# Measurement of Observed Cross Sections for $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow$ hadrons $\left.\right|_{\text {non-D }}$ 

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## Introduction

- The BES-II previously observed an anomalous line shape of the cross section for $e^{+} e^{-} \rightarrow$ hadrons (Phys. Rev. Lett. 101, 102004 (2008)).

- This anomalous line shape more likely indicates that there are new structure(s) in addition to the $\psi(3770)$ around 3.773 GeV or some dynamics affect $\psi(3770)$ and $D \bar{D}$ production and decays.
- S. Dubynskiy and M. B. Voloshin interpret this anomalous line shape of cross section as a "diresonance", "which may arise from existence of both a charmonium state and a 'molecular' DD threshold resonance" (Phys. Rev. D 78, 116014 (2008)).
- However, this observation of the line shape need to be confirmed at other experiment.
- Experimental study of the process of $e^{+} e^{-}$annihilating into hadrons in this energy range with a larger data set at the BESIII experiment will provide some important information to elucidate this anomalous line-shape.


## Data and MC Samples

- Data sets:
$\sim 70 \mathrm{pb}^{-1}$ energy scan data taken in range from 3.64 to 3.87 GeV
- MC samples:

| Process | Generator |
| :--- | :--- |
| $J / \psi \rightarrow$ hadrons | KKMC+BESEVTGEN |
| $\psi(3686) \rightarrow$ hadrons | KKMC+BESEVTGEN |
| $D^{0} \bar{D}^{0}, D^{+} D^{-}$ | KKMC+BESEVTGEN |
| $\psi(3770) \rightarrow$ non- $D \bar{D}$ | KKMC+BESEVTGEN |
| $e^{+} e^{-} \rightarrow q \bar{q}$ | KKMC |
| $e^{+} e^{-} \rightarrow e^{+} e^{-}, \mu^{+} \mu^{-}, \gamma \gamma$ | BABAYAGA |
| $e^{+} e^{-} \rightarrow \tau^{+} \tau^{-}$ | KKMC |
| $e^{+} e^{-} \rightarrow e^{+} e^{-} f \bar{f}$ | BESTwOGAM |

For inclusive decays of the $J / \psi, \psi(3686)$ and $\psi(3770)$, the known decay modes are generated by EVTGEN with branching fractions taken from PDG, while the remaining unknown decay modes are modeled by LUNDCHARM

- Software environment: BOSS 6.6.4.p01


## Event Selection

- At least three good charged tracks
- $\left|V_{r}\right|<1.0 \mathrm{~cm},\left|V_{z}\right|<15.0 \mathrm{~cm},|\cos \theta|<0.93$
- To suppress the Bhabha and dimuon backgrounds, require the ratio of Fox-Wolfram moments $H_{2} / H_{0}<0.85$


- To suppress the two-photon and beam-gas events, require the ratio $\left|p_{z}^{\text {miss }}\right| / E_{\text {vis }}<0.30$
- $p_{z}^{\text {miss }}$ : the $z$ component of the missing momentum of the event
- $E_{\text {vis }}$ : the visible energy of the event


## Event Selection

- In order to separate some beam-gas associated background events and cosmic-ray events, which are suspected to be produced at random $Z$ positions in the beam pipe, we can examine the averaged $Z$ position $\left(Z_{\text {aver }}\right)$ of the charged tracks.
- The numbers of observed hadronic events are extracted from maximum likelihood fits to the $Z_{\text {aver }}$ distributions.


| $E_{\mathrm{cm}}(\mathrm{GeV})$ | $N_{\text {obs }}$ | $E_{\mathrm{cm}}(\mathrm{GeV})$ | $N_{\text {obs }}$ | $E_{\mathrm{cm}}(\mathrm{GeV})$ | $N_{\text {obs }}$ | $E_{\mathrm{cm}}(\mathrm{GeV})$ | $N_{\text {obs }}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 3.6451 | $6905 \pm 102$ | 3.7587 | $74653 \pm 300$ | 3.7873 | $43808 \pm 236$ | 3.8319 | $3791 \pm 76$ |
| 3.7215 | $2208 \pm 52$ | 3.7617 | $76820 \pm 312$ | 3.7915 | $51908 \pm 254$ | 3.8400 | $3560 \pm 68$ |
| 3.7296 | $3533 \pm 69$ | 3.7645 | $58202 \pm 271$ | 3.7952 | $57198 \pm 270$ | 3.8479 | $3555 \pm 68$ |
| 3.7454 | $14932 \pm 138$ | 3.7674 | $44937 \pm 237$ | 3.7989 | $54047 \pm 263$ | 3.8561 | $4102 \pm 76$ |
| 3.7470 | $22267 \pm 167$ | 3.7702 | $38140 \pm 213$ | 3.8030 | $36777 \pm 222$ | 3.8640 | $3895 \pm 75$ |
| 3.7493 | $36395 \pm 216$ | 3.7731 | $35114 \pm 207$ | 3.8068 | $23726 \pm 172$ | 3.6474 | $26947 \pm 200$ |
| 3.7508 | $47770 \pm 244$ | 3.7760 | $35077 \pm 205$ | 3.8099 | $17046 \pm 152$ | 3.6534 | $26396 \pm 191$ |
| 3.7530 | $53721 \pm 258$ | 3.7789 | $36476 \pm 211$ | 3.8128 | $12063 \pm 128$ |  |  |
| 3.7544 | $55601 \pm 261$ | 3.7818 | $38163 \pm 214$ | 3.8160 | $9019 \pm 112$ |  |  |
| 3.7558 | $63800 \pm 278$ | 3.7847 | $42113 \pm 231$ | 3.8240 | $5247 \pm 86$ |  |  |

## Background Subtraction

- Background sources
- QED processes:

$$
e^{+} e^{-} \rightarrow e^{+} e^{-}, \mu^{+} \mu^{-}, \tau^{+} \tau^{-}, \gamma \gamma
$$

- Two-photon interactions:

$$
\begin{aligned}
& e^{+} e^{-} \rightarrow e^{+} e^{-} e^{+} e^{-}, e^{+} e^{-} \mu^{+} \mu^{-}, e^{+} e^{-} \tau^{+} \tau^{-} \\
& e^{+} e^{-} \rightarrow e^{+} e^{-} u \bar{u}, e^{+} e^{-} d \bar{d}, e^{+} e^{-} s \bar{s}
\end{aligned}
$$

- The number of background events

$$
N_{\mathrm{bkg}}=\mathcal{L} \times\left(\sum_{i} \sigma_{\mathrm{bkg}, i} \times \eta_{i}\right)
$$

- $\mathcal{L}$ : luminosity
- $\sigma_{\text {bkg }}$ : theoretical cross section of background process
- $\eta$ : contamination rate


## Background Subtraction — QED Processes

- The cross sections of $e^{+} e^{-} \rightarrow e^{+} e^{-}$, $\mu^{+} \mu^{-}, \gamma \gamma$ are taken from the outputs of MC generators.
- To calculate the $e^{+} e^{-} \rightarrow \tau^{+} \tau^{-}$ cross section, the effects of ISR, FSR, Coulomb interaction, $\psi(3686)$ production and interference are considered (Phys. Rev. D 74, 112003 (2006)).


- The events selection criteria applied to data is also applied to the MC simulated background events.
- The result of fitting a polynomial to these contamination rates is used for further calculations.


## Background Subtraction - Two-Photon Processes



- The cross sections of $e^{+} e^{-} \rightarrow e^{+} e^{-} f \bar{f}$ ( $f=e, \mu, \tau, u, d, s$ ) are obtained from MC simulation.

- The events selection criteria applied to data is also applied to the MC simulated background events.
- The result of fitting a polynomial to these contamination rates is used for further calculations.


## The D $\bar{D}$ Contributions

- The number of $e^{+} e^{-} \rightarrow D \bar{D}$ events is given by $N_{D \bar{D}}=\mathcal{L} \sigma_{D \bar{D}} \varepsilon_{D \bar{D}}$.
- Here $\sigma_{D \bar{D}}$ is measured using singly tagged $D$ events.
- $D^{0} \rightarrow K^{-} \pi^{+}$
- $D^{+} \rightarrow K^{-} \pi^{+} \pi^{+}$

To reduce the statistic, the values expected from the fit are used.



- The efficiencies $\varepsilon_{D \bar{D}}$ are determined from MC simulation. The results of fitting a polynomial to these efficiencies are used for further calculations.


## Efficiencies

- The efficiency for hadronic events is obtained by analyzing the MC simulated signal events.
- We first determine the detection efficiency for each category of hadronic events. To reduce MC statistic, the results of fitting a polynomial to these efficiencies are used for further calculations.


- We then weight these efficiencies according to their cross sections to obtain the overall efficiency.


## Efficiencies



The detection efficiencies for $e^{+} e^{-} \rightarrow$ hadrons $\left.\right|_{\text {non-D }}$ as a function of the CM energy

## Cross Sections

- The observed cross section for $e^{+} e^{-} \rightarrow$ hadrons $\left.\right|_{\text {non-D }}$ is given by

$$
\sigma_{\mathrm{had}-\mathrm{n} D \bar{D}}^{\mathrm{obs}}\left(E_{\mathrm{cm}}\right)=\frac{N_{\mathrm{obs}}\left(E_{\mathrm{cm}}\right)-N_{\mathrm{bkg}}\left(E_{\mathrm{cm}}\right)}{\mathcal{L}\left(E_{\mathrm{cm}}\right) \times \varepsilon\left(E_{\mathrm{cm}}\right)} .
$$



- Analysis of this observed cross section needs the expectedobserved cross sections for this final state, which can be obtained with BW function and ISR sampling function.


## Systematic Uncertainty

| Source | Systematic uncertainty (\%) |
| :--- | :---: |
| Charged track multiplicity | 1.4 |
| $R_{2}$ cut | 0.5 |
| $\left\|p_{Z}^{\text {miss }}\right\| / E_{\text {vis }}$ cut | 0.9 |
| Fit to the $\bar{V}_{z}$ distribution | 0.8 |
| Luminosity | 1.0 |
| Sum in quadrature | 2.2 |

## Expected Cross Section

- The cross section of $e^{+} e^{-} \rightarrow$ hadrons $_{\left.\right|_{D D \bar{D}}}$ can be expressed as

$$
\sigma^{\exp }(s)=\sigma_{J / \psi}^{\exp }(s)+\sigma_{\psi(3686)}^{\exp }(s)+\sigma_{R_{s}(3770)}^{\exp }(s)+\sigma_{q \bar{q}}(s)
$$

where $s \equiv E_{\mathrm{cm}}^{2}$.

- For resonances, one has

$$
\sigma_{R}^{\exp }(s)=\iint \sigma^{0}\left(s^{\prime}(1-x)\right) \mathcal{F}\left(x, s^{\prime}\right) \mathcal{G}\left(s^{\prime}, s\right) d x d s^{\prime}
$$

where $s^{\prime} \equiv s(1-x), x$ is the fraction of the radiative energy to the beam energy, $\mathcal{F}\left(x, s^{\prime}\right)$ is the sampling function, $\mathcal{G}\left(s^{\prime}, s\right)$ is a Gaussian function describing the distribution of the $\mathrm{e}^{+} \mathrm{e}^{-}$collision energy.
We use a Breit-Wigner function to calculate the dressed cross section:

$$
\sigma^{0}(s)=\frac{12 \pi \Gamma^{e e} \Gamma \mathcal{B}_{\mathrm{had}}}{\left(s-M^{2}\right)^{2}+(M \Gamma)^{2}},
$$

where $\Gamma^{e e}$ and $\Gamma$ are respectively the leptonic width and total width of the resonance, $M$ is the mass and $\mathcal{B}_{\text {had }}$ denotes the decay branching fraction.

- The cross section for light hadron production is parameterized as

$$
\sigma_{q \bar{q}}(s)=f_{q \bar{q}} \sigma_{\mu^{+} \mu^{-}}^{\mathrm{B}}(s),
$$

where $\sigma_{\mu^{+} \mu^{-}}^{\mathrm{B}}(s)$ is the Born cross section of $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \mu^{+} \mu^{-}$and $f_{q \bar{q}}$ is a free parameter.

## Fitting Procedure

- We use the least squares method to fit the experimental data. The following function is minimized:

$$
\chi^{2}=\sum_{i}\left(\frac{\sigma^{\mathrm{obs}}\left(E_{\mathrm{cm}, i}\right)-\sigma^{\exp }\left(E_{\mathrm{cm}, i}\right)}{\Delta\left(\sigma^{\mathrm{obs}}\left(E_{\mathrm{cm}, i}\right)\right)}\right)^{2}
$$

where $i$ is the energy point number, $\Delta\left(\sigma^{\text {obs }}\left(E_{\mathrm{cm}, j}\right)\right)$ is the sadistical uncertainty of the measured cross section and $\sigma^{\exp }\left(E_{\mathrm{cm}, i}\right)$ is the predicted cross section.

- The resonance parameters of $J / \psi$ and $\psi(3686)$ are fixed at PDG2016 values in the fit.
- The energy spread is fixed at $\sigma_{E}=1.37 \mathrm{MeV}$.
- For the $R_{s}(3770)$ resonance(s), we use one or two Breit-Wigner amplitude(s) to fit the observed hadronic cross sections.


## Fit Results — Scenario 1

- Assuming that there is only $\psi(3770)$


| $M(\mathrm{MeV})$ | 3773.13 (fixed) |
| :--- | :---: |
| $\Gamma(\mathrm{MeV})$ | 27.2 (fixed) |
| $\Gamma^{e e} \mathcal{B}(\psi(3770) \rightarrow \mathrm{nD} \bar{D})(\mathrm{eV})$ | $39.34 \pm 3.91$ |
| $f_{\mathrm{qa}}$ | $2.794 \pm 0.009$ |
| $\chi^{2} / N_{\text {dof }}$ | $64.6 / 35$ |

## Fit Results — Scenario 1



The observed cross sections after subtracting contributions from $\mathrm{J} / \psi$, $\psi(3686)$ and continuum light hadron production

## Fit Results — Scenario 2

(2) Assuming that there are two structures around 3.773 GeV


| $M_{1}(\mathrm{MeV})$ | $3748.6_{-2.9}^{+1.6}$ |
| :---: | :---: |
| $\Gamma_{1}(\mathrm{MeV})$ | 11.4 $4_{-5.3}^{+11.5}$ |
| $\Gamma_{1}^{\text {ee }} \mathcal{B}\left(R_{1} \rightarrow \mathrm{nD} \overline{\mathrm{D}}\right)(\mathrm{eV})$ | $10.0_{-4}^{+9}$ |
| $M_{2}(\mathrm{MeV})$ | $3773.2_{-1.4}^{+1.5}$ |
| $\Gamma_{2}(\mathrm{MeV})$ | $27.0_{-5.1}^{+6.3}$ |
| $\Gamma_{2}^{e e} \mathcal{B}\left(R_{2} \rightarrow \mathrm{nD} \overline{\mathrm{D}}\right)(\mathrm{eV})$ | $43.3{ }_{-8.9}^{\text {-10.2 }}$ |
| $f_{q \bar{a}}$ | $2.760_{-0.014}^{+0.013}$ |
| $\chi^{2} / N_{\text {dof }}$ | 24.6/30 |

The signal significance of the diresonance is $5.3 \sigma$ (from analyzing the observed cross sections).

## Fit Results — Scenario 2



The observed cross sections after subtracting contributions from $\mathrm{J} / \psi$, $\psi(3686)$ and continuum light hadron production

## Systematic Uncertainties

- We first vary the measured cross sections, parameters of $J / \psi, \psi(3686)$, and energy spread by $\pm 1 \sigma$, then repeat the fit procedure.
- The shifts of the parameters are taken as systematic uncertainties.

| Source | $M_{1}$ ( MeV ) | $\begin{gathered} \Gamma_{1}^{\text {tot }} \\ (\mathrm{MeV}) \end{gathered}$ | $\begin{gathered} \Gamma_{1}^{e e} \mathcal{B}_{R_{1} \rightarrow n D \bar{D}} \\ (\mathrm{eV}) \end{gathered}$ | $\begin{gathered} M_{2} \\ (\mathrm{MeV}) \end{gathered}$ | $\begin{gathered} \Gamma_{2}^{\text {tot }} \\ (\mathrm{MeV}) \end{gathered}$ | $\begin{gathered} \Gamma_{2}^{e e} \mathcal{B}_{R_{2} \rightarrow n D \bar{D}} \\ (\mathrm{eV}) \end{gathered}$ | ${ }^{\prime} q \bar{q}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| $\sigma_{\text {had }}^{\text {obs }}$ measurement | 0.2 | 0.4 | 1.3 | 0.1 | 1.0 | 3.6 | 0.066 |
| $M_{J / \psi}$ | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.000 |
| $\Gamma_{J / \psi}^{\text {tot }}$ | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.000 |
| $\Gamma_{J / \psi}^{e e}$ | 0.0 | 0.1 | 0.0 | 0.0 | 0.1 | 0.3 | 0.004 |
| $M_{\psi(3686)}$ | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.000 |
| $\Gamma_{\psi(3686)}^{\text {tot }}$ | 0.0 | 0.0 | 0.0 | 0.0 | 0.1 | 0.1 | 0.000 |
| $\Gamma_{\psi(3686)}^{e e}$ | 0.1 | 0.3 | 0.8 | 0.0 | 0.9 | 2.2 | 0.002 |
| Energy spread | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.0 | 0.000 |
| Sum in quadrature | 0.2 | 0.5 | 1.5 | 0.1 | 1.4 | 4.2 | 0.066 |

## Comparison




| Parameter | This work | BES-II |
| :--- | :---: | :---: |
| $M_{1}(\mathrm{MeV})$ | $3748.6^{+1.6} \pm 0.2$ | $3762.6 \pm 11.8 \pm 0.5$ |
| $\Gamma_{1}^{\text {tot }}(\mathrm{MeV})$ | $11.4_{-5.3}^{+1.5} \pm 0.5$ | $49.9 \pm 32.1 \pm 0.1$ |
| $M_{2}(\mathrm{MeV})$ | $3773.2_{-1.4}^{+1.5} \pm 0.1$ | $3781.0 \pm 1.3 \pm 0.5$ |
| $\Gamma_{2}^{\text {tot }}(\mathrm{MeV})$ | $27.0_{-5.1}^{+6.3} \pm 1.4$ | $19.3 \pm 3.1 \pm 0.1$ |

- This analysis confirms (at $5.3 \sigma$ ) the BES-II observation of Di-Structure $R_{s}(3770)$ in the range from 3.71 to 3.87 GeV .


## Conclusion

- The cross sections of $e^{+} e^{-} \rightarrow$ hadrons $\left.\right|_{n D \bar{D}}$ are measured in the energy range from 3.64 to 3.87 GeV at the BESIII experiment.
- We observed an enhancement of non- $D \bar{D}$ hadron production in the range from 3.74 to 3.80 GeV , which confirms the anomalous line shape of cross sections for $e^{+} e^{-} \rightarrow$ hadrons observed at the BES-II experiment.
- To well describe the line-shape of these observed cross sections, one more BW amplitude additional to the $\psi(3770)$ resonance is needed.
- By analyzing these observed cross sections, We obtain the parameters of the Di-structures:

| $M_{1}(\mathrm{MeV})$ | $3748.6^{+1.6} \pm 0.2$ |
| :--- | :---: |
| $\Gamma_{1}(\mathrm{MeV})$ | $11.4_{-5.5}^{+1.5} \pm 0.5$ |
| $\Gamma_{1}^{e \mathcal{B}} \mathcal{B}\left(R_{1} \rightarrow\right.$ non $\left.-D \bar{D}\right)(\mathrm{eV})$ | $10.0_{-4.1}^{+9.2} \pm 1.5$ |
| $M_{2}(\mathrm{MeV})$ | $3773.2_{-1.4}^{+1.5} \pm 0.1$ |
| $\Gamma_{2}(\mathrm{MeV})$ | $27.0_{-5.1}^{+6.3} \pm 1.4$ |
| $\Gamma_{2}^{e \Theta} \mathcal{B}\left(R_{2} \rightarrow\right.$ non $\left.-D \bar{D}\right)(\mathrm{eV})$ | $43.3_{-8.9}^{+0.2} \pm 4.2$ |

## Backup Slides

## Fox-Wolfram Moments

For a collection of $N$ particles with momenta $\boldsymbol{p}_{i}$, the $k$-th order FoxWolfram moments $H_{k}$ is defined as

$$
H_{k}=\sum_{i, j}^{N}\left|\boldsymbol{p}_{i}\right|\left|\boldsymbol{p}_{j}\right| P_{k}\left(\cos \theta_{i j}\right),
$$

where $\theta_{i j}$ is the angle between $\boldsymbol{p}_{i}$ and $\boldsymbol{p}_{j}$, and $P_{k}$ is the $k$-th order Legendre polynomial.

## Systematic Uncertainty

- Multiplicity Requirement
- The systematic uncertainty due to the dependence on the charged track multiplicity is estimated by varying the requirement from $N_{\text {ch }} \geq 3$ to $N_{\text {ch }} \geq 2$.
- Additional cuts on $N_{\text {ch }}=2$ events:
- The charged track must not be identified as an electron or a muon:
$E_{\text {EMC }} /$ p $<0.7$ and depth MUC $<30 \mathrm{~cm}$
- The visible energy of event: $E_{\mathrm{vis}}>0.3 \times E_{\mathrm{cm}}$
- The relative changes of $\sigma_{\text {had-n } D \bar{D}}^{\text {obs }}$ are less than $1.4 \%$, which is taken as the systematic uncertainty.
(2) $\mathrm{H}_{2} / \mathrm{H}_{0}$ cut
- To estimate the systematic uncertainty due to the $\mathrm{H}_{2} / \mathrm{H}_{0}$ cut, we vary this cut from its nominal level of 0.85 down to 0.80 .
- The relative changes of $\sigma_{\text {had-n } D \bar{D}}^{\text {obs }}$ are less than $0.5 \%$, which is taken as the systematic uncertainty.
(3) $\left|\rho_{Z}^{\text {miss }}\right| / E_{\mathrm{vis}}$ cut
- To estimate the systematic uncertainty due to the cut on $\left|p_{z}^{\text {miss }}\right| / E_{\mathrm{vis}}$, we vary this cut from its nominal level of 0.30 up to 0.35 .
- The relative changes of $\sigma_{\text {had-nD } \bar{D}}^{\text {obs }}$ are less than $0.9 \%$, which is taken as the systematic uncertainty.


## Systematic Uncertainty

(c) Fit to $\bar{V}_{z}$ distributions

- To estimate the uncertainties due to the fits to the $\bar{V}_{z}$ distributions, we refit the $\bar{V}_{z}$ distributions by varying bin size, fit range ( $[-8,8] \mathrm{cm},[-12,12]$ cm ), background PDF (first, third, fourth order polynomial), and signal PDF (a core Gaussian plus two exponential tails, double-Gaussian).

| Item | Systematic uncertainty (\%) |
| :--- | :---: |
| Bin size | 0.0 |
| Range | 0.2 |
| Background PDF | 0.7 |
| Signal PDF | 0.4 |
| Sum in quadrature | 0.8 |

(6) Luminosity

- 1.0\% (Chin. Phys. C 37, 123001 (2013))

