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MHD-Particle-in-cell Method and its Astrophysical Applications

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Outline

- Motivation and MHD-PIC method
- Application 1: CR acceleration in collisionless shocks
- Application 2: CR streaming instability
- Future directions

How do CRs interact with a thermal plasma?

CRs are collisionless and diffuse by scattering off MHD waves/turbulence:

Galactic CRs' residence time: ~10s Myr in total. Diffusion coefficient: $\kappa \sim R^2/T \sim a$ few x 10²⁸cm²s⁻¹. (e.g., Ginzburg & Syrovatskii 64)

Passive "test particles"

- CRs provide pressure support: $F = -\nabla_{\perp} P_{CR} = -\frac{J_{CR} \times B}{c}$
- When the bulk CRs drift through background plasma faster than the Alfven speed $v_{D,CR} > v_A$, they will drive *streaming instabilities*.

(e.g., Kulsrud & Pearce 1969, Bell 2004)

Need kinetic physics: CRs transfer energy/momentum to gas via Alfven waves.



Simulating CR physics at kinetic level

- Minimum requirement: resolve CR gyro-radii.
- Huge scale separation involved, challenging for conventional PIC codes:
 - Full-PIC: treat all (background+CR) particles as kinetic particles
 - Hybrid-PIC: all ions (background+CR) are kinetic, electrons as massless fluid



Simulating CR physics at kinetic level

- Minimum requirement: resolve CR gyro-radii.
- Overcome this scale separation: skipping over the kinetic scales of the background plasma.
 - MHD-PIC: treat background plasma by MHD, while CRs are kinetic



MHD-PIC approach

- Each computational particle (i.e., superparticle) represents a large collection of real CR particles.
- Each super-particle carries an effective shape, designed to facilitate interpolation from the grid.
- Individual CR particles move under the electro-magnetic field from MHD.
- Total momentum and energy must conserve: particles feedback to MHD cells by depositing changes in momentum and energy locally.



MHD-PIC: formulation and implementation

Equations for the (relativistic) CR particles:

$$\frac{d(\gamma_j \boldsymbol{u}_j)}{dt} = \frac{q_j}{m_j} \left(\boldsymbol{E} + \frac{\boldsymbol{u}_j}{c} \times \boldsymbol{B} \right)$$

Specify the numerical speed of light c >> any velocities in MHD.

Full equations for the gas:

$$\frac{\partial
ho m{v}}{\partial t} +
abla ullet (
ho m{v}m{v} - m{B}m{B} + m{P}^*) = -$$
 Lorentz force on the CRs

 $\frac{\partial E}{\partial t} + \nabla \cdot \left[(E + P^*) \boldsymbol{v} - \boldsymbol{B} (\boldsymbol{B} \cdot \boldsymbol{v}) \right] = - \text{ energy change rate on the CRs}$

Implemented in the Athena MHD code (Bai, Caprioli, Sironi & Spitkovsky 2015).

(See also van Marle et al. 2018, Mignone et al. 2018)

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Non-relativistic shock: MHD-PIC simulation



Fiducial parameters: $M_A \sim 30$, parallel shock $\theta = 0$.

Resolution: 12 ion skin depths per cell (*v.s. 0.5 in hybrid-PIC*) **Particle injection:** artificial (as proof-of-concept) $\eta = 2 \times 10^{-3}$ (to be improved in the near future)

Particle acceleration





Simulation with relativistic particles

Set numerical speed of light *c* a factor ~10-20 larger than $v_{\rm sh}$ to follow particle acceleration to relativistic regime.



Very large box size (4800 $c/\omega_{\rm pi}$ wide), and very long evolution (~10⁵ Ω_c^{-1})

Reduction of shock speed toward later evolution.

Particle acceleration into relativistic regime $t=11088\Omega_c^{-1}$



f(p)~p⁻⁴ through the transition + a drop in normalization.

Bai+ 2015

Evolution of maximum particle energy



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Resonant CR streaming instability (CRSI)

When the bulk drift of the CRs exceed the Alfven velocity, they resonantly trigger the CR streaming instability.

Characteristic growth rate:

$$\Gamma(k) \approx \Omega_c \frac{N_{\rm CR}(p > p_{\rm res}(k))}{n_i} \frac{v_D - v_A}{v_A}$$

Driven primarily by low-energy CRs (the dominant population by number)

(e.g., Kulsrud & Pearce 69, Skilling 75)

Important feedback mechanism to galaxy formation and evolution:

- CR self-confinement
- CR-drive galactic outflows

Matching analytical dispersion relation



Accurately reproduce the linear growth rate over broad spectrum.

Bai et al. 2019

1D simulation: growth and saturation



Simulation performed in the rest frame of the CRs.

Periodic BC.

Gas travels to the left at v_D .

Fiducial parameters:

 $v_{D}=2v_{A}$ $N_{CR}/n_{0}=10^{-4}$ $p_{0}/m=300v_{A}$

2048 ppc, Lx~50 most unstable wavelength.

Bai et al. 2019

Towards saturation: quasi-linear diffusion

wave intensity

$$\frac{\partial f_w}{\partial t} = \frac{\partial}{\partial \mu_w} \left[\frac{1 - \mu_w^2}{2} \nu(\mu_w) \frac{\partial f_w}{\partial \mu_w} \right] + \text{ reflection} \qquad \qquad \text{Scattering frequency:} \\ \nu(\mu_w) = \pi \Omega k_{\text{res}} I(k_{\text{res}})$$

Parameters: $v_D=2v_A$; $N_{CR}/n_0=10^{-3}$



Reduction of CR drift speeds



Future directions

- Incorporate wave damping physics
- Incorporate multi-phase ISM
- Other gyro-resonant instabilities
- CR source problem
- CR escape
- Towards 2D/3D

Summary

- Motivation and development of MHD-PIC method
 - To study kinetic aspect of CRs interacting with background plasmas
 - PIC for CRs, MHD for background plasmas, valid on scales > ion skin depth, implemented to Athena MHD code (fully conservative).
- MHD-PIC simulations of CR acceleration
 - Reproduce hybrid-PIC results using much larger box at much reduced cost, and can follow CR acceleration to relativistic regime.
 - Need (artificial) injection prescription.
- MHD-PIC simulations of resonant CR streaming instability
 - > Overcome the challenges: developed δf method to reduce noise
 - First numerical study: confirmation of linear growth rates, and can follow CR quasi-linear evolution, overcome 90deg problem.
 - Future: more realistic microphysics, CR escape...