



Strong phase at BESIII and possible **new** charm physics at LHCb

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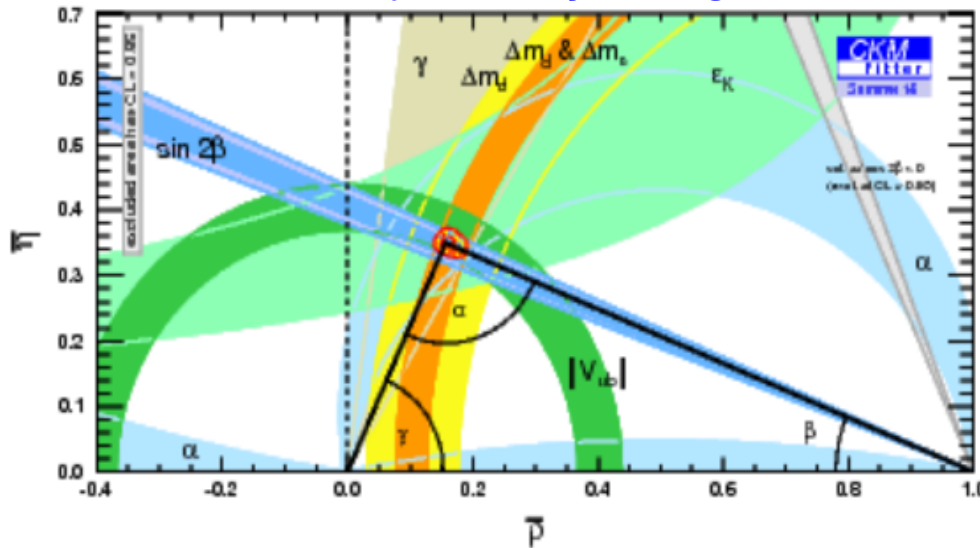
Outline

- Introduction
- Strong phase measurement at BESIII
- Future prospects on Strong phase measurement
- Some new charm physics at LHCb
- Summary and Outlook

Introduction

- γ/ϕ_3 is the only CKM angle that can be measured in tree-level processes, in which the contribution of non-SM effects is expected to be small [JHEP 01(2014)051].
- Measurement of γ provides a benchmark of the SM with negligible theoretical uncertainty.

B_d Unitarity Triangle



Moriond 2018

Phases of CKM elements:

$$\beta = \varphi_1 = \arg\left(-\frac{V_{cd}V_{cb}^*}{V_{td}V_{tb}^*}\right) \quad \beta = (22.0 \pm 0.7)^\circ$$

$$\alpha = \varphi_2 = \arg\left(-\frac{V_{td}V_{tb}^*}{V_{ud}V_{ub}^*}\right) \quad \alpha = (84.9^{+5.1}_{-4.5})^\circ$$

$$\gamma = \varphi_3 = \arg\left(-\frac{V_{ud}V_{ub}^*}{V_{cd}V_{cb}^*}\right) \quad \gamma = (73.5^{+4.2}_{-5.1})^\circ$$

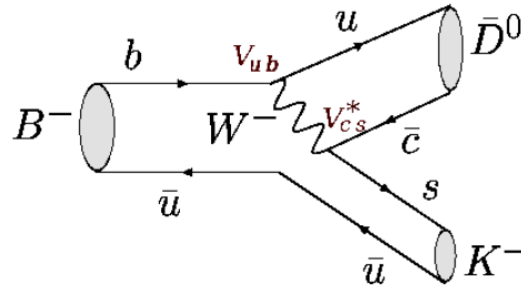
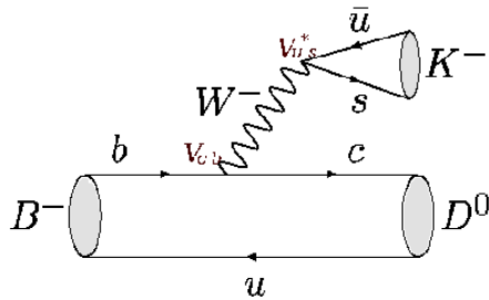
~ an order of magnitude worse than that on β .

$$\gamma \text{ from CKM fitter: } \gamma = (65.8^{+1.0}_{-1.7})^\circ$$

- The current world-average of γ deviates the indirect determination from CKM fitter by $\sim 1.5\sigma$.
- Clearly, an improved knowledge of the measurement of γ is important to further test the SM and probe for new physics.

Strong phase as inputs

- γ/ϕ_3 can be measured by studying the interference between $B^- \rightarrow D^0 K^-$ and $B^- \rightarrow \bar{D}^0 K^-$



$$A(B^- \rightarrow D^0 K^-) = A_B A_D$$

$$A(B^- \rightarrow \bar{D}^0 K^-) = A_B r_B e^{i(\delta_B - \gamma)} A_{\bar{D}}$$

where r_B is ratio of suppressed to favored amplitudes, δ_B is the strong-phase difference between the favoured and suppressed amplitudes.

- Generally, three methods were proposed to measure γ/ϕ_3 :

- ✓ GLW [1]: via $D^0 \rightarrow$ CP eigenstate, $K^+ K^-$, $\pi^+ \pi^-$, $K_S^0 \pi^0$ etc.
- ✓ ADS [2]: via $D^0 \rightarrow$ CF and DCS, such as $K^+ \pi^-$, $K^+ \pi^- \pi^0$, $K^+ \pi^- \pi^- \pi^+$ etc.
- ✓ GGSZ [3]: via with $D^0 \rightarrow$ Multi-body self-conjugate decays, $K_S^0 \pi^+ \pi^-$ etc.

Measurements from strong phases are key inputs.

For GGSZ:

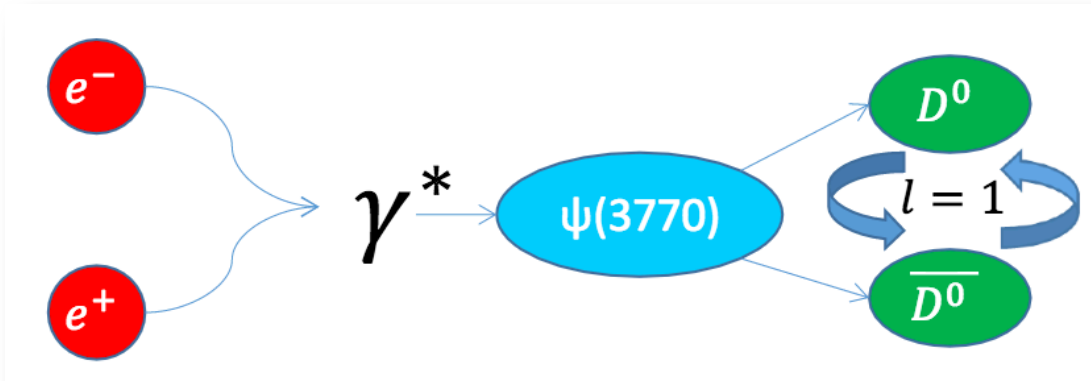
$$d\Gamma(B^\pm \rightarrow DK^\pm) = |\mathcal{A}_D|^2 + r_B^2 |\mathcal{A}_{\bar{D}}|^2 + 2r_B |\mathcal{A}_D| |\mathcal{A}_{\bar{D}}| \times [\cos \Delta \delta_D \cos(\delta_B \pm \phi_3) + \sin \Delta \delta_D \sin(\delta_B \pm \phi_3)]$$

[1] M. Gronau, D. London, *Phys. Lett. B* 253, 483 (1991); M. Gronau, D. Wyler, *Phys. Lett. B* 265, 172 (1991).

[2] D. Atwood, I. Dunietyz and A. Soni, *Phys. Rev. Lett.* 78, 3257 (1997).

[3] A. Giri, Y. Grossman, A. Soffer and J. Zupan, *Phys. Rev. D* 68, 054018 (2003).

Quantum correlated (QC) neutral D state near threshold

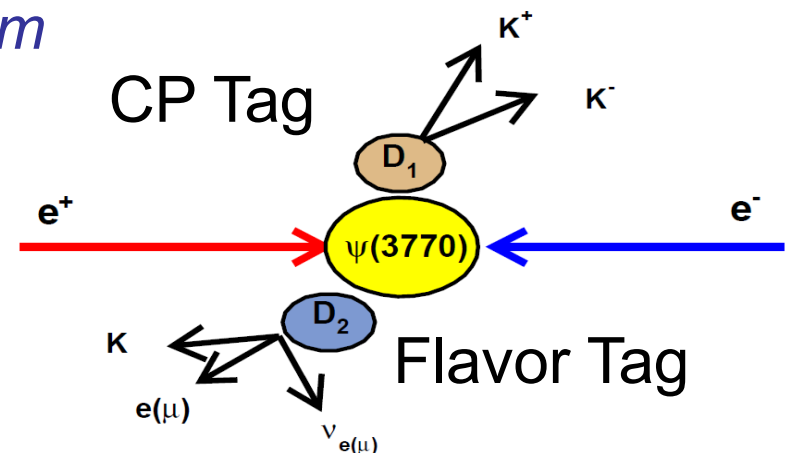


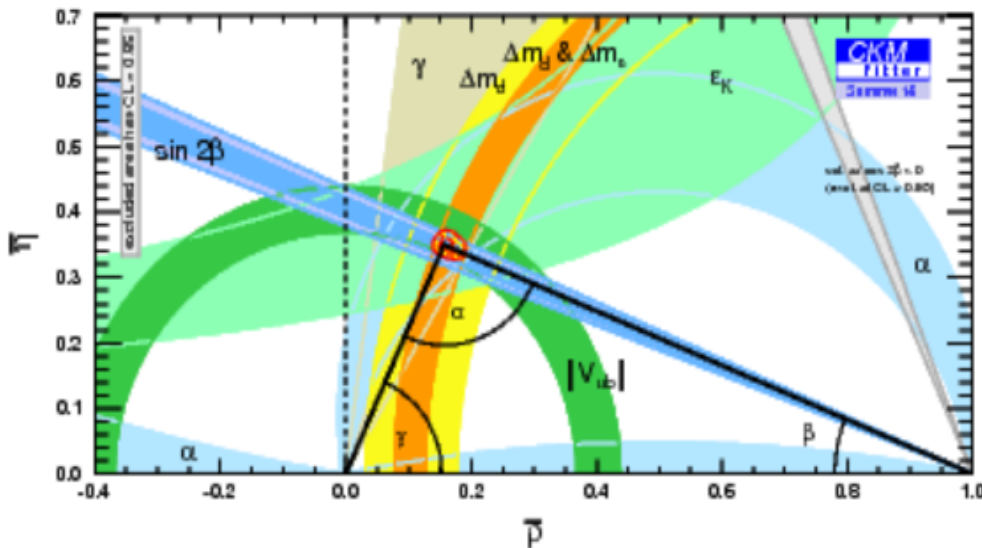
If D^0 in CP eigenstate, \bar{D}^0 must be in opposite CP eigenstate

Quantum Correlations (QC) and CP-tagging are unique

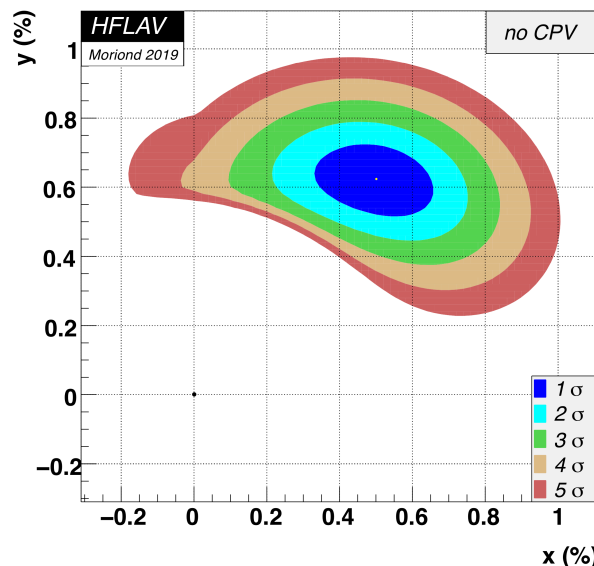
Taking advantage the quantum coherence of $D\bar{D}$ pairs, BESIII can study the charm physics in an unique way

- strong phase in D decays
- D mixing parameters
- direct CP violation
- ...





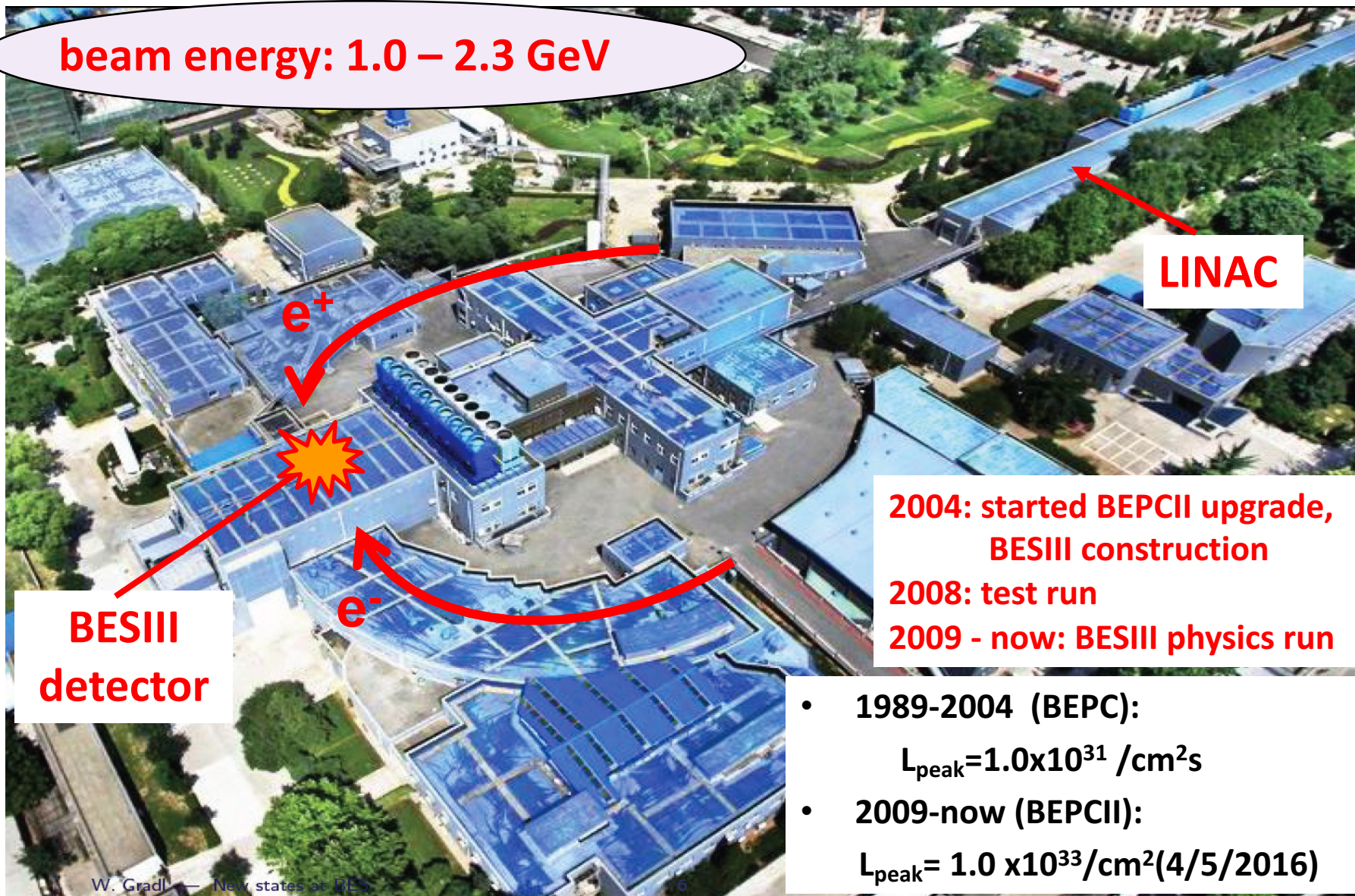
Precision CKM test



Charm Mixing & CP violation

- ◆ **inputs from Quantum Correlated (QC)**
 $\psi(3770) \rightarrow D\bar{D}$ decays
 - ◆ (Averaged) Strong phase difference: δ_D
 - ◆ Coherent factors: R_D
 - ◆ (Averaged) Strong phase in Dalitz bins: c_i, s_i
- ◆ **B factories, LHCb, Super B factories are the customers**

beam energy: 1.0 – 2.3 GeV



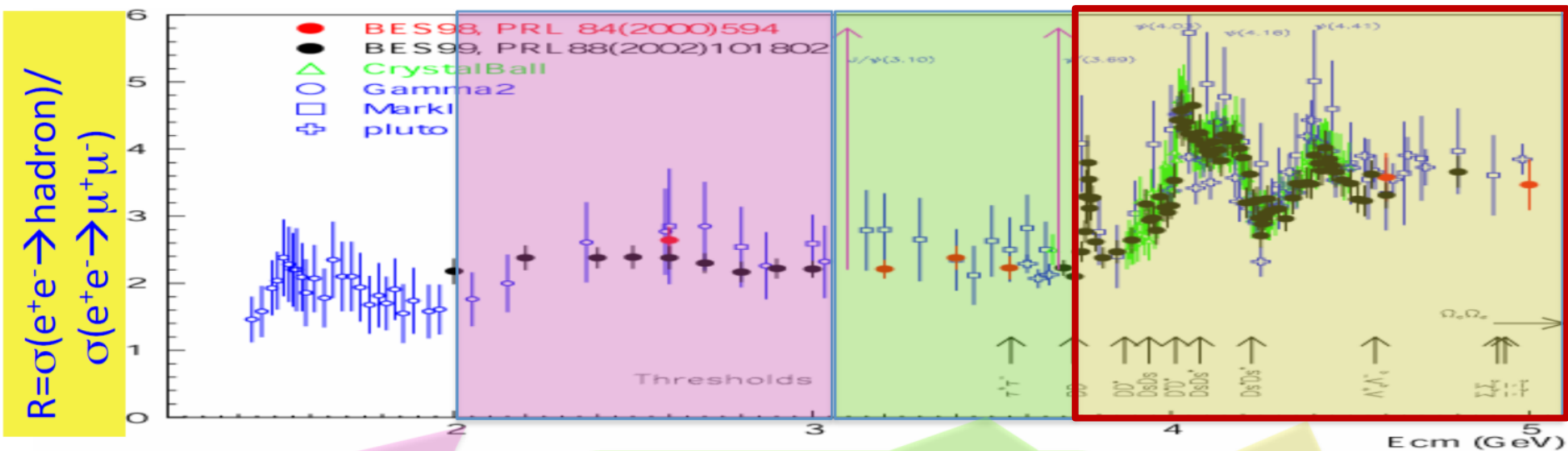
LINAC

BES III
detector

2004: started BEPCII upgrade,
BES III construction
2008: test run
2009 - now: BES III physics run

- 1989-2004 (BEPC):
 $L_{\text{peak}} = 1.0 \times 10^{31} / \text{cm}^2 \text{s}$
- 2009-now (BEPCII):
 $L_{\text{peak}} = 1.0 \times 10^{33} / \text{cm}^2 (4/5/2016)$

W. Gradl — New states of BES



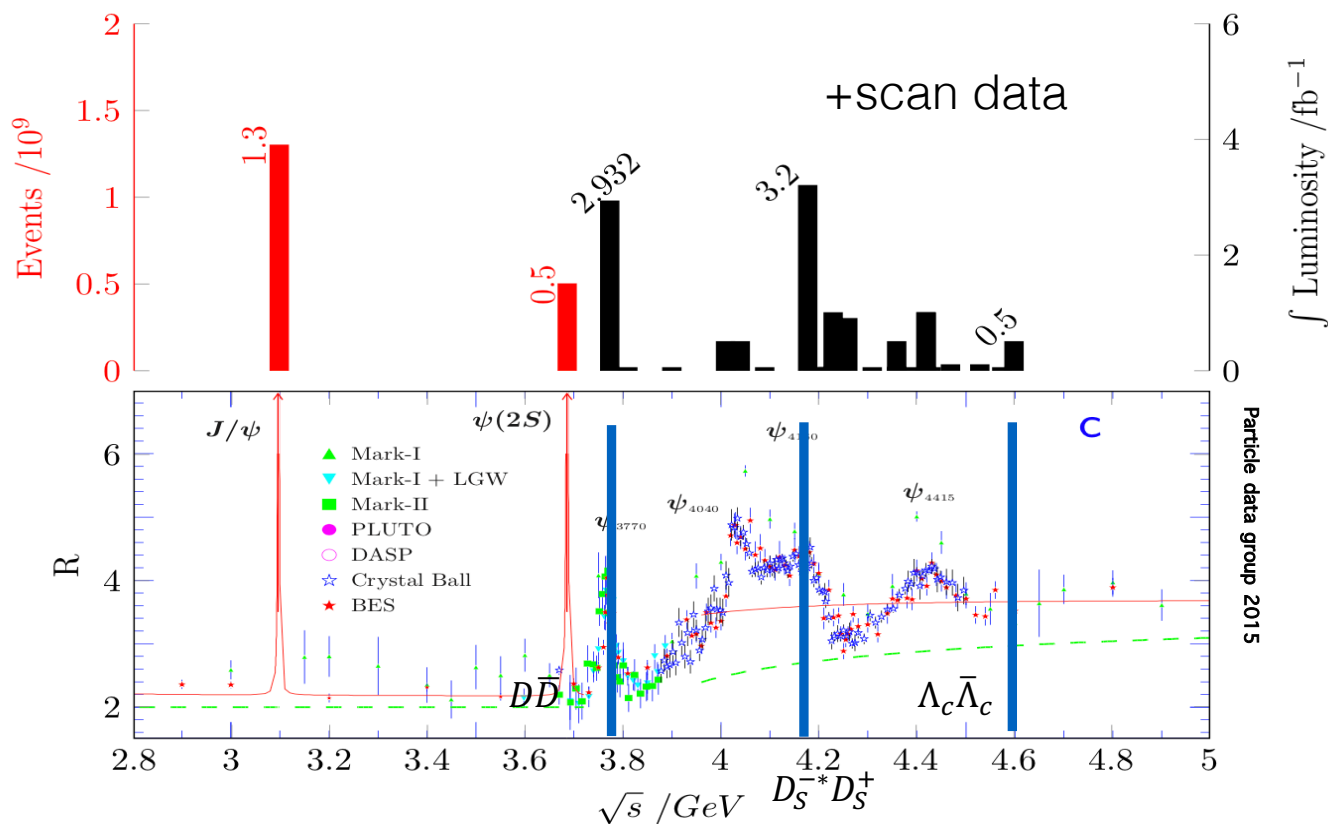
- Hadron form factors
- $Y(2175)$ resonance
- Multiquark states with s quark, Z_s
- MLLA/LPHD and QCD sum rule predictions

- Light hadron spectroscopy
- Gluonic and exotic states
- Process of LFV and CPV
- Rare and forbidden decays
- Physics with τ lepton

- XYZ particles
- D mesons
- f_D and f_{D_s}
- D_0 - \bar{D}_0 mixing
- Charm baryons

$D_{(s)}$ & Λ_c decays:

- (semi-)leptonic decays
- hadronic decays



$\psi(3770)$ is a spin -1 state and therefore the amplitude of $\psi(3770) \rightarrow D^0 \bar{D}^0$:

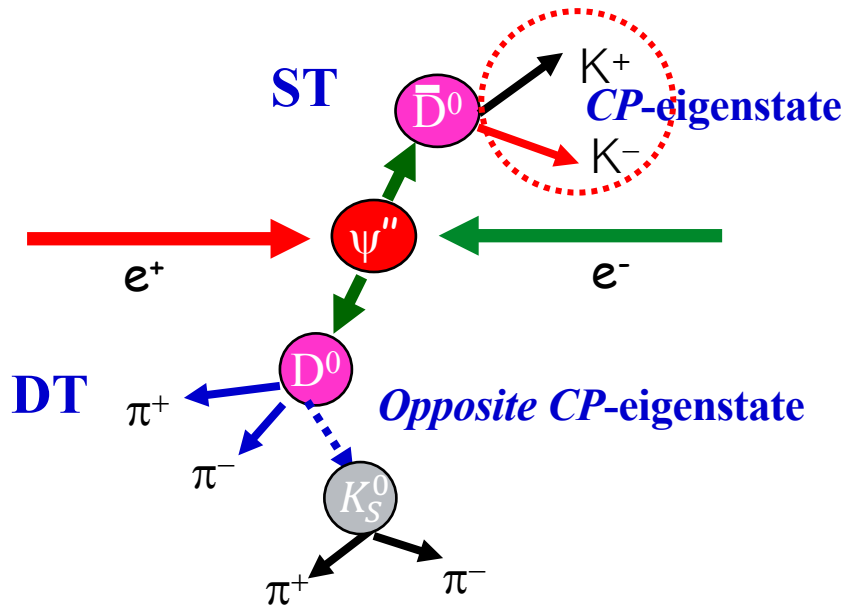
$$(|D^0\rangle|\bar{D}^0\rangle - |\bar{D}^0\rangle|D^0\rangle) / \sqrt{2} \quad [\text{anti-symmetric wave function}]$$

The amplitude for two D mesons to decay to states F and G is

[D. Atwood and A. Soni, PRD68, 033003 (2003)]:

$$\Gamma(F|G) = \Gamma_0 [A_F^2 \bar{A}_G^2 + \bar{A}_F^2 A_G^2 - 2R_F R_G A_F \bar{A}_F A_G \bar{A}_G \cos[\delta_D^F - \delta_D^G]]$$

The coherence factors R_F , the strong-phase difference (or the related parameters) δ_D^F , can be extracted based on the study of the quantum correlated $D\bar{D}$ meson pairs.



The DT mode K^+K^- vs. $K_S^0\pi^+\pi^-$ is selected as an example.

- ✓ **Single tag (ST) samples:**
decay products of only one D meson are reconstructed
- ✓ **Double tag (DT) samples:**
decay products of both D mesons are reconstructed
- ✓ **Some typical reconstructed D decay modes**

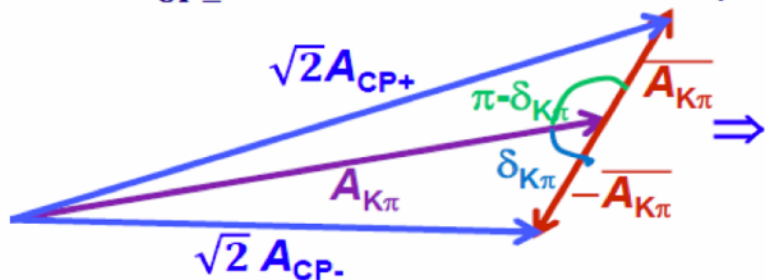
Tag group	
Flavor	$K^+\pi^-, K^+\pi^-\pi^0, K^+\pi^-\pi^-\pi^+, K^+e^-\bar{\nu}_e$
CP-even	$K^+K^-, \pi^+\pi^-, K_S^0\pi^0\pi^0, K_L^0\pi^0, \pi^+\pi^-\pi^0$
CP-odd	$K_S^0\pi^0, K_S^0\eta, K_S^0\omega, K_S^0\eta', K_L^0\pi^0\pi^0$
Mixed-CP	$K_S^0\pi^+\pi^-$

- ✓ **Using this method, CLEO had performed lots of important and excellent measurements related to strong phases.**

2.93 fb⁻¹ @3.773 GeV

Quantum correlation → Interference → access strong phase!

$$\langle K\pi | D_{CP\pm} \rangle = (\langle K\pi | D^0 \rangle \pm \langle K\pi | \bar{D}^0 \rangle) / \sqrt{2} \Rightarrow \sqrt{2} A_{CP\pm} = A_{K\pi} \pm \bar{A}_{K\pi}$$

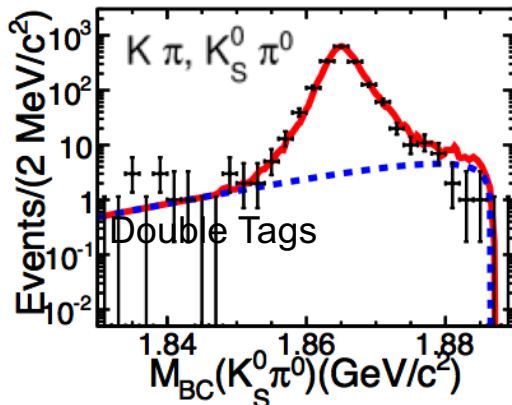
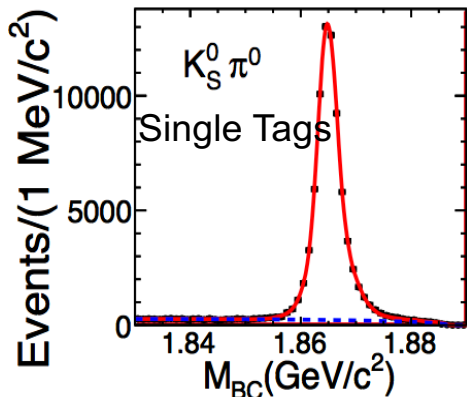


$$2r_{K\pi} \cdot \cos \delta_{K\pi} \approx A_{CP \rightarrow K\pi} \equiv \frac{|A_{CP-}|^2 - |A_{CP+}|^2}{|A_{CP-}|^2 + |A_{CP+}|^2} = \frac{Br(D_{CP-} \rightarrow K\pi) - Br(D_{CP+} \rightarrow K\pi)}{Br(D_{CP-} \rightarrow K\pi) + Br(D_{CP+} \rightarrow K\pi)}$$

Strong phase:

$$\frac{\langle K^-\pi^+ | \bar{D}^0 \rangle^{DCS}}{\langle K^-\pi^+ | D^0 \rangle^{CF}} \equiv -r_{K\pi} e^{-i\delta_{K\pi}}$$

- ◆ Flavor tags: $K^-\pi^+, K^+\pi^-$
- ◆ CP+ tags (5 modes): $K^-K^+, \pi^+\pi^-, K_S^0\pi^0\pi^0, \pi^0\pi^0, \rho^0\pi^0$
- ◆ CP- tags (3 modes): $K_S^0\pi^0, K_S^0\eta, K_S^0\omega$



PLB 734, 227 (2014)

BESIII results:

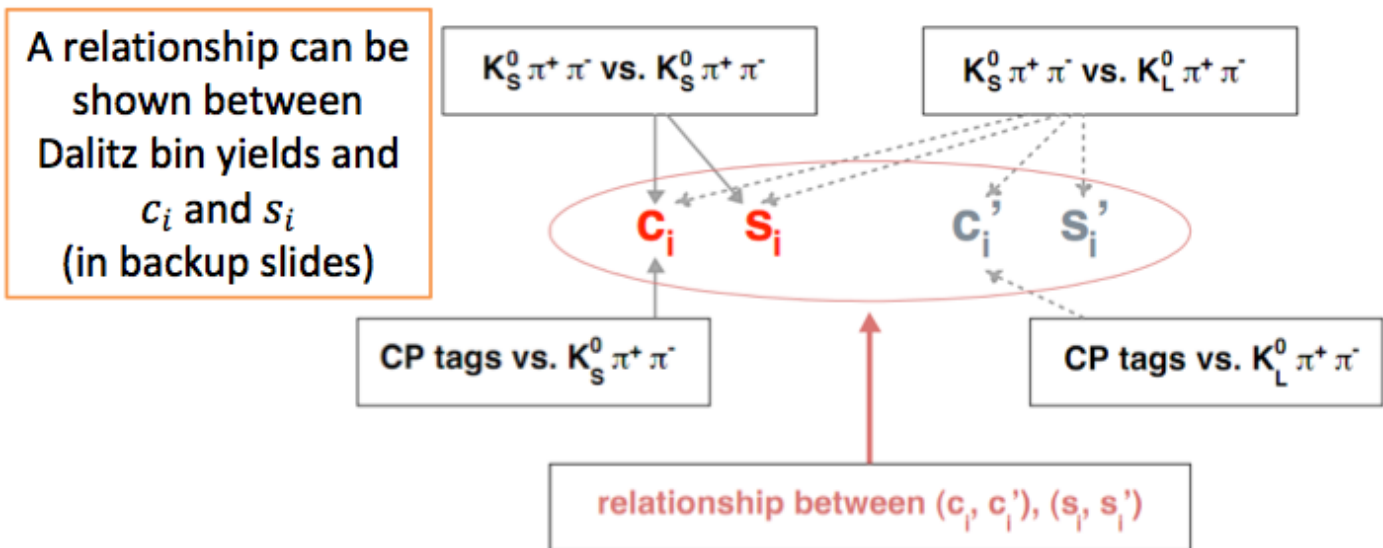
$$\cos \delta_{K\pi} = 1.02 \pm 0.11 \pm 0.06 \pm 0.01$$

- The third error is due to the input parameters
- World best precision
- In 10 /fb BESIII data, precision of $\cos \delta_{K\pi}$ will reach ~ 0.07

$$N_i^\pm = h_B \left[K_{\pm i} + r_b^2 K_{\mp i} + 2\sqrt{K_i K_{-i}} (x_\pm c_i \pm y_\pm s_i) \right]$$

$B^\pm \rightarrow DK^\pm$ yields (points to N_i^\pm)
 from flav.-tagged $D \rightarrow K_S \pi \pi$ (points to $K_{\pm i}$)
 extracted from fit to the B^\pm yields (points to $x_\pm c_i$)
 measured by CLEO [PRD82, 112006 (2010)] (points to $y_\pm s_i$)

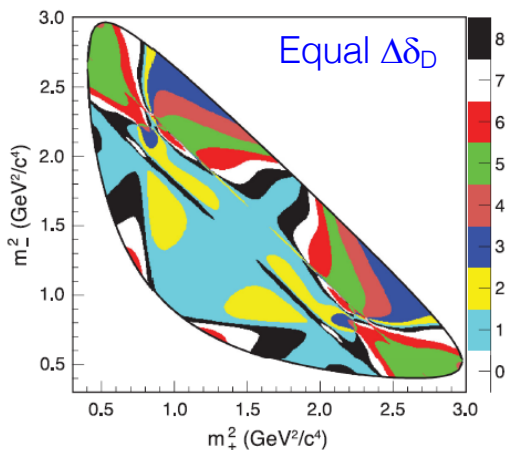
We can calculate c_i and s_i from double tags of $D^0 \rightarrow K_S \pi^+ \pi^-$ vs $D^0 \rightarrow (K_{S,L} \pi^+ \pi^-$ or CP eigenstates)



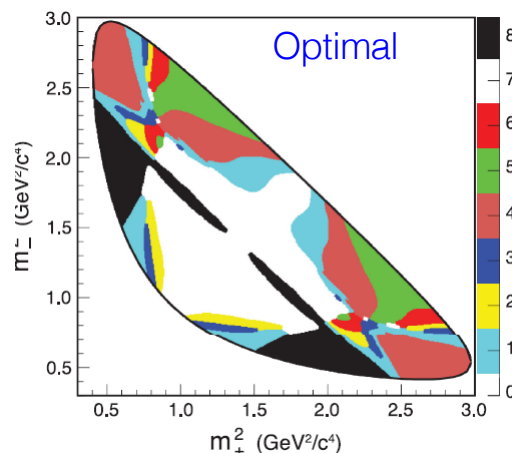
Only c_i, s_i from $K_S \pi^+ \pi^-$ is used to calculate γ .
 However adding in $D^0 \rightarrow K_L \pi^+ \pi^-$ we can calculate c'_i, s'_i and use how they relate to c_i, s_i to further constrain our results in a Global fit.

Three typical binning schemes

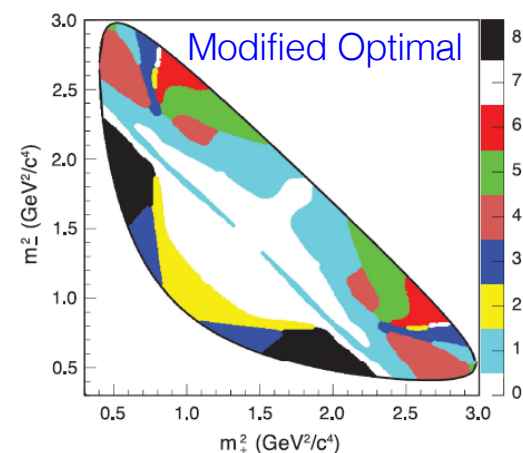
[J. Libby et al. (CLEO Collaboration), Phys. Rev. D 82,112006 (2010)]



DD-mixing¹, β measurements²
[minimum variation in $\Delta\delta_D$]



γ measurements^{3,4}
[Optimized sensitivity]



γ in Low yields
[Optimization including backgrounds]

- ✓ “BaBar K-matrix” $D^0 \rightarrow K_S^0 \pi^+ \pi^-$ model as in Ref. [Phys. Rev. D 78, 034023 (2008)].
- ✓ It should be noted that although the choice of binning is model-dependent, however, a poor choice of model results only in a loss of precision, instead of bias in measuring γ/ϕ_3 .

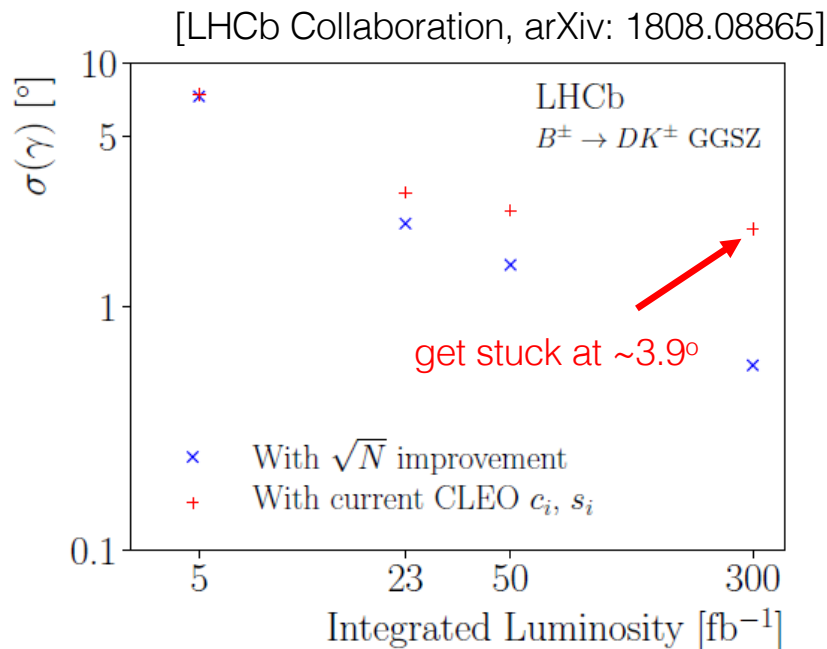
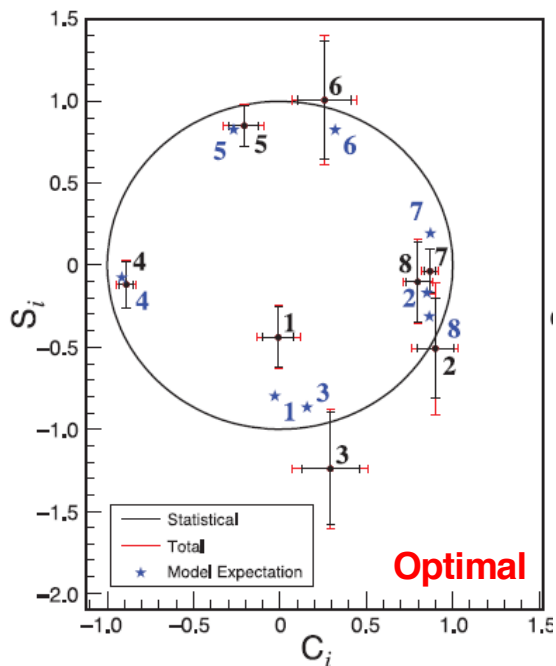
[1] R. Aaij et al. (LHCb Collaboration), Phys. Rev. Lett. 122, 231802 (2019); JHEP 04(2016) 033.

[2] V. Vorobyev et al. (Belle Collaboration), Phys. Rev. D 94, 052004 (2016).

[3] R. Aaij et al. (LHCb Collaboration), Phys. Lett. B 718, 43 (2012); JHEP 10 (2014) 097; JHEP 06 (2016) 131; JHEP 08 (2018) 176.

[4] H. Aihara et al. (Belle Collaboration), Phys. Rev. D 85, 112014 (2012).

- ✓ Results of c_i and s_i in optimal binning from CLEO experiments.



- ✓ The systematic uncertainty in measurement of γ due to input strong-phase parameters is 3.9° for optimal binning. The overall sensitivity is limited to $\sim 3.9^\circ$ for model-independent GGSZ approach.
- ✓ *Therefore, improved measurements in c_i & s_i from BESIII are essential for degree-level precision of measuring γ via model-independent GGSZ approach.*

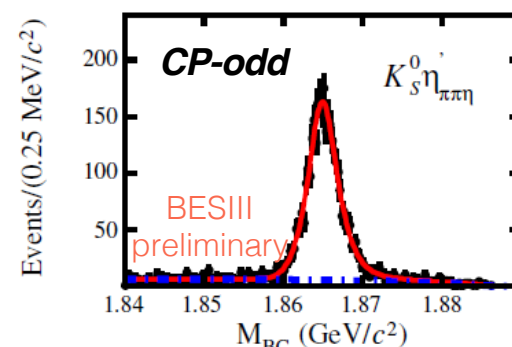
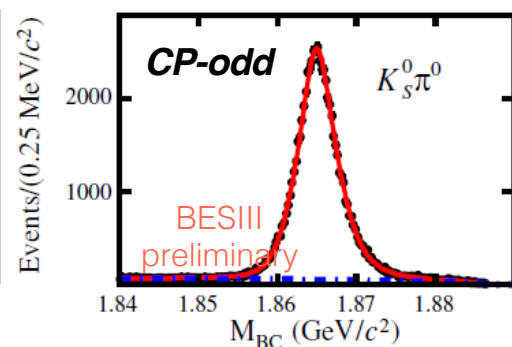
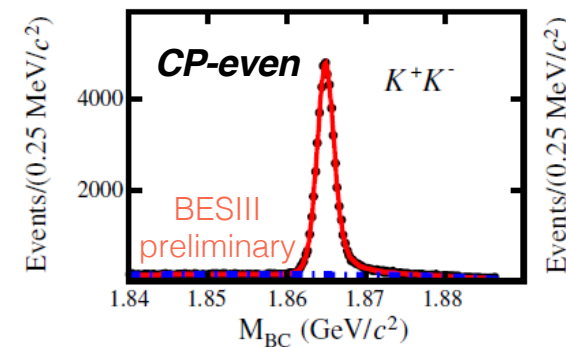
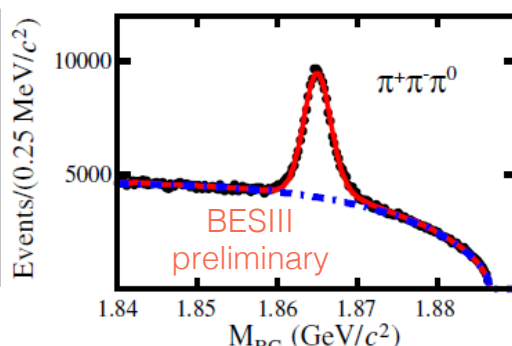
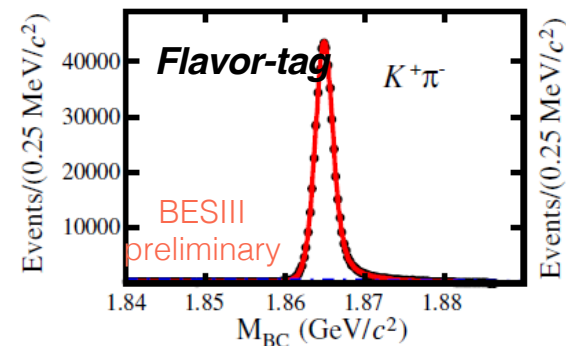
The M_{BC} distributions for ST D mesons

✓ A list of tag decay modes used in this analysis

Tag group	
Flavor	$K^+\pi^-, K^+\pi^-\pi^0, K^+\pi^-\pi^-\pi^+, K^+e^-\bar{\nu}_e$
CP-even	$K^+K^-, \pi^+\pi^-, K_S^0\pi^0\pi^0, K_L^0\pi^0, \pi^+\pi^-\pi^0$
CP-odd	$K_S^0\pi^0, K_S^0\eta, K_S^0\omega, K_S^0\eta', K_L^0\pi^0\pi^0$
Mixed-CP	$K_S^0\pi^+\pi^-$

✓ Beam-constrained mass (M_{BC})

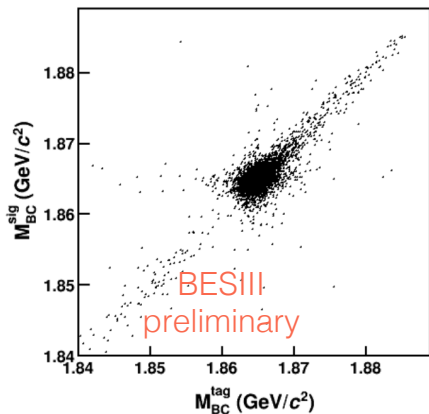
$$M_{BC} = \sqrt{(\sqrt{s}/2)^2 - |\vec{p}_{D_{tag}}|^2}$$



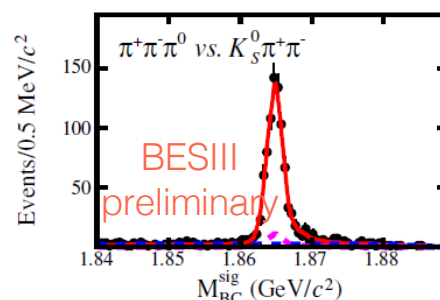
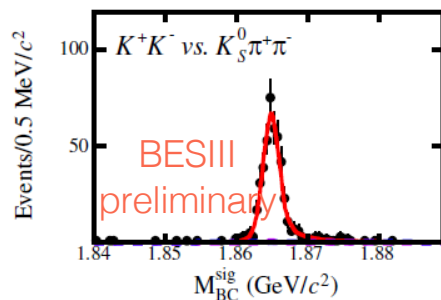
- ✓ $D \rightarrow \pi^+\pi^-\pi^0$ is not fully CP-even and the corrections for the decay is always applied.
- ✓ The fractional CP-even content of $D \rightarrow \pi^+\pi^-\pi^0$: $F_{+\pi^+\pi^-\pi^0} = 0.973 \pm 0.017$ [PLB747, 9 (2015)].

- ✓ The fully reconstructed DT $K_S \pi^+ \pi^-$ events are obtained by searching for the $K_S \pi^+ \pi^-$ signals in the recoiling system of fully-reconstructed ST events.

Scatter plot of M_{BC}^{sig} vs. M_{BC}^{tag}

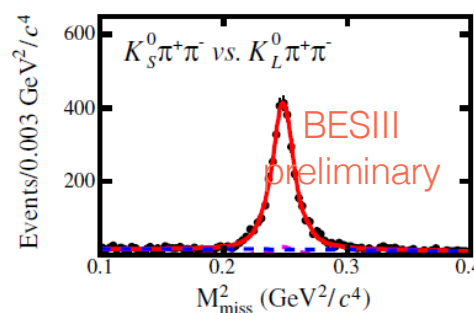
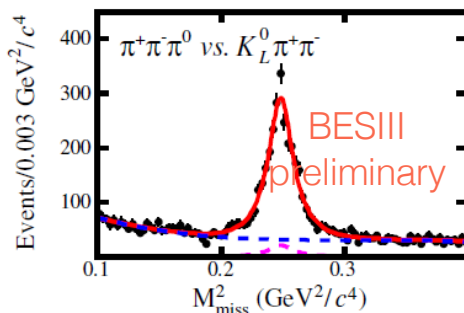
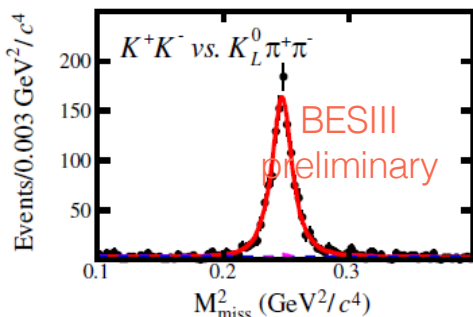


projections of the two-dimensional fits on M_{BC}^{sig}



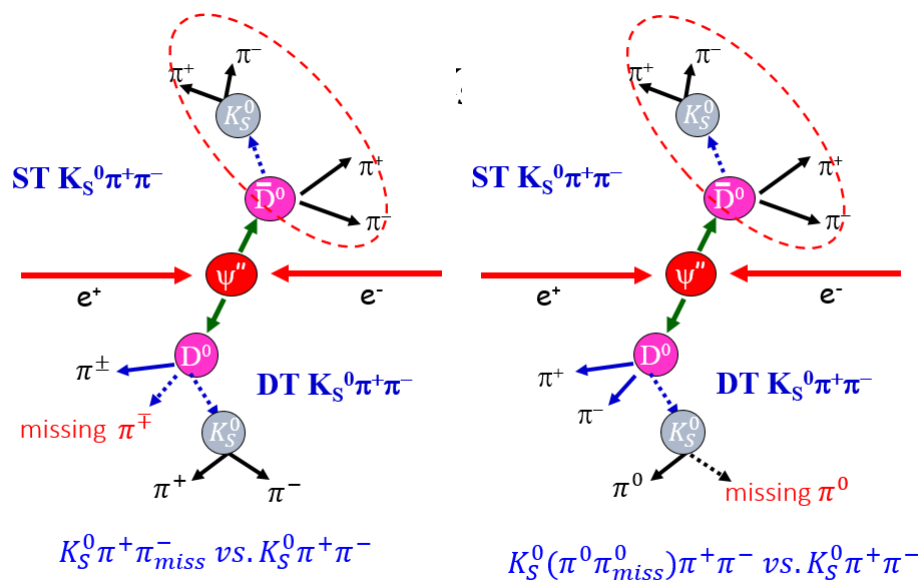
The pink curves show peaking backgrounds dominated by $D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$.

- ✓ DT events containing K_L^0 particles are identified via kinematic variable missing-mass-square (M_{miss}^2).

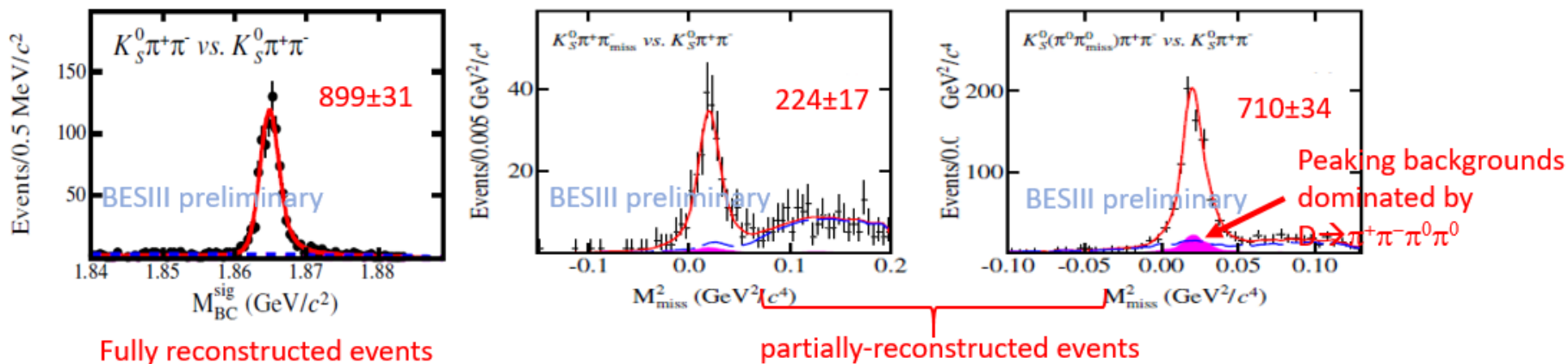


The pink curves show peaking backgrounds dominated by $D \rightarrow K_S \pi^+ \pi^-$ with $K_S^0 \rightarrow \pi^0 \pi^0$.

✓ To increase the sensitivity of measuring s_i , the partially-reconstructed events are introduced into analysis to increase the yield of events .



The yield of DT $K_S^0\pi^+\pi^-$ vs. $K_S^0\pi^+\pi^-$ is doubled by adding the partially-reconstructed samples.



Comparisons of DT events between BESIII and CLEO

Mode	BESIII [signal yields]		CLEO [raw yields]	
	$N_{K_S^0 \pi^+ \pi^-}^{DT}$	$N_{K_L^0 \pi^+ \pi^-}^{DT}$	$N_{K_S^0 \pi^+ \pi^-}^{DT}$	$N_{K_L^0 \pi^+ \pi^-}^{DT}$
Flavor tags				
$K^+ \pi^-$	4740 ± 71	9511 ± 115	1444	2857
$K^+ \pi^- \pi^0$	8899 ± 95	19225 ± 176	2759	5133
$K^+ \pi^- \pi^- \pi^+$	5695 ± 78	11906 ± 132	2240	4100
$K^+ e^- \bar{\nu}_e$	4123 ± 75		1191	
CP-even tags				
$K^+ K^-$	443 ± 22	1289 ± 41	124	357
$\pi^+ \pi^-$	184 ± 14	531 ± 28	61	184
$K_S^0 \pi^0 \pi^0$	198 ± 16	612 ± 35	56	
$\pi^+ \pi^- \pi^0$	790 ± 31	2571 ± 74	--	--
$K_L^0 \pi^0$	913 ± 41		237	
CP-odd tags				
$K_S^0 \pi^0$	643 ± 26	861 ± 46	189	288
$K_S^0 \eta \gamma \gamma$	89 ± 10	105 ± 15	39	43
$K_S^0 \eta_{\pi^+ \pi^-} \pi^0$	23 ± 5	40 ± 9	--	--
$K_S^0 \omega$	245 ± 17	321 ± 25	83	--
$K_S^0 \eta'_{\pi^+ \pi^-} \eta$	24 ± 6	38 ± 8	--	--
$K_S^0 \eta'_{\gamma \pi^+ \pi^-}$	81 ± 10	120 ± 14	--	--
$K_L^0 \pi^0 \pi^0$	620 ± 32		--	
Mixed CP tags				
$K_S^0 \pi^+ \pi^-$	899 ± 31	3438 ± 72	473	1201
$K_S^0 \pi^+ \pi^-_{miss}$	224 ± 17			
$K_S^0 (\pi^0 \pi^0_{miss}) \pi^+ \pi^-$	710 ± 34		--	--

- More tag decay modes are used in BESIII analysis.
- DT events detected at BESIII and comparisons with CLEO [J. Libby et al. (CLEO Collaboration), Phys. Rev. D 82,112006 (2010)]:

CP – eigenstate vs. $K_S^0 \pi^+ \pi^-$: 5.3×CLEO

CP – eigenstate vs. $K_L^0 \pi^+ \pi^-$: 9.2×CLEO

$K_S^0 \pi^+ \pi^-$ vs. $K_S^0 \pi^+ \pi^-$: 3.9×CLEO

$K_L^0 \pi^+ \pi^-$ vs. $K_S^0 \pi^+ \pi^-$: 2.9×CLEO

✓ “--” stands for unused mode in CLEO analysis.

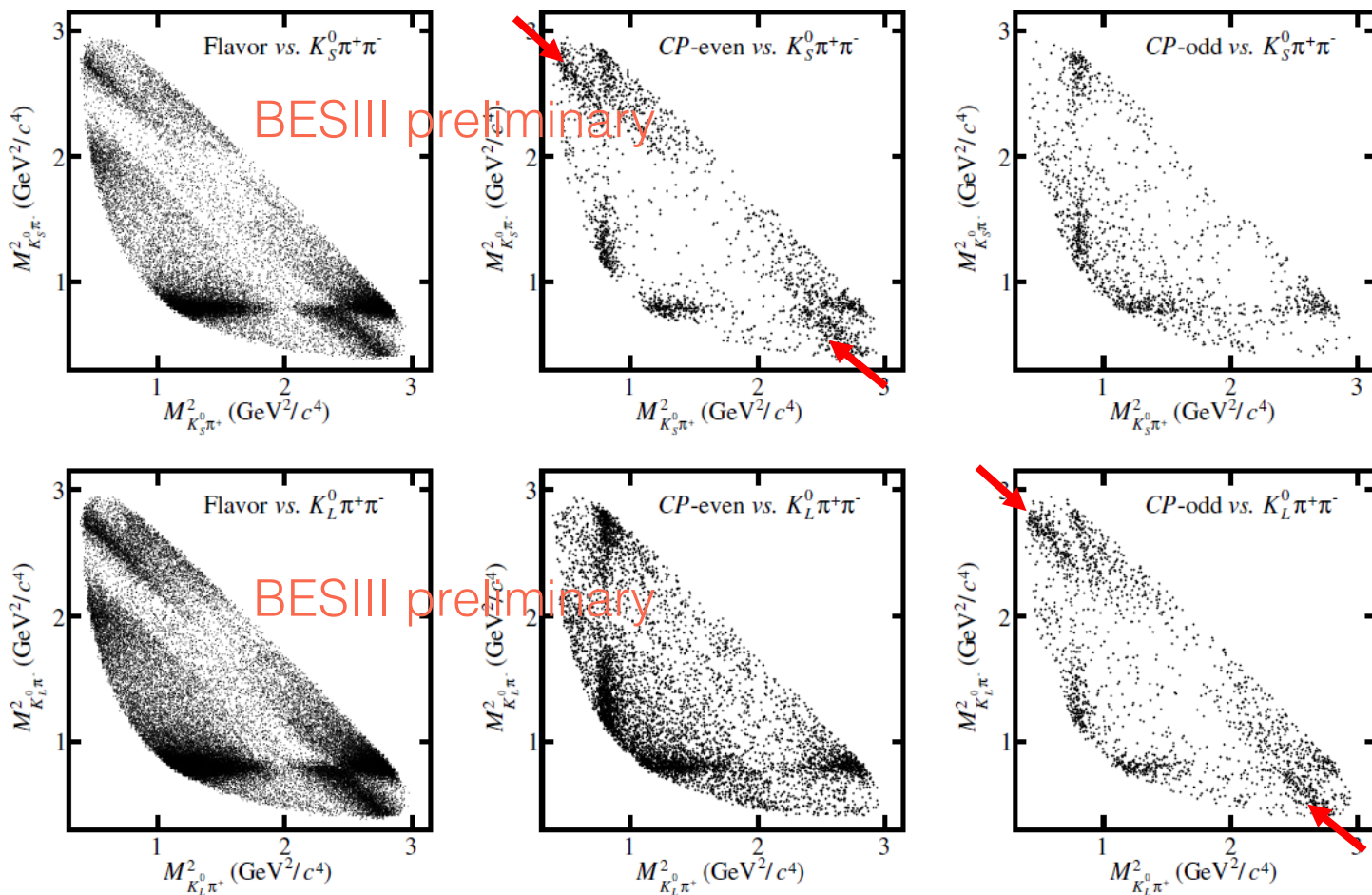
BESIII preliminary

Dalitz plots observed in data

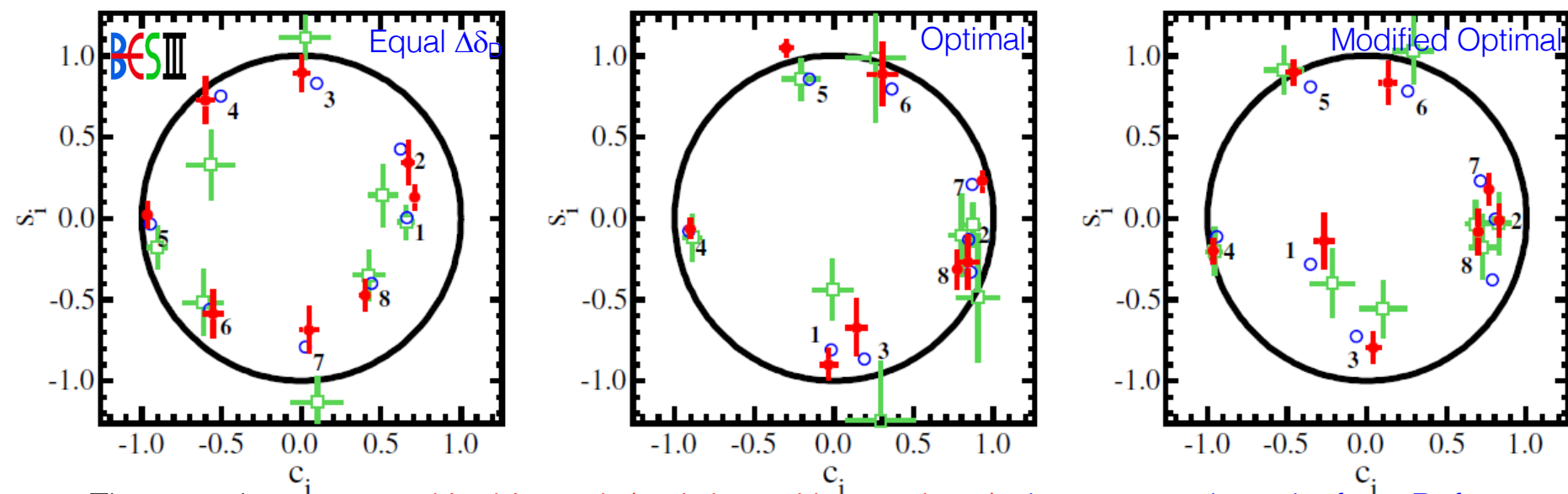


The effect of quantum correlation is immediately seen in Dalitz plots.

- ✓ The CP-odd component $K_S^0 \rho(770)^0$ is visible in CP-even tagged $K_S^0 \pi^+ \pi^-$ decays, but is absent in CP-odd tagged $K_S^0 \pi^+ \pi^-$ decays.

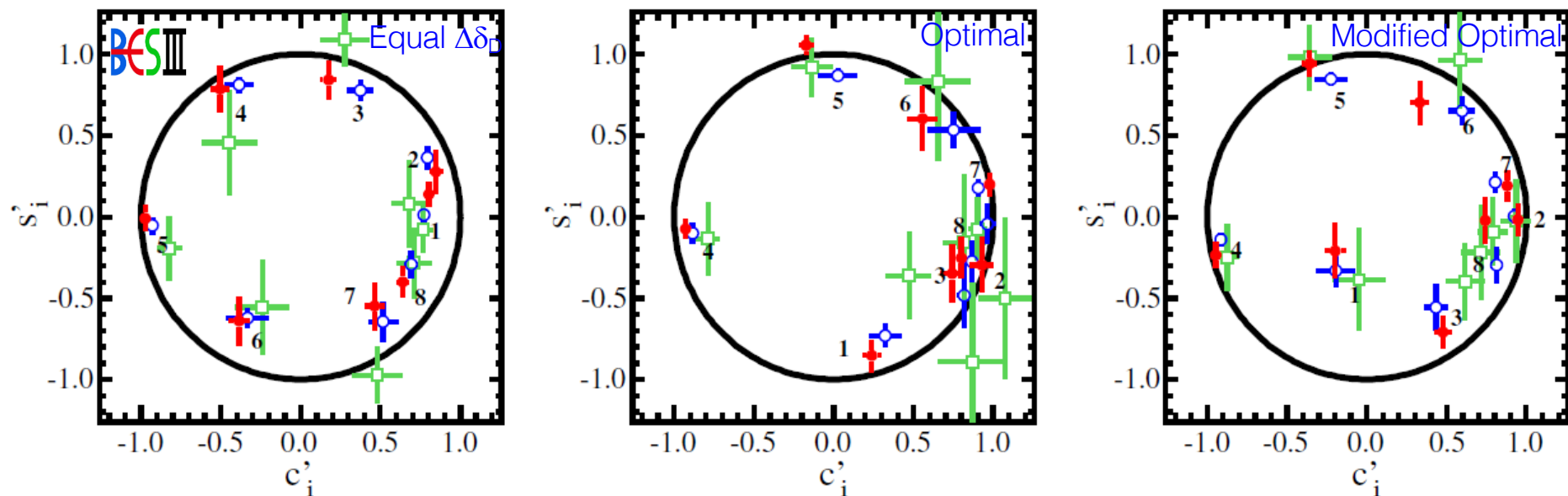


(c_i, s_i)



The c_i and s_i measured in this work (red dots with error bars), the expected results from Ref. [BaBar&Belle, PRD98, 110212 (2018)] (blue open circles) and the CLEO results (green open squares with error bars) [CLEO, PRD 82, 112006 (2010)].

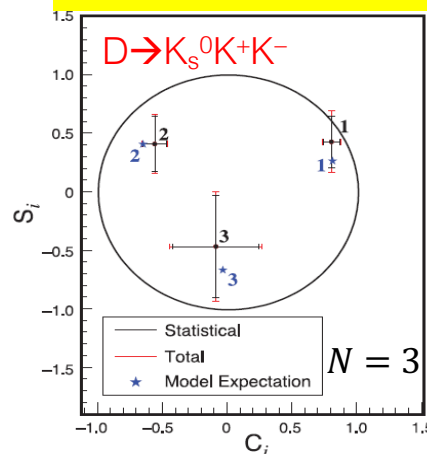
- ✓ The strong-phase parameters are **limited by statistical errors**.
- ✓ *There is no single dominant systematic uncertainty in measurement of c_i & s_i .*
- ✓ On average a factor of ~ 2.5 (2.0) more precise for c_i (s_i) than CLEO measurements.
- ✓ Using BESIII results, the associated uncertainty on γ/ϕ_3 is expected to be approximately **a factor of three** smaller than that from CLEO analysis, if using an analysis of $B^- \rightarrow DK^-$, $D \rightarrow K_S^0 \pi^+ \pi^-$.



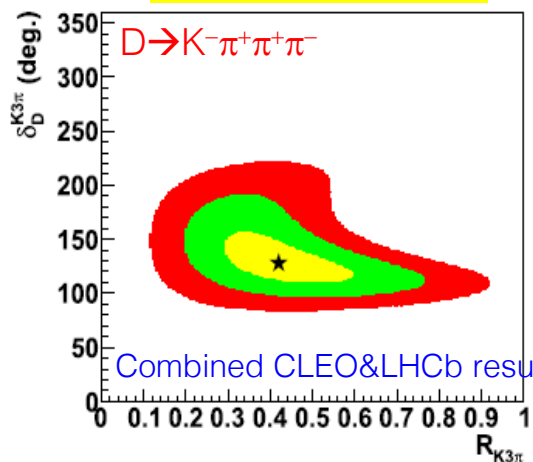
The c'_i and s'_i measured in this work (red dots with error bars), the expected results from Ref. [BaBar&Belle, PRD98, 110212 (2018)] (blue open circles) and the CLEO results (green open squares with error bars) [CLEO, PRD 82, 112006 (2010)].

- ✓ The strong-phase parameters are **limited by statistical errors**.
- ✓ *There is no single dominant systematic uncertainty in measurement of c'_i & s'_i .*
- ✓ on average a factor of ~ 2.8 (2.2) more precise for c'_i (s'_i) than CLEO measurements.
- ✓ The improved precision on c'_i and s'_i are important for Belle-II experiment in γ measurement, if using an analysis of $B^- \rightarrow DK^-$, $D \rightarrow K_L^0 \pi^+ \pi^-$.

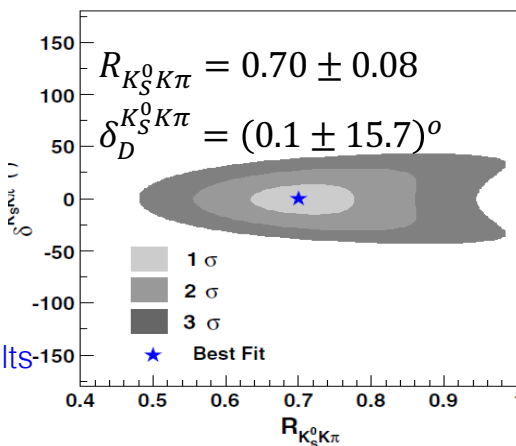
[PRD82, 112006(2010)]



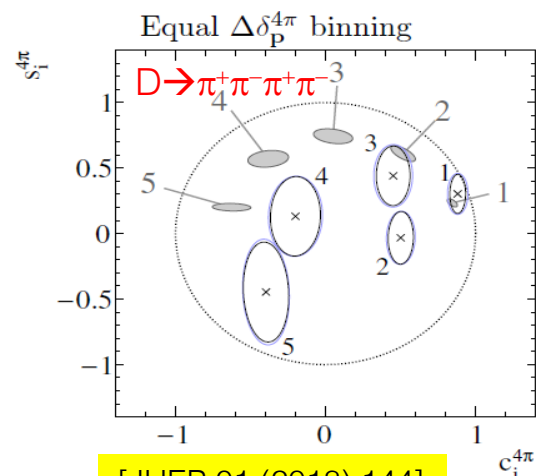
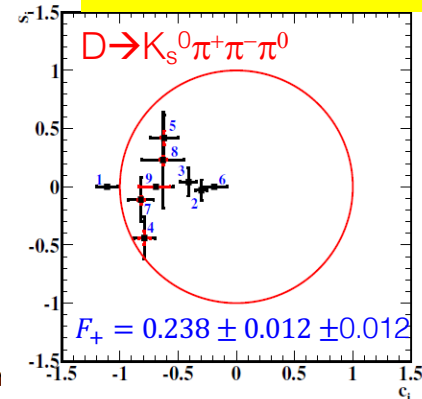
[PLB757, 520(2016)]



[PRD85, 092016(2012)]

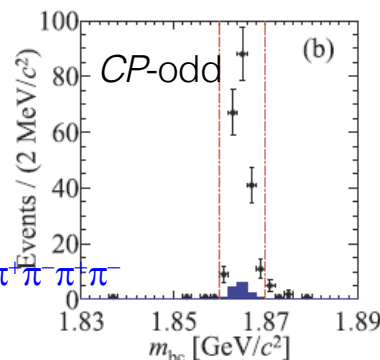
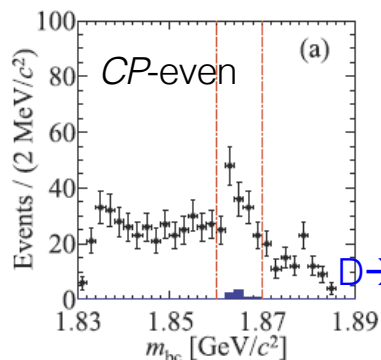


[JHEP 01 (2018) 082]



[JHEP 01 (2018) 144]

✓ CP-even fractions $F_+ = \frac{N_{CP+}}{N_{CP+} + N_{CP-}}$



$D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$

- $D \rightarrow \pi^+ \pi^- \pi^0, K^+ K^- \pi^0$
- $D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$

[PLB740, 1(2015);
 PLB 747, 9 (2015)]

- $\delta_D^{K\pi}$ was also measured by using quantum-correlated data at CLEO, but mixing inputs dominate. [PRD78,012001 & PRL100,221801]

These analyses are ongoing at BESIII

	Decay mode	Quantity of interest	Comments
➤	$D \rightarrow K_s^0 \pi^+ \pi^-$ <i>prel. release</i>	c_i and s_i	Binning schemes as those used in the CLEO-c analysis. With future, very large $\psi(3770)$ data sets, it might be worthwhile to explore alternative binning.
➤	$D \rightarrow K_s^0 K^+ K^-$	c_i and s_i	Binning schemes as those used in the CLEO-c analysis. With future, very large $\psi(3770)$ data sets, it might be worthwhile to explore alternative binning.
➤	$D \rightarrow K^\pm \pi^\mp \pi^+ \pi^-$	R, δ	In bins guided by amplitude models, currently under development by LHCb.
⇨	$D \rightarrow K^+ K^- \pi^+ \pi^-$	c_i and s_i	Binning scheme can be guided by the CLEO model [18] or potentially an improved model from LHCb in the future.
⇨	$D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$	F_+ or c_i and s_i	Unbinned measurement of F_+ . Measurements of F_+ in bins or c_i and s_i in bins could be explored.
➤	$D \rightarrow K^\pm \pi^\mp \pi^0$	R, δ	Simple 2-3 bin scheme could be considered.
⇨	$D \rightarrow K_s^0 K^\pm \pi^\mp$	R, δ	Simple 2 bin scheme where one bin encloses the K^* resonance.
➤	$D \rightarrow \pi^+ \pi^- \pi^0$	F_+	No binning required as $F_+ \sim 1$.
⇨	$D \rightarrow K_s^0 \pi^+ \pi^- \pi^0$	F_+ and c_i and s_i	Unbinned measurement of F_+ required. Additional measurements of F_+ or c_i and s_i in bins could be explored.
➤	$D \rightarrow K^+ K^- \pi^0$	F_+	Unbinned measurement required. Extensions to binned measurements of either F_+ or c_i and s_i possible.
➡	$D \rightarrow K^\pm \pi^\mp$	δ	Of low priority due to good precision available through charm-mixing analyses.

LHCb-PUB-2016-025

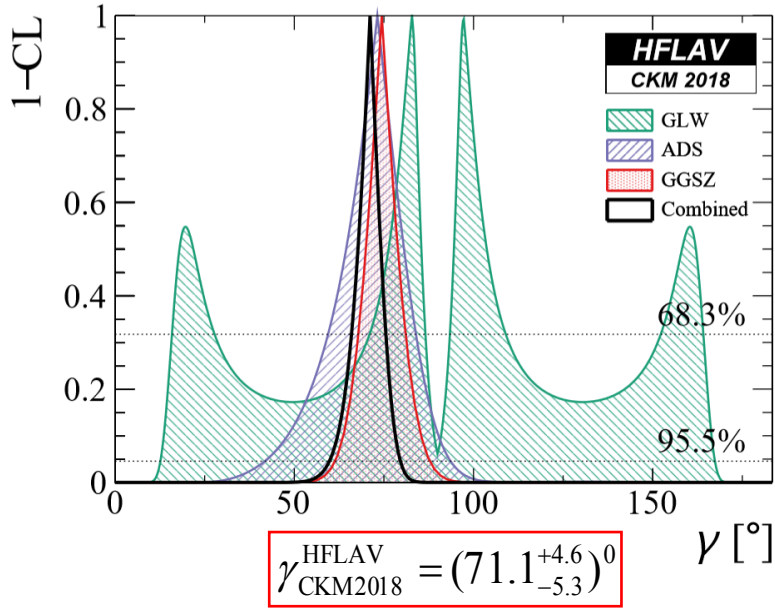
Status at BESIII

➡ published

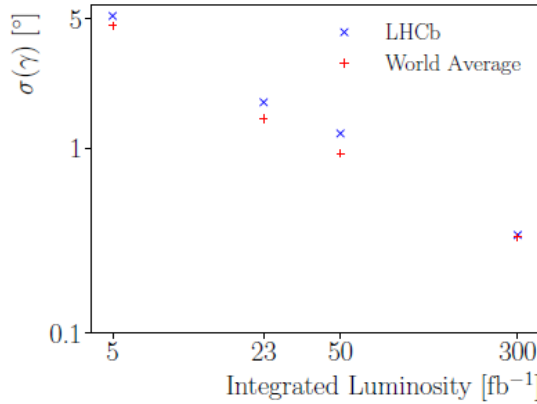
➤ under study

⇨ in plan

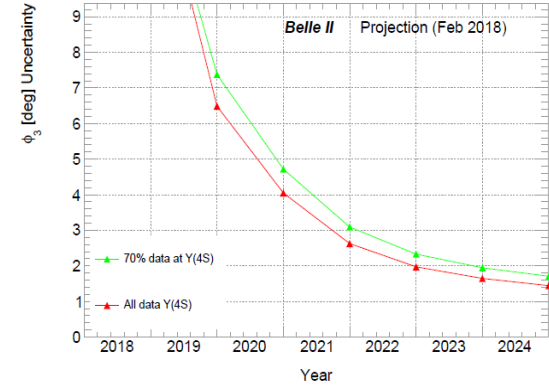




[arXiv: 1808.08865]



[arXiv: 1808.10567]



- ✓ dominated by LHCb measurements
The latest LHCb γ combination gives

$$\gamma = (74.0^{+5.0}_{-5.8})^\circ$$

- ✓ The degree-level precision on γ can be expected in near future.

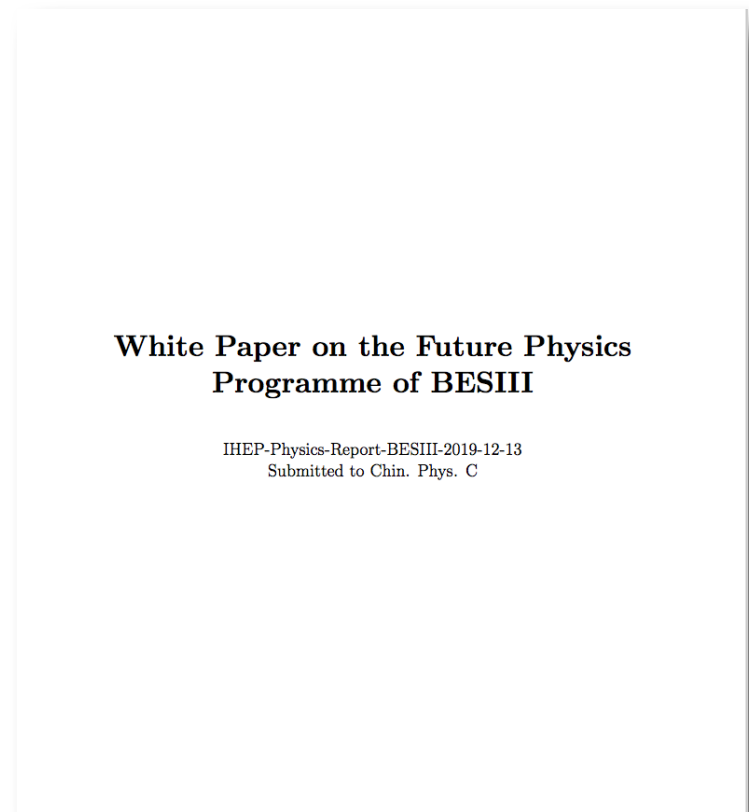
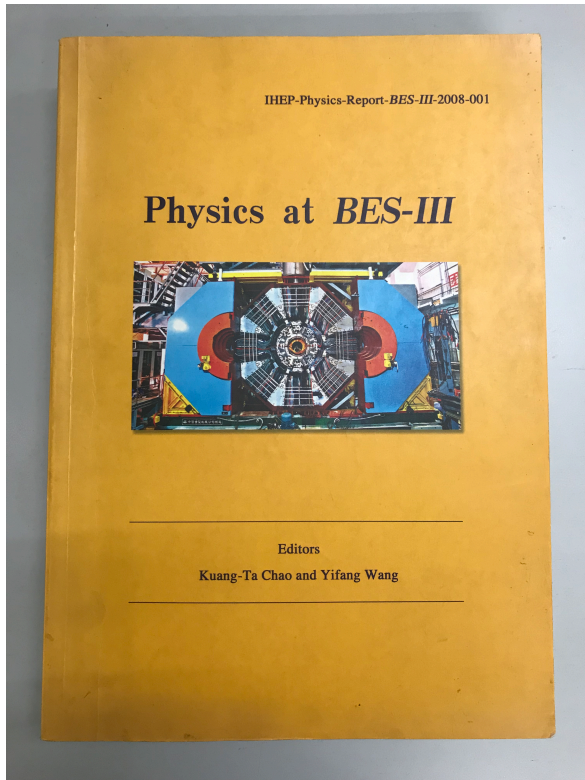
	Collected/ Expected luminosity	Year attained	γ sensitivity
LHCb Run-1 [7, 8 TeV]	3 fb^{-1}	2012	5.5 $^\circ$
LHCb Run-2 [13 TeV]	5 fb^{-1}	2018	2.8 $^\circ$
LHCb phase-1 upgrade [14 TeV]	50 fb^{-1}	2030	0.71 $^\circ$
LHCb phase-2 upgrade [14 TeV]	300 fb^{-1}	2035(?)	0.28 $^\circ$
Belle-II Run	50 ab^{-1}	2025	1.5 $^\circ$

- In above table, sensitivity from LHCb is obtained by scaling Run-I statistical error.

From BESIII physics (yellow) book to BESIII white paper

2008

arXiv:1912.05983



Future BESIII precision



Energy	Physics motivations	Current data	Expected final data
1.8 - 2.0 GeV	R values Nucleon cross-sections	N/A	0.1 fb ⁻¹ (fine scan)
2.0 - 3.1 GeV	R values Cross-sections	Fine scan (20 energy points)	Complete scan (additional points)
J/ψ peak	Light hadron & Glueball J/ψ decays	3.2 fb ⁻¹ (10 billion)	3.2 fb ⁻¹ (10 billion)
$\psi(3686)$ peak	Light hadron & Glueball Charmonium decays	0.67 fb ⁻¹ (0.45 billion)	4.5 fb ⁻¹ (3.0 billion)
$\psi(3770)$ peak	D^0/D^\pm decays	2.9 fb ⁻¹	20.0 fb ⁻¹
3.8 - 4.6 GeV	R values XYZ/Open charm	Fine scan (105 energy points)	No requirement
4.180 GeV	D_s decay XYZ/Open charm	3.2 fb ⁻¹	6 fb ⁻¹

- 2.9/fb $\psi(3770)$ data: strong phase uncertainty to γ is $\sim 1^\circ$
- 20/fb $\psi(3770)$ data: strong phase uncertainty to γ is $\sim 0.4^\circ$
 → sufficient for LHCb upgrade I
- While not enough for LHCb upgrade II

Runs	Collected / Expected integrated luminosity	Year attained	γ/ϕ_3 sensitivity
LHCb Run-1 [7, 8 TeV]	3 fb ⁻¹	2012	8°
LHCb Run-2 [13 TeV]	5 fb ⁻¹	2018	4°
Belle II Run	50 ab ⁻¹	2025	1.5°
LHCb upgrade I [14 TeV]	50 fb ⁻¹	2030	< 1°
LHCb upgrade II [14 TeV]	300 fb ⁻¹	(>)2035	< 0.4°

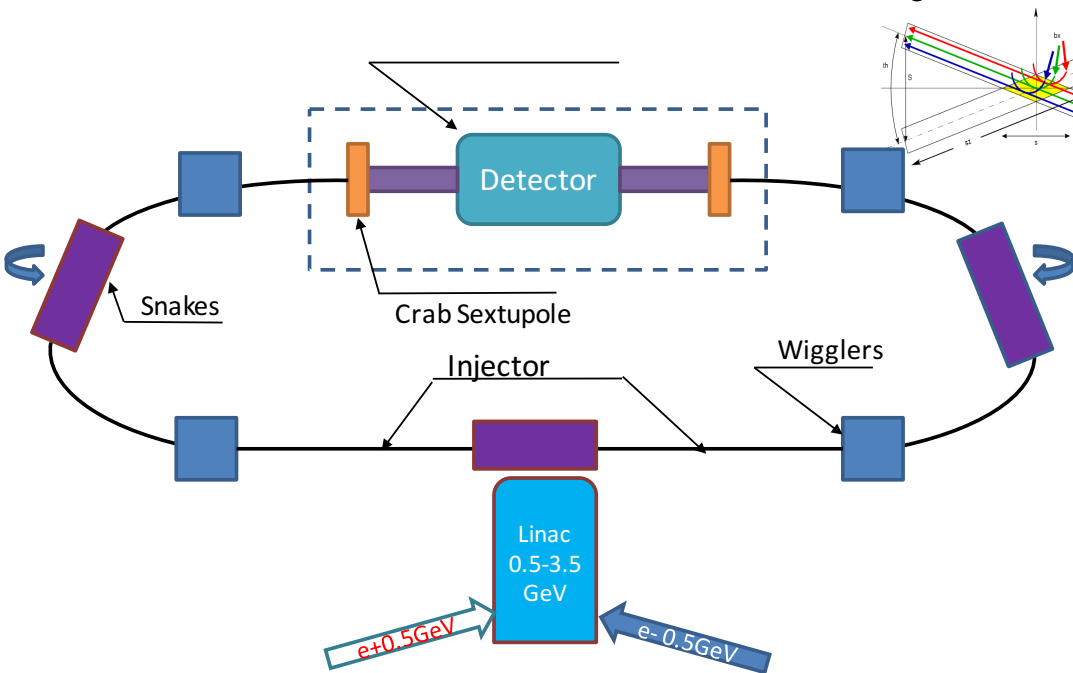
BESIII 20/fb:
 $\sigma(\gamma) \sim 0.4^\circ$

→ STCF is needed!

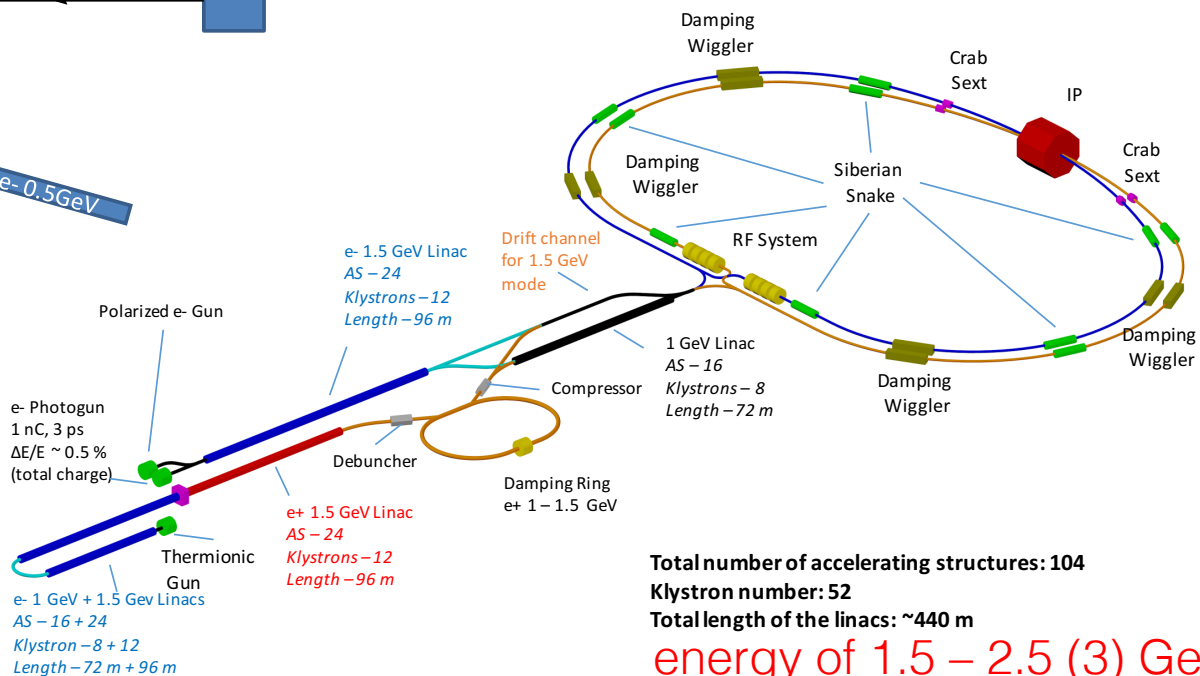
Proposals of the Super Tau-Charm Factory (STCF)



HIEPA in China



Super-CT Project in Russia





Prospects at the STCF

Data samples with 1 ab^{-1} integral luminosity

Data Set	STCF					Belle II		
	process	σ/nb	N	ST eff./%	ST N	σ/nb	N	Tag N
J/ψ	–	–	1.0×10^{12}	–	–	–	–	–
$\psi(2S)$	–	–	3.0×10^{11}	–	–	–	–	–
D^0	$D^0 \bar{D}^0(3.77)$	~ 3.6	3.6×10^9	10.8	0.78×10^9	–	1.4×10^9	–
D^+	$D^+ D^-(3.77)$	~ 2.8	2.8×10^9	9.4	0.53×10^9	–	7.7×10^8	–
D_s	$D_s D_s^*(4.18)$	~ 0.9	0.9×10^9	6.0	0.11×10^9	–	2.5×10^8	–
τ^+	$\tau^+ \tau^-(3.68)$	~ 2.4	2.4×10^9	–	–	0.9	0.9×10^9	–
	$\tau^+ \tau^-(4.25)$	~ 3.6	3.5×10^9	–	–	–	–	–
Λ_c	$\Lambda_c \Lambda_c(4.64)$	~ 0.6	5.5×10^8	5.0	0.55×10^8	–	1.6×10^8	$3.6 \times 10^{4*}$

The luminosity is 1.0 ab^{-1} . * process $e^+e^- \rightarrow D^{(*)-} \bar{p} \pi^+ \Lambda_c^+$.

- Belle-II (50/ab) has 50~100 times more statistics
- STCF is expected to have higher **detection efficiency**
- STCF has low backgrounds for productions at threshold



CKM unitarity triangle

Runs	Collected / Expected integrated luminosity	Year attained	γ/ϕ_3 sensitivity
LHCb Run-1 [7, 8 TeV]	3 fb ⁻¹	2012	8°
LHCb Run-2 [13 TeV]	5 fb ⁻¹	2018	4°
Belle II Run	50 ab ⁻¹	2025	1.5°
LHCb upgrade I [14 TeV]	50 fb ⁻¹	2030	< 1°
LHCb upgrade II [14 TeV]	300 fb ⁻¹	(>)2035	< 0.4°

BESIII 20/fb:
 $\sigma(\gamma) \sim 0.4^\circ$

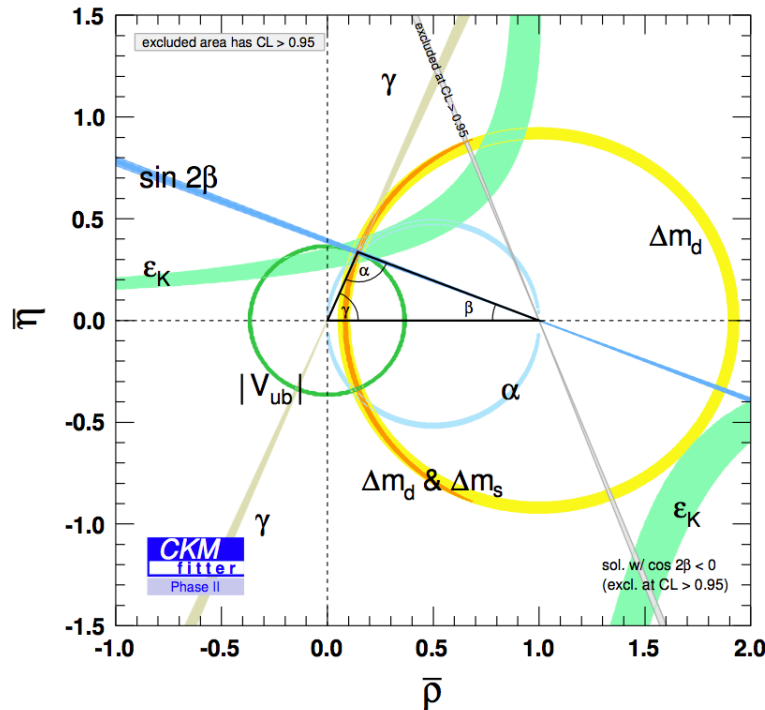
→ STCF is needed!

- ADS method: use *D* doubly Cabibbo-suppressed decays, e.g. $D^0 \rightarrow K^+ \pi^-$
 - With 1 ab⁻¹ @ STCF : $\sigma(\cos\delta_{K\pi}) \sim 0.007$; $\sigma(\delta_{K\pi}) \sim 2^\circ \rightarrow \sigma(\gamma) < 0.5^\circ$
- GGSZ method: use Dalitz plot analysis of 3-body D^0 decays, e.g. $K_S \pi^+ \pi^-$
 - STCF reduces the contribution of *D* Dalitz model to a level of $\sim 0.1^\circ$, since expected precision from future HL-LHCb projects would be $< 0.4^\circ$.
- Cross checks among different methods and modes are crucial to test their consistency and searching for new physics



Scenario beyond 2035

STCF will provide complementary information on the strong phase and allow detailed comparisons in different models



- 300 /fb for LHCb
- 3000 /ab for CMS/ATLAS
- 50 /ab for Belle II
- 20 /fb @ 3773MeV for BESIII

Decay mode	Quantity of interest
$D \rightarrow K_S^0 \pi^+ \pi^-$	c_i and s_i
$D \rightarrow K_S^0 K^+ K^-$	c_i and s_i
$D \rightarrow K^\pm \pi^\mp \pi^+ \pi^-$	R, δ
$D \rightarrow K^+ K^- \pi^+ \pi^-$	c_i and s_i
$D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$	F_+ or c_i and s_i
$D \rightarrow K^\pm \pi^\mp \pi^0$	R, δ
$D \rightarrow K_S^0 K^\pm \pi^\mp$	R, δ
$D \rightarrow \pi^+ \pi^- \pi^0$	F_+
$D \rightarrow K_S^0 \pi^+ \pi^- \pi^0$	F_+, c_i and s_i
$D \rightarrow K^+ K^- \pi^0$	F_+
$D \rightarrow K^\pm \pi^\mp$	δ



Possible **New** Charm physics at LHCb

Opportunities

- ▶ semileptonic charm hadron decays differential rate

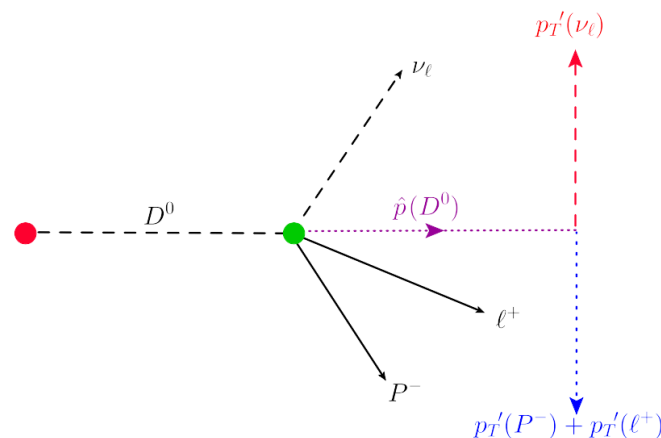
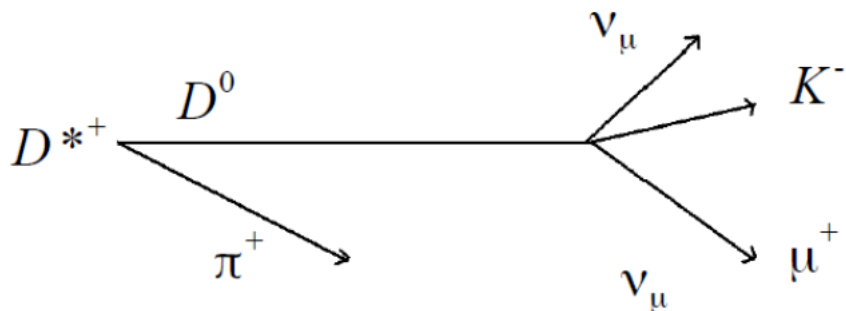
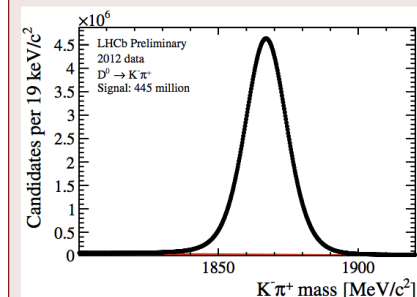
$$\frac{d\Gamma(X_c \rightarrow h\ell\nu)}{dq^2} = (\text{Constants}) \cdot |V_{cy}|^2 \cdot F[\text{form factors}(q^2, m_\ell)], \quad y = d, s$$

$$\frac{d\Gamma(D \rightarrow P\ell\nu)}{dq^2} = \frac{G_F^2}{24\pi^3} \cdot |V_{cy}|^2 \cdot |f_{DP}^+(q^2)|^2 \cdot p^{*3}, \quad y = d, s$$

- ▶ experimental measurements can:
 - ▶ measure $|V_{cs}|$, $|V_{cd}|$ and/or form factors
 - ▶ measure D^0 mixing and CPV
 - ▶ test Lepton Flavour Universality on charm decays

LHCb is a charm factory

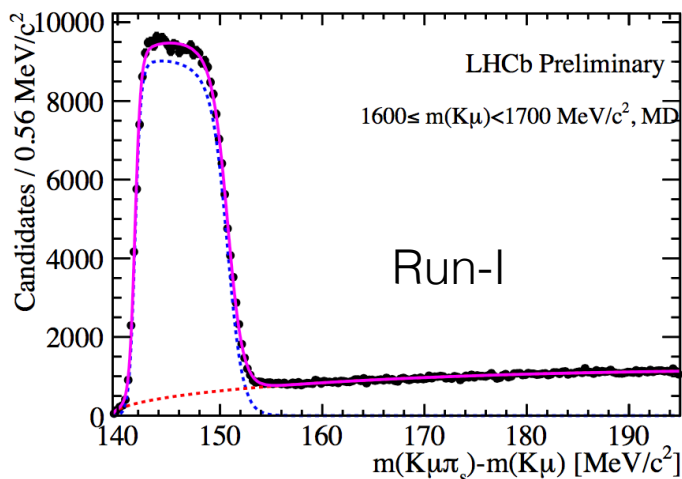
- ▶ 1.4 billion charm hadrons reconstructed in 2010-2015 (LHCb-CONF-2016-005)
- ▶ higher efficiency since 2015



Neutrino reconstruction using D meson flight direction

Many ongoing LHCb charmed-meson SL analyses are using promote production $D^{*+} \rightarrow D^0 \pi^+$

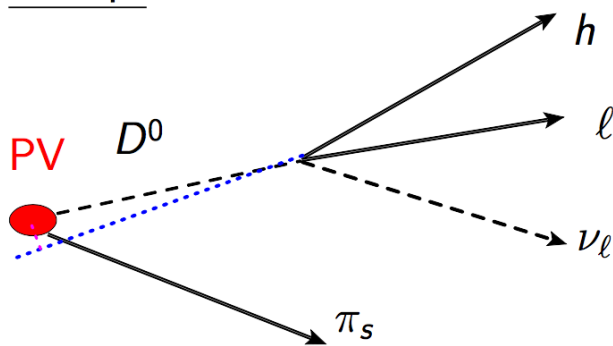
- Pros: high production rate, ...
- Cons:
 - ✓ hard to well control signal shape and non-trivial background treatment in the fit
 - ✓ optimal only for the neutral D meson



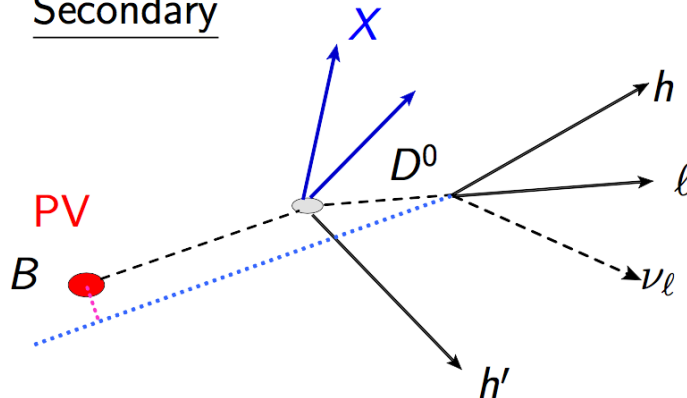
V_{cd} / V_{cs} with world-leading precision	
▶ measure	$\frac{ V_{cd} ^2}{ V_{cs} ^2}$ using $\frac{\mathcal{B}(D^0 \rightarrow \pi^- \mu^+ \nu)}{\mathcal{B}(D^0 \rightarrow K^- \mu^+ \nu)}$
Mixing/CPV measurements on semileptonic D decays	
▶ measure	$\frac{N(D^0 \rightarrow K^+ \mu^- \nu_\mu)}{N(D^0 \rightarrow K^- \mu^+ \bar{\nu}_\mu)}$
Test lepton universality in D meson decays	
▶ measure	$\frac{\mathcal{B}(D \rightarrow h \mu \nu)}{\mathcal{B}(D \rightarrow h e \nu)}$

- SL decays of many other charmed baryons can be studied in a similar fashion: $\Sigma_c^0 \rightarrow \Lambda_c^+ \pi^-$, $\Xi_c(2645) \rightarrow \Xi_c \pi$, Ω_c , Ξ_{cc} etc.

Prompt



Secondary

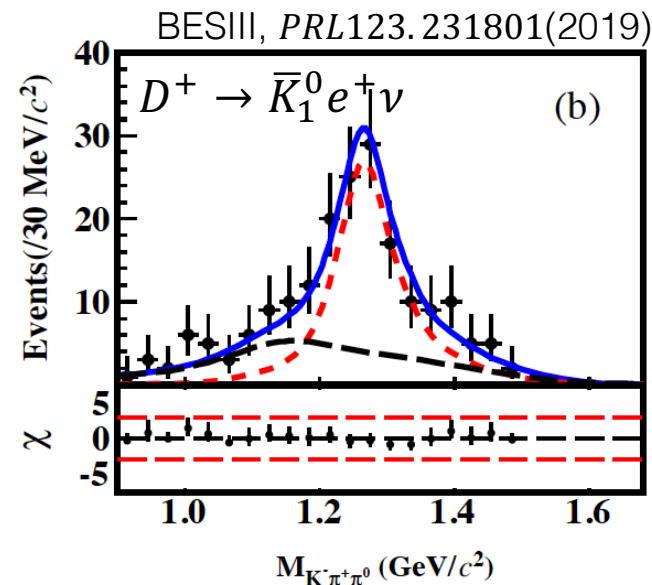
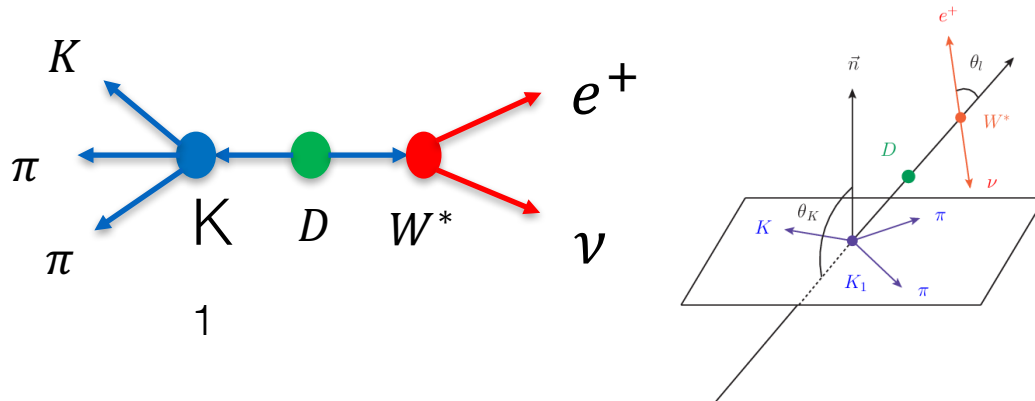


- RUN I+II: keep increasing statistics
- We can study charm SL decays using b -hadron decays, where we expect much lower background level
- More importantly, we will be able to extend to all the charmed hadrons, including **the charged D^+/D_s^+ meson** and many other charged baryons: $\Lambda_c, \Xi_c, \Omega_c$

$$\begin{array}{l} \overline{D^0} \\ D^0 \rightarrow \pi \mu \nu \quad (\mathcal{B} = 0.238 \pm 0.024\%) \\ D^0 \rightarrow K \mu \nu \quad (\mathcal{B} = 3.3 \pm 0.13\%) \\ D^0 \rightarrow K^*(892)^- \mu \nu \quad (\mathcal{B} = 1.92 \pm 0.25\%) \end{array}$$

$$\begin{array}{l} D^+ \\ D^+ \rightarrow K \pi \mu \nu \quad (\mathcal{B} = 3.9 \pm 0.4\%) \\ D^+ \rightarrow K^0 \mu \nu \quad (\mathcal{B} = 9.3 \pm 0.7\%) \\ D^+ \rightarrow K^{*0} \mu \nu \quad (\mathcal{B} = 5.3 \pm 0.15\%) \\ D^+ \rightarrow \eta \mu \nu \quad (\mathcal{B} \sim 1\%) \end{array}$$

$$\begin{array}{l} D_s \\ D_s^+ \rightarrow \phi \mu \nu \quad (\mathcal{B} \sim 2\%) \\ D_s^+ \rightarrow K^0 \mu \nu \quad (\mathcal{B} \sim 0.3\%) \\ D_s^+ \rightarrow \eta^{(\prime)} \mu \nu \quad (\mathcal{B} \sim 3\%) \end{array}$$



ratio of up-down asymmetries

33

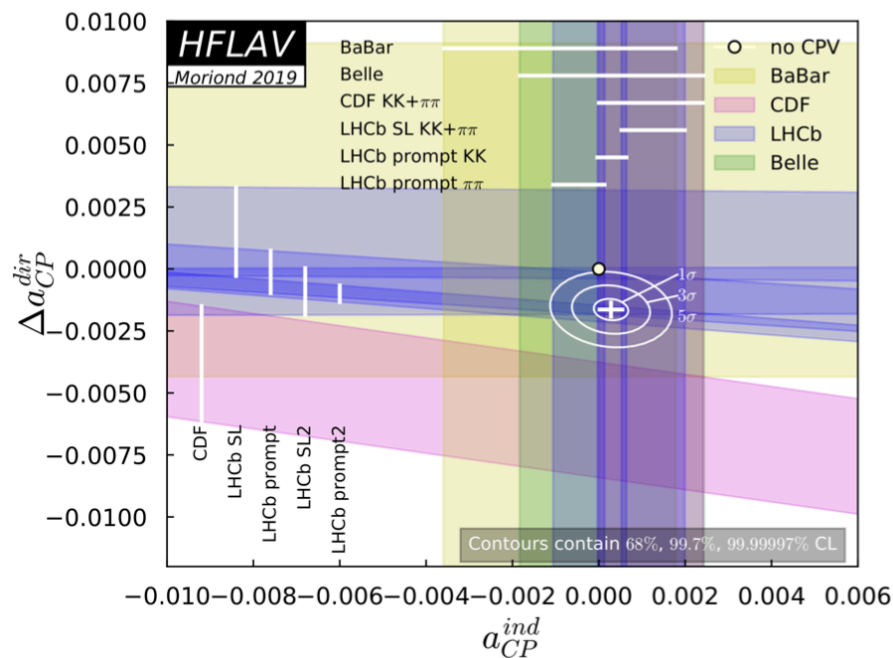
$$\mathcal{A}'_{UD} \equiv \frac{\left[\int_0^1 - \int_{-1}^0 \right] d \cos \theta_K \frac{d\Gamma_{K_1 e^+ \nu}}{d \cos \theta_K}}{\left[\int_0^1 - \int_{-1}^0 \right] d \cos \theta_l \frac{d\Gamma_{K_1 e^+ \nu}}{d \cos \theta_l}} \quad D \rightarrow K_1 (\rightarrow K\pi\pi) e^+ \nu$$

$$\mathcal{A}'_{UD} = \frac{\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]}{|\vec{J}|^2}$$

$$\begin{aligned} \mathcal{A}_{UD} &\equiv \frac{\left[\int_0^1 - \int_{-1}^0 \right] d \cos \theta_K \frac{d\hat{\Gamma}_{K_1 \gamma}}{d \cos \theta_K}}{\left[\int_0^1 + \int_{-1}^0 \right] d \cos \theta_K \frac{d\hat{\Gamma}_{K_1 \gamma}}{d \cos \theta_K}} \\ &= \lambda_\gamma \frac{3 \text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]}{4 |\vec{J}|^2} \quad B \rightarrow K_1 (\rightarrow K\pi\pi) \gamma \end{aligned}$$

- LHCb can measure mounic mode:
 $D^0 \rightarrow K_1^- \mu^+ \nu, K_1^- \rightarrow K^- \pi^+ \pi^-$
- Provides model-independent input to extract photon polarization in $B \rightarrow K_1 \gamma \rightarrow (K\pi\pi) \gamma$
[Wei Wang et al, arXiv:1909.13083]

Charm CPV



→ Updated HFLAV fit

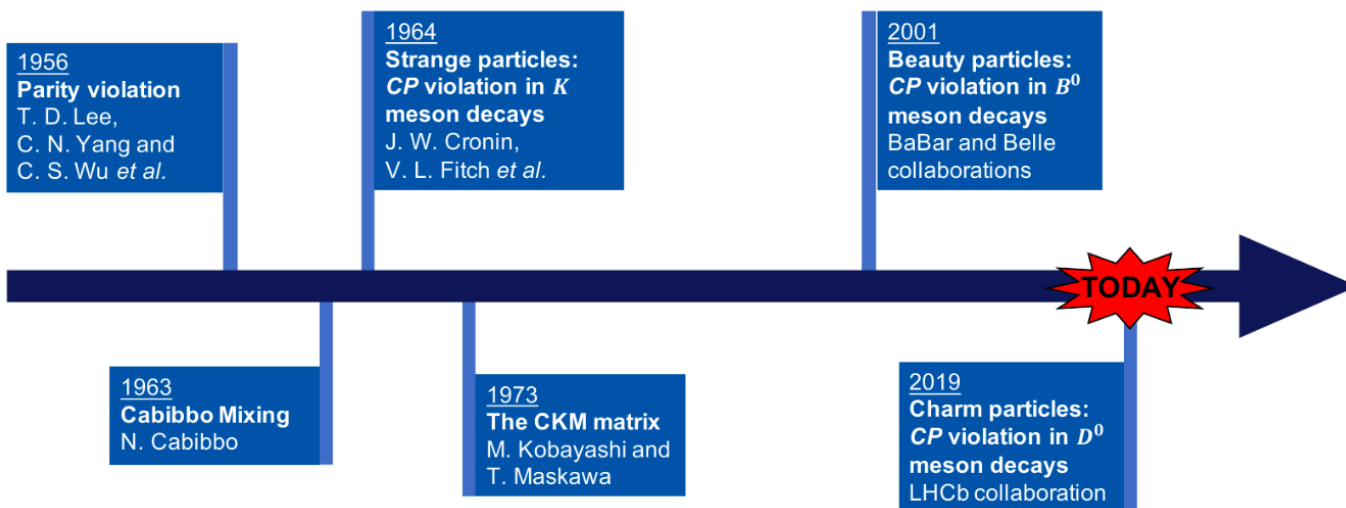
$$\Delta a_{CP}^{dir} = (-16.4 \pm 2.8) \times 10^{-4}$$

$$\Delta a_{CP}^{ind} = (2.8 \pm 2.6) \times 10^{-4}$$

→ Compatible with SM

- Most predictions on $10^{-4} - 10^{-3}$ level

→ Observation in other channels could provide a confirmation of this effect

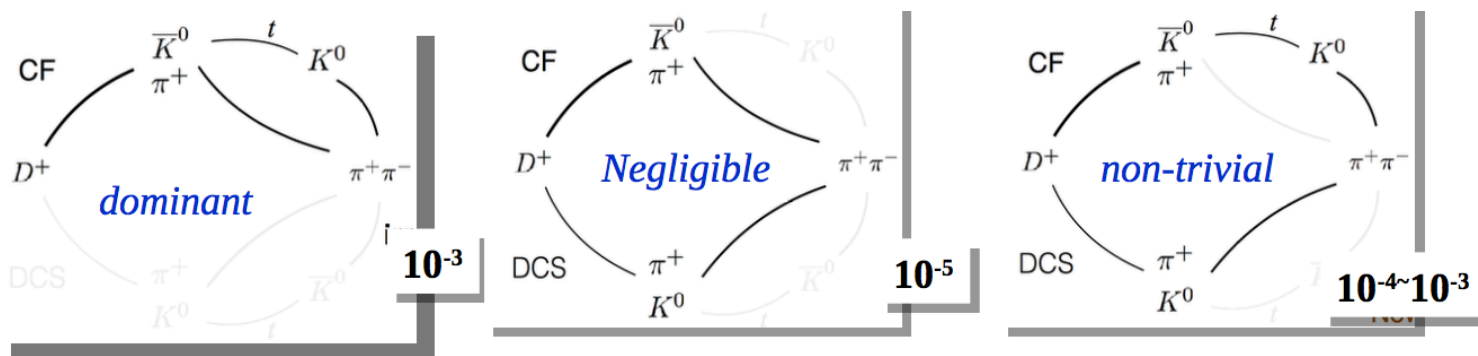


- New CP-violation effect, in order up to 10^{-3} , in charm decays into neutral kaons was proposed [Phys. Rev. Lett. 119, 181802 (2017)].
- This effect emerges from interference between CF and DCS amplitudes with the mixing of final-state kaons by Fu Sheng et al
- Signal channels: $D^+ \rightarrow K_S^0 \pi^+$, $D_s^+ \rightarrow K_S^0 K^+$
- Control channels: $D^+ \rightarrow K^- \pi^+ \pi^+$, $D_s^+ \rightarrow K^+ K^- \pi^+$

$$\Delta A_{CP} = A_{CP}(D_s^+ \rightarrow K_S K^+) - A_{CP}(D^+ \rightarrow K_S \pi^+) =$$

$$[(A_{raw}(D_s^+ \rightarrow K_S K^+) - A_{raw}(D_s^+ \rightarrow K^+ K^- \pi^-)) -$$

$$(A_{raw}(D^+ \rightarrow K_S \pi^+) - A_{raw}(D^+ \rightarrow K^- \pi^+ \pi^+))]$$



- Sensitivity expectation with Run II data (prospect with the full Run II data):

$$D^+ \rightarrow K_S^0 \pi^+ : \delta \mathcal{A}_{CP}^{raw} = 0.03 \%$$

$$D_s^+ \rightarrow K_S^0 K^+ : \delta \mathcal{A}_{CP}^{raw} = 0.07 \%$$

Summary and Outlook

- Quantum correlation of $D^0\bar{D}^0$ pairs produced at threshold is a unique technique to access information of strong-phase parameters in D hadronic decays, which provide the key inputs in ranges of γ measurements, $D^0\bar{D}^0$ mixing and CPV studies.
- A range of quantum-correlated studies are undergoing at BESIII:
 - ✓ Studies in $D \rightarrow K_{S/L}\pi^+\pi^-$ show excellent preliminary results. These results will have important impacts over a wide range studies in flavour physics.
 - ✓ Current BESIII $\psi(3770)$ data will constrain the strong phase uncertainty of γ to be $\sim 1^\circ$
- The future 20 fb^{-1} $\psi(3770)$ data and STCF will be crucial to improve the precision of γ measurement to the level of 0.1°
- Some ideas, like studying semi-leptonic decays and charm CPV searches, at LHCb are discussed

Thank you!

谢谢!

Strong-phase parameters in $D \rightarrow K_{S/L}^0 \pi^+ \pi^-$ at $\psi(3770)$

□ $\psi(3770)$ is a spin -1 state and therefore the amplitude of $\psi(3770) \rightarrow D^0 D^0$:

$$(|D^0\rangle|\bar{D}^0\rangle - |\bar{D}^0\rangle|D^0\rangle)/\sqrt{2} \quad \text{[anti-symmetric wave function]}$$

✓ For CP-tagged $K_S^0 \pi^+ \pi^-$, its amplitude is expressed by:

$$f_{CP\pm} = \frac{1}{\sqrt{2}} [f_D(m_+^2, m_-^2) \pm f_D(m_-^2, m_+^2)]$$

The expected yields in Dalitz Plot (DP) bins:

$$\rightarrow M_i^\pm = h_{CP\pm} (K_i \pm 2c_i \sqrt{K_i K_{-i}} + K_{-i})$$

✓ Similarly, for $K_S^0 \pi^+ \pi^-$ vs. $K_S^0 \pi^+ \pi^-$, the expected yields in DP bins is:

$$\rightarrow M_{ij} = h_{\text{corr}} [K_i K_{-j} + K_{-i} K_j - 2\sqrt{K_i K_{-j} K_{-i} K_j} (c_i c_j + s_i s_j)]$$

c_i & s_i are obtained by studying DT events: $K_S^0 \pi^+ \pi^-$ vs. CP-tag & $K_S^0 \pi^+ \pi^-$ vs. $K_S^0 \pi^+ \pi^-$ tags

Here $h_{CP\pm}$ and h_{corr} are the normalization factors related to yields of single tags and the number of neutral D meson pairs.

From above equations, the precision on s_i is constrained by the mode $K_S^0 \pi^+ \pi^-$ vs. $K_S^0 \pi^+ \pi^-$.

- To improve the precision of measuring \mathbf{s}_i , the $K_S^0\pi^+\pi^-$ vs. $K_L^0\pi^+\pi^-$ events are added, which is dependent on (c_i, s_i, c_i' and s_i'). Due to similarities between the decays, weak model assumptions^[1,2,3] can provide a constraint on the differences between c_i and c_i' , s_i and s_i' .

$$K_S^0\pi^+\pi^- \text{ vs. } K_L^0\pi^+\pi^- \quad M'_{ij} = h'_{\text{corr}} \left[K_i K'_{-j} + K_{-i} K'_j + 2\sqrt{K_i K'_{-j} K_{-i} K'_j} (c_i c'_j + s_i s'_j) \right]$$

$$CP \text{ tag vs. } K_L^0\pi^+\pi^- \quad M'_i{}^{\pm} = h'_{CP\pm} (K'_i \mp 2c'_i \sqrt{K'_i K'_{-i}} + K'_{-i})$$

- The c_i' and s_i' parameters are useful for Belle-II experiment if they use the decay mode $B \rightarrow DK$, with $D \rightarrow K_L^0\pi^+\pi^-$ to measure γ .
- The strong-phase parameters are obtained by minimizing the log-likelihood function constructed by using the observed and expected yields.

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