

# Light nuclei production in ultrarelativistic heavy-ion collisions

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April 8, 2020



## Light nuclei in heavy ion collisions



Deuteron (d)



Tritium (t)



Helium-3 ( ${}^3\text{He}$ )



Hypertriton ( ${}^3_{\Lambda}\text{H}$ )

## Anti-



Deuteron ( $\bar{d}$ )



Tritium ( $\bar{t}$ )



Helium-3 ( ${}^3\bar{\text{He}}$ )



Hypertriton ( ${}^3_{\Lambda}\bar{\text{H}}$ )

These and other nuclei are created in heavy ion collisions

# Anti-helium by Alpha-Magnetic Spectrometer



- Few events (compatible with)  ${}^3\overline{\text{He}}$ ,  ${}^4\overline{\text{He}}$   
*Caveats: hard measurement, 1 event/year, not published*
- Where do they come from?  
*Antimatter clouds? Dark matter annihilations? pp collisions?*

# Understanding anti-helium measurement by AMS

- K. Blum, K. C. Y. Ng, R. Sato and M. Takimoto,  
"Cosmic rays, antihelium, and an old navy spotlight," PRD 96, no. 10, 103021 (2017)

Conclusion:  $\overline{\text{He}}$  production compatible with  $pp$

Use coalescence model for  $pp \rightarrow \overline{\text{He}} + X$

- V. Poulin, P. Salati, I. Cholis, M. Kamionkowski and J. Silk,  
"Where do the AMS-02 antihelium events come from?", PRD 99, no. 2, 023016 (2019)

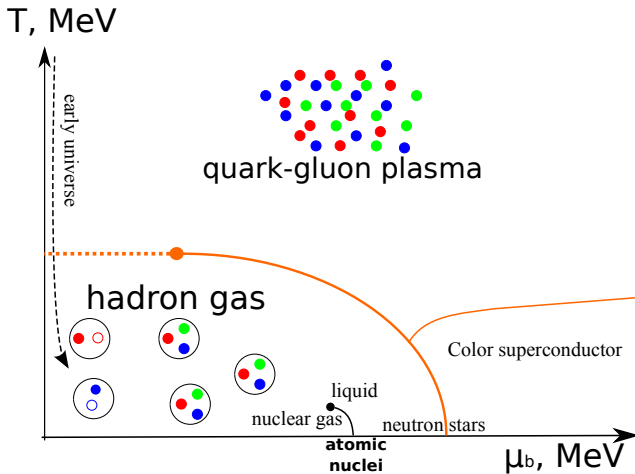
Conclusion:  $pp$  cannot produce that much  $\overline{\text{He}}$

advocate presence of anti-clouds in our Galaxy

Use coalescence model for  $pp \rightarrow \overline{\text{He}} + X$

Both use  $pp$  collisions data from ALICE to calibrate models  
Extrapolation from  $pp \rightarrow \bar{d}$  to  $pp \rightarrow \overline{\text{He}} + X$ ,  $pA \rightarrow \overline{\text{He}} + X$ ,  
 $AA \rightarrow \overline{\text{He}} + X$ , from high to low energies, from midrapidity to  
forward rapidity involved

# Light nuclei and critical fluctuations



Generic critical point feature: **spatial** fluctuations increase

# Nucleon density fluctuations in coordinate space

Kaijia Sun et al., Phys. Lett. B 774, 103 (2017)

Kaijia Sun et al., Phys. Lett. B 781 (2018) 499-504

Proton and neutron density:

$$\rho_n(x) = \langle \rho_n \rangle + \delta \rho_n(x)$$

$$\rho_p(x) = \langle \rho_p \rangle + \delta \rho_p(x)$$

Correlations and fluctuations:

$$C_{np} \equiv \langle \delta \rho_n(x) \delta \rho_p(x) \rangle / (\langle \rho_n \rangle \langle \rho_p \rangle)$$

$$\Delta \rho_n \equiv \langle \delta \rho_n(x)^2 \rangle / \langle \rho_n^2 \rangle$$

From a simple coalescence model

$$N_d \approx \frac{3}{2^{1/2}} \left( \frac{2\pi}{mT} \right)^{3/2} \int d^3x \rho_p(x) \rho_n(x) \sim \langle \rho_n \rangle N_p (1 + C_{np})$$

$$N_t \approx \frac{3^{1/2}}{4} \left( \frac{2\pi}{mT} \right)^3 \int d^3x \rho_p(x) \rho_n^2(x) \sim \langle \rho_n \rangle^2 N_p (1 + 2C_{np} + \Delta \rho_n)$$

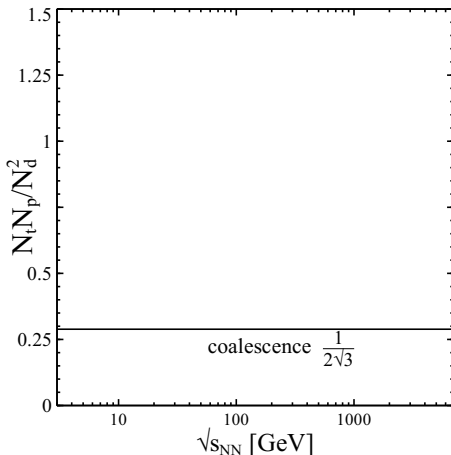
$$\frac{N_t N_p}{N_d^2} = \frac{1}{2\sqrt{3}} \frac{1 + 2C_{np} + \Delta \rho_n}{(1 + C_{np})^2}$$

$$\text{Thermal ratio } \frac{g_t g_p}{g_d^2} \left( \frac{3m \cdot m}{(2m)^2} \right)^{3/2} = \frac{1}{2\sqrt{3}} \quad \text{Fluctuations and correlations}$$

Light nuclei are sensitive to spatial density fluctuations

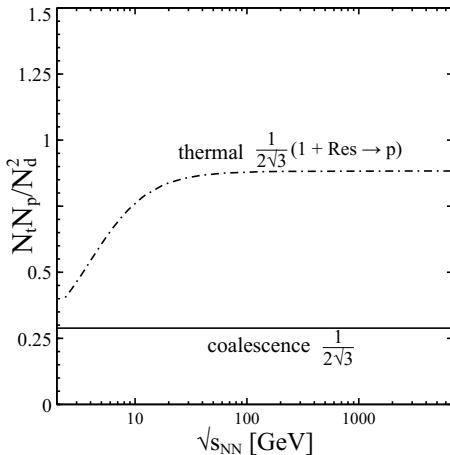
# Comparing the $p$ - $d$ - $t$ ratio to NA49, STAR, and ALICE data

Data: NA49 [Anticic:2010mp,Blume:2007kw,Anticic:2016ckv], STAR [Adam:2019wnb,Zhang:2019wun], ALICE [Adam:2015vda]; model JAM + coalescence [Liu:2019nii]; see DO Quark Matter 2019 proceedings



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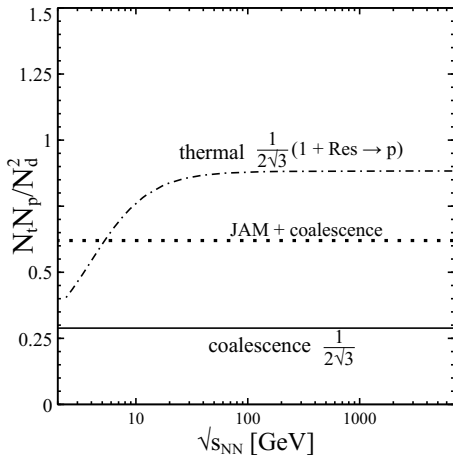


Thermal: first nuclei form, then resonance decays.  
Coalescence: first resonances decay, then nuclei form.



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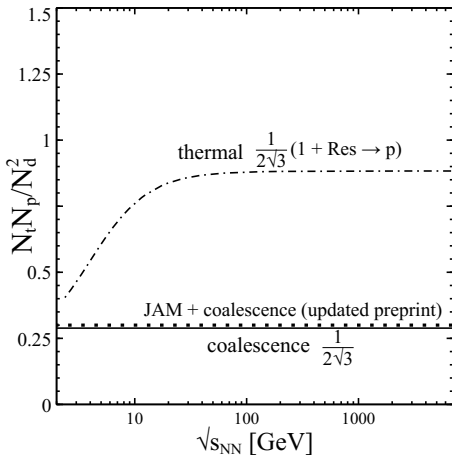
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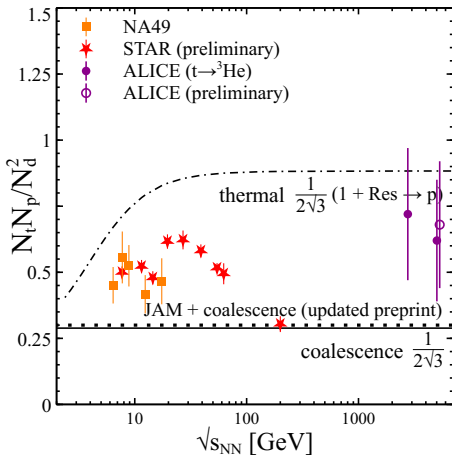
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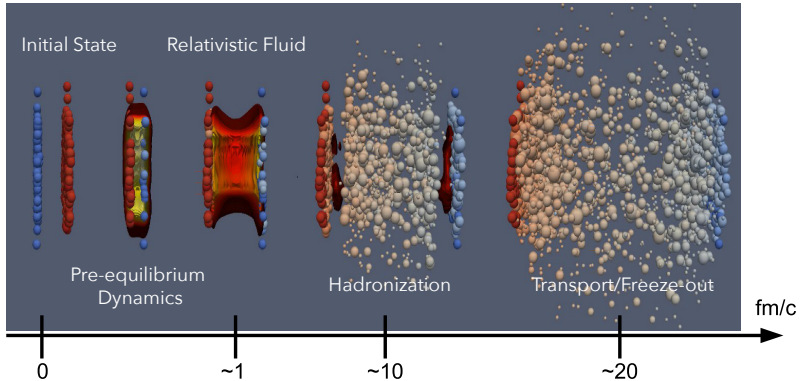


Are the bumps related to fluctuations? Can one generate them without critical point?

## Trying different model, different coalescence

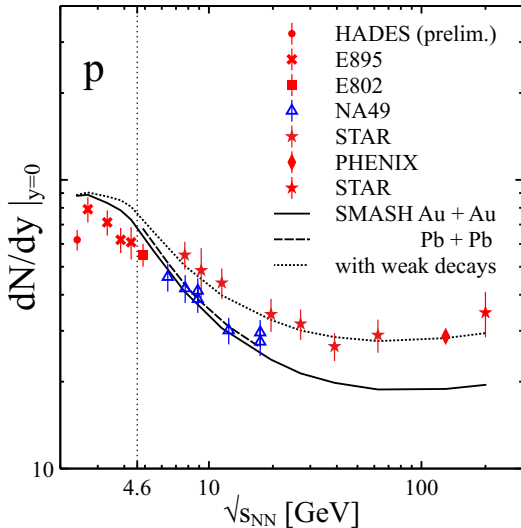
- JAM + coalescence
  - Fixed time coalescence,  $t_{coal} = 50 \text{ fm}/c$
  - $\Delta R_d = 4 \text{ fm}$ ,  $\Delta R_t = 3.4 \text{ fm}$ ,  $\Delta p_d = \Delta p_t = 0.3 \text{ GeV}$
  - Tuned to reproduce d, t  $dN/dy$  at 200 GeV
  - How well does JAM reproduce proton spectra?
- SMASH + coalescence /  
MUSIC hydro + SMASH + coalescence
  - Coalescence at the latest of the last interaction times
  - $pn \rightarrow d$ ,  $dn \rightarrow t$
  - $\Delta p_d = \Delta p_t = 0.42 \text{ GeV}$ ,  $\Delta R_{d,t} = h/\Delta p_d$
  - Rejection with isospin factors 3/8 (d) and 1/4 (t)
  - Tune to 7.7 GeV: STAR and NA49 agree here
  - Control, how well proton spectra are reproduced (still not perfect)

# Hybrid approach

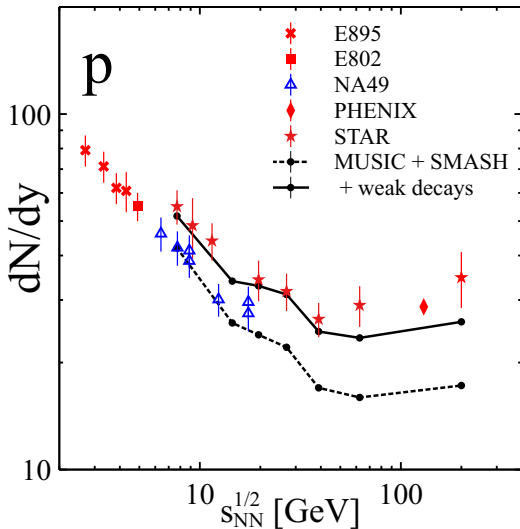


- MUSIC hydro [B. Schenke, S. Jeon, C. Gale, PRC 82, 014903 \(2010\)](#)  
viscous 3+1D hydrodynamics with  $j_B$ , EoS with  $\mu_B, \mu_S, \mu_Q$   
particlization at constant energy density
- SMASH hadronic afterburner [J. Weil et al., PRC 94, no. 5, 054905 \(2016\)](#)

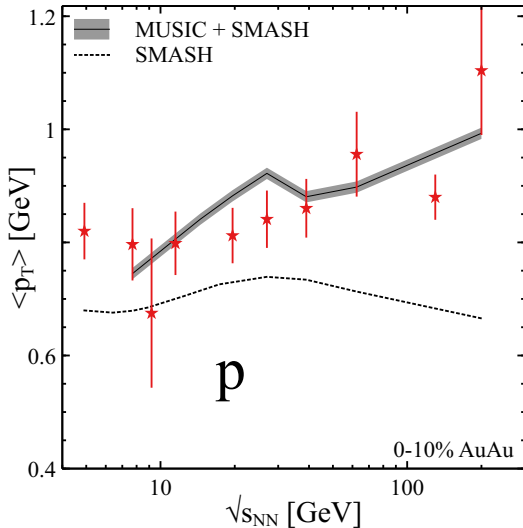
# SMASH or MUSIC + SMASH and proton spectra



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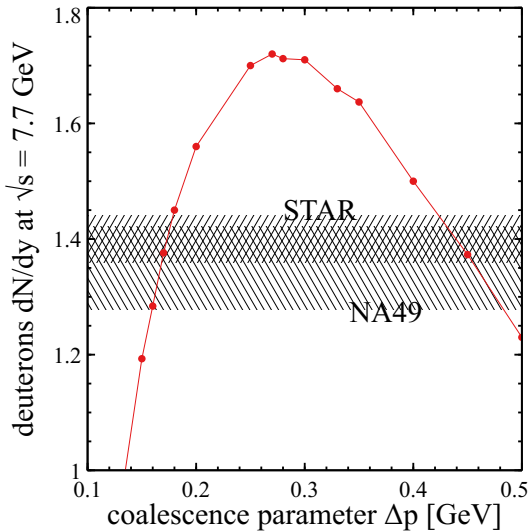


SMASH: underestimates proton mean  $p_T$

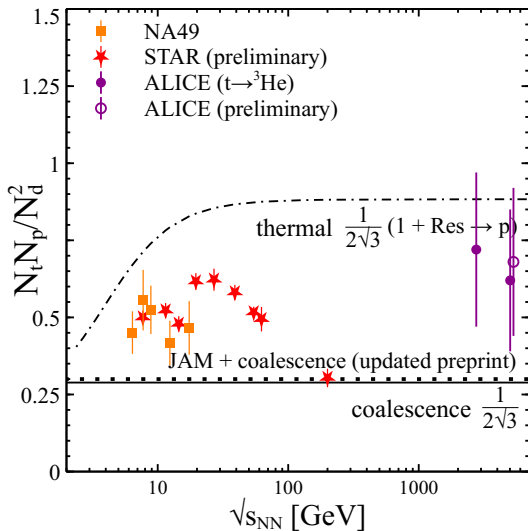
MUSIC + SMASH: reasonable description of proton spectra



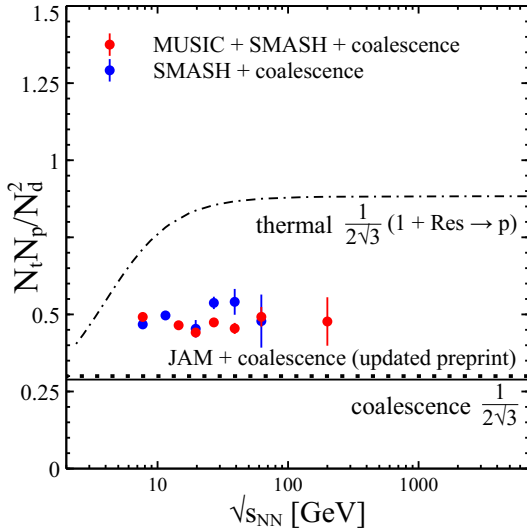
## Tuning the coalescence (there is only one parameter!)



# Coalescence and p-t-d double ratio

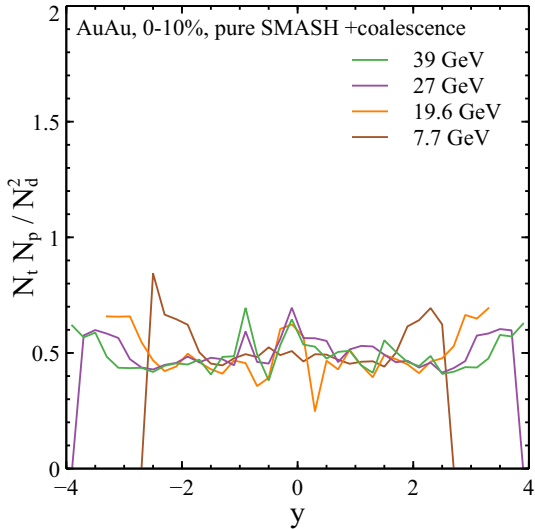


## Coalescence and p-t-d double ratio



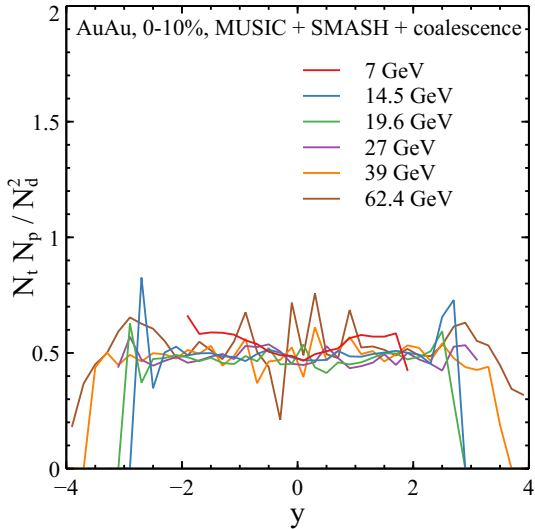
The ratio is flat as a function of energy

## Coalescence and p-t-d double ratio



The ratio is flat as a function of rapidity

## Coalescence and p-t-d double ratio

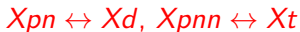


The ratio is flat as a function of rapidity

## Coalescence: conclusions and outlook

- The ratio  $N_t N_p / N_d^2(\sqrt{s}, y)$  is flat regardless of
  - The underlying model: JAM, SMASH, MUSIC + SMASH
  - How well the underlying model describes protons
  - What the specific details of coalescence procedure are
- $N_t N_p / N_d^2$  depends solely on coalescence parameters
- Can we get a non-flat ratio as in experiment?
  - A) Introduce additional fluctuations of nucleon spatial density
    - Nuclear mean-field potentials
    - Critical fluctuations
  - B) Try alternative model of nuclei production

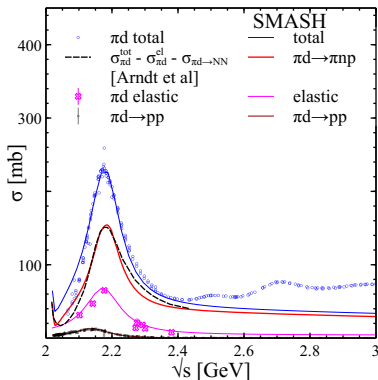
Light nuclei production via explicit catalysis reactions



# Light nuclei production by pion catalysis

- $\pi d \leftrightarrow \pi np$ ,  $\pi t \leftrightarrow \pi nnp$ ,  $\pi^3\text{He} \leftrightarrow \pi npp$
- Disintegration cross section fit to data
- Reverse rates fixed by detailed balance relations
- Put protons and pions in the box, wait until it thermalizes,  
 $N_t N_p / N_d^2 \approx 0.27$

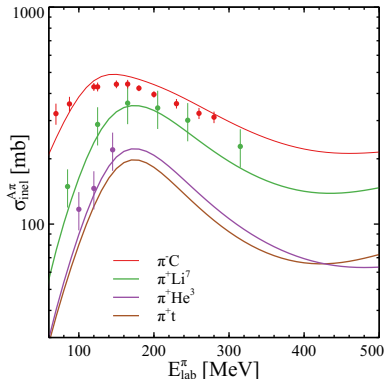
DO, Pang, Elfner, Koch, PRC99 (2019) no.4, 044907



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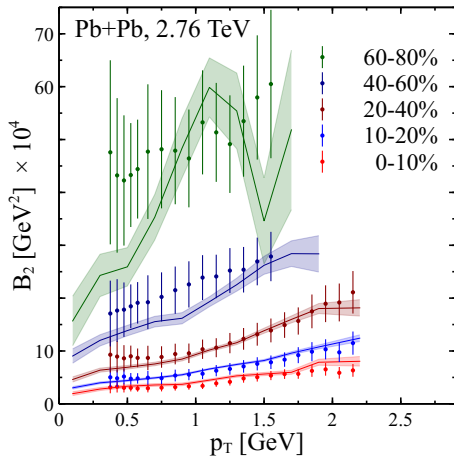
Data: Mihul:1992br, Angelescu:1996ev, Ashery:1981tq, Binon:1970ye





# LHC, deuteron: $B_2(p_T)$ for different centralities

DO, Pang, Elfner, Koch, PRC99 (2019) no.4, 044907

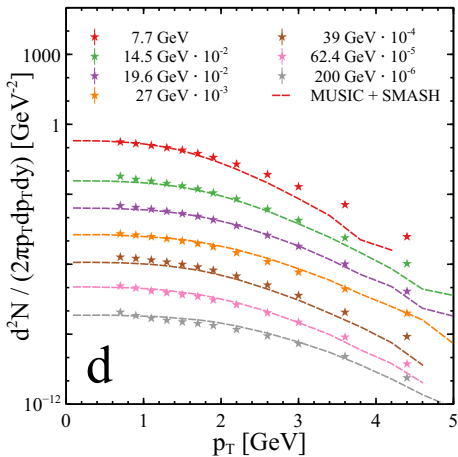


$$B_2(p_T) = \frac{\frac{1}{2\pi} \frac{d^3 N_d}{p_T dp_T dy} \Big|_{p_T^d = 2p_T^p}}{\left( \frac{1}{2\pi} \frac{d^3 N_p}{p_T dp_T dy} \right)^2}$$

No free parameters. Works well for all centralities.

# Deuterons @ STAR by MUSIC + SMASH with catalysis reactions

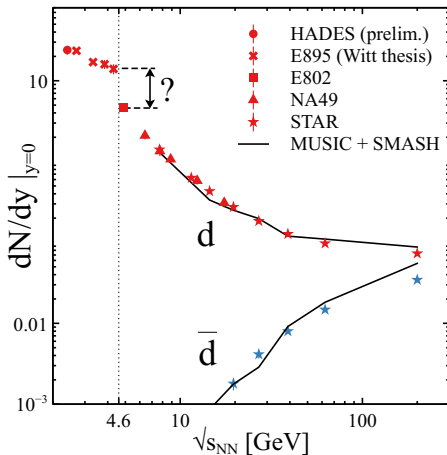
Data: Alt:2006dk, Anticic:2010mp, Adams:2003xp, Adamczyk:2017iwn, Abelev:2009bw, Adcox:2003nr, Klay:2001tf, Ahle:1999in, Adam:2019wnb, Zhang:2019wun



MUSIC + SMASH with explicit catalysis reactions describes deuterons rather well (as well as protons)

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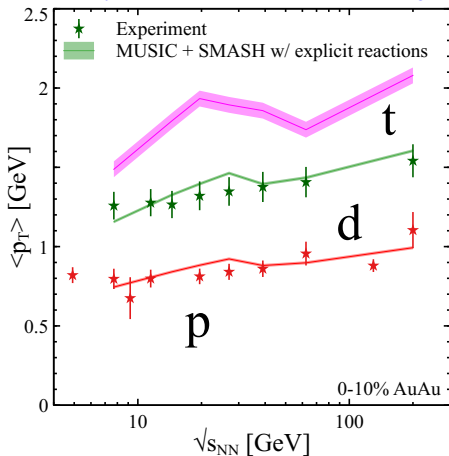
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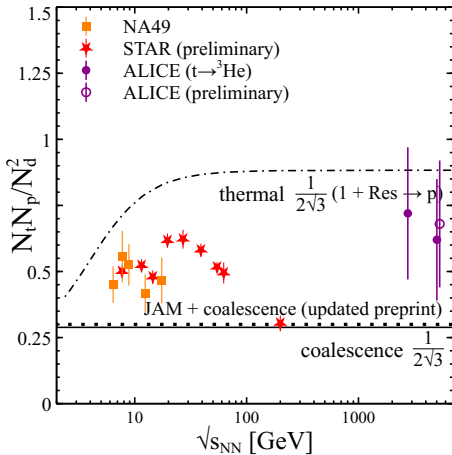
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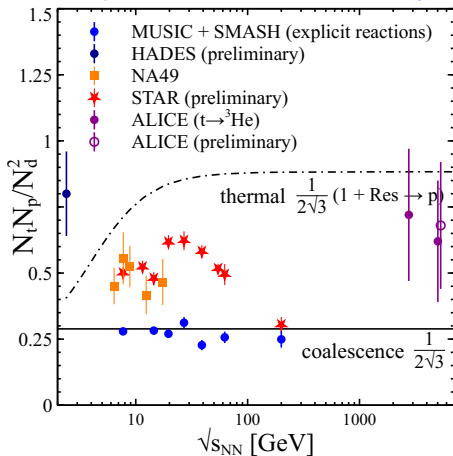
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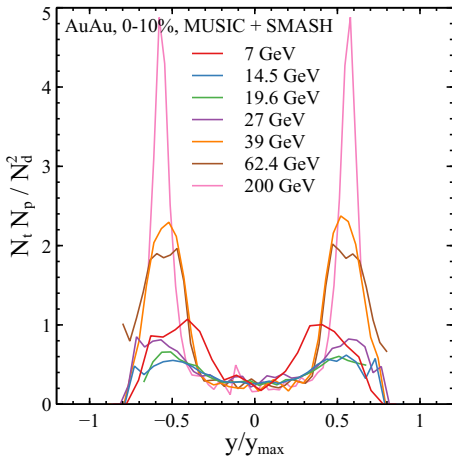
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Do catalysis reactions act as coalescence?

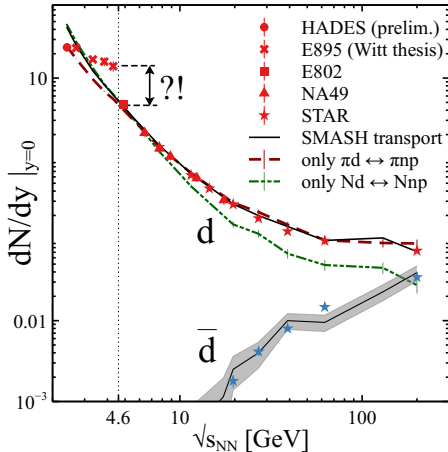
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Do catalysis reactions act as coalescence? Not necessarily!

# Deuterons @ STAR by pure SMASH with catalysis reactions



Pion catalysis dominates at STAR energies

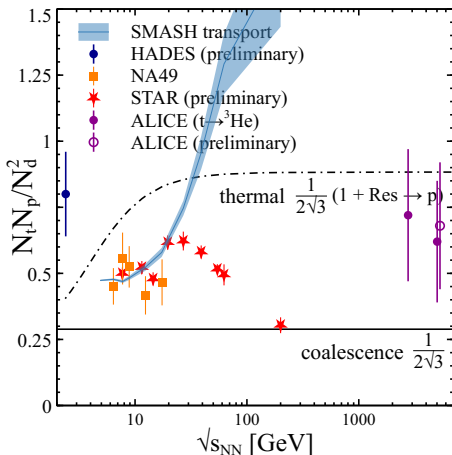
Note the (unpublished) jump around 4.6 GeV,  
no jump there for protons



## Deuterons @ STAR by pure SMASH with catalysis reactions

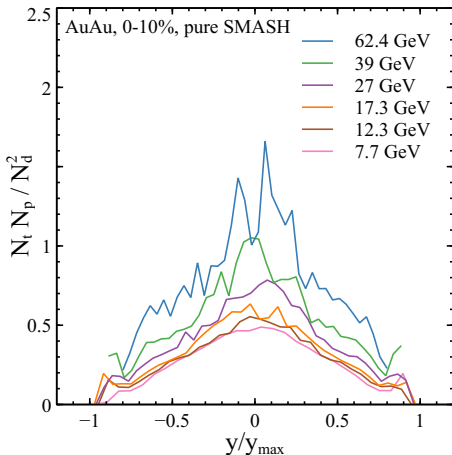
- Pion catalysis dominates at STAR energies
- Deuteron and antideuteron  $dN/dy$  is described very well, because proton  $dN/dy$  is described well too
- **However**, deuteron  $\langle p_T \rangle$  is underestimated, as for protons  $\implies$  need MUSIC + SMASH
- What about  $N_p N_t / N_d^2$ ?

# $N_p N_t / N_d^2$ @ STAR by pure SMASH with catalysis reactions



- Not flat as function of  $\sqrt{s}$  and  $y$
- Triton way above the data at  $\sqrt{s} \geq 20$  GeV
- Same reactions as before!  $\implies$  underlying model matters

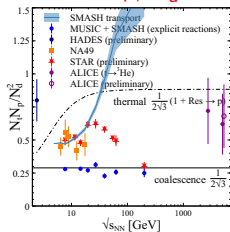
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## Summary and outlook

- The ratio  $N_t N_p / N_d^2$  may be related to the critical point
- In coalescence models  $N_t N_p / N_d^2$  is always flat
- Alternative to coalescence: explicit catalysis reactions
  - Describe deuterons as good as protons in the underlying model
  - $N_t N_p / N_d^2$  may be flat (MUSIC + SMASH), but not necessarily (SMASH)
  - Can we use this knowledge to construct a model, which reproduces experimental  $N_t N_p / N_d^2$  without critical point?

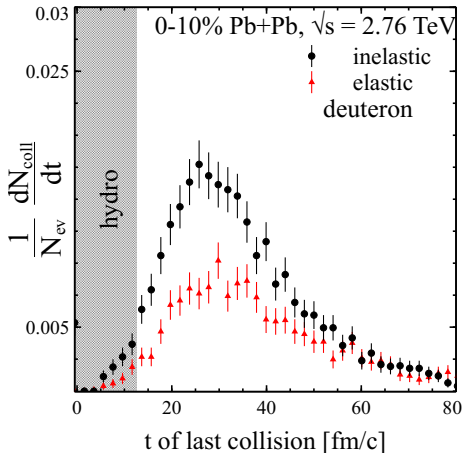


Outlook: understand the difference MUSIC + SMASH vs SMASH, try model with a critical point

Backup

## Does deuteron freeze out at 155 MeV?

Only less than 1% of final deuterons originate from hydrodynamics

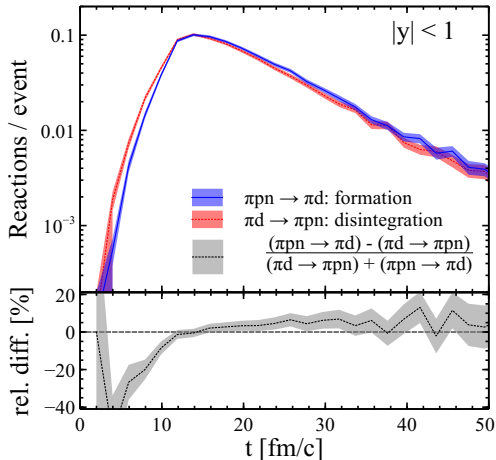


Deuteron freezes out at late time

Its chemical and kinetic freeze-outs roughly coincide

# Is $\pi d \leftrightarrow \pi np$ reaction equilibrated

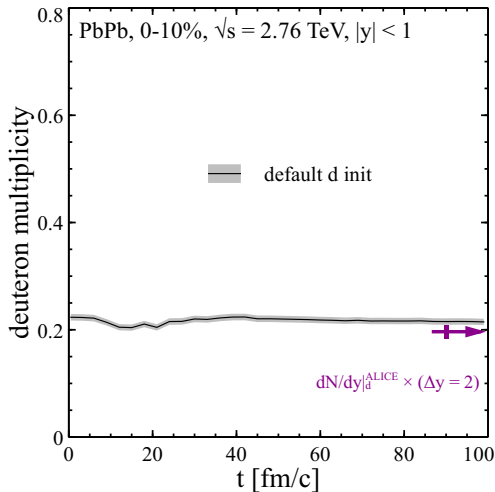
DO, Pang, Elfner, Koch, PRC99 (2019) no.4, 044907



After about 12-15 fm/c within 5%  $\pi d \leftrightarrow \pi np$  is equilibrated

# Deuteron yield

DO, Pang, Elfner, Koch, PRC99 (2019) no.4, 044907

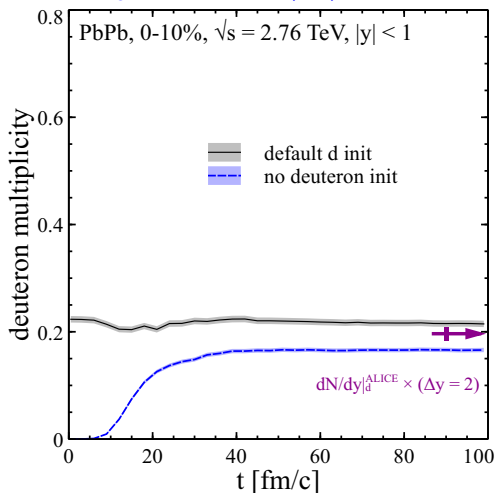


The yield is almost constant. Why? Does afterburner really play any role?



# Deuteron yield

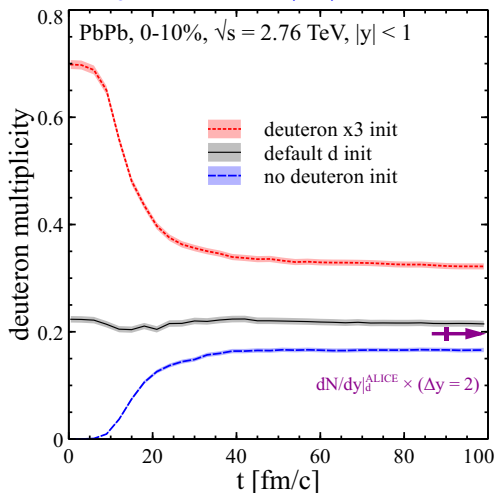
DO, Pang, Elfner, Koch, PRC99 (2019) no.4, 044907



No deuterons at particlization: also possible. Here **all** deuterons are from afterburner.

# Deuteron yield

DO, Pang, Elfner, Koch, PRC99 (2019) no.4, 044907



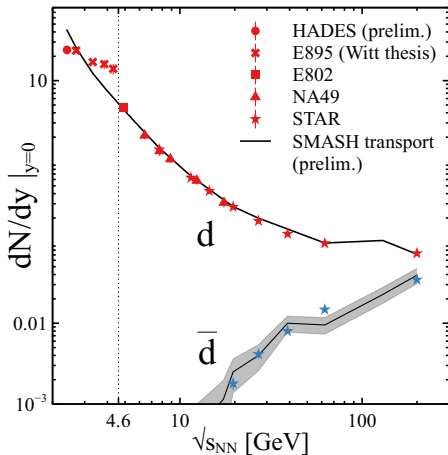
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# Why thermal model describes light nuclei yields at LHC

- Stable hadron yields ( $\pi$ ,  $K$ ,  $N$ ,  $\Lambda$ , ...) comprising resonances are fixed at chemical freeze-out
- Nuclei are kept in partial (relative) equilibrium by huge cross-sections of  $A + h \leftrightarrow A \times N + h$  until kinetic freeze-out
  - Therefore nuclei yields stay constant from hadron chemical freeze-out to kinetic
  - This picture works for all measured nuclei at LHC  
[Xu, Rapp, Eur. Phys. J. A55 \(2019\) no.5, 68](#)  
[Vovchenko et al, arXiv:1903.10024](#)
  - It works even if no nuclei are produced at chemical freeze-out  
[DO, Pang, Elfner, Koch, Phys.Rev. C99 \(2019\) no.4, 044907](#)  
[DO, Pang, Elfner, Koch, MDPI Proc. 10 \(2019\) no.1, 6](#)

# Exactly the same mechanism, lower energies

Data: Alt:2006dk, Anticic:2010mp, Adams:2003xp, Adamczyk:2017iwn,  
Abelev:2009bw, Adcox:2003nr, Klay:2001tf, Ahle:1999in



Still works for deuteron!

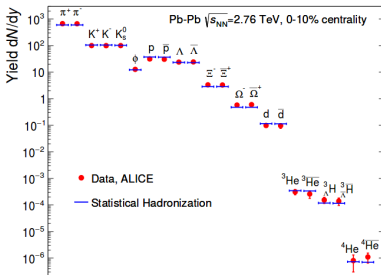
# Thermal model and “snowballs in hell” puzzle

- Nuclei formed early — at hadronic freeze-out

$$N_A \approx g_A V (\pi T m_A / 2)^{3/2} e^{(A\mu_B - m_A)/T}$$

- ALICE fit of yields, Pb+Pb,  $\sqrt{s_{NN}} = 2.76$  TeV:  $T = 155$  MeV
- Nuclei momentum spectra:  $T_{kin} \simeq 110$  MeV
- How can they survive from chemical to kinetic freeze-out?
- Binding energies:  $d, {}^3\text{He}, {}^3_{\Lambda}\text{H}, {}^4\text{He}$  – few MeV

## Snowballs in hell.



Andronic, Braun-Munzinger, Redlich, Stachel, Nature 561 (2018) no.7723, 321-3305

Light nuclei: rapid chemical freeze-out at 155 MeV, like hadrons?