

Photon Helicity in $b \rightarrow s\gamma$ Decays towards New Physics Search

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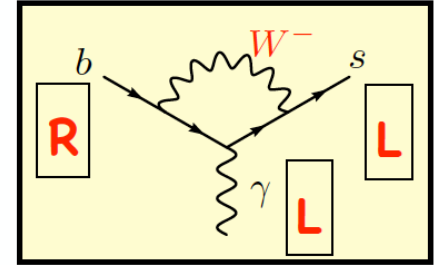
1st LHCb Physics Frontier
14-15 Dec 2019



OUTLINE

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- Photon polarization in $b \rightarrow s\gamma$



- Recent Progresses on photon polarization
- Novel Approach to determine the photon helicity
Wang, Yu, Zhao, 1909.13083
- Summary



BEYOND SM: THREE FRONTIERS

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- High Energy Frontier:
Tevatron, LHC...
Direct search

- High Precision Frontier:
B factories Tau/charm
factory ...
indirect search



BEYOND SM: INTENSITY FRONTIER

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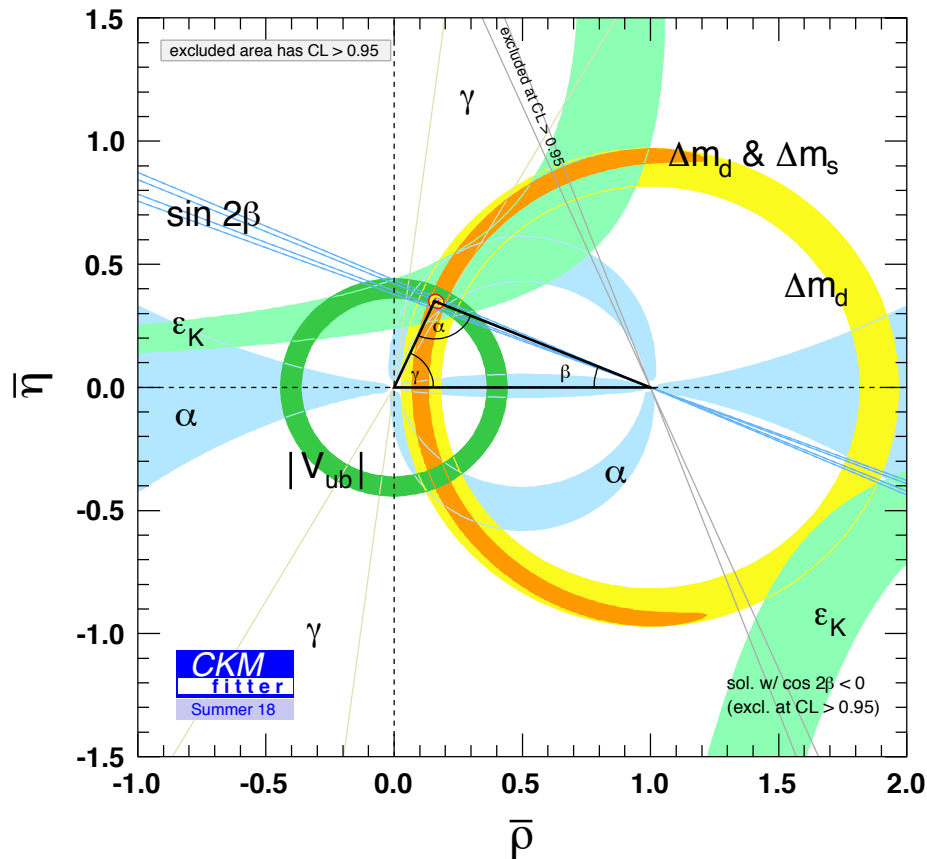
If the LHC/xxx did not discover any new particle beyond SM, precision study becomes an ideal platform to detect NP effects.

If the LHC/xxx discovers new elementary particles beyond SM, then precision physics will be necessary to constrain the underlying framework.



BOTTOM PHYSICS

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➤ Looks great, but can be deceived (tension)

➤ O(10%-15%) NP is still allowed



R_{K^*}/R_K ANOMALY:

$$R_{K^*}[q_{\min}^2, q_{\max}^2] \equiv \frac{\int_{q_{\min}^2}^{q_{\max}^2} dq^2 d\Gamma(B \rightarrow K^* \mu^+ \mu^-)/dq^2}{\int_{q_{\min}^2}^{q_{\max}^2} dq^2 d\Gamma(B \rightarrow K^* e^+ e^-)/dq^2}$$

q^2 :invariant mass of lepton pair

Observable	SM results	Experimental data	
$R_K : q^2 = [1, 6] \text{ GeV}^2$	1.00 ± 0.01	$0.745_{-0.074}^{+0.090} \pm 0.036$	2.6 Sigma
$R_{K^*}^{\text{low}} : q^2 = [0.045, 1.1] \text{ GeV}^2$	$0.920_{-0.006}^{+0.007}$	$0.66_{-0.07}^{+0.11} \pm 0.03$	2.3 Sigma
$R_{K^*}^{\text{central}} : q^2 = [1.1, 6] \text{ GeV}^2$	0.996 ± 0.002	$0.69_{-0.07}^{+0.11} \pm 0.05$	2.5 Sigma

SM: Geng, et.al, 1704.05446

LHCb: PRL, 113, 151601(2014)

LHCb: 1705.05802

Anomalies in heavy flavor:
New Physics or QCD contaminations ?



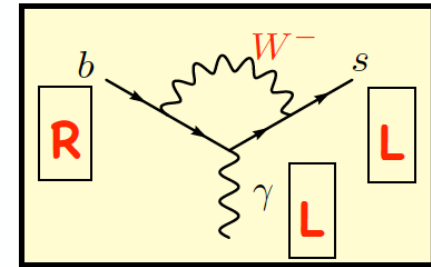
Photon polarization in $b \rightarrow s\gamma$

PHOTON POLARIZATION OF

$$b \rightarrow s\gamma$$

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- The photon polarization of the $b \rightarrow s\gamma$ process has a unique sensitivity to BSM with right-handed couplings.
- However, the photon polarization has never been measured at a **high precision** so far: an important challenge for LHCb and Belle II.



W-boson couples
only left-handed



γ from $b \rightarrow s\gamma$ should
be mostly left-handed

- ✓ $b \rightarrow s\gamma$: left-handed polarization
- ✓ $\bar{b} \rightarrow \bar{s}\gamma$: right-handed polarization



HOW DO WE MEASURE THE POLARIZATION?

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➤ Time-dependent measurements:

$$\checkmark B_d \rightarrow K_S \pi^0 \gamma, B_d \rightarrow \rho \gamma$$

$$\checkmark B_d \rightarrow K_S \pi^+ \pi^- \gamma$$

$$\checkmark B_S \rightarrow K^+ K^- \gamma$$

LHCb(2013):

P_{Λ_b} is “small” :

$(0.06 \pm 0.07 \pm 0.02)$

➤ Angular distribution :

✓ Baryonic decays: $\Lambda_b \rightarrow \Lambda \gamma$, request the polarization of Λ_b or Λ

$$\checkmark B \rightarrow K_{res} (\rightarrow K \pi \pi) \gamma$$



NEW PHYSICS CONTRIBUTIONS IN $b \rightarrow s\gamma$

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$$\mathcal{M}(b \rightarrow s\gamma) \simeq -\frac{4G_F}{\sqrt{2}} V_{ts}^* V_{tb} \left[\underbrace{(C_{7\gamma}^{\text{SM}} + C_{7\gamma}^{\text{NP}}) \langle \mathcal{O}_{7\gamma} \rangle}_{\propto \mathcal{M}_L} + \underbrace{C_{7\gamma}'^{\text{NP}} \langle \mathcal{O}'_{7\gamma} \rangle}_{\propto \mathcal{M}_R} \right]$$

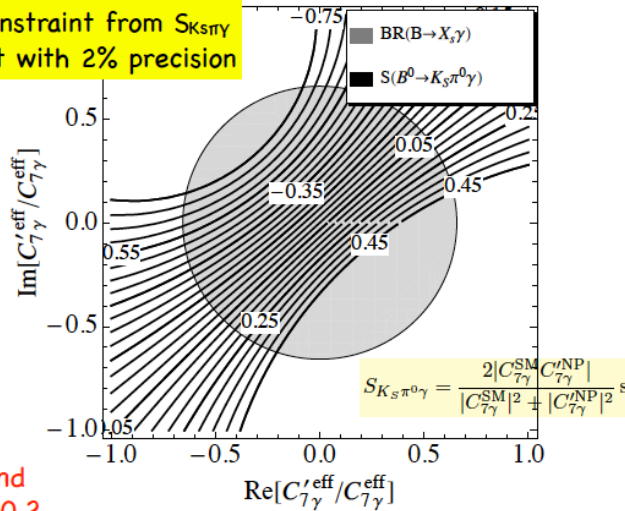
Note:

In NP, $C_{7\gamma}^{\text{NP}}$ and/or $C_{7\gamma}'^{\text{NP}}$ can be complex numbers!
We only consider $C_{7\gamma}'^{\text{NP}}$ in the following.



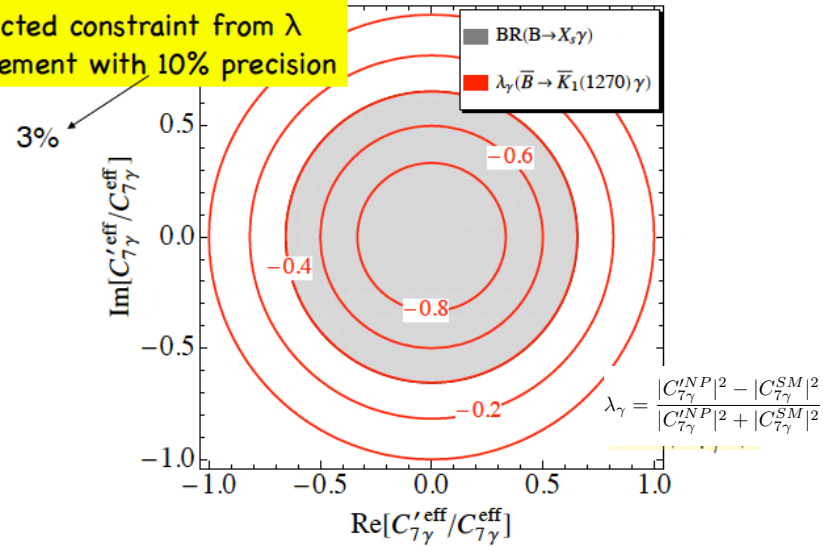
COMPLEMENTARITY

Method I
 Expected constraint from $S_{K_S\pi^0\gamma}$
 measurement with 2% precision



Current bound
 $S_{K_S\pi^0\gamma} = -0.15 \pm 0.2$

Method III
 Expected constraint from λ
 measurement with 10% precision



Discovering NP: Competitive
 Constraining NP: Complementary



Angular distribution of $B \rightarrow K_1 \gamma \rightarrow (K\pi\pi)\gamma$

$$\begin{aligned}\lambda_\gamma &\equiv \frac{|\mathcal{A}(\bar{B} \rightarrow \bar{K}_{1R}\gamma_R)|^2 - |\mathcal{A}(\bar{B} \rightarrow \bar{K}_{1L}\gamma_L)|^2}{|\mathcal{A}(\bar{B} \rightarrow \bar{K}_{1R}\gamma_R)|^2 + |\mathcal{A}(\bar{B} \rightarrow \bar{K}_{1L}\gamma_L)|^2} \\ &= \frac{|C_{7\gamma}^{NP}|^2 - |C_{7\gamma}^{SM}|^2}{|C_{7\gamma}^{NP}|^2 + |C_{7\gamma}^{SM}|^2}\end{aligned}$$

In SM, $\lambda_\gamma \simeq -1$

K₁(1270)

$$I(J^P) = \frac{1}{2}(1^+)$$

Mass $m = 1272 \pm 7$ MeV [1]

Full width $\Gamma = 90 \pm 20$ MeV [1]

K₁(1270) DECAY MODES

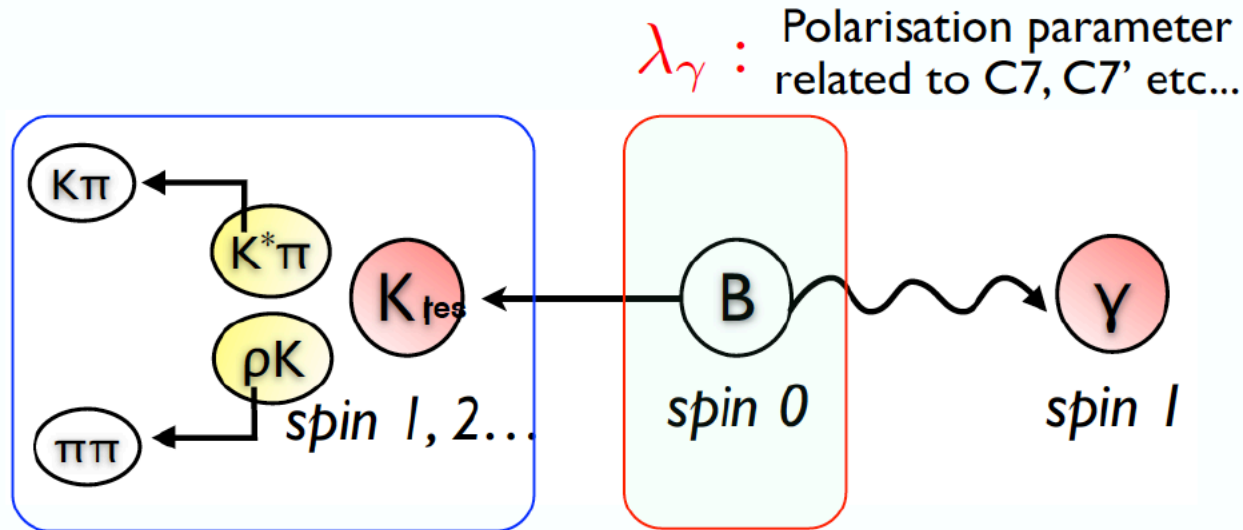
K ₁ (1270) DECAY MODES	Fraction (Γ_i/Γ)	p (MeV/c)
$K \rho$	(42 ± 6) %	46
$K_0^*(1430) \pi$	(28 ± 4) %	†
$K^*(892) \pi$	(16 ± 5) %	302
$K \omega$	(11.0 ± 2.0) %	†
$K f_0(1370)$	(3.0 ± 2.0) %	†
γK^0	seen	539



ANGULAR DISTRIBUTION FOR $b \rightarrow s\gamma$

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B meson is a spin-0 hadron:
Photon polarization is equivalent to the
polarization of Kaon resonance!

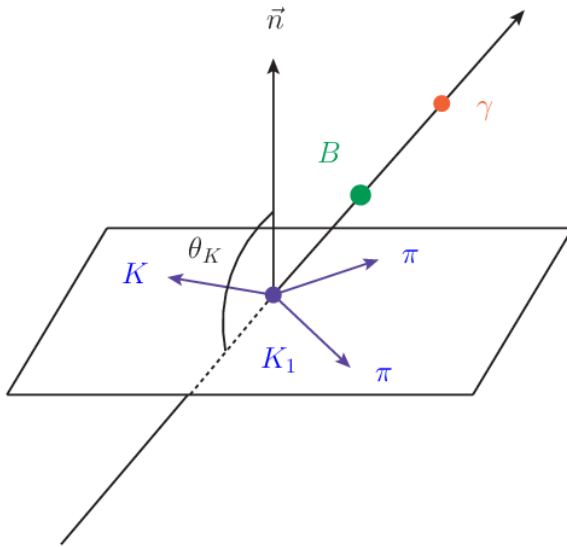


Up-down asymmetry for K_1

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Angular distribution:

Gronau, Grossman, Pirjol, Ryd PRL88('01)



- ✓ $\cos(\theta_K)$: Parity Odd
- ✓ Positive and negative helicities contribute with different signs.

$$\frac{d\Gamma_{K_1\gamma}}{d\cos\theta_K} = \frac{|A|^2|\vec{J}|^2}{4} \times \left[1 + \cos^2\theta_K + 2\lambda_\gamma \cos\theta_K \frac{\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]}{|\vec{J}|^2} \right].$$



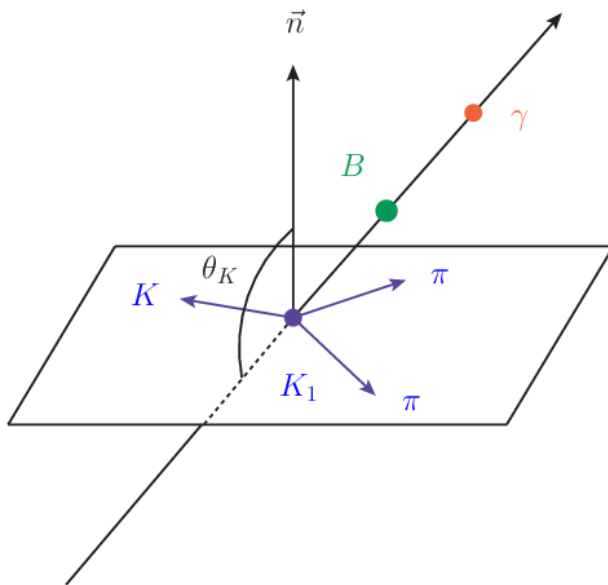
Up-down asymmetry for K_1

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Angular distribution:

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- ✓ To measure λ_γ , we need to know the decay factor $\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]/|\vec{J}|^2$
- ✓ Non-zero decay factor requires imaginary part
- ✓ Source of imaginary part: Breit-Wigner



Up-down asymmetry for K1

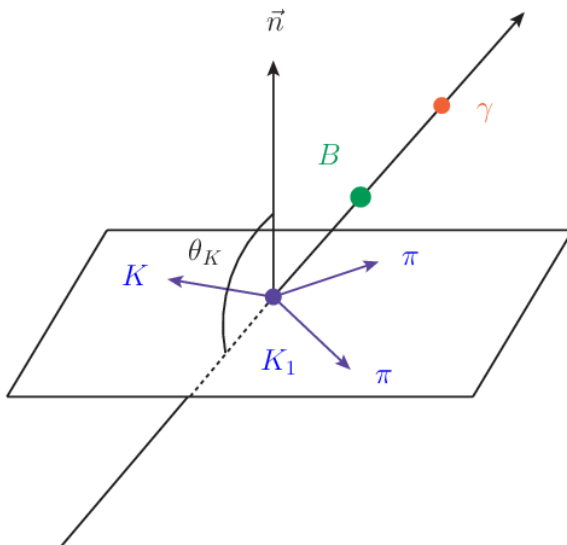
Angular distribution:

Gronau, Grossman, Pirjol, Ryd PRL88('01)

$$\frac{d\Gamma_{K_1\gamma}}{d\cos\theta_K} = \frac{|A|^2|\vec{J}|^2}{4} \times \left[1 + \cos^2\theta_K + 2\lambda_\gamma \cos\theta_K \frac{\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]}{|\vec{J}|^2} \right].$$

Up-down asymmetry for K1

$$\begin{aligned} \mathcal{A}_{\text{UD}} &\equiv \frac{\left[\int_0^1 - \int_{-1}^0 \right] d\cos\theta_K \frac{d\Gamma(B \rightarrow K_1\gamma)}{d\cos\theta_K}}{\left[\int_0^1 + \int_{-1}^0 \right] d\cos\theta_K \frac{d\Gamma(B \rightarrow K_1\gamma)}{d\cos\theta_K}} \\ &= \lambda_\gamma \frac{3}{4} \frac{\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]}{|\vec{J}|^2} \end{aligned}$$



- ✓ To measure λ_γ , we need to know the decay factor $\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]/|\vec{J}|^2$
- ✓ Non-zero decay factor requires imaginary part
- ✓ Source of imaginary part: Breit-Wigner



To extract photon helicity,
we have to reliably understand decay mechanism,
but...

GENERATOR FOR $K_{\text{RES}} \rightarrow K\pi\pi$ DECAYS

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see also M. Gronau, D. Pirjol, Phys.Rev. D96 (2017)

1. $K_{1270}(1+)$ & $K_{1400}(1+)$ decays based on quark model

A.Tayduganov, EK, Le Yaouanc PRD '13

Assume $K_1 \rightarrow K\pi\pi$ comes from quasi-two-body decay, e.g. $K_1 \rightarrow K^*\pi$, $K_1 \rightarrow \rho K$, then, \mathcal{J} function can be written in terms of:

- ▶ 4 form factors (S,D partial wave amplitudes)

2. $K^*_{1410, 1680}(1-)$ and $K^*_{21430}(2+)$

A. Kotenko, B. Knysh talk at Lausanne WS '17

Lesser parameters

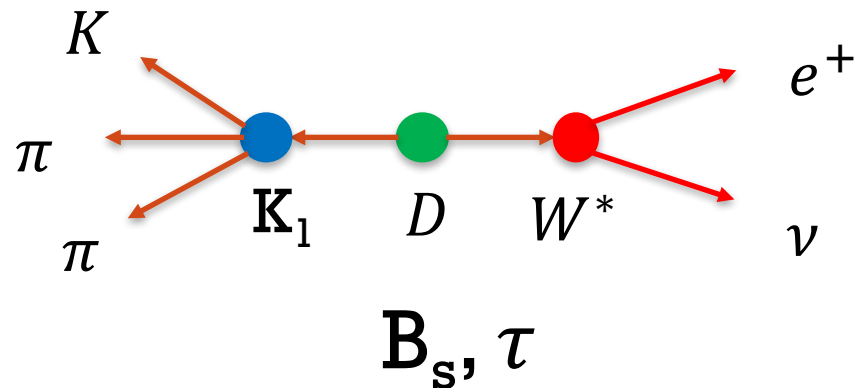
- ▶ Known to decay mainly $K_{\text{res}} \rightarrow K^*\pi$, ρK
- ▶ Only 1 form factor for each resonance

On total 10 complex couplings needed (20 real number)!

From Emi Kou



Semileptonic $D \rightarrow K\pi\pi e^+ \nu$:

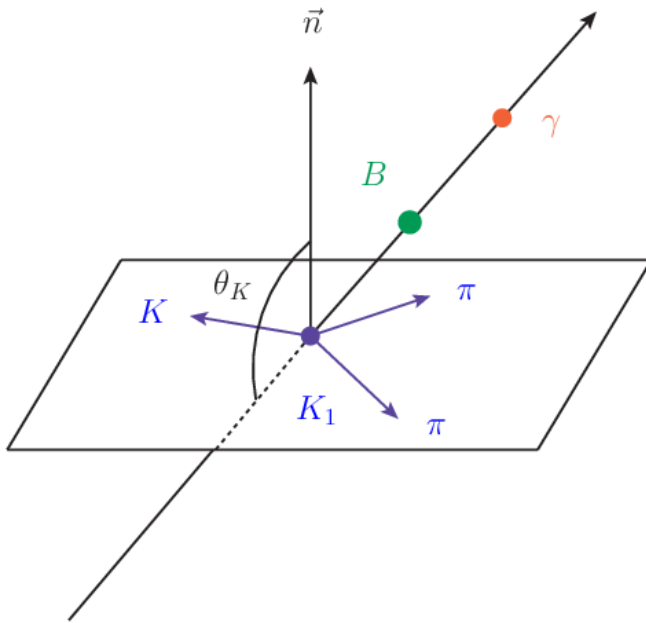


WW, F.S. Yu, Z.X.Zhao, 1909.13083

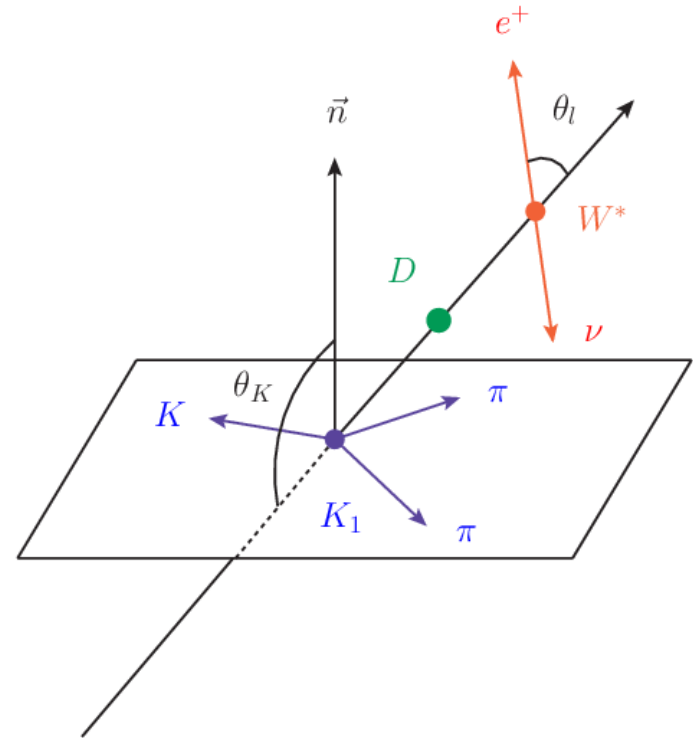


$$B \rightarrow K_1(K\pi\pi)\gamma \quad \text{VS} \quad D \rightarrow K_1(K\pi\pi)e^+\nu$$

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Polarization of γ : +, -



Polarization of W^* : +, -, 0, t

t: timelike, $\sim p_{W^*}$



$$D \rightarrow K_1(\rightarrow K\pi\pi)e^+\nu$$

Angular Distributions:

$$\begin{aligned}\frac{d\Gamma_{K_1 e \nu e}}{d\cos\theta_K d\cos\theta_l} &= d_1[1 + \cos^2\theta_K \cos^2\theta_l] \\ &+ d_2[1 + \cos^2\theta_K] \cos\theta_l \\ &+ d_3 \cos\theta_K [1 + \cos^2\theta_l] \\ &+ d_4 \cos\theta_K \cos\theta_l \\ &+ d_5[\cos^2\theta_K + \cos^2\theta_l].\end{aligned}$$

The angular coefficients are given as:

$$d_1 = \frac{1}{2}|\vec{J}|^2(4|c_0|^2 + |c_-|^2 + |c_+|^2),$$

$$d_2 = -|\vec{J}|^2(|c_-|^2 - |c_+|^2),$$

$$d_3 = -\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)](|c_-|^2 - |c_+|^2),$$

$$d_4 = 2\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)](|c_-|^2 + |c_+|^2),$$

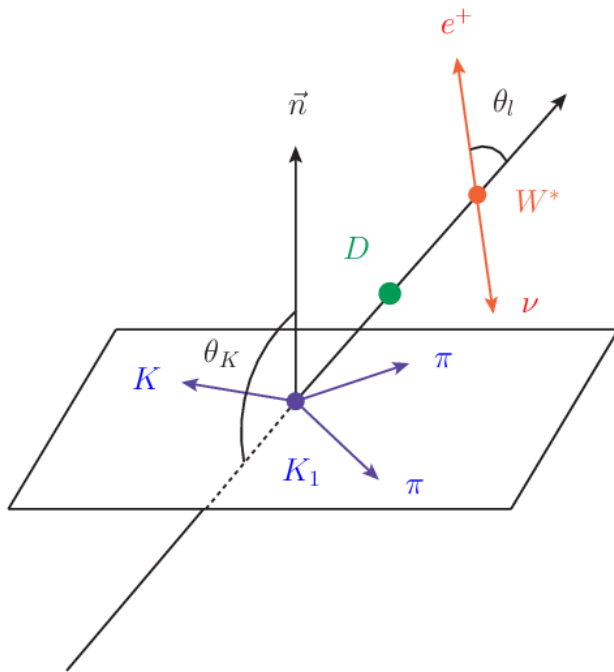
$$d_5 = -\frac{1}{2}|\vec{J}|^2(4|c_0|^2 - |c_-|^2 - |c_+|^2).$$



$D \rightarrow K_1(\rightarrow K\pi\pi)e^+\nu$: RATIO OF UP-DOWN ASYMMETRIES

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Parity Odd: $\cos \theta_K$



$$\cos \theta_K (|c_+|^2 - |c_-|^2) \text{Im}[n \cdot (\vec{J} \times \vec{J}^*)]$$

Parity Violation: $\cos \theta_l$

$$\cos \theta_l (|c_+|^2 - |c_-|^2) |\vec{J}|^2$$



$$\mathcal{A}'_{UD} \equiv \frac{\left[\int_0^1 - \int_{-1}^0 \right] d \cos \theta_K \frac{d\Gamma_{K_1 e \nu e}}{d \cos \theta_K}}{\left[\int_0^1 - \int_{-1}^0 \right] d \cos \theta_l \frac{d\Gamma_{K_1 e \nu e}}{d \cos \theta_l}}$$

$D \rightarrow K_1(\rightarrow K\pi\pi)e^+\nu$

$$\mathcal{A}'_{UD} = \frac{\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]}{|\vec{J}|^2}$$

$$\mathcal{A}_{UD} \equiv \frac{\left[\int_0^1 - \int_{-1}^0 \right] d \cos \theta_K \frac{d\hat{\Gamma}_{K_1 \gamma}}{d \cos \theta_K}}{\left[\int_0^1 + \int_{-1}^0 \right] d \cos \theta_K \frac{d\hat{\Gamma}_{K_1 \gamma}}{d \cos \theta_K}}$$

$B \rightarrow K_1(\rightarrow K\pi\pi)\gamma$

$$= \lambda_\gamma \frac{3}{4} \frac{\text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]}{|\vec{J}|^2}$$

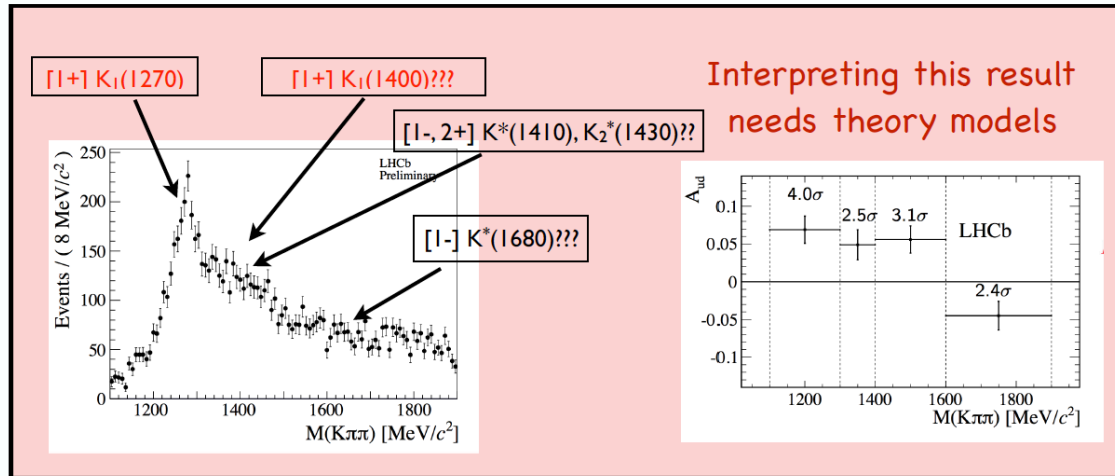


$$\lambda_\gamma = \frac{4}{3} \frac{A_{UD}}{A'_{UD}}$$



LHCb result on up-down asymmetry

LHCb PRL ('14)



Interpreting this result
needs theory models

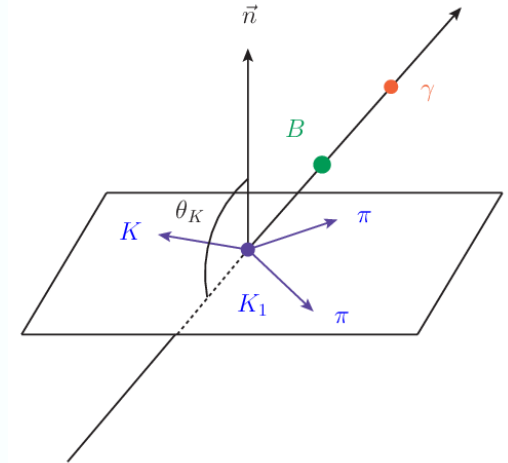


TABLE I. Legendre coefficients obtained from fits to the normalized background-subtracted $\cos \hat{\theta}$ distribution in the four $K^+ \pi^- \pi^+$ mass intervals of interest. The up-down asymmetries are obtained from Eq. (4). The quoted uncertainties contain statistical and systematic contributions. The $K^+ \pi^- \pi^+$ mass ranges are indicated in GeV/c^2 and all the parameters are expressed in units of 10^{-2} . The covariance matrices are given in Ref. [22].

	[1.1,1.3]	[1.3,1.4]	[1.4,1.6]	[1.6,1.9]
c_1	6.3 ± 1.7	5.4 ± 2.0	4.3 ± 1.9	-4.6 ± 1.8
c_2	31.6 ± 2.2	27.0 ± 2.6	43.1 ± 2.3	28.0 ± 2.3
c_3	-2.1 ± 2.6	2.0 ± 3.1	-5.2 ± 2.8	-0.6 ± 2.7
c_4	3.0 ± 3.0	6.8 ± 3.6	8.1 ± 3.1	-6.2 ± 3.2
\mathcal{A}_{ud}	6.9 ± 1.7	4.9 ± 2.0	5.6 ± 1.8	-4.5 ± 1.9

$$\begin{aligned} \mathcal{A}_{\text{UD}} &\equiv \frac{\left[\int_0^1 - \int_{-1}^0 \right] d \cos \theta_K \frac{d\Gamma(B \rightarrow K_1 \gamma)}{d \cos \theta_K}}{\left[\int_0^1 + \int_{-1}^0 \right] d \cos \theta_K \frac{d\Gamma(B \rightarrow K_1 \gamma)}{d \cos \theta_K}} \\ &= \lambda_\gamma \frac{3 \text{Im}[\vec{n} \cdot (\vec{J} \times \vec{J}^*)]}{|\vec{J}|^2}. \end{aligned}$$

[1.1-1.3]GeV:

LHCb:

PRL112.161801(2014)

$$\mathcal{A}_{UD} = (6.9 \pm 1.7) \times 10^{-2}$$

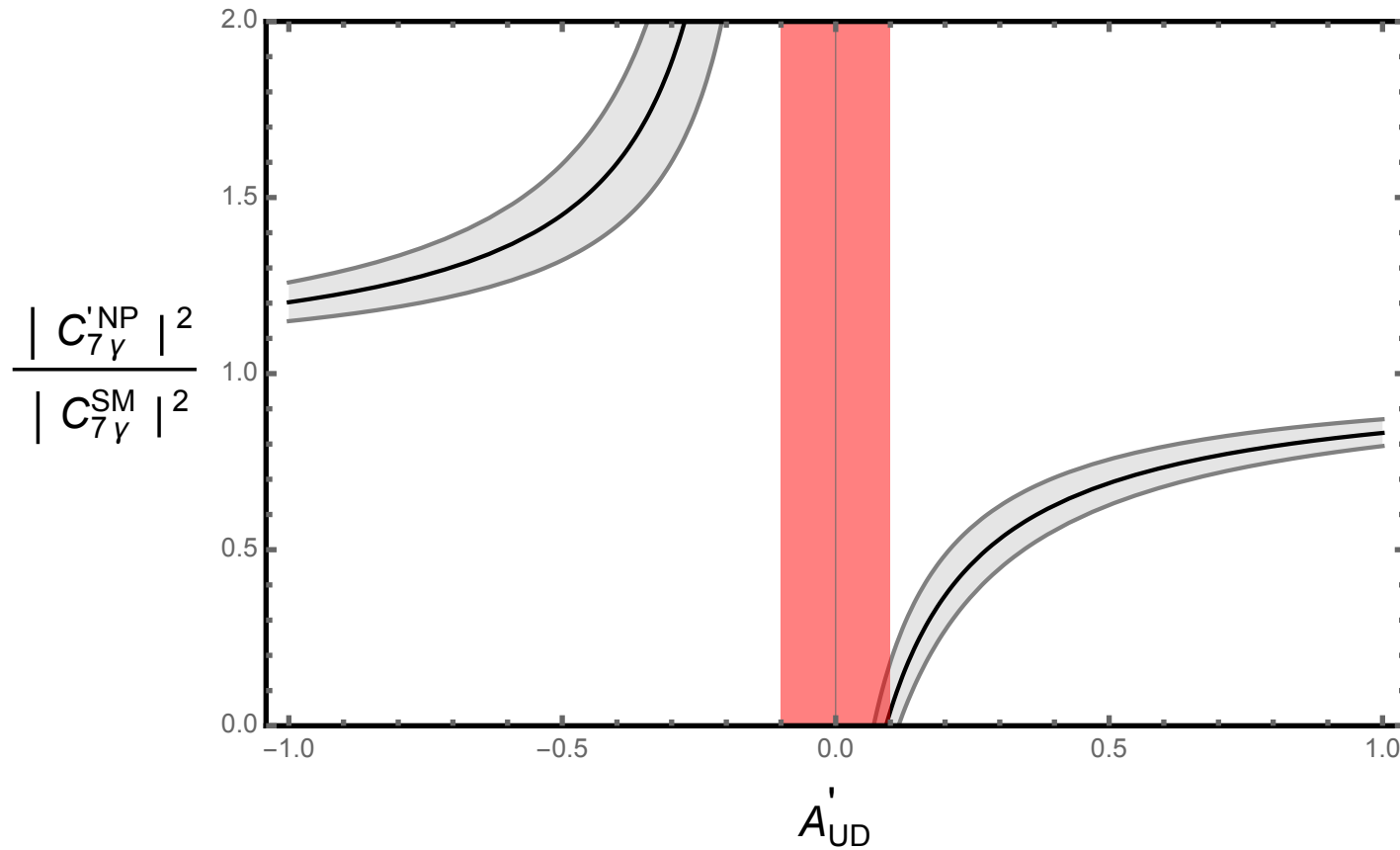
If SM



$$\mathcal{A}'_{UD} = (9.2 \pm 2.3) \times 10^{-2}$$

A significant deviation from the above value would be a clear signal for new physics beyond SM.

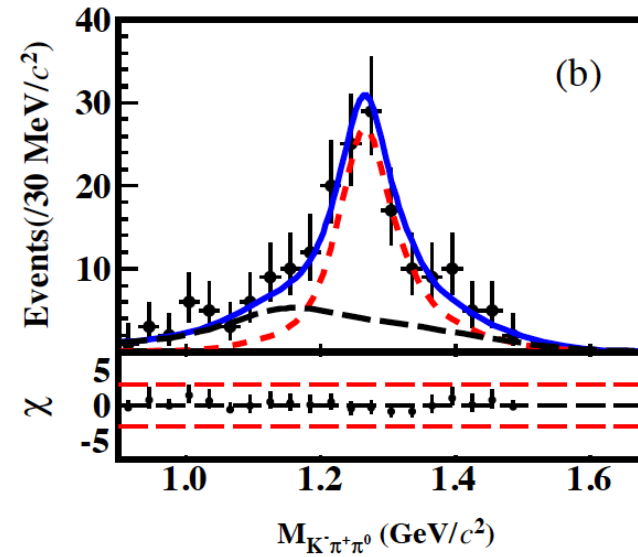
Dependence of $C_{7\gamma}^{\prime NP}$ on ratio of up-down asymmetries



$D \rightarrow K_1(\rightarrow K\pi\pi)e^+\nu$ FROM BESIII

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BESIII: 1907.11370



$$\mathcal{B}(D^+ \rightarrow \bar{K}_1^0 e^+ \nu) = (2.3 \pm 0.26 \pm 0.18 \pm 0.25) \times 10^{-3}.$$

BESIII, BelleII, LHCb, Super Tau-Charm in future?



SUMMARY

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Heavy Flavor Physics: indirect search for NP

Photon polarization in $b \rightarrow s\gamma$: unique to probe right-handed couplings

Model-independent extraction using $D \rightarrow K_1 e^+ \nu$

- ✓ Photon polarization in a model-independent way:
NP?
- ✓ BESIII, BelleII, LHCb, Super tau-charm, CEPC?

Thank you very much!



INCLUDING MORE K_J RESONANCES

The angular distribution for $D \rightarrow K_{res}(\rightarrow K\pi\pi)e^+\nu$

$$\frac{d\hat{\Gamma}}{d\cos\theta_K d\cos\theta_l} = \sum_{K_J=K_1, K_1^*, K_2, K_{12}^I} \frac{d\hat{\Gamma}_{K_J l\nu}}{d\cos\theta_K d\cos\theta_l}$$

$K^*(1410)$

$$\begin{aligned} \frac{d\hat{\Gamma}_{K_1^* l\nu}}{d\cos\theta_K d\cos\theta_l} &= (|c_+''|^2 + |c_-''|^2) \sin^2\theta_K (1 + \cos^2\theta_l) \\ &+ 2(|c_+''|^2 - |c_-''|^2) \sin^2\theta_K \cos\theta_l + 4|c_0''|^2 \cos^2\theta_K \sin^2\theta_l \end{aligned}$$



INCLUDING MORE K_J RESONANCES

$K_2^*(1430)$

$$\begin{aligned} \frac{d\hat{\Gamma}_{K_2 l \nu}}{d \cos \theta_K d \cos \theta_l} &= |c'_0|^2 \frac{3}{2} \sin^2(2\theta_K) \sin^2 \theta_l |\vec{K}|^2 \\ &+ 2|c'_1|^2 \cos^4 \frac{\theta_l}{2} \left\{ |\vec{K}|^2 (\cos^2 \theta_K + \cos^2 2\theta_K) \right. \\ &\quad \left. + 2 \cos \theta_K \cos 2\theta_K \operatorname{Im}[\vec{n} \cdot (\vec{K} \times \vec{K}^*)] \right\} \\ &+ 2|c'_{-1}|^2 \sin^4 \frac{\theta_l}{2} \left\{ |\vec{K}|^2 (\cos^2 \theta_K + \cos^2 2\theta_K) \right. \\ &\quad \left. - 2 \cos \theta_K \cos 2\theta_K \operatorname{Im}[\vec{n} \cdot (\vec{K} \times \vec{K}^*)] \right\} \end{aligned}$$

The $K_1 - K_2$ interference

$$\begin{aligned} &\frac{d\hat{\Gamma}_{K_{12}^l l \nu}}{d \cos \theta_K d \cos \theta_l} \\ &= -4\sqrt{3} \sin^2(\theta_K) \cos \theta_K \sin^2 \theta_l \operatorname{Re}[c_0 (c'_0)^* \vec{J} \cdot \vec{K}^*] \\ &\quad - 8 \cos^4 \frac{\theta_l}{2} \left\{ \frac{1}{2} (3 \cos^2 \theta_K - 1) \operatorname{Im}[c_+ (c'_+)^* \vec{n} \cdot (\vec{J} \times \vec{K}^*)] \right. \\ &\quad \left. + \cos^3 \theta_K \operatorname{Re}[c_1 (c'_1)^* * (\vec{J} \cdot \vec{K}^*)] \right\} \\ &\quad - 8 \sin^4 \frac{\theta_l}{2} \left\{ \frac{1}{2} (1 - 3 \cos^2 \theta_K) \operatorname{Im}[c_- (c'_-)^* \vec{n} \cdot (\vec{J} \times \vec{K}^*)] \right. \\ &\quad \left. + \cos^3 \theta_K \operatorname{Re}[c_{-1} (c'_{-1})^* (\vec{J} \cdot \vec{K}^*)] \right\}. \end{aligned}$$

