

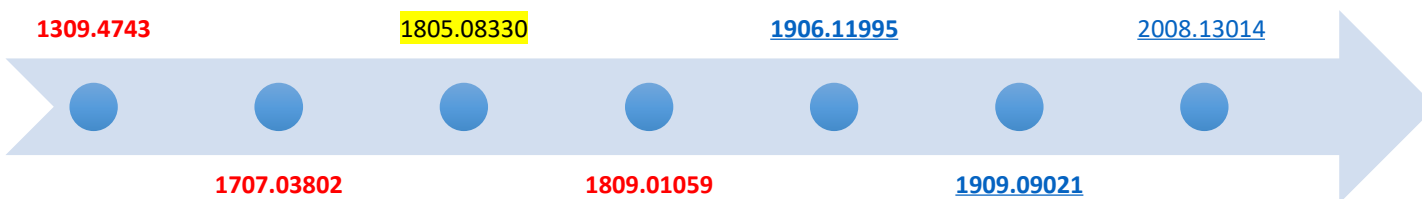
北京航空航天大学  
BEIHANG UNIVERSITY



# 自然界存在新的类原子核物质形态么？

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# Contents

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- Motivation: new types of clusters of color singlets in addition to atomic nuclei—as a nontrivial test of the molecule picture
- $Ds0^*(2317)$  and  $Ds1(2460)$  as **DK/D\*K molecules: theory & lattice**
- An explicit study of the **DDK** system:  $R^{++}(4140)$
- $K^*(4307)$  as excited  $K^*$  with hidden charm
- Summary and outlook

# Contents

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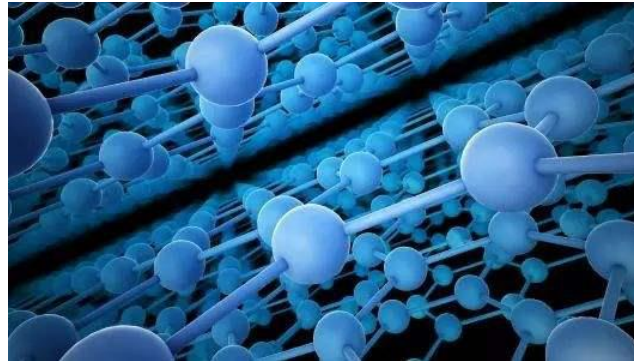
- Motivation: new types of clusters of color singlets in addition to atomic nuclei—as a nontrivial test of the molecule picture
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# Frontiers of Physics

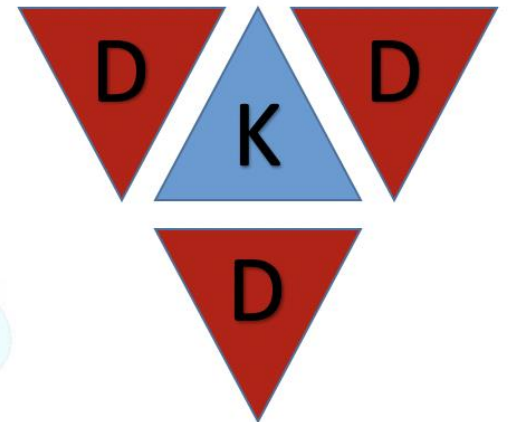
**New Energy**



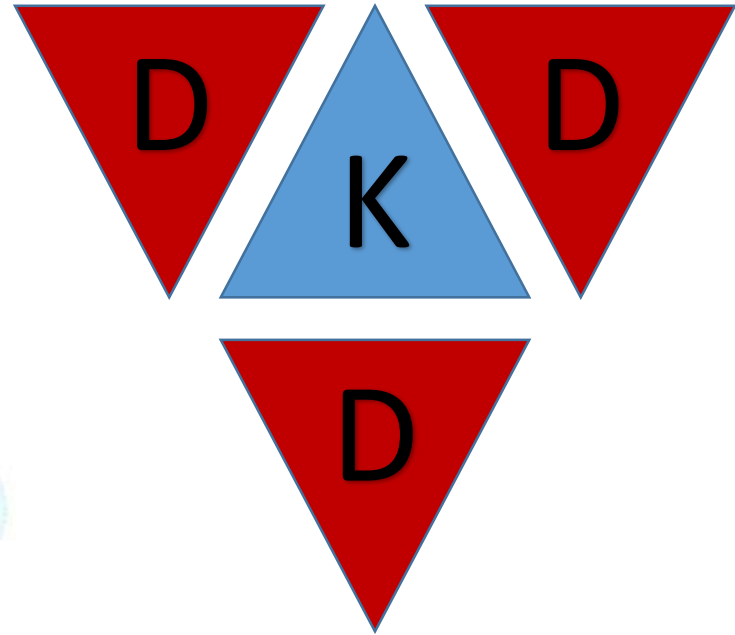
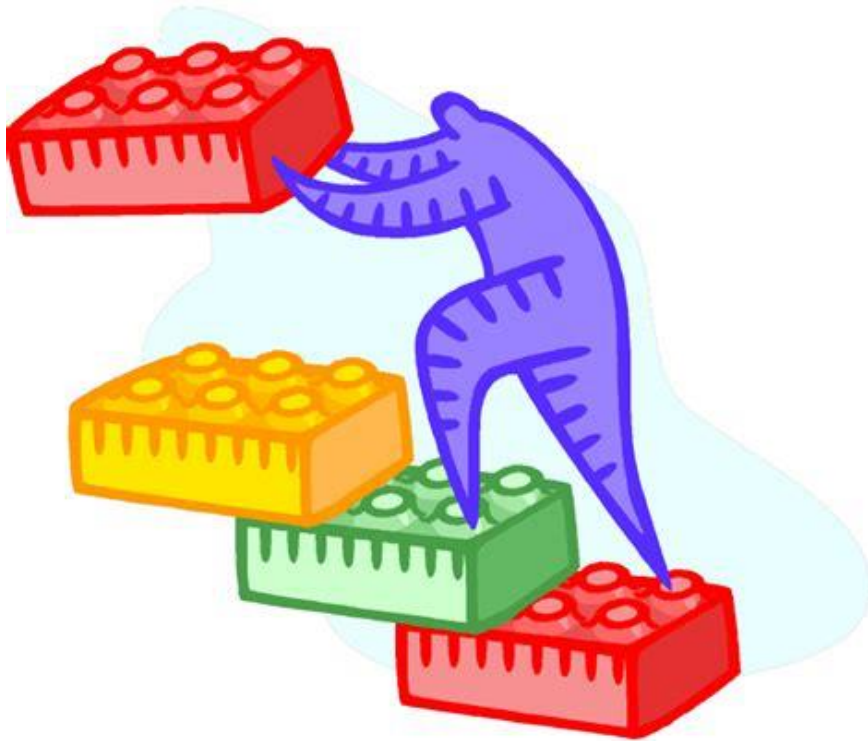
**New Material**

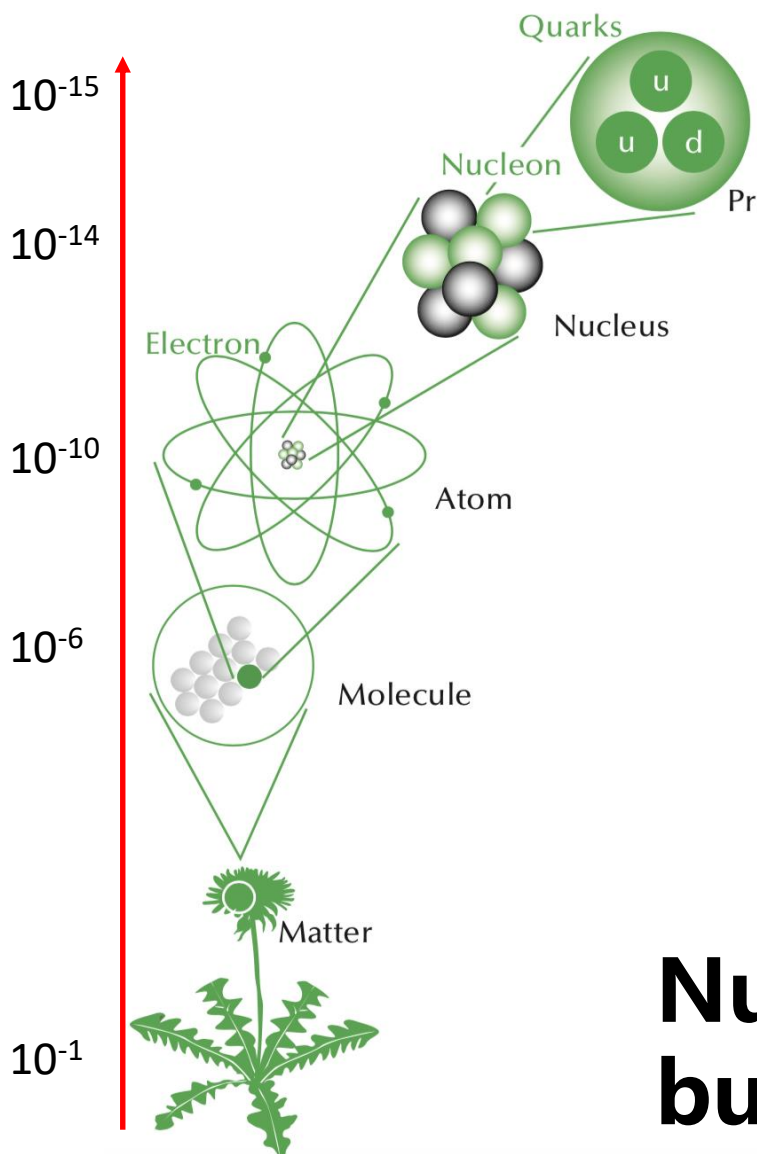


**New Matter**



Build **new forms** of bosonic matter from D & K mesons





**IUPAC Periodic Table of the Elements**

1 <b>H</b> hydrogen 1.008 (1.00784, 1.00824)	2 <b>He</b> helium 4.0026	13 <b>B</b> boron 10.81 (10.806, 10.821)	14 <b>C</b> carbon 12.01 (12.009, 12.012)	15 <b>N</b> nitrogen 14.01 (14.006, 14.008)	16 <b>O</b> oxygen 15.99 (15.999, 16.005)	17 <b>F</b> fluorine 18.998 (18.998, 18.998)	18 <b>Ne</b> neon 20.18 (20.179, 20.183)	19 <b>K</b> potassium 39.098 (39.096, 39.101)	20 <b>Ca</b> calcium 40.078 (40.078, 40.078)	21 <b>Sc</b> scandium 44.956 (44.955, 44.957)	22 <b>Ti</b> titanium 47.867 (47.867, 47.867)	23 <b>V</b> vanadium 50.942 (50.942, 50.942)	24 <b>Cr</b> chromium 51.996 (51.996, 51.996)	25 <b>Mn</b> manganese 54.938 (54.938, 54.938)	26 <b>Fe</b> iron 55.845 (55.845, 55.845)	27 <b>Co</b> cobalt 58.933 (58.933, 58.933)	28 <b>Ni</b> nickel 58.693 (58.693, 58.693)	29 <b>Cu</b> copper 63.546 (63.546, 63.546)	30 <b>Zn</b> zinc 65.38 (65.38, 65.38)	31 <b>Ga</b> gallium 69.723 (69.723, 69.723)	32 <b>Ge</b> germanium 72.63 (72.63, 72.63)	33 <b>As</b> arsenic 74.922 (74.922, 74.922)	34 <b>Se</b> selenium 78.96 (78.96, 78.96)	35 <b>Br</b> bromine 79.904 (79.904, 79.904)	36 <b>Kr</b> krypton 83.796 (83.796, 83.796)	37 <b>Rb</b> rubidium 85.468 (85.468, 85.468)	38 <b>Sr</b> strontium 87.62 (87.62, 87.62)	39 <b>Y</b> yttrium 88.906 (88.906, 88.906)	40 <b>Zr</b> zirconium 91.224 (91.224, 91.224)	41 <b>Nb</b> niobium 92.906 (92.906, 92.906)	42 <b>Mo</b> molybdenum 95.94 (95.94, 95.94)	43 <b>Tc</b> technetium 98 (98, 98)	44 <b>Ru</b> ruthenium 101.07 (101.07, 101.07)	45 <b>Rh</b> rhodium 102.91 (102.91, 102.91)	46 <b>Pd</b> palladium 106.42 (106.42, 106.42)	47 <b>Ag</b> silver 107.87 (107.87, 107.87)	48 <b>Cd</b> cadmium 112.41 (112.41, 112.41)	49 <b>In</b> indium 114.82 (114.82, 114.82)	50 <b>Sn</b> tin 118.71 (118.71, 118.71)	51 <b>Sb</b> antimony 121.76 (121.76, 121.76)	52 <b>Te</b> tellurium 127.6 (127.6, 127.6)	53 <b>I</b> iodine 126.90 (126.90, 126.90)	54 <b>Xe</b> xenon 131.29 (131.29, 131.29)	55 <b>Cs</b> cesium 132.91 (132.91, 132.91)	56 <b>Ba</b> barium 137.33 (137.33, 137.33)	57-71 <b>Lanthanoids</b>	72 <b>Hf</b> hafnium 178.49 (178.49, 178.49)	73 <b>Ta</b> tantalum 180.95 (180.95, 180.95)	74 <b>W</b> tungsten 183.84 (183.84, 183.84)	75 <b>Re</b> rhenium 186.21 (186.21, 186.21)	76 <b>Os</b> osmium 190.23 (190.23, 190.23)	77 <b>Ir</b> iridium 192.22 (192.22, 192.22)	78 <b>Pt</b> platinum 195.08 (195.08, 195.08)	79 <b>Au</b> gold 196.97 (196.97, 196.97)	80 <b>Hg</b> mercury 200.59 (200.59, 200.59)	81 <b>Tl</b> thallium 204.38 (204.38, 204.38)	82 <b>Pb</b> lead 207.2 (207.2, 207.2)	83 <b>Bi</b> bismuth 208.98 (208.98, 208.98)	84 <b>Po</b> polonium 209 (209, 209)	85 <b>At</b> astatine 210 (210, 210)	86 <b>Rn</b> radon 222 (222, 222)	87 <b>Fr</b> francium 223 (223, 223)	88 <b>Ra</b> radium 226 (226, 226)	89-103 <b>actinoids</b>	104 <b>Rf</b> rutherfordium 261 (261, 261)	105 <b>Db</b> dubnium 262 (262, 262)	106 <b>Sg</b> seaborgium 266 (266, 266)	107 <b>Bh</b> bohrium 264 (264, 264)	108 <b>Hs</b> hassium 277 (277, 277)	109 <b>Mt</b> meitnerium 268 (268, 268)	110 <b>Ds</b> darmstadtium 271 (271, 271)	111 <b>Rg</b> roentgenium 272 (272, 272)	112 <b>Cn</b> copernicium 285 (285, 285)	113 <b>Nh</b> nihonium 284 (284, 284)	114 <b>Fl</b> flerovium 289 (289, 289)	115 <b>Mc</b> moscovium 288 (288, 288)	116 <b>Lv</b> livermorium 293 (293, 293)	117 <b>Ts</b> tennessine 289 (289, 289)	118 <b>Og</b> oganesson 294 (294, 294)	57 <b>La</b> lanthanum 138.91 (138.91, 138.91)	58 <b>Ce</b> cerium 140.12 (140.12, 140.12)	59 <b>Pr</b> praseodymium 140.91 (140.91, 140.91)	60 <b>Nd</b> neodymium 144.24 (144.24, 144.24)	61 <b>Pm</b> promethium 145 (145, 145)	62 <b>Sm</b> samarium 150.36 (150.36, 150.36)	63 <b>Eu</b> europium 151.96 (151.96, 151.96)	64 <b>Gd</b> gadolinium 157.25 (157.25, 157.25)	65 <b>Tb</b> terbium 158.93 (158.93, 158.93)	66 <b>Dy</b> dysprosium 162.50 (162.50, 162.50)	67 <b>Ho</b> holmium 164.93 (164.93, 164.93)	68 <b>Er</b> erbium 167.26 (167.26, 167.26)	69 <b>Tm</b> thulium 168.93 (168.93, 168.93)	70 <b>Yb</b> ytterbium 173.05 (173.05, 173.05)	71 <b>Lu</b> lutetium 174.97 (174.97, 174.97)	89 <b>Ac</b> actinium 227.03 (227.03, 227.03)	90 <b>Th</b> thorium 232.04 (232.04, 232.04)	91 <b>Pa</b> protactinium 231.04 (231.04, 231.04)	92 <b>U</b> uranium 238.03 (238.03, 238.03)	93 <b>Np</b> neptunium 237 (237, 237)	94 <b>Pu</b> plutonium 244 (244, 244)	95 <b>Am</b> americium 243 (243, 243)	96 <b>Cm</b> curium 247 (247, 247)	97 <b>Bk</b> berkelium 247 (247, 247)	98 <b>Cf</b> californium 251 (251, 251)	99 <b>Es</b> einsteinium 252 (252, 252)	100 <b>Fm</b> fermium 257 (257, 257)	101 <b>Md</b> mendelevium 258 (258, 258)	102 <b>No</b> nobelium 259 (259, 259)	103 <b>Lr</b> lawrencium 262 (262, 262)
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For notes and updates to this table, see [www.iupac.org](http://www.iupac.org). This version is dated 1 December 2018.  
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# Nucleons are the essential building blocks of Matter!

# The Nuclei of the Three Isotopes of Hydrogen

Protium

Deuterium

Tritium



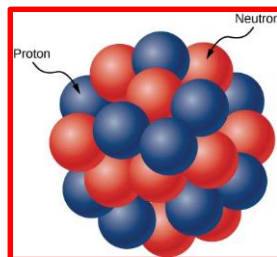
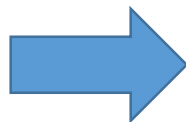
1 proton



1 proton  
1 neutron



1 proton  
2 neutrons



Z, number of protons

proton-rich nuclei

Z=82

<sup>208</sup>Pb

N=126

<sup>100</sup>Sn

N=Z

Z=50

<sup>132</sup>Sn

N=82

<sup>56</sup>Ni

Z=28

<sup>78</sup>Ni

N=50

neutron-rich nuclei

<sup>40</sup>Ca

Z=20

N=28

N=20

Z=8

N=8

- stable
- α decay
- β<sup>+</sup> decay
- β<sup>-</sup> decay
- fission
- magic numbers

N, number of neutrons





# What are D & K mesons?

## Standard Model of Elementary Particles

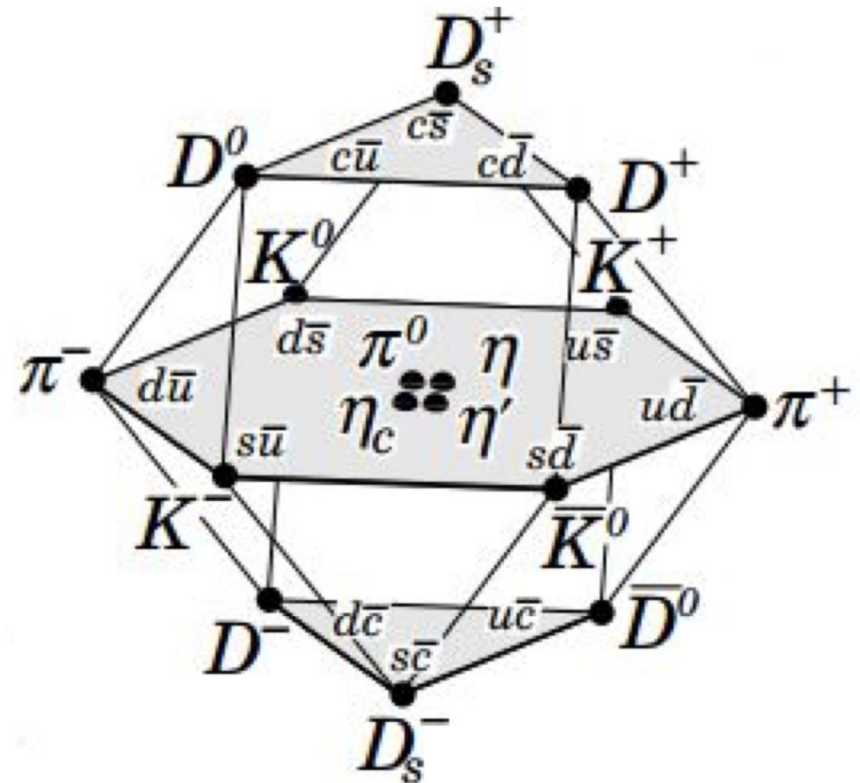
three generations of matter (fermions)						interactions / force carriers (bosons)			
I		II		III					
$\approx 2.2 \text{ MeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	<b>u</b> up	$\approx 1.28 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	<b>c</b> charm	$\approx 173.1 \text{ GeV}/c^2$ $\frac{2}{3}$ $\frac{1}{2}$	<b>t</b> top	0 0 1	<b>g</b> gluon	$\approx 125.09 \text{ GeV}/c^2$ 0 0 0	<b>H</b> higgs
$\approx 4.7 \text{ MeV}/c^2$ $-\frac{1}{3}$ $\frac{1}{2}$	<b>d</b> down	$\approx 1.28 \text{ GeV}/c^2$ $-\frac{1}{3}$ $\frac{1}{2}$	<b>s</b> strange	$\approx 4.18 \text{ GeV}/c^2$ $-\frac{1}{3}$ $\frac{1}{2}$	<b>b</b> bottom	0 0 1	<b><math>\gamma</math></b> photon		
$\approx 0.511 \text{ MeV}/c^2$ -1 $\frac{1}{2}$	<b>e</b> electron	$\approx 1.0566 \text{ MeV}/c^2$ -1 $\frac{1}{2}$	<b><math>\mu</math></b> muon	$\approx 1.7768 \text{ GeV}/c^2$ -1 $\frac{1}{2}$	<b><math>\tau</math></b> tau	0 0 1	<b>Z</b> Z boson		
$< 2.2 \text{ eV}/c^2$ 0 $\frac{1}{2}$	<b><math>\nu_e</math></b> electron neutrino	$< 7 \text{ MeV}/c^2$ 0 $\frac{1}{2}$	<b><math>\nu_\mu</math></b> muon neutrino	$< 15.5 \text{ MeV}/c^2$ 0 $\frac{1}{2}$	<b><math>\nu_\tau</math></b> tau neutrino	$\pm 1$ 1 1	<b>W</b> W boson		

**QUARKS**

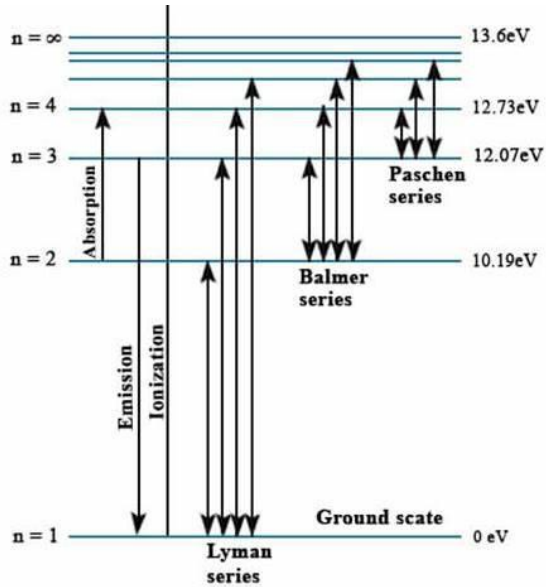
**LEPTONS**

**GAUGE BOSONS VECTOR BOSONS**

**SCALAR BOSONS**

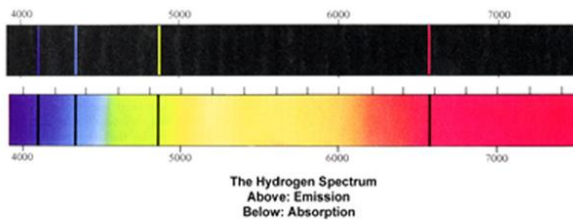
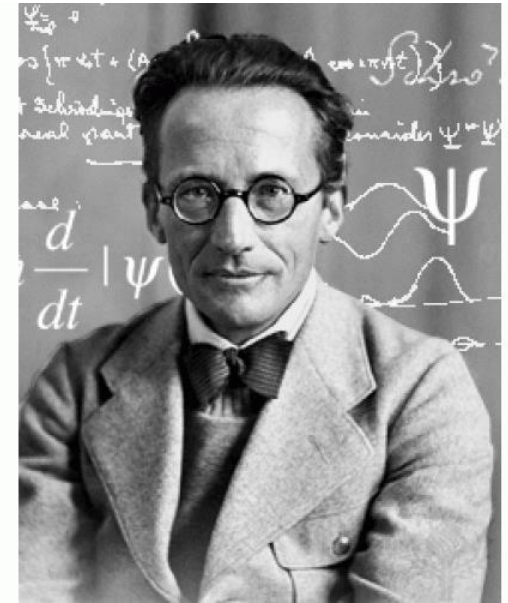






$$\left( \frac{\hat{p}^2}{2\mu} + V \right) \Psi = E\Psi$$

$$V = - \frac{e^2}{4\pi\epsilon_0 r^2}$$

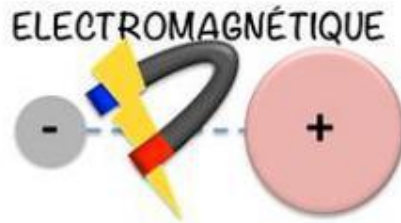


## Schroedinger equation & Hydrogen spectrum

# Strong interaction is much harder

$$F = \frac{kq_1q_2}{r^2}$$

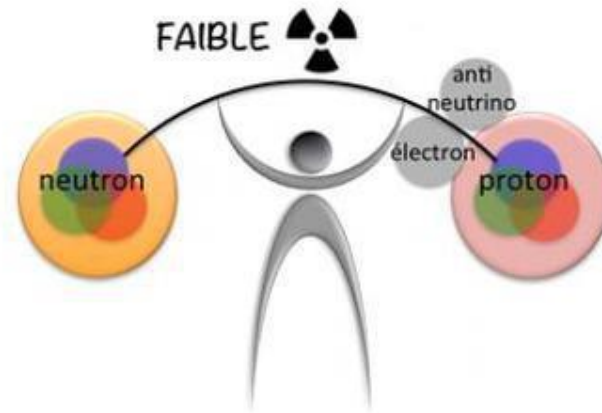
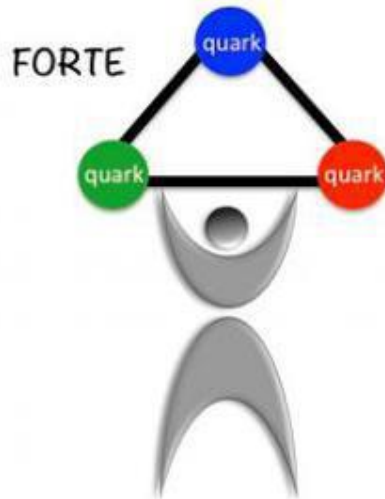
$$F = q\vec{v} \times \vec{B}$$



$$F = \frac{Gm_1m_2}{r^2}$$

$$F = \dots\dots\dots$$

....



$$F = G_F$$

# One of the most difficult problems

in the history of mankind

**SCIENTIFIC AMERICAN**, September 1953

## What Holds the Nucleus Together?

by Hans A. Bethe

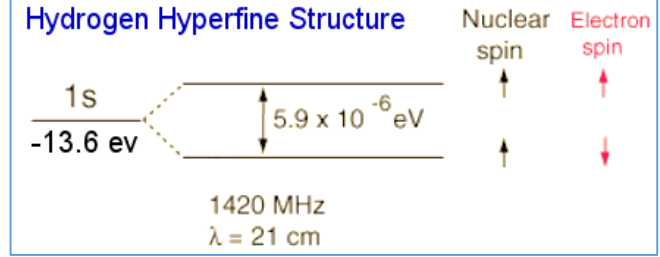
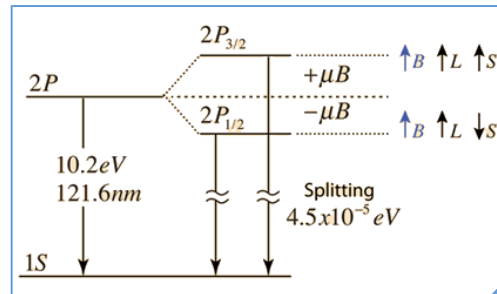
In the past quarter century physicists have devoted a huge amount of experimentation and mental labor to this problem – probably more man-hours than have been given to any other scientific question in the history of mankind.



Hans Bethe  
Nobel Prize in Physics  
1967

高精度协变手征核力—北航

# Why spectroscopy—Atomic



$$\frac{1}{\lambda_{\text{vac}}} = R \left( \frac{1}{n_1^2} - \frac{1}{n_2^2} \right)$$

## Fine structure

## Hyperfine structure



Johannes Rydberg

## Rydberg formula



Niels Henrik David Bohr



Edward W. Morley



Arnold Sommerfeld



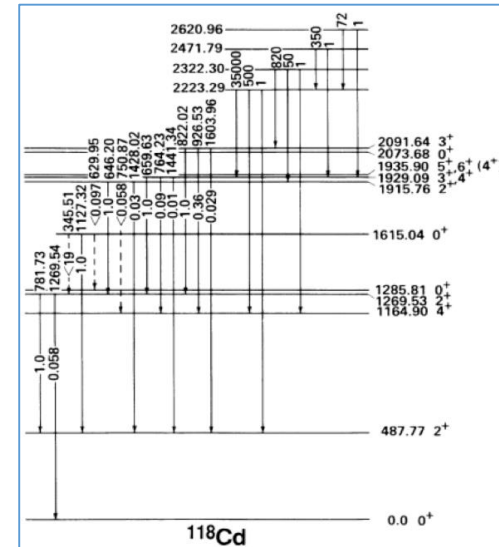
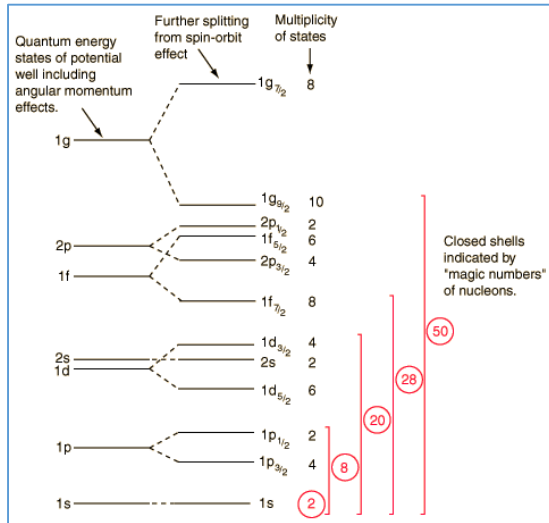
Albert Abraham Michelson



Wolfgang Pauli

# Why spectroscopy—Nuclear

Single  
particle  
motion



Collective  
motion



Eugene Paul Wigner Maria Goeppert Mayer J. Hans D. Jensen

Aage Niels Bohr, Ben Roy Mottelson Leo James Rainwater

# Hadron spectroscopy—QM—QCD

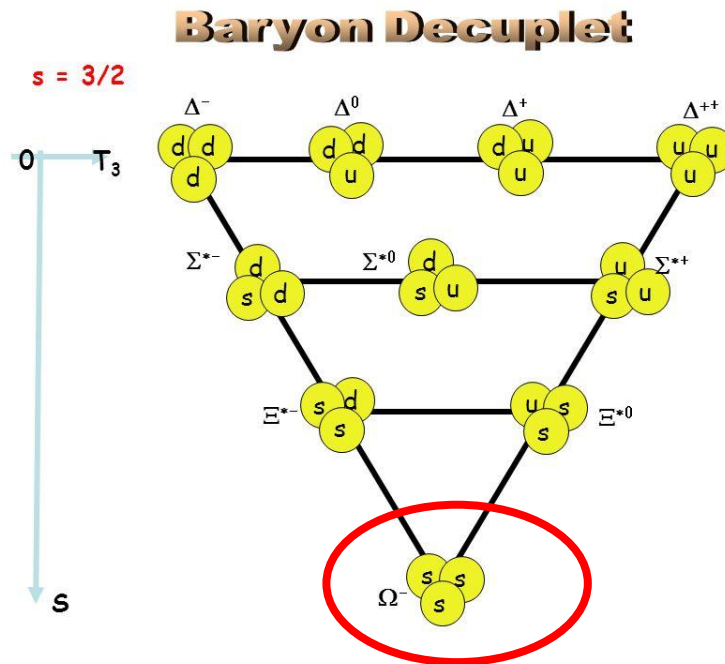
*Put an end to the then chaotic situation*



Murray Gell-Mann



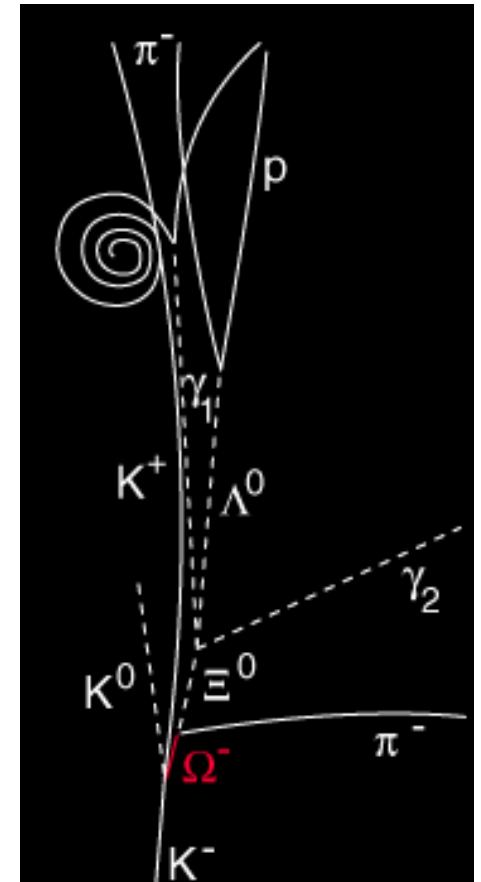
Yuval Ne'eman.



Eightfold  
way-1961

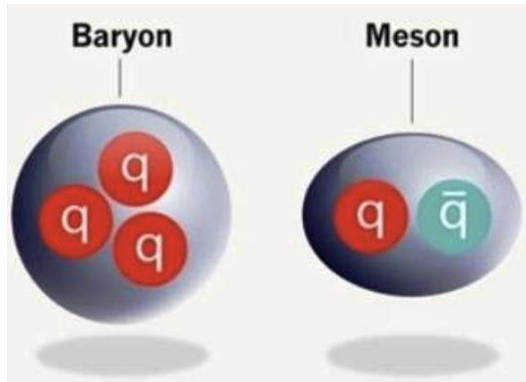
Quark  
model-  
1964

QCD/SM



V. E. Barnes et al., Phys. Rev. Lett. 12, 204 (1964)

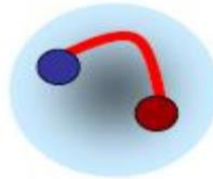
# Beyond QM



In the **naïve quark model**

In principle,  
QCD **allows**

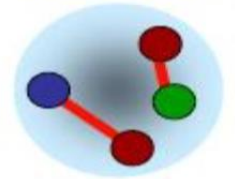
Hybrid



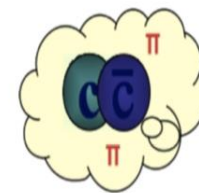
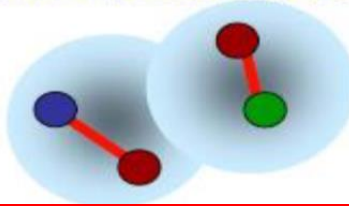
Glueball



Tetraquark

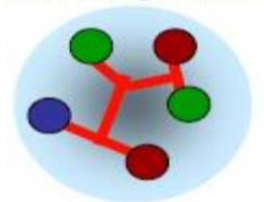


Hadronic molecule



Hadro-  
quarkonium

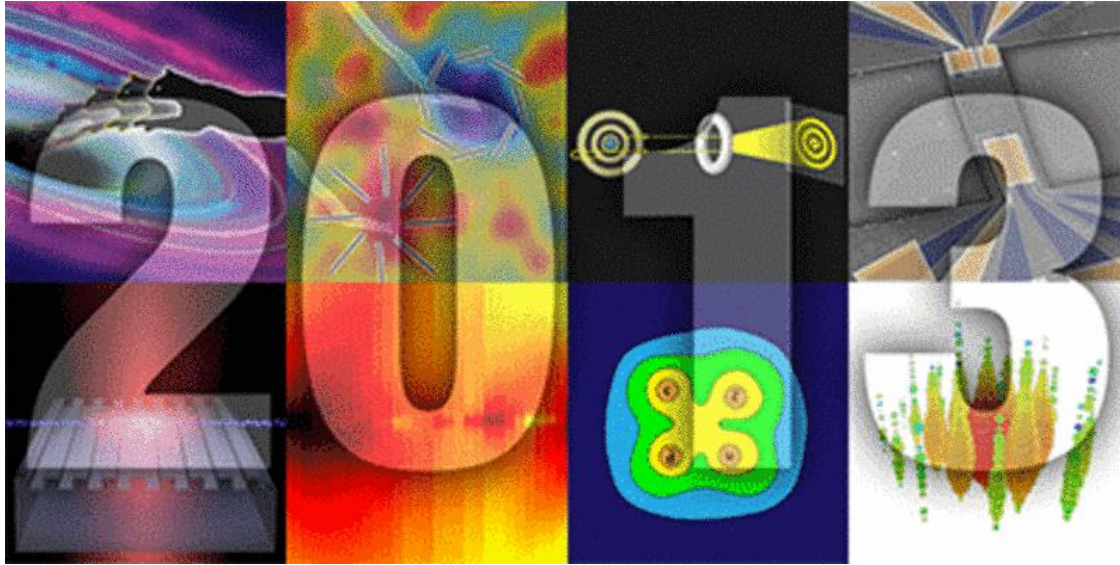
Pentaquark





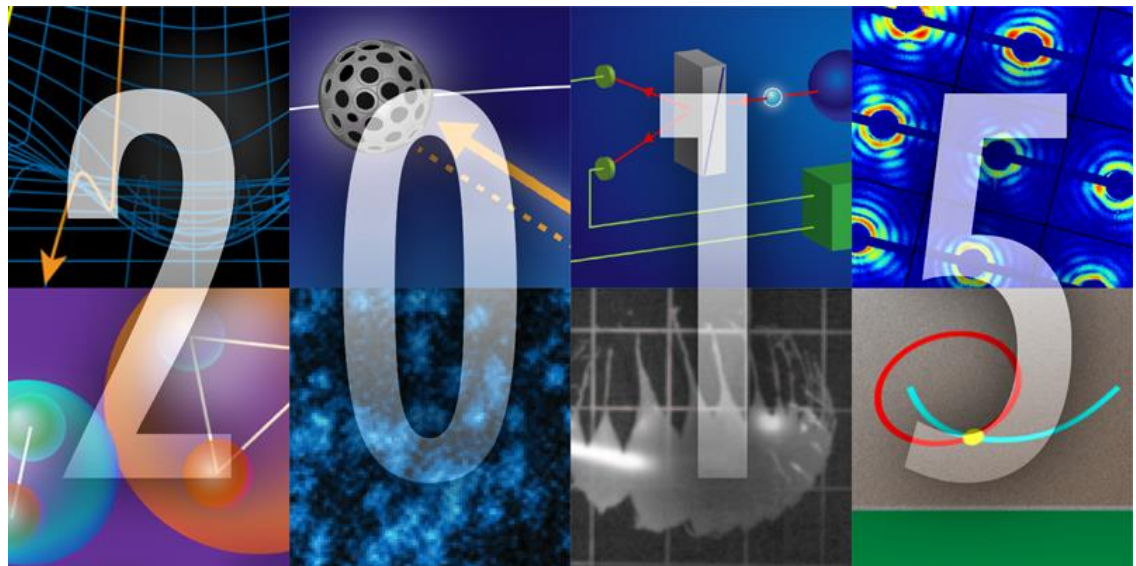
# Highlights of the year

the research covered in Physics  
that **really made waves in and  
beyond the physics community.**

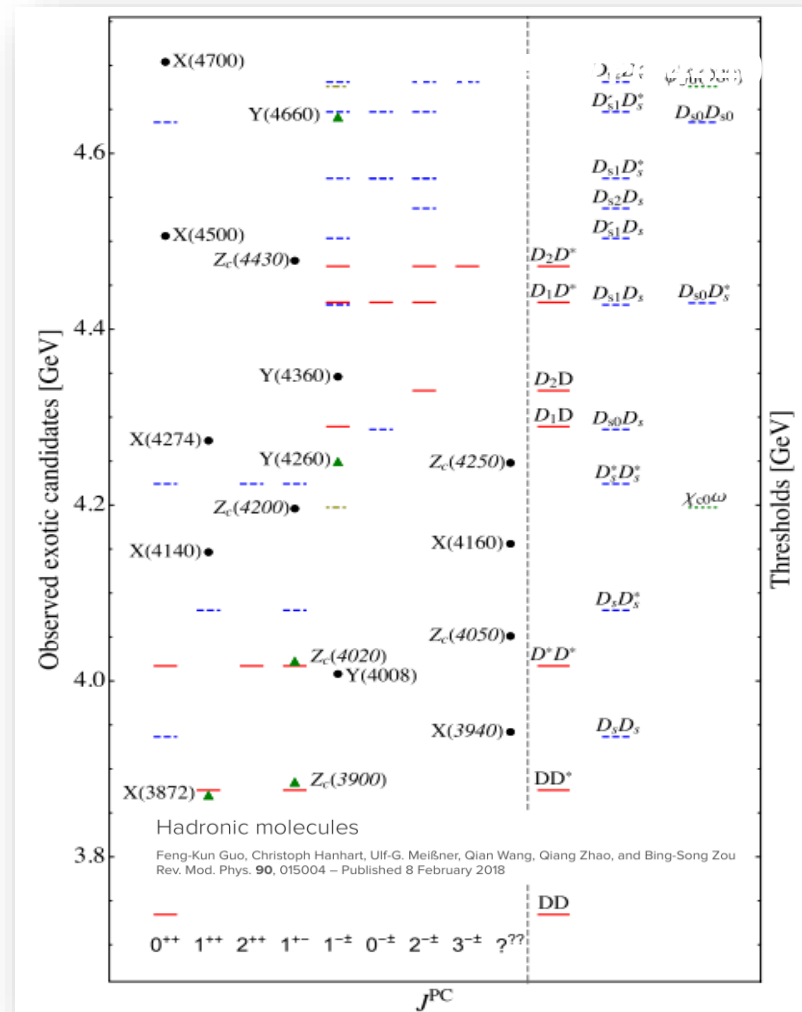


**Four-Quark Matter/BESIII**

**Particle High Five/LHCb**

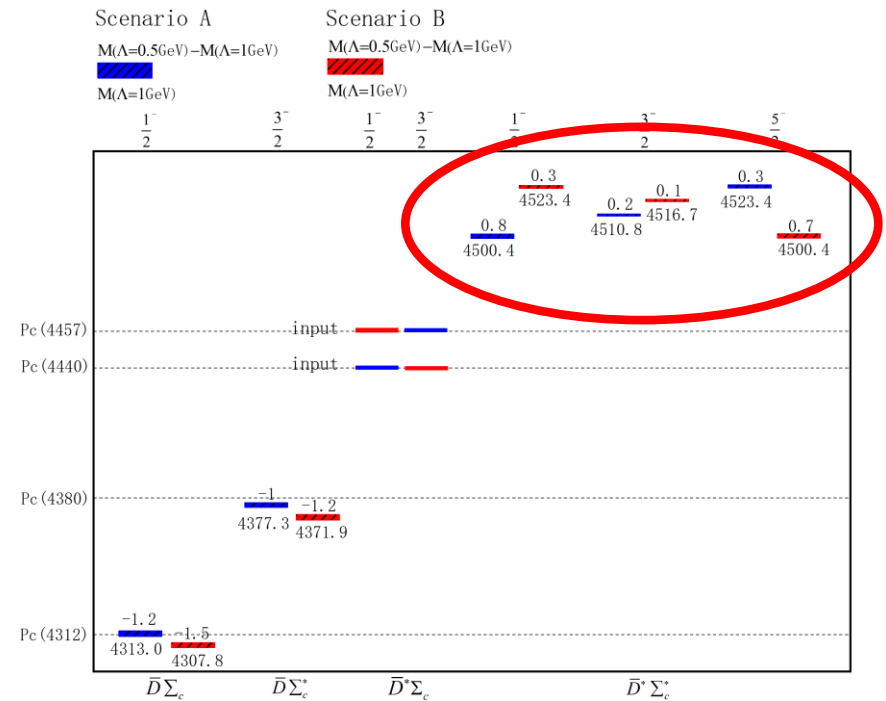
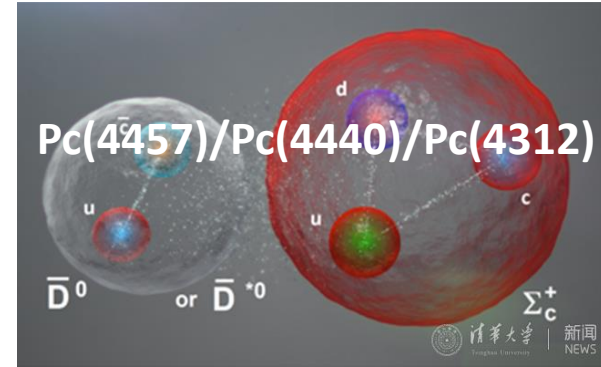
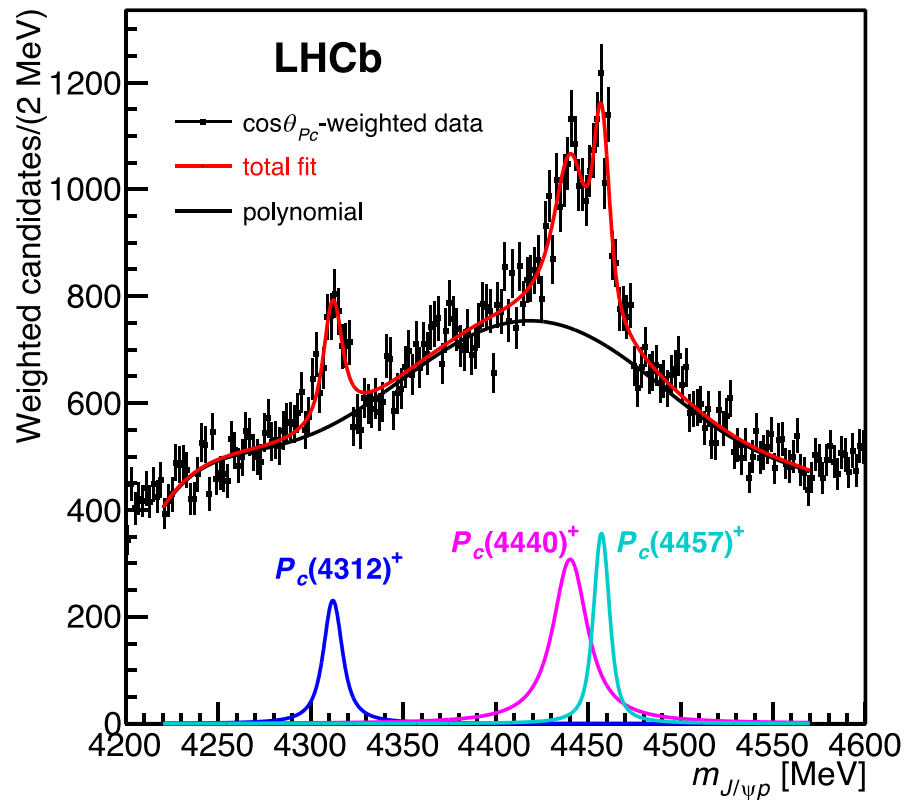


# Most of them close to thresholds



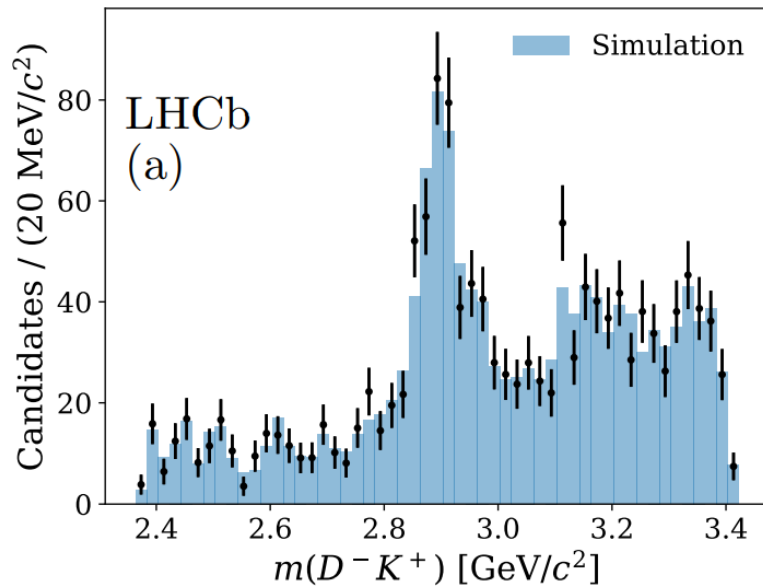
Feng-Kun Guo, Christoph Hanhart, Ulf-G. Meißner, Qian Wang, Qiang Zhao, Bing-Song Zou.  
Rev.Mod.Phys. 90 (2018) 015004.

# Fine structure in Pc

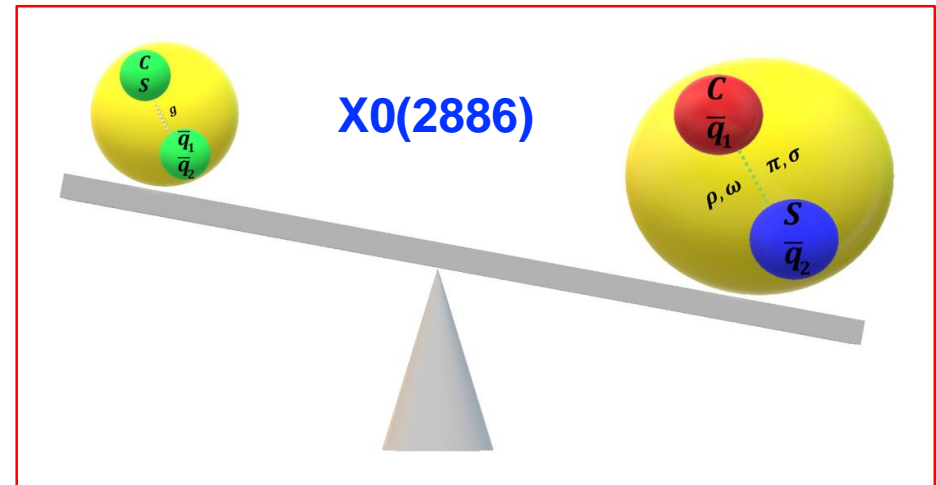


Phys.Rev.Lett. 122 (2019) 242001

# Another new state



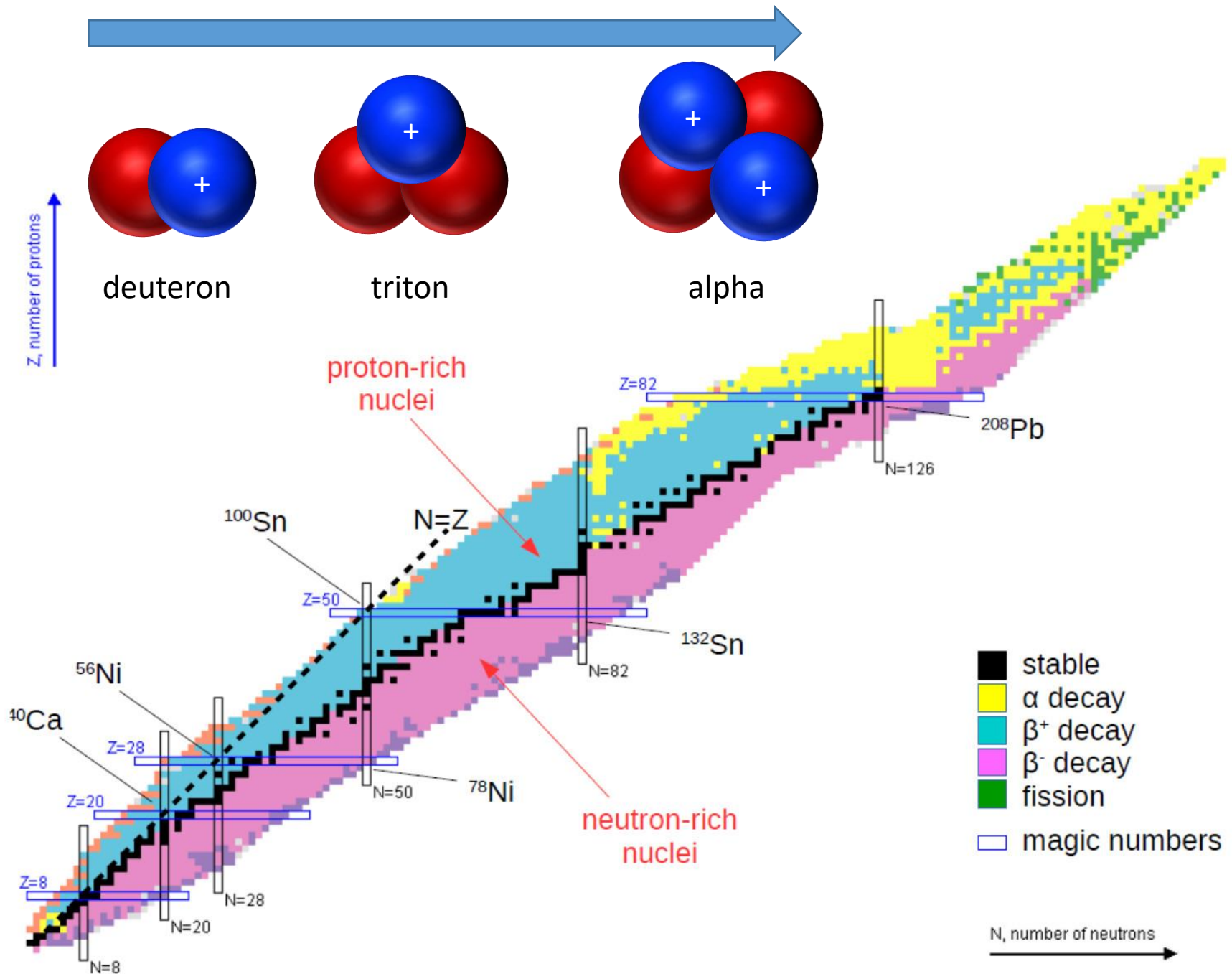
2009.00025



[2008.07389](#), PRD rapid communication

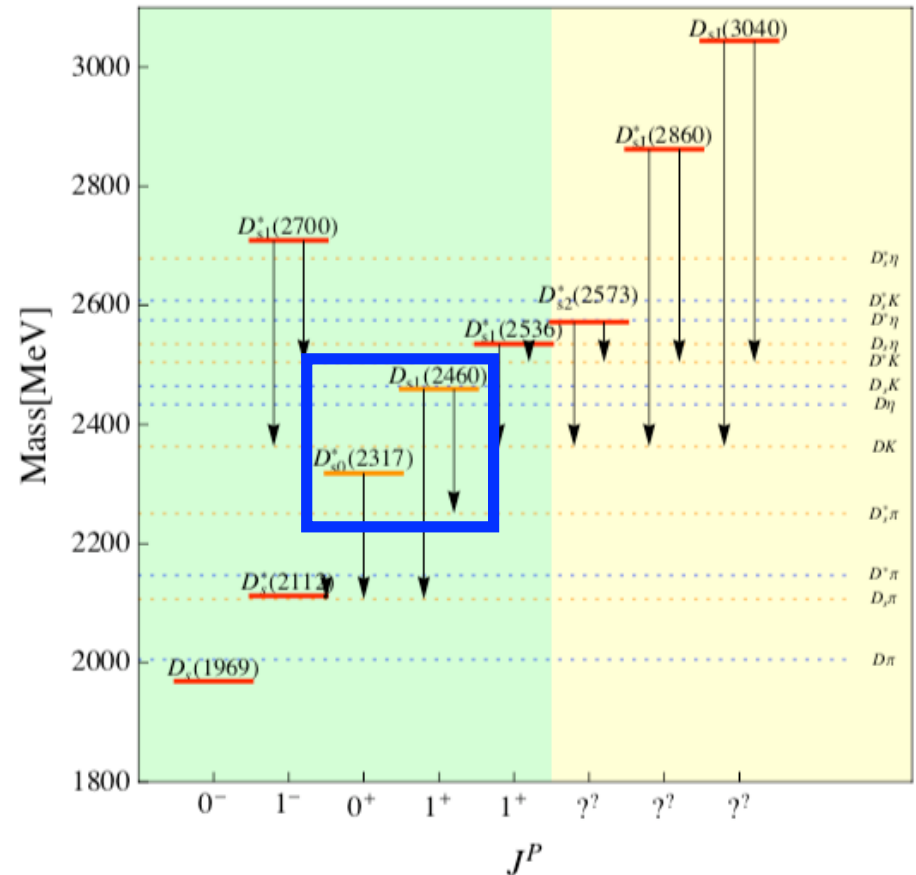
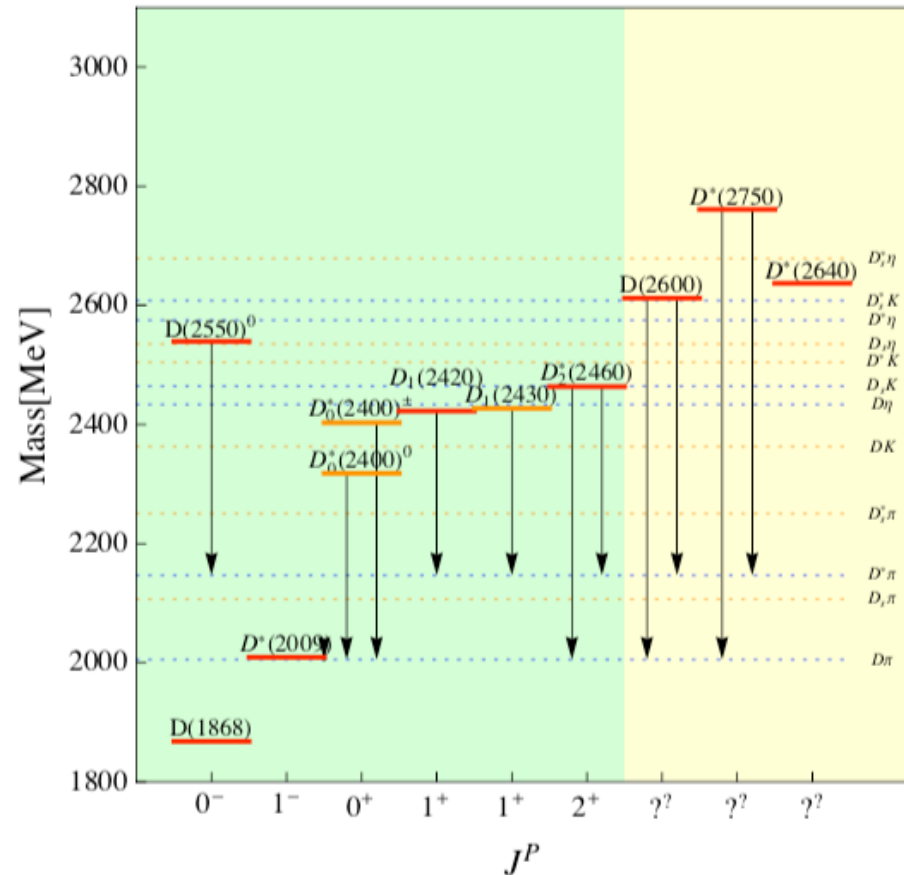
How to check the molecular picture?  
our naïve answer—go to many body







# Next best two-body molecule candidates



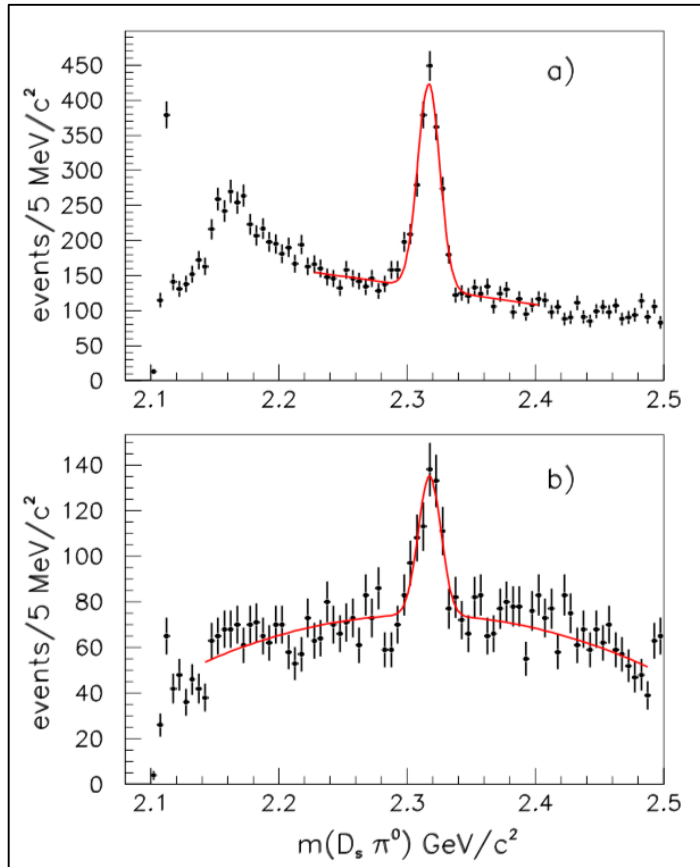
**$D_{s0}^*(2317)$**

**$D_{s1}(2460)$**



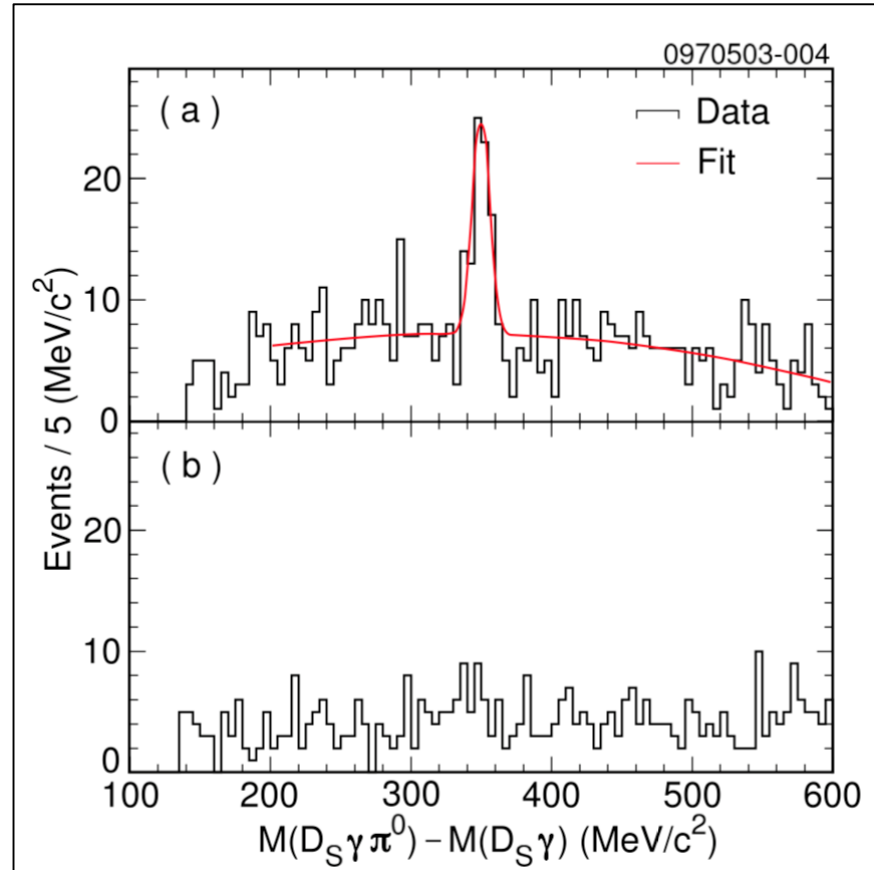
# Discovery channels

## Ds0\*(2317)



BaBar PRL90,242001(2003)

## Ds1(2460)



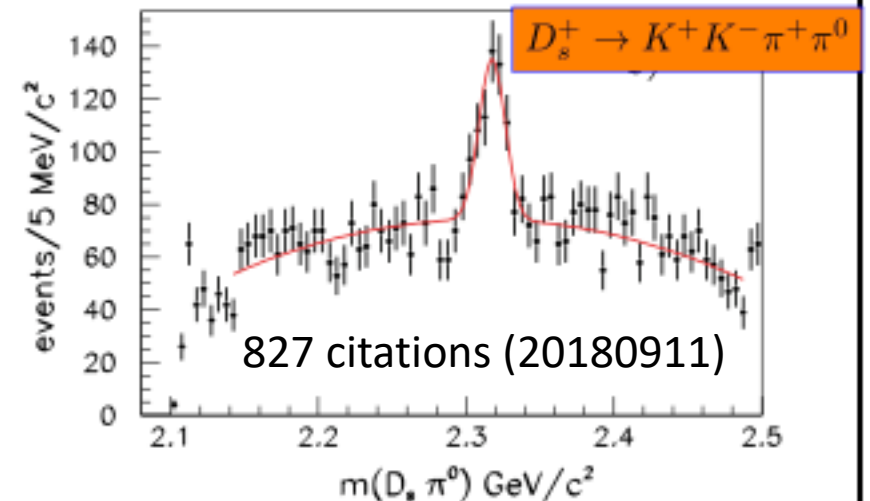
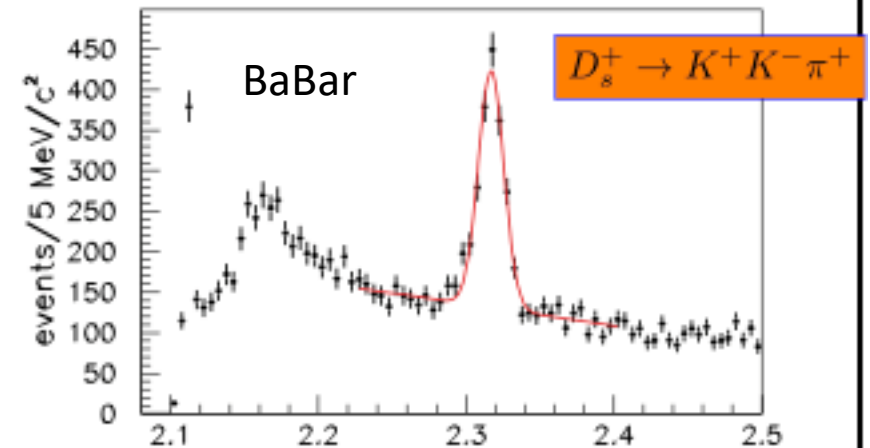
CLEO PRD68,032002(2003)

# What are special about these two states

- $D_{s0}^*(2317)$ ,  $D_{s1}(2460)$
- 160/70 MeV lower than quark model predictions--difficult to be understood as conventional  $c\bar{s}$ bar states.
- “Dynamically generated” from strong DK interaction
  - ✓ E. E. Kolomeitsev 2004, [SEP]
  - ✓ F. K. Guo 2006,
  - ✓ D. Gamermann 2007

$$m_{D_{s1}(2460)} - m_{D_{s0}^*(2317)} \approx m_{D^*} - m_D$$

$M = 2317.8 \pm 0.6$  and  $\Gamma < 3.8$  MeV



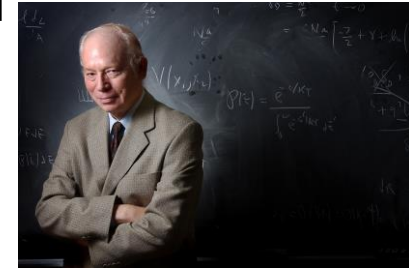
PRL90,242001(2003)

# Contents

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- Motivation: new types of clusters of color singlets in addition to nuclei—as a nontrivial test of the molecule picture
- $Ds0^*(2317)$  and  $Ds1(2460)$  as  $DK/D^*K$  molecules: theory&lattice
- An explicit study of the  $DDK$  system:  $R^{++}(4140)$
- $K^*(4307)$  as excited  $K^*$  with hidden charm
- Summary and outlook

# UChPT in Bethe-Salpeter equation

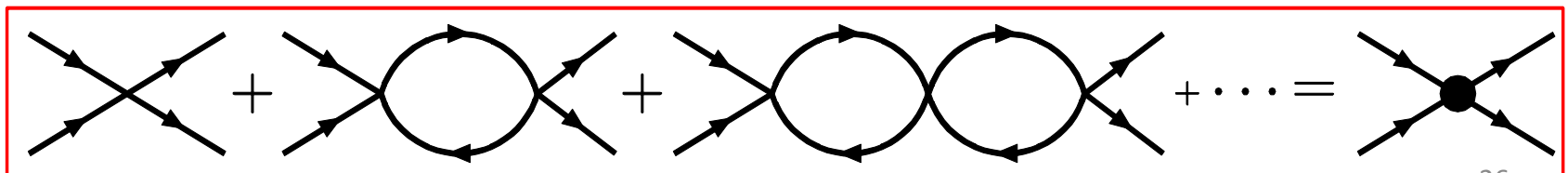


## □ Model independent DK interaction from ChPT

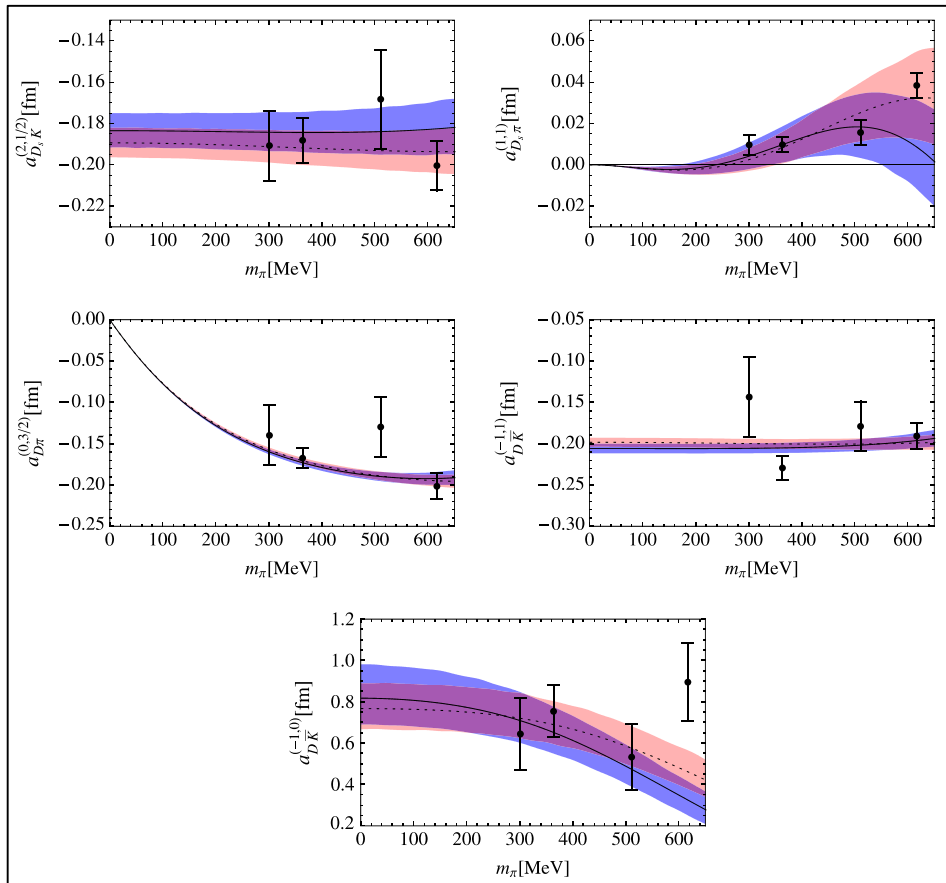
$$\mathcal{V}_{\text{WT}}(P(p_1)\phi(p_2) \rightarrow P(p_3)\phi(p_4)) = \frac{1}{4f_0^2} \mathcal{C}_{\text{LO}} (s - u) \quad \text{Weinberg-Tomazawa}$$

$$\begin{aligned} \mathcal{V}_{\text{NLO}}(P(p_1)\phi(p_2) \rightarrow P(p_3)\phi(p_4)) = & -\frac{8}{f_0^2} C_{24} \left( c_2 p_2 \cdot p_4 - \frac{c_4}{m_P^2} (p_1 \cdot p_4 p_2 \cdot p_3 + p_1 \cdot p_2 p_3 \cdot p_4) \right) \\ & -\frac{4}{f_0^2} C_{35} \left( c_3 p_2 \cdot p_4 - \frac{c_5}{m_P^2} (p_1 \cdot p_4 p_2 \cdot p_3 + p_1 \cdot p_2 p_3 \cdot p_4) \right) \\ & -\frac{4}{f_0^2} C_6 \frac{c_6}{m_P^2} (p_1 \cdot p_4 p_2 \cdot p_3 - p_1 \cdot p_2 p_3 \cdot p_4) \\ & -\frac{8}{f_0^2} C_0 c_0 + \frac{4}{f_0^2} C_1 c_1, \end{aligned} \quad (11)$$

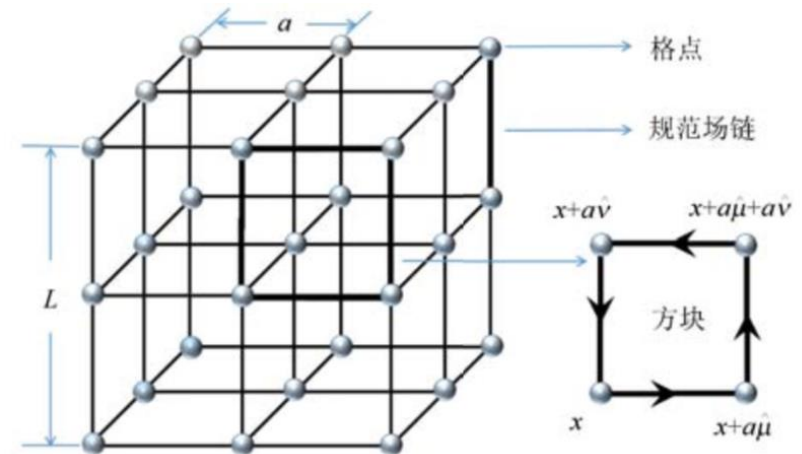
## □ Resummed in the Bethe-Salpeter equation (two-body elastic unitarity)



# Fixing the LECs using latest LQCD\* data



- NLO ChPT kernel: 5 LECs
- A quite good description of the 20 Lattice **scattering lengths of pseudoscalar mesons and D mesons** (I=0 DK excluded) can be achieved.



# Ds0 and Ds1 dynamically generated

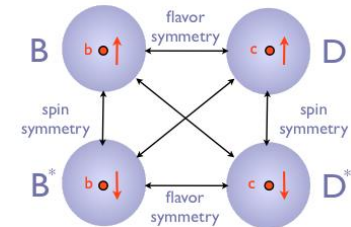
## ● Charm sector

“Post-diction”

$D_{s0}^*(2317)$ ,  $D_{s1}(2460)$

TABLE V. Pole positions  $\sqrt{s} = M - i\frac{\Gamma}{2}$  (in units of MeV) of charm mesons dynamically generated in the HQS UChPT.

$(S, I)$	$J^P = 0^+$	$J^P = 1^+$
$(1, 0)$	$2317 \pm 10$	$2457 \pm 17$
$(0, 1/2)$	$(2105 \pm 4) - i(103 \pm 7)$	$(2248 \pm 6) - i(106 \pm 13)$



## ● Bottom Sector

TABLE VI. Pole positions  $\sqrt{s} = M - i\frac{\Gamma}{2}$  (in units of MeV) of bottom mesons dynamically generated in the HQS UChPT.

$(S, I)$	$J^P = 0^+$	$J^P = 1^+$
$(1, 0)$	$5726 \pm 28$	$5778 \pm 26$
$(0, 1/2)$	$(5537 \pm 14) - i(118 \pm 22)$	$(5586 \pm 16) - i(124 \pm 25)$

# Predicted $B_s0$ and $B_s1$ states

Physics Letters B 750 (2015) 17–21



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Physics Letters B

[www.elsevier.com/locate/physletb](http://www.elsevier.com/locate/physletb)



## Predicting positive parity $B_s$ mesons from lattice QCD



C.B. Lang<sup>a</sup>, Daniel Mohler<sup>b,\*</sup>, Sasa Prelovsek<sup>c,d</sup>, R.M. Woloshyn<sup>e</sup>

<sup>a</sup> Institute of Physics, University of Graz, A-8010 Graz, Austria

<sup>b</sup> Fermi National Accelerator Laboratory, Batavia, IL 60510-5011, USA

<sup>c</sup> Department of Physics, University of Ljubljana, 1000 Ljubljana, Slovenia

<sup>d</sup> Jozef Stefan Institute, 1000 Ljubljana, Slovenia

<sup>e</sup> TRIUMF, 4004 Wesbrook Mall, Vancouver, BC V6T 2A3, Canada

**Table 5**

Comparison of masses from this work to results from various model based calculations; all masses in MeV.

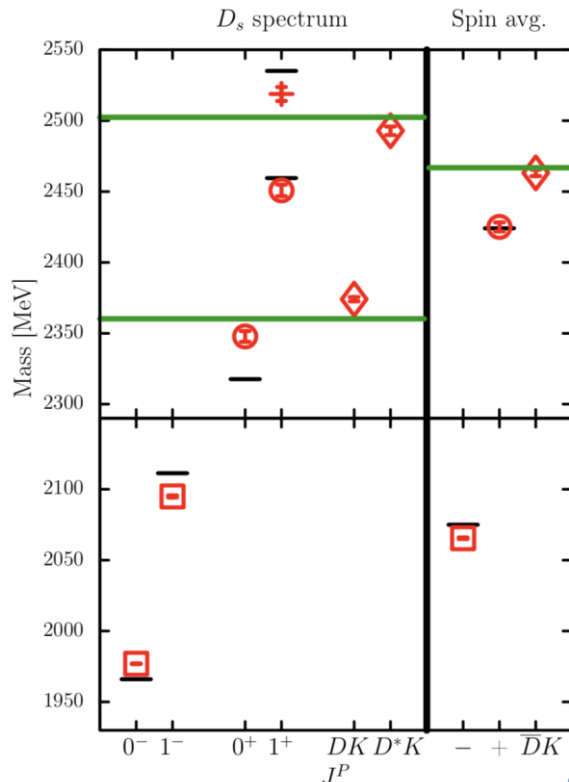
$J^P$	$0^+$	$1^+$
Covariant (U)ChPT [24]	5726(28)	5778(26)
NLO UHMChPT [19]	5696(20)(30)	5742(20)(30)
LO UChPT [17,18]	5725(39)	5778(7)
LO $\chi$ -SU(3) [16]	5643	5690
HQET + ChPT [20]	5706.6(1.2)	5765.6(1.2)
Bardeen, Eichten, Hill [15]	5718(35)	5765(35)
rel. quark model [5]	5804	5842
rel. quark model [22]	5833	5865
rel. quark model [23]	5830	5858
HPQCD [30]	5752(16)(5)(25)	5806(15)(5)(25)
this work	5713(11)(19)	5750(17)(19)

In agreement with LQCD



# More support from recent IQCD studies

- [G.K.C. Cheung et al., arXiv:2008.06432\[hep-lat\].](#)
- [G. S. Bali et al., arXiv:1706.01247 \[hep-lat\].](#)
- C. B. Lang et al., arXiv:1403.8103 [hep-lat].
- D. Mohler et al., arXiv:1308.3175 [hep-lat].



**“DK components substantial”**

FIG. 12. On the left, our final results for the lower lying  $D_s$  spectrum as detailed in Table VII. The short horizontal black lines indicate the corrected experimental values (see Section II) while the green horizontal lines give the positions of the  $DK$  and  $D^*K$  non-interacting thresholds. Our lattice results for the finite volume thresholds are labelled  $DK$  and  $D^*K$ , respectively. The errors indicated are statistical only. On the right, the negative parity spin-averaged  $1S$  mass  $m_- = \frac{1}{4}(m_{0^-} + 3m_{1^-})$  is shown and denoted  $-$ , while the same spin-average of the positive parity  $0^+$  and  $1^+$  states is labelled with  $+$  and the weighted average of the threshold is labelled as  $\overline{DK}$ .

# Further **tests** of the DK interaction

---

- Experiments, theory, and lattice QCD all show that  $DK$  or  $D^*K$  interaction is strong enough to form  $Ds0^*(2317)$  or  $Ds1(2460)$
- A natural question is: if we add one more  $D(\bar{D})$  or  $D^*(\bar{D}^*)$ , can they form molecules of three hadrons?
- This seems to **be a rather straightforward and naive question**, but **remains unexplored** until quite recently

# Contents

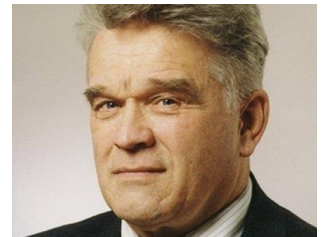
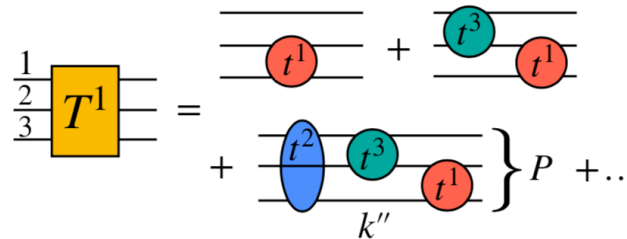
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- Motivation: new types of clusters of color singlets in addition to nuclei—as a nontrivial test of the molecule picture
- $Ds0^*(2317)$  and  $Ds1(2460)$  as DK/ $D^*K$  molecules
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# An **explicit three-body** study of DDK

- Coupled-three-channel problem:  $D(\text{DK} - D_S\pi - D_S\eta)$
- Three-body scattering matrix (Faddeev)

$$T = \sum_{i=1}^3 T^i$$



$$T^i = t^i \delta^3(\vec{k}'_i - \vec{k}_i) + \sum_{j \neq i=1}^3 T_R^{ij}, \quad i = 1, 2, 3,$$

$$T_R^{ij} = t^i g^{ij} t^j + t^i \left[ G^{iji} T_R^{ji} + G^{ijk} T_R^{jk} \right],$$

**A. Martínez Torres**, K. P. Khemchandani, and E. Oset PRC **77**, 042203(R)

**A. Martínez Torres**, K.P. Khemchandani, **LSG**, M. Napsuciale, E. Oset, PRD78 (2008) 074031

# Two-body inputs

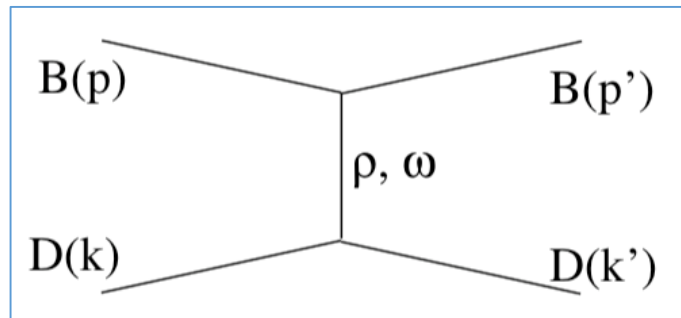
- **DK: leading order UChPT**  $DK$ ,  $D_s\eta$  and  $D_s\pi$

$$V_{ij} = -\frac{C_{ij}}{4f^2}(s - u)$$

$$a(\mu) = -1.846, \mu = 1000 \text{ MeV} \Rightarrow \text{Pole} = 2318 \text{ MeV}$$

*F.-K. Guo, P.-N. Shen, H.-C. Chiang, R.-G. Ping, and B.-S. Zou, PL B641, 278 (2006).*

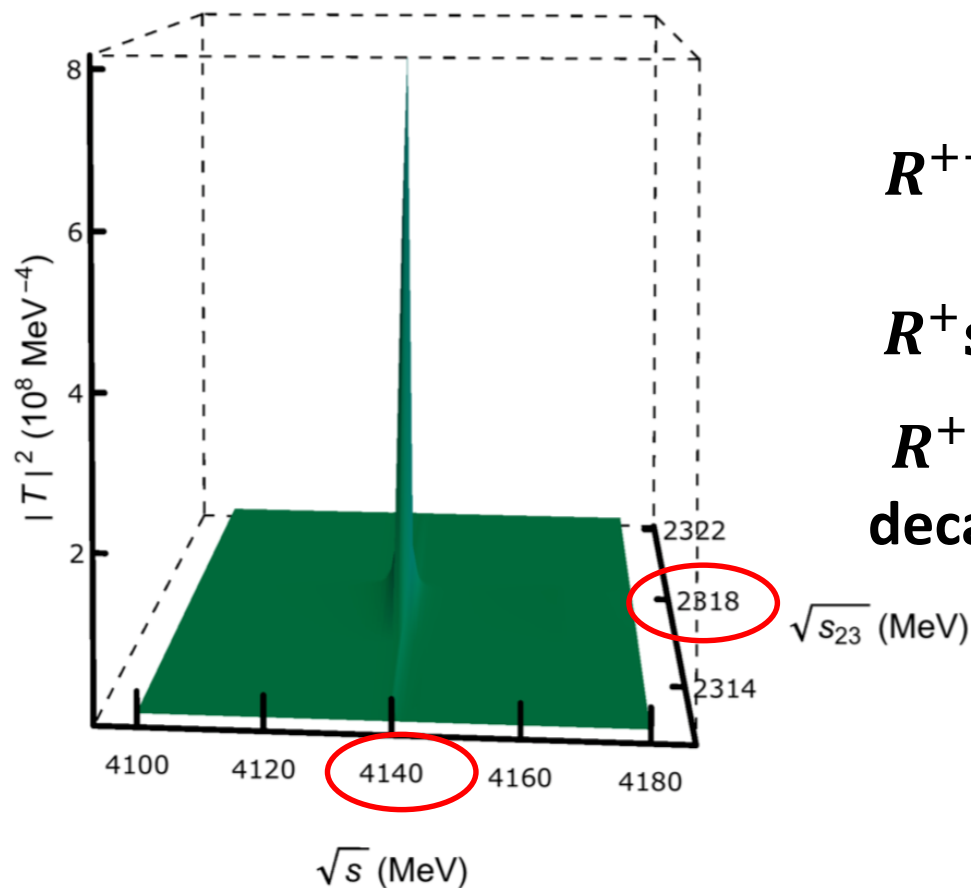
- **DD(Ds): local hidden gauge theory**



$$a(\mu) = -1.3 \sim -1.5, \mu = 1500 \text{ MeV} \Leftarrow \text{fixed from } D\bar{D}/D\bar{D}^* \rightarrow \chi(3700)/\chi(3872)$$

*S. Sakai, L. Roca, and E. Oset, PRD96, 054023 (2017).*

# Three-body amplitudes

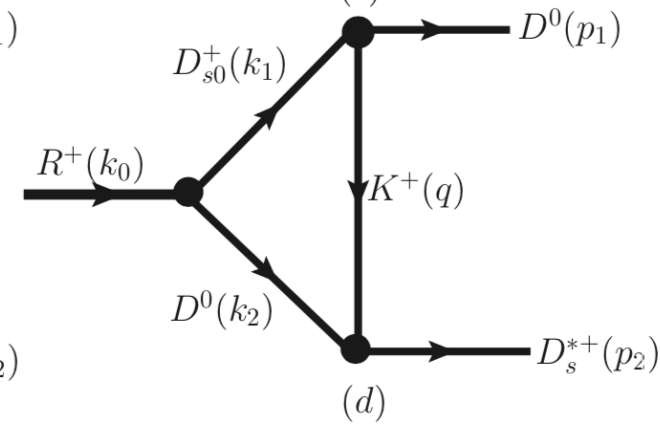
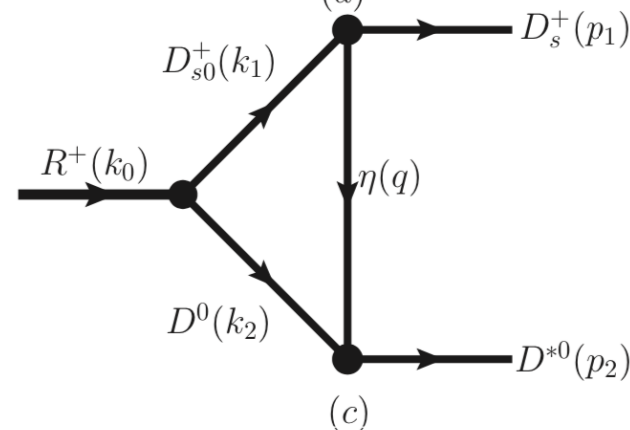
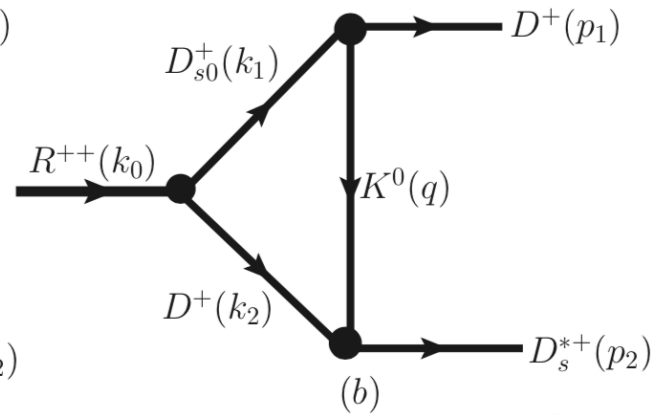
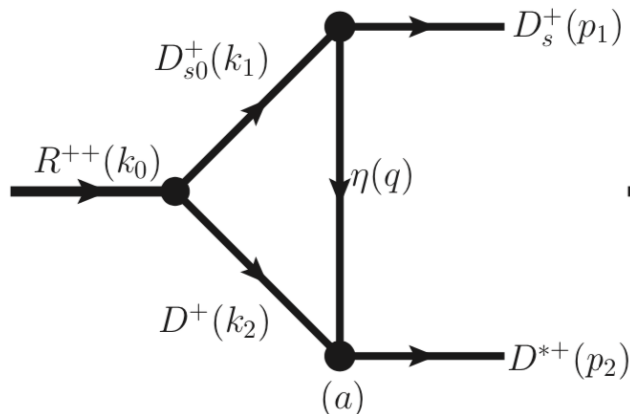
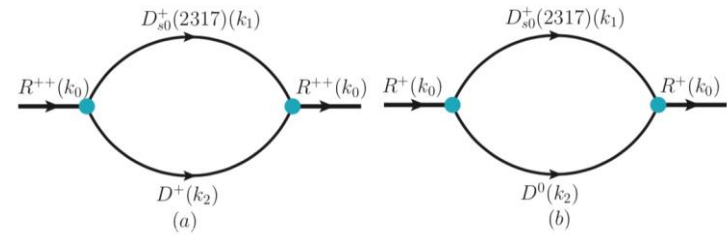


$$R^{++} = (I, I_{23}) = \left(\frac{1}{2}, 0\right)$$

$R^+$  should also exist

$R^{++}$  is a bound state, but can decay strongly

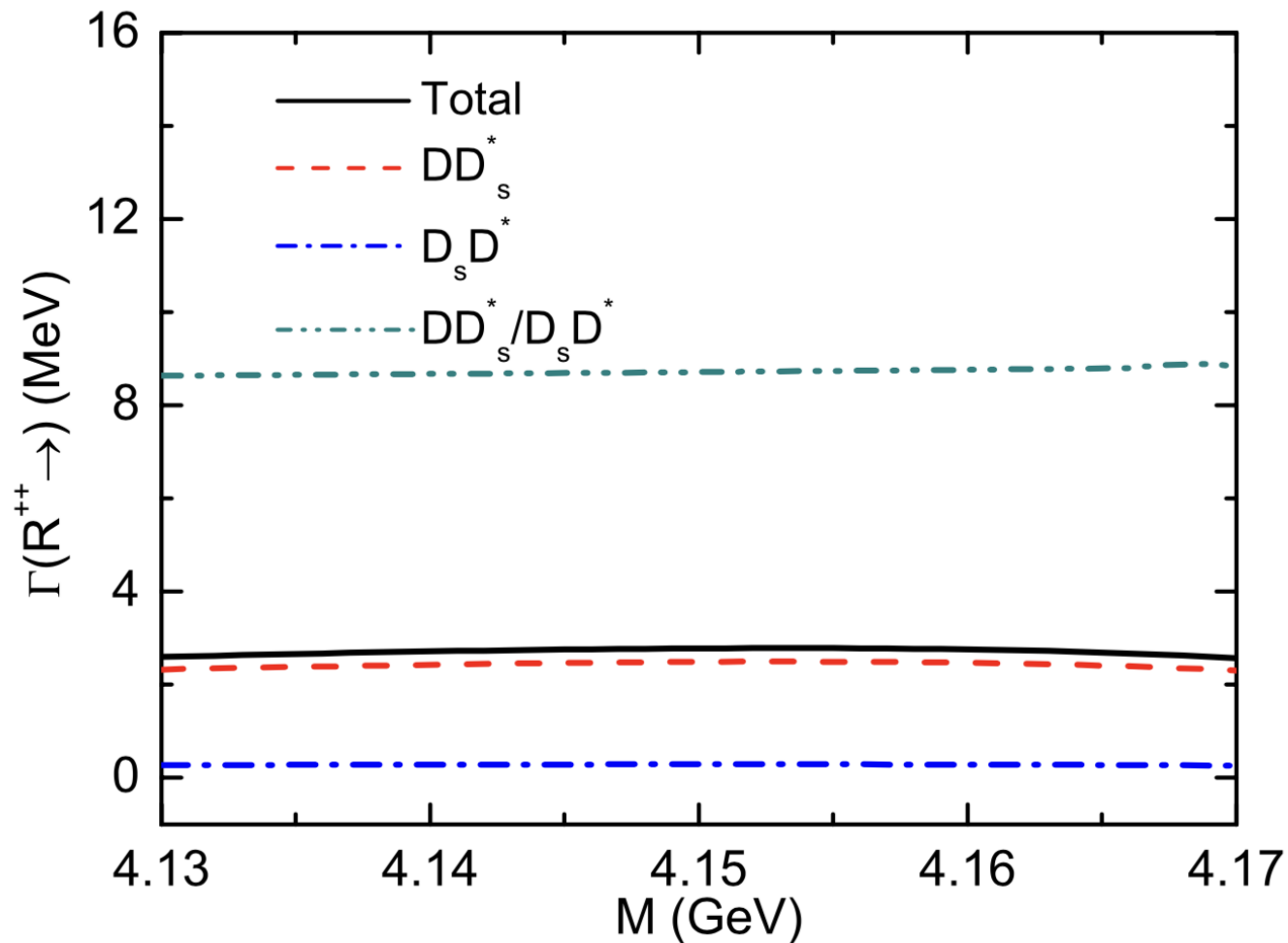
# Two-body decay width





# Two-body decay width

Kaon-Exchange Dominant

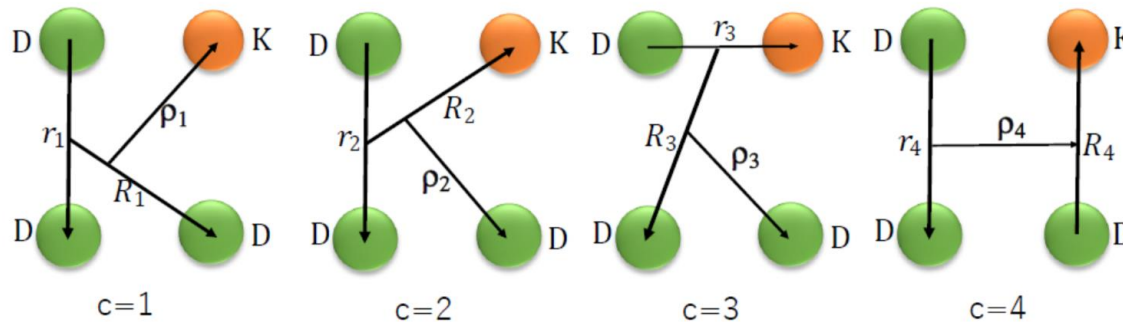
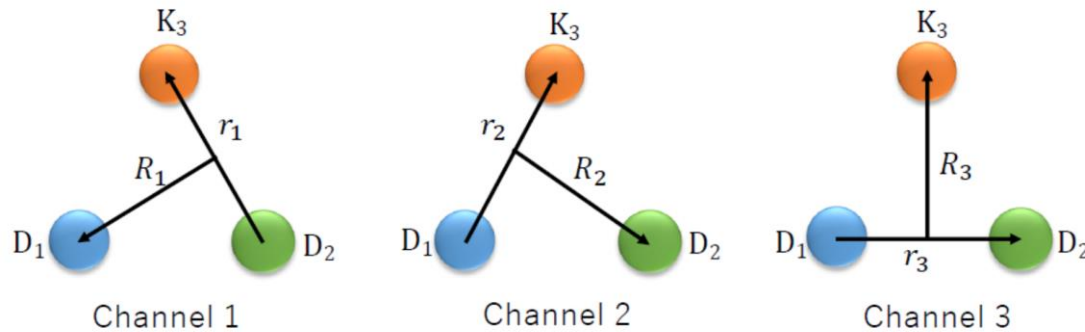


# A DDDK state

 $1(0^+)$ 

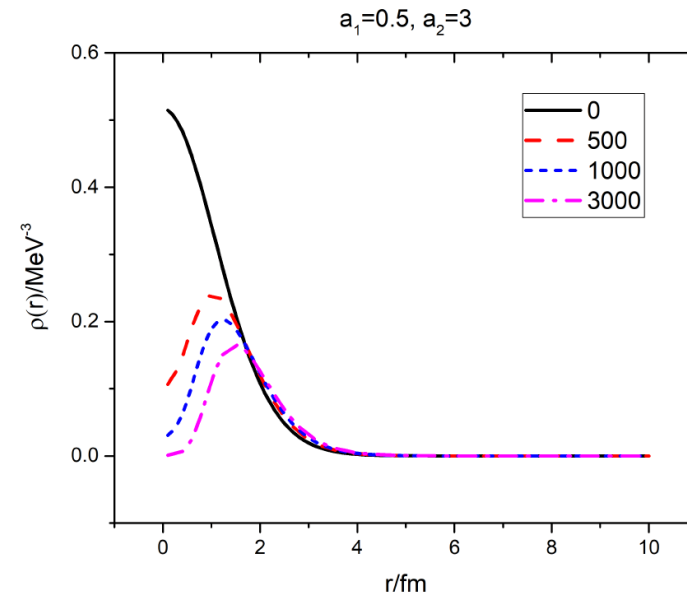
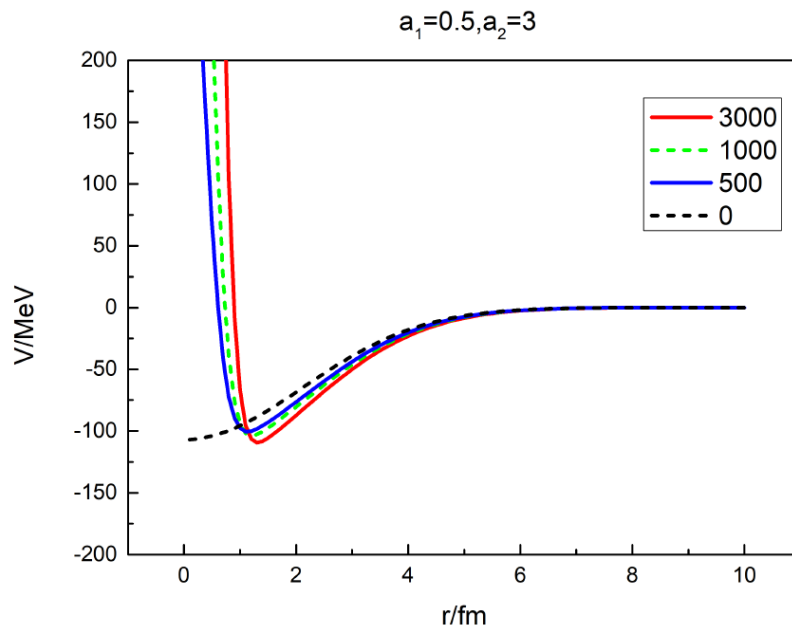
Gaussian Expansion Method

What if we add one more D?



# A DDDK state $1(0^+)$

What if we add one more D? **Our study** shows that such a state exists as well



Uncertainties are at **the order of 10-20 MeV**

$$V_{DK}(r) = C_1 e^{-r^2/a_1^2} + C_2 e^{-r^2/a_2^2}$$

	DK*	DDK	DDDK
Binding	45 MeV	(67-71) MeV	91-107 MeV

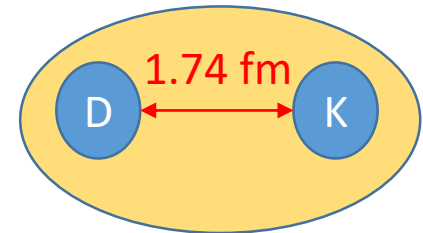
# DD interactions play a minor role

$\frac{C_S}{\pi R_S^3}$	$\frac{C(R_c)}{\pi R_c^3}$	$E_2$	$E_3(\text{only } V_{DK})$	$E_3(V_{DK} + V_{DD})$	$E_4(\text{only } V_{DK})$	$E_4(V_{DK} + V_{DD})$
$R_S = 0.5\text{fm}$			$R_c = 1\text{fm}$			
0	-320.1	-45.0	-65.8	-71.2	-89.4	-106.8
500	-455.4	-45.0	-65.8	-70.4	-89.2	-103.5
1000	-562.6	-45.0	-65.7	-69.7	-88.8	-101.4
3000	-838.7	-45.0	-65.0	-68.4	-87.0	-97.3
$R_S = 0.5\text{fm}$			$R_c = 2\text{fm}$			
0	-149.1	-45.0	-66.0	-68.8, -45.1	-88.7, -66.3	-97.6, -70.7
500	-178.4	-45.0	-65.9	-68.2, -45.5	-88.5, -66.7	-95.5, -70.9
1000	-195.0	-45.0	-65.8, -45.2	-67.9, -45.8	-88.2, -66.9	-94.5, -71.2
3000	-225.9	-45.0	-65.3, -45.6	-67.2, -46.6	-87.0, -67.0	-92.6, -71.7
$R_S = 0.5\text{fm}$			$R_c = 3\text{fm}$			
0	-107.0	-45.0	-66.2, -47.3	-68.0, -48.3	-88.8, -70.2	-94.4, -74.3
500	-119.4	-45.0	-66.2, -48.2	-67.7, -49.3	-88.7, -71.0	-93.2, -74.8
1000	-125.6	-45.0	-66.1, -48.7	-67.5, -49.8	-88.4, -71.3	-92.5, -75.2
3000	-136.2	-45.0	-65.8, -49.4	-67.1, -50.7	-87.6, -71.7	-91.4, -75.7

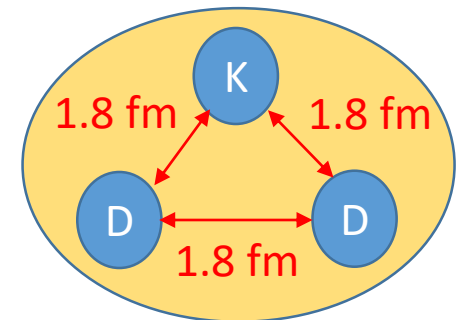
# Spatial distributions

$\frac{C_S}{\pi R_S^3}$	$\frac{C(R_c)}{\pi R_c^3}$	$r_2(DK)$	$r_3(DK)$	$r_3(DD)$	$\langle T \rangle$	$\langle V_{DK} \rangle$	$\langle V_{DD} \rangle$
$R_S = 0.5\text{fm } R_c = 1\text{fm}$							
0	-320.1	1.28	1.32	1.36	124.37	-189.61	-5.98
500	-455.4	1.39	1.44	1.47	99.51	-164.83	-5.03
1000	-562.6	1.46	1.53	1.54	91.43	-156.67	-4.51
3000	-838.7	1.61	1.69	1.68	93.24	-157.80	-3.82
$R_S = 0.5\text{fm } R_c = 2\text{fm}$							
0	-149.1	1.74	1.80	1.80	60.20	-125.74	-3.23
500	-178.4	1.91	1.98	1.96	51.00	-116.59	-2.64
1000	-195.0	1.99	2.07	2.04	50.63	-116.12	-2.43
3000	-225.9	2.13	2.22	2.15	53.61	-118.59	-2.24
$R_S = 0.5\text{fm } R_c = 3\text{fm}$							
0	-107.0	2.13	2.19	2.17	39.49	-105.35	-2.13
500	-119.4	2.31	2.38	2.34	34.80	-100.73	-1.77
1000	-125.6	2.37	2.47	2.42	34.90	-100.77	-1.65
3000	-136.2	2.53	2.61	2.53	36.66	-102.24	-1.54

Ds0\*(2317)



R(4140)



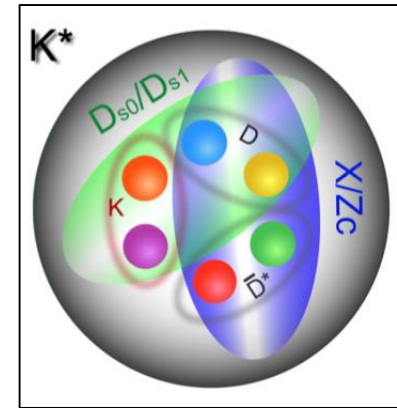
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# $K^*(4307)$

Instead of a  $D$ , what happens if we add a  $\bar{D}^*$  to the DK pair



Physics Letters B 785 (2018) 112–117



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$K^*$  mesons with hidden charm arising from  $KX(3872)$  and  $KZ_c(3900)$  dynamics



Xiu-Lei Ren<sup>a</sup>, Brenda B. Malabarba<sup>b</sup>, Li-Sheng Geng<sup>c,d</sup>, K.P. Khemchandani<sup>e,c</sup>,  
A. Martínez Torres<sup>b,c,\*</sup>

<sup>a</sup> Institut für Theoretische Physik II, Ruhr-Universität Bochum, D-44780 Bochum, Germany

<sup>b</sup> Instituto de Física, Universidade de São Paulo, C.P. 66318, 05389-970 São Paulo, São Paulo, Brazil

<sup>c</sup> School of Physics and Nuclear Energy Engineering & Beijing Key Laboratory of Advanced Nuclear Materials and Physics, Beihang University, Beijing 100191, China

<sup>d</sup> Beijing Advanced Innovation Center for Big Data-Based Precision Medicine, Beihang University, Beijing 100191, China

<sup>e</sup> Universidade Federal de São Paulo, C.P. 01302-907, São Paulo, Brazil

# $K^*(4307)$

- Fixed center approximation (FCA):

$$K(D\bar{D}^* + \bar{D}D^*) \sim KX(3872)/Zc(3900)$$

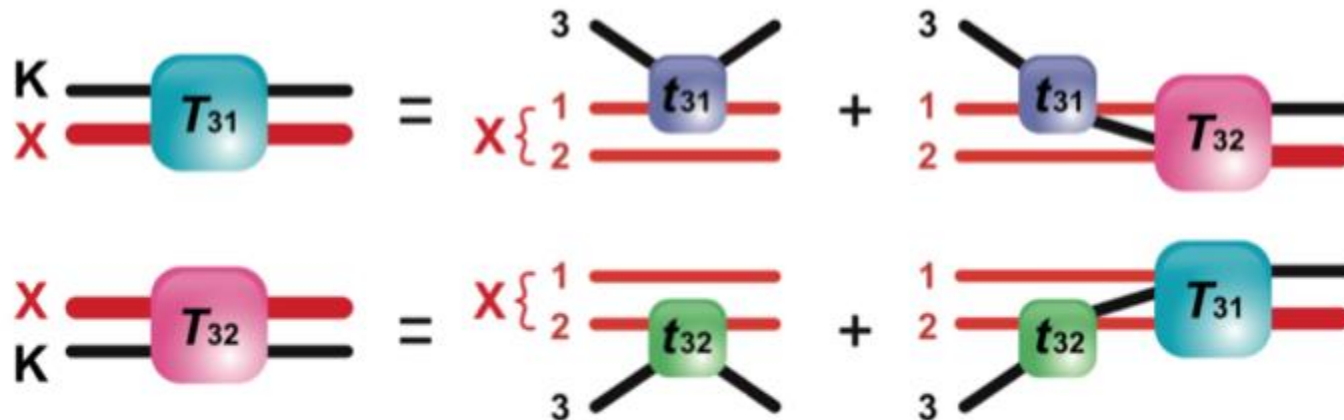
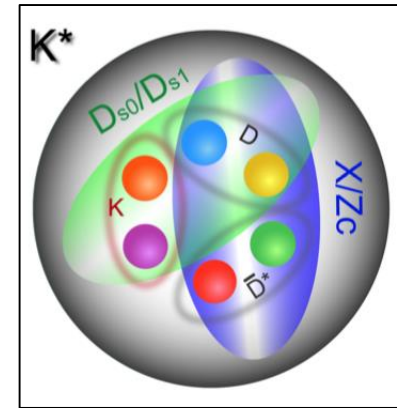
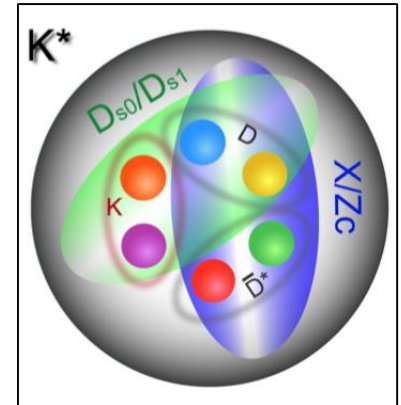
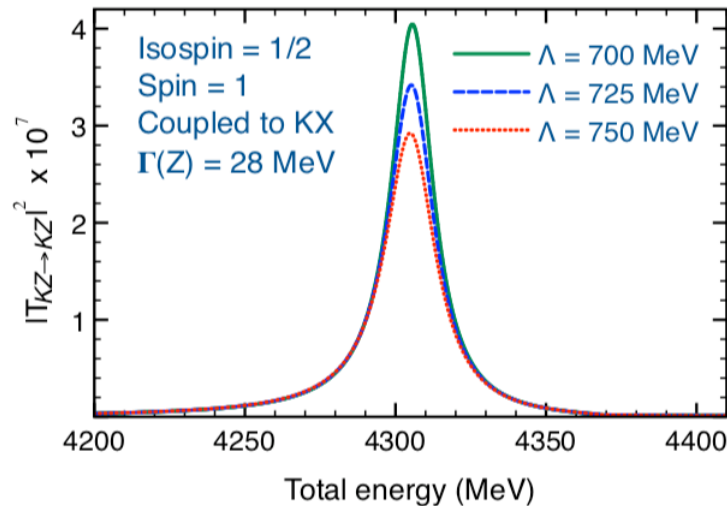
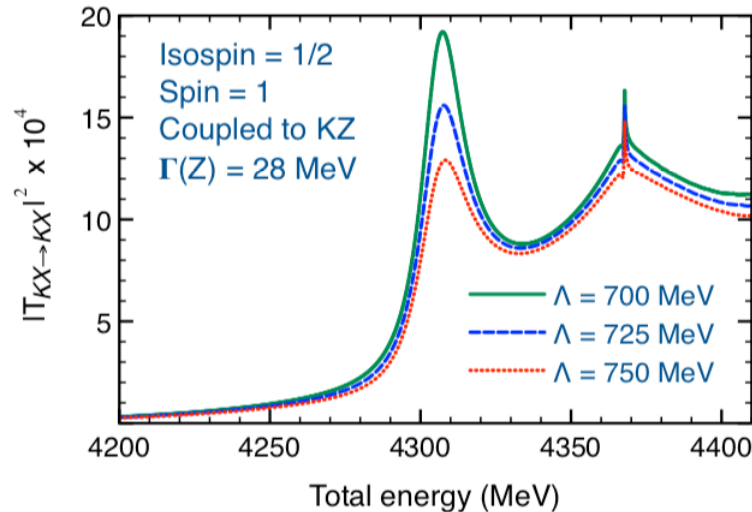


Figure 2: Diagrams showing the scattering of the particle labeled “3” ( $K$ ) on a cluster ( $X$ ) made of particles 1 ( $D$ ) and 2 ( $\bar{D}^*$ ).



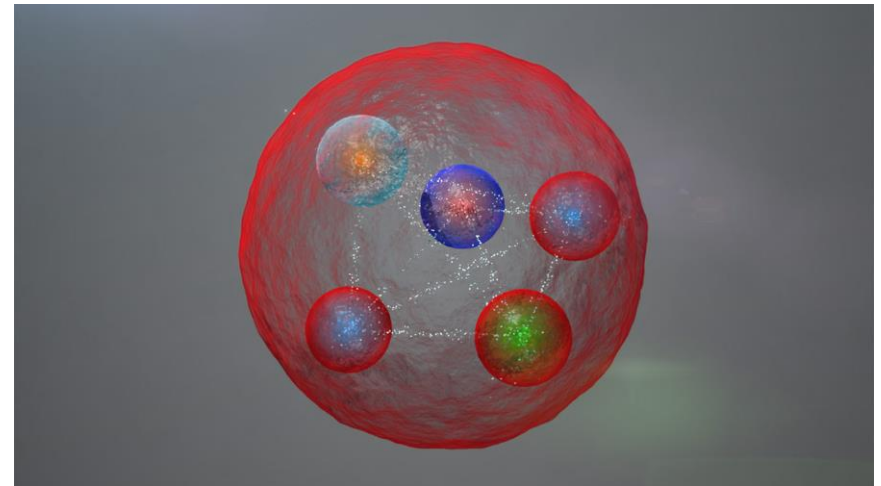
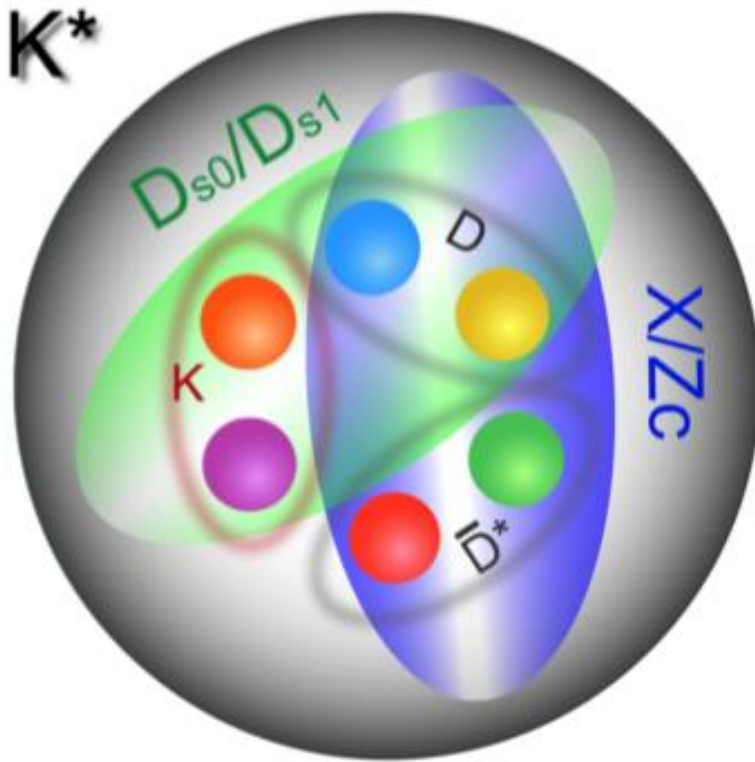
# $K^*(4307)$



- Treating KX and KZ as coupled channel systems
- A resonance with  $M=(4307 \pm 2) - i(9 \pm 2) \text{ MeV}$  with  $I(J^P) = 1/2(1^-)$

*In agreement with Li Ma, Qian Wang, Ulf-G. Meißner, 1711.06143, but with completely different dynamics*

# $K^*(4307)$ —**bosonic counterpart** of $P_c$



**Pentaquark ( $N^*$ ) by LHCb**

*Phys.Rev.Lett.* 115 (2015) 072001

Prediction of narrow  $N^*$  and  $\Lambda^*$  resonances with hidden charm above 4 GeV,  
Jia-Jun Wu, R. Molina, E. Oset, B.S. Zou, 1007.0573

# Analogy between $KD$ and $\bar{K}N$

$D_{s0}^*(2317)$

- DK bound state
- **Dynamically generated**--  
Unitary heavy hadron  
chiral perturbation theory

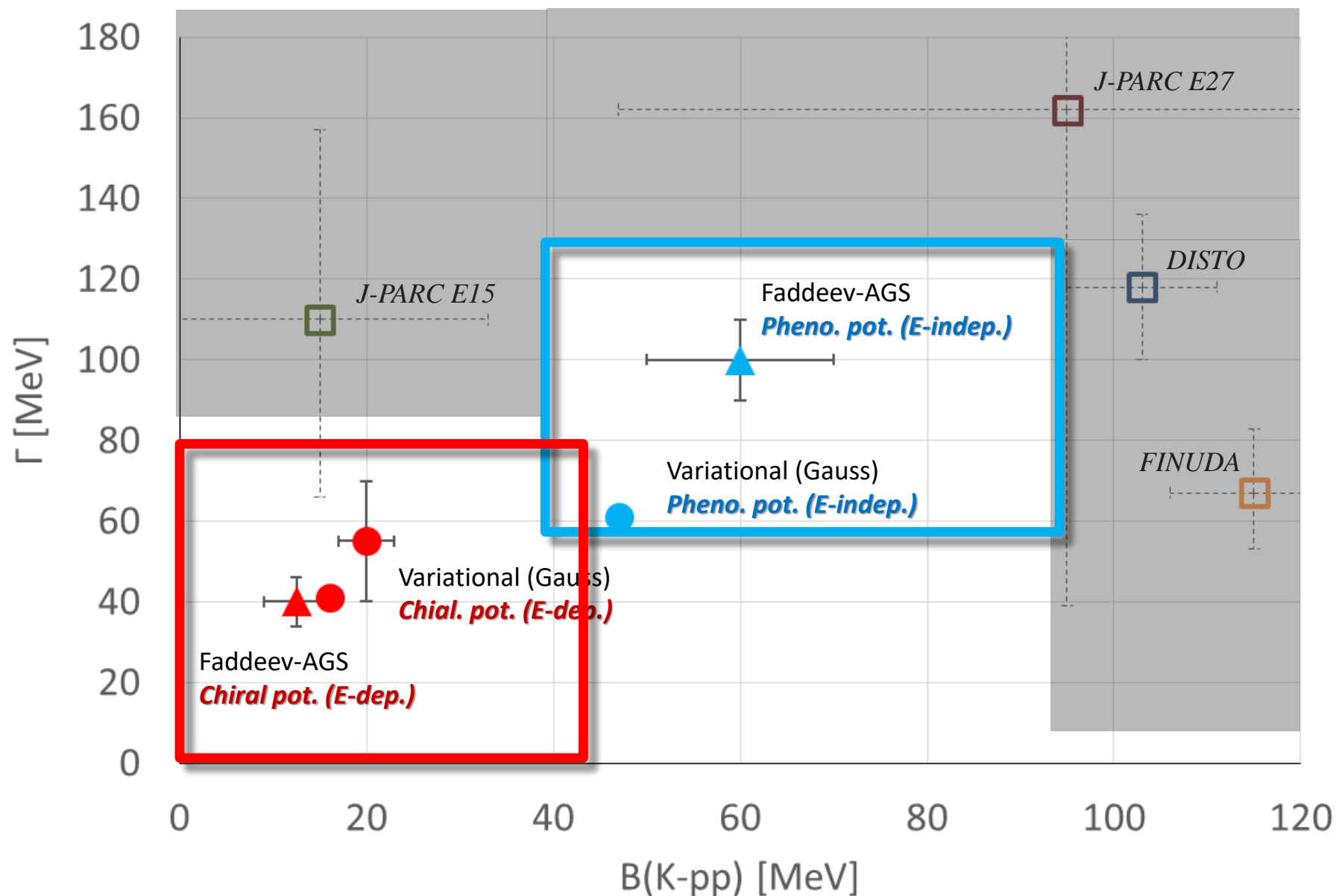
$\Lambda(1405)$

- N-Kbar bound state
- **Dynamically generated**--  
Unitary baryon chiral  
perturbation theory

The interaction between a **kaon** and a **heavy particle** seems to play an important role

# Current status on “ $K^-pp$ ”

A. Dote, Menu2019



# Experimental search by the Belle Collaboration

2008.13341

Search for a doubly-charged  $DDK$  bound state in  $\Upsilon(1S, 2S)$  inclusive decays and via direct production in  $e^+e^-$  collisions at  $\sqrt{s} = 10.520, 10.580$ , and  $10.867$  GeV

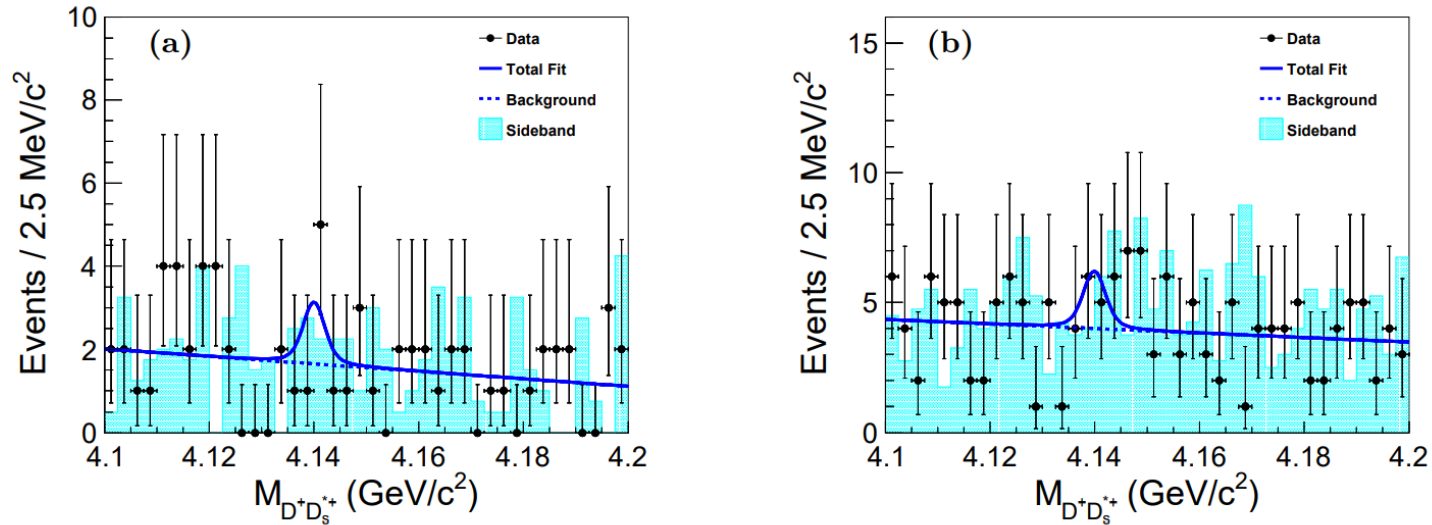


FIG. 4: The invariant-mass spectra of  $D^+D_s^{*+}$  in the (a)  $\Upsilon(1S)$  and (b)  $\Upsilon(2S)$  data samples. The cyan shaded histograms are from the normalized  $M_{D^+}$  and  $M_{D_s^{*+}}$  sideband events. The blue solid curves show the fitted results with the  $R^{++}$  mass fixed at  $4.14 \text{ GeV}/c^2$  and width fixed at  $2 \text{ MeV}$ , and the blue dashed curves are the fitted backgrounds.

# Experimental search by the Belle Collaboration

2008.13341

Search for a doubly-charged  $DDK$  bound state in  $\Upsilon(1S, 2S)$  inclusive decays and via direct production in  $e^+e^-$  collisions at  $\sqrt{s} = 10.520, 10.580, \text{ and } 10.867 \text{ GeV}$

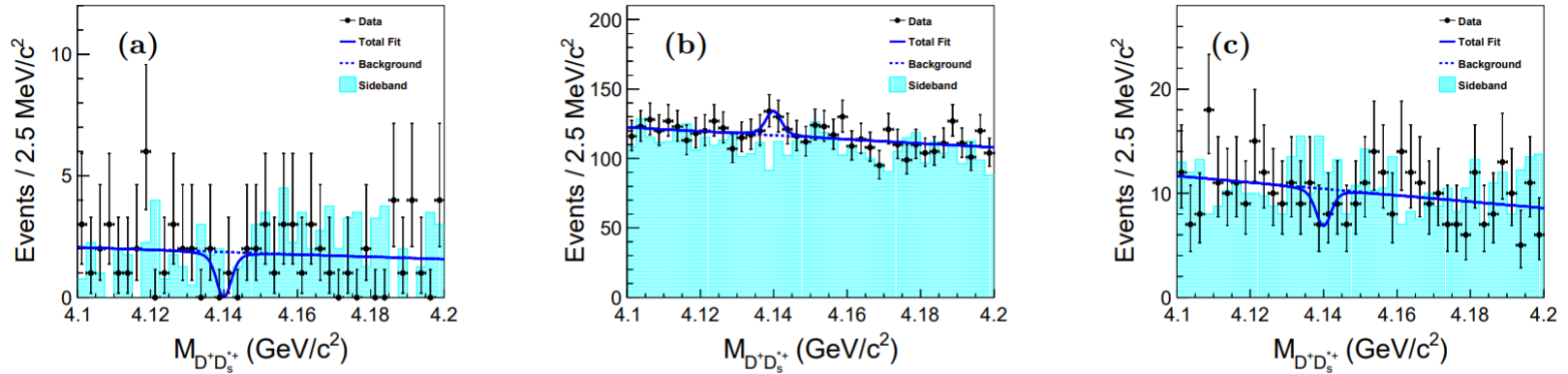


FIG. 10: The invariant-mass spectra of the  $D^+D_s^{*+}$  from  $e^+e^-$  annihilations at (a)  $\sqrt{s} = 10.520 \text{ GeV}$ , (b)  $\sqrt{s} = 10.580 \text{ GeV}$ , and (c)  $\sqrt{s} = 10.867 \text{ GeV}$  data samples. The cyan shaded histograms are from the normalized  $M_{D^+}$  and  $M_{D_s^{*+}}$  sideband events. The blue solid curves show the fitted results with the  $R^{++}$  mass fixed at  $4.14 \text{ GeV}/c^2$  and width fixed at  $2 \text{ MeV}$ , and the blue dashed curves are the fitted backgrounds.

# Looking into future

- ❑ We may have a better chance with Belle II, whose statistics will be larger by 50 times
- ❑ How about LHCb?



## Invited theory talk (vidyo pin: 2018)

📅 星期三 2018年10月10日 下午2:00 → 下午4:00 Europe/Zurich

📍 32/1-A24 (CERN)

视频会议室

 DK\_DDK\_DDDK\_molecules

Join

下午2:00 → 下午4:00 DK/DDK/DDDK molecules

🕒 2h

The  $Ds_0^*(2317)$  is widely accepted as a DK molecule. Its existence indicates that the DK interaction is attractive and strong enough to form a bound state. A natural question is then whether the system will still bind with one or more D mesons added. In a series of recent works, we explored such possibilities and showed that the DDK three-body molecular state exists, with a mass around 4140 MeV and a width of about 10 MeV. Due to the doubly charmed and doubly charged nature, such a state is apparently exotic. We have also performed a preliminary study of the DDDK system. In this talk, I will report on these studies and a few new binding mechanisms which have led to rather model independent predictions on the existence of hadronic molecules.

References:  
1309.4743, 1704.06123, 1705.00516, 1707.03802, 1805.08330, 1809.01059

报告人: Lisheng Geng (Beihang University)

 Geng-LHCb-1010.pdf

# Summary and outlook

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- Many interpretations exist for newly discovered exotic hadrons and are difficult to be distinguished from each other
- We proposed a **novel way** to **validate the molecule picture**.  
Taking the  $Ds_0^*(2317)$  as an example, if it **is indeed a molecule of DK**, then new forms of matter may be built upon them, similar to the build up of the periodic table.
- **We have performed explicit few-body studies**—demonstrating that indeed both  $DDK/DD\bar{b}^*K$  and  $DDDK$  states **bind**
- Now we need experimental or lattice QCD confirmations and further theoretical studies on their production and decay mechanisms



道生一，一生二，二生三

三生万物

**Thanks for your attention!**

# DDK system in finite volume

*Jin-Yi Pang, Jia-Jun Wu, and Li-Sheng Geng, 2008.13014*

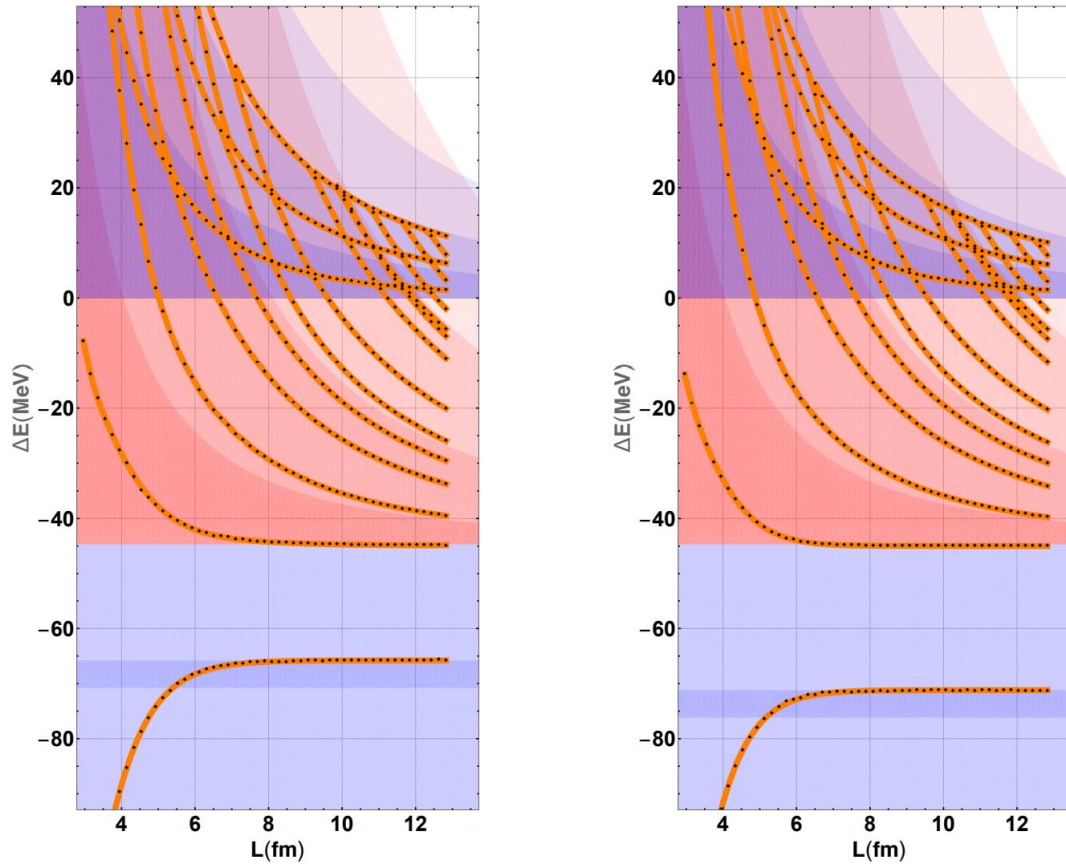


Figure 5: *DDK* states in finite volume. Left: only the *DK* interaction is considered. Right: both *DK* and *DD* interactions are taken into account. The upper blue regions indicate the case of 3 free particles in finite volume. The red regions indicate the case of free  $D_{s0}^*(2317)$  and  $D$ . The lower blue regions indicate the *DDK* bound state below the  $DD_{s0}^*(2317)$  threshold.