



Automated BSM at NLO

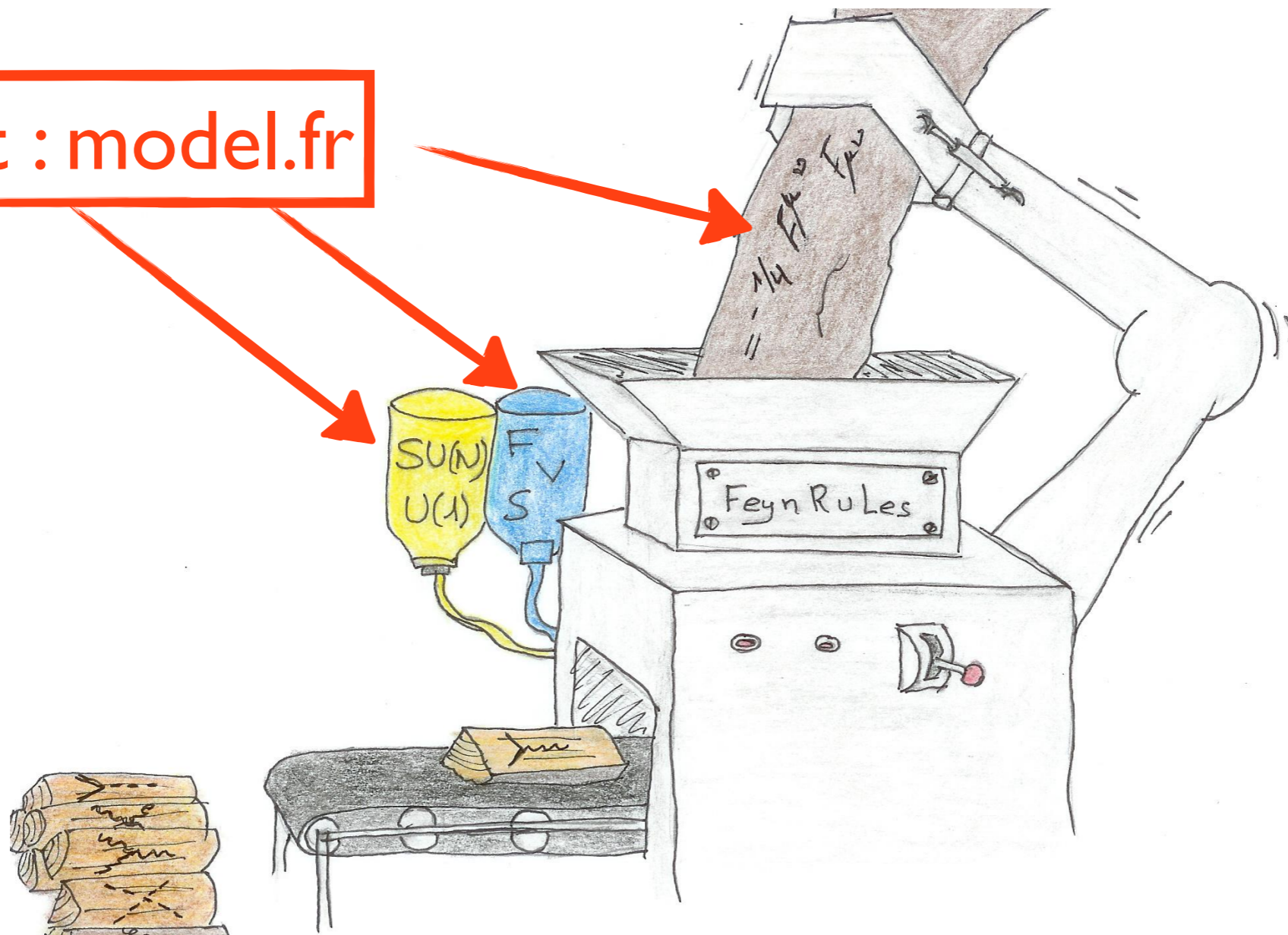
Celine Degrande (IPPP, Durham University)
MC4BSM 2016
Beijing

Plan

- FeynRules in a nutshell
- BSM@NLO
 - Ingredients
 - How does it work?
- Examples :
 - Charged Higgs production
 - SUSY QCD
 - ...

FeynRules

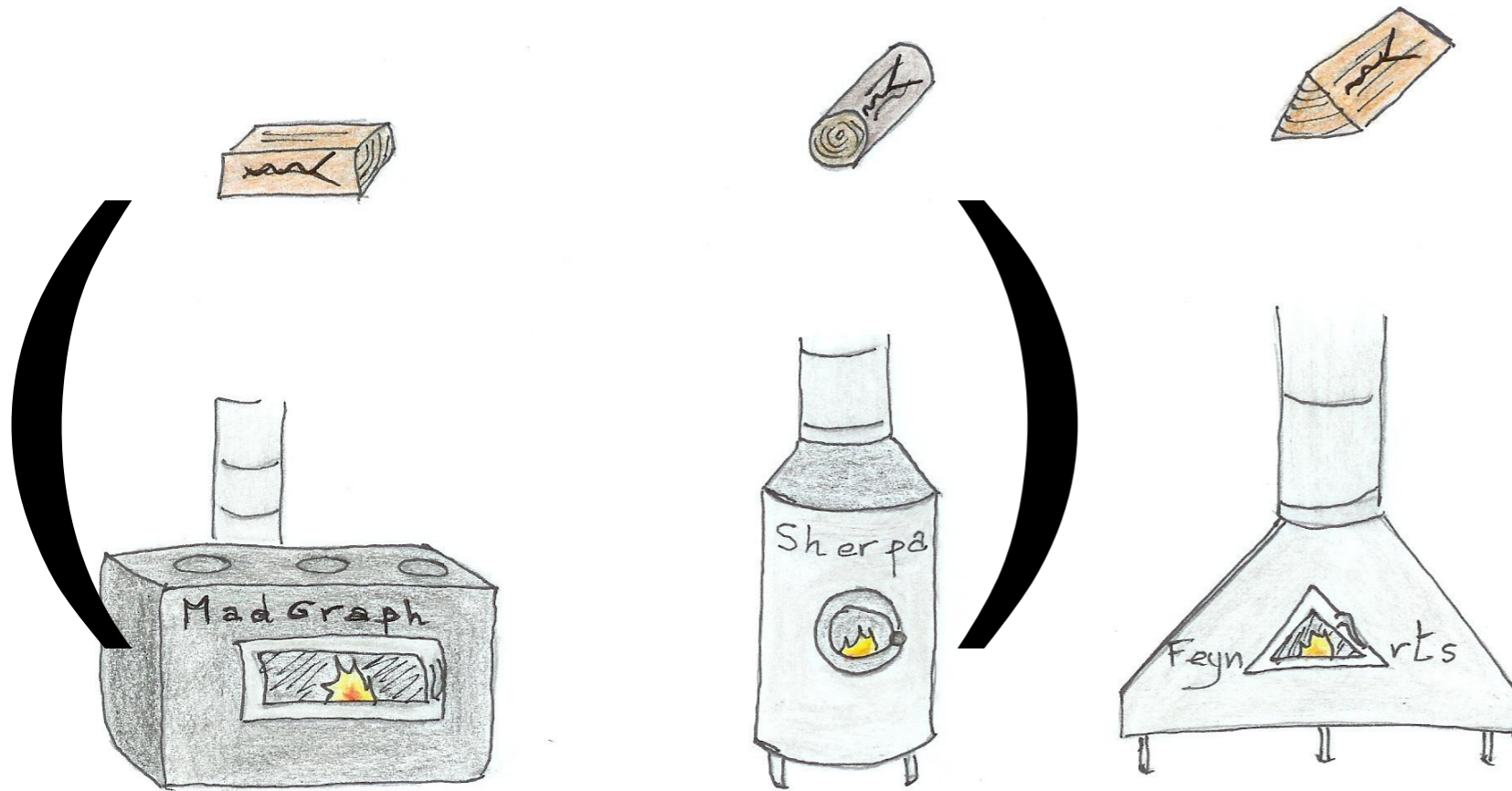
Input : model.fr



Output : vertices

A. Alloul, N. D. Christensen, CD, C. Duhr and B. Fuks, CPC185 (2014) 2250

FeynRules outputs



FeynRules outputs
can be used
directly by event
generators

UFO : output with the
full information
used by several
generators



UFO

- Generator independent output with full model information
- Contains the list of particles, parameters, vertices, decays (1 to 2), coupling orders
- vertices are split into **Lorentz structures**, **colours** and **couplings** and all are included in the model!

$$-ig_s T_{ij}^a \gamma_\mu$$

- Used in MG5, Herwig, Gosam, Sherpa

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Why NLO/Loop BSM?

- Discovery :
 - Loop-induced production/decay
 - NLO : Refine search strategies
- Measurement of properties/couplings : NLO corrections
 - QCD corrections are large at the LHC
- Quantification of the constraints NP should not be limited by the th. error on EFT

Madgraph5_aMC@NLO

Wide BSM support

+

Automated NLO computation

- Computation of the born

- Computation of the real

- Computation of the loop

- Matching with parton shower 'à la' MC@NLO

MG5

MadFKS (IR)

MadLoop

MadLoop

$$\mathcal{A}^{1-loop} = \sum_i d_i \text{Box}_i + \sum_i c_i \text{Triangle}_i + \sum_i b_i \text{Bubble}_i \\ + \sum_i a_i \text{Tadpole}_i + R$$

- Box, Triangle, Bubble and Tadpole are known scalar integrals
- Loop computation = find the coefficients
 - Tensor reduction (OPP)
- R : rational terms should be partially provided
- UV counterterm vertices have to be provided

To be provided : R_2

$$\bar{A}(\bar{q}) = \frac{1}{(2\pi)^4} \int d^d \bar{q} \frac{\bar{N}(\bar{q})}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}}, \quad \bar{D}_i = (\bar{q} + p_i)^2 - m_i^2$$

$$\bar{N}(\bar{q}) = N(q) + \tilde{N}(\tilde{q}, q, \epsilon)$$

The diagram shows the decomposition of the numerator $\bar{N}(\bar{q})$ into two terms: $N(q)$ and $\tilde{N}(\tilde{q}, q, \epsilon)$. Each term is enclosed in a red circle. Red arrows point from each circle to a variable below it: d for $\bar{N}(\bar{q})$, 4 for $N(q)$, and ϵ for $\tilde{N}(\tilde{q}, q, \epsilon)$.

$$R_2 \equiv \lim_{\epsilon \rightarrow 0} \frac{1}{(2\pi)^4} \int d^d \bar{q} \frac{\tilde{N}(\tilde{q}, q, \epsilon)}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}}$$

Finite set of vertices that can be computed once
for all

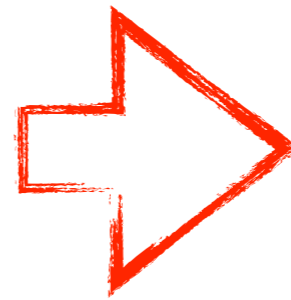
Computed in MadLoop :R₁

Due to the ϵ dimensional parts of the denominators

Like for the 4 dimensional part but with a different set of integrals

$$\int d^n \bar{q} \frac{\tilde{q}^2}{\bar{D}_i \bar{D}_j} = -\frac{i\pi^2}{2} \left[m_i^2 + m_j^2 - \frac{(p_i - p_j)^2}{3} \right] + \mathcal{O}(\epsilon),$$
$$\int d^n \bar{q} \frac{\tilde{q}^2}{\bar{D}_i \bar{D}_j \bar{D}_k} = -\frac{i\pi^2}{2} + \mathcal{O}(\epsilon),$$
$$\int d^n \bar{q} \frac{\tilde{q}^4}{\bar{D}_i \bar{D}_j \bar{D}_k \bar{D}_l} = -\frac{i\pi^2}{6} + \mathcal{O}(\epsilon).$$

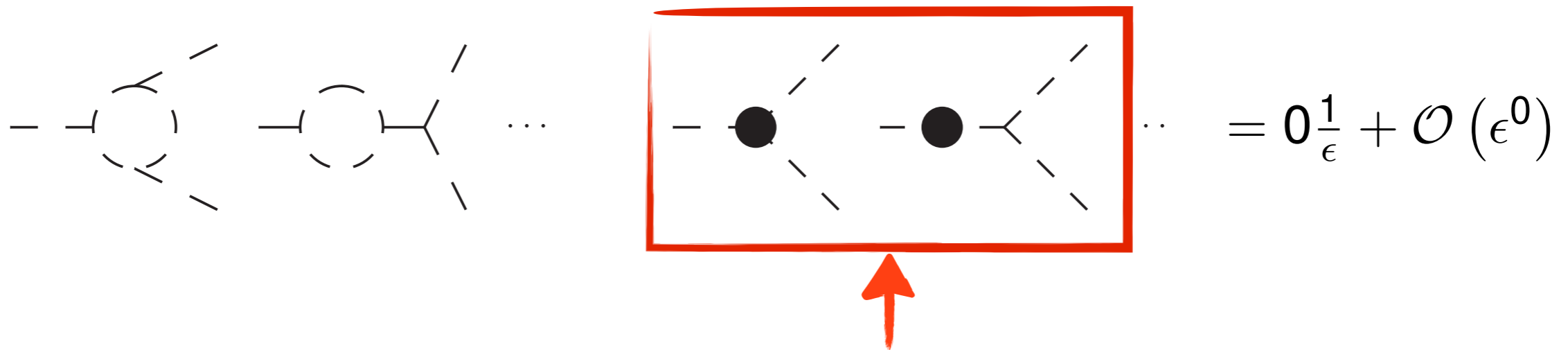
Only $R = R_1 + R_2$ is gauge invariant



Check

UV

$$\bar{A}(\bar{q}) = \frac{1}{(2\pi)^4} \int d^d \bar{q} \frac{\bar{N}(\bar{q})}{\bar{D}_0 \bar{D}_1 \dots \bar{D}_{m-1}} = K \frac{1}{\epsilon} + \mathcal{O}(\epsilon^0)$$



Relations fixed by the Lagrangian (finite part)

Finite set of vertices that can be computed once
for all

Renormalization

External parameters

$$\begin{aligned}x_0 &\rightarrow x + \delta x, \\ \phi_0 &\rightarrow \left(1 + \frac{1}{2}\delta Z_{\phi\phi}\right)\phi + \sum_x \frac{1}{2}\delta Z_{\phi\chi}\chi.\end{aligned}$$

Same for the conjugate field

Internal parameters are renormalised by replacing the external parameters in their expressions

$$\begin{aligned}gg & (1 + \delta Z_{gg}) TL \\ ggg & \left(1 + \frac{1}{2}\delta\alpha_s + \frac{3}{2}\delta Z_{gg}\right) TL \\ gggg & \left(1 + \delta\alpha_s + 2\delta Z_{gg}\right) TL\end{aligned}$$

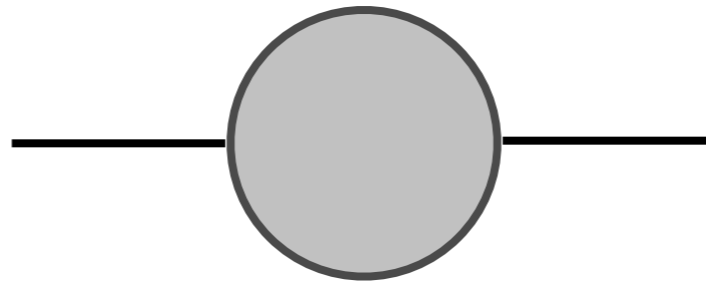
Fixed by

Renormalization conditions

On-shell scheme (or **complex mass** scheme):

Renormalized mass = Physical mass

Two-point function vanishes on-shell (No external bubbles)



$$i\delta_{ij} (\not{p} - m_i) + i [f_{ij}^L(p^2) \not{p}\gamma_- + f_{ij}^R(p^2) \not{p}\gamma_+ + f_{ij}^{SL}(p^2) \gamma_- + f_{ij}^{SR}(p^2) \gamma_+]$$

$$\cancel{\tilde{\mathcal{L}}} [f_{ij}^L(p^2) m_i + f_{ij}^{SR}(p^2)] \Big|_{p^2=m_i^2} = 0$$

$$\cancel{\tilde{\mathcal{L}}} [f_{ij}^R(p^2) m_i + f_{ij}^{SL}(p^2)] \Big|_{p^2=m_i^2} = 0$$

$$\cancel{\tilde{\mathcal{L}}} \left[2m_i \frac{\partial}{\partial p^2} [(f_{ii}^L(p^2) + f_{ii}^R(p^2)) m_i + f_{ii}^{SL}(p^2) + f_{ii}^{SR}(p^2)] + f_{ii}^L(p^2) + f_{ii}^R(p^2) \right] \Big|_{p^2=m_i^2} = 0$$

Similar for the vectors and scalars

How does it work?

FeynRules
Renormalize the Lagrangian

model.mod
model.gen

FeynArts
Write the amplitudes

NLOCT.m
Compute the NLO vertices

model.nlo



CD, CPC 197 (2015) 239

Restrictions/Assumptions

- Renormalizable Lagrangian, maximum dimension of the operators is 4
- Feynman Gauge
- $\{\gamma_\mu, \gamma_5\} = 0$
- 't Hooft-Veltman scheme
- On-shell scheme for the masses and wave functions
- $\overline{\text{MS}}$ by default for everything else (zero-momentum possible for fermion gauge boson interaction)

R2 : Validation

- tested* on the SM (QCD:P. Draggiotis et al. +QED:M.V. Garzelli et al)
- tested* on MSSM (QCD:H.-S. Shao, Y.-J. Zhang) : test the Majorana

*Analytic comparison of the expressions

UV Validation

- SM QCD : tested* (W. Beenakker, S. Dittmaier, M. Kramer, B. Plumper)
- SM EW : tested* (expressions given by H.-S. Shao from A. Denner)

*Analytic comparison of the expressions

Tests in event generators

- aMC@NLO
- The SM QCD has been tested by V. Hirschi (Comparison with the built-in version)
- SM EW (MZ scheme): comparison to published results for ME by H.-S. Shao and V. Hirschi
- Various BSM
 - gauge invariance
 - pole cancelation

Test EW

== a a > t t~ ['QED'] ==
== a a > t t~ a ['QED'] ==
== a a > w+ w- ['QED'] ==
== a b > t w- ['QED'] ==
== d~ d > w+ w- ['QCD'] ==
== d~ d > w+ w- ['QED'] ==
== d~ d > z z ['QCD'] ==
== d~ d > z z ['QED'] ==
== e+ e- > t t~ a ['QED'] ==
== e+ e- > t t~ g ['QED'] ==
== g b > t w- ['QED'] ==
== g g > h h ['QCD'] ==
== g g > t t~ ['QED'] ==
== g g > t t~ g ['QED'] ==
== g g > t t~ h ['QCD'] ==
== g g > t t~ h ['QED'] ==
== h h > h h ['QED'] ==
== h h > h h h ['QED'] ==
== t t~ > w+ w- ['QED'] ==

== u b > t d ['QED'] ==
== u d~ > t b~ ['QED'] ==
== u g > t d b~ ['QED'] ==
== u u~ > a a ['QED'] ==
== u u~ > e+ e- ['QED'] ==
== u u~ > g a ['QCD QED'] ==
== u u~ > u u~ ['QCD QED'] ==
== u u~ > u u~ a ['QCD QED'] ==
== u u~ > u u~ g ['QCD QED'] ==
== u u~ > w+ w- ['QED'] ==
== u u~ > z a ['QED'] ==
== u u~ > z z ['QED'] ==
== u~ d > w- z ['QCD'] ==
== u~ d > w- z ['QED'] ==
== u~ u > w+ w- ['QCD'] ==
== u~ u > w+ w- ['QED'] ==
== u~ u > z z ['QCD'] ==
== u~ u > z z ['QED'] ==
== ve ve~ > e+ e- ['QED'] ==
== w+ w- > h h ['QED'] ==

Massive and massless b

Future development

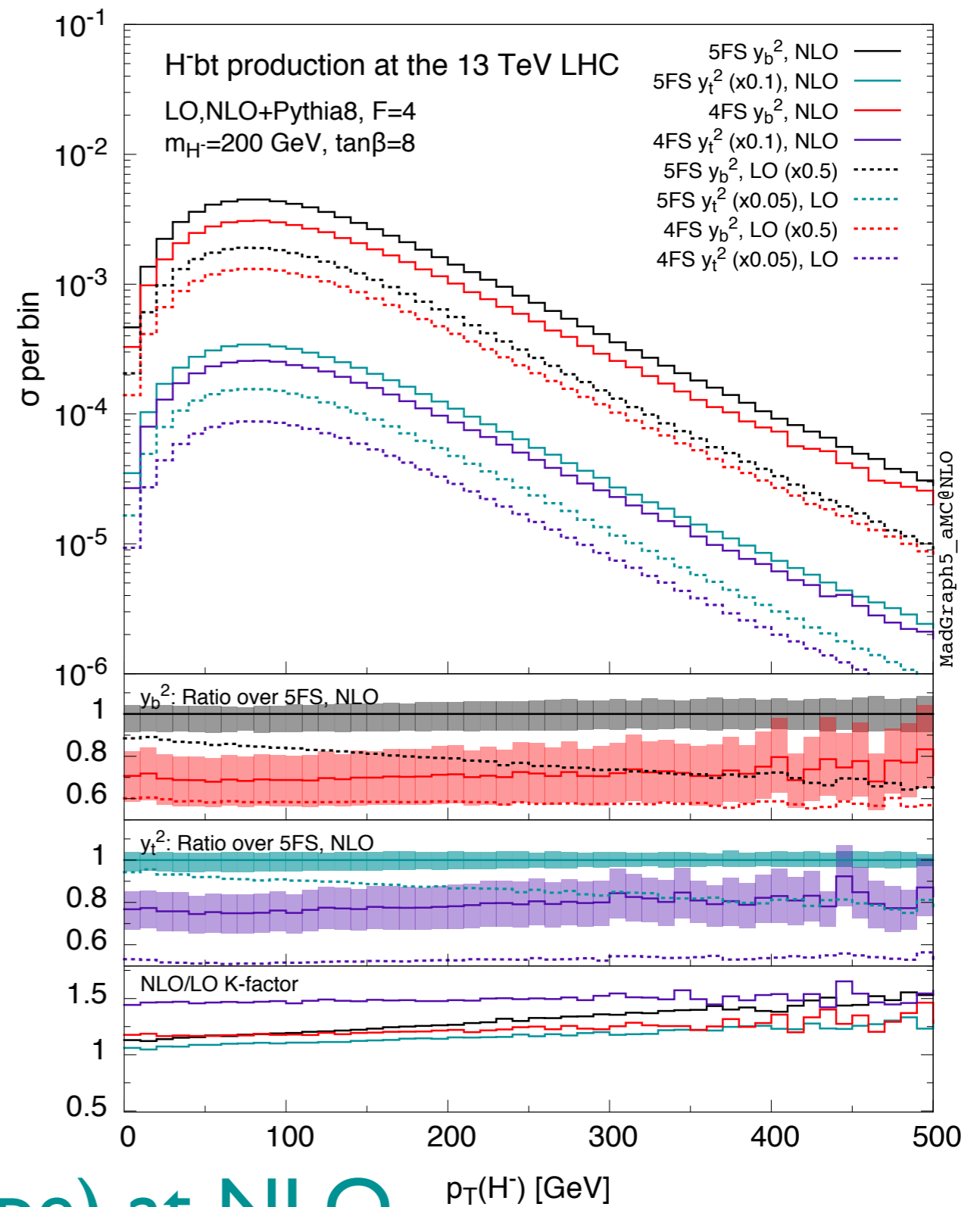
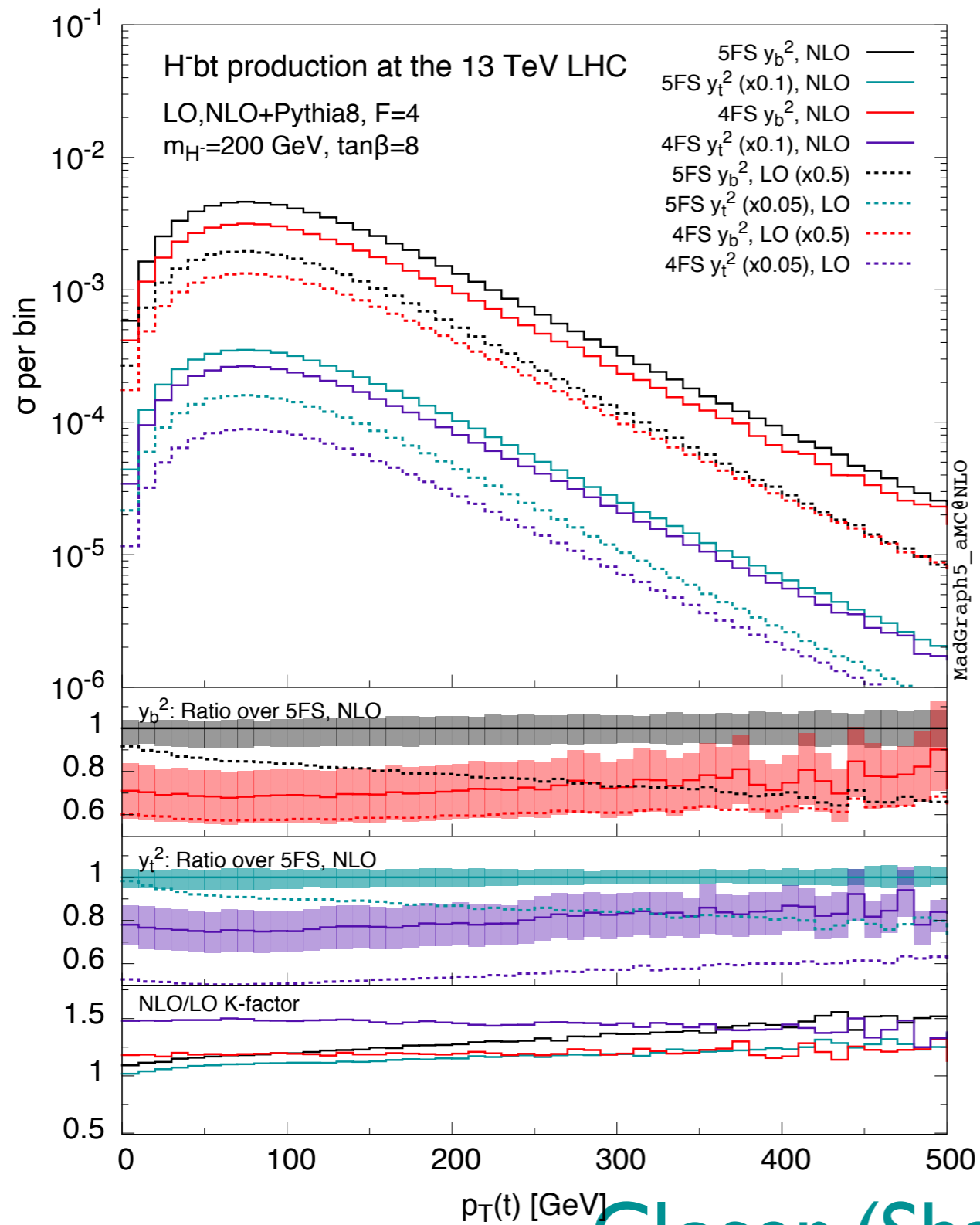
- EFT : done but 4F operators (in progress)
 - any gauge
- UFO@NLO in Gosam (N. Greiner)
- DRED (asked by Gosam)
- UFO 2.0

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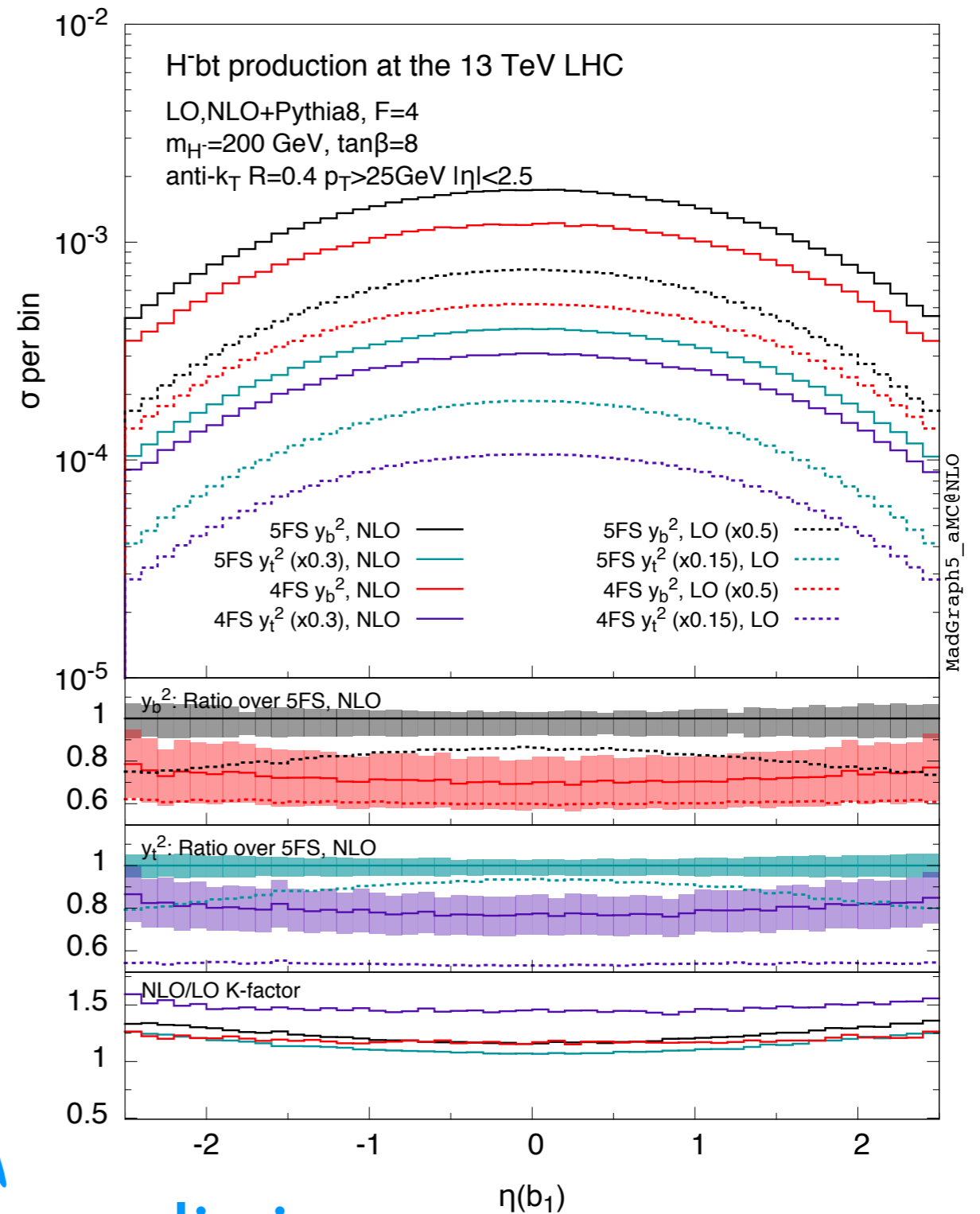
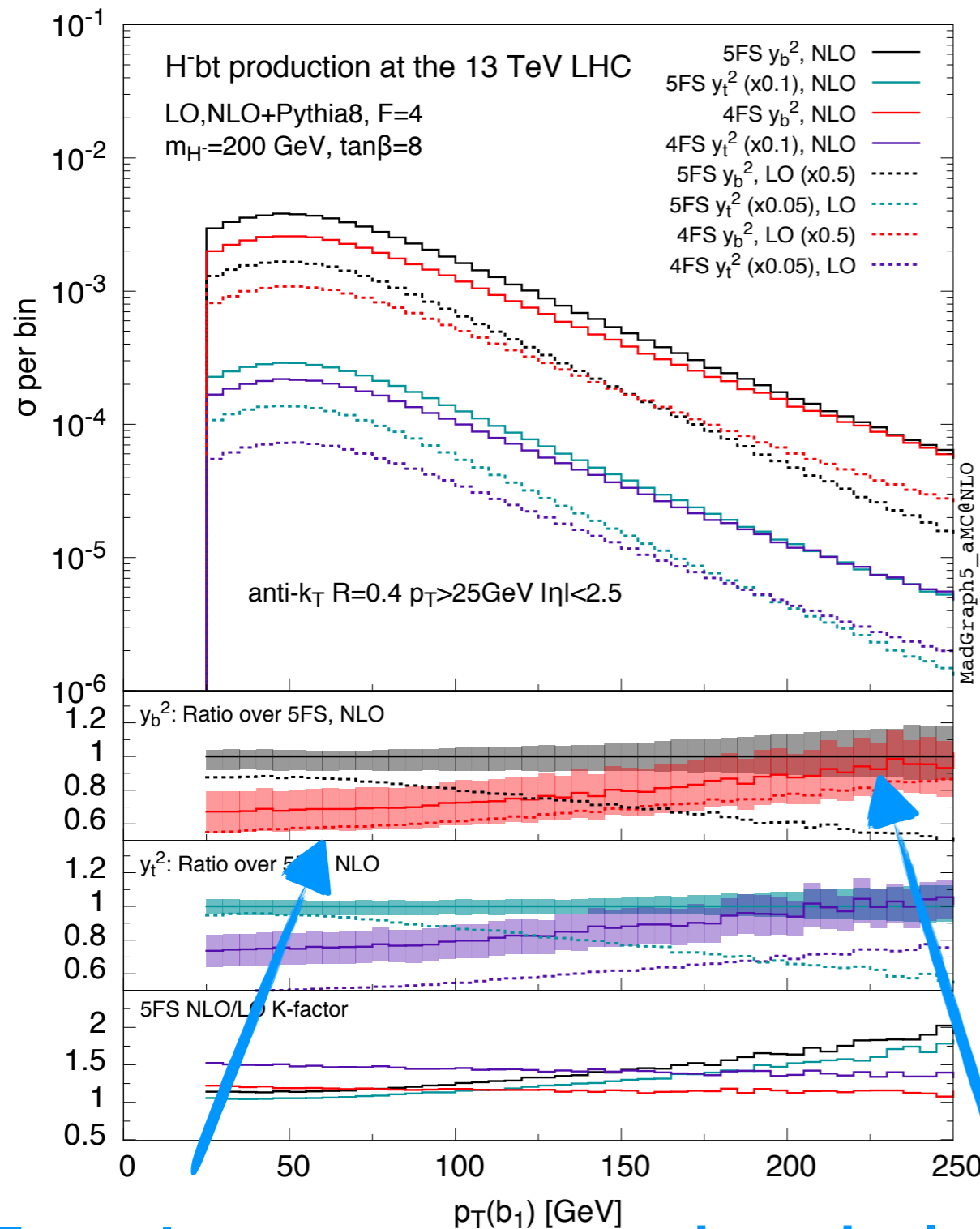
H⁺ production : 4F vs 5F

C. D., M. Ubiali, M. Wiesemann and M. Zaro, JHEP 1510 (2015) 145



Closer (Shape) at NLO

H⁺ production : 4F vs 5F

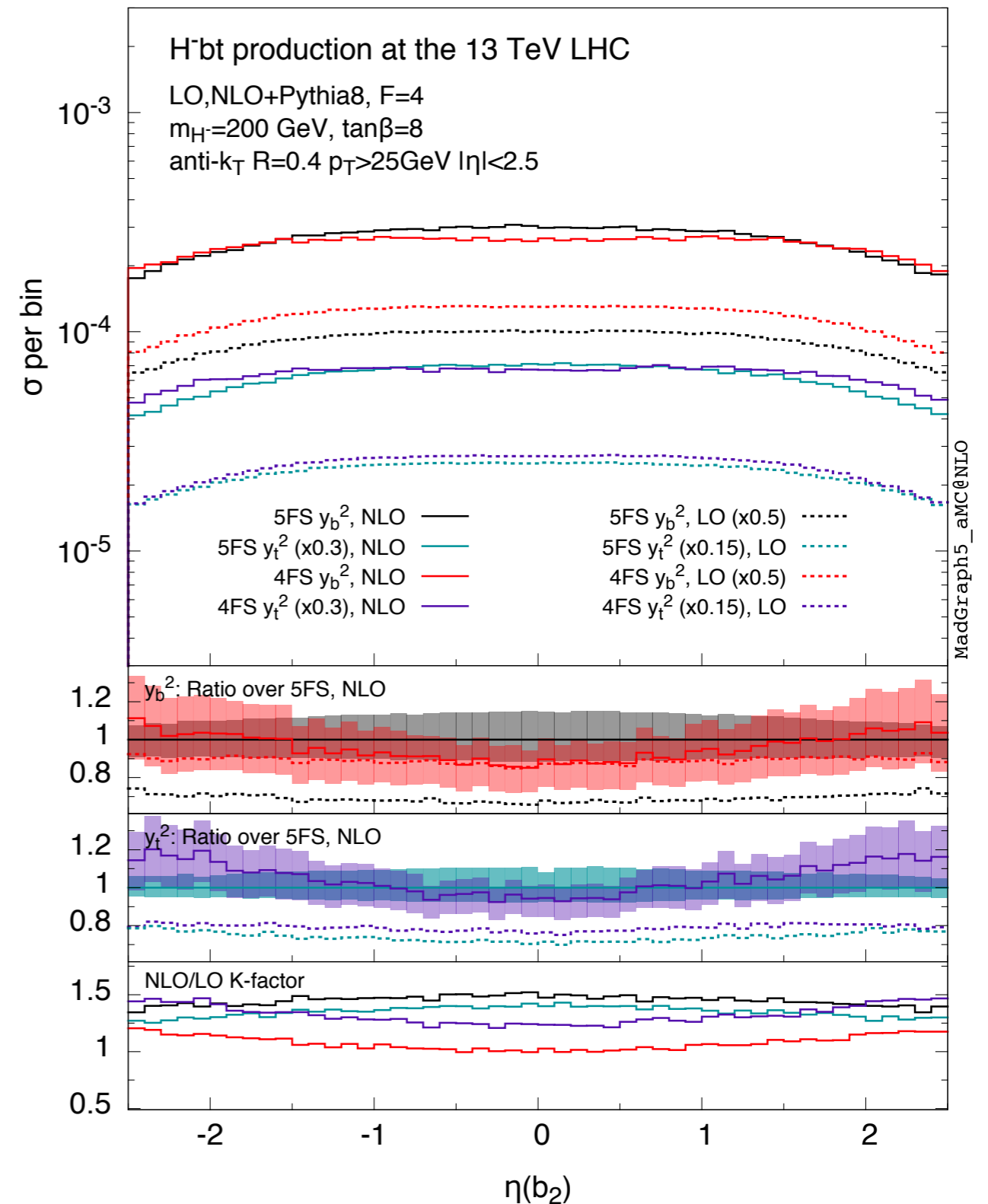
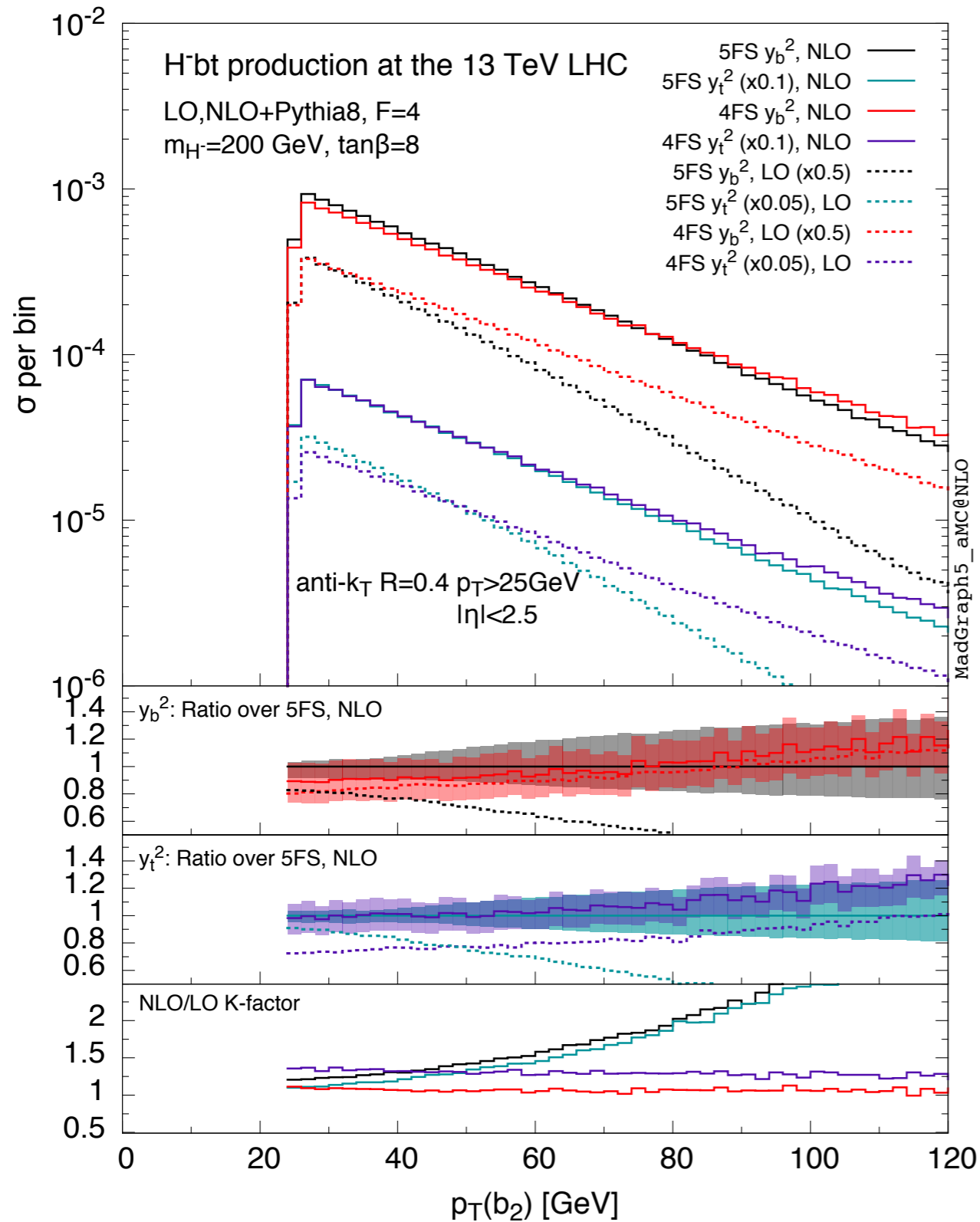


Top decay

hard gluon splitting

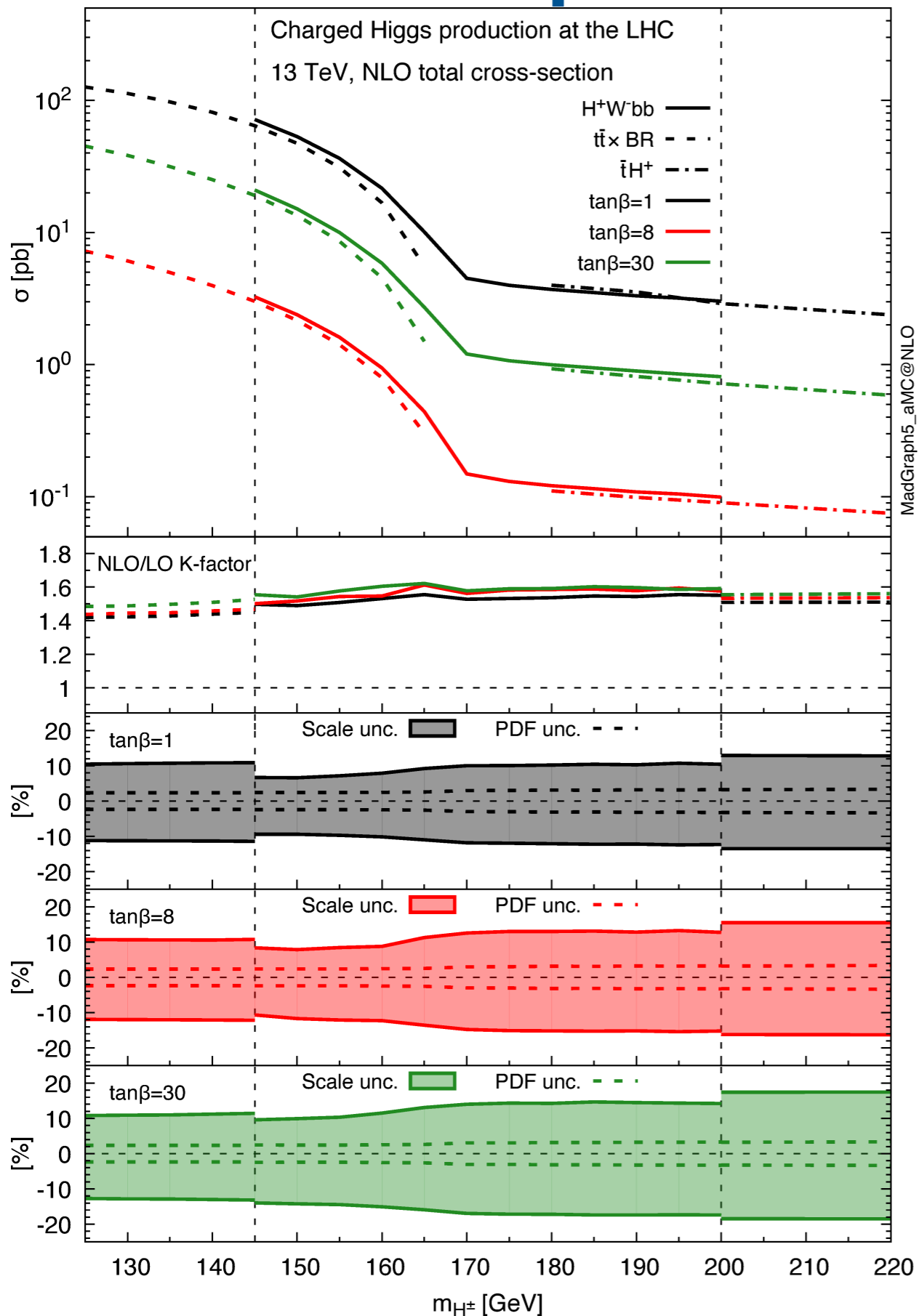
more exclusive in b more different

H⁺ production : 4F vs 5F



Only LO in 5F

H^+ production : $m_H \sim m_t$



1607.05291 : C. D., R.
Frederix, V. Hirschi, M. Ubiali,
M. Wiesemann and M. Zaro

Example II: Gluino pair production

CD, B. Fuks, V. Hirschi, J. Proudom and H. S. Shao, PLB 755, 82 (2016)

12 scalar triplets
no mixing

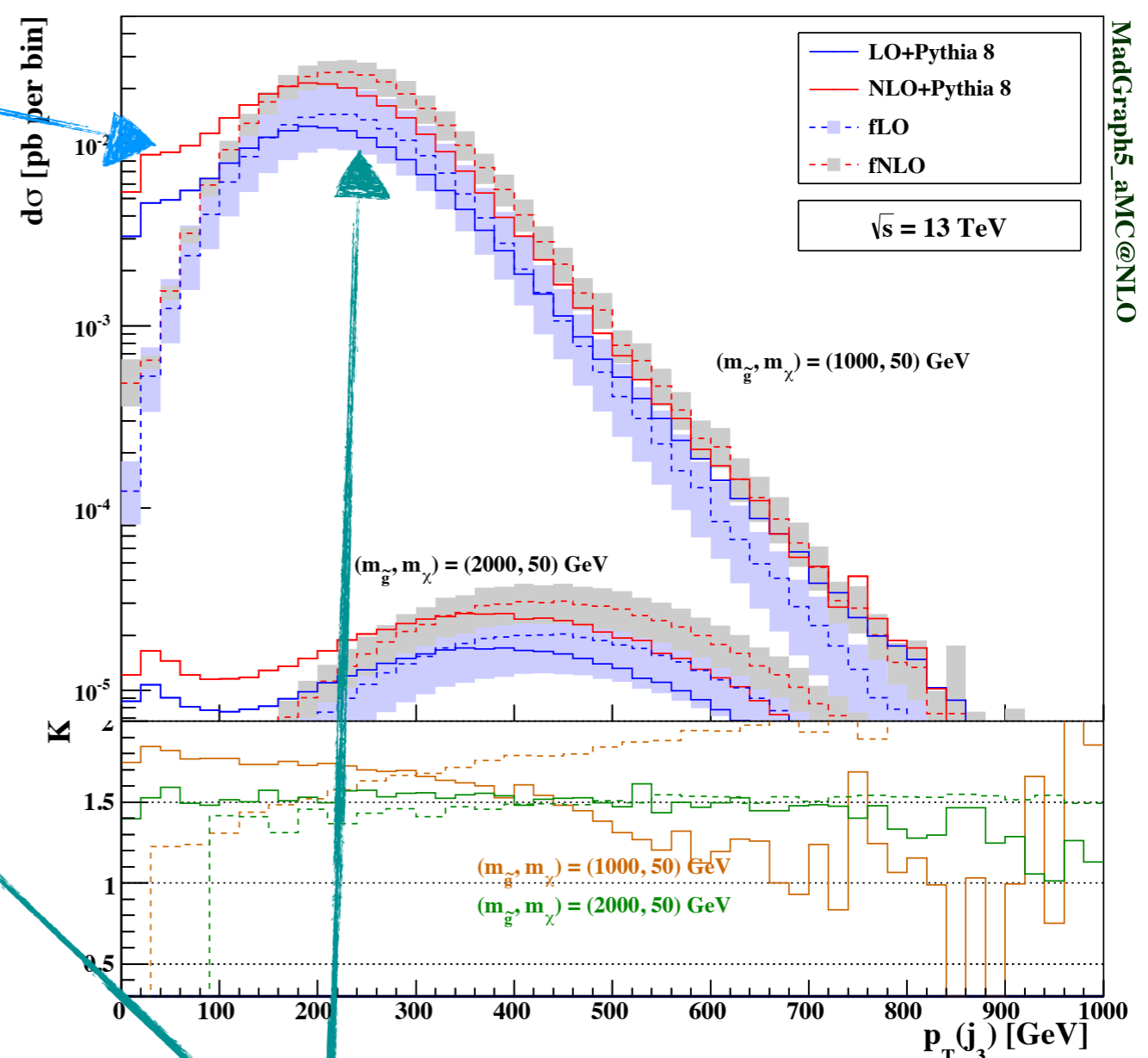
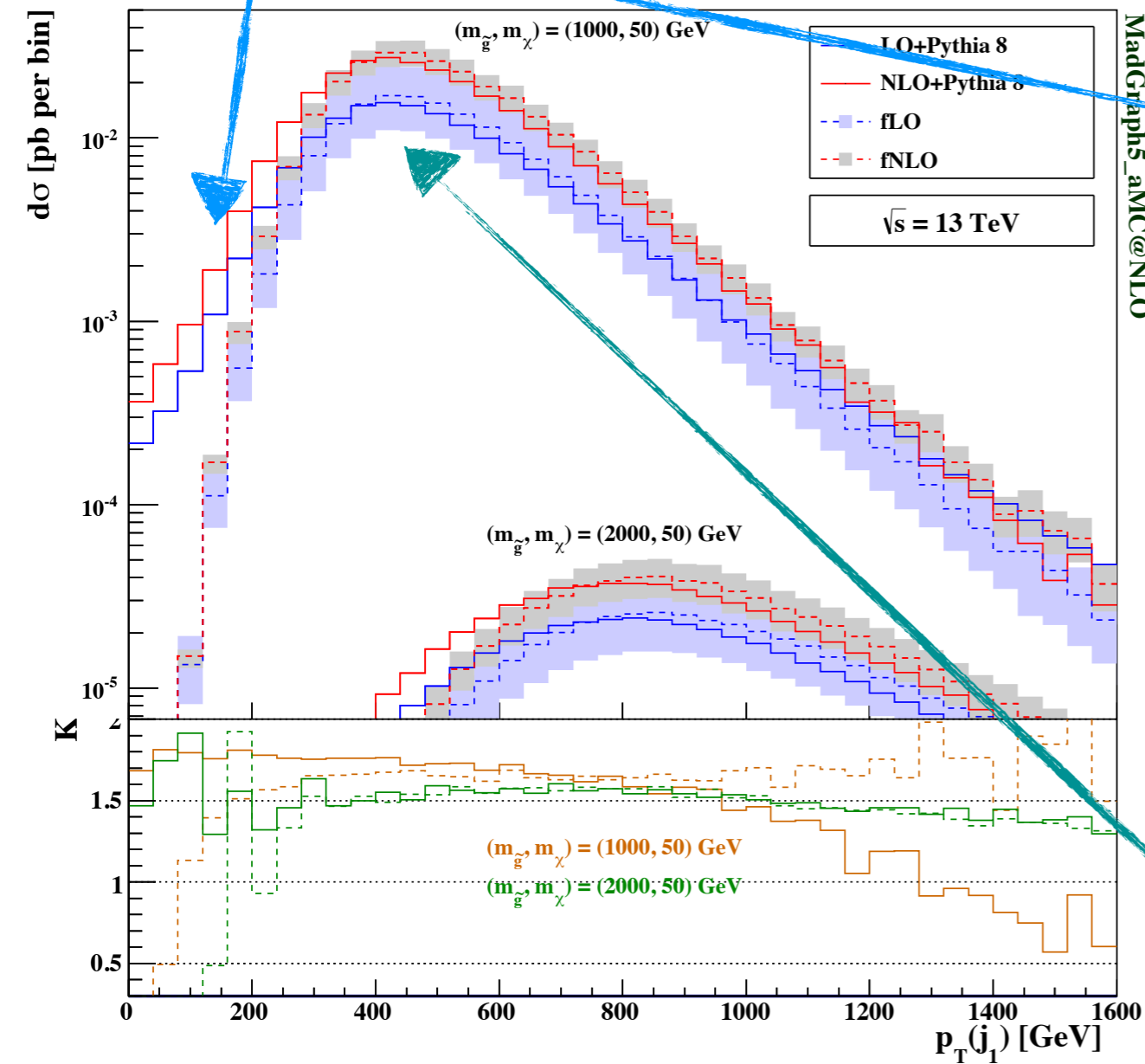
1 Majorana octet

$$\begin{aligned}\mathcal{L}_{\text{SQCD}} = & D_\mu \tilde{q}_L^\dagger D^\mu \tilde{q}_L + D_\mu \tilde{q}_R^\dagger D^\mu \tilde{q}_R + \frac{i}{2} \bar{g} \not{D} \tilde{g} \\ & - m_{\tilde{q}_L}^2 \tilde{q}_L^\dagger \tilde{q}_L - m_{\tilde{q}_R}^2 \tilde{q}_R^\dagger \tilde{q}_R - \frac{1}{2} m_{\tilde{g}} \bar{g} \tilde{g} \\ & + \sqrt{2} g_s \left[- \tilde{q}_L^\dagger T (\bar{g} P_L q) + (\bar{q} P_L \tilde{g}) T \tilde{q}_R + \text{h.c.} \right] \\ & - \frac{g_s^2}{2} \left[\tilde{q}_R^\dagger T \tilde{q}_R - \tilde{q}_L^\dagger T \tilde{q}_L \right] \left[\tilde{q}_R^\dagger T \tilde{q}_R - \tilde{q}_L^\dagger T \tilde{q}_L \right],\end{aligned}$$

Full QCD sector of the MSSM

Example II: Gluino pair production

Softer



Non flat K-factor

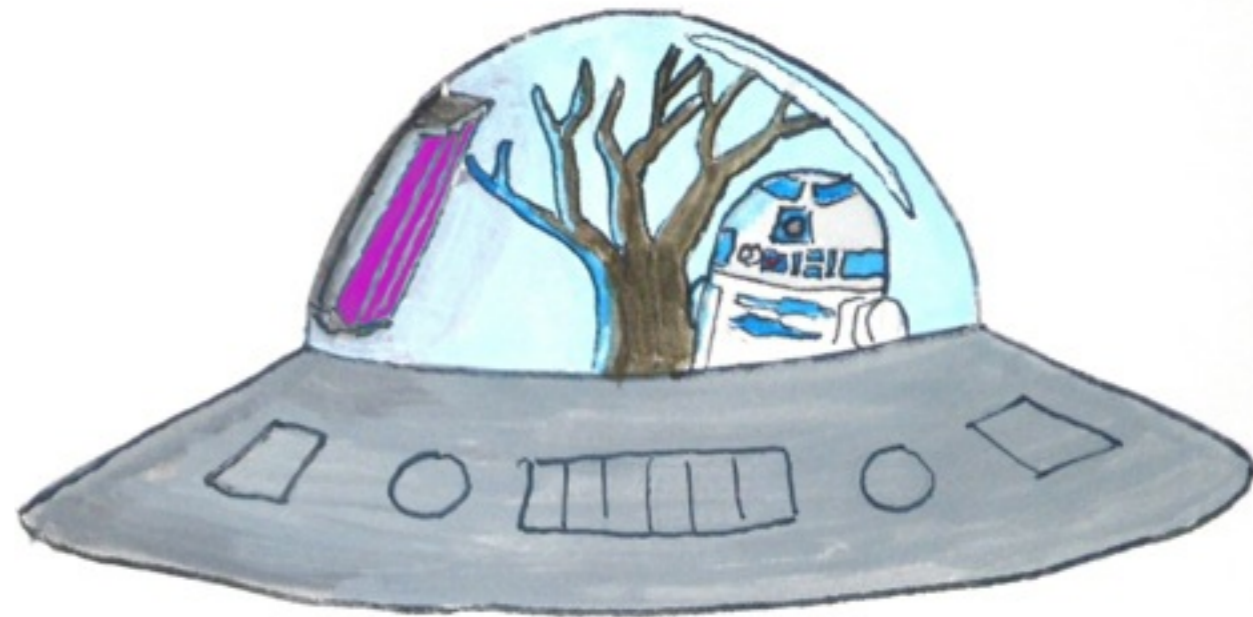
Gluino decay

More examples

- Squark and sgluon production [CD, B. Fuks, V. Hirschi, J. Proudome and H. S. Shao, PRD 91 (2015) no.9, 094005]
- GM [CD, K. Hartling, H. E. Logan, A. D. Peterson and M. Zaro, PRD 93 (2016) no.3, 035004]
- Heavy neutrino [CD, O. Mattelaer, R. Ruiz and J. Turner, arXiv:1602.06957]
- Spin-2 (dim-5 operators) [G. Das, CD, V. Hirschi, F. Maltoni and H. S. Shao, arXiv:1605.09359]
- Top dim-6
 - FCNC [CD F. Maltoni, J. Wang and C. Zhang, PRD 91 (2015) 034024]
 - Pair production [D. Buarque Franzosi and C. Zhang, PRD 91 (2015) no. 11, 114010]
 - tth [F. Maltoni, E. Vryonidou and C. Zhang, arXiv:1607.05330]
- ...

Final remarks

- Automatic BSM@NLO
 - Renormalizable (Public)
- Pheno
 - Spin 2
 - Charged Higgs threshold
 - Top EFT
 - ...
- Jointly by FeynRules and Madgraph_aMC@NLO teams



Back-up

Example I: Charged Higgs production

$$m_{H^\pm} = 200 \text{ GeV} \quad \text{and} \quad m_{H^\pm} = 600 \text{ GeV}$$

$$\tan \beta = 8 \quad \text{but} \quad y_b^2, y_t^2 \text{ and } y_b y_t$$

NNPDF2.3 at NLO/ NNPDF3.0 at LO with 4/5F

$$\alpha_s(M_Z) = 0.118 \quad \text{(5F)} \quad \alpha_s(M_Z) = 0.1226 \quad \text{(4F)}$$

$$m_b^{\text{pole}} = 4.75 \text{ GeV} \quad m_t^{\text{pole}} = 172.5 \text{ GeV}$$

$$\bar{m}_b(\bar{m}_b) = 4.3377 \text{ GeV}$$

$$\mu_{R,F} = H_T/3 \equiv \frac{1}{3} \sum_i \sqrt{m(i)^2 + p_T(i)^2}$$

$$\text{Anti-}k_T \quad \Delta R = 0.4 \quad p_T(j) \geq 25 \text{ GeV}, \quad |\eta(j)| \leq 2.5.$$

Renormalization conditions

Zero momentum scheme available for the gauge couplings

$$\Gamma_{FFV}^\mu(p_1, p_2) = igT^a \delta_{f_1, f_2} \left[\gamma^\mu \left(\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{1}{2} \delta Z_{FF}^R + \frac{1}{2} \delta Z_{FF}^L + \frac{g'_V}{2g} \delta Z_{V'V} \right) \right. \\ \left. + \gamma^\mu \gamma_5 \left(\frac{1}{2} \delta Z_{FF}^R - \frac{1}{2} \delta Z_{FF}^L + \frac{g'_A}{2g} \delta Z_{V'V} \right) \right. \\ \left. + \left(\gamma^\mu h^V(k^2) + \gamma^\mu \gamma_5 h^A(k^2) + \frac{(p_1 - p_2)^\mu}{2m} h^S(k^2) + \frac{k_\mu}{2m} h^P(k^2) \right) \right]$$

$$\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{1}{2} \delta Z_{FF}^R + \frac{1}{2} \delta Z_{FF}^L + \frac{g'_V}{2g} \delta Z_{V'V} + h^V(0) + h^S(0) = 0 \\ \frac{1}{2} \delta Z_{FF}^R - \frac{1}{2} \delta Z_{FF}^L + \frac{g'_A}{2g} \delta Z_{V'V} + h^A(0) = 0.$$

By gauge invariance

$$\frac{\delta g}{g} + \frac{1}{2} \delta Z_{VV} + \frac{g'_V}{2g} \delta Z_{V'V} + \frac{g'_A}{2g} \delta Z_{V'V} = 0$$

Only from
two-point
functions

$\overline{\text{MS}}$ scheme for everything else (option for all)

How does it work?

FeynRules :

```
...  
Lren = OnShellRenormalization[ LSM , QCDOOnly -> True];  
WriteFeynArtsOutput[ Lren , Output -> "SMrenoL",  
GenericFile -> False]
```

FeynArts / NLOCT :

```
WriteCTI[ "SMrenoL/SMrenoL" , "Lorentz", Output->  
"SMQCDreno", QCDOOnly -> True]
```

FeynRules :

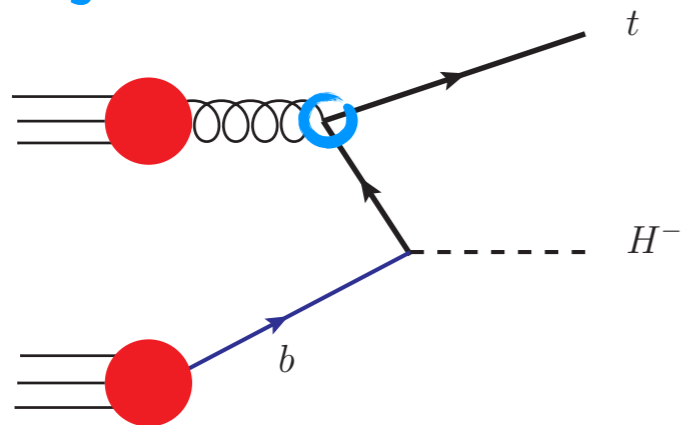
```
...  
Get["SMQCDreno.nlo"];  
WriteUFO[ LSM , UVCounterterms -> UV$vertlist ,  
R2Vertices -> R2$vertlist]
```

Example 1: Charged Higgs production

5 Flavours

- $m_b=0$ (but $m_b^y>0$)
- In the PDF
- In the running of α_s
- Handle collinear logarithms

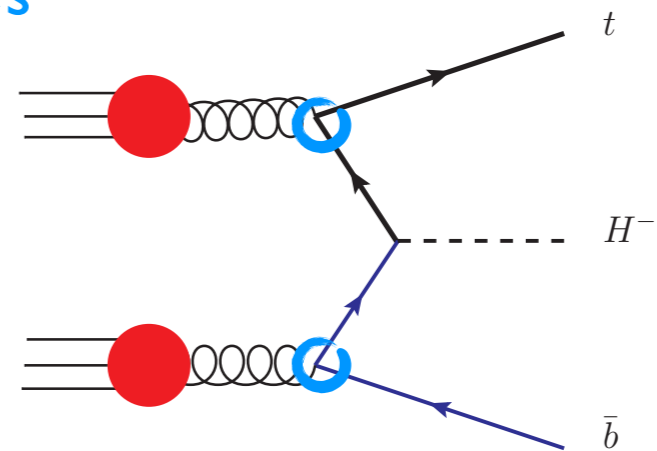
● α_s



4 Flavours

- $m_b>0$
- Not in the PDF
- Not in the running of α_s
- Contribution to b observable at LO

● α_s^2



Example I: Charged Higgs production

Type-II 2HDM

$$V_{t\bar{b}H^-} = -i \left(y_t P_R \frac{1}{\tan \beta} + y_b P_L \tan \beta \right)$$

$$P_{R/L} = (1 \pm \gamma_5)/2$$
$$y_{t/b} \equiv \sqrt{2} \frac{m_{t/b}^y}{v}$$

$$\delta y_{t/b} = \sqrt{2} \frac{\delta m_{t/b}}{v} \xrightarrow{\text{On-shell sc.}} \delta m_{t/b} = -\frac{g_s^2}{12\pi^2} m_{t/b} \left(\frac{3}{\bar{\epsilon}} + 4 - 6 \log \frac{m_{t/b}}{\mu_R} \right)$$
$$\xrightarrow{\overline{\text{MS}} \text{ sc.}} \delta y_b = -\frac{\sqrt{2}}{v} \frac{g_s^2 m_b^y}{4\pi^2 \bar{\epsilon}}$$

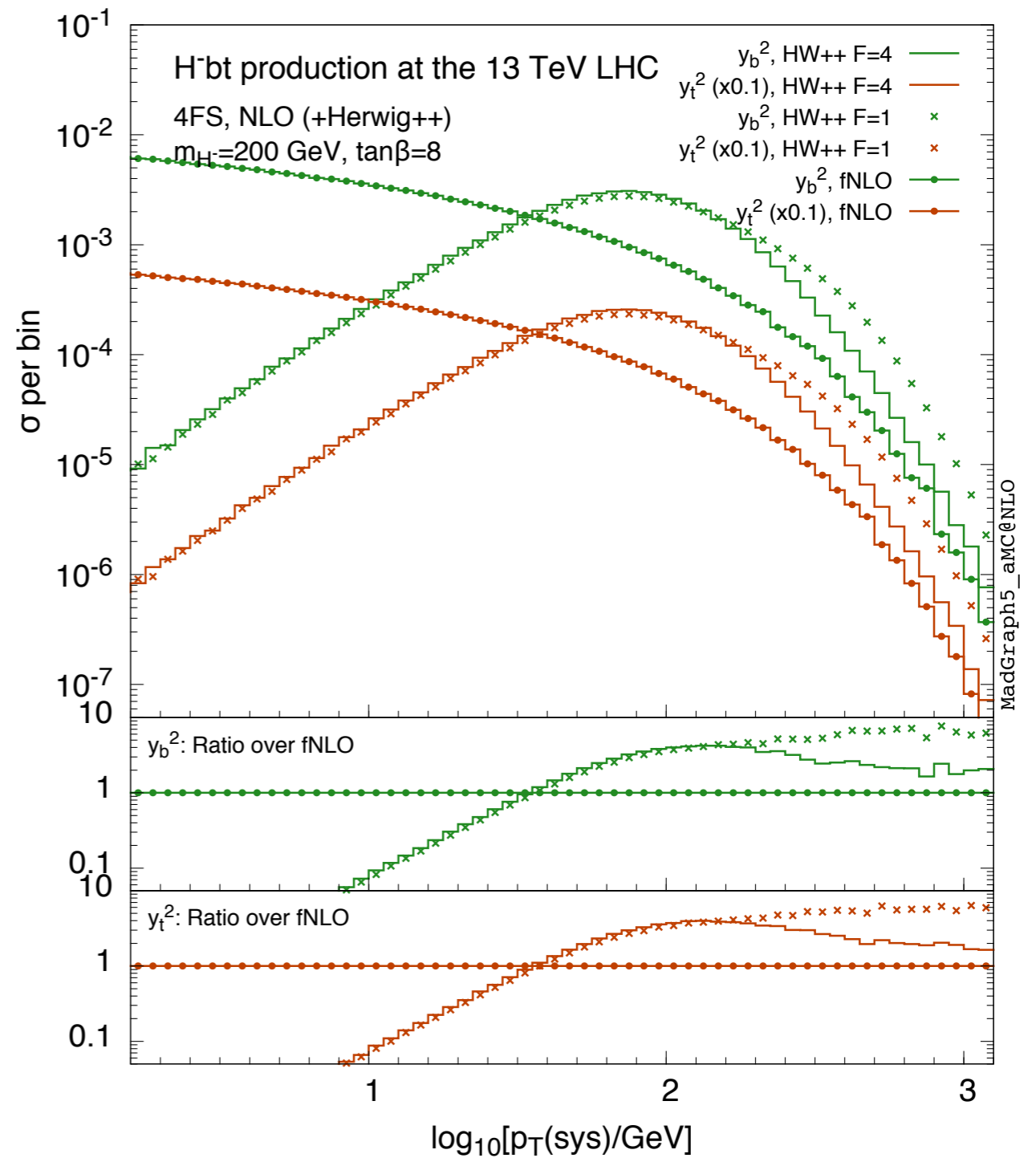
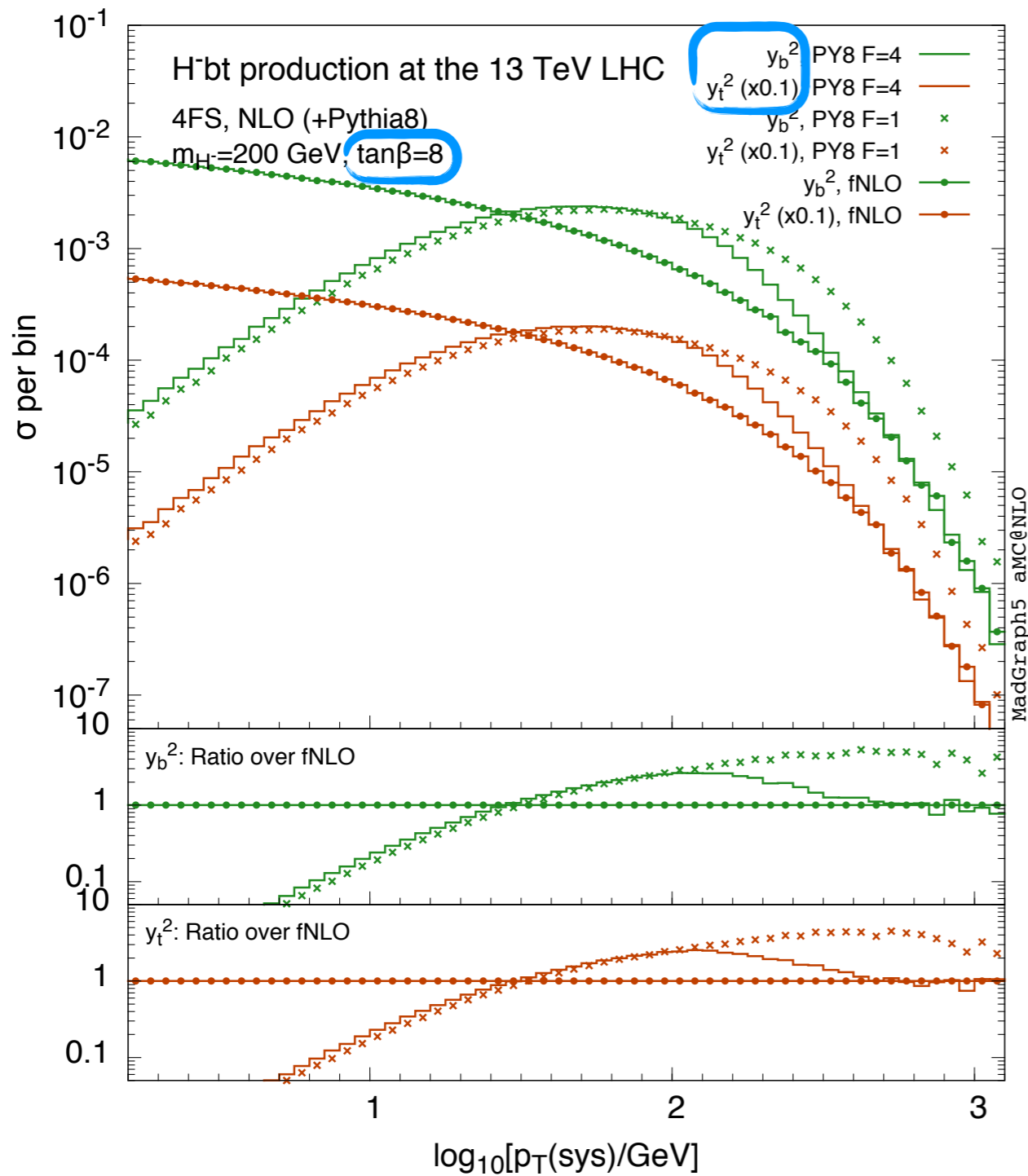
Input : -FR model

-running of the b yukawa mass

Validation : -Comparison with S. Dittmaier, M. Kramer, M. Spira
and M. Walser, PRD 83 (2011) 055005

-Recover ttH

Example I: Charged Higgs production



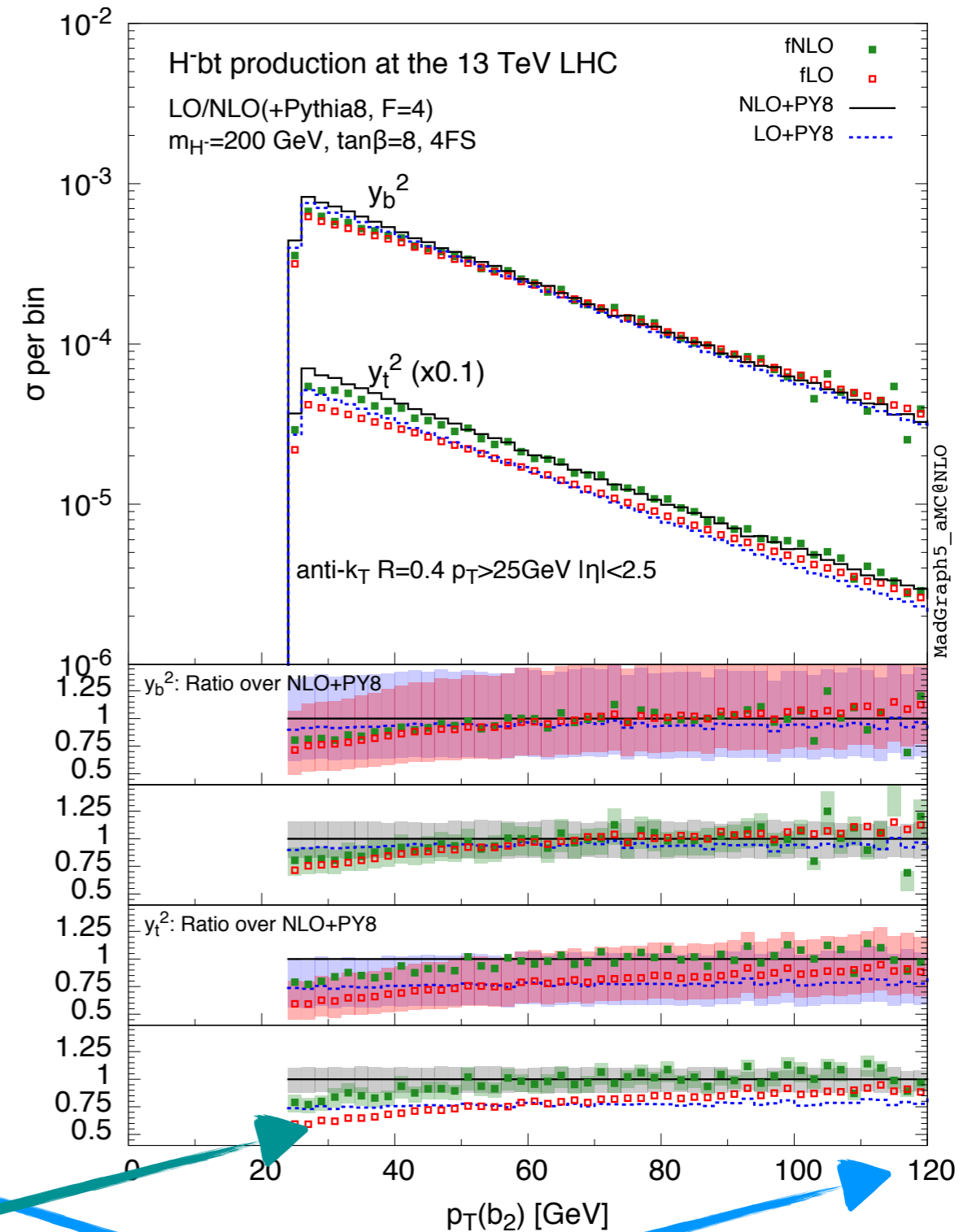
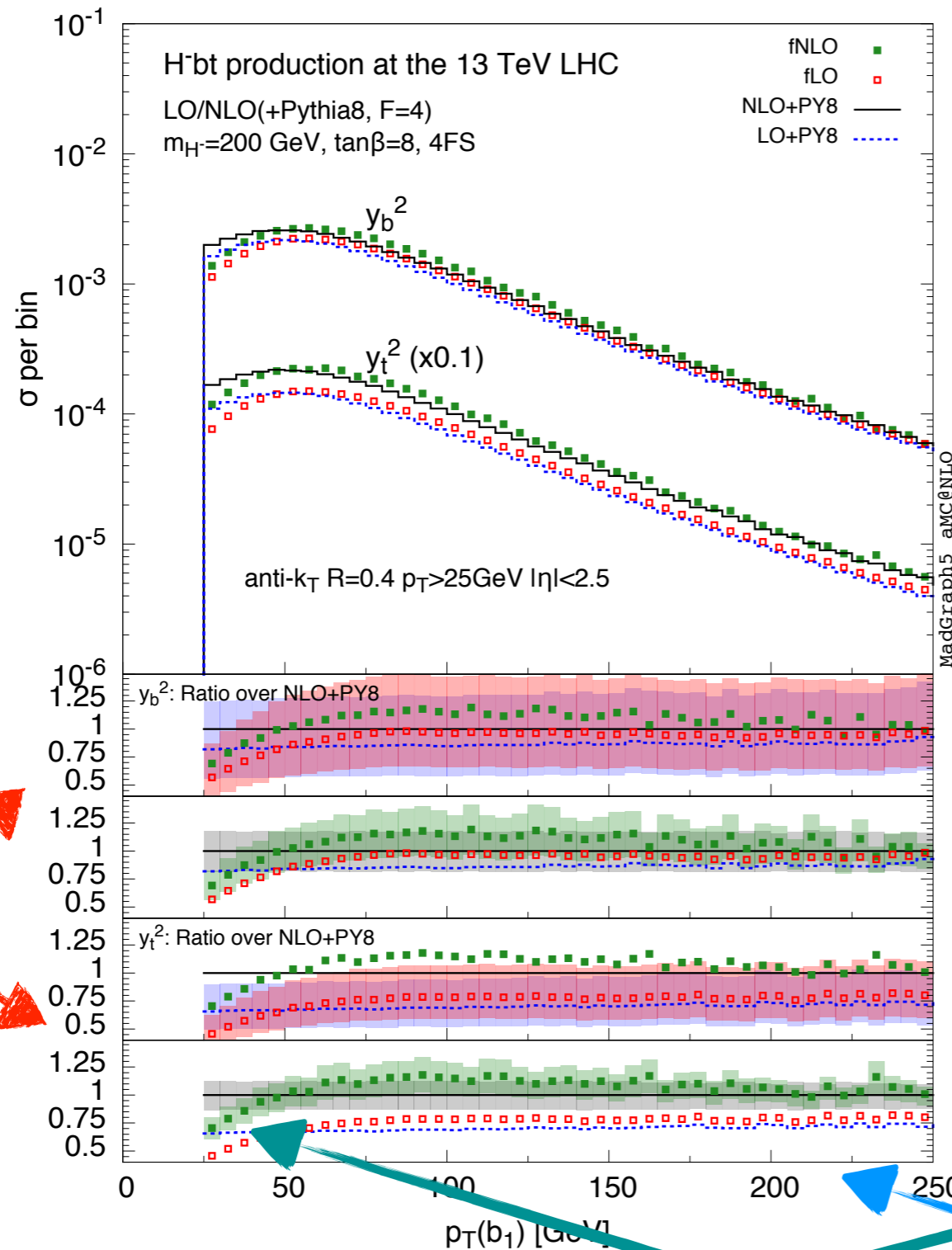
$$\mu_{R,F} = H_T/3 \equiv \frac{1}{3} \sum_i \sqrt{m(i)^2 + p_T(i)^2}$$

$$\frac{0.1}{F} \hat{s} \leq \mu_{\text{sh}} \leq \frac{1}{F} \hat{s}$$

Example 1: Charged Higgs production

$\sigma(m_{H^-} = 200 \text{ GeV})$ [fb]		NLO		LO		Reduced Top leptonic decay
		y_b^2	y_t^2	y_b^2	y_t^2	
$\geq 1j_b$	Inclusive	$50.40^{+17.8\%}_{-18.6\%}$	$42.43^{+12.4\%}_{-13.1\%}$	$42.12^{+52.2\%}_{-31.9\%}$	$28.68^{+36.3\%}_{-24.7\%}$	
	F.O.	$45.47^{+17.5\%}_{-18.4\%}$	$38.31^{+12.2\%}_{-13.0\%}$	$38.26^{+51.9\%}_{-31.8\%}$	$26.09^{+36.1\%}_{-24.6\%}$	
	Pythia8	$43.44^{+17.4\%}_{-18.4\%}$	$36.67^{+12.0\%}_{-13.0\%}$	$36.81^{+52.0\%}_{-31.8\%}$	$25.09^{+36.1\%}_{-24.7\%}$	
	Herwig++	42.64	36.04	36.08	24.61	
$\geq 2j_b$	F.O.	$11.55^{+10.9\%}_{-15.4\%}$	$9.76^{+6.5\%}_{-10.0\%}$	$11.22^{+50.4\%}_{-31.2\%}$	$7.79^{+35.0\%}_{-24.1\%}$	
	Pythia8	$12.55^{+15.3\%}_{-17.4\%}$	$10.67^{+10.4\%}_{-12.1\%}$	$11.73^{+51.2\%}_{-31.5\%}$	$8.12^{+35.6\%}_{-24.4\%}$	
	Herwig++	11.03	9.33	10.09	7.00	
				$\tan \beta = 8$	\neq K-factor	
$\sigma(m_{H^-} = 600 \text{ GeV})$ [fb]		NLO		LO		
		y_b^2	y_t^2	y_b^2	y_t^2	
$\geq 1j_b$	Inclusive	$2.400^{+20.3\%}_{-20.1\%}$	$2.117^{+13.1\%}_{-14.2\%}$	$1.794^{+54.9\%}_{-33.0\%}$	$1.339^{+40.1\%}_{-26.5\%}$	
	F.O.	$2.187^{+19.9\%}_{-19.9\%}$	$1.925^{+12.6\%}_{-14.0\%}$	$1.649^{+54.7\%}_{-32.9\%}$	$1.232^{+39.9\%}_{-26.5\%}$	
	PYTHIA8	$2.115^{+19.9\%}_{-19.9\%}$	$1.865^{+12.5\%}_{-14.0\%}$	$1.601^{+54.8\%}_{-32.9\%}$	$1.197^{+40.0\%}_{-26.5\%}$	
	HERWIG++	2.077	1.836	1.570	1.175	
$\geq 2j_b$	F.O.	$0.630^{+12.6\%}_{-17.0\%}$	$0.548^{+5.9\%}_{-10.8\%}$	$0.548^{+53.8\%}_{-32.6\%}$	$0.413^{+39.2\%}_{-26.2\%}$	
	PYTHIA8	$0.697^{+16.7\%}_{-18.6\%}$	$0.611^{+9.6\%}_{-12.6\%}$	$0.588^{+54.3\%}_{-32.8\%}$	$0.443^{+39.6\%}_{-26.3\%}$	
	HERWIG++	0.602	0.532	0.498	0.376	

Example I: Charged Higgs production



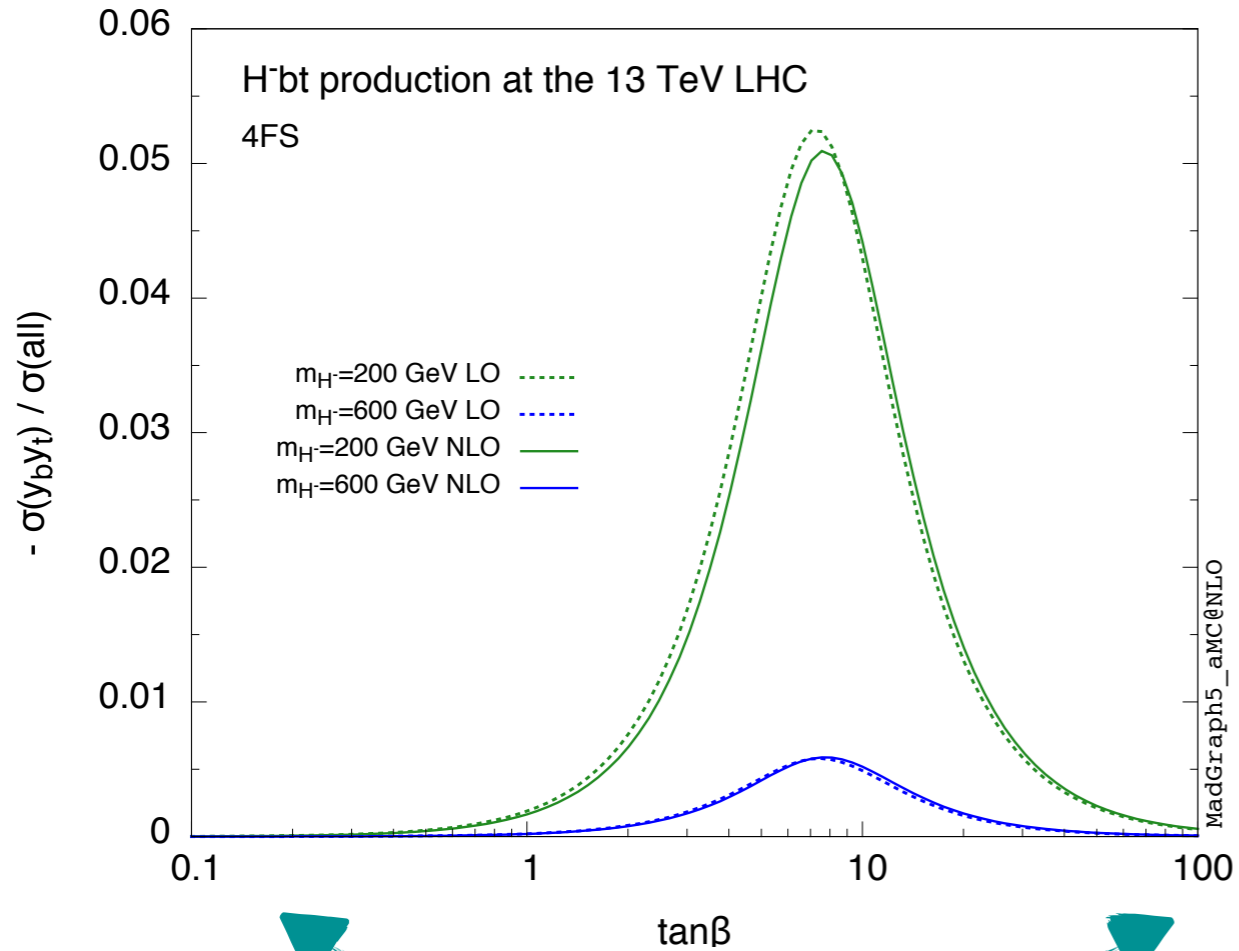
PS effects

FO~PS

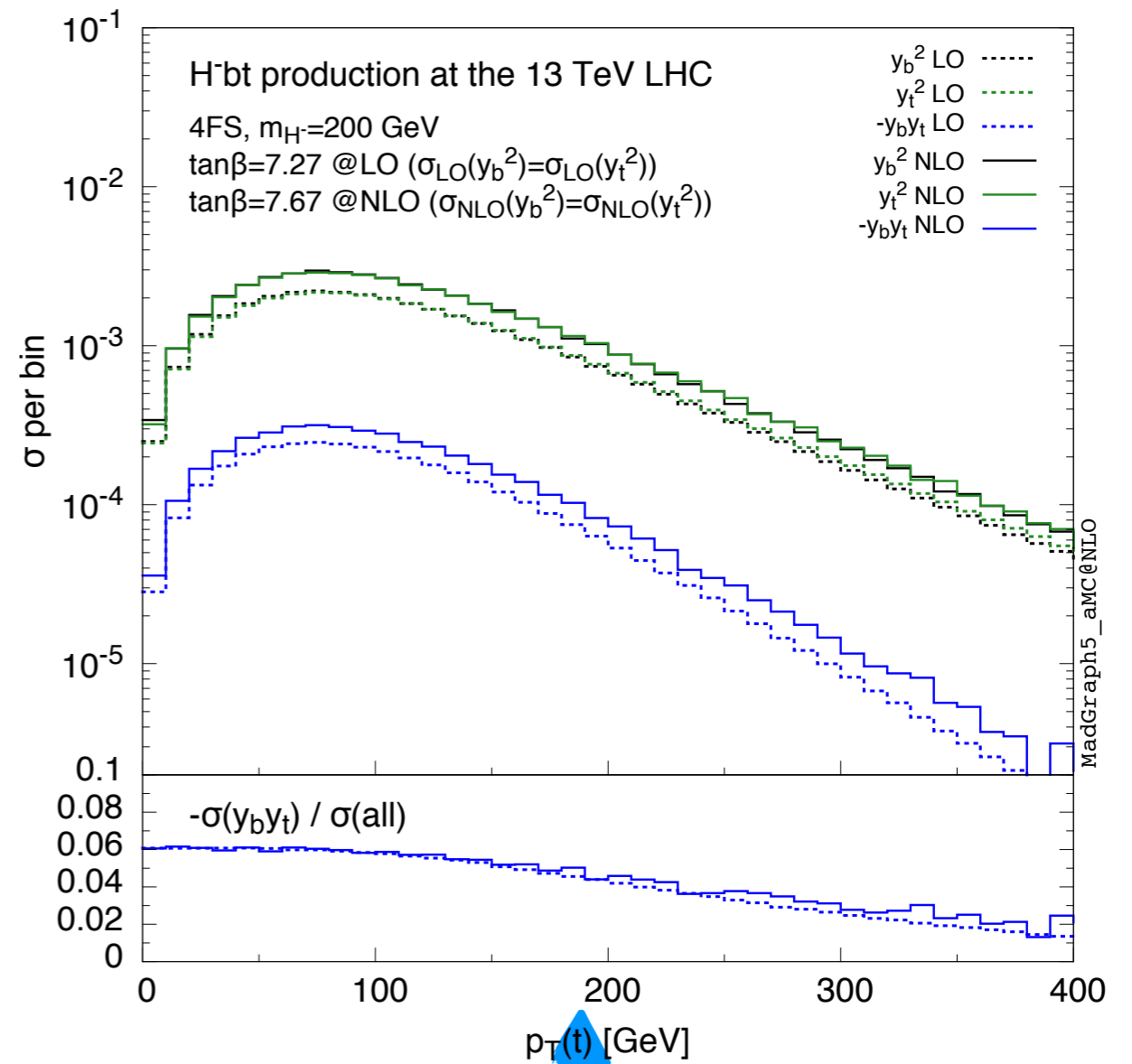
C. Degrande

Example I: Charged Higgs production

Interference : $\propto m_b$

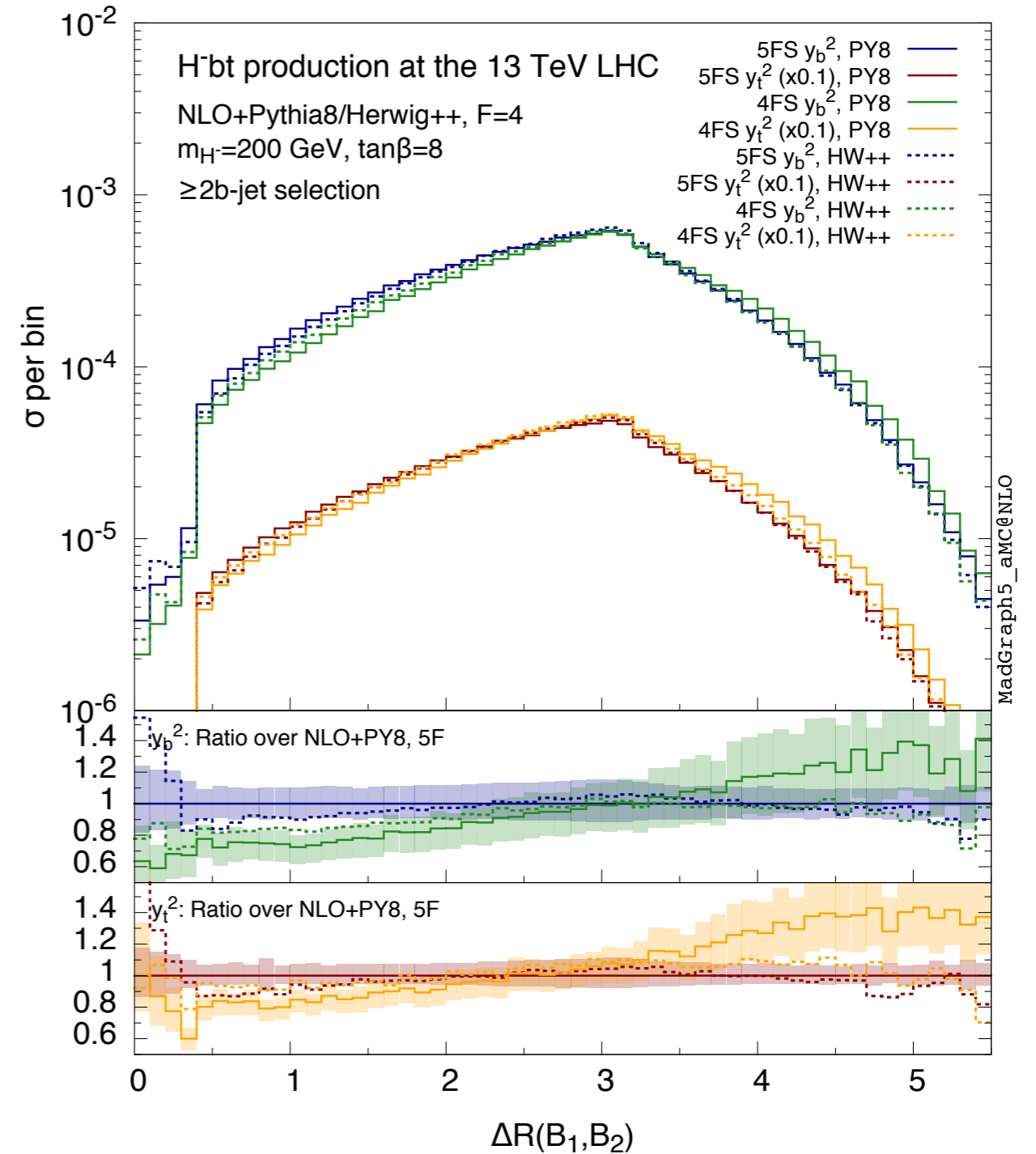
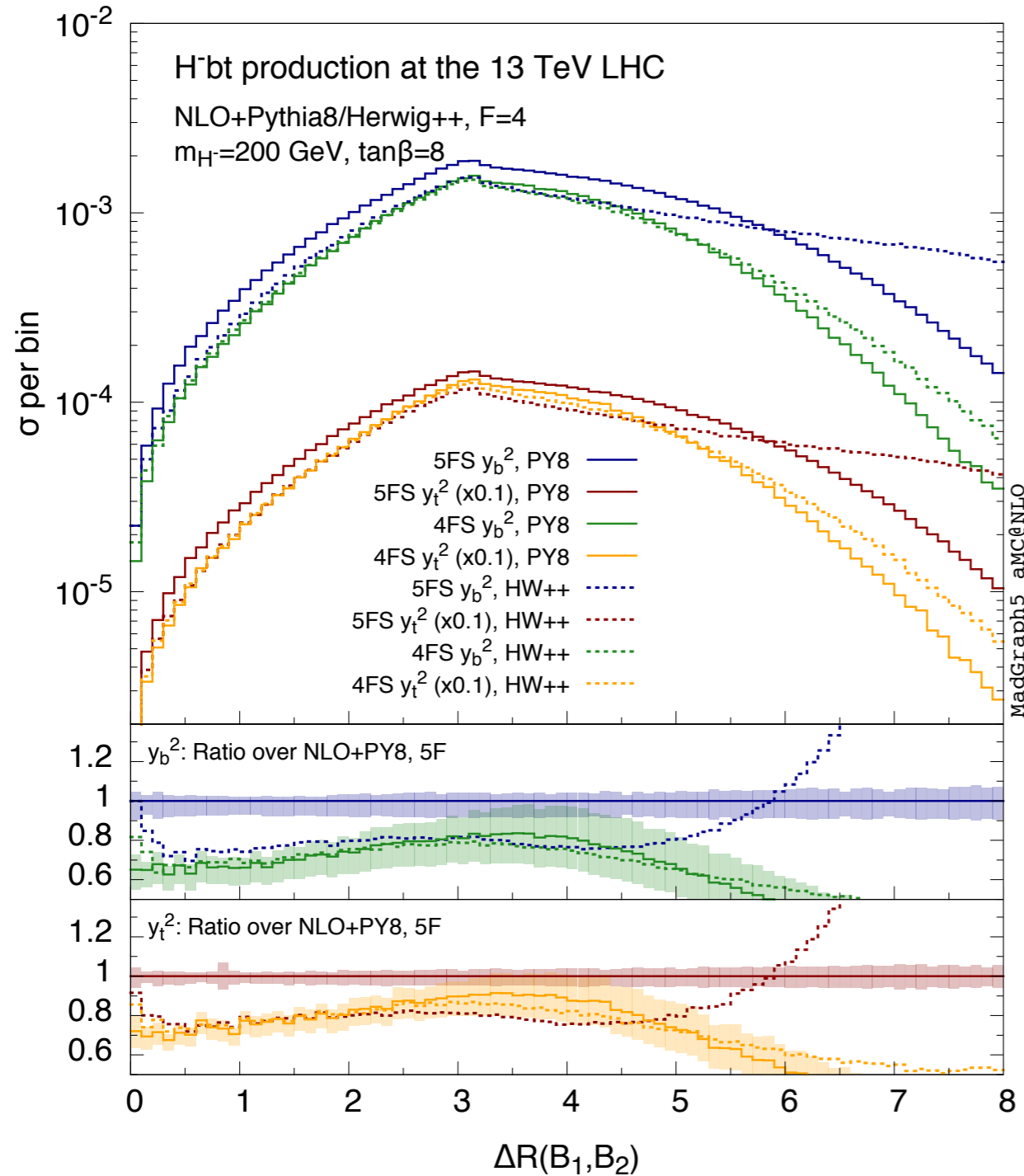


Other terms
are enhanced



More important
at low energy

Example I: Charged Higgs production




Shower dependent in 5F

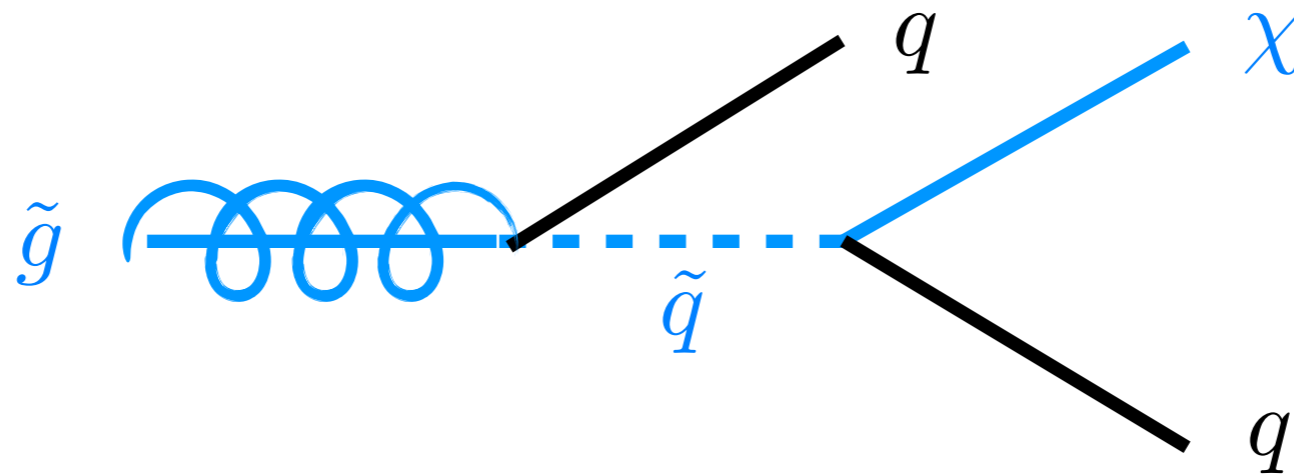
Example II: Gluino pair production

I Majorana gauge singlet (bino-like)

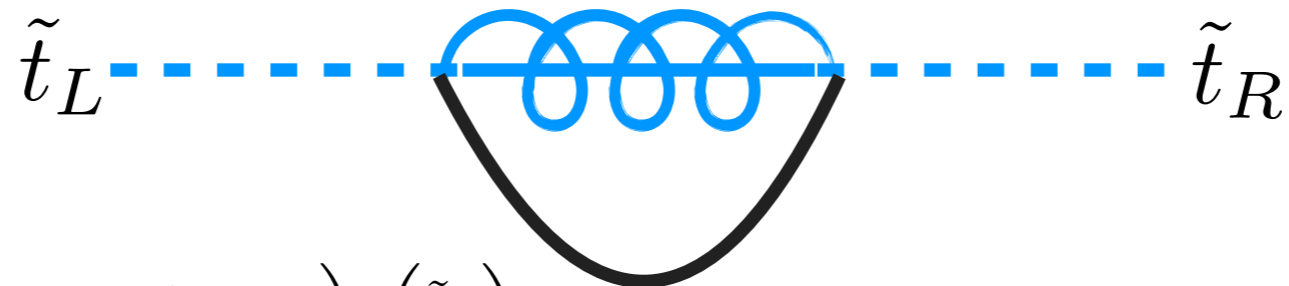
$$\mathcal{L}_{\text{decay}} = \frac{i}{2} \bar{\chi} \not{\partial} \chi - \frac{1}{2} m_{\chi} \bar{\chi} \chi$$



$$+ \sqrt{2} g' \left[- \tilde{q}_L^\dagger Y_q (\bar{\chi} P_L q) + (\bar{q} P_L \chi) Y_q \tilde{q}_R + \text{h.c.} \right]$$



Example II: Gluino pair production



$$\begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \end{pmatrix} \rightarrow \begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \end{pmatrix} + \frac{1}{2} \begin{pmatrix} \delta Z_{\tilde{t}_L} & \delta Z_{\tilde{t},LR} \\ \delta Z_{\tilde{t},RL} & \delta Z_{\tilde{t}_R} \end{pmatrix} \begin{pmatrix} \tilde{t}_L \\ \tilde{t}_R \end{pmatrix} \quad \delta \mathcal{L}_{\text{off}} = -\delta m_{\tilde{t},LR}^2 (\tilde{t}_L^\dagger \tilde{t}_R + \tilde{t}_R^\dagger \tilde{t}_L)$$

Zero-momentum for the massive d.o.f.

$$\begin{aligned} \frac{\delta \alpha_s}{\alpha_s} &= \frac{\alpha_s}{2\pi\bar{\epsilon}} \left[\frac{n_f}{3} - \frac{11n_c}{6} \right] + \frac{\alpha_s}{6\pi} \left[\frac{1}{\bar{\epsilon}} - \log \frac{m_t^2}{\mu_R^2} \right] \\ &+ \frac{\alpha_s n_c}{6\pi} \left[\frac{1}{\bar{\epsilon}} - \log \frac{m_{\tilde{g}}^2}{\mu_R^2} \right] + \frac{\alpha_s}{24\pi} \sum_{\tilde{q}} \left[\frac{1}{\bar{\epsilon}} - \log \frac{m_{\tilde{q}}^2}{\mu_R^2} \right] \end{aligned}$$

Finite CT to remove SUSY artificial breaking

Example II: Gluino pair production

Validation : Comparison with Prospino (degenerate spectrum)

$$m_{\tilde{t}_R} = 17\text{TeV}, \quad m_{\tilde{t}_L} = 16\text{TeV}, \quad m_{\tilde{q}_{R/L}} = 15\text{TeV}$$

No resonant squark

$m_{\tilde{g}}$ [GeV]	σ^{LO} [pb]	σ^{NLO} [pb]
200	$2104^{+30.3\% +14.0\%}_{-21.9\% -14.0\%}$	$3183^{+10.8\% +1.8\%}_{-11.6\% -1.8\%}$
500	$15.46^{+34.7\% +19.5\%}_{-24.1\% -19.5\%}$	$24.90^{+12.5\% +3.7\%}_{-13.4\% -3.7\%}$
750	$1.206^{+35.9\% +23.5\%}_{-24.6\% -23.5\%}$	$2.009^{+13.5\% +5.5\%}_{-14.1\% -5.5\%}$
1000	$1.608 \cdot 10^{-1}^{+36.3\% +26.4\%}_{-24.8\% -26.4\%}$	$2.743 \cdot 10^{-1}^{+14.4\% +7.3\%}_{-14.8\% -7.3\%}$
1500	$6.264 \cdot 10^{-3}^{+36.2\% +29.4\%}_{-24.7\% -29.4\%}$	$1.056 \cdot 10^{-2}^{+16.1\% +11.3\%}_{-15.8\% -11.3\%}$
2000	$4.217 \cdot 10^{-4}^{+35.6\% +29.8\%}_{-24.5\% -29.8\%}$	$6.327 \cdot 10^{-4}^{+17.7\% +17.8\%}_{-16.6\% -17.8\%}$