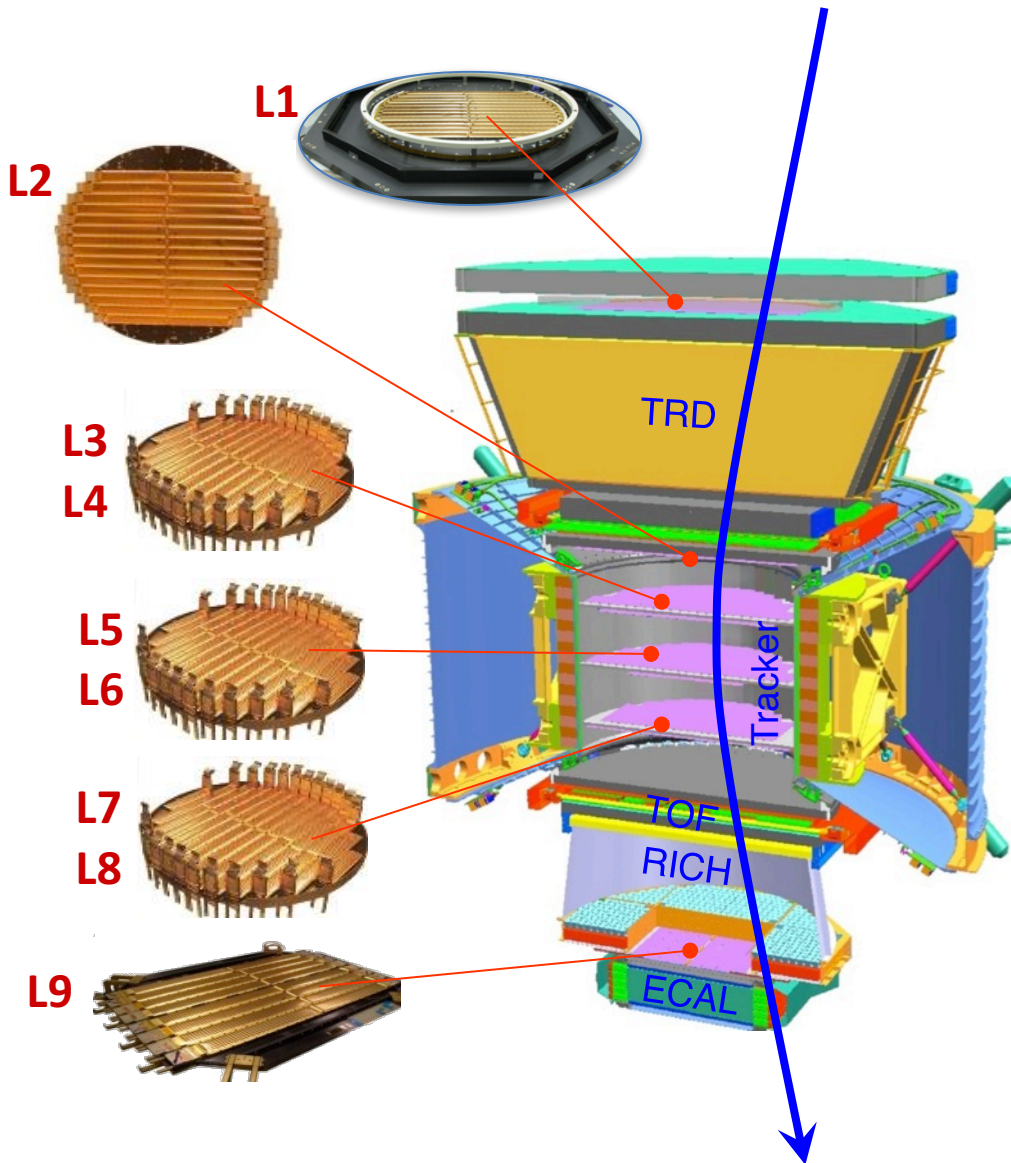


Strip Silicon Charge Detector Performance

A. Oliva

Istituto Nazionale di Fisica Nucleare, Sezione di Bologna

AMS-02 Silicon Tracker



9 layers of **300 μm** thickness **double-sided** silicon sensors arranged in 192 ladders.

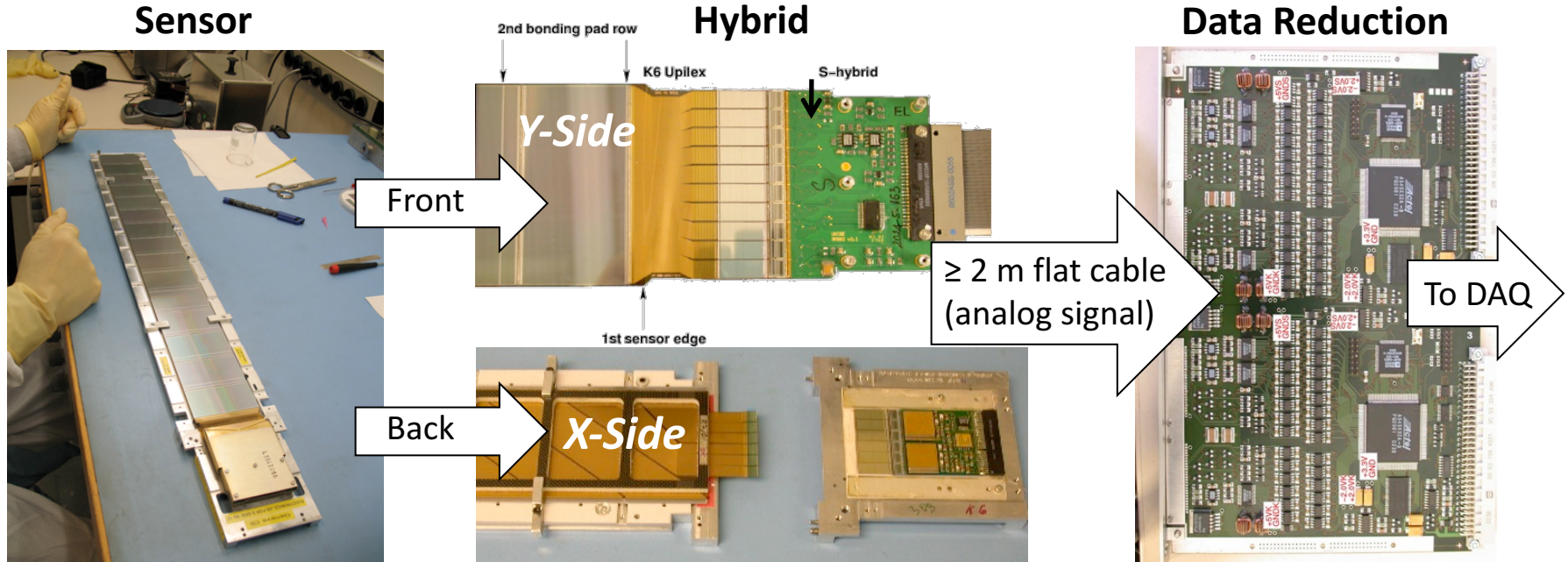
6 honeycomb carbon fiber planes. Overall detector material of about $0.04 X_0$.

A total of **196k channels**.

About **10 μm** for $Z=1$ and **5 μm** for $Z>1$ bending coordinate spatial resolution.

High dynamic range front end for charge measurement (Z measurement up to Iron).

AMS-02 Silicon Tracker Readout



300 μm , 7 \times 4 cm^2
27.5(104) μm strip impl. pitch
110(208) μm readout pitch
Charge sharing
Capacitive coupling (1 pF/cm)

640(384) readout channels
Amplification (100 MIP range)
Shaping (4 μs)
Sample-and-Hold
Each channel 0.7 mW power
10(6) VA_hdr64a

3 ADC
Pedestal/Noise eval.
Common noise sub.
Cluster search
Comp. factor of \sim 1000

AMS-02 Tracker Charge Calibration

Use of the first highest 5 strips of the cluster separately.

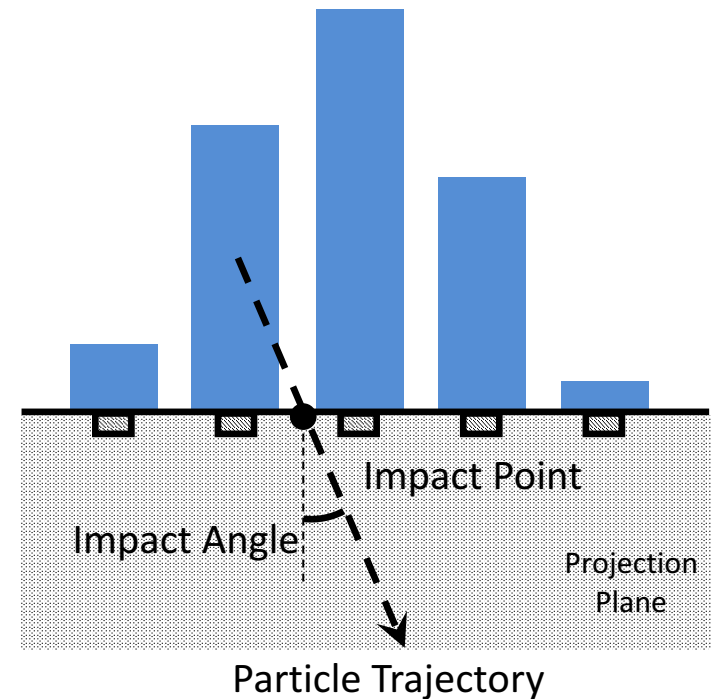
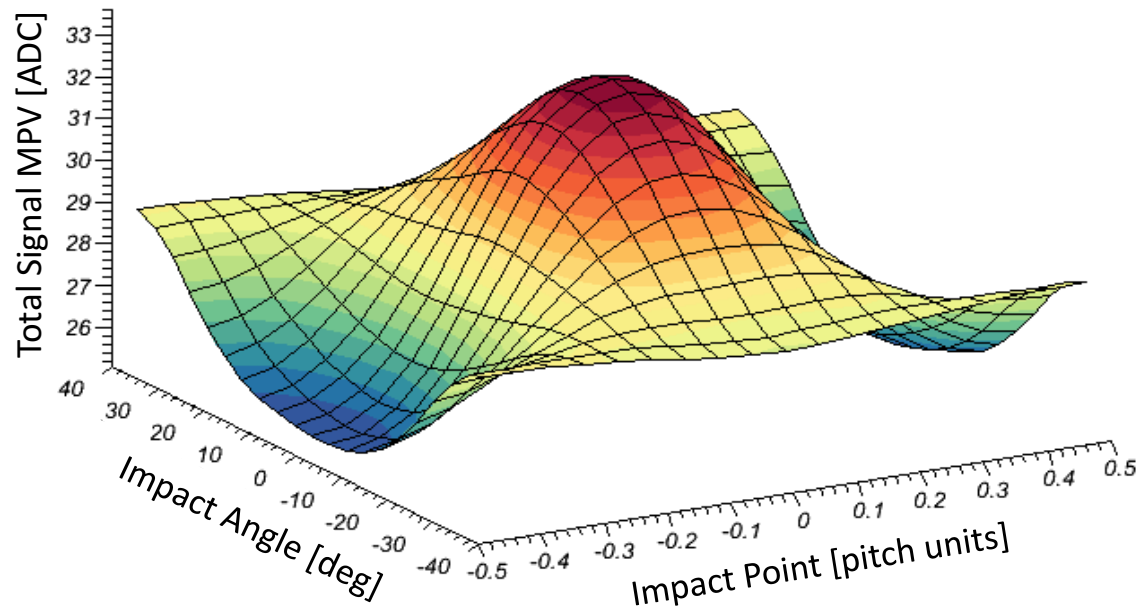
Equalization of 3072 VAs as function of charge (non-linearities).

Path-length correction.

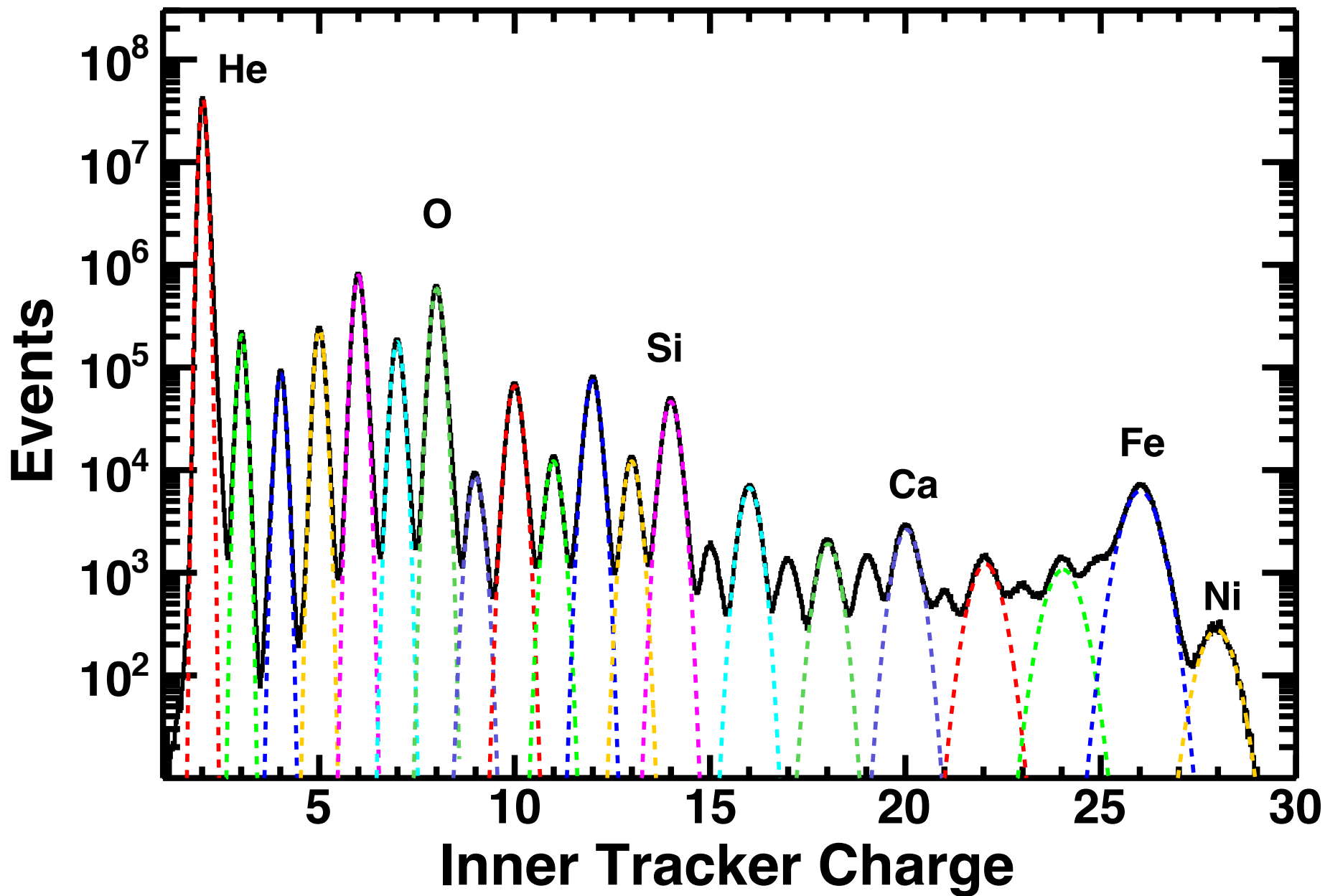
Correction of charge collection effects (coordinate and angular dependences).

Description of the dependence with energy ($\sim 1/\beta^2$).

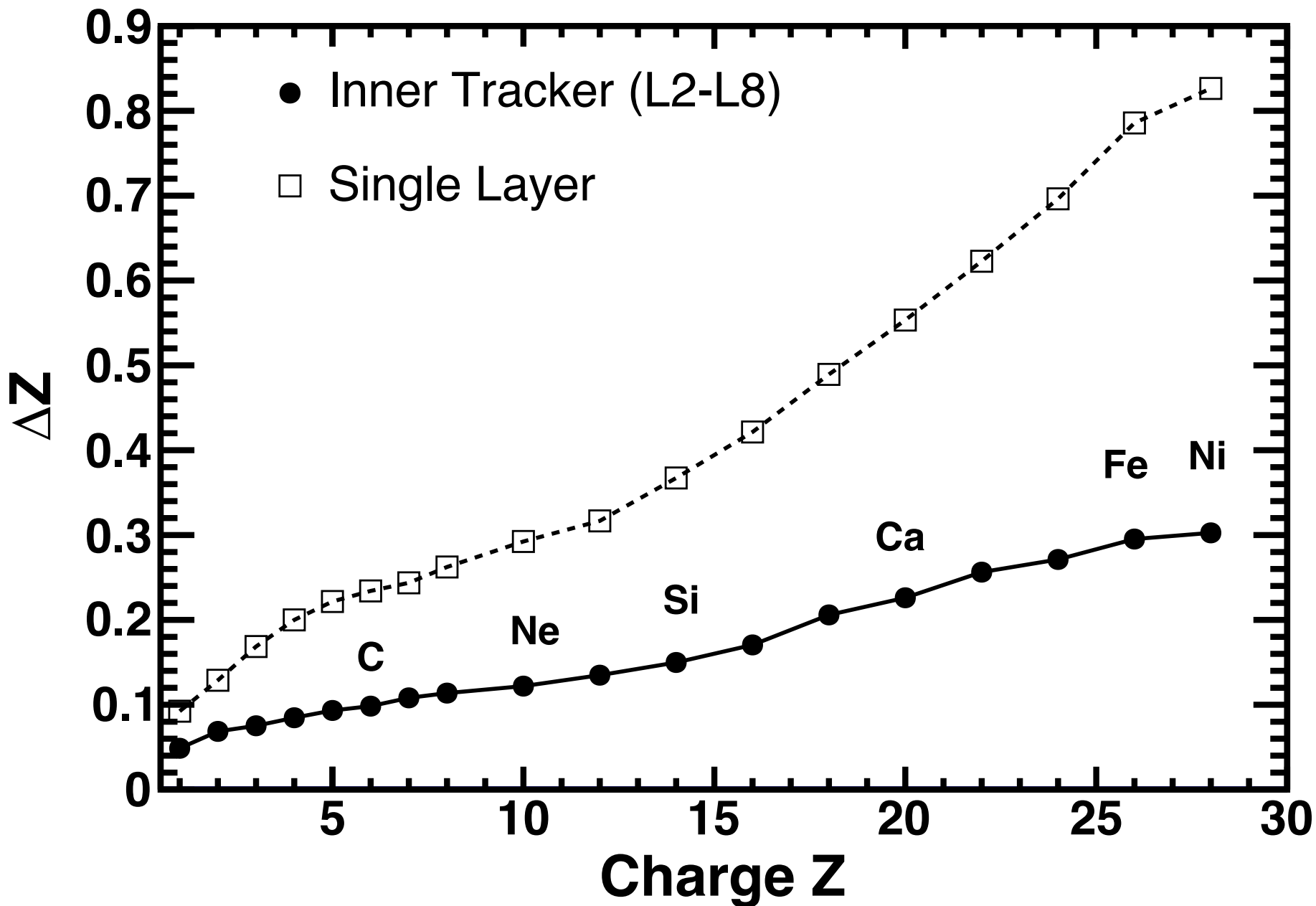
Time variation correction.



AMS-02 Inner Tracker Charge (L2-L8)

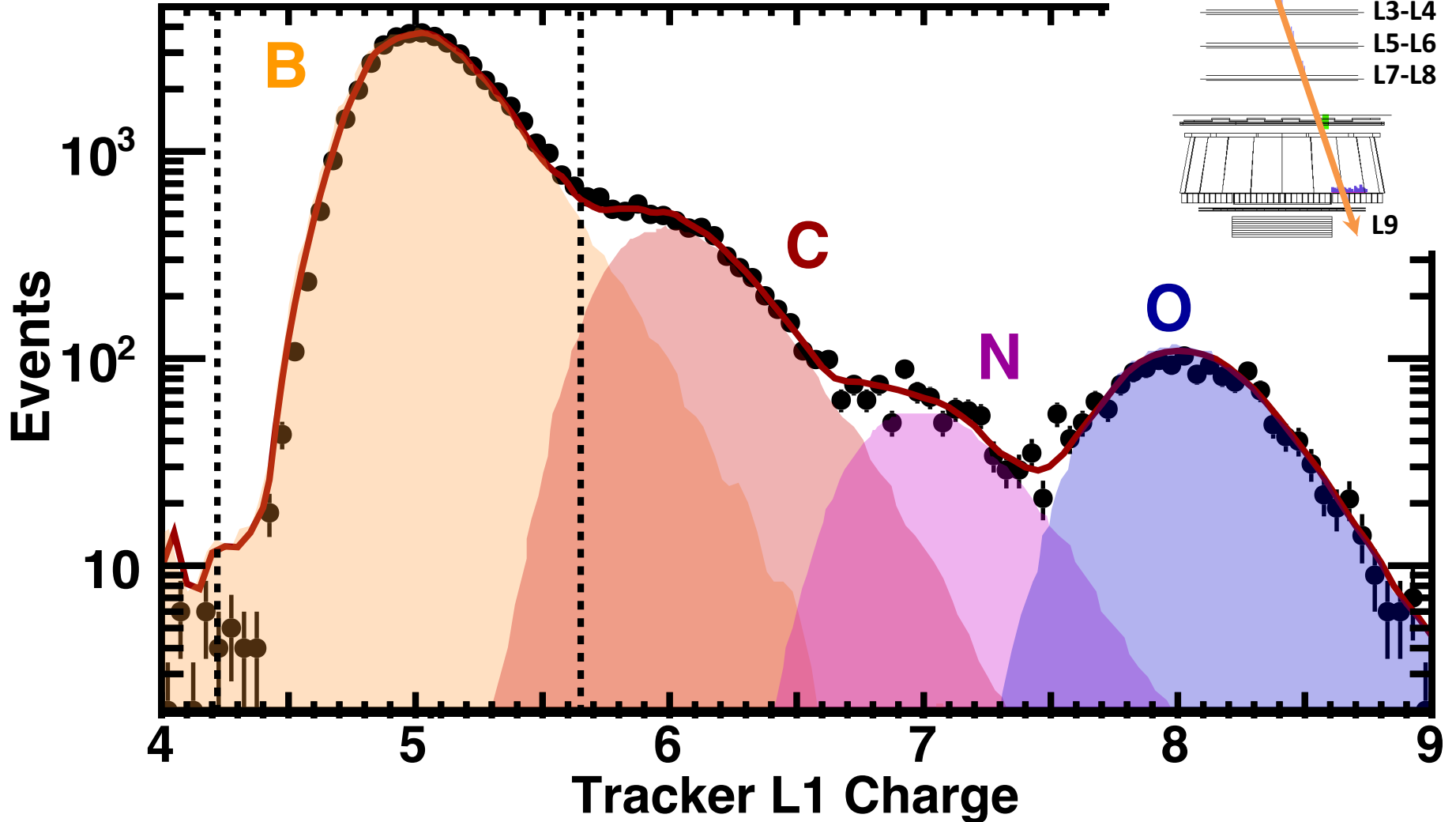


AMS-02 Tracker Charge Resolution



The Importance of Measuring Charge at the Top-of-the-Instrument

L1 is used for *background estimation* (interaction inside AMS), and for *MC interaction tuning* (selecting samples of incident nuclei).



Energy Loss (from Particle Data Book)

$$k = \frac{\xi}{W_{\max}}$$

Vavilov parameter

$$\xi = \frac{K}{2} \frac{Z}{A} \frac{z^2}{\beta^2} \rho x \quad \xi[\text{MeV}] = 0.1535 \frac{K}{2} \frac{Z}{A} \frac{z^2}{\beta^2} X[\text{g/cm}^2]$$

Landau width

$$W_{\max} = \frac{2m_e c^2 \beta^2 \gamma^2}{1 + 2\gamma \frac{m_e}{M} + \left(\frac{m_e}{M}\right)^2} \stackrel{m_e \gamma \ll M}{\approx} 2m_e c^2 \beta^2 \gamma^2$$

Maximum energy transfer to e^-

For Silicon below 1 mm Silicon, $\beta \approx 1$, and impinging nuclei, **k is below 0.1**, corresponding to **Landau regime**.

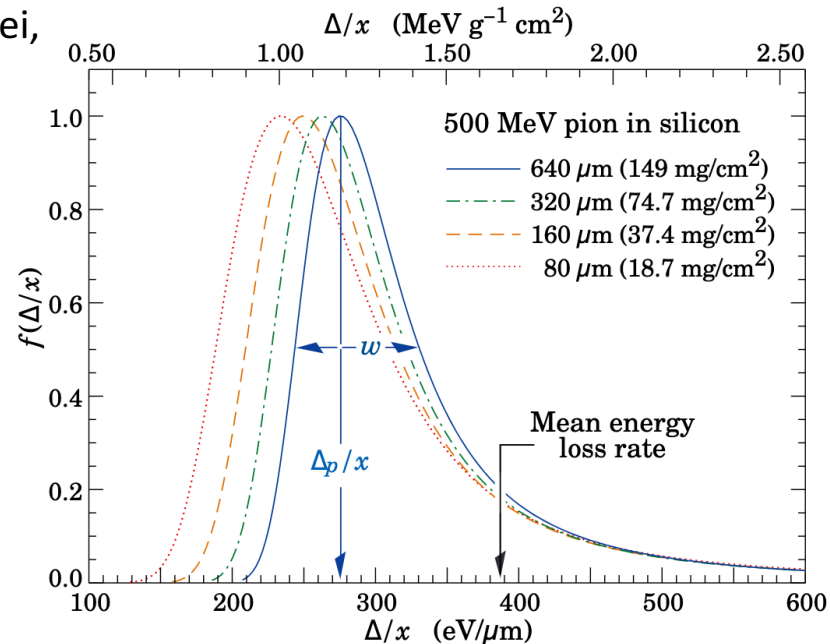
Landau most probable energy loss:

$$\Delta_p = \xi \left[\ln \frac{2m_e c^2 \beta^2 \gamma^2}{I} + \ln \frac{\xi}{I} + 0.2 - \beta^2 - \delta(\beta\gamma) \right]$$

Landau full-width at half maximum (FWHM):

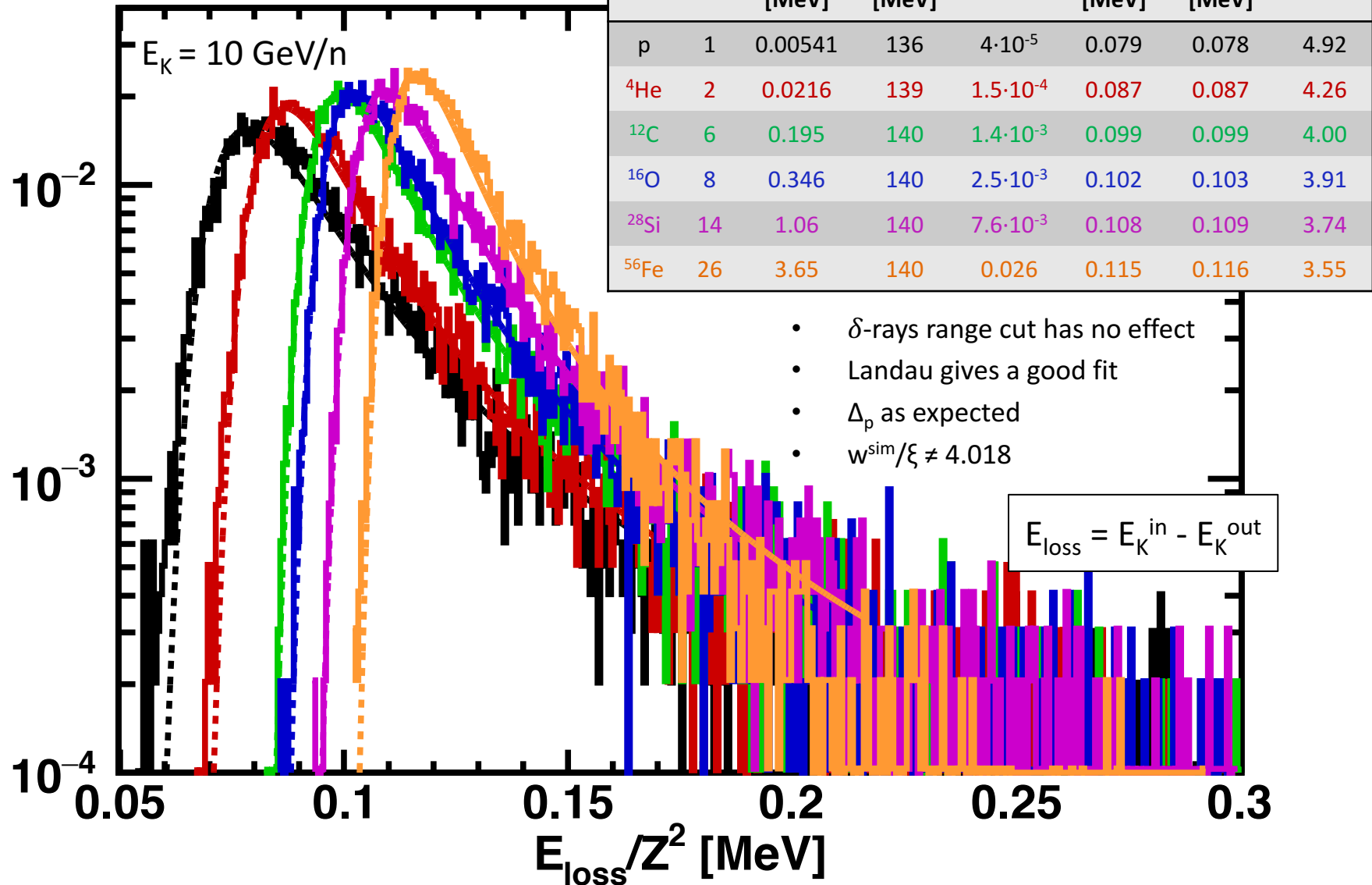
$$w_L \sim 4.018\xi$$

→ Verification with MC

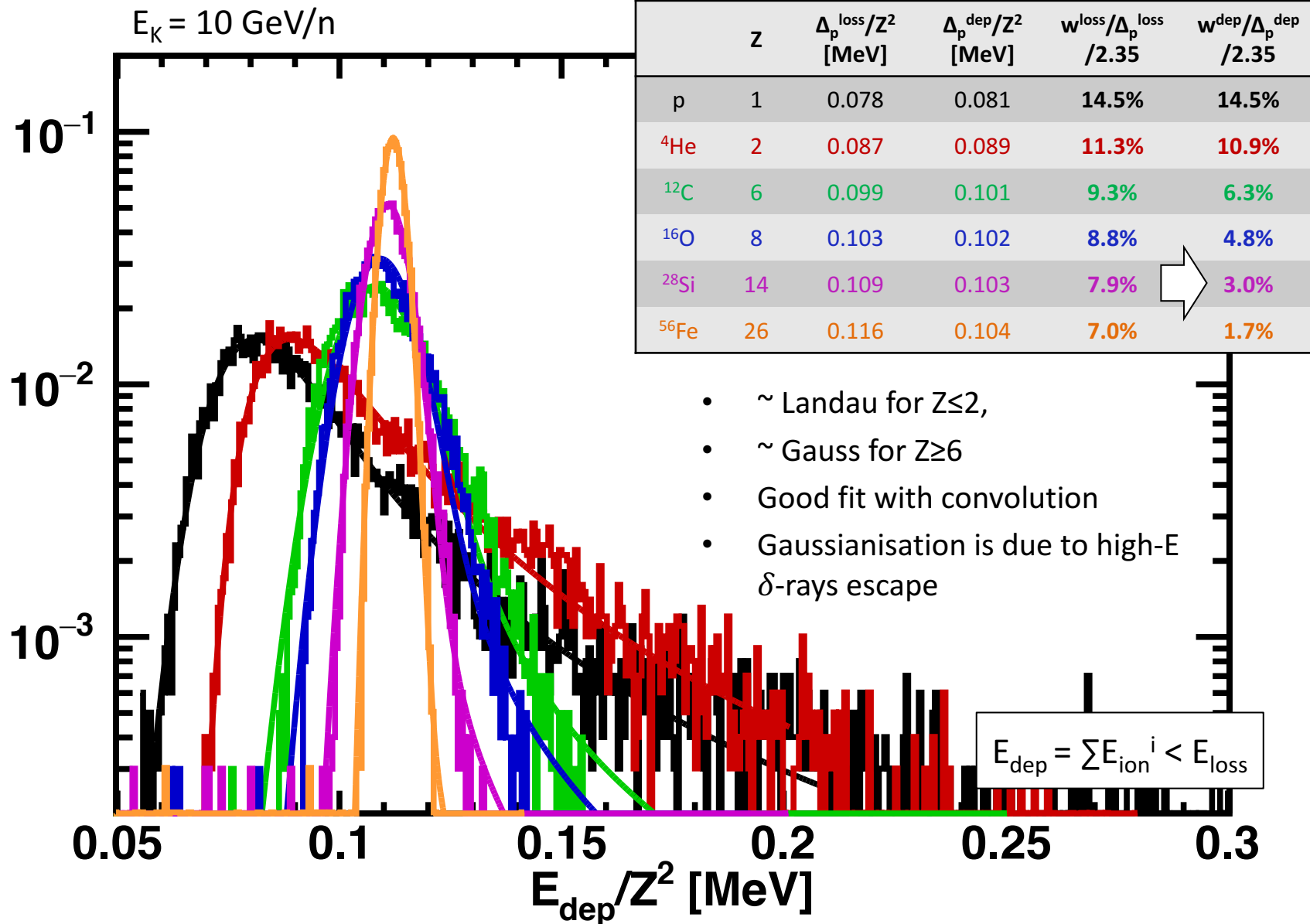


Energy Loss in 300 μm Silicon

From a simple GEANT4 simulation.

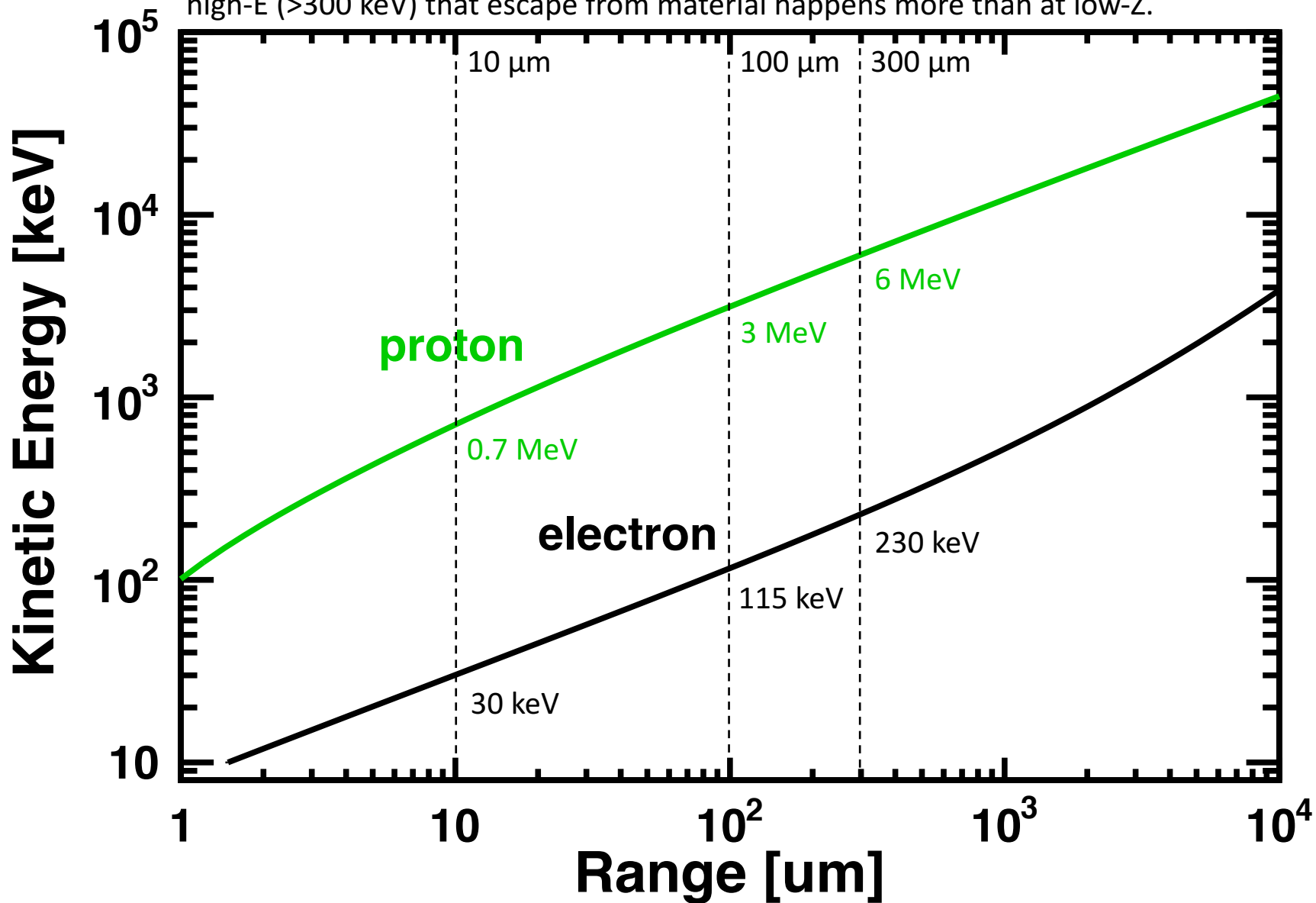


Energy Deposition in 300 μm Silicon



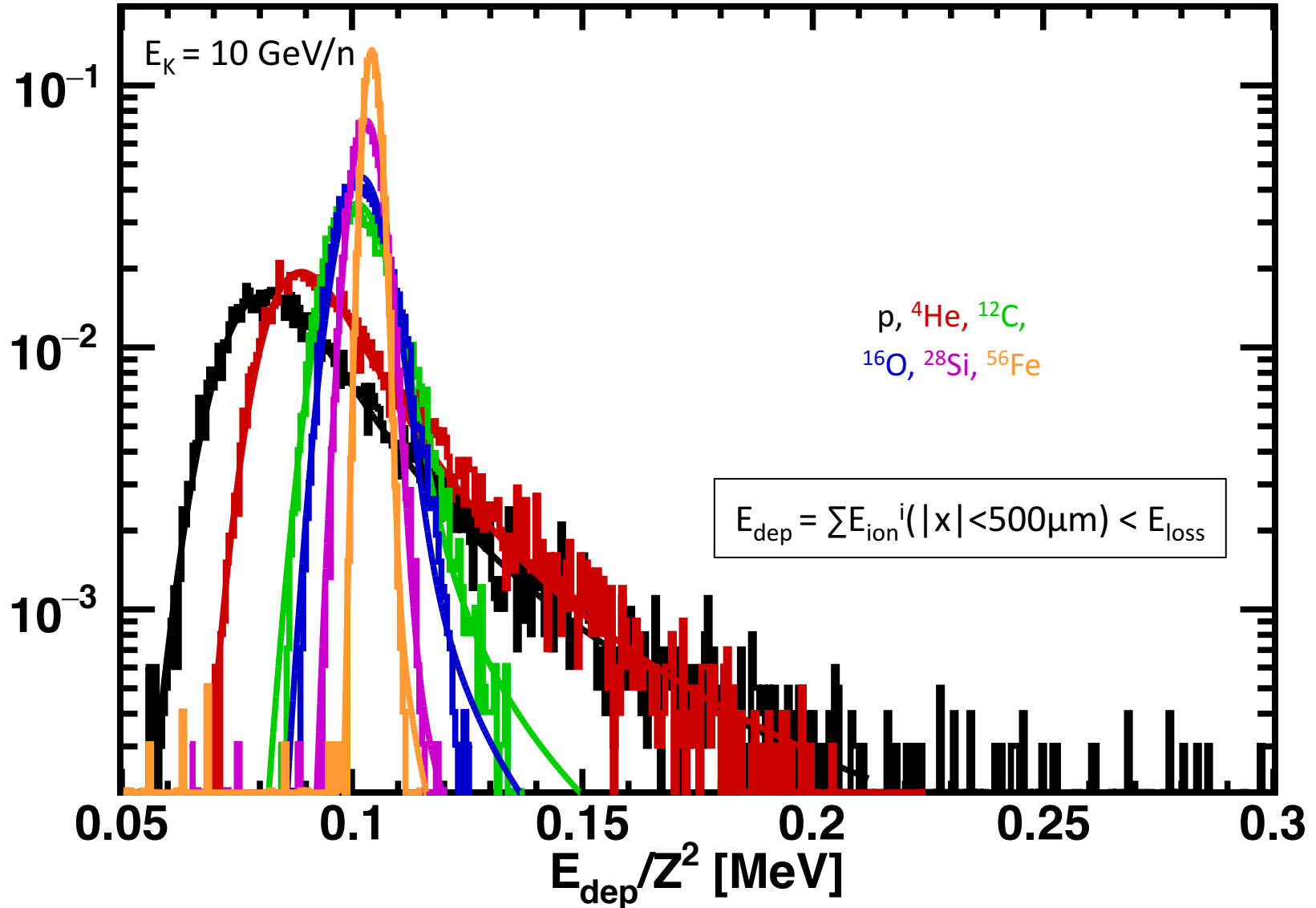
δ -Rays Escape

δ -rays production cross section $d\sigma/dE \approx kZ^2/E^2$. At high-Z emission of δ -rays of high-E (>300 keV) that escape from material happens more than at low-Z.



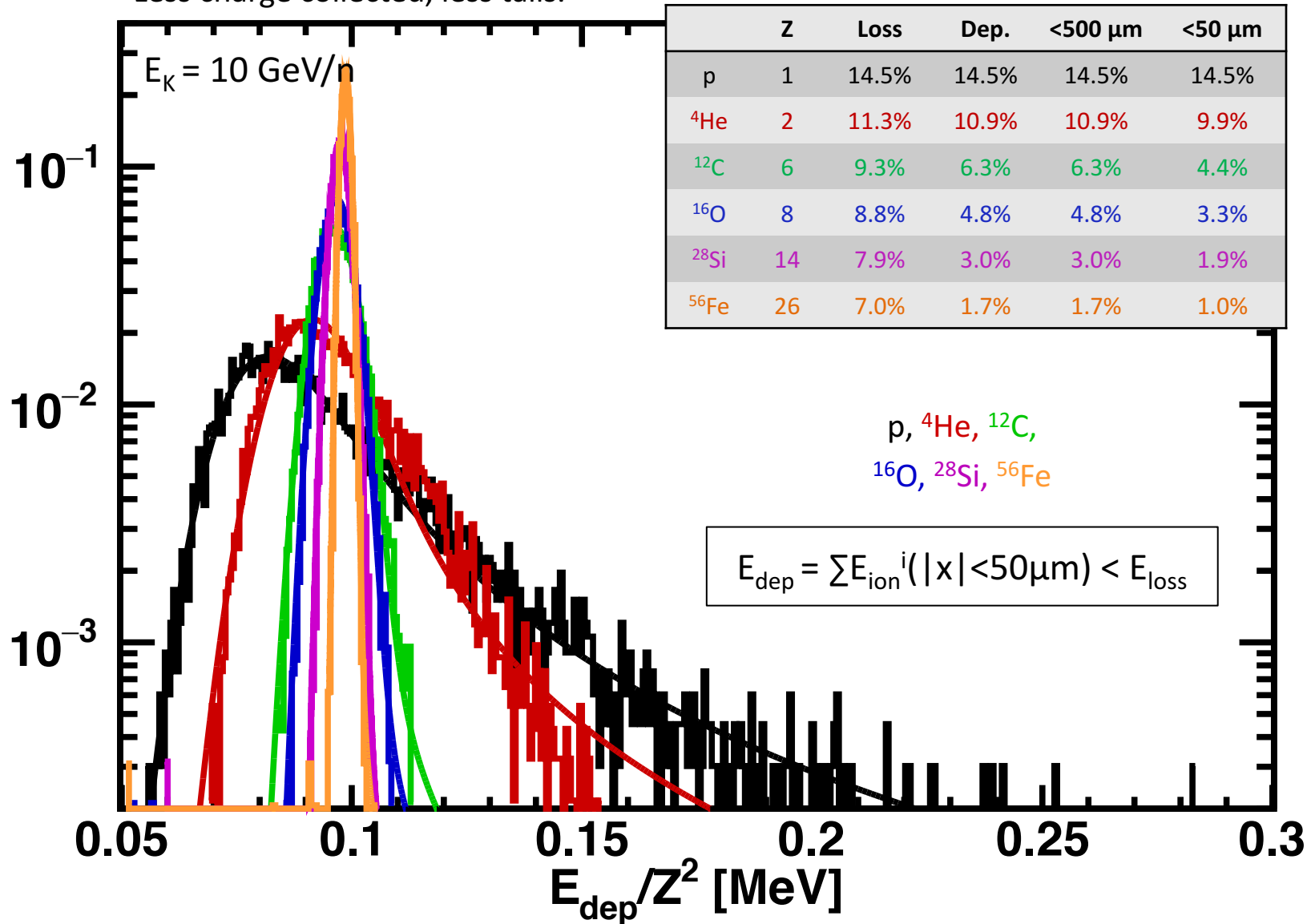
Energy Loss in 300 μm Silicon (“cluster”)

Pretty similar to previous.



Energy Loss in 300 μm Silicon (“strip”)

Less charge collected, less tails.



Charge Measurement and Resolution

$$E = k \frac{z^2}{\beta^2} f(\beta\gamma) = K z^2$$

$$\sigma_E = \frac{dE}{dz} \sigma_z = 2K z \sigma_z$$

$$\frac{\sigma_z}{z} = \frac{1}{2} \frac{\sigma_E}{E}$$

$$\left(\frac{\sigma_E}{E}\right)_{\text{stra}} = \frac{1}{2.35} \frac{w_L}{\Delta_p} \propto \frac{1}{\sqrt{z}}$$

Straggling
(from simulation)

$$\left(\frac{\sigma_E}{E}\right)_{\text{stat}} = \sqrt{\frac{WF}{E}} \propto \frac{1}{z}$$

E/W = number of carriers (W = 3.6 eV for Si)
F = Fano's factor (0.12 for Si)

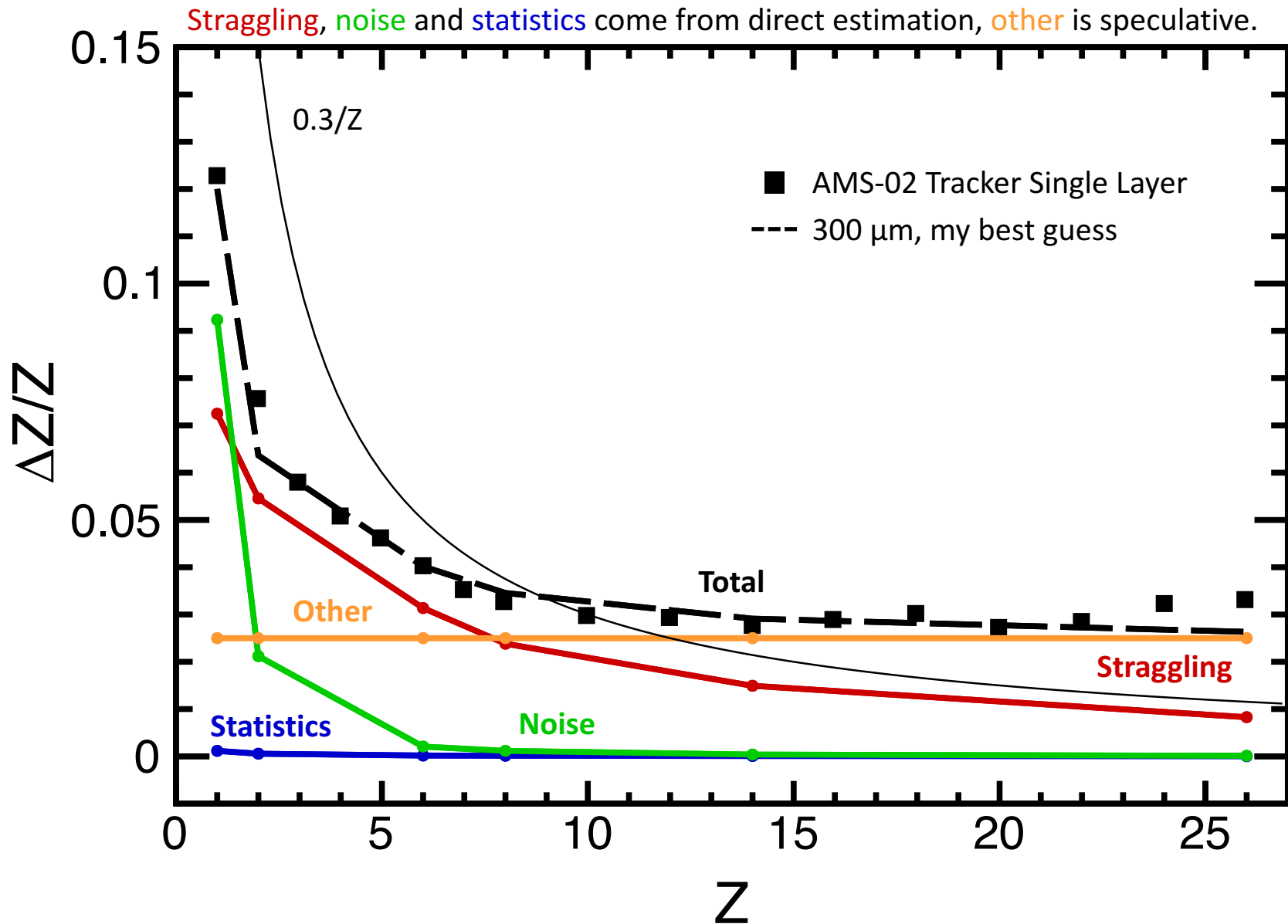
Statistical
(below % even with 50 μm)

$$\left(\frac{\sigma_E}{E}\right)_{\text{noise}} = \frac{c_{\text{noise}}}{E} \propto \frac{1}{z^2}$$

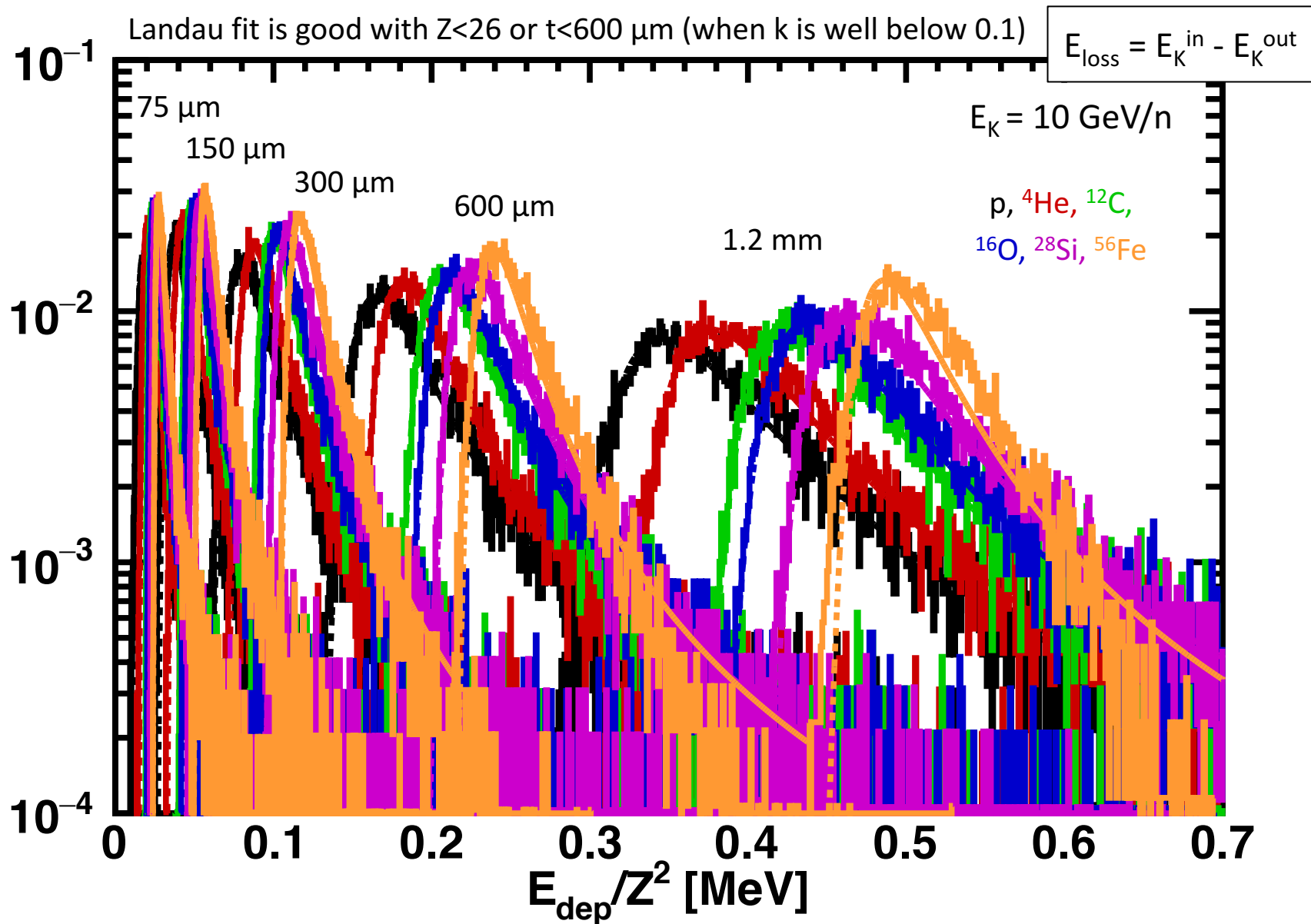
Noise
(depends on S/N)

Other: systematics due *detector effects*, that can be *irreducible* (saturation, space charge, sample-and-hold) or that can be *improved* with better understanding of the detector (amplification non-linearities, angular and spatial dependencies, ...). For the latter, large statistics at high-Z may help.

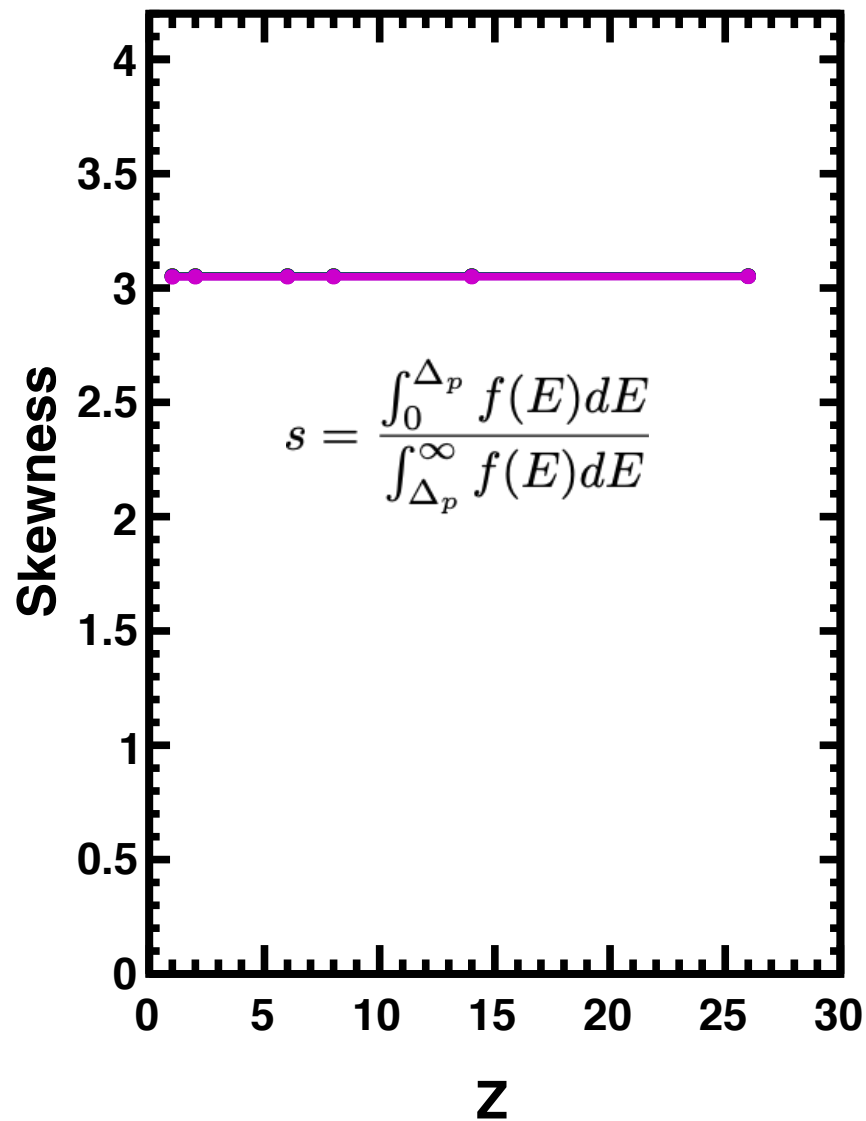
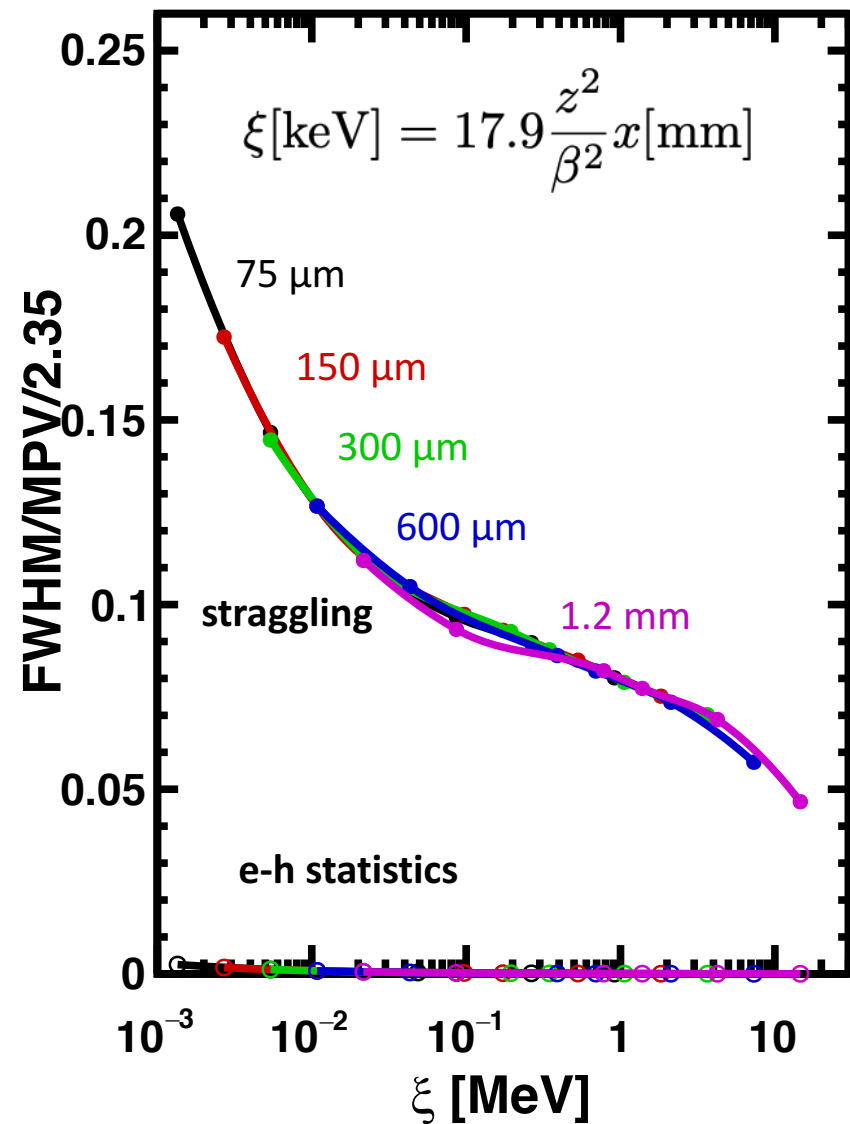
AMS-02 Single Layer Charge Resolution



Single Layer Energy Loss for Different Thicknesses

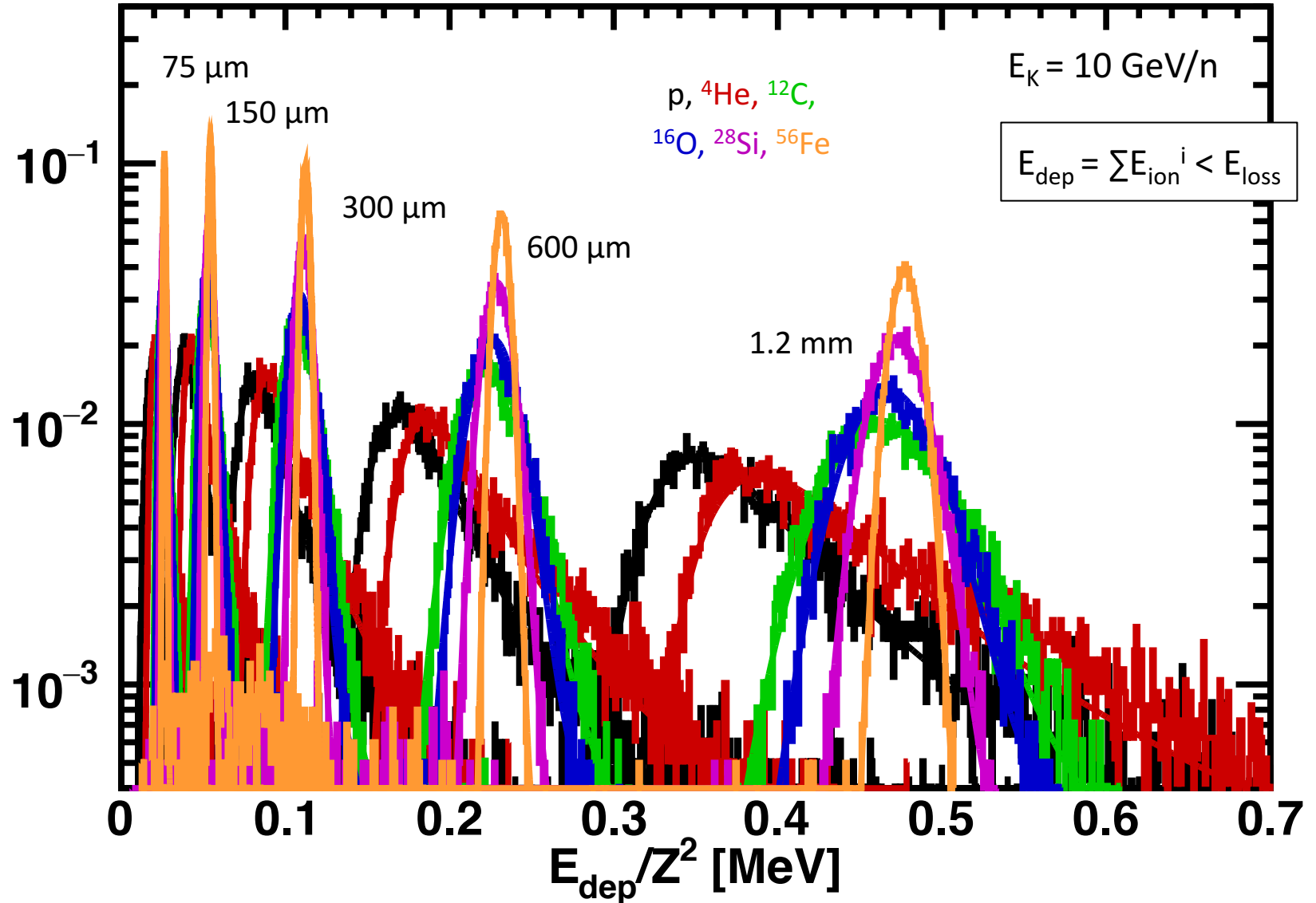


Single Layer Energy Loss for Different Thicknesses



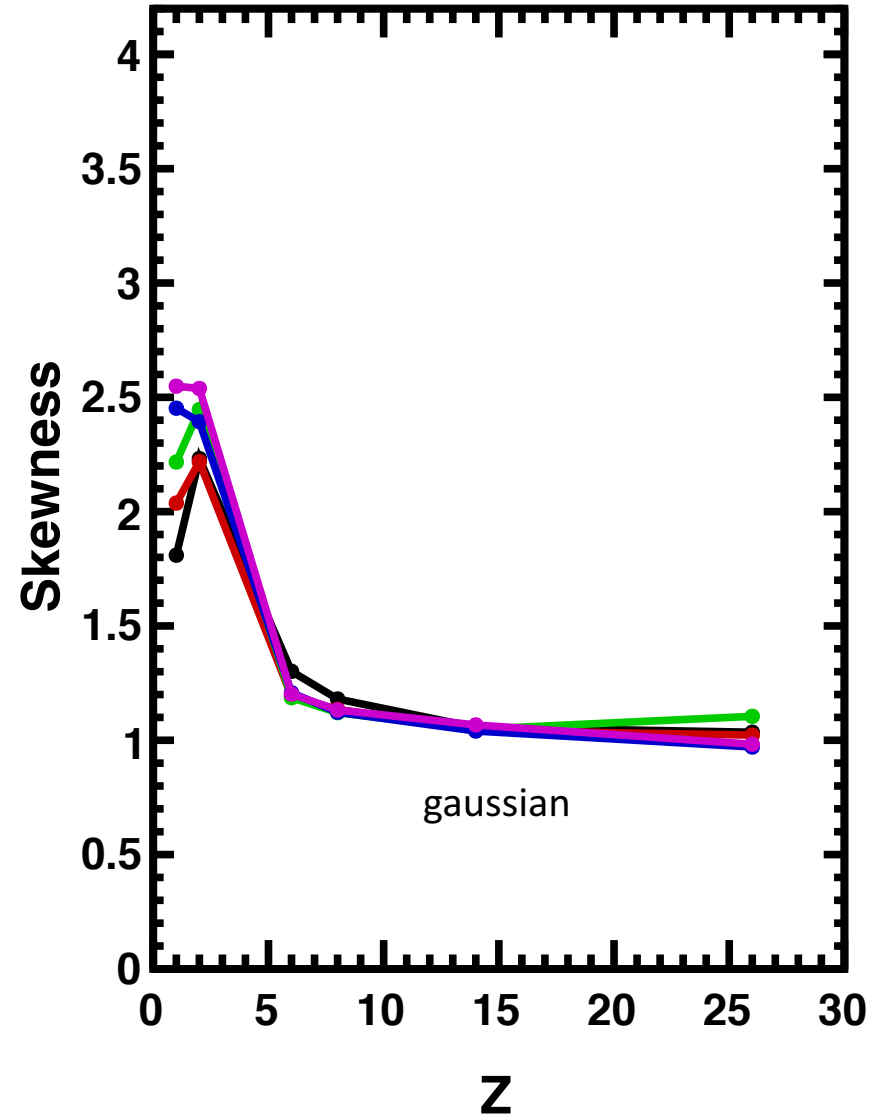
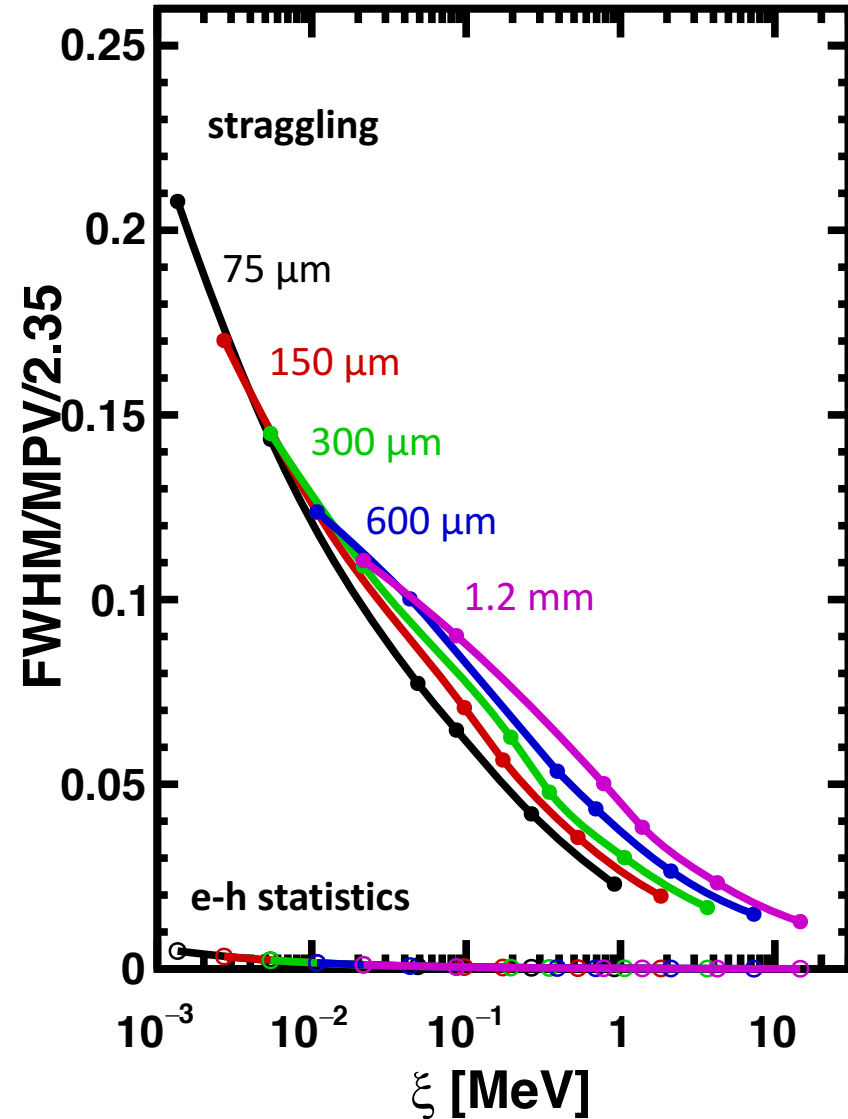
Single Layer Energy Deposit inside 500 μm for Different Thicknesses

Gaussian above $Z=6$, independently from thickness (high-E δ -rays escaping).



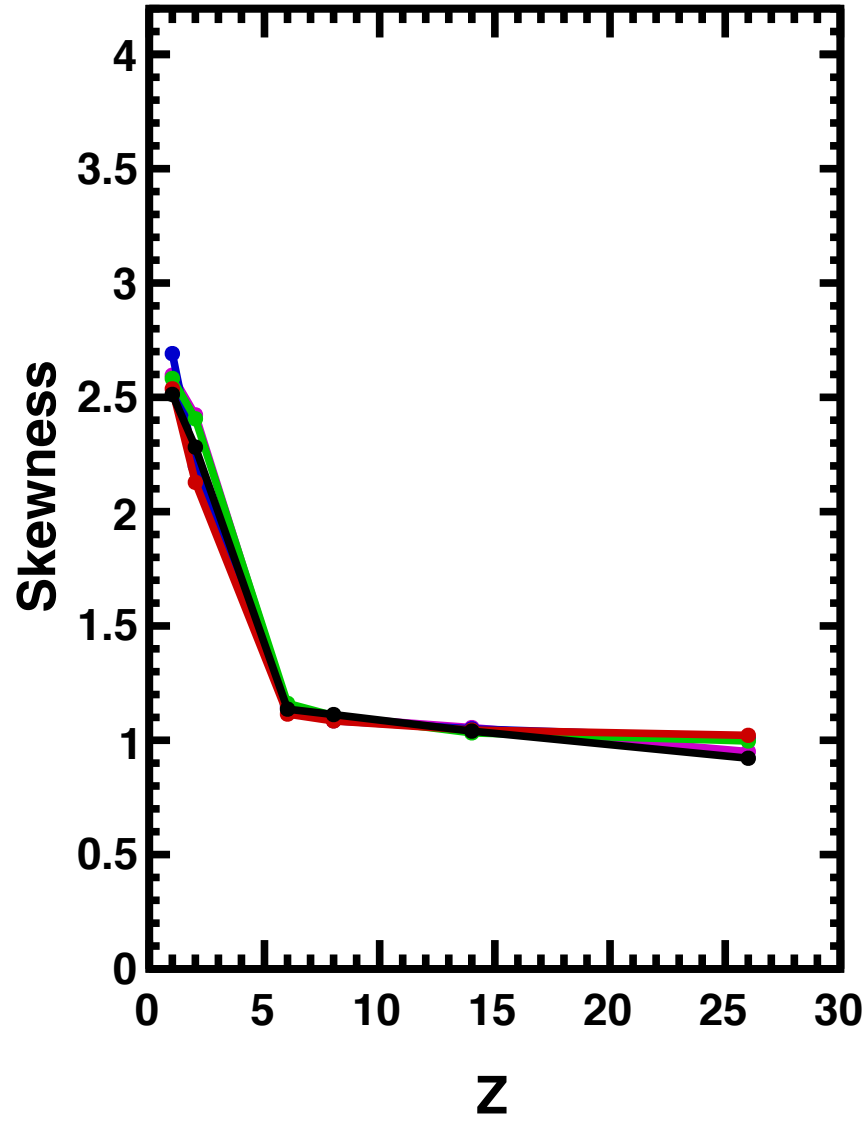
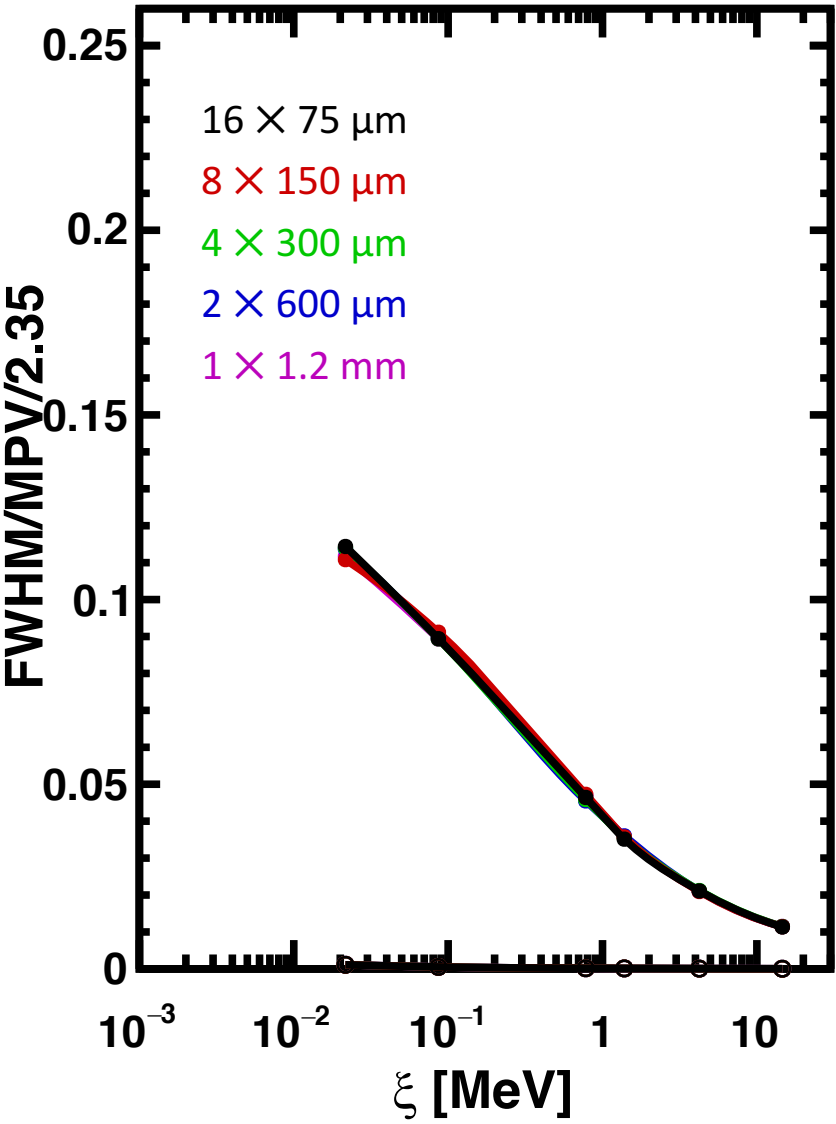
Single Layer Energy Deposit inside 500 μm for Different Thicknesses

Similar things can be seen for integral and $<50 \mu\text{m}$ energy deposit.



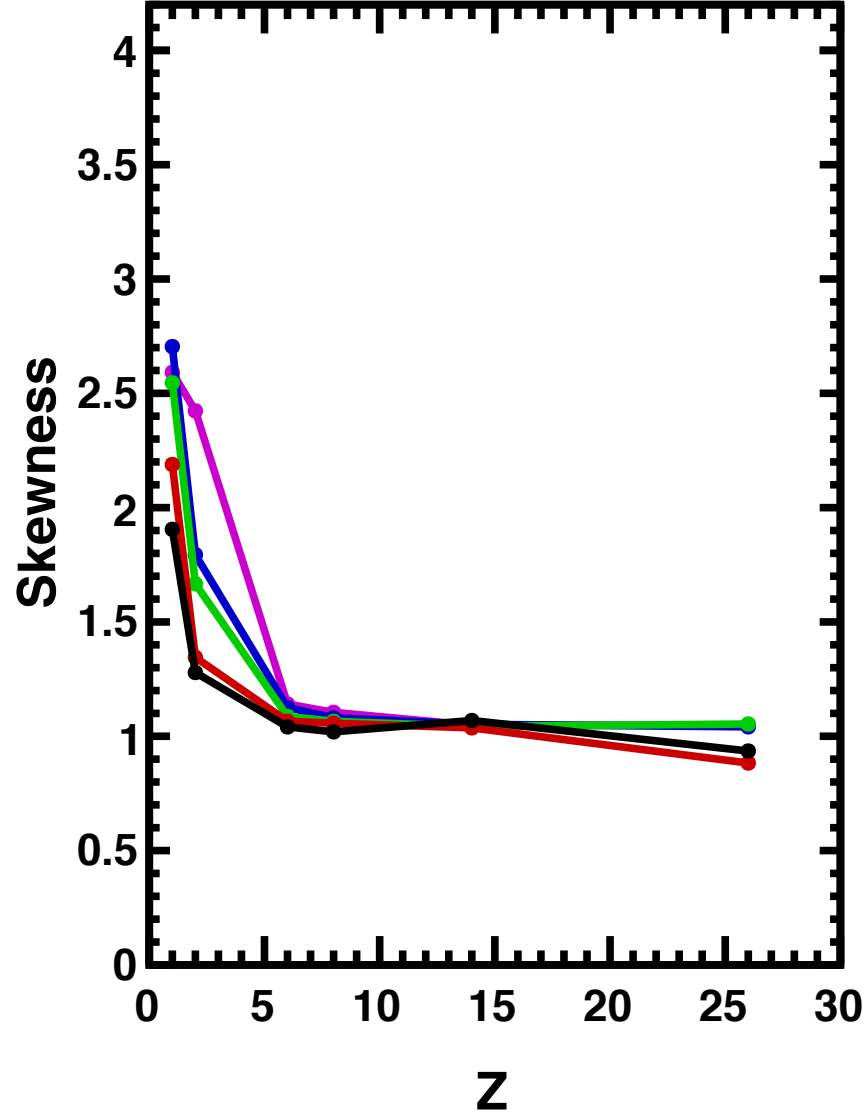
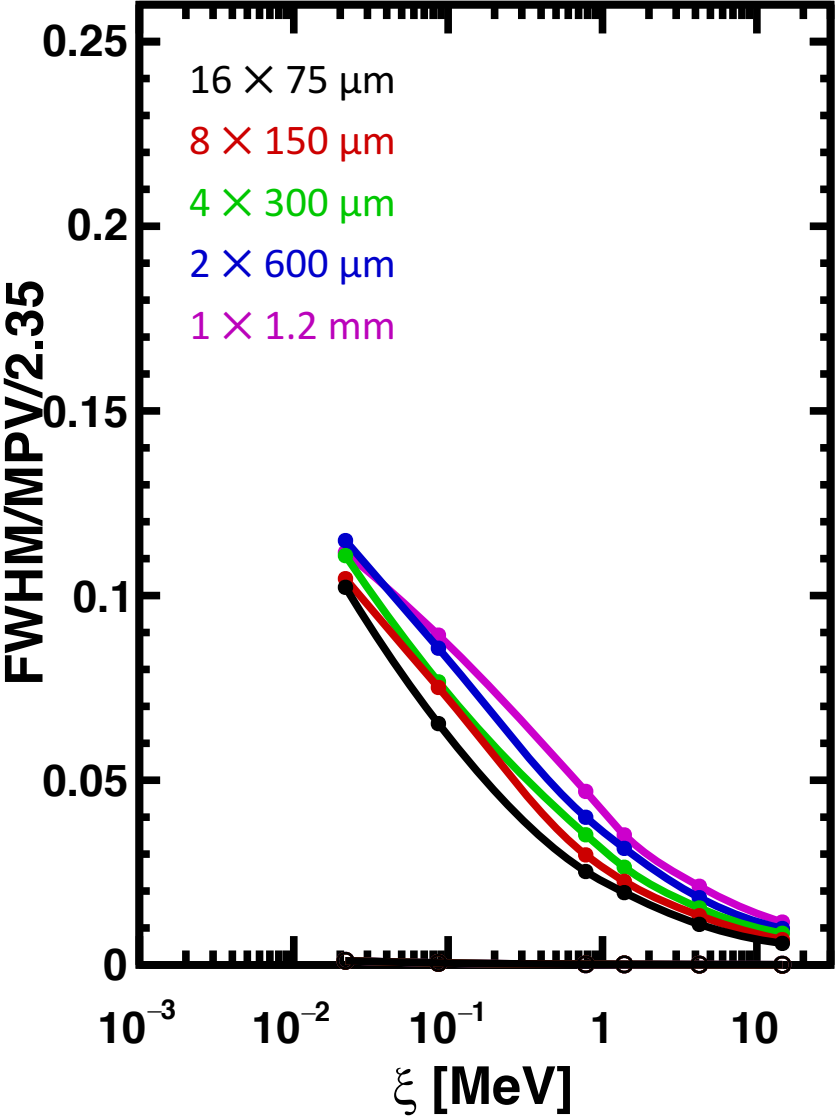
Energy Deposit inside 500 μm for the Same Overall Thickness of 1.2 mm

As expected, just by simple segmentation, there is no difference for any energy deposition/loss.



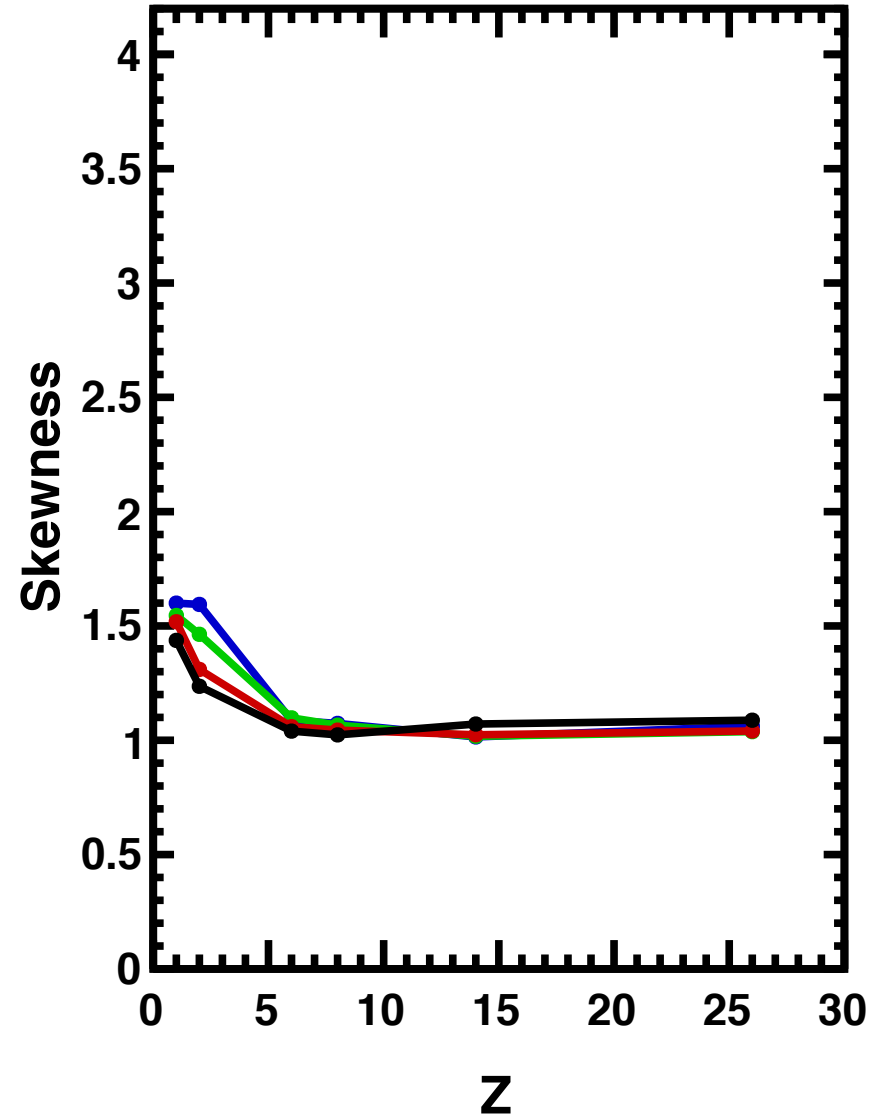
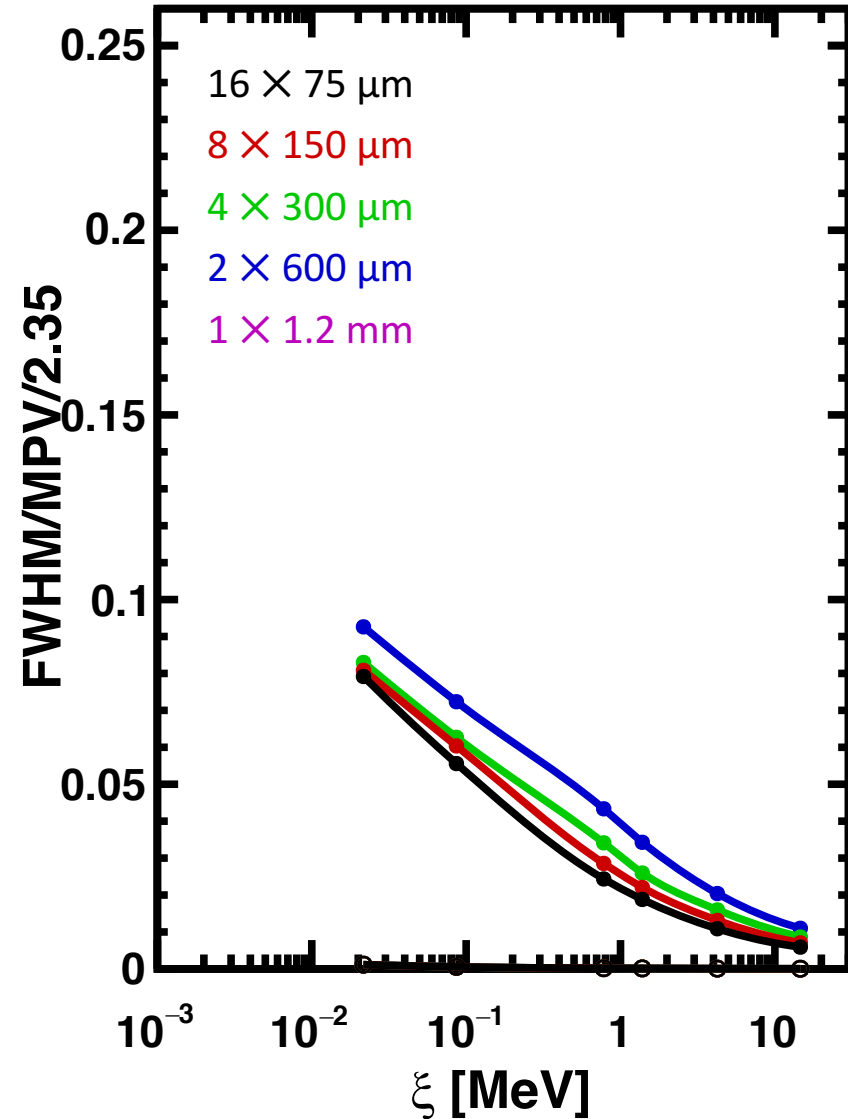
Energy Deposit inside 500 μm for the Same Overall Thickness of 1.2 mm

Some improvement arise if there is a gap, here of 1 cm (reduce δ -rays correlations).



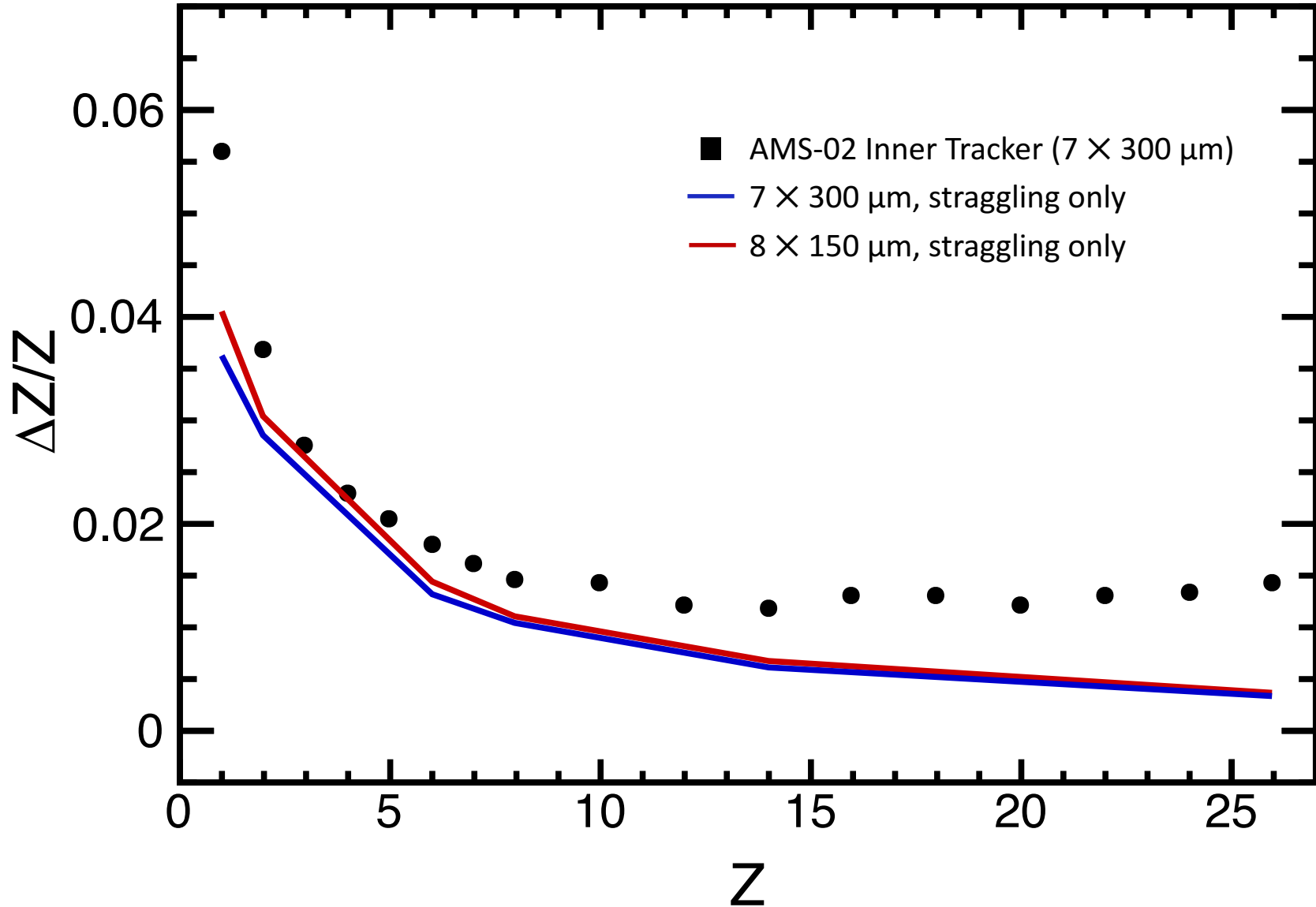
Energy Deposit inside 500 μm for the Same Overall Thickness of 1.2 mm

Additional improvement comes from the use of truncated mean (use of likelihood ...).



AMS-02 Inner Tracker Charge Resolution

Even if the use of $8 \times 150 \mu\text{m}$ has an intrinsic lower performance, it seems adequate.



Conclusions

A simple simulation able to give a baseline charge performance has been developed and verified with AMS-02 data.

At low-Z important contribution from straggling and noise.

At high-Z systematics could be dominant. It is important to plan ion beam test campaigns to check the linearity of the system and understand charge dependence from impact angle and coordinate. Also a uniformity of the response of the front-end electronics can be important for an accurate calibration.

A thin SCD with 8 layers of 150 μm has been studied and has an intrinsic resolution that is adequate (to be confirmed in test beam).

A more detailed performance evaluation of different scenarios will be carried out with HERD simulation.