NEUTRINO FLOOR IN DM DIRECT DETECTIONS

WEI CHAO BEIJING NORMAL UNIVERSITY

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We study impacts of non-standard neutrino interactions to the
neutrino floor

NSI	Enhancement	Estimated values
Vector	\checkmark	~several times
Axial-vector	×	×
Tensor	×	×
Scalar	\checkmark	~several times
Pseudo-scalar	\checkmark	~30%

Wei Chao, J. Zhang, X. Wang and X. Zhang, JCAP1908,010

Evidence of DM





DM incidents, 1907.06674

Death by Dark Matter

Jagjit Singh Sidhu¹, Robert J. Scherrer², Glenn Starkman¹ Physics Department/CERCA/ISO Case Western Reserve University Cleveland, Ohio 44106-7079, USA and ²Department of Physics & Astronomy, Vanderbilt University, Nashville, TN 37235

Macroscopic dark matter refers to a variety of dark matter candidates that would be expected to (elastically) scatter off of ordinary matter with a large geometric cross-section. A wide range of macro masses M_X and cross-sections σ_X remain unprobed. We show that over a wide region within the unexplored parameter space, collisions of a macro with a human body would result in serious injury or death. We use the absence of such unexplained impacts with a well-monitored subset of the human population to exclude a region bounded by $\sigma_X > 10^{-8} - 10^{-7}$ cm² and $M_X < 50$ kg. Our results open a new window on dark matter: the human body as a dark matter detector.

What is dark matter

We do not exactly know!

W boson

neutrino

neutrino



10-39

Black Hole Remnant

wimpzilla

 $10^{-33}10^{-30}10^{-27}10^{-24}10^{-21}10^{-18}10^{-15}10^{-12}10^{-9}10^{-6}10^{-3}10^{0}10^{3}10^{6}10^{9}10^{12}10^{15}10^{18}10^{18}10^{15}10^{18}10^{15}10^{18}10^{15}10^{18}10^{15}10^{18}10^{15}10^{18}10^{15}10^{18}10^{15}10^{18}10^$ mass (GeV)

Ways of probing WIMP



Detecting technologies



DM direct detections



Where to go for Direct detections



Why neutrinos relevant?



Sterile neutrino can be DM candidate There might be neutrino portals to the DM Their experiments can detect both neutrino&DM

Two relevant issues

Precision calculations of the direct detection cross section.

Understanding the neutrino floor.
Neutrino flux
Neutrino interactions

Neutrino flux from the universe



Sun as the source of neutrinos



Supernova and atmosphere neutrino



Neutrino flux on the earth



Neutrino-nuclei scattering

Charged currents coupling to electroweak gauge boson

Neutral currents coupling to electroweak gauge boson

$$\sum_{\substack{\alpha=e,\mu,\tau\\\alpha=e,\mu,\tau}} W_{\mu}^{+}(\bar{\nu}_{\alpha}\gamma^{\mu}\alpha) + h.c.$$

Coherent neutrino-nucleus scattering in the SM

Number of expected events

$$\frac{d\sigma_{\nu}}{dE_R} = \frac{G_F^2}{4\pi} Q_{\nu N}^2 m_N \left(1 - \frac{m_N E_R}{2E_{\nu}^2}\right) F^2(E_R)$$

Weak hyper-charge of target nucleus

Nuclear form factor

$$N = \frac{\varepsilon}{m_N} \int_{E_T}^{E_{max}} dE_R \int dE_\nu \frac{d\phi_\nu}{dE_\nu} \frac{d\sigma_\nu}{dE_R}$$

Neutrino floor in the SM

Neutrino floor in the SM

Neutrino floor in the SM

Billard, et al., PRD89,023524

Non-standard Neutrino interactions

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Neutrino oscillations

Neutral current NSI: Propagation of neutrinos in matter

Charged current NSI: Production and detection

$$i\frac{d}{dx}\begin{pmatrix}\nu_e\\\nu_\mu\\\nu_\tau\end{pmatrix} = H\begin{pmatrix}\nu_e\\\nu_\mu\\\nu_\tau\end{pmatrix}$$

$$H = H_{\rm vac} + H_{\rm matt}$$

$$H_{\text{vac}} = U \text{Diag} \left(\frac{m_1^2}{2E}, \frac{m_2^2}{2E}, \frac{m_3^2}{2E} \right) U^{\dagger}$$
$$H_{\text{matt}} = \sqrt{2} G_F N_e \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$H_{\text{matt}} = \sqrt{2}G_F N_e \begin{pmatrix} 1 & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} + \sqrt{2}G_F \sum_{f=e,u,d} \begin{pmatrix} \varepsilon_{ee}^f & \varepsilon_{e\mu}^f & \varepsilon_{e\tau}^f \\ \varepsilon_{\mu e}^f & \varepsilon_{\mu\mu}^f & \varepsilon_{\mu\tau}^f \\ \varepsilon_{\tau e}^f & \varepsilon_{\tau\mu}^f & \varepsilon_{\tau\tau}^f \end{pmatrix}$$

 $-0.008 < \varepsilon_{ee}^{uV} < 0.618 \qquad -0.111 < \varepsilon_{\mu\mu}^{uV} < 0.402$ $-0.012 < \varepsilon_{ee}^{dV} < 0.361 \qquad -0.103 < \varepsilon_{\mu\mu}^{dV} < 0.361$ Altmannshofer, et al., 1812.02778

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CONHERENT

GOAL:Measure N² dependence of CEvNS process

Beam ON coincidence window	547 counts
Anticoincidence window	405 counts
Beam-on bg: prompt beam neutrons	7.0 ± 1.7
Beam-on bg: NINs (neglected)	4.0 ± 1.3
Signal counts, single-bin counting	136 ± 31
Signal counts, 2D likelihood fit	134 ± 22
Predicted SM signal counts	173 ± 48

Confirm CEvNS at 6.7 sigma

CHARM

$$R_e = \frac{\sigma(\nu_e N \to \nu_e X) + \sigma(\bar{\nu}_e N \to \bar{\nu}_e X)}{\sigma(\nu_e N \to e^- X) + \sigma(\bar{\nu}_e N \to e^+ X)} = 0.406 \pm 0.140$$

CHARM, PLB180,303

$$R_{\mu} = \frac{\sigma(\nu_{\mu}N \to \nu_{\mu}X)}{\sigma(\nu_{\mu}N \to \mu^{-}X)} = 0.3093 \pm 0.0031$$

CHARM, Z. Phys. C36,611

 $R_e^{\rm SM} = 0.3221 \pm 0.0006$

 $R_{\nu_{\mu}}^{\rm SM} = 0.3156 \pm 0.0006$

Falkowski, et al., 1706.03783

$$R_e^{\rm NSI} = R_e^{\rm SM} + \frac{\Delta \sigma_{\rm NSI}}{\sigma_{\rm CC}}$$

$$R_{\nu_{\mu}}^{\rm NSI} = R_{\nu_{\mu}}^{\rm SM} + \frac{\Delta \sigma_{\rm NSI}}{\sigma_{\rm CC}^{\nu_{\mu}}}$$

Combined constraints

Couplings	Constraints	Couplings	Constraints	Couplings	Constraints	Couplings	Constraints
$\zeta_{u,S}^{eX}$	0.051	$\zeta^{\mu X}_{u,S}$	0.035	$\zeta_{u,P}^{eX}$	4.863	$\zeta^{\mu X}_{u,P}$	0.484
$\zeta_{d,S}^{eX}$	0.051	$\zeta^{\mu X}_{d,S}$	0.034	$\zeta_{d,P}^{eX}$	6.256	$\zeta^{\mu X}_{d,P}$	0.686
$\zeta^{eX}_{s,S}$	0.866	$\zeta^{\mu X}_{s,S}$	0.579	$\zeta_{s,P}^{eX}$	11.87	$\zeta^{\mu X}_{s,P}$	1.603
$\zeta_{u,T}^{eX}$	0.632	$\zeta^{\mu X}_{u,T}$	0.064	$\zeta_{u,A}^{eX}$	0.996	$\zeta^{\mu X}_{u,A}$	0.178
$\zeta_{d,T}^{eX}$	0.866	$\zeta^{\mu X}_{d,T}$	0.093	$\zeta_{d,A}^{eX}$	0.996	$\zeta^{\mu X}_{d,A}$	0.250
$\zeta_{s,T}^{eX}$	1.680	$\zeta^{\mu X}_{s,T}$	0.215	$\zeta_{s,A}^{eX}$	2.123	$\zeta^{\mu X}_{s,A}$	0.500
$\zeta^{eX}_{u,V}$	0.123	$\zeta^{\mu X}_{u,V}$	0.084				
$\zeta_{d,V}^{eX}$	0.112	$\zeta^{\mu X}_{d,V}$	0.072				
$\zeta^{eX}_{s,V}$	2.123	$\zeta^{\mu X}_{s,V}$	0.566				

Neutrino floor with exotic neutrino interactions

Quark level	Nucleon level	Matching conditions
$\frac{G_F}{\sqrt{2}}\zeta_{q,S}\bar{\nu}_{\alpha}P_L\nu_{\beta}\bar{q}q$	$\frac{G_F}{\sqrt{2}}\zeta_{N,S}\bar{\nu}_{\alpha}P_L\nu_{\beta}\bar{N}N$	$\zeta_{N,S} = \sum_{q=u,d} \zeta_{q,S} rac{m_N}{m_q} f_{T_q}^N$
$rac{G_F}{\sqrt{2}}\zeta_{q,P}ar{ u}_{lpha}P_L u_{eta}ar{q}i\gamma^5q$	$\frac{G_F}{\sqrt{2}}\zeta_{N,P}\bar{\nu}_{\alpha}P_L\nu_{\beta}\bar{N}i\gamma^5N$	$\zeta_{N,P} = \sum_{q=u,d} \zeta_{q,P} \frac{m_N}{m_q} \left(1 - \frac{\bar{m}}{m_q}\right) \Delta_q^N$
$\frac{G_F}{\sqrt{2}}\zeta_{q,V}\bar{\nu}_{\alpha}\gamma_{\mu}P_L\nu_{\beta}\bar{q}\gamma^{\mu}q$	$\frac{G_F}{\sqrt{2}}\zeta_{N,V}\bar{\nu}_{\alpha}\gamma_{\mu}P_L\nu_{\beta}\bar{N}\gamma^{\mu}N$	$\zeta_{p,V} = 2\zeta_{u,V} + \zeta_{d,V}; \zeta_{n,V} = \zeta_{u,V} + 2\zeta_{d,V}$
$\frac{G_F}{\sqrt{2}}\zeta_{q,A}\bar{\nu}_{\alpha}\gamma_{\mu}P_L\nu_{\beta}\bar{q}\gamma^{\mu}\gamma^5q$	$\frac{G_F}{\sqrt{2}}\zeta_{N,A}\bar{\nu}_{\alpha}\gamma_{\mu}P_L\nu_{\beta}\bar{N}\gamma^{\mu}\gamma^5N$	$\zeta_{N,A} = \sum_q \zeta_{q,A} \Delta_q^N$
$\frac{G_F}{\sqrt{2}}\zeta_{q,T}\bar{\nu}_{\alpha}\sigma_{\mu\nu}P_L\nu_{\beta}\bar{q}\sigma^{\mu\nu}q$	$\left \frac{G_F}{\sqrt{2}} \zeta_{N,T} \bar{\nu}_{\alpha} \sigma_{\mu\nu} P_L \nu_{\beta} \bar{N} \sigma^{\mu\nu} N \right.$	$\zeta_{N,T} = \sum_{q} \zeta_{q,T} \delta_q^N $

Numerical results-1: vector interactions with Xe131

Numerical results-1: vector interactions with X131

Numerical results-2: axial-vector interactions with Xe131

Numerical results-2: axial-vector interactions with Xe131

Numerical results-3: scalar interactions with Xe131

Numerical results-3: scalar interactions with X131

Numerical results-4: pseudo-scalar interactions with Xe131

Numerical results-4: pseudo-scalar interactions with X131

Numerical results-5: tensor interactions with Xe131

Numerical results-5: tensor interactions with X131

Numerical results-6: vector interactions with Ge72

Numerical results-6: scalar interactions with Ge72

Conclusions

Impact of non-standard neutrino interactions to the neutrino floor was studied

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