



郑州大学

# The Collins asymmetry in electroproduction of Kaon at the EICs within TMD factorization

薛世晨

第五届“强子谱和强子结构研讨会”2021.01.23-25

Collaborated with Xiaoyu Wang, De-Min Li(Zhengzhou U.) and Zhun Lu(Southeast U., China)

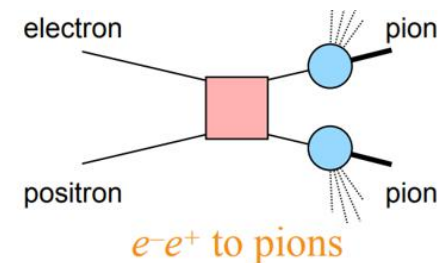
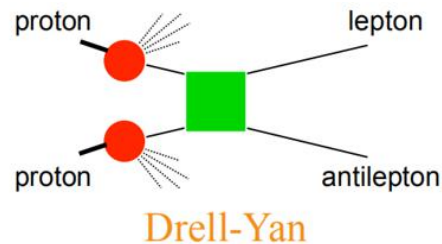
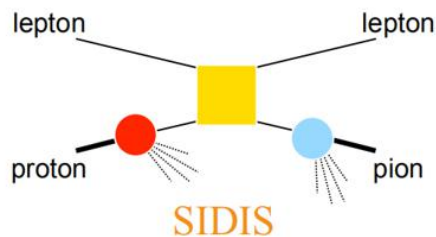
Based on S. C. Xue, X. Wang, D. M. Li and Z. Lu, arXiv: 2003.05679, accepted by EPJC

# OUTLINE



- Introduction
- Framework
- Numerical Estimate
- Conclusion

- Understanding the structure of the nucleon is one of the main tasks in hadronic physics.
- Based on QCD, the internal structure of a polarized nucleon explored in high energy process is described by TMD PDFs.
- These TMD PDFs can be extracted from the asymmetries arising in high energy scattering processes.



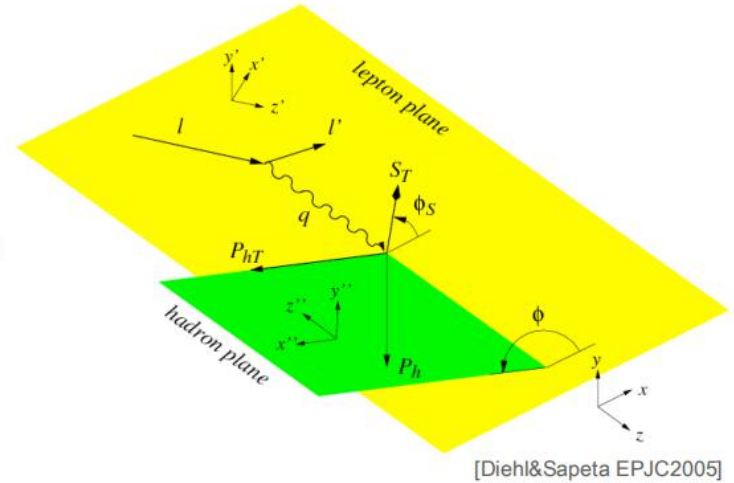
## Structure Functions

### SIDIS differential cross section

18 structure functions  $F(x, z, Q^2, P_T)$ ,  
model independent. (one photon exchange approximation)



Bacchetta et al, JHEP0702, 093 (2007)

$$\begin{aligned}
 & \frac{d\sigma}{dx dy dz dP_T^2 d\phi_h d\psi} \\
 &= \frac{\alpha^2}{xyQ^2} \frac{y^2}{2(1-\epsilon)} \left( 1 + \frac{\gamma^2}{2x} \right) \\
 & \times \left\{ F_{UU,T} + \epsilon F_{UU,L} + \sqrt{2\epsilon(1+\epsilon)} F_{UU}^{\cos\phi_h} \cos\phi_h + \epsilon F_{UU}^{\cos 2\phi_h} \cos 2\phi_h + \lambda_e \sqrt{2\epsilon(1-\epsilon)} F_{LU}^{\sin\phi_h} \sin\phi_h \right. \\
 & + S_L \left[ \sqrt{2\epsilon(1+\epsilon)} F_{UL}^{\sin\phi_h} \sin\phi_h + \epsilon F_{UL}^{\sin 2\phi_h} \sin 2\phi_h \right] + \lambda_e S_L \left[ \sqrt{1-\epsilon^2} F_{LL} + \sqrt{2\epsilon(1-\epsilon)} F_{LL}^{\cos\phi_h} \cos\phi_h \right] \\
 & + S_T \left[ (F_{UT,T}^{\sin(\phi_h-\phi_S)} + \epsilon F_{UT,L}^{\sin(\phi_h-\phi_S)}) \sin(\phi_h - \phi_S) + \epsilon F_{UT}^{\sin(\phi_h+\phi_S)} \sin(\phi_h + \phi_S) + \epsilon F_{UT}^{\sin(3\phi_h-\phi_S)} \sin(3\phi_h - \phi_S) \right. \\
 & \quad \left. + \sqrt{2\epsilon(1+\epsilon)} F_{UT}^{\sin\phi_S} \sin\phi_S + \sqrt{2\epsilon(1+\epsilon)} F_{UT}^{\sin(2\phi_h-\phi_S)} \sin(2\phi_h - \phi_S) \right] \\
 & + \lambda_e S_T \left[ \sqrt{1-\epsilon^2} F_{LT}^{\cos(\phi_h-\phi_S)} \cos(\phi_h - \phi_S) \right. \\
 & \quad \left. + \sqrt{2\epsilon(1-\epsilon)} F_{LT}^{\cos\phi_S} \cos\phi_S + \sqrt{2\epsilon(1-\epsilon)} F_{LT}^{\cos(2\phi_h-\phi_S)} \cos(2\phi_h - \phi_S) \right] \left. \right\}
 \end{aligned}$$



## TMD PDFs at leading-twist

		Quark polarization		
		Unpolarized (U)	Longitudinally polarized (L)	Transversely polarized (T)
Nucleon polarization	U	$f_1 = \text{Nucleon spin} \uparrow$		$h_1^\perp = \text{Quark spin} \downarrow - \text{Quark spin} \uparrow$ Boer-Mulder
	L		$g_1 = \text{Nucleon spin} \rightarrow - \text{Quark spin} \rightarrow$ Helicity	$h_{1L}^\perp = \text{Nucleon spin} \rightarrow - \text{Quark spin} \nearrow$
	T	$f_{1T}^\perp = \text{Nucleon spin} \uparrow - \text{Quark spin} \downarrow$ Sivers	$g_{1T}^\perp = \text{Nucleon spin} \rightarrow - \text{Quark spin} \rightarrow$	$h_{1T}^\perp = \text{Nucleon spin} \uparrow - \text{Quark spin} \downarrow$ Transversity $h_{1T}^\perp = \text{Nucleon spin} \uparrow - \text{Quark spin} \nearrow$

 Nucleon spin    
  Quark spin

Unpolarized distribution function, helicity, transversity: Survive at collinear limit  
 Boer-Mulders, Sivers: Time reversal odd

## Transversity distribution function

- **Fundamental distribution to encode the nucleon structure**
- **Chiral-odd**
- **Hard to access compared to helicity and unpolarized distribution function**

## How to access transversity

- **(TMD) factorization frame in SIDIS ——Collins function**  
J. C. Collins, NPB 396, 161 (1993)
- **Collinear factorization in SIDIS—twist-3 fragmentation function**  
X.Wang, Z. Lu PRD 93, 074009 (2016)
- **Collinear factorization in SIDIS—dihadron fragmentation function**  
A. Bacchetta et al., PRL107, 012001(2011)
- **Drell-Yan process—the antiquark transversity**  
V. Barone et al., Phys.Rept. 359 (2002) 1-168

Although progress has been made, sea quark transversity is almost unknown.

## Collins function

- **Correlation between fragmenting quark transverse spin and unpolarized hadron transverse momentum**
- **Analyzer of quark spin in nucleon**
- **Kaon Collins function has been extracted from the  $e^+e^-$  annihilation data from BaBar**

*M. Anselmino et al., PRD 93, 034025 (2016)*

The Collins asymmetry in Kaon production SIDIS process makes an ideal tool to access sea quark transversity due to its strange constituent.



What have we done?

**We predict the Collins asymmetry for charged Kaon production at EICs.**

- **apply the TMD factorization formalism**
- **Approach-TMD evolution formalism**
- **adopt the parametrization of the non-perturbative Sudakov form factor**
- **NLL accuracy**

# OUTLINE



- Introduction
- Framework
- Numerical Estimate
- Conclusion

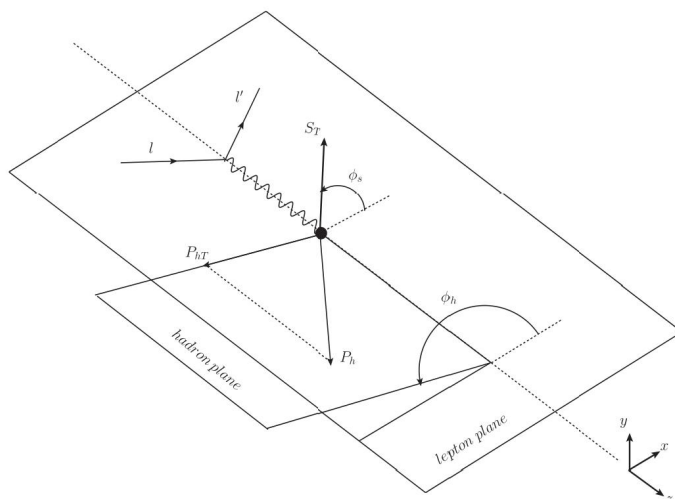
- We consider the SIDIS process

$$e(\ell) + p^\uparrow(P) \longrightarrow e(\ell') + K(P_h) + X(P_X)$$

- The invariants defined as

$$x_B = \frac{Q^2}{2P \cdot q}, \quad y = \frac{P \cdot q}{P \cdot \ell} = \frac{Q^2}{x_{BS}}, \quad z_h = \frac{P \cdot P_h}{P \cdot q},$$
$$Q^2 = -q^2, \quad s = (P + \ell)^2,$$

- The definition of azimuthal angles in SIDIS



- The 5-fold differential cross section with a transversely polarized target has the following general form

$$\frac{d^5\sigma(S_T)}{dx_B dy dz_h d^2P_{hT}} = \sigma_0(x_B, y, Q^2) \left[ F_{UU} + \sin(\phi_h + \phi_s) \frac{2(1-y)}{1+(1-y)^2} F_{UT}^{\sin(\phi_h + \phi_s)} + \dots \right],$$

- The Collins asymmetry can be written in terms of the structure functions

$$A_{UT}^{\sin(\phi_h + \phi_s)} = \frac{\sigma_0(x_B, y, Q^2)}{\sigma_0(x_B, y, Q^2)} \frac{2(1-y)}{1+(1-y)^2} \frac{F_{UT}^{\sin(\phi_h + \phi_s)}}{F_{UU}},$$

$$F_{UU}(Q; P_{hT}) = \mathcal{C}[f_1 D_1],$$

$$F_{UT}^{\sin(\phi_h + \phi_s)}(Q; P_{hT}) = \mathcal{C} \left[ \frac{-\hat{h} \cdot \mathbf{k}_T}{M_h} h_1 H_1^\perp \right].$$

$$\mathcal{C}[\omega f D] = x \sum_q e_q^2 \int d^2\mathbf{p}_T d^2\mathbf{k}_T \delta^{(2)}(\mathbf{p}_T - \mathbf{k}_T - \mathbf{P}_{h\perp}/z) \omega(\mathbf{p}_T, \mathbf{k}_T) f^q(x, p_T^2) D^q(z, k_T^2).$$

## ➤ Performing the Fourier Transformation

$$\begin{aligned}
 F_{UU}(Q; P_{hT}) &= \mathcal{C}[f_1 D_1] \\
 &= \sum_q e_q^2 \int d^2 \mathbf{p}_T d^2 \mathbf{k}_T \delta^{(2)}(\mathbf{p}_T - \mathbf{k}_T - \mathbf{P}_{hT}/z_h) f_1^q(x_B, p_T^2) D_1^q(z_h, k_T^2) \\
 &= \frac{1}{z_h^2} \sum_q e_q^2 \int d^2 \mathbf{p}_T d^2 \mathbf{K}_T \delta^{(2)}(\mathbf{p}_T + \mathbf{K}_T/z_h - \mathbf{P}_{hT}/z_h) f_1^q(x_B, p_T^2) D_1^q(z_h, K_T^2) \\
 &= \frac{1}{z_h^2} \sum_q e_q^2 \int d^2 \mathbf{p}_T d^2 \mathbf{K}_T \int \frac{d^2 b}{(2\pi)^2} e^{-i(\mathbf{p}_T + \mathbf{K}_T/z_h - \mathbf{P}_{hT}/z_h) \cdot \mathbf{b}} f_1^q(x_B, p_T^2) D_1^q(z_h, K_T^2) \\
 &= \frac{1}{z_h^2} \sum_q e_q^2 \int \frac{d^2 b}{(2\pi)^2} e^{i\mathbf{P}_{hT}/z_h \cdot \mathbf{b}} \tilde{f}_1^q(x_B, b) \tilde{D}_1^q(z_h, b).
 \end{aligned}$$

## ➤ introducing the TMDs definition

$$\begin{aligned}
 \int d^2 \mathbf{p}_T e^{-i\mathbf{p}_T \cdot \mathbf{b}} f_1^q(x_B, p_T^2) &= \tilde{f}_1^q(x_B, b) \\
 \int d^2 \mathbf{K}_T e^{-i\mathbf{K}_T/z_h \cdot \mathbf{b}} D_1^q(z_h, K_T^2) &= \tilde{D}_1^q(z_h, b)
 \end{aligned}$$

## ➤ The Collins structure function

$$\begin{aligned}
 F_{UT}^{\sin(\phi_h + \phi_s)}(Q; P_{hT}) &= C \left[ \frac{-\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} h_1 H_1^\perp \right] \\
 &= \sum_q e_q^2 \int d^2 \mathbf{p}_T d^2 \mathbf{k}_T \delta^{(2)}(\mathbf{p}_T - \mathbf{k}_T - \mathbf{P}_{hT}/z_h) \frac{-\hat{\mathbf{h}} \cdot \mathbf{k}_T}{M_h} h_1^q(x_B, p_T^2) H_1^{\perp, q}(z_h, k_T^2) \\
 &= \sum_q e_q^2 \frac{1}{z_h^2} \int d^2 \mathbf{p}_T d^2 \mathbf{K}_T \delta^{(2)}(\mathbf{p}_T + \mathbf{K}_T/z_h - \mathbf{P}_{hT}/z_h) \frac{\hat{\mathbf{h}} \cdot \mathbf{K}_T}{z_h M_h} h_1^q(x_B, p_T^2) H_1^{\perp, q}(z_h, K_T^2) \\
 &= \sum_q e_q^2 \frac{1}{z_h^2} \int d^2 \mathbf{p}_T d^2 \mathbf{K}_T \int \frac{d^2 b}{(2\pi)^2} e^{-i(\mathbf{p}_T + \mathbf{K}_T/z_h - \mathbf{P}_{hT}/z_h) \cdot \mathbf{b}} \frac{\hat{\mathbf{h}} \cdot \mathbf{K}_T}{z_h M_h} h_1^q(x_B, p_T^2) H_1^{\perp, q}(z_h, K_T^2) \\
 &= \sum_q e_q^2 \frac{1}{z_h^2} \frac{1}{z_h} \int \frac{d^2 b}{(2\pi)^2} e^{i\mathbf{P}_{hT}/z_h \cdot \mathbf{b}} \hat{h}_\alpha \tilde{h}_1^q(x_B, b) \tilde{H}_1^{\perp, \alpha, q}(z_h, b). \tag{10}
 \end{aligned}$$

## ➤ The transversity distribution function and Collins function in $\mathbf{b}$ space are defined as

$$\begin{aligned}
 \int d^2 \mathbf{p}_T e^{-i\mathbf{p}_T \cdot \mathbf{b}} h_1^q(x_B, p_T^2) &= \tilde{h}_1^q(x_B, b); \\
 \int d^2 \mathbf{K}_T e^{-i\mathbf{K}_T/z_h \cdot \mathbf{b}} \frac{K_T^\alpha}{M_h} H_1^{\perp, q}(z_h, K_T^2) &= \tilde{H}_1^{\perp, \alpha, q}(z_h, b).
 \end{aligned}$$

- **The evolution of the TMDs is usually performed in  $\mathbf{b}$  space.**
- **The solution of the energy dependence for TMDs has the general form as**

$$\tilde{F}(x_B, b; Q) = \mathcal{F} \times e^{-S} \times \tilde{F}(x_B, b; \mu)$$

$$\tilde{D}(z_h, b; Q) = \mathcal{D} \times e^{-S} \times \tilde{D}(z_h, b; \mu)$$

- **The Sudakov-like form factor can be separated into a perturbatively calculable part and a nonperturbative part**

$$S(Q; b) = S_{\text{pert}}(Q; b_*) + S_{\text{NP}}(Q; b)$$

$$S_{\text{pert}}(Q; b_*) = \int_{\mu_b^2}^{Q^2} \frac{d\bar{\mu}^2}{\bar{\mu}^2} \left[ A(\alpha_s(\bar{\mu})) \ln\left(\frac{Q^2}{\bar{\mu}^2}\right) + B(\alpha_s(\bar{\mu})) \right]$$

$$S_{\text{NP}}^{\text{pdf}}(Q; b) = \frac{g_2}{2} \ln\left(\frac{Q}{Q_0}\right) b^2 + g_1^{\text{pdf}} b^2,$$

$$S_{\text{NP}}^{\text{ff}}(Q; b) = \frac{g_2}{2} \ln\left(\frac{Q}{Q_0}\right) b^2 + g_1^{\text{ff}} b^2.$$

## ➤ Unpolarized structure function

$$\begin{aligned}
 F_{UU}(Q; P_{hT}) &= \frac{1}{z_h^2} \int_0^\infty \frac{db b}{(2\pi)} J_0(P_{hT}/z_h b) \tilde{F}_{UU}(Q; b) \\
 &= \frac{1}{z_h^2} \sum_q e_q^2 \int_0^\infty \frac{db b}{(2\pi)} J_0(P_{hT}/z_h b) e^{-S_{\text{pert}}(Q; b_*) - S_{\text{NP}}^{\text{SIDIS}}(Q; b)} \\
 &\quad \left( \sum_i C_{q \leftarrow i}^{(\text{SIDIS})} \otimes f_1^{i/p}(x_B, \mu_b) \right) \times \left( \sum_j \hat{C}_{j \leftarrow q}^{(\text{SIDIS})} \otimes D_1^{K/j}(z_h, \mu_b) \right).
 \end{aligned}$$

## ➤ Collins structure function

$$\begin{aligned}
 F_{UT}(Q; P_{hT}) &= \frac{-1}{2z_h^3} \sum_q e_q^2 \int_0^\infty \frac{db b^2}{(2\pi)} J_1(P_{hT}/z_h b) e^{-S_{\text{pert}}(Q; b_*) - S_{\text{NP}}^{\text{SIDIS Collins}}(Q; b)} \\
 &\quad \left( \sum_i \delta C_{q \leftarrow i}^{(\text{SIDIS})} \otimes h_1^{i/p}(x_B, \mu_b) \right) \left( \sum_j \delta \hat{C}_{j \leftarrow q}^{(\text{SIDIS})} \otimes \hat{H}_{K/j}^{(3)}(z_h, \mu_b) \right).
 \end{aligned}$$



# OUTLINE



- Introduction
- Framework
- Numerical Estimate
- Conclusion

## ➤ Parametrization for the collinear transversity distribution function

**Valence quark** 
$$h_1^q(x, Q_0) = N_q^h x^{a_q} (1-x)^{b_q} \frac{(a_q + b_q)^{a_q + b_q}}{a_q^{a_q} b_q^{b_q}} \frac{1}{2} (f_1^q(x, Q_0) + g_1^q(x, Q_0))$$

Z. B. Kang et al, PRD 93, 014009 (2016)

**Sea quark** 
$$h_1^q(x, Q_0) = N_s \frac{1}{2} (f_1^q(x, Q_0) + g_1^q(x, Q_0))$$

## ➤ Parametrization for Collins function

$$\Delta^N D_{h/q^\uparrow}(z_h, p_\perp) = \tilde{\Delta}^N D_{h/q^\uparrow}(z_h) h(p_\perp) \frac{e^{-p_\perp^2 / \langle p_\perp^2 \rangle}}{\pi \langle p_\perp^2 \rangle} \quad \text{M. Anselmino et al, PRD 93, 034025 (2016)}$$

$$\hat{H}_{h/j}^{(3)}(z_h) = \frac{\sqrt{2}e}{M_C} \mathcal{N}_q^C(z_h) D_{h/q}(z_h) \left( \frac{M_C^2}{M_C^2 + \langle p_\perp^2 \rangle} \right)^2 \langle p_\perp^2 \rangle$$

## ➤ Kinematical region for EIC

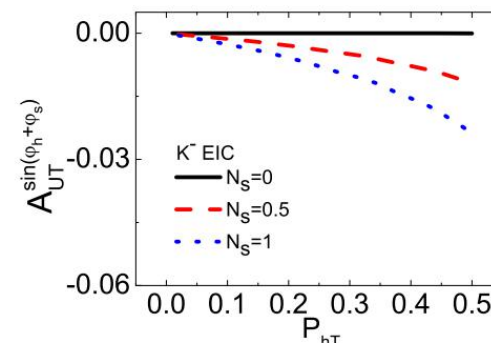
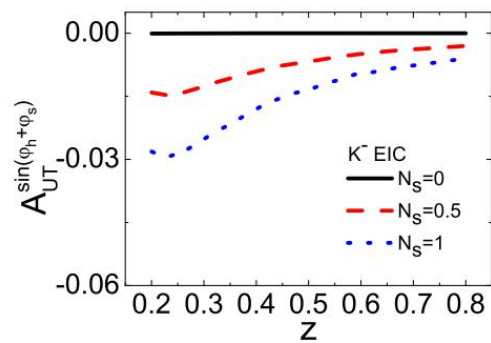
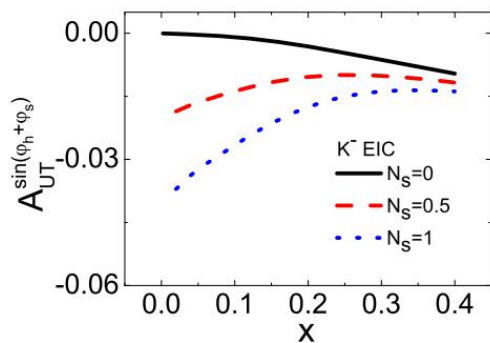
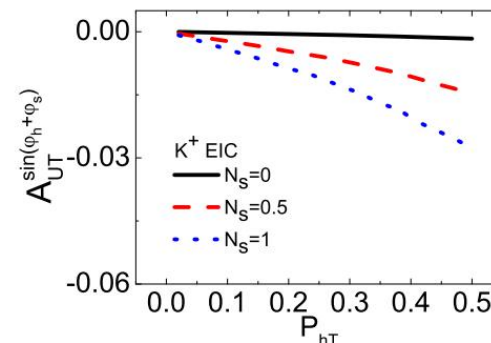
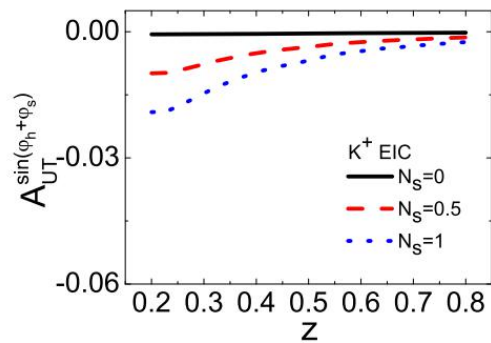
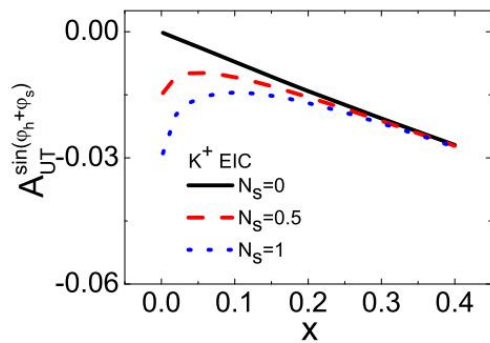
$$0.001 < x < 0.4, \quad 0.07 < y < 0.9, \quad 0.2 < z < 0.8,$$
$$1 \text{ GeV}^2 < Q^2, \quad W > 5 \text{ GeV}, \quad \sqrt{s} = 100 \text{ GeV}, \quad P_{hT} < 0.5 \text{ GeV}.$$

## ➤ For EicC

$$0.005 < x < 0.5, \quad 0.07 < y < 0.9, \quad 0.2 < z < 0.7,$$
$$1 \text{ GeV}^2 < Q^2 < 200 \text{ GeV}^2, \quad W > 2 \text{ GeV}, \quad \sqrt{s} = 16.7 \text{ GeV}, \quad P_{hT} < 0.5 \text{ GeV}$$

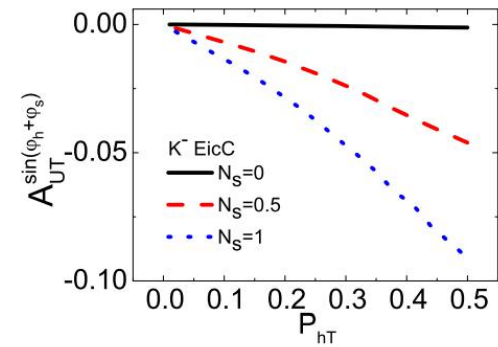
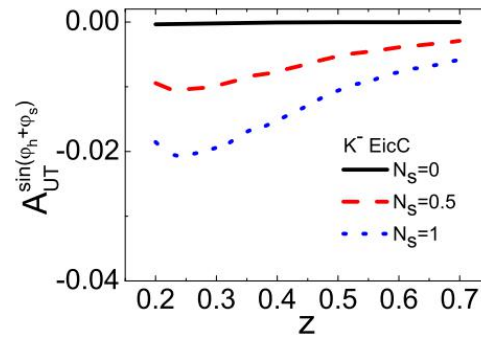
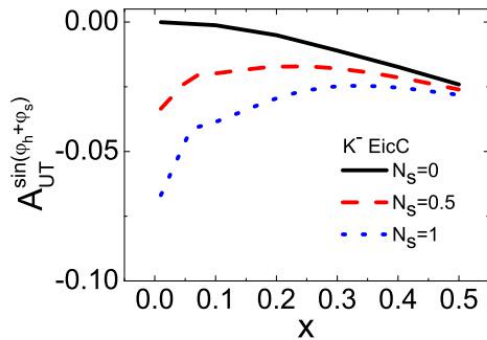
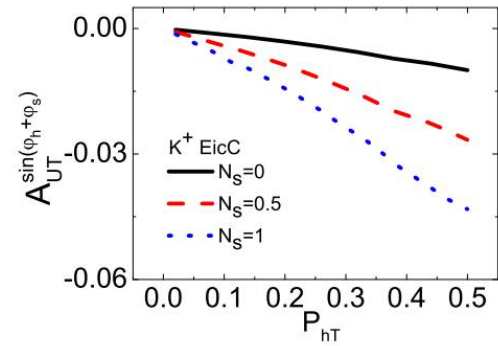
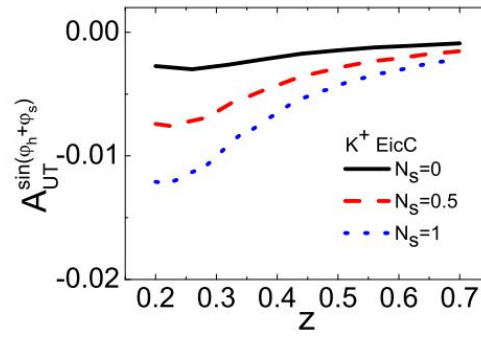
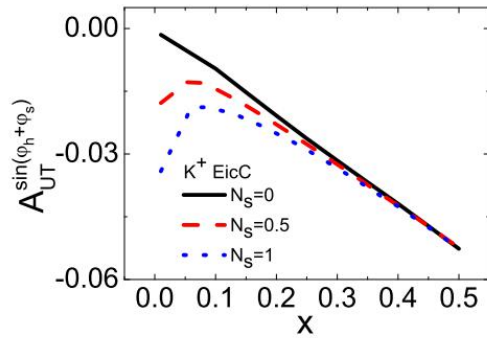
# Numerical Estimate

## ➤ For EIC



# Numerical Estimate

## ➤ For EicC



- **Both sizable at EIC and EicC**
- **The magnitude of the asymmetry increases with increasing  $N_s$  of the collinear sea quark transversity function.**
- **The effect of the transversity of the sea quarks turns out to be smaller in the  $K^+$  production process than that in the  $K^-$  process.**
- **For the asymmetry as the function of  $x$ , there is a clear peak at  $x \approx 0.05$  at EicC when considering the non-zero sea quark transversity, while the peak vanishes with zero sea contribution of transversity. Although the peak turns to be vague at EIC, the tendency still remains.**

- Introduction
- Framework
- Numerical Estimate
- Conclusion

- **The measurement on the Collins asymmetry of semi-inclusive Kaon production at future electron ion colliders can provide useful constraints on the sea quark transversity.**
- **We note that there are large errors in the extraction of the Kaon Collins function, which indicates the importance of more precision  $e^+e^-$  data in order to constrain the Kaon Collins function.**





Thank you!