Detecting pure triangle singularity effect through $\psi(2S) \rightarrow \overline{p}p\eta$ process

Qi Huang arXiV: 2011.14590 Collaborators: Chao-Wei Shen, Jia-Jun Wu UCAS

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Outline

- Motivation
- Results
- Summary

Triangle Singularity (TS): L. D. Landau, Nucl. Phys. 13, 181-192 (1959) ⇒ Concept.



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BESIII, 2012: Isospin breaking process $\eta(1405) \rightarrow \pi^0 f_0(980)$.

Nature of $a_1(1420)$: Triangle singularity caused by $K^*\overline{K}K$ loop.

Lots of theoretical works.

A recent review: F. K. Guo, et. al., Prog. Part. Nucl. Phys. 112, 103757 (2020)

Still not confirmed experimentally !

Meaning:

- 1. Understand triangle singularity itself.
- 2. Confirm hadron loop mechanism.
- 3. Help studying properties of hadrons.





1. Threshold

Example:

 $Z_c(3900) - \overline{D}D^*$ threshold ~ 25 MeV $P_c \sim 4.45 \ GeV - \chi_{c1}p$ threshold ~ O(10) MeV X.-H. Liu, M. Oka and Q. Zhao, Phys. Lett. B753, 297 (2016):

If a threshold enhancement falls into the TS kinematic region, distinguishing it from TS will be complicated.

2. Width of internal particles

Example:

 $\eta(1405) \rightarrow \pi^0 f_0(980) \rightarrow 3\pi$

3. Information of interaction vertex

Example:

 $P_c(4450)$: $\Lambda_b \to J/\psi pK$ via $\chi_{c1}p\Lambda^*$ loop. \checkmark No experimental data to constraint the $\Lambda_b \chi_{c1}\Lambda^*$ vertex \Rightarrow Line shape only.

N. N. Achasov, et.al., Phys. Rev. D92, 036003 (2015):





 $m_{TS} @ m_{p\eta}$: 1.56387 GeV

1. Far away from the threshold of the relative channel.

Position of TS is 80 MeV away from $p\eta$ threshold. TS and threshold can be separated clearly.

2. Narrow widths of all internal particles.

Peak caused by triangle singularity must be very sharp. Distinguishing it from N^* states is very easy.

3. The well known three vertices in the triangle loop.

All three vertices can be constrained by experimental data. Strength of TS can be predicted precisely.

TS in $\psi(2S) \rightarrow p\bar{p}\eta$ process



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Main decay mechanism:



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Main decay mechanism:



Loop diagram:

 $p\eta \rightarrow p\eta$ vertex: s-channel contribution: $p\eta \rightarrow N^* \rightarrow p\eta$ N^* near m_{TS} : N(1440), N(1520), N(1535), N(1650). Considerable $B(N^* \rightarrow p\eta)$: N(1535), N(1650). $\Rightarrow N^* = N(1535), N(1650)$.

$$\mathcal{F}(q,m,\Lambda)=rac{\Lambda^4}{(q^2-m^2)^2+\Lambda^4}$$
 , $\Lambda=m+lpha\,\Lambda_{QCD}$.

Tree diagram: Dominant! BESIII: PRD 88, 032010(2013)



Purpose:

To check if the narrow peak caused by the triangle loop diagram is visible after mixing it with the tree diagram.

TS in $\psi(2S) \rightarrow p\bar{p}\eta$ process

Loop diagram only



- 1. The peak caused by the triangle singularity is very sharp, its width is only about 1 MeV.
- 2. The contribution of N(1650) is negligible.
- 3. The peak at the triangle singularity is little dependent on the α that appears in the expression of form factor.

Stable!

TS in $\psi(2S) \rightarrow p\bar{p}\eta$ process

Interference between tree and loop diagrams



Loop diagram: $\alpha = 1, N(1535) + N(1650)$.

- 1. There exists a visible enhancement at the right shoulder of N(1535).
- 2. The width of the peak caused by the triangle singularity is enlarged to 5 MeV.
- 3. The enhancement of the peak comparing to the tree diagram is about 10%.

TS in $\psi(2S) \rightarrow p\bar{p}\eta$ process

Effect of the tree diagram $\psi(2S) \rightarrow \eta(J/\psi \rightarrow p\bar{p})$



Schmid theorem can't be applied directly here:

1. $p\eta \rightarrow p\eta$ is not purely elastic process, $p\pi$ channel must couples with it, which is effectively included in the imaginary part of N^* propagator.

2. $\psi(2S) \rightarrow \bar{p}(N^* \rightarrow p\eta)$ will modify the amplitude to $|t_{J/\psi}^{\text{Tree}} + t_{N^*}^{\text{Loop}}|^2$, where $t_{\text{elastic}}^{\text{Loop}} \neq t_{\text{elastic}}^{\text{Loop}}$.

3. Contribution of $\psi(2S) \rightarrow \eta(J/\psi \rightarrow p\overline{p})$ can be ----removed by applying a cut $m_{p\overline{p}} < m_{J/\psi}$.



TS in $\psi(2S) \rightarrow p\bar{p}\eta$ process

Interference between tree and loop diagrams



Loop diagram: $\alpha = 1, N(1535) + N(1650)$.

Black solid: No $\psi(2S) \rightarrow \eta(J/\psi \rightarrow p\bar{p})$, no cut. Red dashed: $\psi(2S) \rightarrow \eta(J/\psi \rightarrow p\bar{p})$, $m_{p\bar{p}} < 3.067$ GeV.

- 1. The contribution of $\psi(2S) \rightarrow \eta(J/\psi \rightarrow p\bar{p})$ can be neglected after applying the $m_{p\bar{p}}$ cut, thus there will be no effect of Schmid theorem.
- 2. The line changes little, our previous conclusions on the triangle singularity that appears in the $m_{p\eta}$ of $\psi(2S) \rightarrow p\bar{p}\eta$ process are still valid.

TS in $\psi(2S) \rightarrow p\bar{p}\eta$ process

Discussion with experimentalists:

- 1. High statistics: 4 billion $\psi(2S) \Rightarrow$ about 120 events.
- 2. High resolution: The resolution of the detector should be less than 2-3 MeV (BESIII: ~4.3 MeV).

An unavoidable weakness:

Phase angle between tree and loop diagrams: In our calculation, the phase angle is set to 0.

Only assumption but as a weak point with no solution: Enhancement might be invisible at some specific values.

We still hope the future experiments such as STCF could do precise analysis on this process.

Improvement: Selection of internal particles

A lesson: Widths of internal particles should not be too narrow !

Or:

1. Peak of TS is too sharp to be detected by experiment. \implies Requirement of the resolution is too high.





- We predict a detectable pure TS effect in $\psi(2S) \rightarrow p\overline{p}\eta$ process.
- A precise calculation on the strength of this TS is done.
- A visible enhancement shows up on the $m_{p\eta}$ invariant mass spectrum.
- We suggest experiments do a precise analysis on $\psi(2S) \rightarrow \overline{p}p\eta$ process.

Thank you ~