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## Search for the Chiral Magnetic Effect from STAR Beam Energy Scan-II data

HENPIC seminar, Sept 19, 2024

### Supported in part by



### **Zhiwan Xu** University of California, Los Angeles





## Outline

- Introduction to CME
- Previous results at STAR

- Summary

### • New Method: Event Shape Selection

### • Measurement with RHIC BES-II data



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## ~ Chirality





## **Discovery of Parity Violation in Weak Interaction**



the weak interaction.

**1957:** Nobel Prize for Yang and Lee





Foundation archive. Chen Ning Yang Prize share: 1/2

Photo from the Nobel Foundation archive.

Prize share: 1/2

### 1956: Co60 experiment (by Chien-Shiung Wu et al) discovered Parity symmetry breaking in







## **Parity Violation in Strong Interaction?**



### **QCD** vacuum is not empty





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### Image courtesy of Brookhaven National Laboratory

### http://www.physics.adelaide.edu.au/theory/staff/leinweber/VisualQCD/Nobel/

https://www.science.smith.edu/~jbrady/petrology/igrocks-diagrams/unary/H2O.php

### Phase change to Plasma



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## **Relativistic Heavy-Ion Collider**



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## Heavy-Ion Collisions at RHIC



• Participants: • At high T: quarks are liberated • Excited gluons can probe QCD vacuum





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• Spectator protons carry "+" charges • Create B field

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## The STAR detector



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## **Quark Gluon Plasma: the Small Bang**

### The Big Bang Theory



### "Baryogengesis"

### **Matter > Anti-matter**

https://www.bnl.gov/newsroom/news.php?a=11795

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### the "Small Bang" at RHIC



### "Chirogenesis"

### More LH>RH (RH>LH) in local domain



## **QCD Vacuum Topology and chirogenesis**

### **QCD** Lagrangian



$$^{2}+(\partial^{\mu}b^{a})(D_{\mu}c)^{a}$$

$$\nu - \Lambda_{QCD}^3$$



## **Strongest Magnetic Field on Earth**



Newsroom Media & Communications Office



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Phys. Rev. X 14, 011028



## **Chiral Magnetic Effect**



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Magnetic field (B) can induce charge separation (current J) for quarks at chirality imbalance ( $\mu_5$ ): CME.

Manifestly violate local **P** and **CP** symmetry.

### **3 conditions of CME**

- Chiral Symmetry Restoration
- **Topological Vacuum Transition**  $\mu_5 \neq 0$

• A strong B field The key condition





## **CME Observables in HEP**



 $\propto \mu_5 |\vec{B}|$ 

Parity Odd, can not directly observe

Parity Even, sensitive to charge separation

### **Common CME observables:** • **v**<sup>112</sup> correlator

S.A. Voloshin, Phys. Rev. C,70, 057901 (2004)

### • **R** correlator

N. N. Ajitanand *et al.*, Phys. Rev. C83, 011901(R) (2011)

### Signed balance functions

A.H. Tang, Chin. Phys. C,44, No.5 054101 (2020)

### Similar sensitivities to the CME signal and to the background. (Best Paper Award 2023)

S. Choudhury et al.(STAR), Chin. Phys. C46(2022)014101





## Early CME v<sup>112</sup> Measurements



The positively finite  $\Delta \gamma_{112}$  meets the CME expectation, but could contain contributions from: • Flow-related background  $\propto v_2$  (elliptic flow) Nonflow-related background (di-jets)

ALICE, Phys. Rev. Lett. 110(2013)012301 STAR, Phys. Rev. Lett. 113(2014)52302

 $\Delta \gamma^{OS} > \Delta \gamma^{SS}$ 

$$\Delta \gamma^{112} = \Delta \gamma^{CME} + k \frac{v_2}{N} + \Delta \gamma^{nor}$$

Signal?

**Background?** 



nflow





# **Rough Background Estimation**



- o very low beam energies: chiral symmetry breaking?
- o very high energies: no duration of the magnetic field?
- At energies in-between: AMPT could underestimate the background.

ALICE, Phys. Rev. Lett. 110(2013)012301 STAR, Phys. Rev. Lett. 113(2014)52302

Compared with a pure-background model, the CME signal seems to disappear at 7.7 GeV and 2.76 TeV.





## **Isobar Collisions**

### One approach is to look for signal difference in controlled experiment of two isobars:



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## **Isobar Results**



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(b)

0.2

n₅/s

0.15

20

**Need to investigate the BKG contribution** Should go back to large collisions system

## "Traditional" Event Shape Engineering

### • Three sub-events are used: one for POI, one for event plane, and one for event shape variable, $q_2$ , the modulus of the flow vector.



$$q_x \equiv \frac{1}{\sqrt{N}} \sum_{i}^{N} \cos(2\phi_i) \qquad \qquad q_y \equiv \frac{1}{\sqrt{N}} \sum_{i}^{N} \sin(2\phi_i)$$

- Measure  $\Delta y^{112}$  vs  $q_2$  and  $v_2$  vs  $q_2$ , then plot  $\Delta y^{112}$  vs v<sub>2</sub> to extrapolate zero-v<sub>2</sub> intercept.
- At LHC energies, all the ESE results are consistent with zero. (no duration of the magnetic field?)
- Since POI are excluded from q<sub>2</sub>, the 0 extrapolation is long and unstable.



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- At LHC energies, all the ESE results are consistent with zero. (no duration of the magnetic field?)
- Since POI are excluded from q<sub>2</sub>, the 0 extrapolation is long and unstable – works if the signal is very large, while systematic uncertainties too large for small signals

0.1

(d)



## **Schematic Diagram of Event Shape**

### Ideally, if we control eccentricity, we control flow for everything. But large event-by-event fluctuations could dominate the observable.

- participant zone geometry expected to be long range in rapidity
- emission pattern fluctuations more localized, less correlated over rapidity



H. Petersen and B. Müller, Phys. Rev. C 88, 044918

Event shape variables based on particles of interest (**POI**) are sensitive to both geometry and emission pattern.

**CME** background comes from combined eccentricity and emission patterns





## **Event Shape Selection and v<sub>2</sub> Control**

**(a)** 

**(b)** 

### **Event shape variable**

single q<sup>2</sup> (POI)

pair q<sup>2</sup> (POI)

 $q_2^2 = \frac{1}{N} \left[ \left( \sum_{i=1}^N \sin 2\varphi_i \right)^2 + \left( \sum_{i=1}^N \cos 2\varphi_i \right)^2 \right]$  $= 1 + \frac{1}{N} \sum_{i \neq j} \cos[2(\varphi_i - \varphi_j)],$  $\langle q_2^2 \rangle \approx 1 + N v_2^2 \{2\}$ 

$$q_2^2 = \frac{\left(\sum_{i=1}^N \sin 2\varphi_i\right)^2 + \left(\sum_{i=1}^N \cos 2\varphi_i\right)^2}{N(1 + N\langle v_2 \rangle^2)}$$

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**Elliptic flow variable** 

(C) single v<sub>2</sub> (POI) pair v<sub>2</sub> (POI) <u>~(d)</u>

- ESS recipes (a) and (b) involve direct event-by-event correlations between  $q_2^2$  and  $v_2$ , which will cause under-subtraction of background.
- We should use "mixed" recipes, (c) or (d).
- Redefine  $q_2^2$  with an extra normalization.
- Pair  $q_2^2$  and pair  $v_2$  are based on  $\varphi_p$ .



- $v_2\{\pi\} = 0; v_2\{\rho\} > 0$   $\Delta \gamma^{112} = 1$  (pure-BKG)
- $q^2_{POI} = 0$  $0 q^{2}_{PPOI} = 1$



### Simulation



- In AVFD, the optimal ESS recipe (c) accurately matches the input true CME signal.

• Mixed combinations further suppress residual BKG: intercepts follow an ordering (a)>(b)>(c)>(d) • With AMPT, all ESS schemes seem to over-estimate the BKG (same ordering as AVFD).









### **ESS procedure**

### • A novel method to control emission pattern: utilize event shapes of POI kinematic region

**1. Categorize events** Z. Xu et al, PLB 848(2024)138367 Flow vector with higher-order normalization  $q_{2}^{2} = \frac{\left(\sum_{i=1}^{N} \sin 2\varphi_{i}\right)^{2} + \left(\sum_{i=1}^{N} \cos 2\varphi_{i}\right)^{2}}{N(1 + N\langle v_{2} \rangle)}$ 

**2. Measure the Δγ Observable & v<sub>2</sub> flow**. Optimal Solution pair q<sub>2</sub> (PPOI) ...... single v<sub>2</sub> (POI)

• adding momenta of two POI particles

~ mimic resonance decay.

**3.** Plot  $\Delta \gamma$  against v<sub>2</sub> to extrapolate  $\Delta \gamma_{ESS}^{112}$ 

$$\Delta \gamma_{ESS}^{112} = Intercept \times (1 - v_2)^2$$

Non-interdependent Flow, Z.Xu et al Phys. Rev. C 107, L061902



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## Beam Energy Scan at RHIC

![](_page_23_Figure_1.jpeg)

### **Beam Energy Scan**

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![](_page_23_Picture_5.jpeg)

![](_page_23_Picture_8.jpeg)

# Beam Energy Scan at RHIC

	BES-I	
$\sqrt{S_{NN}}$ (GeV)	Events (10 <sup>6</sup> )	Year
62.4	46	2010
39	86	2010
27	30	2011
19.6	15	2011
14.6	13	2014
11.5	7	2010
9.2	0.3	2008
7.7	4	2010

![](_page_24_Figure_2.jpeg)

![](_page_24_Picture_4.jpeg)

BES-II			
$\sqrt{S_{NN}}$ (GeV)	Events (10 <sup>6</sup> )	Yea	
27	555	201	
19.6	478	201	
14.6	324	201	
11.5	230	202	
9.2	160	202	
7.7	101	202	

Baryon Chemical Potential -  $\mu_{\rm B}$ (MeV)

### **BES-II Statistics:** • 10-20 times higher.

### **Detector Upgrades:**

• 2018 EPD : high EP resolution into spectator region (2.1<n<5.1)  $\eta > y_{\text{beam}}$ : Forward spectators

![](_page_24_Figure_12.jpeg)

![](_page_24_Figure_15.jpeg)

![](_page_24_Picture_16.jpeg)

## **The Event Plane Detector at STAR**

![](_page_25_Figure_1.jpeg)

- Higher resolution, BES-II new detector (EPD)
- The inner EPD detects first-order spectator plane
  - Targeting the spectator regions for B field

![](_page_25_Picture_9.jpeg)

![](_page_25_Picture_10.jpeg)

![](_page_25_Picture_11.jpeg)

## ESS applied to Au+Au at 19.6 GeV

![](_page_26_Figure_1.jpeg)

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![](_page_26_Figure_3.jpeg)

**Event Shape Selection** Spectator  $\Psi_1$  $\Delta \gamma^{112} = \Delta \gamma^{\text{CME}} + \frac{k^2}{k} + \Delta \gamma^{\text{nonflow}}$ Signal Background

- ESS using POI allows much shorter extrapolation to zero v<sub>2</sub>.
- The ordering of y-intercepts follows predictions from both AVFD and AMPT
- The y-intercept requires a small correction to restore the unbiased CME signal:

$$\Delta \gamma_{ESS}^{112} = Intercept \times (1 - v_2)^2$$

Z.Xu et al Phys. Rev. C 107, L061902

![](_page_26_Picture_12.jpeg)

![](_page_26_Picture_13.jpeg)

![](_page_26_Picture_14.jpeg)

![](_page_26_Picture_15.jpeg)

### Au+Au at 19.6 GeV

![](_page_27_Figure_1.jpeg)

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![](_page_27_Figure_3.jpeg)

- The ESS is applied to different centralities. 0
- The ordering of four intercept  $\Delta \gamma_{ESS}^{112}$  follows prediction from both AMPT and AVFD model.

### Au+Au at 19.6 GeV

![](_page_28_Figure_1.jpeg)

![](_page_28_Picture_10.jpeg)

## **Beam Energy Scan II - Event Shape Selection**

![](_page_29_Figure_1.jpeg)

- $\Delta \gamma_{ESS}^{112}$  from the optimal ESS (c), pair q2 and single v2:
  - At 200 GeV, using ZDC-SMD planes, no signal is observed.

  - At 9.2 and 7.7 GeV, data favor the zero-CME scenario.
- $\Delta \gamma_{ESS}^{132}$  is consistent with zero.

• At 19.6, 14.6 and 11.5 GeV, a finite  $\Delta \gamma_{ESS}^{112}$  (3 $\sigma$  significance) in the 20-50% centrality.

## **Beam Energy Dependence of CME observable**

![](_page_30_Figure_1.jpeg)

- **BKG-indicator**  $\Delta \gamma_{ESS}^{132}$  **consistent with zero**
- ° At least 80% of  $\langle \Delta \gamma^{112} \rangle$  is from the background.
- At 200 GeV, ratio is (-2 ± 5.1 ± 1.6)%
  - upper limit of fCME~10% in Au+Au
  - upper limit of fCME~ 5% in isobars using participant planes: 0.7% difference, too small to detect
- Combine three points at 19.6, 14.6 and 11.5 GeV, the literal average of the ESS results reaches an over 5σ significance (assuming similar physics conditions between 10 and 20 GeV).
- The ESS results approach zero around 9.2 and 7.7 GeV.

![](_page_30_Figure_13.jpeg)

![](_page_30_Figure_14.jpeg)

![](_page_30_Picture_15.jpeg)

## **Connection from ESS to H-correlator**

![](_page_31_Figure_1.jpeg)

• In the BES-I data, the H correlator is introduced to subtract the flow BKG:

 $\delta = B + H$ 

$$H(\kappa_{bg}) \equiv (\kappa_{bg} v_2 \delta - \gamma^{112}) / (1 + \kappa_{bg} v_2)$$
$$\Delta \bar{H} \equiv H_{SS} - H_{OS}$$
$$\gamma = \kappa V_2 B - H \qquad \delta = \cos(\phi_1 - \phi_2)$$

Bzdak, V. Koch and J. Liao, Lect. Notes Phys. 871, 503 (2013)

- $\circ \kappa_{bq}$  is an adjustable parameter, unknown priori. It quantifies the coupling between elliptic flow and other mechanisms manifested in the twoparticle correlation.
- With  $\kappa_{bg}$  set to 2.5,  $\Delta H$  agrees with the ESS result at all beam energies under study.
- The flow background can be reasonably well described by a universal coupling between  $v_2$ and the two-particle correlation.

## A sweet-zone: 10-30 GeV?

![](_page_32_Figure_1.jpeg)

 $\circ$  dv<sub>1</sub>/dy is sensitive to the EOS: "softest point collapse" of flow.

- Near critical point region, the topological fluctuation will be enhanced.
- STAR net-proton cumulant measured a significant deviation from model.
- At large B, the chiral symmetry breaking (split from deconfiment)?

![](_page_32_Figure_9.jpeg)

A. J. Mizher, M. N. Chernodub, and E. S. Fraga, PRD 82 (2010) 105016

![](_page_32_Figure_13.jpeg)

![](_page_32_Picture_14.jpeg)

![](_page_32_Picture_15.jpeg)

## Summary

- The search for the CME in heavy-ion collision probes the intrinsic properties of QCD.
- STAR latest CME searches use the novel Event Shape Selection to effectively suppress flowrelated backgrounds.
  - At 200 GeV, upper limit of  $f_{CME} \sim 10\%$ .
  - ° At each of 11.5, 14.6 and 19.6 GeV, a positively finite  $\Delta \gamma_{ESS}^{112}$  (>3 $\sigma$ ). Over 5 $\sigma$  if combined. Around 7.7 GeV, approaches zero CME with large uncertainties.
- More theoretical insights are needed:
  - The remaining B field may be too weak at 200 GeV?
  - Chiral symmetry breaking/QGP disappering around 7.7 GeV?
  - The chance of the CME occurrence is enhanced near the critical point?

![](_page_33_Picture_13.jpeg)

# Thank you.

### Acknowledgement: Thank all the collaborators at STAR, especially folks from the CME-focus group.

![](_page_34_Picture_2.jpeg)