



基于CDEX-1实验的亚GeV暗物质探 测与调制效应研究

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Outline

Introduction of CDEX-1B experiment

> Annual modulation analysis

Detection of sub-GeV dark matter





Elastic Scattering



Lower Background
Lower Energy threshold
Long-time stability

- ➤ 1 kg-scale-mass HPGe detector, cooled by cold finger.
- ➤ A NaI(Tl) detector is used as active shielding to veto the gamma-ray induced background events.
- The detector has been under stable data taking conditions since March 27th, 2014.





RUN-1













➤ The data used for the TI analysis is from March 2014 to July 2017, with a total exposure of 737.1 kg·day.
 ➤ The energy threshold is 160 eVee, with a combined efficiency of 7019/12/17%.

- Exclusion plots of spin-independent and spin-dependent χ-N scattering at 90% confidence level. (derived by binned poisson method)
 - L. T. Yang et al., Chin. Phys. C 42, 23002 (2018)









 $v_E = 232 + 30 \times 0.51 \cos(2\pi/T \times (t - \Phi)) \text{ km/s}$

Lower Background

Long-time stability

t = 0 h

- AM provide smoking-gun signatures for WIMPs independent of background modeling, while only requires background at relevant energy range is stable with time;
- The expected χN rates have distinctive AM features with maximum intensity in June and a period of one year due to Earth's motion relative to the galactic WIMP-halo.





Compton contribution from High-energy gammas

L-Shell X-ray contribution – time varying



days,

and 428.1 live

Best-fit of modulation amplitude w/ phase=152.5day SI Limits at 90% C.L. from AM



✓ CDEX-1B excludes DAMA/LIBRA phase-1's interpretation with the spin-independent WIMP interaction with Standard Halo model in Germanium crystal.

PRL 123.221301 (2019)



➢ In dark matter direct detection experiments, it is usually assumed that the extranuclear electrons around the recoil nucleus immediately follow the motion of the nucleus.



✓ nuclear recoil without atomic excitation $\chi + N \rightarrow \chi + N(E_R)$

The recoil "atom" loses its energy through scattering with adjacent atoms in the medium where the inelastic scatterings involve the ionization and excitation of the atoms.

➢ However, it takes some time for the electrons to catch up, which causes ionization and excitation of the recoil atom. The ionization and the excitation result in extra electronic energy injections into the detectors.



χ-N inelastic scattering

M.J. Dolan, et al. Phys. Rev. Lett. 121, 101801 (2018)

Single electron transition rates:

$$\sum_{F} |Z_{FI}|^{2} = |Z_{II}|^{2} + \sum_{n,\ell,n',\ell'} p_{q_{e}}^{d} (n\ell \to n'\ell') + \sum_{n,\ell} \int \frac{dE_{e}}{2\pi} \frac{d}{dE_{e}} p_{q_{e}}^{c} (n\ell \to E_{e})$$

excitation ionization

The excitation probabilities are much smaller than the ionization probabilities for a given initial n.

$$\blacktriangleright \text{ The event rates:} \qquad \frac{dR}{dE_R dE_e dv_{DM}} \simeq \frac{dR_0}{dE_R dv_{DM}} \times \frac{1}{2\pi} \sum_{n,\ell} \frac{d}{dE_e} p_{q_e}^c (n\ell \to E_e)$$
$$\frac{dR_0}{dE_R dv_{DM}} \simeq \frac{1}{2} \frac{\rho_{DM}}{m_{DM}} \frac{1}{\mu_N^2} \tilde{\sigma}_N(q_A) \times \frac{\tilde{f}(v_{DM})}{v_{DM}}.$$

M. Ibe, et al. J. High Energy Phys. 03 (2018) 194.

2019/12/7

$$E_R \simeq \frac{q_A^2}{2m_A}, \quad q_e \simeq \frac{m_e}{m_A} q_A.$$

The total electronic energy released at the ionization is given by :

19



The most important is dp^{c}/dE_{e} M. Ibe, et al. J. High Energy Phys. 03 (2018)
 The dp^{c}/dE_{e} for a different q_{e} is obtained by multiplying q_{e}^{2}/q_{ref}^{2}

The atomic structure of Ge



The energy HPGe detector will detect is :
$$E_{det} = E_{EM} + Q_{nr}E_R$$
So the spectrum is given by :
$$\frac{dR}{dE_{det}dv_{DM}} = \int dE_R dE_{EM} \delta(E_{det} - Q_{nr}E_R - E_{EM}) \frac{dR}{dE_R dE_{EM} dv_{DM}}$$

➢ parameter setting:

- $\rho_{DM} = 0.3 GeV/cm^3$, $v_E = 232 \ km/s$
- Maxwell velocity distribution $v_0 = 220 \ km/s$
- the Galactic escape velocity $v_{esc} = 544 \ km/s$
- Helm form factor

Lindhard formula (semi-rational)



$$\frac{Q_{nr}}{1 + \kappa g(\varepsilon)}$$
The quenching factor in Ge

 $\kappa g(\varepsilon)$

- The quenching factor in Ge used is calculated by Lindhard formula.
- The value of κ in Lindhard formula is fitted by experiment data under 2 keVnr with 30% uncertailty.



Expected measureable spectra of the χ -N elastic SI-scattering and χ -N inelastic SI-scattering with Migdal effect incorporated.





TI analysis

Experiment data : from CDEX-1B \succ Exposure : 737.1 kg \cdot day

The ionization of N shell electrons were not considered.





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AM analysis



The solid lines are the expected spectra due to Migdal effect in June (red solid line), September (black solid line) and December (cyan solid line).

The AM-amplitude due to Migdal effect (red line) in 0.25-0.3 keVee energy range at different m_{χ}





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Summary

- CDEX-1B excludes DAMA/LIBRA phase-1's interpretation with the spin-independent WIMP interaction with Standard Halo model in Germanium crystal.
- ➤ Analysis on time-integrated (TI) and annual modulation (AM) effects on CDEX-1B data are performed, with 737.1 kg•day exposure and 160 eVee threshold for TI analysis, 1107.5 kg•day exposure and 250 eVee threshold for AM analysis.
- The sensitive windows in m_χ are expanded by an order of magnitude with Migdal effect incorporated. New limits on σ^{SI}_{χN} at 90% confidence level at m_χ ~ 0.05-0.18 GeV/c² and 0.075-3.0 GeV/c² in the TI and AM analysis are derived, respectively.





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