

# Hadron production by equal-velocity quark combination mechanism in high energy collisions

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Refs:

- [arxiv:2008.03017v2](#), [arxiv:2007.14588v1](#),
- [Phys.RevC.102.014911\(2020\)](#), [Chin.Phys. C44, 014101\(2020\)](#),
- [Eur.Phys.J. C78, 344 \(2018\)](#), [Phys.Rev. C97, 064915 \(2018\)](#),
- [Phys. Lett. B774, 516\(2017\)](#), [Phys. Rev. D96,094010\(2017\)](#), [Phys. Rev. C95, 064911\(2017\)](#).

# outline

1. Quark (re-)combination/coalescence mechanism
2. Quark Number Scaling (**QNS**) property in  $p_T$  spectra of hadrons in pp and p-Pb collisions at LHC
3. Equal-Velocity Combination (**EVC**) of quarks and antiquarks at hadronization
4. EVC of light-flavor quarks
5. EVC of light-flavor quarks and charm quarks
6. Apply EVC to hadronic elliptic flow  $v_2$  in heavy-ion collisions
7. Energy-scan study for hadronic  $p_T$  spectra in Au+Au collisions at  
 $\sqrt{s_{NN}} = 7.7\text{--}200$  GeV
8. Summary and outlook

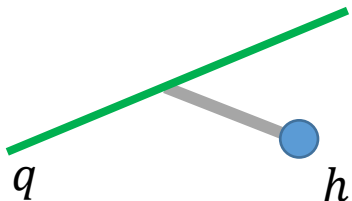
# 1. Hadronization

- the process of the formation of hadrons out of final-state quarks and/or gluons produced in high energy reactions
- Non-perturbative QCD process
- Currently modeled and/or parameterized in phenomenological methods.

## Two pictures:

### Fragmentation

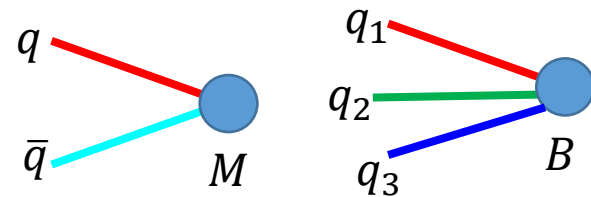
String fragmentation,  
cluster fragmentation,  
etc.



Probability  $D_{q \rightarrow h}(z)$

### Combination

quark (re-)combination,  
parton coalescence,  
etc.



## Combination models on market

### ➤ quark recombination model

R. J. Fries, B. Muller, C. Nonaka, S. A. Bass,  
e.g., Phys. Rev. C **68**, 044902 (2003)

### ➤ parton coalescence model

V. Greco, C. M. Ko, P. Lévai, L.W. Chen, et al.  
e.g., Phys. Rev. C **68**, 034904 (2003)

### ➤ quark recombination model

R. C. Hwa and C. B. Yang, e.g, Phys. Rev. C **70**, 024904 (2004)

### ➤ resonance recombination model

L. Ravagli, R. Rapp, e.g., Phys.Lett. B 655,126 (2007)  
M. He, R.J. Fries, R. Rapp, e.g., Phys.Rev.C 82, 034907 (2010)

### ➤ quark combination model (Shandong Group)

Q.B.Xie, F.L.Shao, et al., e.g., Phys. Rev. C71,044903 (2005)  
phenomenological combine rule

### ➤ transport model

P.F.Zhuang, et al, e.g., Phys. Rev. C76, 014907(2009)

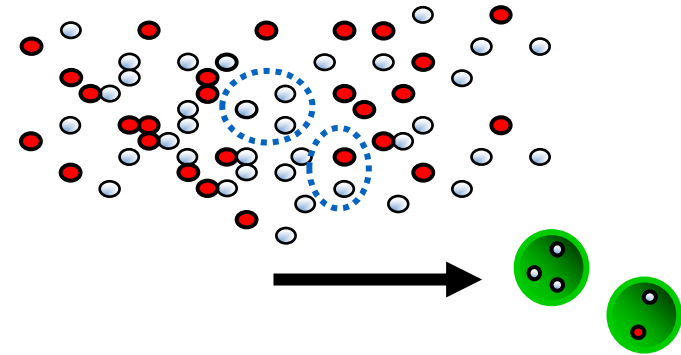
### ➤ quark molecular dynamics model

M. Hofmann et al, e.g., Nucl. Phys. B 478,161(2000)

### ➤ variational model

A.Alala, etal, e.g., Phys. Rev. C77,044901(2009)

....



what kind of quarks?

how to combine?

how to test in experiments?

## 2.

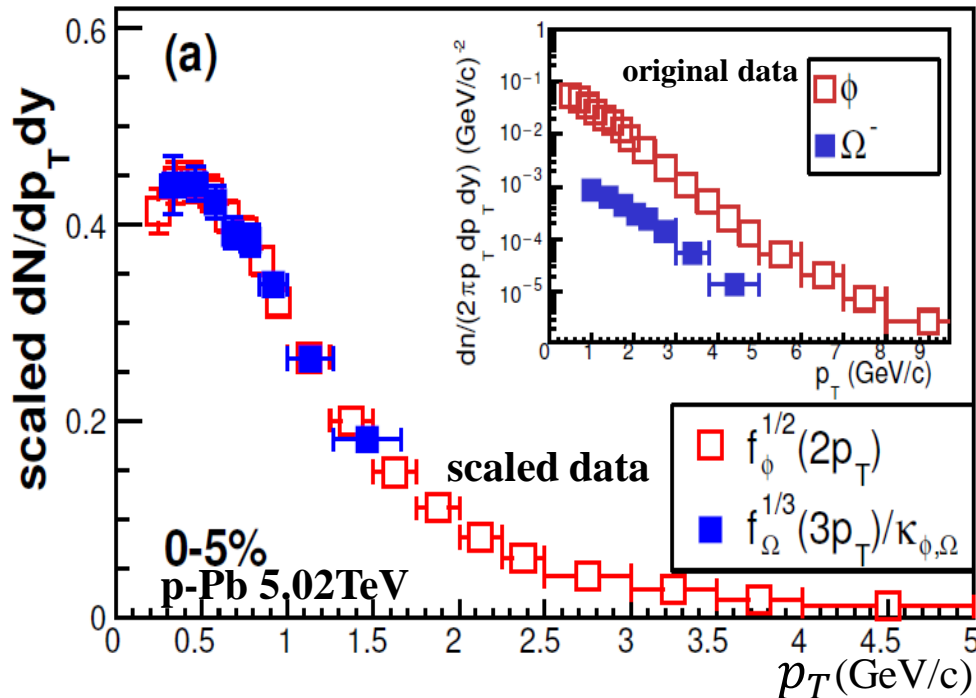
# Quark Number Scaling (**QNS**) property in $p_T$ spectra of hadrons in pp and p-Pb collisions at LHC

Refs:

- Jian-wei Zhang, Hai-hong Li, Feng-lan Shao, and Jun Song, [Chin.Phys. C44, 014101\(2020\)](#).
- Jun Song, Xing-rui Gou, Feng-lan Shao, Zuo-tang Liang, [Phys. Lett. B774, 516\(2017\)](#).
- Xing-rui Gou, Feng-lan Shao, Rui-qin Wang, Hai-hong Li, Jun Song, [Phys. Rev. D96,094010\(2017\)](#).

## Manipulate exp data of $p_T$ spectra for $\Omega(sss)$ and $\phi(s\bar{s})$ in p-Pb collisions at 5.02 TeV

- (1) divide  $p_T$  bin by quark number,  $p_{T_\Omega}/3$ ,  $p_{T_\phi}/2$ ,
- (2) take the inverse quark number power of density  $dN_h/dp_T dy$ ,  
i.e.  $dN_\Omega^{1/3}/dp_T dy$  and  $dN_\phi^{1/2}/dp_T dy$
- (3) Divide the scaled  $\Omega$  data by a constant  $\kappa$  to keep the same magnitude with that of  $\phi$



ALICE data : Phys. Lett. B 758, 389(2016).  
Eur. Phys. J. C76, 245 (2016)

## mathematic relationship

$$f_{\Omega}^{1/3}(3p_T) = \kappa_{\phi,\Omega} f_{\phi}^{1/2}(2p_T)$$

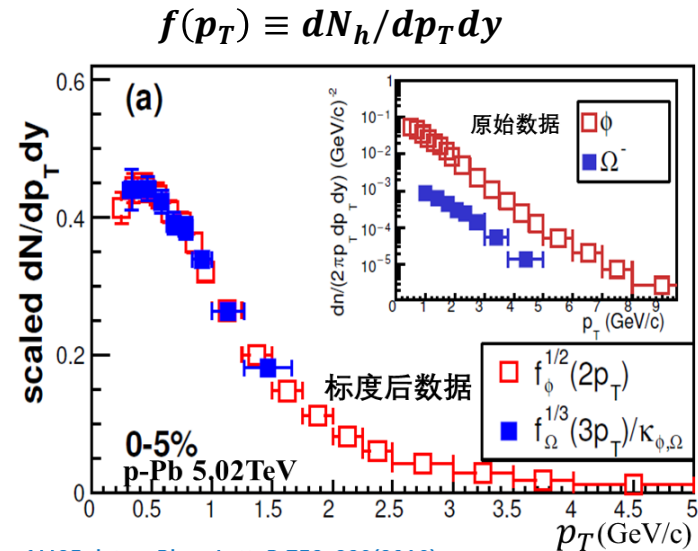
## in another form

$$f_{\Omega}(3p_T) = \kappa_{\Omega} f_s^3(p_T)$$

$$f_{\phi}(2p_T) = \kappa_{\phi} f_s^2(p_T)$$

Coefficients  $\kappa_{\phi}$ ,  $\kappa_{\Omega}$ ,  $\kappa_{\phi,\Omega}$  are independent of  $p_T$

A clear signal of quark combination at hadronization!

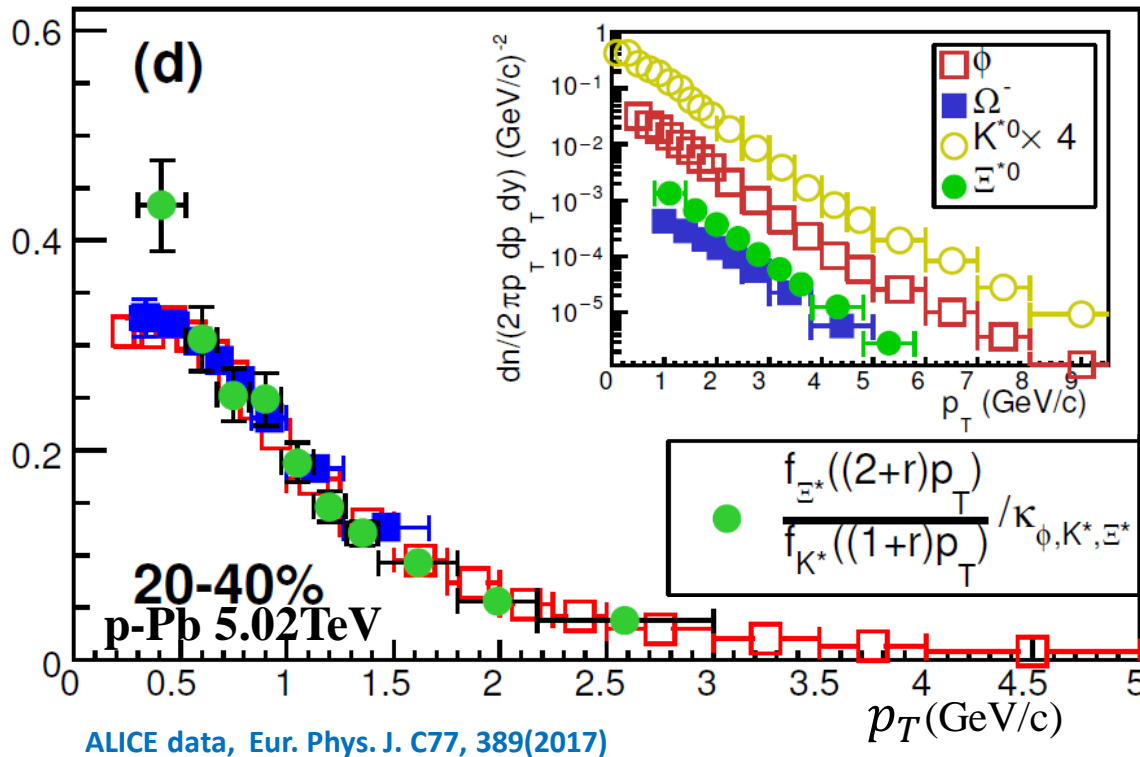


ALICE data : Phys. Lett. B 758, 389(2016).  
Eur. Phys. J. C76, 245 (2016)

For  $\Xi^{*0}(uss)$  and  $K^{*0}(d\bar{s})$ , we also find

$$\frac{f_{\Xi^{*0}}((2+r)p_T)}{f_{K^{*0}}((1+r)p_T)} = \kappa_{K^*,\Xi^*} f_s(p_T)$$

where  $r \approx 0.67$

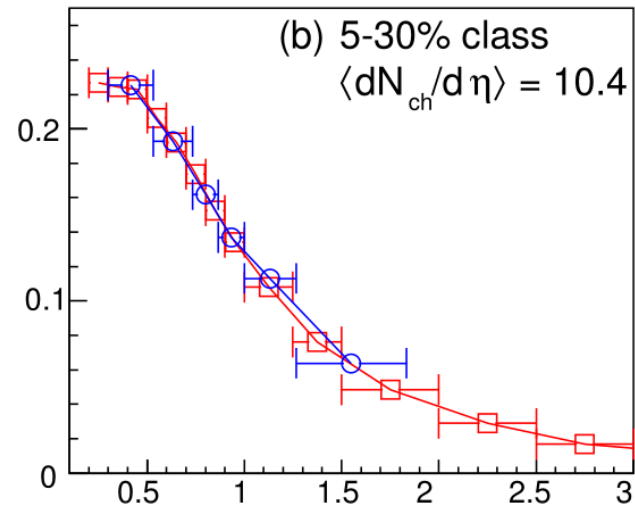
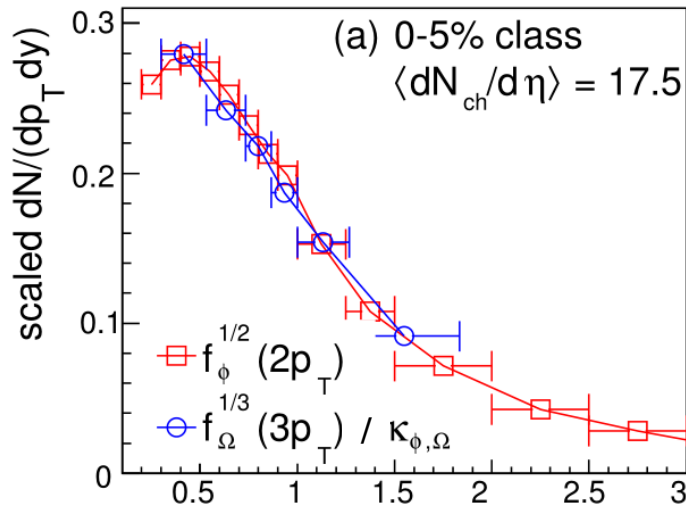


$$m_s \approx 500 \text{ MeV}$$

$$m_u \approx 330 \text{ MeV}$$

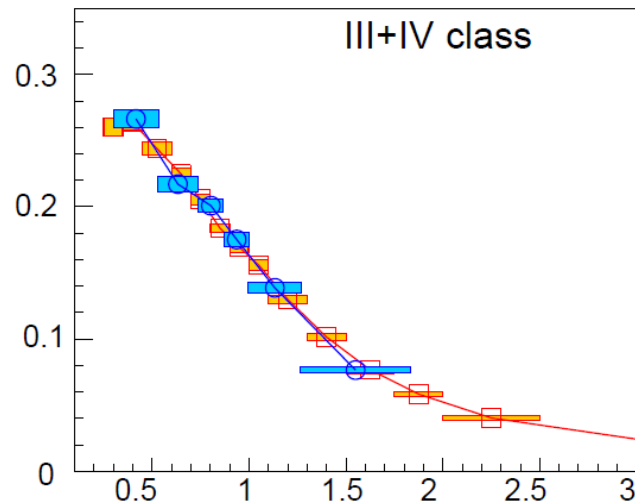
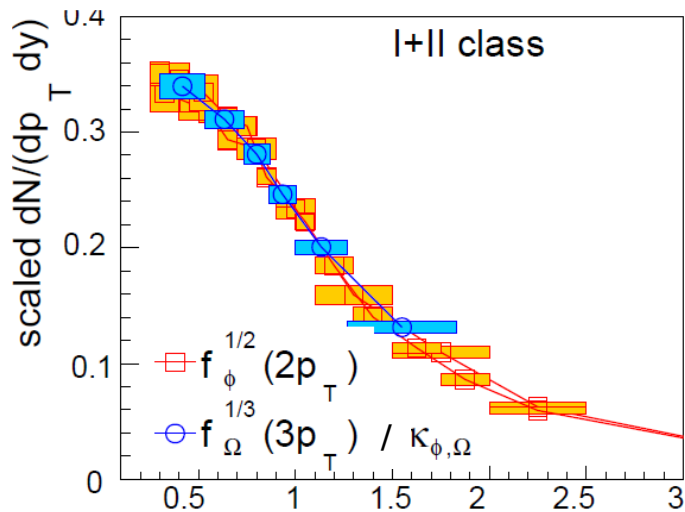
$$\frac{m_u}{m_s} \approx r$$





pp at  $\sqrt{s} = 7\text{TeV}$

Zhang,Shao,Song,CPC(2020)



pp at  $\sqrt{s} = 13\text{TeV}$

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3.

**Equal-Velocity Combination (EVC) of quarks  
and antiquarks at hadronization**

Start from general formula in momentum space

$$f_{B_j}(p_B) = \int dp_1 dp_2 dp_3 R_{B_j}(p_1, p_2, p_3; p_B) f_{q_1 q_2 q_3}(p_1, p_2, p_3)$$
$$f_{M_j}(p_M) = \int dp_1 dp_2 R_{M_j}(p_1, p_2; p_M) f_{q_1 \bar{q}_2}(p_1, p_2)$$

Assume independent distribution of (anti-)quarks

$$f_{q_1 q_2 q_3}(p_1, p_2, p_3) = f_{q_1}(p_1) f_{q_2}(p_2) f_{q_3}(p_3)$$
$$f_{q_1 \bar{q}_2}(p_1, p_2) = f_{q_1}(p_1) f_{\bar{q}_2}(p_2)$$

## equal velocity combination (EVC) approximation

$$R_{B_j}(p_1, p_2, p_3; p_B) = \kappa_{B_j} \prod_{i=1}^3 \delta(p_i - x_i p_B)$$
$$R_{M_j}(p_1, p_2; p_M) = \kappa_{M_j} \prod_{i=1}^2 \delta(p_i - x_i p_M)$$

**momentum fraction**

for meson  $\mathbf{x}_{1,2} = \frac{m_{1,2}}{m_1+m_2}$ , for baryon  $\mathbf{x}_{1,2,3} = \frac{m_{1,2,3}}{m_1+m_2+m_3}$ ,

$$m_s = 500 \text{ MeV}, m_u = m_d = 330 \text{ MeV}.$$

**We obtain**

$$f_{B_j}(p_B) = \kappa_{B_j} f_{q_1}(x_1 p_B) f_{q_2}(x_2 p_B) f_{q_3}(x_3 p_B)$$

$$f_{M_j}(p_M) = \kappa_{M_j} f_{q_1}(x_1 p_M) f_{\bar{q}_2}(x_2 p_M)$$

## For combination of $s$ and $\bar{s}$

$$f_{\Omega}(3p_T) = \kappa_{\Omega} f_s^3(p_T)$$

$$f_{\phi}(2p_T) = \kappa_{\phi} f_s^2(p_T)$$



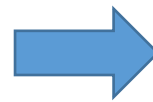
$$f_{\Omega}^{\frac{1}{3}}(3p_T) = \kappa_{\phi, \Omega} f_{\phi}^{\frac{1}{2}}(2p_T)$$

$f_s(p_T) = f_{\bar{s}}(p_T)$  is taken at LHC

For combination of  $u(d)$  and  $s$ , denote  $\frac{x_u}{x_s} = \frac{m_u}{m_s} = r$

$$f_{\Xi^{*0}}((2+r)p_T) = \kappa_{\Xi^{*0}} f_s^2(p_T) f_u(r p_T)$$

$$f_{K^{*0}}((1+r)p_T) = \kappa_{K^{*0}} f_s(p_T) f_{\bar{d}}(r p_T)$$



$$\frac{f_{\Xi^{*0}}((2+r)p_T)}{f_{K^{*0}}((1+r)p_T)} = \kappa_{\phi, K^*, \Xi^*} f_{\phi}^{\frac{1}{2}}(2p_T)$$

Consider stochastic combination and flavor-blind of strong interaction

$$\frac{\kappa_{Mj}}{A_{Mj}} = C_{Mj} P_{q\bar{q} \rightarrow M}$$

$$\frac{\kappa_{Bj}}{A_{Bj}} = C_{Bj} N_{iter} P_{qqq \rightarrow B}$$

$A_{Bj} = 1/\int dp_T \prod_{i=1}^3 f_{q_i}^{(n)}(x_i p_T)$ ,  $A_{Mj} = 1/\int dp_T f_{q_1}^{(n)}(x_1 p_T) f_{\bar{q}_2}^{(n)}(x_2 p_T)$  are normalization coefficients of jointed quark distributions.

$$P_{q\bar{q} \rightarrow M} \approx \frac{2}{x(1-z^2)} \left[ 1 - z \frac{(1+z)^a + (1+z)^a}{(1+z)^a - (1+z)^a} \right], \quad \text{averaged probability of } q\bar{q} \rightarrow M$$

$$P_{qqq \rightarrow B} \approx \frac{8}{3x^2(1+z)^3} \frac{(1+z)^a}{(1+z)^a - (1+z)^a}, \quad \text{averaged probability of } qqq \rightarrow B$$

Song, Shao, PRC 88, 027901(2013)

$N_{iter}=1,3,6$  for three identical, two identical and three different flavors

$$x = N_q + N_{\bar{q}}, \quad z = (N_q - N_{\bar{q}})/x, \quad a = 1 + \frac{1}{3} \left( \overline{N}_M / \overline{N}_B \right)_{z=0} \approx 4.86 \pm 0.1 \text{ light-flavor sector}$$

$N_q$  number of all quarks;  $N_{\bar{q}}$  that of all antiquarks

$C_{M_j}$  and  $C_{B_j}$  are fine-tune parameters

$$C_{M_j} = \begin{cases} \frac{1}{1+R_{V/P}} & \text{for } J^P = 0^- \text{ mesons} \\ \frac{R_{V/P}}{1+R_{V/P}} & \text{for } J^P = 1^- \text{ mesons} \end{cases} \quad C_{B_j} = \begin{cases} \frac{R_{O/D}}{1+R_{O/D}} & \text{for } J^P = (1/2)^+ \text{ baryons} \\ \frac{1}{1+R_{O/D}} & \text{for } J^P = (3/2)^+ \text{ baryons} \end{cases} \quad \begin{array}{l} R_{V/P} \approx 0.5 \\ R_{O/D} \approx 2.0 \end{array}$$

Include decay contributions

$$f_{h_j}^{(final)}(p) = f_{h_j}(p) + \sum_{i \neq j} \int dp' f_{h_i}(p') D_{ij}(p', p)$$

decay function  $D_{ij}(p', p)$  is determined by the decay kinematics and decay branch ratios in PDG

**Model inputs:**  $f_{q_i}(p)$  fixed by experimental data

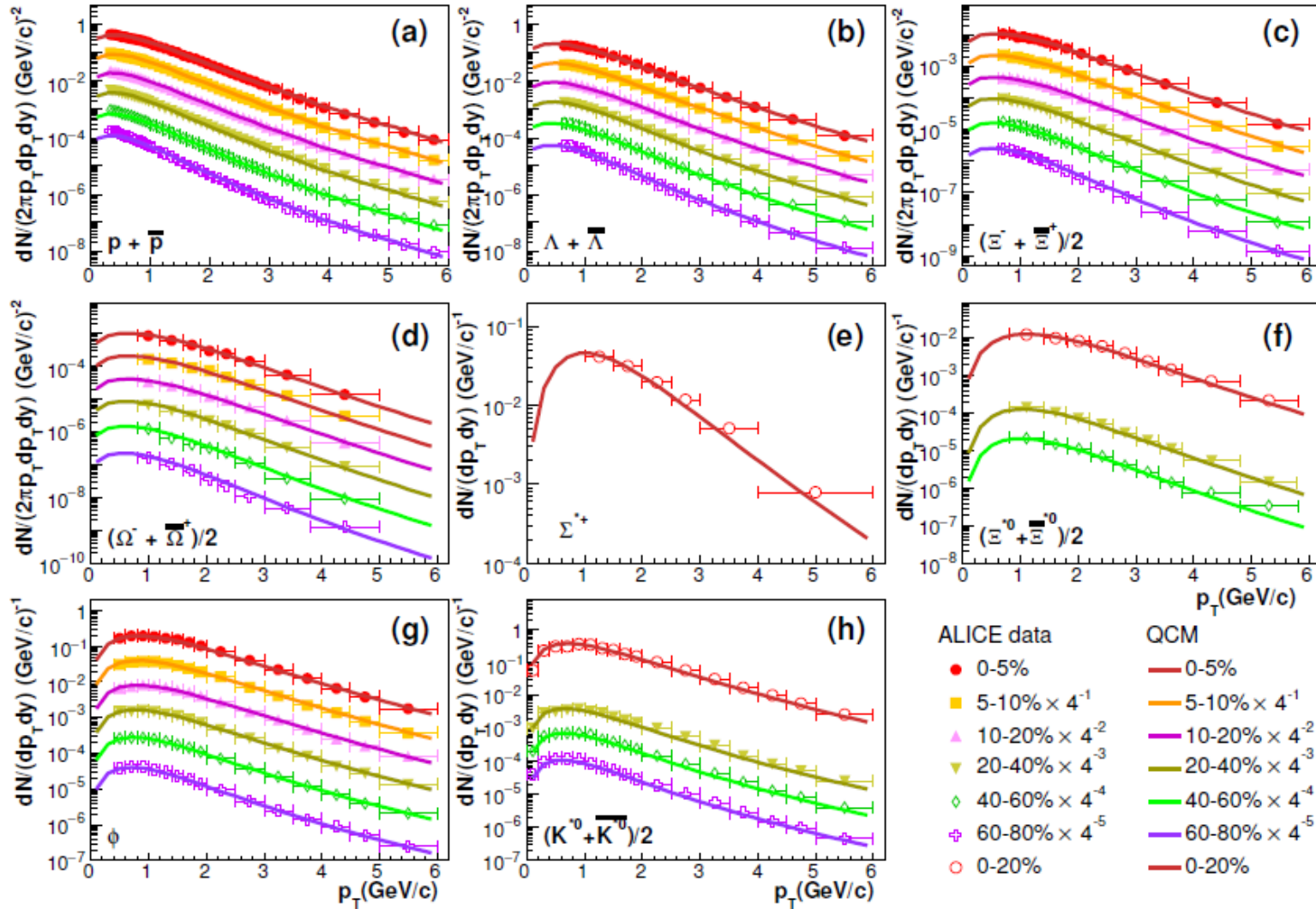
# 4.

## EVC of light-flavor quarks

Refs:

- Jian-wei Zhang, Hai-hong Li, Feng-lan Shao, and Jun Song, [Chin.Phys. C44, 014101\(2020\)](#).
- Jun Song, Xing-rui Gou, Feng-lan Shao, Zuo-tang Liang, [Phys. Lett. B774, 516\(2017\)](#).
- Xing-rui Gou, Feng-lan Shao, Rui-qin Wang, Hai-hong Li, Jun Song, [Phys. Rev. D96,094010\(2017\)](#).
- Feng-lan Shao, Guo-jing Wang, Rui-qin Wang, Hai-hong Li, Jun Song, [Phys. Rev. C95, 064911\(2017\)](#).

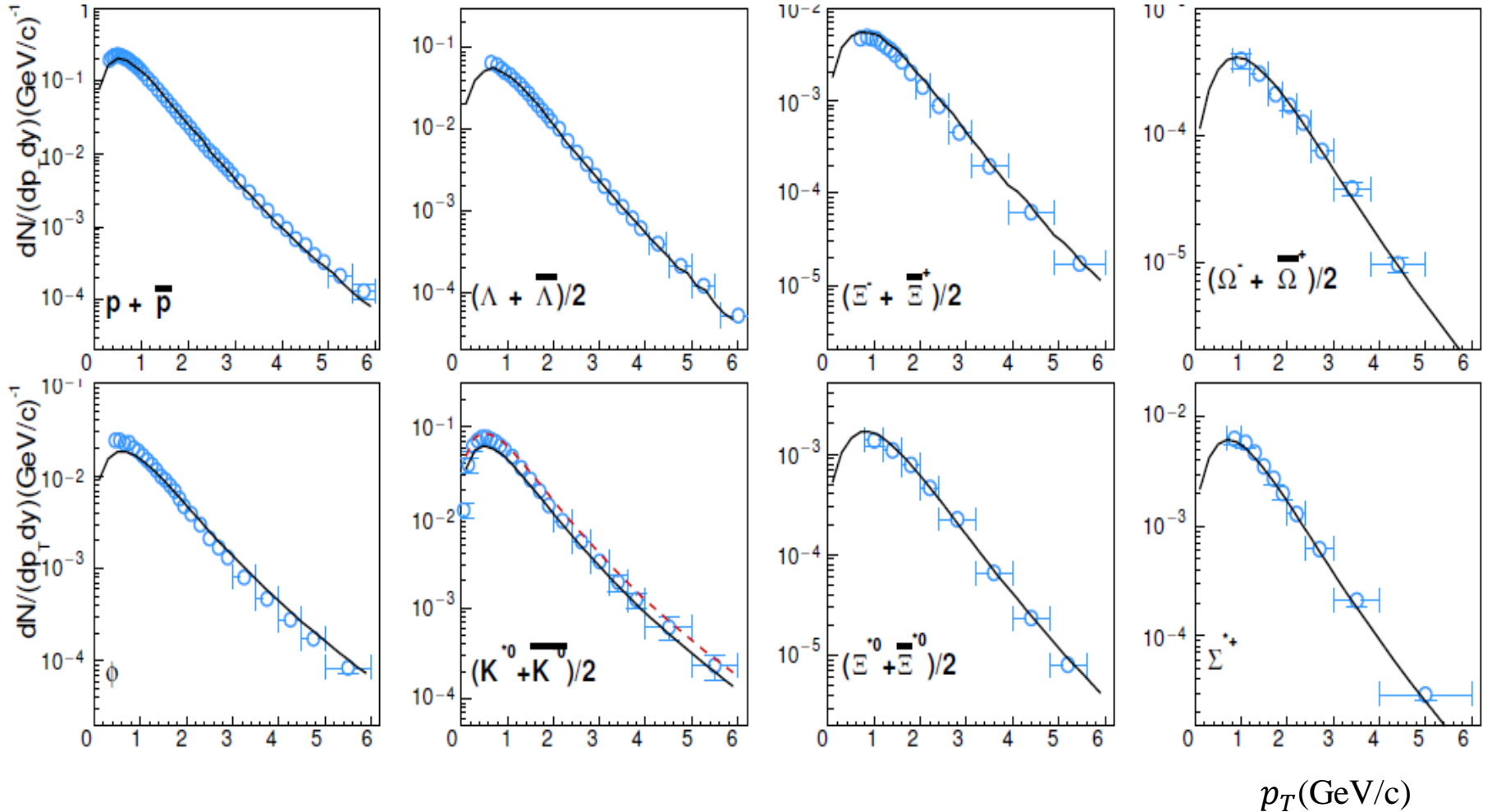




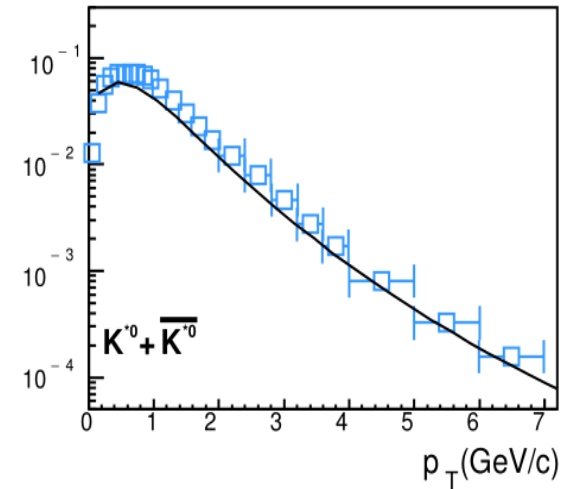
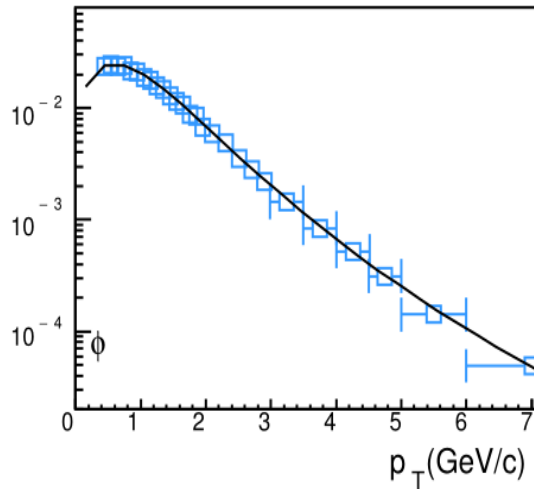
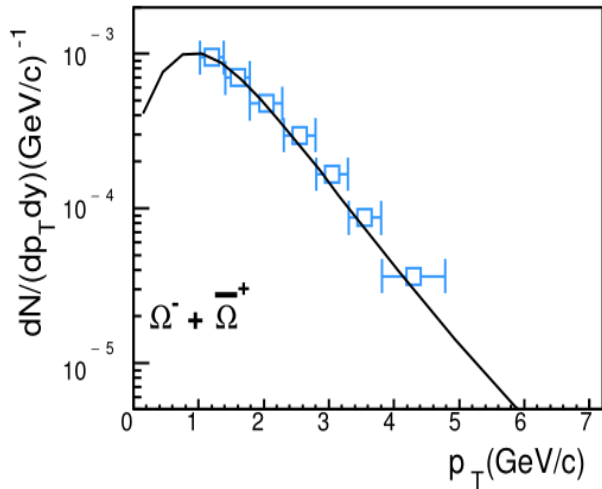
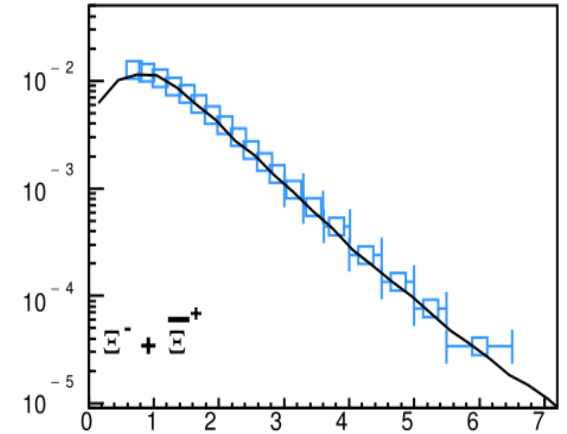
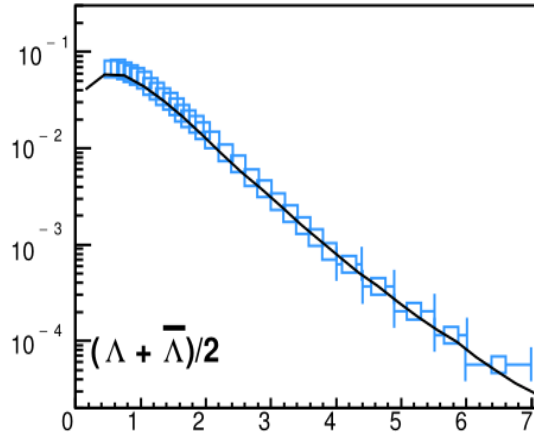
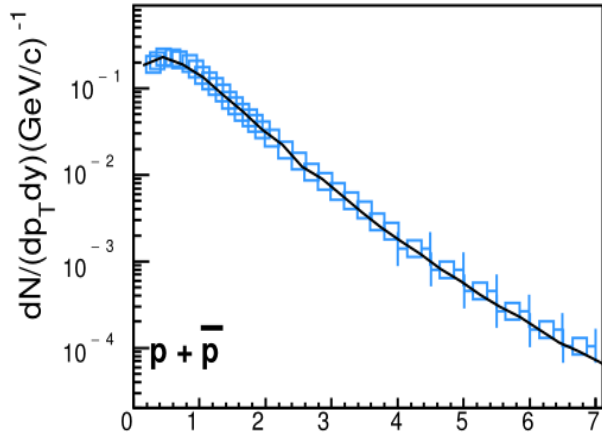
Two inputs:  $f_u(p_T) = f_d(p_T) = f_{\bar{u}}(p_T) = f_{\bar{d}}(p_T)$   
 $f_s(p_T) = f_{\bar{s}}(p_T)$

## pp collisions at 7 TeV

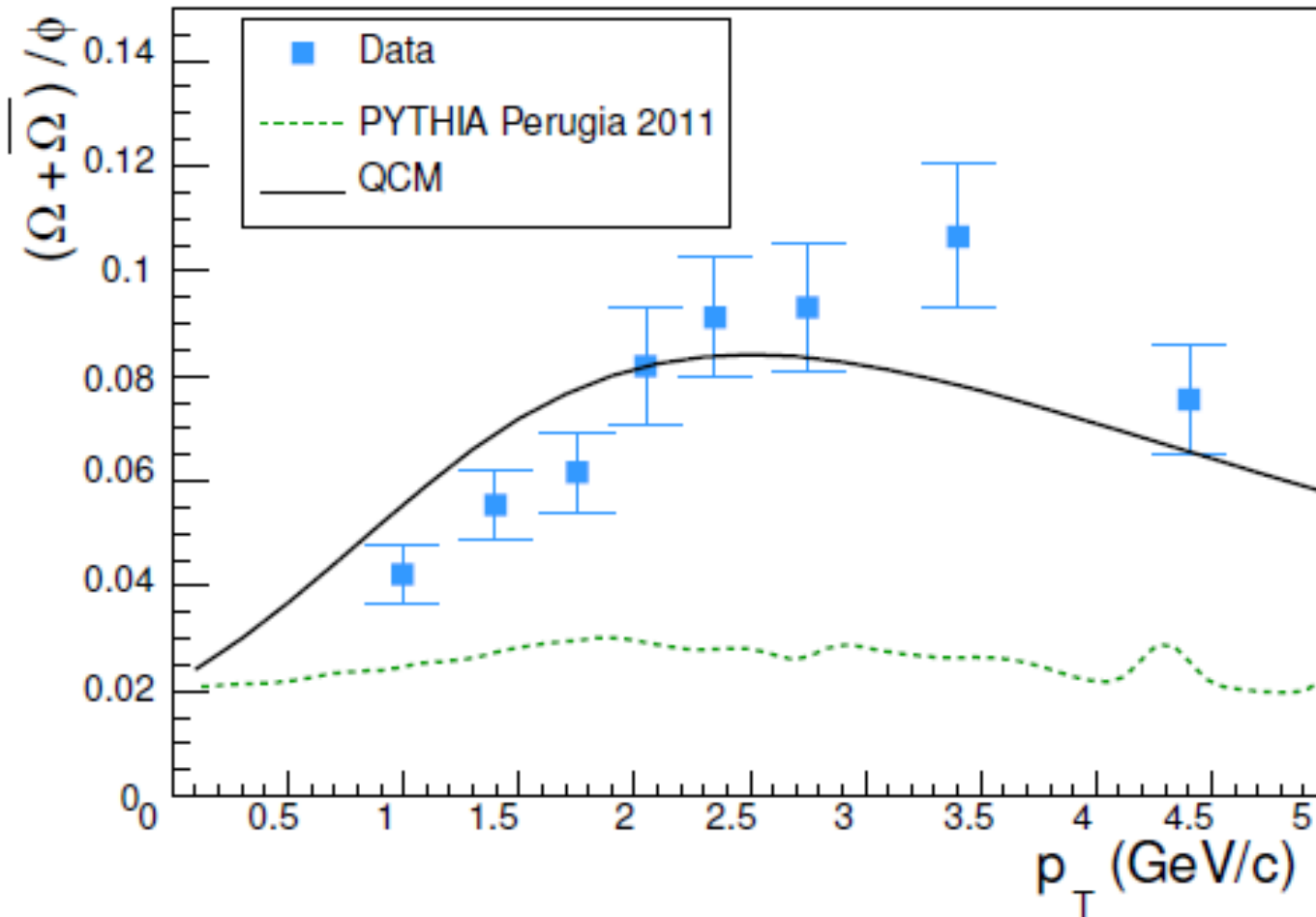
## Minimum bias events



## pp collisions at 13 TeV



## Baryon/Meson ratio



Minimum bias events in pp collisions at 7 TeV

# Kaon and pion production in EVC

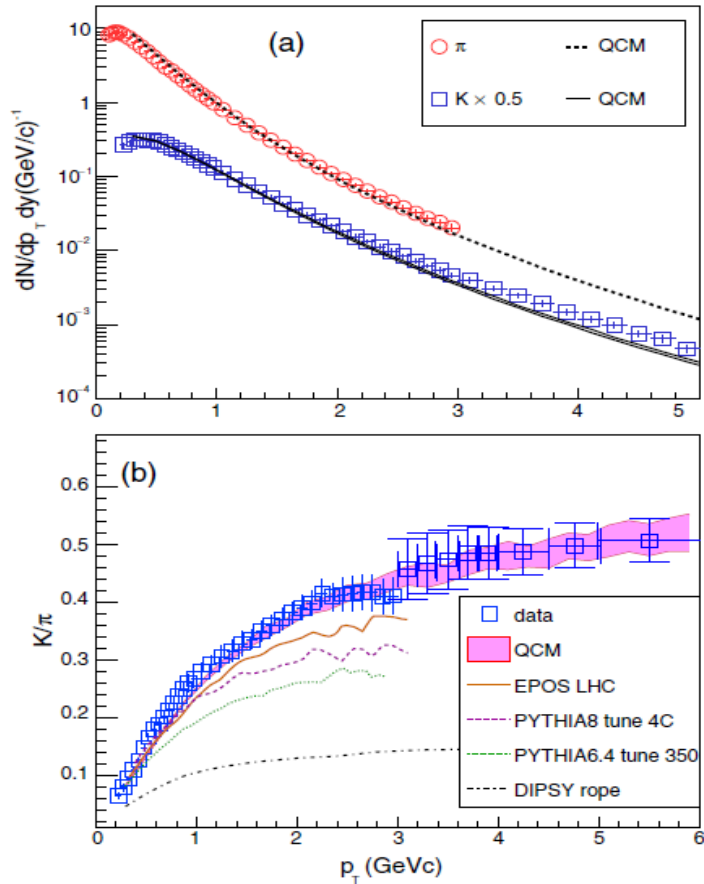


FIG. 10. Midrapidity  $p_T$  spectra of kaon and pion in minimum-bias  $pp$  collisions at  $\sqrt{s} = 7$  TeV and the ratio between them. Symbols are experimental data [22,58] and results of other models and/or event generators are taken from [22,51].

reconcile off-shell issue

$$u + \bar{d} \rightarrow \pi^+ + X$$

$$u + \bar{s} \rightarrow K^+ + X$$

$X$  is identified as pions.

# 5.

## EVC of charm and light-flavor quarks

Refs:

- Jun Song, Hai-hong Li, Feng-lan Shao, [Eur.Phys.J. C78, 344 \(2018\)](#).
- Hai-hong Li, Feng-lan Shao, Jun Song, Rui-qin Wang, [Phys.Rev. C97, 064915 \(2018\)](#).

## single-charmed hadrons in EVC

$$f_D(p_T) = \kappa_D f_c \left( \frac{r_{cu}}{1 + r_{cu}} p_T \right) f_{\bar{u}} \left( \frac{1}{1 + r_{cu}} p_T \right)$$

$$f_{D_s}(p_T) = \kappa_{D_s} f_c \left( \frac{r_{cs}}{1 + r_{cs}} p_T \right) f_{\bar{s}} \left( \frac{1}{1 + r_{cs}} p_T \right)$$

$$f_{\Lambda_c}(p_T) = \kappa_{\Lambda_c} f_c \left( \frac{r_{cu}}{2 + r_{cu}} p_T \right) f_u^2 \left( \frac{1}{1 + r_{cu}} p_T \right)$$

0.75-0.85

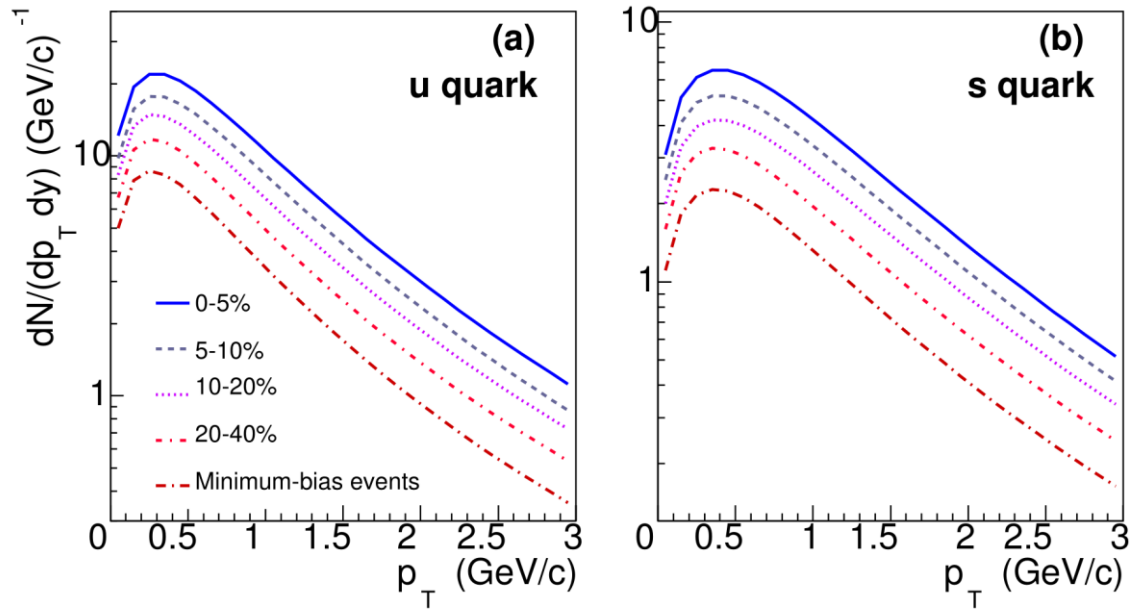
0.15-0.25

$$r_{cu} = \frac{m_c}{m_u}$$
$$r_{cs} = \frac{m_c}{m_s}$$

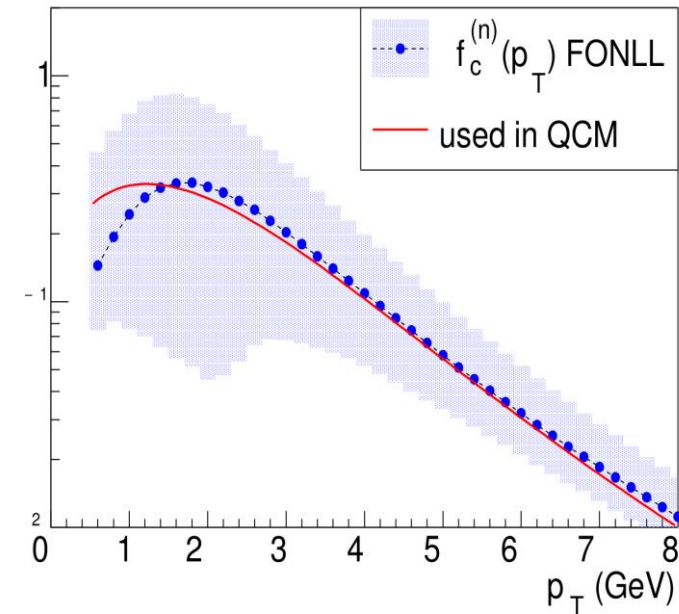
$$m_c = 1.5 \text{ GeV}$$

## Quark spectra at hadronization are known.

e.g., in p-Pb collisions at 5.02 TeV



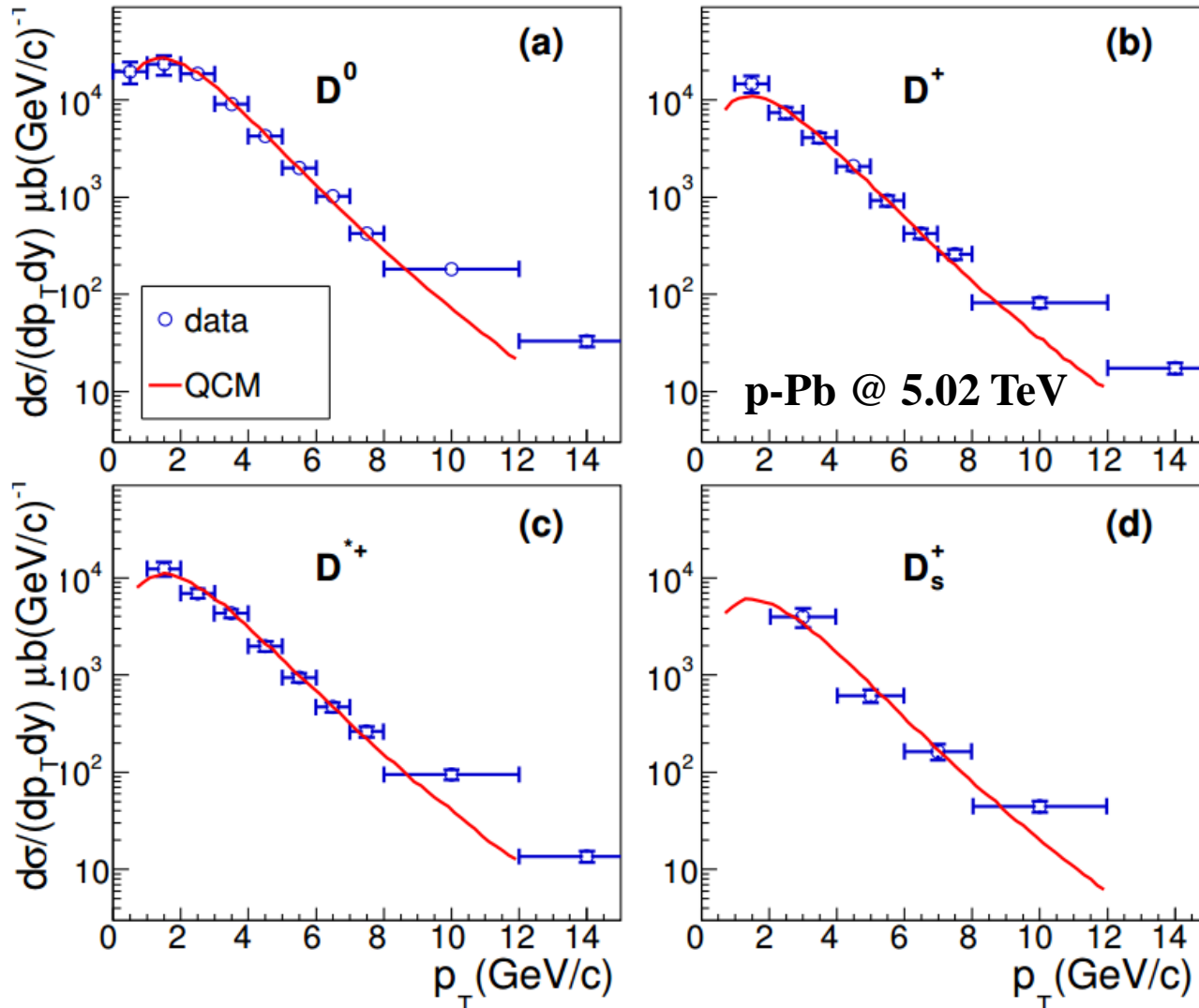
obtained from light-flavor hadrons



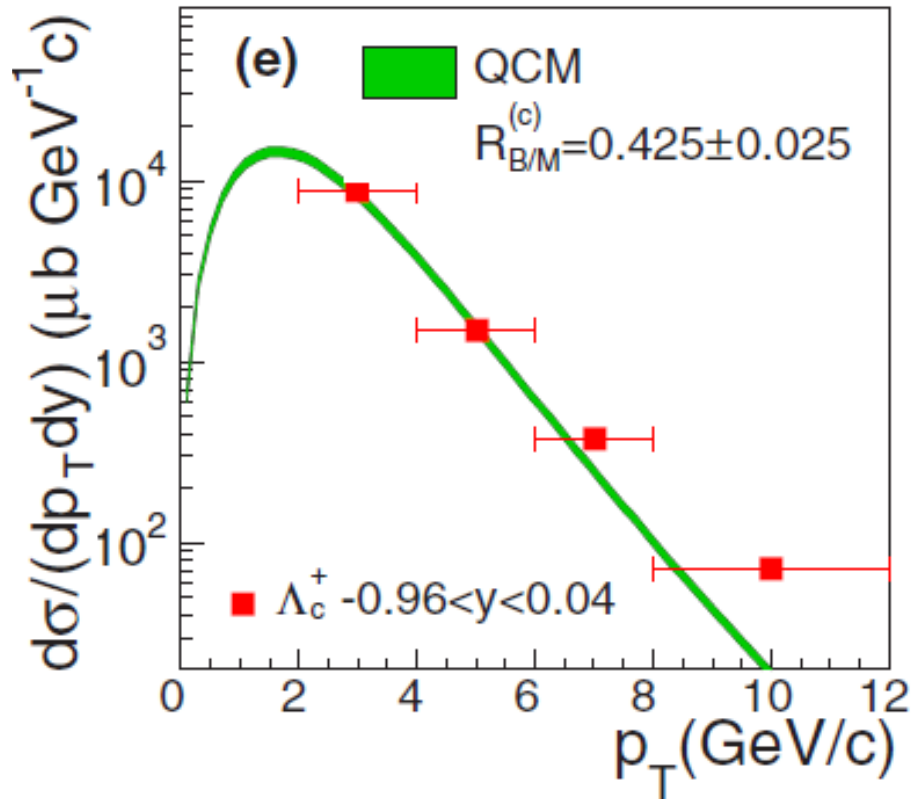
consistent with pQCD calculations



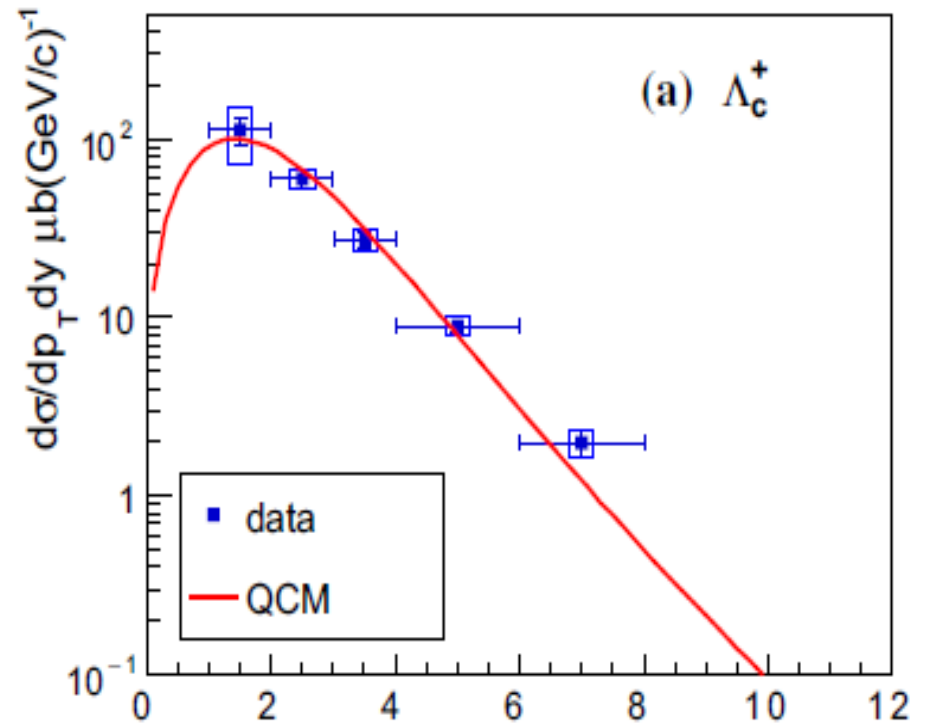
**$c$  quark of  $p_{T,c} \lesssim 6$  +  $l$  quark of  $p_{T,l} \lesssim 2$  GeV/c**



## $p_T$ spectrum of $\Lambda_c^+$

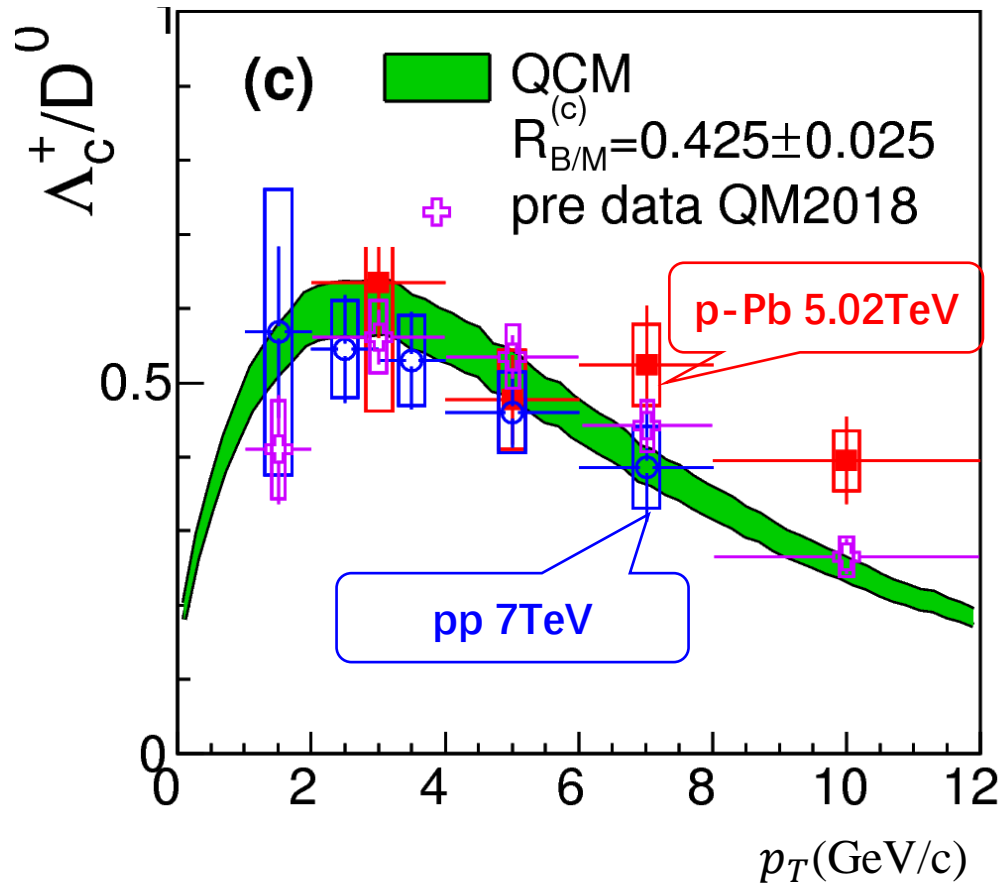


pPb @ 5.02 TeV

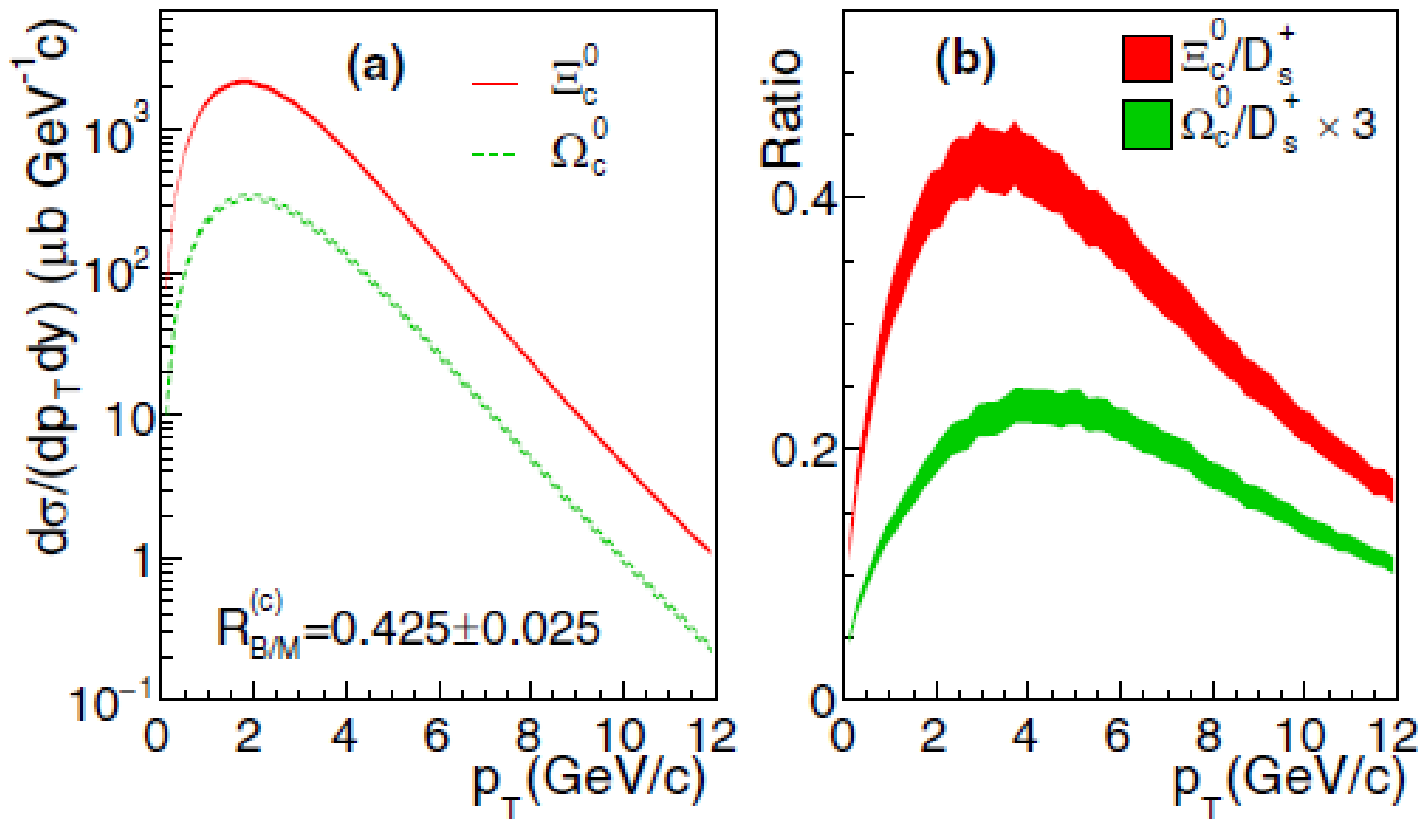


pp @ 7 TeV

# $p_T$ dependence of charmed Baryon/Meson ratio



## Prediction of $\Xi_c^0$ , $\Omega_c^0$



pPb @ 5.02 TeV

# 6.

## Apply EVC to elliptic flow $v_2$ of hadrons

Refs:

- Jun Song, Hai-hong Li, and Feng-lan Shao, [arxiv:2008.03017v2](https://arxiv.org/abs/2008.03017v2).

## Applying EVC to $(p_T, \varphi)$ plane

$$f_{M_i}(p_T, \varphi) = \kappa_{M_i} f_{q_1}(x_1 p_T, \varphi) f_{\bar{q}_2}(x_2 p_T, \varphi)$$

$$f_{B_i}(p_T, \varphi) = \kappa_{B_i} f_{q_1}(x_1 p_T, \varphi) f_{q_2}(x_2 p_T, \varphi) f_{q_3}(x_3 p_T, \varphi)$$

## with quark distribution

$$f_q(p_T, \varphi) = f_q(p_T) \left[ 1 + 2 \sum_{n=1}^{\infty} v_{n,q}(p_T) \cos(n\varphi) \right]$$

## Hadron flow is

$$v_{n,h}(p_T) = \frac{\int f_h(p_T, \varphi) \cos(n\varphi) d\varphi}{\int f_h(p_T, \varphi) d\varphi}$$

We neglect  $v_{1,q}$  ( $\lesssim 10^{-3}$  at mid-rapidity in HIC) and  $v_{n,q}^{3,4,\dots}$  and obtain

$$v_{2,M_i}(p_T)$$

$$= v_{2,q_1}(x_1 p_T) \left[ 1 + \sum_{n=2}^{\infty} \frac{v_{n,q_1}(x_1 p_T)}{v_{2,q_1}(x_1 p_T)} v_{n+2,\bar{q}_2}(x_2 p_T) \right]$$

$$+ v_{2,\bar{q}_2}(x_2 p_T) \left[ 1 + \sum_{n=2}^{\infty} \frac{v_{n,\bar{q}_2}(x_2 p_T)}{v_{2,\bar{q}_2}(x_2 p_T)} v_{n+2,q_1}(x_1 p_T) \right]$$

$$v_{2,B_i}(p_T)$$

$$= v_{2,q_1}(x_1 p_T) \left\{ 1 + \sum_{n=2}^{\infty} \frac{v_{n,q_1}(x_1 p_T)}{v_{2,q_1}(x_1 p_T)} [v_{n+2,q_2}(x_2 p_T) + v_{n+2,q_3}(x_3 p_T)] \right\}$$

$$+ v_{2,q_2}(x_2 p_T) \left\{ 1 + \sum_{n=2}^{\infty} \frac{v_{n,q_2}(x_2 p_T)}{v_{2,q_2}(x_2 p_T)} [v_{n+2,q_1}(x_1 p_T) + v_{n+2,q_3}(x_3 p_T)] \right\}$$

$$+ v_{2,q_3}(x_3 p_T) \left\{ 1 + \sum_{n=2}^{\infty} \frac{v_{n,q_3}(x_3 p_T)}{v_{2,q_3}(x_3 p_T)} [v_{n+2,q_1}(x_1 p_T) + v_{n+2,q_2}(x_2 p_T)] \right\}$$

**Finally, we obtain the simplest form**

$$v_{2,M_i}(p_T) = v_{2,q_1}(x_1 p_T) + v_{2,\bar{q}_2}(x_2 p_T),$$

$$v_{2,B_i}(p_T) = v_{2,q_1}(x_1 p_T) + v_{2,q_2}(x_2 p_T) + v_{2,q_3}(x_3 p_T).$$

$$v_{2,p}(p_T) = 3v_{2,u}(p_T/3), \quad (13)$$

$$v_{2,\Lambda}(p_T) = 2v_{2,u}\left(\frac{1}{2+r}p_T\right) + v_{2,s}\left(\frac{r}{2+r}p_T\right), \quad (14)$$

$$v_{2,\Xi}(p_T) = v_{2,u}\left(\frac{1}{1+2r}p_T\right) + 2v_{2,s}\left(\frac{r}{1+2r}p_T\right), \quad (15)$$

$$v_{2,\Omega}(p_T) = 3v_{2,s}(p_T/3), \quad (16)$$

$$v_{2,\phi}(p_T) = 2v_{2,s}(p_T/2) \quad r = \frac{x_s}{x_u} = \frac{m_s}{m_u} \quad (17)$$



**Express quark  $v_2$  via hadron's  $v_2$**

$$v_{2,u}(p_T) = \frac{1}{3} v_{2,p}(3p_T), \quad (18)$$

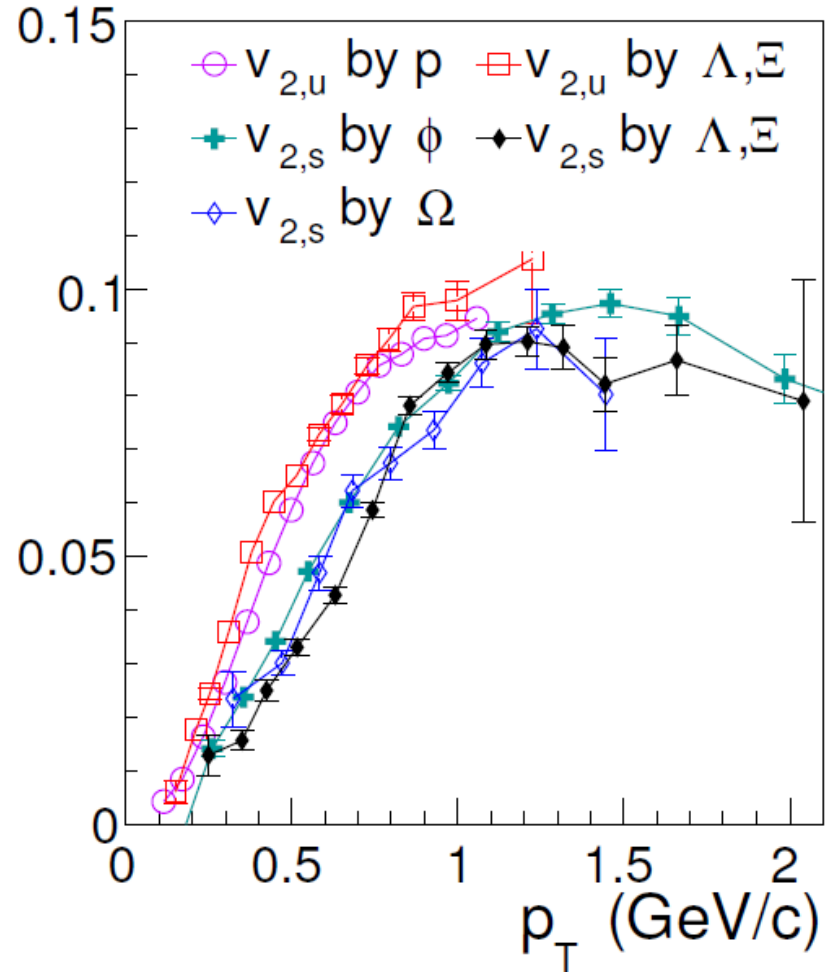
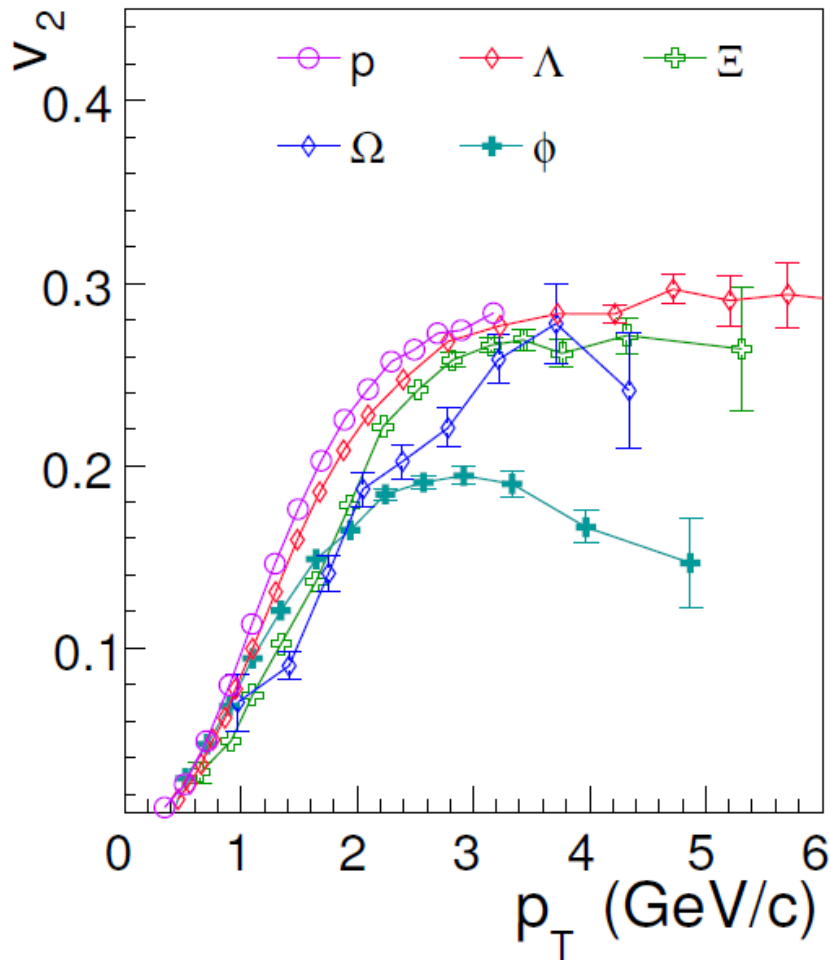
$$v_{2,u}(p_T) = \frac{1}{3} [2v_{2,\Lambda}((2+r)p_T) - v_{2,\Xi}((1+2r)p_T)], \quad (19)$$

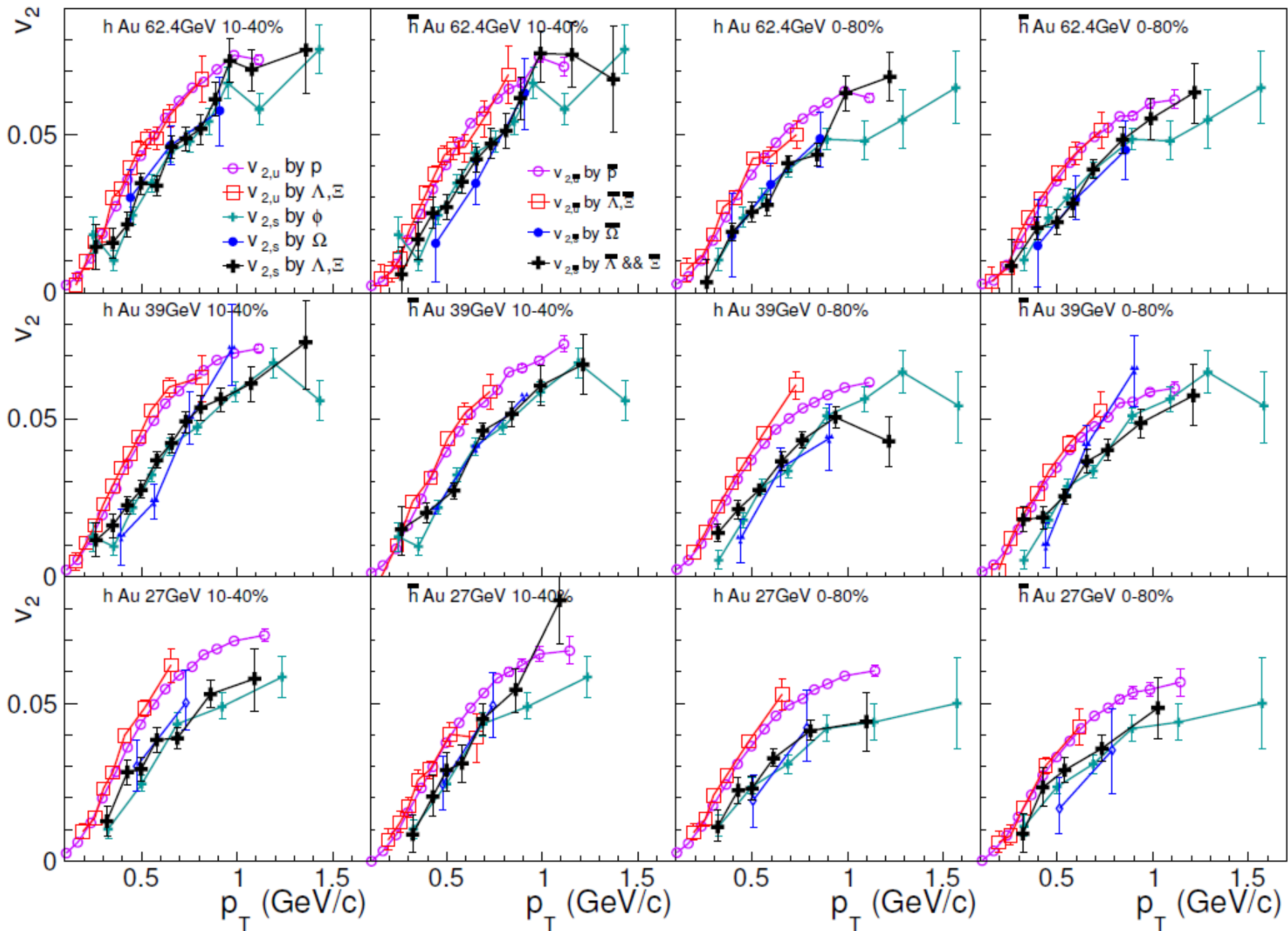
$$v_{2,s}(p_T) = \frac{1}{3} v_{2,\Omega}(3p_T), \quad (20)$$

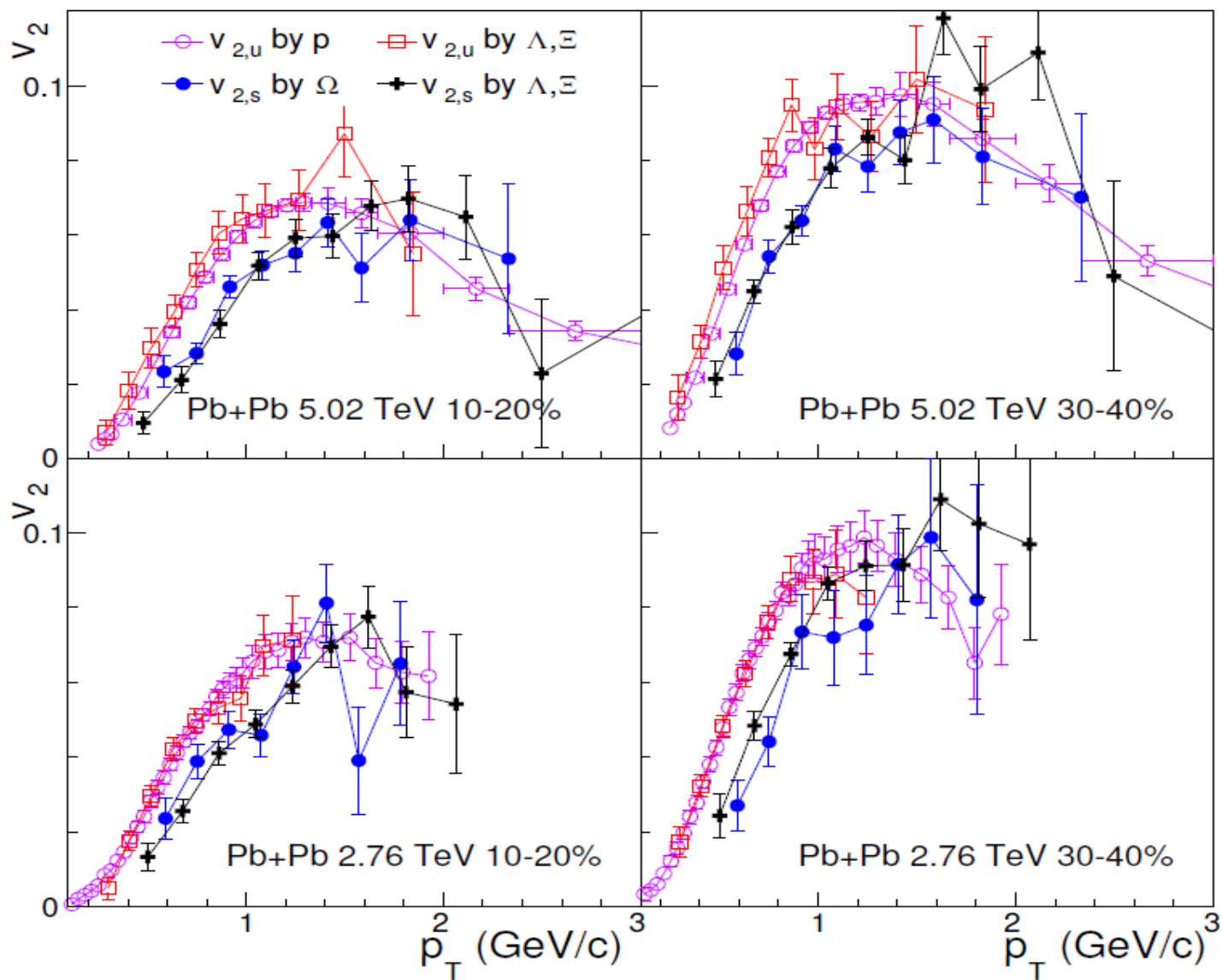
$$v_{2,s}(p_T) = \frac{1}{2} v_{2,\phi}(2p_T), \quad (21)$$

$$v_{2,s}(p_T) = \frac{1}{3} \left[ 2v_{2,\Xi} \left( \frac{1+2r}{r} p_T \right) - v_{2,\Lambda} \left( \frac{2+r}{r} p_T \right) \right]. \quad (22)$$

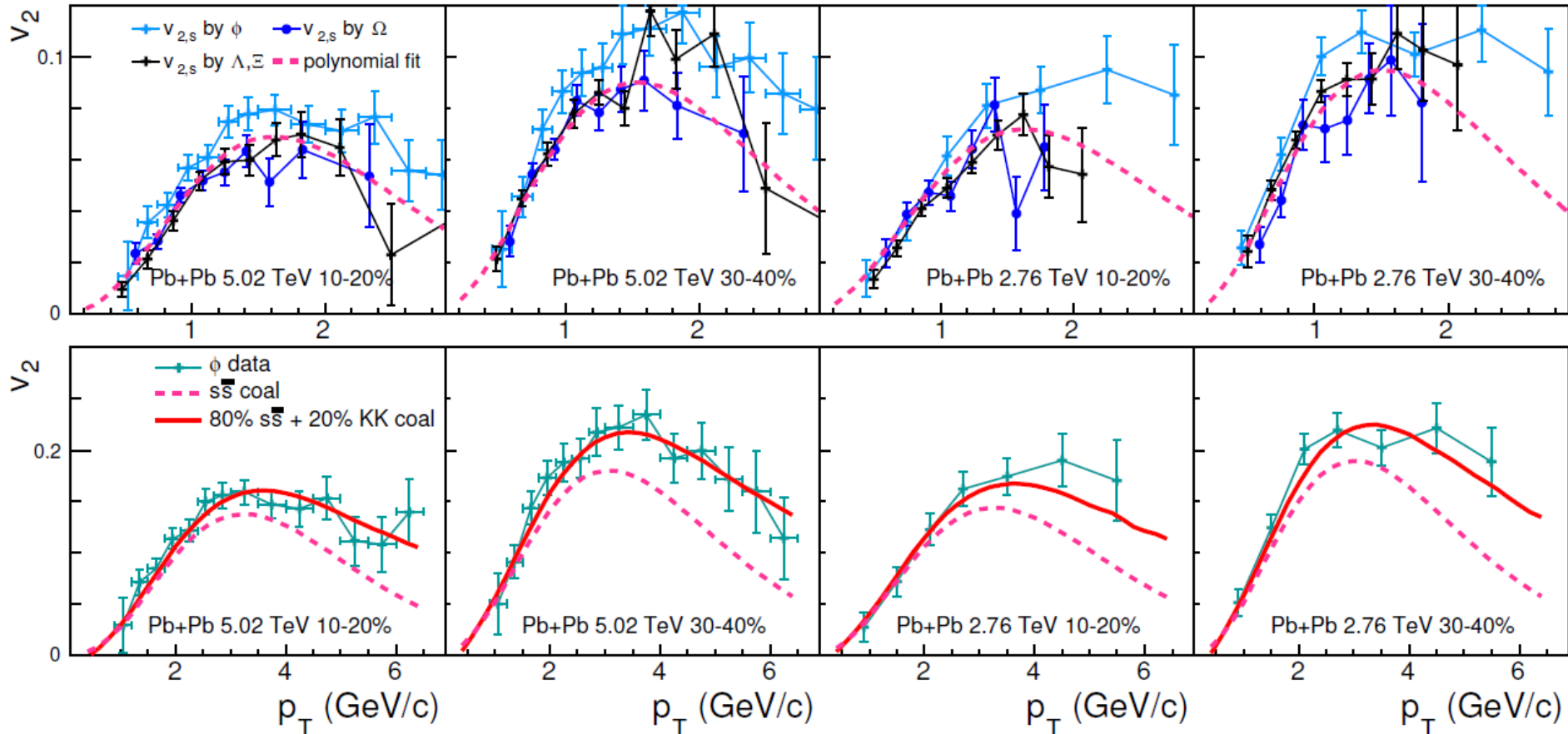
## AuAu 200 GeV for 30-80% centrality







# $\phi$ production at LHC



$$v_{2,\phi}^{(final)}(p_T) \approx 2 \left[ \underbrace{\left( (1-z) v_{2,s} \left( \frac{p_T}{2} \right) \right)}_{s\bar{s}} + \underbrace{z v_{2,K} \left( \frac{p_T}{2} \right)}_{KK \text{ coal}} \right]$$

## $v_2$ of charm quark

$D^0(c\bar{u})$  meson  $v_2$  in EVC QCM

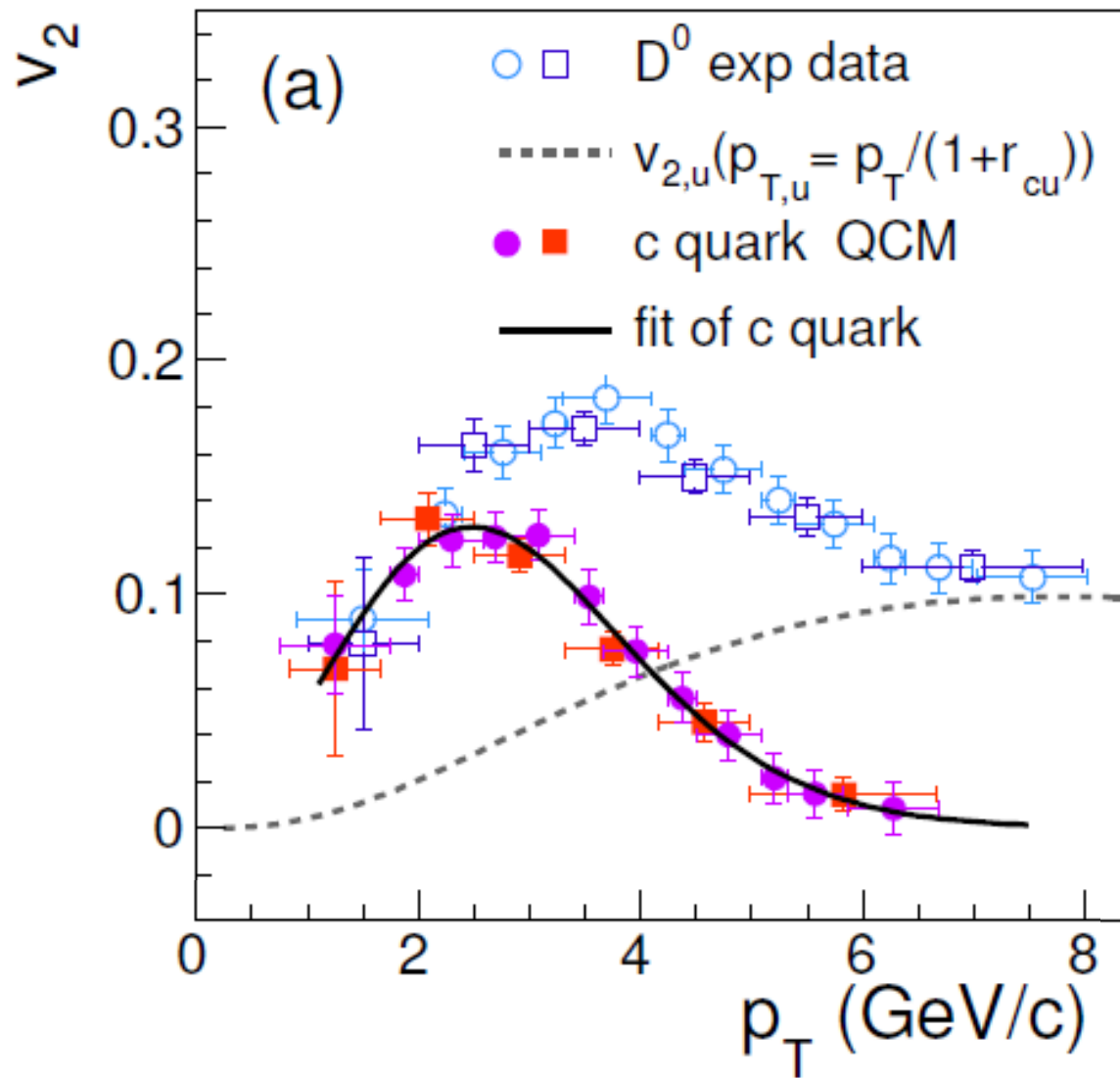
$$v_{2,D}(p_T) = v_{2,u} \left( \frac{1}{1+r_{cu}} p_T \right) + v_{2,c} \left( \frac{r_{cu}}{1+r_{cu}} p_T \right) \quad (27)$$

From  $D^0$  data, we get charm quark  $v_2$

$$v_{2,c}(p_T) = v_{2,D} \left( \frac{1+r_{cu}}{r_{cu}} p_T \right) - v_{2,u} \left( \frac{1}{r_{cu}} p_T \right). \quad (28)$$

where  $r_{cu} = m_c/m_u = 5$

# In Pb+Pb collisions at 5.02 TeV for 30-50% centrality



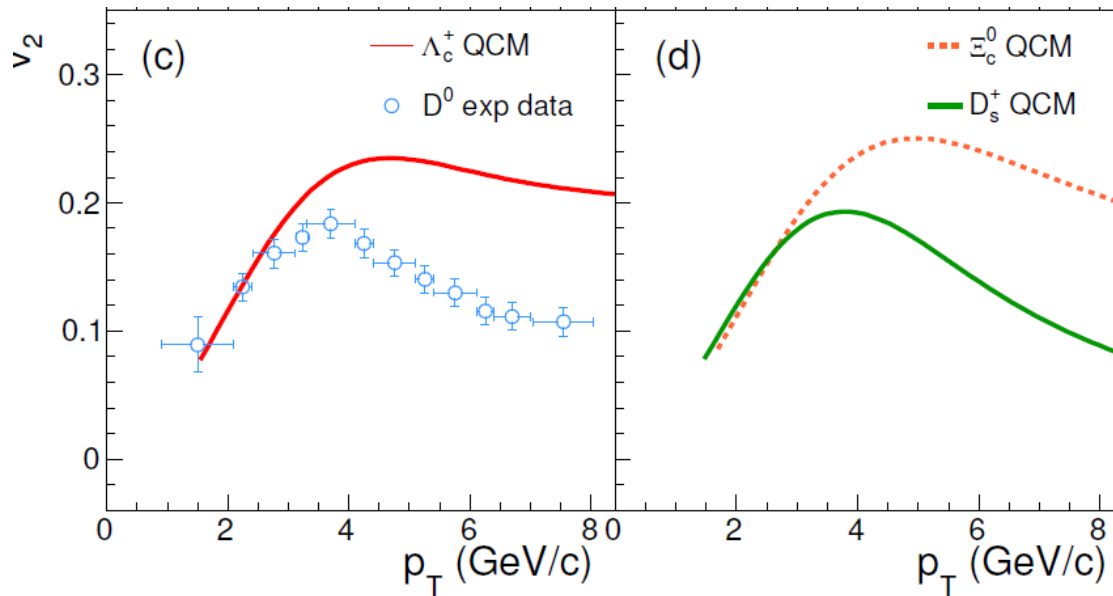
$$v_{2,D_s}(p_T) = v_{2,s} \left( \frac{1}{1+r_{cs}} p_T \right) + v_{2,c} \left( \frac{r_{cs}}{1+r_{cs}} p_T \right)$$

$$r_{cs} = m_c/m_s$$

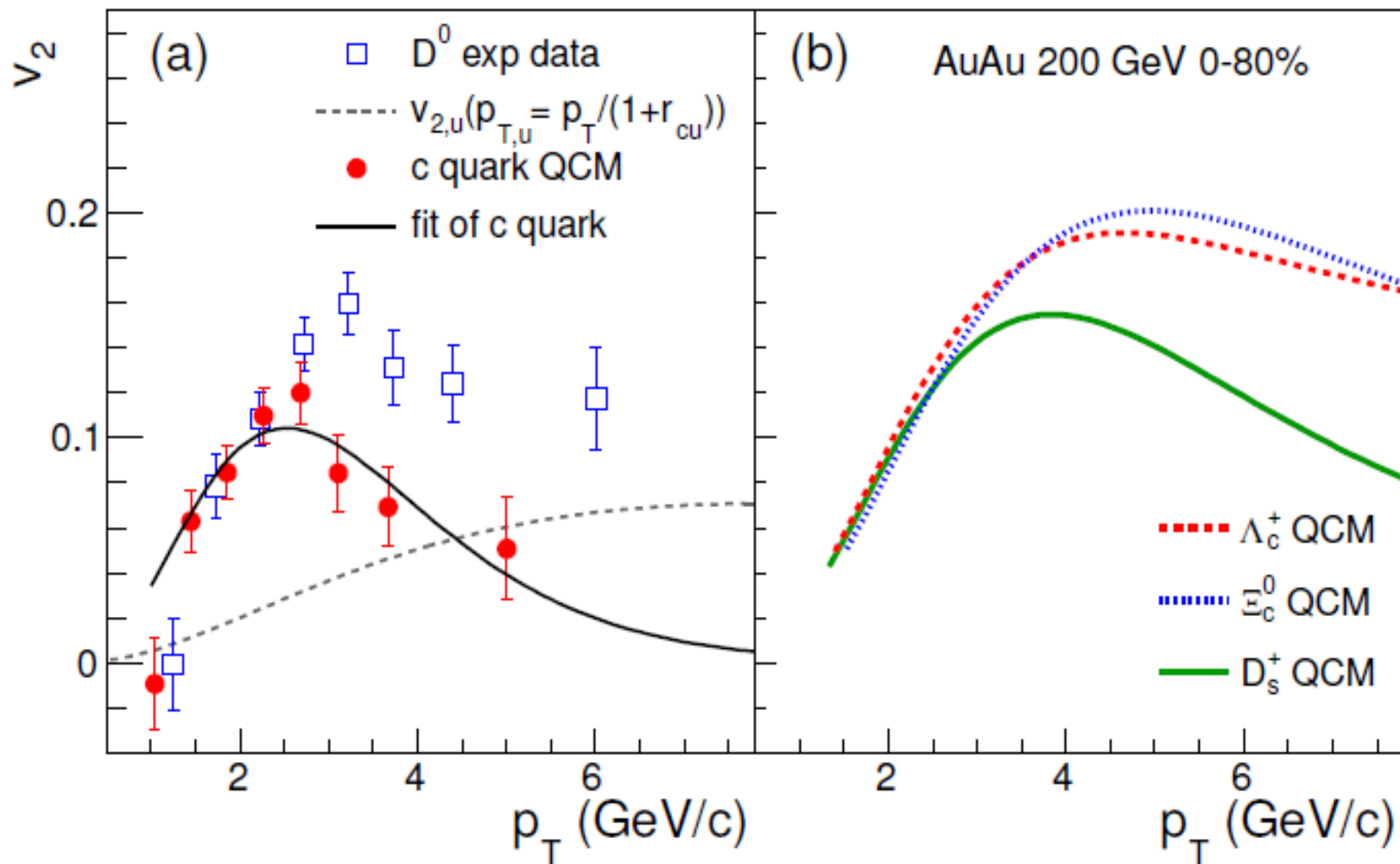
$$r_{cu} = m_c/m_u$$

$$v_{2,\Lambda_c}(p_T) = 2v_{2,u} \left( \frac{1}{2+r_{cu}} p_T \right) + v_{2,c} \left( \frac{r_{cu}}{2+r_{cu}} p_T \right)$$

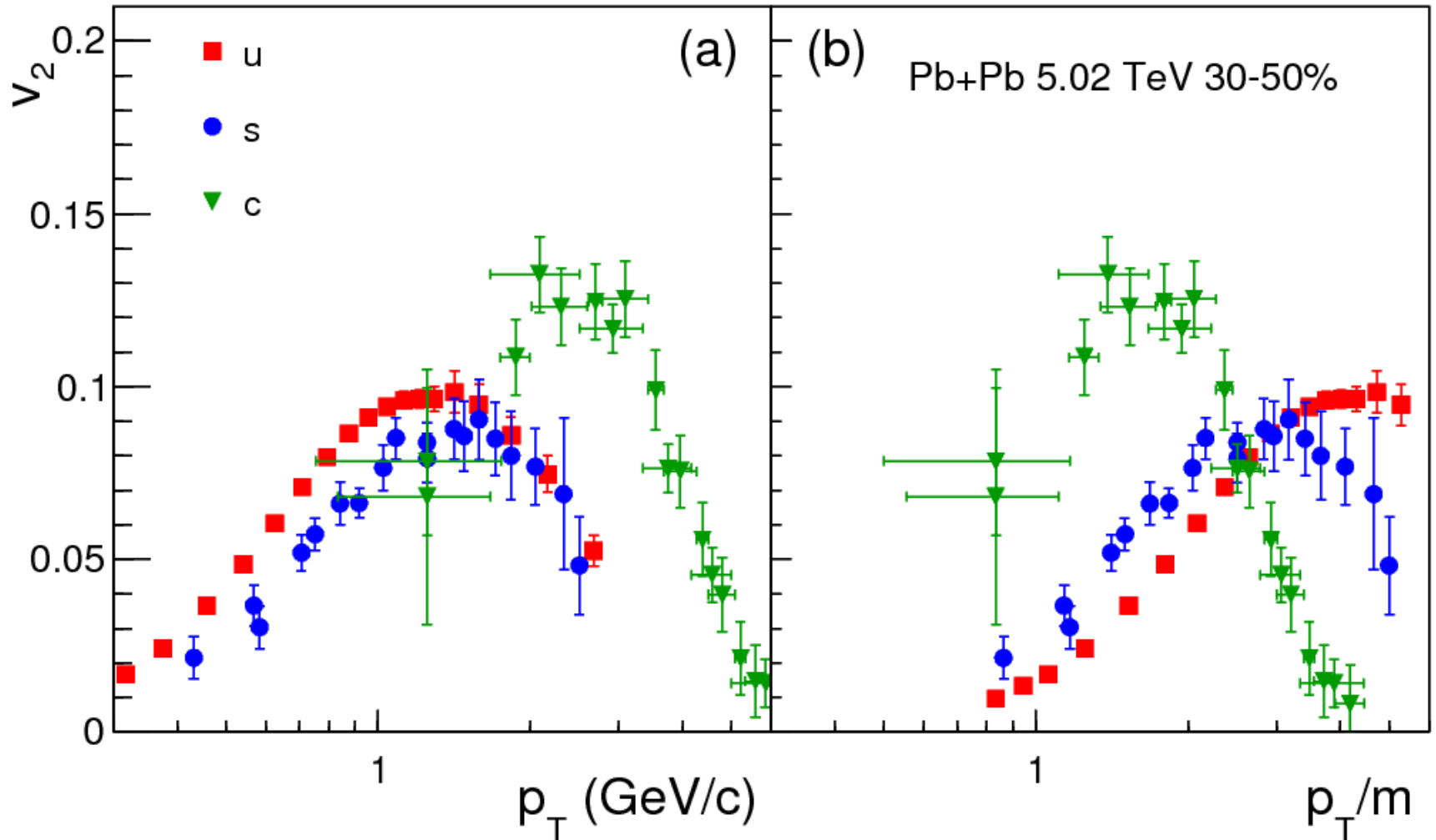
$$v_{2,\Xi_c}(p_T) = v_{2,u} \left( \frac{1}{1+r+r_{cu}} p_T \right) + v_{2,s} \left( \frac{r}{1+r+r_{cu}} p_T \right) + v_{2,c} \left( \frac{r_{cu}}{1+r+r_{cu}} p_T \right)$$







# Compare $v_2$ of up, strange and charm quarks



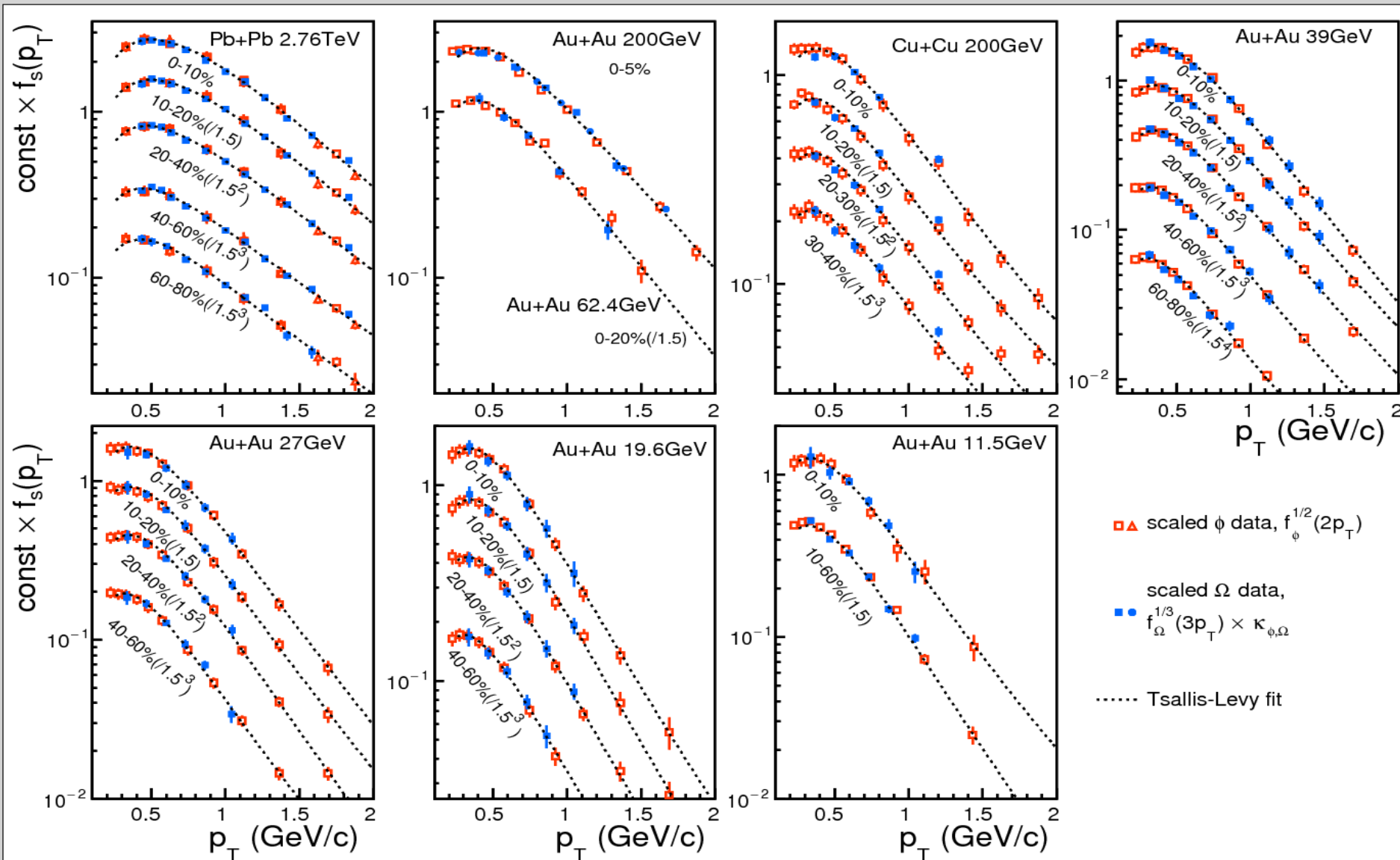
# 7.

## **Energy-scan study for hadronic $p_T$ spectra in Au+Au collisions at $\sqrt{s_{NN}} = 7.7 - 200 \text{ GeV}$**

Refs:

- Jun Song, xiao-feng Wang, Hai-hong Li, Rui-qin Wang, Feng-lan Shao, [arxiv:2007.14588v1](#).
- Jun Song, Feng-lan Shao, and Zuo-tang Liang, [Phys.RevC.102.014911\(2020\)](#).

# QNS in heavy-ion collisions

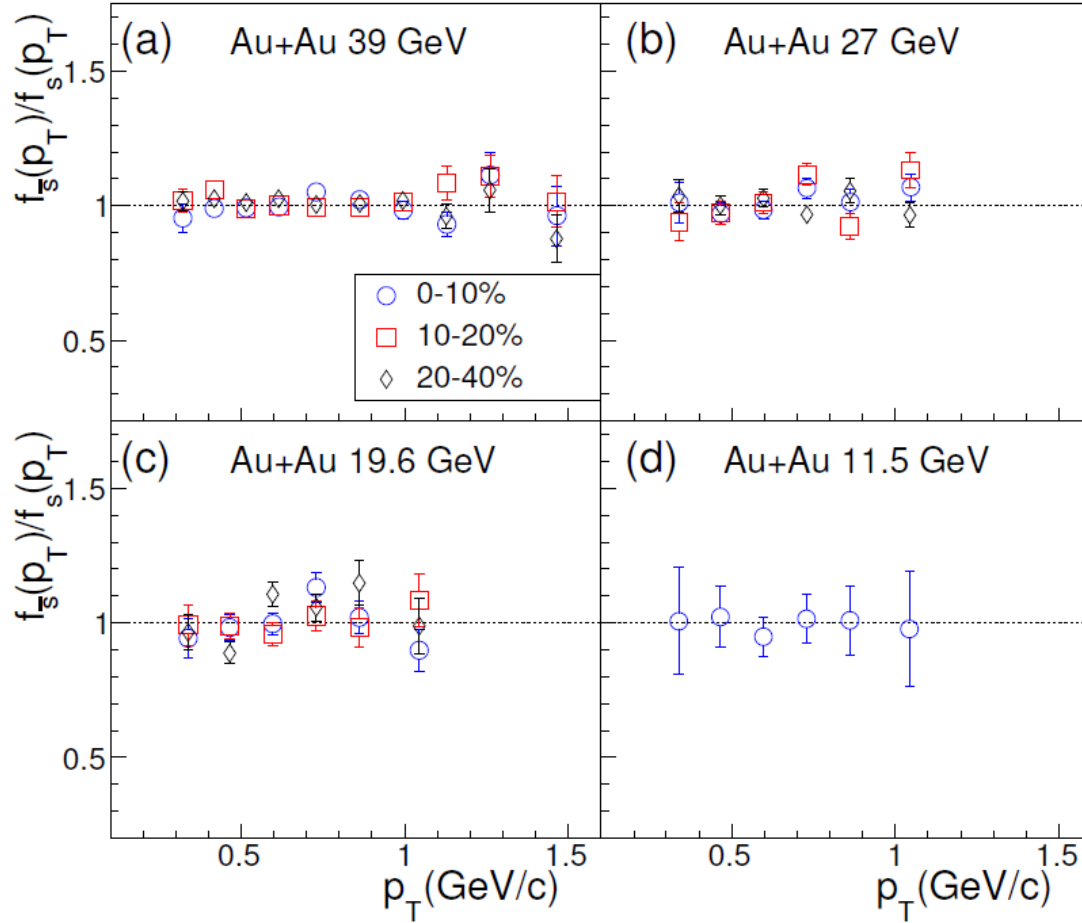


# Strangeness neutralization

$$\text{Asymmetry factor : } z_S = \frac{N_{\bar{S}} - N_S}{N_{\bar{S}} + N_S}$$

Table I. Strangeness asymmetry factor  $z_S$  calculated by yield data of strange hadrons and anti-hadrons in central Au+Au collisions [29, 41–45, 48].

$\sqrt{s_{NN}}$ (GeV)	$K^{+,-}$	$\Lambda (\bar{\Lambda})$	$\Xi^- (\bar{\Xi}^+)$	$\Omega^- (\bar{\Omega}^+)$	$\phi$	$z_S$																																																														
200	$48.9 \pm 6.3$	$16.7 \pm 1.1$	$2.17 \pm 0.2$	$0.53 \pm 0.06$	$7.95 \pm 0.11$	$-0.004 \pm 0.06$																																																														
	$45.7 \pm 5.2$	$12.7 \pm 0.9$	$1.83 \pm 0.2$				62.4	$37.6 \pm 2.7$	$15.7 \pm 2.3$	$1.63 \pm 0.2$	$0.212 \pm 0.028$	$3.52 \pm 0.08$	$-0.019 \pm 0.04$	$32.4 \pm 2.3$	$8.3 \pm 1.1$	$1.03 \pm 0.11$	$0.167 \pm 0.027$	39	$32.0 \pm 2.9$	$11.02 \pm 0.03$	$1.54 \pm 0.01$	$0.191 \pm 0.006$	$3.38 \pm 0.03$	$-0.002 \pm 0.05$	$25.0 \pm 2.3$	$3.82 \pm 0.01$	$0.78 \pm 0.01$	$0.139 \pm 0.004$	27	$31.1 \pm 2.8$	$11.67 \pm 0.04$	$1.57 \pm 0.01$	$0.154 \pm 0.008$	$3.01 \pm 0.04$	$-0.006 \pm 0.05$	$22.6 \pm 2.0$	$2.75 \pm 0.01$	$0.598 \pm 0.006$	$0.0972 \pm 0.0049$	19.6	$29.6 \pm 2.9$	$12.58 \pm 0.04$	$1.62 \pm 0.02$	$0.155 \pm 0.01$	$2.57 \pm 0.04$	$-0.0002 \pm 0.05$	$18.8 \pm 1.9$	$1.858 \pm 0.009$	$0.421 \pm 0.005$	$0.0811 \pm 0.0048$	11.5	$25.0 \pm 2.5$	$14.17 \pm 0.08$	$1.35 \pm 0.02$	$0.082 \pm 0.012$	$1.72 \pm 0.04$	$-0.004 \pm 0.05$	$12.3 \pm 1.2$	$0.659 \pm 0.009$	$0.169 \pm 0.004$	$0.0356 \pm 0.0052$	7.7	$20.8 \pm 1.7$	$15.3 \pm 0.11$	$1.19 \pm 0.03$	$0.0271 \pm 0.0048$	$1.21 \pm 0.06$	$-0.021 \pm 0.04$
62.4	$37.6 \pm 2.7$	$15.7 \pm 2.3$	$1.63 \pm 0.2$	$0.212 \pm 0.028$	$3.52 \pm 0.08$	$-0.019 \pm 0.04$																																																														
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$$\frac{f_{\bar{s}}(p_T)}{f_s(p_T)} = \kappa_{\bar{\Omega}, \Omega} \left[ \frac{f_{\bar{\Omega}}(3p_T)}{f_{\Omega}(3p_T)} \right]^{1/3}$$

We have

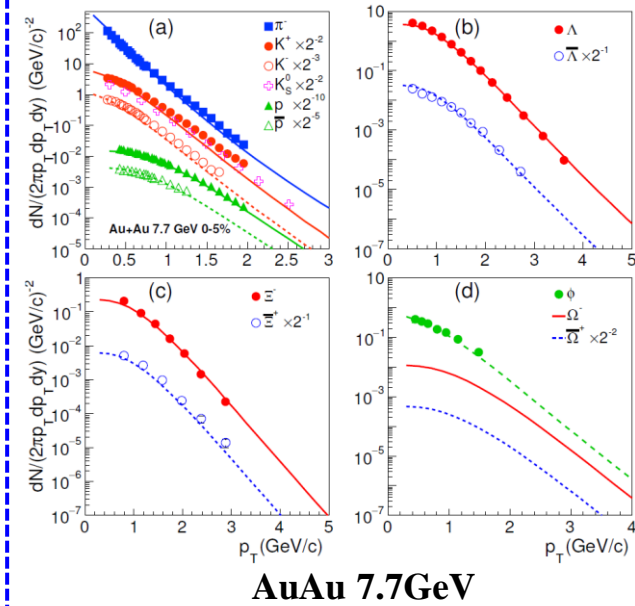
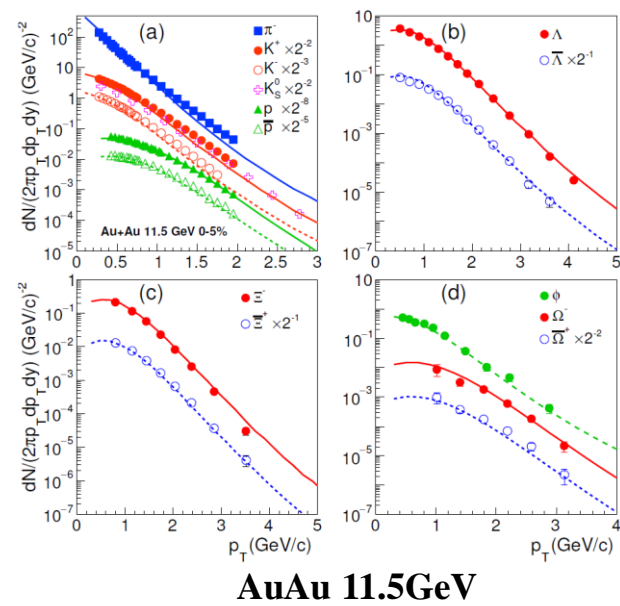
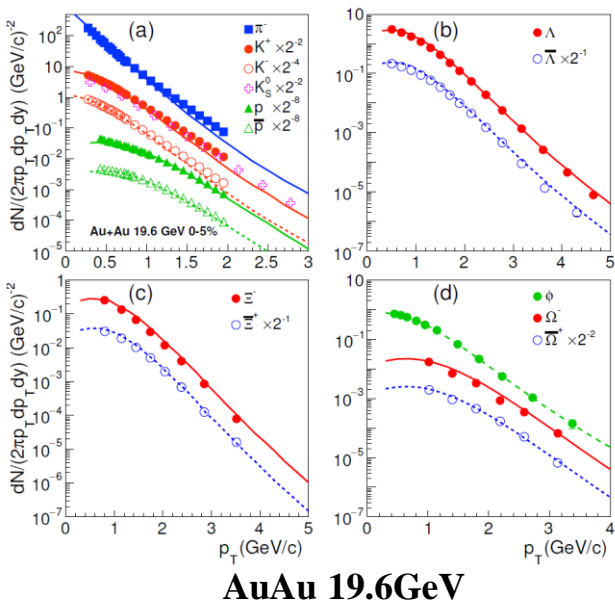
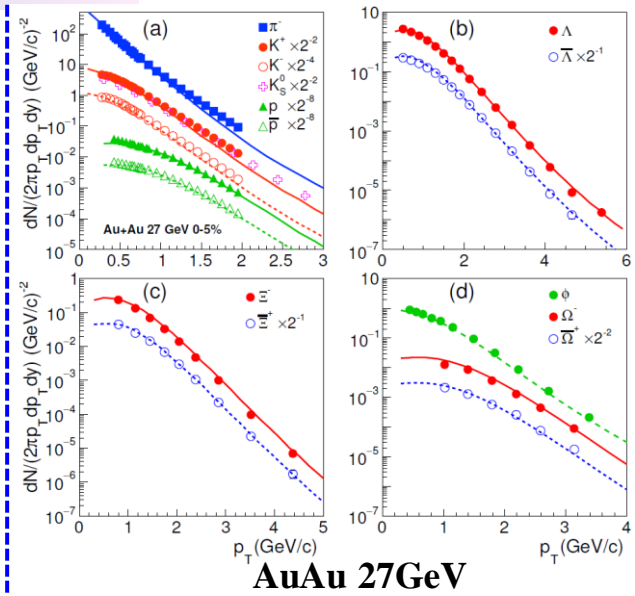
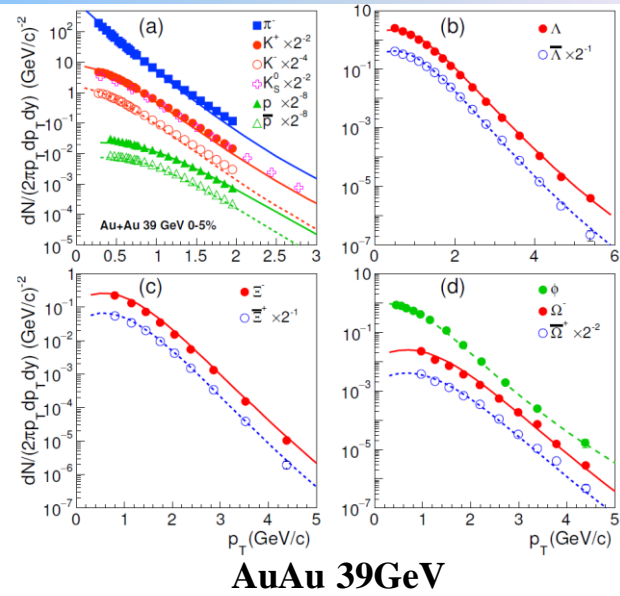
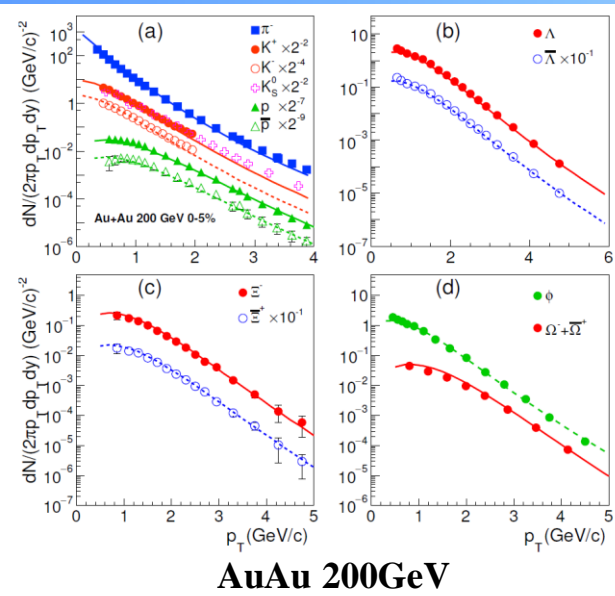
$$f_s(p_T) \approx f_{\bar{s}}(p_T)$$

at midrapidity in HIC

Figure 1. Ratio  $f_{\bar{s}}(p_T)/f_s(p_T)$  in Au+Au collisions at different collision energies obtained from experimental data of  $\Omega^-$  and  $\bar{\Omega}^+$  in central and semi-central collisions [29, 48] by Eq.

# $p_T$ spectra at STAR BES energies

Song, Shao, et al., arxiv:2007.14588v1



Three inputs  $f_u(p_T) \approx f_d(p_T)$ ,  $f_{\bar{u}}(p_T) \approx f_{\bar{d}}(p_T)$ ,  $f_s(p_T) \approx f_{\bar{s}}(p_T)$  at each energy.

# Summary

EVC is an effective mechanism.

- ✓ **quark number scaling for hadronic  $p_T$  spectra.**
- ✓ **self-consistent description for  $p_T$  spectra of light-flavor hadrons as well as single-charmed hadrons in pp and pPb collisions at LHC energies.**
- ✓  **$v_2$  of hadrons in heavy-ion collisions.**
- ✓ **energy-scan test of hadronic  $p_T$  spectra in Au+Au collisions at  $\sqrt{s_{NN}} = 7.7 - 200$  GeV.**

## Outlook

- **Systematic analysis on the obtained quark distribution functions at hadronization.**
- **The underlying physics of EVC.**

**Thanks**

