

Structure of 5-quark Systems in Chiral Quark Model

Gang Yang & Jia-lun Ping

Nanjing Normal University

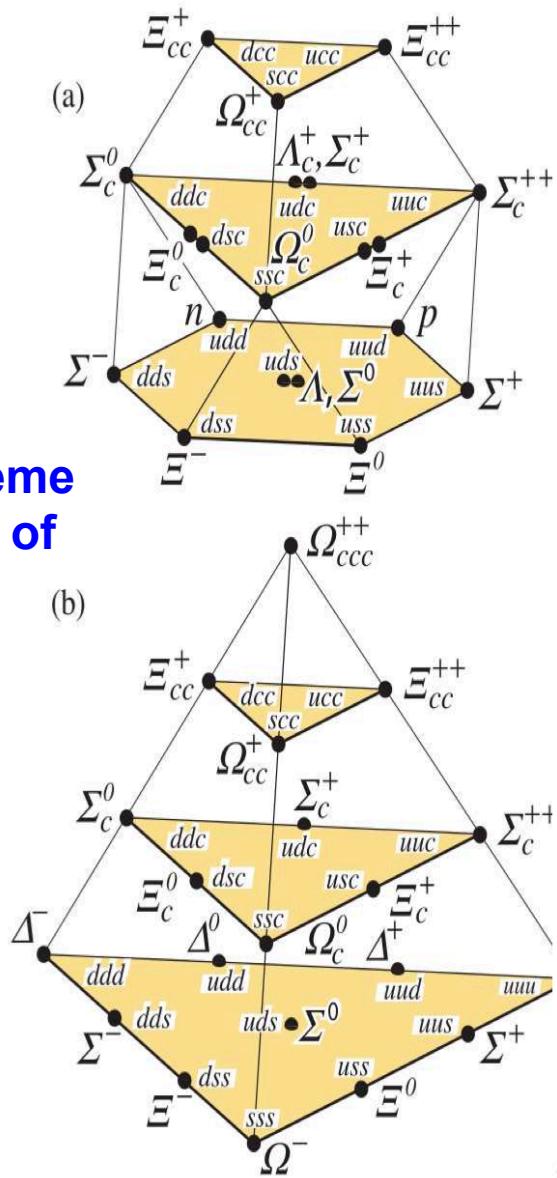
Outline:

- **Introductions**
- **ChQM and calculation methods**
- **Review on P_c^+ and Ω_c^0 from the LHCb collaboration**
- **5-quark study in nucleon sector**
- **Summary**

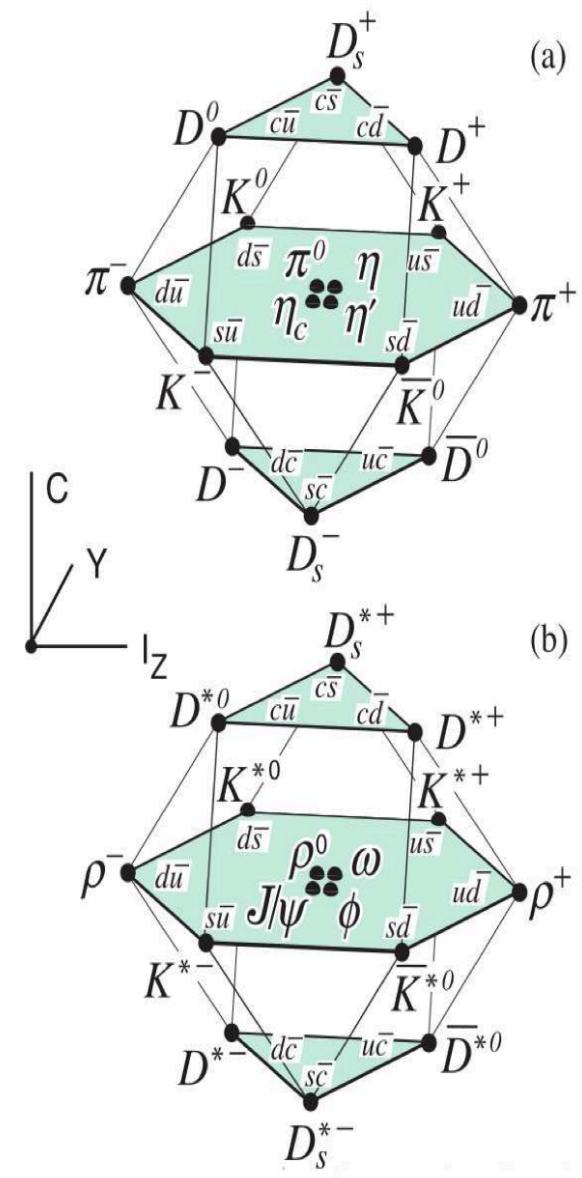
Introduction

**Successful classification scheme
organizing the large number of
conventional hadrons**

Baryons



Mesons



$$3 \otimes 3 \otimes 3 = 10_S \oplus 8_M \oplus 8_M \oplus 1_A$$

$$3 \otimes \bar{3} = 8 \oplus 1$$

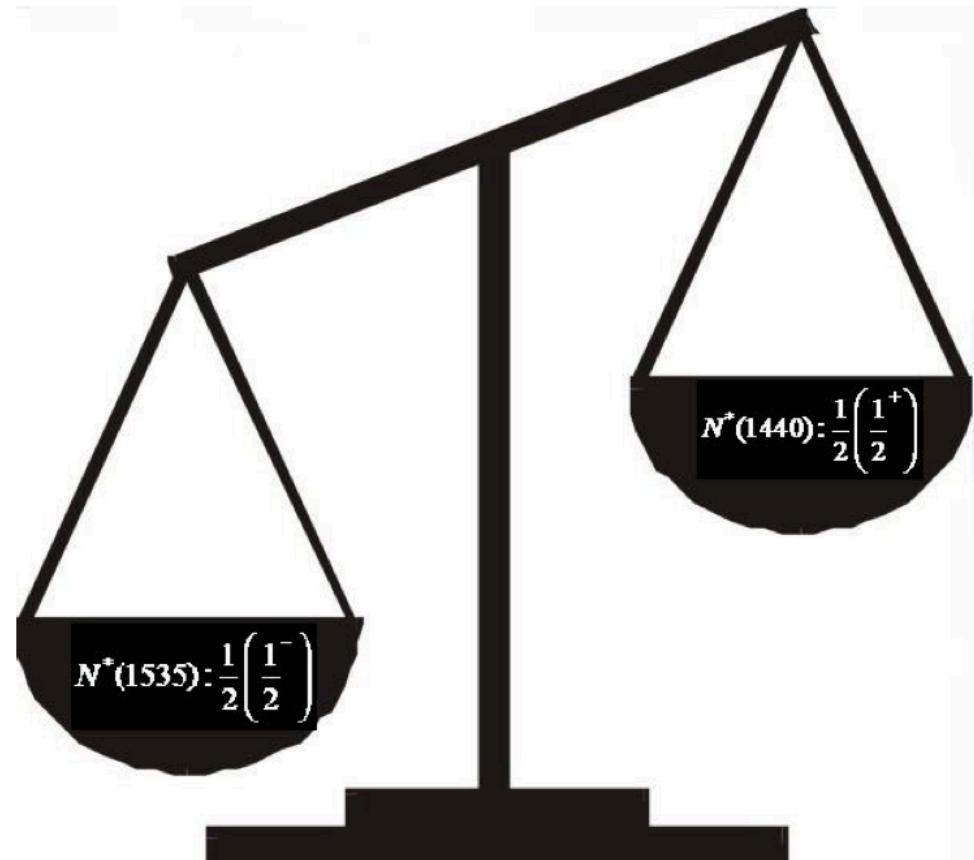
- Mass order reverse

V. D. Burkert, C. D. Roberts,
arXiv: 1710.02549 [nucl-ex]

$$E_{SHO} = \left(2n + L + \frac{3}{2} \right) \hbar\omega$$

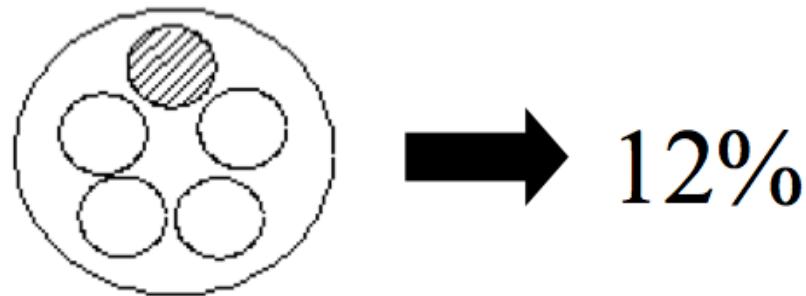
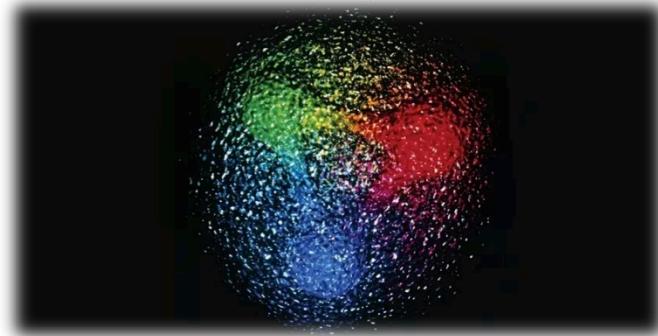
$$E_{n=0, L=1} < E_{n=1, L=0}$$

Roper resonance
puzzle



- \bar{u} and \bar{d} asymmetry

$$\bar{d} - \bar{u} \approx 0.12$$



Garvey G T, Peng J C., Prog. Part. Nucl. Phys., 2001, 47: 203.

- More exotic states

$\Lambda(1405)$ uds (Pwave) ; [ud][su] \bar{u}

$\Lambda(1670)$ uds (Pwave) ; [us][ds] \bar{s}

Zou *et al*, NPA835 (2010) 199 ; CLAS, PRC 87(2013) 035206

Dibaryon: H, d*, N- Ω , $\Omega-\Omega$...

J. L. Ping *et al*, PRC 79(2009) 024001 ; PRC 83(2011) 015202

R. L. Jaffe, PRL 38(1977) 195

- **Pentaquark:** Θ^+ , Ξ^- , P_c^+ ...

T. Nakano *et al*, PRL 91(2003) 012002

R. Aaij *et al*, PRL 115(2015) 072001

Theoretical interpretation:

- **Couple-channel unitary approach:**
B. S. Zou *et al*, PRL 105(2010) 232001; PRC 84(2011) 015202
- **One-boson-exchange model:**
X. Liu , S. L. Zhu *et al*, CPC 36(2012) 6
R. Chen, X. Liu, X. Q. Li and S. L. Zhu, PRL 115(2015) 132002
- **Bethe-Salpeter equation method:**
J. He , PLB 753(2016) 547
- **QCD sum rule approach:**
H. X. Chen, X. Liu and S. L. Zhu *et al*, PRL 115(2015) 172001
Z. G. Wang, EPJC 76(2016) 70

- **Diquark-triquark model:**

R. L. Zhu, C. F. Qiao, PLB 756(2016) 259

- **Non-resonance explanations:**

Q. Zhao *et al*, PLB 757(2016) 231

F. K. Guo *et al*, PRD 92(2015) 071502

- **Tetraquark: X(3872), Y(3940), Z(3900)...**

The stage of multi-quark systems may coming!

ChQM and calculation methods

- ChQM has acquired great achievement both in describing the hadron spectra and hadron-hadron interaction.

A. Valcarce *et al*, Rep. Prog. Phys. 68(2005) 965-1041

A. Valcarce *et al*, J. Phys. G: Nucl. Part. Phys. 31(2005) 481-506

- General form of Hamiltonian

$$H = \sum_{i=1}^5 \left(m_i + \frac{p_i^2}{2m_i} \right) - T_{CM} + \sum_{j>i=1}^5 V_{ij}^C(\vec{r}_{ij})$$

- **The central part of potential**

- *Screened potential:*

$$V_{CON}(\mathbf{r}_{ij}) = \boldsymbol{\lambda}_i^c \cdot \boldsymbol{\lambda}_j^c \left[-a_c(1 - e^{-\mu_c r_{ij}}) + \Delta \right],$$

- *OGE:*

$$V_{OGE}(\mathbf{r}_{ij}) = \frac{1}{4}\alpha_s \boldsymbol{\lambda}_i^c \cdot \boldsymbol{\lambda}_j^c \left[\frac{1}{r_{ij}} - \frac{1}{6m_i m_j} \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j \frac{e^{-r_{ij}/r_0(\mu)}}{r_{ij} r_0^2(\mu)} \right], \quad r_0(\mu) = \hat{r}_0/\mu,$$

- *Quark-quark interaction:*

$$V_\chi(\mathbf{r}_{ij}) = v_\pi(\mathbf{r}_{ij}) \sum_{a=1}^3 \boldsymbol{\lambda}_i^a \cdot \boldsymbol{\lambda}_j^a + v_K(\mathbf{r}_{ij}) \sum_{a=4}^7 \boldsymbol{\lambda}_i^a \cdot \boldsymbol{\lambda}_j^a + v_\eta(\mathbf{r}_{ij}) [\boldsymbol{\lambda}_i^8 \cdot \boldsymbol{\lambda}_j^8 \cos \theta_P - \boldsymbol{\lambda}_i^0 \cdot \boldsymbol{\lambda}_j^0 \sin \theta_P],$$

$$v_\chi(\mathbf{r}_{ij}) = \frac{g_{ch}^2}{4\pi} \frac{m_\chi^2}{12m_i m_j} \frac{\Lambda_\chi^2}{\Lambda_\chi^2 - m_\chi^2} m_\chi \left[Y(m_\chi r_{ij}) - \frac{\Lambda_\chi^3}{m_\chi^3} Y(\Lambda_\chi r_{ij}) \right] \boldsymbol{\sigma}_i \cdot \boldsymbol{\sigma}_j, \quad \chi = \pi, K, \eta,$$

$$V_s(\mathbf{r}_{ij}) = v_\sigma(\mathbf{r}_{ij})(\boldsymbol{\lambda}_i^0 \cdot \boldsymbol{\lambda}_j^0) + v_{a_0}(\mathbf{r}_{ij}) \sum_{a=1}^3 \boldsymbol{\lambda}_i^a \cdot \boldsymbol{\lambda}_j^a + v_\kappa(\mathbf{r}_{ij}) \sum_{a=4}^7 \boldsymbol{\lambda}_i^a \cdot \boldsymbol{\lambda}_j^a + v_{f_0}(\mathbf{r}_{ij})(\boldsymbol{\lambda}_i^8 \cdot \boldsymbol{\lambda}_j^8),$$

$$v_s(\mathbf{r}_{ij}) = -\frac{g_{ch}^2}{4\pi} \frac{\Lambda_s^2}{\Lambda_s^2 - m_s^2} m_s \left[Y(m_s r_{ij}) - \frac{\Lambda_s}{m_s} Y(\Lambda_s r_{ij}) \right], \quad s = \sigma, a_0, \kappa, f_0.$$

• Calculation method

In solving the Schrodinger equation by variation method

$$\langle \Psi | H - E | \Psi \rangle = 0$$

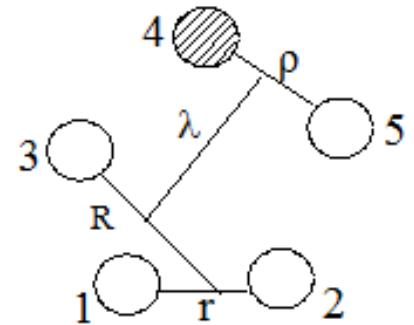
The complete wave function for 5-body

$$\Psi = \sum_{g=1}^{N_{total}} C_g A \phi_{n_1 l_1 m_1} (\vec{r}) \psi_{n_2 l_2 m_2} (\vec{R}) \omega_{n_3 l_3 m_3} (\vec{\rho}) \eta_{n_4 l_4 m_4} (\vec{\lambda}) \chi_i^\sigma \chi_j^f \chi_k^c$$

$g = \{ n_1, l_1, m_1, n_2, l_2, m_2, n_3, l_3, m_3, n_4, l_4, m_4, i, j, k \}$

Antisymmetry operator:

$$A = 1 - (1\ 3) - (1\ 5) - (2\ 3) - (2\ 5) - (3\ 5) \\ + (1\ 3\ 5) + (1\ 5\ 3) + (2\ 3\ 5) + (2\ 5\ 3) + (1\ 3)(2\ 5) + (1\ 5)(2\ 3)$$



• Gaussian basis function

E. Hiyama et al., Prog. Part. Nucl. Phys. 51(2003) 223-307

The spatial part is constructed by an expansion of a series of different size of Gaussian basis.

$$\begin{aligned}\phi_{nlm}(\vec{r}) &= \phi_{nl}^G(r) Y_{lm}(\hat{r}) \\ &= N_{nl} r^l e^{-\nu_n r^2} Y_{lm}(\hat{r}) \\ &= \sqrt{\frac{2^{l+2} (2\nu_n)^{\frac{l+3}{2}}}{\sqrt{\pi} (2l+1)!!}} r^l e^{-\nu_n r^2} Y_{lm}(\hat{r})\end{aligned}$$

Transform into the ISG form,

$$\phi_{nlm}(\vec{r}) = N_{nl} \lim_{\varepsilon \rightarrow 0} \left(\frac{l}{4\nu_n \varepsilon} \right)^{l \left| \frac{l-m}{2} \right|} \sum_{j=0}^p A_{lm,j} \sum_{s=0}^p \sum_{t=0}^q \sum_{u=0}^j C_p^s C_q^t C_j^u e^{-\nu_n (\vec{r} - \varepsilon \vec{D})^2}$$

$$A_{lm,j} = \sqrt{\frac{(2l+1)(l-m)!}{4\pi(l+m)!}} \frac{(l+m)!}{2^m} \frac{(-1)^j}{4^j j! (m+j)! (l-m-2j)!} \quad \begin{array}{l} p = l-m-2j \\ q = m+j \end{array}$$

The parameter of Gaussian is taken as the geometric progression numbers.

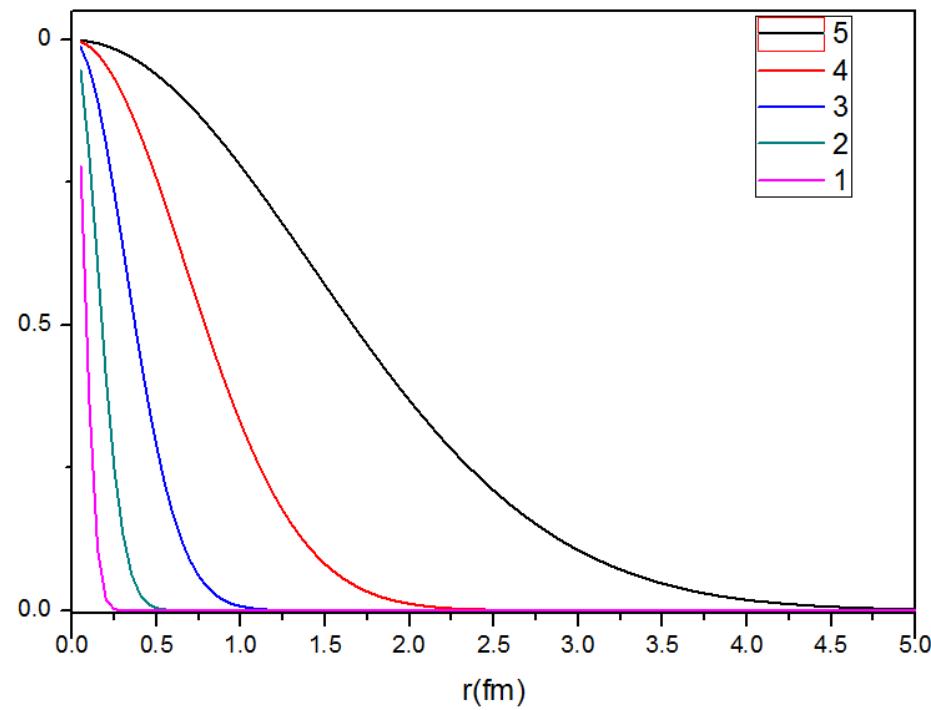
$$\nu_n = [r_{\max}^{N-1} \cdot r_{\min}^{\frac{(n-1)(N-2)}{N-1}}]^{-2}$$

e.g.

$$r_{\min} = 0.1 \text{ fm}; r_{\max} = 2.0 \text{ fm}; N = 5.$$

$$\nu_1 = 100; \nu_2 = 22.4; \nu_3 = 5; \nu_4 = 1.1; \nu_5 = 0.3;$$

$$\varphi_n = e^{-\nu_n \cdot r^2}$$

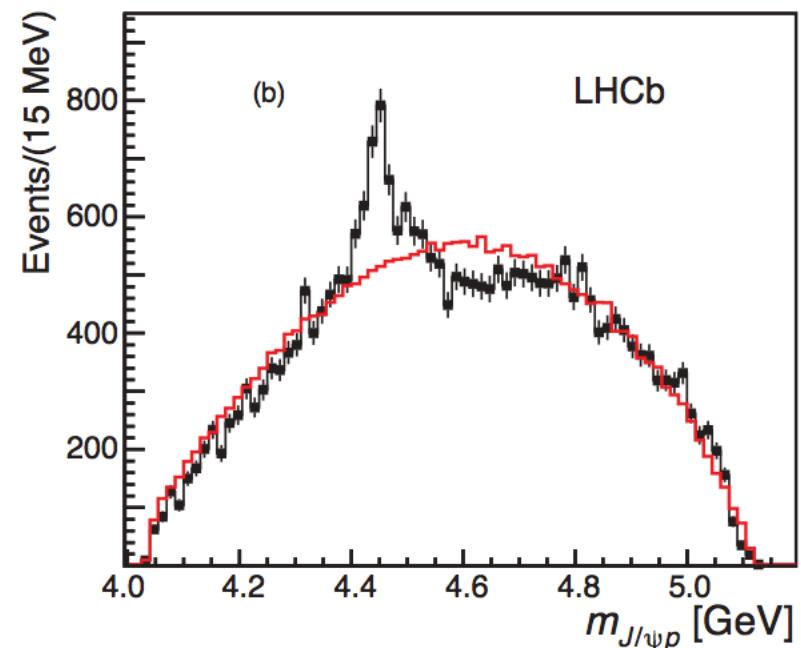
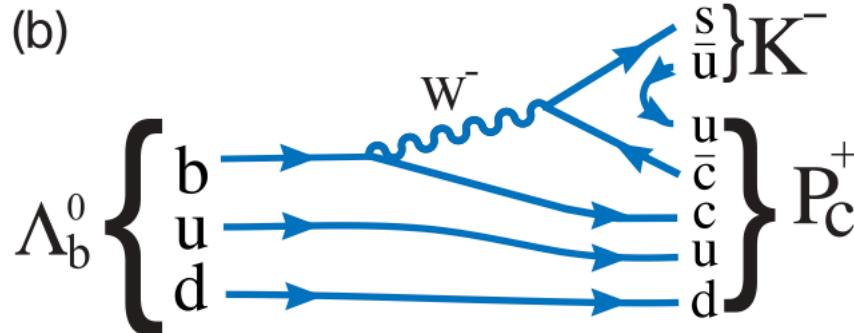


The results on P_c^+ and Ω_c^0 from LHCb data

- Two pentaquark states have been discovered in the $J/\psi p$ invariant mass spectrum.

$P_c^+(4380) : 4380 \pm 8 \pm 29 \text{ MeV}; \Gamma = 205 \pm 18 \pm 86 \text{ MeV}.$

$P_c^+(4449.8) : 4449.8 \pm 1.7 \pm 2.5 \text{ MeV}; \Gamma = 39 \pm 5 \pm 19 \text{ MeV}.$



R. Aaij et al, PRL 115(2015) 072001

- Flavor wave-function of P_c^+ in $I=1/2$.

$$\begin{aligned}\chi_1^f &= \sqrt{\frac{2}{3}}B_{11}M_{\frac{1}{2},-\frac{1}{2}} - \sqrt{\frac{1}{3}}B_{10}M_{\frac{1}{2},\frac{1}{2}}, \\ \chi_2^f &= B_{00}M_{\frac{1}{2},\frac{1}{2}}, \\ \chi_3^f &= B_{\frac{1}{2},\frac{1}{2}}^1 M_{00}, \\ \chi_4^f &= B_{\frac{1}{2},\frac{1}{2}}^2 M_{00}.\end{aligned}$$

5q

Sub-cluster

$$\left\{ \begin{array}{l} B_{11} = uuc, \quad B_{10} = \frac{1}{\sqrt{2}}(ud + du)c, \quad B_{1-1} = ddc, \\ B_{00} = \frac{1}{\sqrt{2}}(ud - du)c, \\ B_{\frac{1}{2},\frac{1}{2}}^1 = \frac{1}{\sqrt{6}}(2uud - udu - duu), \\ B_{\frac{1}{2},\frac{1}{2}}^2 = \frac{1}{\sqrt{2}}(ud - du)u, \\ M_{\frac{1}{2},\frac{1}{2}} = \bar{c}u, \quad M_{\frac{1}{2},-\frac{1}{2}} = \bar{c}d, \quad M_{00} = \bar{c}c. \end{array} \right. \quad (4)$$

- Spin wave-function of P_c^+ .

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 1}(5) = \sqrt{\frac{1}{6}} \chi_{\frac{3}{2}, -\frac{1}{2}}^{\sigma}(3) \chi_{11}^{\sigma} - \sqrt{\frac{1}{3}} \chi_{\frac{3}{2}, \frac{1}{2}}^{\sigma}(3) \chi_{10}^{\sigma} + \sqrt{\frac{1}{2}} \chi_{\frac{3}{2}, \frac{3}{2}}^{\sigma}(3) \chi_{1-1}^{\sigma}$$

Sub-cluster

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 2}(5) = \sqrt{\frac{1}{3}} \chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 1}(3) \chi_{10}^{\sigma} - \sqrt{\frac{2}{3}} \chi_{\frac{1}{2}, -\frac{1}{2}}^{\sigma 1}(3) \chi_{11}^{\sigma}$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 3}(5) = \sqrt{\frac{1}{3}} \chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 2}(3) \chi_{10}^{\sigma} - \sqrt{\frac{2}{3}} \chi_{\frac{1}{2}, -\frac{1}{2}}^{\sigma 2}(3) \chi_{11}^{\sigma}$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 4}(5) = \chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 1}(3) \chi_{00}^{\sigma}$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 5}(5) = \chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 2}(3) \chi_{00}^{\sigma}$$

$$\chi_{\frac{3}{2}, \frac{3}{2}}^{\sigma 1}(5) = \sqrt{\frac{3}{5}} \chi_{\frac{3}{2}, \frac{3}{2}}^{\sigma}(3) \chi_{10}^{\sigma} - \sqrt{\frac{2}{5}} \chi_{\frac{3}{2}, \frac{1}{2}}^{\sigma}(3) \chi_{11}^{\sigma}$$

$$\chi_{\frac{3}{2}, \frac{3}{2}}^{\sigma 2}(5) = \chi_{\frac{3}{2}, \frac{3}{2}}^{\sigma}(3) \chi_{00}^{\sigma}$$

$$\chi_{\frac{3}{2}, \frac{3}{2}}^{\sigma 3}(5) = \chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 1}(3) \chi_{11}^{\sigma}$$

$$\chi_{\frac{3}{2}, \frac{3}{2}}^{\sigma 4}(5) = \chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 2}(3) \chi_{11}^{\sigma}$$

$$\chi_{\frac{5}{2}, \frac{5}{2}}^{\sigma 1}(5) = \chi_{\frac{3}{2}, \frac{3}{2}}^{\sigma}(3) \chi_{11}^{\sigma}$$

Total part

$$\chi_{\frac{3}{2}, \frac{3}{2}}^{\sigma}(3) = \alpha \alpha \alpha,$$

$$\chi_{\frac{3}{2}, \frac{1}{2}}^{\sigma}(3) = \frac{1}{\sqrt{3}} (\alpha \alpha \beta + \alpha \beta \alpha + \beta \alpha \alpha),$$

$$\chi_{\frac{3}{2}, -\frac{1}{2}}^{\sigma}(3) = \frac{1}{\sqrt{3}} (\alpha \beta \beta + \beta \alpha \beta + \beta \beta \alpha),$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 1}(3) = \frac{1}{\sqrt{6}} (2 \alpha \alpha \beta - \alpha \beta \alpha - \beta \alpha \alpha),$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{\sigma 2}(3) = \frac{1}{\sqrt{2}} (\alpha \beta \alpha - \beta \alpha \alpha),$$

$$\chi_{\frac{1}{2}, -\frac{1}{2}}^{\sigma 1}(3) = \frac{1}{\sqrt{6}} (\alpha \beta \beta - \alpha \beta \beta - 2 \beta \beta \alpha),$$

$$\chi_{\frac{1}{2}, -\frac{1}{2}}^{\sigma 2}(3) = \frac{1}{\sqrt{2}} (\alpha \beta \beta - \beta \alpha \beta),$$

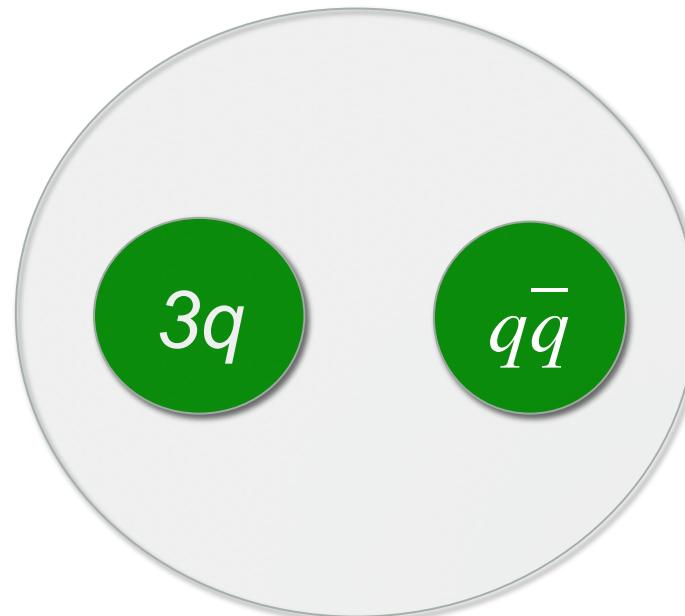
$$\chi_{11}^{\sigma} = \alpha \alpha, \quad \chi_{10}^{\sigma} = \frac{1}{\sqrt{2}} (\alpha \beta + \beta \alpha), \quad \chi_{1-1}^{\sigma} = \beta \beta,$$

$$\chi_{00}^{\sigma} = \frac{1}{\sqrt{2}} (\alpha \beta - \beta \alpha).$$

- Color wave-function of P_c^+ .

I. *Color singlet channel*

$$\chi_1^c = \frac{1}{\sqrt{18}} (rgb - rbg + gbr - grb + brg - bgr) \\ \times (\bar{r}r + \bar{g}g + \bar{b}b),$$



$[1] \otimes [1]$

II. Hidden color channel

[8] \otimes [8]

$$\chi_k^c = \frac{1}{\sqrt{8}} \left(\chi_{3,1}^k \chi_{2,8}^k - \chi_{3,2}^k \chi_{2,7}^k - \chi_{3,3}^k \chi_{2,6}^k + \chi_{3,4}^k \chi_{2,5}^k + \chi_{3,5}^k \chi_{2,4}^k - \chi_{3,6}^k \chi_{2,3}^k - \chi_{3,7}^k \chi_{2,2}^k + \chi_{3,8}^k \chi_{2,1}^k \right)$$

$$\chi_{3,1}^2 = \frac{1}{\sqrt{6}} (2rrg - rgr - grr), \quad \chi_{3,1}^3 = \frac{1}{\sqrt{2}} (rgr - grr),$$

$$\chi_{3,2}^2 = \frac{1}{\sqrt{6}} (rgg + grg - 2ggr), \quad \chi_{3,2}^3 = \frac{1}{\sqrt{2}} (rgg - grg),$$

$$\chi_{3,3}^2 = \frac{1}{\sqrt{6}} (2rrb - rbr - brr), \quad \chi_{3,3}^3 = \frac{1}{\sqrt{2}} (rbr - brr),$$

$$\chi_{3,4}^2 = \frac{1}{\sqrt{12}} (2rgb - rbg + 2grb - gbr - brg - bgr),$$

$$\chi_{3,4}^3 = \frac{1}{\sqrt{4}} (rbg + gbr - brg - bgr),$$

$$\chi_{2,1} = \bar{b}r, \quad \chi_{2,2} = \bar{b}g$$

$$\chi_{2,3} = -\bar{g}r, \quad \chi_{2,4} = \frac{1}{\sqrt{2}} (\bar{r}r - \bar{g}g),$$

$$\chi_{2,5} = \frac{1}{\sqrt{6}} (2\bar{b}b - \bar{r}r - \bar{g}g), \quad \chi_{2,6} = \bar{r}g$$

$$\chi_{2,7} = -\bar{g}b, \quad \chi_{2,8} = \bar{r}b.$$

Sub-cluster

$$\chi_{3,5}^2 = \frac{1}{\sqrt{4}} (rbg - gbr + brg - bgr),$$

$$\chi_{3,5}^3 = \frac{1}{\sqrt{12}} (2rgb + rbg - 2grb - gbr - brg + bgr),$$

$$\chi_{3,6}^2 = \frac{1}{\sqrt{6}} (2ggb - gbg - bgg), \quad \chi_{3,6}^3 = \frac{1}{\sqrt{2}} (gbg - bgg),$$

$$\chi_{3,7}^2 = \frac{1}{\sqrt{6}} (rbb + brb - 2bbr), \quad \chi_{3,7}^3 = \frac{1}{\sqrt{2}} (rbb - brb),$$

$$\chi_{3,8}^2 = \frac{1}{\sqrt{6}} (gbb + bgb - 2bbg), \quad \chi_{3,8}^3 = \frac{1}{\sqrt{2}} (gbb - bgb),$$

TABLE IV: The lowest eigen-energies of the $udcc\bar{u}$ system with $J^P = \frac{3}{2}^-$ (unit: MeV).

Channel	E	E_{th}^{Theo}	E_B	E_{th}^{Exp}	E'
$J^P = 3/2^-$					
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 3, 4, j = 3, 4, k = 1$	3841	3841	0	4036(NJ/ψ)	4036
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 3, 4, j = 3, 4, k = 2, 3$	4722				
color-singlet+hidden color	3841				
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 3, 4, j = 2, k = 1$	4115	4115	0	4293($\Lambda_c \bar{D}^*$)	4293
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 3, 4, j = 2, k = 2, 3$	4680				
color-singlet+hidden color	4115				
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 3, 4, j = 1, k = 1$	4518	4520	-2	4462($\Sigma_c \bar{D}^*$)	4460
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 3, 4, j = 1, k = 2, 3$	4961				
color-singlet+hidden color	4517	4520	-3	4462	4459
percentage(S;H): 96.3%; 3.7%					
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 2, j = 1, k = 1$	4444	4447	-3	4385($\Sigma_c^* \bar{D}$)	4382
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 2, j = 1, k = 2, 3$	4754				
color-singlet+hidden color	4432	4447	-15	4385	4370
percentage(S;H): 82.6%; 17.4%					
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 1, j = 1, k = 1$	4564	4566	-2	4527($\Sigma_c^* \bar{D}^*$)	4525
$\chi_{3/2}^{\sigma i} \chi_j^f \chi_k^c$ $i = 1, j = 1, k = 2, 3$	4623				
color-singlet+hidden color	4549	4566	-17	4527	4510
percentage(S;H): 61.1%; 38.9%					
mixed (only color-singlet)	3841				
mixed (color-singlet+hidden color)	3841				

P_c^+ (4380)

unbound

JP=1/2-

Channel	E	E_{th}^{Theo}	E_B	E_{th}^{Exp}	E'
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 4, 5, \quad j = 3, 4, \quad k = 1$	3745	3745	0	$3919(N\eta_c)$	3919
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 4, 5, \quad j = 3, 4, \quad k = 2, 3$	4714				
color-singlet+hidden color	3745				
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 2, 3, \quad j = 3, 4, \quad k = 1$	3841	3841	0	$4036(NJ/\psi)$	4036
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 2, 3, \quad j = 3, 4, \quad k = 2, 3$	4964				
color-singlet+hidden color	3841				
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 4, 5, \quad j = 2, \quad k = 1$	3996	3996	0	$4151(\Lambda_c \bar{D})$	4151
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 4, 5, \quad j = 2, \quad k = 2, 3$	4663				
color-singlet+hidden color	3996				
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 2, 3, \quad j = 2, \quad k = 1$	4115	4115	0	$4293(\Lambda_c \bar{D}^*)$	4293
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 2, 3, \quad j = 2, \quad k = 2, 3$	4599				
color-singlet+hidden color	4115				
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 4, 5, \quad j = 1, \quad k = 1$	4398	4402	-4	$4320(\Sigma_c \bar{D})$	4316
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 4, 5, \quad j = 1, \quad k = 2, 3$	4835				
color-singlet+hidden color	4394	4402	-8	4320	4312
		percentage(S;H):	91.0%; 7.0%		
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 2, 3, \quad j = 1, \quad k = 1$	4518	4520	-2	$4462(\Sigma_c \bar{D}^*)$	4460
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 2, 3, \quad j = 1, \quad k = 2, 3$	4728				
color-singlet+hidden color	4479	4520	-41	4462	4421
		percentage(S;H):	67.4%; 32.6%		
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 1, \quad j = 1, \quad k = 1$	4563	4566	-3	$4527(\Sigma_c^* \bar{D}^*)$	4524
$\chi_{1/2}^{\sigma i} \chi_j^f \chi_k^c \quad i = 1, \quad j = 1, \quad k = 2, 3$	4476				
color-singlet+hidden color	4461	4566	-105	4527	4422
		percentage(S;H):	23.0%; 77.0%		
mixed (only color singlet)	3745				
mixed (color singlet+hidden color)	3745				

Scattering state

$J^P = 5/2^-$						
$\chi_{5/2}^{\sigma i} \chi_j^f \chi_k^c$	$i = 1, j = 1, k = 1$	4563	4566	-3	4527($\Sigma_c^* \bar{D}^*$)	4524
$\chi_{5/2}^{\sigma i} \chi_j^f \chi_k^c$	$i = 1, j = 1, k = 2, 3$	5002				
color-singlet+hidden color		4477	4566	-89	4527	4438
		percentage(S;H): 66.2%; 33.8%				

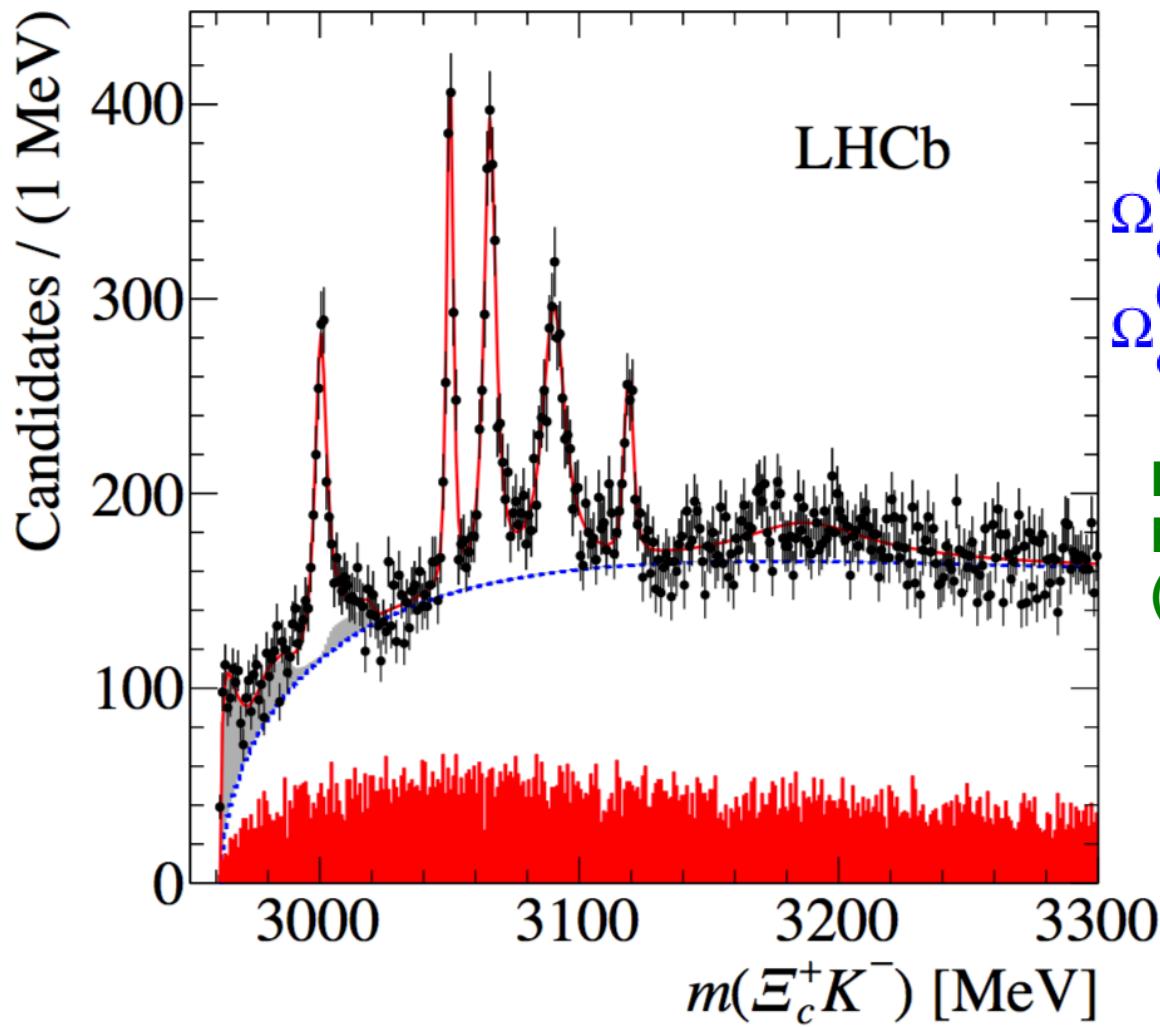
TABLE V: Distances between any two quarks (unit: fm).

J^P	Channel	r_{12}	r_{13}	r_{14}	r_{34}
$\frac{1}{2}^-$	$\chi_{1/2}^{\sigma i} \chi_j^f \ i = 4, 5, j = 1, k = 1 \ (\Sigma_c \bar{D})$	0.8	0.7	2.1	2.1
	$\chi_{1/2}^{\sigma i} \chi_j^f \ i = 4, 5, j = 1, k = 2, 3$	1.0	0.8	0.8	0.4
	$\chi_{1/2}^{\sigma i} \chi_j^f \ i = 2, 3, j = 1, k = 1 \ (\Sigma_c \bar{D}^*)$	0.8	0.7	2.2	2.1
	$\chi_{1/2}^{\sigma i} \chi_j^f \ i = 2, 3, j = 1, k = 2, 3$	0.9	0.8	0.8	0.4
	$\chi_{1/2}^{\sigma i} \chi_j^f \ i = 1, j = 1, k = 1 \ (\Sigma_c^* \bar{D}^*)$	0.9	0.8	2.1	2.0
	$\chi_{1/2}^{\sigma i} \chi_j^f \ i = 1, j = 1, k = 2, 3$	0.9	0.8	0.8	0.4
$\frac{3}{2}^-$	$\chi_{3/2}^{\sigma i} \chi_j^f \ i = 3, 4, j = 1, k = 1 \ (\Sigma_c \bar{D}^*)$	0.8	0.7	2.4	2.3
	$\chi_{3/2}^{\sigma i} \chi_j^f \ i = 3, 4, j = 1, k = 2, 3$	1.1	0.9	0.9	0.5
	$\chi_{3/2}^{\sigma i} \chi_j^f \ i = 2, j = 1, k = 1 \ (\Sigma_c^* \bar{D})$	0.9	0.8	2.2	2.2
	$\chi_{3/2}^{\sigma i} \chi_j^f \ i = 2, j = 1, k = 2, 3$	1.0	0.9	0.9	0.5
	$\chi_{3/2}^{\sigma i} \chi_j^f \ i = 1, j = 1, k = 1 \ (\Sigma_c^* \bar{D}^*)$	0.9	0.8	2.6	2.4
	$\chi_{3/2}^{\sigma i} \chi_j^f \ i = 1, j = 1, k = 2, 3$	0.9	0.9	0.8	0.4
$\frac{5}{2}^-$	$\chi_{5/2}^{\sigma i} \chi_j^f \ i = 1, j = 1, k = 1 \ (\Sigma_c^* \bar{D}^*)$	0.9	0.8	2.4	2.3
	$\chi_{5/2}^{\sigma i} \chi_j^f \ i = 1, j = 1, k = 2, 3$	1.3	1.4	1.3	0.8

udccu

~P_c⁺ (4380)

- Five new Ω_c^0 states are discovered in the $\Xi_c^+ K^-$ mass spectrum of pp collision.



$\Omega_c^0(3000), \Omega_c^0(3050), \Omega_c^0(3066)$
 $\Omega_c^0(3090), \Omega_c^0(3119)$

R. Aaij et al.
Phys. Rev. Lett. 118, 182001
(2017)

- **Theoretical interpretation**

H. X. Chen, X. Liu, S. L. Zhu *et al.* Phys. Rev. D 95, 094008 (2017)

Q. Zhao, X. H. Zhong *et al.* Phys. Rev. D 95, 116010 (2017)

Z. G. Wang Eur. Phys. J. C77 (2017) 325

X. Liu *et al.* arXiv: 1704.02583 [hep-ph]

H. X. Huang, J. L. Ping *et al.* arXiv: 1704.01421 [hep-ph]

...

- Attempt in the $3q$ configuration

Masses of baryons and mesons in ChQM with two sets of parameters.

Too heavy

P	$N(939)$	$\Delta(1232)$	$\Omega(1672)$	$\Lambda(1116)$	$\Sigma(1189)$	$\Xi(1315)$
set I						
+	936	1208	1643	1154	1173	1362
-	1575	1625	2203	1772	1777	1981
set II						
+	939	1231	1671	1187	1209	1408
-	1661	1716	2301	1889	1895	2098
$\Sigma^*(1383)$	$\Xi^*(1532)$	$\Omega_c(2695)$	$\Omega_c(2765)$	$\Xi_c(2467)$	$\Xi_c^*(2645)$	
set I						
+	1342	1488	2675	2748	2541	2603
-	1805	1999	3257	3282	3086	3093
set II						
+	1393	1539	2748	2818	2629	2727
-	1928	2119	3378	3389	3145	3166
set I						
P	$\pi(140)$	$\rho(775)$	$\eta(548)$	$\omega(782)$	$K(495)$	$K^*(892)$
-	93	800	611	705	326	965
$\eta'(958)$	$\phi(1019)$	$D^0(1865)$	$D^*(2007)$			
-	914	1056	1842	2043		

A test on our results with E. Hiyama
et al. PRD 92(2015) 114029.



TABLE XVI: P-wave results. (unit: MeV)

$J^P = \frac{1}{2}^-$	Λ_c	Σ_c	Ω_c	Ξ_{cc}
I	2632	2813	3037	3952
HIYAMA	2628	2802	3030	3947

- The calculated masses of Ω_c^0 in ground states and low-lying excited states

nL	$\Omega_c(2695)$	$\Omega_c(2765)$
set I		
$1S$	2675	2748
$2S$	3074	3112
$1P$	3257	3282
set II		
$1S$	2748	2818
$2S$	3167	3201
$1P$	3378	3389

- Flavor wave-function of 5-quark systems.

$$B_{00}^1 = ssc, \quad B_{00}^2 = sss,$$

$$B_{\frac{1}{2}, \frac{1}{2}}^1 = \frac{1}{\sqrt{6}}(sus + uss - 2ssu),$$

$$B_{\frac{1}{2}, -\frac{1}{2}}^1 = \frac{1}{\sqrt{6}}(sds + dss - 2ssd),$$

$$B_{\frac{1}{2}, \frac{1}{2}}^2 = \frac{1}{\sqrt{2}}(us - su)s, \quad B_{\frac{1}{2}, -\frac{1}{2}}^2 = \frac{1}{\sqrt{2}}(ds - sd)s,$$

$$B_{\frac{1}{2}, \frac{1}{2}}^3 = \frac{1}{\sqrt{3}}(ssu + sus + uss),$$

$$B_{\frac{1}{2}, -\frac{1}{2}}^3 = \frac{1}{\sqrt{3}}(ssd + sds + dss),$$

$$B_{\frac{1}{2}, \frac{1}{2}}^4 = \frac{1}{\sqrt{2}}(us + su)c, \quad B_{\frac{1}{2}, -\frac{1}{2}}^4 = \frac{1}{\sqrt{2}}(ds + sd)c,$$

$$B_{\frac{1}{2}, \frac{1}{2}}^5 = \frac{1}{\sqrt{2}}(us - su)c, \quad B_{\frac{1}{2}, -\frac{1}{2}}^5 = \frac{1}{\sqrt{2}}(ds - sd)c.$$

$$M_{\frac{1}{2}, \frac{1}{2}}^1 = \bar{d}c, \quad M_{\frac{1}{2}, -\frac{1}{2}}^1 = -\bar{u}c,$$

$$M_{\frac{1}{2}, \frac{1}{2}}^2 = \bar{d}s, \quad M_{\frac{1}{2}, -\frac{1}{2}}^2 = -\bar{u}s,$$

$$M_{00}^1 = \frac{1}{\sqrt{2}}(\bar{u}u + \bar{d}d), \quad M_{00}^2 = \bar{s}s, \quad M_{00}^3 = \bar{s}c.$$

$$\begin{aligned}\chi_1^f &= \sqrt{\frac{1}{2}}(B_{\frac{1}{2}, \frac{1}{2}}^1 M_{\frac{1}{2}, -\frac{1}{2}}^1 - B_{\frac{1}{2}, -\frac{1}{2}}^1 M_{\frac{1}{2}, \frac{1}{2}}^1), \\ \chi_2^f &= \sqrt{\frac{1}{2}}(B_{\frac{1}{2}, \frac{1}{2}}^2 M_{\frac{1}{2}, -\frac{1}{2}}^1 - B_{\frac{1}{2}, -\frac{1}{2}}^2 M_{\frac{1}{2}, \frac{1}{2}}^1), \\ \chi_3^f &= \sqrt{\frac{1}{2}}(B_{\frac{1}{2}, \frac{1}{2}}^3 M_{\frac{1}{2}, -\frac{1}{2}}^1 - B_{\frac{1}{2}, -\frac{1}{2}}^3 M_{\frac{1}{2}, \frac{1}{2}}^1), \\ \chi_4^f &= \sqrt{\frac{1}{2}}(B_{\frac{1}{2}, \frac{1}{2}}^4 M_{\frac{1}{2}, -\frac{1}{2}}^1 - B_{\frac{1}{2}, -\frac{1}{2}}^4 M_{\frac{1}{2}, \frac{1}{2}}^1), \\ \chi_5^f &= \sqrt{\frac{1}{2}}(B_{\frac{1}{2}, \frac{1}{2}}^5 M_{\frac{1}{2}, -\frac{1}{2}}^1 - B_{\frac{1}{2}, -\frac{1}{2}}^5 M_{\frac{1}{2}, \frac{1}{2}}^1), \\ \chi_6^f &= B_{00}^1 M_{00}^1, \quad \chi_7^f = B_{00}^1 M_{00}^2, \quad \chi_8^f = B_{00}^2 M_{00}^3.\end{aligned}$$

5q part

Sub-cluster

• Possible channels

TABLE III: The channels with $IJ^P = 0\frac{1}{2}^-$.

index	$\chi_{1/2}^{\sigma_i}$	χ_j^f	χ_k^c	physical channel
1	$i = 1$	$j = 3$	$k = 1$	$\Xi^* \bar{D}^*$
2	$i = 1$	$j = 3$	$k = 3$	
3	$i = 1$	$j = 4$	$k = 1$	$\Xi_c^* \bar{K}^*$
4	$i = 1$	$j = 4, 5$	$k = 2, 3$	
5	$i = 1$	$j = 6$	$k = 1$	$\Omega_c^* \omega$
6	$i = 1$	$j = 6$	$k = 3$	
7	$i = 2, 3$	$j = 1, 2$	$k = 1$	$\Xi \bar{D}^*$
8	$i = 2, 3$	$j = 1, 2$	$k = 2, 3$	
9	$i = 2, 3$	$j = 4, 5$	$k = 1$	$\Xi_c \bar{K}^*$
10	$i = 2, 3$	$j = 4, 5$	$k = 2, 3$	
11	$i = 2$	$j = 6$	$k = 1$	$\Omega_c \omega$
12	$i = 2, 3$	$j = 6$	$k = 2, 3$	
13	$i = 4, 5$	$j = 1, 2$	$k = 1$	$\Xi \bar{D}$
14	$i = 4, 5$	$j = 1, 2$	$k = 2, 3$	
15	$i = 4, 5$	$j = 4, 5$	$k = 1$	$\Xi_c \bar{K}$
16	$i = 4, 5$	$j = 4, 5$	$k = 2, 3$	
17	$i = 4$	$j = 6$	$k = 1$	$\Omega_c \eta$
18	$i = 4, 5$	$j = 6$	$k = 2, 3$	

TABLE IV: The channels with $IJ^P = 0\frac{3}{2}^-$.

index	$\chi_{3/2}^{\sigma_i}$	χ_j^f	χ_k^c	physical channel
1	$i = 1$	$j = 3$	$k = 1$	$\Xi^* \bar{D}^*$
2	$i = 1$	$j = 3$	$k = 3$	
3	$i = 1$	$j = 4$	$k = 1$	$\Xi_c^* \bar{K}^*$
4	$i = 1$	$j = 4, 5$	$k = 2, 3$	
5	$i = 1$	$j = 6$	$k = 1$	$\Omega_c^* \omega$
6	$i = 1$	$j = 6$	$k = 3$	
7	$i = 2$	$j = 3$	$k = 1$	$\Xi^* \bar{D}$
8	$i = 2$	$j = 3$	$k = 3$	
9	$i = 2$	$j = 4$	$k = 1$	$\Xi_c^* \bar{K}$
10	$i = 2$	$j = 4, 5$	$k = 2, 3$	
11	$i = 2$	$j = 6$	$k = 1$	$\Omega_c^* \eta$
12	$i = 2$	$j = 6$	$k = 3$	
13	$i = 3, 4$	$j = 1, 2$	$k = 1$	$\Xi \bar{D}^*$
14	$i = 3, 4$	$j = 1, 2$	$k = 2, 3$	
15	$i = 3, 4$	$j = 4, 5$	$k = 1$	$\Xi_c \bar{K}^*$
16	$i = 3, 4$	$j = 4, 5$	$k = 2, 3$	
17	$i = 3$	$j = 6$	$k = 1$	$\Omega_c \omega$
18	$i = 3, 4$	$j = 6$	$k = 2, 3$	

TABLE V: The channels with $IJ^P = 0\frac{5}{2}^-$.

index	$\chi_{5/2}^{\sigma_i}$	χ_j^f	χ_k^c	physical channel
1	$i = 1$	$j = 3$	$k = 1$	$\Xi^* \bar{D}^*$
2	$i = 1$	$j = 3$	$k = 3$	
3	$i = 1$	$j = 4$	$k = 1$	$\Xi_c^* \bar{K}^*$
4	$i = 1$	$j = 4, 5$	$k = 2, 3$	
5	$i = 1$	$j = 6$	$k = 1$	$\Omega_c^* \omega$
6	$i = 1$	$j = 6$	$k = 3$	

- Calculation in the 5-quark configuration

TABLE VII: The lowest eigen-energies of the $ss\bar{c}u\bar{u} + ss\bar{c}d\bar{d}$ system with $J^P = \frac{1}{2}^-$ (unit: MeV). The percentages of color-singlet (S) and hidden-color (H) channels are also given.

Channel	E	E_{th}^{Theo}	E_B	E_{th}^{Exp}	E'
1	3526	3531	-5	3539($\Xi^* \bar{D}^*$)	3534
2	4016				
1+2	3525		-6		3533
				percentage(S;H): 99.8%; 0.2%	
3	3566	3568	-2	3537($\Xi_c^* \bar{K}^*$)	3535
4	3616				
3+4	3564		-4		3533
				percentage(S;H): 96.3%; 3.7%	
5	3453	3453	0	3548($\Omega_c^* \omega$)	3548
6	3404				
5+6	3402		-51		3497
				percentage(S;H): 0.2%; 99.8%	
7	3374	3405	-31	3322($\Xi \bar{D}^*$)	3291
8	3672				
7+8	3373		-32		3290
				percentage(S;H): 99.8%; 0.2%	
9	3495	3506	-11	3359($\Xi_c \bar{K}^*$)	3348
10	3613				
9+10	3472		-34		3325
				percentage(S;H): 85.2%; 14.8%	

$ss\bar{c}u\bar{u}; ss\bar{c}d\bar{d}$

11		3380	3380	0	3477($\Omega_c \omega$)	3477	
12		3608					
11+12		3380					
13		3175	3204	-29	3185($\Xi \bar{D}$)	3156	
14		3811					
13+14		3175		-29		3156	
					percentage(S;H): 100.0%; 0.0%		
15		2867	2867	0	2961($\Xi_c \bar{K}$)	2961	
16		3807					
15+16		2855		-12		2949	
					percentage(S;H): 96.7%; 3.3%		
17		3286	3286	0	3243($\Omega_c \eta$)	3243	
18		3828					
17+18		3286					
mixed (singlet)		2771	2867	-96	2961($\Xi_c \bar{K}$)	2865	
mixed (full)		2675	2867	-192	2961($\Xi_c \bar{K}$)	2769	

TABLE IX: The lowest eigen-energies of the $ssc\bar{u}u+ssc\bar{d}d$ system with $\frac{3}{2}^-$ (unit: MeV).

Channel	E	E_{th}^{Theo}	E_B	E_{th}^{Exp}	E'				
1	3521	3531	-10	3539($\Xi^*\bar{D}^*$)	3529				
2	4026								
1+2	3521		-10		3529				
				percentage(S;H): 100.0%; 0.0%					
3	3565	3568	-3	3537($\Xi_c^*\bar{K}^*$)	3534				
4	3617								
3+4	3562		-6		3531				
				percentage(S;H): 94.0%; 6.0%					
5	3453	3453	0	3548($\Omega_c^*\omega$)	3548				
6	3477								
5+6	3453								
7	3309	3330	-21	3397($\Xi^*\bar{D}$)	3376				
8	4145								
7+8	3309		-21		3376				
				percentage(S;H): 100.0%; 0.0%					
9	2929	2929	0	3139($\Xi_c^*\bar{K}$)	3139				
10	3782								
9+10	2928		-1		3138				
				percentage(S;H): 99.7%; 0.3%					

$\sim \Omega_c^0(3066)$

TABLE X: The lowest eigen-energies of the $ssc\bar{u}u+ssc\bar{d}d$ system with $\frac{5}{2}^-$ (unit: MeV).

Channel	E	E_{th}^{Theo}	E_B	E_{th}^{Exp}	E'
1	3508	3531	-23	3539($\Xi^* \bar{D}^*$)	3516
2	4042				
1+2	3507		-24		3515
				percentage(S;H): 99.8%; 0.2%	
3	3568	3568	0	3537($\Xi_c^* \bar{K}^*$)	3537
4	3646				
3+4	3532		-36		3501
				percentage(S;H): 80.0%; 20.0%	
5	3453	3453	0	3548($\Omega_c^* \omega$)	3548
6	3563				
5+6	3453				
mixed (singlet)	3453				
mixed (full)	3453				

unbound

TABLE XI: Distances between quarks, q is for u, d quark and Q is for c quark (unit: fm).

J^P	Channel	r_{qq}	r_{qQ}	$r_{q\bar{q}}$	$r_{Q\bar{q}}$
$\frac{1}{2}^-$	$\Omega_c^0(2769)$	1.3	1.1	1.4	1.2
$\frac{3}{2}^-$	$\Omega_c^0(3067)$	1.4	1.1	1.2	1.4

compact

5-quark study in N, P sector

- Flavor wave-function of 5-quark systems.

$$B_{\frac{3}{2}, \frac{3}{2}}^1 = uuu, \quad B_{\frac{3}{2}, -\frac{3}{2}}^1 = ddd,$$

$$B_{\frac{3}{2}, \frac{1}{2}}^1 = \frac{1}{\sqrt{3}}(uud + udu + duu), \quad B_{\frac{3}{2}, -\frac{1}{2}}^1 = \frac{1}{\sqrt{3}}(udd + dud + ddu),$$

$$B_{\frac{1}{2}, \frac{1}{2}}^1 = \frac{1}{\sqrt{6}}(2uud - udu - duu), \quad B_{\frac{1}{2}, \frac{1}{2}}^2 = \frac{1}{\sqrt{2}}(ud - du)u,$$

$$B_{\frac{1}{2}, -\frac{1}{2}}^1 = \frac{1}{\sqrt{6}}(udd + dud - 2ddu), \quad B_{\frac{1}{2}, -\frac{1}{2}}^2 = \frac{1}{\sqrt{2}}(ud - du)d,$$

$$M_{1,1}^1 = \bar{d}u, \quad M_{1,-1}^1 = -\bar{u}d,$$

$$M_{10}^1 = \frac{1}{\sqrt{2}}(\bar{d}\bar{d} - \bar{u}\bar{u}), \quad M_{00}^1 = -\frac{1}{\sqrt{2}}(\bar{u}u + \bar{d}d).$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{f1} = \sqrt{\frac{1}{6}}B_{\frac{3}{2}, -\frac{1}{2}}^1 M_{11} - \sqrt{\frac{1}{3}}B_{\frac{3}{2}, \frac{1}{2}}^1 M_{10} + \sqrt{\frac{1}{2}}B_{\frac{3}{2}, \frac{3}{2}}^1 M_{1-1}$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{f2} = \sqrt{\frac{1}{3}}B_{\frac{1}{2}, \frac{1}{2}}^1 M_{10} - \sqrt{\frac{2}{3}}B_{\frac{1}{2}, -\frac{1}{2}}^1 M_{11}$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{f3} = \sqrt{\frac{1}{3}}B_{\frac{1}{2}, \frac{1}{2}}^2 M_{10} - \sqrt{\frac{2}{3}}B_{\frac{1}{2}, -\frac{1}{2}}^2 M_{11}$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{f4} = B_{\frac{1}{2}, \frac{1}{2}}^1 M_{00}$$

$$\chi_{\frac{1}{2}, \frac{1}{2}}^{f5} = B_{\frac{1}{2}, \frac{1}{2}}^2 M_{00}$$

$$\chi_{\frac{3}{2}, \frac{3}{2}}^{f1} = \sqrt{\frac{3}{5}}B_{\frac{3}{2}, \frac{3}{2}}^1 M_{10} - \sqrt{\frac{2}{5}}B_{\frac{3}{2}, \frac{1}{2}}^1 M_{11}$$

$$\chi_{\frac{3}{2}, \frac{3}{2}}^{f2} = B_{\frac{3}{2}, \frac{3}{2}}^1 M_{00}$$

$$\chi_{\frac{3}{2}, \frac{3}{2}}^{f3} = B_{\frac{1}{2}, \frac{1}{2}}^1 M_{11}$$

$$\chi_{\frac{3}{2}, \frac{3}{2}}^{f4} = B_{\frac{1}{2}, \frac{1}{2}}^2 M_{11}$$

$$\chi_{\frac{5}{2}, \frac{5}{2}}^{f1} = B_{\frac{3}{2}, \frac{3}{2}}^1 M_{11}.$$

- Possible channels in negative parity.

		$I = \frac{1}{2}$		$I = \frac{3}{2}$		$I = \frac{5}{2}$	
J^P	index	$\chi_J^{\sigma_i}; \chi_I^{f_j}; \chi_k^c$ [$i; j; k$]	phys. channel	$\chi_J^{\sigma_i}; \chi_I^{f_j}; \chi_k^c$ [$i; j; k$]	phys. channel	$\chi_J^{\sigma_i}; \chi_I^{f_j}; \chi_k^c$ [$i; j; k$]	phys. channel
$\frac{1}{2}^-$	1	[4, 5; 4, 5; 1]	$P\eta$	[4, 5; 3, 4; 1]	$P\pi$	[1; 1; 1]	$\Delta\rho$
	2	[4, 5; 4, 5; 2, 3]		[4, 5; 3, 4; 2, 3]		[1; 1; 3]	
	3	[4, 5; 2, 3; 1]	$P\pi$	[2, 3; 3, 4; 1]	$P\rho$		
	4	[4, 5; 2, 3; 2, 3]		[2, 3; 3, 4; 2, 3]			
	5	[2, 3; 4, 5; 1]	$P\omega$	[1; 2; 1]	$\Delta\omega$		
	6	[2, 3; 4, 5; 2, 3]		[1; 2; 3]			
	7	[2, 3; 2, 3; 1]	$P\rho$	[1; 1; 1]	$\Delta\rho$		
	8	[2, 3; 2, 3; 2, 3]		[1; 1; 3]			
	9	[1; 1; 1]	$\Delta\rho$				
	10	[1; 1; 3]					
$\frac{3}{2}^-$	1	[3, 4; 4, 5; 1]	$P\omega$	[3, 4; 3, 4; 1]	$P\rho$	[2; 1; 1]	$\Delta\pi$
	2	[3, 4; 4, 5; 2, 3]		[3, 4; 3, 4; 2, 3]		[2; 1; 3]	
	3	[3, 4; 2, 3; 1]	$P\rho$	[2; 2; 1]	$\Delta\eta$	[1; 1; 1]	$\Delta\rho$
	4	[3, 4; 2, 3; 2, 3]		[2; 2; 3]		[1; 1; 3]	
	5	[2; 1; 1]	$\Delta\pi$	[2; 1; 1]	$\Delta\pi$		
	6	[2; 1; 3]		[2; 1; 3]			
	7	[1; 1; 1]	$\Delta\rho$	[1; 2; 1]	$\Delta\omega$		
	8	[1; 1; 3]		[1; 2; 3]			
	9			[1; 1; 1]	$\Delta\rho$		
	10			[1; 1; 3]			
$\frac{5}{2}^-$	1	[1; 1; 1]	$\Delta\rho$	[1; 2; 1]	$\Delta\omega$	[1; 1; 1]	$\Delta\rho$
	2	[1; 1; 3]		[1; 2; 3]		[1; 1; 3]	
	3			[1; 1; 1]	$\Delta\rho$		
	4			[1; 1; 3]			

- Dynamic study in ChQM for the light quark sector(u,d) with three sets of parameters.

TABLE II: Masses of baryon and meson in ChQM with the three sets of parameters (unit: MeV).

set	$N(939)$	$\Delta(1232)$	$\Omega(1672)$	$\Lambda(1116)$
I	939	1231	1671	1178
II	749	1116	1623	1033
III	833	1178	1704	1124
	$\Sigma(1189)$	$\Xi(1315)$	$\Sigma^*(1383)$	$\Xi^*(1532)$
I	1201	1408	1381	1527
II	1058	1260	1341	1428
III	1147	1406	1350	1525
	$\pi(140)$	$\rho(775)$	$\eta(548)$	$\omega(782)$
I	305	817	646	693
II	139	742	551	573
III	145	783	561	653
	$K(495)$	$K^*(892)$	$\eta'(958)$	$\phi(1019)$
I	517	968	771	1070
II	483	886	958	1027
III	483	955	857	1093

*qqq
systems*

**Total mass error for
baryons and mesons in 3
sets of parameters.**

set	ΔE_b	ΔE_m
I	176	632
II	770	272
III	373	405

TABLE V: The lowest eigen-energies of system with $IJ^P = \frac{1}{2} \frac{3}{2}^-$ in three sets of parameters (unit: MeV). The percentages of color-singlet (S) and hidden-color (H) channels are also given.

Channel	E_{th}^{Exp}	set I				set II				set III			
		E	E_{th}^{Theo}	E_B	E'	E	E_{th}^{Theo}	E_B	E'	E	E_{th}^{Theo}	E_B	E'
1	1722($P\omega$)	1632	1632	0	1722	1322	1322	0	1722	1486	1486	0	1722
2		2063				1878				1964			
1+2		1632		0	1722	1322		0	1722	1486		0	1722
percentage (S; H):			100%; 0.0%				100%; 0.0%				100%; 0.0%		
3	1714($P\rho$)	1756	1756	0	1714	1491	1491	0	1714	1616	1616	0	1714
4		2143				1991				2049			
3+4		1756		0	1714	1491		0	1714	1616		0	1714
percentage (S; H):			100%; 0.0%				100%; 0.0%				100%; 0.0%		
5	1371($\Delta\pi$)	1533	1536	-3	1368	1252	1255	-3	1368	1321	1323	-2	1369
6		2556				2568				2478			
5+6		1533		-3	1368	1252		-3	1368	1320		-3	1368
percentage (S; H):			100%; 0.0%				100%; 0.0%				100%; 0.0%		
7	2007($\Delta\rho$)	2048	2048	0	2007	1858	1858	0	2007	1961	1961	0	2007
8		2516				2517				2431			
7+8		2048		0	2007	1858		0	2007	1961		0	2007
percentage (S; H):			100%; 0.0%				100%; 0.0%				100%; 0.0%		
mixed (singlet)	1371($\Delta\pi$)	1514	1536	-22	1349	1223	1255	-32	1339	1301	1323	-22	1349
mixed (full)	1371($\Delta\pi$)	1497	1536	-39	1332	1201	1255	-54	1317	1283	1323	-40	1331

TABLE VII: The percentage of each component within the full channel-coupling of the lowest energy for $IJ^P = \frac{1}{2}\frac{3}{2}^-$ state in three set of parameters.

	$(P\omega)^1$	$(P\omega)^8$	$(P\rho)^1$	$(P\rho)^8$
<i>components</i> (set I)	27.1%	1.5%	11.1%	0.7%
	$(\Delta\pi)^1$	$(\Delta\pi)^8$	$(\Delta\rho)^1$	$(\Delta\rho)^8$
	<u>51.7%</u>	1.5%	4.7%	1.7%
<i>components</i> (set II)	$(P\omega)^1$	$(P\omega)^8$	$(P\rho)^1$	$(P\rho)^8$
	26.9%	0.7%	8.4%	0.8%
	$(\Delta\pi)^1$	$(\Delta\pi)^8$	$(\Delta\rho)^1$	$(\Delta\rho)^8$
	<u>57.6%</u>	1.9%	1.5%	2.2%
<i>components</i> (set III)	$(P\omega)^1$	$(P\omega)^8$	$(P\rho)^1$	$(P\rho)^8$
	23.1%	1.1%	9.5%	0.2%
	$(\Delta\pi)^1$	$(\Delta\pi)^8$	$(\Delta\rho)^1$	$(\Delta\rho)^8$
	<u>59.3%</u>	1.5%	3.7%	1.6%

TABLE VI: The lowest eigen-energies of system with $IJ^P = \frac{1}{2} \frac{5}{2}^-$ (unit: MeV) in three sets of parameters. The percentages of color-singlet (S) and hidden-color (H) channels are also given.

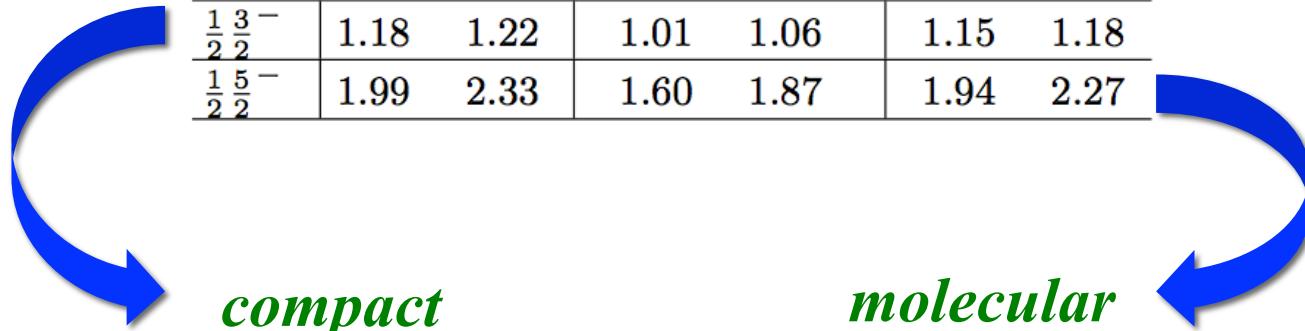
Channel	E	E_{th}^{Theo}	E_B	E_{th}^{Exp}	E'
<i>set I</i>					
1	2043	2048	-5	2007($\Delta\rho$)	2002
2	2633				
1+2	2042		-6		2001
percentage (S; H): 99.3%; 0.7%					
<i>set II</i>					
1	1849	1858	-9	2007($\Delta\rho$)	1998
2	2674				
1+2	1847		-11		1996
percentage (S; H): 99.3%; 0.7%					
<i>set III</i>					
1	1954	1961	-7	2007($\Delta\rho$)	2000
2	2563				
1+2	1953		-8		1999
percentage (S; H): 99.3%; 0.7%					

TABLE VIII: Distances between quarks and quark-antiquark in three sets of parameters (unit: fm).

	set I		set II		set III	
IJ^P	r_{qq}	$r_{q\bar{q}}$	r_{qq}	$r_{q\bar{q}}$	r_{qq}	$r_{q\bar{q}}$
$\frac{1}{2}\frac{3}{2}^-$	1.18	1.22	1.01	1.06	1.15	1.18
$\frac{1}{2}\frac{5}{2}^-$	1.99	2.33	1.60	1.87	1.94	2.27

compact

molecular



- May be a model-independent results for the nature of pentaquark states of $\Delta\pi$, $\Delta\rho$ with the spin-parity of 3/2- and 5/2-, respectively, and $I=1/2$.

$$\frac{1}{2}\left(\frac{3}{2}\right)^{-} : N^{*}(1331)$$

$$\frac{1}{2}\left(\frac{5}{2}\right)^{-} : N^{*}(2000)$$

Conclusions and prospects

- The pentaquark state P_c^+ (4380) can be interpreted as the molecular baryons of $\Sigma_c^* D$.
- One of the newly discovered Ω_c^0 excited states may be the pentaquark state of $\Xi_c^* K$.
- Two higher excited states of Ω_c^0 can be explained as the 2S states.
- The high spin N* states may exist the tightly bound states of pentaquark.
- *The coupling between 3q and 5q in unquenched quark model should be developed in future.*
- *The spin-orbit, tensor force may be considered in future.*

Thanks!