

Constraints on the maximum mass of neutron stars with strangeness

Ang Li
Xiamen Univ.
liang@xmu.edu.cn

Miao & Li, [2107.07979](#)

Miao, Jiang, Li+ 2021 ApJL, [2107.13997](#)

Li+, 2021 ApJ, [2103.15119](#)

Li+, 2021 MNRAS, [2009.12571](#)

Miao, Li+, 2020 ApJ, [2006.00839](#)



Outline

- Intro. on neutron star (NS) and dense matter equation of state (EOS)
- NS maximum mass from LIGO/Virgo and NICER with microscopic EOSs
- Take-home message

NSs: Densest and smallest stars observed in the Universe

- Radius: $R \sim 10 - 15 \text{ km}$; Mass $M \sim 1 - 3 M_{\odot}$;
 - For $R = 10 \text{ km}$, $M = 1.4 M_{\odot}$, average number density $\sim 0.4 \text{ fm}^{-3}$; average (energy) density: $\sim 6.9 \times 10^{14} \text{ g/cm}^3 \sim (2-3)\rho_0$, **exceeding** the ordinary nuclear density;
 - **Equation of state (EOS)**, mainly $p(\epsilon)$, informative of the composition and inner structure of a NS;
- Such extreme conditions make it **impossible** to attain EOS by experimental methods only!

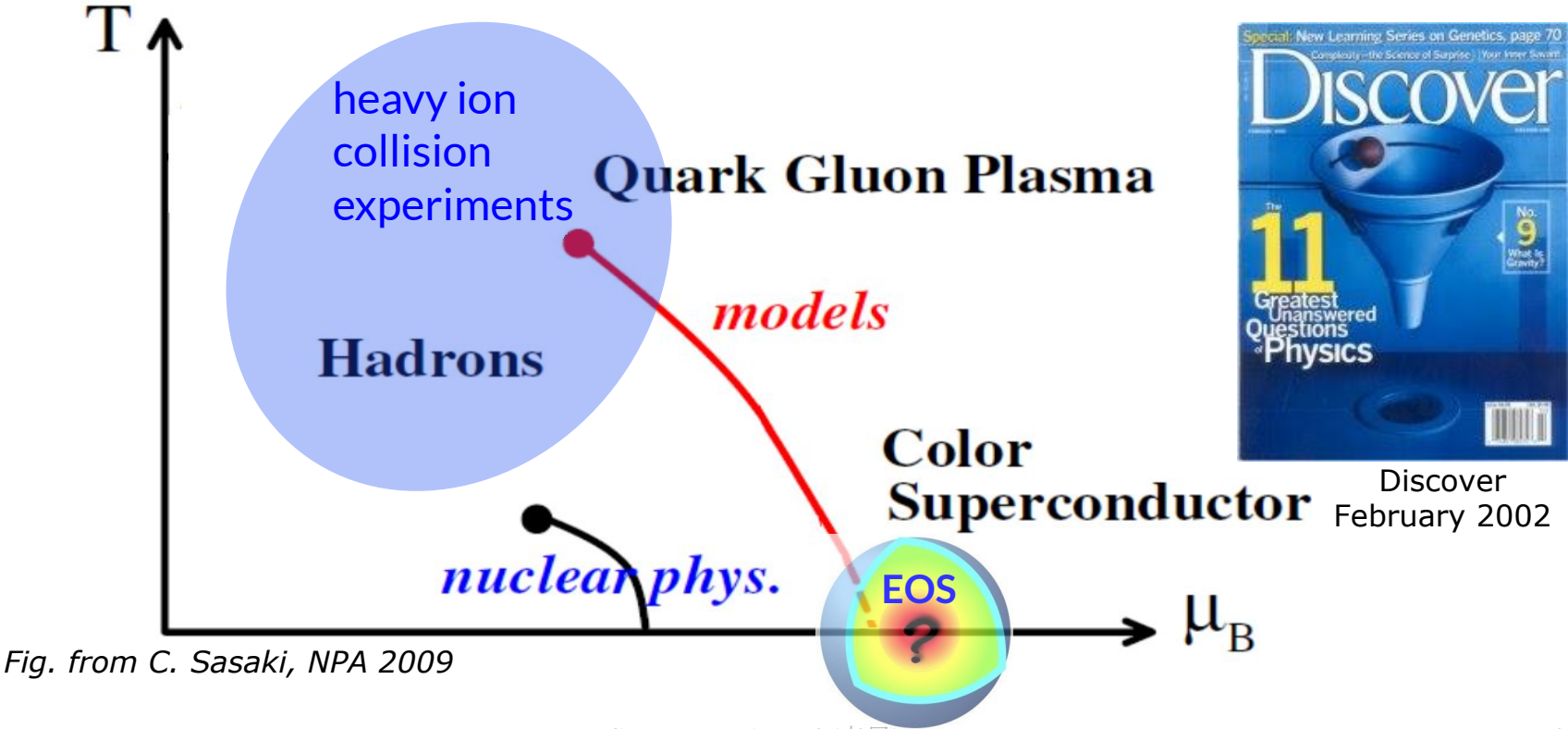


Fig. from C. Sasaki, NPA 2009

Discover
February 2002

NSs: Densest and smallest stars observed in the Universe

- Dense matter EOS: **One of 11 unanswered questions** of Physics whose resolutions could provide a new era in science!
- Electromagnetic (EM) and gravitational wave (GW) observations of NSs are ideal probe of dense QCD;
- **Macroscopic** properties (e.g., M, R) of NSs have an intrinsic correlation with **microphysical EOS**.

1. What is dark matter?
2. What is dark energy?
3. How were the heavy elements from iron to uranium made?
4. Do neutrinos have mass?
5. Where do ultra-energy particles come from?
6. Is a new theory of light and matter needed to explain what happens at very high energies and temperatures?
- 7. Are there new states of matter at ultrahigh temperatures and densities?**
8. Are protons unstable?
9. What is gravity?
10. Are there additional dimensions?
11. How did the Universe begin?



Discover
February 2002

NSs described in terms of General Relativity

EOS,
mainly $p(\varepsilon)$

+

$$\frac{dy}{dr} = -\frac{1}{r} [y^2 + yF(p, \varepsilon) + r^2 Q(p, \varepsilon)]$$

$$\frac{dp}{dr} = -\frac{(\varepsilon + p)(m + 4\pi r^3 p)}{r(r - 2m)}$$

$$\frac{dm}{dr} = 4\pi r^2 \varepsilon$$

Tolman-Oppenheimer-Volkoff
(TOV) equation:
GR version of the hydrostatic equilibrium

M_{TOV} ; M-R
relations

$$\begin{matrix} M \\ R \\ y_R \equiv y(R) \end{matrix}$$

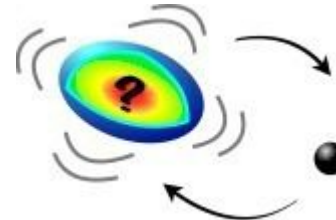
$$\begin{aligned} k_2 = & \frac{1}{20} \left(\frac{R_s}{R}\right)^5 \left(1 - \frac{R_s}{R}\right)^2 \left[2 - y_R + (y_R - 1) \frac{R_s}{R}\right] \left\{ \frac{R_s}{R} \left(6 - 3y_R + \frac{3R_s}{2R} (5y_R - 8)\right) \right. \\ & + \frac{1}{4} \left(\frac{R_s}{R}\right)^3 \left[26 - 22y_R + \frac{R_s}{R} (3y_R - 2) + \left(\frac{R_s}{R}\right)^2 (y_R + 1)\right] \\ & \left. + 3 \left(1 - \frac{R_s}{R}\right)^2 \left[2 - y_R + (y_R - 1) \frac{R_s}{R}\right] \ln \left(1 - \frac{R_s}{R}\right) \right\}^{-1} \end{aligned}$$

$$F(r) = \frac{r - 4\pi r^3 [\mathcal{E}(r) - P(r)]}{r - 2M(r)},$$

$$Q(r) = \frac{4\pi r \left[5\mathcal{E}(r) + 9P(r) + \frac{\mathcal{E}(r) + P(r)}{c_s^2} - \frac{6}{4\pi r^2}\right]}{r - 2M(r)}$$

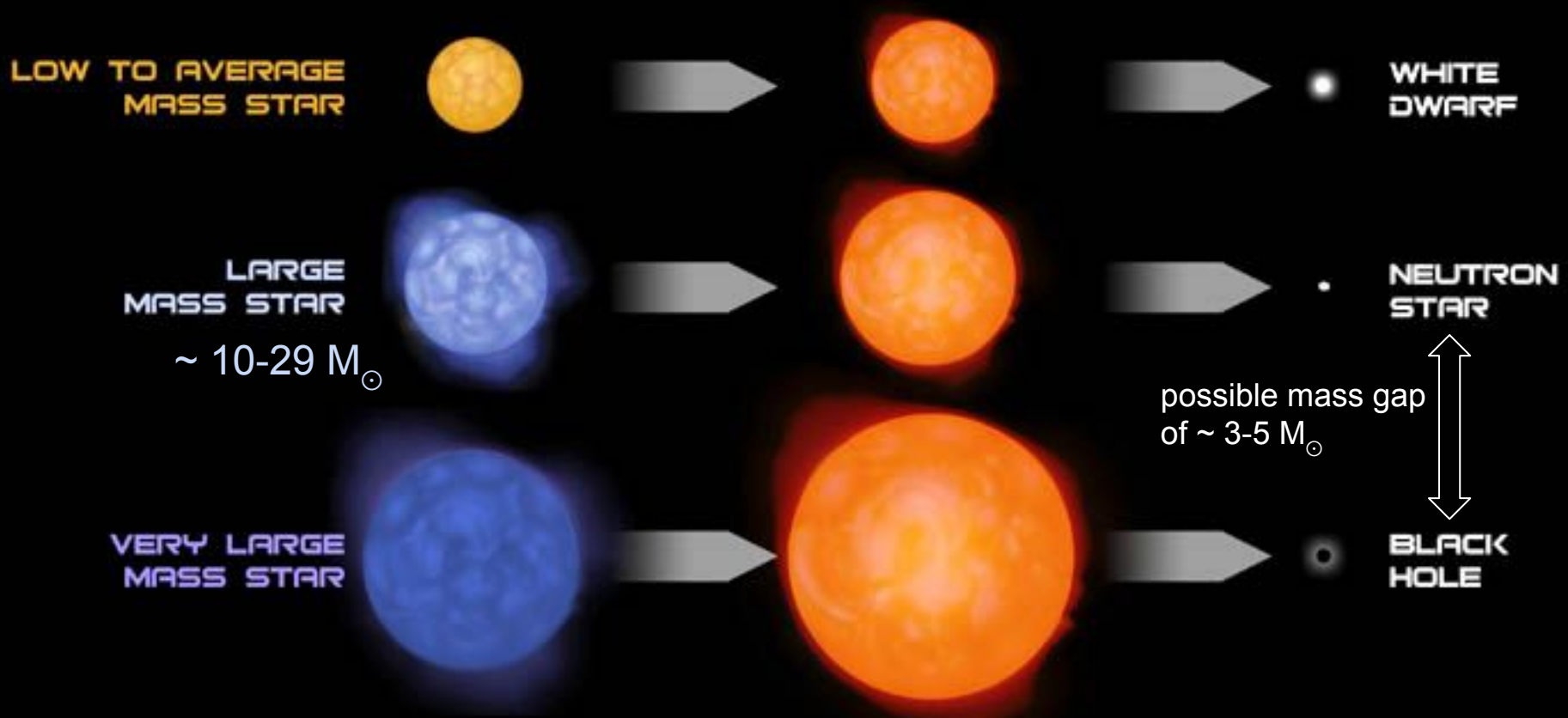
$$-4 \left\{ \frac{M(r) + 4\pi r^3 P(r)}{r[r - 2M(r)]} \right\}^2.$$

tidal
deformability

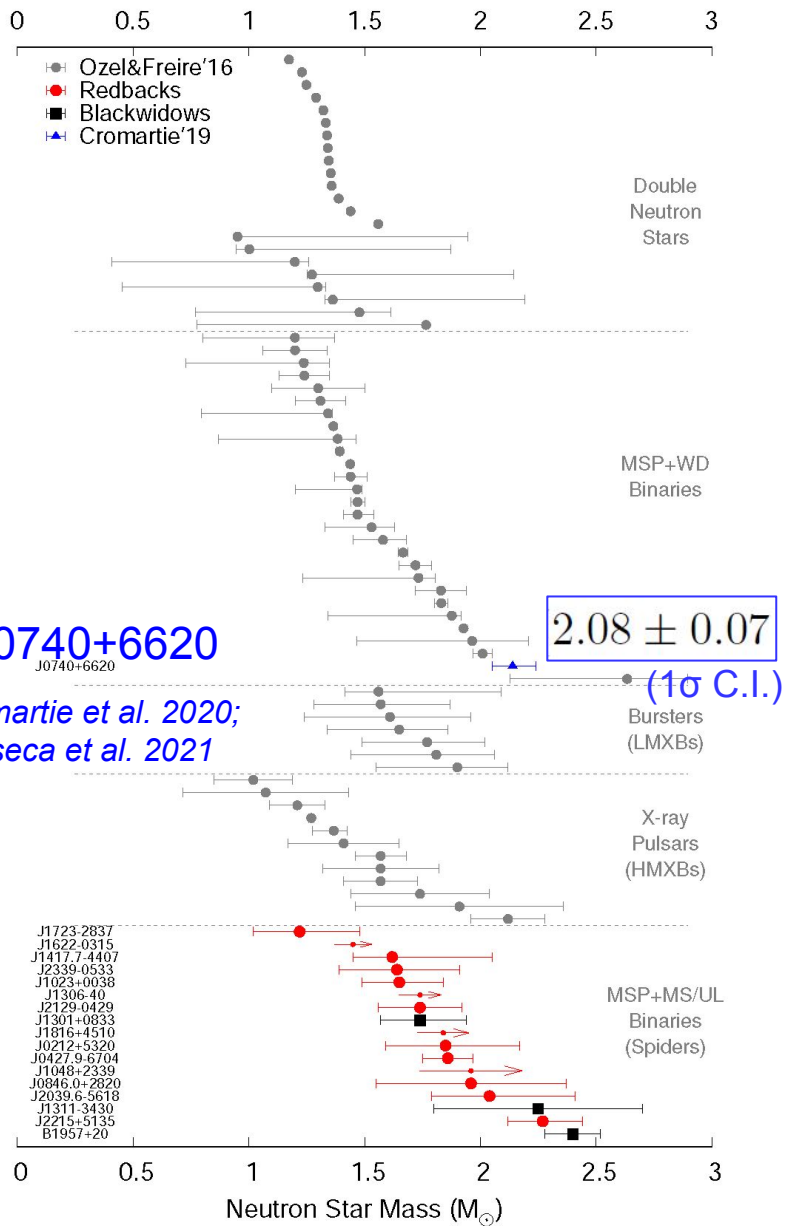


$$\lambda = \frac{2}{3} k_2 R^5$$

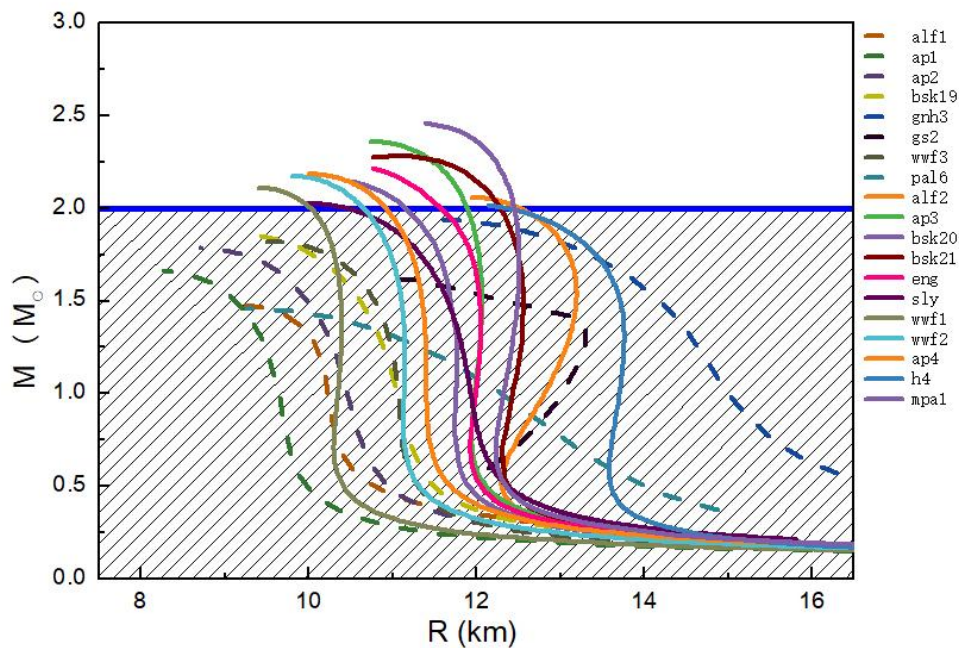
NS maximum mass and the mass gap between NS and BH



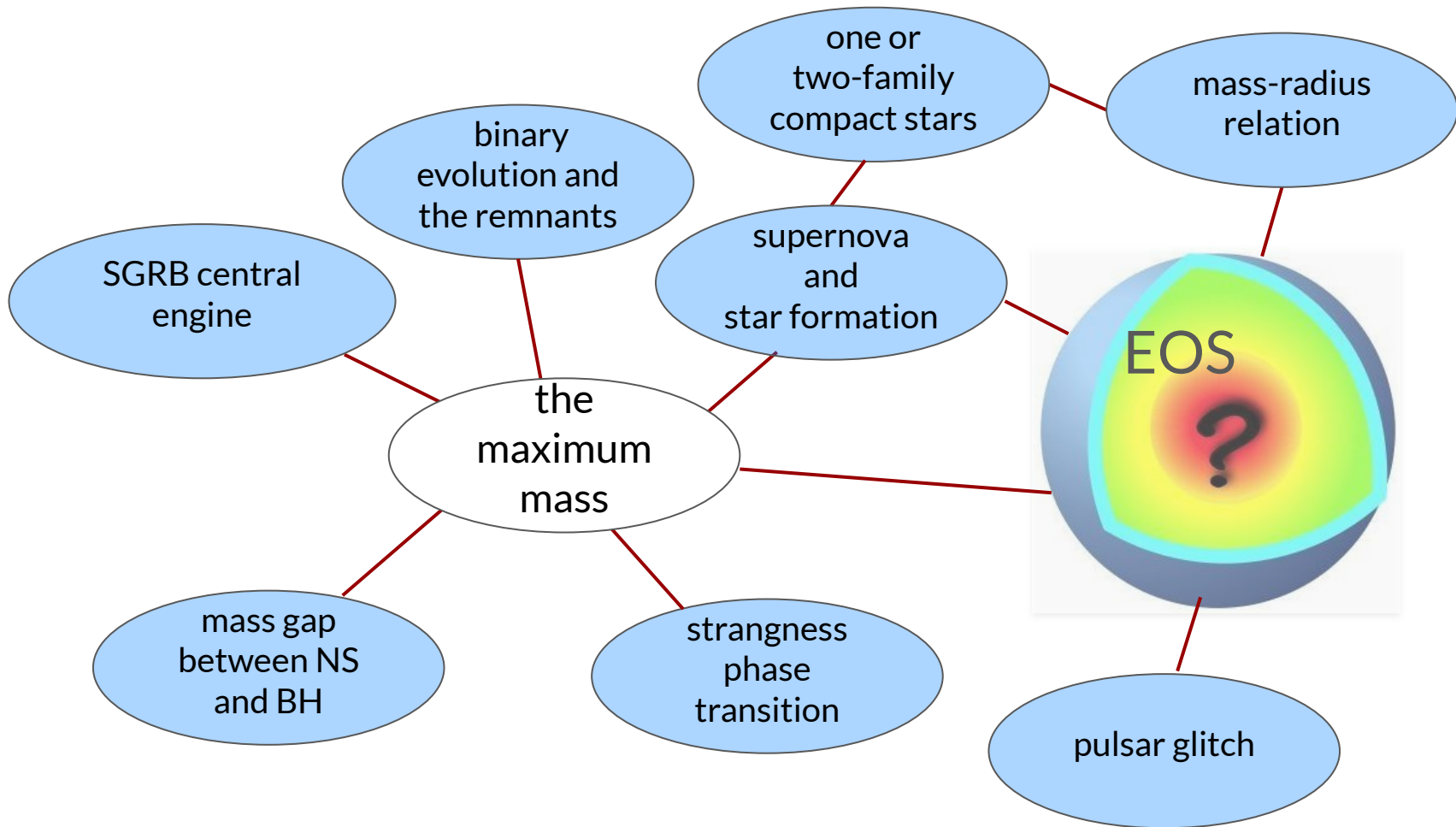
Neutron star maximum mass: Observational



$M_{\text{TOV}}(\text{EOS}) \geq$ all measured neutron star masses.



Why is understanding the NS EOS important?



EOS insensitive constraints on the maximum mass $\sim 3.2M_{\odot}$

A general treatment of M_{TOV} was given by Rhoades and Ruffini (1974), under the following set of assumptions:

1. General relativity is the correct theory of gravity. In particular, this means that the OV equation determines the equilibrium structure.
2. The equation of state satisfies the “microscopic stability” condition,

$$\frac{dP}{d\rho} \geq 0. \quad (9.5.1)$$

If this condition were violated, small elements of matter would spontaneously collapse.

3. The equation of state satisfies the causality condition

$$\frac{dP}{d\rho} \leq c^2; \quad (9.5.2)$$

that is, the speed of sound is less than the speed of light.

4. The equation of state below some “matching density” ρ_0 is known.

extreme causal EOS (1974)

the EOS is uncertain above a fiducial density of ρ_0 , the region in which they assumed that the equation of state is the **stiffest possible**, producing a sound velocity equal to the velocity of light,

$$P = P_0 + (\rho - \rho_0)c^2, \quad \rho \geq \rho_0.$$

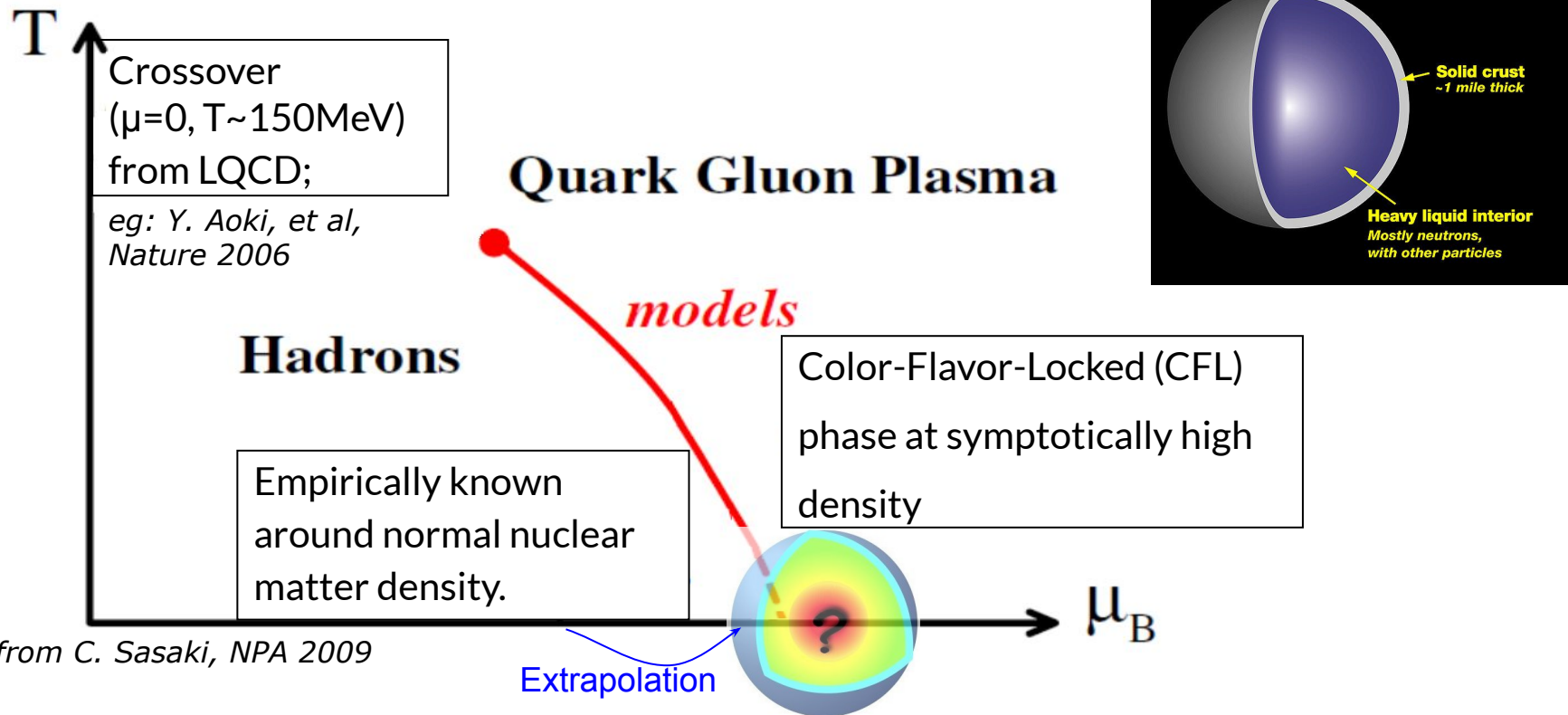
$$M_{\text{max}} \approx 3.2 \left(\frac{\rho_0}{4.6 \times 10^{14} \text{ g cm}^{-3}} \right)^{-1/2} M_{\odot}$$

Outline

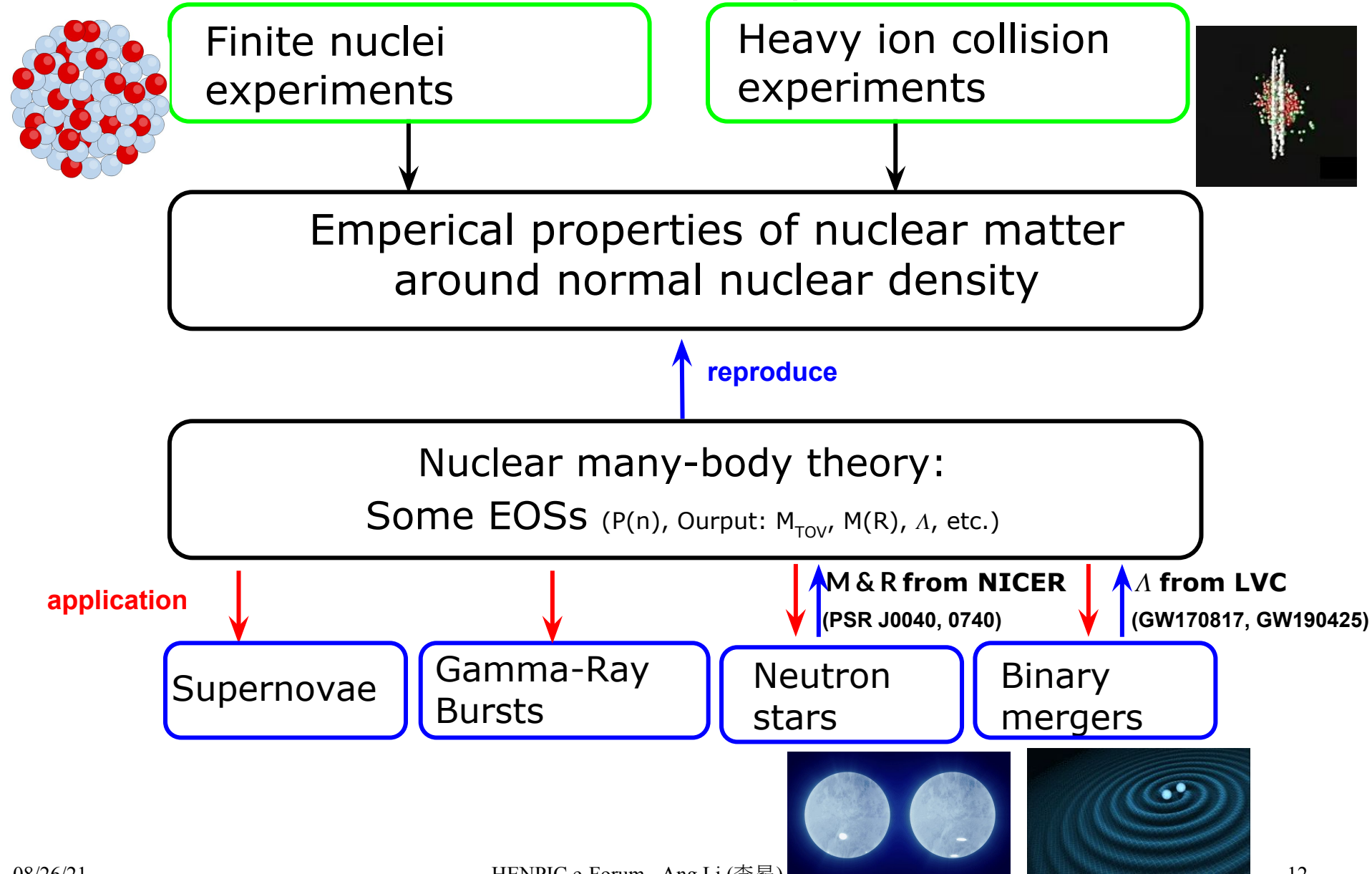
- Intro. on neutron star (NS) and dense matter equation of state (EOS)
- **NS maximum mass from LIGO/Virgo and NICER with microscopic EOSs**
- Take-home message

Maximum mass exact value is determined by the EOS

- M_{TOV} depends upon the strong **interaction** part of the EOS, e.g, Oppenheimer & Volko (1939) used an EOS for degenerate Fermi gas, obtaining a maximum mass of only $0.7 M_{\odot}$;
- The problem is to find the EOS in a regime where laboratory measurements of particle interactions are **inadequate** and the necessary theories of multi-body interactions are still **incomplete**.

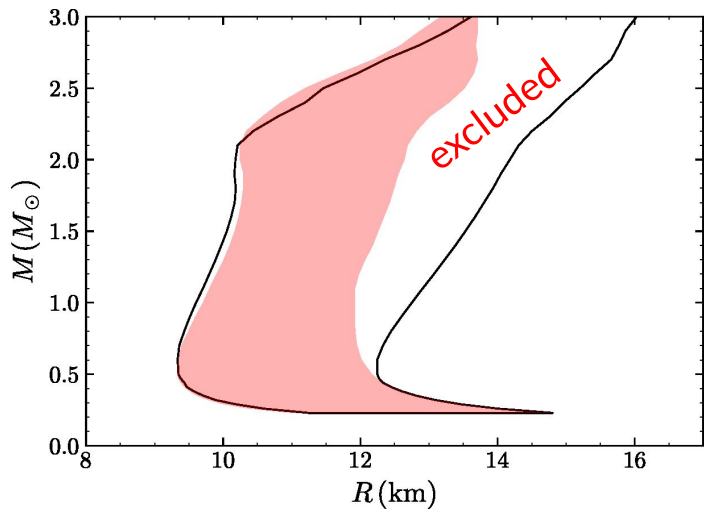


Connecting neutron star observations and nuclear experiment consistently

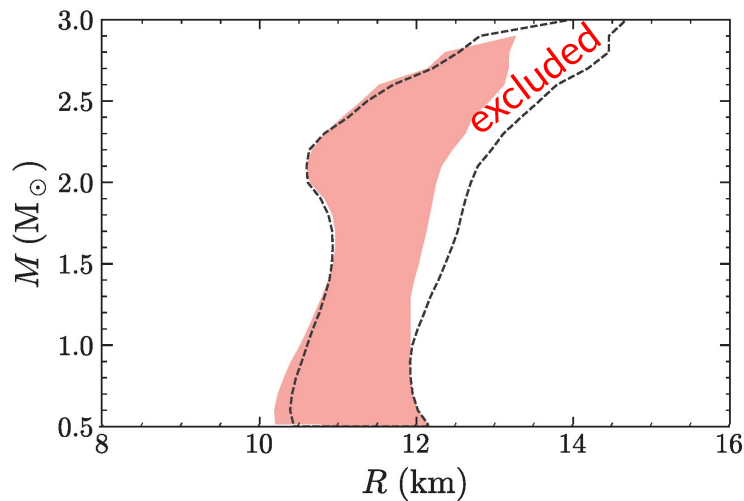


Important EOS constraints: Λ , I measurements

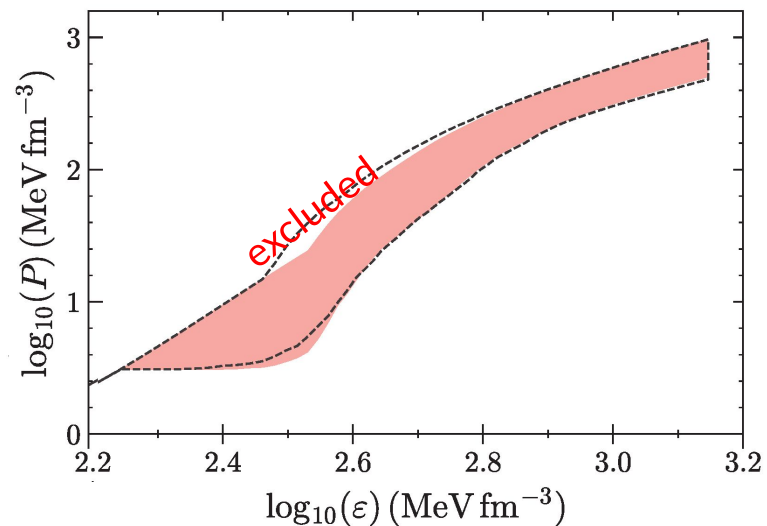
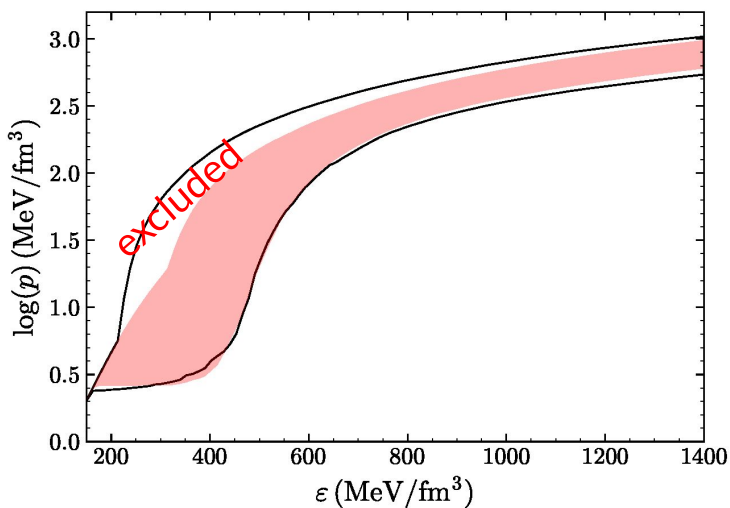
GW170817 Λ measurement



mocked J0737's I measurement
(from e.g., SKA)



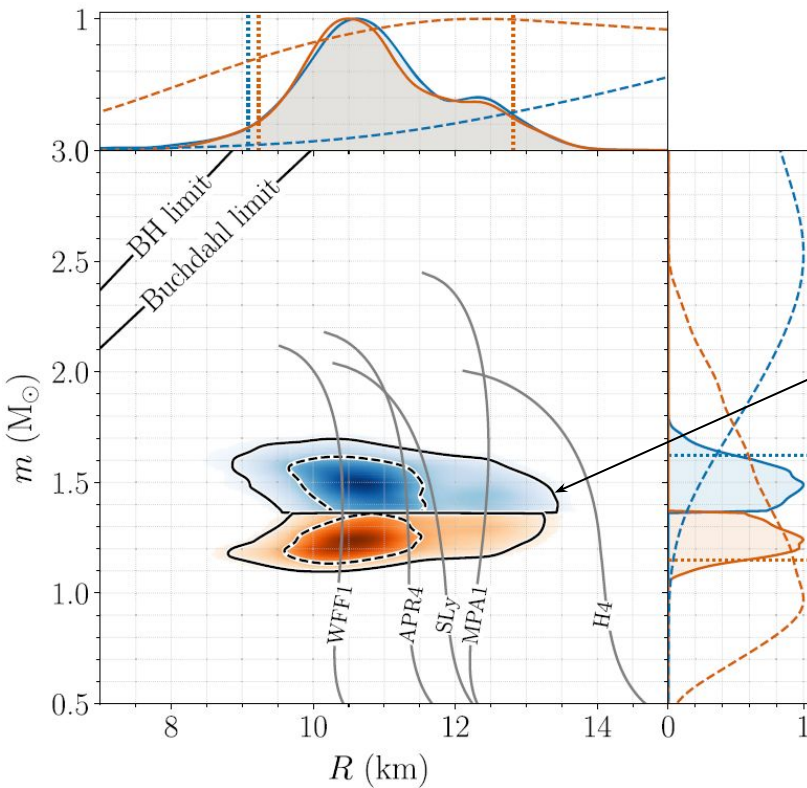
Ska



GW170817: Measurements of Neutron Star Radii and Equation of State

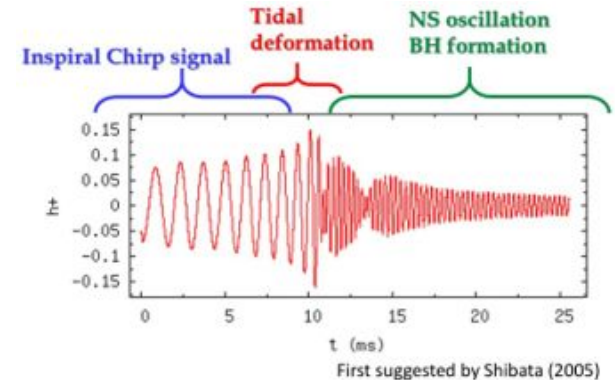
B. P. Abbott *et al.**

(The LIGO Scientific Collaboration and the Virgo Collaboration)



Our results are comparable and consistent with studies that use the tidal measurement from [5] to obtain bounds on NS radii. Using our bound of $\Lambda_{1.4} < 800$ (the only tidal parameter in [5], which assumed a common EOS for both NSs) and different EOS parametrizations, several studies found $R_{1.4} \lesssim 13.5$ km [56,58,62,64].

$$\Lambda = \frac{2}{3} k_2 \left(\frac{R}{M} \right)^5$$



[56] E.-P. Zhou, X. Zhou, and A. Li, *Phys. Rev. D* **97**, 083015 (2018).

(1.36, 1.60) M_{\odot}

(1.16, 1.36) M_{\odot}

Exemplary quark mean-field (QMF) NS EOS in the light of GW170817

THE ASTROPHYSICAL JOURNAL, 862:98 (9pp), 2018 August 1
 © 2018. The American Astronomical Society. All rights reserved.

<https://doi.org/10.3847/1538-4357/aacc28>

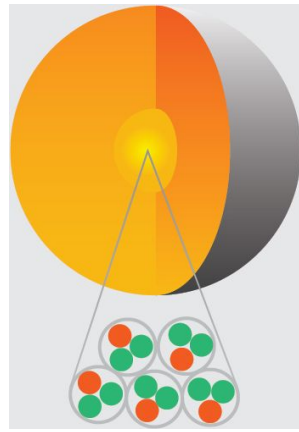


Neutron Star Equation of State from the Quark Level in Light of GW170817

Zhen-Yu Zhu¹, En-Ping Zhou², and Ang Li¹

¹Department of Astronomy, Xiamen University, Xiamen, Fujian 361005, People's Republic of China; liang@xmu.edu.cn

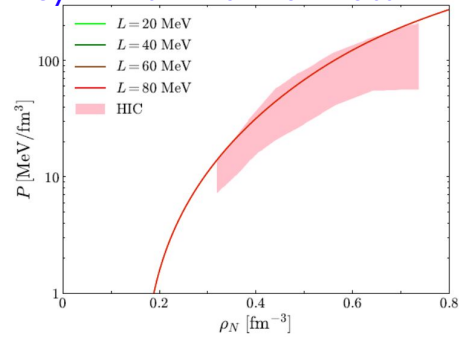
²State Key Laboratory of Nuclear Science and Technology and School of Physics, Peking University, Beijing 100871, People's Republic of China



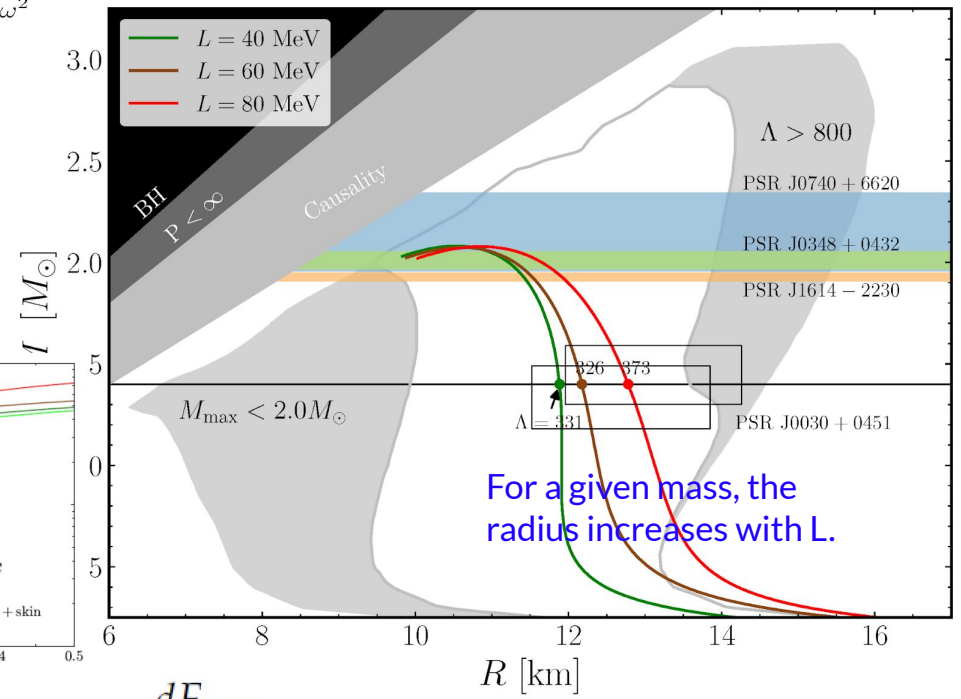
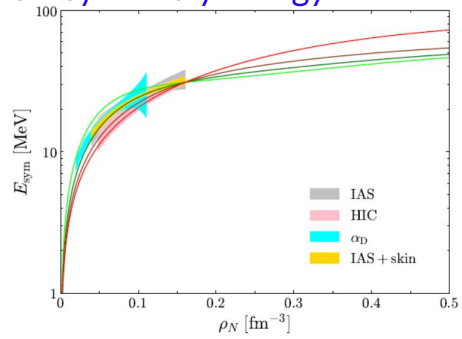
$$\begin{aligned} \mathcal{L} = & \bar{\psi} (i\gamma_{\mu}\partial^{\mu} - M_N^* - g_{\omega N}\omega\gamma^0 - g_{\rho N}\rho\tau_3\gamma^0) \psi \\ & - \frac{1}{2}(\nabla\sigma)^2 - \frac{1}{2}m_{\sigma}^2\sigma^2 - \frac{1}{3}g_2\sigma^3 - \frac{1}{4}g_3\sigma^4 \\ & + \frac{1}{2}(\nabla\rho)^2 + \frac{1}{2}m_{\rho}^2\rho^2 + \frac{1}{2}(\nabla\omega)^2 + \frac{1}{2}m_{\omega}^2\omega^2 \\ & + \frac{1}{2}g_{\rho N}^2\rho^2\Lambda_v g_{\omega N}^2\omega^2 \\ & [\gamma^0(\epsilon_q - g_{\omega q}\omega - \tau_{3q}g_{\rho q}\rho) \\ & - \vec{\gamma} \cdot \vec{p} - (m_q - g_{\sigma q}\sigma) \\ & - U(r)]\psi_q(\vec{r}) = 0 \end{aligned}$$

Microscopic NS EOS from the quark level, connecting **consistently** nuclear experiments and GW+EM observations.

symmetric nuclear matter EOS



symmetry energy



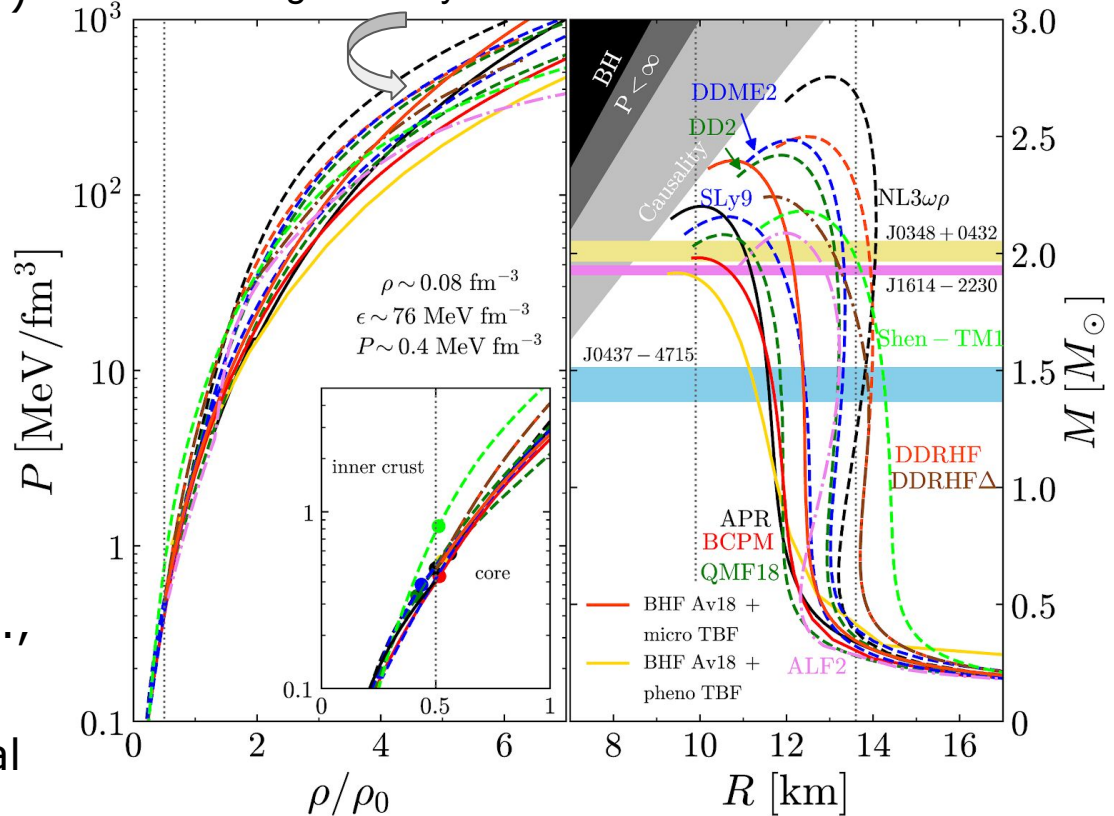
08/26/21

$$E(\rho, \delta) = E_{\text{SNM}}(\rho) + E_{\text{sym}}(\rho)\delta^2, \quad L \equiv 3\rho_0 \frac{dE_{\text{sym}}}{d\rho}(\rho_0)$$

Many nuclear many-body models can be employed

- Green's Function Monte Carlo
- Chiral Perturbation Theory (ChPT)
- Variational Many-Body (VMB; e.g., APR)
- V_{lowk} + Renormalization Group
- Brueckner-Hartree-Fock (BHF)
- Dirac-Brueckner-Hartree-Fock (DBHF)
- Quark mean-field (QMF)
- Quark Meson Coupling (QMC)
- Relativistic mean-field (RMF; e.g., DD2, NL3, TM1)
- Skyrme energy density functional (e.g., BSk20, Sly)...

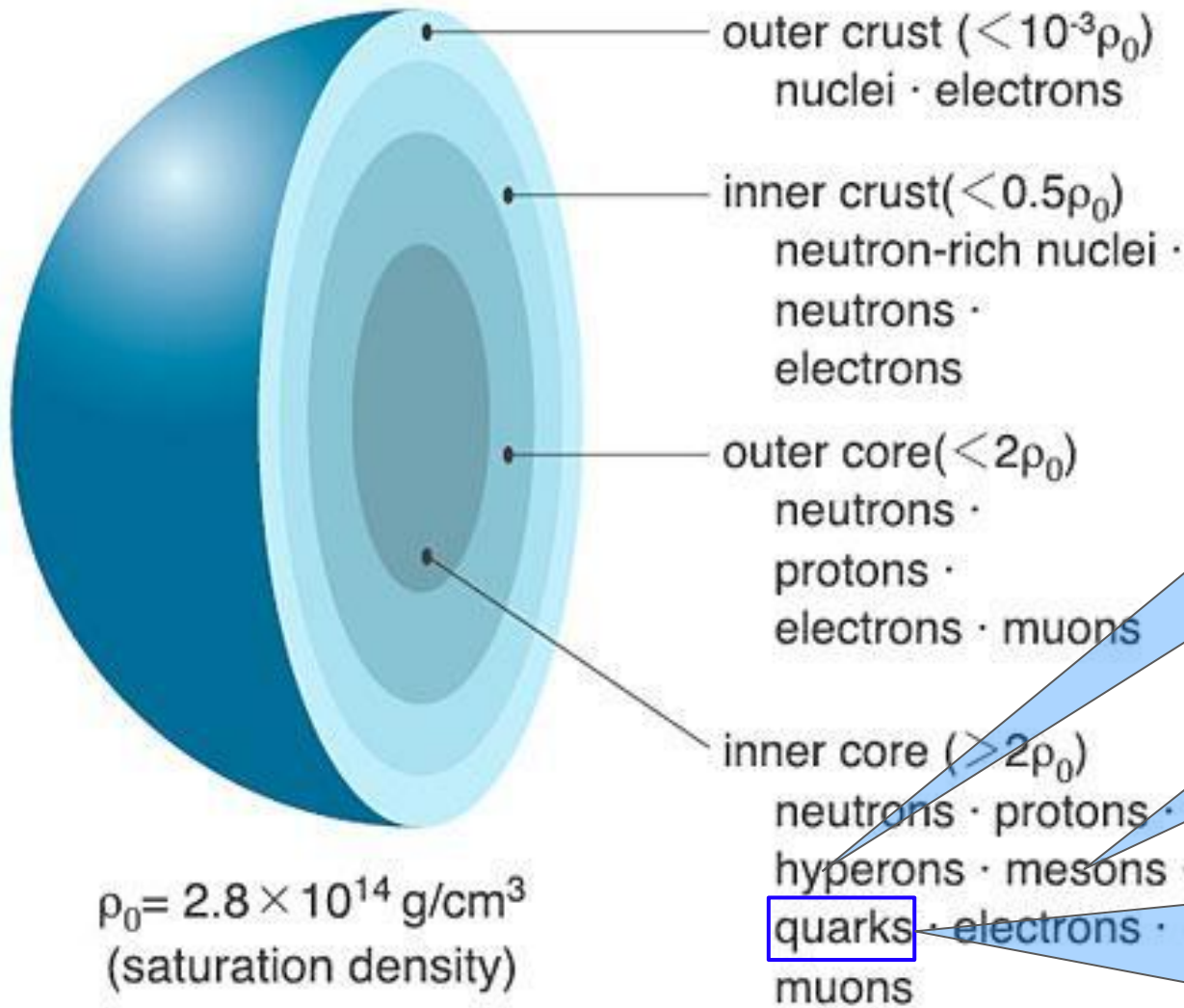
a fan of **different** prediction on high-density EOS



《致密物质状态方程：中子星与奇异星》李昂等，2019 原子核物理评论

<http://www.npr.ac.cn/article/doi/10.11804/NuclPhysRev.36.01.001>

Hyperon puzzle; hyperon/kaon/quark competition



《Neutron star equation of state: Modeling and applications》

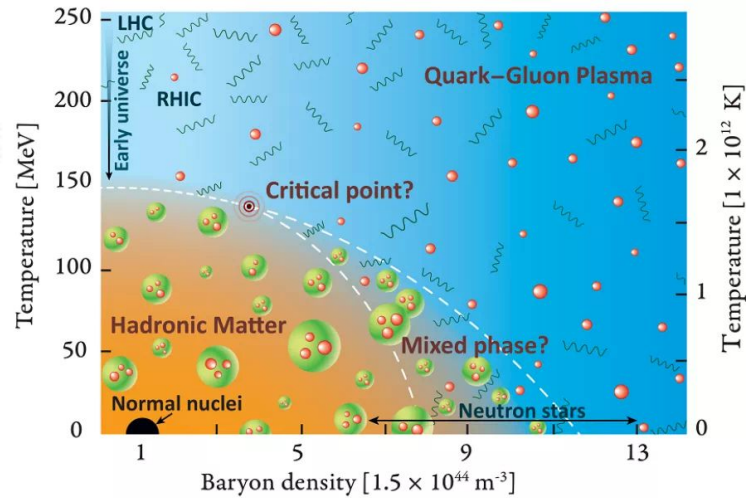
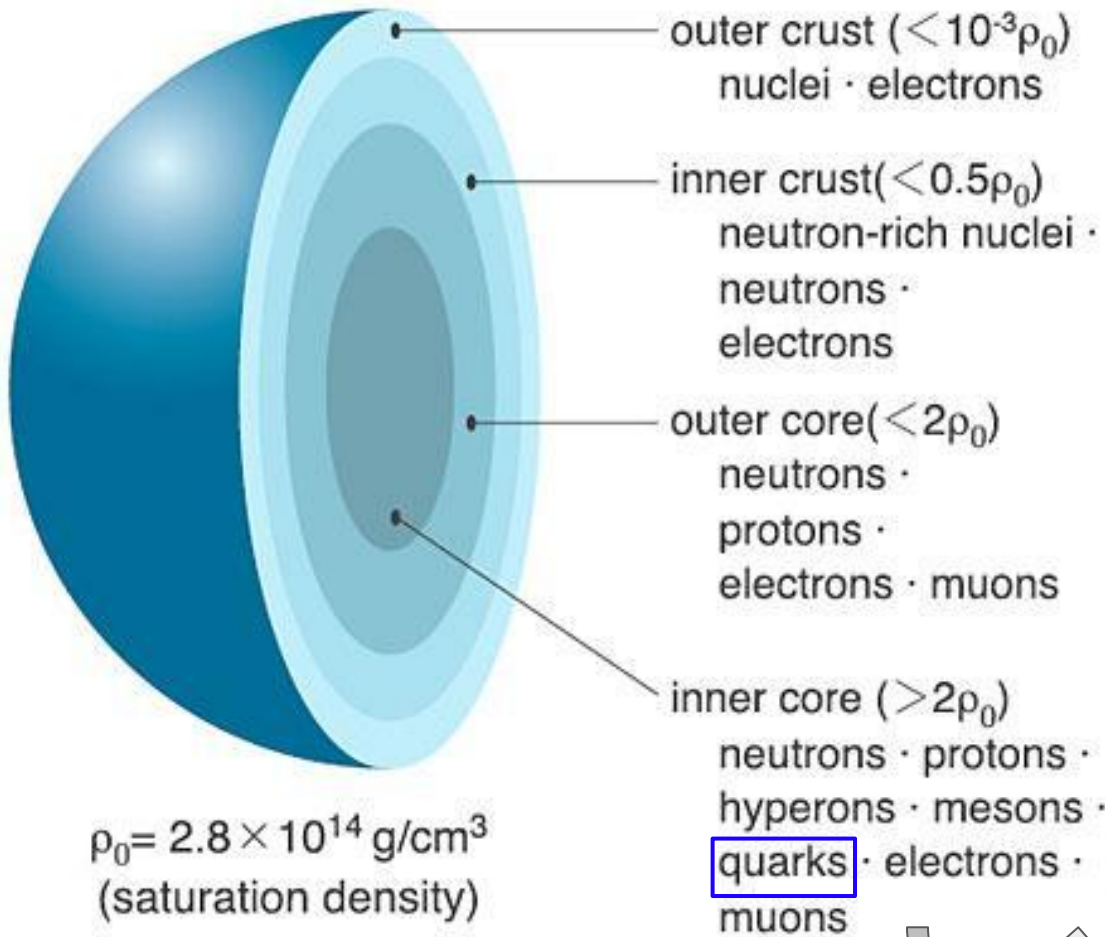
Ang Li et al. 2020 Journal of High Energy Astrophysics 2007.05116

~~should appear, but unlikely due to hyperon puzzle (Li et al. 2007; Burgio, Schuzle, Li 2011);~~

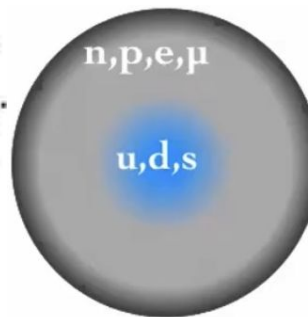
~~high threshold; likely not relevant to compact stars (Li et al. 2006, 2010);~~

✓ should appear early; hinder other strangeness; massive stars still possible (Li et al. 2015; Miao, Li et al. 2020);

Phase transition and hybrid stars



Assuming **1st-order** hadron-quark phase transition: Most promising scenario to be tested or distinguished from pure hadronic matter by future observations.

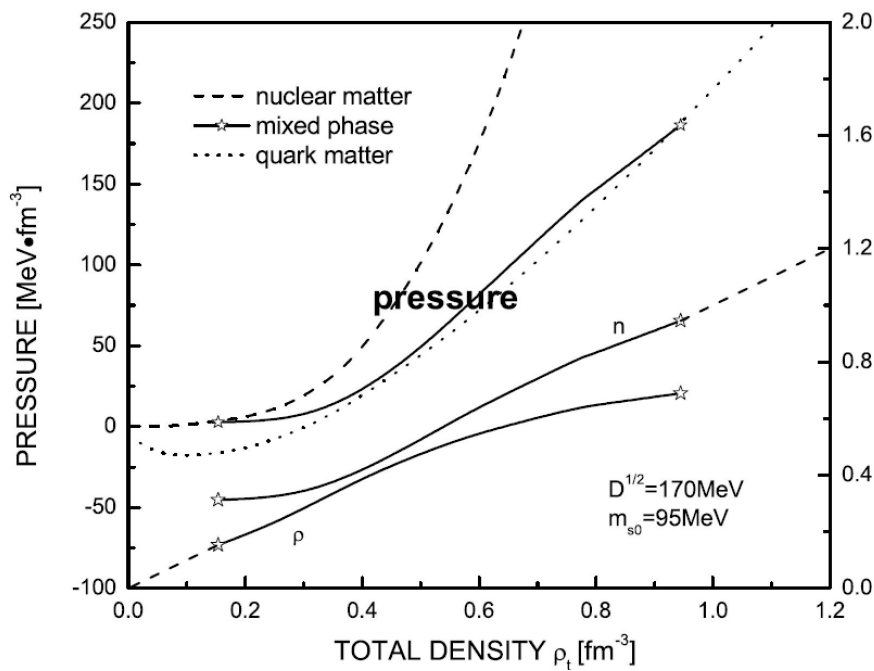


Massive stars could be hybrid stars with a stiff quark-matter core

PHYSICAL REVIEW C **91**, 035803 (2015)

Massive hybrid stars with a first-order phase transition

A. Li,^{1,2,*} W. Zuo,^{2,3} and G. X. Peng^{4,5}

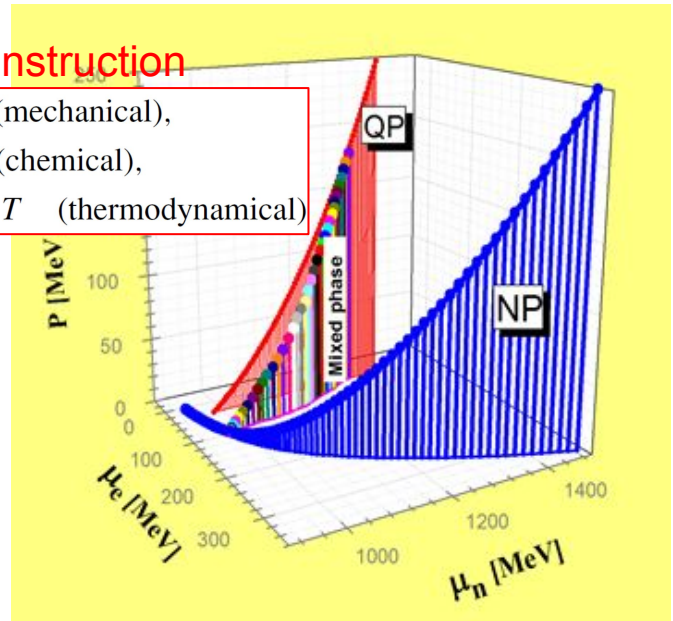


Gibbs construction

$$P_N = P_q \quad (\text{mechanical}),$$

$$\mu_N = \mu_q \quad (\text{chemical}),$$

$$T_N = T_q \equiv T \quad (\text{thermodynamical})$$



Many models of quark matter do exist, but they all contain a high degree of uncertainty; Here CDDM used;

-> To **constrain** M_{max} (EOS) and **quark deconfinement phase transition** parameters from current observations, also to stimulate **new** oriented observations.



Constraining Hadron-quark Phase Transition Parameters within the Quark-mean-field Model Using Multimessenger Observations of Neutron Stars

Zhiqiang Miao¹ , Ang Li^{1,5} , Zhenyu Zhu^{1,2} , and Sophia Han^{3,4}

¹ Department of Astronomy, Xiamen University, Xiamen, Fujian 361005, People's Republic of China; liang@xmu.edu.cn

² Institute for Theoretical Physics, D-60438 Frankfurt am Main, Germany

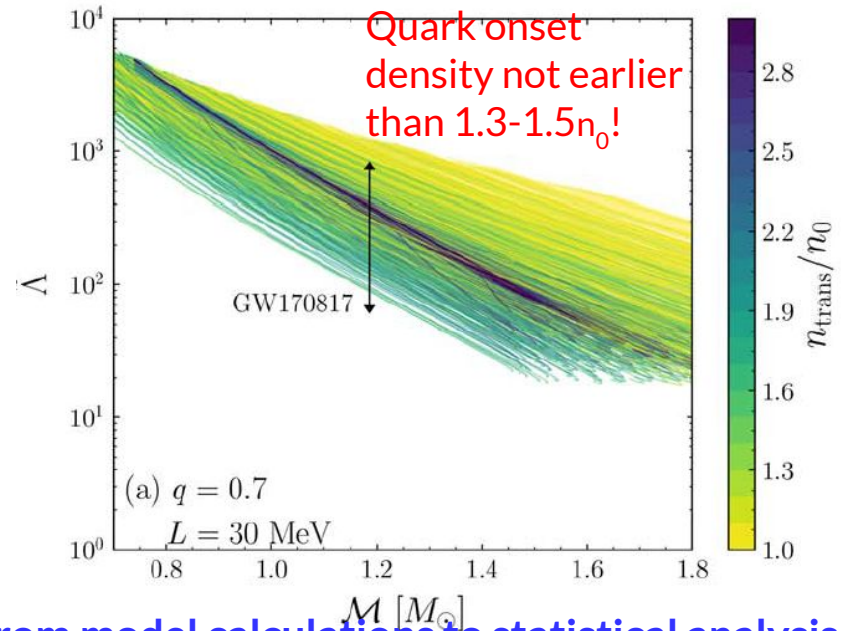
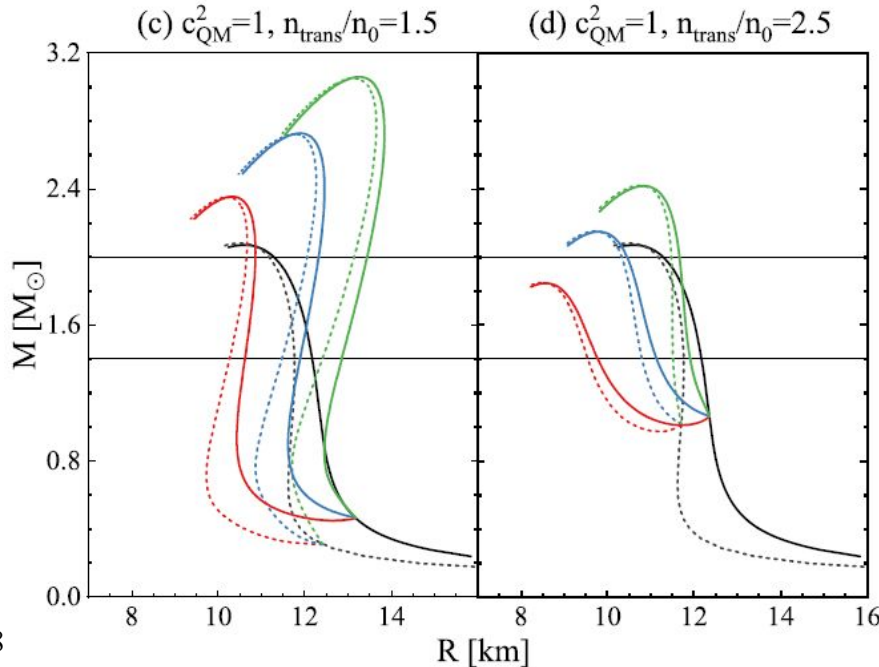
³ Department of Physics and Astronomy, Ohio University, Athens, OH 45701, USA

⁴ Department of Physics, University of California Berkeley, Berkeley, CA 94720, USA

QMF hadronic EOS plus CSS ($n_{\text{trans}}, \Delta\epsilon, c_{\text{QM}}$) characterize the high-density (quark matter) phase:

$$\epsilon(p) = \begin{cases} \epsilon_{\text{HM}}(p), & p < p_{\text{trans}} \\ \epsilon_{\text{HM}}(p_{\text{trans}}) + \Delta\epsilon + c_{\text{QM}}^{-2}(p - p_{\text{trans}}), & p > p_{\text{trans}} \end{cases}$$

Constant-speed-of-sound (CSS) scheme is a general parametrization suitable for expressing experimental constraints in a **model-independent** way.



from model calculations to statistical analysis->

Bayesian inference of the NS EOS parameter space from LIGO/Virgo and NICER

- Assuming the sources are NSs;
- Considering quark deconfinement phase transition in the EOS prior;
- Limiting EOS by the lower bound on M_{TOV} from heaviest MSP J0740+6620.

The Bayes's theorem

$$p(\boldsymbol{\theta} | \mathbf{d}, \mathbb{M}) = \frac{p(\boldsymbol{\theta} | \mathbb{M})p(\mathbf{d} | \boldsymbol{\theta}, \mathbb{M})}{p(\mathbf{d} | \mathbb{M})} \propto p(\boldsymbol{\theta} | \mathbb{M})p(\mathbf{d} | \boldsymbol{\theta}, \mathbb{M})$$

full parameter space from $1/\sqrt{3}$ (the conformal limit in perturbative QCD matter) to 1 (the causal limit)

\mathbb{M} : The QMF+CSS/DD2+CSS model

$\boldsymbol{\theta}$: parameters, including EOS parameters $\boldsymbol{\theta}_{\text{EOS}} = \{n_{\text{trans}}/n_0, \Delta\epsilon/\epsilon_{\text{trans}}, c_{\text{QM}}^2\}$ and $\boldsymbol{\theta}_{\text{GW}}$





\mathbf{d} : observational data, including three measurements: the mass of MSP J0740+6620, the tidal deformability from GW170817 and mass-radius of PSR J0030+0451

$p(\mathbf{d} | \boldsymbol{\theta}, \mathbb{M})$: likelihood, which can be expressed as $p(\mathbf{d} | \boldsymbol{\theta}, \mathbb{M}) = \mathcal{L}_{M_s} \times \mathcal{L}_{\text{GW}} \times \mathcal{L}_{\text{PSR}}$

$p(\boldsymbol{\theta} | \mathbb{M})$: prior for the parameters



Constraints on the Maximum Mass of Neutron Stars with a Quark Core from GW170817 and NICER PSR J0030+0451 Data

Ang Li¹ , Zhiqiang Miao¹ , Sophia Han^{2,3} , and Bing Zhang⁴ 

¹ Department of Astronomy, Xiamen University, Xiamen, Fujian 361005, People's Republic of China; liang@xmu.edu.cn

² Institute for Nuclear Theory, University of Washington, Seattle, WA 98195, USA

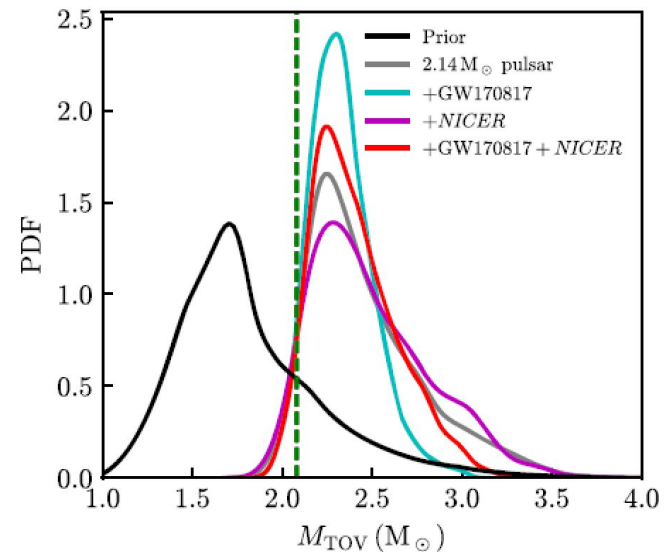
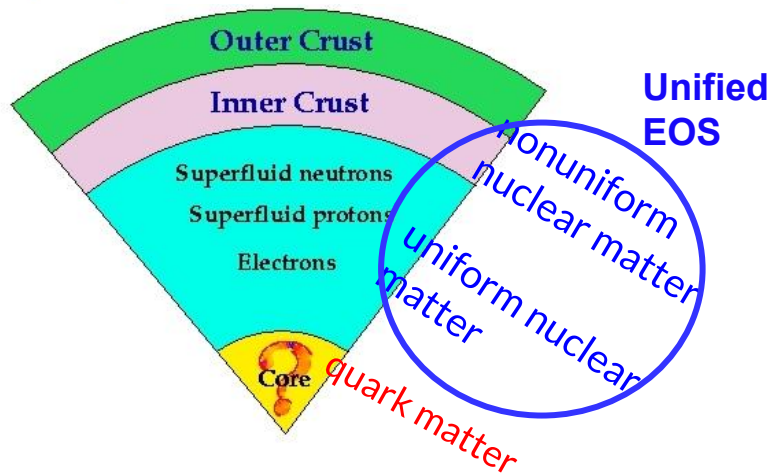
³ Department of Physics, University of California Berkeley, Berkeley, CA 94720, USA

⁴ Department of Physics and Astronomy, University of Nevada Las Vegas, Las Vegas, NV 89154, USA

Received 2021 January 24; revised 2021 March 26; accepted 2021 March 28; published 2021 May 21





Abstract

We perform a Bayesian analysis of the maximum mass M_{TOV} of neutron stars with a quark core, incorporating the observational data from tidal deformability of the GW170817 binary neutron star merger as detected by LIGO/Virgo and the mass and radius of PSR J0030+0451 as detected by the Neutron Star Interior Composition Explorer. The analysis is performed under the assumption that the hadron–quark phase transition is of first order, where the low-density hadronic matter described in a unified manner by the soft QMF or the stiff DD2 equation of state (EOS) transforms into a high-density phase of quark matter modeled by the generic “constant-sound-speed” parameterization. The mass distribution measured for the $2.14 M_{\odot}$ pulsar MSP J0740+6620 is used as the lower limit on M_{TOV} . We find the most probable values of the hybrid star maximum mass are $M_{\text{TOV}} = 2.36^{+0.49}_{-0.26} M_{\odot}$ ($2.39^{+0.47}_{-0.28} M_{\odot}$) for QMF (DD2), with an absolute upper bound around $2.85 M_{\odot}$, to the 90% posterior credible level. Such results appear robust with respect to the uncertainties in the hadronic EOS. We also discuss astrophysical implications of this result, especially on the postmerger product of GW170817, short gamma-ray bursts, and other likely binary neutron star mergers.





Constraints on the Maximum Mass of Neutron Stars with a Quark Core from GW170817 and NICER PSR J0030+0451 Data

Ang Li¹ , Zhiqiang Miao¹ , Sophia Han^{2,3} , and Bing Zhang⁴ 

¹ Department of Astronomy, Xiamen University, Xiamen, Fujian 361005, People's Republic of China; liang@xmu.edu.cn

² Institute for Nuclear Theory, University of Washington, Seattle, WA 98195, USA

³ Department of Physics, University of California Berkeley, Berkeley, CA 94720, USA

⁴ Department of Physics and Astronomy, University of Nevada Las Vegas, Las Vegas, NV 89154, USA

Received 2021 January 24; revised 2021 March 26; accepted 2021 March 28; published 2021 May 21

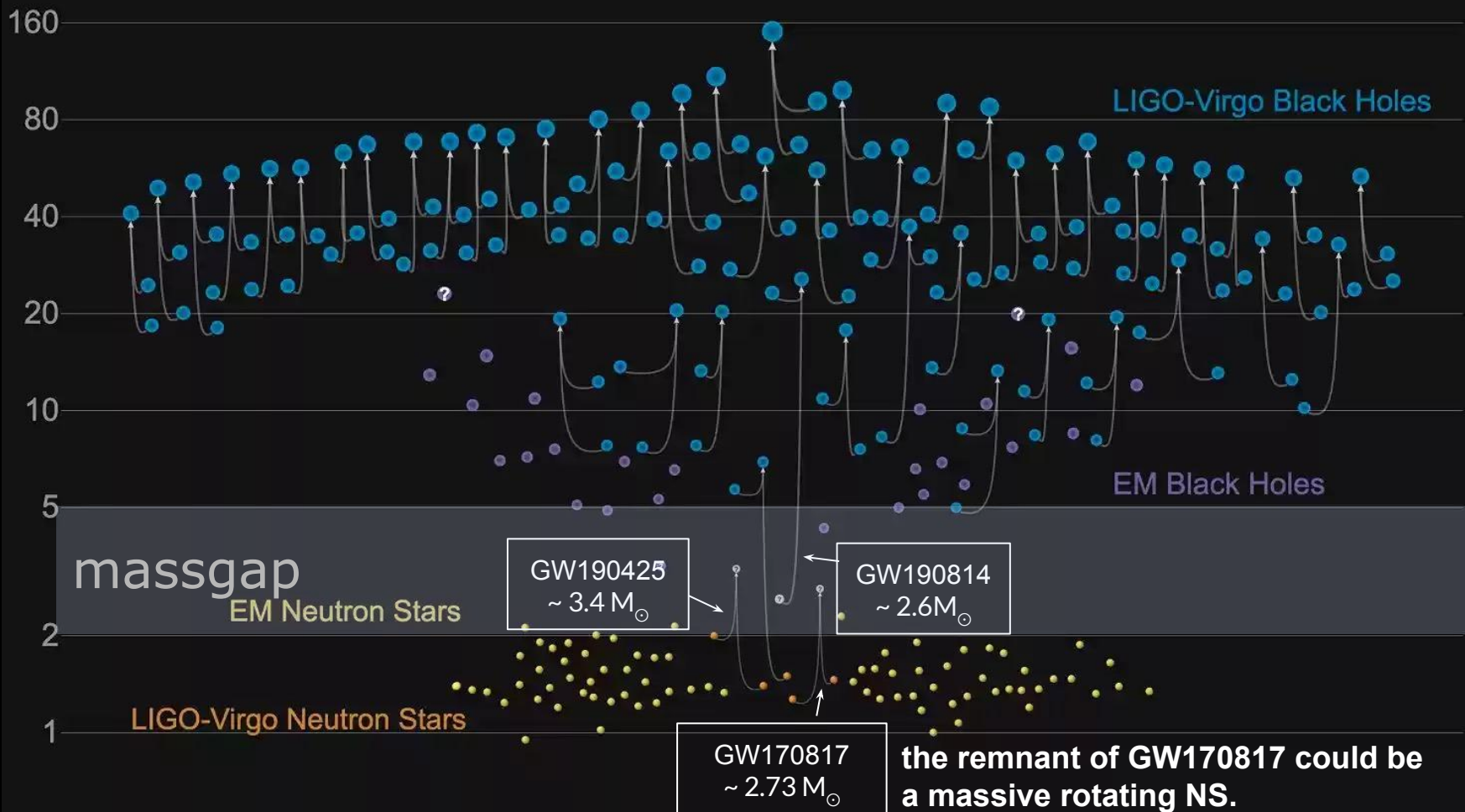
Abstract

We perform a Bayesian analysis of the maximum mass M_{TOV} of neutron stars with a quark core, incorporating the observational data from tidal deformability of the GW170817 binary neutron star merger as detected by LIGO/Virgo and the mass and radius of PSR J0030+0451 as detected by the Neutron Star Interior Composition Explorer. The analysis is performed under the assumption that the hadron–quark phase transition is of first order, where the low-density hadronic matter described in a unified manner by the soft QMF or the stiff DD2 equation of state (EOS) transforms into a high-density phase of quark matter modeled by the generic “constant-sound-speed” parameterization. The mass distribution measured for the $2.14 M_{\odot}$ pulsar MSP J0740+6620 is used as the lower limit on M_{TOV} . We find the most probable values of the hybrid star maximum mass are $M_{\text{TOV}} = 2.36^{+0.49}_{-0.26} M_{\odot}$ ($2.39^{+0.47}_{-0.28} M_{\odot}$) for QMF (DD2), with an absolute upper bound around $2.85 M_{\odot}$, to the 90% posterior credible level. Such results appear robust with respect to the uncertainties in the hadronic EOS. We also discuss astrophysical implications of this result, especially on the postmerger product of GW170817, short gamma-ray bursts, and other likely binary neutron star mergers.

- The **general** requirements adopted here (e.g., causality) should also apply to any alternative hadron–quark phase transition scenarios or other types of strangeness phase transitions;
- Our conclusions are **valid and useful** for identifying compact objects' nature with their mass falling into the possible mass gap.

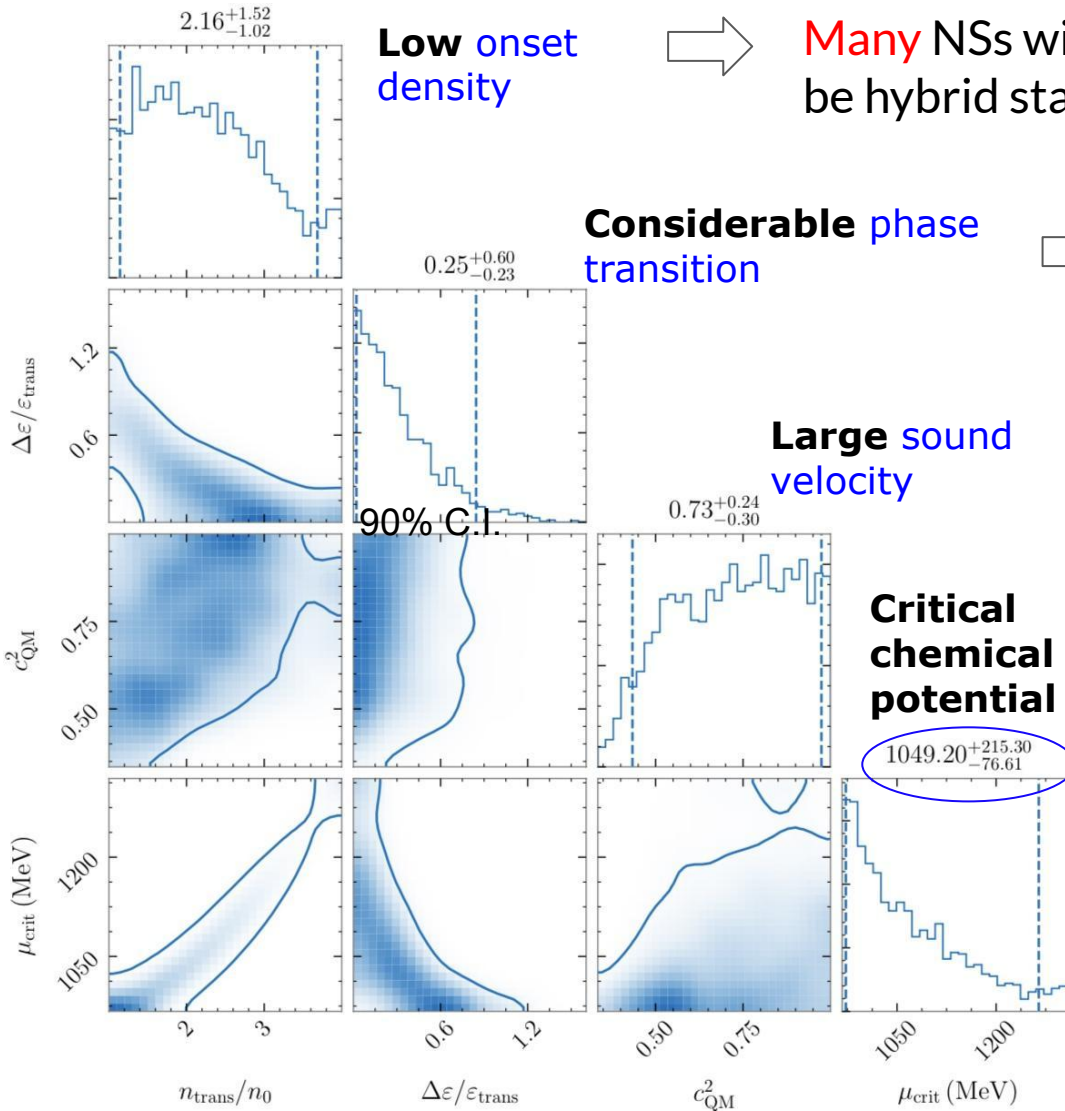
NS maximum mass $\sim 2.4M_{\odot}$: Help to identify the nature of compact objects with the mass falling into the NS-BH gap

Masses in the Stellar Graveyard *in Solar Masses*



GWTC-2 plot v1.0

general results on the hadron-quark phase transition described by the CSS parametrization



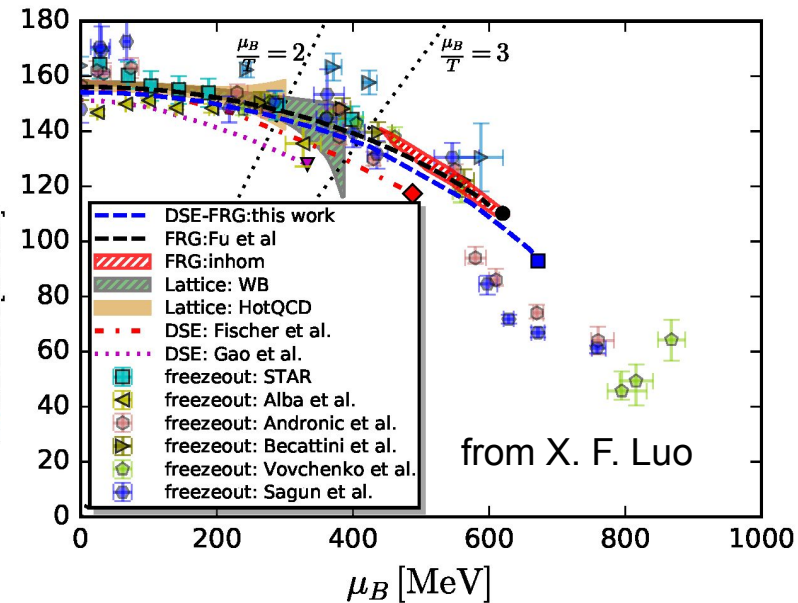
Many NSs with mass above $\sim 1.4M_{\odot}$ should be hybrid stars!



Subsequent support on the critical point



Large NS mass **needs** large sound speed!

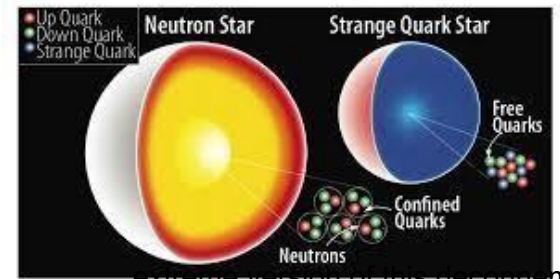


GW170817's tidal deformability are compatible with a binary-quark-star merger!

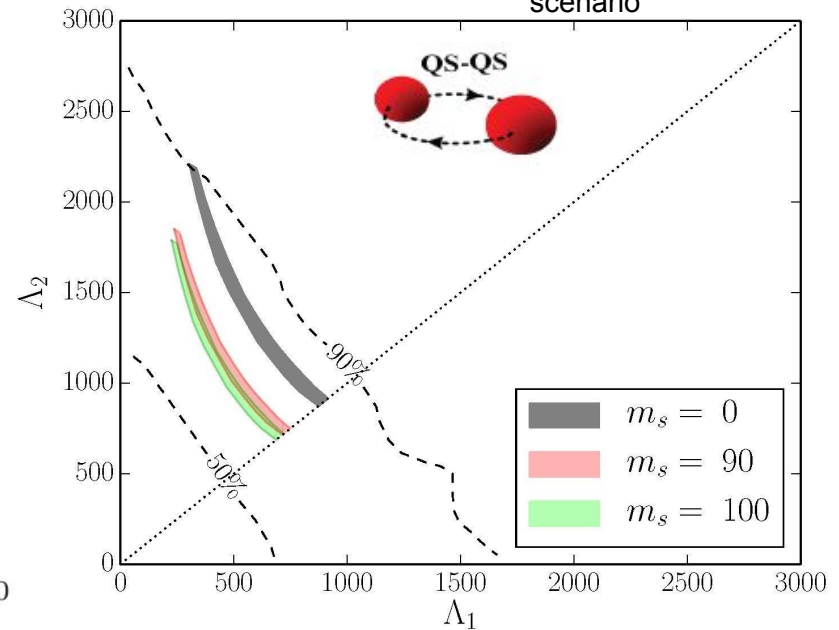
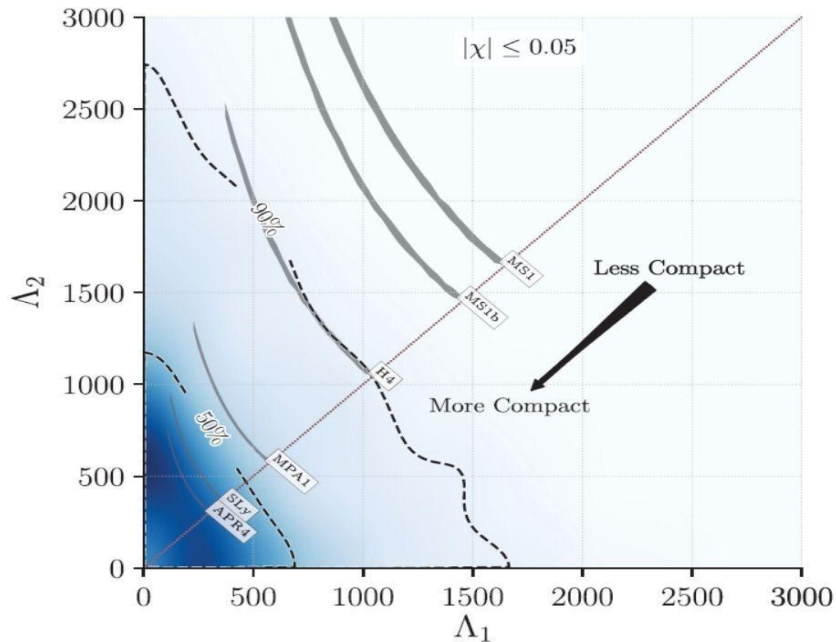
PHYSICAL REVIEW D **97**, 083015 (2018)

Constraints on interquark interaction parameters with GW170817 in a binary strange star scenario

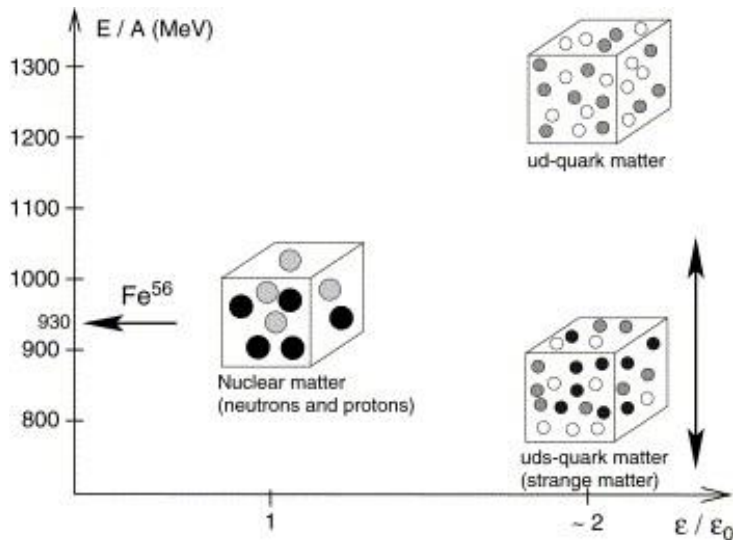
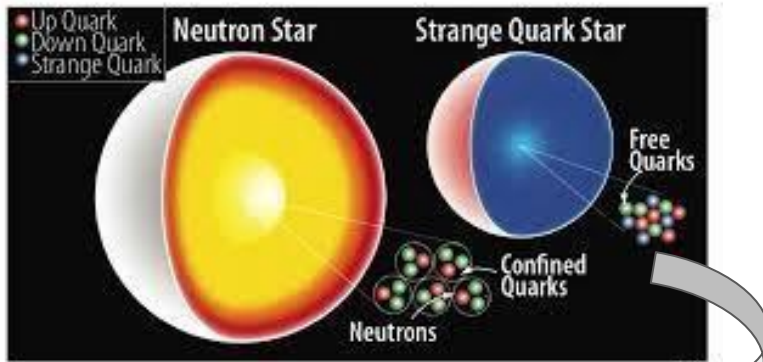
En-Ping Zhou,^{1,2} Xia Zhou,³ and Ang Li^{4,*}



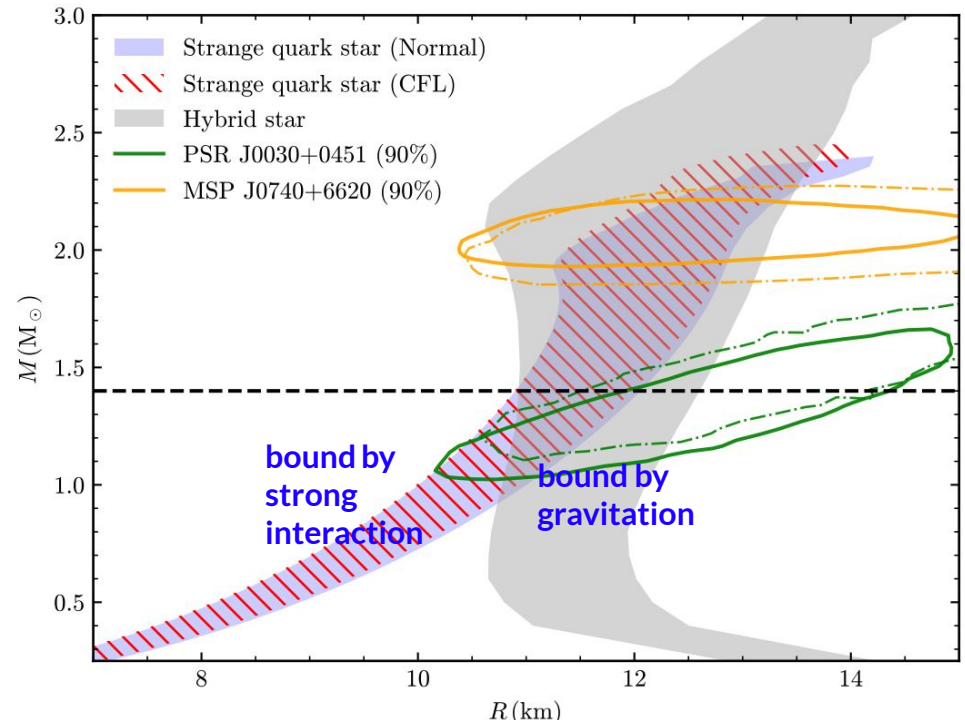
extreme version of this decomposition scenario



One or two-family scenario for compact stars



NS+QS **two**-families scenario or
NS/QS **one** family scenario?



Strange quark matter being the true ground state of matter?

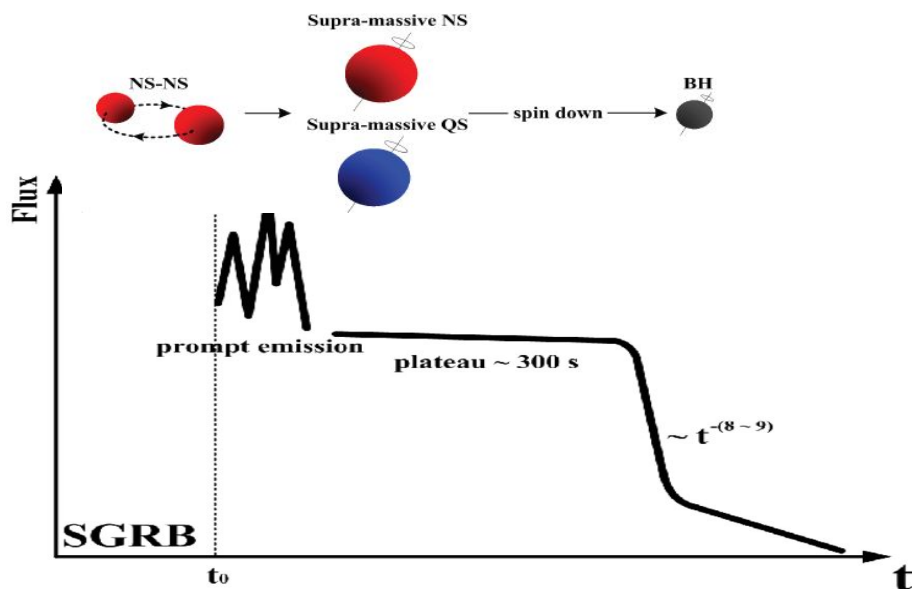
—**Bodmer-Witten conjecture**

Bodmer 1971; Witten 1984

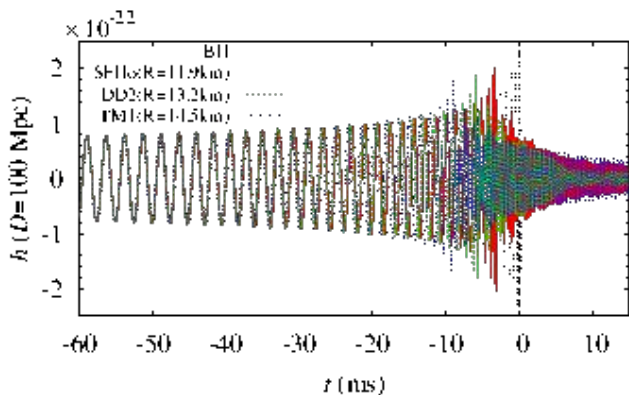
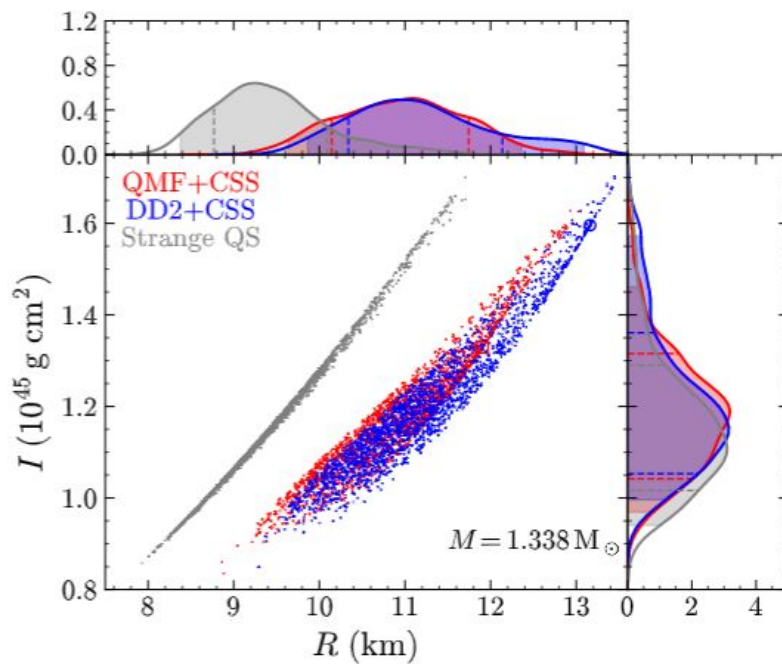
How to distinguish?

How to distinguish?

combined analysis of **SGRBs** and **kilonovae** events (Li et al. 2016, 2017);



< ~1km accuracy **radius (+Mol)** measurement (Miao & Li 2021);

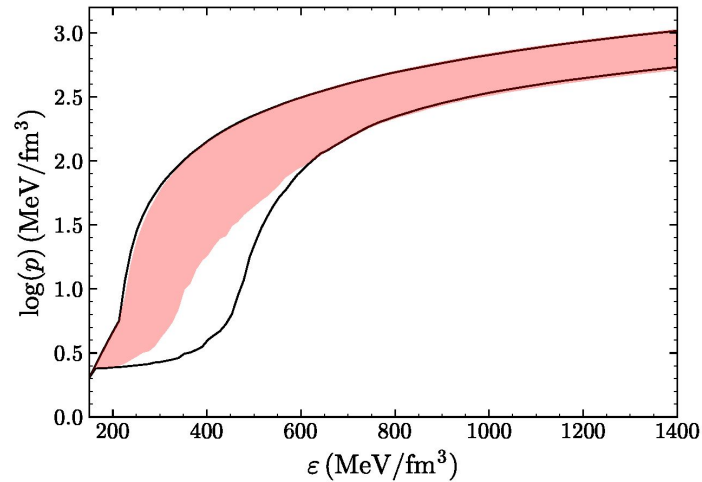
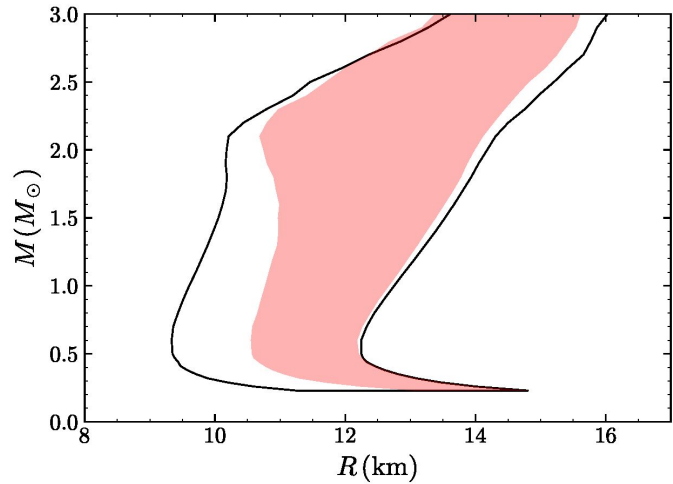


post-merger **GW** signals (possibly from O5 until maybe 2030) with the help of **merger simulation** (in progress).

Radius measurement of massive NSs is vital

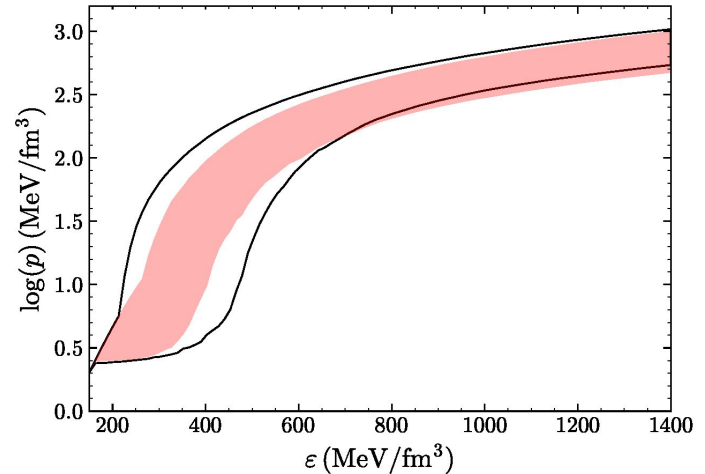
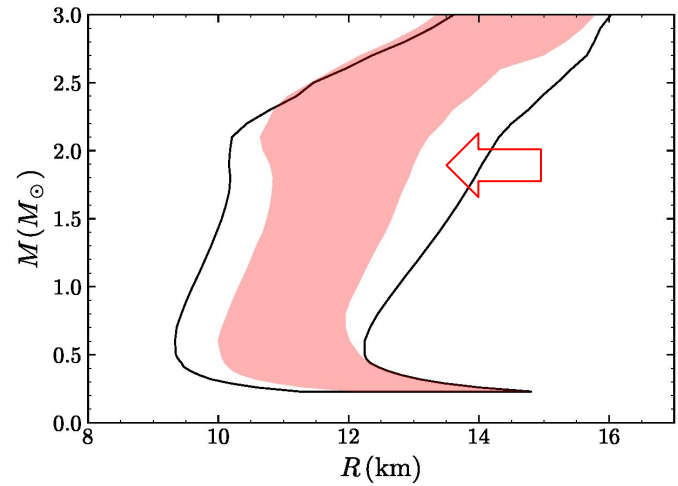
~1.4 M_{\odot} J0030's M,R measurement

$1.34^{+0.15}_{-0.16} M_{\odot}$ and $12.71^{+1.14}_{-1.19}$ km



~2 M_{\odot} J0740's M,R measurement

$2.072^{+0.067}_{-0.066} M_{\odot}$ and $12.39^{+1.30}_{-0.98}$ km



Radius measuring: one of primary goal of next generation of hard x-ray timing instruments

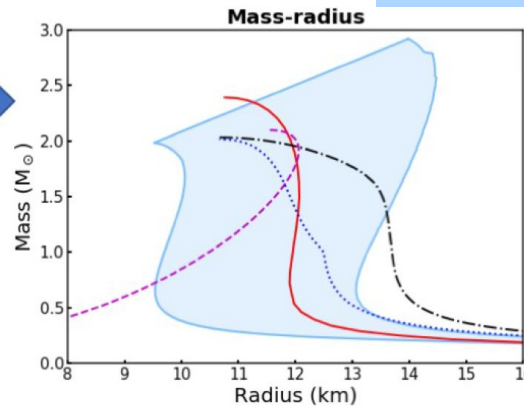
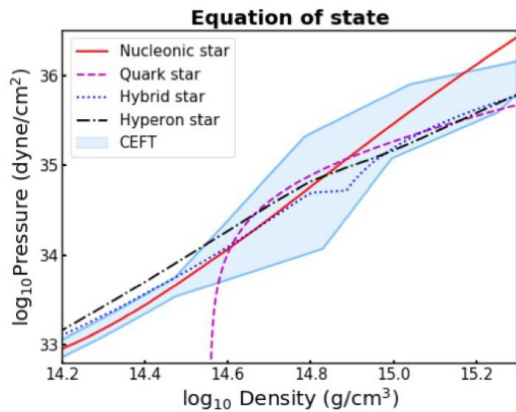
Astro2020 Science White Paper

Determining the Equation of State of Cold, Dense Matter with X-ray Observations of Neutron Stars

- Thematic Areas:
- Planetary Systems
 - Star and Planet Formation
 - Formation and Evolution of Compact Objects
 - Cosmology and Fundamental Physics
 - Stars and Stellar Evolution
 - Resolved Stellar Populations and their Environments
 - Galaxy Evolution
 - Multi-Messenger Astronomy and Astrophysics

Principal Author:

Name: Slavko Bogdanov
 Institution: Columbia University
 Email: slavko@astro.columbia.edu



SCIENCE CHINA

Physics, Mechanics & Astronomy



• Invited Review •
 Special Issue: The X-ray Timing and Polarimetry Frontier with eXTP

February 2019 Vol. 62 No. 2: 029503
<https://doi.org/10.1007/s11433-017-9188-4>

Dense matter with eXTP

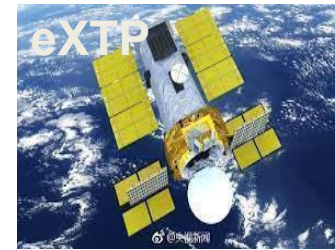
Anna L. Watts^{1*}, WenFei Yu², Juri Poutanen^{3,4}, Shu Zhang⁵, Sudip Bhattacharyya⁶,

STROBE-X

STROBE-X: X-ray Timing and Spectroscopy on Dynamical Timescales from Microseconds to Years

White Paper Submitted to Astro 2020 Decadal Survey

Paul S. Ray
 U.S. Naval Research Laboratory
paul.ray@nrl.navy.mil



launch by 2025
 (IHEP)

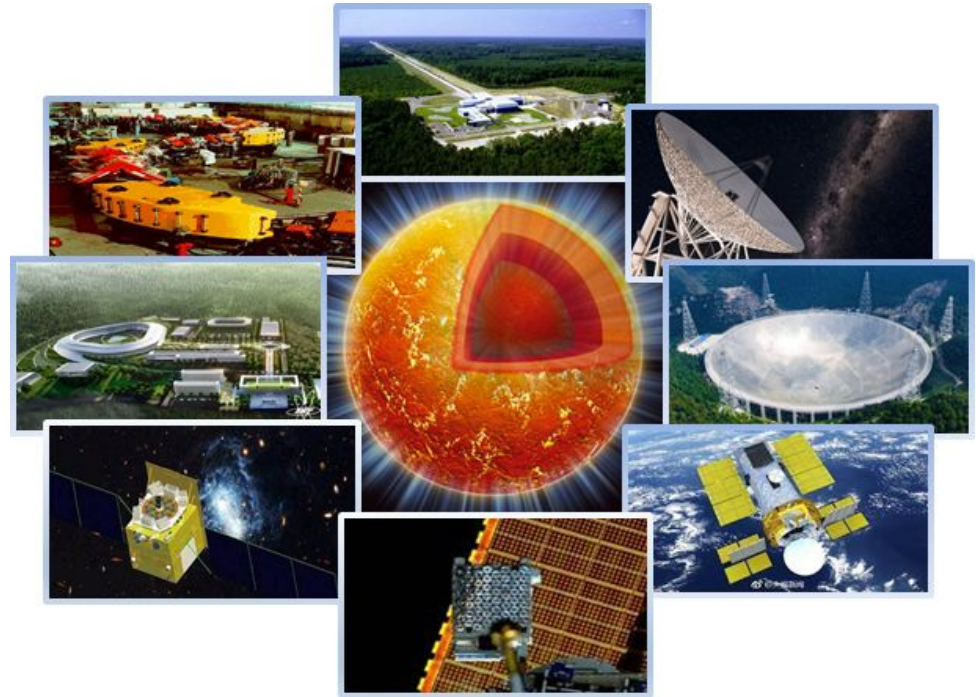


ready for
 construction in the
 2020s (NASA)

Magenta - quark star, composed entirely of quark matter from Li et al., 2016

Take-home message

- Connect consistently nuclear exp. and LIGO/Virgo+NICER obs. for “quantitative” studies of neutron stars;
- Current limits on M_{\max} (EOS):
~2.08(observational)
~2.4(theoretical);
- Joint efforts from nuclear and astro. for probing the phase state of dense QCD matter: theory + simulation + data!



Thank you and Q&A !