

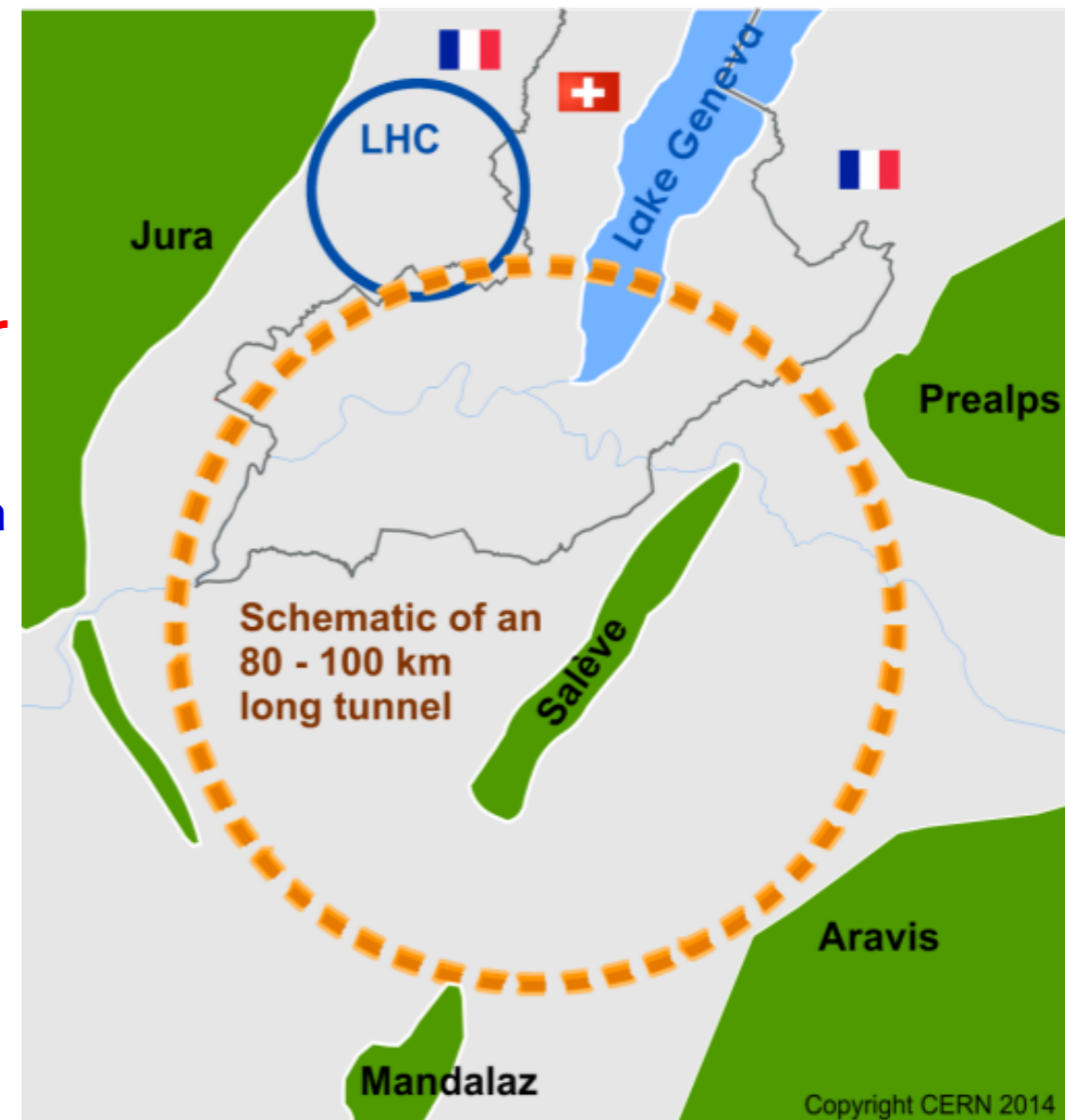
Beam Optics & Dynamic Aperture of FCC-ee

Nov. 1, 2017 @ DA Workshop

K. Oide (KEK) on behalf of the FCC-ee Collaboration

The FCC Design Study

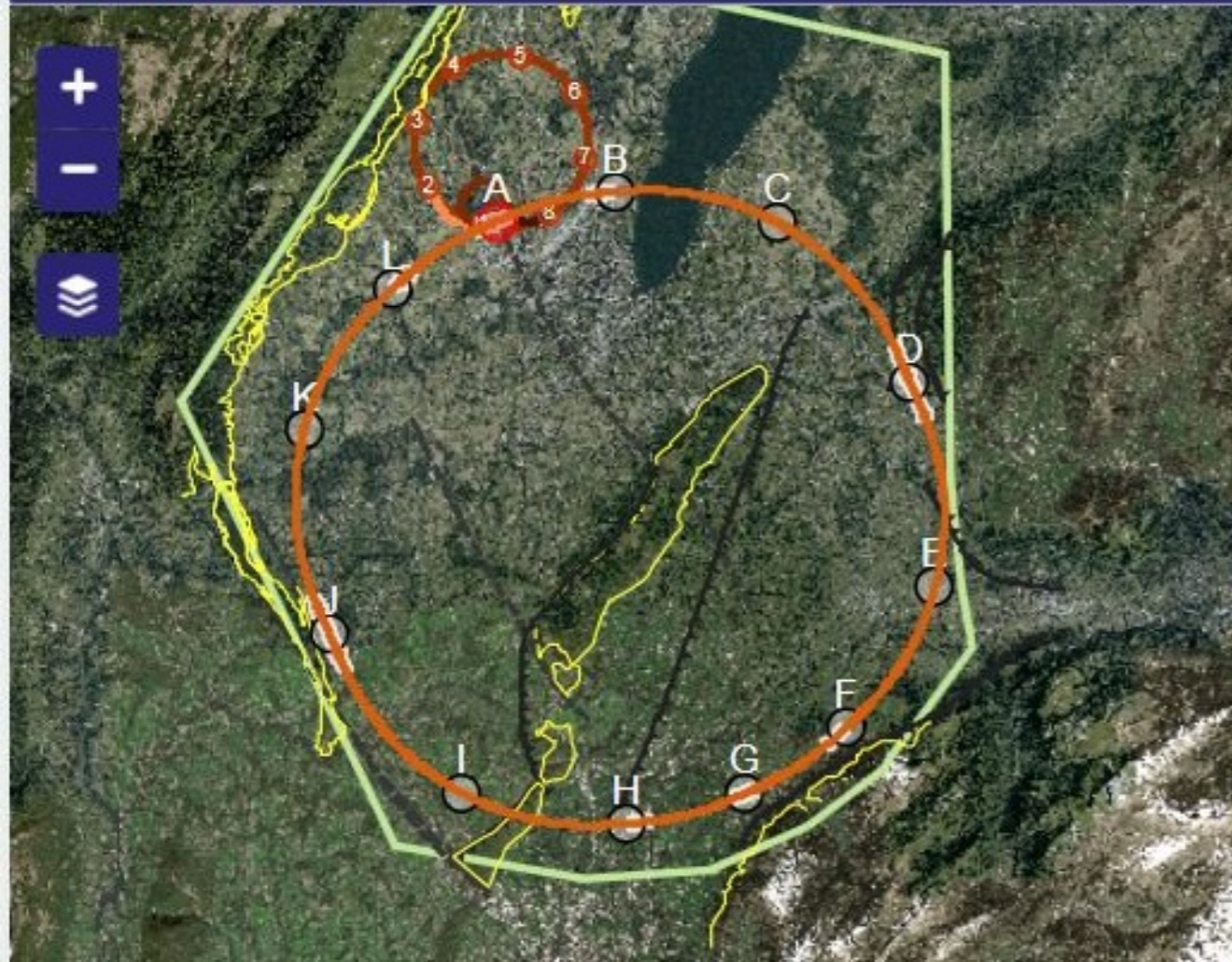
- **Requested from European Strategy (2013)**
 - ◆ “Ambitious post-LHC accelerator project”
 - Study kicked off in Geneva in Feb. 2014
- **International collaboration to study circular colliders (111 institutes)**
 - ◆ Fitting in a new 100 km infrastructure, in the Geneva area
- **Ultimate goal: 100 TeV pp collider (FCC-hh)**
 - ◆ Requires R&D for 16T magnets
 - ◆ Defines the infrastructure
- **Possible first steps**
 - ◆ e^+e^- collider (FCC-ee) at the intensity frontier
 - High luminosity, $\sqrt{s} = 90\text{-}400$ GeV
 - ◆ pp collider (HE-LHC) in the LEP/LHC tunnel
 - With FCC-hh technology (16T \rightarrow 28 TeV)
- **Possible add-on**
 - ◆ e-p option (FCC-eh)



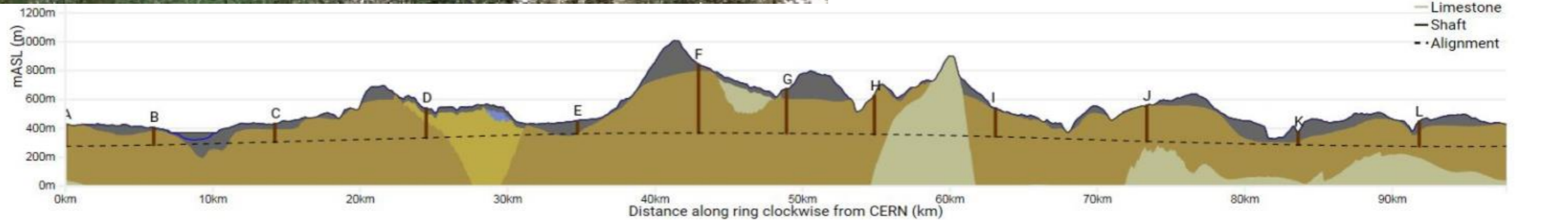
- **European Strategy update (2019)**
 - ◆ Conceptual design report (CDR)
 - ◆ Cost review for tunnel and each collider
 - ◆ Schedules and operation models

The FCC Home

Alignment Location



- **Optimized length: 97.5 km**
 - ◆ Accessibility, rock type, shaft depth, etc.
 - ◆ Tried different options from 80 to 100 km
- **Tunneling**
 - ◆ Molasse 90% (easy to dig)
 - ◆ Limestone 5%, Moraines 5% (tougher)
- **Shallow implementation**
 - ◆ 30m below Lemman lakebed
 - ◆ Only one very deep shaft (F, 476m)
 - Alternatives studied (e.g. inclined access)



- Quaternary
- Lake
- Wildflysch
- Molasse subalpine
- Molasse
- Limestone
- Shaft
- - Alignment

Geology Intersected by Tunnel

Geology Intersected by Section

84.6%

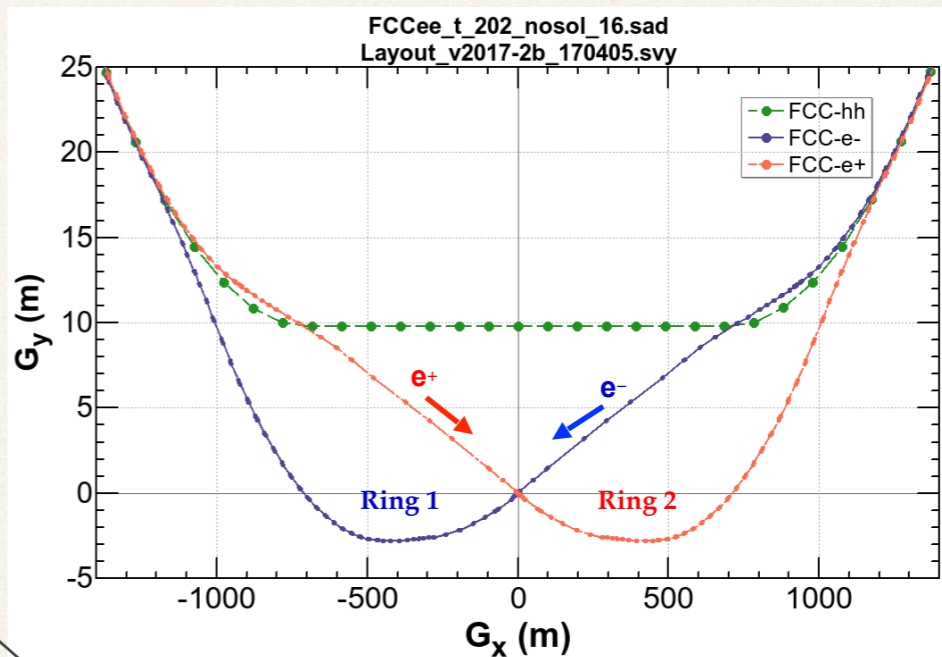
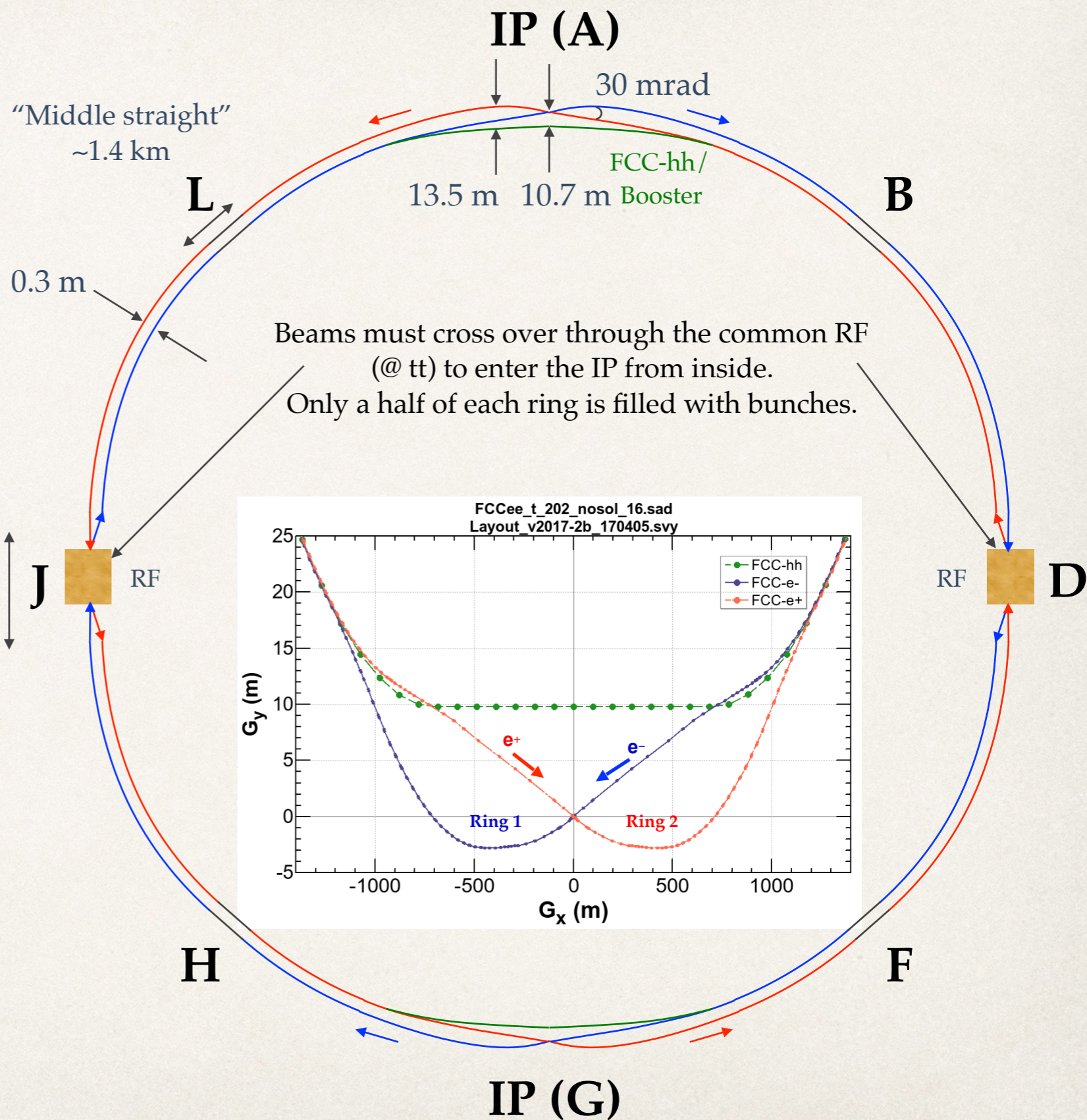
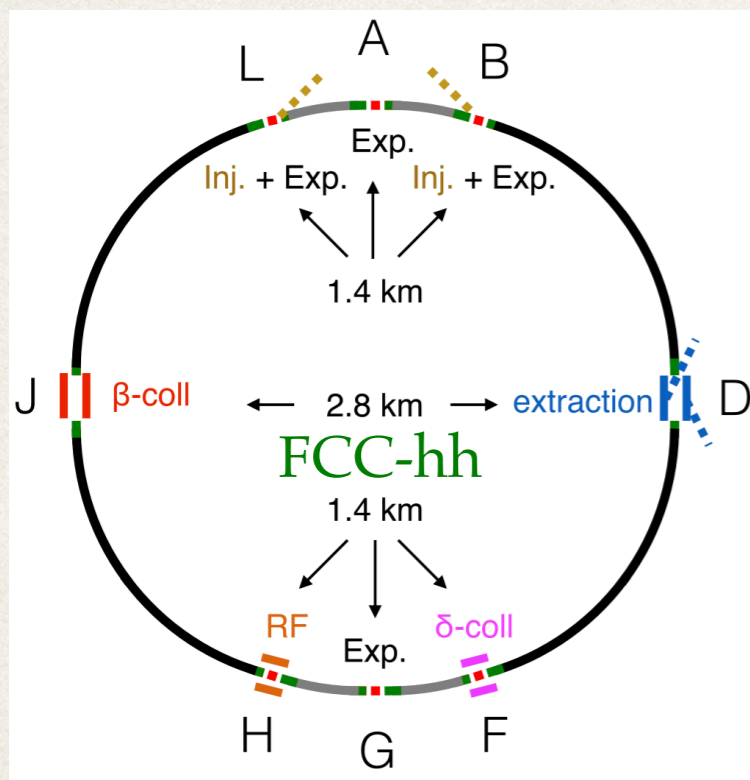
5.2%

5.5%

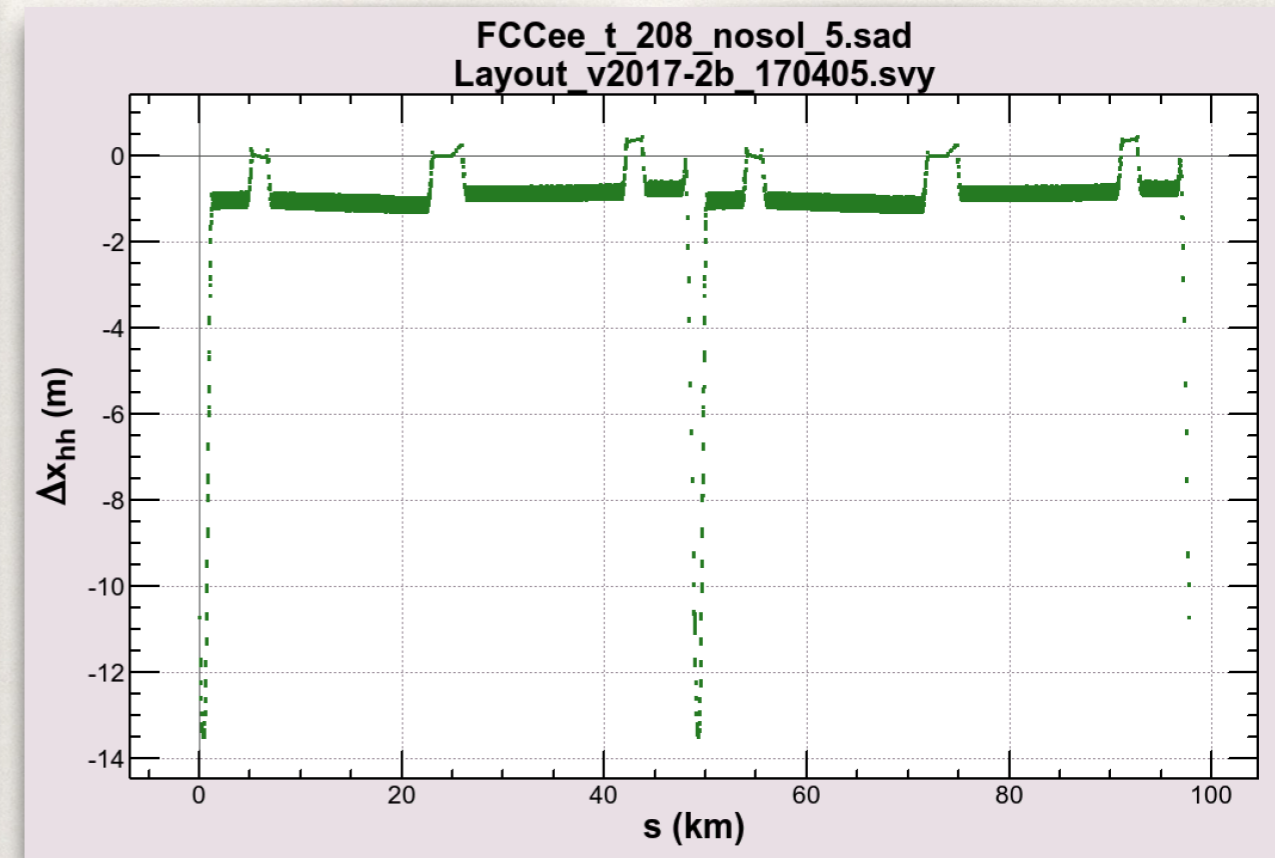
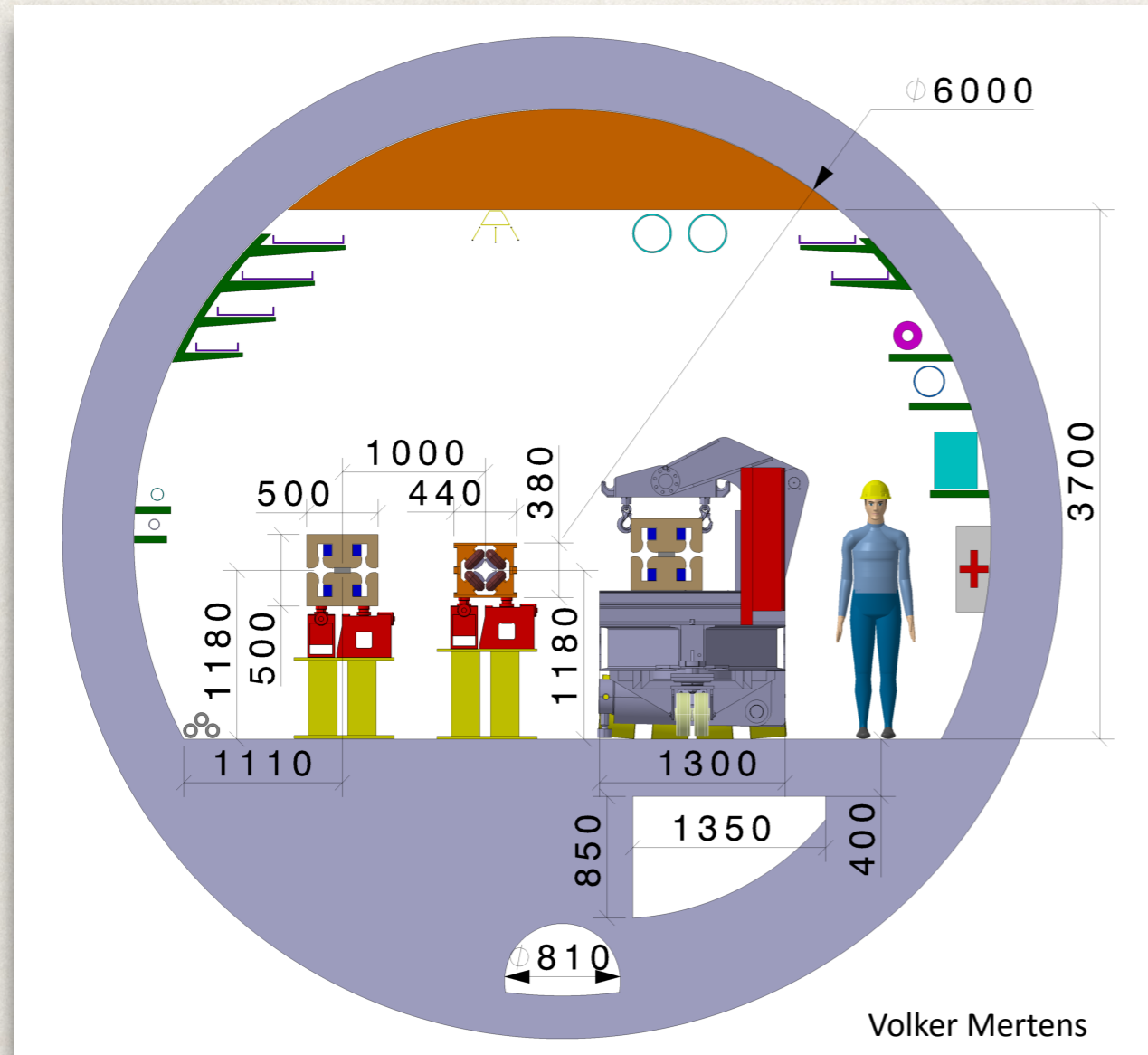
4.7%

- ❖ A baseline optics* for FCC-ee was once established in Oct. 2016 characterized by:
 - ❖ 100 km circumference, 2 IP / ring
 - ❖ Common lattice for all energies
 - ❖ 90°/90° FODO cell in the arc with non-interleaved sextupole pairs (294 families)
 - ❖ 30 mrad crossing angle at the IP, with the crab-waist scheme
 - ❖ Local chromaticity correction for y -plane, incorporated with crab sextupoles
 - ❖ 100 MW total SR power for all energies
 - ❖ Limit the SR toward the IP below 100 keV at 175 GeV, up to 450 m upstream
 - ❖ Tapering of magnets along the ring to compensate the effects of SR on orbit/optics
 - ❖ Sufficient dynamic aperture for beamstrahlung and top-up injection
- ❖ Motivations for changes in 2017:
 - ❖ Mitigation of the coherent beam-beam instability at Z
 - ❖ *Smaller β_x^* at Z, W, Zh.*
 - ❖ *60°/60° cell in the arc at Z*
 - ❖ *Longer bunch lengths by SR and beamstrahlung.*
 - ❖ The baseline beam energy for ttbar is set at 182.5 GeV.
 - ❖ Adopt the “Twin Aperture Quadrupole” scheme for arc quadrupoles
 - ❖ Fit the footprint to a new FCC-hh layout

Layout of FCC-ee



Layout in the tunnel

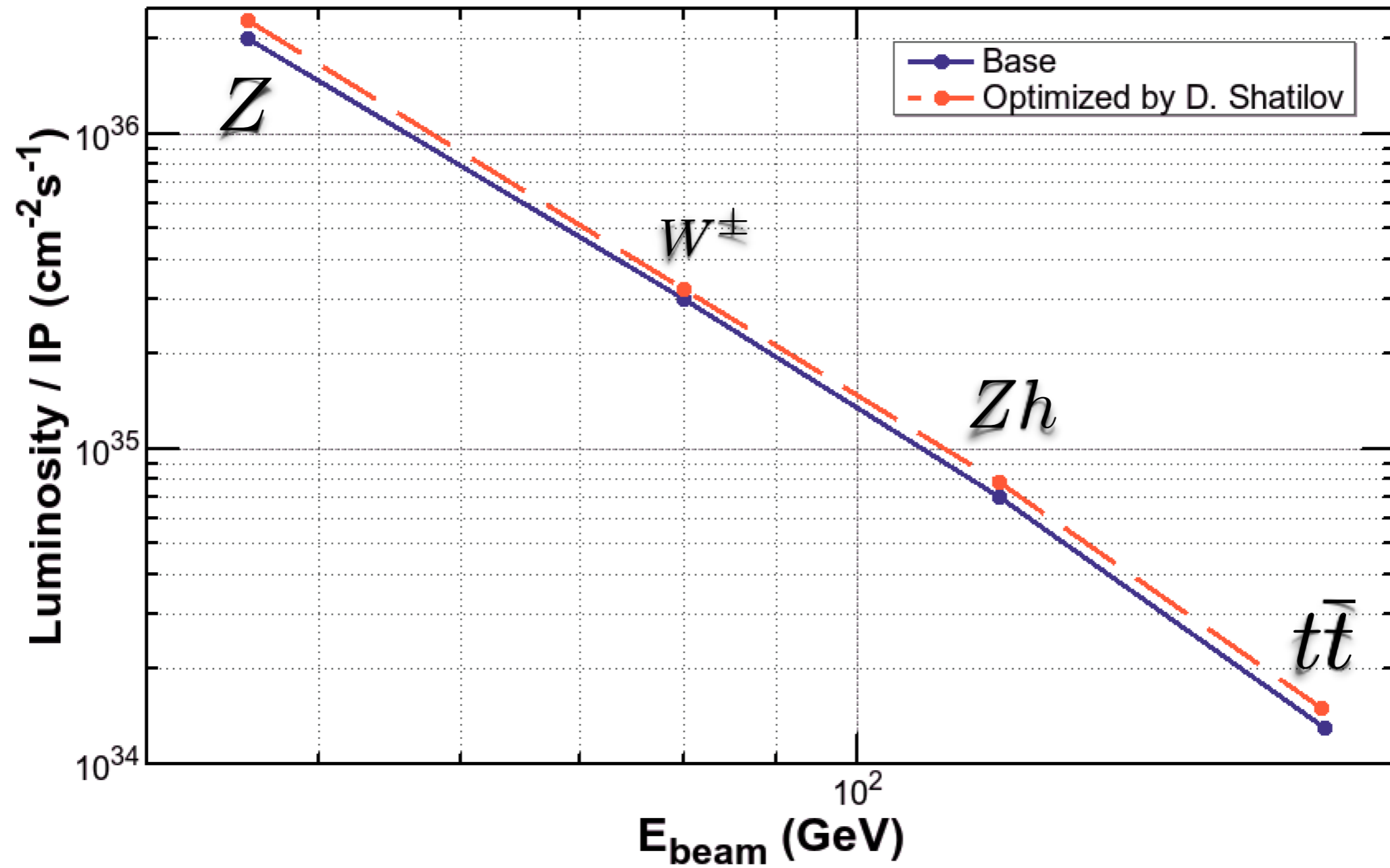


- ❖ A new layout of the FCC-ee collider and booster has been proposed by V. Mertens as above.
- ❖ The main IP sits right on the FCC-hh footprint.
- ❖ The collider comes outside of the booster/FCC-hh by 1 m.
- ❖ The collider optics has been modified to fit the new layout.
- ❖ The shift of the ee's IP from the hh's IP is 10.7 m.
- ❖ The layouts in the straight sections are subject to change.

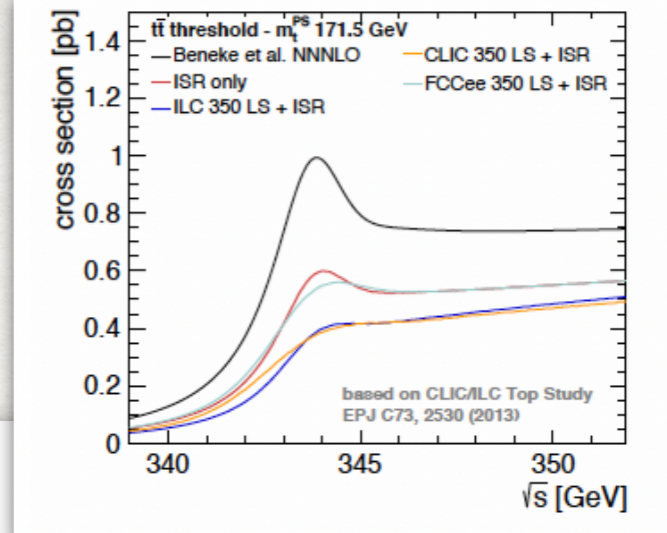
Parameters 2017



parameter	Z	W	H (ZH)	ttbar
beam energy [GeV]	45.6	80	120	182.5
Circumference [km]	97.756			
arc cell optics	60/60	90/90		
momentum compaction [10^{-5}]	1.48	0.73		
horizontal emittance [nm]	0.27	0.28	0.63	1.43
vertical emittance [pm]	1.0	1.0	1.3	2.9
horizontal beta* [m]	0.15	0.2	0.3	1
vertical beta* [mm]	0.8	1	1	2
length of interaction area [mm]	0.42	0.5	0.9	2.0
tunes, half-ring (x, y)	(0.569, 0.60)	(0.577, 0.61)	(0.565, 0.60)	(0.552, 0.588)
synchrotron tune (half ring)	0.0124	0.0114	0.0180	0.0340
long. damping time [ms]	414	77	23	6.6
SR energy loss / turn [GeV]	0.0360	0.334	1.69	9.05
total RF voltage [GV]	0.096	0.43	1.96	10.6
RF acceptance [%]	1.9	1.9	2.3	4.7
energy acceptance [%]	± 1.3	± 1.3	± 1.5	-2.8/+2.4
energy spread (SR / BS) [%]	0.038 / 0.132	0.066 / 0.153	0.099 / 0.151	0.153 / 0.195
bunch length (SR / BS) [mm]	3.5 / 12.1	3.3 / 7.65	3.15 / 4.9	2.5 / 3.3
Piwnski angle (SR / BS)	8.2 / 28.5	6.6 / 15.3	3.4 / 5.3	1.0 / 1.3
bunch intensity [10^{11}]	1.7	1.5	1.5	2.8
number of bunches / beam	16640	2030	401	40
beam current [mA]	1390	150	29.6	5.5
luminosity [$10^{34} \text{ cm}^{-2}\text{s}^{-1}$]	> 200	> 30	> 7	> 1.3
beam-beam parameter (x / y)	0.004 / 0.133	0.0065 / 0.118	0.016 / 0.108	0.090 / 0.148
luminosity lifetime [min]	70	50	42	39
time between injections [sec]	122	44	31	32
allowable asymmetry [%]	± 5	± 3	± 3	± 3
required lifetime by BS [min]	29	16	11	12
actual lifetime (w) by BS [min]	> 200	20	20	25

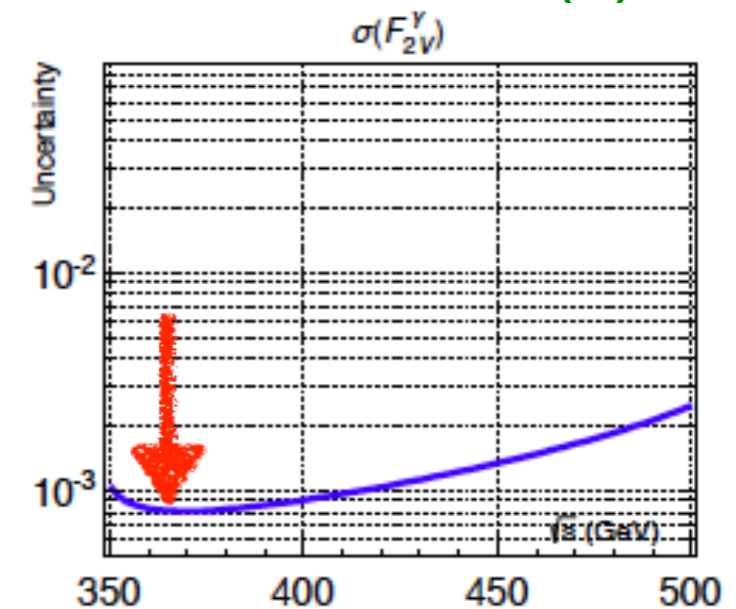
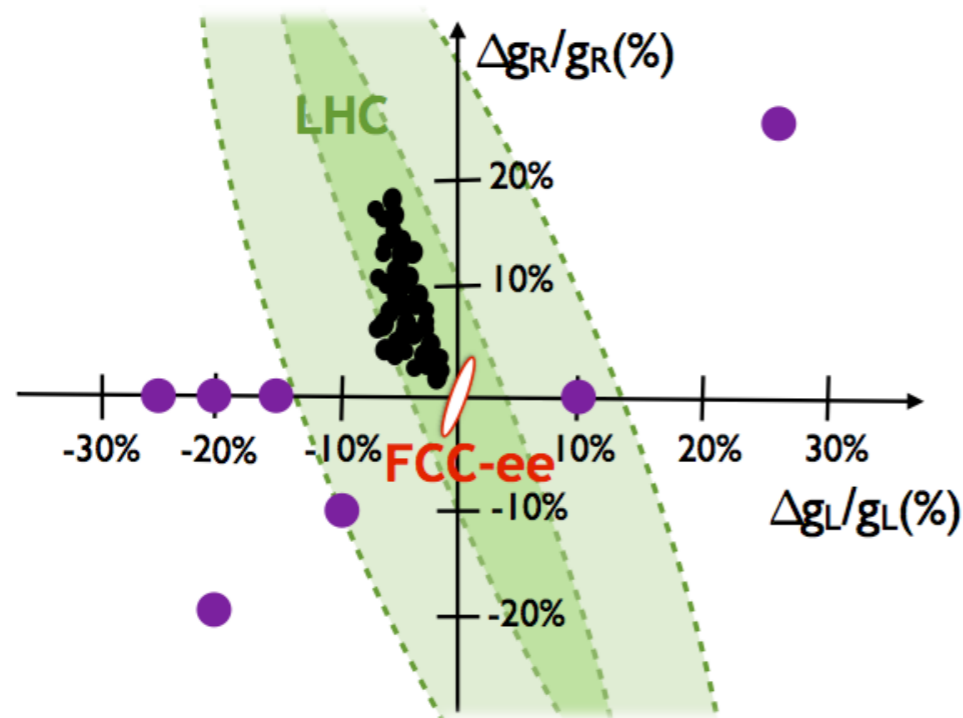


Why 365 GeV CM for $t\bar{t}$?



ELECTROWEAK COUPLINGS OF THE TOP QUARK(2)

- Large statistics and final state polarization allow a full separation of the $t\bar{t}Z/\gamma$ couplings with **NO need for polarization in the initial state.**
- Optimal $\sqrt{s} = 365\text{-}370$ GeV



- Fit includes conservative assumptions detector performance
- Theory uncertainty on production mechanism dominates

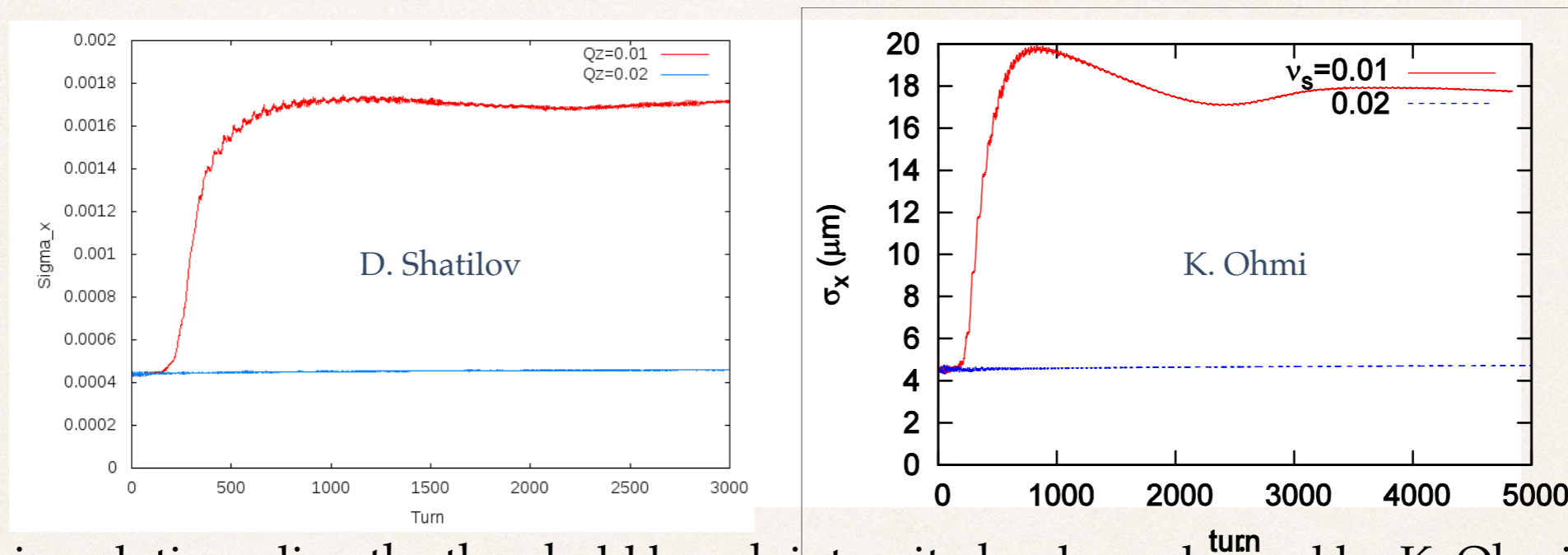
FCC-ee expected precision of order 10^{-2} to 10^{-3}

conservative revised operation model

working point	luminosity/IP [$10^{34} \text{ cm}^{-2}\text{s}^{-1}$]	total luminosity (2 IPs)/ year	physics goal	run time [years]
Z first 2 years	100	26 $\text{ab}^{-1}/\text{year}$	150 ab^{-1}	4
Z later	200	52 $\text{ab}^{-1}/\text{year}$		
<i>W</i>	30	7.8 $\text{ab}^{-1}/\text{year}$	10 ab^{-1}	~ 1
<i>H</i>	7.0	1.8 $\text{ab}^{-1}/\text{year}$	5 ab^{-1}	3
top	1.3	0.34 $\text{ab}^{-1}/\text{year}$	1.5 ab^{-1}	5

total run time: 12-13 years

- ❖ A new coherent instability in x-z plane was first found by K. Ohmi by FCC Week 2016 with a strong-strong beam-beam simulation.
- ❖ D. Shatilov confirmed their phenomenon by a completely independent simulation, with a turn-by-turn alternating quasi-strong-strong method. The results of these two agree to each other very well.



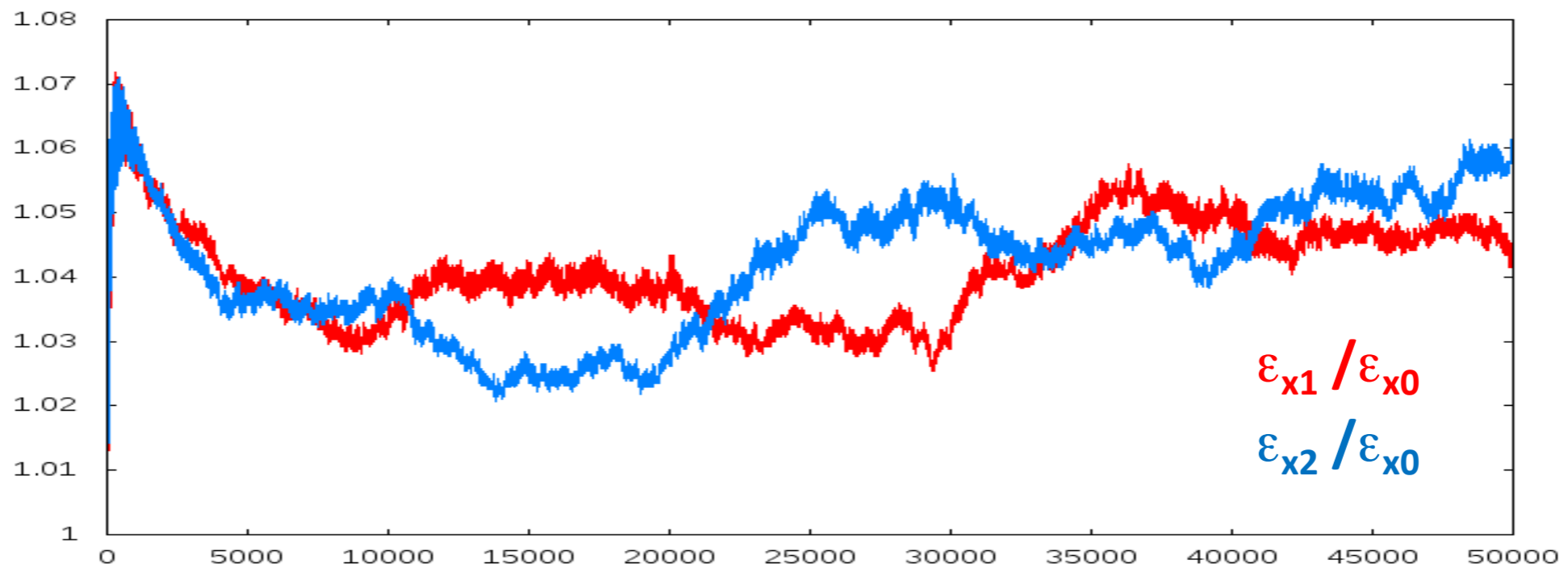
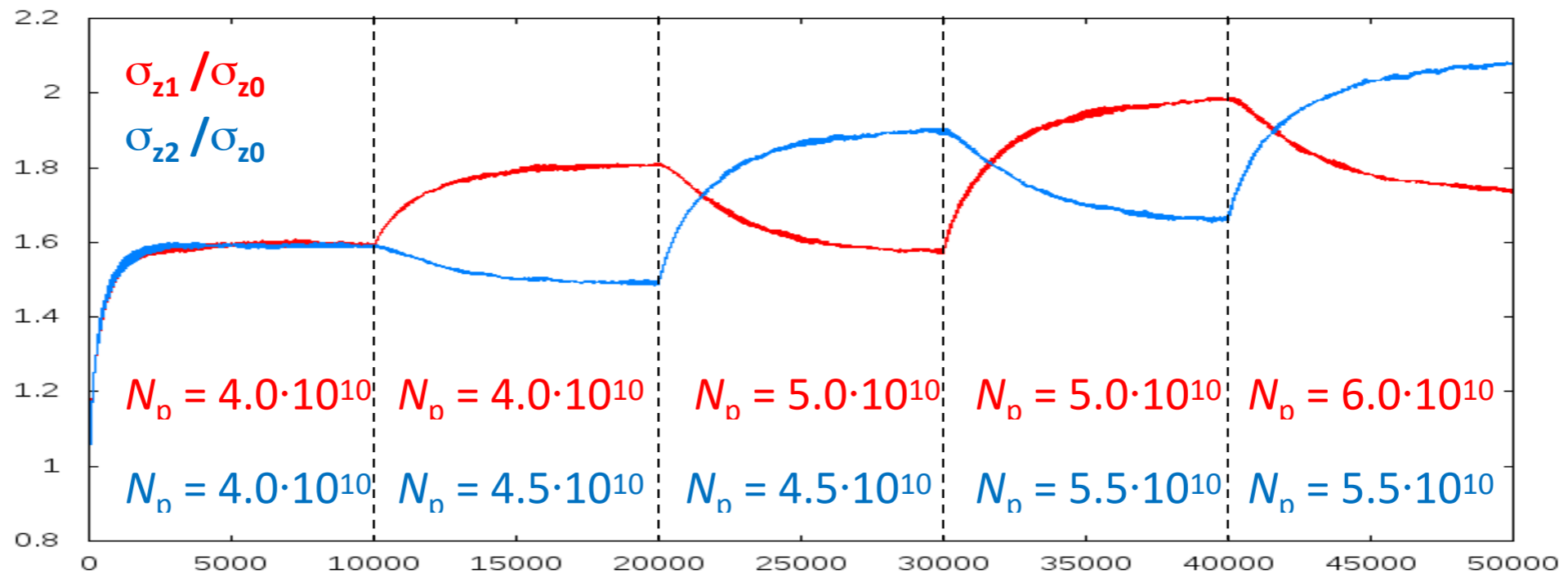
- ❖ A semi-analytic scaling the threshold bunch intensity has been derived by K. Ohmi:

$$N_{\text{th}} \propto \frac{\alpha_p \sigma_\delta \sigma_z}{\beta_x^*}$$

- ❖ Thus a smaller β_x^* and a larger momentum compaction α_p are favorable. The latter can be achieved by changing the phase advances to $60^\circ / 60^\circ$ of the arc at Z.
- ❖ Longer bunch lengths by SR as well as beamstrahlung stabilizes the instability.
- ❖ This instability appears up to Zh (120 GeV). A longer bunch length by a strong beamstrahlung can mitigate it, then an imbalance in the charges of the two colliding bunches may cause flip-flop: a “bootstrap” injection is necessary (D. Shatilov).

Bootstrapping

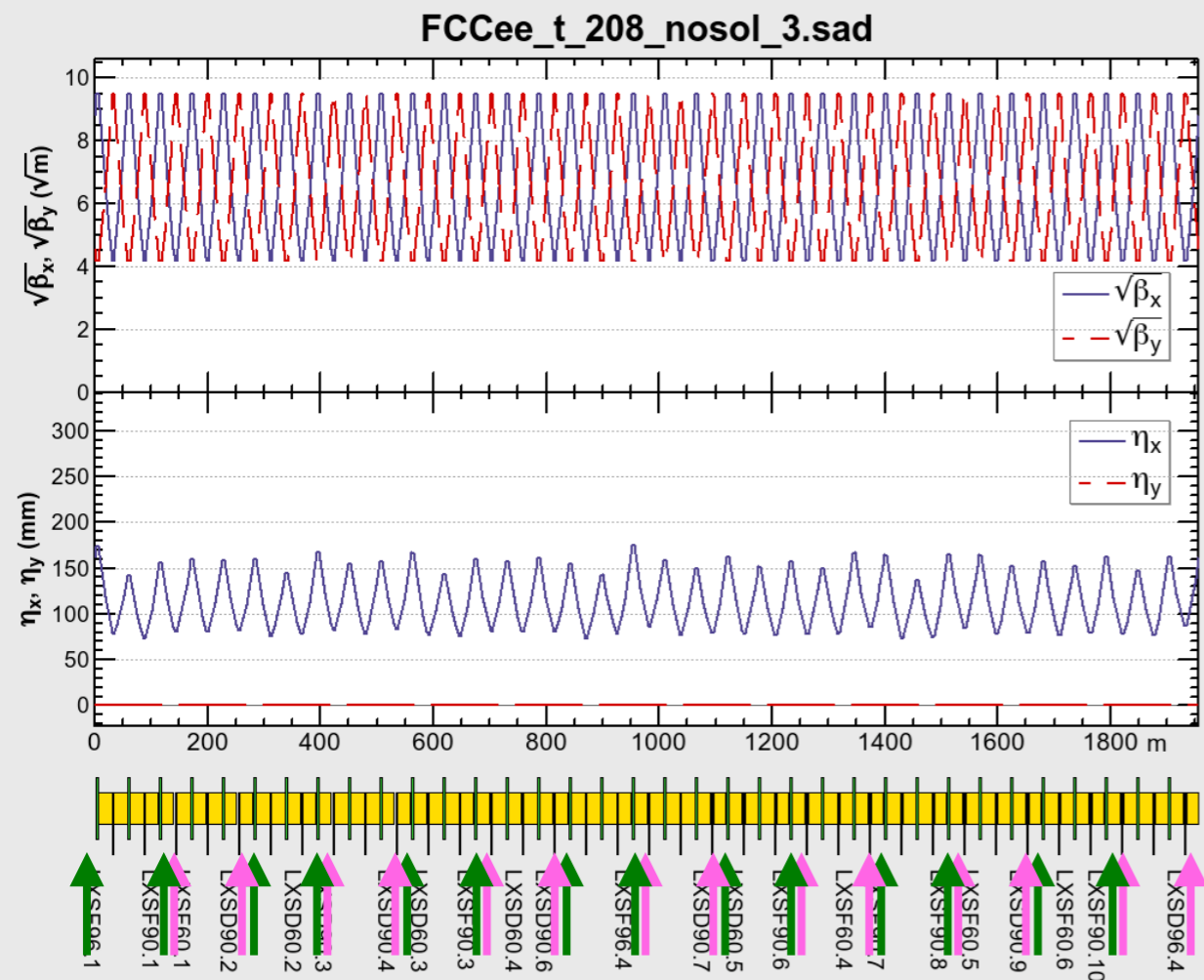
D. Shatilov



- The maximum bunch charge is determined considering the balance of beamstrahlung, momentum acceptance, and the capability of injector.

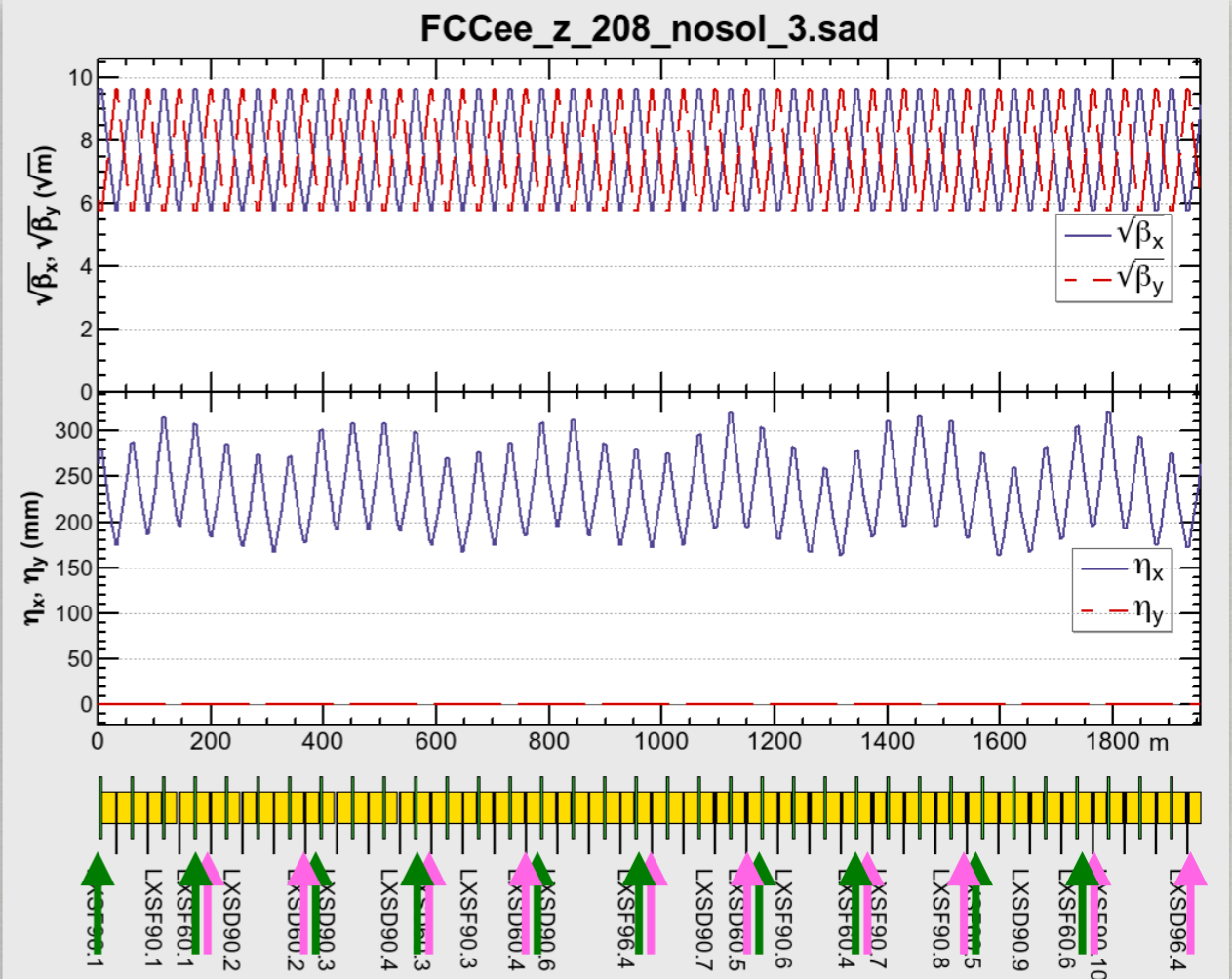
The arc cells

90°/90° (tt, Zh, WW)



14 sext pairs / 28 FODOs

60°/60° (Z)



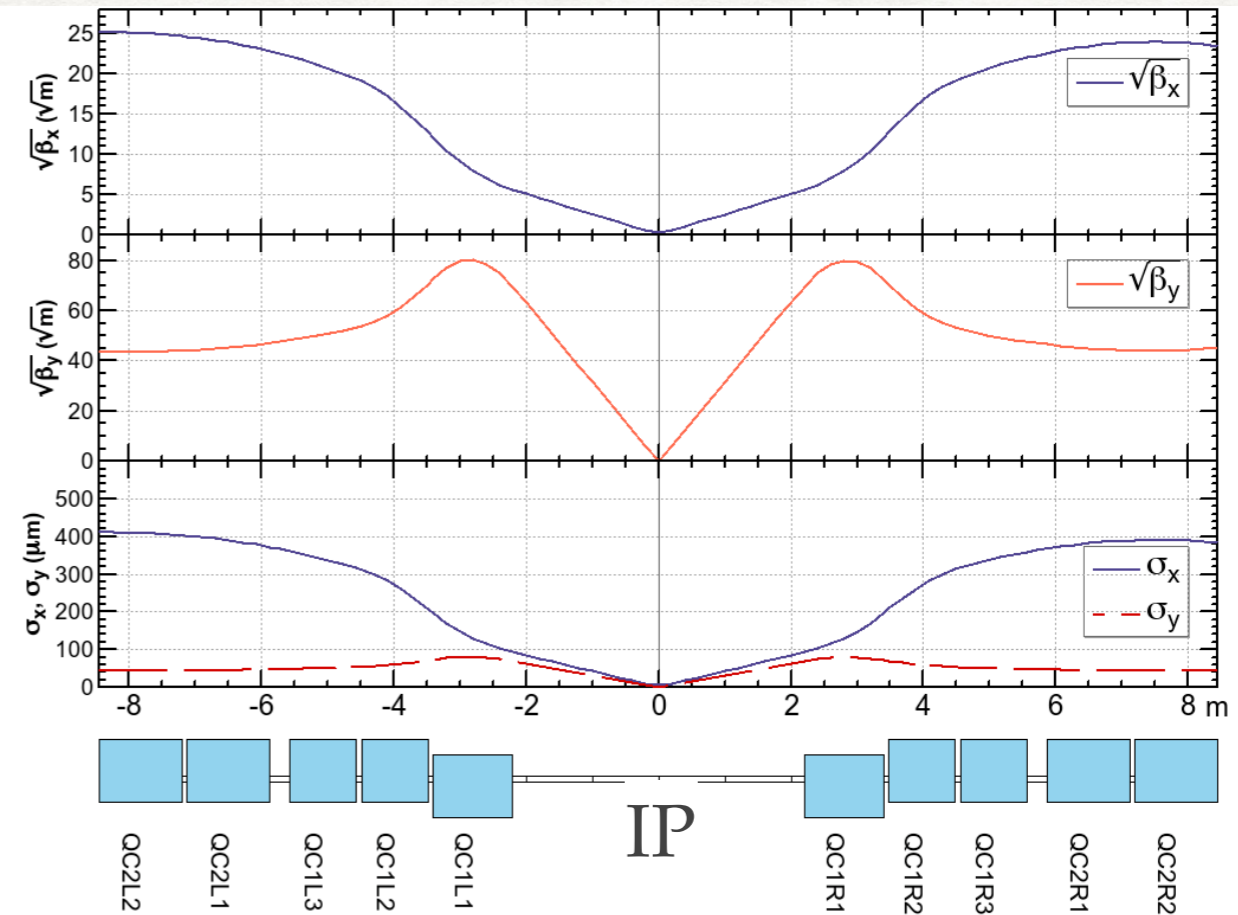
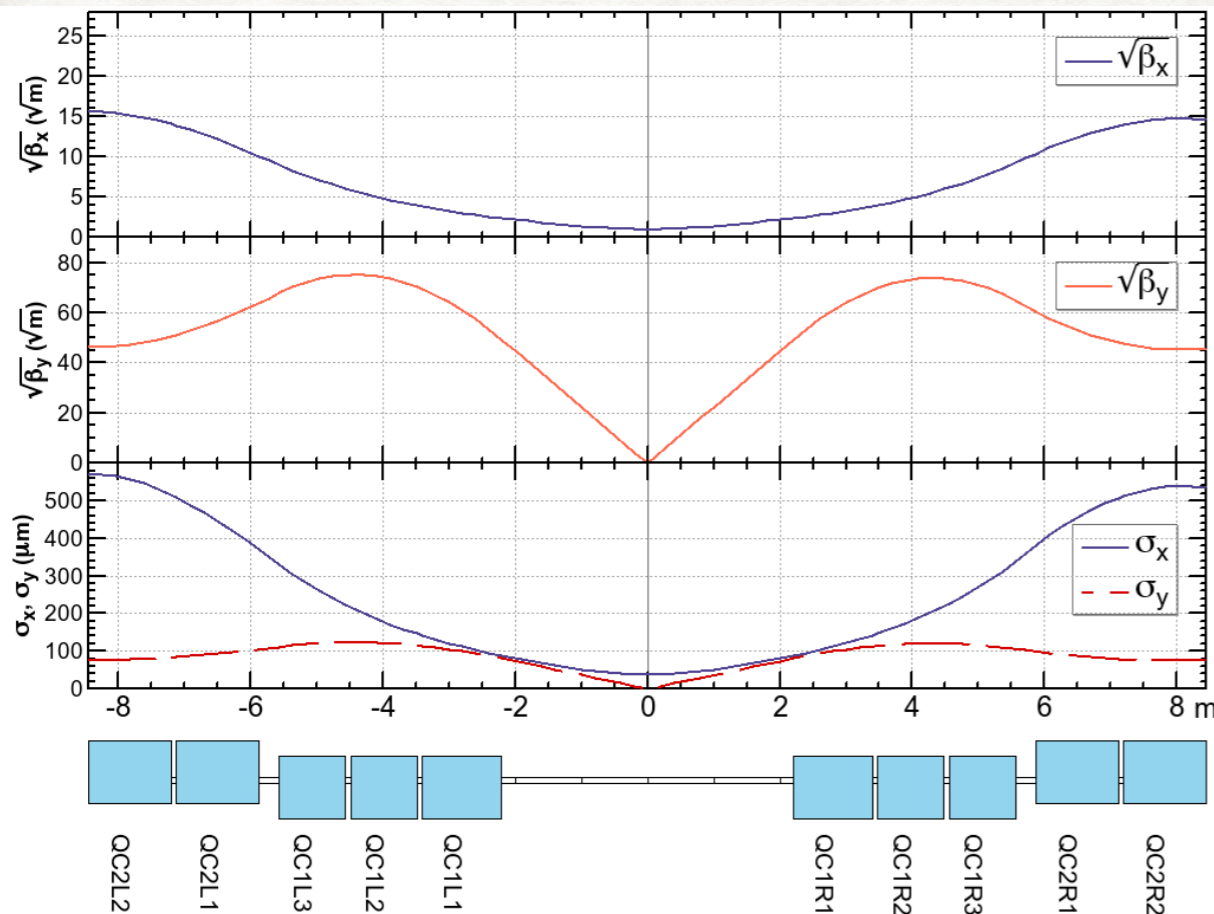
10 sext pairs / 28 FODOs

- The length of dipoles are adjusted to fill spaces, with three lengths of dipoles (21.84 m, 23.54 m, 24.44 m) with equal bending radius.
- While the β -functions are more or less uniform, there appear some wiggling in the horizontal dispersion.
- Sextupoles are made longer, from 0.7 m to 0.8 m (Z), and from 2.1 m to 2.4 m (tt).
- The resulting packing factor of the dipoles in the arc has been improved from 81.7% to 83.2%.

Reduction of β_x^* at Z, W, Zh

$$\beta_{x,y}^* = (1 \text{ m}, 2 \text{ mm}) @ \text{tt}$$

$$\beta_{x,y}^* = (15 \text{ cm}, 1 \text{ mm}) @ \text{Z}$$

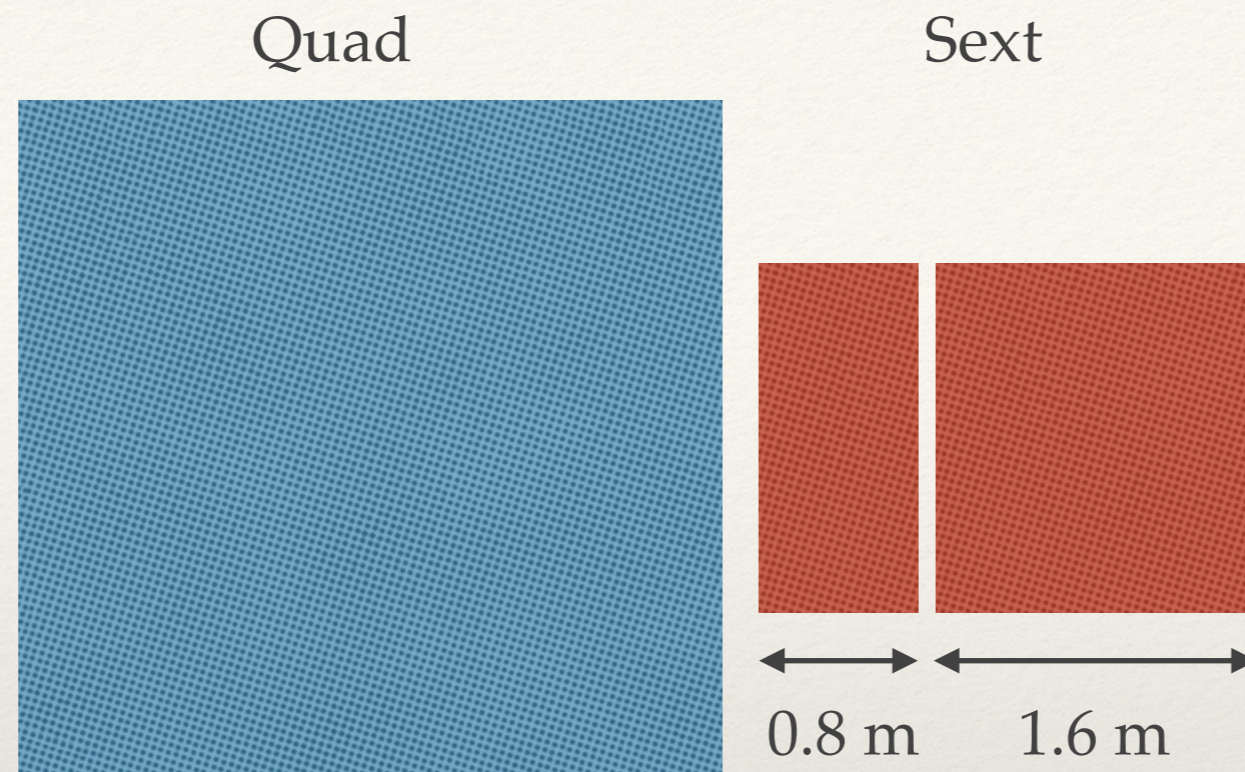


- ❖ Divide QC1 into three independent pieces, reverse the polarity at Z.

	L (m)	B' @ tt (T/	B' @ Z (T/
QC1L1	1.2	-94.4	-96.3
QC1L2	1	-92.6	+50.3
QC1L3	1	-96.7	+9.8
QC2L1	1.25	+45.8	+6.7
QC2L2	1.25	+74.0	+3.2

	L (m)	B' @ tt (T/	B' @ Z (T/
QC1R1	1.2	-99.9	-97.2
QC1R2	1	-99.9	+51.2
QC1R3	1	-99.9	+12.0
QC2R1	1.25	+78.6	+7.3
QC2R2	1.25	+76.2	+7.2

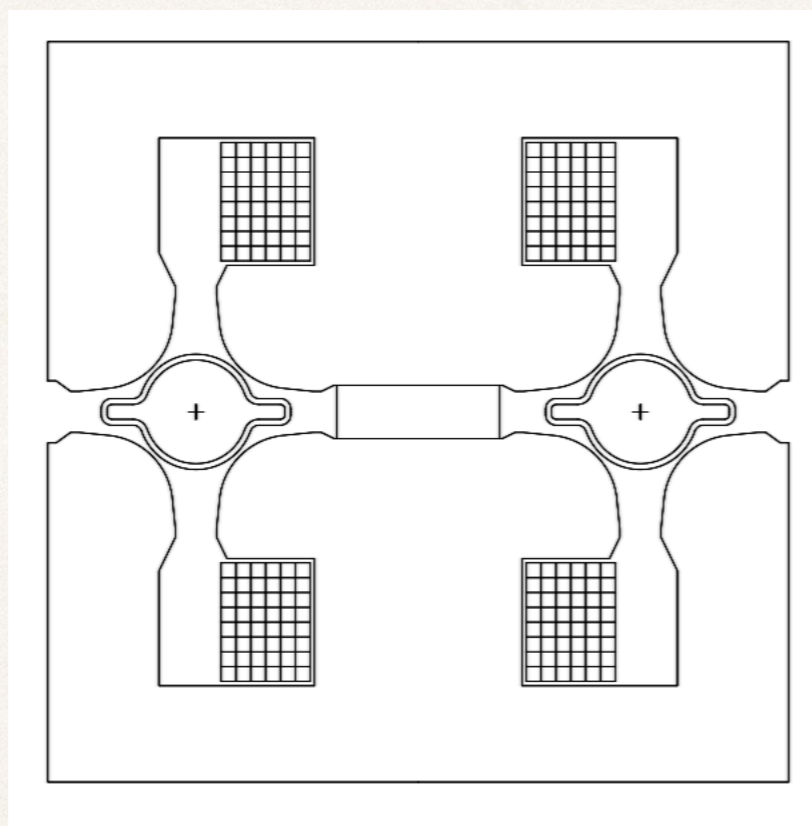
- ❖ By this split the chromaticity and the peaks of $\beta_{x,y}$ around the IP are suppressed for the reduction of $\beta_{x,y}^*$ at Z to (1/7, 1/2) at tt.



- ❖ A sext consists of two pieces, 0.8 m and 1.6 m long with same field strength.
- ❖ Only the thinner piece of the sext is used at Z.

Twin Aperture Quadrupole

- ❖ An idea of “twin aperture quadrupole” has been developed by A. Milanese to save the power consumption of quadrupole magnets.
- ❖ The currents in the magnet are always surrounded by iron to maximize the usage.



An example of the cross section of a twin aperture quadrupole for FCC-ee (A. Milanese).

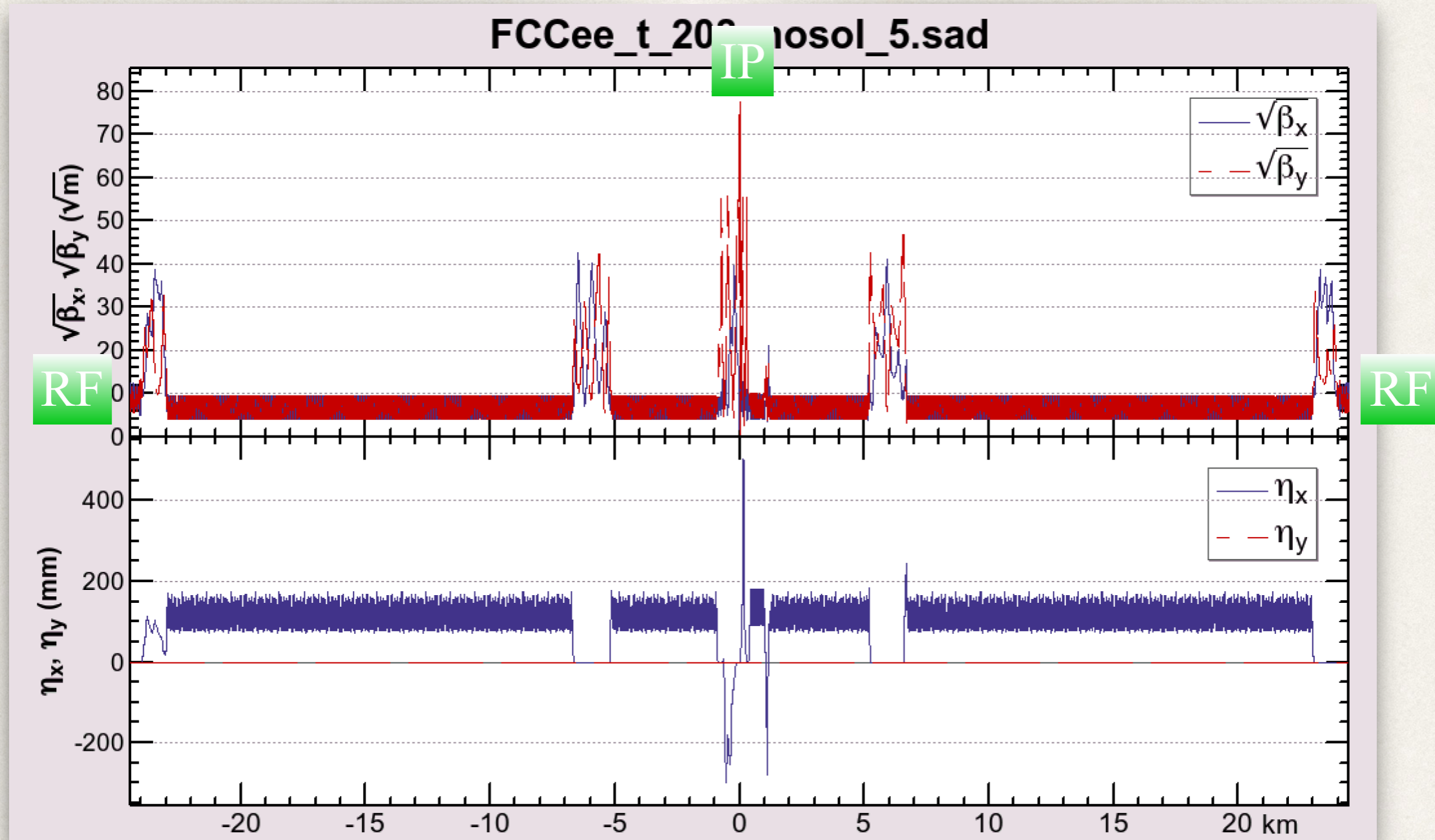
The separation between two beams is 30 cm.

- ❖ The power consumption of the twin aperture quad: 22 MW at 175 GeV with Cu coil = half of single-aperture quads.
- ❖ Dipoles are also “twin”: power consumption = 17 MW at 175 GeV with Al bus bar.



- ❖ Measurement of field quality is going on.
- ❖ A prototype of quadrupole comes soon.

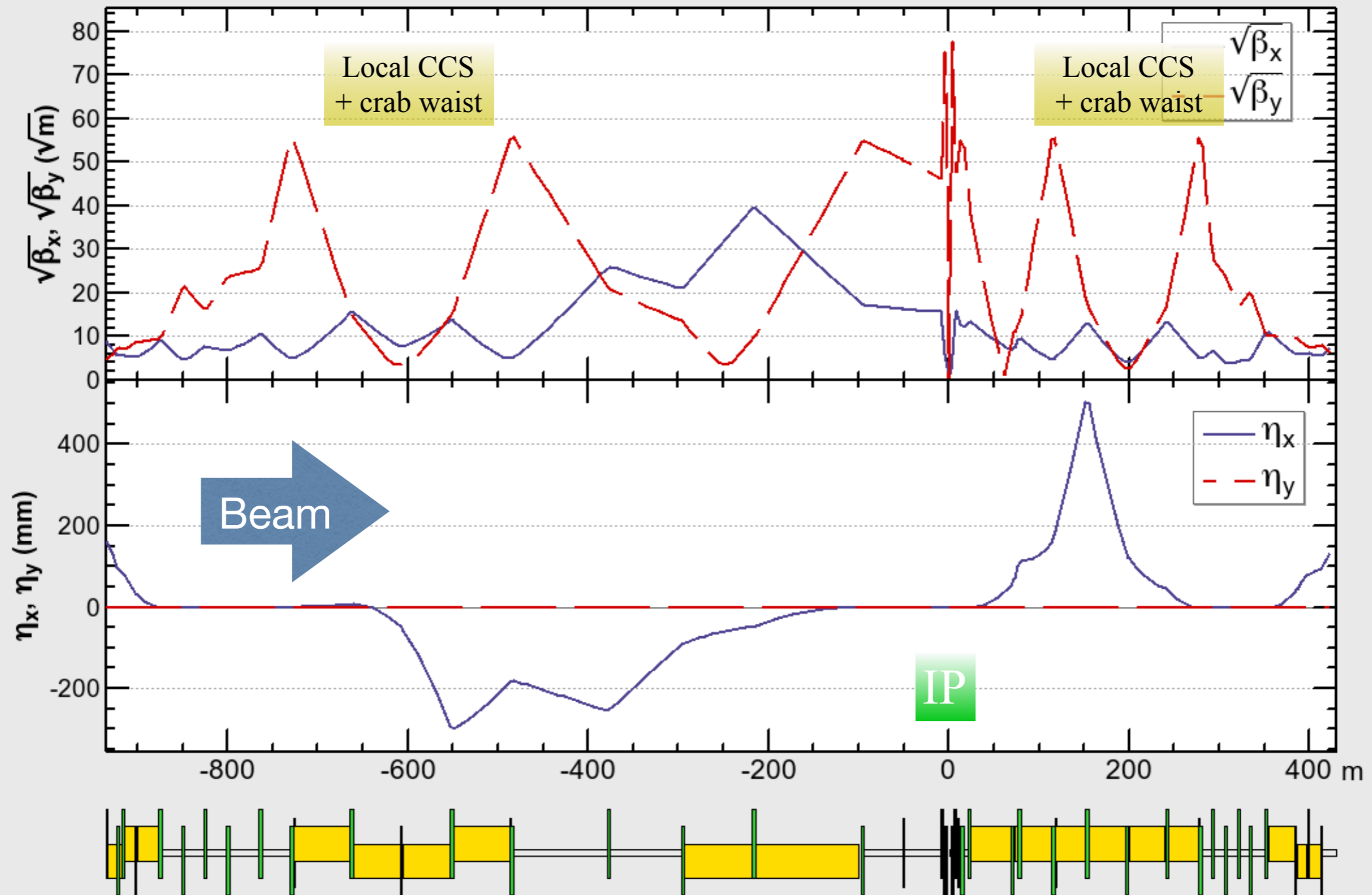
A. Milanese



- Above are the half optics $\beta^*_{x/y} = 1 \text{ m} / 2 \text{ mm}$ at ttbar.
- 2 IPs / ring.
- The optics for straight sections except for the IR are tentative, to be customized for injection/extraction/collimation/polarimeter/beam instrumentation, etc.

Synchrotron radiation toward the IP @ 182.5 GeV

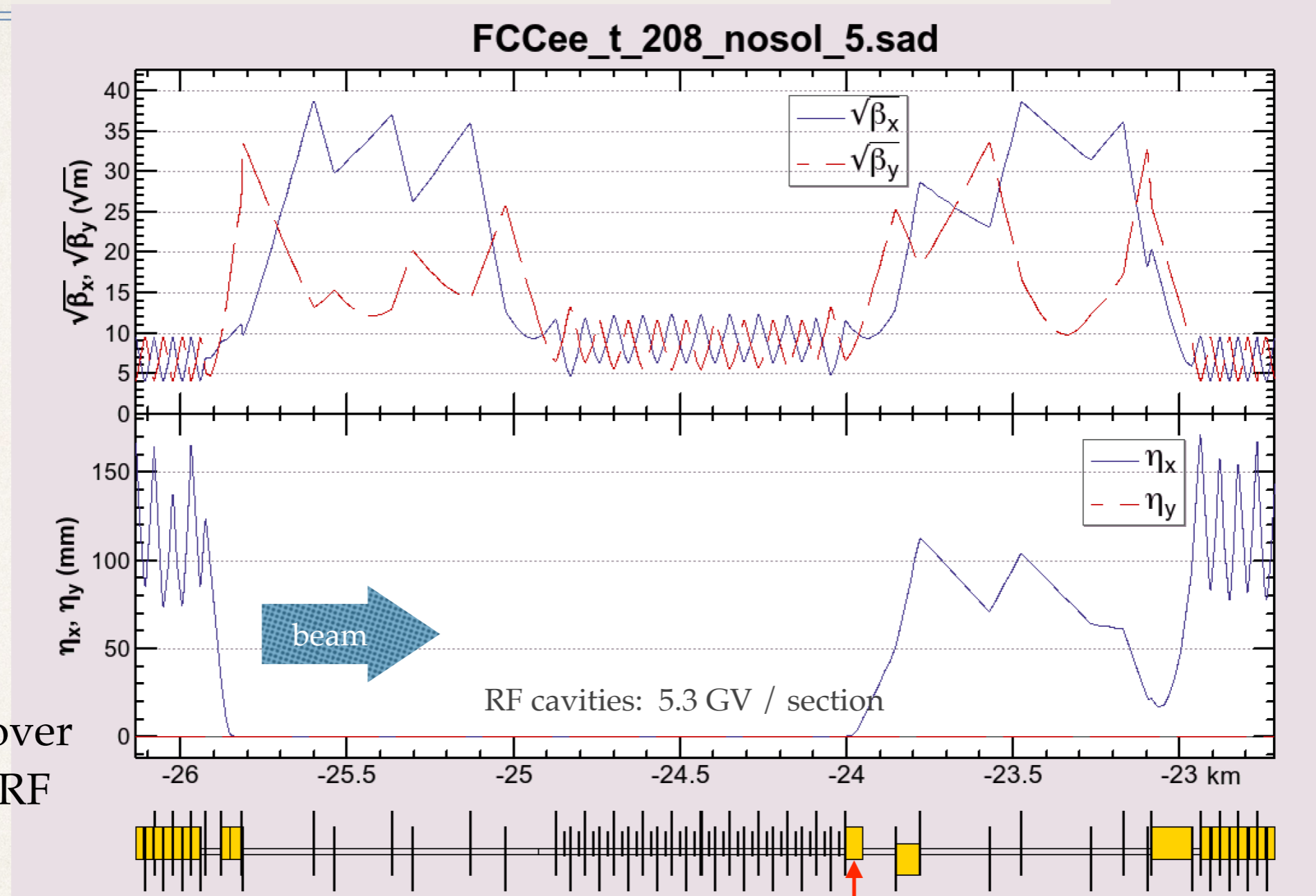
FCCee_t_208_nosol_3.sad



u_c (keV)	1460	571	472	572	226	88.9	93.2	741	648	1004	175	917	...
P_{SR} (kW)	24.2	7.9	7.9	7.9	1.5	0.28	0.47	11.5	0.75	15.9	0.47	0.16	...

$u_c < 100$ keV up to 480 m from the IP.

The RF section (FCC-ee @ 182.5 GeV)

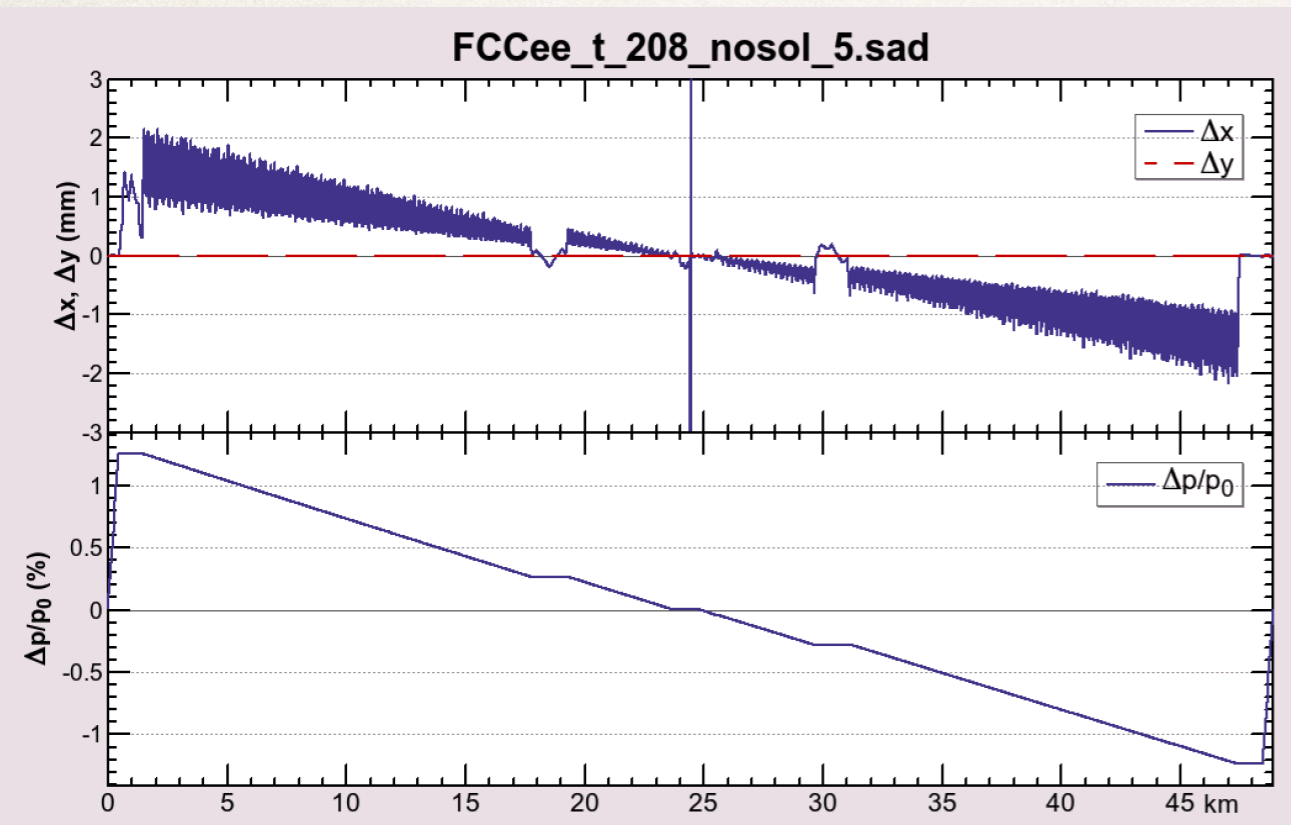


Beams cross over through the RF section.

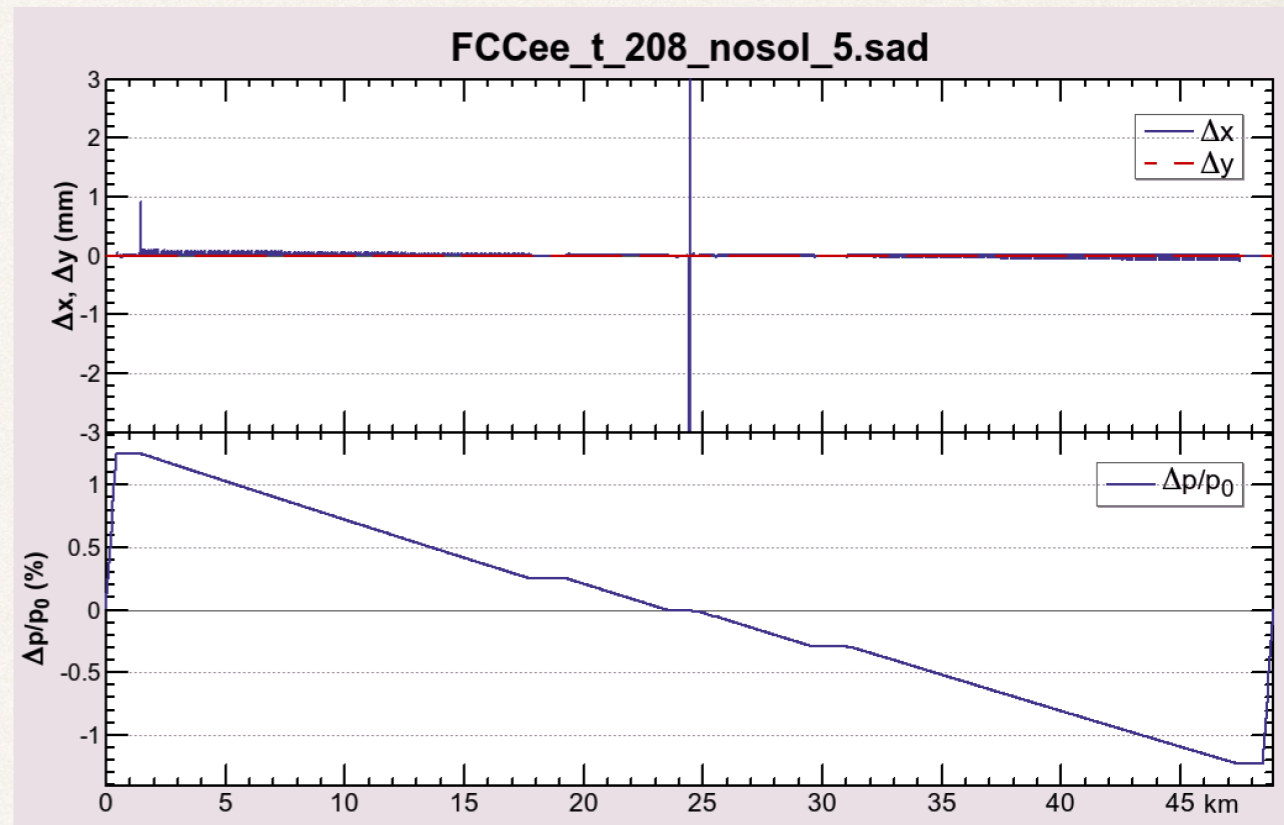
An electrostatic separator, combined with a dipole magnet

- ❖ The usage of the straights on the both sides of the RF is to be determined.
- ❖ If the nominal strengths of quads are symmetrical in the common section, it matches to the optics of both beam.
- ❖ This section is compatible with the RF staging scenario. For lower energy, the common RF and cross over will not be necessary.

No Taper



Tapered

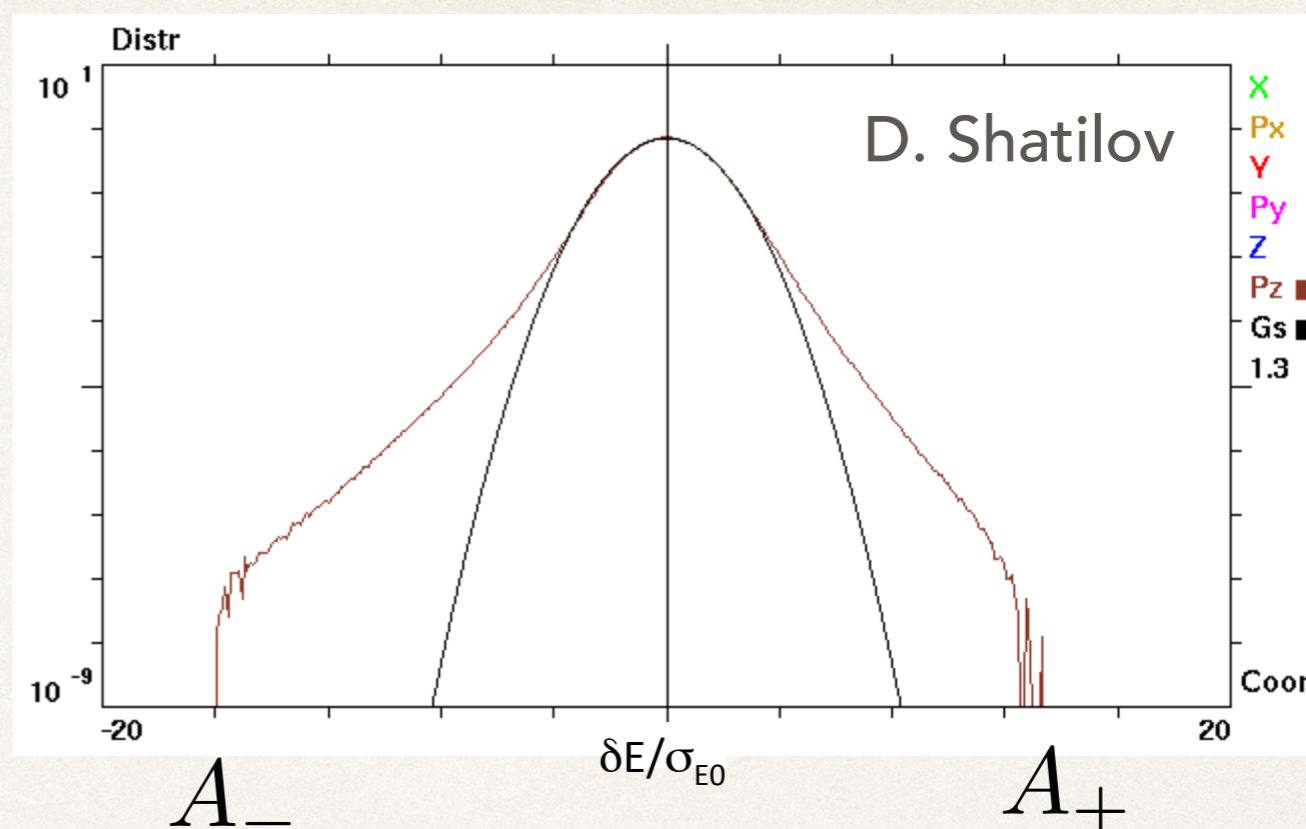


- ❖ The change of the orbit due to energy loss along the arc causes serious deformation on the optics, causing the loss of the dynamic aperture.
- ❖ Everything can be cured almost completely by “tapering”, i.e. scaling the strengths of all magnets along the local energy of the beam: this is one of the best merits of a double-ring collider (F. Zimmermann).

E = 182.5 GeV

$\sigma_{E0} = 0.00153$, $\sigma_E = 0.00193$,
Black line: Gauss with $\sigma_E = 1.3 \sigma_{E0}$

Energy acceptance: 2.5% = $16.3 \sigma_{E0}$



- ❖ The expected energy distribution of the beam has asymmetric tail due to beamstrahlung (D. Shatilov, as above).
- ❖ Thus the required momentum acceptance should be asymmetric: Wider aperture in the negative side.
- ❖ The aperture of the positive side can be expressed as the summation of damping and diffusion terms in a half synchrotron period:

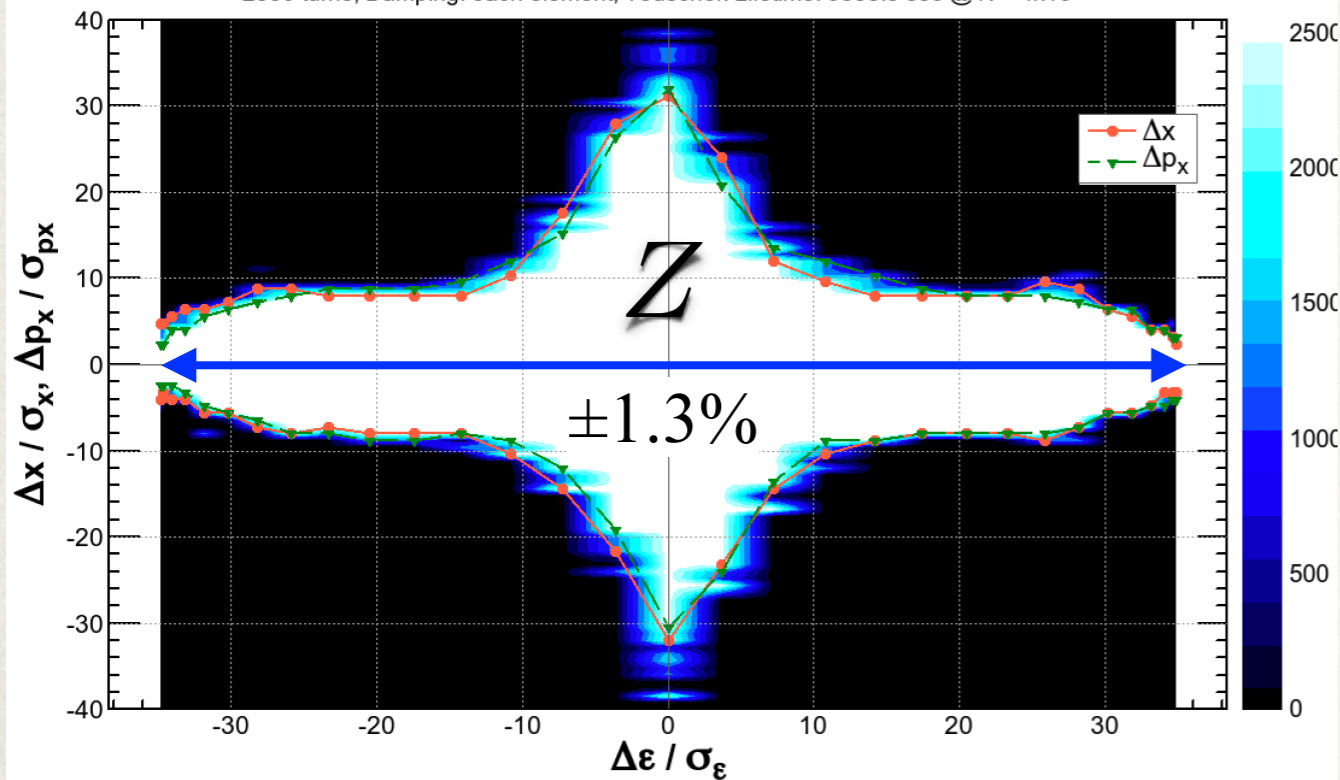
$$A_+ \approx -A_- \exp(-\alpha_z/2\nu_s) + 3\sigma_{\delta,BS} \sqrt{1 - \exp(-\alpha_z/\nu_s)}$$

with the damping rate α_z .

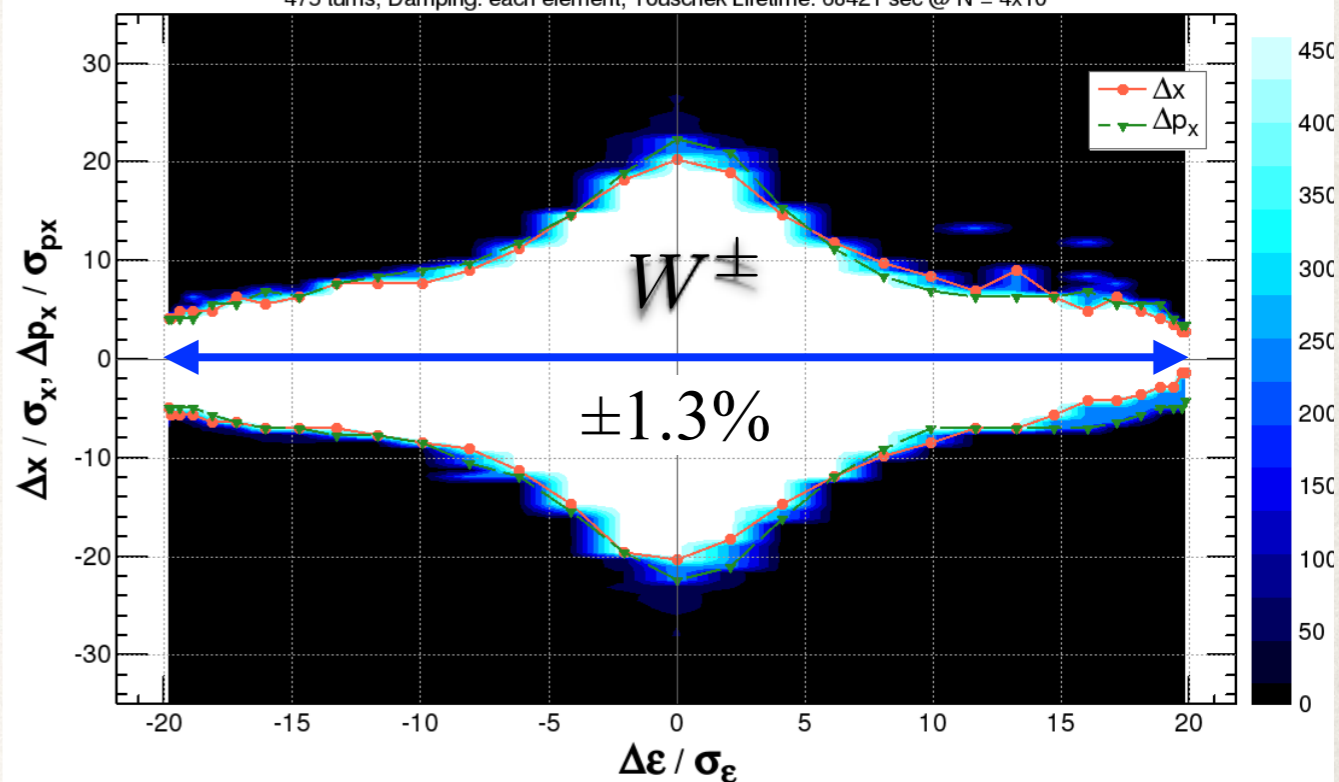
Dynamic Aperture



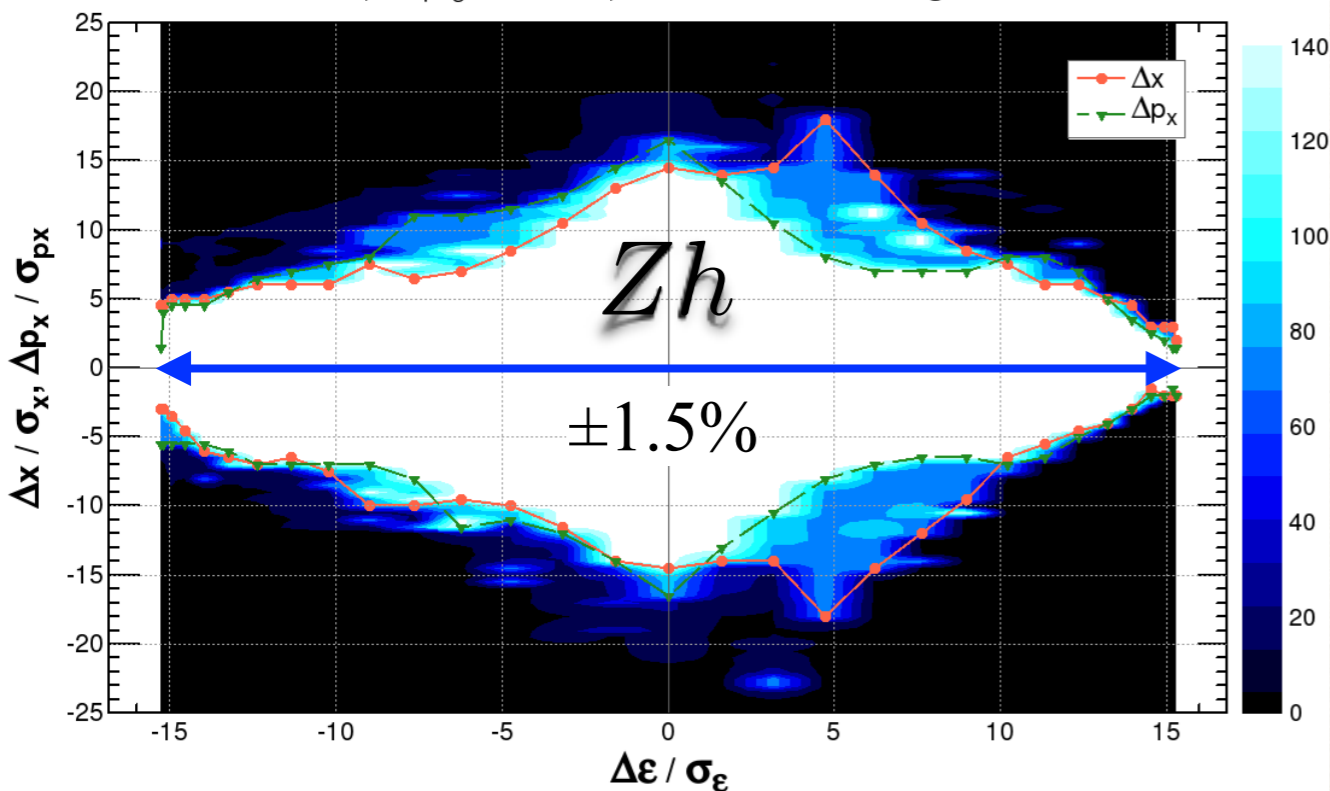
FCCee_z_208_nosol_4.sad: $\epsilon_x = .26$ nm, $\epsilon_y/\epsilon_x = 0.38\%$, $\sigma_\epsilon = 0.037\%$, $\sigma_z = 3.5$ mm,
 $\beta_{x,y} = \{.15$ m, .8 mm $\}$, $v_{x,y,z} = \{273.1380, 271.2199, -0.0245\}$, Crab Waist = 97%
 2550 turns, Damping: each element, Touschek Lifetime: 9538.5 sec @ $N = 4 \times 10^{10}$



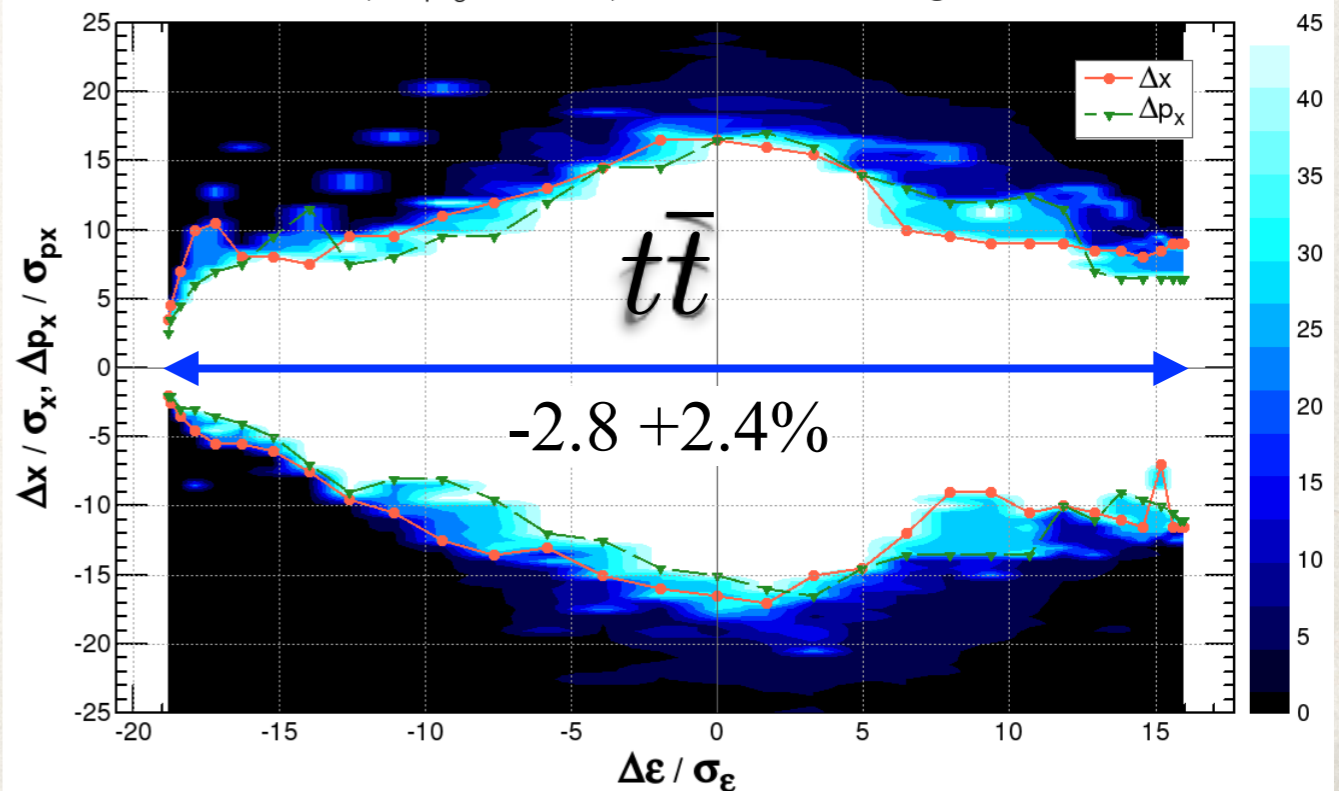
FCCee_w_208_nosol_6.sad: $\epsilon_x = .28$ nm, $\epsilon_y/\epsilon_x = 0.36\%$, $\sigma_\epsilon = 0.065\%$, $\sigma_z = 3.3$ mm,
 $\beta_{x,y} = \{.2$ m, 1 mm $\}$, $v_{x,y,z} = \{389.1538, 389.2195, -0.0227\}$, Crab Waist = 90%
 475 turns, Damping: each element, Touschek Lifetime: 68421 sec @ $N = 4 \times 10^{10}$



FCCee_h_208_nosol_13.sad: $\epsilon_x = .62$ nm, $\epsilon_y/\epsilon_x = 0.16\%$, $\sigma_\epsilon = 0.098\%$, $\sigma_z = 3.2$ mm,
 $\beta_{x,y} = \{.3$ m, 1 mm $\}$, $v_{x,y,z} = \{389.1293, 389.1985, -0.0354\}$, Crab Waist = 85%
 145 turns, Damping: each element, Touschek Lifetime: 1.11E6 sec @ $N = 4 \times 10^{10}$



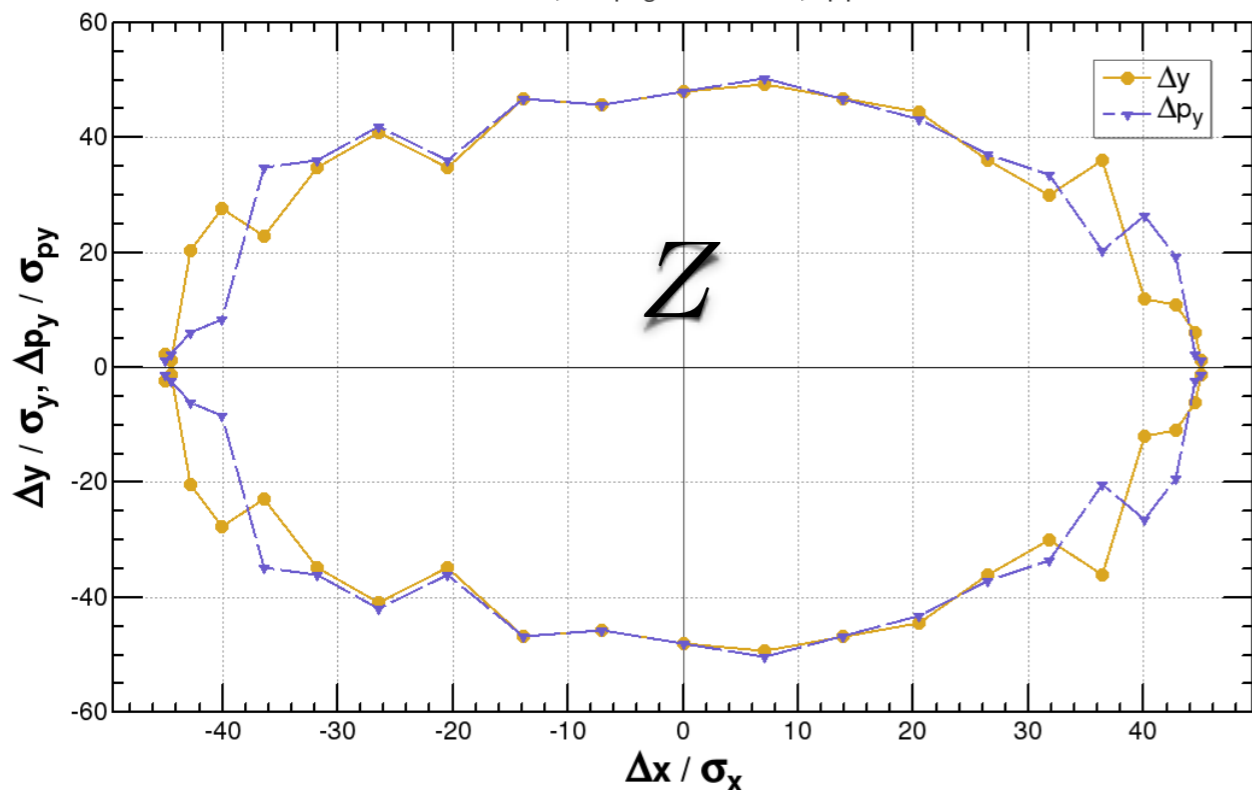
FCCee_t_208_nosol_5.sad: $\epsilon_x = 1.43$ nm, $\epsilon_y/\epsilon_x = 0.20\%$, $\sigma_\epsilon = 0.149\%$, $\sigma_z = 2.5$ mm,
 $\beta_{x,y} = \{1$ m, 1.98 mm $\}$, $v_{x,y,z} = \{389.1038, 389.1761, -0.0680\}$, Crab Waist = 50%
 45 turns, Damping: each element, Touschek Lifetime: 1.02E8 sec @ $N = 1 \times 10^{10}$



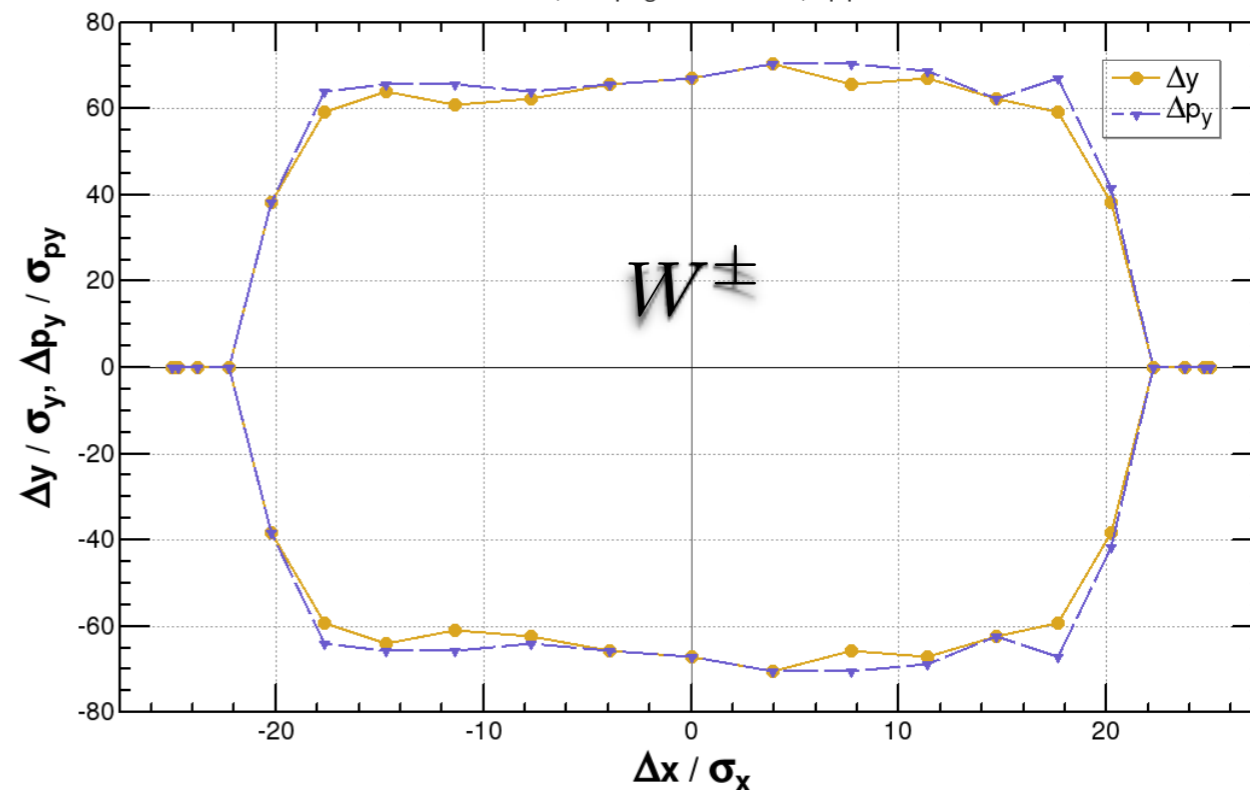
Dynamic Aperture (cont'd : on-momentum transverse)



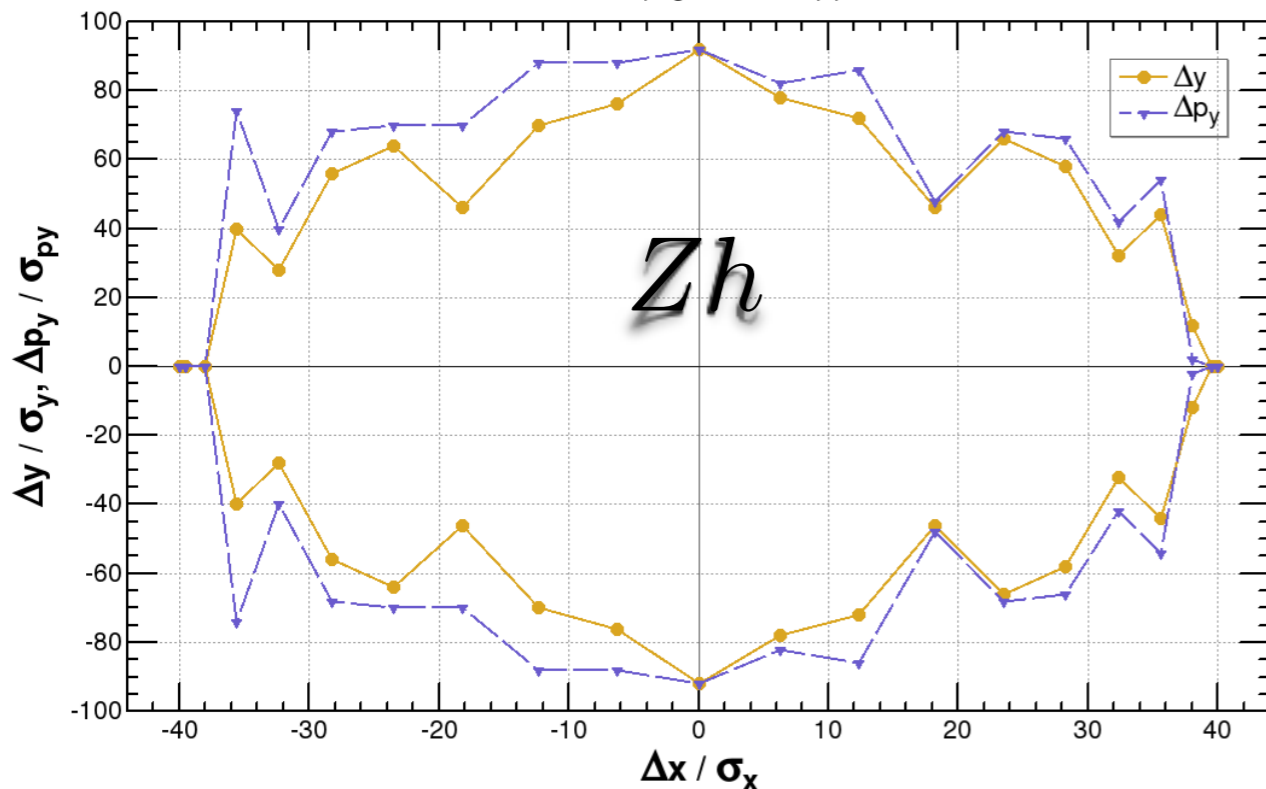
FCCee_z_208_nosol_4.sad: $\epsilon_x = .26$ nm, $\epsilon_y/\epsilon_x = 0.38\%$, $\sigma_\epsilon = 0.037\%$, $\sigma_z = 3.5$ mm,
 $\beta_{x,y} = \{.15$ m, .8 mm $\}$, $v_{x,y,z} = \{273.1380, 271.2199, -0.0245\}$, Crab Waist = 97%
 2550 turns, Damping: each element, $\Delta p/p = 0\%$



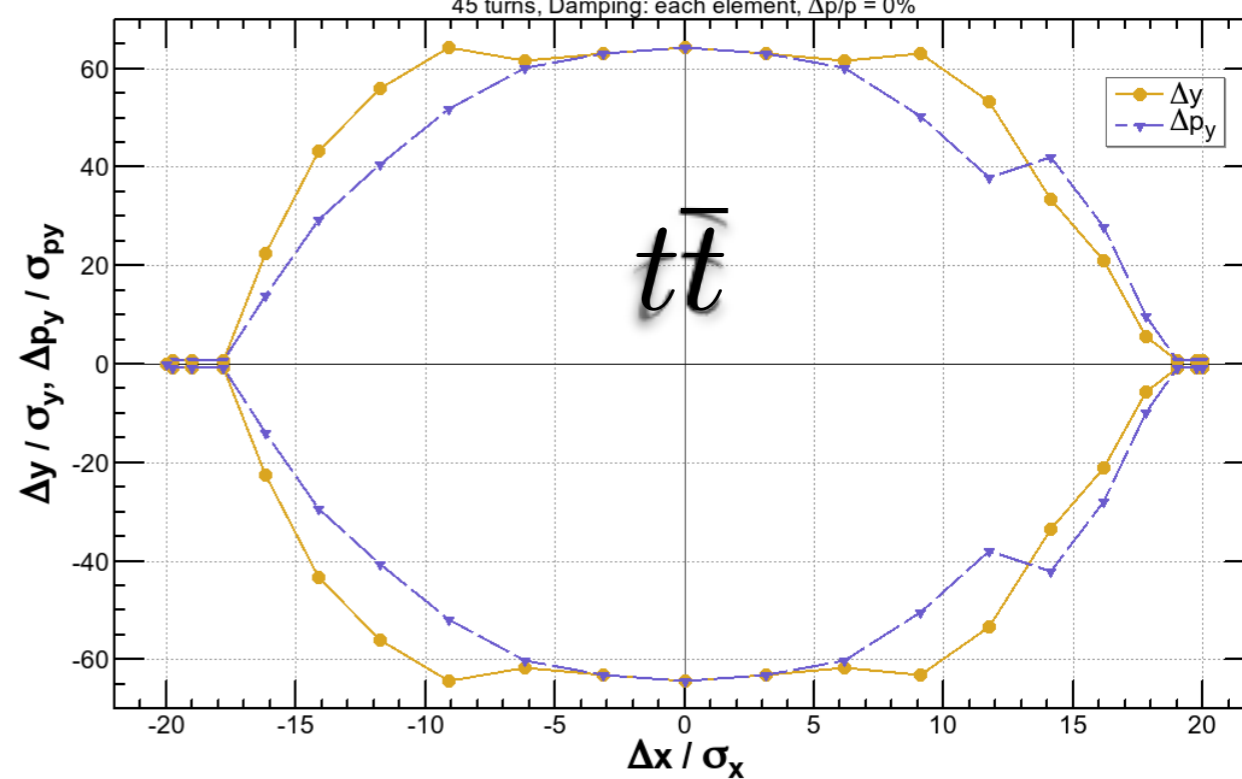
FCCee_w_208_nosol_6.sad: $\epsilon_x = .28$ nm, $\epsilon_y/\epsilon_x = 0.36\%$, $\sigma_\epsilon = 0.065\%$, $\sigma_z = 3.3$ mm,
 $\beta_{x,y} = \{.2$ m, 1 mm $\}$, $v_{x,y,z} = \{389.1538, 389.2195, -0.0227\}$, Crab Waist = 90%
 475 turns, Damping: each element, $\Delta p/p = 0\%$



FCCee_h_208_nosol_13.sad: $\epsilon_x = .62$ nm, $\epsilon_y/\epsilon_x = 0.40\%$, $\sigma_\epsilon = 0.098\%$, $\sigma_z = 3.2$ mm,
 $\beta_{x,y} = \{.3$ m, 1 mm $\}$, $v_{x,y,z} = \{389.1300, 389.2000, -0.0354\}$, Crab Waist = 85%
 145 turns, Damping: each turn, $\Delta p/p = 0\%$



FCCee_t_208_nosol_5.sad: $\epsilon_x = 1.43$ nm, $\epsilon_y/\epsilon_x = 0.20\%$, $\sigma_\epsilon = 0.149\%$, $\sigma_z = 2.5$ mm,
 $\beta_{x,y} = \{1$ m, 1.98 mm $\}$, $v_{x,y,z} = \{389.1038, 389.1761, -0.0680\}$, Crab Waist = 50%
 45 turns, Damping: each element, $\Delta p/p = 0\%$



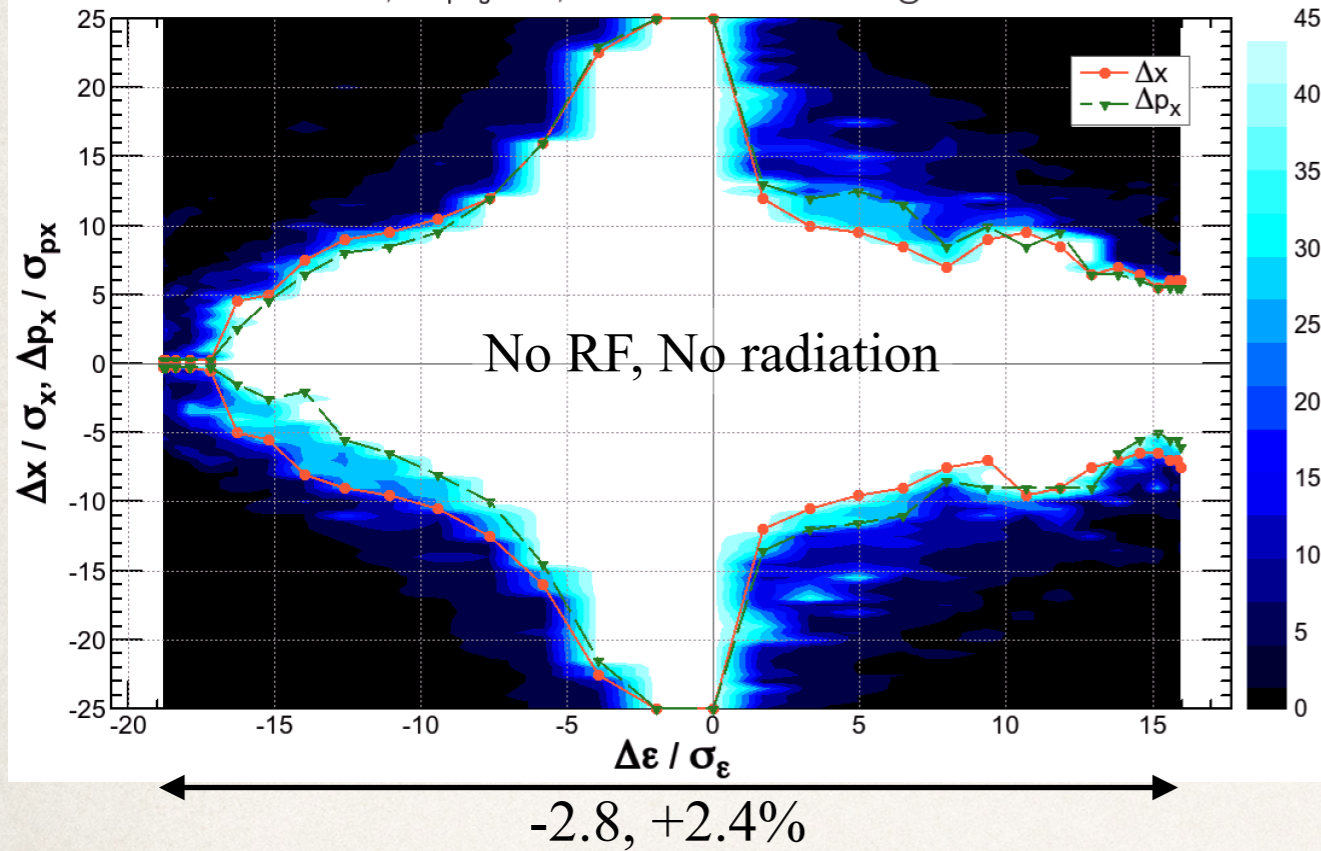
Effects included in the dynamic aperture survey



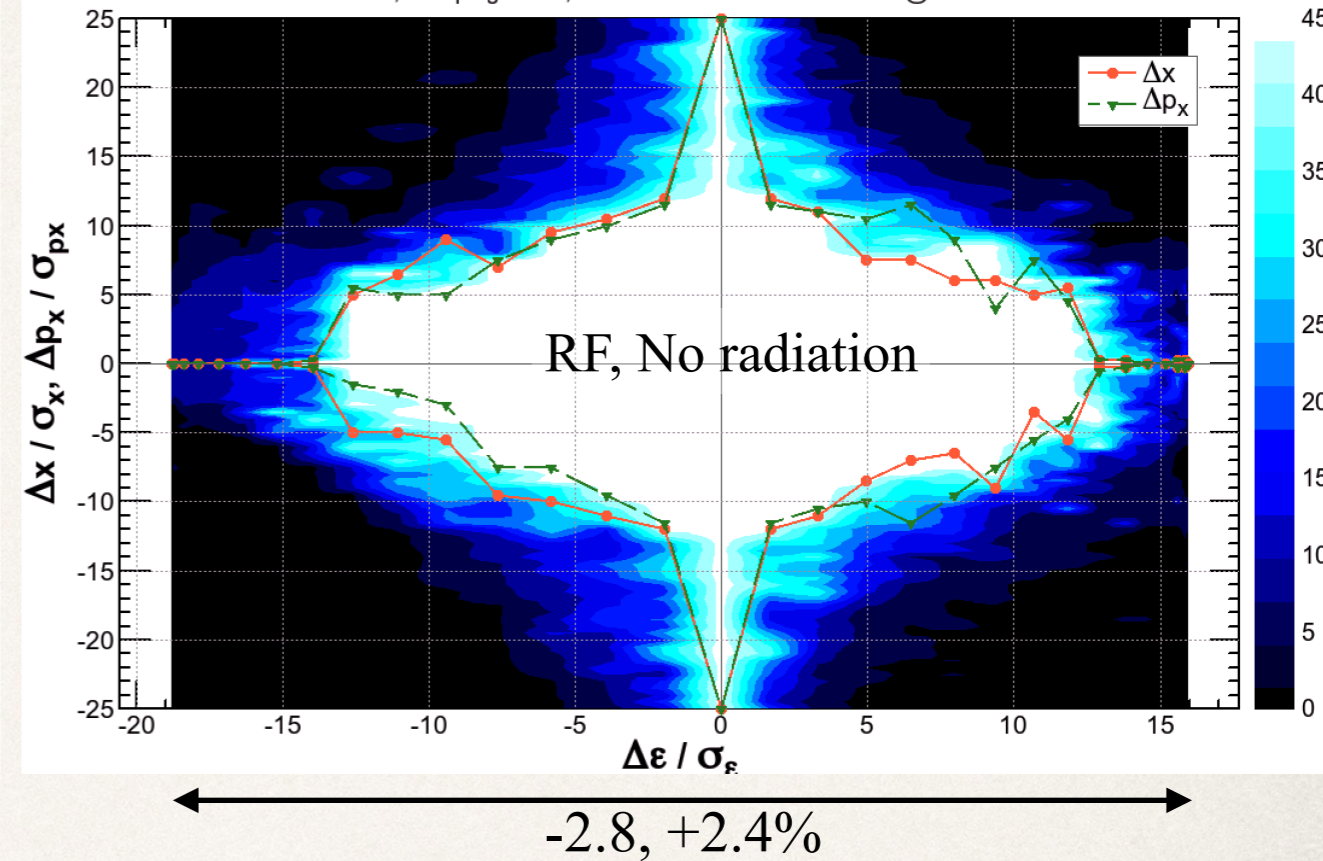
Effects	Included?	Significance
Synchrotron motion	Yes	Essential
Radiation loss in dipoles	Yes	Essential – improves the aperture
Radiation loss in quadrupoles	Yes	Essential – reduces the aperture esp. at $t\bar{t}$
Radiation fluctuation	after optimization	Essential
Tapering	Yes	Essential
Crab waist	Yes	transverse aperture is reduced by $\sim 20\%$
Maxwellian fringes	Yes	small
Kinematical terms	Yes	small
Solenoids	Evaluated separately	minimal, if locally compensated
Beam-beam effects and beamstrahlung for stored beam	after optimization (D. Zhou, D. El Kechen)	affects the lifetime for $\beta_y^* = 1 \text{ mm}$ at $t\bar{t}$
Beam-beam effects for injected beam	on going	
Higher order fields / errors / misalignments	on going	Essential , development of correction/tuning scheme is necessary

Several effects on the dynamic aperture (ttbar)

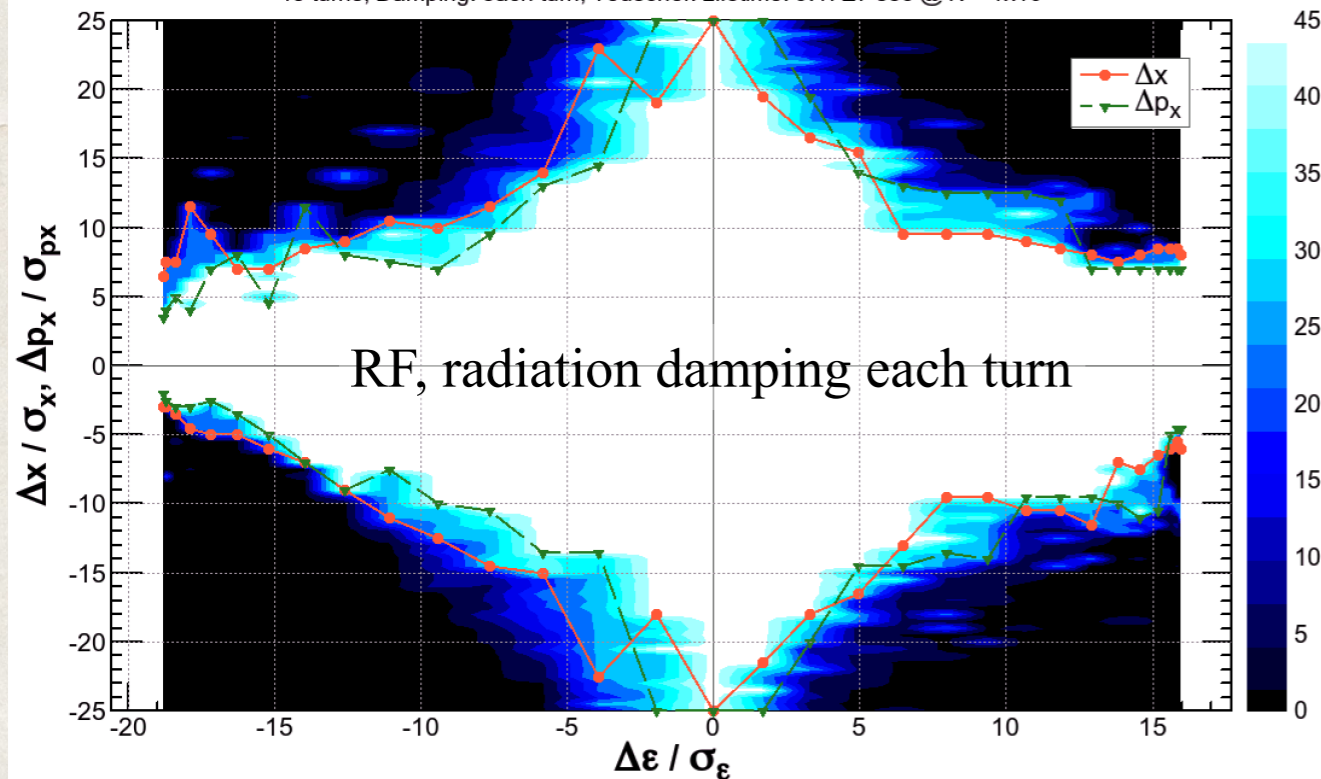
FCCee_t_208_nosol_5.sad: $\epsilon_x = 1.43$ nm, $\epsilon_y/\epsilon_x = 0.20\%$, $\sigma_\epsilon = 0.149\%$, $\sigma_z = 2.5$ mm,
 $\beta_{x,y} = \{1$ m, 2 mm $\}$, $v_{x,y,z} = \{389.1060, 389.1800, -0.0680\}$, Crab Waist = 50%
 45 turns, Damping: none, Touschek Lifetime: $8.16E7$ sec @ $N = 1 \times 10^{10}$



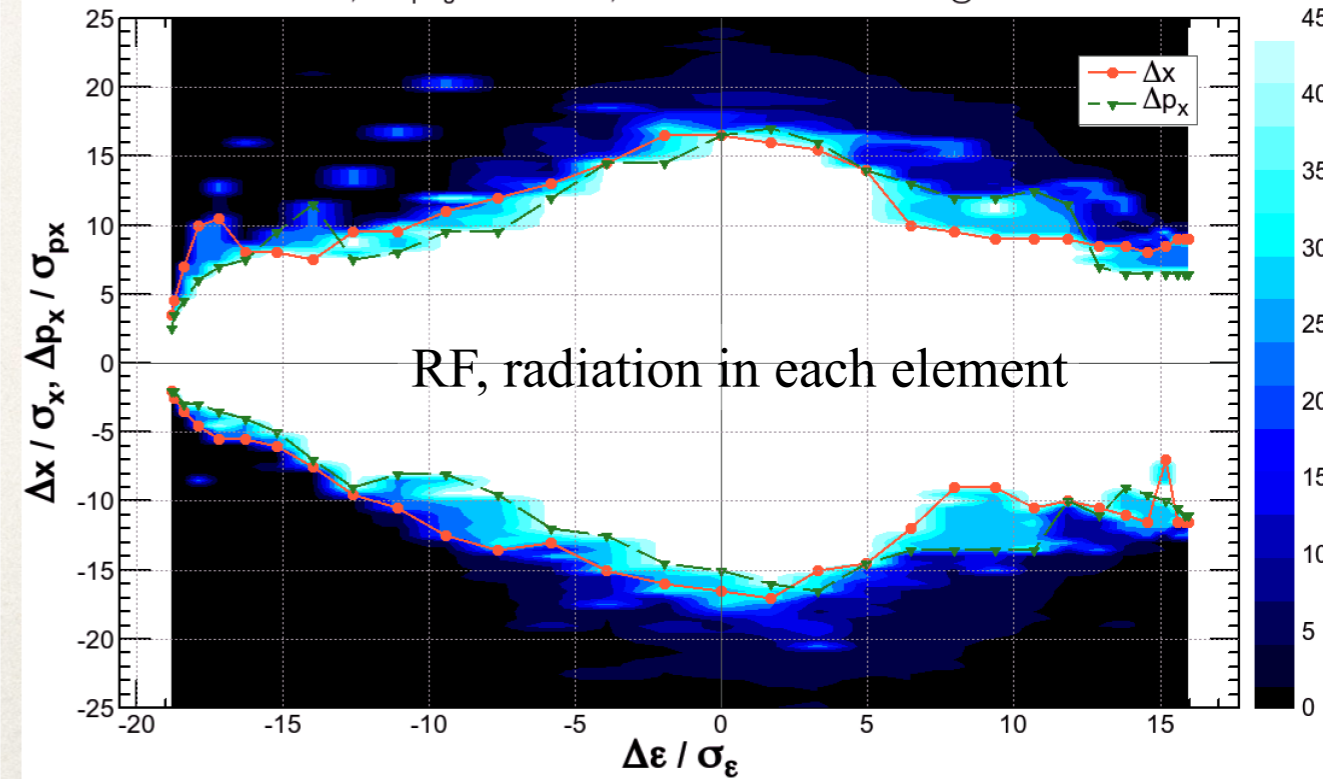
FCCee_t_208_nosol_5.sad: $\epsilon_x = 1.43$ nm, $\epsilon_y/\epsilon_x = 0.20\%$, $\sigma_\epsilon = 0.149\%$, $\sigma_z = 2.5$ mm,
 $\beta_{x,y} = \{1$ m, 2 mm $\}$, $v_{x,y,z} = \{389.1060, 389.1800, -0.0680\}$, Crab Waist = 50%
 45 turns, Damping: none, Touschek Lifetime: $4.88E7$ sec @ $N = 1 \times 10^{10}$



FCCee_t_208_nosol_5.sad: $\epsilon_x = 1.43$ nm, $\epsilon_y/\epsilon_x = 0.20\%$, $\sigma_\epsilon = 0.149\%$, $\sigma_z = 2.5$ mm,
 $\beta_{x,y} = \{1$ m, 2 mm $\}$, $v_{x,y,z} = \{389.1060, 389.1800, -0.0680\}$, Crab Waist = 50%
 45 turns, Damping: each turn, Touschek Lifetime: $9.47E7$ sec @ $N = 1 \times 10^{10}$

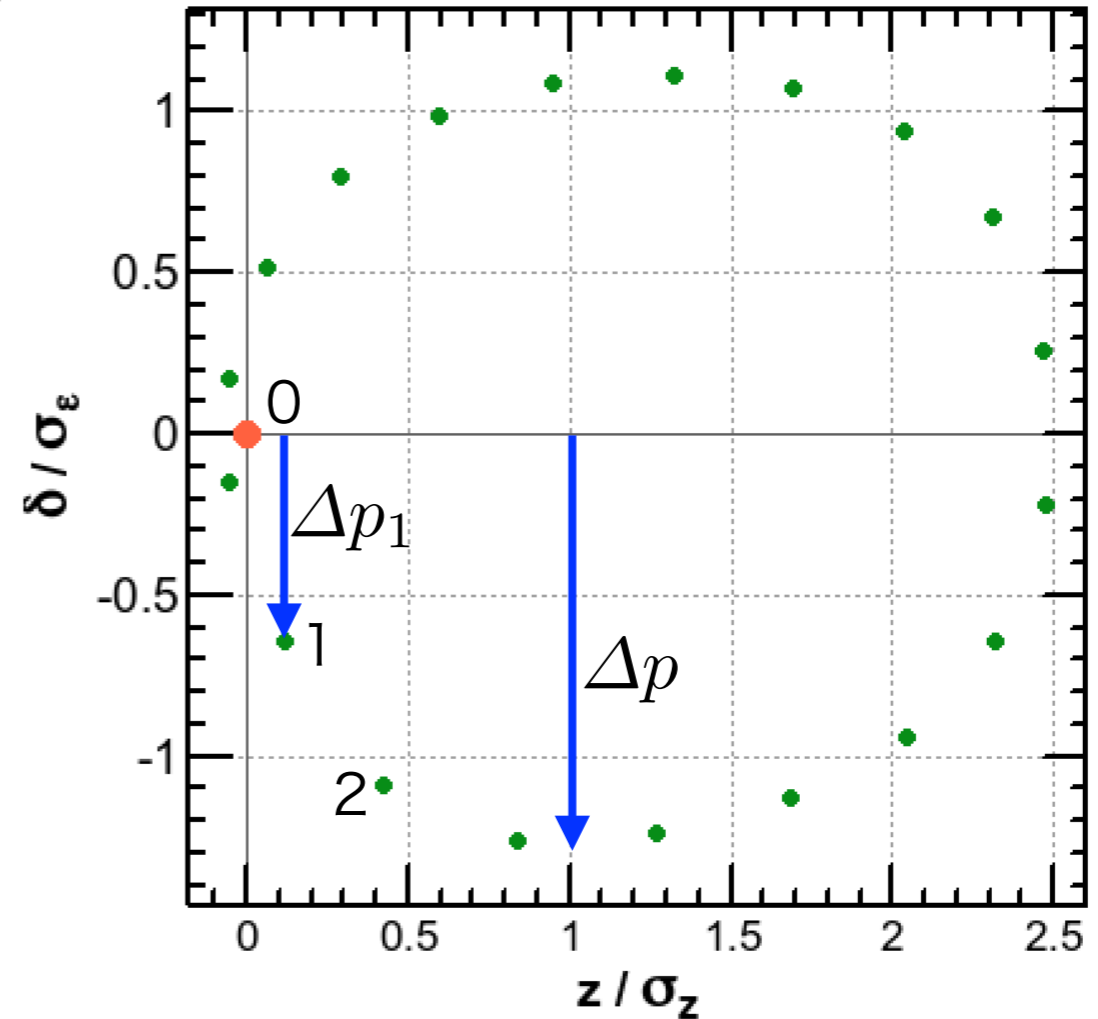
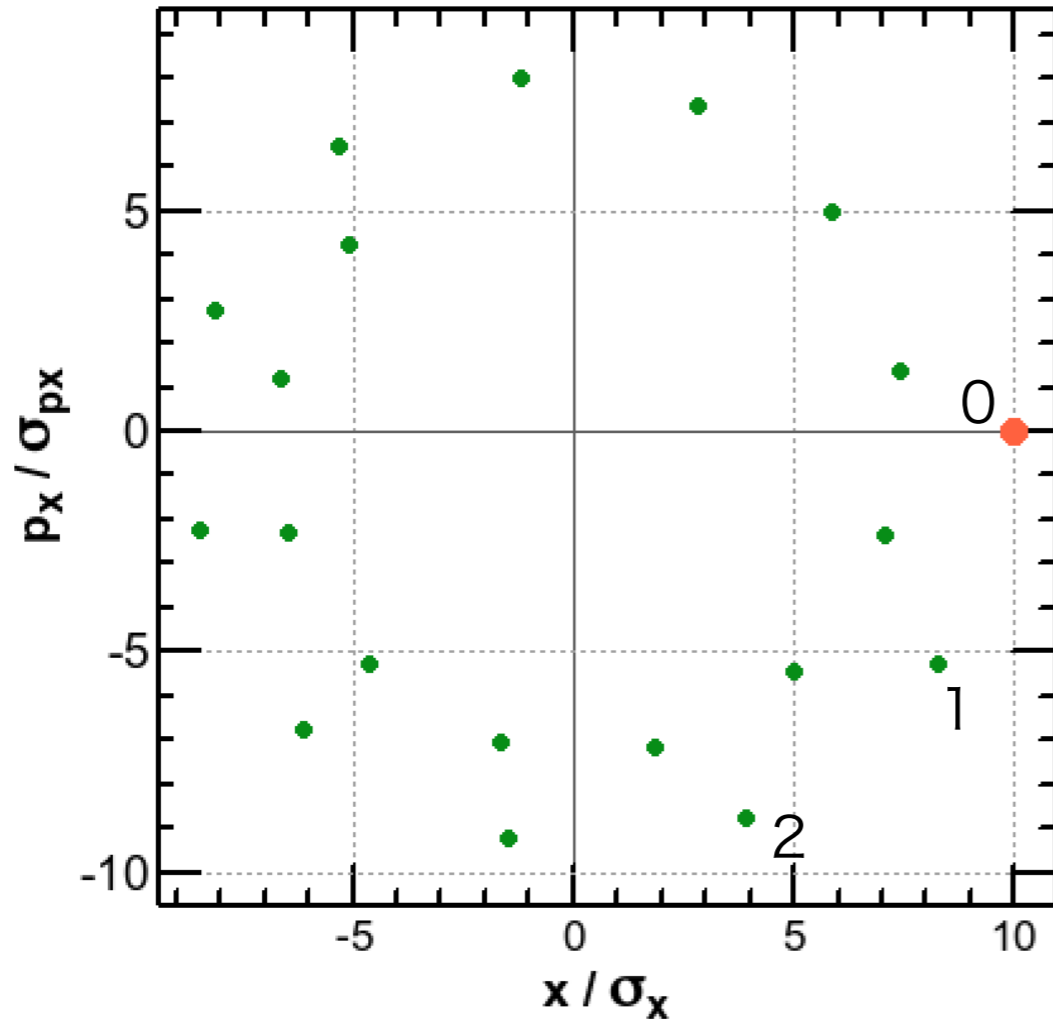


FCCee_t_208_nosol_5.sad: $\epsilon_x = 1.43$ nm, $\epsilon_y/\epsilon_x = 0.20\%$, $\sigma_\epsilon = 0.149\%$, $\sigma_z = 2.5$ mm,
 $\beta_{x,y} = \{1$ m, 1.98 mm $\}$, $v_{x,y,z} = \{389.1038, 389.1761, -0.0680\}$, Crab Waist = 50%
 45 turns, Damping: each element, Touschek Lifetime: $1.02E8$ sec @ $N = 1 \times 10^{10}$



Synchrotron radiation in quadrupoles

FCCee_t_82_symls_by2_1a.sad: $\varepsilon_x = 1.26$ nm, $\varepsilon_y/\varepsilon_x = 0.2\%$, $\sigma_\varepsilon = 0.141\%$, $\sigma_z = 2.9$ mm,
 $\beta_{x,y} = \{1$ m, 2 mm $\}$, $\nu_{x,y,z} = \{387.0800, 387.1400, -0.0658\}$, Crab Waist = 100%
 50 turns, Damping: each element



- ❖ Horizontal betatron oscillation (left) causes a synchrotron motion (right) due to the energy loss by the synchrotron radiation in arc quadrupoles.

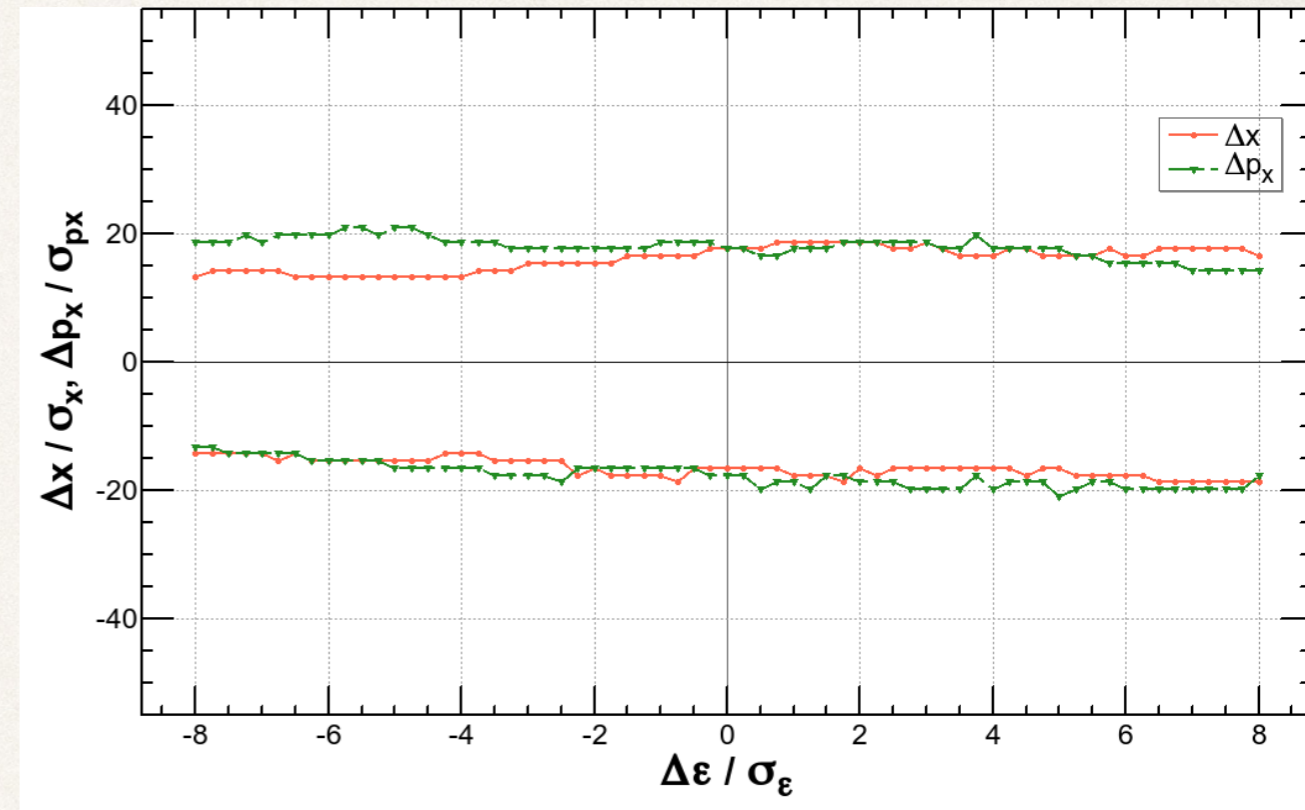
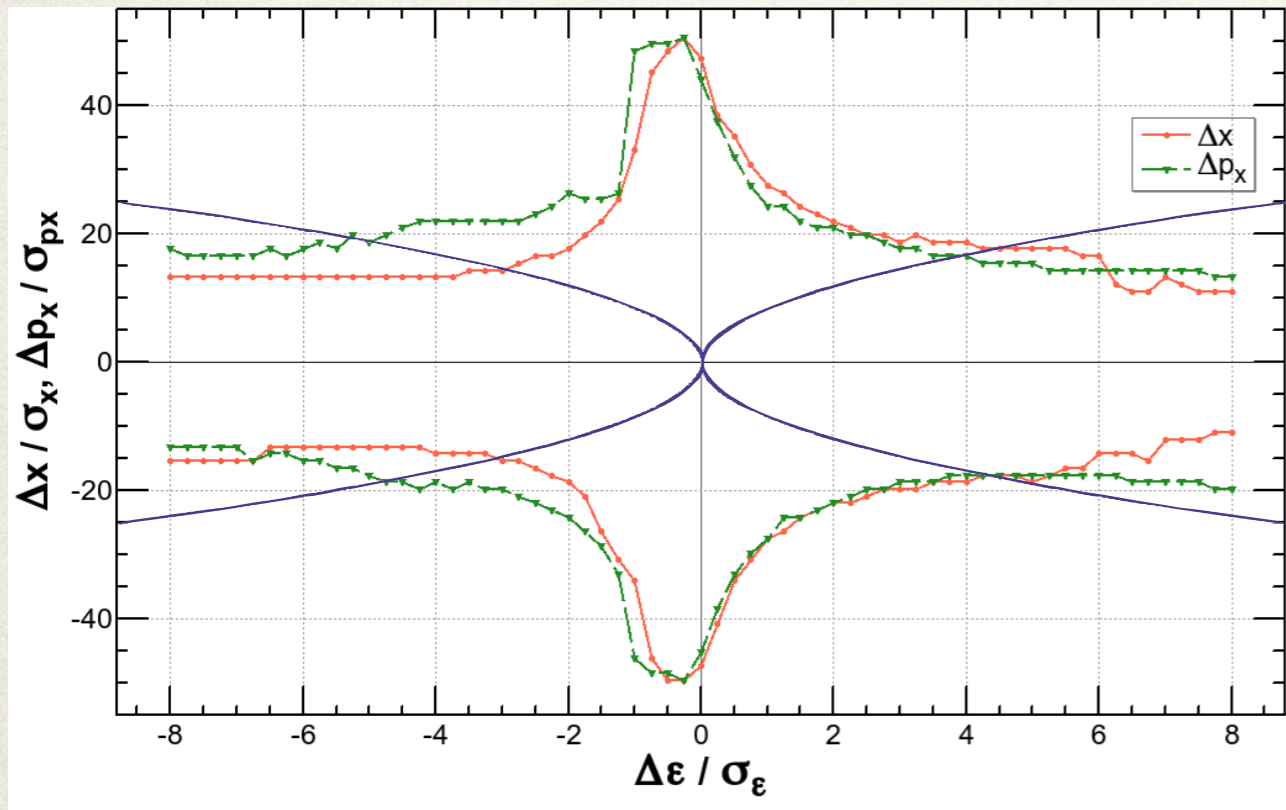
$$\Delta p = \frac{\Delta p_1}{2\pi\nu_s} \exp(-\alpha_z/4\nu_s) = \frac{\alpha_z}{\pi\nu_s J_z} R_Q n^2 \varepsilon_x \exp(-\alpha_z/4\nu_s) ,$$

$$R_Q = \frac{2\sqrt{2}}{\theta_c^2} \left(\frac{\sqrt{2}+1}{l_{QF}} + \frac{\sqrt{2}-1}{l_{QD}} \right) \quad (90^\circ \text{ FODO})$$

- ❖ Such particles can not stay on momentum: reduction of the dynamic aperture.

Synchrotron radiation in quadrupoles (cont'd)

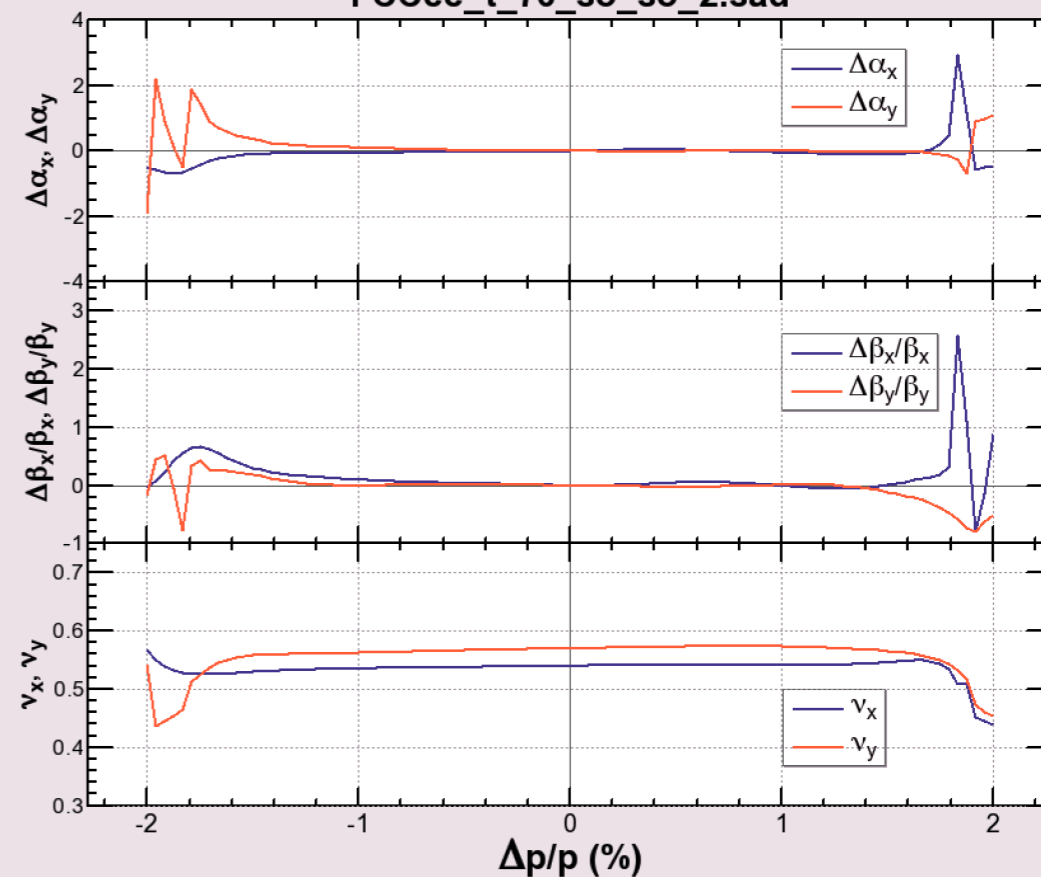
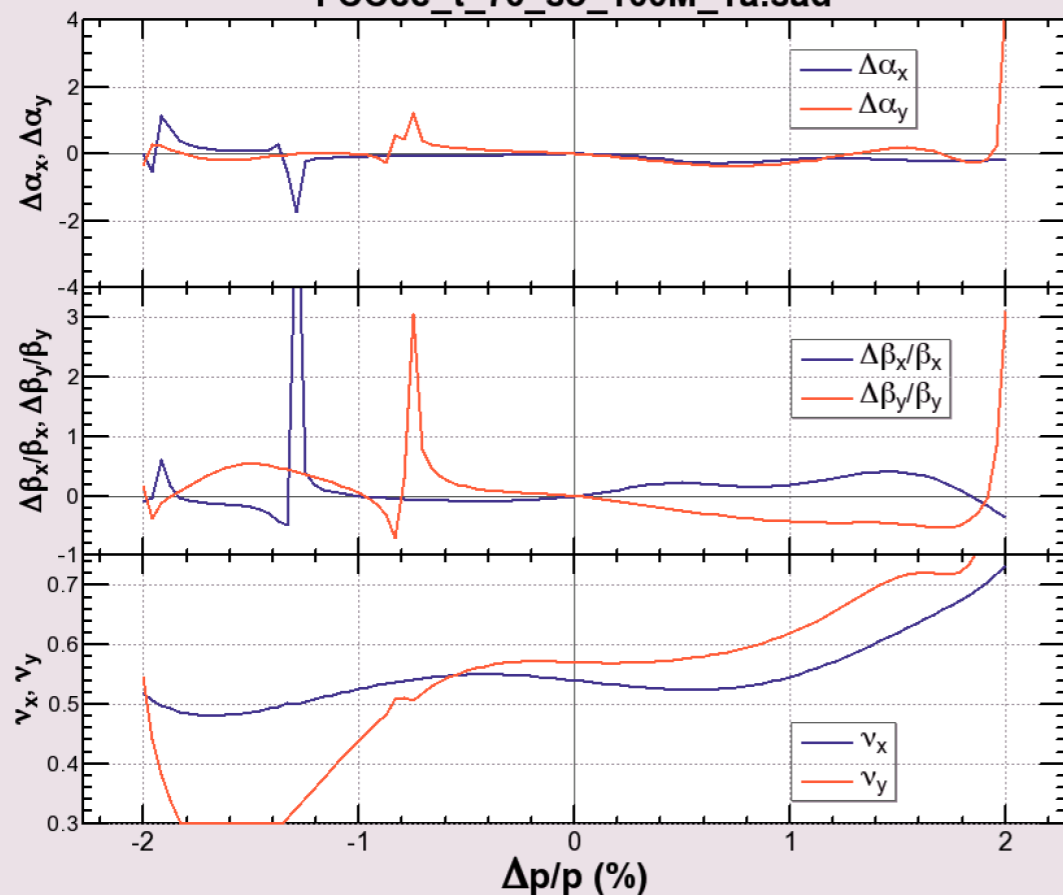
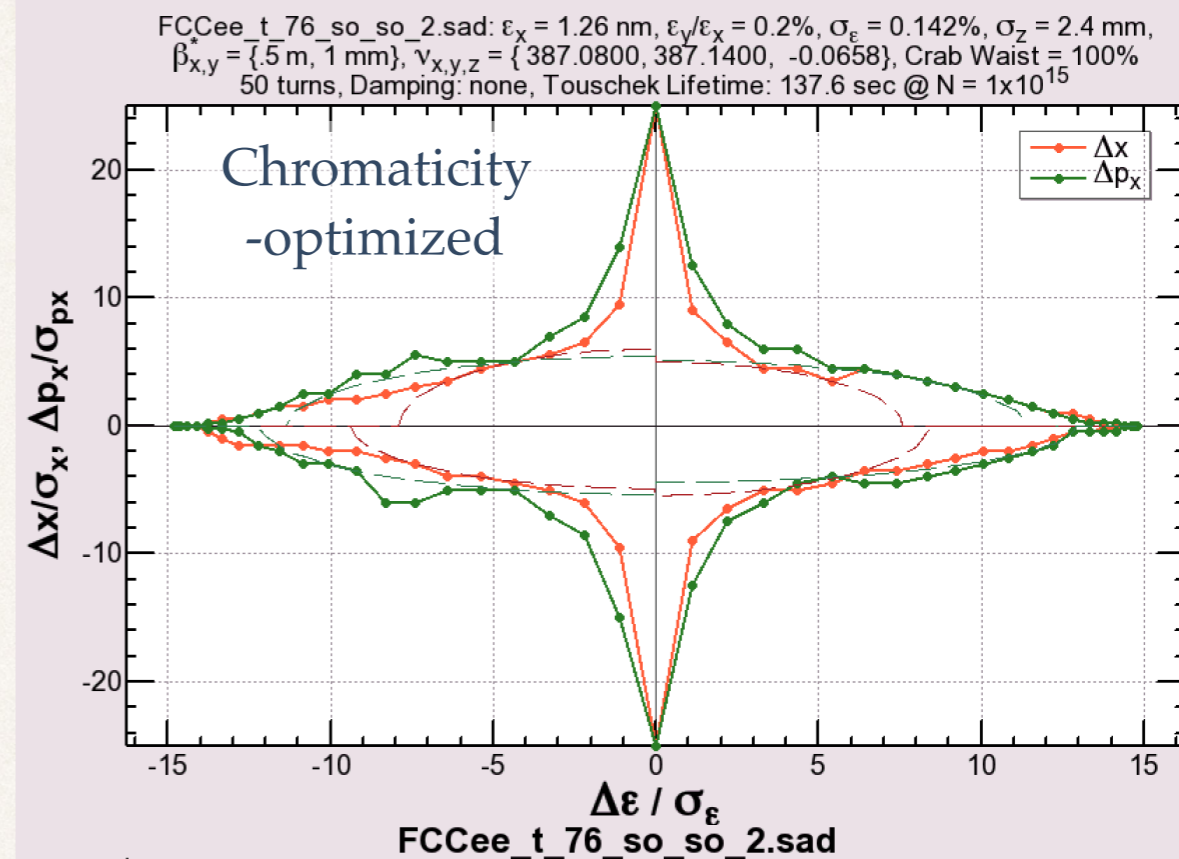
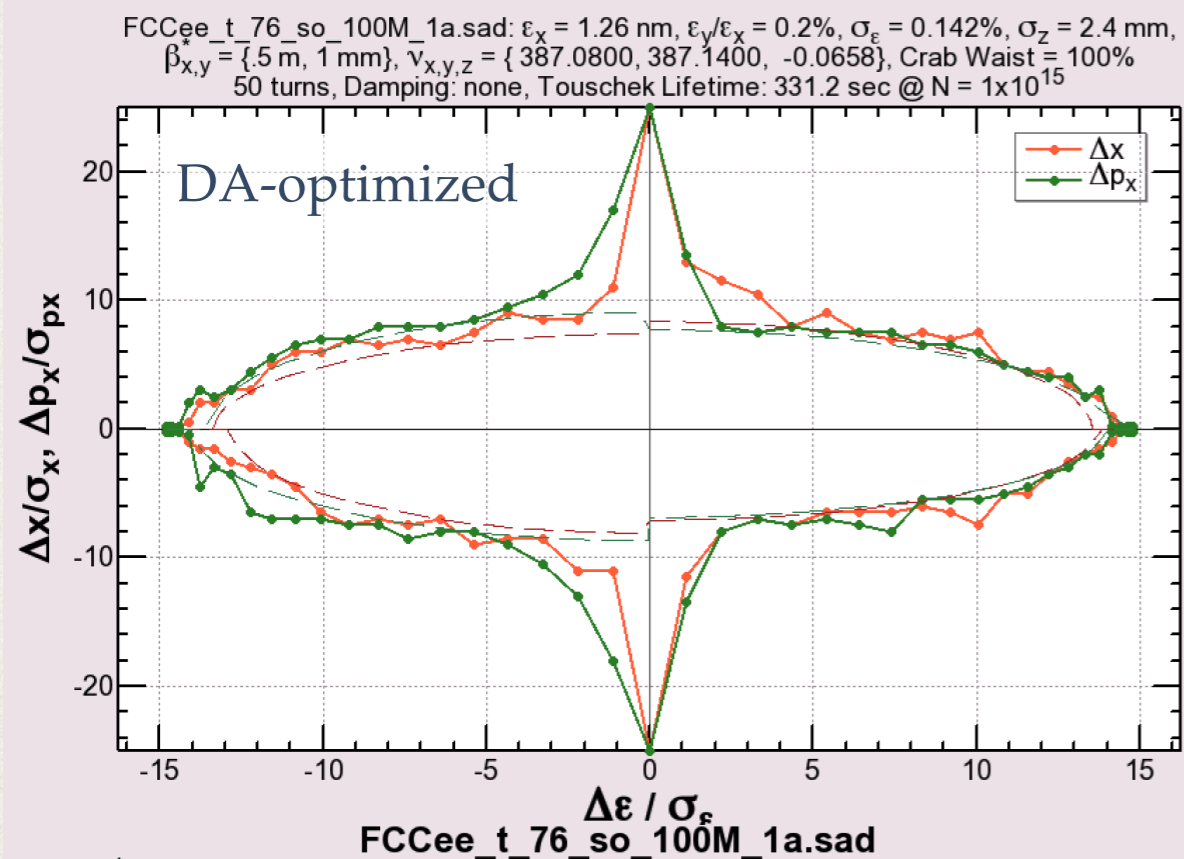
$$E = 175 \text{ GeV}, \beta_{x,y} = (1 \text{ m}, 2 \text{ mm})$$



- ❖ The dynamic aperture without radiation loss in quadrupoles (left) has a sharp peak at on momentum.
- ❖ The peak is destroyed if the radiation in quads is turned on (right).
- ❖ The parabolas on the left show the amplitude of the synchrotron motion due to the radiation in the quadrupole. For a given transverse amplitude, if the parabola is beyond the DA, the particle with that amplitude will be lost.
- ❖ More effects will be presented by A. Bogomyagkov in this workshop.

Less chromaticity \neq better dynamic aperture

$\beta^*_{x,y} = (0.5 \text{ m}, 1 \text{ mm}),$ no radiation damping



Effect of Radiation Fluctuation

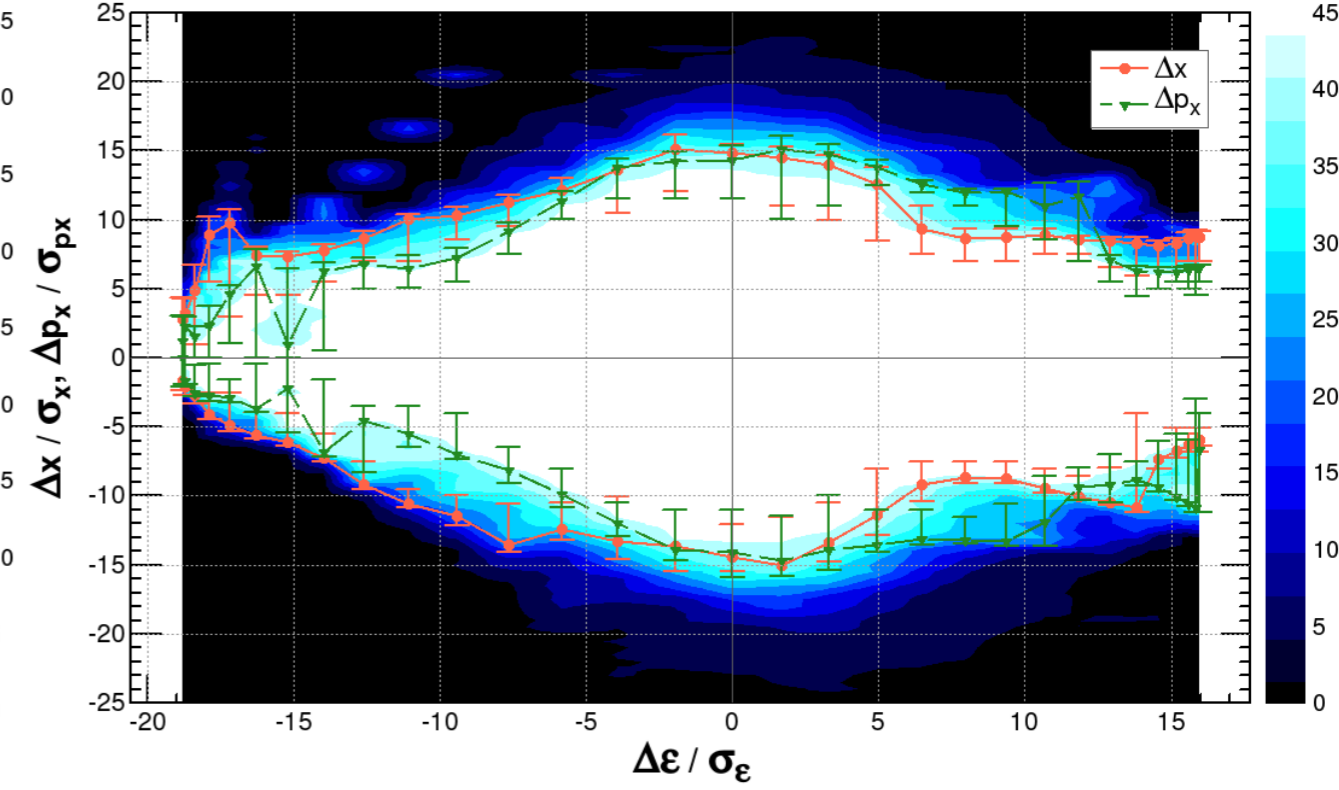
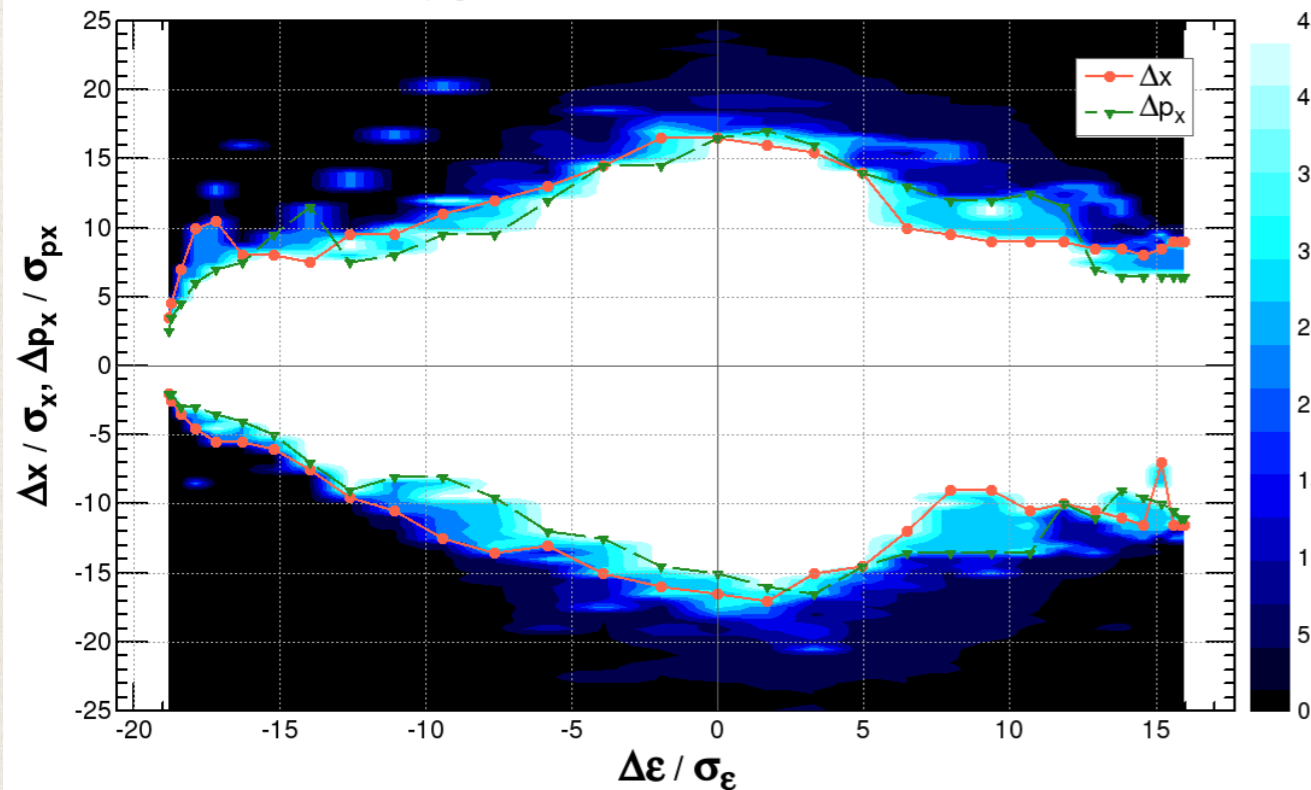
$E = 182.5 \text{ GeV}, \beta_{x,y} = (1 \text{ m}, 2 \text{ mm})$

Radiation damping only

Radiation damping + fluctuation

FCCee_t_208_nosol_5.sad: $\epsilon_x = 1.43 \text{ nm}, \epsilon_y/\epsilon_x = 0.20\%, \sigma_\epsilon = 0.149\%, \sigma_z = 2.5 \text{ mm},$
 $\beta_{x,y} = \{1 \text{ m}, 1.98 \text{ mm}\}, v_{x,y,z} = \{389.1038, 389.1761, -0.0680\}, \text{Crab Waist} = 50\%$
 45 turns, Damping: each element, Touschek Lifetime: $1.02\text{E}8 \text{ sec @ } N = 1 \times 10^{10}$

FCCee_t_208_nosol_5.sad: $\epsilon_x = 1.43 \text{ nm}, \epsilon_y/\epsilon_x = 0.20\%, \sigma_\epsilon = 0.149\%, \sigma_z = 2.5 \text{ mm},$
 $\beta_{x,y} = \{1 \text{ m}, 1.98 \text{ mm}\}, v_{x,y,z} = \{389.1038, 389.1761, -0.0680\}, \text{Crab Waist} = 50\%$
 45 turns, Damping: each element, Touschek Lifetime: $9.11\text{E}7 \text{ sec @ } N = 1 \times 10^{10}, \text{Fluctuation: } 100 \text{ samples}$



$-2.8, +2.4\%$

- (Right figure) 100 samples are taken to evaluate the dynamic aperture with radiation fluctuation.
 - Within the lines: particles survive for 75% of the samples.
 - Error bars correspond to the range of survival between 50% and 100% of the samples.
- It may reasonable that the 50% loss corresponds to the original aperture.
- The thickness between 50% and 100% survival can be attributed to the fractal structure of unstable orbits or resonances in the phase space.

- ❖ Modification of the beam optics for FCC-ee has been performed over the base line optics 2016:
 - ❖ Mitigation of the coherent beam-beam instability at Z
 - ❖ *By achieving smaller β_x^* at Z, W, Zh.*
 - ❖ *Applying 60°/60° cell in the arc, only at Z, compatible with 90°/90° cell at higher energies.*
 - ❖ The base energy for tuba is raised to 182.5 GeV, keeping the SR to IP below 100 keV.
 - ❖ Adopt the “Twin Aperture Quadrupole” scheme for arc quadrupoles.
 - ❖ Fit the footprint to a new FCC-hh layout.
- ❖ The resulting dynamic aperture is sufficient for the beamstrahlung and top-up injection.