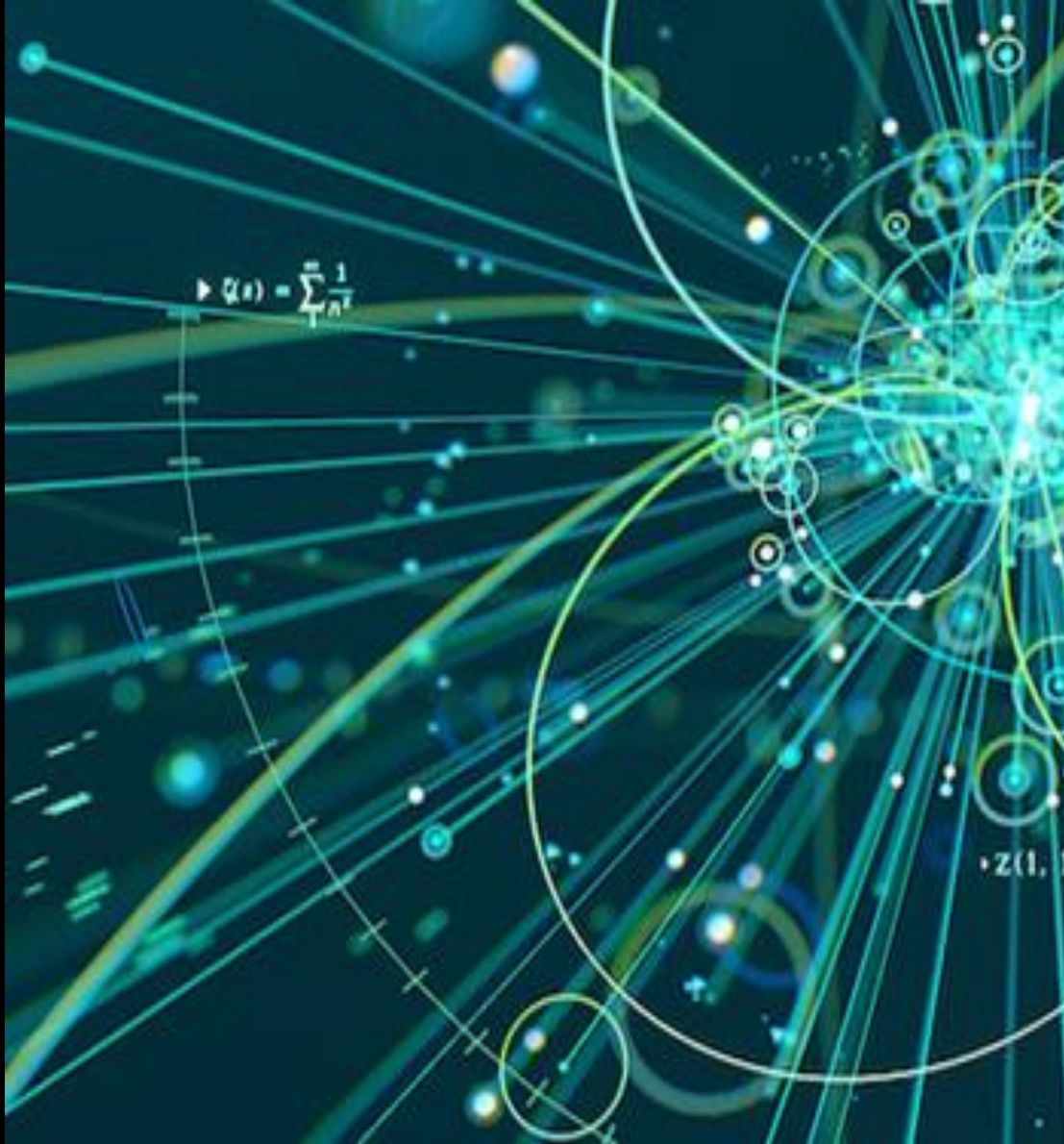




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Measurement of the relative phase between EM and strong amplitudes in $\psi(2S) \rightarrow p\bar{p}$

B. Passalacqua, M. Maggiora, M. Destefanis



OUTLINE

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- Cross Section
- Relative Phase

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INTRODUCTION

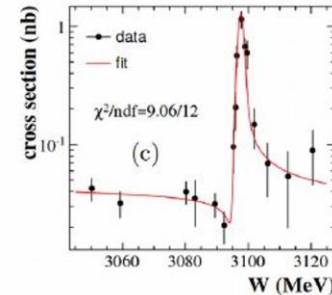
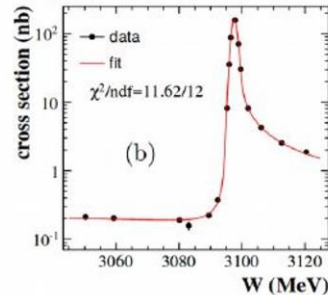
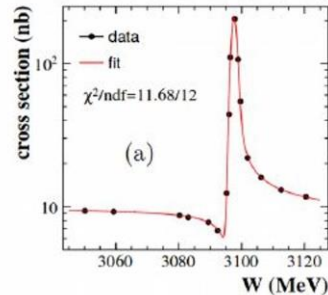
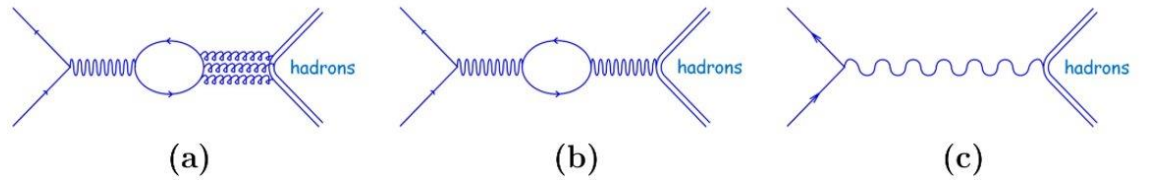


Phase Measurement

The relative phase measurement by means of the interference pattern of the e^+e^- reaction cross section as a function of the center of mass energy (W) near the resonance.

Process $e^+e^- \rightarrow \text{hadrons}$ around Charmonia

pQCD regime \rightarrow all amplitudes are expected to be almost real



$$\sigma_0 \cong |A_g(W) + A_\gamma(W)e^{i\phi_{g,EM}} + A_{cont}(W)e^{i\phi_{g,EM}}|^2$$

But experiments pointing to another direction for the J/ψ

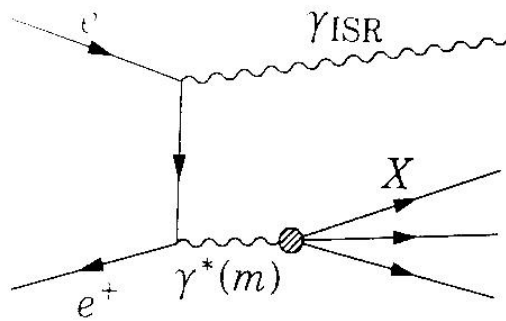
- e.g. $J/\psi \rightarrow p\bar{p}$ $\Phi = 89^\circ \pm 8^\circ$
- e.g. $J/\psi \rightarrow \rho\pi$ $\Phi = 106^\circ \pm 10^\circ$

A possible explanation:

- Quarkonium OZI breaking decay as Freund and Nambu

Initial State Radiation

In a e^+e^- pair collision one or both leptons can eventually radiate one or more photons:
the radiated energy reduces the effective CM energy of the e^+e^- annihilation.



The probability of radiating an ISR photon is described by the *radiator function* $W(s, x, \theta_\gamma)$

- x is the fraction of the beam energy carried away by the ISR photon
- θ_γ is the angle of the photon.
- ISR photon energy $\sim 50\text{-}100$ MeV
- ISR correction factor $1+\delta \equiv \int_0^1 \frac{\sigma(x)}{\sigma_0} W(x) dx$, where $x = 1 - \frac{E^2}{E_0^2}$

Cross section:

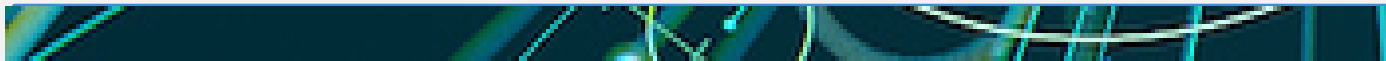
$$\sigma = \frac{N}{L\epsilon'(1+\delta)}$$

Check with ad-hoc generator

→ $p(k)dk = \beta k^{\beta-1}$ probability distribution of the ISR photon
The factor $\beta \cong 0.07$ is parametrized as:

$$\beta = 2 \frac{\alpha}{\pi} \left[\ln \left(\frac{Q^2}{m^2} \right) - 1 \right]$$

DATA ANALYSIS



Event Selection

Data collected during the 2018 run around the $\psi(2S)$ resonance (3.4 - 3.8 GeV)

→ $BR(\psi(2S) \rightarrow p\bar{p}) = (2.94 \pm 0.08) \times 10^{-4}$

Beam Energies:

Nominal E [MeV]	E [MeV]	σ_E [MeV]	L [pb^{-1}]
3580.0	3581.543	0.060	85.7
3670.0	3670.158	0.063	84.7
3681.0	3680.144	0.061	84.8
3683.0	3682.752	0.115	28.7
3684.0	3684.224	0.119	28.7
3685.5	3685.264	0.105	26.0
3686.6	3686.496	0.120	25.1
3690.0	3691.363	0.075	69.4
3710.0	3709.755	0.074	70.3

Kinematic cuts for the proton tracks:

- $|R_{xy}| < 1 \text{ cm}$, $|R_z| < 10 \text{ cm}$
- $P \leq 2 \text{ GeV}/c$
- $|\cos\theta| < 0.8$
- $E_{show}/P < 0.5$ for protons

Cuts for both the proton and the antiproton tracks:

- $178^\circ < \theta_{p\bar{p}} < 180^\circ$, $\theta_{p\bar{p}}$ is the polar angle between the two tracks $p\bar{p}$ in the CM frame
- PID tags selecting proton and antiproton
- $1.4 \text{ GeV}/c < P_{p\bar{p}} < 1.7 \text{ GeV}/c$

Selections optimization:

- Barrel region
- Back to back and charged tracks



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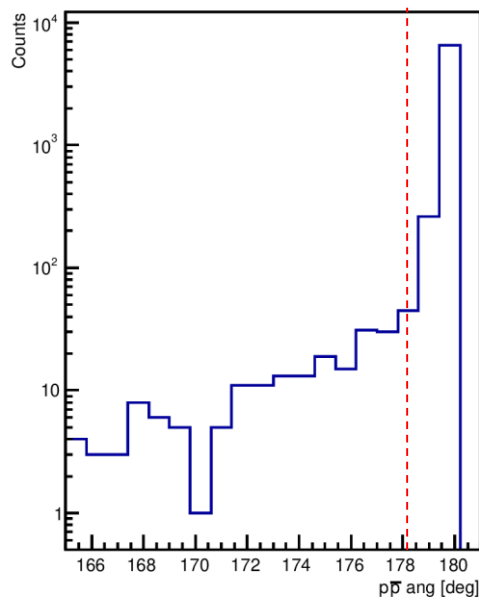
MonteCarlo Simulations $e^+e^- \rightarrow \psi(2S) \rightarrow p\bar{p}$

N of generated
event: 10000

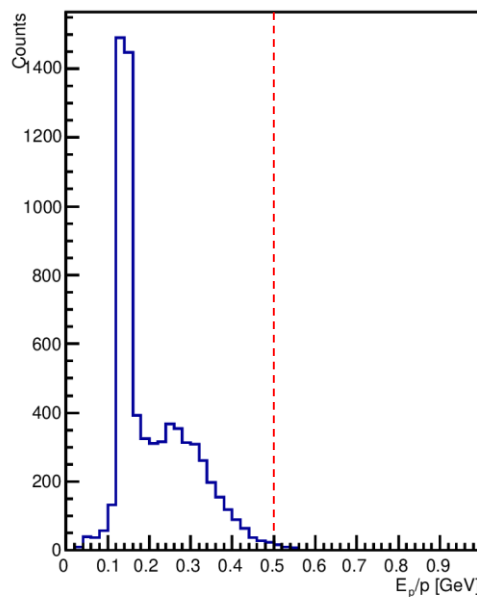
Generator:
BesEvtGen

Transport:
Geant4

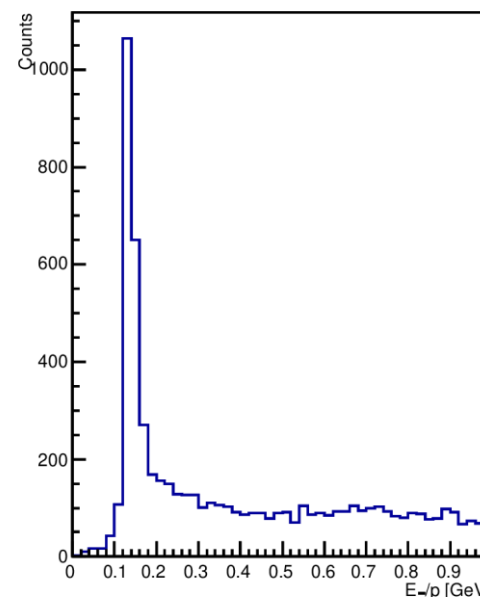
$p\bar{p}$ angle
E = 3.580 GeV



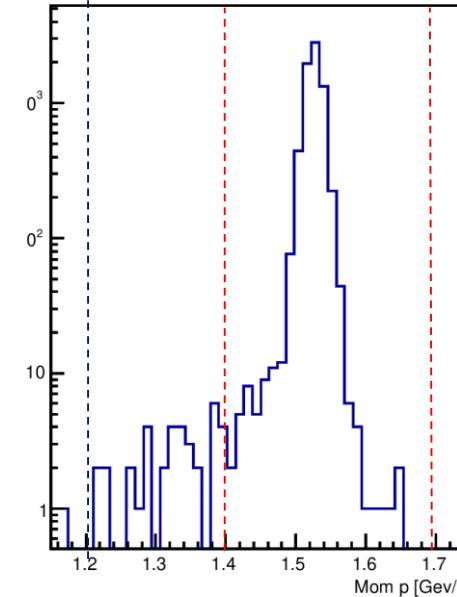
$\frac{E_{show}}{p}$ for proton
E = 3.580 GeV



$\frac{E_{show}}{p}$ for antiproton
E = 3.580 GeV



p momentum
E = 3.580 GeV



J2BB1 Model: $\frac{d|M|^2}{d\cos\theta} \propto (1 + \alpha \cos^2\theta) \quad \alpha = 0.68$

Monte Carlo Generators for Tau-Charm Physics at BESIII, Rong-Gang Ping and Cai Ying Pang, BESIII internal note, BAD522 V6 (2006)

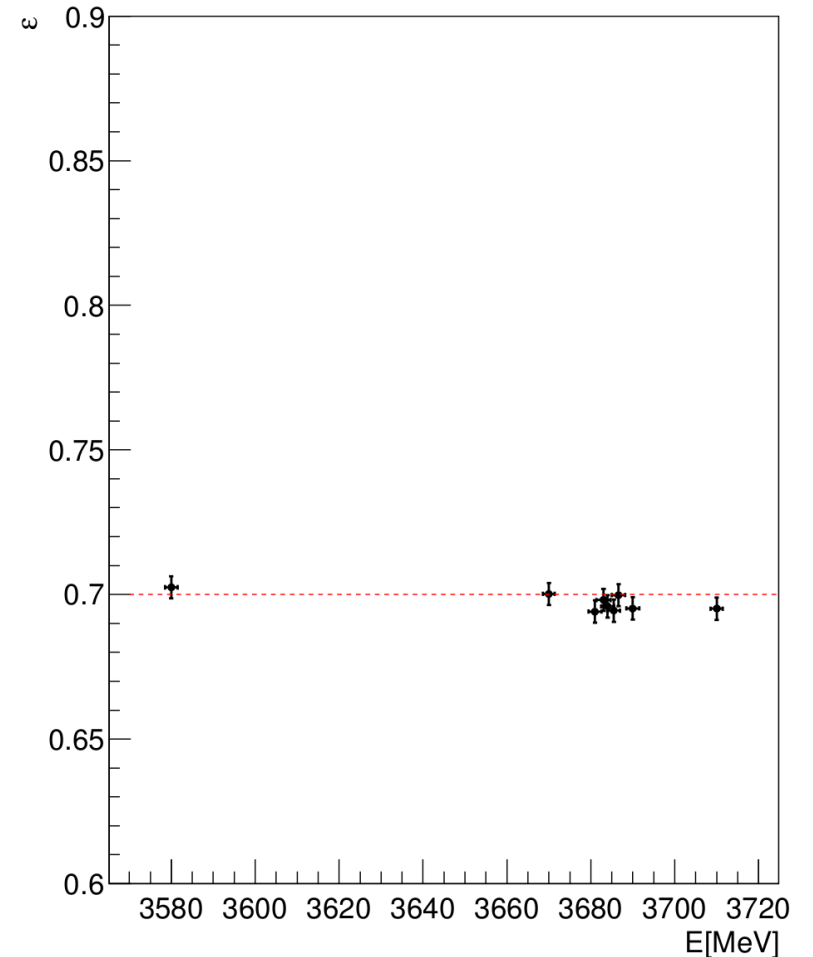
Efficiency

Values of reconstructed events obtained from the Montecarlo simulations

Nominal Energy [MeV]	$N_{reconstructed}$	Efficiency	Error Efficiency
3580.0	110	0.7025	0.0038
3670.0	180	0.7002	0.0038
3681.0	257	0.6941	0.0038
3683.0	304	0.6981	0.0038
3684.0	1408	0.6959	0.0038
3685.5	3113	0.6944	0.0038
3686.6	2955	0.6998	0.0038
3690.0	622	0.6952	0.0038
3710.0	300	0.6951	0.0038

The statistical uncertainty is estimated as binomial:

$$\frac{\sigma_{\varepsilon}}{\varepsilon} = \sqrt{\frac{1 - \varepsilon}{N_{gen}}}$$





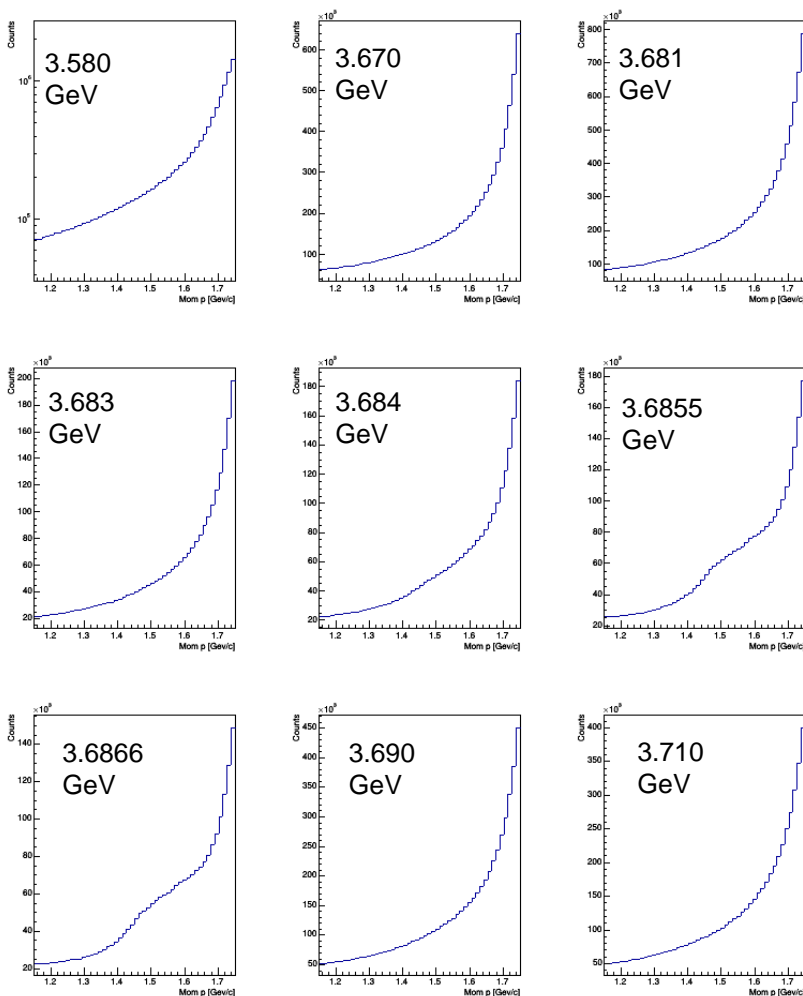
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Event Selection - Real data $e^+e^- \rightarrow \psi(2S) \rightarrow p\bar{p}$

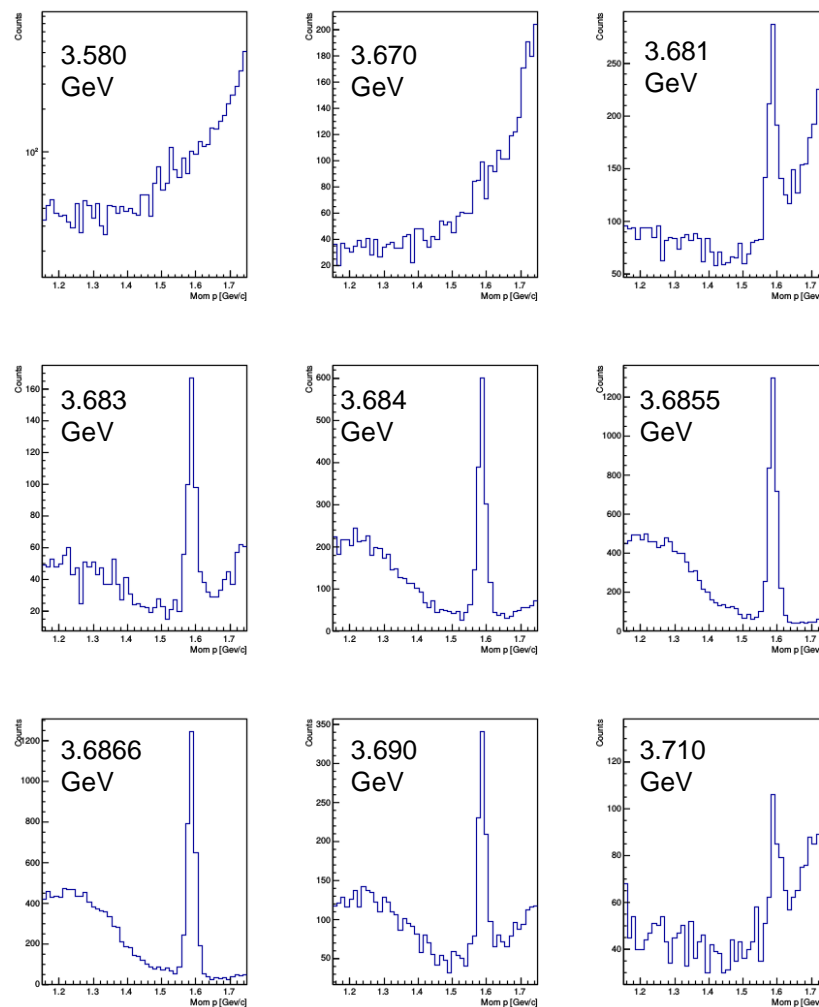
Proton momentum spectra:

no cuts applied



Proton momentum spectra:

the cut on momentum is not applied



Radiative Corrections

In central production process:

$$\beta = 4 \frac{\alpha}{\pi} \left[\ln \left(\frac{W_1}{m_e} \right) - 0.5 \right]$$

According to Touscheck,
the correction factor is:

$$C = |1 - E_n^{(1-\beta)} + 0.5E_n^{(2-\beta)}|$$

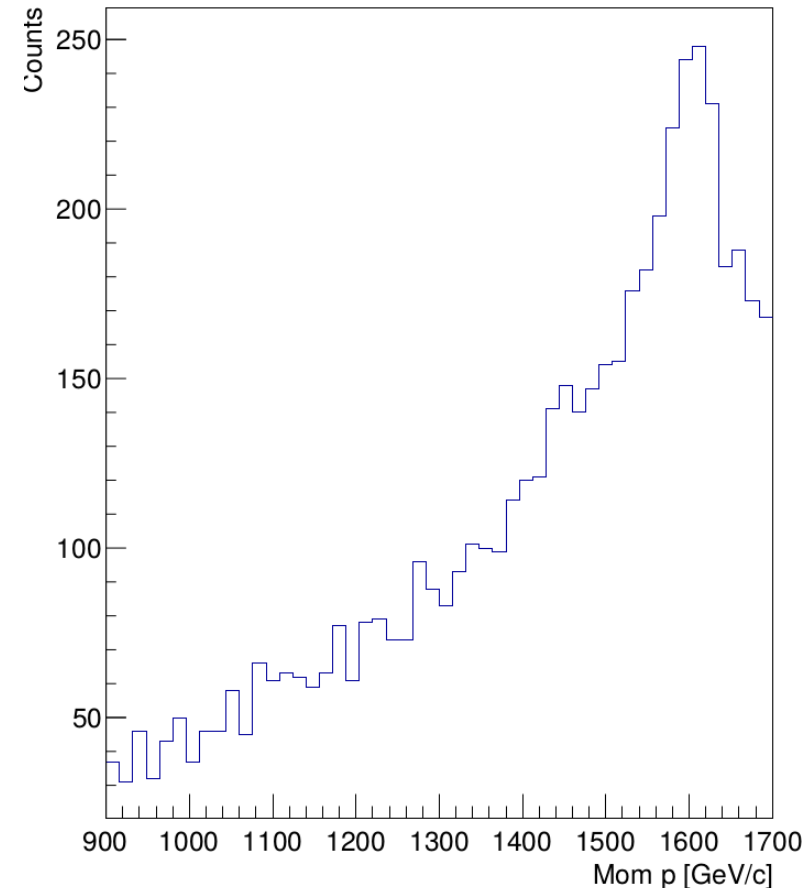
Where $E_n = k/R$

The energy after ISR:

$$W_2 = \sqrt{W_1^2 - 2kW_1}$$

- Simulated angular distribution $\propto 1 + \alpha \cos^2 \theta$ where $\alpha = 0.68$
- Photon ISR energy $\langle E_\gamma \rangle \sim 100 \text{ keV}$
- Collinearity: usually $\theta_{DIFF} \sim 4^\circ$,
where $\theta_{DIFF} = 180^\circ - \theta_{afterISR}$

Preliminary distribution of proton momentum after
ISR, private algorithm for simulation at 3.710 GeV



Number of event
generated: 10000

From Strange to Charm: Meson production in electron-positron collisions, PhD Thesis, Petterson J., Uppsala Universitet, 2019.



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Signal distribution:

Crystal Ball

Background:

Polynomial function

↓

$$N_{ev} = I \cdot N_{sig}$$

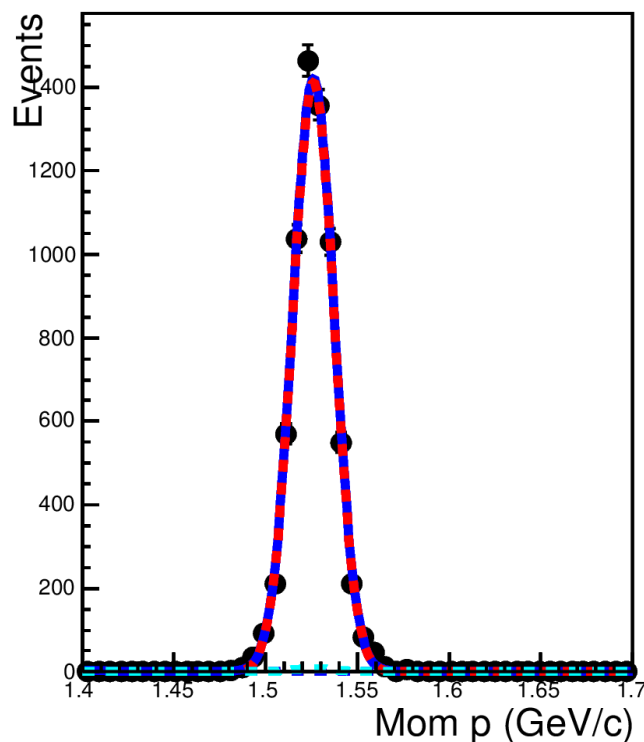
$$\sigma = \sqrt{N_{ev}}$$

Simultaneous Fit

p momentum

E=3.580 GeV

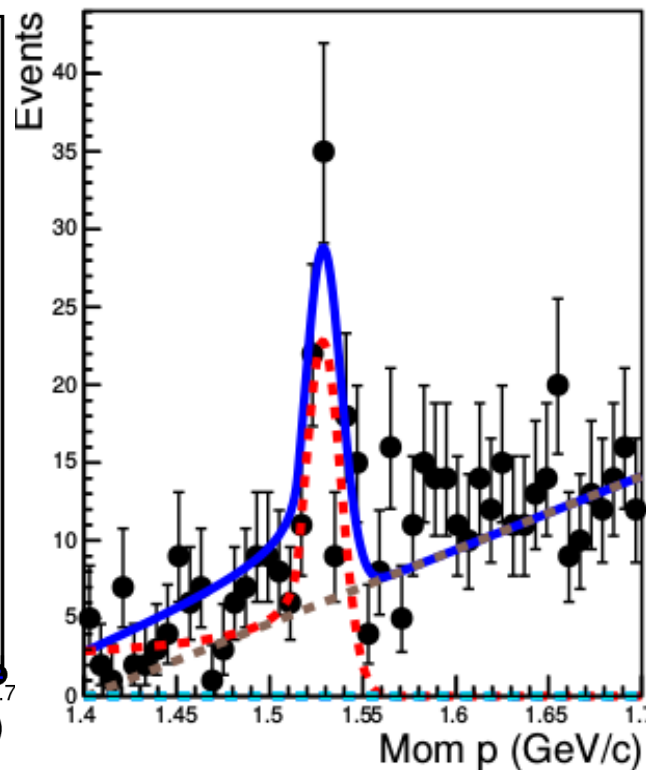
MonteCarlo Simulations



p momentum

E=3.580 GeV

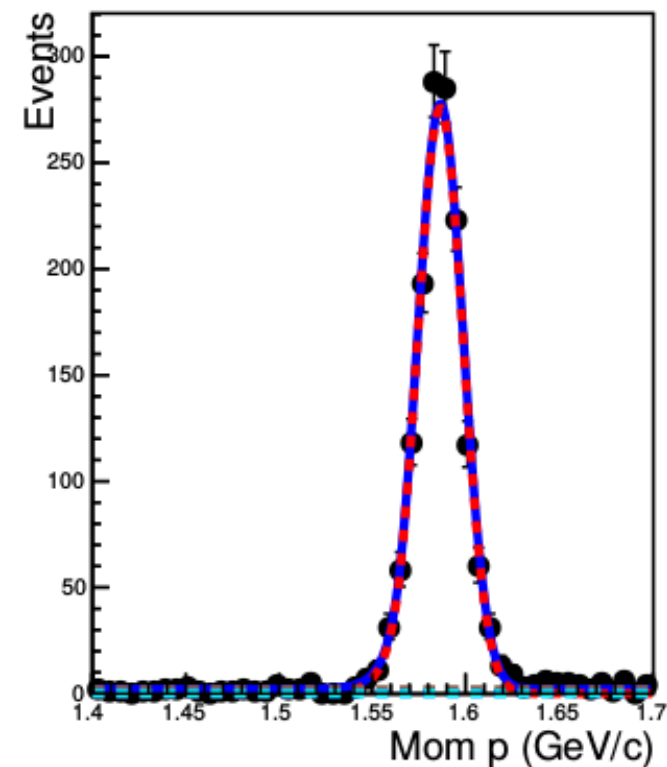
Real dataset



p momentum

E=3.6866 GeV

Real dataset



Background Studies

$$e^+e^- \rightarrow e^+e^-$$

E = 3.684 GeV

Simulations – PHSP Model

$$e^+e^- \rightarrow \mu^+\mu^-$$

E = 3.684 GeV,

Simulations – PHSP Model

$$e^+e^- \rightarrow J/\psi \rightarrow p\bar{p}$$

Invariant mass

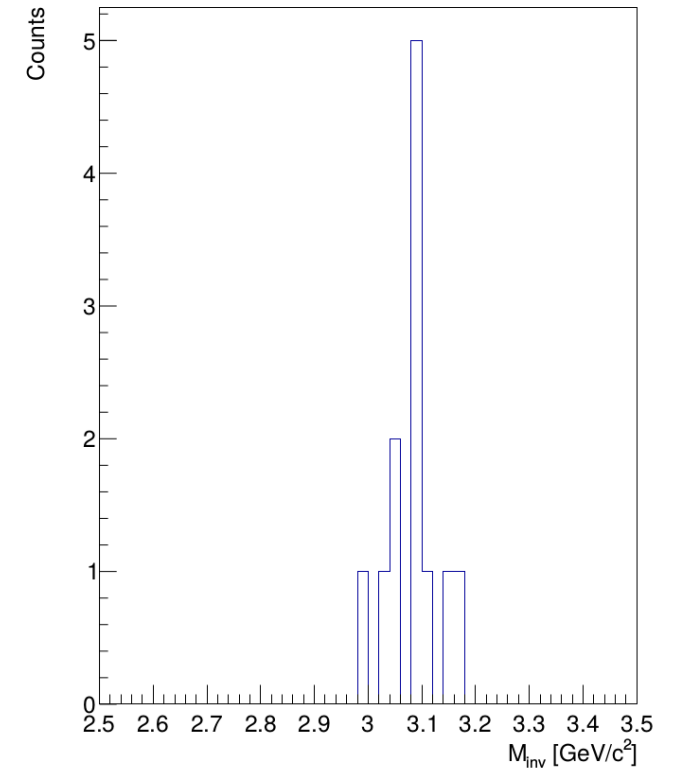
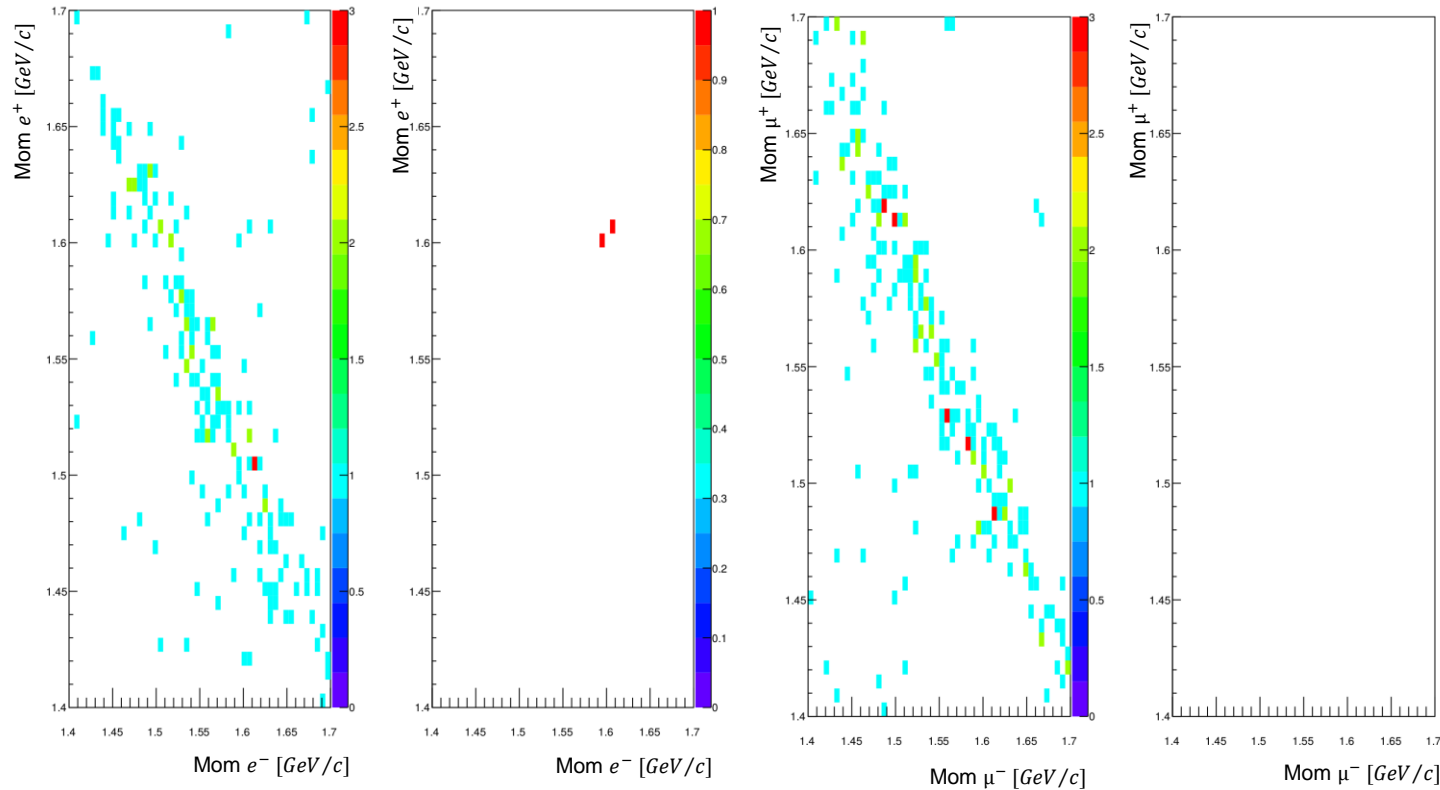
E = 3.670 GeV

Real dataset

N of event
generated: 10000

Generator:
KKMC

Transport:
Geant4



Sistematic Uncertanties

Variations on the selection criteria:

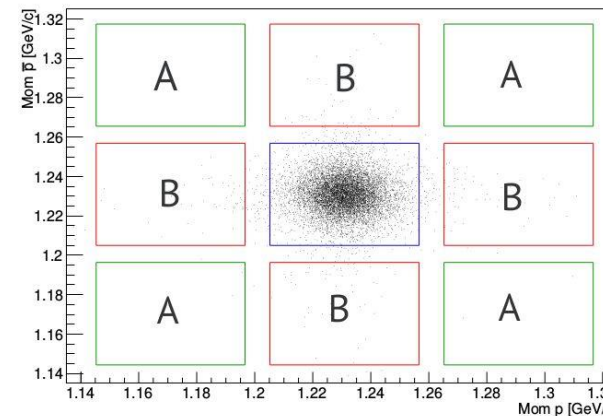
Cut	Value	Variation
E_{show}/p	0.5	± 0.05
$\theta_{p\bar{p}}$	$178^\circ < \quad < 180^\circ$	-0.5° and $+1^\circ$
fit	$-3\sigma < \quad < 3\sigma$	± 0.5
PID	0.00	$+ 0.001$

Fit routine:

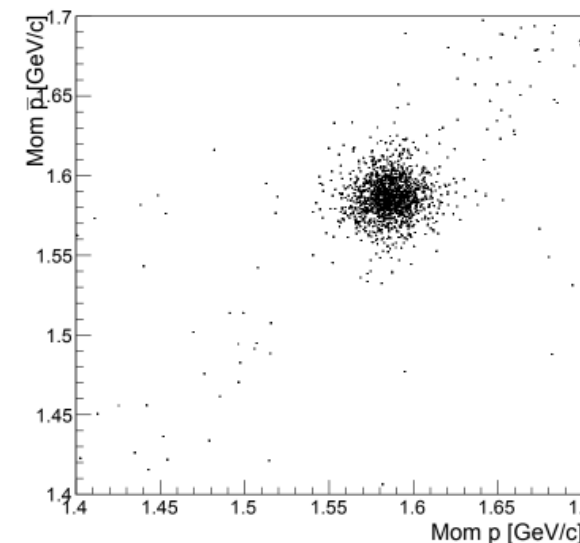
- Simultaneous fit
- Sideband method

Considering the variables as uncorrelated:

$$\sigma_{syst} = \sqrt{\sum_i \sigma_i^2}$$



$$N_0 = N - \frac{1}{2}B + \frac{1}{4}A$$



$p\bar{p}$

$E = 3.684 \text{ GeV}$



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Sistematic Uncertanties

Number of events and their error for each number of $p\bar{p}$ pairs variation and total systematic error:

Energy [MeV]	N_E	σ_E	N_T	σ_T	N_{PID}	σ_{PID}	N_F	σ_F	σ_{tot}
3580.0	80.03	5.19	6.05	0.94	80.55	8.97	93.51	5.36	9.84
3670.0	62.77	4.57	10.94	2.54	62.77	7.92	122.54	5.23	9.28
3681.0	537.38	13.38	27.02	3.29	537.38	23.18	531.39	13.11	25.11
3683.0	349.42	10.80	12.73	2.38	350.28	18.72	349.87	10.71	20.31
3684.0	1381.20	21.54	5.17	1.66	1396.93	37.38	1405.44	21.23	40.20
3685.5	3076.45	32.09	45.88	4.06	3097.80	55.66	3126.30	31.87	60.09
3686.6	2938.12	31.29	31.07	3.83	2094.12	54.17	3125.32	31.58	58.82
3690.0	736.47	15.68	21.02	2.83	740.49	27.21	752.39	15.43	29.37
3710.0	236.59	8.87	53.67	3.44	236.59	15.38	259.07	8.91	16.97

$$E \equiv \frac{E_{show}}{p} \quad T \equiv \theta_{p\bar{p}} \quad F \equiv \text{fit}$$

RESULTS



Cross Section

$$\sigma = \frac{N_{p\bar{p}}}{L \varepsilon}$$

N number of $p\bar{p}$ pairs

L integrated luminosity

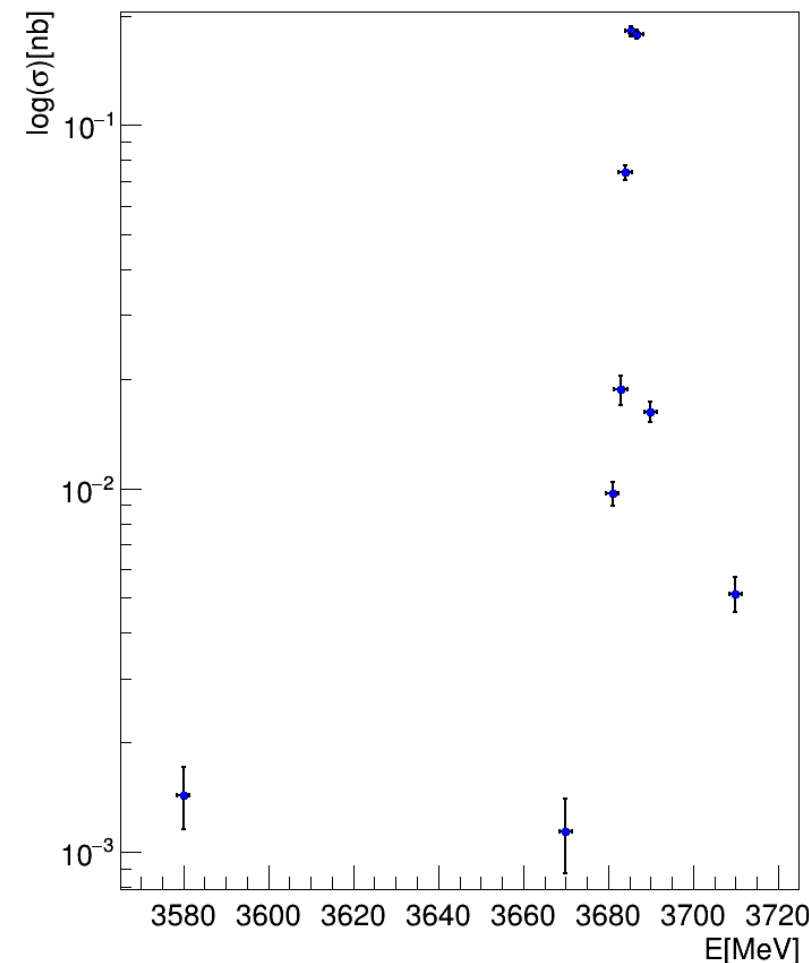
ε efficiency

Cross section for each CM energy with their statistical and systematic error

Nominal Energy [MeV]	σ [pb]	σ_{stat} [pb]	σ_{syst} [pb]
3580.0	1.43	0.16	0.23
3670.0	1.14	0.14	0.22
3681.0	9.66	0.42	0.58
3683.0	18.68	1.00	1.40
3684.0	74.01	1.98	2.79
3685.5	181.42	3.26	4.59
3686.6	177.80	3.28	4.66
3690.0	16.23	0.59	0.84
3710.0	5.13	0.33	0.47

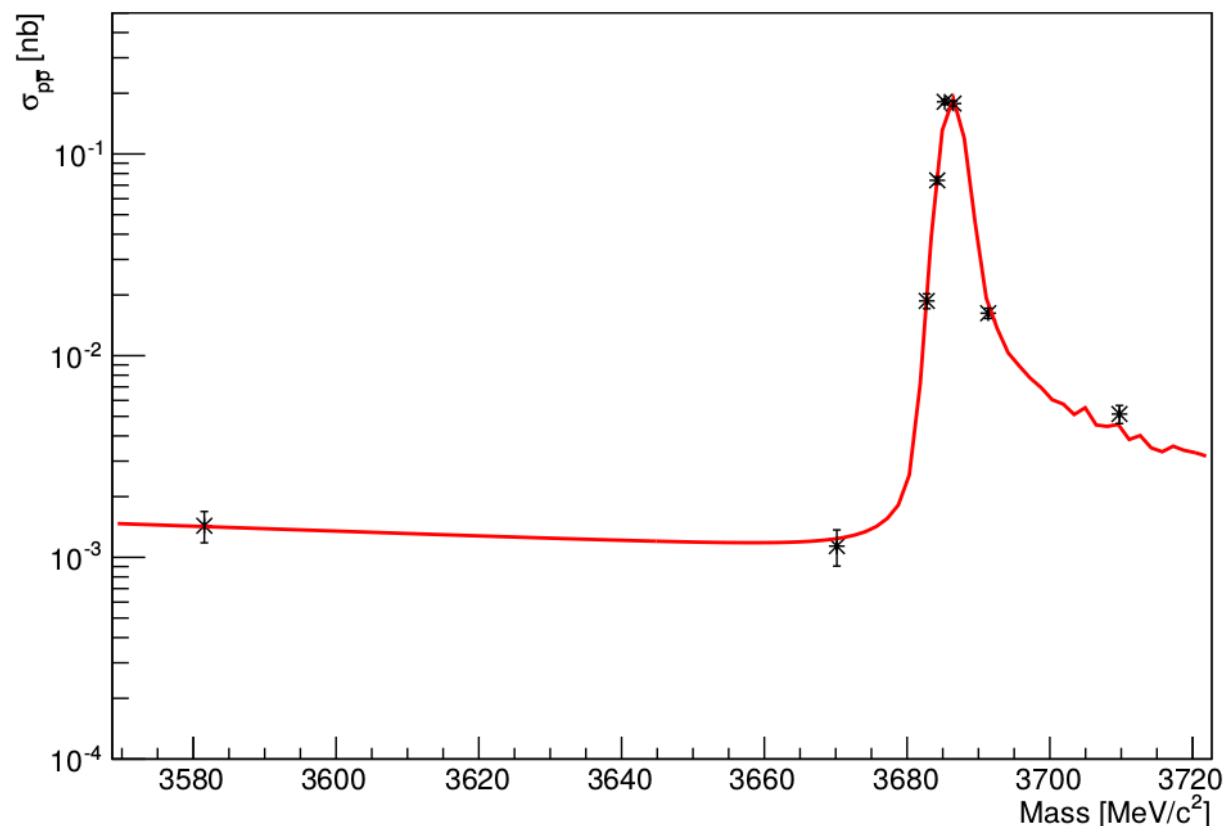
Observed cross section for the $p\bar{p}$ final states

Error bars include both statistical and systematic uncertainties



Relative Phase

Fit of the $p\bar{p}$ cross section



The cross section can be written as:

$$\sigma[nb] = \left| \sqrt{12\pi B_{in} B_{out}} \left[\frac{\hbar c}{W} \right]^2 \cdot 10^7 \frac{C_1 + C_2 e^{i\phi}}{M_\psi - W - i\frac{\Gamma}{2}} + C_3 e^{i\phi} \right|^2$$

Where C_1, C_2 and C_3 are the three parameters which correspond to the A_{3g}, A_γ and A_{EM}

- Multiple extraction to simulate ISR effects
- Cross section calculated at each extraction

$$\sigma = \frac{1}{N_{ext}} \sum \sigma_i$$

First generation measurement

Relative Phase: $\phi = (89.05 \pm 14.70)^\circ$

Branching Ratio: $B_{out} = (3.06 \pm 0.07) \times 10^{-4}$

$B_{PDG} = (2.94 \pm 0.08) \times 10^{-4}$

Cross section at the continuum: $\sigma_c = (7.54 \pm 1.12) pb$

Summary

01

$e^+e^- \rightarrow \psi(2S) \rightarrow p\bar{p}$ Event Selection

02

Simultaneous fit of momentum spectra

03

Background studies and Systematic uncertainties

04

Cross section

05

Relative Phase:

$$\phi = (89.05 \pm 14.70)^\circ$$

and Branching Fraction

$$B_{out} = (3.06 \pm 0.07) \times 10^{-4}$$

06

Next steps:

- Energy optimization
- New data



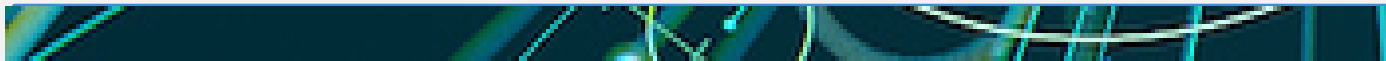
BES III



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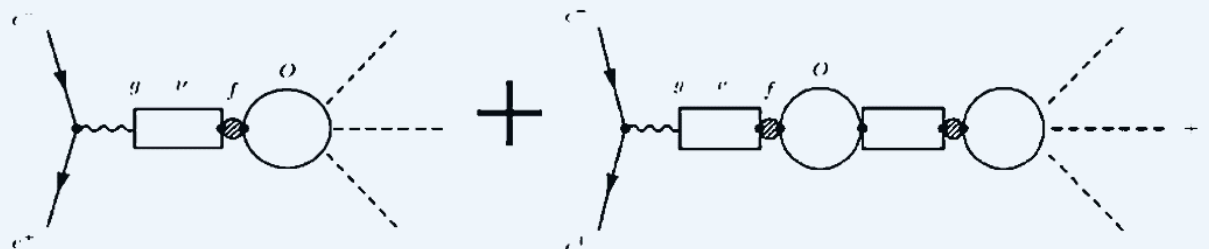
Thanks!

BACKUP SLIDES



A proposal: Quarkonium OZI breaking decay as Freund and Nambu

Considering quarkonium as a superposition of a narrow resonance ν_O , not directly decay into hadrons, and a wide resonance, a glueball O , not coupled to leptons but strongly coupled to hadrons:



Scheme of the process iterated in f, where f is the coupling between ν and O

$$A_{strong} = \frac{\sqrt{\Gamma_{ee}} M_V M_O f \sqrt{\Gamma_O}}{(M_V^2 - W^2 - i M_V \Gamma_V)(M_O^2 - W^2 - i M_O \Gamma_O) - M_V M_O f^2}$$

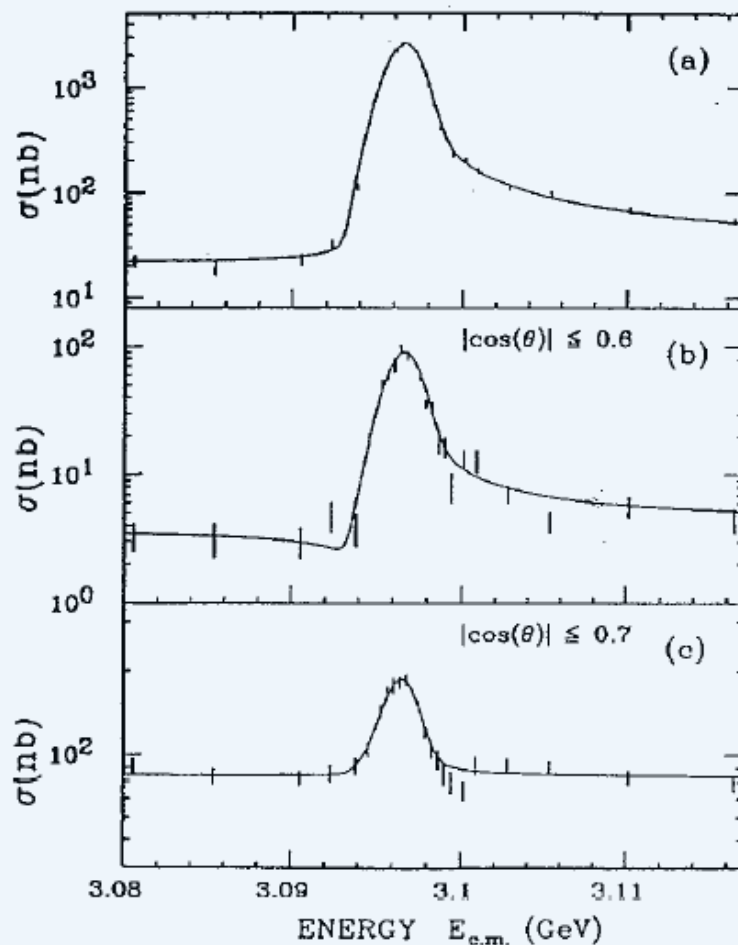
assuming $\Gamma_O \gg \Gamma_{J/\psi}$, $f^2 \sim \Gamma_O (\Gamma_{J/\psi} - \Gamma_V)$

$$A_{strong} \sim \frac{(i) \sqrt{B_{ee}} M_V f \sqrt{B_h}}{M_{J/\psi}^2 - W^2 - i M_{J/\psi} \Gamma_{J/\psi}} \quad A_{em} = \frac{\sqrt{B_{ee}} M_V \Gamma_{J/\psi} \sqrt{B_{em}}}{M_{J/\psi}^2 - W^2 - i M_{J/\psi} \Gamma_{J/\psi}}$$

Cross section of J/ψ reproduced with
 $|f| \sim 0.012 \text{ GeV}$
 $M_O \sim M_{J/\psi} \cong 3.096 \text{ GeV}$
 $\Gamma_O \sim 0.5 \text{ GeV}$

Dynamics of the Zweig- Izuka Rule and a New Vector Meson below $2 \text{ GeV}/c^2$, Peter G. O. Freund and Yoichiro Nambu Phys. Rev. Lett. 34, 1645
 R. Baldini, C. Bini, E. Luppi, Phys. Lett. B404, 362 (1997)

Was an interference already seen?



$$e^+e^- \rightarrow \text{hadrons}$$

$$e^+e^- \rightarrow \mu^+\mu^-$$

$$e^+e^- \rightarrow e^+e^-$$

Yes,
without the strong contribution

J.Z. Bai et al., Phys. Lett. B 355, 374-380 (1995)

Radiator function

$$W(s, x) = \Delta \beta x^{\beta-1} - \frac{\beta}{2} (2 - x) + \frac{\beta^2}{8} ((2 - x)(3 \ln(1 - x) - 4 \ln x) - 4 \frac{\ln(1 - x)}{x} - 6 + x)$$

where

$$L = 2 \ln \frac{\sqrt{s}}{m_e}$$

$$\Delta = 1 + \frac{\alpha}{\pi} \left(\frac{3}{2} L + \frac{1}{3} \pi^2 - 2 \right) + \left(\frac{\alpha}{\pi} \right)^2 \delta_2$$

$$\delta_2 = L^2 \left(\frac{9}{8} - 2 \xi_2 \right) - L \left(\frac{45}{16} - \frac{11}{2} \xi_2 - 3 \xi_3 \right) - \frac{6}{5} \xi_2^2 - \frac{9}{2} \xi_3 - 6 \xi_2 \ln 2 + \frac{57}{12}$$

$$\beta = \frac{2\alpha}{\pi} (L - 1), \quad \xi_2 = 1.64493407, \quad \xi_3 = 1.2020569$$

The angular distribution of the ISR photon is described by:

$$P(\theta_\gamma) = \frac{\sin^2 \theta_\gamma - \frac{x^2 \sin^4 \theta_\gamma}{2(x^2 - 2x + 2)}}{(\sin^2 \theta_\gamma + \frac{m^2}{E^2} \cos^2 \theta_\gamma)^2} - \frac{\frac{m^2(1-2x) \sin^2 \theta_\gamma - x^2 \cos^4 \theta_\gamma}{E^2 (x^2 - 2x + 2)}}{(\sin^2 \theta_\gamma + \frac{m^2}{E^2} \cos^2 \theta_\gamma)}$$

Crystal Ball function

$$f(x, \alpha, n, \bar{x}, \sigma) = N \begin{cases} \exp\left(-\frac{(x - \bar{x})^2}{2\sigma^2}\right) & \text{for } \frac{x - \bar{x}}{\sigma} > -\alpha \\ A \left(B - \frac{x - \bar{x}}{\sigma}\right)^n & \text{for } \frac{x - \bar{x}}{\sigma} \leq -\alpha \end{cases}$$

where

$$A = \left(\frac{n}{|\alpha|}\right)^n \exp\left(-\frac{|\alpha|^2}{2}\right)$$

$$B = \frac{n}{|\alpha|} - |\alpha|$$

$$N = \frac{1}{\sigma(C + D)}$$

$$C = \frac{n}{|\alpha|} \frac{1}{n-1} \exp\left(-\frac{|\alpha|^2}{2}\right)$$

$$D = \sqrt{\frac{\pi}{2}} \left(1 + \operatorname{erf}\left(\frac{|\alpha|}{\sqrt{2}}\right)\right)$$

M. J. Oreglia, *A Study of the Reaction $\psi' \rightarrow \gamma\gamma\psi$* , Ph. D. Thesis, SLAC-R-236, 1980

The cross section can be written as:

$$\sigma[nb] = \left| \sqrt{12\pi B_{in} B_{out}} \left[\frac{\hbar c}{W} \right]^2 \cdot 10^7 \frac{C_1 + C_2 e^{i\phi}}{M_\psi - W - i\frac{\Gamma}{2}} + C_3 e^{i\phi} \right|^2$$

Where C_1, C_2 and C_3 are the three parameters which correspond to the A_{3g}, A_γ and A_{EM}

The Real and the Imaginary part of the cross section AA and BB respectively, can be defined as:

$$AA = \sqrt{C_0} \frac{(C_1 + C_2 \cos\phi) - (M_\psi - W) + C_2 \Gamma/2 \sin\phi}{(M_\psi - W)^2 + (\Gamma/2)^2} + C_3 \cos\phi$$

$$BB = \sqrt{C_0} \frac{(C_1 + C_2 \cos\phi) \Gamma/2 + C_2 \sin\phi}{(M_\psi - W)^2 + (\Gamma/2)^2} + C_3 \sin\phi$$

For each extraction the cross section is:

$$\sigma_i = AA^2 + BB^2$$

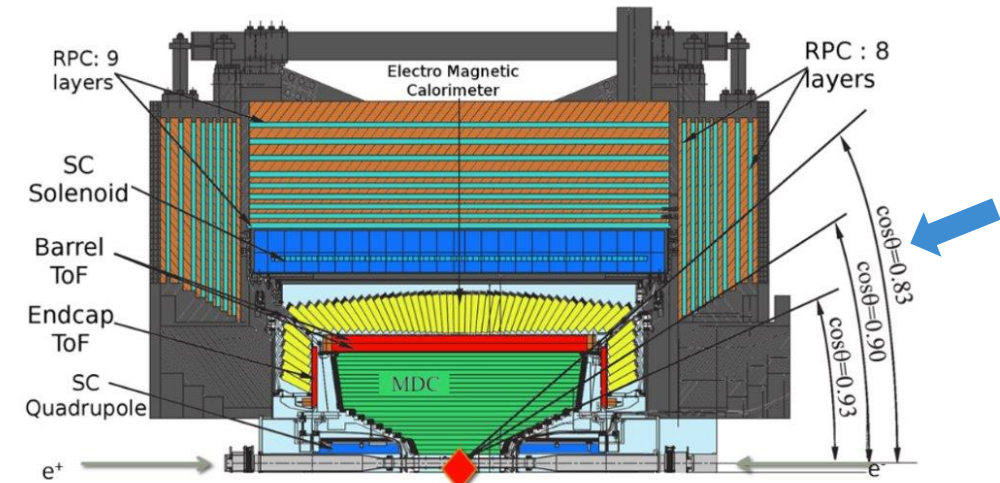
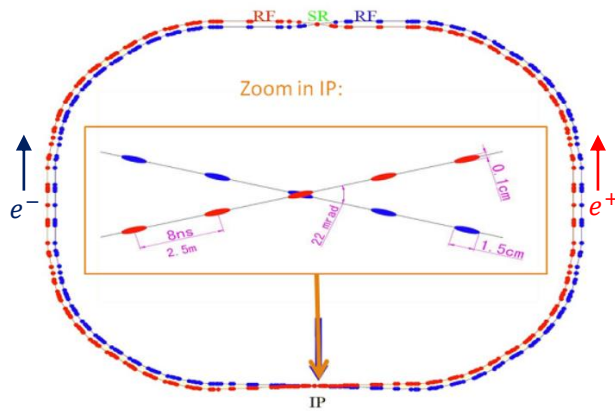
The final value of the simulated cross section is:

$$\sigma = \frac{1}{N_{est}} \sum \sigma_i$$

The BESIII Experiment

Where? Beijing in People's Republic of China (PRC)

BESIII Collaboration now has ~500 members from 72 institution
from 15 countries, including IHEP and INFN



Beijing Electron Positron Collider II (BEPCII)

- Beam energy: 1.0 – 2.3 GeV/c
- Design Luminosity: $10^{33} \text{ cm}^{-2}\text{s}^{-1}$
- Center of mass energy: ranging 2.0 - 4.6 GeV
- Circumference: 237 m

Physics of BESIII

Charmonium, D, τ , Light Hadron Spectroscopy and search for New Hadronic states

BEijing Spectrometer III (BESIII)

- **Drift chamber** (MDC), momentum resolution for charged particles is 0.5% at 1 GeV
- **Electromagnetic calorimeter** EMC, energy resolution* 2.5% and position resolution* 6 mm
- **Time of Flight system** (TOF), time resolution* 80 ps
- **Solenoid magnet** providing a 1.0 Tesla magnetic field
- **Muon Chamber System** (MUC) made of Resistive Plate Chamber
- **Geometrical acceptance** 93 % of 4π

*in the barrel

Physics at BESIII, Asner D. M. et al. Int. J. Mod. Phys A24 (2009) S1-794 arXiv: 0809.1869 [hep – ex]