# Deciphering the nature of X(3872) in heavy ion collisions 

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Based on collaboration with Hui Zhang, Jinfeng Liao, Enke Wang, Qian Wang arXiv: 2004.00024, 2005.xxxxx

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## Outline

# - Introduction 

What is $\mathrm{X}(3872)$ ?

## $\uparrow X(3872)$ in heavy ion collisions

$\uparrow$ X(3872) in jet
$\uparrow$ Summary

## Exotic State XYZ

- Particle physics textbooks tell us that hadrons appear in two modes:
- mesons $(q \bar{q})$
- baryons (qqq)
- Many other types of color singlet compound hadrons, the so-called exotics, could exist


Glueball

tetraquark

pentaquark

X - unknown
Y - the vector exotic states $1^{--}$
$Z$ - charged quarkoniumlike states

## Exotic State X(3872)

- First observed by Belle collaboration (2003)

$$
B \rightarrow J / \psi \pi^{+} \pi^{-} K
$$

- Mass PDG 2012

$$
m_{X}=3871.68 \pm 0.17 M e V
$$

- Quantum numbers

CDF PRL 98, 132002(2007)
LHCb PRL 110, 222001 (2013)


$$
J^{P C}=1^{++}
$$

- Decay pattern PDG 2012
$J / \psi \rho\left(\pi^{+} \pi^{-}\right), J / \psi \omega\left(\pi^{+} \pi^{-} \pi^{0}\right), D^{0} \bar{D}^{* 0} / \bar{D}^{0} D^{* 0} / D \bar{D} \pi, J / \psi \gamma$



## Remaining mystery

## $\uparrow$ The internal structure of X(3872)

Figs from Yen-Jie Lee


Tetraquark
Hadronic molecule

Hybrid


Charmonium

No conclusive statement yet about the internal structure of X(3872).

## The inner structure of X(3872)

## - Loosely bound molecule state

- $\mathbf{x}(3872)$ is a loosely bound state of $D^{0} \bar{D}^{* 0} / \bar{D}^{0} D^{* 0}$
- The mass, quantum number and the large isospin violation can be understood naturally.
- The large production rate seems to be questionable Bignamini et al, PRL 09

$$
\sigma_{C D F}^{t h}<0.085 n b \quad \sigma_{C D F}^{e x}=3.1 \pm 0.7 n b
$$

- Rescattering effects may enhance the rate, if the upper bound of the relative momentum of the molecule state is as large as 3m_pi
Artoisenet and Braaten, PRD 10


## The inner structure of X(3872)

## $\uparrow$ Compact tetraquark state

- $\mathrm{X}(3872)$ is a compact four quark state
- A tetraquark system with two quarks arrange their color in a diquark before interacting with the antiquarks
- The mass, quantum number and the large isospin violation can be understood naturally.
- Stimulated the discovery of charged exotic states, e.g., Zc(3900)



## The inner structure of X(3872)

$\downarrow$ Quantum mixture of $\chi_{c 1}(2 p)-D^{0} \bar{D}^{*} 0$

- $\mathrm{X}(3872)$ is a mixed state of $\chi_{c 1}(2 p)$ and $D^{0} \bar{D}^{* 0} / \bar{D}^{0} D^{* 0}$

$$
|X\rangle=\sqrt{Z_{c \bar{c}}}\left|\chi_{c 1}(2 p)\right\rangle+\sqrt{Z_{m o l}}\left|D \bar{D}^{*}\right\rangle
$$

Meng, Gao and Chao PRD 87(2013)074035

- Different number of 'valence' quarks are superimposed
- Both the two components are substantial:
$\checkmark \chi_{c 1}$ component controls the short distance production
$\checkmark D \bar{D}^{*}$ components is mainly in charge of the hadronic decays


Meng, Han and Chao
PRD 96(2017)074014


Butenschoen, He and Kniehl PRL 123(2019)032001

## $X(3872)$ production in heavy ion collisions

$\uparrow X(3872)$ is usually studied in leptonic or hadronic collisions
$\uparrow \mathrm{HI}$ is very different with pp, which could provide a unique opportunity to explore the nature of $\mathrm{X}(3872)$


## $X(3872)$ production in heavy ion collisions

## $\downarrow$ Rich quark/gluon environment in HI




PRL 116 (2016) 222302

## X(3872) production in heavy ion collisions

## $\downarrow$ First experimental evidence of X(3872) in HI




Primary Vertex
$X(3872) \rightarrow J / \psi \pi^{+} \pi^{-}$was used for reconstruction.

## X(3872) production in heavy ion collisions

## $\downarrow$ Theoretical estimation of X(3872) in HI



Orders of magnitude difference indicates the advantage of HI in identifying the inner structure of $\mathbf{X ( 3 8 7 2 )}$.

ExHIC collaboration
PRL 106 (2011) 212001

## $X(3872)$ production in heavy ion collisions

-A "realistic" simulation by AMPT z-W. Lin etal, PRC 72 (2005) 064901


## X(3872) simulation by AMPT

## $\downarrow$ Calibration of the baseline




- AMPT does not have spin degrees of freedom, we distribute the yield into different spin state according to thermal model approximation

$$
R \equiv \frac{\mathrm{Yield}(\mathrm{~A})}{\mathrm{Yield}(\mathrm{~B})}=\exp \left(\frac{M_{B}-M_{A}}{T}\right)
$$

- 30\% from D* and 70\% from D
- 35\% for spin triplet, 65\% from spin singlet diquark


## X(3872) simulation by AMPT

- X(3872) coalescence


1. Coalescence of $D$ mesons
2. The average size:

$$
R_{D \bar{D}^{*}} \sim 5-7 \mathrm{fm}
$$

3. mass: $2 M_{D}<M_{X}<2 M_{D^{*}}$

Zhang, Liao, Wang, Wang, Xing arXiv: 2004.00024


1. Partonic coalescence of diquark and anti-diquark
2. The relative distance between diquark pairs $R_{[c q][\bar{c} \bar{q}]}<1 \mathrm{fm}$
3. mass:
$2 M_{|00\rangle_{0}}<M_{X}<2 M_{|11\rangle_{0}}$

## $X(3872)$ production in heavy ion collisions

## $\uparrow$ Total yields in 1M events

220k for hadronic molecule and 880 for compact tetraquark state.
$\uparrow \mathrm{P}_{\mathrm{T}}$ and rapidity dependence



Orders of magnitude difference between hadronic molecule and compact tetraquark scenarios, an unique opportunity for HI collisions.

## $X(3872)$ production in heavy ion collisions

## - Centrality dependence



- Strongly decreasing for hadronic molecule
- Mild change for compact tetraquark
- System size dependence could be a good probe to X(3872) inner structure.
more differential = more power


## $X(3872)$ production in heavy ion collisions

## $\uparrow$ Elliptic flow



- Elliptic flow is the key observable for collective property of bulk medium
- This is the first estimation of elliptic flow for exotic states
- Quark number scaling of tetraquark state?


## Puzzling result from CMS




- Energy loss leads to suppression in large pt
- Disassociation leads to suppression in large pt
- What caused enhancement in large pt? Strong coalescence/ regeneration?


## $\mathrm{X}(3872)$ in large pt

## $\downarrow$ Quantum mixture of $\chi_{c 1}(2 p)-D^{0} \bar{D}^{* 0}$

Butenschoen, He and Kniehl
PRL 123(2019)032001

$$
|X\rangle=\sqrt{Z_{c \bar{c}}}\left|\chi_{c 1}(2 p)\right\rangle+\sqrt{Z_{m o l}}\left|D \bar{D}^{*}\right\rangle
$$



- NRQCD

$$
\begin{aligned}
& d \sigma\left(p p \rightarrow \chi_{c 1}^{\prime}\right)=\sum_{n} d \hat{\sigma}\left((c \bar{c})_{n}\right) \frac{\left\langle\mathcal{O}_{n}^{\gamma_{c 1}^{\prime}}\right\rangle}{m_{c}^{2 L_{n}}}=\sum_{i, j, n} \int d x_{1} d x_{2} G_{i / p} G_{j / p} d \hat{\sigma}\left(i j \rightarrow(c \bar{c})_{n}\right)\left\langle\mathcal{O}_{n}^{\gamma_{c 1}^{\prime}}\right\rangle \\
& n={ }^{3} S_{1}^{8}, 3 P_{1}^{1}
\end{aligned}
$$

- LDMEs

|  | ${ }^{3} S_{1}^{8}\left(\mathrm{GeV}^{3}\right)$ | ${ }^{3} P_{1}^{1}\left(\mathrm{GeV}^{5}\right)$ |
| :---: | :---: | :---: |
| Kniehl | $0.83_{-0.16}^{+0.12} \times 10^{-4}$ | $0.34_{-0.15}^{+0.12} \times 10^{-2}$ |
| Chao | $0.87_{-0.51}^{+0.71} \times 10^{-4}$ | $0.75_{-0.32}^{+0.32} \times 10^{-3}$ |

Meng, Han and Chao
PRD 96(2017)074014

- Both loosely bound hadronic molecule and compact tetraquark state have problems to describe large $p_{T} \mathbf{X}$ (3872)
- Quantum mixture scenario is successful in large $p_{T}$ region, confirmed by two groups from NLO NRQCD, but with different LDMEs.


## X(3872) production in jet

## $\downarrow$ Jet substructure



## Jet substructure

## $\uparrow$ Light hadron production in jet

Xing et al., JHEP (2016) Kang et al., JHEP

$$
F\left(z_{h}, p_{T}\right)=\frac{d \sigma^{h}}{d p_{T} d \eta d z_{h}} / \frac{d \sigma^{h}}{d p_{T} d \eta}
$$





NLO + LL can describe the light hadron data very well.

## Jet substructure

## $\uparrow$ Open heavy flavor production in jet



Anderle et al., PRD (2017)


NLO + LL failed to describe the open heavy flavor data, eventually leads to new FFs global fit.

## Jet substructure

## $\checkmark$ Heavy quarkonium production in jet

TABLE I. $\quad J / \psi$ NRQCD LDMEs from four different groups.

|  | $\left.\mathcal{O}\left({ }^{3} S_{1}^{[1]}\right)\right\rangle$ <br> $\mathrm{GeV}^{3}$ | $\left\langle\mathcal{O}\left({ }^{1} S_{0}^{[8]}\right)\right\rangle$ <br> $10^{-2} \mathrm{GeV}^{3}$ | $\left\langle\mathcal{O}\left({ }^{3} S_{1}^{[8]}\right)\right\rangle$ <br> $10^{-2} \mathrm{GeV}^{3}$ | $\left\langle\mathcal{O}\left({ }^{3} 0^{-2} P_{0}^{[8]}\right)\right\rangle$ |
| :--- | :---: | :---: | :---: | :---: |
|  | $\mathrm{GeV}^{5}$ |  |  |  |
| Bodwin | $0^{\mathrm{a}}$ | 9.9 | 1.1 | 1.1 |
| Butenschoen | 1.32 | 3.04 | 0.16 | -0.91 |
| Chao | 1.16 | 8.9 | 0.30 | 1.26 |
| Gong | 1.16 | 9.7 | -0.46 | -2.14 |



- Both four sets of LDMEs can describe inclusive $J / \psi$ production in pp at high $p_{T}$.
- Significant difference in the prediction for JFFs.
- $J / \psi$ in jet is a sensitivity observable to probe the $J / \psi$ production mechanism.


## Jet substructure

## $\uparrow J / \psi$ polarization in jet

$$
\lambda_{F}\left(z_{h}, p_{T}\right)=\frac{F_{T}^{J / \psi}-F_{L}^{J / \psi}}{F_{T}^{J / \psi}+F_{L}^{J / \psi}}= \begin{cases}+1, & \text { Transverse } \\ -1, & \text { Longitudinal }\end{cases}
$$





## Jet substructure - X(3872)

## $\uparrow$ X(3872) production in jet




JFFs for $\mathrm{X}(3872)$ is a powerful observable to test the quantum mixture scenario.
$\uparrow \mathrm{HI}$ collisions provide a unique opportunity to differentiate hadronic molecule and compact tetraquark scenarios for $\mathrm{X}(3872)$.
$\checkmark$ X(3872) in jet is a rigorous observable to further test the picture of quantum mixture of $\chi_{c 1}(2 p)-D^{0} \bar{D}^{* 0}$.
$\checkmark$ Please stay tuned for further simulations in HI and precision pQCD (NRQCD) calculations for $X(3872)$ in high $p_{T}$.

## Thanks for your attention!

