called Φ

Questions addressed in the note

1. If Φ (=X750) is real, what would the implications be for the program of the ILC?
2. Will the ILC be able to shed light on this resonance or on accompanying new physics?

Caution

*It might turn out that **the Φ is a relatively minor player** in a new sector of physics that the LHC will begin to uncover in the next few years.*

For this reason, ***it is premature to discuss a new accelerator intended specifically to target the Φ*** or any other new particle that turns up in the early 13 TeV LHC data.

Properties of X750 (hereafter called Φ)

$\Phi \rightarrow \gamma\gamma$ means $\Phi = \text{color singlet}$ with $J \neq 1$

we assume $J = 0$

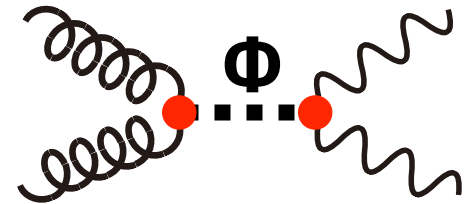
Φ is seen at 13 TeV but is much less apparent at 8 TeV

	gg	$b\bar{b}$	$d\bar{d}$	$u\bar{u}$	$\gamma\gamma$
$r_{p\bar{p}}$	4.8	5.7	2.7	2.6	1.9

$\mathcal{L}_{13\text{TeV}}/\mathcal{L}_{8\text{TeV}}$ prefers production via gg-fusion (or bb annihilation)

Assume production via gg:

$$\sigma(pp \rightarrow \Phi \rightarrow \gamma\gamma) = 5 \text{ fb} \longrightarrow \Gamma(\Phi \rightarrow \gamma\gamma) \geq 0.5 \text{ MeV}$$



Effective Lagrangian

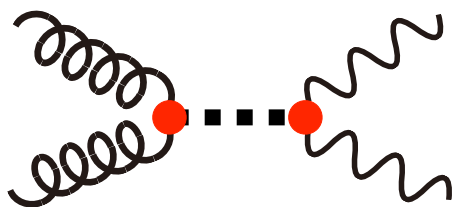
Underlying BSM physics must respect $SU(3) \times SU(2) \times U(1)$

$\rightarrow \gamma$ must be a mixture of neutral $SU(2)$ and $U(1)$ gauge bosons

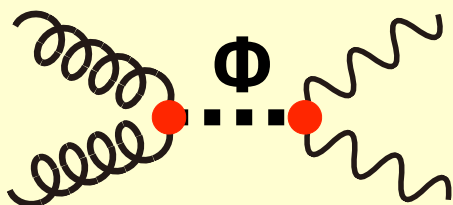
$$\mathcal{L} = \frac{\alpha_s}{4} A_3 \Phi G_{\mu\nu} G^{\mu\nu} + \frac{\alpha_w}{4} A_2 \Phi W_{\mu\nu} W^{\mu\nu} + \frac{\alpha'}{4} A_1 \Phi B_{\mu\nu} B^{\mu\nu}$$

Typical Models and Effects

Effective Couplings



Φ =RS radion

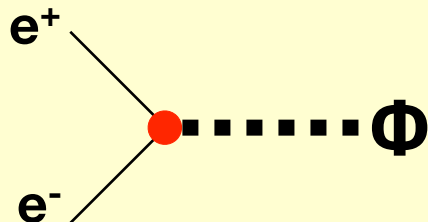


KK-loop correction
→ hWW, hZZ

~8% deviation expected for 5 TeV
KK gluon.

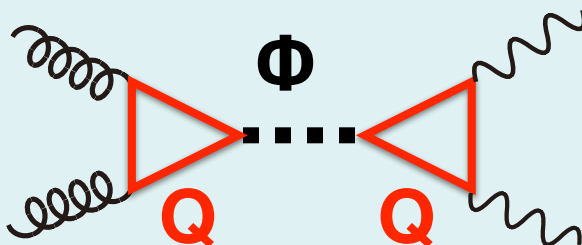
Φ =RS graviton

$J=2$

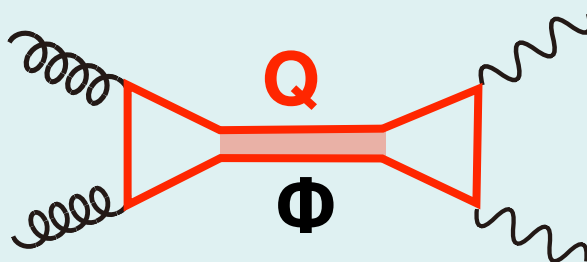
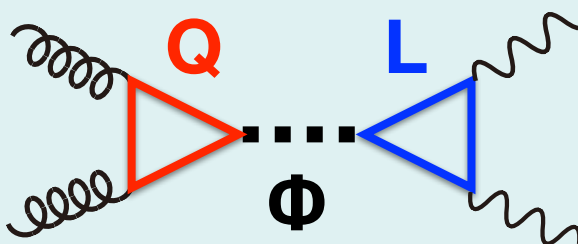


direct coupling to e^+e^-
→ s -channel Φ production
still not completely excluded.

Elementary Scalar



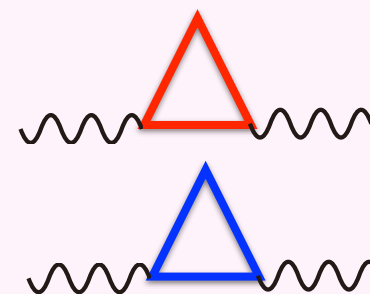
New vector-like fermions
in the loops



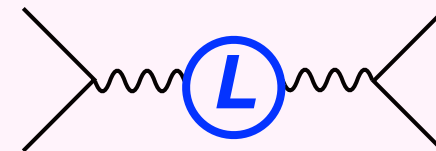
Resonance/pNGB

might be accompanied by
DM within ILC's reach

Oblique Corrections



→ 2-to-2 processes



with $\delta\sigma/\sigma=0.1\%$, ILC
sensitivity exceeds LHC

Mixings

Q-t mixing → ttZ

L- τ mixing

Φ -h mixing

→ $h\gamma\gamma, hgg$

→ hWW, hZZ

a few % deviation expected
→ well within H20 target

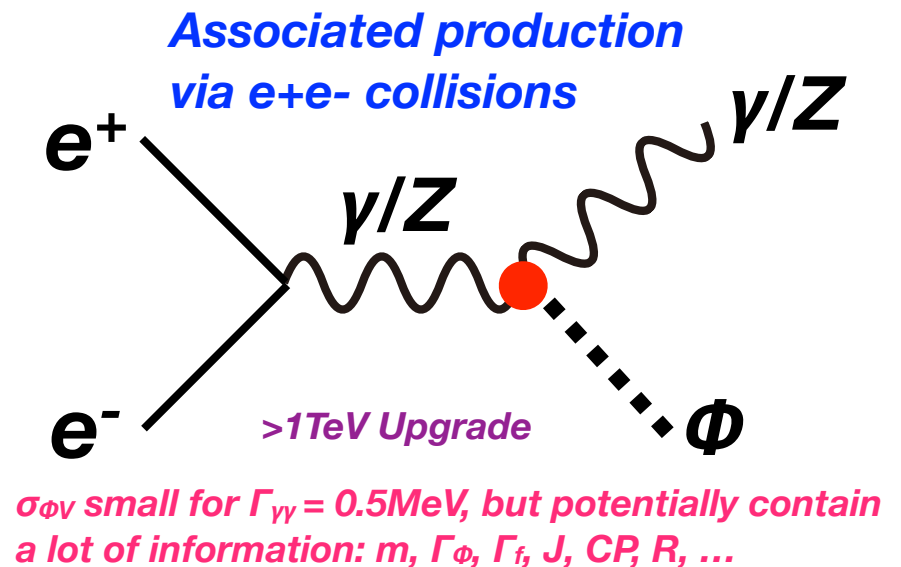
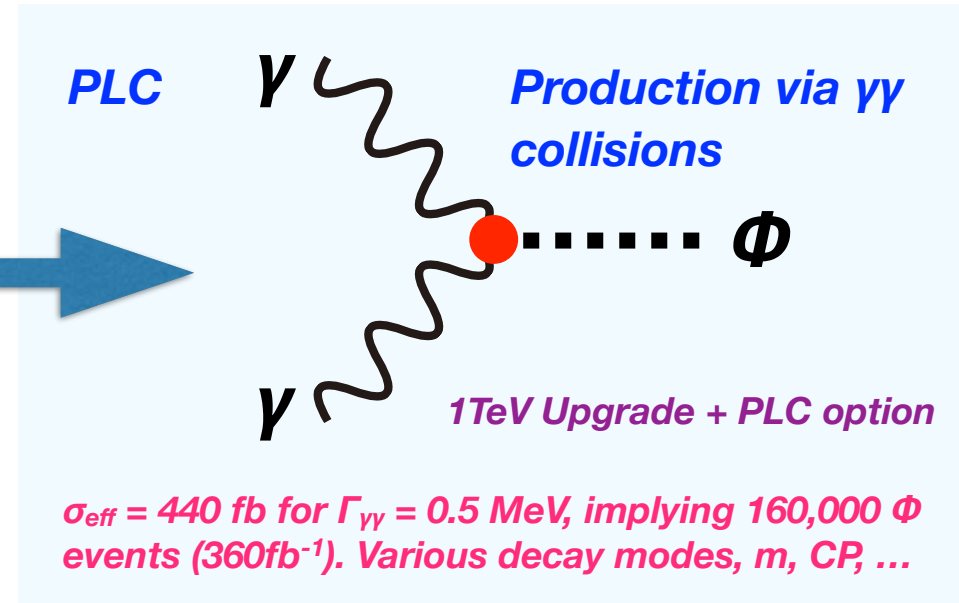
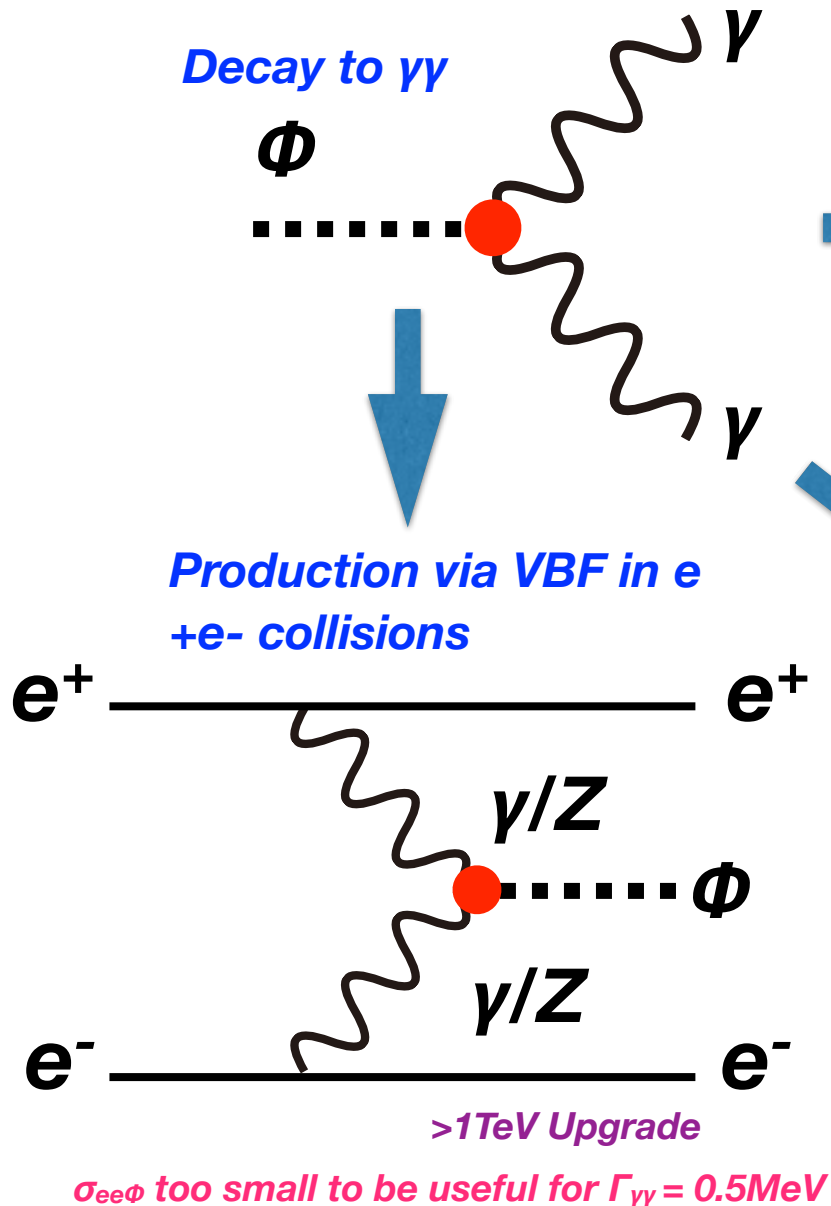
	hWW hZZ	$hb\bar{b}$ $h\tau\tau$	$h\gamma\gamma$ hgg	$ht\bar{t}$	$h \rightarrow$ invis.	$h\tau\mu$	$t\bar{t}Z$	$ee \rightarrow$ $ee, \mu\mu$	$ee \rightarrow$ $\gamma +$ invis.
Vectorlike fermions		X	X	X			X	X	
2 Higgs doublet	X	X	X	X					
Higgs singlet	X	X		X			X		
NMSSM	X	X	X	X	X				X
Flavored Higgs	X	X	X			X			
NR bound state		X		X				X	
Pion of new forces		X	X	X	X		X	X	X
RS radion	X	X	X	X			X		
RS graviton	X	X		X			X		

Table 2: Anomalies in precision measurements expected to be visible at the ILC for the models of the Φ discussed in this section.

***Direct observation of the Φ
in e^+e^- and $\gamma\gamma$ collisions***

Observation of Φ in e^+e^- and $\gamma\gamma$ collisions

$\Phi \rightarrow \gamma\gamma$ means there is a **$\Phi\gamma\gamma$ coupling** which implies



PLC: Production via $\gamma\gamma$ collisions

	$qq + gg$	bb	tt	$ee/\mu\mu/\tau\tau$	$\gamma\gamma$	$Z\gamma$	ZZ	hh	WW	Zh
σ (fb)	46	2	760	40	20	20	20	< 0.4	7600	1
BR 5σ	0.4%	0.1%	2%	0.4%	0.3%	0.3%	0.3%	0.04%	5%	0.06%

Table 3: Standard Model background cross sections for the observation of decays of a spin 0 resonance Φ at the PLC. The second line gives the branching ratio, relative to $BR(\Phi \rightarrow gg)$, for a 5σ observation with a 360 fb^{-1} data set as described in the text.

The capability for direct observation of the gg decay and the sensitivity to bb , tt , and Higgs modes far exceed what will be possible at the LHC.

Properties to be measured at PLC

$\Gamma_{\gamma\gamma} \times BR(\Phi \rightarrow gg)$

J (from decay angular distribution)

CP (from transversely polarized initial photons)

$R := A_2/A_1$ (from ratios of rates to different decay modes)

Γ_Φ (directly from the mass spectrum if it is 10s of GeV as suggested by ATLAS)

1. *The note is intended to show*
 - ***The 500 GeV ILC has a lot to say about X750 through precision measurements plus possible discovery of NPs associated with X750.***
 - ***Possible energy upgrade with PLC option will open up even greater opportunities to uncover the new physics operating behind X750 together with LHC.***
2. ***Our strategy stated in the ICFA letter to MEXT's ILC Advisory Panel is intact:***

While performing precision studies of the Higgs boson and the top quark, we will prepare for the energy upgrade of the ILC taking advantage of energy expandability enabled by its linear shape.

Caution

*It might turn out that **the Φ is a relatively minor player** in a new sector of physics that the LHC will begin to uncover in the next few years.*

For this reason, ***it is premature to discuss a new accelerator intended specifically to target the Φ*** or any other new particle that turns up in the early 13 TeV LHC data.

ILC Project Status

MEXT's ILC Review

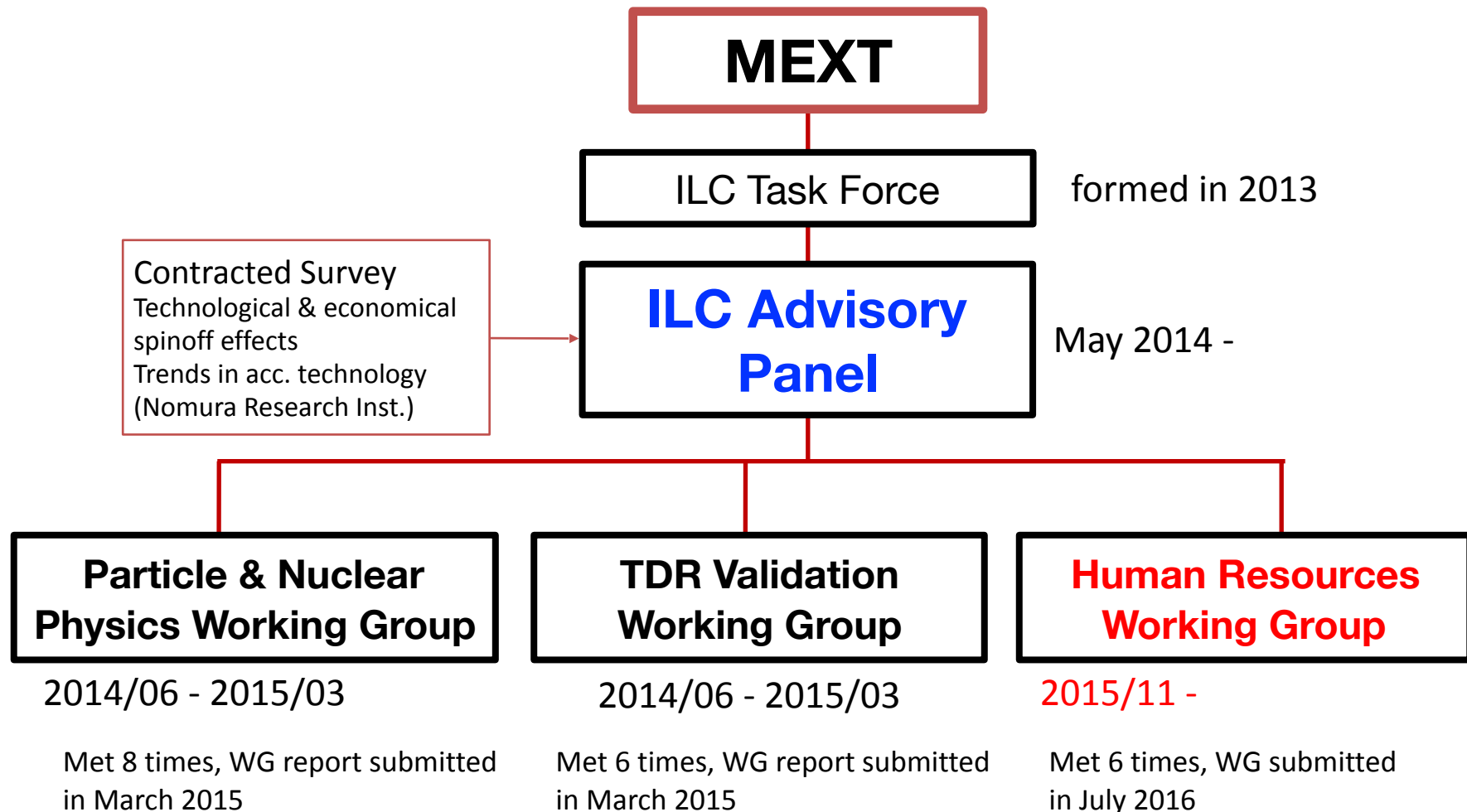
MEXT

=

Japan's
**Ministry of
Education,
Culture, Sports, Science and
Technology**

MEXT's ILC Review

In May, 2014 MEXT setup the ILC Advisory Panel for discussion on various issues concerning ILC construction in Japan. **ILC is now officially being studied.**



Interim Report

ILC Advisory Panel released “Summary of Discussions” in Aug., 2015 and clarified issues to be addressed in their recommendations:

Recommendation 1: The ILC project requires huge investment that is so huge that a single country cannot cover, thus *it is indispensable to share the cost internationally*. From the viewpoint that the huge investments in new science projects must be weighed based upon the scientific merit of the project, *a clear vision on the discovery potential of new particles as well as that of precision measurements of the Higgs boson and the top quark has to be shown* so as to bring about novel development that goes beyond the Standard Model of the particle physics.

[ICFA letter & its followup](#)

Recommendation 2: Since the specifications of the performance and the scientific achievements of the ILC are considered to be designed based on the results of LHC experiments, which are planned to be executed through the end of 2017, it is necessary to *closely monitor, analyze and examine the development of LHC experiments*. Furthermore, it is necessary to *clarify how to solve technical issues and how to mitigate cost risk* associated with the project.

[X750 note](#)

Recommendation 3: While presenting the total project plan, including not only the plan for the accelerator and related facilities but also the plan for other infrastructure as well as efforts pointed out in Recommendations 1 & 2, it is important to *have general understanding on the project by the public and science communities*.

[ILC brochure](#)

Recent International Move

February 11 and 12, 2016, Washington DC

The 4th visit by representatives of Federation of Diet Members for ILC (Formed in the Japanese Diet in 2008, More than 150 members).

The forum is hosted by the Hudson Institute and Advanced Accelerator Association Promoting Science and Technology of Japan (AAA), and attended by Scientists, Industrial people, and DOE and MEXT officials.

- 1st day: “Forum on Enhancing US-Japan Alliance Through Science and Technology”
- 2nd day: “US-Japan ILC Technical Session”
- Japanese Diet members visited several key persons in US Congress and Senate to discuss US-Japan cooperation for the ILC project.
- Diet members plan to continue this forum to establish stronger cooperation with US Congresspersons and Senators



Feb 11, 2016
Rayburn House Office Building



Feb 12, 2016
Hudson Institute

AsiaHEP/ACFA Statement on the ILC & Circular Electron Positron Collider

- AsiaHEP and ACFA reassert their **strong endorsement of the ILC**, which is in a mature state of technical development. The aim of ILC is to explore physics beyond the Standard Model by unprecedented precision measurements of the Higgs boson and top quark, as well as searching for new particles which are difficult to discover at LHC. The Higgs studies at higher energies are especially important for measurement of WW fusion process, to fix the full Higgs decay width, and to measure the Higgs self-coupling. In continuation of decades of world-wide coordination, **we encourage redoubled international efforts at this critical time to make the ILC a reality in Japan.** The past few years have seen growing interest in a large radius circular collider, first focused as a “Higgs factory”, and ultimately for proton-proton collisions at the high energy frontier. **We encourage the effort lead by China in this direction, and look forward to the completion of the technical design in a timely manner.**

June 1, 2016: Executive Meeting of Federation of Diet Members for ILC

Attending executive members:

7 Diet members (executive members)

From MEXT:

Yayoi KOMATSU (Director, Research Promotion Bureau),

Hiroshi IKUKAWA (Deputy Director-General, Research Promotion Bureau),

Masami WATANABE (Director, Basic Research Promotion Division, Research Promotion Bureau),

Sadahiro HAGIWARA (Director for Particle and Nuclear Physics Promotion Office, Basic Research Promotion Division, Research Promotion Bureau) +others

Scientists: Sakue Yamada, Toshinori Mori, Tadashi Ishikawa, Tomohiko Tanabe

AAA: Jun-ichi NISHIYAMA



Agenda

- **Report by MEXT on MEXT-DOE meeting on May 25**
- Discussions and recommendations by Diet members
- Remarks by scientist (Toshinori Mori, as P5 committee member)

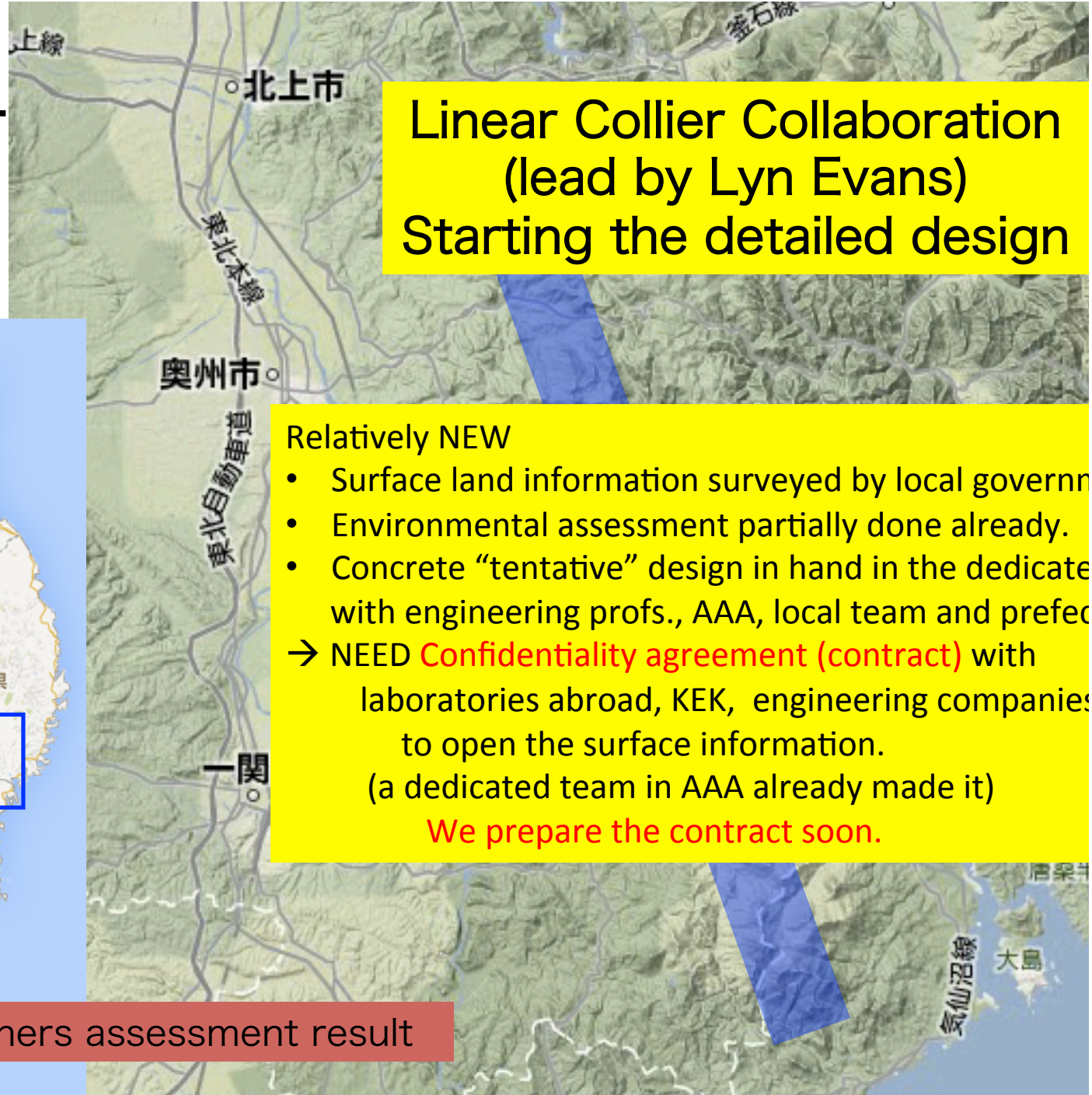
Confirming and Conclusions of the meeting

- **The next meeting of the MEXT-DOE Discussion Group should be held in July/August.**
- **By October this year, items on joint research should be identified and concluded, in time for the next round of budget request (Japan JFY2017, US FY2018)**
- **Discussion Group's planning should include researchers and industry for concrete R&D planning.**

Local Support

4

Kitakami



Linear Collier Collaboration
 (lead by Lyn Evans)
 Starting the detailed design

Relatively NEW

- Surface land information surveyed by local government
- Environmental assessment partially done already.
- Concrete “tentative” design in hand in the dedicated with engineering profs., AAA, local team and prefecture

→ NEED Confidentiality agreement (contract) with laboratories abroad, KEK, engineering companies to open the surface information.
 (a dedicated team in AAA already made it)
 We prepare the contract soon.



Researchers assessment result



Oshu City



Ichinoseki Station



Morioka



Tohoku tourism ad seen on Tokyo Metro



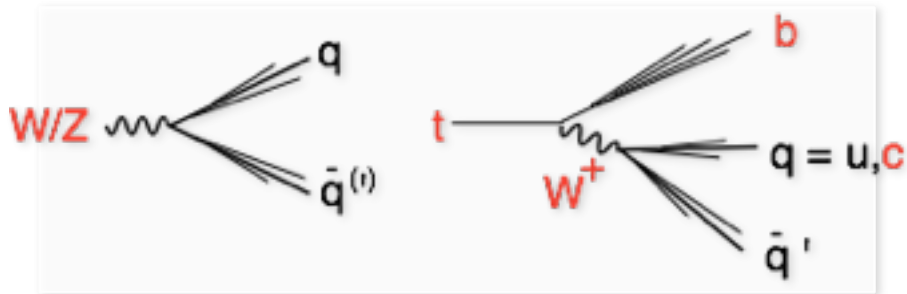
Posters and "Toy ILC" by school children of Oshu City welcoming international workshop on ILC

ILC Detectors

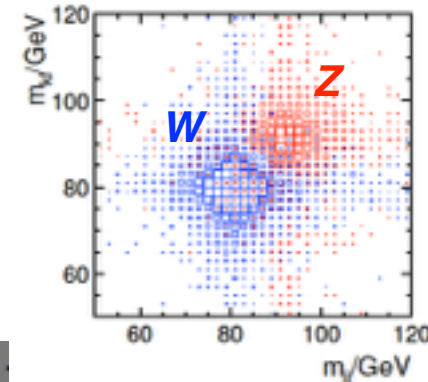
New Paradigm :

View events as viewing a Feynman diagram

Reconstruct final states in terms of fundamental particles (quarks, leptons, gauge bosons, and Higgs bosons)



Identify W/Z/top/Higgs with their jet invariant mass: M_{jets}

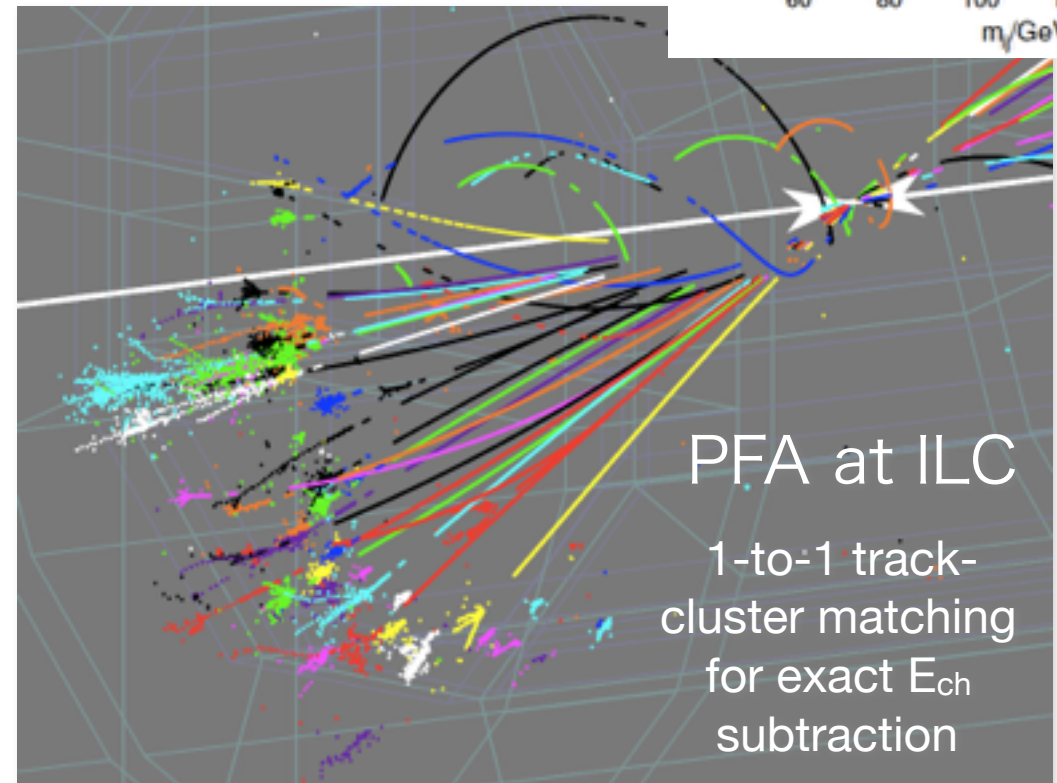


Particle Flow Analysis

PFA is the key to achieve excellent jet invariant mass resolution comparable to the natural width of the weak boson:

$$\sigma_{M_{\text{jets}}} \simeq \Gamma_Z$$

Use tracker for charged particles, **use CAL only for neutral particles**, removing energy deposits by charged particles (E_{ch}) in CAL by **1-to-1 track to CAL cluster matching**



1-to-1 matching requires
High resolution tracking
High granularity calorimetry

Detailed **B**aseline Design (TDR vol.4)

arXiv: 1306.6329

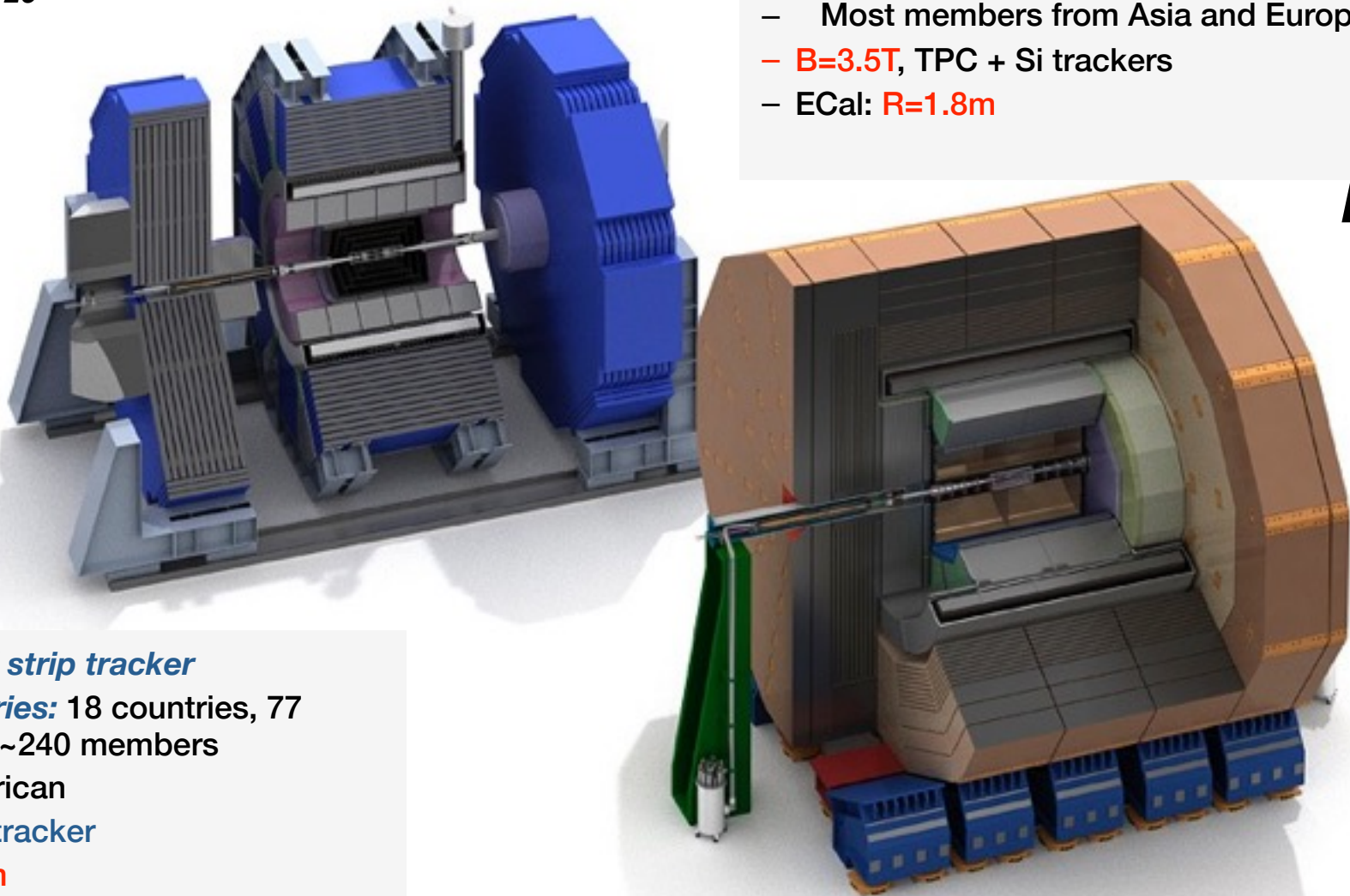
- Large R with **TPC tracker**
- **LOI signatories:** 32 countries, 151 institutions, ~700 members
- Most members from Asia and Europe
- **B=3.5T**, TPC + Si trackers
- ECal: **R=1.8m**

ILD

SiD

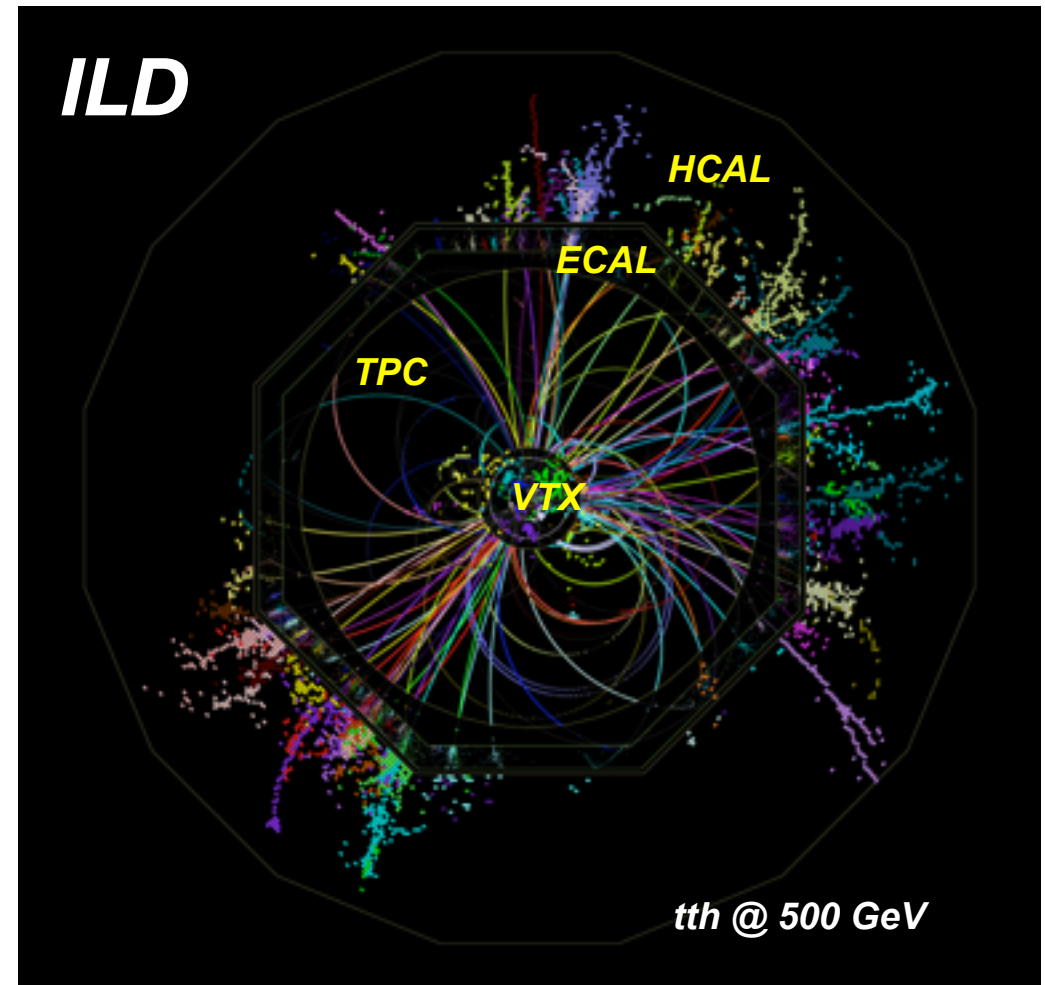
- High B with **Si strip tracker**
- **LOI signatories:** 18 countries, 77 institutions, ~240 members
- Mostly American
- **B=5T**, Si only tracker
- ECal: **R=1.27m**

Both detector concepts are optimized for
Particle Flow Analysis



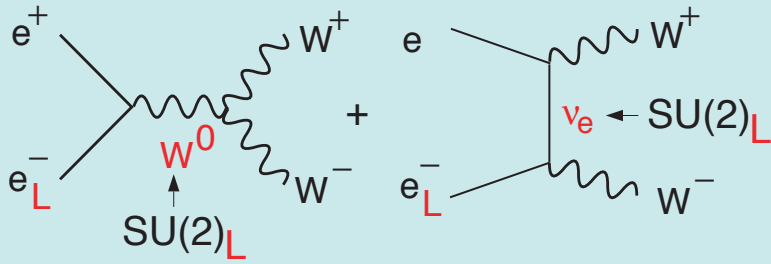
Features of ILC Detectors

- Compared with LHC detectors, ILC detectors have **~10 times better momentum resolution** and **100~1000 times finer granularity**.
- This performance **can be achieved only in the clean environment of the ILC**, and cannot be achieved in the LHC environment.



Power of Beam Polarization

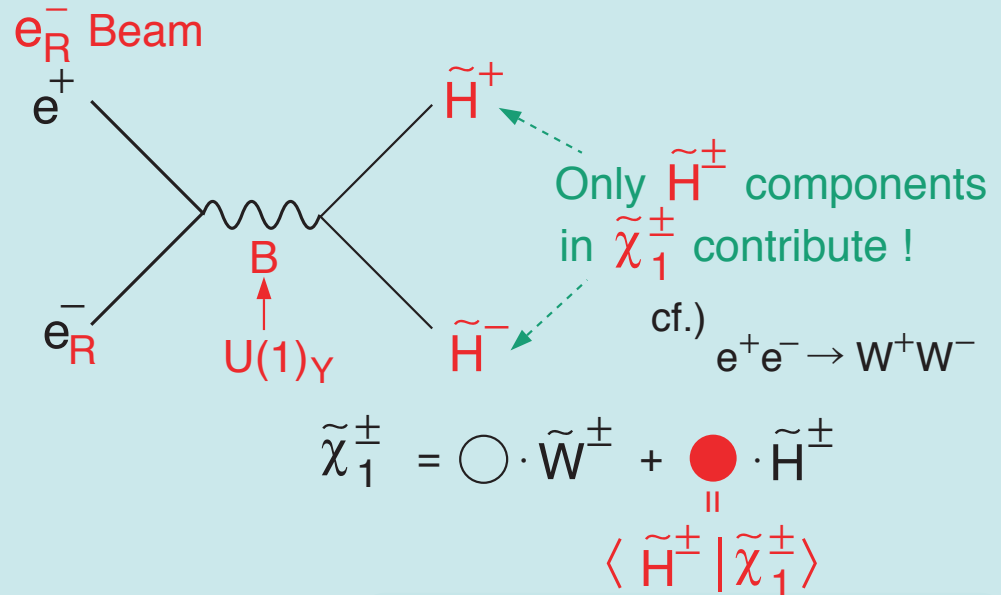
$W^+ W^-$ (Largest SM BG in SUSY searches)



In the symmetry limit, $\sigma_{WW} \rightarrow 0$ for e_R^- !

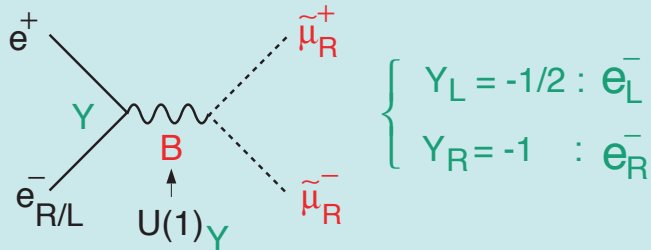
BG Suppression

Chargino Pair



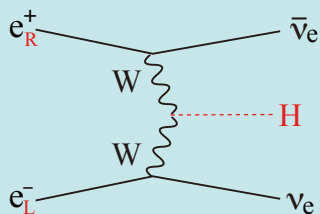
Decomposition

Slepton Pair



In the symmetry limit, $\sigma_R = 4 \sigma_L$!

WW-fusion Higgs Prod.



	ILC
Pol (e ⁻)	-0.8
Pol (e ⁺)	+0.3
$(\sigma/\sigma_0)_{\nu H}$	$1.8 \times 1.3 = 2.34$

Signal Enhancement

Higgs

Why 500 GeV?

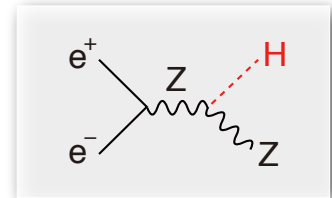
Higgs-related Physics at $E_{cm} \approx 500 \text{ GeV}$

Three well know thresholds

ZH @ 250 GeV ($\sim M_Z + M_H + 20 \text{ GeV}$) :

- Higgs mass, width, J^{PC}
- Gauge quantum numbers
- Absolute measurement of HZZ coupling (**recoil mass**)
- BR($h \rightarrow VV, qq, ll, \text{invisible}$) : $V=W/Z(\text{direct}), g, \gamma(\text{loop})$

\rightarrow Higgs couplings (other than top)

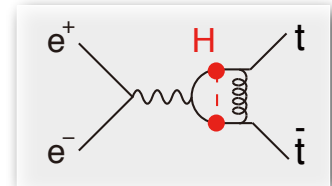


$t\bar{t}$ @ 340-350 GeV ($\sim 2m_t$) : ZH meas. Is also possible

- Threshold scan \rightarrow theoretically clean m_t measurement:
 - \rightarrow test stability of the SM vacuum
 - \rightarrow indirect meas. of top Yukawa coupling
- A_{FB} , Top momentum measurements
- Form factor measurements

$$\Delta m_t(\overline{MS}) \simeq 100 \text{ MeV}$$

$\gamma\gamma \rightarrow HH$ @ 350 GeV possibility



$v\bar{v}H$ @ 350 - 500 GeV :

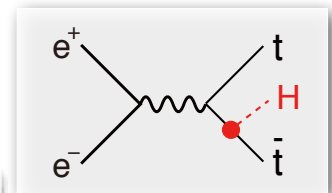
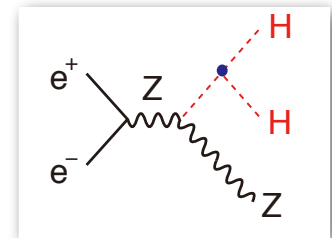
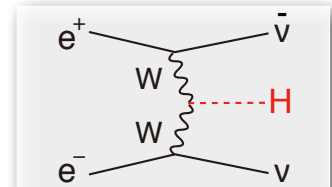
- **HWW coupling** \rightarrow total width \rightarrow absolute normalization of Higgs couplings

ZHH @ 500 GeV ($\sim M_Z + 2M_H + 170 \text{ GeV}$) :

- Prod. cross section attains its maximum at around 500 GeV \rightarrow Higgs self-coupling

$t\bar{t}H$ @ 500 GeV ($\sim 2m_t + M_H + 30 \text{ GeV}$) :

- Prod. cross section becomes maximum at around 800 GeV.
- QCD threshold correction enhances the cross section \rightarrow top Yukawa measurable at 500 GeV concurrently with the self-coupling



We can access all the relevant Higgs couplings at $\sim 500 \text{ GeV}$ for the mass-coupling plot!

Higgs Physics at Higher Energy

Self-coupling with WBF, top Yukawa at xsection max., other higgses, ...

vvH @ $\sqrt{s} > 1\text{TeV}$: $> 1\text{ab}^{-1}$ (pol e^+, e^-)=(+0.2,-0.8)

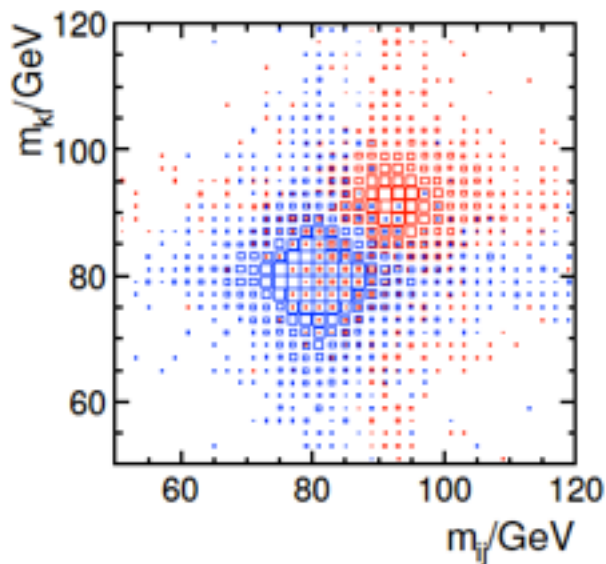
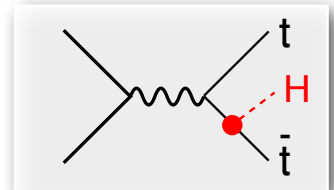
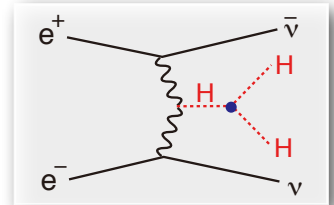
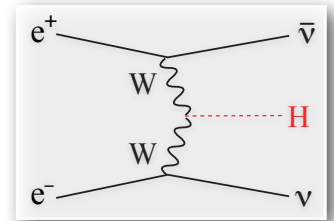
- allows us to measure rare decays such as $H \rightarrow \mu^+ \mu^-$, ...
- further improvements of coupling measurements

vvHH @ 1TeV or higher : 2ab^{-1} (pol e^+, e^-)=(+0.2,-0.8)

- cross section increases with E_{cm} , which compensates the dominance of the background diagrams at higher energies, thereby giving a better precision for the self-coupling.
- If possible, we want to see the running of the self-coupling (very very challenging).

ttbarH @ 1TeV : 1ab^{-1}

- Prod. cross section becomes maximum at around 800GeV.
- CP mixing of Higgs can be unambiguously studied.



Obvious but most important advantage of higher energies in terms of Higgs physics is, however, its **higher mass reach to other Higgs bosons** expected in extended Higgs sectors and **higher sensitivity to $W_L W_L$ scattering** to decide whether the Higgs sector is strongly interacting or not.

In any case we can improve the mass-coupling plot by including the data at 1TeV!

Model-independent Global Fit for Couplings

33 σ_{BR} measurements (Y_i) and σ_{ZH} ($Y_{34,35}$)

$$\chi^2 = \sum_{i=1}^{35} \left(\frac{Y_i - Y'_i}{\Delta Y_i} \right)^2$$

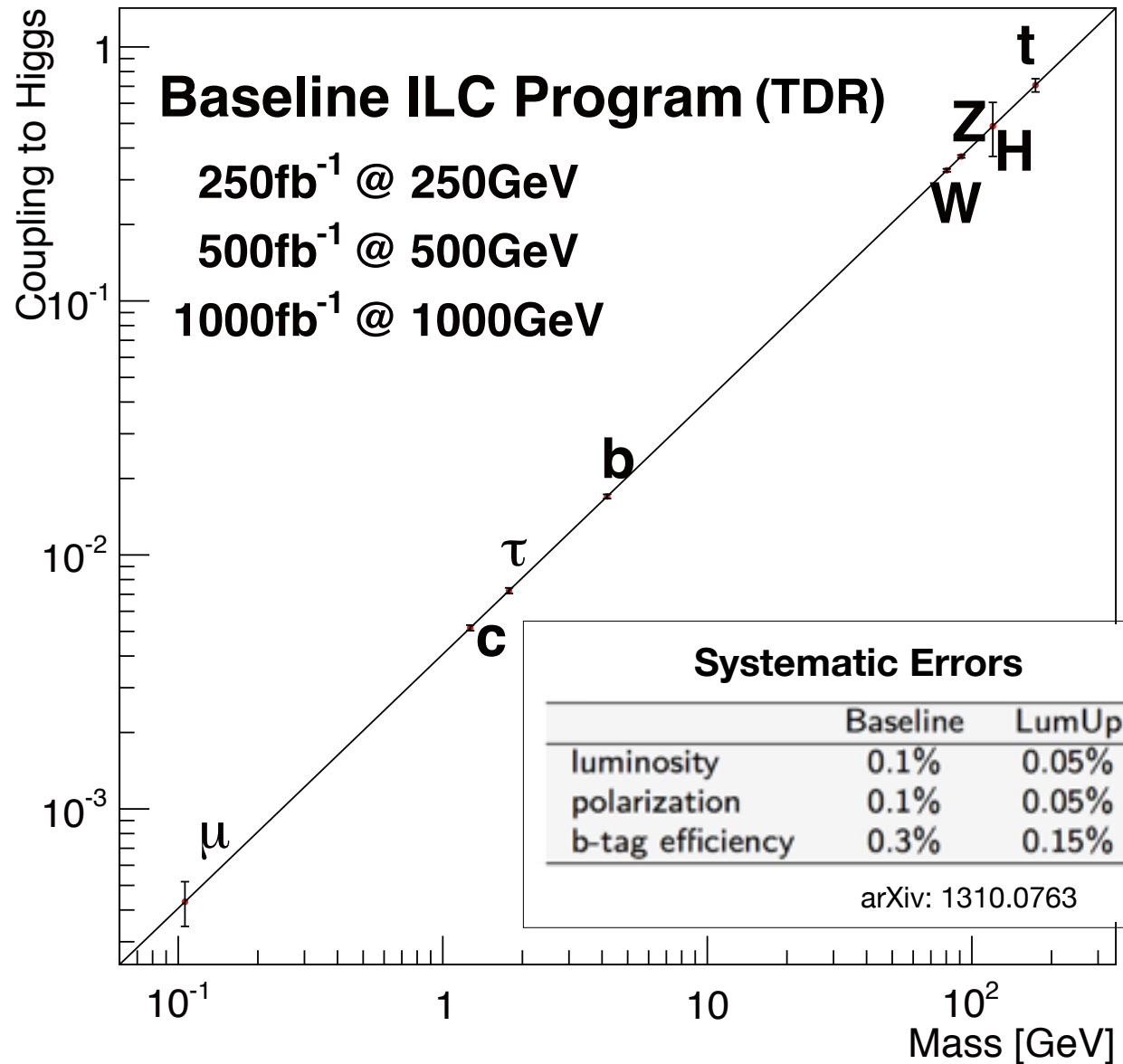
$$Y'_i = F_i \cdot \frac{g_{HA_i A_i}^2 \cdot g_{HB_i B_i}^2}{\Gamma_0}$$

($i = 1, \dots, 33$)
 ($A_i = Z, W, t$)
 ($B_i = b, c, \tau, \mu, g, \gamma, Z, W$: decay)

$$F_i = S_i G_i$$

$G_i = \left(\frac{\Gamma_i}{g_i^2} \right)$

$$S_i = \left(\frac{\sigma_{ZH}}{g_{HZZ}^2} \right), \left(\frac{\sigma_{\nu\bar{\nu}H}}{g_{HWW}^2} \right), \text{ or } \left(\frac{\sigma_{t\bar{t}H}}{g_{Htt}^2} \right)$$



ILC's precisions will eventually reach sub-% level!

Independent Higgs Measurements at ILC

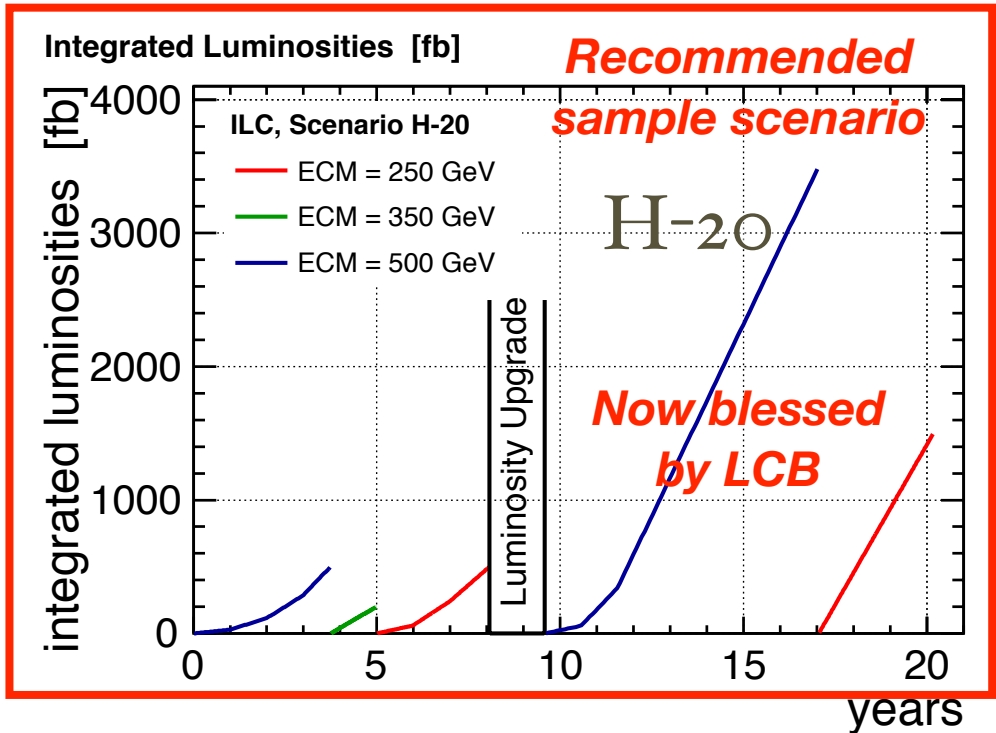
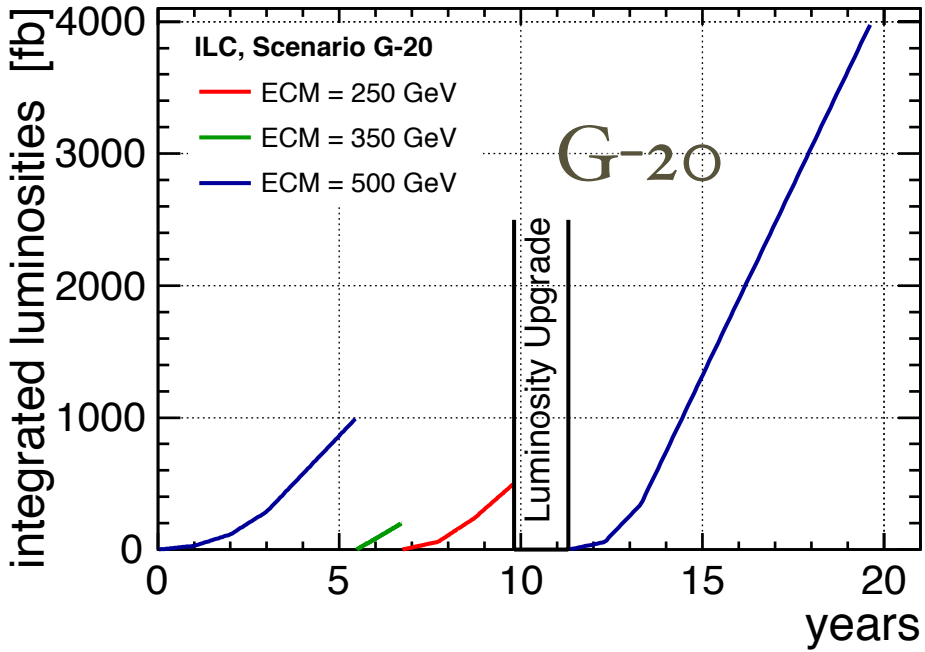
Baseline (=TDR) ILC program

250 GeV: 250 fb⁻¹
 500 GeV: 500 fb⁻¹
 1 TeV: 1000 fb⁻¹

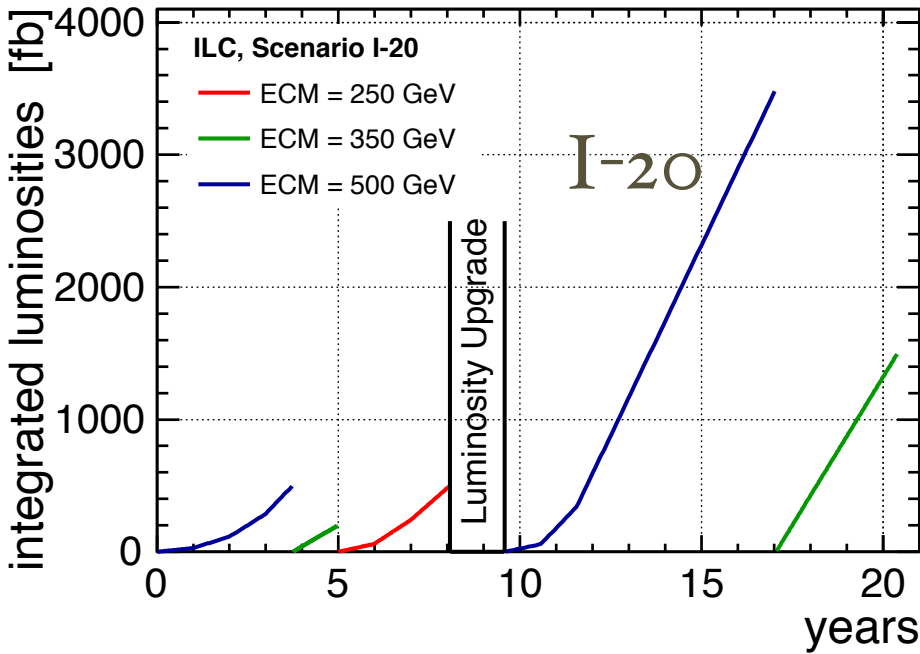
(M_H = 125 GeV)

Ecm	250 GeV		500 GeV		1 TeV
luminosity [fb ⁻¹]	250		500		1000
polarization (e ⁻ ,e ⁺)	(-0.8, +0.3)		(-0.8, +0.3)		(-0.8, +0.2)
process	ZH	vvH(fusion)	ZH	vvH(fusion)	vvH(fusion)
cross section	2.6%	-	3%	-	
	σ·Br	σ·Br	σ·Br	σ·Br	σ·Br
H→bb	1.2%	10.5%	1.8%	0.66%	0.32%
<u>H→cc</u>	8.3%		13%	6.2%	3.1%
H→gg	7%		11%	4.1%	2.3%
H→WW*	6.4%		9.2%	2.4%	1.6%
H→ττ	3.2%		5.4%	9%	3.1%
H→ZZ*	19%		25%	8.2%	4.1%
H→γγ	34%		34%	19%	7.4%
H→μμ	72%	-	88%	72%	31%
tth/H→bb	-		28% (12%@550GeV)		6.2%

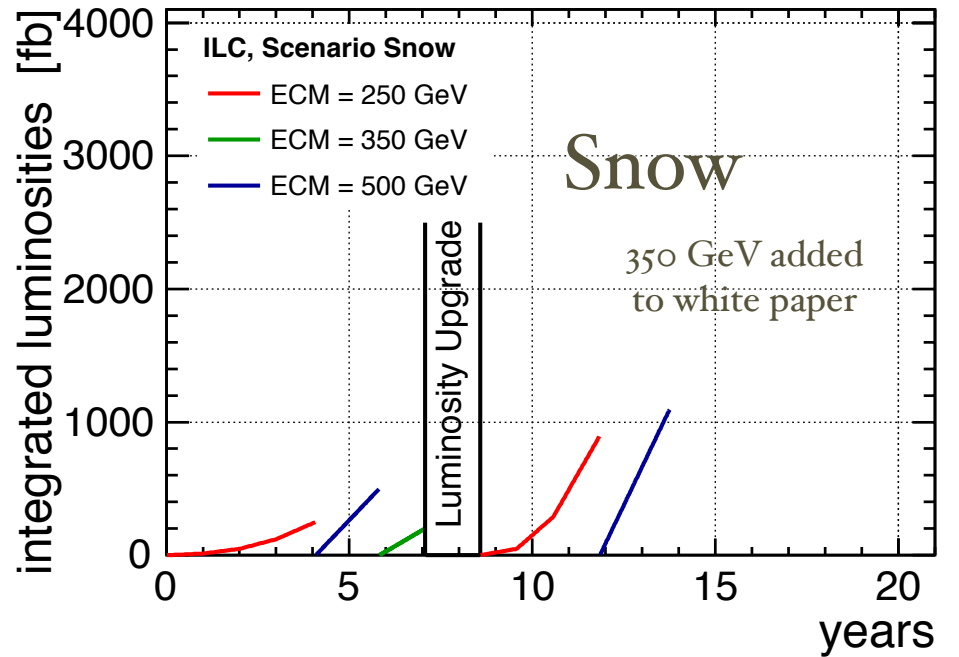
Integrated Luminosities [fb]



Integrated Luminosities [fb]

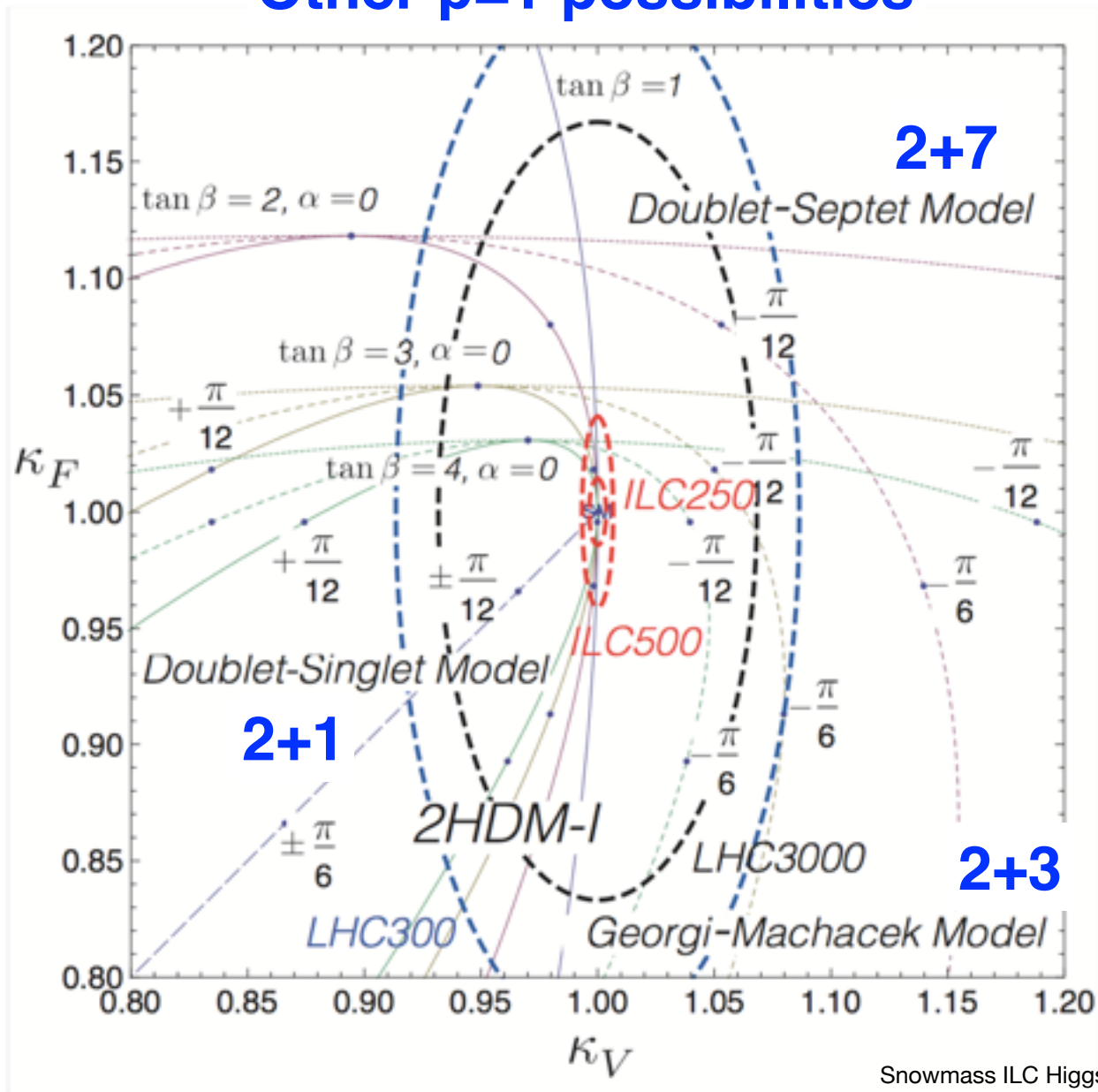


Integrated Luminosities [fb]



Multiplet Structure

Other $\rho=1$ possibilities



Snowmass ILC Higgs White Paper (arXiv: 1310.0763)
 Kanemura et al (arXiv: 1406.3294)

Figure 1.18. The scaling factors in models with universal Yukawa coupling constants.

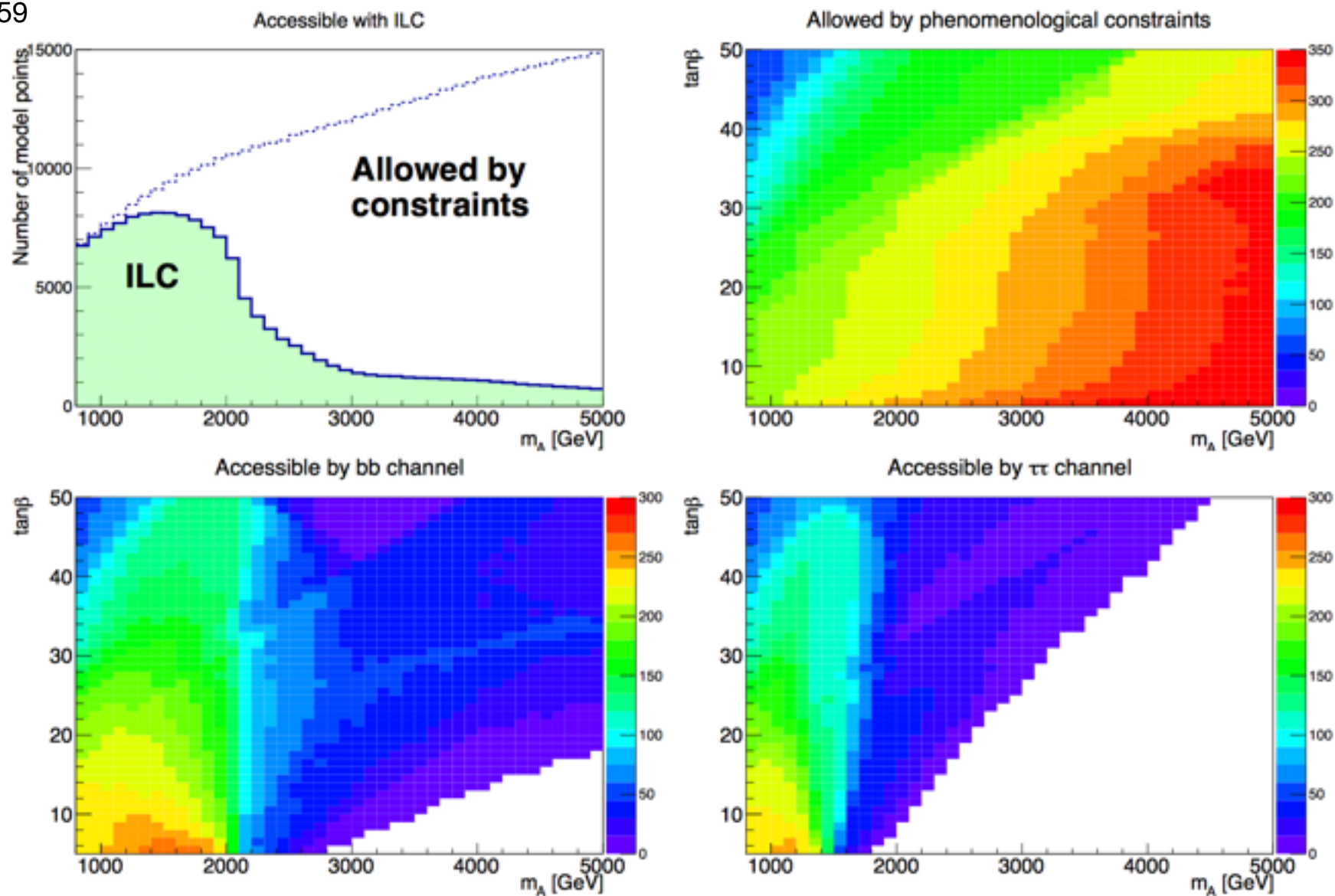
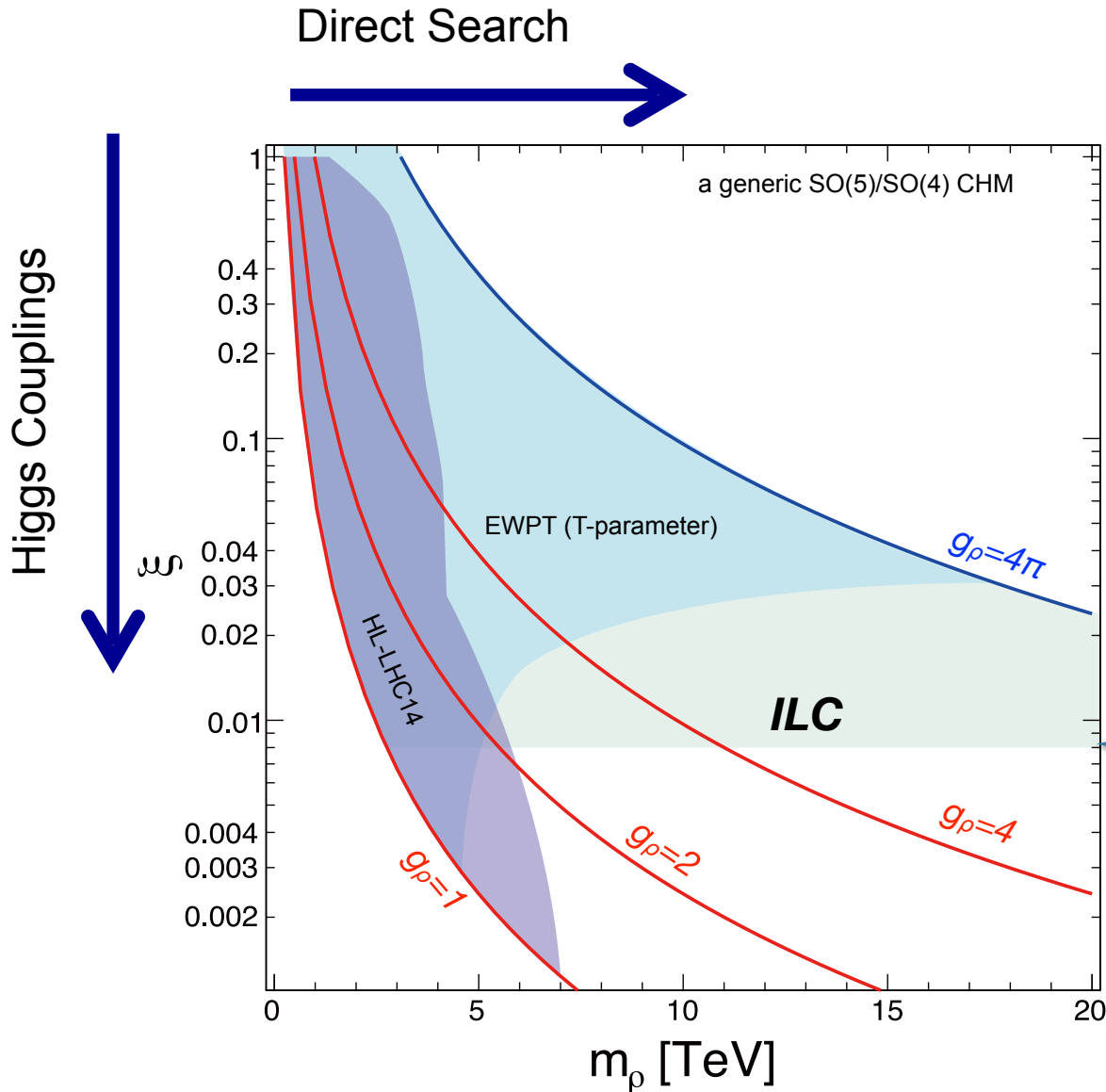
Motoi Endo^(a,b), Takeo Moroi^(a,b), and Mihoko M. Nojiri^(b,c,d)

Figure 8: Upper-left: The number of model points accessible with ILC by at least one decay mode of h as a function of m_A (green histogram), as well as that of model points allowed by the phenomenological constraints (dotted histogram). Upper-right: The number of model points allowed by the phenomenological constraints on m_A vs. $\tan\beta$ plane. Lower-left: The number of model points accessible with ILC by $h \rightarrow \bar{b}b$. Lower-right: The number of model points accessible with ILC by $h \rightarrow \bar{\tau}\tau$.

Composite Higgs: Reach

Complementary approaches to probe composite Higgs models

- Direct search for heavy resonances at the LHC
 - Indirect search via Higgs couplings at the ILC
- Comparison depends on the coupling strength (g_*)



Based on Contino, et al, JHEP 1402 (2014) 006
Torre, Thamm, Wulzer 2014
Grojean @ LCWS 2014

$$\xi = \frac{g_\rho^2}{m_\rho^2} v^2 = \frac{v^2}{f^2}$$

$$\frac{g_{hVV}}{g_{h_{SM}VV}} = \sqrt{1 - \xi}$$

ILC (250+500 LumiUP)

$$\Delta \frac{g_{hVV}}{g_{h_{SM}VV}} = 0.4\%$$

New resonance scale and fingerprint identification in minimal composite Higgs models

Shinya Kanemura,¹ Kunio Kaneta,² Naoki Machida,¹ and Tetsuo Shindou³

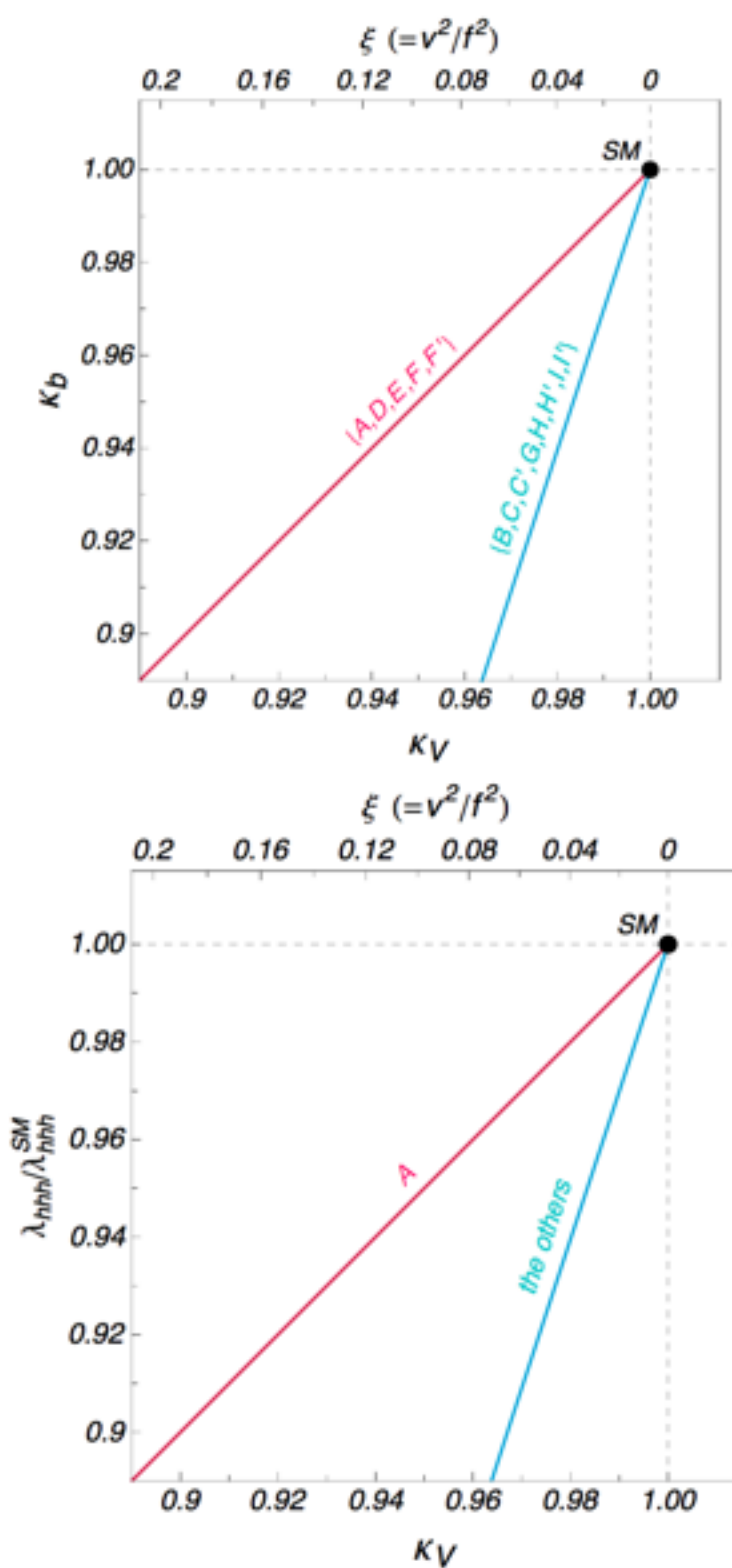
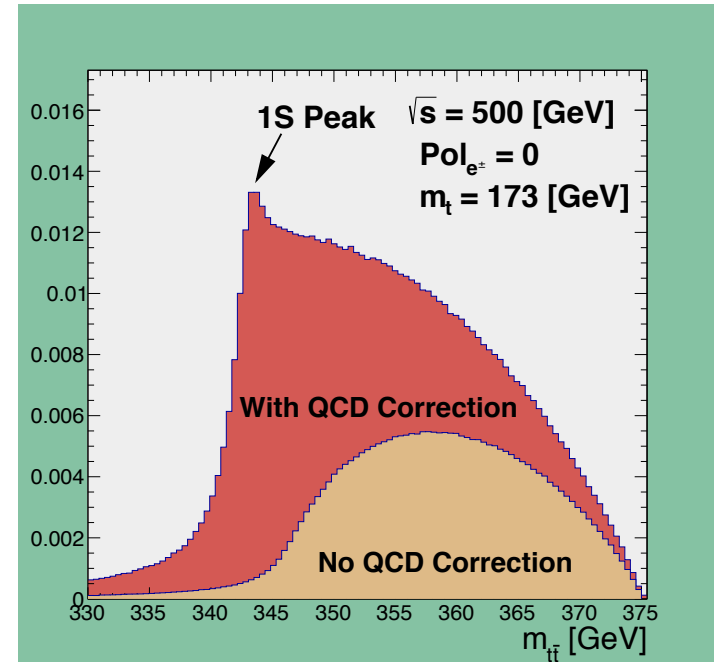
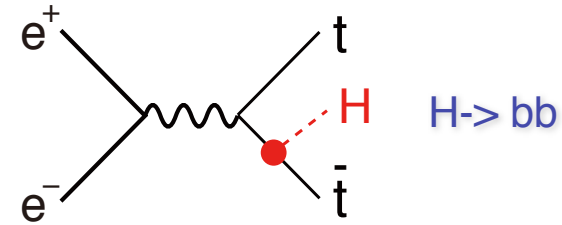
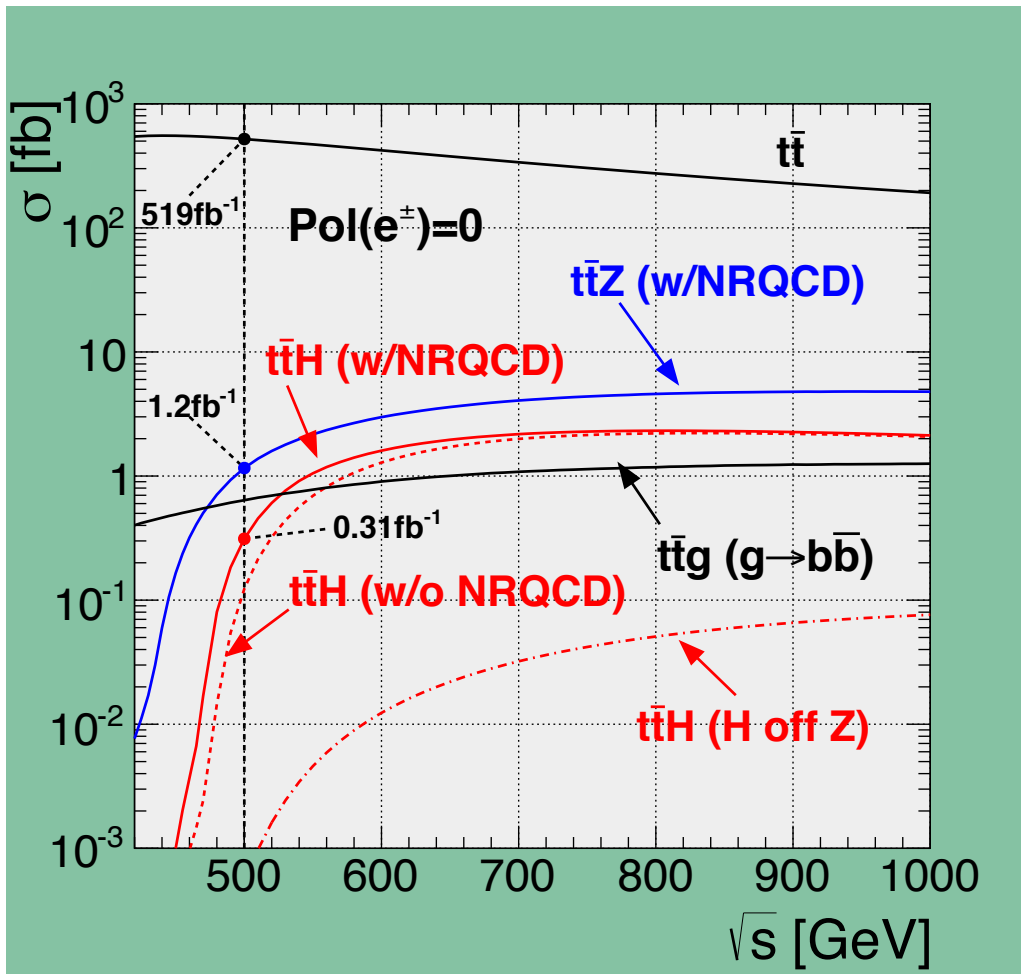


TABLE I: Scale factors for MCHMs with various matter representations. The labels are used in Fig. 7, where C, H and I are the case of $M_2^2 \rightarrow 0$, and C', H' and I' are the case of $M_2^2 \rightarrow 0$.

Label	Model	κ_V	κ_{AVV}	κ_{AAA}	κ_{AAA}	κ_t	κ_b	κ_{hh}	κ_{hhh}
A	MCHM ₄	$\sqrt{1-\xi}$	$1-2\xi$	$\sqrt{1-\xi}$	$1-\frac{1}{3}\xi$	$\sqrt{1-\xi}$	$\sqrt{1-\xi}$	$-\xi$	$-\xi$
B	MCHM ₅	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(3+2\xi^2/3)}{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
B	MCHM ₁₀	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(3+2\xi^2/3)}{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
C, C'	MCHM ₁₄	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_2	F_3	$\frac{1-2\xi}{\sqrt{1-\xi}}$	F_6	-4ξ
D	MCHM _{5,5-10}	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(3+2\xi^2/3)}{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\sqrt{1-\xi}$	-4ξ	$-\xi$
E	MCHM _{5,10-10}	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(3+2\xi^2/3)}{1-\xi}$	$\sqrt{1-\xi}$	$\sqrt{1-\xi}$	$-\xi$	$-\xi$
F, F'	MCHM _{5,14-10}	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_2	F_3	$\sqrt{1-\xi}$	F_6	$-\xi$
G	MCHM _{10,5-10}	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(3+2\xi^2/3)}{1-\xi}$	$\sqrt{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$-\xi$	-4ξ
B	MCHM _{10,14-10}	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_2	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
B	MCHM _{14,5-10}	$\sqrt{1-\xi}$	$1-2\xi$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi(3+2\xi^2/3)}{1-\xi}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
H, H'	MCHM _{14,5-10}	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_2	F_3	$\frac{1-2\xi}{\sqrt{1-\xi}}$	F_7	-4ξ
B	MCHM _{14,10-10}	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_2	$\frac{1-2\xi}{\sqrt{1-\xi}}$	$\frac{1-2\xi}{\sqrt{1-\xi}}$	-4ξ	-4ξ
I, I'	MCHM _{14,14-10}	$\sqrt{1-\xi}$	$1-2\xi$	H_1	H_2	F_3	$\frac{1-2\xi}{\sqrt{1-\xi}}$	F_6	-4ξ

Top Yukawa Coupling

The largest among matter fermions, but not yet directly observed



A factor of 2 enhancement from QCD bound-state effects

Cross section maximum at around $E_{cm} = 800\text{GeV}$

Philipp Roloff, LCWS12

Tony Price, LCWS12

DBD Full Simulation

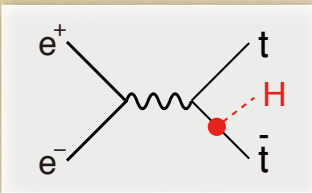
$$1 \text{ ab}^{-1} @ 500 \text{ GeV} \quad m_H = 125 \text{ GeV}$$

$$\Delta g_Y(t)/g_Y(t) = 9.9\%$$

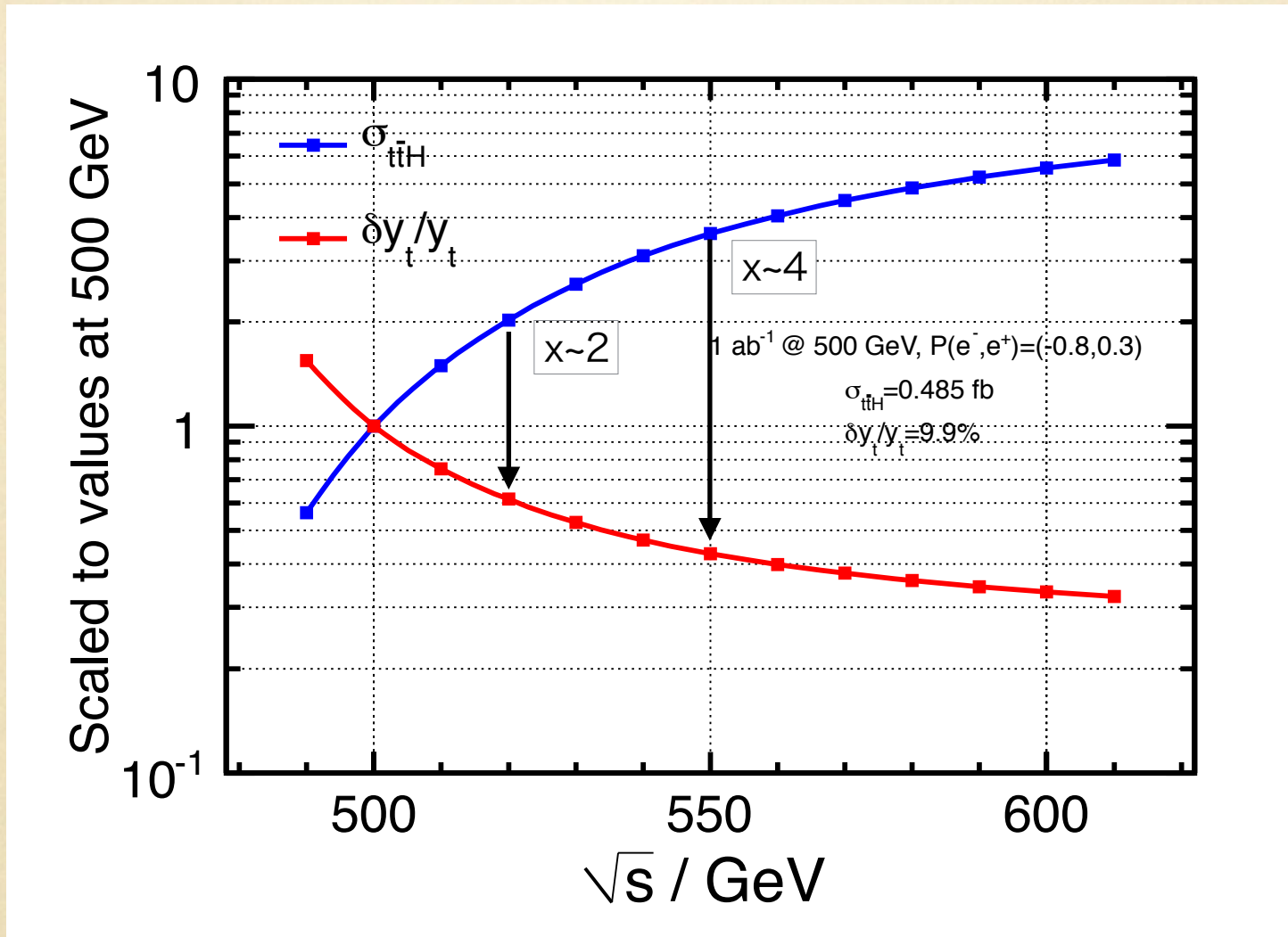
Tony Price, LCWS12

scaled from $m_H=120 \text{ GeV}$

Notice $\sigma(500+20\text{GeV})/\sigma(500\text{GeV}) \sim 2$
Moving up a little bit helps significantly!



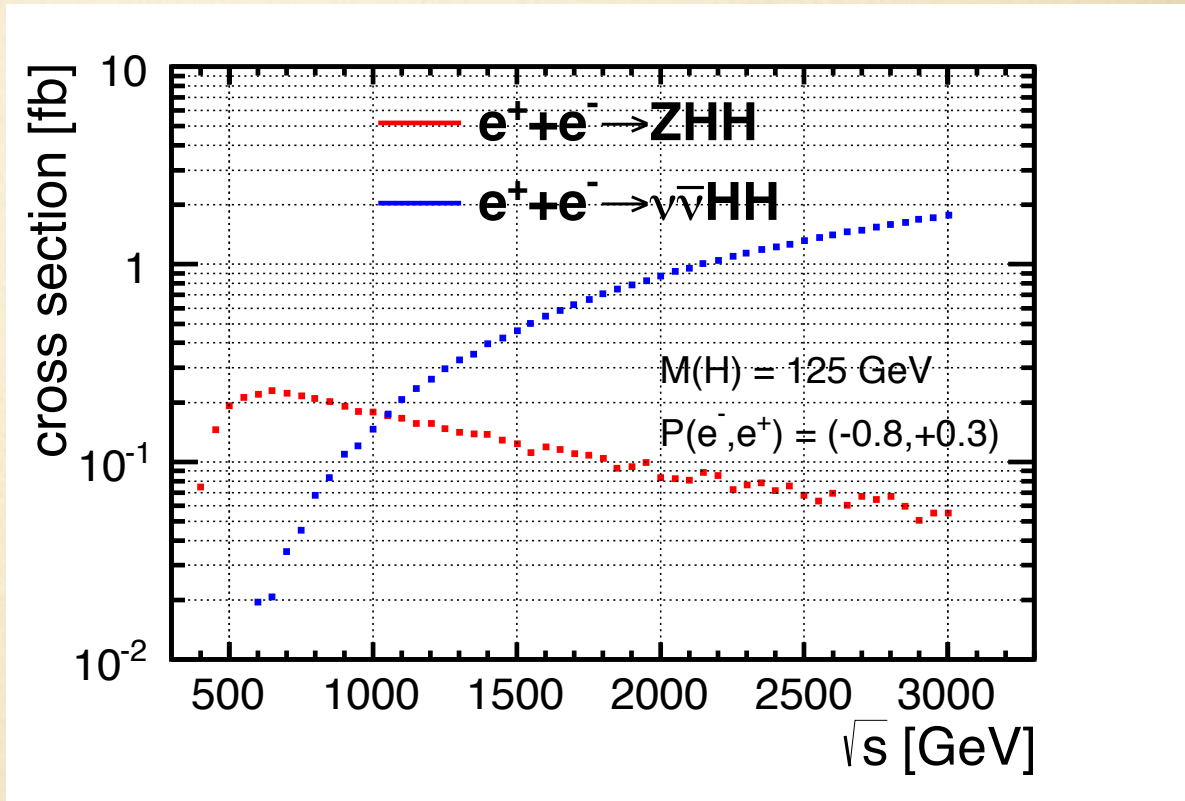
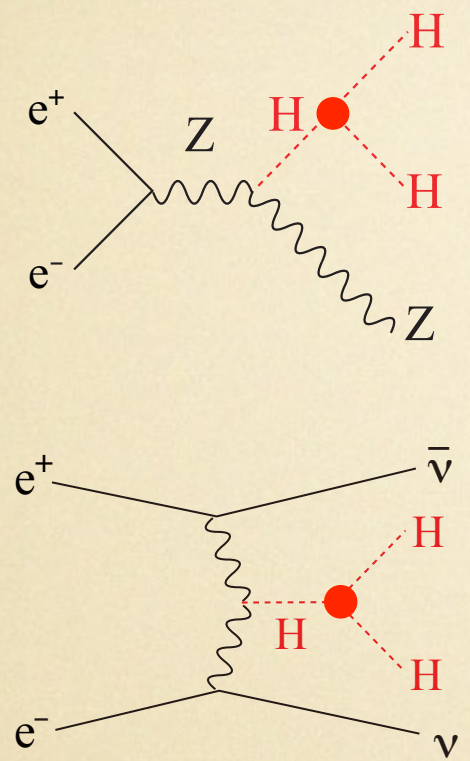
Top Yukawa coupling



Y. Sudo

Slight increase of E_{max} is very beneficial!

prospects of Higgs self-coupling @ linear colliders



prospects from full simulation studies:

ILC	$\Delta\lambda_{HHH}/\lambda_{HHH}$	500 GeV	+ 1 TeV
	Snowmass	46%	13%
	H20	27%	10%

(ref. H20 arXiv: 1506.07870)

J. Tian, LC-REP-2013-003

CLIC	1.4 TeV (1.5 ab ⁻¹)	+3 TeV (2 ab ⁻¹)
	21%	10%

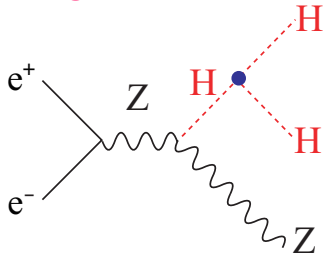
(arXiv: 1307.5288)

M. Kurata, LC-REP-2014-025

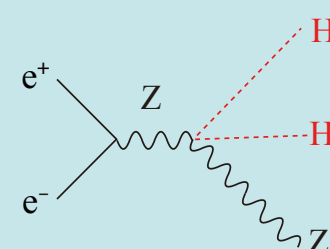
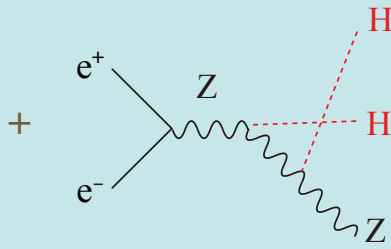
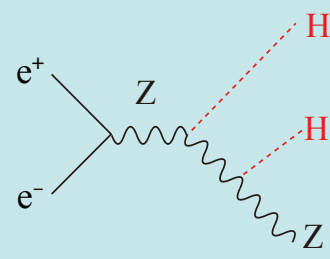
C. Dürig @ ALCW15

The Problem : BG diagrams dilute self-coupling contribution

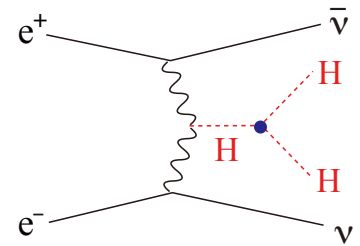
Signal diagram



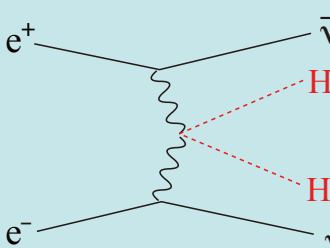
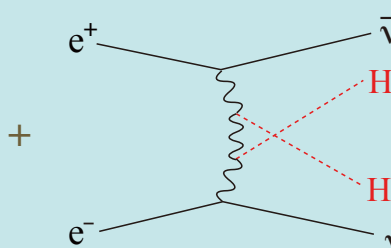
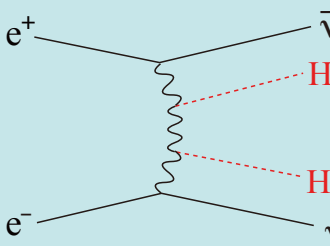
Irreducible BG diagrams



Signal diagram



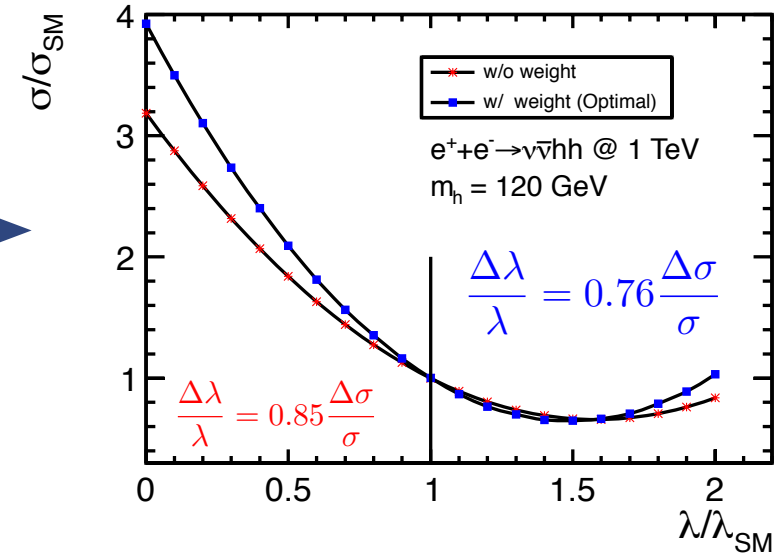
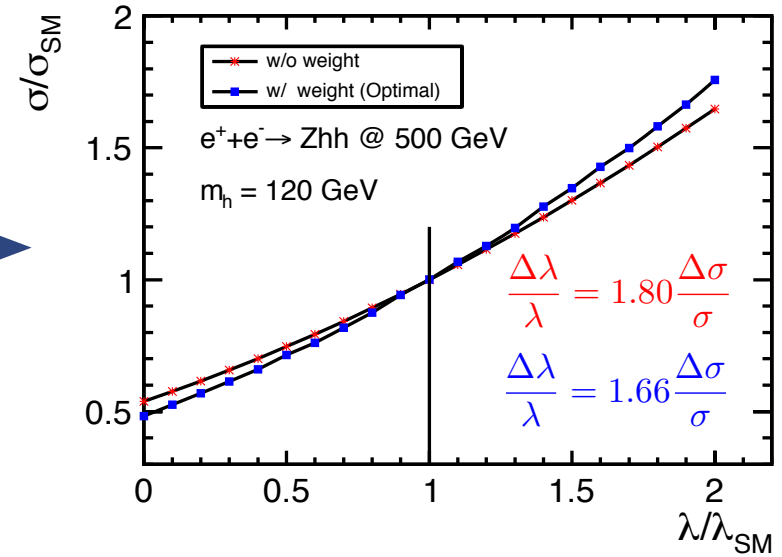
Irreducible BG diagrams



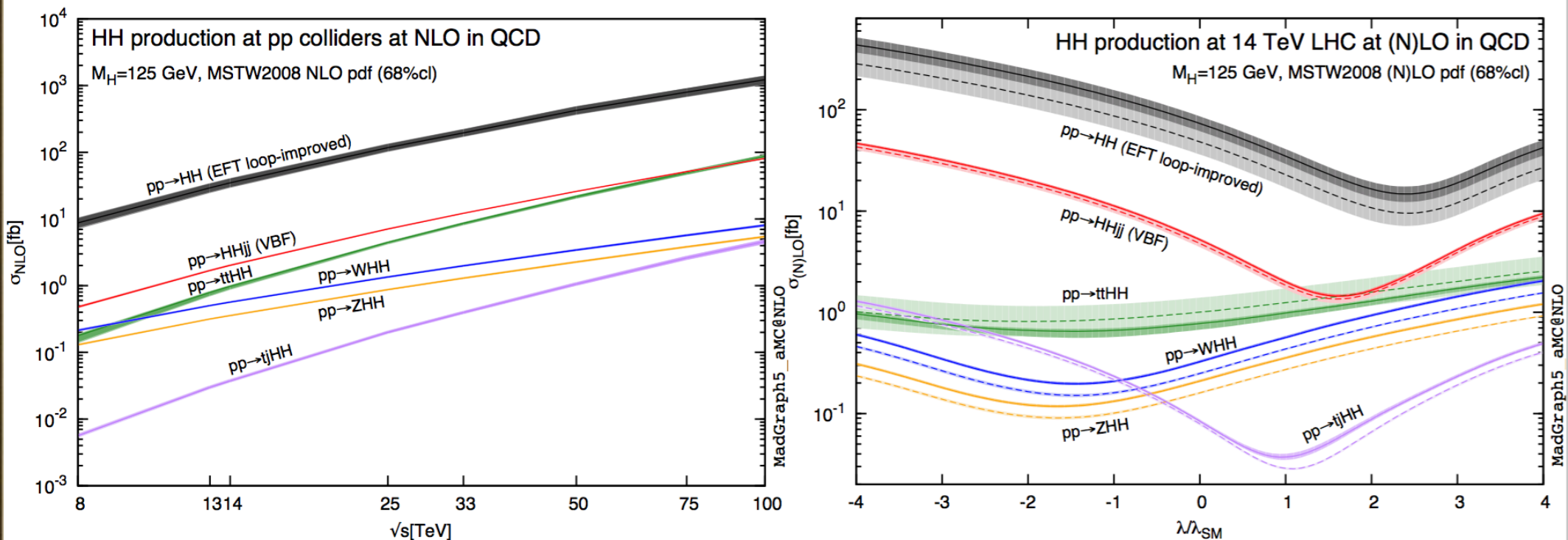
$$\sigma = \lambda^2 S + \lambda I + B$$

$$\frac{\Delta\lambda}{\lambda} = F \cdot \frac{\Delta\sigma}{\sigma}$$

$F=0.5$ if no BG diagrams



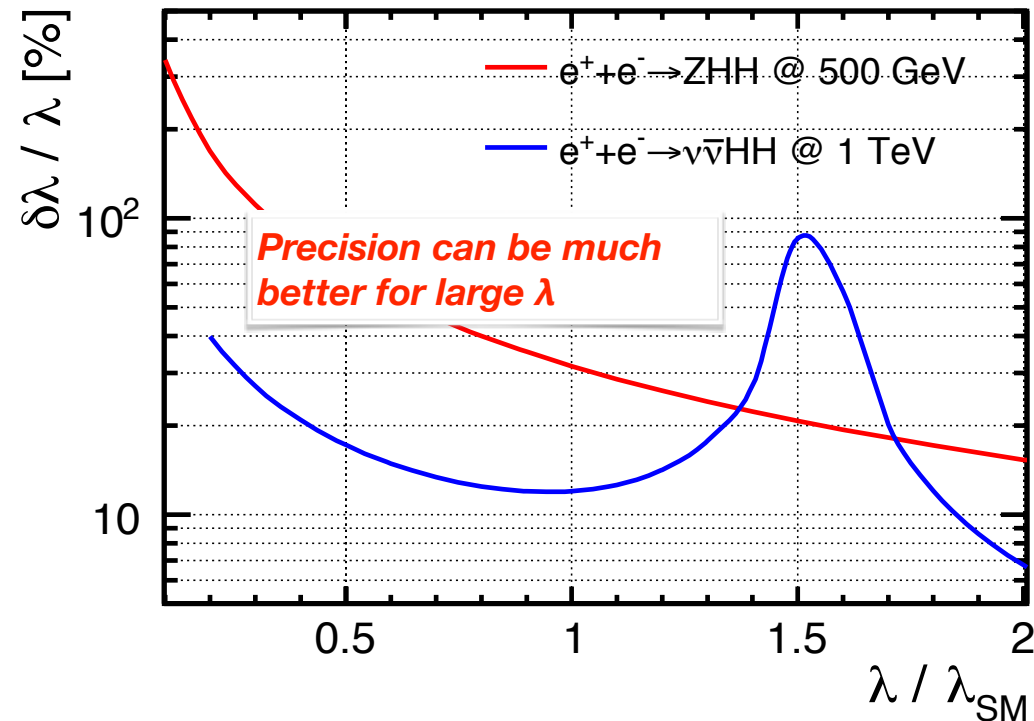
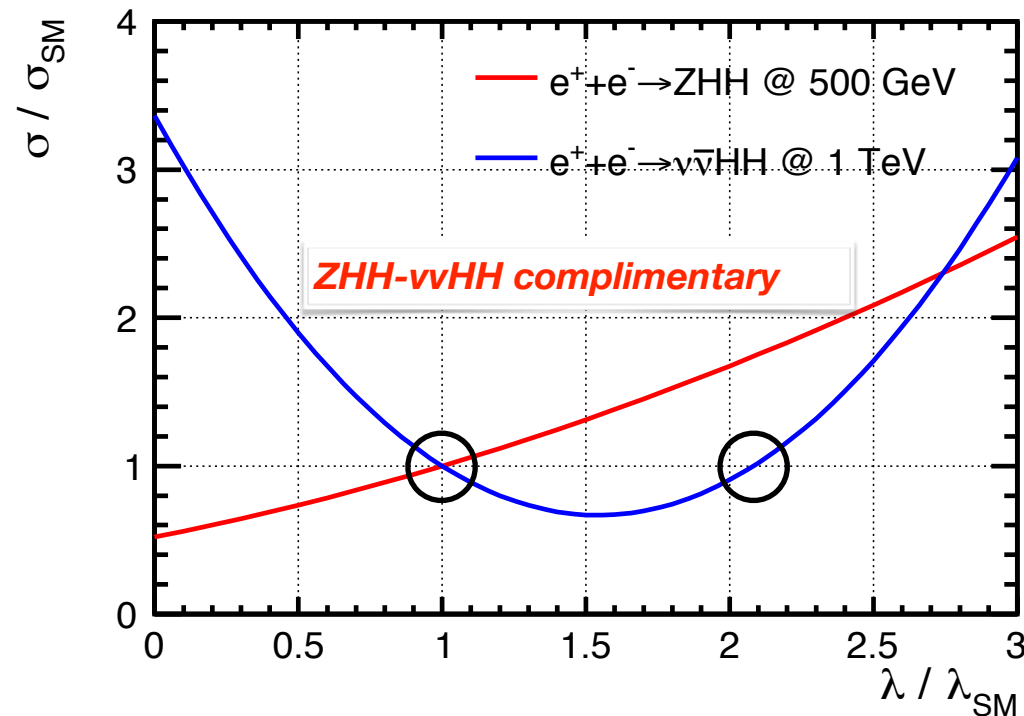
What if $\lambda \neq \lambda_{SM}$? @ LHC



arXiv:1401.7304

- interference is destructive, σ minimum at $\lambda \sim 2.5\lambda_{SM}$; if λ is enhanced, it's going to be very difficult (from snowmass study by 3000 fb⁻¹ @ 14 TeV, significance of double Higgs production is only $\sim 2\sigma$, if cross section decreases by a fact of 2~3, very challenging to observe $pp \rightarrow HH$)

What if $\lambda \neq \lambda_{\text{SM}}$? @ LCs



- for ZHH, interference is constructive, enhanced λ will increase σ , and improve sensitivity factor as well, e.g. if $\lambda = 2\lambda_{\text{SM}}$, σ increases by 60%, F reduced by 1/2, $\delta\lambda/\lambda \sim 15\%$
 \rightarrow we may finish the λ story at 500 GeV ILC!

In EWSB models with classical conformal symmetry
 (Hashino, Kanemura, Orikasa, arXiv:1508.03245)

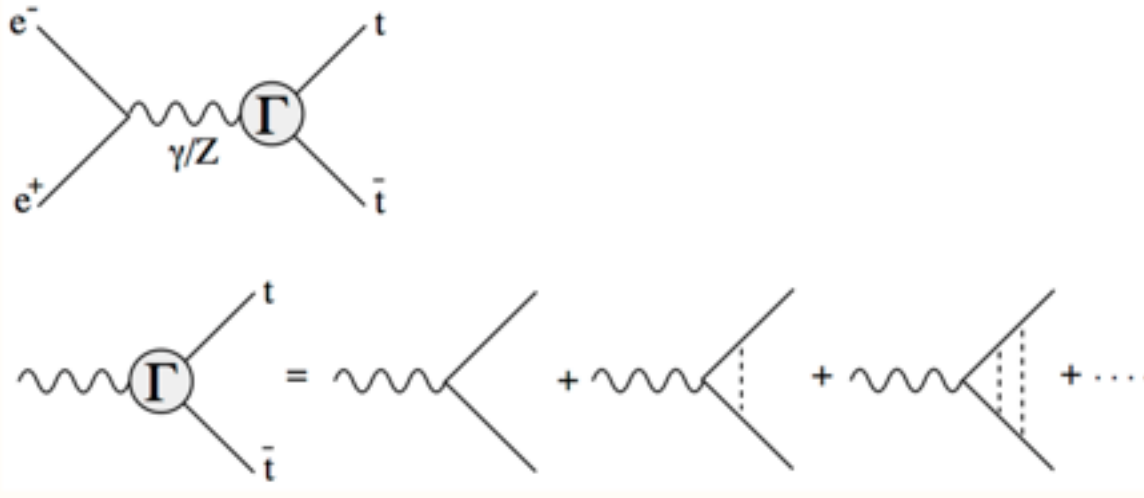
$$\Gamma_{hhh}^{\text{CSI}} = \frac{5m_h^2}{v} = \frac{5}{3} \times \Gamma_{hhh}^{\text{SMtree}}.$$

- for $\nu\bar{\nu}\text{HH}$, interference is destructive, enhanced λ will decrease σ , minimum when $\lambda \sim 1.5\lambda_{\text{SM}}$, $\delta\lambda/\lambda$ degrades significantly if $\lambda/\lambda_{\text{SM}} \in (1.3, 1.7)$
- but if $\lambda < \lambda_{\text{SM}}$, more difficult to use ZHH, have to rely more on $\nu\bar{\nu}\text{HH}$
- two channels are complementary in terms of λ measurement in BSM

Top

Top Quark

Threshold Region



At threshold both the top quark and the anti-top quark are slow and stay close to each other, allowing multiple exchange of Coulombic gluons.

⇒ **Leading contribution**

The threshold correction factor (bound-state effect) denoted by Γ satisfies the Bethe-Salpeter equation which reduces to Schroedinger's equation:

$$\left[H - \left(E + \frac{i}{2} \Gamma_{\Theta} \right) \right] G = 1$$

in the non-relativistic limit. The operator G is related to Γ through

$$\Gamma_V^k \simeq - \left(\frac{1}{D_t} + \frac{1}{D_{\bar{t}}} \right) \cdot \tilde{G}(\mathbf{p}; E) \cdot \gamma^k$$

$$\tilde{G}(\mathbf{p}; E) \equiv \langle \mathbf{p} | G | \mathbf{x} = \mathbf{0} \rangle$$

for vector part

$$\Gamma_A^k \simeq - \left(\frac{1}{D_t} + \frac{1}{D_{\bar{t}}} \right) \cdot \left(\frac{\tilde{F}^l(\mathbf{p}; E)}{m_t} \right) \cdot \sigma^{kl} \gamma^5$$

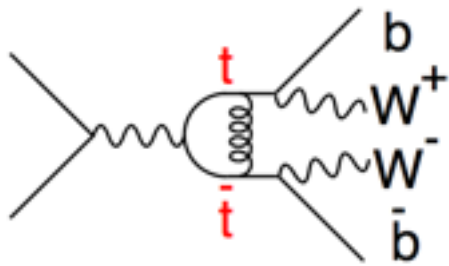
$$\tilde{F}^l(\mathbf{p}; E) \equiv \langle \mathbf{p} | G \cdot \hat{p}^l | \mathbf{x} = \mathbf{0} \rangle$$

for axial vector part

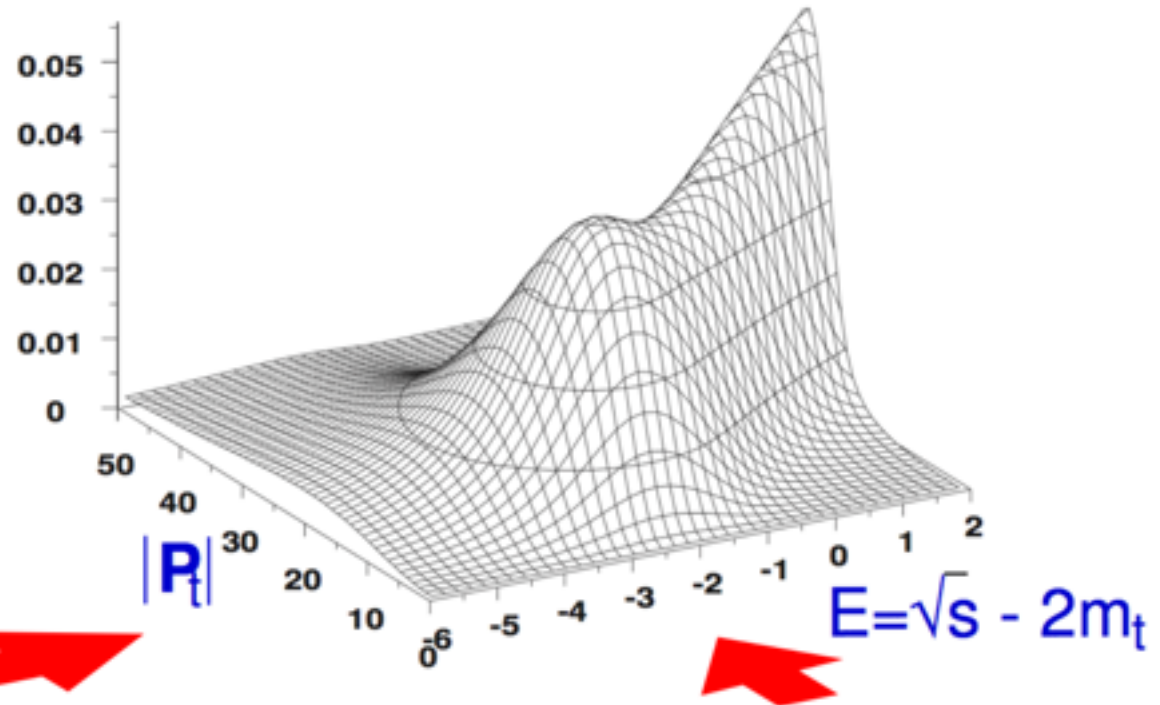
Top Quark

Threshold Region

How to access G experimentally



$$p_{top} = p_{bW} = p_{3jets}$$



Momentum Dist.

$$\frac{d\sigma_{t\bar{t}}}{d|\mathbf{p}|} \propto |\langle \mathbf{p} | G | \mathbf{x} = \mathbf{0} \rangle|^2$$

$$\simeq \left| \sum_n \frac{\phi_n(\mathbf{p}) \Psi_n^*(\mathbf{0})}{E - E_n + i\Gamma_n/2} \right|^2$$

momentum space wave fun.

Threshold Scan

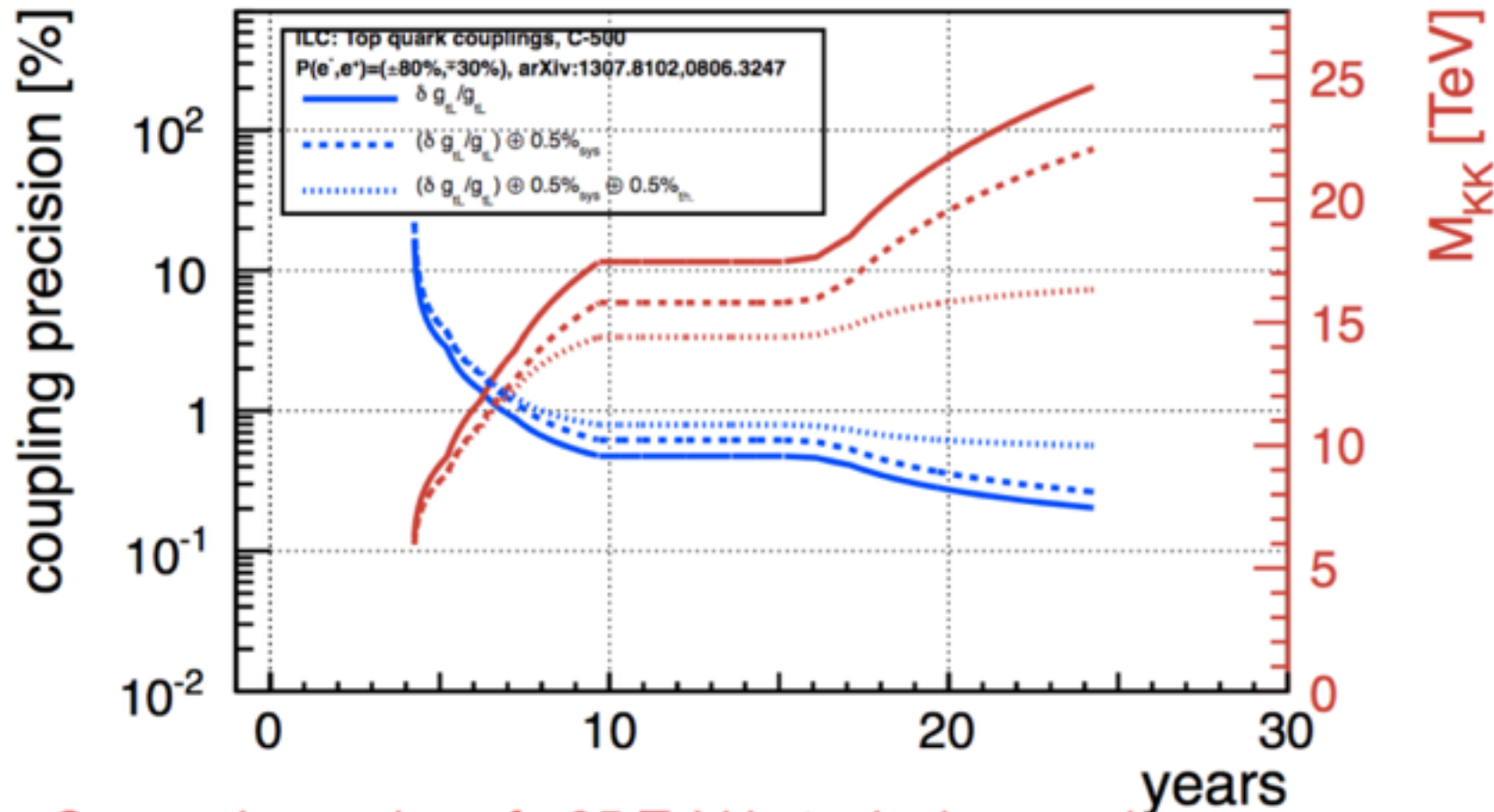
$$\sigma_{t\bar{t}} \propto \text{Im} \langle \mathbf{x} = \mathbf{0} | G | \mathbf{x} = \mathbf{0} \rangle$$

$$\simeq \text{Im} \sum_n \frac{|\Psi_n(\mathbf{0})|^2}{E - E_n + i\Gamma_n/2}$$

wave function at origin

New physics reach for typical BSM scenarios with composite Higgs/Top and or extra dimensions

Based on phenomenology described in Pomerol et al. arXiv:0806.3247



Can probe scales of ~25 TeV in typical scenarios

(... and up to 80 GeV for extreme scenarios)

=> Important guidance for e.g. 100 TeV pp-collider

Comparison to FCC-ee

Recent publication assesses potential of FCC-ee

P. Janot, arXiv:1503.01325, arXiv:1510.09056

- run right above threshold; study assumes 2.4 ab^{-1} at $\sqrt{s} = 365 \text{ GeV}$

(theory systematics close to threshold to be evaluated)

- no beam polarization, use final-state polarization instead

(ILC beam polarization expected to be known to 10^{-3} , can one understand final state polarization to that level?)

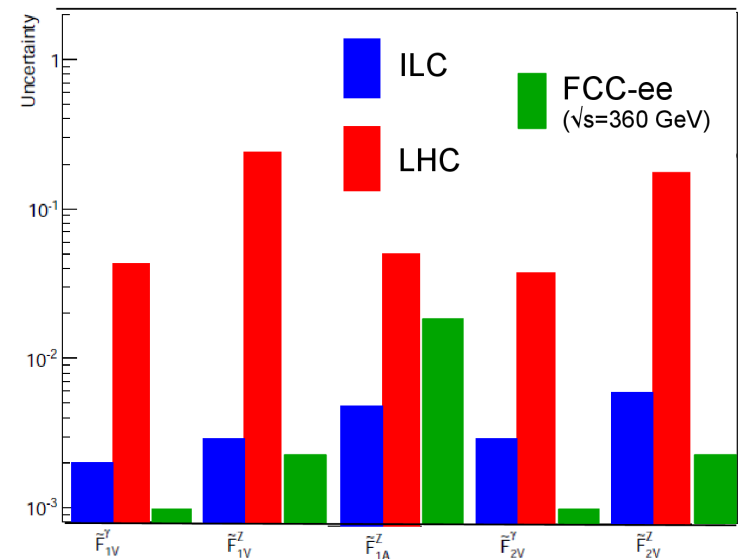
Fast simulation analysis based on lepton energy and angle yields:

- similar precision to ILC for Z couplings, except F1AZ

- significantly better than ILC for photon couplings

Complementarity

Good to see interest in this measurement
Full study needed to understand systematics



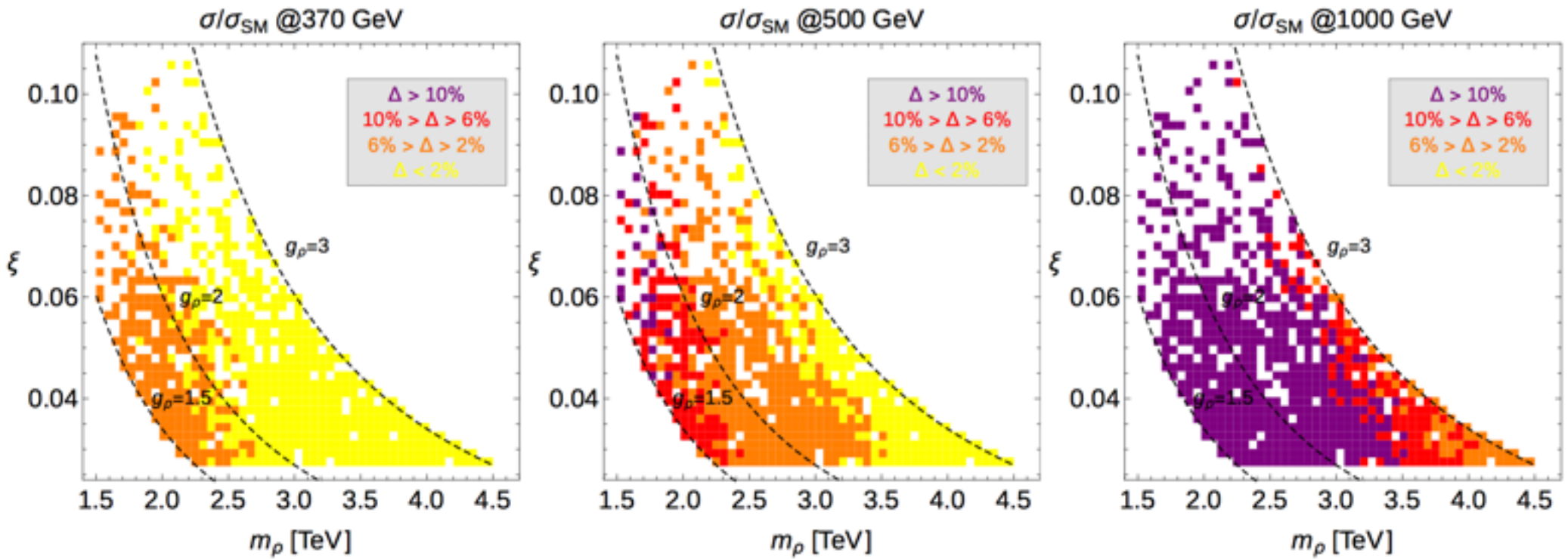


Figure 5. Predicted deviations for the cross section of the process $e^+e^- \rightarrow t\bar{t}$ at 370, 500, 1000 GeV in the 4DCHM compared with the SM as functions of $m_\rho = fg_\rho$ and $\xi = v^2/f^2$. For each point we have selected the configuration yielding the maximal deviation defined as $\Delta = (\sigma^{4\text{DCHM}} - \sigma^{\text{SM}})/\sigma^{\text{SM}}$. The points correspond to $f = 0.75 - 1.5$ TeV, $g_\rho = 1.5 - 3$. Bounds on the masses of the extra fermions are the same as in figure 2.

DM

Slepton decays to DM with small mass differences

Study of stau pair production at the ILC

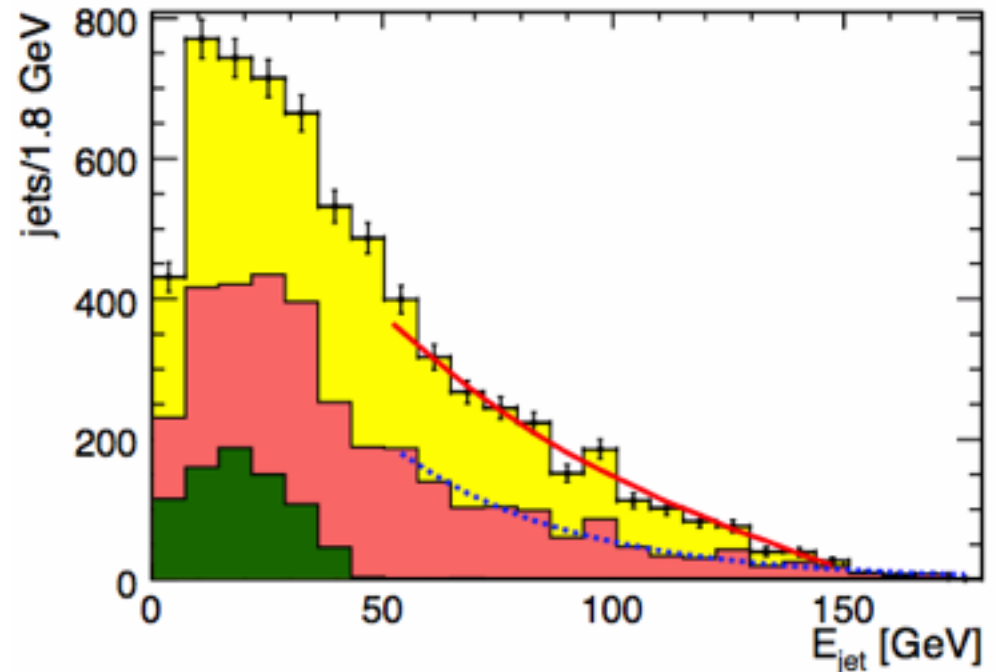
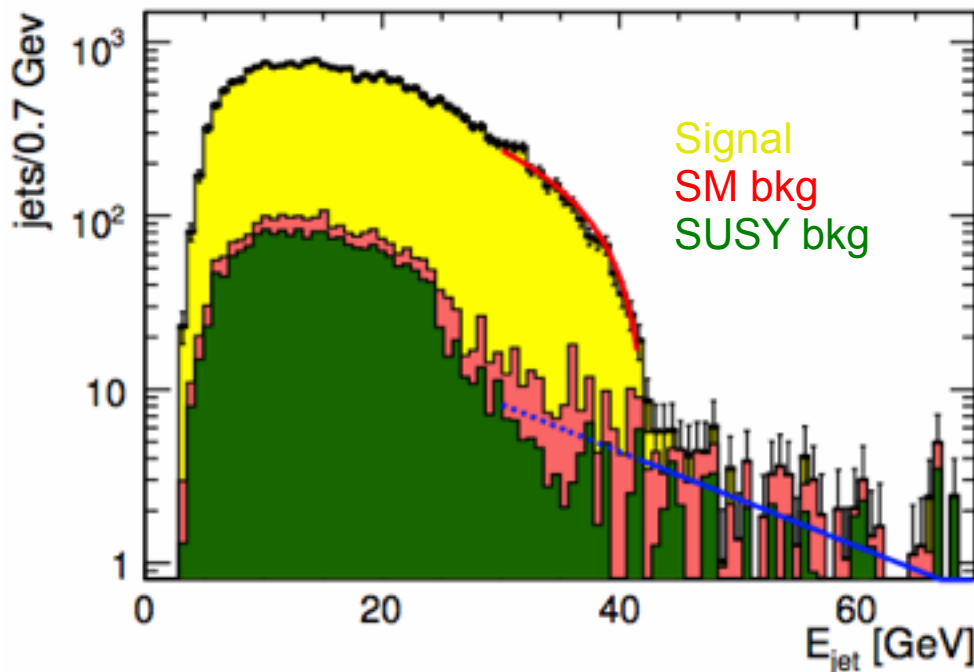
Observation of lighter and heavier stau states with decay to DM + hadronic tau

Benchmark point: $m(\text{LSP}) = 98 \text{ GeV}$, $m(\text{stau1}) = 108 \text{ GeV}$, $m(\text{stau2}) = 195 \text{ GeV}$

$$\sigma(e^+e^- \rightarrow \tilde{\tau}_1^+ \tilde{\tau}_1^-) = 158 \text{ fb}$$

$$\sigma(e^+e^- \rightarrow \tilde{\tau}_2^+ \tilde{\tau}_2^-) = 18 \text{ fb}$$

Bechtle, Berggren, List, Schade, Stempel, arXiv:0908.0876, PRD82, 055016 (2010)



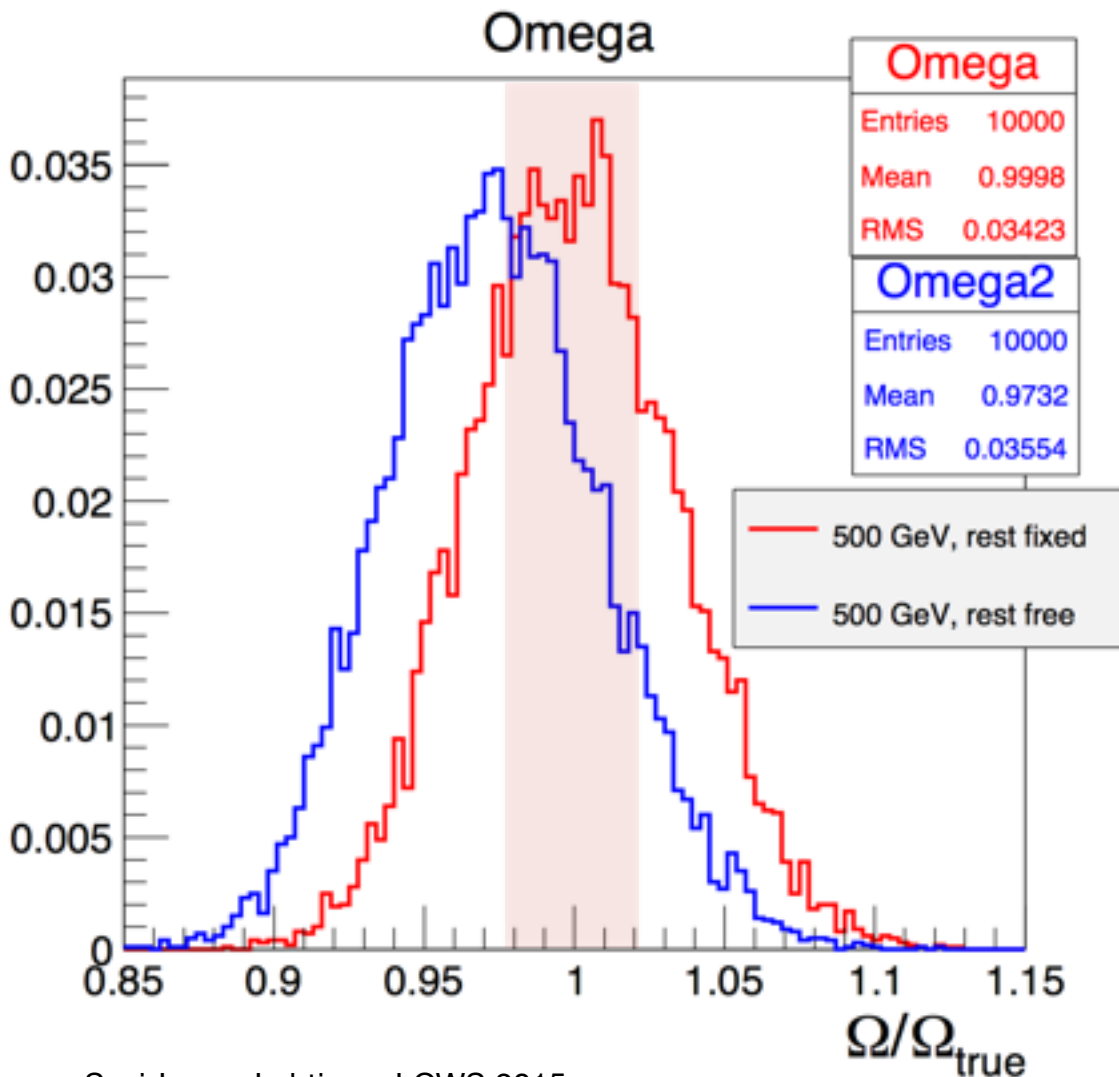
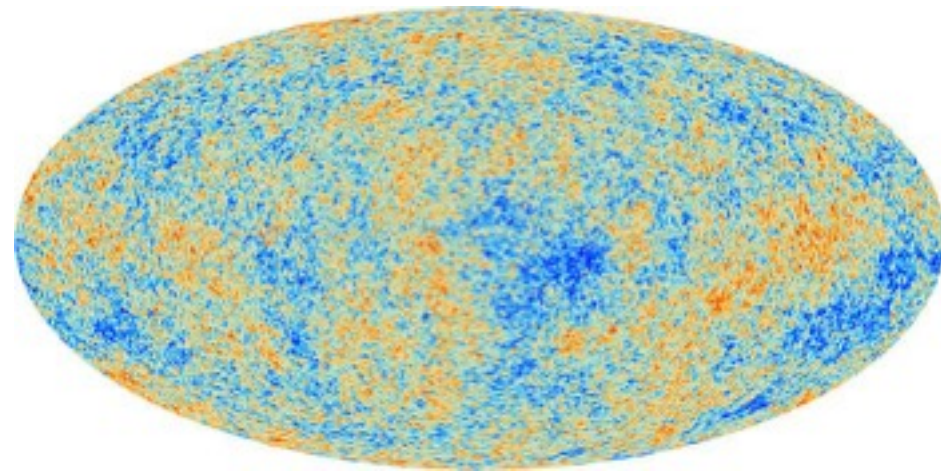
$\sqrt{s}=500 \text{ GeV}$, $\text{Lumi}=500 \text{ fb}^{-1}$, $P(e^-,e^+)=(+0.8,-0.3)$
 Stau1 mass $\sim 0.1\%$, Stau2 mass $\sim 3\%$ \rightarrow LSP mass $\sim 1.7\%$

DM Relic Abundance

WMAP/Planck (68% CL)

$$\Omega_c h^2 = 0.1196 \pm 0.0027$$

ESA/Planck



Once a DM candidate is discovered, crucial to check the consistency with the measured DM relic abundance.

Mass and couplings measured at ILC

→ **DM relic density to compare with the CMB data**

Other Probes

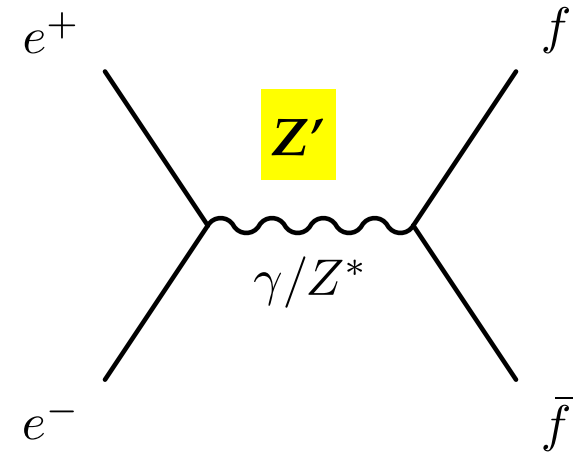
Z'

Z' : Heavy Neutral Gauge Bosons

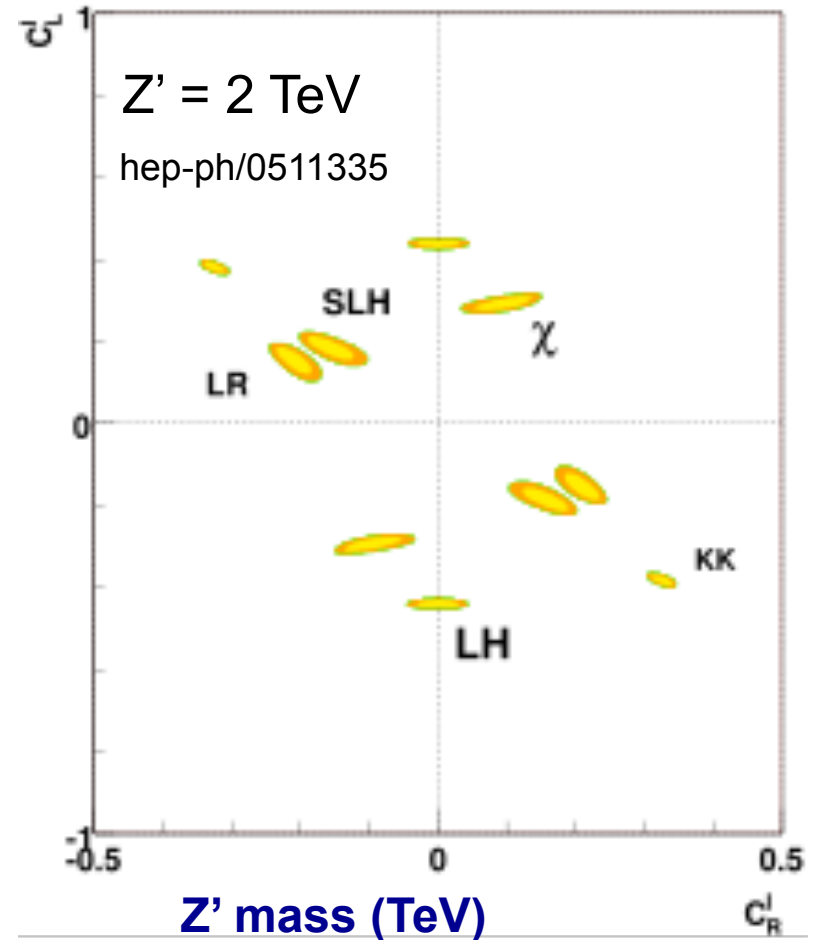
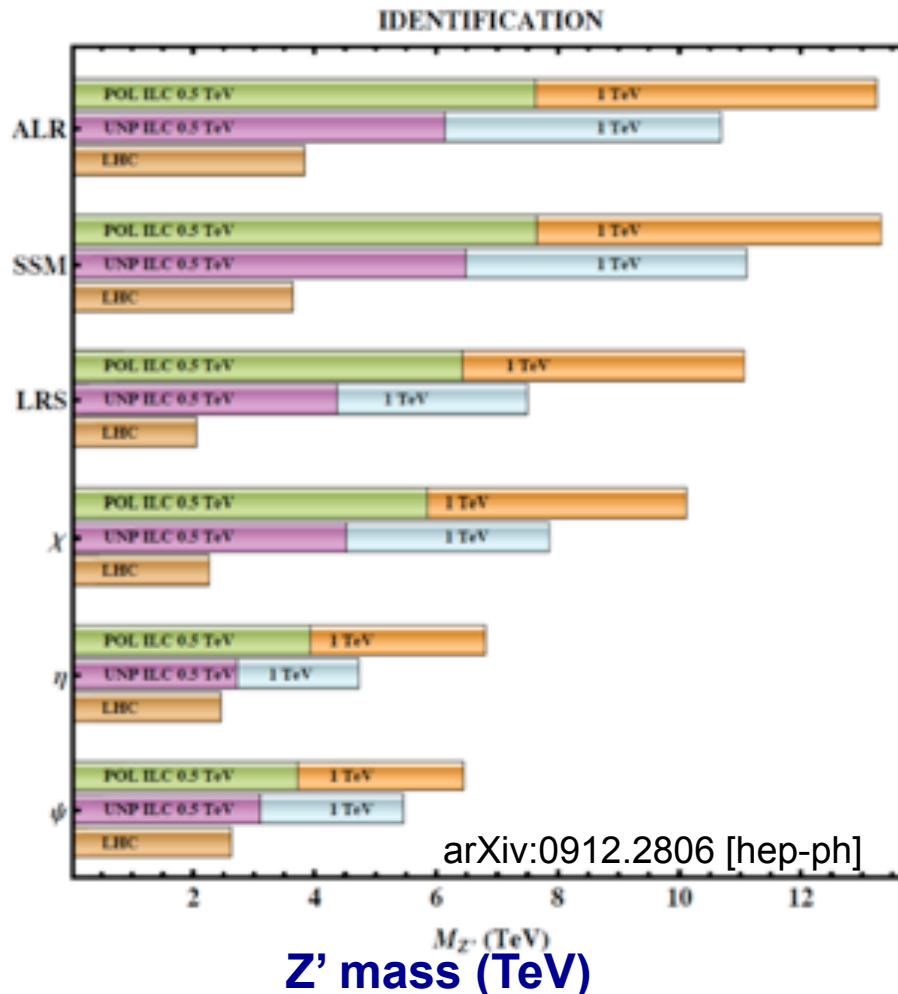
New gauge forces imply existence of heavy gauge bosons (Z')

Complementary approaches LHC/ILC

- LHC: Direct searches for Z' (mass determination)
- ILC: Indirect searches via interference effects (coupling measurements and model discrimination) – beam polarizations improve reach and discrimination power



Models with Z' boson



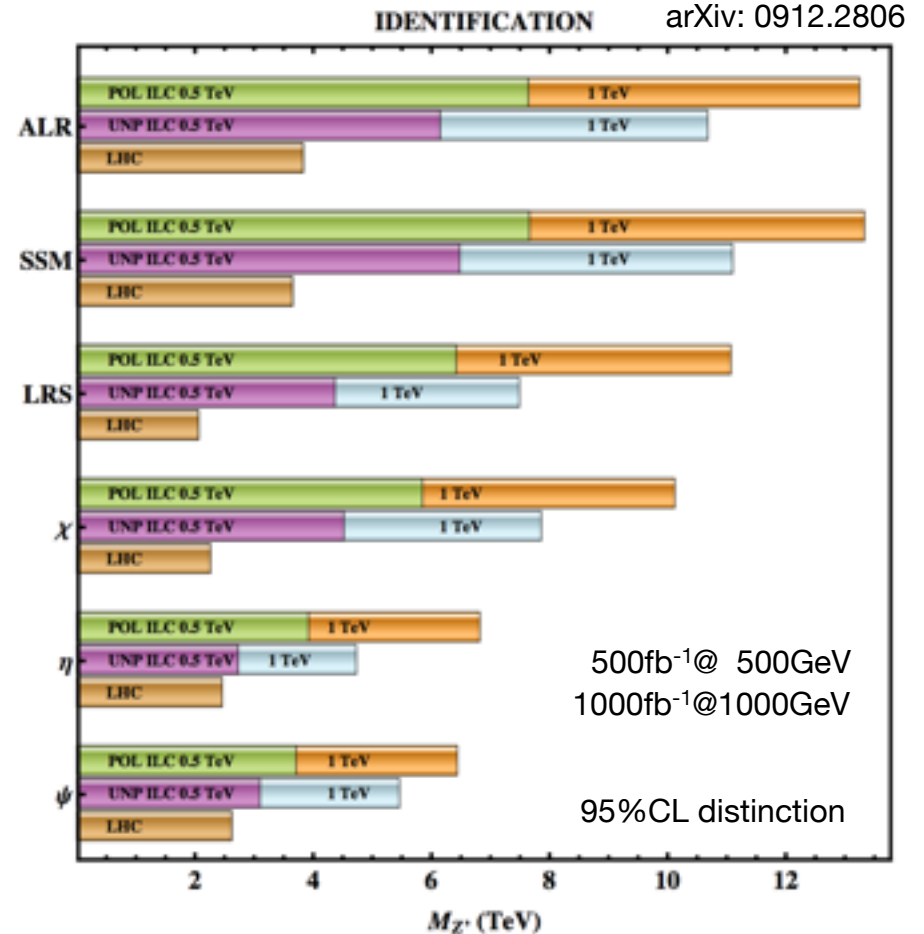
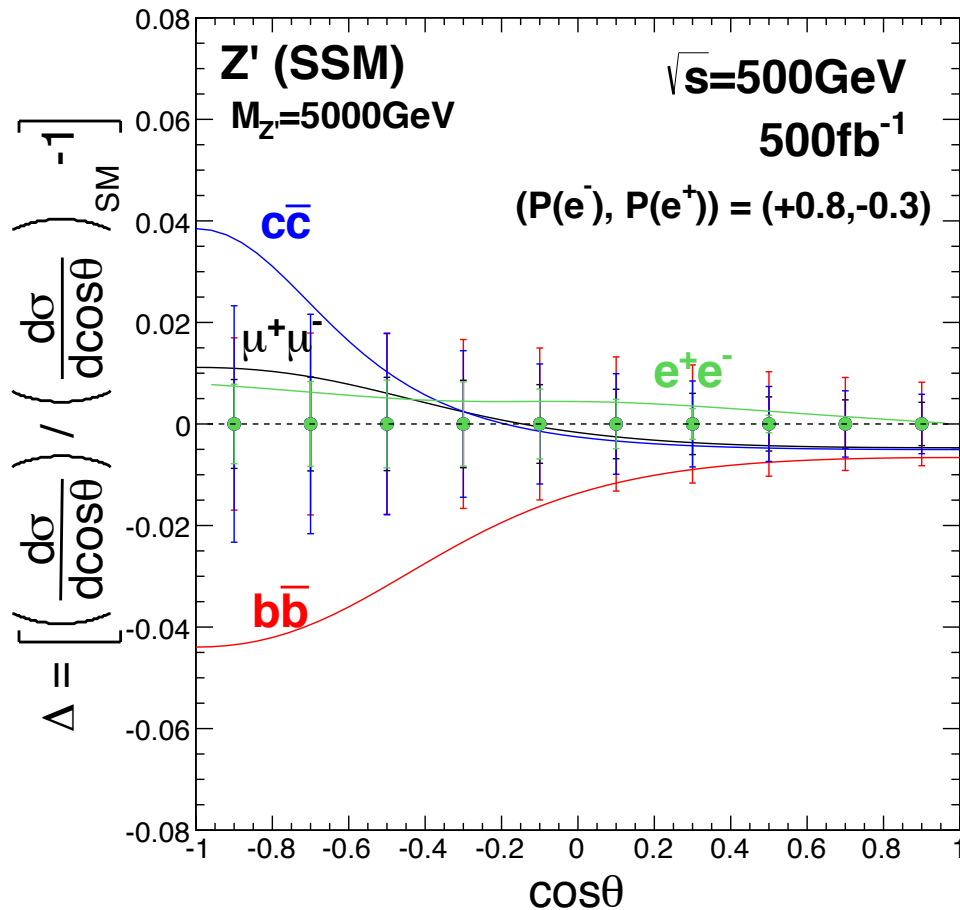
Two-Fermion Processes

Z' Search / Study

Observables: $d\sigma(P^-, P^+)/d \cos\theta$

$$\chi^2 = \sum_f \sum_{P^-, P^+} \sum_{i \in \text{bins}} \frac{|n_i(SM + Z') - n_i(SM)|^2}{\Delta n_i} \quad (f=e, \mu, \tau, c, b)$$

Example: Sequential SM-like Z'



Two-Fermion Processes

Z' Search / Study

arXiv:0912.2806 [hep-ph]

hep-ph/0511335

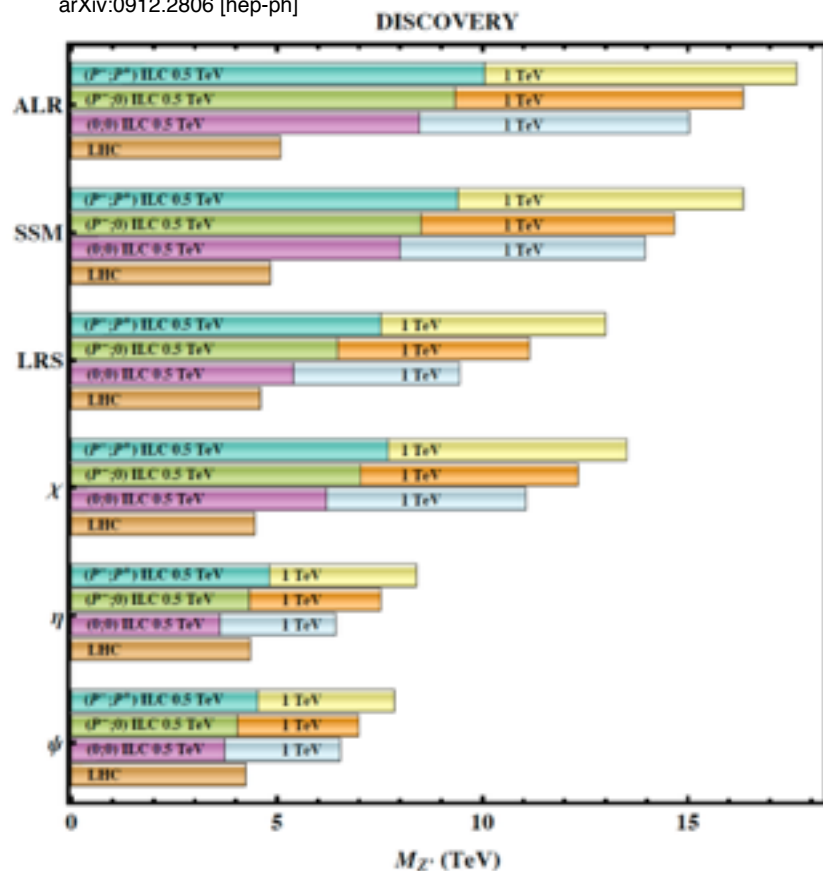
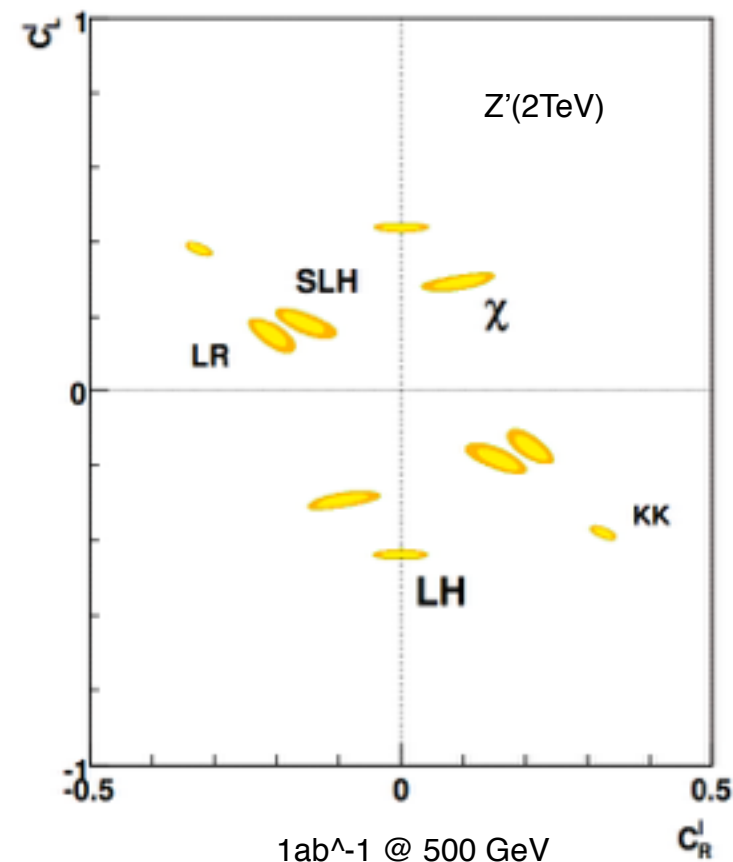


Figure 23: Sensitivity of the ILC to various candidate Z' bosons, quoted at 95% conf., with $\sqrt{s} = 0.5$ (1.0) TeV and $\mathcal{L}_{int} = 500$ (1000) fb^{-1} . The sensitivity of the LHC-14 via Drell-Yan process $pp \rightarrow \ell^+\ell^- + X$ with 100 fb^{-1} of data are shown for comparison. For details, see [14].



ILC's Model ID capability is expected to exceed that of LHC even if we cannot hit the Z' pole.

Beam polarization is essential to sort out various possibilities.

Two-Fermion Processes

Compositeness

S. Riemann, LC-TH-2001-007

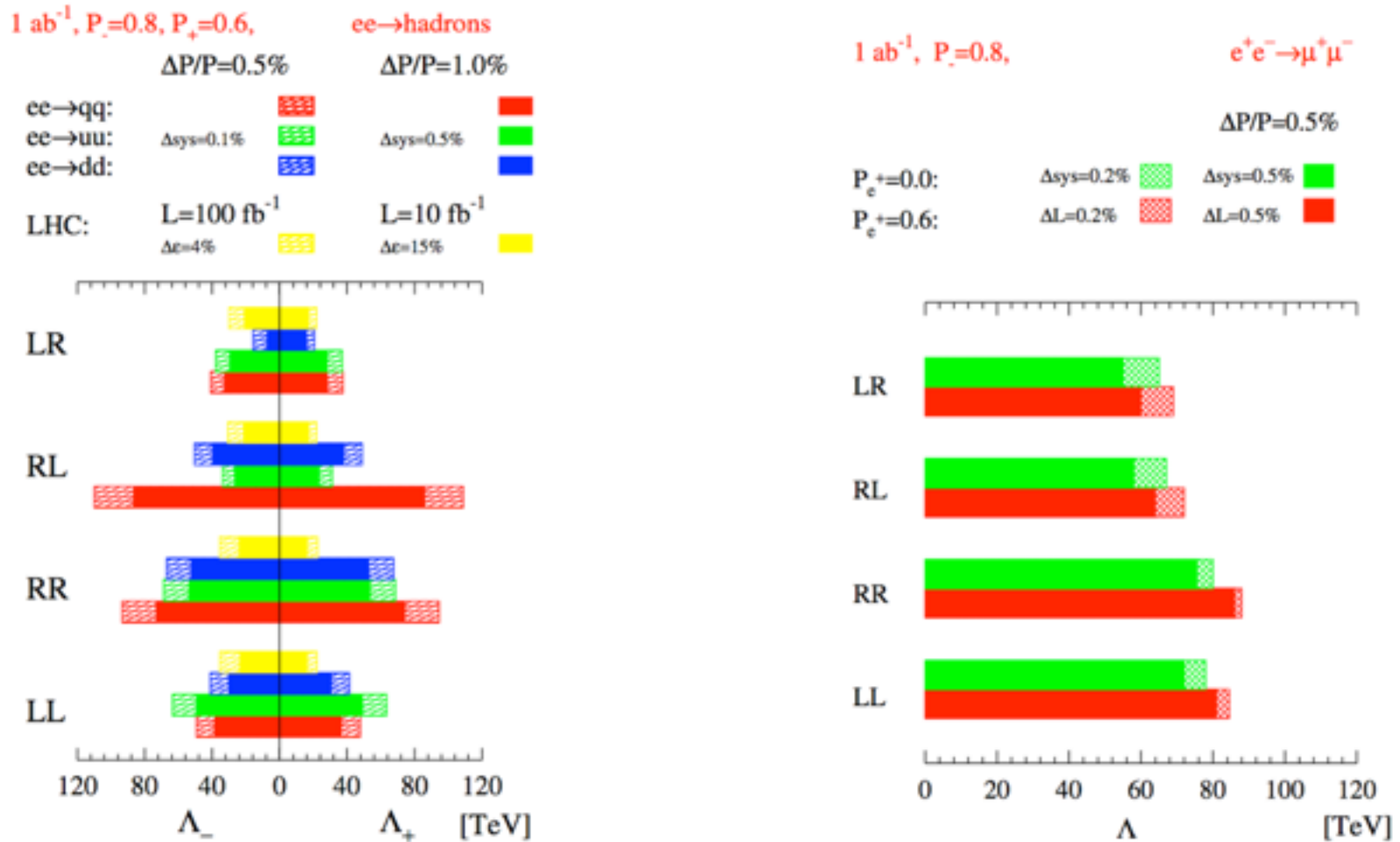


Figure 26: Sensitivities (95% c.l.) of a 500 GeV ILC to contact interaction scales Λ for different helicities in $e^+e^- \rightarrow \text{hadrons}$ (left) and $e^+e^- \rightarrow \mu^+\mu^-$ (right), including beam polarization [18].

Beam polarization is essential to sort out various possibilities.