



X(5568): understandings and status

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- Experimental status
- Theoretical understandings:
 - (1) molecule picture
 - (2) tetraquark picture
 - (3) more discussions
- Spectra in CMI model
- Production problem
- Puzzles to be answered
- Summary

Hadrons with exotic properties

- 2003: X(3872), D_{sJ}(2317), D_{sJ}(2460), ~~Θ(1540)~~
- 2004: X(3915), ~~D_{sJ}(2632)~~
- 2005: Y(4260)
- 2007: Y(4008), Y(4360), Y(4660), X(4160), Z(4430)
- 2008: Z₁(4050), Z₂(4250)
- 2009: X(4140)
- 2010: X(4274), X(4350), X(4315)?
- 2011: Z_b(10610), Z_b(10650)
- 2013: Z_c(3900), Z_c(3885), Z_c(4020), Z_c(4025)
- 2014: Z_c(4200)
- 2015: P_c(4380), P_c(4450)
- 2016: X(4500), X(4700), X(5568) ?

Narrow X(5568)

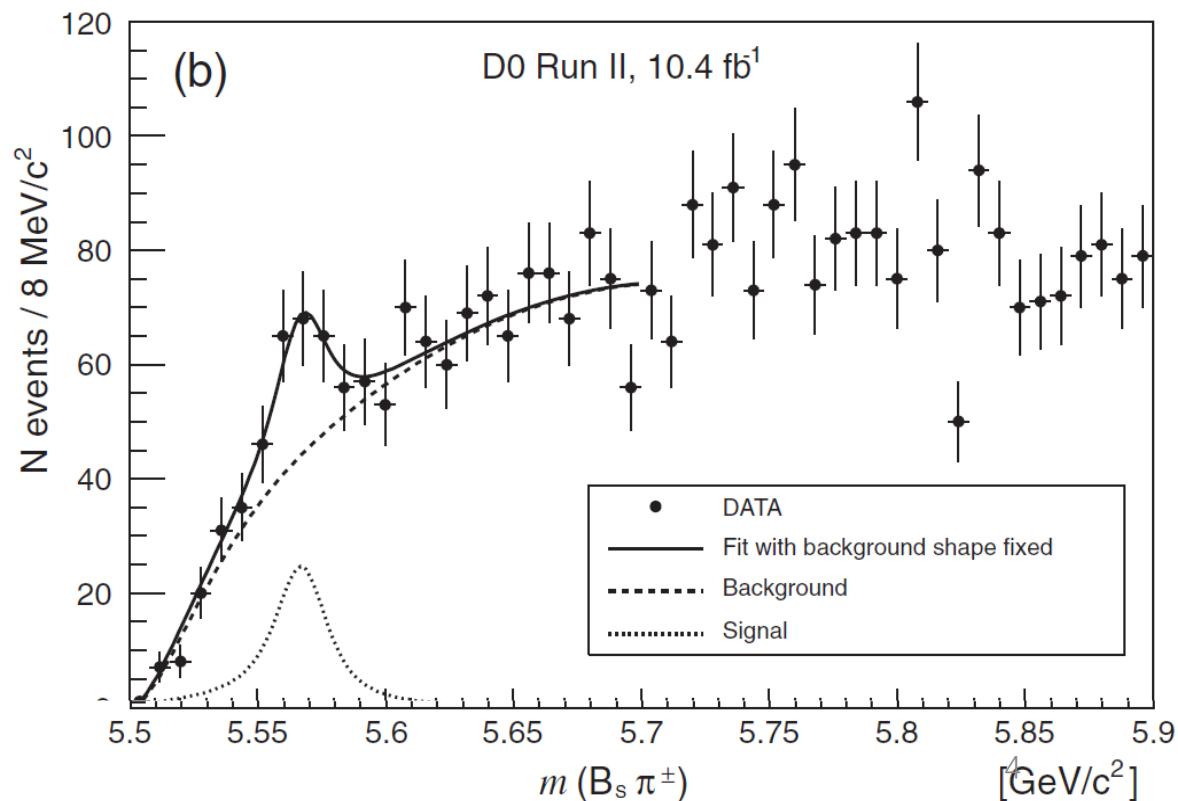
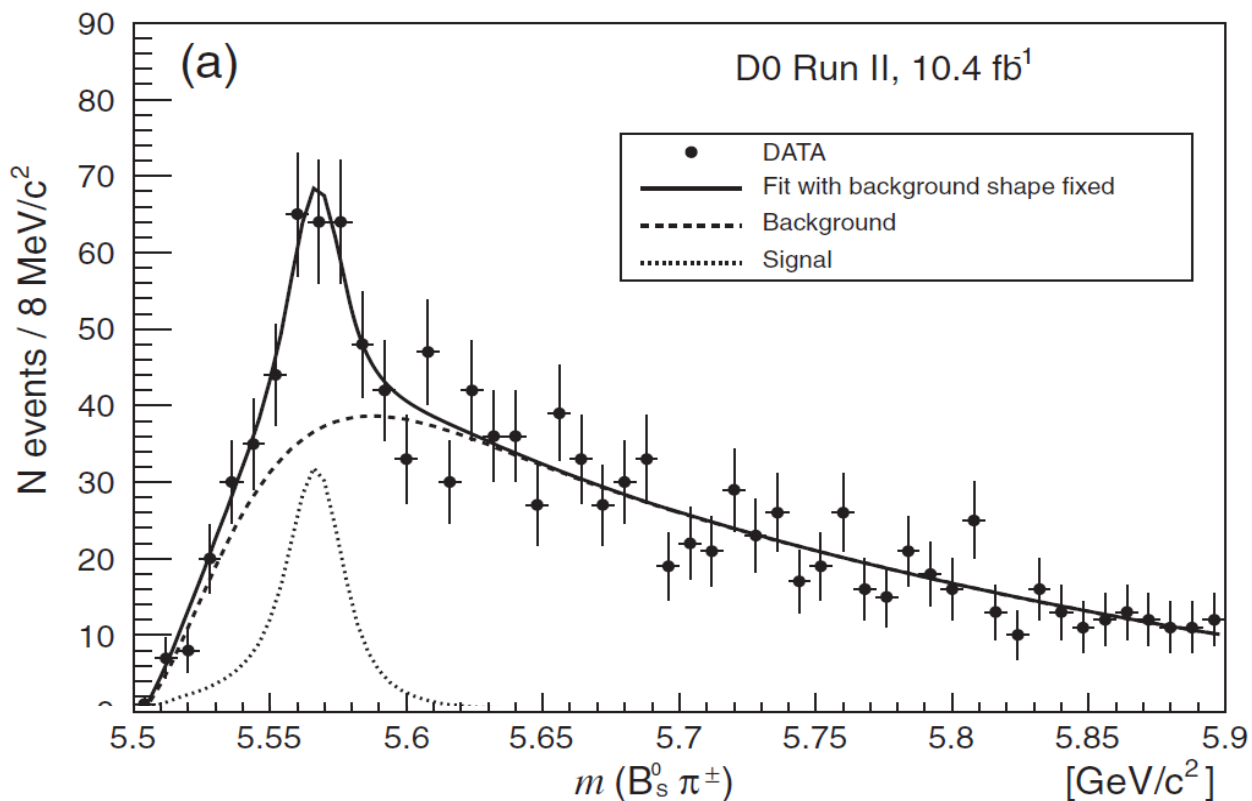
$$J^P = 0^+$$

$$J^P = 1^+$$

- 1602.07588 (D0): $X(5568) \rightarrow B_s^0 \pi^\pm$ [$X(5568 + 49) \rightarrow B_s^* \pi^\pm \rightarrow B_s^0 \gamma \pi^\pm$]

$$m = 5567.8 \pm 2.9_{-1.9}^{+0.9} \text{ MeV} \quad \Gamma = 21.9 \pm 6.4_{-2.5}^{+5.0} \text{ MeV}$$

- Four different flavors: $b\bar{s}u\bar{d}$ or $b\bar{s}\bar{u}d$



Narrow X(5568)

- 1608.00435 (LHCb): No significant excess is found
- 1712.06144 (CMS): No statistically significant peaks
“central kinematic region of B_s^0 candidates similar to D0”
- 1712.09620 (CDF): No evidence for this state is found
“the same collision conditions and similar kinematic range as in the original observation of this state by D0”
- 1712.10176 (D0): Evidence for a narrow structure, again
 $B_s^0 \rightarrow J/\psi_{(\rightarrow\mu^+\mu^-)} + \phi_{(\rightarrow K^+K^-)}$
 $B_s^0 \rightarrow \mu^\mp + D_s^\pm_{(\rightarrow\phi\pi^\pm)} + X$
 $m = 5566.9_{-3.1}^{+3.2}{}_{-1.2}^{+0.6} \text{ MeV}$
 $\Gamma = 18.6_{-6.1}^{+7.9}{}_{-3.8}^{+3.5} \text{ MeV}$
- 1802.01840 (ATLAS): No significant signal was found

Narrow $X(5568)$: cone cut/line shape

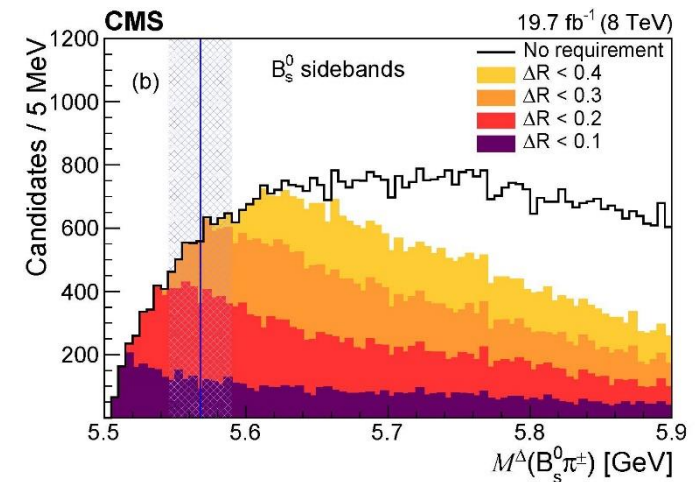
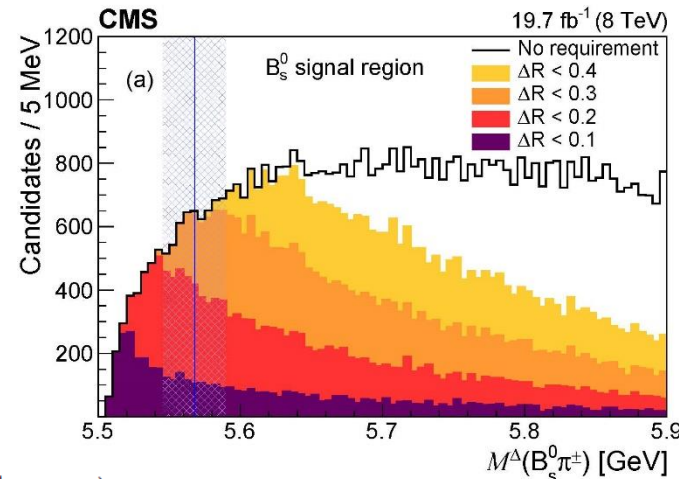
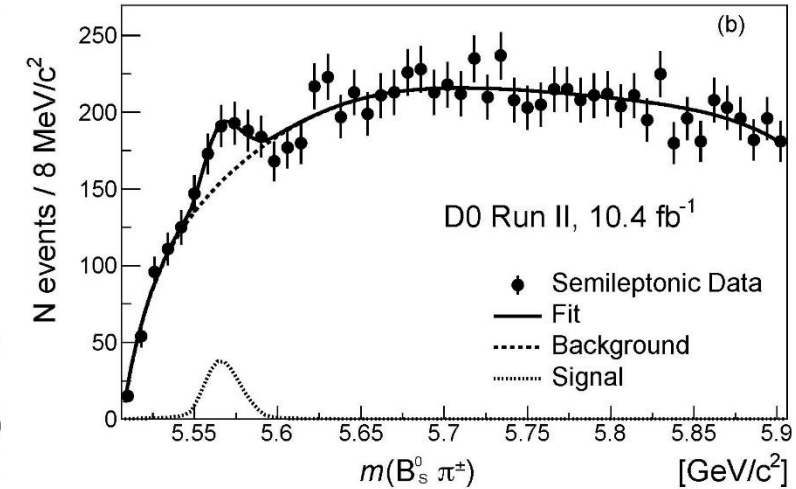
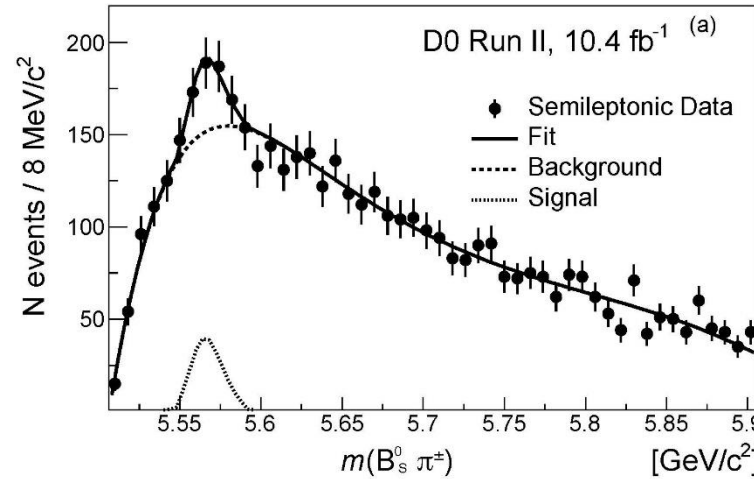
$$\Delta R = \sqrt{\Delta\eta^2 + \Delta\phi^2} < 0.3$$

η : pseudorapidity
 ϕ : azimuthal angle

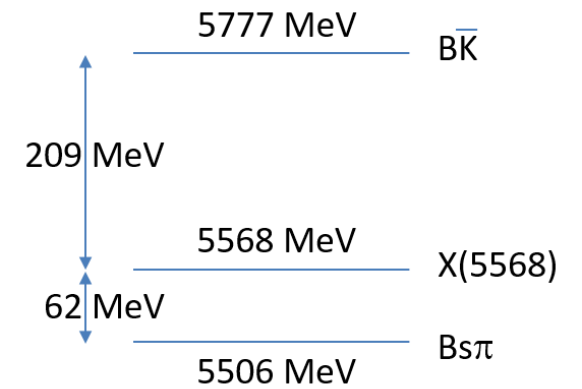
D0: w./w.o. cone cut: 5.1/3.9 σ

CMS: cone cut changes the $M^\Delta(B_s^0\pi^\pm)$ distribution and can even produce a peaking shape. Not used

$$M^\Delta(B_s^0\pi^\pm) = M(J/\psi K^+ K^- \pi^\pm) - M(J/\psi K^+ K^-) + m_{B_s^0}$$



X(5568): molecule?



- Agaev, Azizi, Sundu, [1603.02708](#) [QSR, $0^+ B\bar{K}$]

$$J^{X_b}(x) = \left[\bar{d}^a(x) \gamma_5 s^a(x) \right] \left[\bar{b}^b(x) \gamma_5 u^b(x) \right]$$

⇒ mass, decay constant, width: **NO**

- C.J. Xiao, D.Y. Chen, [1603.00228](#) [assume]

⇒ decay: **unnatural** molecule

- Albaladejo, Nieves, Oset, Z.F. Sun, X. Liu, [1603.09230](#) [CC]

⇒ mass, width: **NO**

- X.W. Kang, Oller, [1606.06665](#) [EFT: P-wave 1^-]

⇒ **NO**

X(5568): molecule?

- Lang, Mohler, Prelovsek, [1607.03185](#) [lat.: $0^+ B_s \pi - B \bar{K}$]
⇒ scattering: **NO**
- R, Chen, X. Liu, [1607.05566](#) [OBE: $0^+, 1^+$]
⇒ mass: **NO**
- J.X. Lu, X.L. Ren, L.S. Geng, [1607.06327](#) [sym., fin.vol. χ PT, comp.]
⇒ **NO**
- B.X. Sun, F.Y. Dong, J.L. Ping, [1609.04068](#) [0^+ CC BSE]
⇒ pole: **OK**
- Z.Y. Wang, J.J. Qi, C. Wang, X.H. Guo, [1802.04566](#) [BSE: 0^+]
⇒ mass, width: **OK**
- H.W. Ke, X.Q. Li, [1801.00675](#) [assume 0^+ in LFQM]; [1802.08823](#) [mole-tetra]
⇒ decay: **OK**

X(5568): tetraquark?

- Agaev, Azizi, Sundu, [1602.08642](#), [1603.00290](#) [QSR: diquark]
 $0^{++}: J^{X_b}(x) = \varepsilon^{ijk} \varepsilon^{imn} [s^j(x) C \gamma_\mu u^k(x)] [\bar{b}^m(x) \gamma^\mu C \bar{d}^n(x)]$ (5584±137) MeV
⇒ mass, decay constant, width: OK [not 1^+ in [1608.04785](#)]
- W. Chen, H.X. Chen, X. Liu, T.G. Steele, S.L. Zhu, [1602.08916](#) [QSR]
 $0^+:$ $J_3 = s_a^T C \gamma_5 u_b (\bar{b}_a \gamma_5 C \bar{d}_b^T - \bar{b}_b \gamma_5 C \bar{d}_a^T)$ 5.58±0.14 GeV
 $1^+:$ $J_{3\mu} = s_a^T C \gamma_5 u_b (\bar{b}_a \gamma_\mu C \bar{d}_b^T - \bar{b}_b \gamma_\mu C \bar{d}_a^T)$ 5.59±0.15 GeV
⇒ mass: OK
- Z.G. Wang, [1602.08711](#), [1603.02498](#) [QSR, diquark]
 $0^+:$ $J(x) = \varepsilon^{ijk} \varepsilon^{imn} u^j(x) C \gamma_5 s^k(x) \bar{d}^m(x) \gamma_5 C \bar{b}^n(x)$ (5.57±0.12) GeV
⇒ mass,width: OK

X(5568): tetraquark?

- M. Nielsen et. al., [1602.09041](#), [1603.02249](#) [QSR]
 0^+ : $j_S = \epsilon_{abc}\epsilon_{dec}(u_a^T C\gamma_5 s_b)(\bar{d}_d\gamma_5 C\bar{b}_e^T) \quad \sim 6.4 \text{ GeV}$
 \Rightarrow mass: **NO**, width: OK
- L. Tang, C.F. Qiao, [1603.04761](#) [QSR, 8c*8c]
 0^+ : $j_A(x) = [i\bar{b}^j(x)\gamma_5(t^a)_{jk}s^k(x)][i\bar{d}^m(x)\gamma_5(t^a)_{mn}u^n(x)] \quad 5.57\pm 0.15 \text{ GeV}$
 0^+ : $j_C(x) = [\bar{b}^j(x)\gamma^\mu(t^a)_{jk}s^k(x)][\bar{d}^m(x)\gamma_\mu(t^a)_{mn}u^n(x)] \quad 5.58\pm 0.15 \text{ GeV}$
 \Rightarrow mass: **OK**
- J.R. Zhang, J.L. Zou, J.Y. Wu, [1703.09043](#) [QSR: 0^+ diquark-diquark]
scalar-scalar: $5.57_{-0.23}^{+0.35} \text{ GeV}$; axial-axial: $5.77_{-0.33}^{+0.44} \text{ GeV}$
 \Rightarrow mass: **OK**

X(5568): tetraquark?

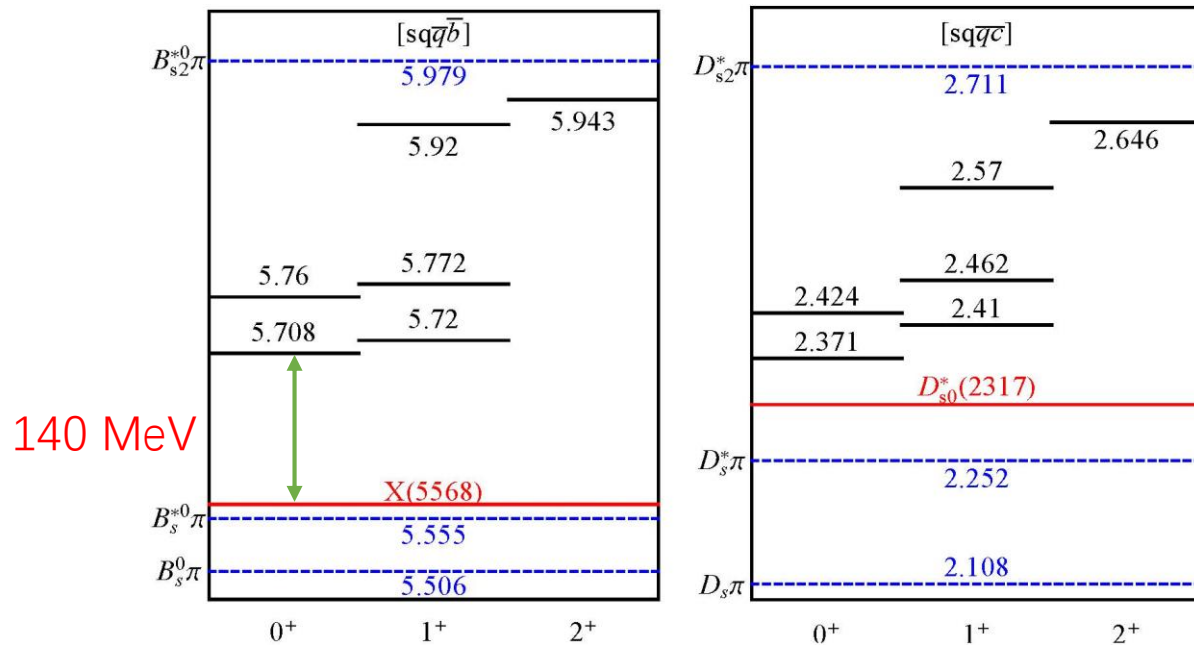
- Albuquerque, Narison, Rabemananjara, Rabetiarivony, [1604.05566](#) [QSR]

| <i>Molecule</i> | | | <i>Molecule</i> | | |
|-------------------|-------|--|---|-------|---------------|
| BK | 0^+ | $(\bar{b} i\gamma_5 u)(\bar{d} i\gamma_5 s)$ | B^*K | 1^+ | 5186 ± 13 |
| $B_s\pi$ | 0^+ | $(\bar{b} i\gamma_5 s)(\bar{d} i\gamma_5 u)$ | BK | 0^+ | 5195 ± 15 |
| B^*K | 1^+ | $(\bar{b} i\gamma_\mu u)(\bar{d} i\gamma_5 s)$ | $B_s^*\pi$ | 1^+ | 5200 ± 18 |
| $B_s^*\pi$ | 1^+ | $(\bar{b} i\gamma_\mu s)(\bar{d} i\gamma_5 u)$ | $B_s\pi$ | 0^+ | 5199 ± 24 |
| <i>Four-quark</i> | | | <i>Four-quark (su)($\bar{b}\bar{d}$)</i> | | |
| | 1^- | $(s^T C \gamma_5 u)(\bar{b} \gamma_\mu \gamma_5 C \bar{d}^T) + k(s^T C u)(\bar{b} \gamma_\mu C \bar{d}^T)$ | A_b | 1^+ | 5186 ± 16 |
| | 1^+ | $(s^T C \gamma_5 u)(\bar{b} \gamma_\mu C \bar{d}^T) + k(s^T C u)(\bar{b} \gamma_\mu \gamma_5 C \bar{d}^T)$ | S_b | 0^+ | 5196 ± 17 |

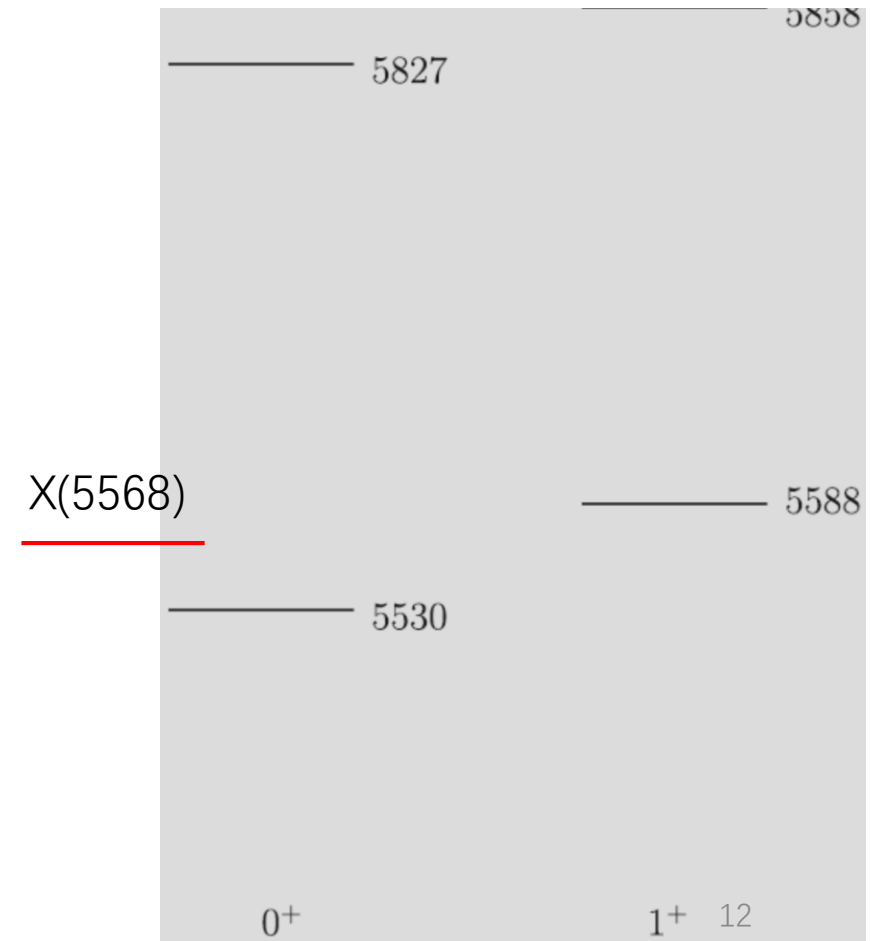
\Rightarrow **not** pure molecule or 4-quark state, can **mix** $B\bar{K}$ & S_b ($\sin 2\theta=0.15$)

X(5568): tetraquark?

- W. Wang, R.L. Zhu, [1602.08806](#) [CMI: diquark]
 \Rightarrow mass: **not favored**



- Fl. Stancu, [1603.03322](#) [CMI]
 \Rightarrow mass: **OK**
- Y.R.Liu, X.Liu, S.L.Zhu, [1603.01131](#) [CMI]



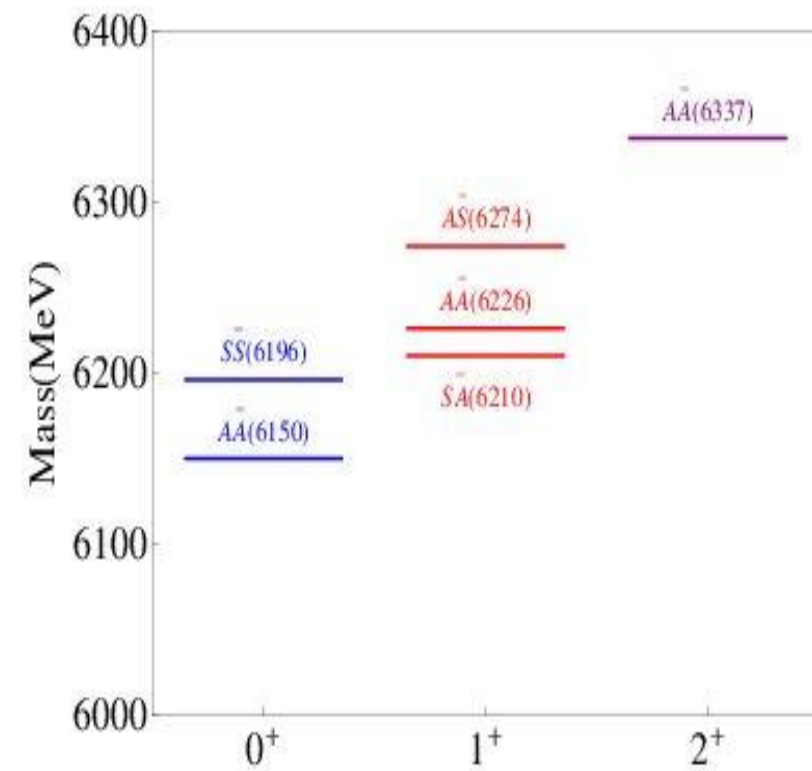
X(5568): tetraquark?

- Q.F. Lv, Y.B. Dong, [1603.06316](#) [GI model: diquark]
⇒ mass: **NO**
- Espositoa, Pilloni, Polosa, [1603.07667](#)
[hybridized tetraquark]
⇒ mass: **NO**
- Ali, Maiani, Polosa, Riquer, [1604.01731](#) [diquark]
⇒ mass: **NO**
- X.Y. Chen, J.L. Ping, [1604.05651](#) [CQM, 0^+]

| $qq - \bar{q}\bar{q}$ | $E_{\bar{3} \otimes 3}$ | $E_{6 \otimes \bar{6}}$ | E_{cc} |
|-----------------------|-------------------------|-------------------------|----------|
| $sud\bar{b}$ | 6397.6 | 6466.4 | 6351.0 |

⇒ mass: diquark **too high**, **no** mole

- Goerke et.al., [1608.04656](#) [4-quark]
⇒ can fit width with 5568 or 5771



X(5568): more discussions

- X.H. Liu, G. Li, [1603.00708](#) [triangle singularity]
⇒ not rescattering effects
- Y. Jin, S.Y. Li, S.Q. Li, [1603.03250](#) [production]
⇒ difficult to understand large production rate
- T.J. Burns, E.S. Swanson, [1603.04366](#) [th., cusp, mole., tetra.]
“none of the interpretations seems a natural fit for X(5568)”
- F.K. Guo, U.-G. Meissner, B.S. Zou, [1603.06316](#) [tetra,mole,th.]
⇒ NO
- Z. Yang, Q. Wang, U.-G. Meißner, [1609.08807](#)
⇒ kinematic reflection

X(5568): more related discussions

- X.G. He, P. Ko, [1603.02915](#) [EFT]: partners, decays, mass sum rules
- X.H. Liu, Ulf-G. Meißner, [1703.09043](#) [TS]
“Generating a resonance-like structure in the reaction $B_c \rightarrow B_s \pi \pi$ ”
- W. Chen, H.X. Chen, X. Liu, T.G. Steele, S.L. Zhu, [1705.10088](#) [QSR]
 $\Rightarrow sq\bar{q}\bar{c}, qq\bar{q}\bar{c}, sq\bar{q}\bar{b}, qq\bar{q}\bar{b}$
- F.S. Yu, [1709.02571](#)
 \Rightarrow detectable weakly decaying $bs\bar{u}\bar{d}$ [X(5568): $b\bar{s}u\bar{d}$ or $b\bar{s}\bar{u}d$]

Content

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 - (1) molecule picture
 - (2) tetraquark picture
 - (3) more discussions
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- Production problem
- Puzzles to be answered
- Summary

} ~40

HEP 找到 1 笔记录

1. **Evidence for a $B_s^0\pi^\pm$ state**
D0 Collaboration (V.M. Abazov (Dubna, JINR) *et al.*). Feb 24, 2016. 8 pp.
Published in *Phys.Rev.Lett.* **117** (2016) no.2, 022003
FERMILAB-PUB-16-038-E
DOI: [10.1103/PhysRevLett.117.022003](https://doi.org/10.1103/PhysRevLett.117.022003)
e-Print: [arXiv:1602.07588](https://arxiv.org/abs/1602.07588) [hep-ex] | [PDF](#)
[References](#) | [BibTeX](#) | [LaTeX\(US\)](#) | [LaTeX\(EU\)](#) | [Harvmac](#) | [EndNote](#)
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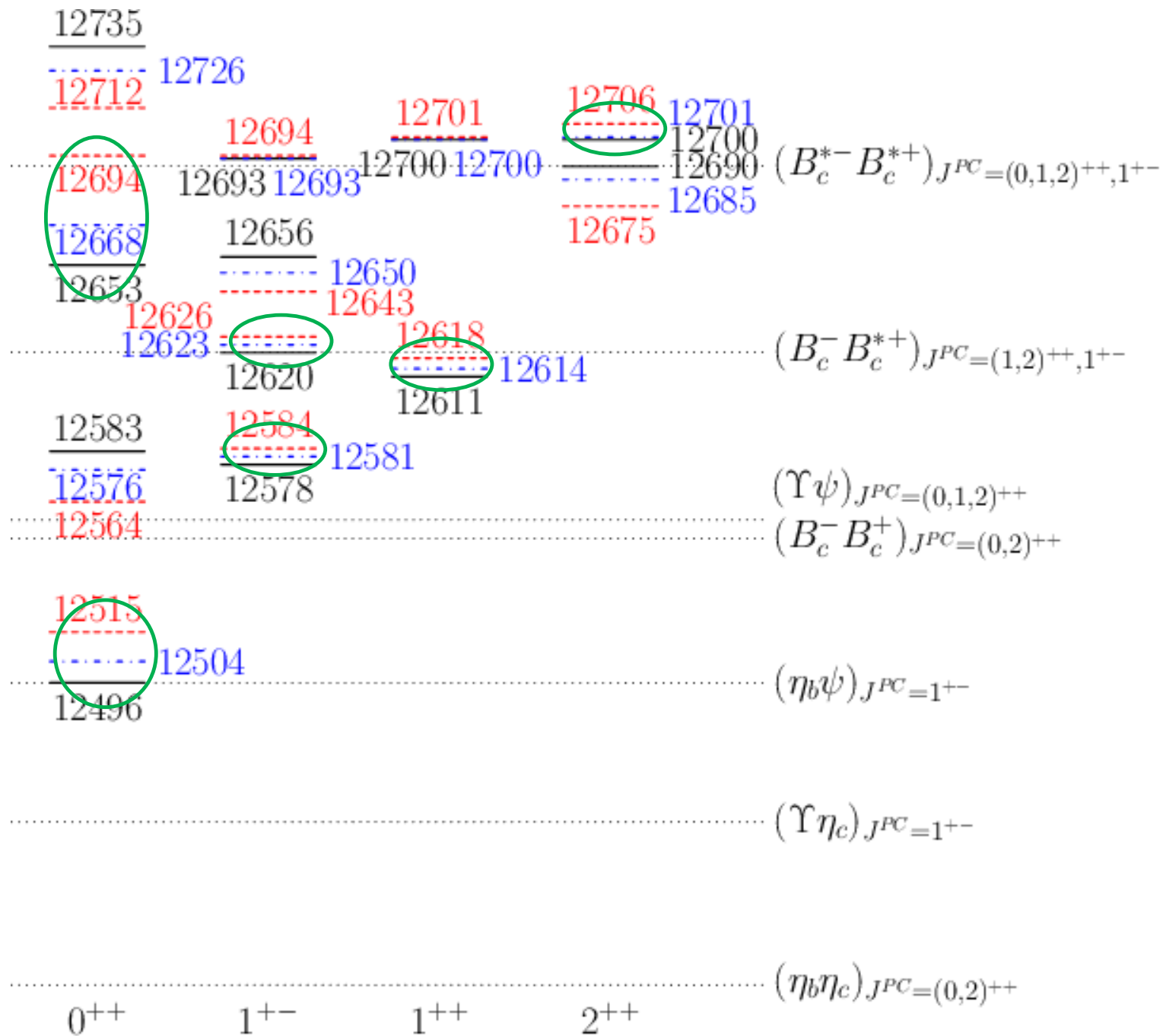
Decay ?

Spectra in CMI model

- $H = \sum_i m_i^{eff} + \sum_{i<j} C_{ij} \lambda_i \cdot \lambda_j \sigma_i \cdot \sigma_j$
simple, basic features, mass differences
- Stancu, 1603.03322:
“difficulty of extracting effective quark masses, from mesons and baryons, to be used in multiquark systems” ,
“large value of $C_{u\bar{d}}$ and **color-mixing** play key role”
- Even can be used: mass **uncertainty**?
- Ours: “X(5568)&partners” where low quark mass to recover DsJ(2632)?
- **Improvements** in applying the model to multiquark states!

Spectra in CMI model

- Our strategy:
 - (1) Important color-mixing (basis independent results?)
 - (2) $M = \sum_i m_i^{eff} + \langle H_{CMI} \rangle$ as upper limit
 - (3) $M = M_{ref} - \langle H_{CMI} \rangle_{ref} + \langle H_{CMI} \rangle$, find lower limit and roughly **reasonable masses**
 - (4) Effective force within diquark: **stable or narrow** multiquarks
- With these considerations, non-existence of X(5568) is expected.



Even if X(5568) does not exist, more interesting exotics are worthwhile search for

[Y.R. Liu, X. Liu, S.L. Zhu, 1603.01131 and following CMI studies]

Production problem [Y. Jin, S.Y. Li, 1603.03250]

- Most general hadrons can be produced both in **decay and directly** in multiproduction process of high energy collisions
- **Exotic** hadrons produced in decays of heavier hadrons, but **difficult in direct production in multiproduction processes**
- High mass: X(5568) in decay hard.
- Production in D0: multiproduction process
production mechanism must be quite special if exist.

Difficulty of production of multiquark states (**Unitarity**)

ALL quarks go into hadrons (confinement--Tous les hommes sont mortels !)

Formal descriptions: by a unitary time evolution operator

$$\sum_h |\langle h|U|q \rangle|^2 = \langle q|U^\dagger U|q \rangle = \sum |q \rangle \langle q| = \sum |h \rangle \langle h| = 1 \quad \sum_{h=B,B,M} |\langle h|U|q \rangle|^2 \sim 1 - \varepsilon, \varepsilon \rightarrow 0^+$$

Total probability to become any state of hadron system is 1

'Reversal employment' of the Gell-Mann–Zweig constituent quark model, where all the quark states and the hadron states are different bases of the *same* Hilbert space of states/static \rightarrow production

Most of the limited space has been occupied by general hadrons, not any more left for the populated exotics---they really of large population even only considering the heavy

See, Unitarity and Entropy Change in Exclusive Quark Combination Models, arXiv:1005.4664 [hep-ph] ; Exotic hadron production in quark combination model, Phys.Rev. C80 (2009) 035202; talks here at Moriond QCD'11(unitarity and exotic) , and some recent ones to repeat this observation

Difficulty of production of multiquark states (**colour connection**)

- **Complexity of the colour structure of the multiquark hadrons**
- **Eg, qqbar qqbar four quark system**

$$(3_1 \otimes 3_3) \otimes (3_2^* \otimes 3_4^*) = (3_{13}^* \oplus 6_{13}) \otimes (3_{24} \oplus 6_{24}^*) = (3_{13}^* \otimes 3_{24}) \oplus (6_{13} \otimes 6_{24}^*) \oplus \dots$$

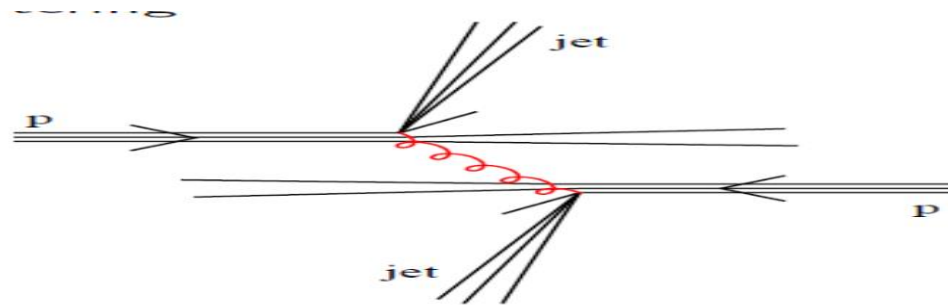
$$(3_1 \otimes 3_2^*) \otimes (3_3 \otimes 3_4^*) = (1_{12} \oplus 8_{12}) \otimes (1_{34} \oplus 8_{34}) = (1_{12} \otimes 1_{34}) \oplus (8_{12} \otimes 8_{34}) \oplus \dots,$$

$$(3_1 \otimes 3_4^*) \otimes (3_3 \otimes 3_2^*) = (1_{14} \oplus 8_{14}) \otimes (1_{32} \oplus 8_{32}) = (1_{14} \otimes 1_{32}) \oplus (8_{14} \otimes 8_{32}) \oplus \dots,$$

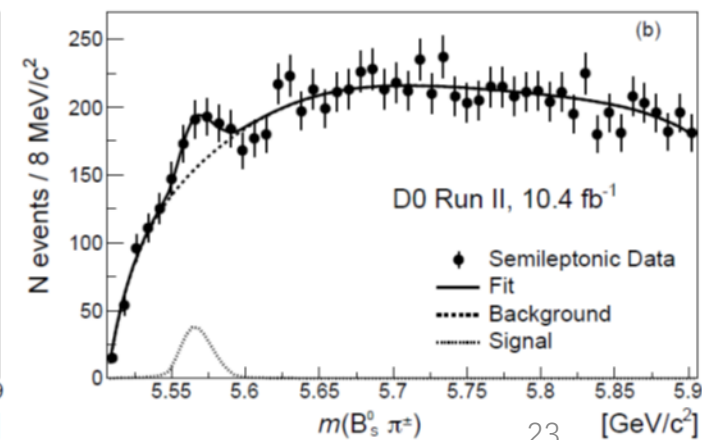
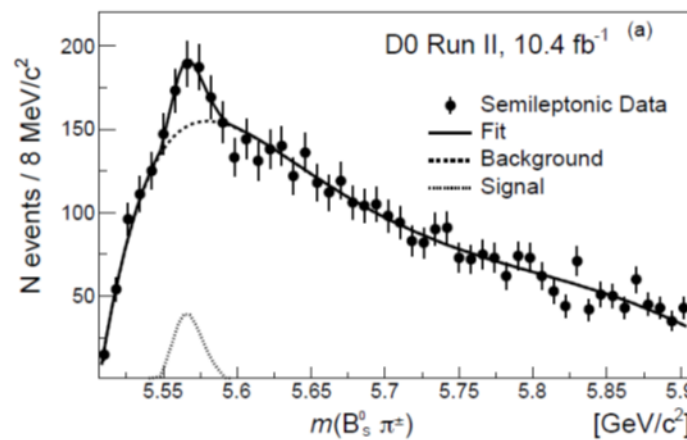
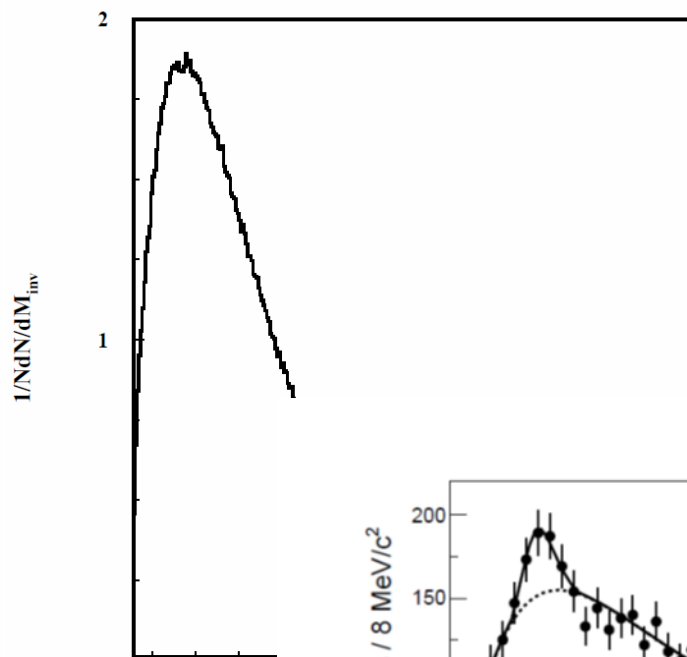
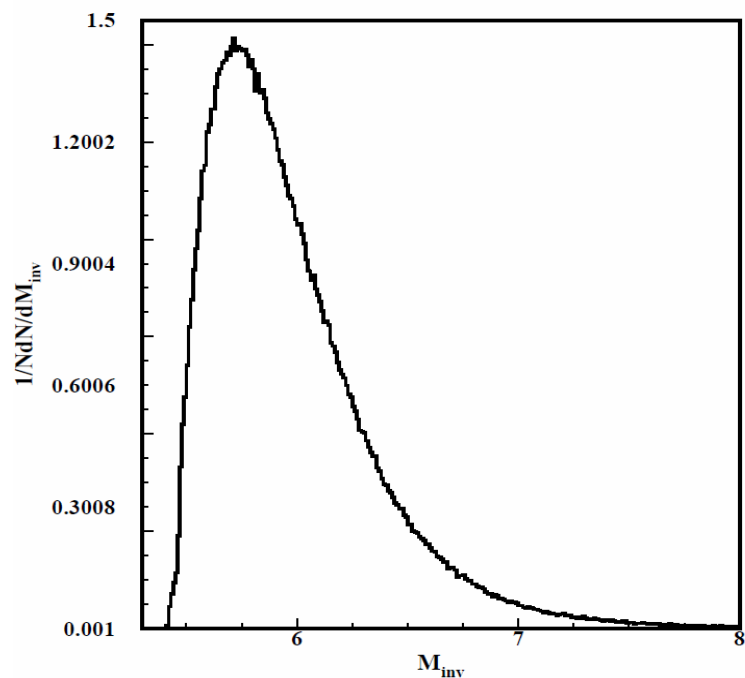
- **--->N-body problem? sub-clusters? This is a problem both in hadron (static properties) as well as in production**

- Colour reconnection, colour recombination..., total probability again 1
- Probability of rare connection is small
- See, e.g, Studying color connection effects of $e^+e^- \rightarrow c\bar{c}c\bar{c} \rightarrow \Xi c\bar{c} + X$ process within Quark Combination Model, P.R. D91 (2015) no.11, 114017, Search for a doubly charmed hadron at B factories, P.R.D89 (2014) no.9, 094006, Colour connections of four quark $QQ^+Q^-Q^-$ system and doubly heavy baryon production in $e^+e^-e^+e^-$ annihilation, P.L.B727 (2013) 468-473

- $e+e- \rightarrow \text{jet} + \text{jet},$



Line shape: peak around threshold, I-Bpi,R-Dpi



Puzzles to be answered

- Is $X(5568)$ a resonance?

If Yes: $D_{sJ}(2632)$? other exotic hadrons?

production mechanism? why not in other experiments?

is there any decisive channel?

similar to Ξ_{cc} (SELEX: 3519 MeV LHCb: 3621.40 MeV)?

...

If No: why is there a signal at D0? How to apply theoretical models?

[Recall $\Theta(1540)$ and $D_{sJ}(2632)$.]

- Whether $X(5568)$ exists or not:

inconsistency in theories? challenging methods? need improvements? ...

Summary

- Most studies disfavor molecule, tetraquark, or non-resonance interpretations for $X(5568)$.
- Inconsistent results (ex. & th.) need to be understood.
- $X(5568)$ motivates other exotic multiquarks which are worthwhile search in hadron decays or e^+e^- collisions.

Thank you very much!