

Heavy Quarkonium Associated Production and Multi-parton Interaction

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- 2 The frame of Calculation
- 3 Loop Induced Contributions
- 4 Numerical Result of $\Upsilon + J/\psi$
- 5 Numerical result of $\Upsilon + J/\psi + \phi$ and triple parton scattering
- 6 Summary

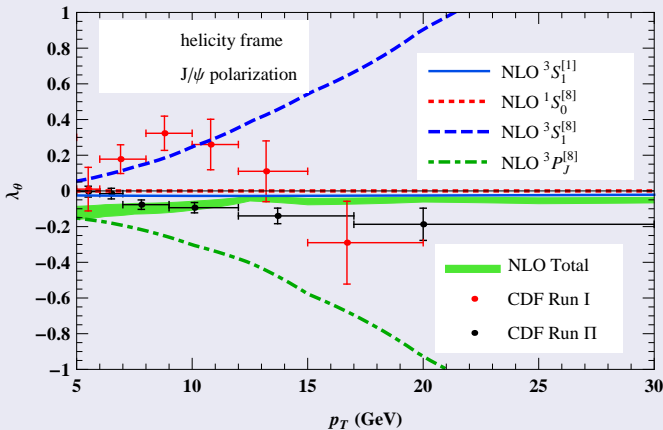
Introduction

Quarkonium productions

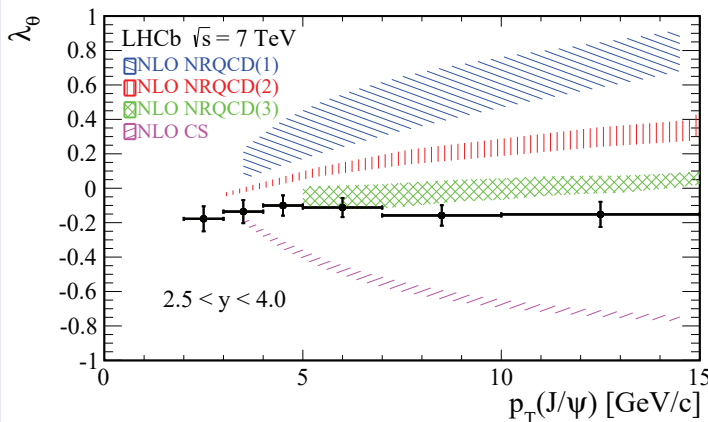
Quarkonium production have been studied by

- 1 Kuang-Ta Chao group
- 2 Yu Jia group
- 3 B. A. Kniehl group
- 4 Cong-Feng Qiao group
- 5 Jian-Xiong Wang group
- 6 ...

Quarkonium productions

NLO J/ψ polarization at CDF, arXiv:1201.2675

Quarkonium productions

NLO J/ψ at LHCb, Chao/Wang/Kniehl, 1506.03981

Quarkonium production and double parton scattering

Many quarkonium associated production processes seems to be dominant by Double-Parton Scattering (DPS).

- 1 $J/\psi + W$ and $J/\psi + Z$, (ATLAS, arXiv:1401.2831, 1412.6428)

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- 2 $J/\psi + charm$ and $\Upsilon + charm$ (LHCb, arXiv:1205.0975, 1510.05949)

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- 3 $J/\psi + J/\psi$ (D0, arXiv:1406.2380; CMS, arXiv:1406.0484)

Quarkonium production and double parton scattering

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- 3 $J/\psi + J/\psi$ (D0, arXiv:1406.2380; CMS, arXiv:1406.0484)
- 4 $\Upsilon + J/\psi$ (D0, arXiv:1511.02428)

Multi parton scattering

The inclusive cross section to produce n hard particles in hadronic colliders is a convolution of generalized n -parton distribution functions (PDF) and elementary partonic cross sections summed over all involved partons,

$$\begin{aligned}
 \sigma_{hh' \rightarrow a_1 \dots a_n}^{\text{NPS}} = & \\
 & \left(\frac{m}{n!}\right) \sum_{i_1, \dots, i_n, i'_1, \dots, i'_n} \int \Gamma_h^{i_1 \dots i_n}(\mathbf{x}_1, \dots, \mathbf{x}_n; \mathbf{b}_1, \dots, \mathbf{b}_n; Q_1^2, \dots, Q_n^2) \\
 & \times \hat{\sigma}_{a_1}^{i_1 i'_1}(\mathbf{x}_1, \mathbf{x}'_1, Q_1^2) \cdots \hat{\sigma}_{a_n}^{i_n i'_n}(\mathbf{x}_n, \mathbf{x}'_n, Q_n^2) \quad (1) \\
 & \times \Gamma_{h'}^{i'_1 \dots i'_n}(\mathbf{x}'_1, \dots, \mathbf{x}'_n; \mathbf{b}_1 - \mathbf{b}, \dots, \mathbf{b}_n - \mathbf{b}; Q_1^2, \dots, Q_n^2) \\
 & \times d\mathbf{x}_1 \dots d\mathbf{x}_n d\mathbf{x}'_1, \dots, d\mathbf{x}'_n d^2b_1, \dots, d^2b_n d^2b.
 \end{aligned}$$

Double parton scattering and Single parton scattering

SPS and DPS

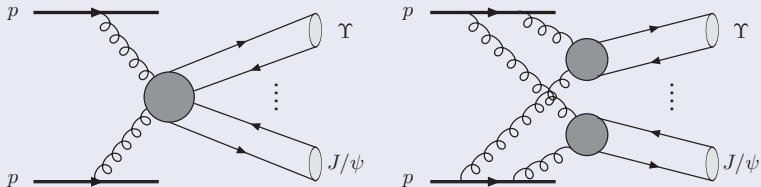


Figure: SPS and DPS of $pp \rightarrow J/\psi + \Upsilon + X$.

Triple parton scattering

TPS

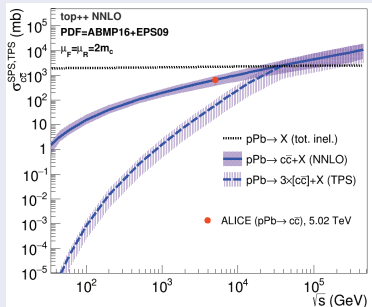
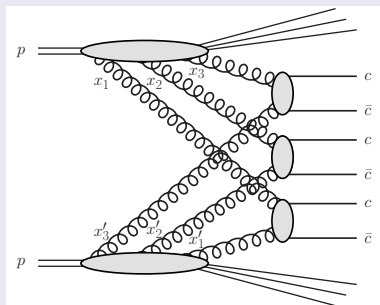


Figure: TPS of $pp \rightarrow c\bar{c} + c\bar{c} + c\bar{c}$ (PRL118, 122001).

The n -parton distribution function (1708.07519)

It encodes all the 3D structure information of the hadron.

- 1 Assumption 1: the n -PDF are factored in terms of longitudinal and transverse components,

$$\Gamma_h^{i_1 \dots i_n} = D_h^{i_1 \dots i_n}(\mathbf{x}_1, \dots, \mathbf{x}_n; Q_1^2, \dots, Q_n^2) f(\mathbf{b}_1) \dots f(\mathbf{b}_n) \quad (2)$$

- 2 We can get hadron-hadron overlap function
 $T(\mathbf{b}) = \int f(\mathbf{b}_1) f(\mathbf{b}_1 - \mathbf{b}) d^2 b_1$, where $1 = \int T(\mathbf{b}) d^2 b$.

- 3 Assumption 2: the longitudinal components reduce to the product of independent single PDF

$$D_h^{i_1 \dots i_n}(\mathbf{x}_1, \dots, \mathbf{x}_n; Q_1^2, \dots, Q_n^2) = D_h^{i_1}(\mathbf{x}_1; Q_1^2) \dots D_h^{i_n}(\mathbf{x}_n; Q_n^2) \quad (3)$$

The cross sections and $\sigma_{\text{eff}}^{nPS}$ (1708.07519)

The cross sections of n -particle associated production

Then we can get

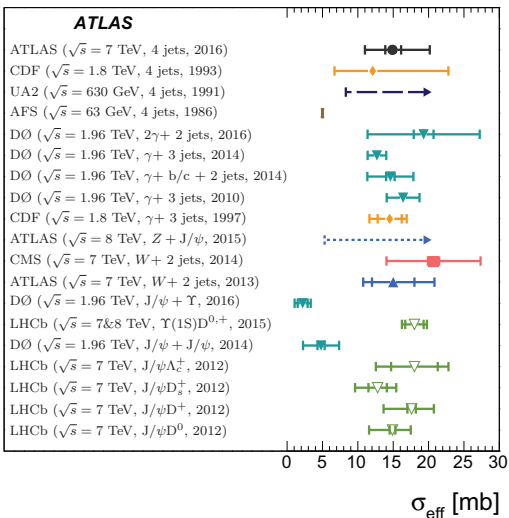
$$\sigma_{hh' \rightarrow a_1 \dots a_n}^{nPS} = \left(\frac{m}{n!} \right) \frac{\sigma_{hh' \rightarrow a_1}^{SPS} \dots \sigma_{hh' \rightarrow a_n}^{SPS}}{(\sigma_{\text{eff}}^{nPS})^{n-1}}, \quad (4)$$

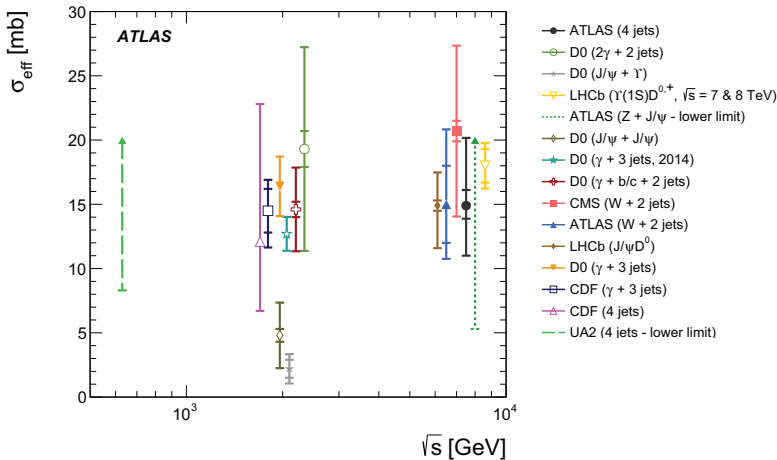
$\sigma_{\text{eff}}^{nPS}$

$$\left(\frac{1}{\sigma_{\text{eff}}^{nPS}} \right)^{n-1} = \int d^2b T^n(\mathbf{b}) \quad (5)$$

$\sigma_{\text{eff}}^{\text{DPS}}$ (arXiv:1608.01857)

Experiment (energy, final state, year)



$\sigma_{\text{eff}}^{\text{DPS}}$ (arXiv:1608.01857)


Prompt $J/\psi + \Upsilon$ @ D0

Prompt $J/\psi + \Upsilon(1S, 2S, 3S)$ @ D0 (arXiv:1511.02428)

$$\sigma_{D0}^{J/\psi+\Upsilon} = 27 \pm 9 \pm 7 \text{ fb} \quad (6)$$

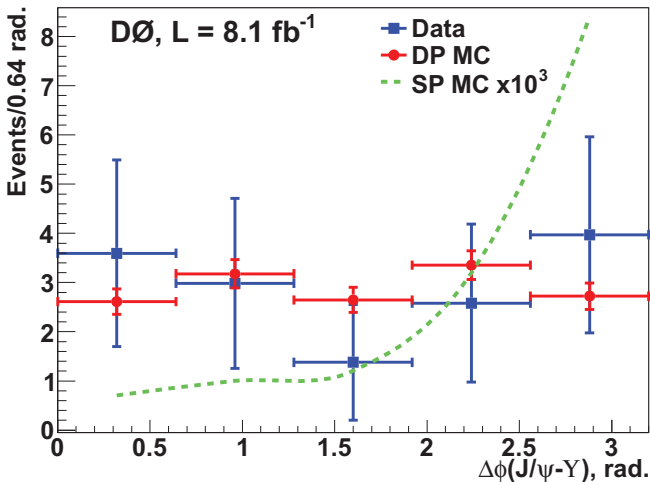
Ignore the SPS contribution

$$\sigma_{DPS}^{J/\psi+\Upsilon} = \sigma_{D0}^{J/\psi+\Upsilon} = \frac{\sigma^{J/\psi} \sigma^{\Upsilon}}{\sigma_{eff}} \quad (7)$$

σ_{eff}

$$\sigma_{eff} = 2.2 \pm 0.7 \pm 0.9 \text{ mb} \quad (8)$$

The distribution of the azimuthal angle between the $J/\psi + \Upsilon$



Color-Singlet contributions of $J/\psi + \Upsilon$

Color-Singlet contributions

Unlike J/ψ -pair or Υ -pair production, neither $\mathcal{O}(\alpha_S^4)$ nor $\mathcal{O}(\alpha_S^5)$ contributions survive in Color-Singlet Model (CSM).

The approximated Loop-Induced (LI) contribution

The approximated Loop-Induced (LI) contribution in CSM at $\mathcal{O}(\alpha_S^6)$ was estimated in Ref. (arXiv:1503.00246) with in the specific limit $\hat{s} \gg |\hat{t}| \gg m_{\psi, \Upsilon}^2$, where \hat{s} and \hat{t} are the Mandelstam variables.

Color-Octet contributions of $J/\psi + \Upsilon$

Color-Octet contributions

The process is a golden observable to probe the so-called Color-Octet Mechanism (COM) (arXiv:1007.3095)

Color-Octet contributions at $\sqrt{s} = 115$ GeV

The Color Octet (CO) contribution were predicted for AFTER@LHC energies $\sqrt{s} = 115$ GeV (arXiv:1504.06531) with HELAC-Onia (arXiv:1212.5293, 1507.03435).

Hadroproduction of $\Upsilon + J/\psi$

SPS contributions were absence

However, the exact calculations of the complete SPS contributions were absence in the literature.

First complete study of $\Upsilon + J/\psi$

We present the first complete study of the simultaneous production of prompt ψ and Υ mesons by including all leading contributions, at order $\mathcal{O}(\alpha_S^6)$ or equivalent.

The frame of Calculation

Cross sections

Hadron and Parton level cross sections

$$\sigma(h_1 h_2 \rightarrow \mathcal{C} + \mathcal{B} + X) = \sum_{a,b} f_{a/h_1} \otimes f_{b/h_2} \otimes \hat{\sigma}(ab \rightarrow \mathcal{C} + \mathcal{B} + X). \quad (9)$$

Parton level cross section

$$d\hat{\sigma}(ab \rightarrow \mathcal{C} + \mathcal{B} + X) = \sum_{n_1, n_2} \hat{\sigma}(ab \rightarrow c\bar{c}[n_1] + b\bar{b}[n_2] + X) \langle O^{\mathcal{C}}(n_1) \rangle \langle O^{\mathcal{B}}(n_2) \rangle \quad (10)$$

Matrix elements

Fock States Of J/ψ

$$\begin{aligned}
 |J/\psi\rangle &= \mathcal{O}(1)|c\bar{c}(^3S_1^{[1]})\rangle + \mathcal{O}(v_c^2)|c\bar{c}(^3S_1^{[8]})gg\rangle \\
 &+ \mathcal{O}(v_c^2)|c\bar{c}(^3P_J^{[1,8]})g\rangle + \mathcal{O}(v_c^2)|c\bar{c}(^1S_0^{[8]})g\rangle + \dots
 \end{aligned}$$

 v^2

$$v_b^2 \sim v_c^2 \sim 0.1 - 0.3$$

$$\alpha_S \sim 0.2$$

$$\alpha_S \sim v_c^2 \sim v_b^2 \tag{11}$$

QED

J^{PC} Of J/ψ or Υ are 1^{--}

QED contributions may be important too.

α

$$\alpha \sim 0.008$$

$$\alpha_S \sim \sqrt{\alpha} \tag{12}$$

$$\mathcal{O}(\alpha_S^6)$$

Color Singlet

The $\mathcal{O}(\alpha_S^4)$ and $\mathcal{O}(\alpha_S^5)$ contributions to $\Upsilon + \psi$ direct production in CSM vanish because of P-parity and C-parity conservation.

Color Octet

$$\mathcal{O}(\alpha_S^4 v_c^i v_b^j) \leq \mathcal{O}(\alpha_S^6) \text{ with } i + j \geq 4$$

EW

$$\mathcal{O}(\alpha_S^2 \alpha^2) \leq \mathcal{O}(\alpha_S^6) \text{ with } i + j \geq 4$$

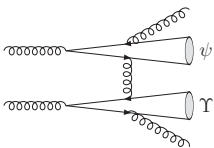
Feeddown for $\chi_{c,b}$

$$\mathcal{O}(\alpha_S^4 v_c^i v_b^j) \leq \mathcal{O}(\alpha_S^6) \text{ with } i + j \geq 4$$

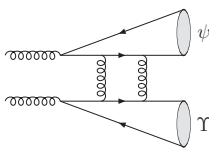
Order of SPS

Label	HELAC-ONIA 2.0 syntax	First order
DR	$g g > cc\sim(3S11) bb\sim(3S11) g g$	$\mathcal{O}(\alpha_S^6)$
LI	addon 8	$\mathcal{O}(\alpha_S^6)$
EW	$p p > cc\sim(3S11) bb\sim(3S11)$	$\mathcal{O}(\alpha_S^2 \alpha^2)$
INTER	addon 8	$\mathcal{O}(\alpha_S^4 \alpha)$
COM	$g g > jpsi y(1s)$	$\mathcal{O}(\alpha_S^4 v_c^i v_b^j), i + j \geq 4$

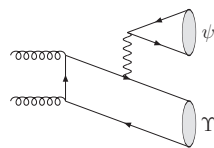
Feynman Diagram of SPS



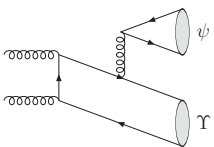
DR



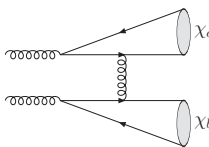
LI



EW



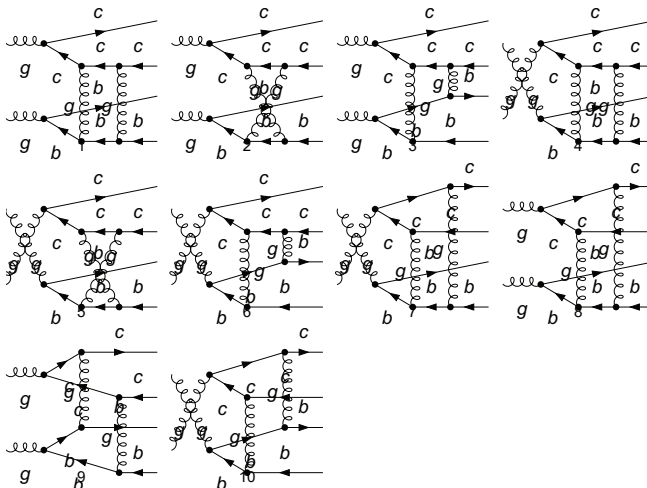
NRQCD



FD

Loop Induced Contributions

Feynman Diagram



Loop induced contributions of $J/\psi + \Upsilon$

Tree contributions is 0

Unlike J/ψ -pair or Υ -pair production, neither $\mathcal{O}(\alpha_S^4)$ nor $\mathcal{O}(\alpha_S^5)$ contributions survive in Color-Singlet Model (CSM).

The Loop-Induced (LI) contributions are UV and IR finite

- 1 The amplitude can be calculated in $D = 4$ directly.
- 2 The gluon mass is introduced to test IR divergence.
- 3 The momentum and polarization vector can be written in $D = 4$ directly.
- 4 ...

Momentum for IR and UV finite amplitude

Momentum of $g(k_1)g(k_2) \rightarrow J/\psi(p_1) + \Upsilon(p_2)$

$$k_1 = \left\{ \frac{\sqrt{s}}{2}, 0, 0, \frac{\sqrt{s}}{2} \right\}$$

$$k_2 = \left\{ \frac{\sqrt{s}}{2}, 0, 0, -\frac{\sqrt{s}}{2} \right\}$$

$$p_1 = \{E_1, 0, \mathbf{p} \times \sin\theta, \mathbf{p} \times \cos\theta\}$$

$$p_2 = \{E_2, 0, -\mathbf{p} \times \sin\theta, -\mathbf{p} \times \cos\theta\} \quad (13)$$

Momentum and polarization vector of $g(k_1)g(k_2)$

Momentum of $g(k_1)g(k_2)$

$$\begin{aligned}k_1 &= \left\{ \frac{\sqrt{s}}{2}, 0, 0, \frac{\sqrt{s}}{2} \right\} \\k_2 &= \left\{ \frac{\sqrt{s}}{2}, 0, 0, -\frac{\sqrt{s}}{2} \right\}\end{aligned}\quad (14)$$

Polarization of $g(k_1)g(k_2)$

$$\begin{aligned}\epsilon_1(k_1) &= \epsilon_2(k_2) = \{0, 1, 0, 0\} \\ \epsilon_2(k_1) &= \epsilon_1(k_2) = \{0, 0, 1, 0\}\end{aligned}\quad (15)$$

Momentum and polarization of $J/\psi(p_1) + \Upsilon(p_2)$

Momentum of $J/\psi(p_1) + \Upsilon(p_2)$

$$\begin{aligned}
 p_1 &= \{E_1, 0, \mathbf{p} \times \sin\theta, \mathbf{p} \times \cos\theta\} \\
 p_2 &= \{E_2, 0, -\mathbf{p} \times \sin\theta, -\mathbf{p} \times \cos\theta\}
 \end{aligned}
 \tag{16}$$

Polarization of $J/\psi(p_1) + \Upsilon(p_2)$

$$\begin{aligned}
 \epsilon_L(p_1) &= 1/m_J \{\mathbf{p}, 0, E_1 \sin\theta, E_1 \cos\theta\} \\
 \epsilon_L(p_2) &= 1/m_\Upsilon \{\mathbf{p}, 0, -E_2 \sin\theta, -E_2 \cos\theta\} \\
 \epsilon_{T1}(p_1) &= \epsilon_{T1}(p_2) = \{0, 1, 0, 0\} \\
 \epsilon_{T2}(p_1) &= \epsilon_{T2}(p_2) = \{0, 0, -\cos\theta, \sin\theta\}
 \end{aligned}
 \tag{17}$$

Amplitude

Amplitude of $g(k_1, 1)g(k_2, 1) \rightarrow J/\psi(p_1, L) + \Upsilon(p_2, L)$

- Color factor can be calculated diagram by diagram. It can be considered as a global factor.

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- Spin projector operator can be used directly.
- The scalar product of k_1 , k_2 , p_1 , p_2 and polarization vector can be expressed by s , m_J , m_Υ , E_1 , p , θ .

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- Loop integrate.

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- Spin projector operator can be used directly.
- The scalar product of k_1, k_2, p_1, p_2 and polarization vector can be expressed by $s, m_J, m_\Upsilon, E_1, p, \theta$.
- Loop integrate.
- Amplitude can be expressed by $s, m_J, m_\Upsilon, E_1, p, \theta$.
- Simplify the amplitude

Numerical Result of $\Upsilon + J/\psi$

Direct SPS cross sections @ D0 in fb

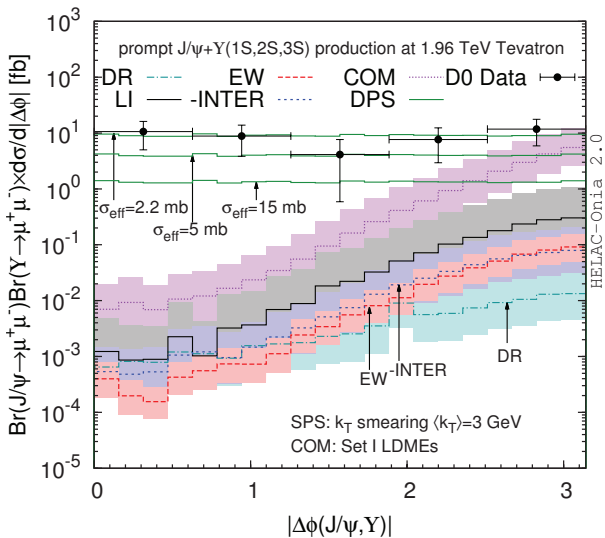
		J/ψ	$\psi(2S)$
DR	$\Upsilon(1S)$	$3.58^{+233\%}_{-66.4\%} \pm 4.4\%$	$2.34^{+233\%}_{-66.4\%} \pm 4.4\%$
	$\Upsilon(2S)$	$1.78^{+233\%}_{-66.4\%} \pm 4.4\%$	$1.17^{+233\%}_{-66.4\%} \pm 4.4\%$
	$\Upsilon(3S)$	$1.36^{+233\%}_{-66.4\%} \pm 4.4\%$	$0.894^{+233\%}_{-66.4\%} \pm 4.4\%$
LI	$\Upsilon(1S)$	$56.2^{+264\%}_{-70.2\%} \pm 4.7\%$	$36.8^{+264\%}_{-70.2\%} \pm 4.7\%$
	$\Upsilon(2S)$	$28.0^{+264\%}_{-70.2\%} \pm 4.7\%$	$18.4^{+264\%}_{-70.2\%} \pm 4.7\%$
	$\Upsilon(3S)$	$21.4^{+264\%}_{-70.2\%} \pm 4.7\%$	$14.0^{+264\%}_{-70.2\%} \pm 4.7\%$
EW	$\Upsilon(1S)$	$15.8^{+75.4\%}_{-46.4\%} \pm 4.6\%$	$10.4^{+75.4\%}_{-46.4\%} \pm 4.6\%$
	$\Upsilon(2S)$	$7.90^{+75.4\%}_{-46.4\%} \pm 4.6\%$	$5.18^{+75.4\%}_{-46.4\%} \pm 4.6\%$
	$\Upsilon(3S)$	$6.04^{+75.4\%}_{-46.4\%} \pm 4.6\%$	$3.96^{+75.4\%}_{-46.4\%} \pm 4.6\%$
INTER	$\Upsilon(1S)$	$-16.6^{+162\%}_{-62.0\%} \pm 4.8\%$	$-10.9^{+162\%}_{-62.0\%} \pm 4.8\%$
	$\Upsilon(2S)$	$-8.29^{+162\%}_{-62.0\%} \pm 4.8\%$	$-5.43^{+162\%}_{-62.0\%} \pm 4.8\%$
	$\Upsilon(3S)$	$-6.34^{+162\%}_{-62.0\%} \pm 4.8\%$	$-4.15^{+162\%}_{-62.0\%} \pm 4.8\%$
COM	$\Upsilon(1S)$	$409^{+138\%}_{-56.7\%} \pm 4.4\%$	$174^{+138\%}_{-56.8\%} \pm 4.4\%$
	$\Upsilon(2S)$	$135^{+139\%}_{-57.0\%} \pm 4.4\%$	$57.6^{+139\%}_{-57.1\%} \pm 4.4\%$
	$\Upsilon(3S)$	$197^{+137\%}_{-56.6\%} \pm 4.4\%$	$84.1^{+138\%}_{-56.7\%} \pm 4.4\%$

SPS cross sections @ D0 & LHCb

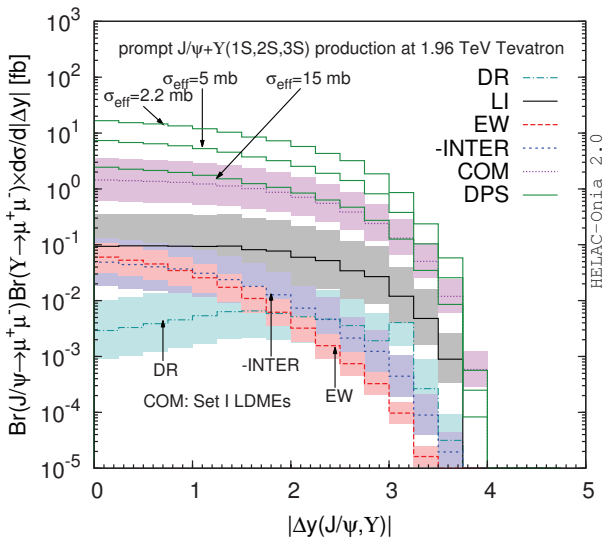
Experiment	CSM				COM			
	DR	LI	EW	INTER	Set I	Set II	Set III	Set IV
D0: $27 \pm 42.2\%$	$0.0146^{+233\%}_{-66.6\%}$	$0.229^{+264\%}_{-70.4\%}$	$0.065^{+75.5\%}_{-46.6\%}$	$-0.068^{+162\%}_{-62.2\%}$	$2.96^{+135\%}_{-56.2\%}$	$1.41^{+160\%}_{-77.6\%}$	$1.80^{+143\%}_{-58.0\%}$	$0.418^{+144\%}_{-58.3\%}$
LHCb	$0.255^{+391\%}_{-79.7\%}$	$6.05^{+436\%}_{-82.2\%}$	$1.71^{+135\%}_{-65.2\%}$	$-3.23^{+262\%}_{-75.9\%}$	$38.8^{+238\%}_{-73.0\%}$	$21.2^{+243\%}_{-73.6\%}$	$28.1^{+243\%}_{-73.8\%}$	$6.57^{+243\%}_{-73.9\%}$

TABLE III: Cross sections $\sigma(pp(\bar{p}) \rightarrow J/\psi\Upsilon) \times \text{Br}(J/\psi \rightarrow \mu^+\mu^-)\text{Br}(\Upsilon \rightarrow \mu^+\mu^-)$ (in units of fb) of prompt J/ψ and $\Upsilon(1S, 2S, 3S)$ simultaneous production at the Tevatron in the D0 fiducial region [10] and at $\sqrt{s} = 13$ TeV LHC in the LHCb acceptance $2 < y_{J/\psi, \Upsilon} < 4.5$, where we have also included feeddown contributions from higher-excited quarkonia decay.

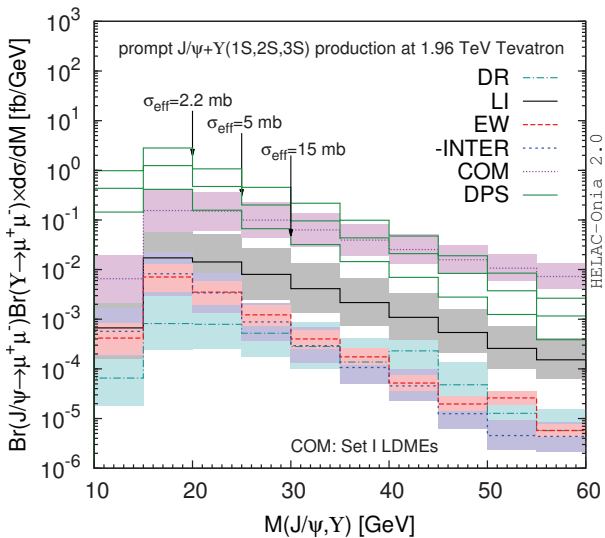
dphi @ D0



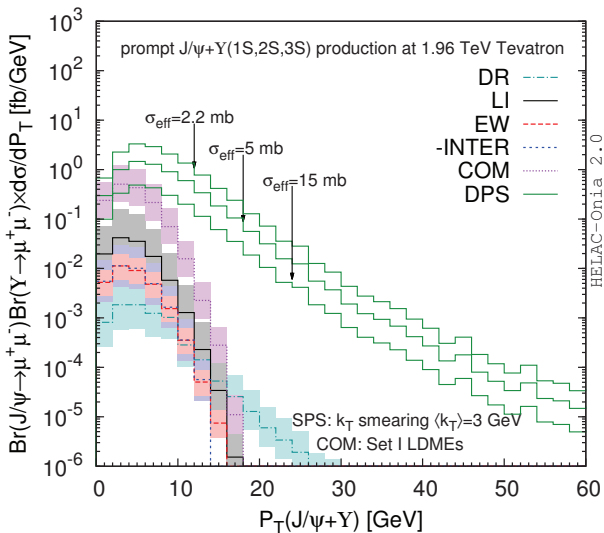
dy @ D0



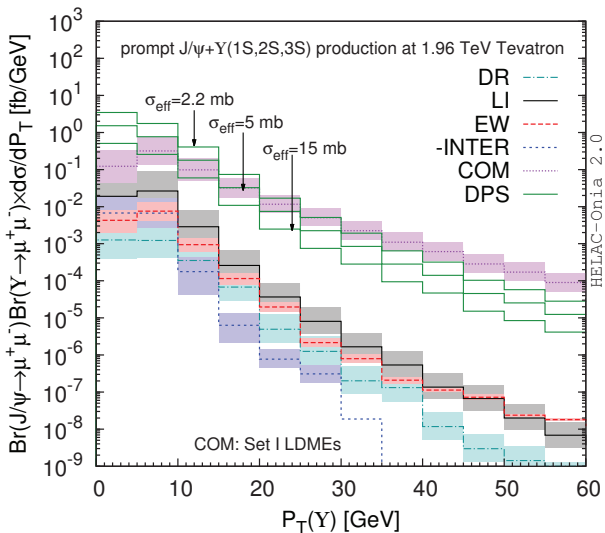
dM @ D0



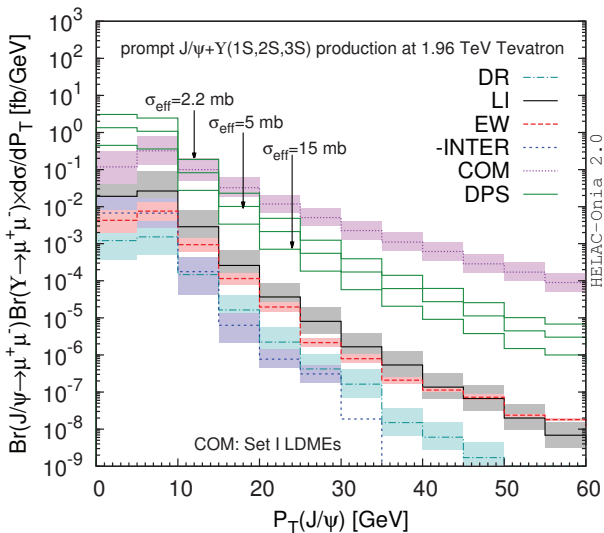
dPt @ D0



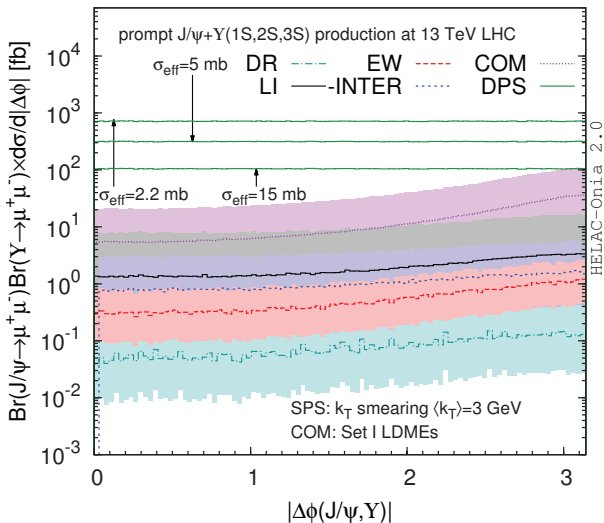
dptY @ D0



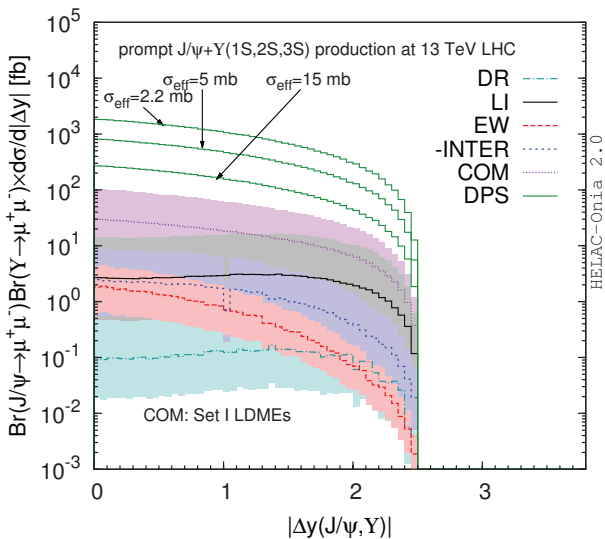
dptpsi @ D0



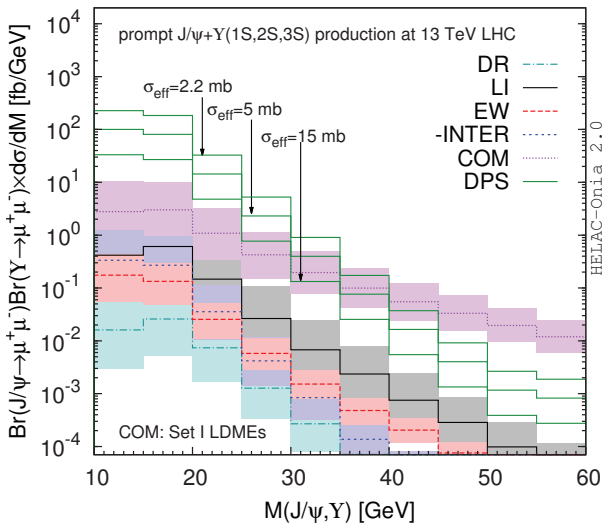
dphi @ LHCb



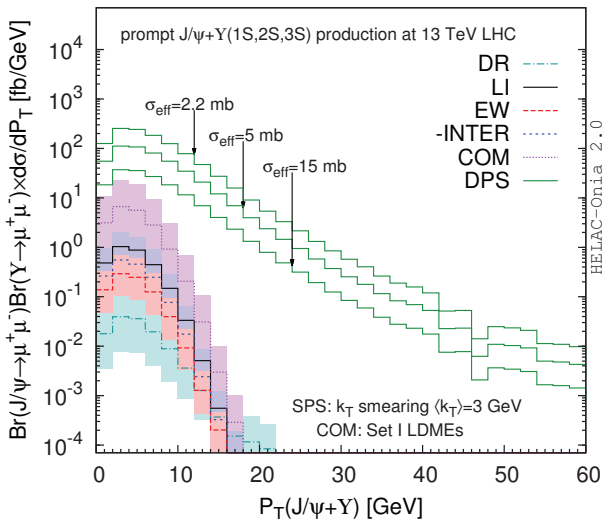
dy @ LHCb



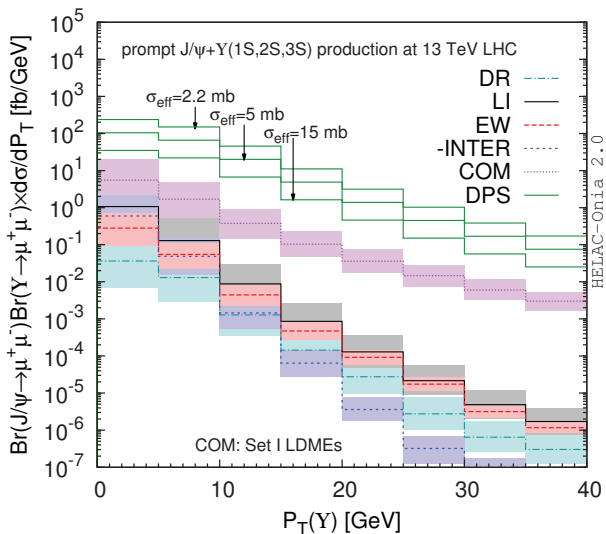
dM @ LHCB



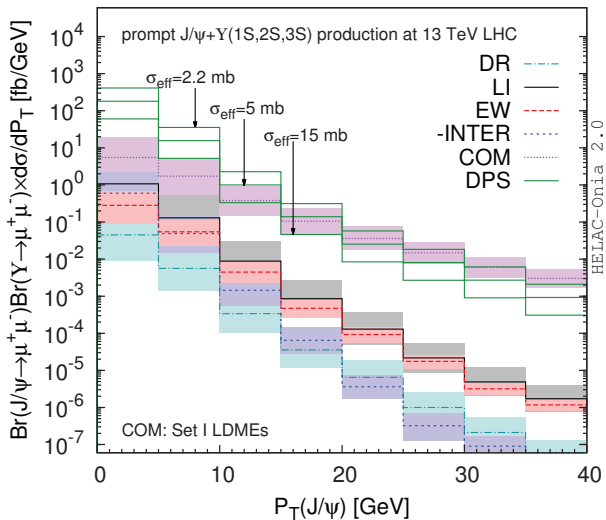
dPt @ LHCb



dptY @ LHCb



dptpsi @ LHCb



Numerical Result of $\Upsilon + J/\psi + \phi$

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SPS cross section of $\Upsilon, J/\psi, \phi$ at LHCb

- 1 We can get the inclusive cross sections of $\Upsilon, J/\psi, \phi$ at $\sqrt{s} = 13$ TeV at LHCb, it $0.2(15) \mu\text{b}$ for $\Upsilon(J/\psi)$, and the cross sections is 0.6 mb for $p_T(\phi) > 2 \text{ GeV}$.

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- 5 TPS $\Upsilon + J/\psi + \phi$: about $\frac{\sigma[\Upsilon]\sigma[J/\psi]\sigma[\phi]}{(\sigma_{\text{eff}}^{\text{TPS}})^2} \sim 28 \text{ pb}$ for $p_T(\phi) > 2 \text{ GeV}$ and $\sigma_{\text{eff}}^{\text{TPS}} \sim 8 \text{ mb}$.

Search for TPS in $\Upsilon + J/\psi + \phi$ at LHCb

Estimate the number of events

- 1 $Br[\Upsilon(J/\psi) \rightarrow \mu^+ \mu^-] = 0.024(0.06)$ and
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- 3 Number of events for $\Upsilon(\mu^+\mu^-) + J/\psi(\mu^+\mu^-) + \phi(K^+K^-)$ with $p_T(\phi) > 2 \text{ GeV}$ is about **80**.

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- 4 We can introduce cut to suppress SPS and DPS contributions.

Numerical Result for $\Upsilon + J/\psi + \phi$ at CMS/Atlas

SPS cross section of $\Upsilon, J/\psi, \phi$ at CMS/Atlas

- 1 We can get the inclusive cross sections of $\Upsilon, J/\psi, \phi$ at $\sqrt{s} = 13$ TeV at CMS/Atlas, it 0.4, 30, 1200 μb for $p_T(\phi) > 2$ GeV.

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- 2 SPS and DPS are very small.
- 3 $Br[\Upsilon(J/\psi) \rightarrow \mu^+\mu^-] = 0.024(0.06)$ and $Br[\phi \rightarrow K^+K^-] = 0.5$.
- 5 Integrated luminosity of LHCb is about 40 fb^{-1} at $\sqrt{s} = 13$ TeV.

Numerical Result for $\Upsilon + J/\psi + \phi$ at CMS/Atlas

SPS cross section of Υ , J/ψ , ϕ at CMS/Atlas

- 1 We can get the inclusive cross sections of Υ , J/ψ , ϕ at $\sqrt{s} = 13$ TeV at CMS/Atlas, it 0.4, 30, 1200 μb for $p_T(\phi) > 2$ GeV.
- 2 SPS and DPS are very small.
- 3 TPS $\Upsilon + J/\psi + \phi$: about $\frac{\sigma[\Upsilon]\sigma[J/\psi]\sigma[\phi]}{(\sigma_{\text{eff}}^{\text{TPS}})^2} \sim 200$ pb for $p_T(\phi) > 2$ GeV and $\sigma_{\text{eff}}^{\text{TPS}} \sim 8$ mb.
- 4 $Br[\Upsilon(J/\psi) \rightarrow \mu^+\mu^-] = 0.024(0.06)$ and $Br[\phi \rightarrow K^+K^-] = 0.5$.
- 5 Integrated luminosity of LHCb is about 40 fb^{-1} at $\sqrt{s} = 13$ TeV.
- 6 Number of events for $\Upsilon(\mu^+\mu^-) + J/\psi(\mu^+\mu^-) + \phi(K^+K^-)$ with $p_T(\phi) > 2$ GeV is about **6000**.

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Summary

We have performed the first complete analysis of simultaneous production of prompt ψ and Υ mesons including all leading SPS contributions.

Our work shows that it is in fact most probably dominated by DPS contributions for D0 data.

Finally, we show that $\Upsilon + J/\psi + \phi$ at LHC is dominated by TPS. It may be studied by experimenters.