

μ SR Applications in the Research of Diluted Magnetic Semiconductors

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(Material synthesis + μ SR + NMR)

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In collaboration with:

Y.J. Uemura' Group (Columbia)

Changqing Jin's Group (IOP)

Experiments:

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Cui Ding (now a postdoc at Tsinghua with Prof. Qikun Xue)

Shengli Guo

Licheng Fu

Yilun Gu

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Jianhua Zhao(Semi), Kai Chang (Semi), Kaiyou Wang(semi), Igor Zutic (Buffalo), S. Maekawa (JAEA), Bo Gu(KITP), Igor Mazin (Navy), Ravin Bhatt (Princeton), Xin Wan (ZJU), Chao Cao (HZNU), Ni Ni (UCLA)

Outline

a. Introduction of some newly fabricated DMSs

“111”, “122”, “1111”, “42622”, “32522”

----- $\text{Li}_{1+y}(\text{Zn}_{1-x}\text{Mn}_x)\text{P}$

----- n-type $\text{Ba}(\text{Zn}_{1-x}\text{Co}_x)_2\text{As}_2$

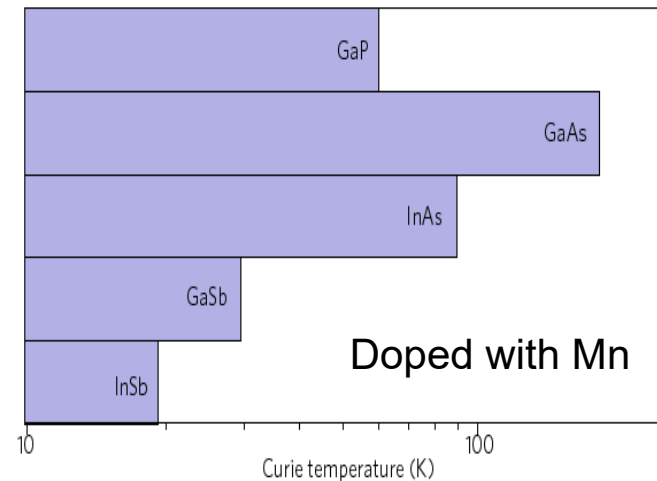
b. Microscopic characterization

----- μSR results on bulk form DMSs

----- NMR results on $\text{Li}_{1.1}(\text{Cd}_{1-x}\text{Mn}_x)\text{P}$ and $\text{Li}_{1.15}(\text{Zn}_{1-x}\text{Mn}_x)\text{P}$

Why bulk form DMS?

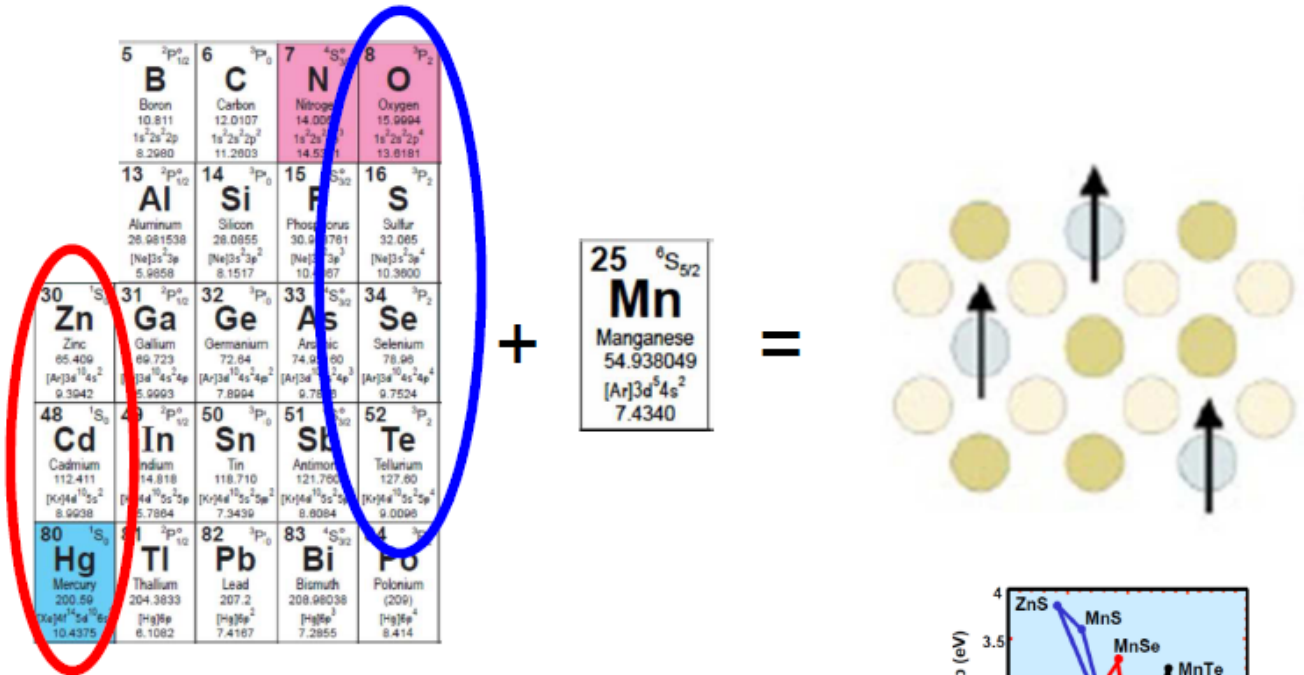
- We want to create “spintronics”, and use not only charge but also spin information
- $\text{Ga}_{1-x}\text{Mn}_x\text{As}$, $T_C \sim 200\text{K}$, limited for practical applications; holes and spins are all introduced by Mn atoms (**Jianhua Zhao's group in Institute of Semiconductor, Beijing**)
- Mismatch of valences for Mn^{2+} and Ga^{3+} atoms, chemical solubility is less than 1% **in bulk form**
- Only low dimensional thin film can be fabricated with Molecular Beam Epitaxy
- Mn^{2+} substitution for Ga^{3+} introduces holes and spins
- **NMR, μSR and neutron scattering** are limited to investigate thin film.
- **No NMR work ever done on DMS**
- DMS in bulk form is **highly expected!**



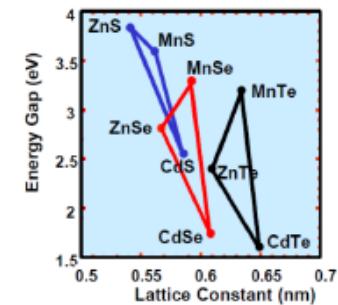
Currently achieved highest T_C in various III-V DMS thin films

(Dietl et al, Nature Materials, 2010, (9)965)

II-VI Diluted Magnetic Semiconductors



Spin Glass Ordering, saturation moment $\sim 0.01\mu_B$



0.003 μ_B/Mn

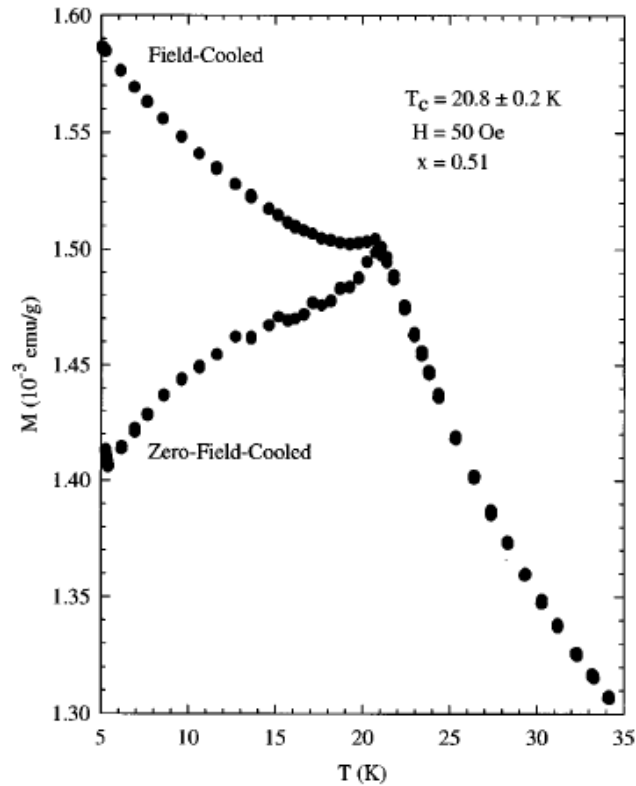


FIG. 6. Zero-field-cooled and field-cooled magnetization vs temperature T for $\text{Zn}_{0.49}\text{Mn}_{0.51}\text{Te}$. The cusp in the magnetization at $T = 20.8 \pm 0.2$ K for the zero-field-cooled case is characteristic of a transition to the spin-glass phase.

Phys Rev B, 58, 12876 (1998)

$\text{Zn}_{0.49}\text{Mn}_{0.51}\text{Te}$

- Zn and Mn valences of “+2”,
- No solid solution limit
the doping can be as high as ~70%
- Mn ions do not provide carriers,
only provide local moment
- Difficult to control density and type
of carriers

ZnSe/ZnTe \rightarrow FeSe
(“11” family of Fe-based SC)

Dilute Moment n -Type Ferromagnetic Semiconductor Li(Zn,Mn)AsJ. Mašek,¹ J. Kudrnovský,¹ F. Máca,¹ B. L. Gallagher,² R. P. Campion,² D. H. Gregory,³ and T. Jungwirth^{4,2}¹*Institute of Physics ASCR, Na Slovance 2, 182 21 Praha 8, Czech Republic*²*School of Physics and Astronomy, University of Nottingham, Nottingham NG7 2RD, United Kingdom*³*Department of Chemistry, University of Glasgow, Glasgow G12 8QQ, United Kingdom*⁴*Institute of Physics ASCR, Cukrovarnická 10, 162 53 Praha 6, Czech Republic*

(Received 6 September 2006; published 7 February 2007)

We propose to replace Ga in (Ga,Mn)As with Li and Zn as a route to high Curie temperature, carrier mediated ferromagnetism in a dilute moment n -type semiconductor. Superior material characteristics, rendering Li(Zn,Mn)As a realistic candidate for such a system, include high solubility of the isovalent substitutional Mn impurity and carrier concentration controlled independently of Mn doping by adjusting Li-(Zn,Mn) stoichiometry. Our predictions are anchored by *ab initio* calculations and comparisons with the familiar and directly related (Ga,Mn)As, by the physical picture we provide for the exchange interaction between Mn local moments and electrons in the conduction band, and by analysis of prospects for the controlled growth of Li(Zn,Mn)As materials.

DOI: [10.1103/PhysRevLett.98.067202](https://doi.org/10.1103/PhysRevLett.98.067202)

PACS numbers: 75.50.Pp, 73.61.Ey, 75.30.Hx

“III-V”



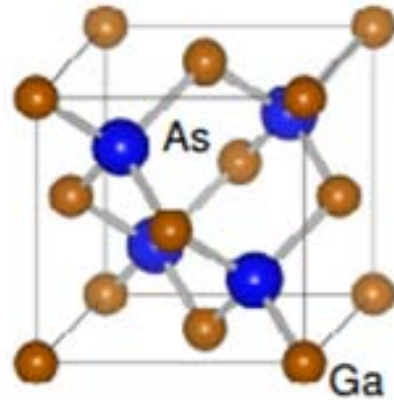
“I-II-V”



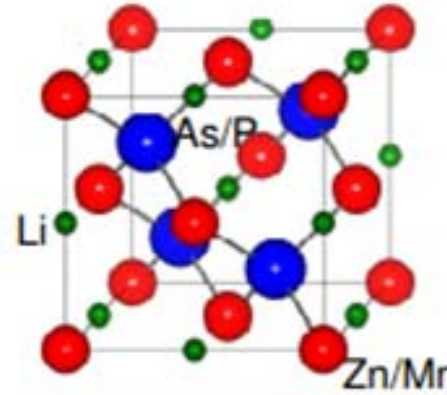
“II-VI”

$(\text{Li}^+\text{Zn}^{2+}) = \text{Ga}^{3+}$
 $\rightarrow \text{LiZnAs} = \text{GaAs}$

$(\text{Li}^+\text{As}^{3-}) = \text{Se}^{2-}$
 $\rightarrow \text{Zn}(\text{LiAs}) = \text{ZnSe}$



GaAs (1.42ev)



LiZnAs (1.6ev)
LiZnP (2.04ev)
LiCdP (1.30ev)

$\text{Ga}_{1-x}\text{Mn}_x\text{As}$: Mn dope both spins and holes

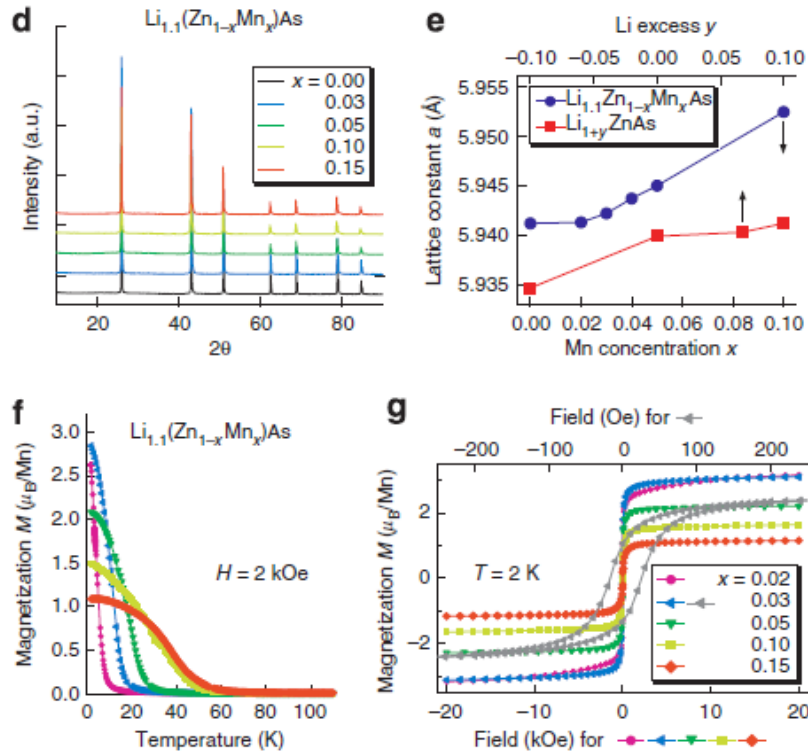
$\text{Li}_{1+y}(\text{Zn}^{+2}_{1-x}\text{Mn}^{+2}_x)\text{As}$, $\text{Li}_{1+y}(\text{Zn}^{+2}_{1-x}\text{Mn}^{+2}_x)\text{P}$, $\text{Li}_{1+y}(\text{Cd}^{+2}_{1-x}\text{Mn}^{+2}_x)\text{P}$ ($T_C \sim 40\text{-}50\text{K}$)
: decouple spins and carriers

- Mn doping induces
- Li extra or deficient

LiZnAs \rightarrow LiFeAs
(“111” family of Fe-based SC)

(C.Q. Jin's group in JET, Li Dong et al, Nature Comm, 2014, PRB 2016)

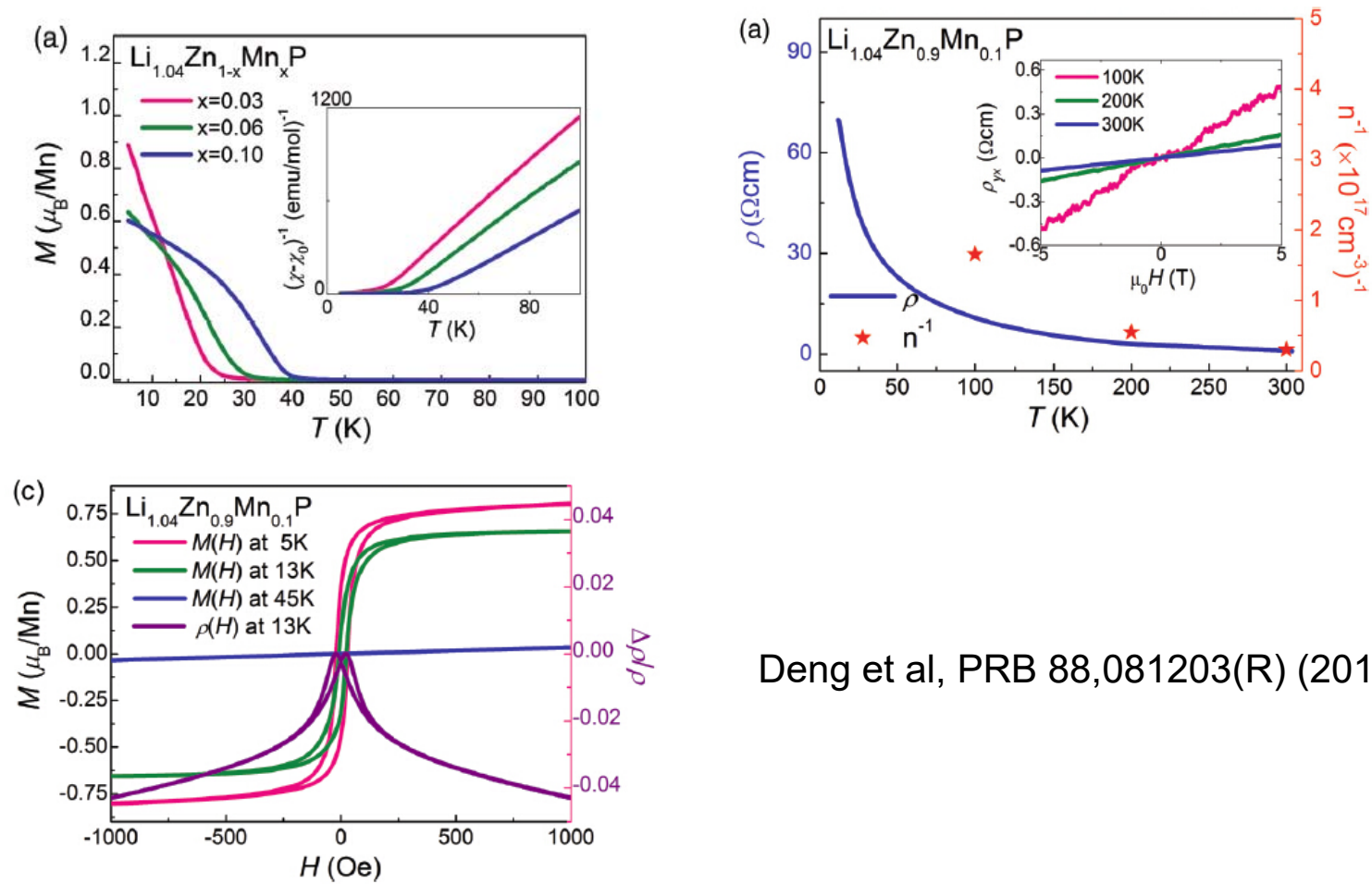
“111”- $\text{Li}_{1+y}(\text{Zn}_{1-x}\text{Mn}_x)\text{As}$ ($T_C \sim 50$ K)



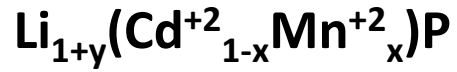
- Saturation moments continuously decrease with more Mn
- Effective moment is $5\sim 6\mu_B/\text{Mn}$, high spin state
- Coercive field ~ 50 Oe, for magnetic manipulation

Z. Deng et al, Nature Communications (2011)

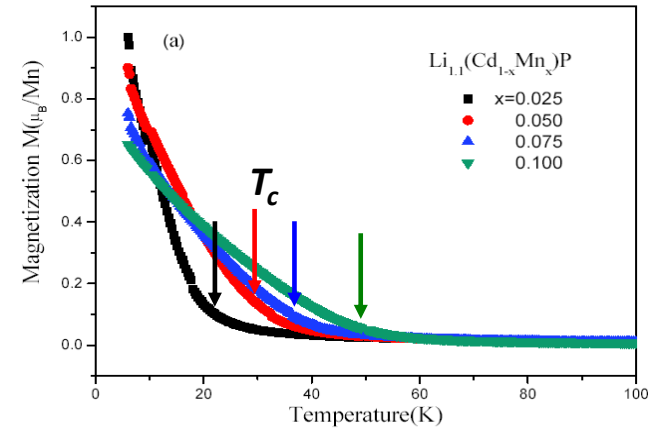
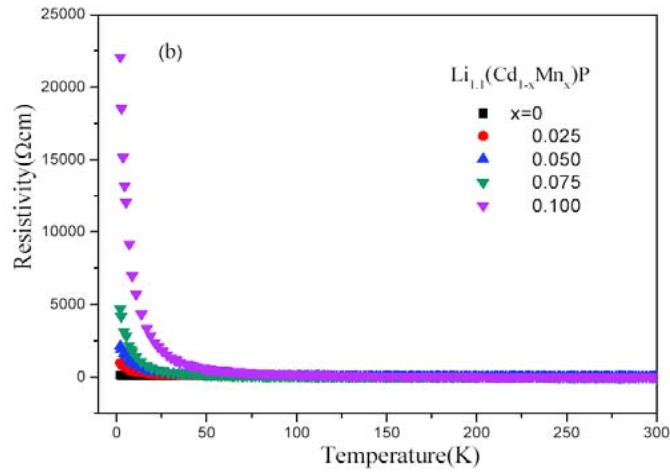
$\text{Li}_{1+y}(\text{Zn}^{2+}_{1-x}\text{Mn}^{2+}_x)\text{P}$: Max $T_C \sim 40$ K



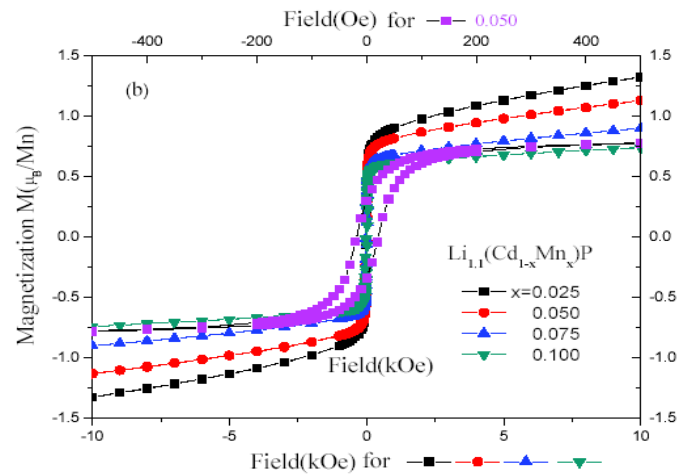
Deng et al, PRB 88,081203(R) (2013)



Wei Han et al, Scientific Reports 9, 7490 (2019)



Magnetizations of Li1.1(Cd,Mn)P as function of Mn content

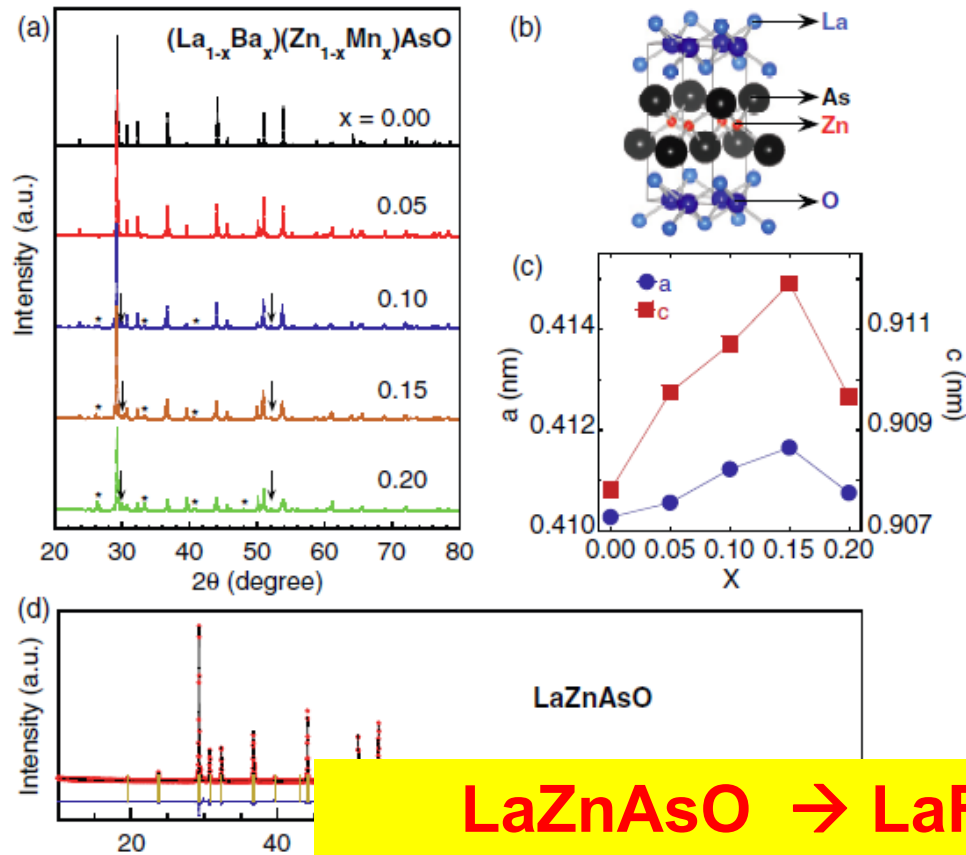


The MH curves of Li1.1(Cd,Mn)P

T_c and moment size at 6K

$x^*(100\%)$	T_c	Mn(6K)
2.5	21	1.00
5.0	28	0.90
7.5	35	0.75
10	43	0.65

$(\text{La}_{1-x}\text{Ba}_x)(\text{Zn}_{1-x}\text{Mn}_x)\text{AsO}$ $T_C \sim 40\text{K}$

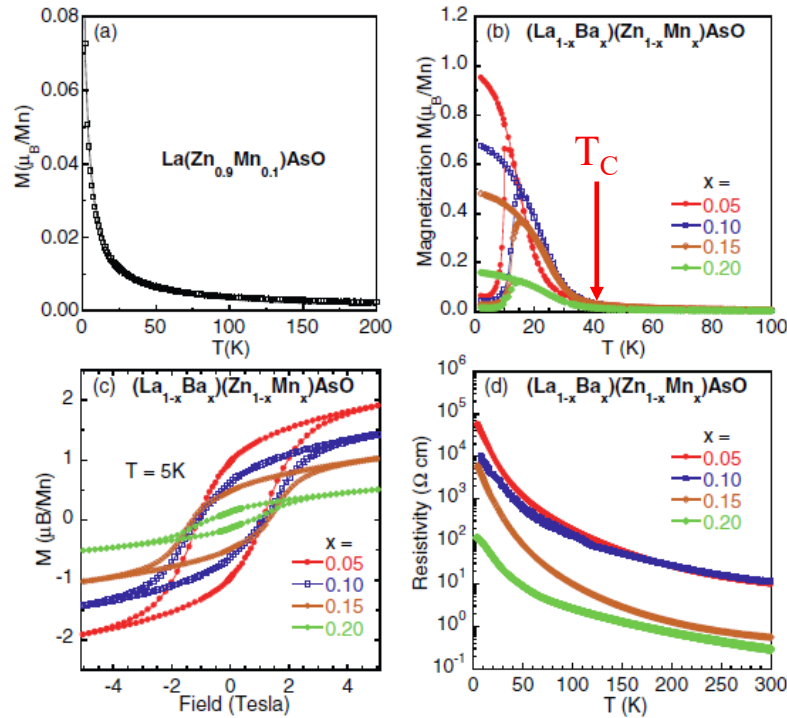


**• Isostructural to Fe-based SC
“1111”- $\text{LaFeAsO}_{1-x}\text{F}_x$**

**$\text{LaZnAsO} \rightarrow \text{LaFeAsO}$
 (“1111” family of Fe-based SC)**

C. Ding and F.L. Ning, PRB 88, 041102 (2013)

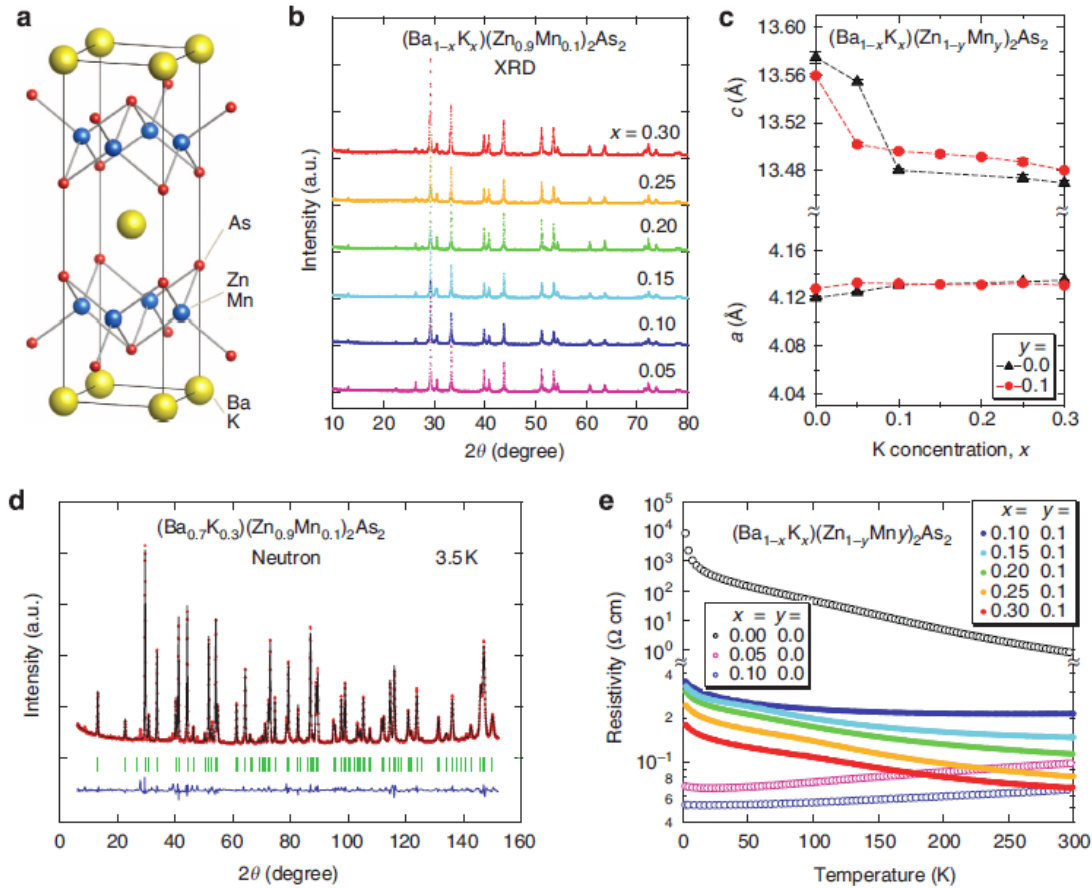
(La_{1-x}Ba_x)(Zn_{1-x}Mn_x)AsO $T_C \sim 40K$



- doping Mn only, paramagnetic state
- Only when carriers are doped via Ba substitution for La, FM ordering is observed
- From resistivity, semiconducting behavior
- Large coercive field ~ 1 Tesla

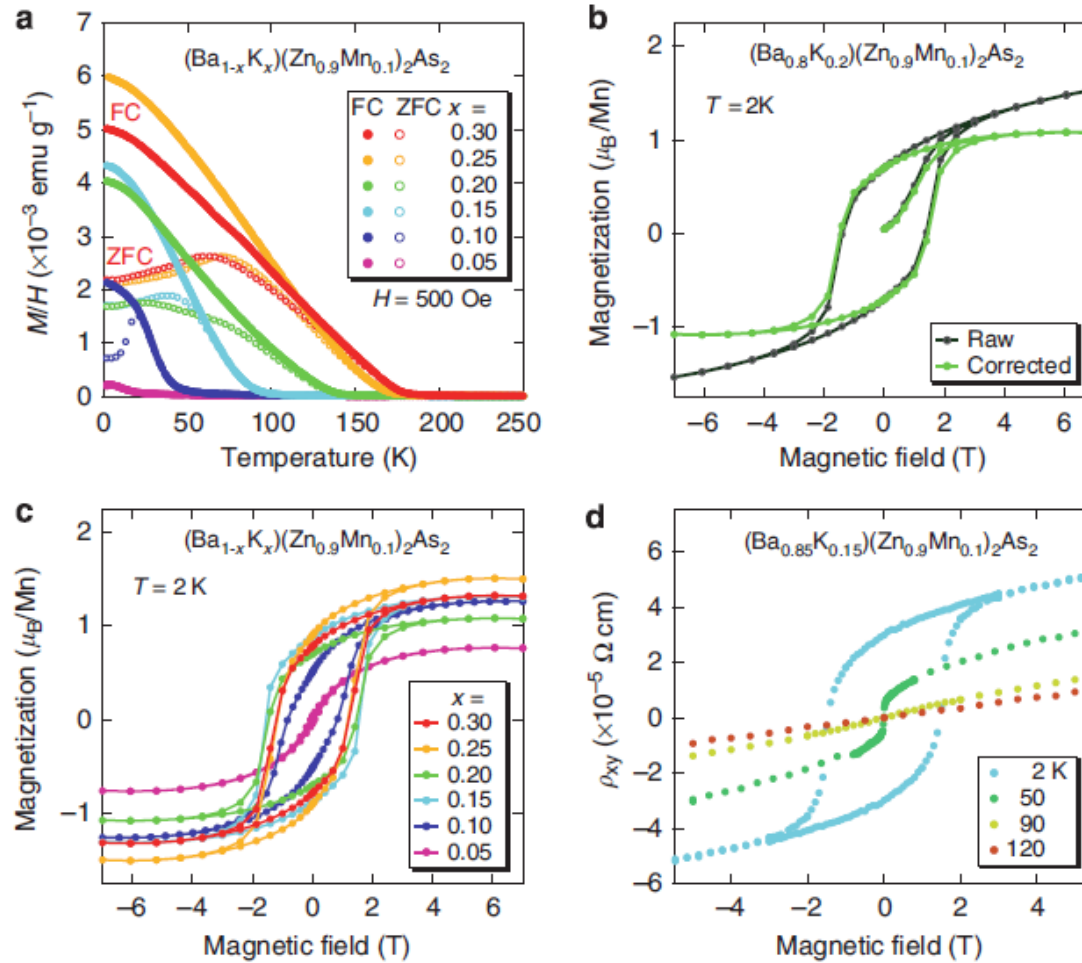
C. Ding and F.L. Ning, PRB 88, 041102 (2013)

$(\text{Ba}_{1-x}\text{K}_x)(\text{Zn}_{1-y}\text{Mn}_y)_2\text{As}_2$ $T_C \sim 180\text{K}$



K. Zhao et al, Nature Communications, 4, 1442 (2013)

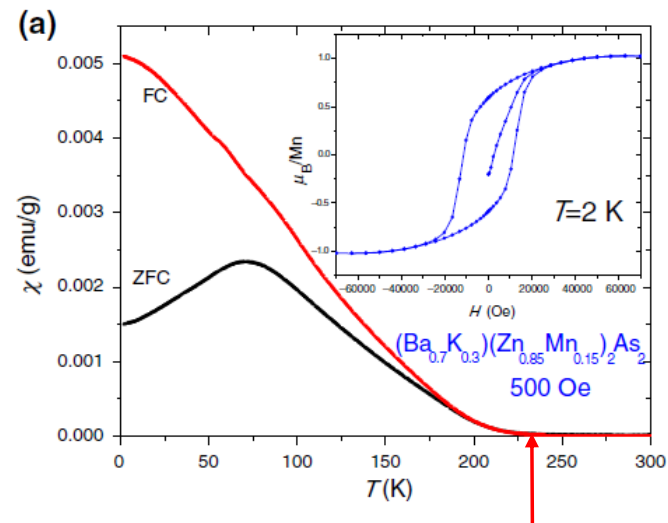
$(\text{Ba}_{1-x}\text{K}_x)(\text{Zn}_{1-y}\text{Mn}_y)_2\text{As}_2$ $T_C \sim 180\text{K}$



K. Zhao et al, Nature Communications, 4, 1442 (2013)

Ferromagnetism at 230 K in $(\text{Ba}_{0.7}\text{K}_{0.3})(\text{Zn}_{0.85}\text{Mn}_{0.15})_2\text{As}_2$ diluted magnetic semiconductor

Kan Zhao · Bijuan Chen · Guoqiang Zhao ·
Zhen Yuan · Qingqing Liu · Zheng Deng ·
Jinlong Zhu · Changqing Jin



BaZn₂As₂ → BaFe₂As₂
(“122” family of Fe-based SC)

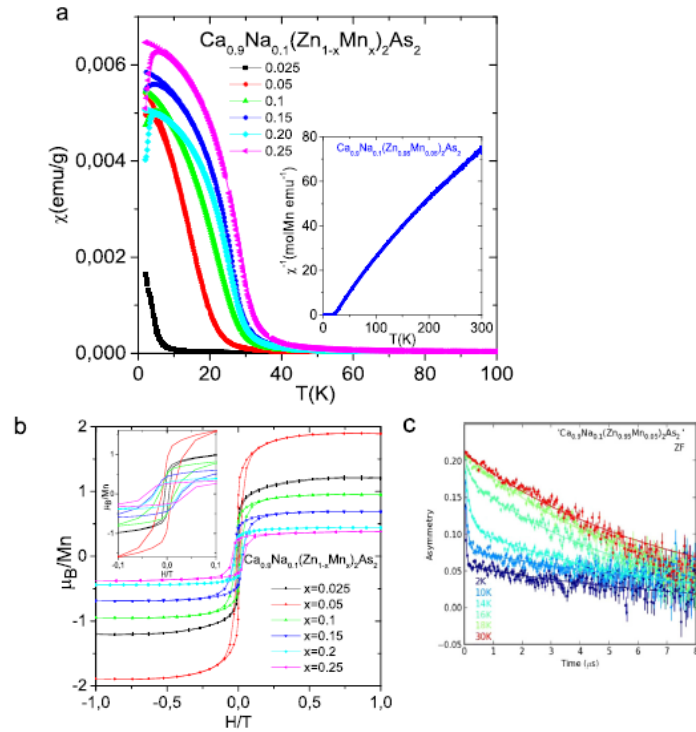
$T_c \sim 230$ K

Fabricate new DMS and what we have

GaAs	----- Ga _{1-x} Mn _x As -----	MnAs
	FM, $T_c \sim 210\text{K}(\text{film})$	FM, $T_c \sim 320\text{K}$
LiZnAs	----- LiZn _{1-x} Mn _x As -----	LiMnAs
“111”	FM, $T_c \sim 55\text{K}(\text{bulk})$	AFM, $T_N \sim 400\text{K}$
I-II-V	(Nature Comm. 2011, IOP)	
BaZn ₂ As ₂	---- Ba _{1-y} K _y (Zn _{1-x} Mn _x) ₂ As ₂ -----	BaMn ₂ As ₂
“122”	FM, $T_c \sim 230\text{K}(\text{bulk})$	AFM, $T_N \sim 620\text{K}$
II-II-V	(Nature Comm. 2013, IOP)	
LaZnAsO	---- La _{1-y} Ba _y Zn _{1-x} Mn _x AsO -----	LaMnAsO
“1111”	FM, $T_c \sim 40\text{K}(\text{bulk})$	AFM, $T_N \sim 320\text{K}$
	(PRB 2013, Zhejiang)	
“32522”	---- (Sr ₃ La ₂ O ₅)(Zn _{1-x} Mn _x) ₂ As ₂ -----	AFM(?)
	FM, $T_c \sim 40\text{K}(\text{bulk})$	
	(EPL 2014, Zhejiang)	
“42622”	---- (Sr _{1-x} K _x) ₄ Ti ₂ O ₆ (Zn _{1-x} Mn _x) ₂ As ₂ -----	AFM(?)
	FM, $T_c \sim 25\text{K}(\text{bulk})$	
	(To be submitted, Zhejiang)	

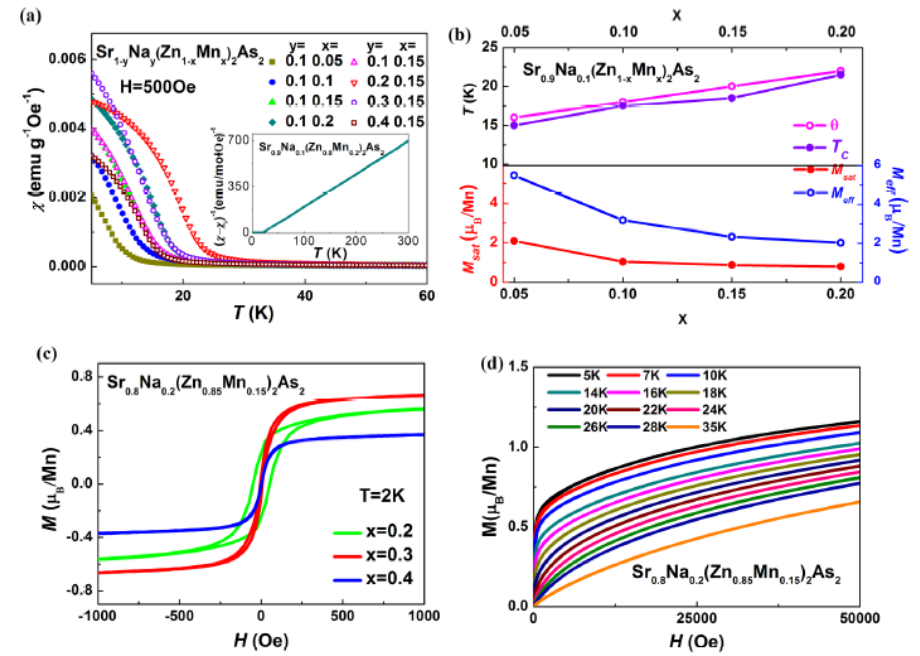
(Ca,Na)(Zn,Mn)₂As₂: A new spin and charge doping decoupled diluted ferromagnetic semiconductor

K. Zhao,¹ B. J. Chen,¹ Z. Deng,¹ W. Han,^{1,2} G. Q. Zhao,¹ J. L. Zhu,¹ Q. Q. Liu,¹ X. C. Wang,¹ B. Frandsen,³ L. Liu,³ S. Cheung,³ F. L. Ning,⁴ T. J. S. Munsie,⁵ T. Medina,⁵ G. M. Luke,⁵ J. P. Carlo,⁶ J. Munevar,⁷ G. M. Zhang,⁸ Y. J. Uemura,³ and C. Q. Jin^{1,3,a}



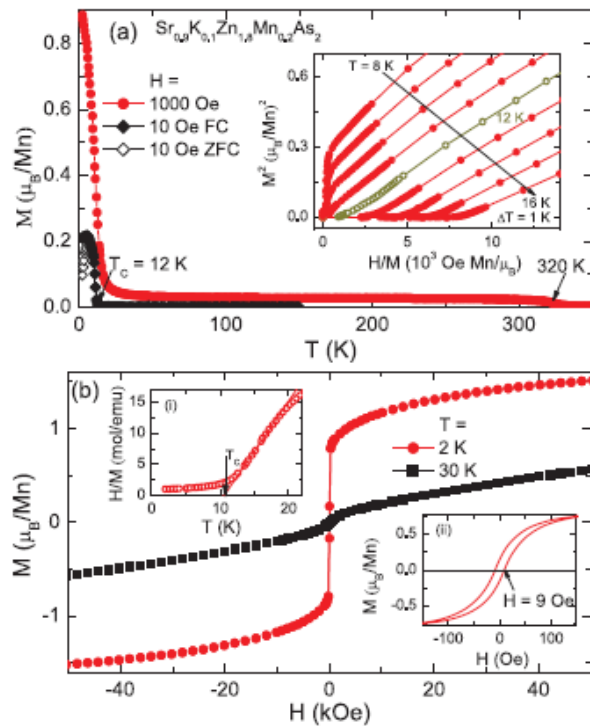
(Sr,Na)(Zn,Mn)₂As₂: A diluted ferromagnetic semiconductor with the hexagonal CaAl₂Si₂ type structure

B. J. Chen,¹ K. Zhao,¹ Z. Deng,¹ W. Han,¹ J. L. Zhu,¹ X. C. Wang,¹ Q. Q. Liu,¹ B. Frandsen,² L. Liu,² S. Cheung,² F. L. Ning,³ T. J. S. Munsie,⁴ T. Medina,⁴ G. M. Luke,⁴ J. P. Carlo,⁵ J. Munevar,⁶ Y. J. Uemura,² and C. Q. Jin¹



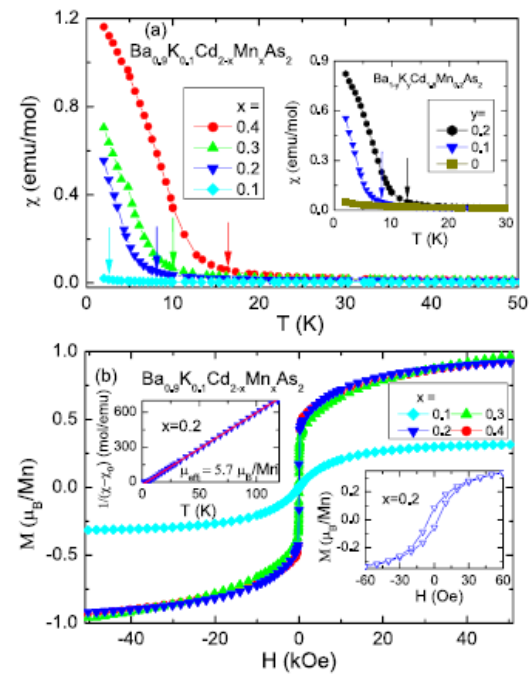
Sr_{0.9}K_{0.1}Zn_{1.8}Mn_{0.2}As₂: A ferromagnetic semiconductor with colossal magnetoresistance

XIAOJUN YANG¹, QIAN CHEN¹, YUPENG LI¹, ZHEN WANG¹, JINKE BAO¹, YUKE LI², QIAN TAO¹, GUANGHAN CAO¹ and ZHU-AN XU^{1(a)}



K and Mn co-doped BaCd₂As₂: A hexagonal structured bulk diluted magnetic semiconductor with large magnetoresistance

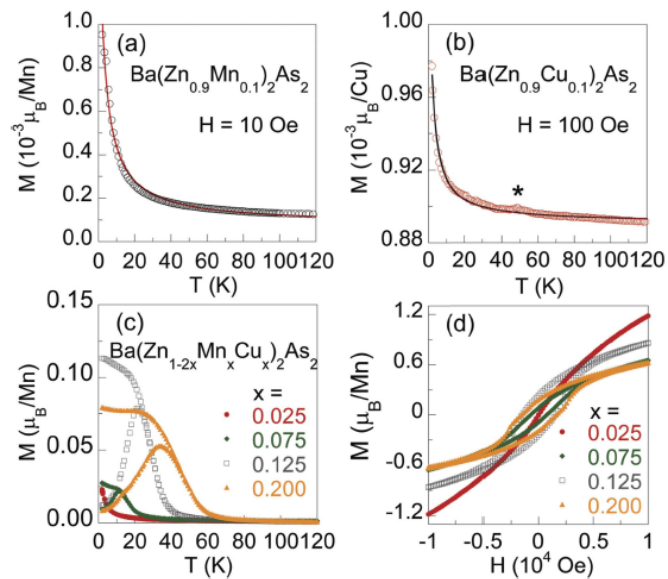
Xiaojun Yang,¹ Yuke Li,² Pan Zhang,¹ Hao Jiang,¹ Yongkang Luo,¹ Qian Chen,¹ Chunmu Feng,¹ Chao Cao,² Jianhui Dai,² Qian Tao,¹ Guanghan Cao,¹ and Zhu-An Xu^{1(a)}
¹Department of Physics and State Key Laboratory of Silicon Materials, Zhejiang University, Hangzhou 310027, China
²Department of Physics, Hangzhou Normal University, Hangzhou 310036, China



OPEN **Ba(Zn_{1-2x}Mn_xCu_x)₂As₂: A Bulk Form Diluted Ferromagnetic Semiconductor with Mn and Cu Codoping at Zn Sites**

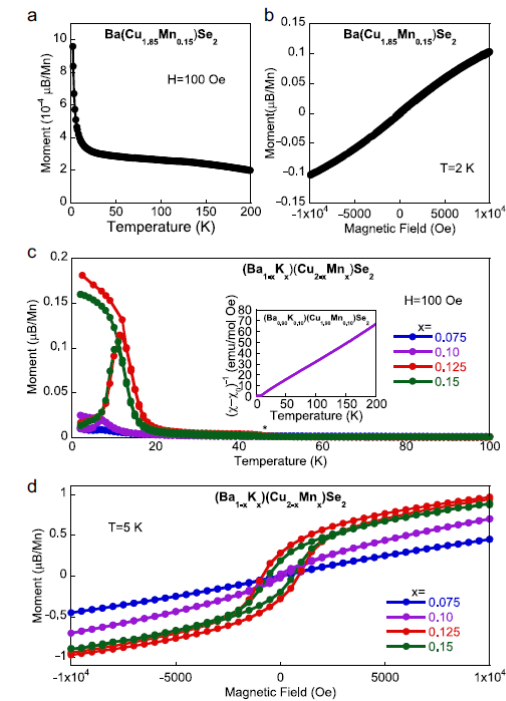
Received: 07 April 2015
Accepted: 28 September 2015
Published: 23 October 2015

Huiyuan Man^{1,2}, Shengli Guo^{1,2}, Yu Su³, Yang Guo⁴, Bin Chen⁴, Hangdong Wang⁴, Cui Ding^{1,2} & F.L. Ning^{1,2}

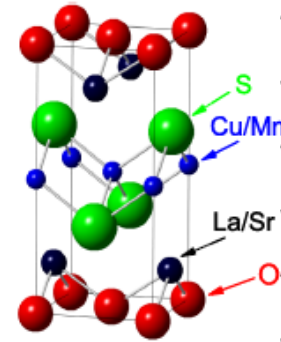
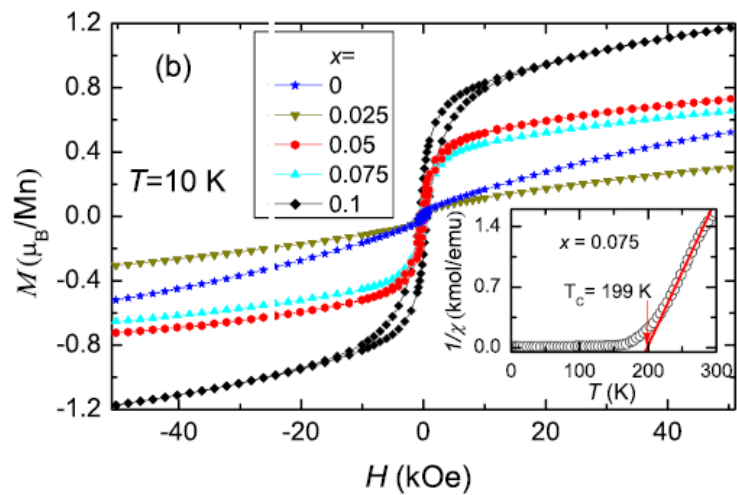
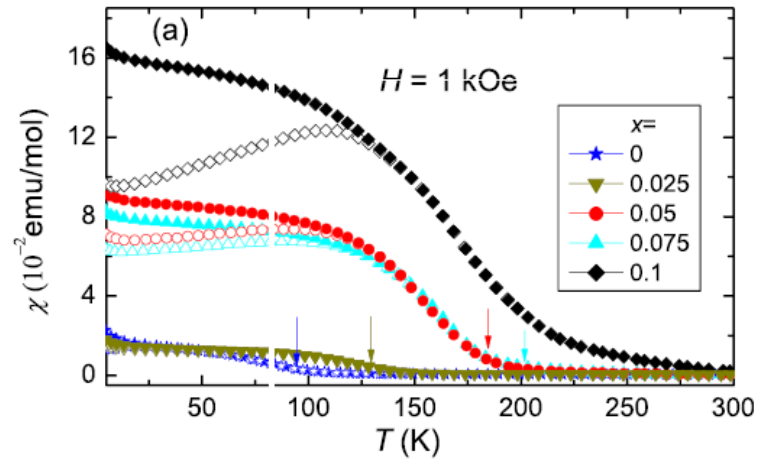


(Ba_{1-x}K_x)(Cu_{2-x}Mn_x)Se₂: A copper-based bulk form diluted magnetic semiconductor with orthorhombic BaCu₂S₂-type structure

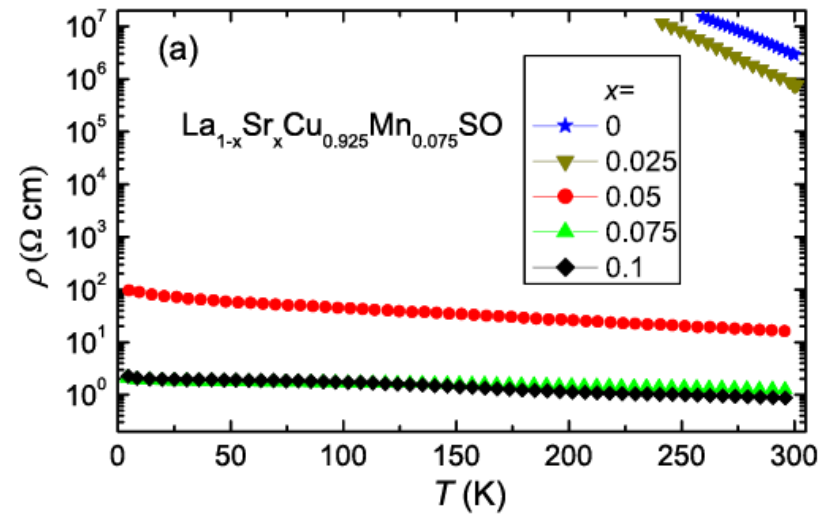
Shengli Guo^{a,b}, Huiyuan Man^{a,b}, Xin Gong^{a,b}, Cui Ding^{a,b}, Yao Zhao^{a,b}, Bin Chen^c, Yang Guo^c, Hangdong Wang^c, F.L. Ning^{a,b,*}



Sr and Mn co-doped LaCuSO₄: “1111”- type DMS

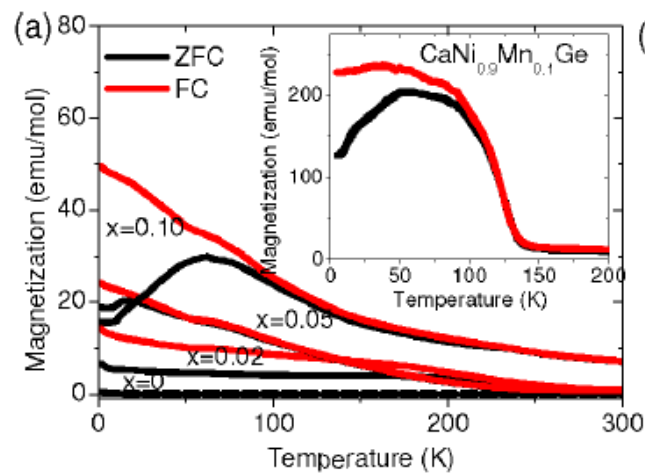


$T_C \sim 200 \text{ K}$

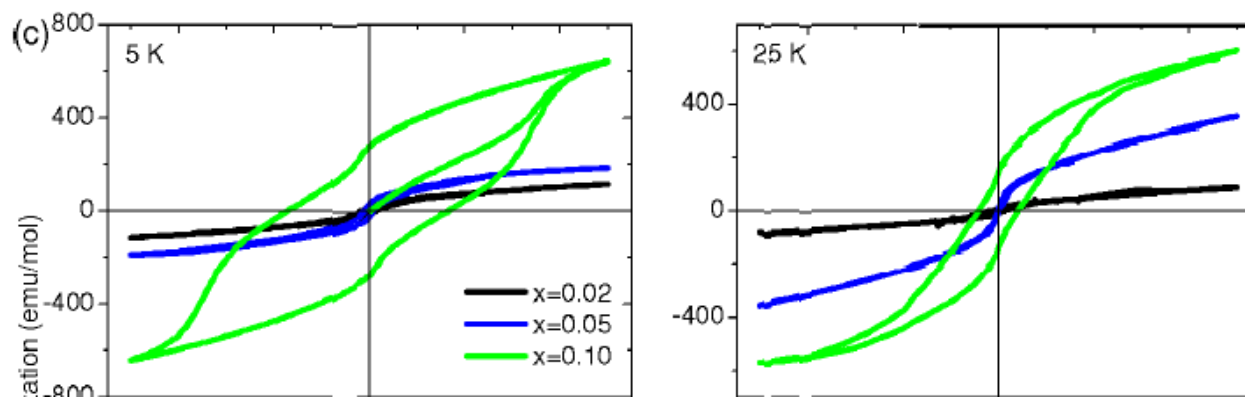
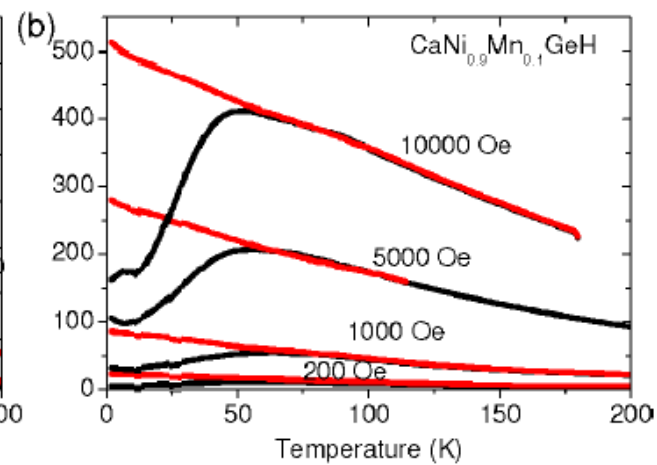


Z.A. Xu et al, Applied Physics Letters 103, 022410 (2013)

CaNi_{1-x}Mn_xGe



CaNi_{1-x}Mn_xGeH



H. Hosono et al, JACS, 134, 11687 (2012)

H. Hosono et al, Physical Review B **84**, 214439 (2011)

Ga_{1-x}Mn_xAs T_c ~ 200K(film)

LiZn_{1-x}Mn_xAs T_c ~ 55K(bulk)-----"111"

Ba_{1-y}K_y(Zn_{1-x}Mn_x)₂As₂ T_c ~ 230K(bulk)---"122"



P-type DMS

La_{1-y}Ba_yZn_{1-x}Mn_xAsO T_c ~ 40K(bulk) ---"1111"

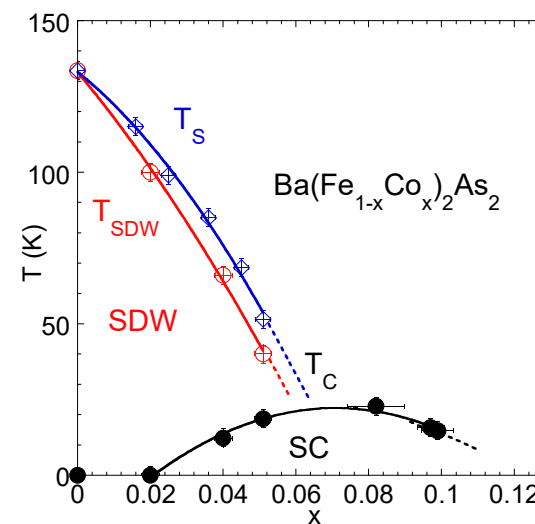
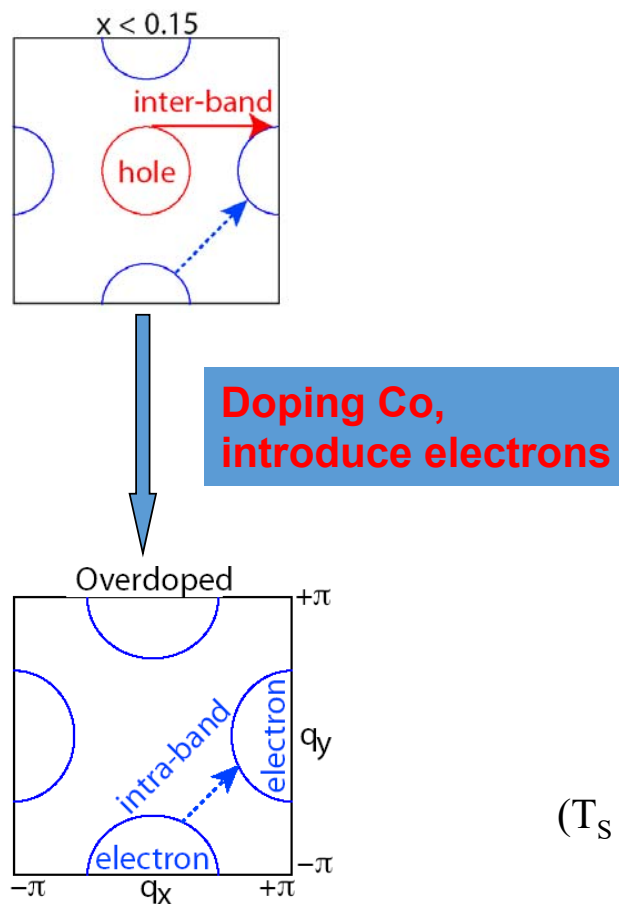
(Sr₃La₂O₅)(Zn_{1-x}Mn_x)₂As₂ T_c ~ 40K(bulk)---"32522"

(Sr_{1-x}K_x)₄Ti₂O₆(Zn_{1-x}Mn_x)₂As₂ T_c ~ 25K(bulk)---"42622"



For making junctions, n-type DMS are expected!

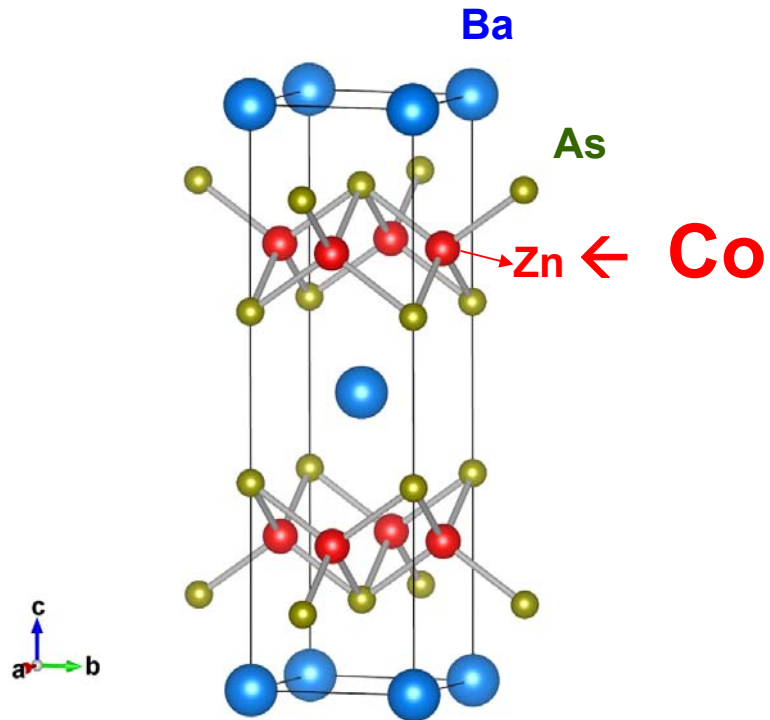
**“122” – $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ - SC
 $T_c = 25 \text{ K}$**



(F.L. Ning et al, JPSJ 78, 013711(2009))

(T_S is from Fisher et al, PRB 79, 014506(2009)).

“122” - BaZn_2As_2

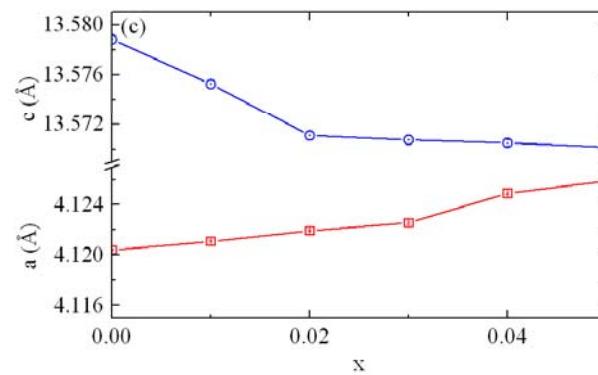
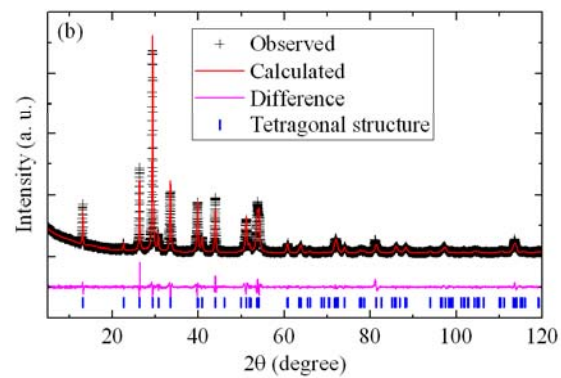
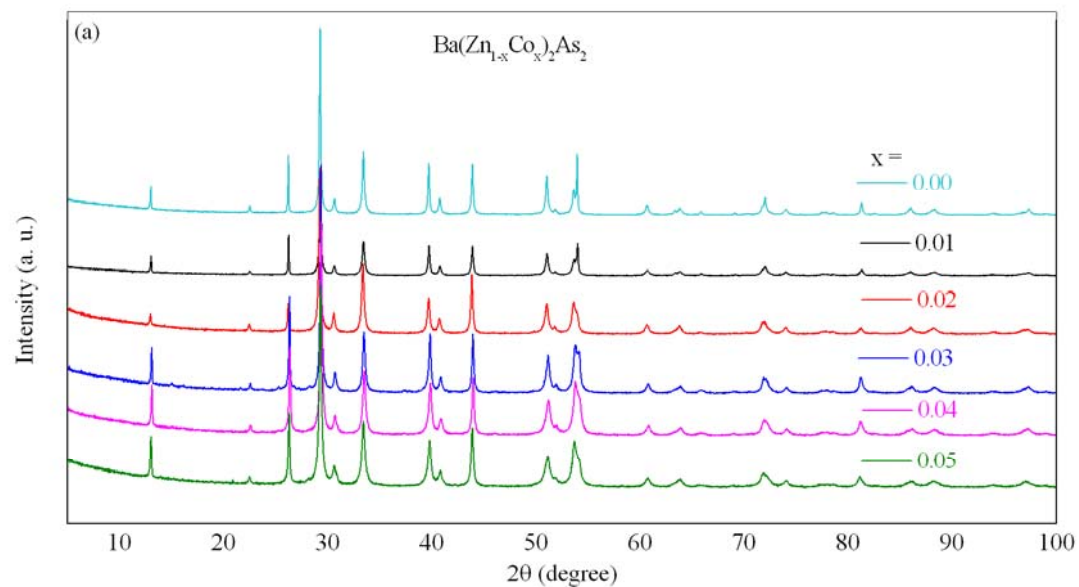
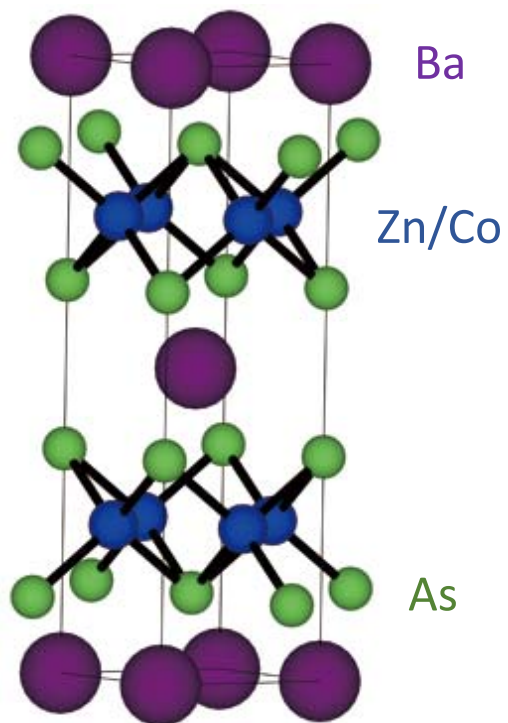


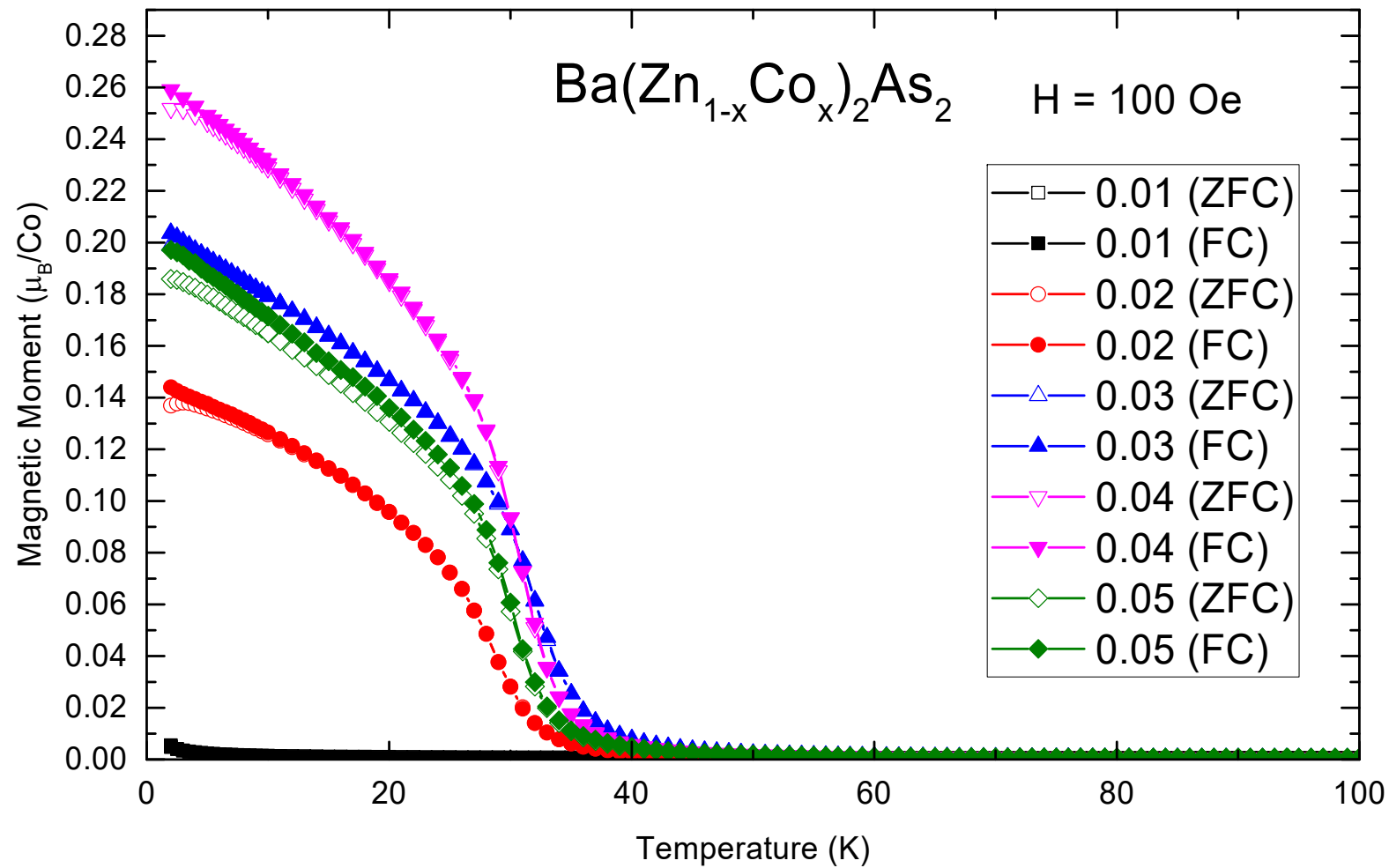
The idea here is:

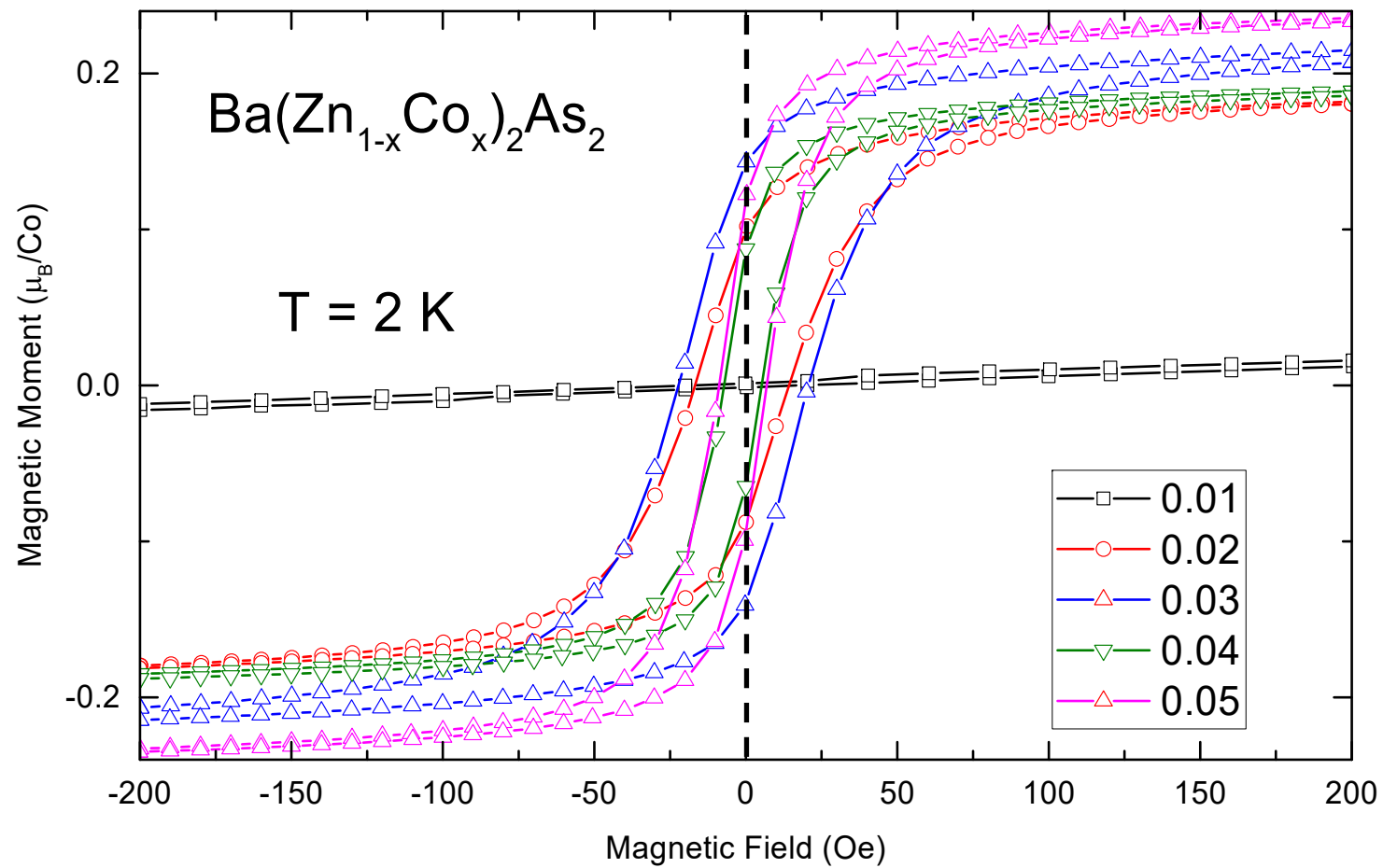
---Co introduce both moments and electrons

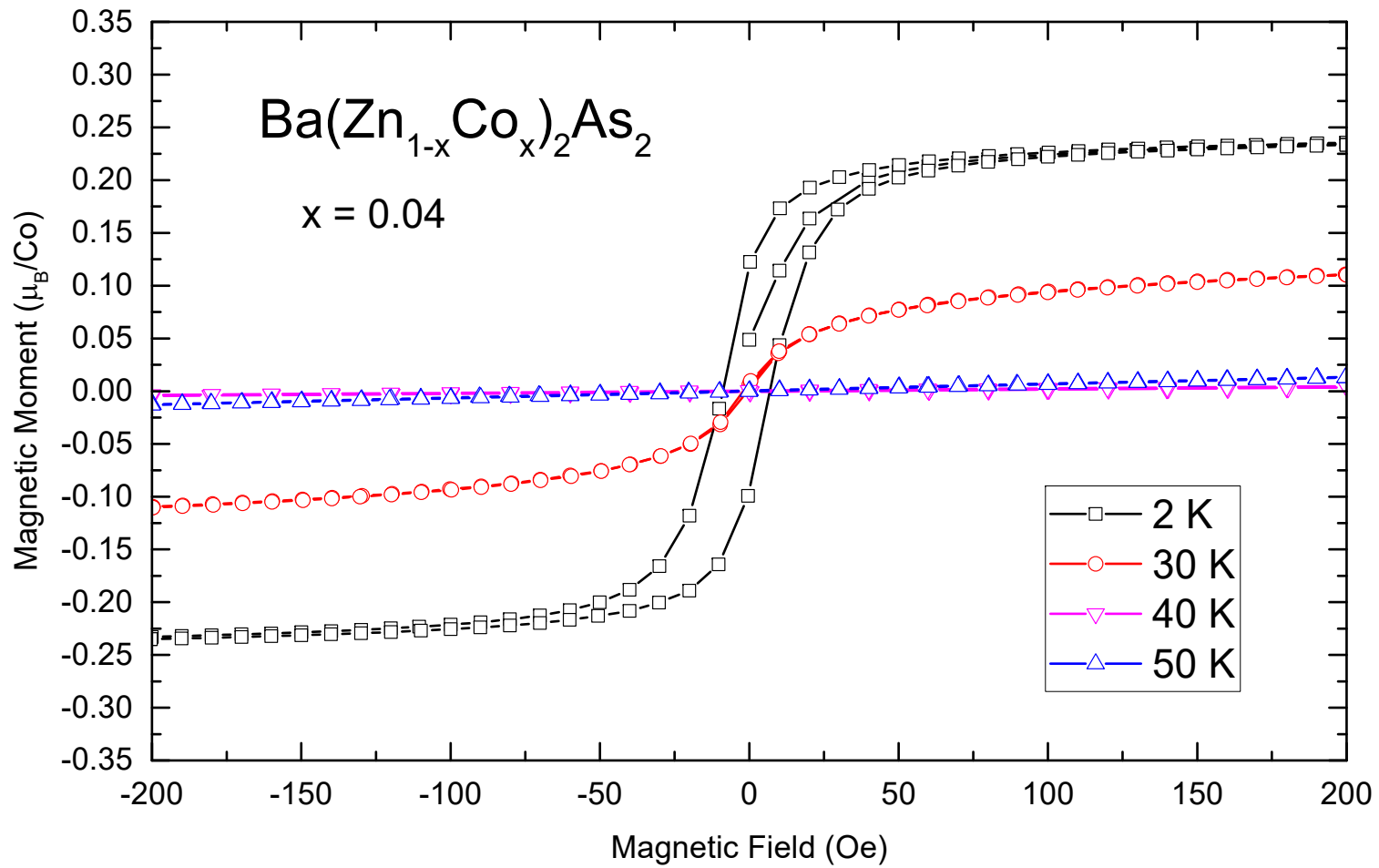
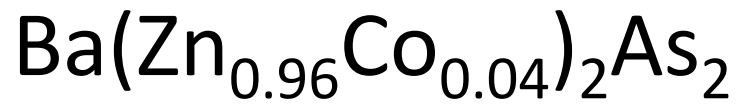
(Started the work in 2013)

Until 2016: Shengli Guo obtained pure phase

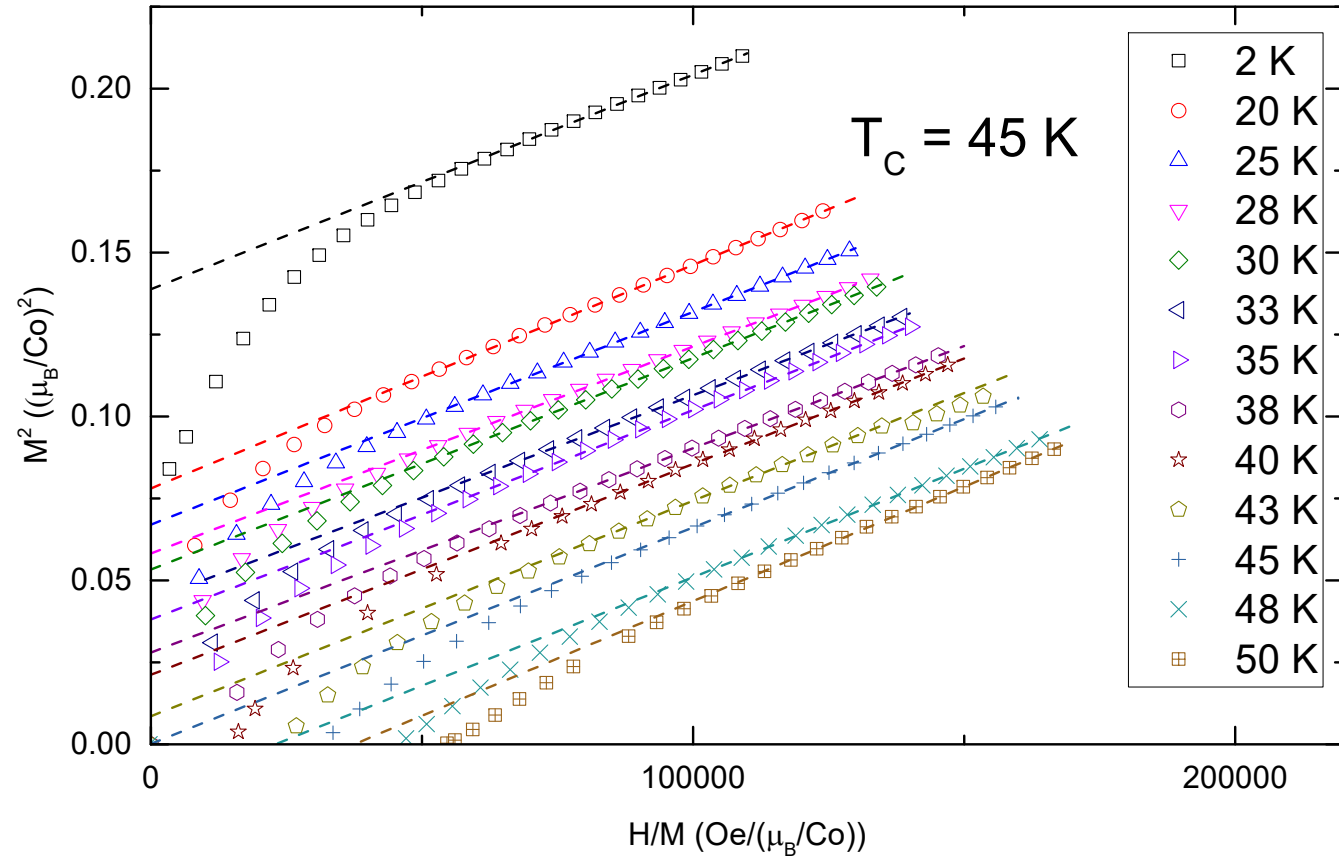
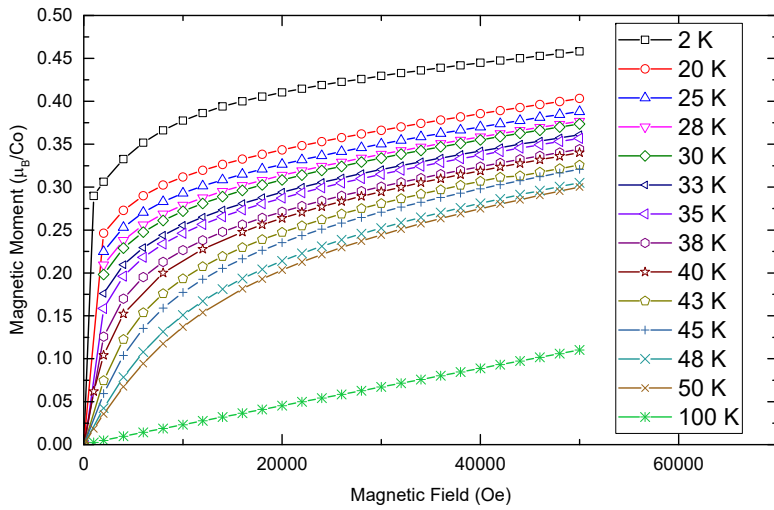


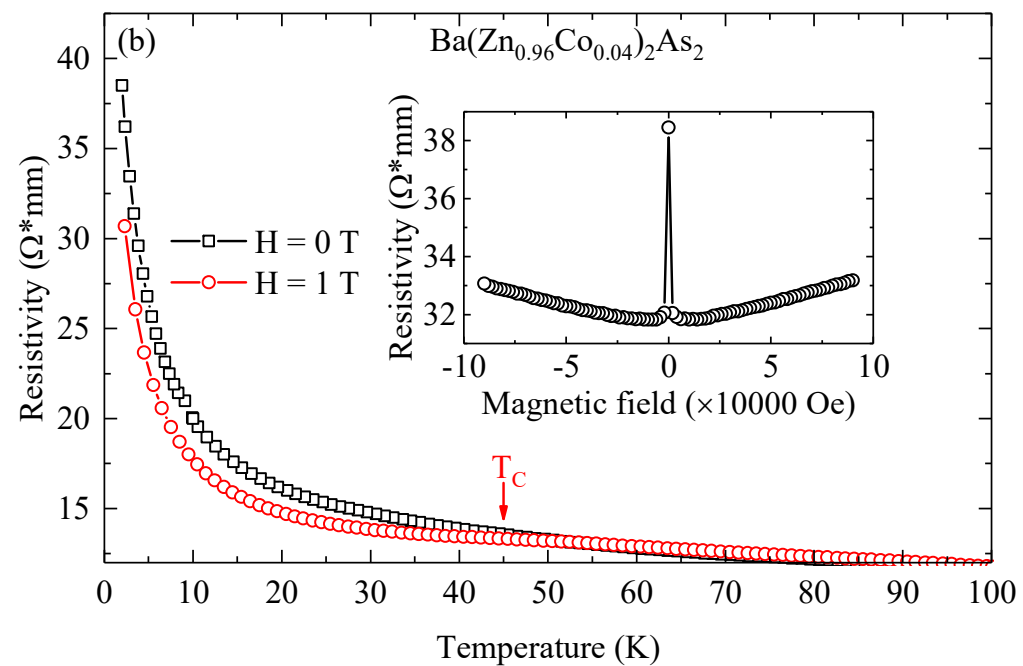
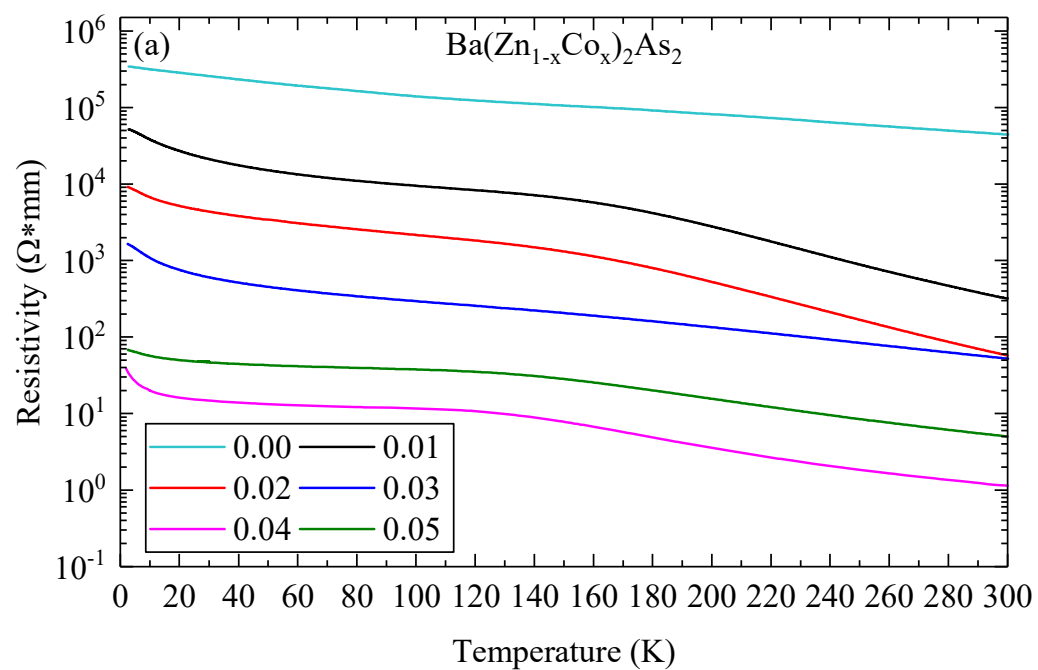


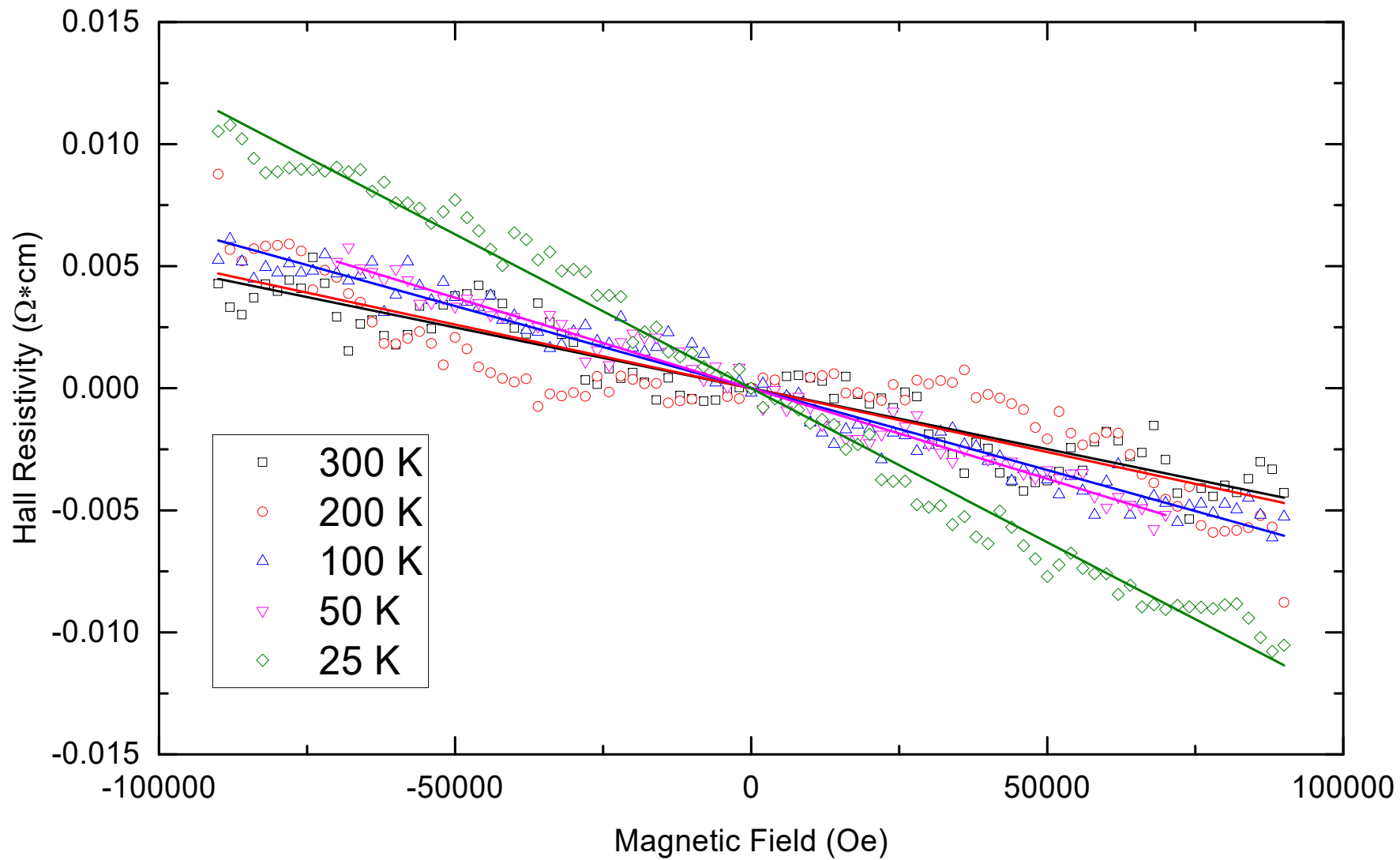




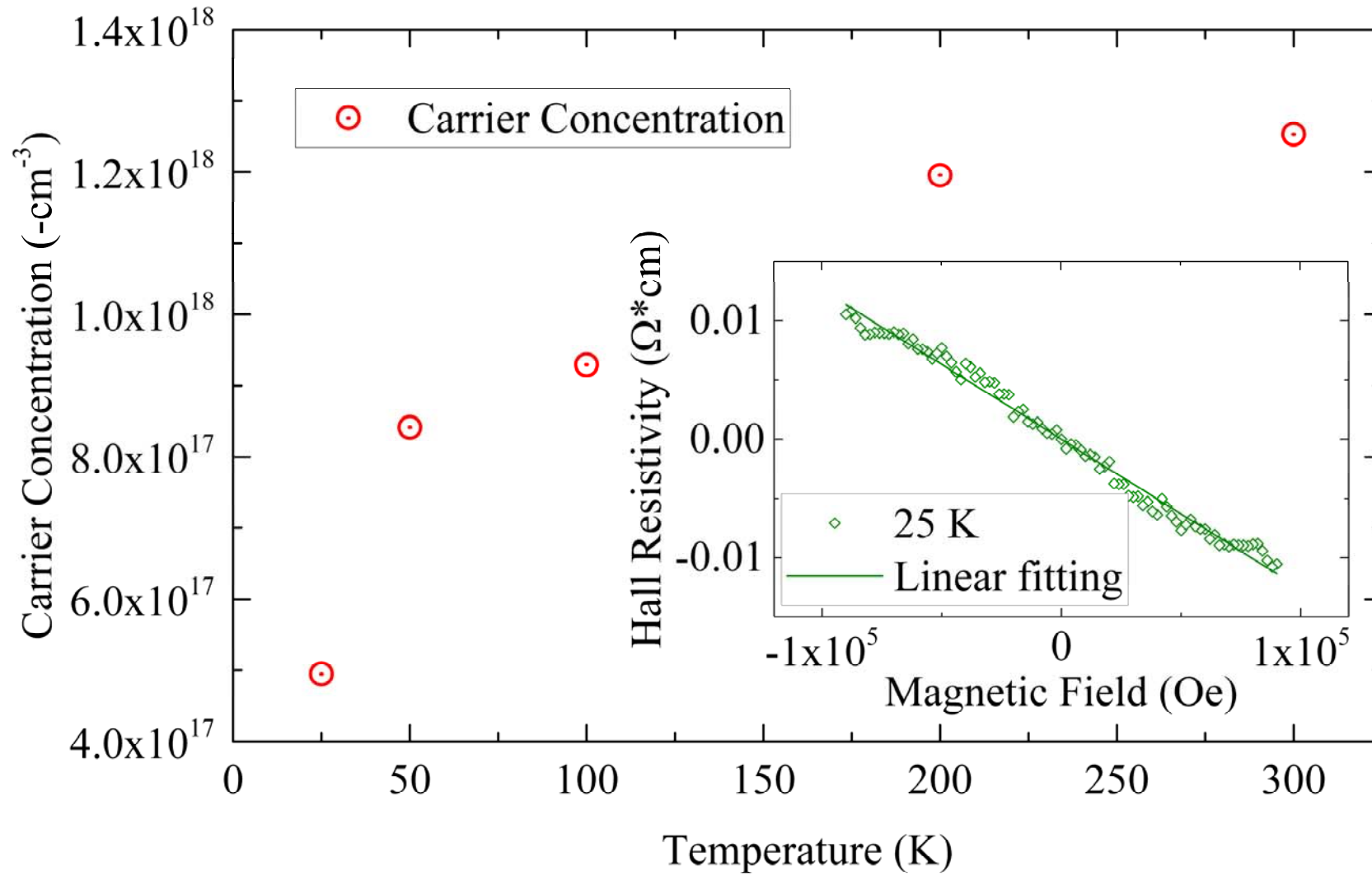
Ba(Zn_{0.96}Co_{0.04})₂As₂



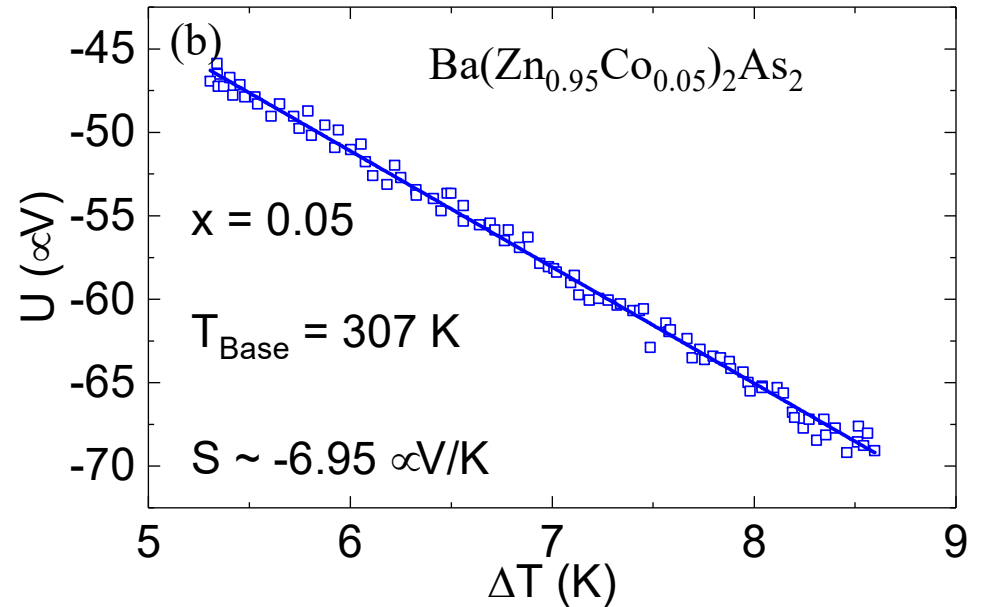
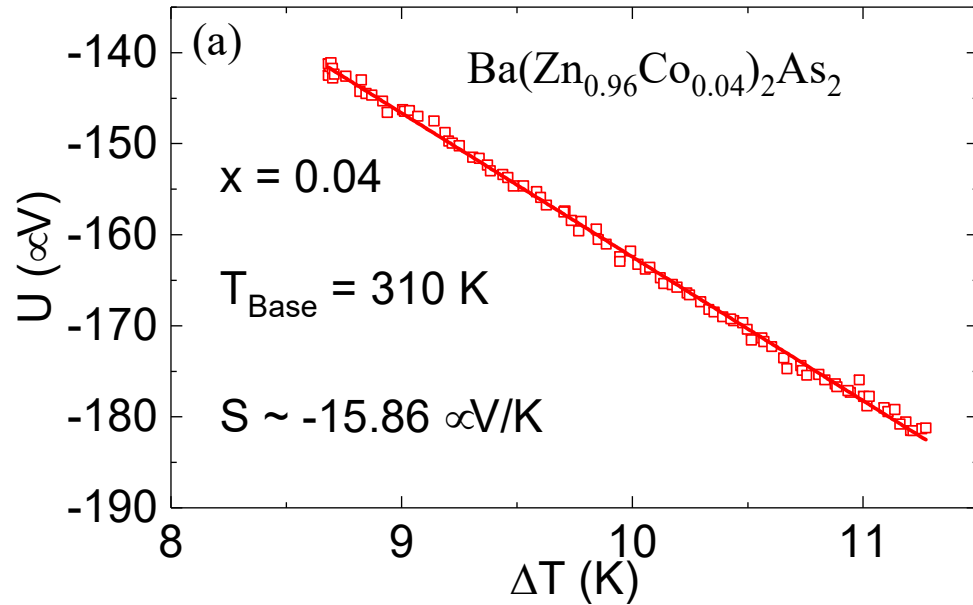




-n VS. T plot



Negative Seebeck coefficient



New Diluted Ferromagnetic Semiconductors (Zhejiang U)

1. $(\text{La}_{1-x}\text{Ba}_x)(\text{Zn}_{1-x}\text{Mn}_x)\text{AsO}$ ($T_C = 40\text{K}$) (PRB, 2013)
2. $(\text{SrLa}_2\text{O}_5)(\text{Zn}_{1-x}\text{Mn}_x)_2\text{As}_2$ ($T_C = 40\text{K}$) (EPL, 2014)
3. $\text{Li}_{1+y}(\text{Zn}_{1-x}\text{Cr}_x)\text{As}$ ($T_C = 210\text{K}$) (JAP, 2014)
4. $\text{Ba}(\text{Zn}_{1-2x}\text{Mn}_x\text{Cu}_x)_2\text{As}_2$ ($T_C = 70\text{K}$) (SR, 2015)
5. $(\text{Ba}_{1-x}\text{K}_x)(\text{Cu}_{1-x}\text{Mn}_x)\text{Se}_2$ ($T_C = 18\text{K}$) (JMMM, 2016)
6. $\text{Li}(\text{Zn}_{1-2x}\text{Mn}_x\text{Cu}_x)\text{As}$ ($T_C = 33\text{K}$) (JPCM, 2016)
7. $\text{La}(\text{Zn}_{1-2x}\text{Mn}_x\text{Cu}_x)\text{AsO}$ ($T_C = 8\text{K}$) (EPL, 2016)
8. $(\text{La}_{1-x}\text{Ca}_x)(\text{Zn}_{1-x}\text{Mn}_x)\text{AsO}$ ($T_C = 30\text{K}$) (JPCM, 2016)
9. $\text{La}(\text{Zn}_{1-2x}\text{Mn}_x\text{Cu}_x)\text{SbO}$ ($T_C = 15\text{K}$) (EPL, 2017)
10. $\text{Ba}(\text{Zn}_{1-x}\text{Co}_x)_2\text{As}_2$ ($T_C = 40\text{K}$) (N-type, PRB, 2019)
11. $(\text{Sr}_{1-x}\text{K}_x)\text{F}(\text{Zn}_{1-x}\text{Mn}_x)\text{Sb}$ ($T_C = 40\text{K}$) (JMMM, 2019)
12. $\text{Cu}_2(\text{Zn}_{1-x}\text{Mn}_x)(\text{Sn}_{1-y}\text{Al}_y)\text{Se}_4$ ($T_C = 5\text{K}$) (Submitted, 2019)
13. $(\text{Ba}_{1-x}\text{K}_x)(\text{Zn}_{1-y}\text{Mn}_y)_2\text{Sb}_2$ (In preparation, 2019)
14. $(\text{Ca}_{1-x}\text{Na}_x)(\text{Zn}_{1-y}\text{Mn}_y)_2\text{Sb}_2$ ($T_C = 10\text{K}$) (In preparation, 2019)

**In past several years, our collaboration team (Uemura & Jin & Ning)
published ~ 40 papers on DMS**

- 1, S.L. Guo et al, PRB 99, 155201 (2019)
- 2, L.C. Fu et al, JMMM 483, 95 (2019)
- 3, Y.L. Gu et al, JOS 40, 081506 (2019)
- 4, G.X. Zhi et al, Cond Matt 3, 42(2018)
- 5, S.L. Guo et al, CPB 27, 097502(2018)
- 6, Yao Zhao et al, EPL 120, 47005 (2017)
- 7, F.F. Zhu et al, APL 111, 062106 (2017)
- 8 S.L. Guo et al, EPL 114, 57008 (2016)
- 9, S.L. Guo et al, JMMM 400, 295 (2016)
- 10, Cui Ding et al, JPCM 28, 026003 (2016)
- 11, S.L. Guo et al, JPCM 28, 366001 (2016)
- 12, H.Y. Man et al, Scientific Reports 5, 15507 (2015)
- 13, Cui Ding et al, Science China (2014)
- 14, Cui Ding et al, EPL 107, 17004 (2014)
- 15, F.L. Ning et al, PRB 90, 185023 (2014)
- 16, H.Y. Man et al, EPL 105, 67004 (2014)
- 17, Q. Wang et al, JAP 115, 083917 (2014)
- 18, J.C. Lu et al, EPL 103, 67001 (2013)
- 19, C. Ding et al PRB 88, 041102(R) (2013)
- 20, C. Ding et al PRB 88, 041108(R) (2013)
- 21, W. Han et al, Scientific Reports 9, 7490 (2019)
- 22, Y. Peng et al, CPB 28, 057501 (2019)
- 23, G.Q. Zhao et al, JPCM 30, 254001 (2018)
- 24, G.X. Gu et al, APL 112, 032402 (2018)
- 25, Sun Fei et al, CPL 34, 0256 (2017)
- 26, R. Wang et al, AIP Advances 7, 045017 (2017)
- 27, Sun Fei et al, PRB 95, 094412 (2017)
- 28, B.J.Chen et al, Scientific Reports 6, 36578 (2016)
- 29, B.Frandsen and Uemura et al, PRB 94, 094102 (2016)
- 30, F. Sun et al, PRB 93, 224403 (2016)
- 31, H. Suzuki and Fujimori et al, PRB 92, 215320 (2015)
- 32, H. Suzuki and Fujimori et al, PRB 91, 140401 (2015)
- 33, K. Zhao et al, Chin. Sci. Bull. (2014)
- 34, B.J. Chen et al PRB 90, 155202 (2014)
- 35, B.J. Chen et al JAP 116,163906 (2014)
- 36, W. Han et al, Science China (2013)
- 37, Z. Deng et al, PRB 88, 081203(R) (2013)
- 38, K. Zhao et al, Nature Comm. 4, 1442 (2013)
- 39, Z. Deng et al, JPCS 400, 032033 (2012)
- 40, Z. Deng et al, Nature Comm. 2, 422 (2011)

**Photoemission and X-ray Absorption Studies of the Diluted Magnetic Semiconductor
Ba_{1-y}K_y(Zn_{1-x}Mn_x)₂As₂ Isostructural to Fe-based Superconductors**

H. Suzuki¹, K. Zhao², G. Shibata¹, Y. Takahashi¹, S. Sakamoto¹, K. Yoshimatsu¹, B. J. Chen², H. Kumigashira³, F.-H. Chang⁴, H.-J. Lin⁴, D. J. Huang⁴, C. T. Chen⁴, Bo Gu⁵, S. Maekawa⁵, Y. J. Uemura⁶, C. Q. Jin² and A. Fujimori¹

Theory of Mn-doped I-II-V Semiconductors

J. K. Glasbrenner,¹ I. Žutić,² and I. I. Mazin³

¹ *National Research Council/Code 6393, Naval Research Laboratory, Washington, DC 20375, USA*

² *Department of Physics, University at Buffalo, State University of New York, NY 14260, USA and*

³ *Code 6393, Naval Research Laboratory, Washington, DC 20375, USA*

(Dated: May 13, 2014)

Re: APS Invited Symposium Nomination Notification

From: Jacek K. Furdyna <furdyna@nd.edu>
Time: 22:02:31 Sep 14, 2014 (Sunday)
To: ningfl@zju.edu.cn, furdyna.1@nd.edu

Dear Fanlong,

It was a real pleasure to be able to nominate your beautiful and ground-breaking work at the APS March APS accepts this nomination. It should be a very strong

And I hope we get a chance to meet sometime in the future.

With my best regards,
Jacek

Re: APS Invited Symposium Nomination Notification

From: zigor <zigor@buffalo.edu>
Time: 01:11:15 Sep 10, 2014 (Wednesday)
To: ningfl@zju.edu.cn

My pleasure Fanlong,

there will be lots of politics involved, but it was certainly worth trying for such an exciting work.

Regards,
Igor

P- and N-type diluted magnetic semiconductors with narrow band gaps

Bo Gu ^{1*} and Sadamichi Maekawa^{1,2}

¹ *Advanced Science Research Center, Japan Atomic Energy Agency, Tokai 319-1195, Japan*

² *ERATO, Japan Science and Technology Agency, Sendai 980-8577, Japan*

(Dated: January 19, 2016)

We propose a method to realize diluted magnetic semiconductors (DMS) with p- and n-type carriers by choosing host semiconductors with a narrow band gap. By employing a combination of the density function theory and quantum Monte Carlo simulation, we demonstrate such semiconductors using Mn-doped BaZn₂As₂, which has a band gap of 0.2 eV. In addition, we found a new non-toxic DMS Mn-doped BaZn₂Sb₂, of which the Curie temperature T_c is predicted to be higher than that of Mn-doped BaZn₂As₂, the T_c of which was up to 230 K in the recent experiment.

PACS numbers: 75.50.Pp, 75.30.Hx, 02.70.Ss

After the discovery of ferromagnetism in (Ga,Mn)As, diluted magnetic semiconductors (DMS) have received considerable attention owing to potential applications based on the use of both their charge and spin degrees of freedom in electronic devices [1, 2]. Thus far, the highest Curie temperature of (Ga,Mn)As has been $T_c = 190$ K [3]. The substitution of divalent Mn atoms into trivalent Ga sites introduces hole carriers; thus, (Ga,Mn)As is a p-type DMS. The valence mismatch between Mn and Ga leads to severely limited chemical solubility for Mn in GaAs. Moreover, owing to simultaneous doping of charge and spin induced by Mn substitution, it is difficult to individually optimize charge and spin densities.

To overcome these difficulties, a new type of DMS, i.e., Li(Zn,Mn)As was proposed [4] and later fabricated with $T_c = 50$ K [5]. It is based on LiZnAs, a I–II–V semiconductor. Spin is introduced by isovalent (Zn²⁺, Mn²⁺) substitution, which is decoupled from carrier doping with excess/deficient Li concentration. Although Li(Zn,Mn)As was proposed as a promising p-type DMS

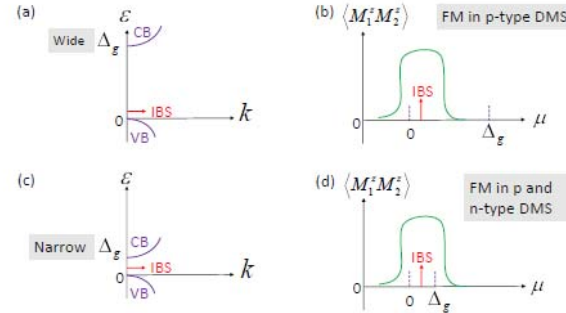


FIG. 1. (Color online) Schematic pictures of impurity bound state (IBS) and ferromagnetic (FM) coupling in diluted magnetic semiconductors (DMS). (a) Host bands $\epsilon(k)$ with a wide band gap Δ_g between the valence band (VB) and the conduction band (CB). The position of the IBS ω_{IBS} (arrow) is close to the top of the VB owing to strong mixing between the impurity and the VB, and usually no IBS appears below the top of the CB. (b) Plot of the magnetic moment $\langle M_i M_i \rangle$ vs chemical potential μ for FM in p-type DMS. (c) Host bands $\epsilon(k)$ with a narrow band gap Δ_g between the VB and the CB. The IBS (arrow) is located near the top of the VB. (d) Plot of the magnetic moment $\langle M_i M_i \rangle$ vs chemical potential μ for FM in p and n-type DMS.

Diluted Ferromagnetic Semiconductors in Bulk Form with decoupling spins and carriers

REVIEWS OF MODERN PHYSICS, VOLUME 86, JANUARY–MARCH 2014

Dilute ferromagnetic semiconductors: Physics and spintronic structures

By Tomasz Dietl and Hideo Ohno

In addition to current interest in various magnetically doped semiconductors, oxides, and organic materials, a lot of attention will be devoted to four emerging families of compounds: (i) high Curie temperature ferrimagnetic spinel oxides and Heusler compounds such as Mn_2CoAl (Ouardi *et al.*, 2013), awaiting for mastering of defect and carrier control; (ii) high Néel temperature semiconductors, e.g., LiMnAs , for antiferromagnetic spintronics (Jungwirth *et al.*, 2011); (iii) topological insulators, in which ferromagnetism might be mediated by Dirac electrons, e.g., $(\text{Bi}, \text{Mn})_2\text{Te}_3$ (Checkelsky *et al.*, 2012); and (iv) derivatives of FeAs

superconductors, such as $(\text{K}, \text{Ba})(\text{Zn}, \text{Mn}, \text{Fe})_2\text{As}_2$ compounds for studies of interplay between p - d Zener ferromagnetism, antiferromagnetic superexchange, and superconductivity (Zhao *et al.*, 2013). Here nanocharacterization protocols, elaborated over the recent years for DMSs (Bonanni, 2011), will play an essential role in the meaningful development of new materials.

Uemura, Jin, Ning started the research topic.

Dietl and Ohno, RMP, 86, 187 (2014): “a lot of attention will be devoted to four emerging families of compounds, ..., **IV: FeAs derived compound**, ..., for studies of interplay between p - d Zener ferromagnetism, antiferromagnetic, superexchange, and superconductivity.”

Roadmap for Emerging Materials for Spintronic Device Applications

Atsufumi Hirohata¹, Hiroaki Sukegawa², Hideto Yanagihara³, Igor Žutić⁴, Takeshi Seki⁵, Shigemi Mizukami⁶, and Raja Swaminathan⁷

¹Department of Electronics, University of York, York YO10 5DD, U.K.

²Magnetic Materials Unit, National Institute for Materials Science, Tsukuba 305-0047, Japan

³Graduate School of Pure and Applied Sciences, University of Tsukuba, Tsukuba 305-8577, Japan

⁴Department of Physics, University at Buffalo–The State University of New York, Buffalo, NY 14260 USA

⁵Institute for Materials Research, Tohoku University, Sendai 980-8577, Japan

⁶WPI Advanced Institute for Materials Research, Tohoku University, Sendai 980-8577, Japan

⁷Intel Corporation, Chandler, AZ 85226 USA

The Technical Committee of the IEEE Magnetism Society has selected seven research topics to develop their roadmaps, where major developments should be listed alongside expected timelines: 1) hard disk drives; 2) magnetic random access memories; 3) domain-wall devices; 4) permanent magnets; 5) sensors and actuators; 6) magnetic materials; and 7) organic devices. Among them, magnetic materials for spintronic devices have been surveyed as the first exercise. In this roadmap exercise, we have targeted magnetic tunnel and spin-valve junctions as spintronic devices. These can be used, for example, as a cell for a magnetic random access memory and a spin-torque oscillator in their vertical form as well as a spin transistor and a spin Hall device in their lateral form. In these devices, the critical role of magnetic materials is to inject spin-polarized electrons efficiently into a nonmagnet. We have accordingly identified two key properties to be achieved by developing new magnetic materials for future spintronic devices: 1) half-metallicity at room temperature (RT) and 2) perpendicular anisotropy in nanoscale devices at RT. For the first property, five major magnetic materials are selected for their evaluation for future magnetic/spintronic device applications: 1) Heusler alloys; 2) ferrites; 3) rutiles; 4) perovskites; and 5) dilute magnetic semiconductors. These alloys have been reported or predicted to be half-metallic ferromagnets at RT. They possess a bandgap at the Fermi level E_F only for its minority spins, achieving 100% spin polarization at E_F . We have also evaluated $L1_0$ alloys and $D0_{22}$ -Mn alloys for the development of a perpendicularly anisotropic ferromagnet with large spin polarization. We have listed several key milestones for each material on their functionality improvements, property achievements, device implementations, and interdisciplinary applications within 35 years time scale. The individual analyses and the projections are discussed in the following sections.

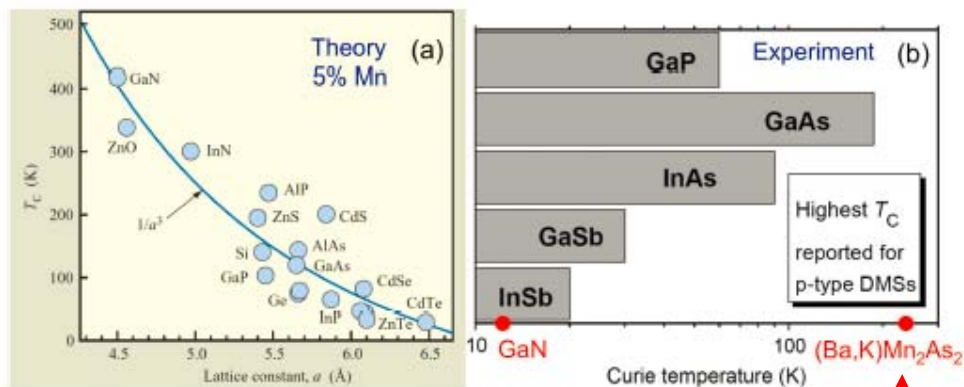
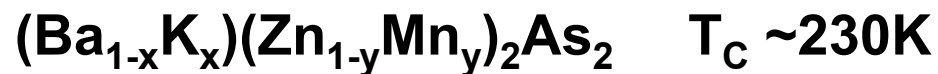


Fig. 8. (a) Theoretical predictions for T_C in DMS [41], adapted from [45]. (b) Reliable highest experimental T_C reported for Mn-doped DMS, adapted from [36].

some of the previous limitations are overcome: the absence of carriers in II–VI and the low-Mn solubility in III–Vs. With 30% K and 15% Mn doping, their $T_C \sim 230$ K [48] exceeds the value in (Ga, Mn)As. Selected highest reliable experimental T_C reported for the Mn-doped DMS is shown in Fig. 8(b). Circles are given for GaN, which has about 30 times smaller T_C than predicted in Fig. 8(a), and (Ba, K)Zn₂As₂, a current record for DMS. The *ab initio* studies predict a further increase in T_C [49]. We expect that the tunable RT carrier-mediated ferromagnetism will be realized in II–II–V DMS within 5 years.

Regarding (m3.2), while DMS is often viewed as the

II-II-V DMS



Uemura, Jin, Ning started the research topic.

Several advantages of bulk form DMS materials:

- Synthesized in thermally equilibrium condition, **less defects** are expected;
- **Bulk form, single crystals** are available, for NMR, neutron, μ SR measurements;
- Carriers can be doped **at different ionic sites**, for example for BaZn_2As_2 , K substitution for Ba introduces holes, while Mn substitution for Zn introduces electrons;
- Charge and spin doping are **decoupled**, each can be precisely controlled and tuned;
- **Iso-structural** to its superconductor, antiferromagnet variant, junctions via As layer can be made;
- **Thin films** of parent compound LaZnAsO and BaZn_2As_2 have been successfully made;
- $(\text{Ba}_{1-y}\text{K}_y)(\text{Zn}_{1-x}\text{Mn}_x)_2\text{As}_2$: $T_C = 230$ K for now, thin film may achieve **higher T_C ?** (FeSe-SC)
- Share **the same origin of ferromagnetism** as $(\text{Ga,Mn})\text{As}$, helpful to understand the general mechanism;

Outline

a. Introduction of some newly fabricated DMSs

“111”, “122”, “1111”, “42622”, “32522”

----- $\text{Li}_{1+y}(\text{Zn}_{1-x}\text{Mn}_x)\text{P}$

----- n-type $\text{Ba}(\text{Zn}_{1-2x}\text{Co}_x\text{Mn}_x)_2\text{As}_2$

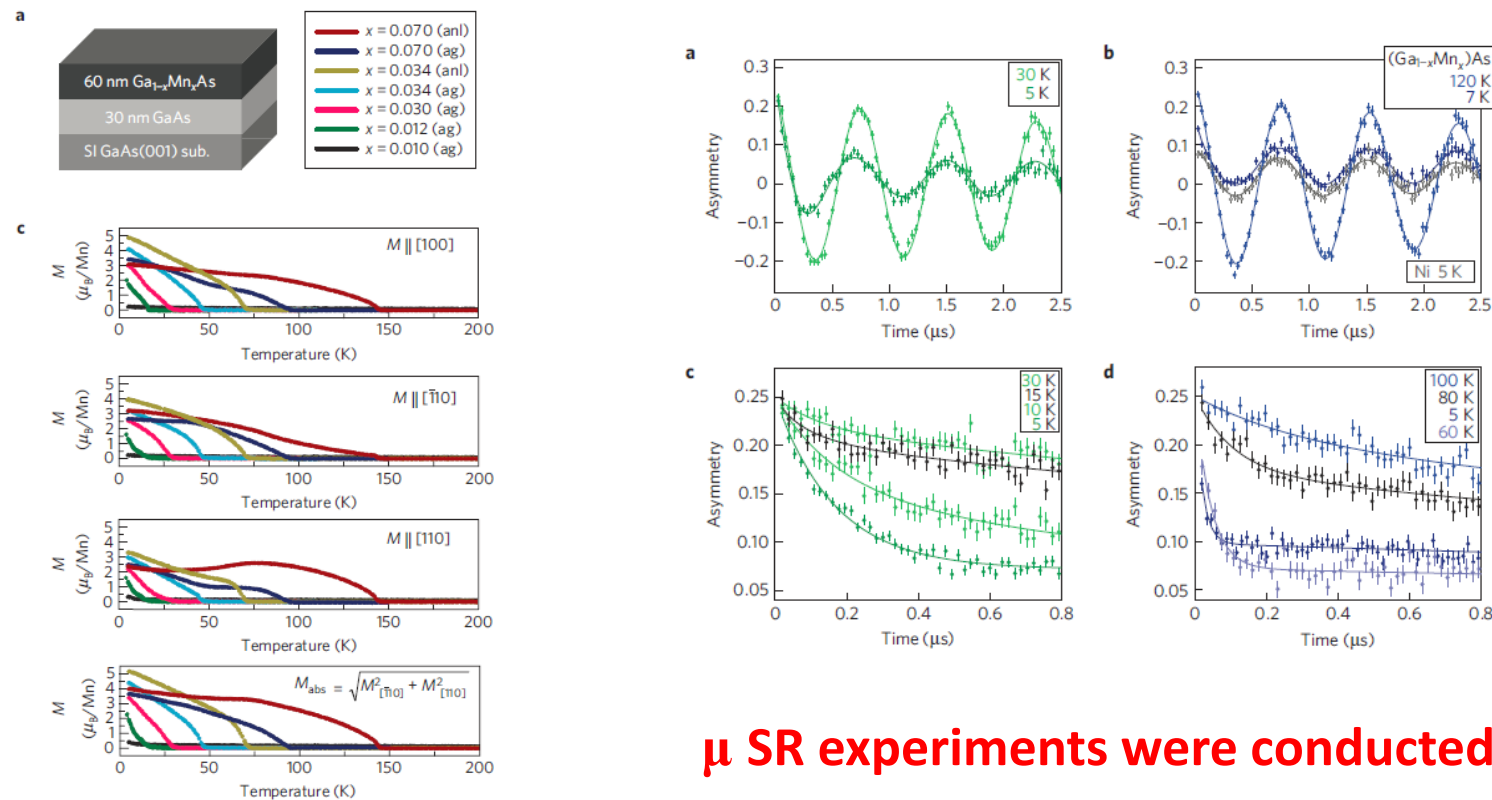
b. Microscopic characterization

----- μSR results on $(\text{Ga},\text{Mn})\text{As}$ and $\text{Li}_{1.15}(\text{Zn}_{1-x}\text{Mn}_x)\text{P}$

----- NMR results on $\text{Li}_{1.1}(\text{Cd}_{1-x}\text{Mn}_x)\text{P}$ and $\text{Li}_{1.15}(\text{Zn}_{1-x}\text{Mn}_x)\text{P}$

Spatially homogeneous ferromagnetism of (Ga, Mn)As

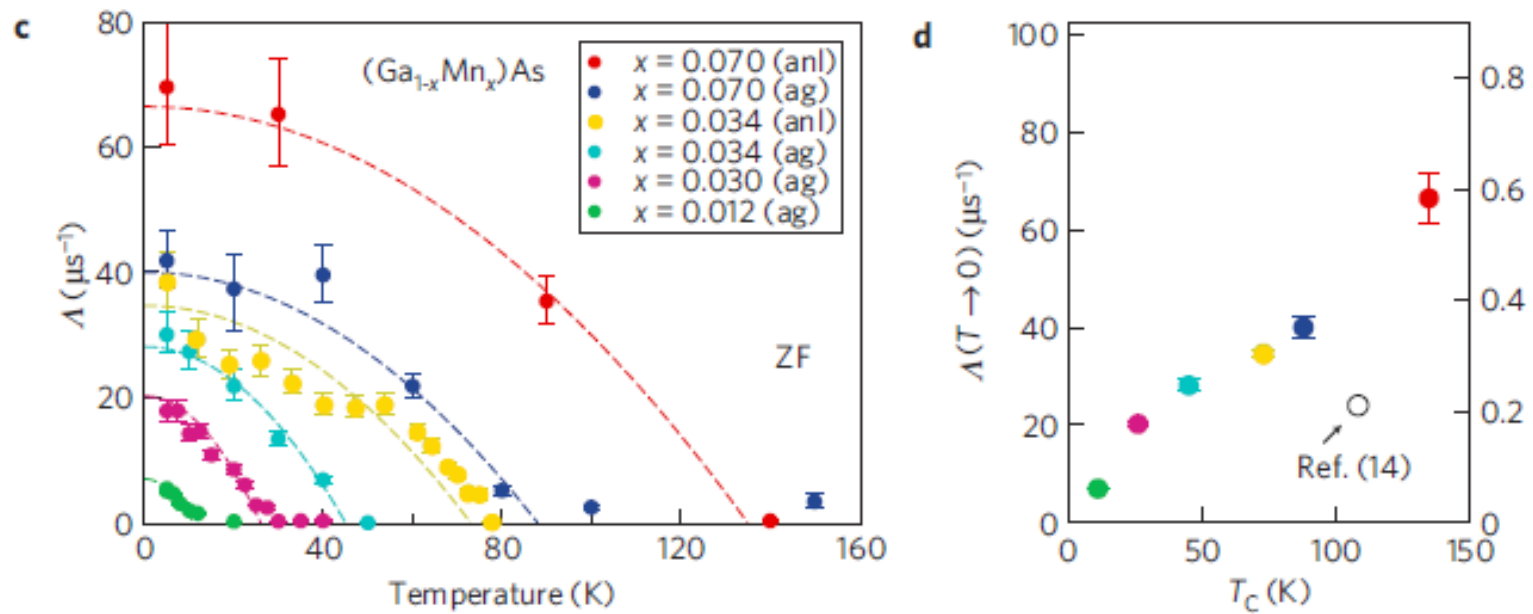
S. R. Dunsiger^{1,2}, J. P. Carlo¹, T. Goko^{1,3}, G. Nieuwenhuys⁴, T. Prokscha⁴, A. Suter⁴, E. Morenzoni⁴, D. Chiba^{5,6}, Y. Nishitani⁶, T. Tanikawa^{5,6}, F. Matsukura^{5,6}, H. Ohno^{5,6}, J. Ohe^{7,8}, S. Maekawa^{7,8} and Y. J. Uemura^{1*}



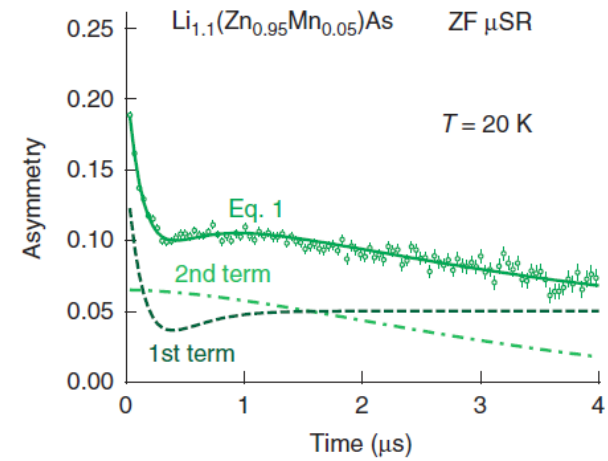
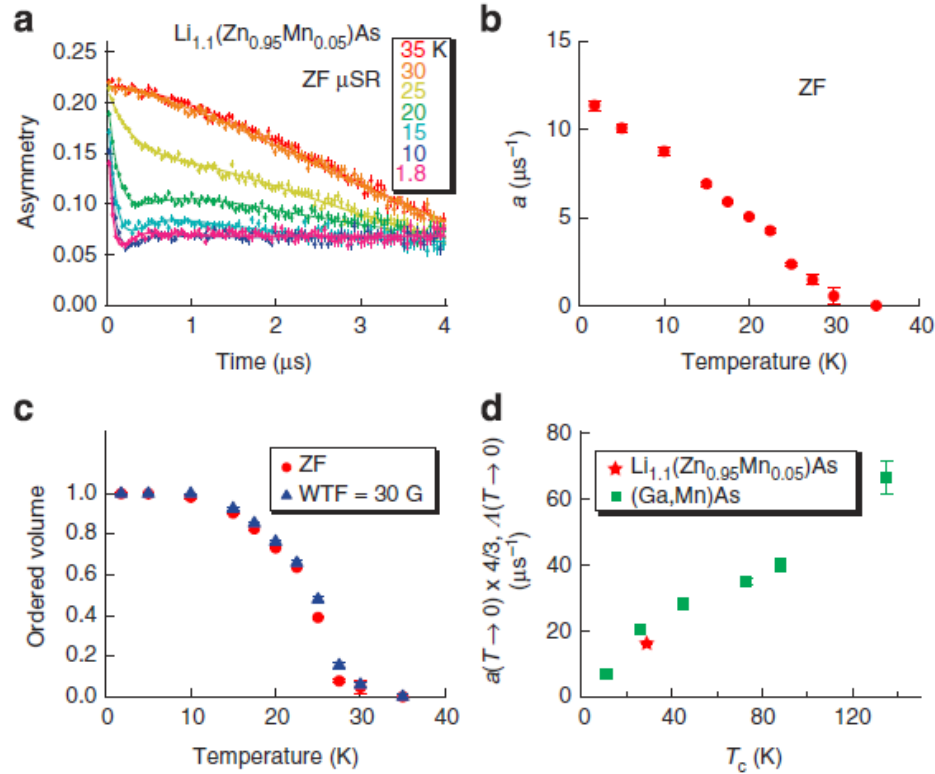
μ SR experiments were conducted in PSI

Spatially homogeneous ferromagnetism of (Ga, Mn)As

S. R. Dunsiger^{1,2}, J. P. Carlo¹, T. Goko^{1,3}, G. Nieuwenhuys⁴, T. Prokscha⁴, A. Suter⁴, E. Morenzoni⁴, D. Chiba^{5,6}, Y. Nishitani⁶, T. Tanikawa^{5,6}, F. Matsukura^{5,6}, H. Ohno^{5,6}, J. Ohe^{7,8}, S. Maekawa^{7,8} and Y. J. Uemura^{1*}



Li_{1.1}(Zn_{0.95}Mn_{0.05})As (T_C ~ 30 K)

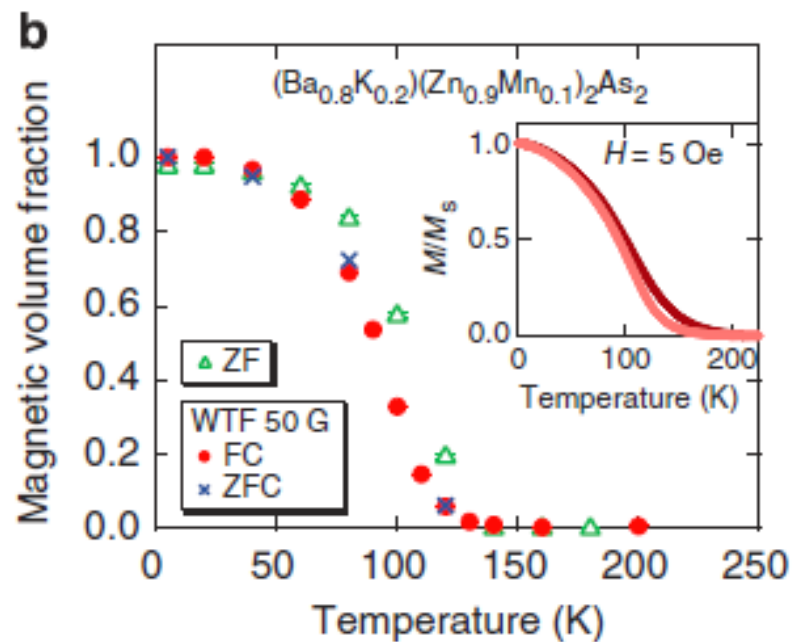
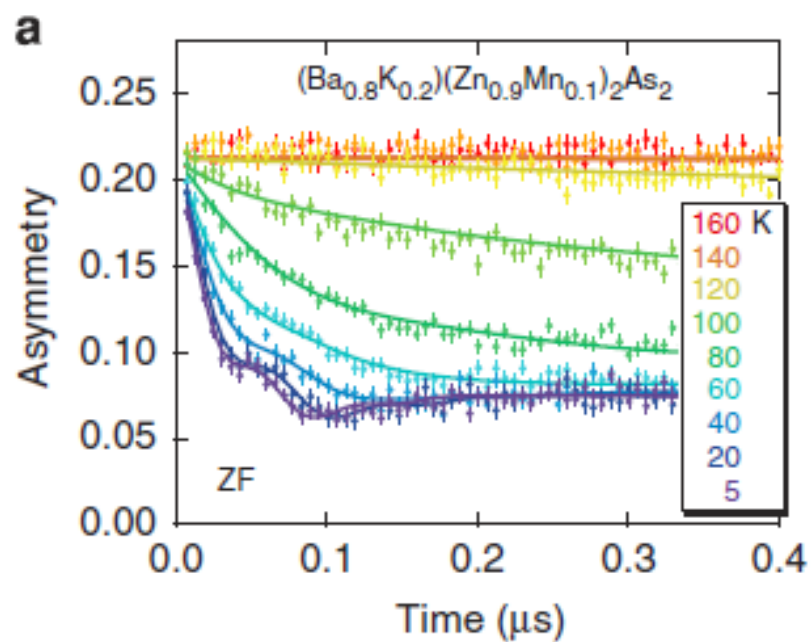


$$A_{\text{mag}} G_z^L(t) + A_{\text{para}} \exp(-(\lambda t)^\beta).$$

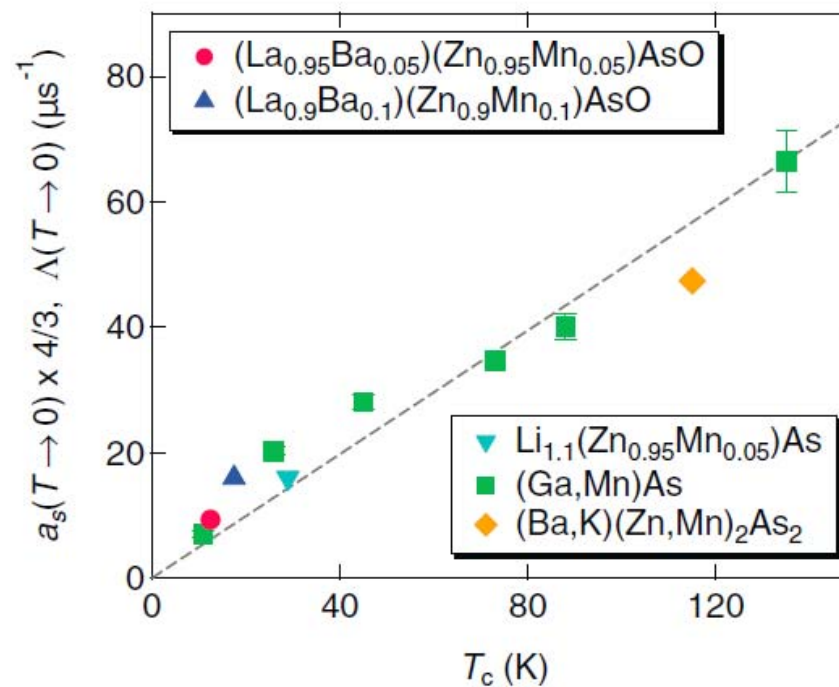
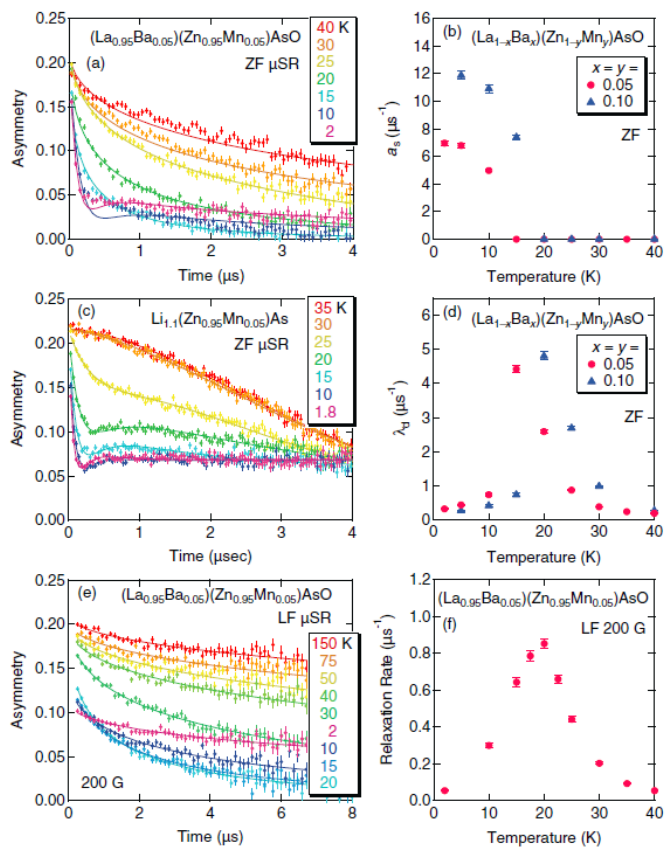
$$G_z^L(T) = \frac{1}{3} + \frac{2}{3}(1 - at) \exp(-at),$$

a_s --- relaxation rate (static local field amplitude)

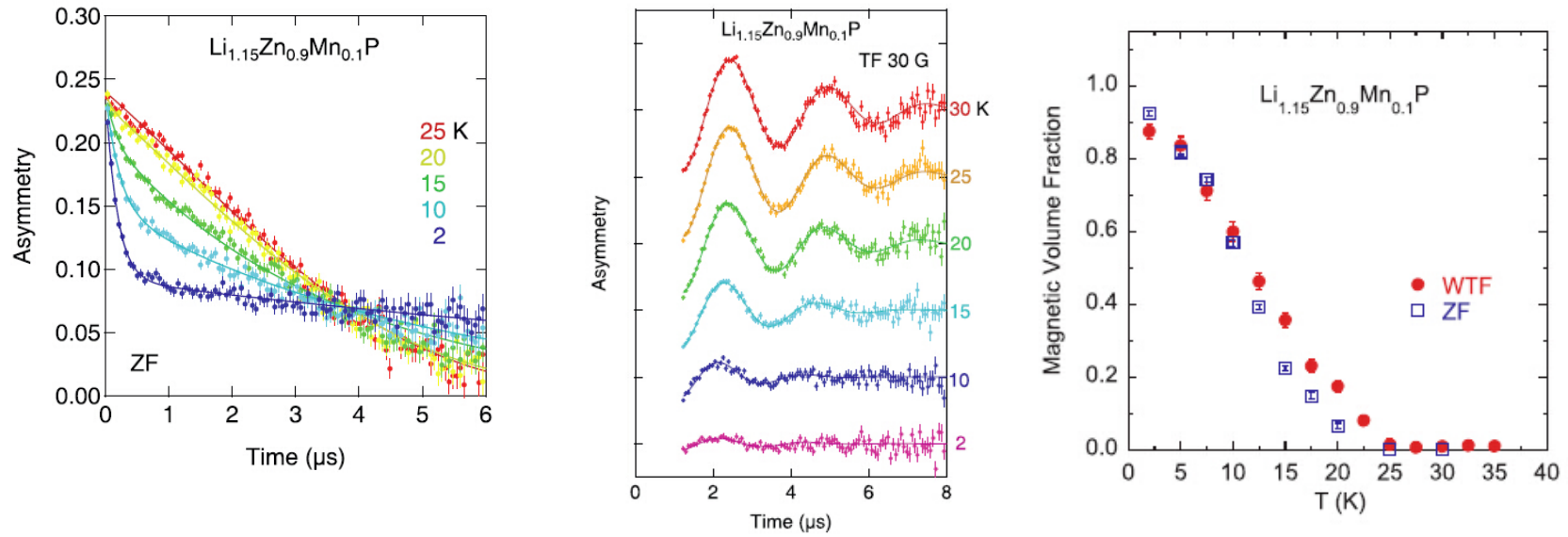
$(\text{Ba}_{1-x}\text{K}_x)(\text{Zn}_{1-y}\text{Mn}_y)_2\text{As}_2$ ($T_C \sim 180$ K)



$(\text{La}_{1-x}\text{Ba}_x)(\text{Zn}_{1-x}\text{Mn}_x)\text{AsO}$ $T_c \sim 40\text{K}$



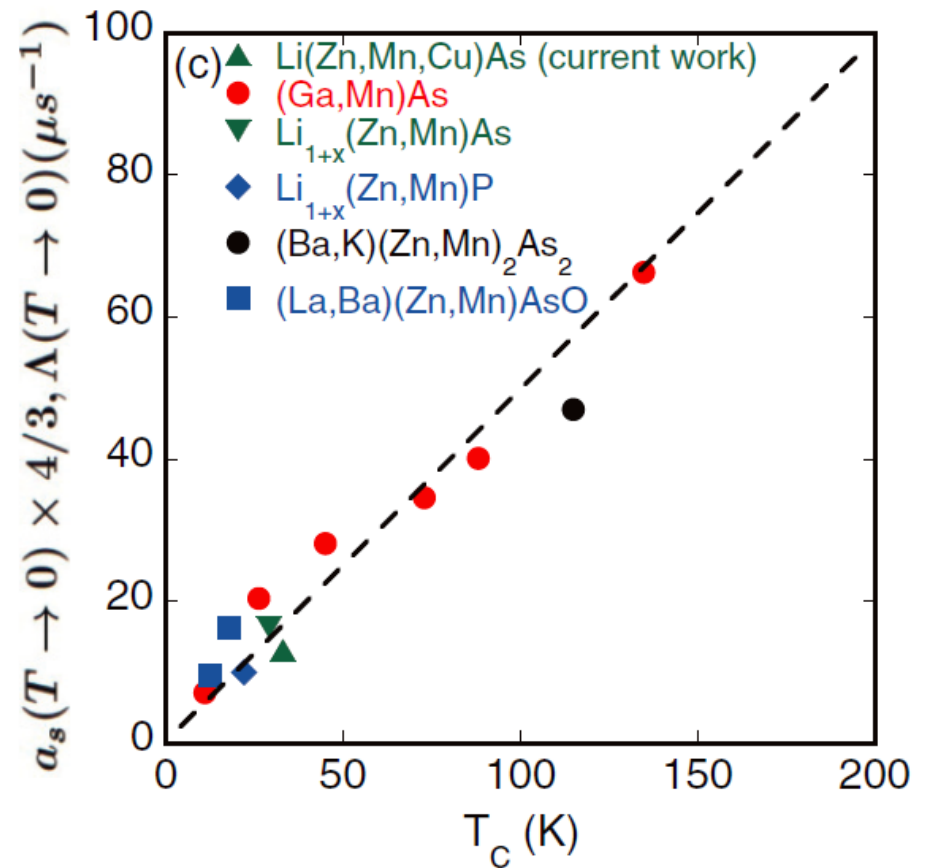
$\text{Li}_{1.15}(\text{Zn}_{0.9}\text{Mn}_{0.1})\text{P}$: $T_C \sim 22 \text{ K}$



F.L. Ning et al, PRB 90,085123 (2014)

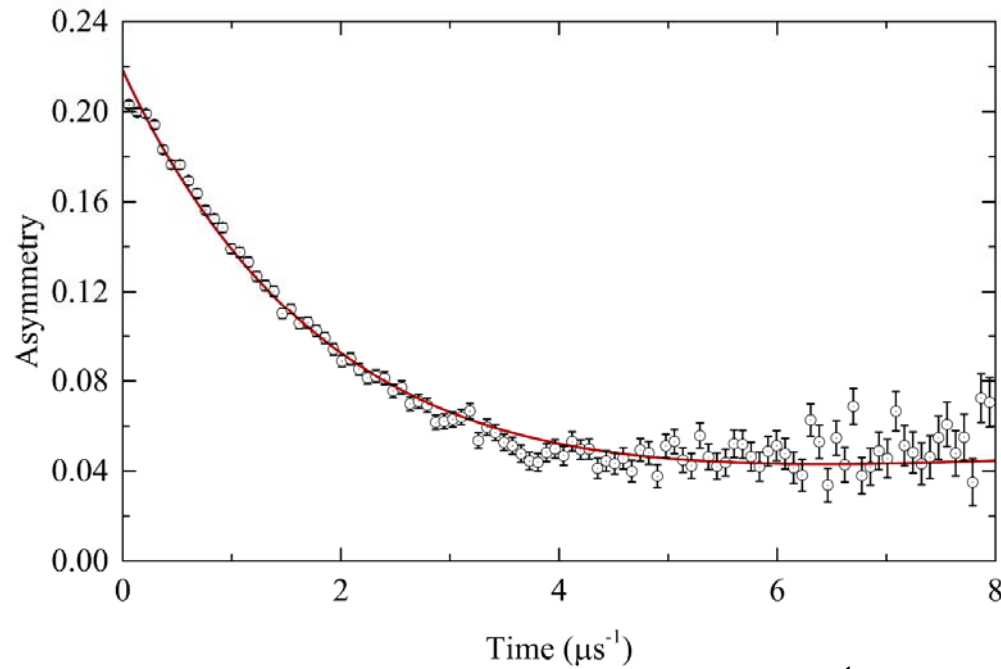
**Static local field amplitude a_s :
which is proportional to the
individual ordered moment
size multiplied by the moment
concentration**

S. Dunsiger, Nature Material, 2009
C. Ding, Physical Review B, 2013
K. Zhao, Nature Communications, 2013
Z. Deng, Nature Communications, 2011
F.L. Ning, Physical Review B, 2014



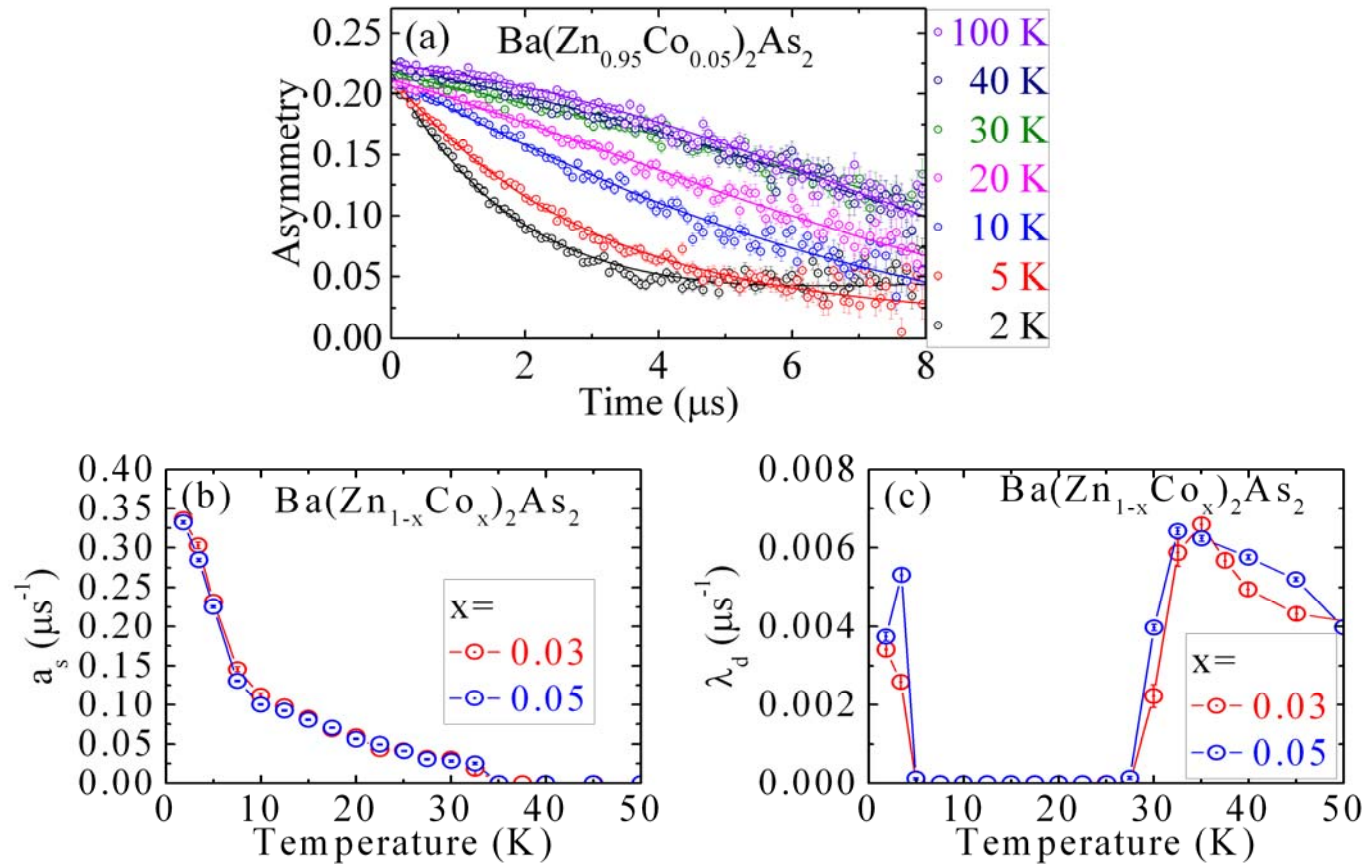
**T_c : which is a measure of the effective
average ferromagnetic interaction**

Zero field (ZF-) MuSR asymmetry spectra of Ba(Zn_{0.95}Co_{0.05})₂As₂ at 2 K



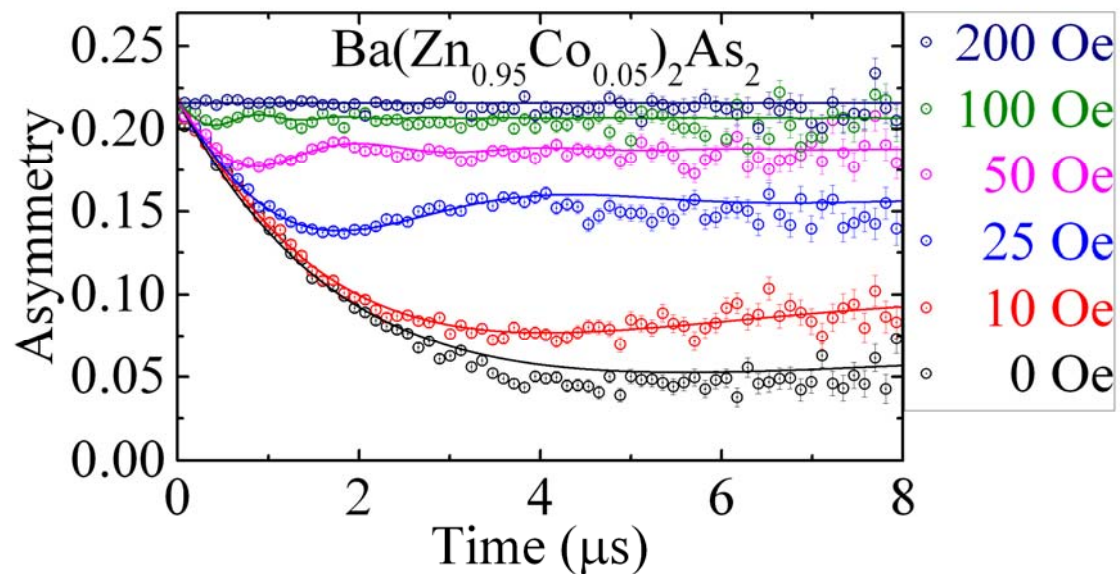
Fitting Function for dilute spin systems: $G_z(t) \times g_z^G(t)$, $G_z(t) = \frac{1}{3} \exp\left(-(\lambda_d t)^{\frac{1}{2}}\right) + \frac{2}{3} \left[1 - \frac{a_s^2 t^2}{(\lambda_d t + a_s^2 t^2)^{\frac{1}{2}}}\right] \exp\left[-(\lambda_d t + a_s^2 t^2)^{\frac{1}{2}}\right]$, $g_z^G(t) = \frac{1}{3} + \frac{2}{3} (1 - \Delta_0^2 t^2) \exp\left(-\frac{1}{2} \Delta_0^2 t^2\right)$. Uemura *et al.* *Phys. Rev. B* **31**, 546 (1985).

ZF-MuSR results for $\text{Ba}(\text{Zn}_{1-x}\text{Co}_x)_2\text{As}_2$



Static field arises below T_C

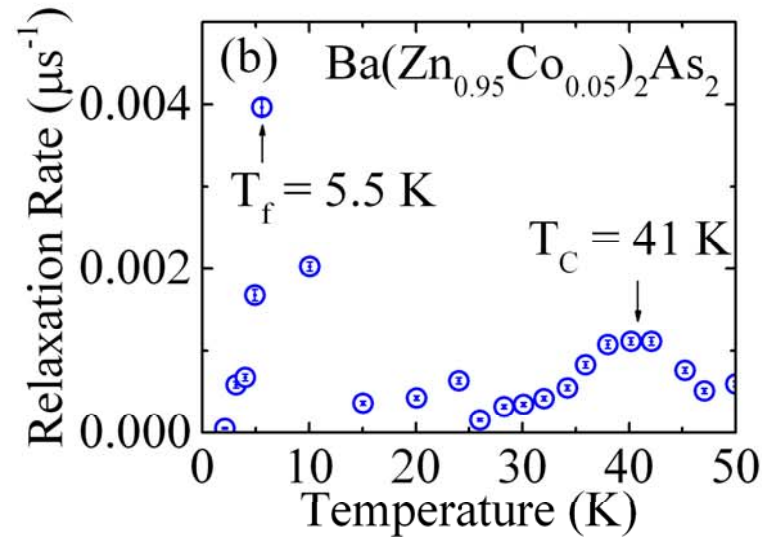
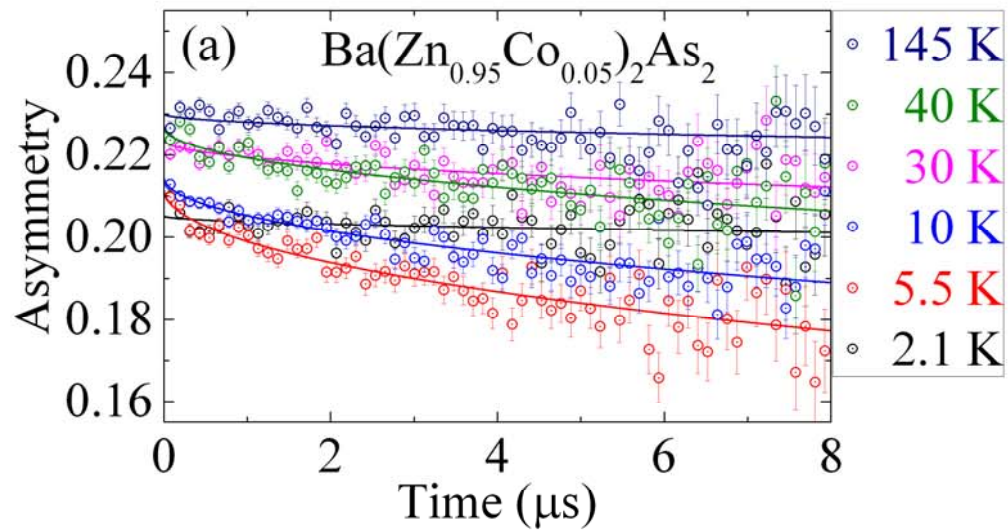
Longitudinal field (LF-) MuSR results for $\text{Ba}(\text{Zn}_{1-x}\text{Co}_x)_2\text{As}_2$ at 2 K

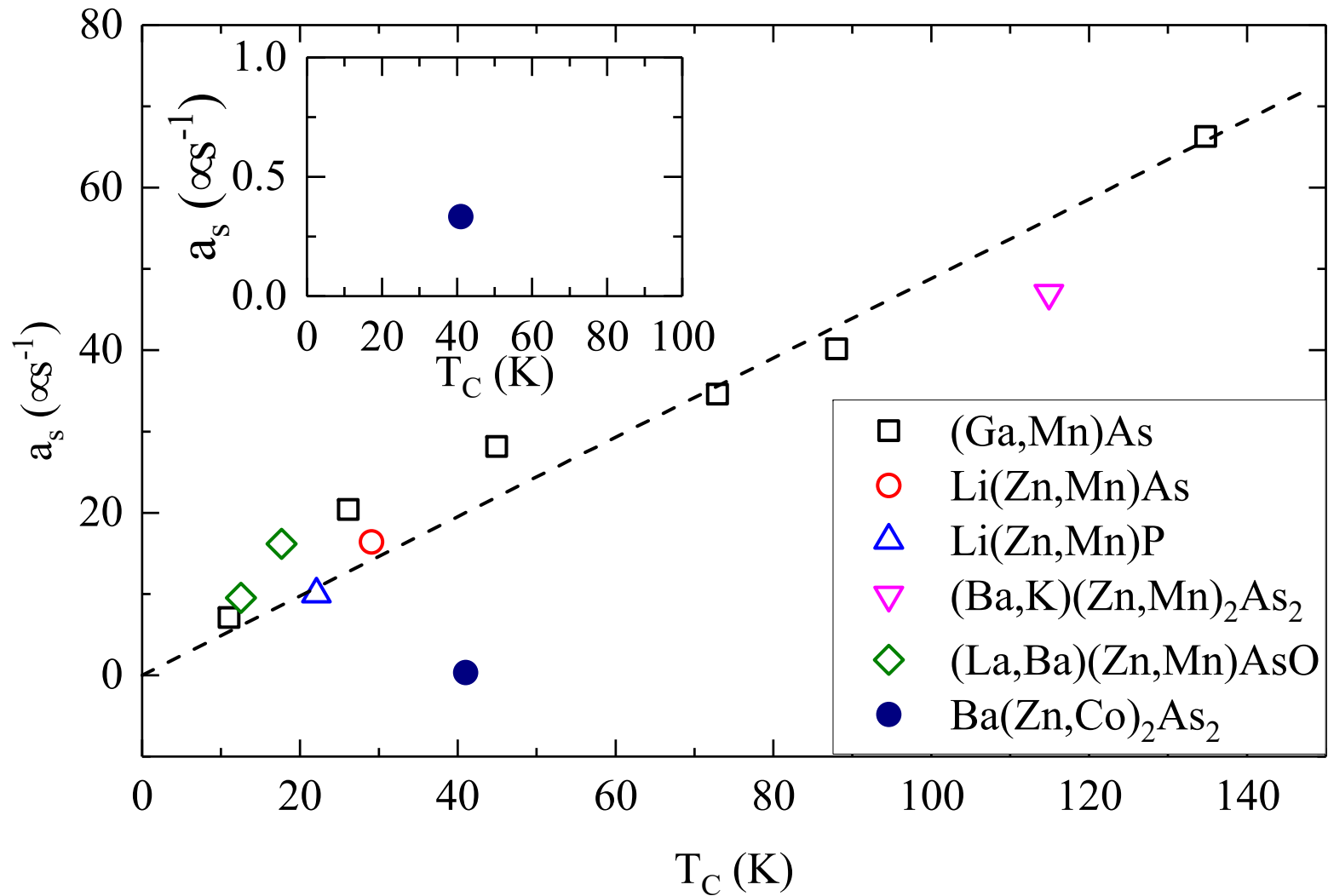


The asymmetry is fully decoupled at ~ 100 Oe.

Conclusion: the internal field at 2 K muon suffered is **fully static** and magnitude is ~ 10 Oe.

Temperature scan 100 Oe LF-MuSR results for $\text{Ba}(\text{Zn}_{1-x}\text{Co}_x)_2\text{As}_2$





Ferromagnetic exchange coupling is **much larger** in this system compared these other p-type systems.

Outline

a. Introduction of some newly fabricated DMSs

“111”, “122”, “1111”, “42622”, “32522”

----- $\text{Li}_{1+y}(\text{Zn}_{1-x}\text{Mn}_x)\text{P}$

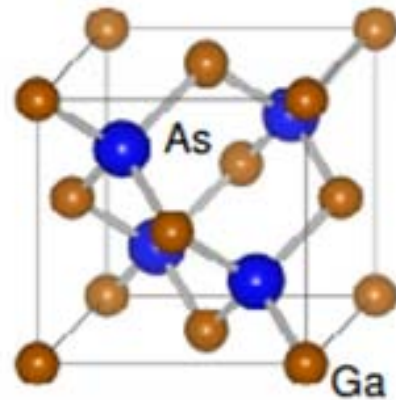
----- n-type $\text{Ba}(\text{Zn}_{1-2x}\text{Co}_x\text{Mn}_x)_2\text{As}_2$

b. Microscopic characterization

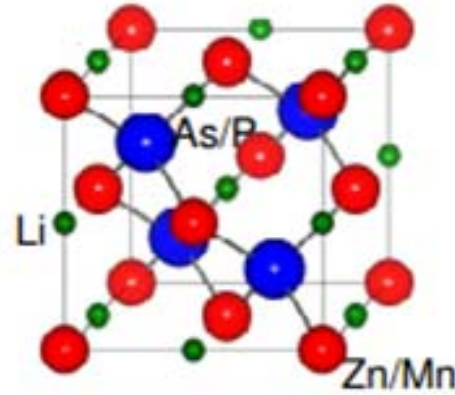
----- μSR on $\text{Li}_{1.15}(\text{Zn}_{1-x}\text{Mn}_x)\text{P}$

----- **NMR results on $\text{Li}_{1.1}(\text{Cd}_{1-x}\text{Mn}_x)\text{P}$ and $\text{Li}_{1.15}(\text{Zn}_{1-x}\text{Mn}_x)\text{P}$**

“III-V” versus “I-II-V” semiconductors



GaAs (1.42eV)



LiZnAs (1.6eV)
LiZnP (2.04eV)
LiCdP (1.30eV)



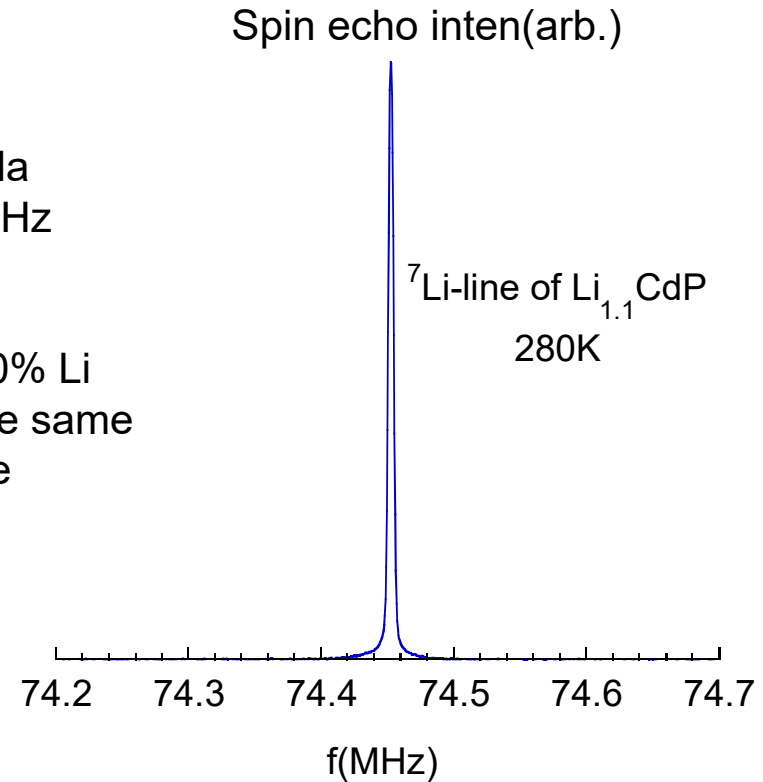
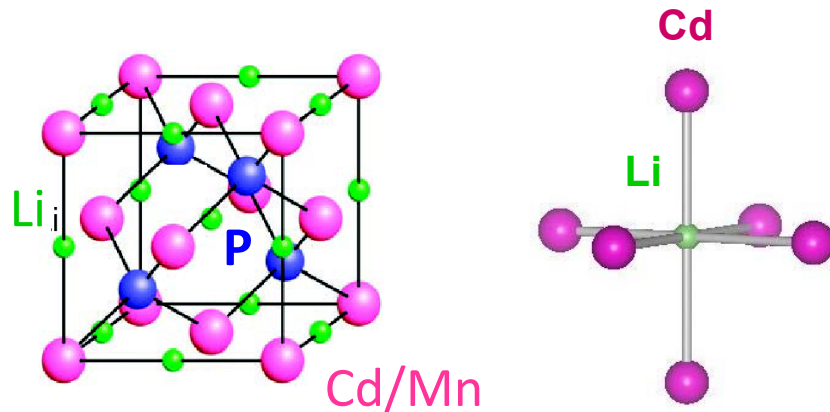
- decouple spins and carriers
- Mn doping induces local moments
- Li extra or deficiencies create itinerant or localized carriers
- use Li to control the type of carriers

^7Li NMR line shape ---- $\text{Li}_{1.1}\text{CdP}$

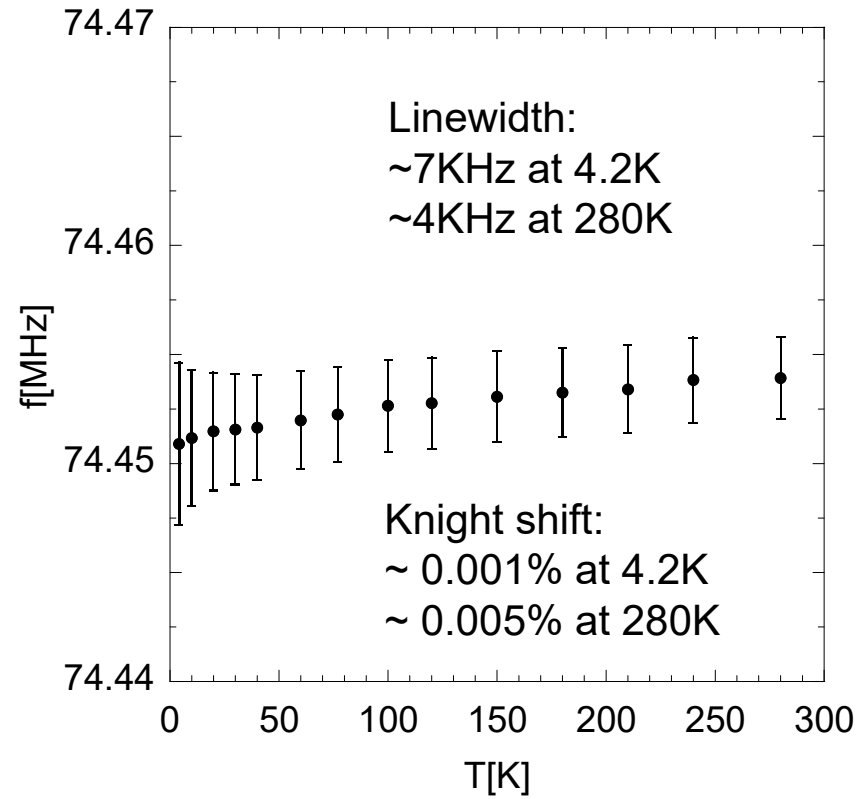
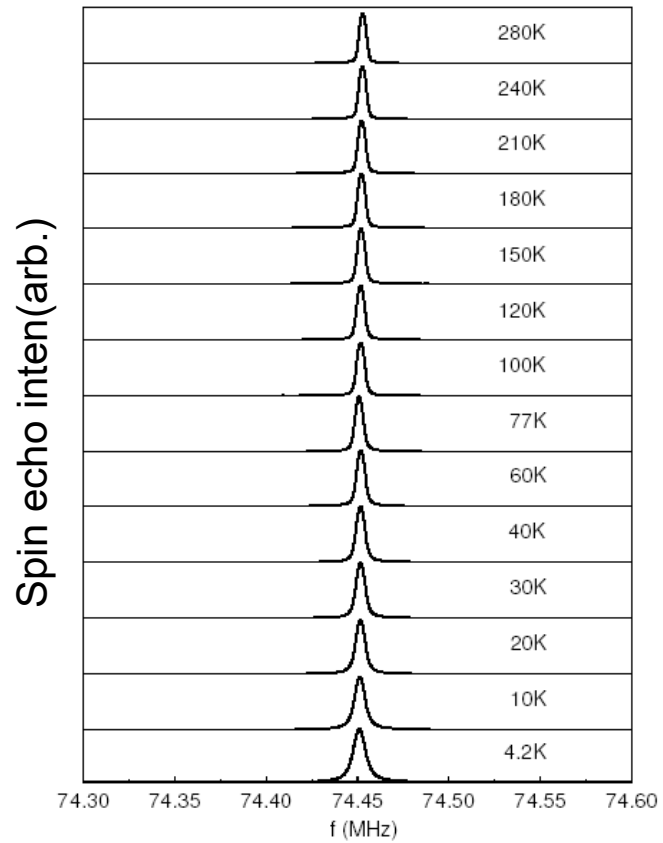
For ^7Li :

- Nuclear spin = $3/2$
- Gyromagnetic ratio = 16.546MHz/Tesla
- $B_{\text{ext}} = 4.5\text{ Tesla} \rightarrow f_{\text{resonance}} = 74.457\text{MHz}$

- Single line with line width = 4KHz
- No additional Li line found for extra 10% Li
- 6 Cd atoms at N.N. site, all Li have the same environment -----define it as $\text{Li}(0)$ site

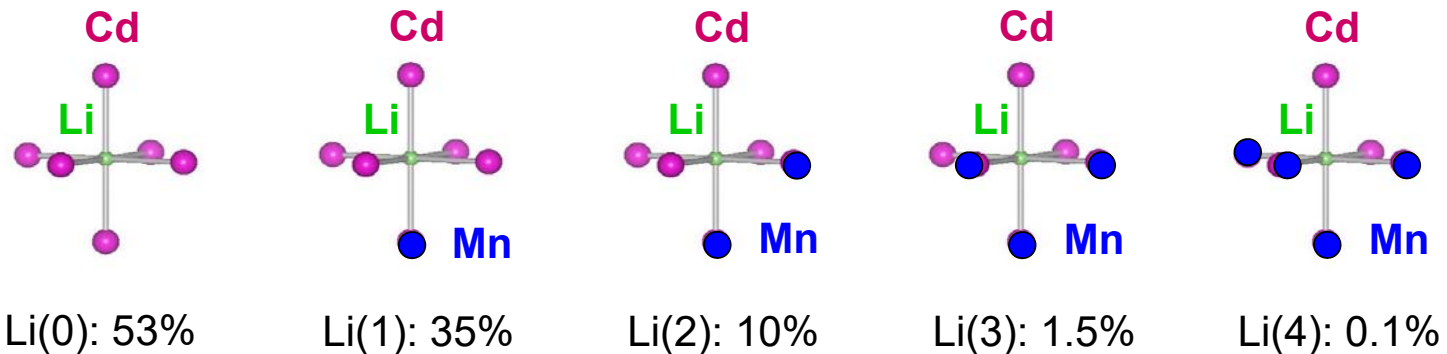


^7Li NMR line shape ---- $\text{Li}_{1.1}\text{CdP}$



^7Li NMR line shape ---- $\text{Li}_{1.1}\text{Cd}_{0.9}\text{Mn}_{0.1}\text{P}$

10% Mn doped into Cd sites, the probability to find $n = 0, 1, 2, 3, 4, 5, 6$ Mn atoms at N.N. sites are $P(n) = C_6^n(0.1)^n(0.9)^{6-n}$

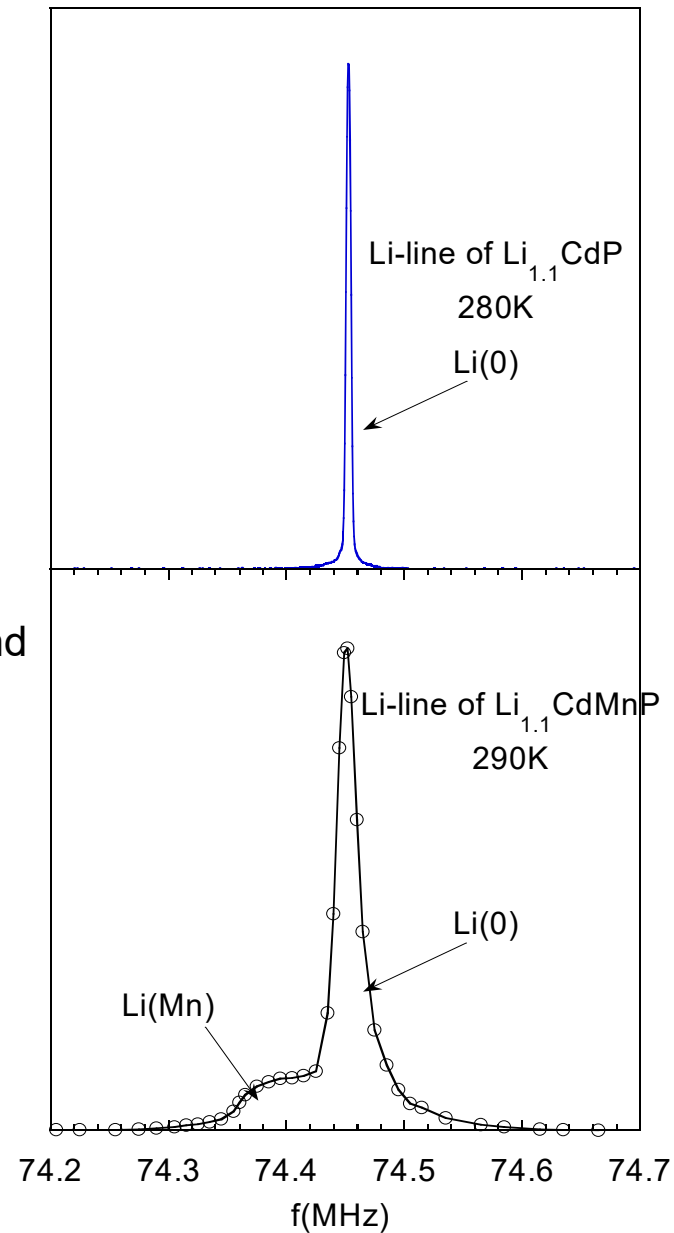


Li(0):Li(Mn) ~ 50%:50%

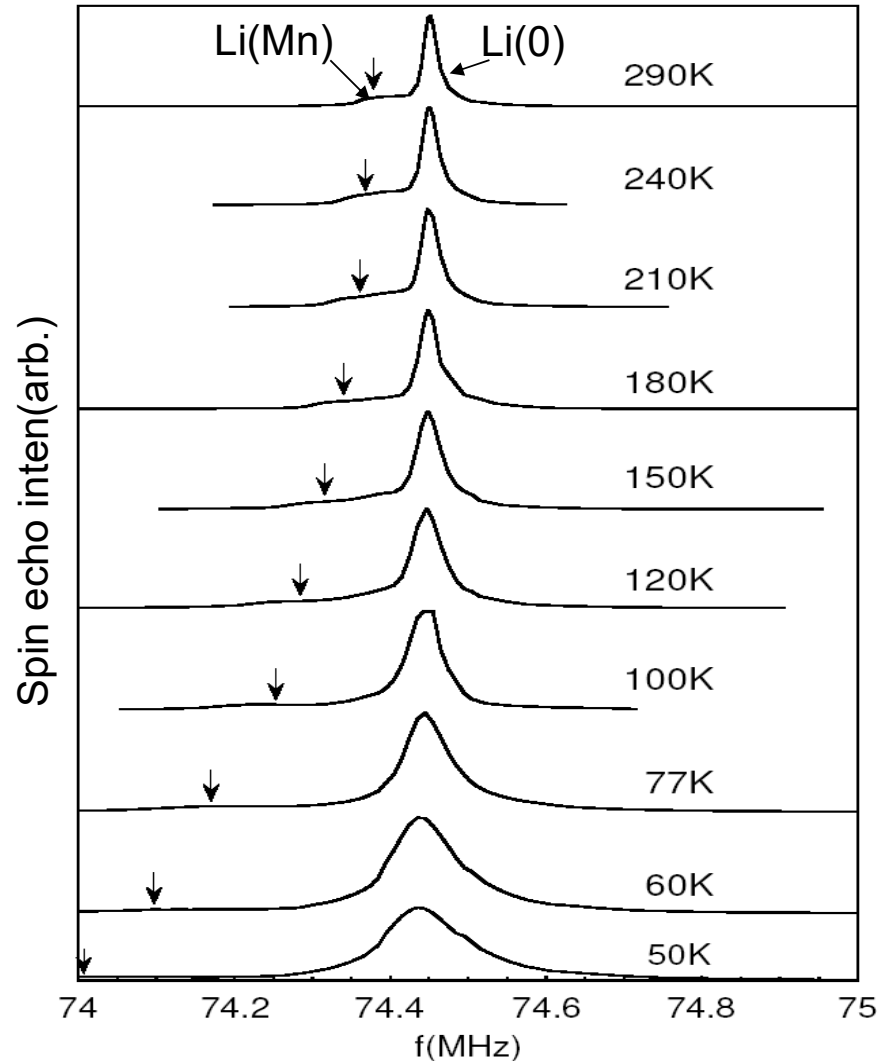
^7Li NMR line shape ---- $\text{Li}_{1.1}\text{Cd}_{0.9}\text{Mn}_{0.1}\text{P}$

Key features:

- Li(0) line preserve but broaden
- Li(Mn) line appears, broad peak at left hand
- Li(Mn) line has a large negative shift
- Integrated area, **Li(Mn):Li(0) = 50:50**



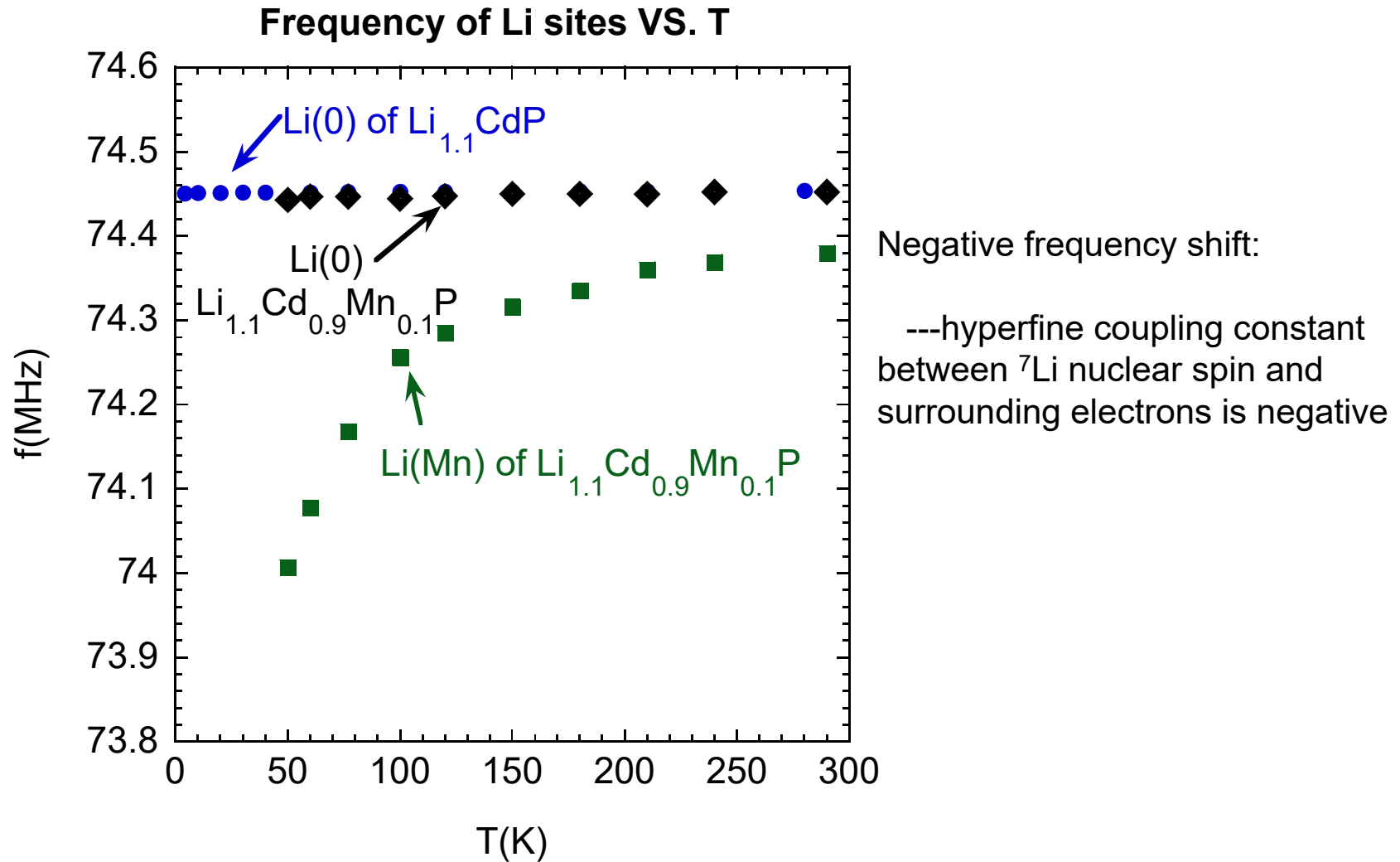
^7Li NMR line shape ---- $\text{Li}_{1.1}\text{Cd}_{0.9}\text{Mn}_{0.1}\text{P}$



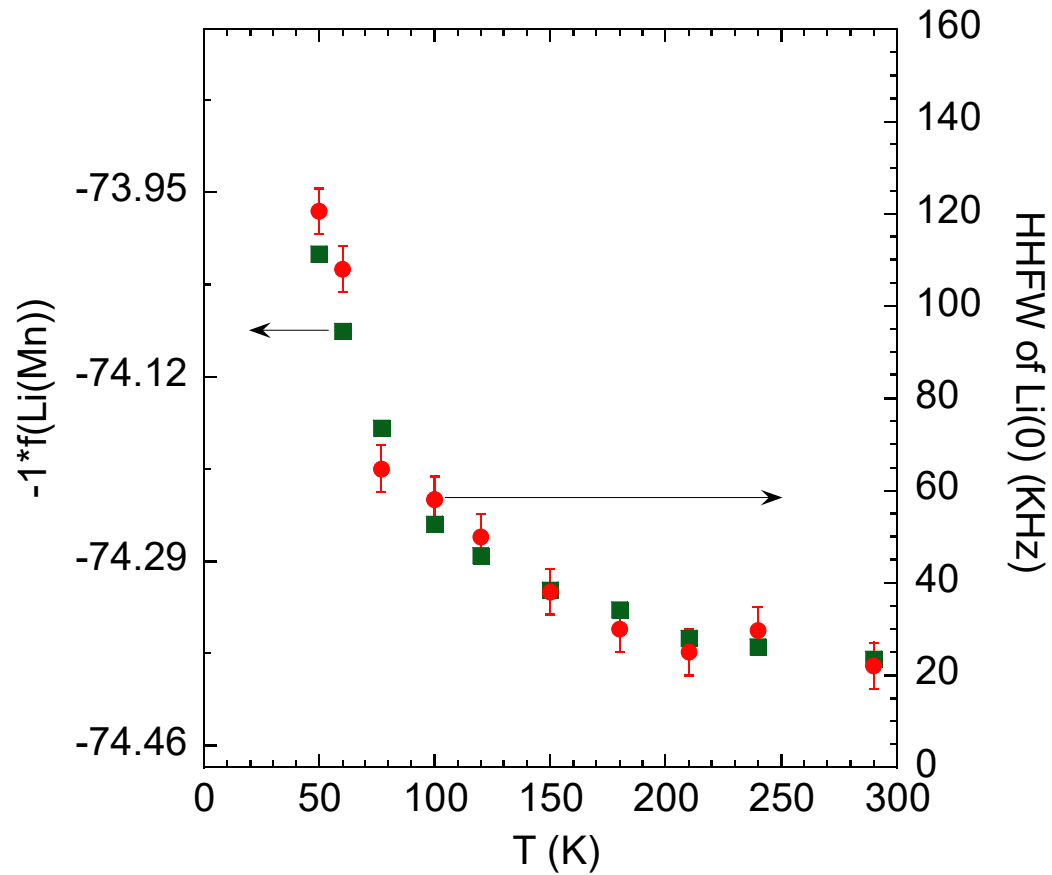
Key features:

- Frequency shifts of Li(Mn) follows a Curie-Weiss law
- Li(0) continuously broaden from 20KHz (at 290K) to 240KHz (at 4.2K) ~ 10 times broadening. Line width follow a Curie-Weiss law
- For Li(0) site in the parent compound $\text{Li}_{1.1}\text{CdP}$, only increase from 4KHz (at 280K) to 7KHz (at 4.2K).

Frequency VS. Temperature



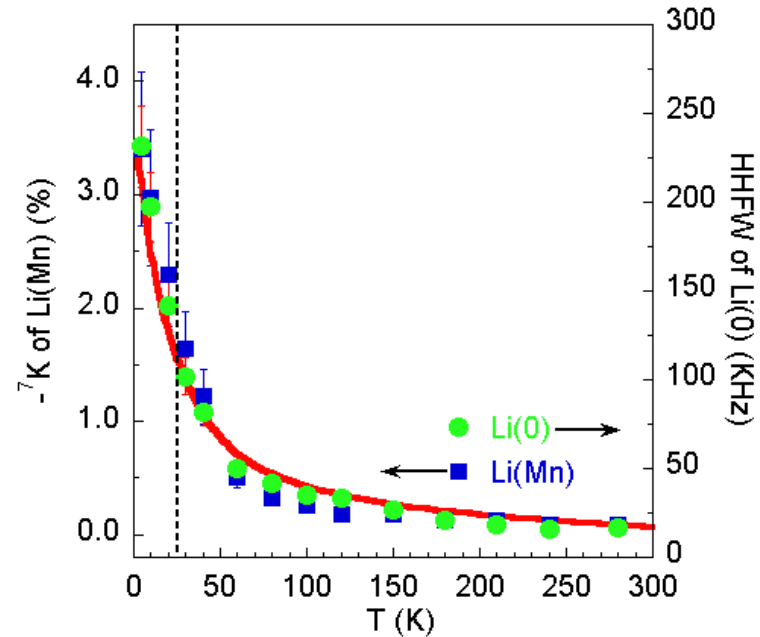
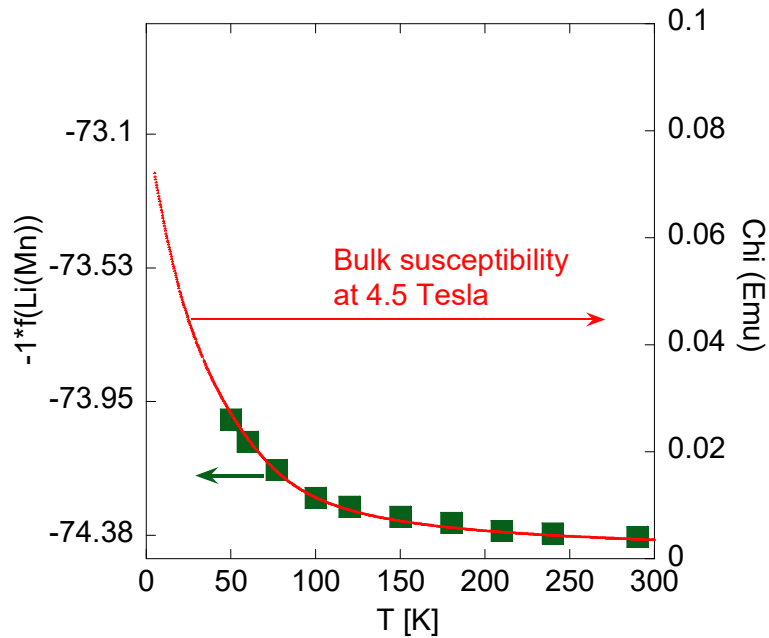
$\text{Li}_{1.1}\text{Cd}_{0.9}\text{Mn}_{0.1}\text{P}$



The line broadening of Li(0) which do not have Mn at N.N. is the same as the frequency shifts of Li(Mn).



Li(0) and Li(Mn) sites are electronically coupled!

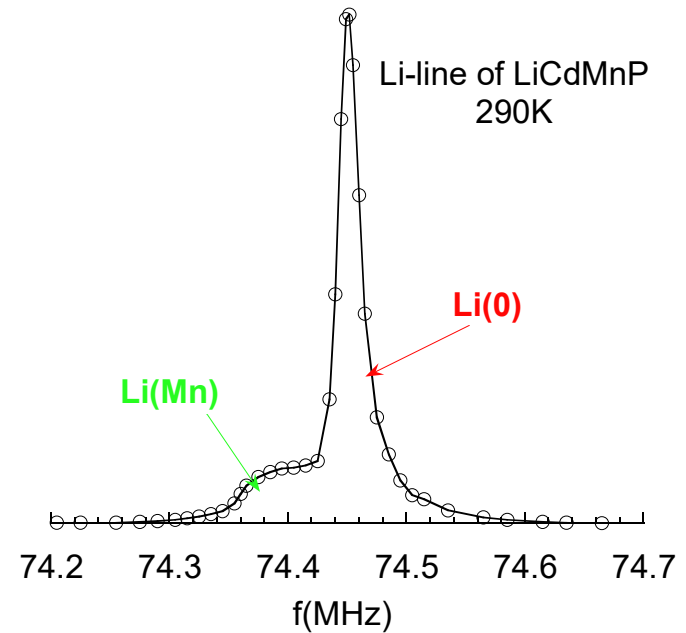
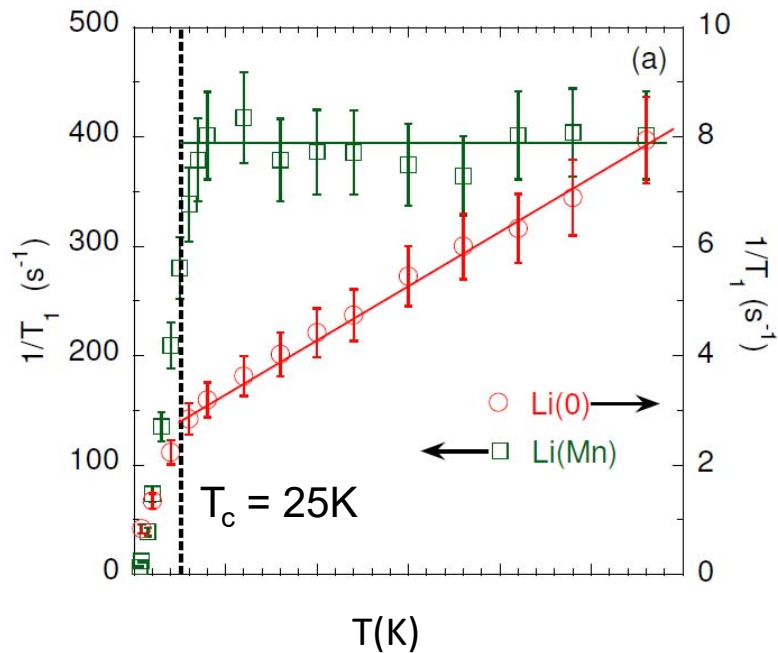


Mn atoms are homogenously distributed in the compound, and no Mn clusters exist.

Consistent with 100% magnetic ordered volume from μ SR experiment.

The ferromagnetic ordering is intrinsic, from the Mn atoms in the ionic sites in I-II-V DMS! **It is NOT from spurious phase.**

Spin dynamics--- $\text{Li}_{1.15}\text{Zn}_{0.9}\text{Mn}_{0.1}\text{P}$ and $\text{Li}_{1.1}\text{Cd}_{0.9}\text{Mn}_{0.1}\text{P}$

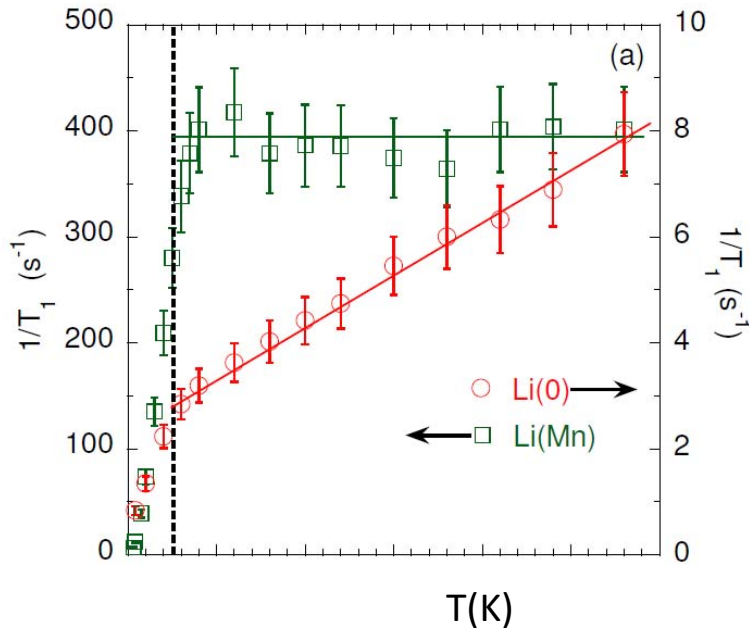


T_1 of Li(Mn) is 50 times shorter than T_1 of Li(0)

Below T_c , for both Li(Mn) and Li(0) sites, T_1 decrease

C. Ding et al., Physics Review B 88, 041108 (R) (2013).

Spin dynamics----Li_{1.15}Zn_{0.9}Mn_{0.1}P



For Li(Mn) site, $1/T_1 \sim \text{constant}$
Spin fluctuations are locked
 at T range ~ 100 K, $J \sim 100$ K

For Li(0) site, $1/T_1 \sim a + bT$
 Korringa relation,
 Related to information of **carriers**

Li(Mn) represent “**spins**”
 Li(0) represent “**carries**”

Carrier mediated ferromagnetism!

$$\frac{1}{T_1} = \sqrt{2\pi} \frac{S(S+1)}{3\omega_e} \left(\frac{A_0}{\hbar}\right)^2 \sim 400 \text{ s}^{-1}$$

Moriya's exchange narrowed regime

$$\omega_e \sim 1.1 \times 10^{15} \text{ rad/s}$$

Characteristic frequency of the Mn spin fluctuations

$$\omega_e^2 = \frac{2}{3} z S(S+1) \left(\frac{J}{\hbar}\right)^2$$

Moriya's Gaussian approximation for the
 spin-spin correlation function

$Z \sim 10^3$ is the number of Mn sites within the
 range of Mn-Mn interactions

Summary

- Series of bulk DMSs with **(highest $T_C = 230$ K)** have been fabricated, and these DMSs have common mechanism with (Ga,Mn)As, and can serve as model systems to study the ferromagnetism in DMS
- A **n-type DMS** $\text{Ba}(\text{Zn}_{1-x}\text{Co}_x)_2\text{As}_2$ ($T_C = 45\text{K}$) have been successfully synthesized
- NMR and μSR : Mn atoms are **homogenously** doped, the ferromagnetic ordering is from Mn in the ionic Zn site
- NMR results are consistent with a **carrier mediated ferromagnetism** picture.
- Now with the same 122 tetragonal structure, junctions could be made **via As layer between following materials:**

N-type DMS $\text{Ba}(\text{Zn}_{1-x}\text{Co}_x)_2\text{As}_2$ ($T_C = 45\text{K}$)

P-type DMS $(\text{Ba}_{1-y}\text{K}_y)(\text{Zn}_{1-x}\text{Mn}_x)_2\text{As}_2$ ($T_C=230\text{K}$)

Superconductor $\text{Ba}(\text{Fe}_{1-x}\text{Co}_x)_2\text{As}_2$ ($T_c=25\text{K}$)

Antiferromagnet BaMn_2As_2 ($T_N=625\text{K}$)

Paramagnetic BaCo_2As_2

Thank you for your attention!