

BSM phenomenology at the next-to-leading order

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Mini-workshop on precision physics and future colliders

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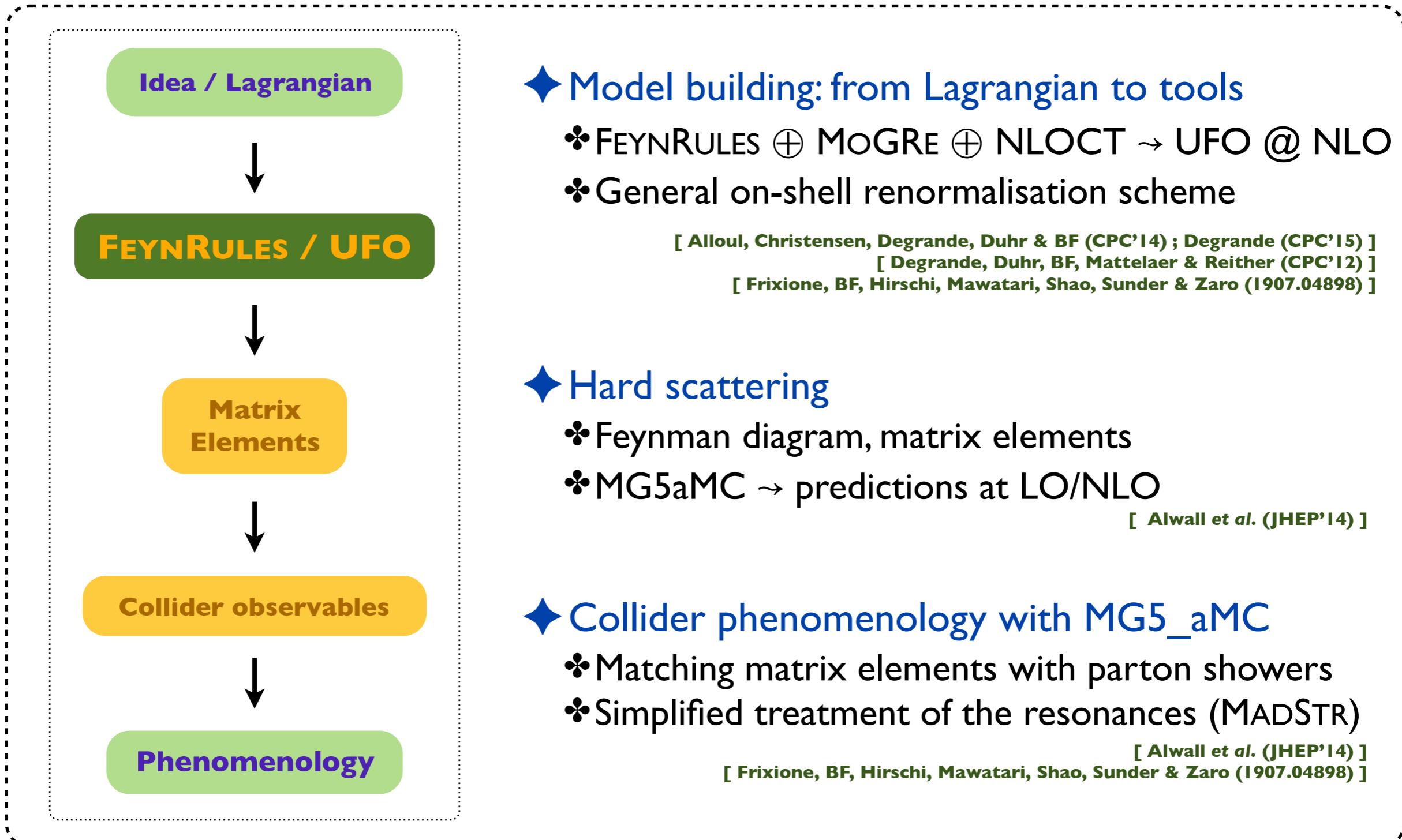
A need for precision predictions for BSM?

- ◆ Final words on any potential new physics at the LHC
 - ❖ Accurate measurements + precision predictions (NLO QCD + PS)

- ◆ New physics is standard in the simulation tools
 - ❖ 20-25 years of developments
 - ❖ Simulations at the NLO accuracy in QCD can be easily achieved
 - ★ For any model → the `MADGRAPH5_aMC@NLO` framework

A comprehensive approach to new physics calculations

[Christensen, de Aquino, Degrande, Duhr, BF, Herquet, Maltoni & Schumann (EPJC'11)]



Outline

1. Automating NLO calculations in QCD for new physics
2. Examples: supersymmetry, dark matter & vector-like quarks
3. Summary - conclusions

Automating NLO calculations in QCD for new physics

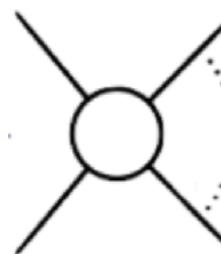
NLO calculations in a nutshell

◆ Dissecting an NLO calculation in QCD

- ❖ Three ingredients: the Born, virtual loop and real emission contributions

$$\sigma_{NLO} = \int d^4\Phi_n \mathcal{B} + \int d^4\Phi_n \int_{\text{loop}} d^d\ell \mathcal{V} + \int d^4\Phi_{n+1} \mathcal{R}$$

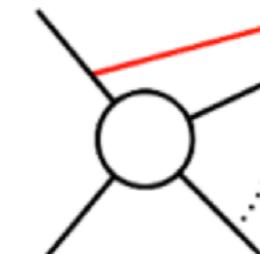
Born



Virtuels: one extra power
of α_s and divergent



Reals: one extra power
of α_s and divergent



Loop calculations

◆ Dimensional regularisation: calculations in $d = 4 - 2\epsilon$

- ❖ Divergences explicit ($1/\epsilon^2, 1/\epsilon$)
- ❖ Reduction of tensor loop-integrals to scalar integrals

◆ The reduction must be performed in a d -dimensional space

- ❖ Numerical methods work in **4 dimensions** $\rightarrow R_1$ and R_2 terms

$$\int d^d \ell \frac{N(\ell, \tilde{\ell})}{\bar{D}_0 \bar{D}_1 \cdots \bar{D}_{m-1}} \quad \text{with } \bar{\ell} = \ell + \tilde{\ell}$$

D-dim 4-dim (-2 ϵ)-dim

[Ossala, Papadopoulos, Pittau (NPB'07; JHEP'08)]

◆ The R_1 terms originate from the denominators

MADLOOP

$$\frac{1}{\bar{D}} = \frac{1}{D} \left(1 - \frac{\tilde{\ell}^2}{\bar{D}} \right)$$

- ❖ 3 **generic** non-vanishing integrals

[Hirschi et al. (JHEP'11)]

◆ The R_2 terms originate from the numerator

NLOCT

- ❖ Process-dependent contributions proportional to $\tilde{\ell}^2$
- ❖ Renormalisable theory: finite number of R_2 's $\rightarrow R_2$ **Feynman rules**

[Degrade (CPC'15)]

Matching fixed order with parton showers

[Frederix, Frixione, Maltoni & Stelzer (JHEP'09); Frixione & Webber (JHEP'02)]

◆ Subtracting the poles

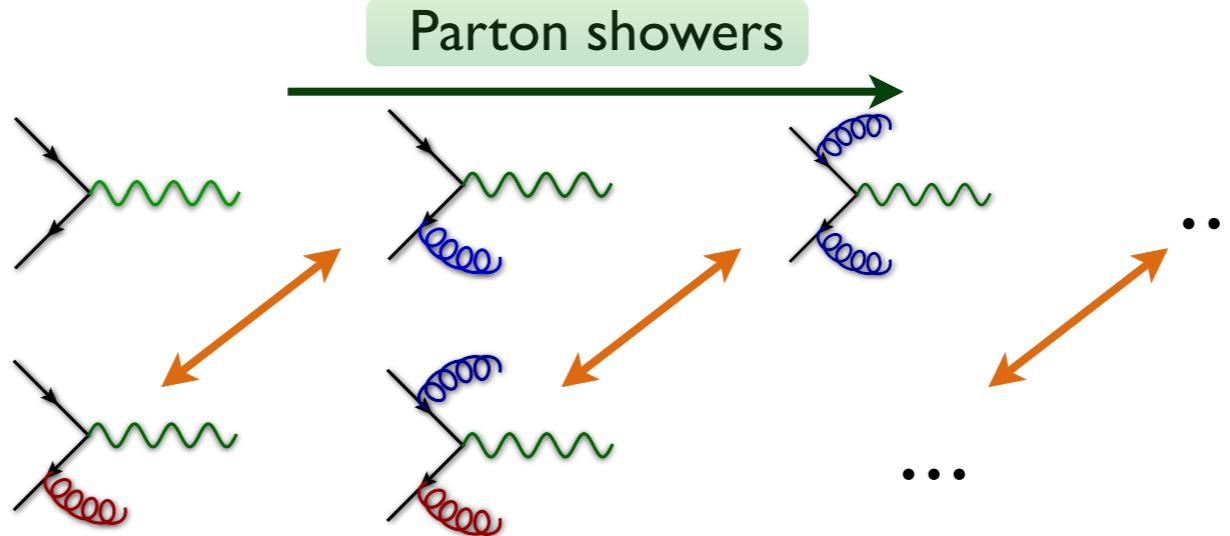
MADFKS

- ❖ The structure of the poles is known \rightarrow subtraction
 - ★ \mathcal{C} subtracted from the reals \rightarrow finite
 - ★ \mathcal{C} integrated and added back to the virtuals \rightarrow finite
 - ★ Integrals can be calculated numerically (and in 4D)

$$\sigma_{NLO} = \int d^4\Phi_n \mathcal{B} + \int d^4\Phi_{n+1} [\mathcal{R} - \mathcal{C}] + \int d^4\Phi_n \left[\int_{\text{loop}} d^d\ell \mathcal{V} + \int d^d\Phi_1 \mathcal{C} \right]$$

◆ Matching with parton showers

MC@NLO

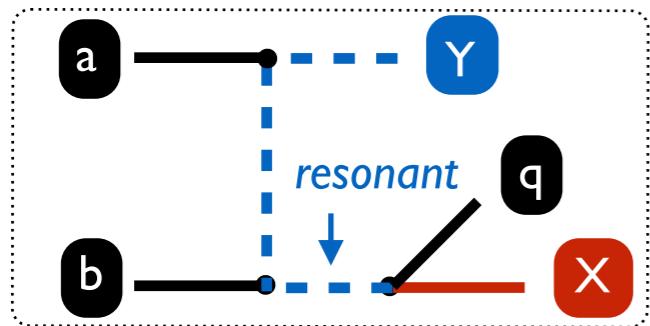


- ❖ Two sources of double counting
 - ★ Radiation: reals vs. shower
 - ★ No radiation: virtuals vs. no-emission probability

Intermediate resonances

[Frixione, BF, Hirschi, Mawatari, Shao, Sunder & Zaro (1907.04898)]

◆ Resonant contribution could appear at NLO (real emission)



❖ Overlap

- ★ YY @ LO \otimes Y \rightarrow Xq decay
- ★ YX @ NLO (real emission)

❖ Possible (huge) enhancement w.r.t. LO (if YY dominates over XY)

- ★ Spoiling the perturbative expansion for the original process

❖ All potential subprocesses may need to be considered separately

- ★ Resonances must be subtracted

◆ Resonance subtraction/removal from the squared matrix element

$$|\mathcal{A}|^2 = |\mathcal{A}^{(\text{non-res.})}|^2 + 2\Re(\mathcal{A}^{(\text{non-res.})}\mathcal{A}^{(\text{res.})\dagger}) + |\mathcal{A}^{(\text{res.})}|^2$$

❖ DR: the resonant diagrams are removed

❖ DR+I: diagram removal while keeping the interferences

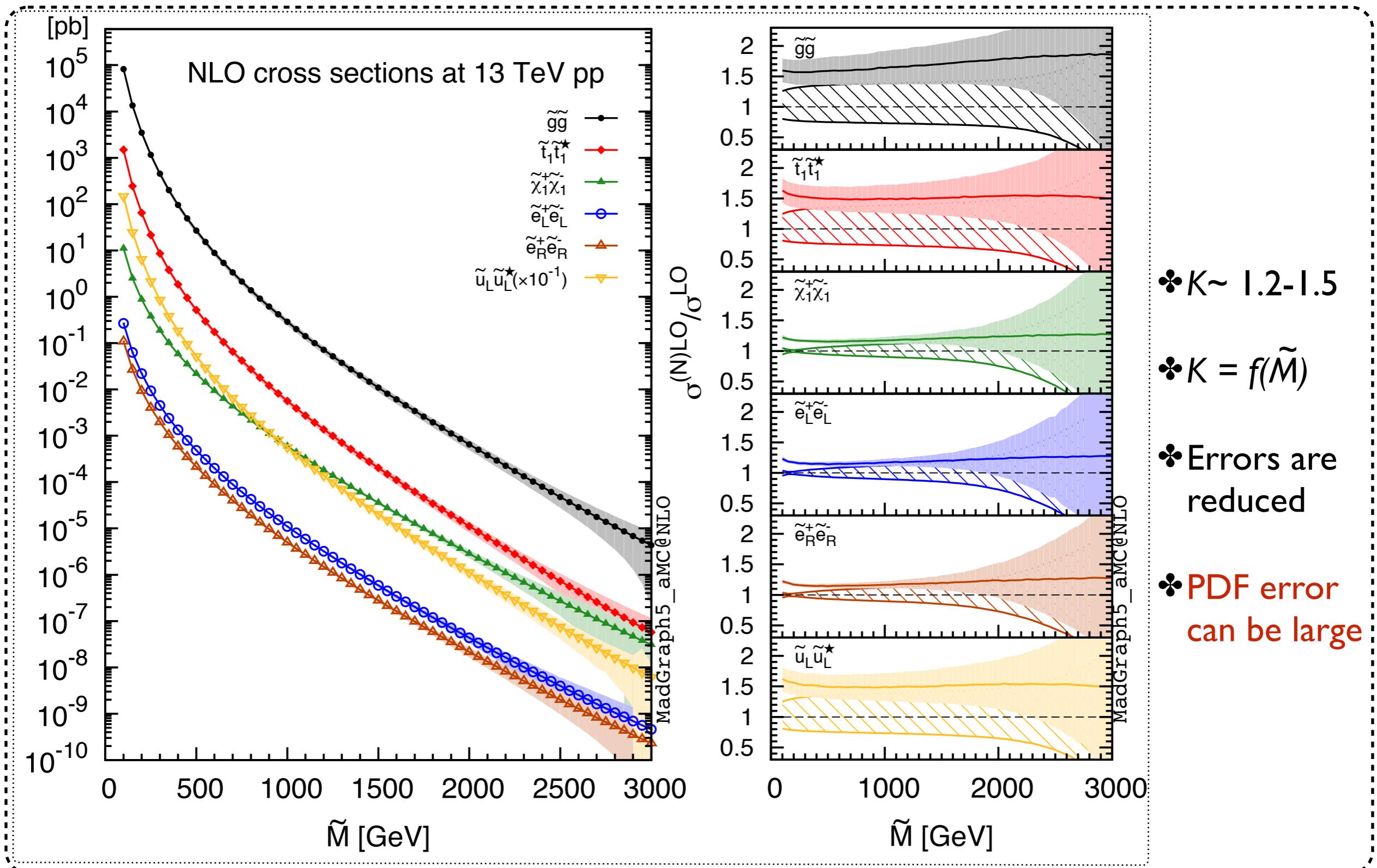
❖ DS: the purely resonant part is subtracted from the last term

- ★ There are different ways to handle this (momenta projections)

Supersymmetry @ NLO

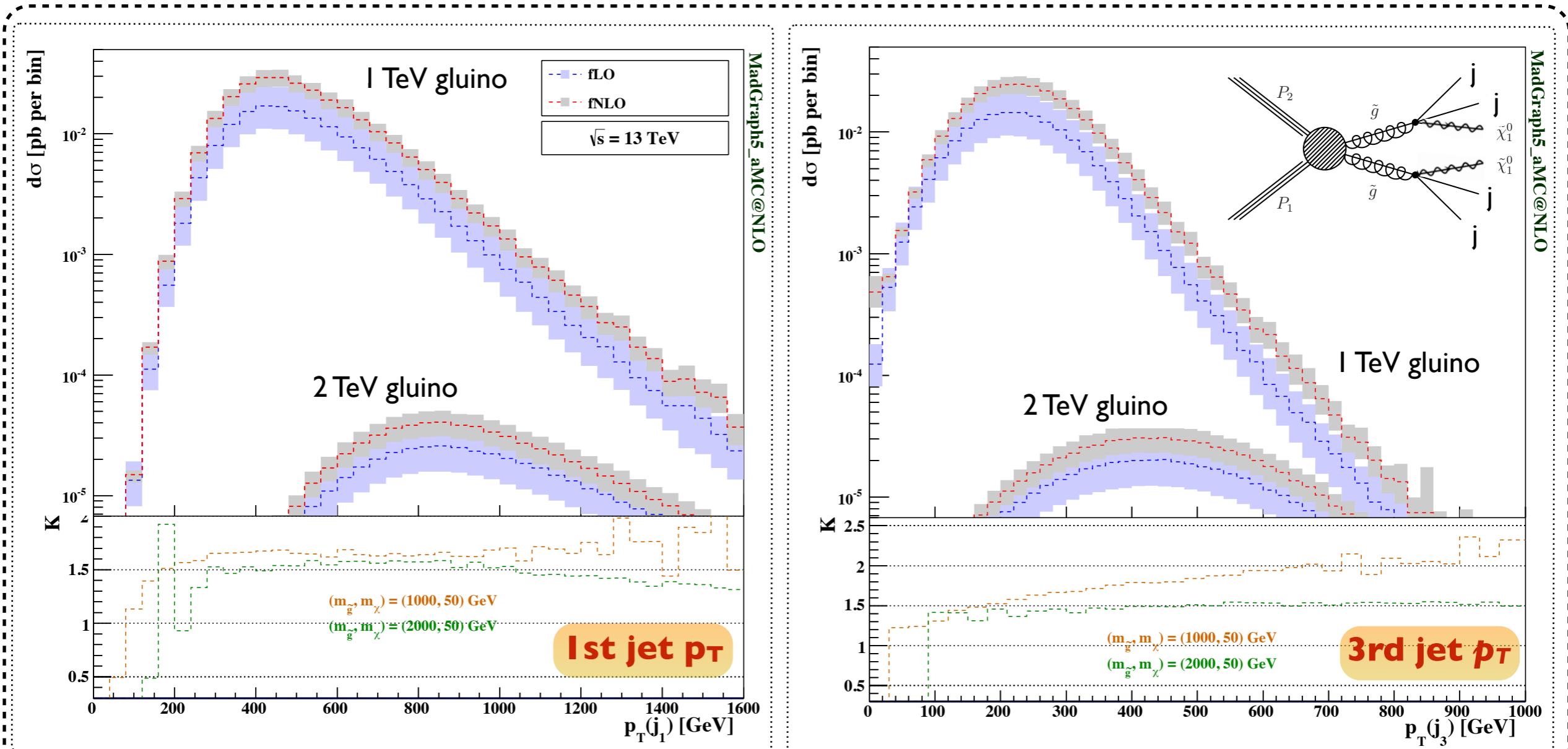
SUSY rates at 13 TeV (simplified models)

[Frixione, BF, Hirschi, Mawatari, Shao, Sunder & Zaro (1907.04898)]



Fixed-order distributions: jet properties

[Degrade, BF, Hirschi, Proudom & Shao (PRD'15; PLB'16)]

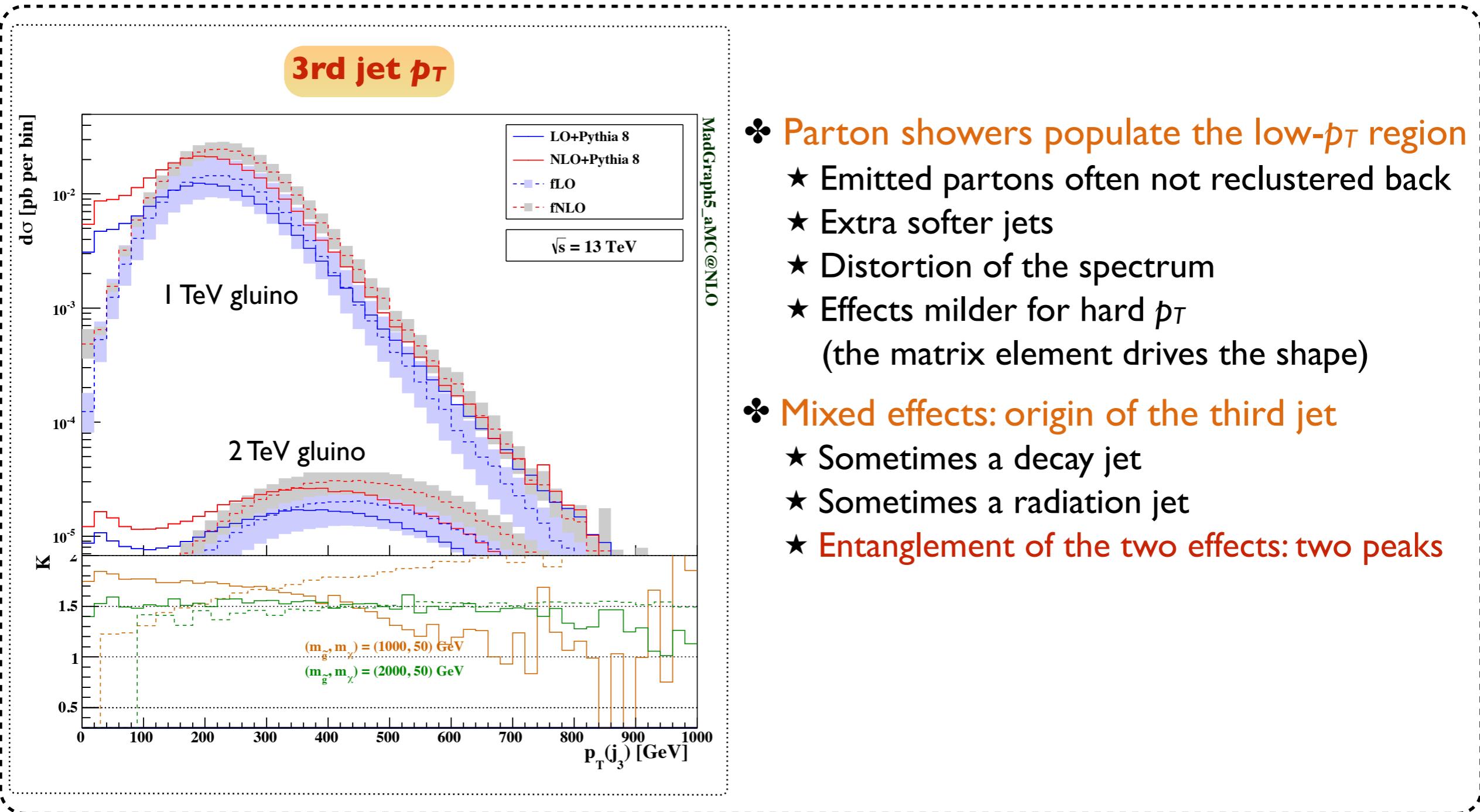


- ❖ Two potential jet origins
 - ★ Decay jet (hard)
 - ★ Radiation jet (soft, not for the 1st/2nd jets)

- ❖ Constant K-factors not accurate
 - ★ Normalisation modification
 - ★ Distortion of the shapes
 - ★ Reduction of the theoretical uncertainties

NLO+PS distributions: jet properties

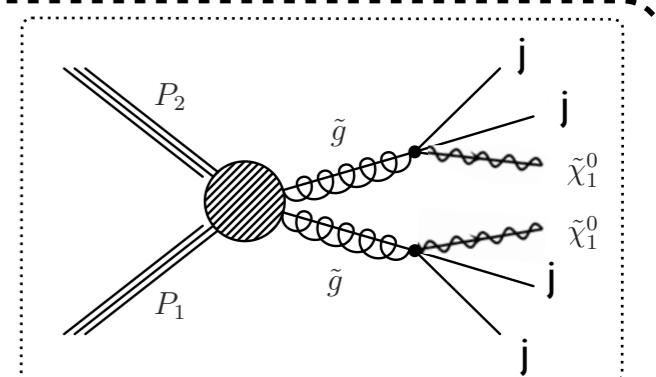
[Degrade, BF, Hirschi, Proudom & Shao (PRD'15; PLB'16)]



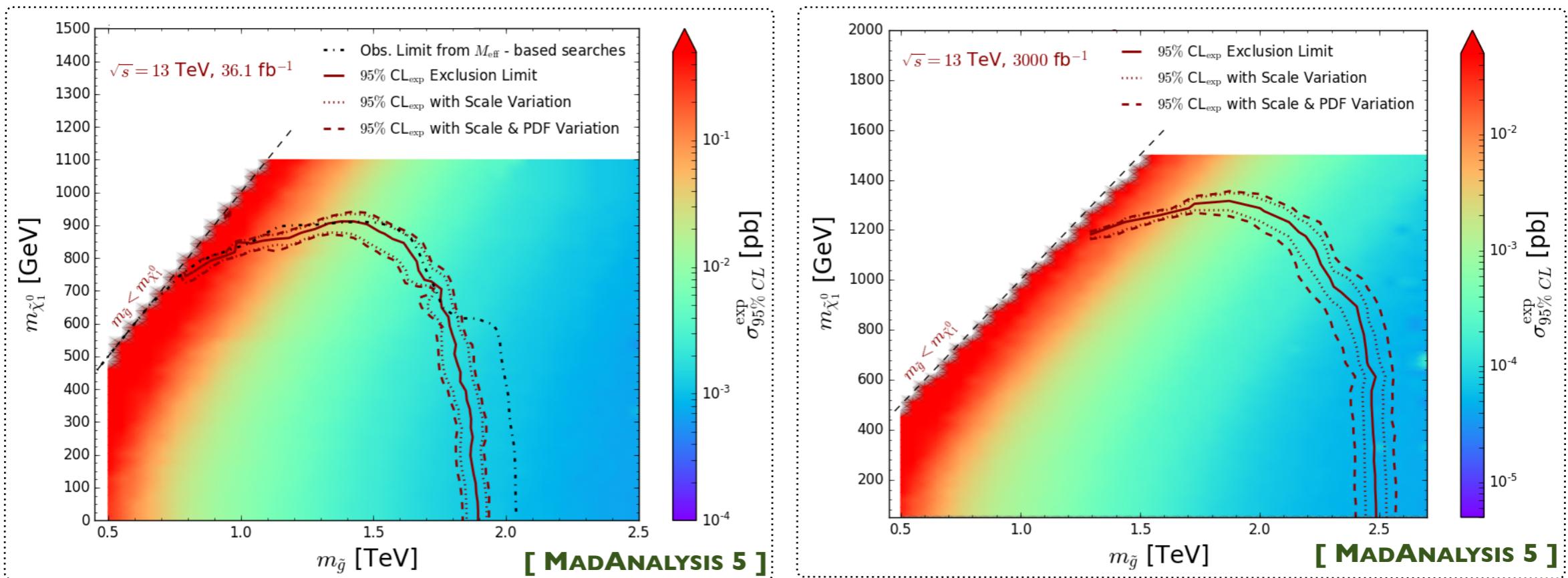
Impact of the uncertainties → future colliders

[Araz, Frank & BF (to appear)]

- ◆ Constraining gluino pair production and decay @ LHC
 - ❖ NLO impact on the shapes of the distributions
 - ❖ Impact on the limits?
 - ❖ Impact of the theory uncertainties?



- ◆ Recasting ATLAS multijet + MET analysis (ATLAS SUSY 2016-07)
 - ❖ Left: reproduction of the ATLAS results (LO-merged; $\sigma_{\text{NLL/NLO}}$) with NLO signals
 - ❖ Right: extrapolation for HL-LHC → impact of the errors



Treatment of the resonances

[Frixione, BF, Hirschi, Mawatari, Shao, Sunder & Zaro (1907.04898)]

	[fb]	DR	DR + I	DS			LO
$\tilde{g}\tilde{g}$	$\sigma_{\text{inclusive}}$	0.331	$0.330^{+19\%}_{-18\%} \pm 28\%$	0.327	0.322	0.330	$0.187^{+44\%}_{-29\%} \pm 27\%$
	σ_{fiducial}	0.228	$0.227^{+19\%}_{-18\%} \pm 28\%$	0.225	0.222	0.228	$0.128^{+44\%}_{-29\%} \pm 27\%$
$\tilde{g}\tilde{q}$	$\sigma_{\text{inclusive}}$	8.42	$8.39^{+12\%}_{-14\%} \pm 6.9\%$	8.38	8.35	8.41	$5.49^{+38\%}_{-25\%} \pm 7.0\%$
	σ_{fiducial}	5.93	$5.91^{+12\%}_{-14\%} \pm 6.9\%$	5.90	5.87	5.93	$3.86^{+38\%}_{-26\%} \pm 7.0\%$
$\tilde{q}\tilde{q}$	$\sigma_{\text{inclusive}}$	20.4	$20.4^{+7.8\%}_{-10\%} \pm 2.2\%$	20.4	20.4	20.4	$14.9^{+30\%}_{-22\%} \pm 2.2\%$
	σ_{fiducial}	14.8	$14.8^{+7.8\%}_{-9.9\%} \pm 2.2\%$	14.8	14.7	14.8	$10.8^{+30\%}_{-21\%} \pm 2.2\%$

❖ Benchmark (allowed by data)

- ★ Multi-TeV squarks and gluinos
- ★ 50 GeV lightest neutralino (decays into jets and missing energy)
- ★ Typical H_T/MET selection (+ N_{jets} requirement)

❖ NLO impact

- ★ Large K-factors (especially for $\tilde{g}\tilde{g}$), reduction of the theory errors
- ★ 50 GeV lightest neutralino (decays into jets and missing energy)
- ★ Results compatible regardless of how resonances are treated

No double counting

Dark matter @ NLO

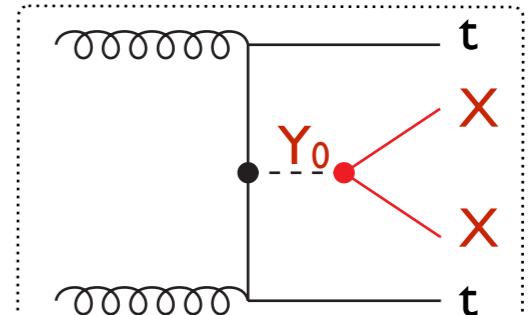
Top-philic dark matter

[Arina, Backovic, Conte, BF, Guo, Heisig, Hespel, Krämer, Maltoni, Martini, Mawatari, Pellen & Vryonidou (JHEP'16)]

◆ A simplified model for top-philic dark matter

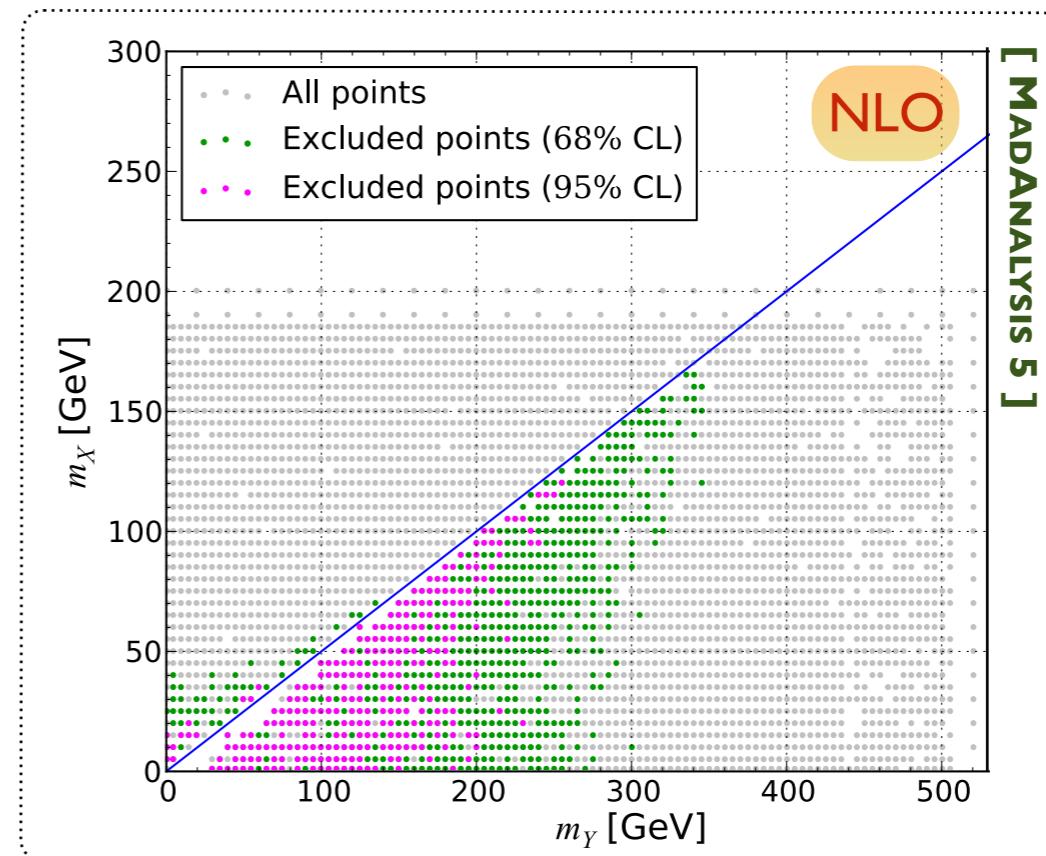
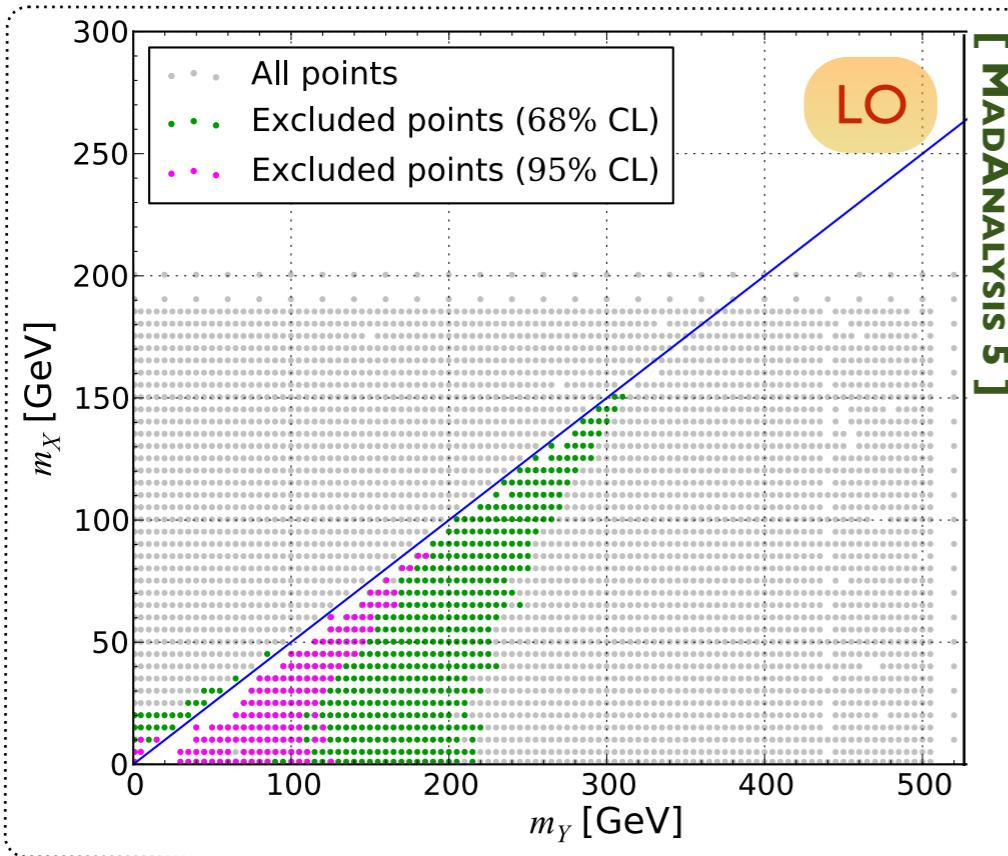
- ◆ A dark sector with a fermionic dark matter candidate X
- ◆ A (scalar) mediator Y_0 linking the dark sector and the top

$$\mathcal{L}_{t,X}^{Y_0} = - \left(g_t \frac{y_t}{\sqrt{2}} \bar{t}t + g_X \bar{X}X \right) Y_0$$



- ◆ Could be probed with $t\bar{t}$ +MET events (CMS-B2G-14-004)

◆ For central scales: mild (but visible) NLO effects on the exclusions



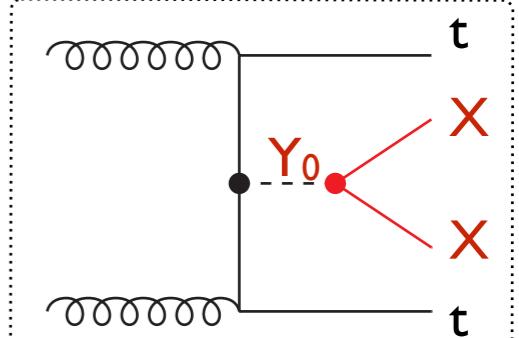
- ◆ How is the picture changing when including scale variations?

NLO effects on a CLs

[Arina, Backovic, Conte, BF, Guo, Heisig, Hespel, Krämer, Maltoni, Martini, Mawatari, Pellen & Vryonidou (JHEP'16)]

◆ There are theoretical uncertainties on a CLs number

(m_Y, m_X)	σ_{LO} [pb]	CL_{LO} [%]	σ_{NLO} [pb]	CL_{NLO} [%]
I (150, 25) GeV	$0.658^{+34.9\%}_{-24.0\%}$	$98.7^{+0.8\%}_{-13.0\%}$	$0.773^{+6.1\%}_{-10.1\%}$	$95.0^{+2.7\%}_{-0.4\%}$
II (40, 30) GeV	$0.776^{+34.2\%}_{-24.1\%}$	$74.7^{+19.7\%}_{-17.7\%}$	$0.926^{+5.7\%}_{-10.4\%}$	$84.2^{+0.4\%}_{-14.4\%}$
III (240, 100) GeV	$0.187^{+37.1\%}_{-24.4\%}$	$91.6^{+6.4\%}_{-18.1\%}$	$0.216^{+6.7\%}_{-11.4\%}$	$86.5^{+8.6\%}_{-5.5\%}$

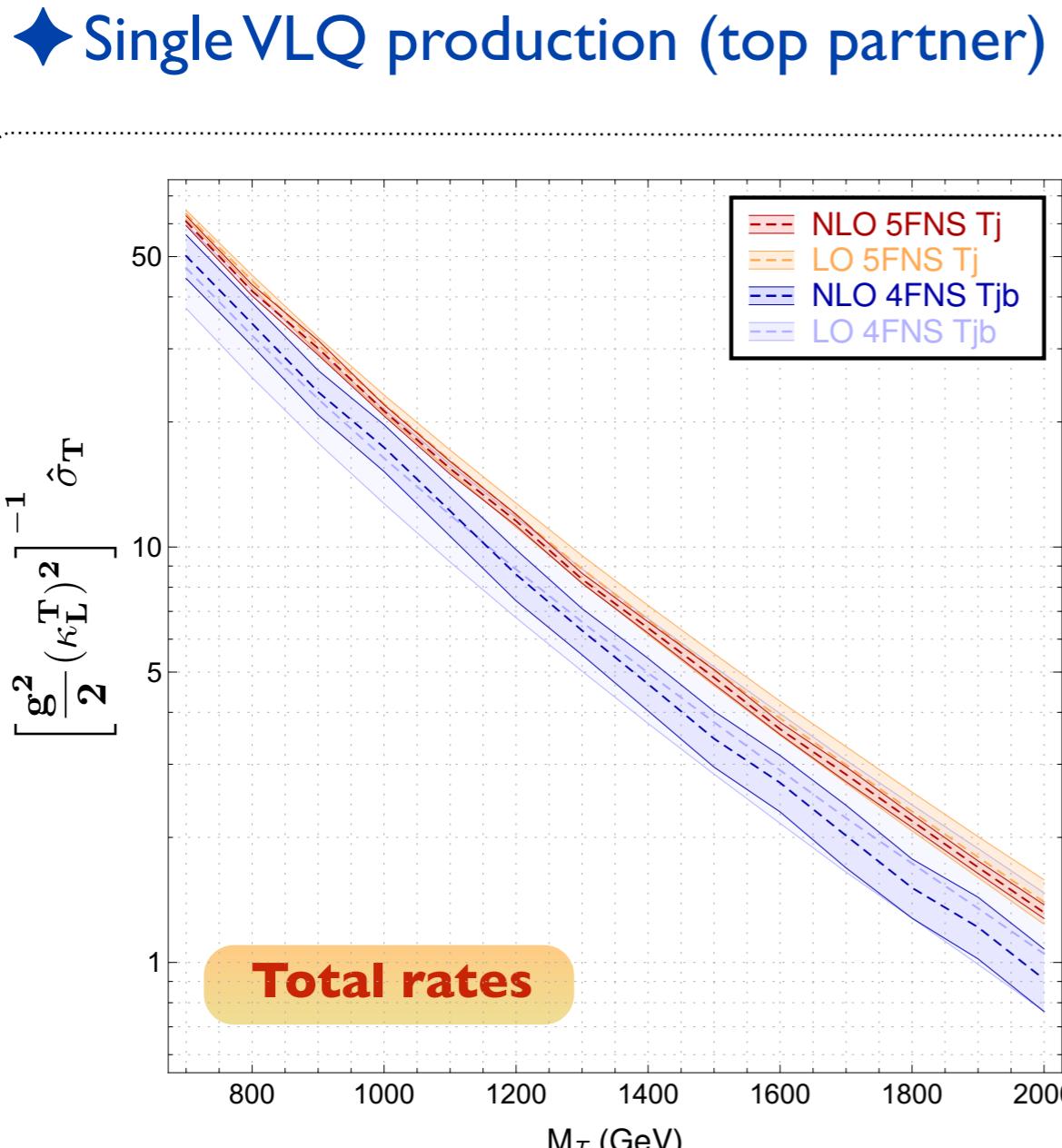


- ❖ An excluded point may not be excluded when accounting for uncertainties
- ❖ The CLs number can increase / decrease at NLO
- ❖ The error band is reduced

3rd generation VLQ @ NLO

Single VLQ production: third generation

[Cacciapaglia, Carvalho, Deandrea, Flacke, BF, Majumder, Panizzi & Shao (PLB'19)]

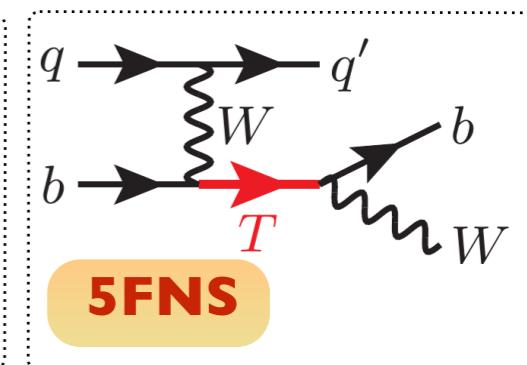
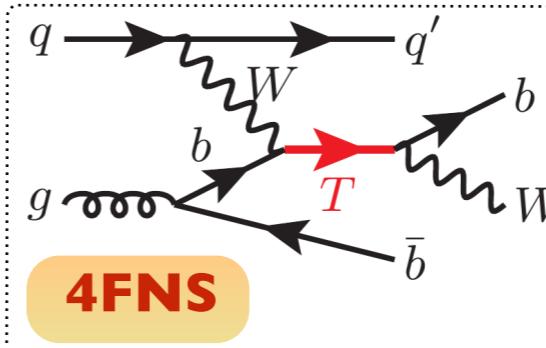


◆ Lagrangian and diagrams

♣ Production through W-couplings

$$\mathcal{L}_{\text{VLQ}} = i\bar{T}\not{D}T - m_T\bar{T}T + \frac{\sqrt{2}g}{2} \kappa_L^T [\bar{T}W P_L q_d + \text{h.c.}]$$

♣ Diagrams



◆ Total rates at NLO (4 and 5 FNS)

♣ 5FNS: $K < I$ (virtuals)

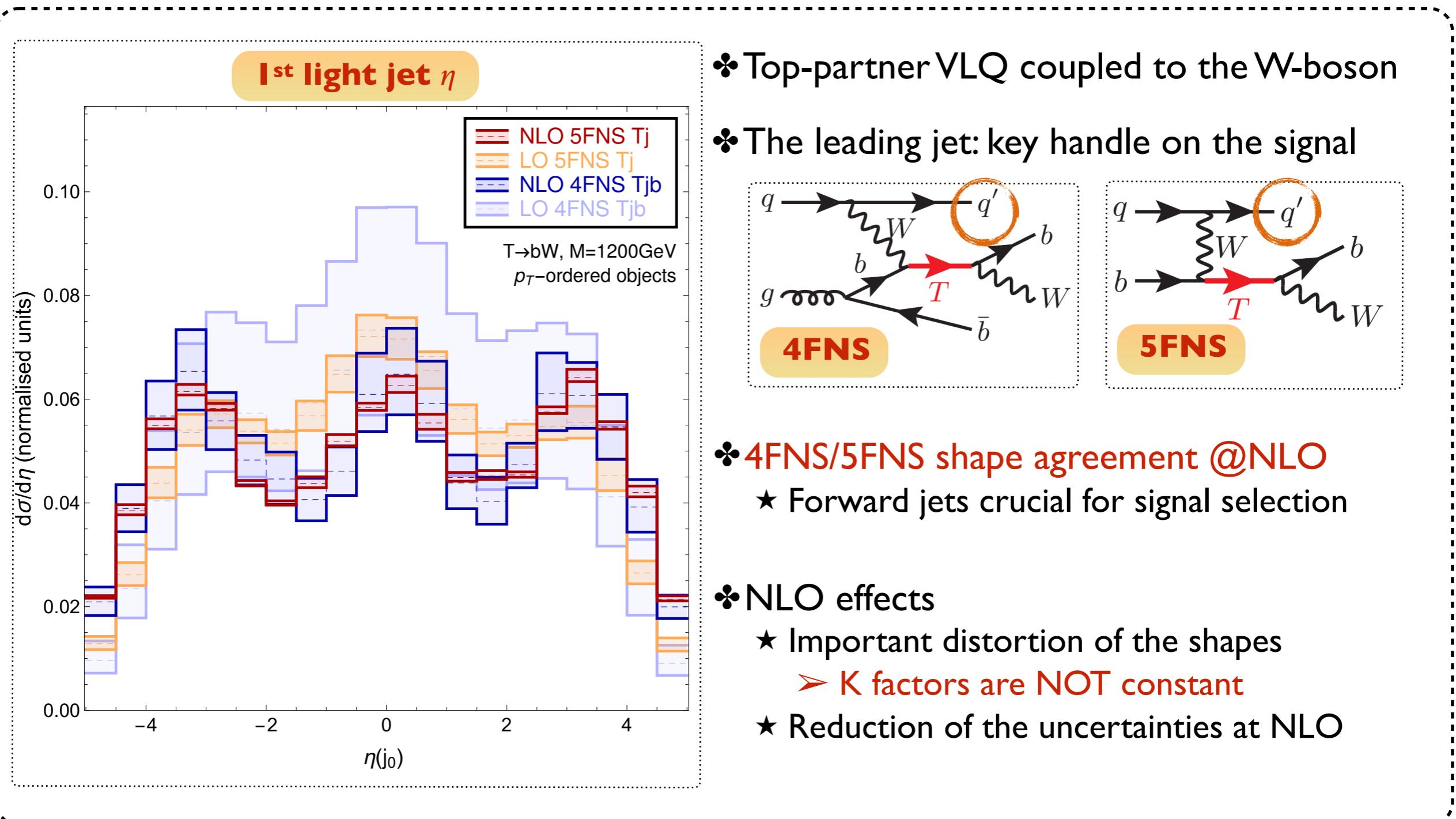
♣ 4FNS: $K = f(M_T)$

♣ Reduction of the uncertainties

♣ Log Q/m_b resummation (5FNS)
(differences at NLO for large masses)

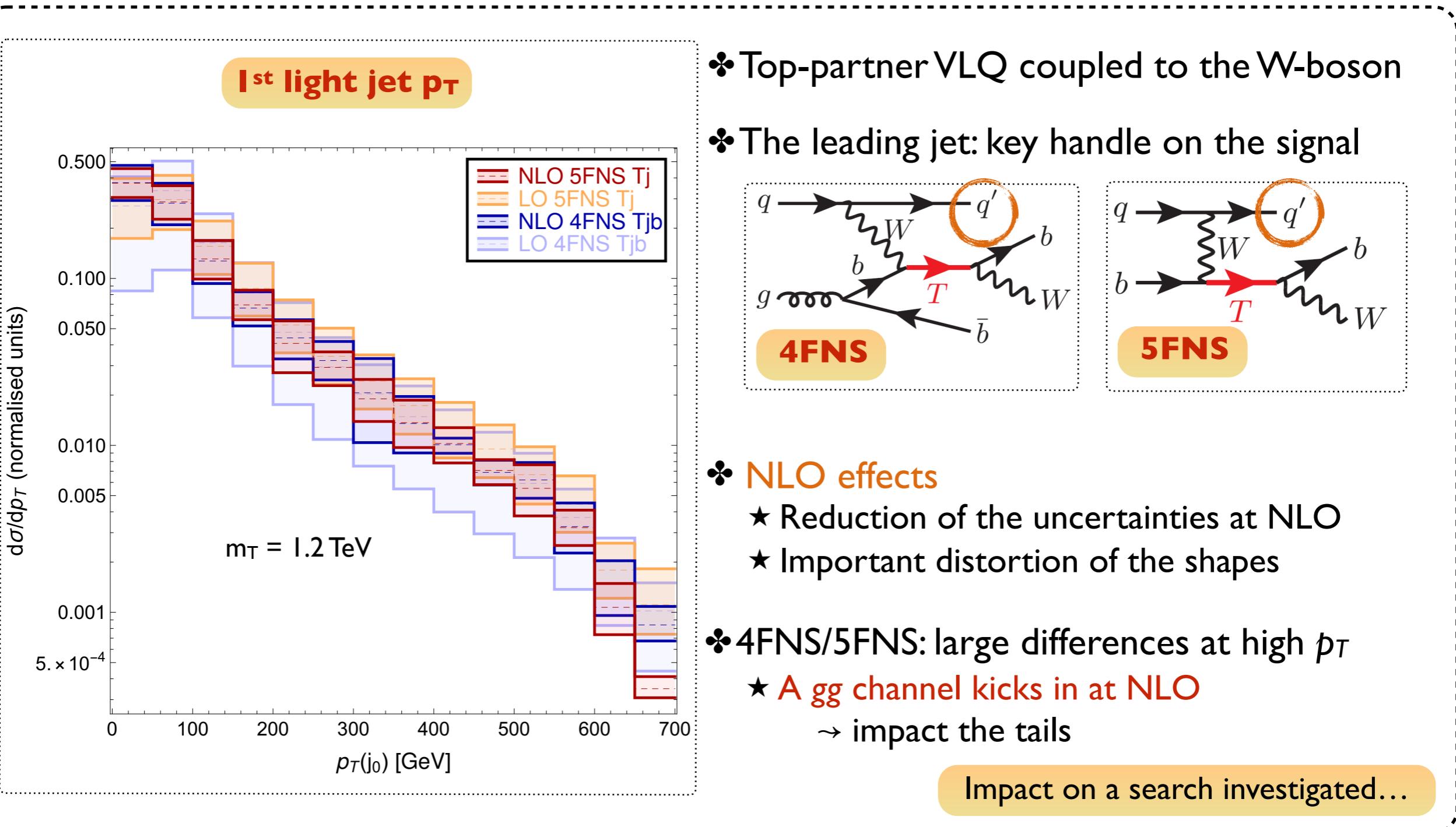
Leading jet pseudorapidity

[Cacciapaglia, Carvalho, Deandrea, Flacke, BF, Majumder, Panizzi & Shao (PLB'19)]



Leading jet transverse momentum

[Cacciapaglia, Carvalho, Deandrea, Flacke, BF, Majumder, Panizzi & Shao (PLB'19)]



Summary

Summary

◆ NLO-QCD simulations for new physics are easy to handle

- ❖ In particular via a joint use of FeynRules and MADGRAPH5_aMC@NLO
- ❖ Many models are publicly available
 - ★ Supersymmetric (simplified or not) models
 - ★ BSM Higgs models
 - ★ Dark matter simplified models
 - ★ Higgs and top effective field theories
 - ★ Vector-like quark models
 - ★ Extra gauge bosons

[<http://feynrules.irmp.ucl.ac.be/wiki/NLOModels>]

◆ Impact

- ❖ NLO effects are important and should be accounted for
- ❖ Shape distortion, large K-factors
- ❖ Uncertainties under better control
- ❖ More robust predictions