

# Measurement and Monte Carlo

or "how to make a *useful* measurement"

## Part II

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\*Thanks to Jon Butterworth and Emily Nurse for the use of their slides



Monte Carlo net

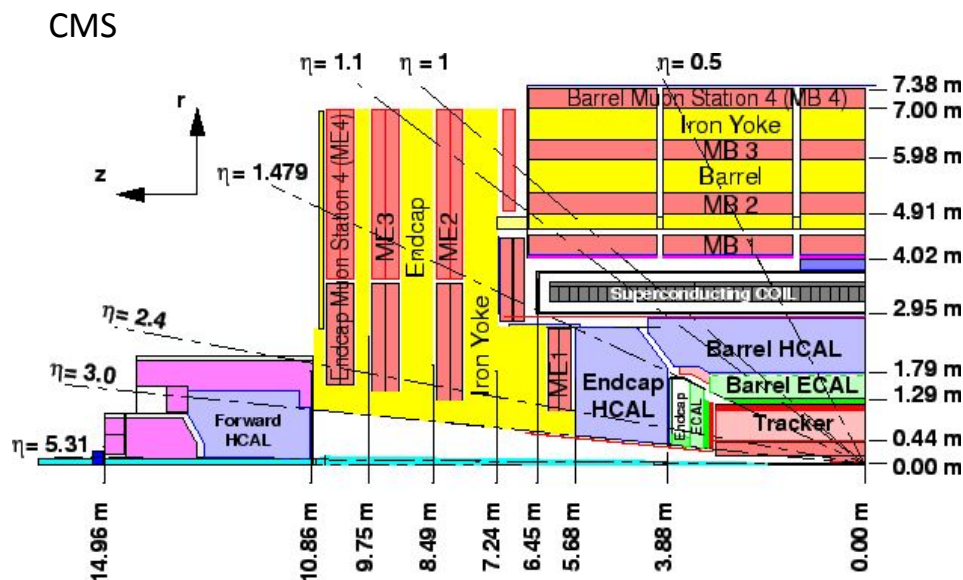
$$\frac{1}{\sigma_0} \frac{d\sigma}{dx_1 dx_2} = \frac{\alpha_s C_F}{2\pi e} \left\{ \frac{x_1^2 + x_2^2}{(1-x_1)(1-x_2)} - \frac{2m_s}{s} \left( \frac{1}{(1-x_1)^2} + \frac{1}{(1-x_2)^2} \right) \right\}$$

$$\int_{x_{\min}}^{x_{\max}} f(x) dx = R \int_{x_{\min}}^{x_{\max}} f(x) dx$$
$$x_{\min}) + R(F(x_{\max}) - F(x_{\min}))$$

# Fiducial phase-space

Regardless of detector efficiencies and resolutions, there are uninstrumented kinematic regions that we don't measure *at all*:

**We do not have  $4\pi$  detectors, and can't reconstruct particles down to zero  $p_T$ !**



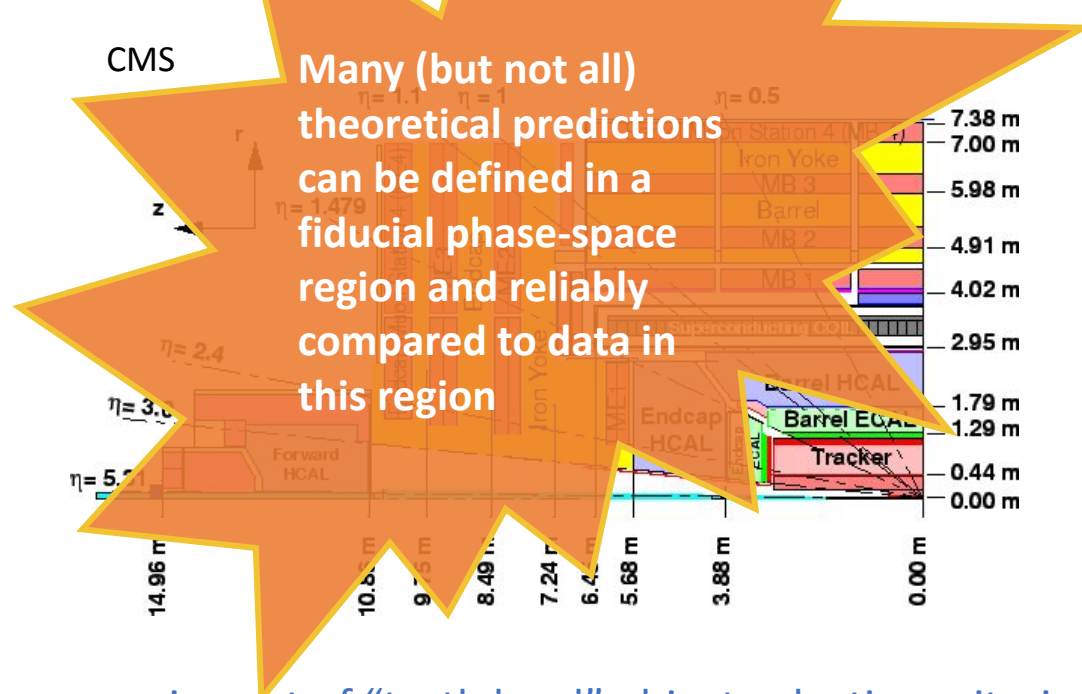
A fiducial phase-space is a set of “truth level” object-selection criteria that align well with the sensitivity of the real detector + reco:

e.g.: Select events with one (and only one) muon with  $p_T > 25$  GeV,  $|\eta| < 2.4$  and  $p_T^{\text{miss}} > 30$  GeV.

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# Fiducial phase-space

Regardless of detector efficiencies and ... are uniformly instrumented kinematic regions that would not be measured

We do not have  $4\pi$  detector

If you have a theory prediction that cannot be calculated in a fiducial phase space (e.g. using resummation techniques) then provide a separate acceptance factor, but *also* publish the fiducial result!

A fiducial phase-space is a set of “truth level” object-selection criteria that align well with the sensitivity of the real detector + reco:

e.g.: Select events with one (and only one) muon with  $p_T > 25$  GeV,  $|\eta| < 2.4$  and  $p_T^{\text{miss}} > 30$  GeV.



# What is a final-state particle?

Remember...

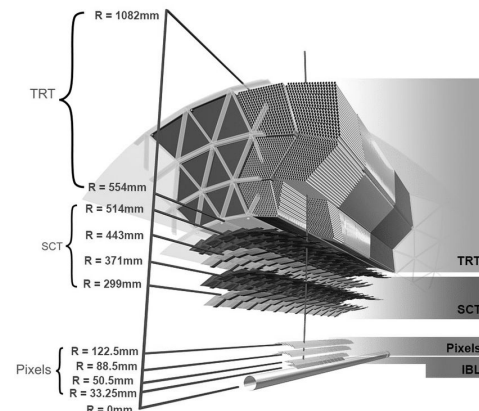
We only “see” *stable final-state* particles :

- electrons: stable
- muons: stable (  $\tau_0 = 2.2 \mu\text{s}$ , mean decay after  $\sim 1 \text{ km}$  at 20 GeV!!)
- taus: unstable (  $\tau_0 = 0.3 \text{ ps}$ , mean decay after  $\sim 1 \text{ mm}$  at 20 GeV!!)
- neutrinos: stable (but invisible)
- photons: stable
- hadrons: unstable
  - more leptons (charged and neutral), photons, more hadrons...
- quarks and gluons: unstable — are they even real?
  - hadrons → jets
- W, Z, Higgs, top: unstable, varying degrees of objective reality!
  - everything!

Common choice :  $\tau_0 > 30 \text{ ps}$ , after hadronization.

Lots of useful advice and discussion here:

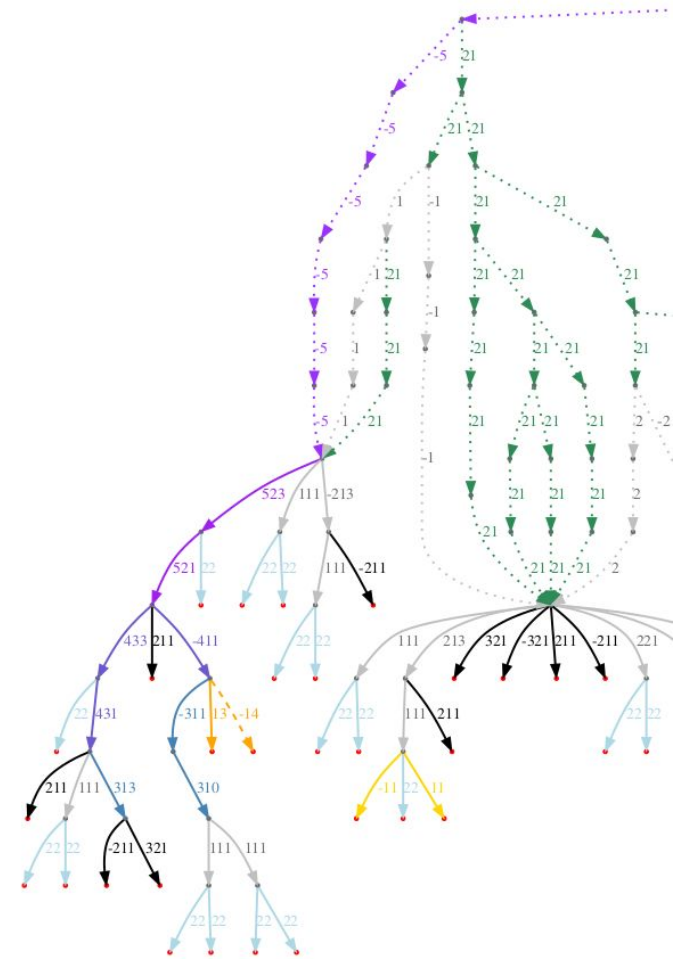
<https://cds.cern.ch/record/2022743>



# What is promptness / directness?

Not all final-state particles are equally significant:

- Hadrons are ten-a-penny! Even c and b hadrons can be produced fairly easily from semi-soft  $g \rightarrow b\bar{b}$  splittings
- Similar for photons: a high-energy photon direct from the hard process is a significant EM-interaction event, but *lots* of low-energy ones from  $\pi^0 \rightarrow \gamma\gamma$
- It's harder to make electrons and muons come out of hadron collisions! Implies an electroweak process internally  $\Rightarrow$  weak coupling. Can still happen in semi-leptonic (esp. heavy) hadron decays, with neutrinos.
- Leptonic taus are indistinguishable from  $e/\mu + \nu$
- Define *directness* “backwards”: recursively eliminate anything from a hadron. Only leptons and photons can be unambiguously direct



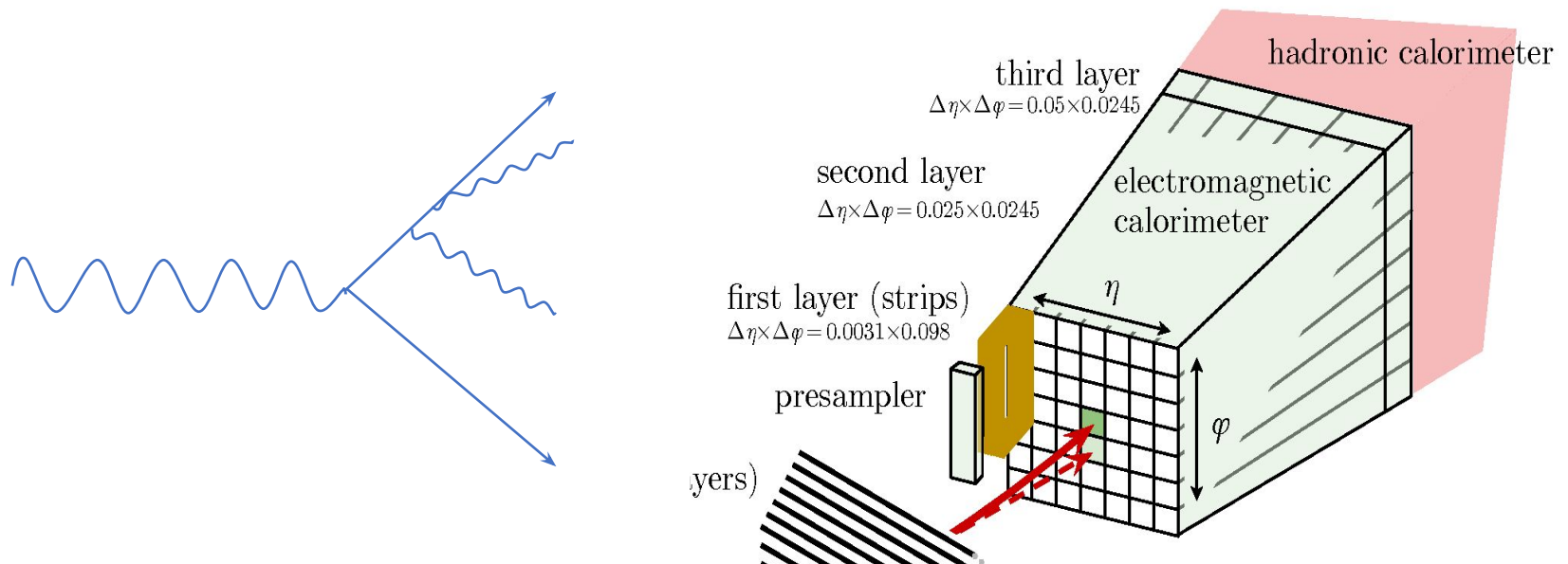
Often see “prompt” used to mean direct... but misleading since not a timing or displacement thing, unlike b-physics usage

# What is a final-state electron/muon?

- Electrons/muons from hadron decays are typically removed in the data analysis by isolation cuts / fake removal:
- ✓ Can define **direct leptons** via the **not-from-hadron-decays rule** and only consider these: more robust and model-independent than asking that the lepton comes from a certain propagator in the hard process
- ✓ **Well-defined in Rivet**, but you may need to also implement it in your experiment's software
- We don't usually define particle-level isolation, but rather correct for inefficiencies of these requirements
  - It might be worth reconsidering this in specific analyses where proximity to jets has a large effect on results

# What is a final-state electron/muon?

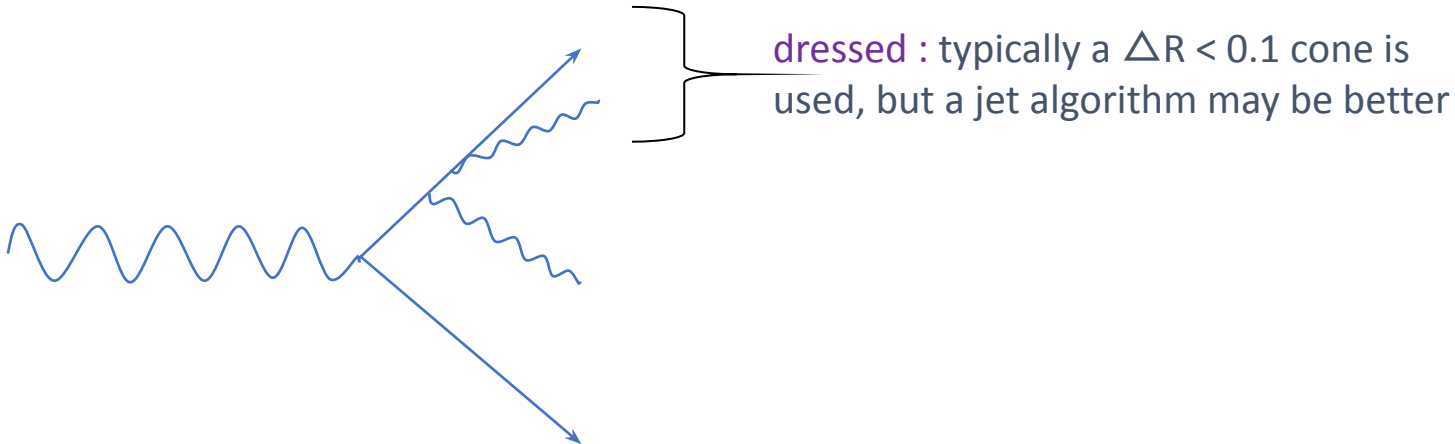
- Electrons and muons emit FSR photon radiation
  - (and lots of it, especially in the collinear limit, especially for electrons)
- For muons we measure the charged particle track, photon energy is not included
- For electrons we cluster calorimeter cells, and most collinear radiation will be included in the energy measurement





# What is a final-state electron/muon?

- Electrons and muons emit FSR photon radiation (and lots of it, especially in the collinear limit, especially for electrons)
- We can define lepton momenta as:
  - **Born leptons** – as if FSR never happened  $\Rightarrow$  not what we measure
  - **Bare leptons** – after all FSR  $\Rightarrow$  closest to muon measurement
  - **Dressed leptons** – with the momenta of close-by photons “clustered” into the lepton momenta  $\Rightarrow$  closest to electron measurement

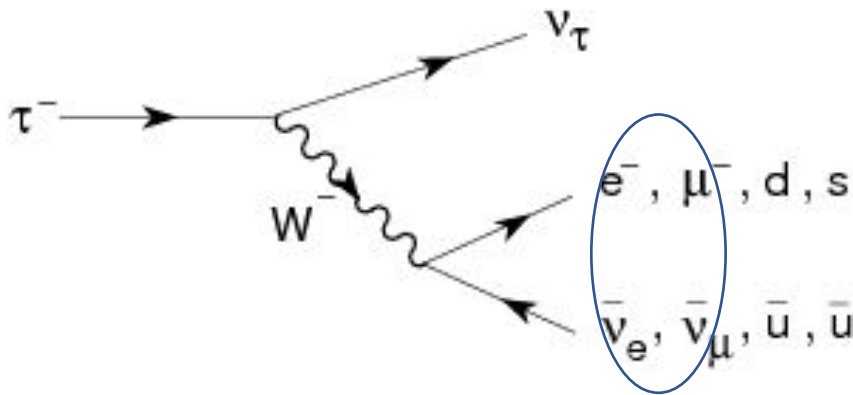


# What is a final-state electron/muon?

- Electron and muon final states can be very different for bare leptons, but much closer for Born and dressed leptons: [see the Rivet tutorial](#)
- It is often argued that dressed should be used for both to allow for *easy combination of final states*. Also bare versus dressed is much closer for muons than bare versus dressed for electrons
- Similarly, fiducial phase space cuts often harmonized for the two, requiring a small extrapolation in phase space for one
- But electrons  $\neq$  muons
- We may want to retain sensitivity to differences
  - Especially cf. lepton (non-)universality anomalies!
- It may be better to measure both, publish correlations between uncertainties, and make choices that are best for each individual channel...

# What is a final-state tau?

Recall: unstable ( $c\tau_0=0.1$  mm, decays after  $\sim 1$  mm at 20 GeV)

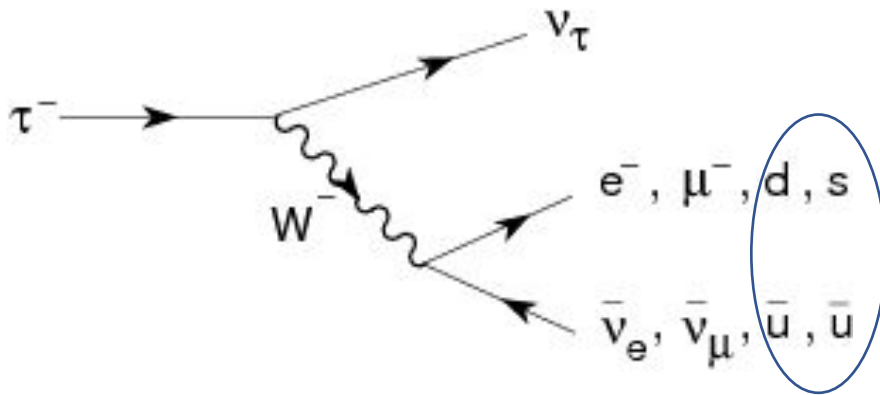


## Leptonic decays

- The final state particles are electrons/muons and neutrinos
- Define fiducial phase-space with those (but we need to be careful to check lepton efficiencies as, e.g. impact-parameter cuts can be less efficient for leptons from taus)

# What is a final-state tau?

Recall: unstable ( $c\tau_0=0.1$  mm, decays after  $\sim 1$  mm at 20 GeV)



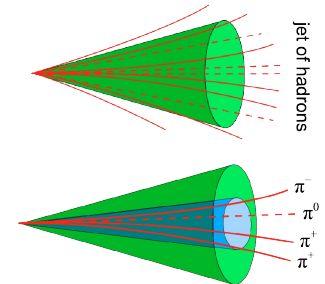
$$\tau^\pm \rightarrow \pi^\pm \pi^0 \nu$$

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$$\tau^\pm \rightarrow \pi^\pm \pi^\pm \pi^\mp \nu$$

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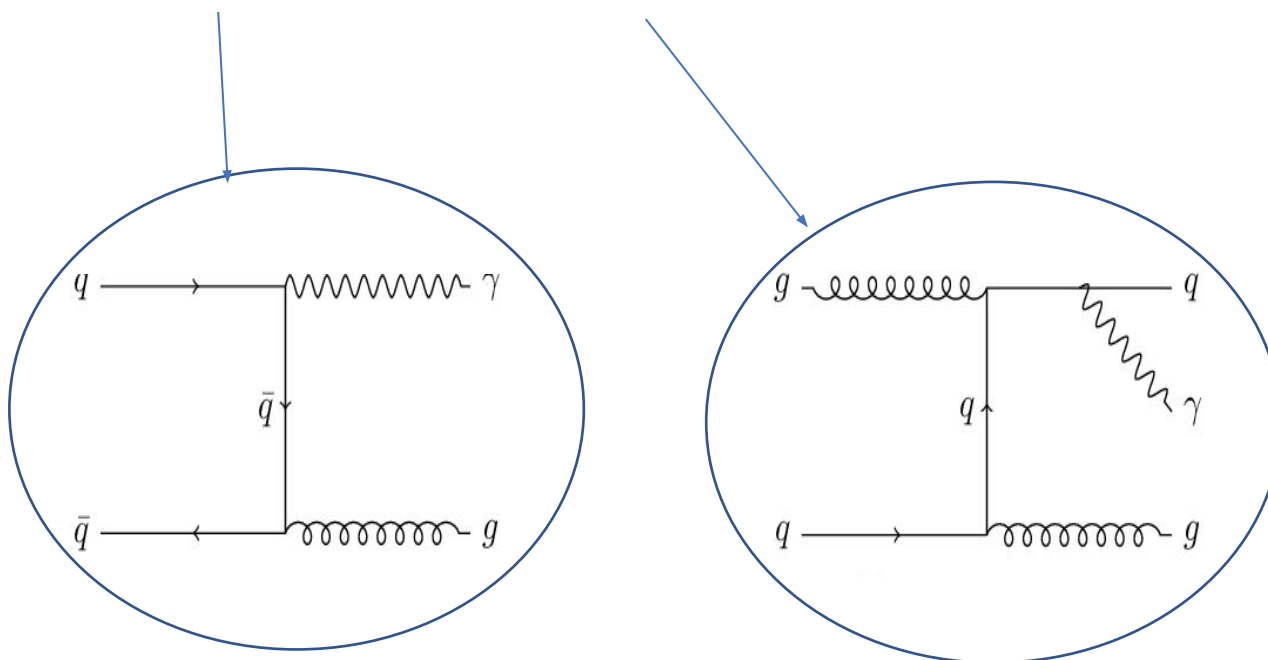


## Hadronic decays

- Final-state particles are hadrons ( $\rightarrow$  jets) and neutrinos
- Complicated by the large number of hadrons *not* from taus. Experimental cuts reject backgrounds based on features of the jets: hard to replicate at particle-level
- A compromise might be best: require a hadron in the jet to have come from a direct tau (this is final-state based *in principle*).
- Not much experience! More detailed studies would be interesting

# What is a final-state photon?

- Analyses usually measure direct/prompt, isolated photons
- Recall: direct or prompt mean “not from a hadron decay”
- But photons can be further divided into those from the hard scatter and those from parton fragmentation & non-perturbative hadronisation



A particle-level isolation criteria is necessary to replicate the isolation applied at reco-level.  
Maybe jet-based. See also problems at fixed-order, cf. Frixione and soft-drop isolation

Note in principle this could also be done for prompt leptons, but it is much less important ... *why?*



# What is a final-state neutrino/invisible?

Invisible in the detector and existence inferred by  $p_T^{\text{miss}}$

Recall:

$$\mathbf{E}_T^{\text{miss}} = - \underbrace{\sum_{\text{selected electrons}} \mathbf{p}_T^e}_{\mathbf{E}_T^{\text{miss}, e}} - \underbrace{\sum_{\text{accepted photons}} \mathbf{p}_T^\gamma}_{\mathbf{E}_T^{\text{miss}, \gamma}} - \underbrace{\sum_{\text{accepted } \tau\text{-leptons}} \mathbf{p}_T^{\tau_{\text{had}}}}_{\mathbf{E}_T^{\text{miss}, \tau_{\text{had}}}} - \underbrace{\sum_{\text{selected muons}} \mathbf{p}_T^\mu}_{\mathbf{E}_T^{\text{miss}, \mu}} - \underbrace{\sum_{\text{accepted jets}} \mathbf{p}_T^{\text{jet}}}_{\mathbf{E}_T^{\text{miss}, \text{jet}}} - \underbrace{\sum_{\text{unused tracks}} \mathbf{p}_T^{\text{track}}}_{\mathbf{E}_T^{\text{miss}, \text{soft}}}$$

hard term soft term

- Sometimes the momenta of (prompt?) invisible\* particles are summed
- An alternative is to take – the sum of all the visible particles within detector acceptance, which is closer to what we measure but can be a bit complicated. E.g. what  $p_T$  of hadrons are we actually sensitive to? And what about reco calibrations? (More on this later)

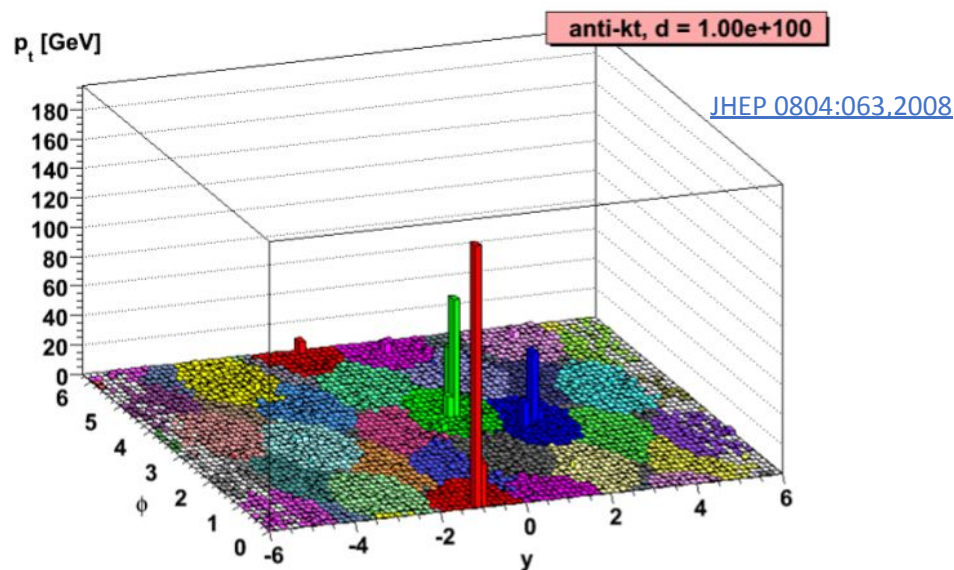
\* neutrinos are indistinguishable from BSM invisible particles, and multiple invisibles are indistinguishable from single invisibles

# What is a final-state parton?

- Partons radiate more partons, recursively. They all eventually hadronize
- Run a jet algorithm on the final-state particles
  - Form a list of particles (this would be clusters / tracks at reco-level)
  - Merge the smallest pair according to a “distance” parameter
  - Iterate (until a stopping condition, cf. the “beam distance”)
- Algorithms assign each hadron to a jet. The energy/momentum of the jet represents the energy/momentum of the parton from the hard scatter
- Think carefully about what is included as inputs: Muons? Neutrinos?

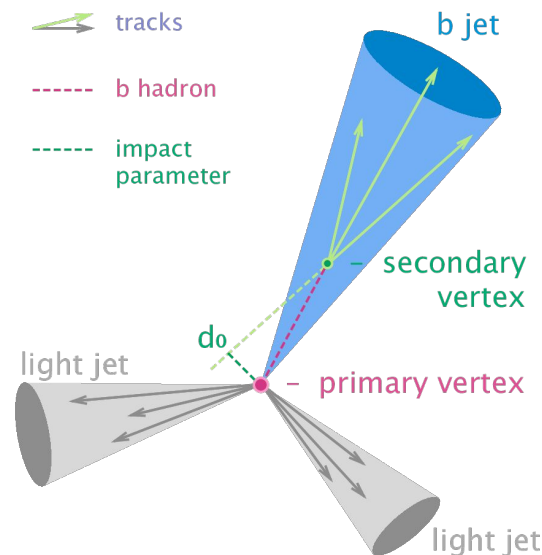
Note: In principle an electron will always form a jet experimentally.

We remove these jets using overlap removal at both reco- and truth-level (e.g. remove any jets with  $\Delta R < 0.4$  from a prompt electron)



# What is a final-state $b$ -quark/jet?

- Recall decay length for a 20 GeV  $b$ -hadrons  $\sim 2$  mm, they are therefore unstable and not included as final-state particles
- But we select them experimentally by making displaced-vertex selections



- Common “compromise” is to associate the *non-final state*  $b$ -hadrons to jets
  - More “in-principle final-state”: it *could* have been reconstructed
- If a jet contains a  $b$ -hadron it is considered a particle-level  $b$ -jet
  - Maybe with a  $p_T$  cut on the  $b$ -hadron, and careful definition of “contains”

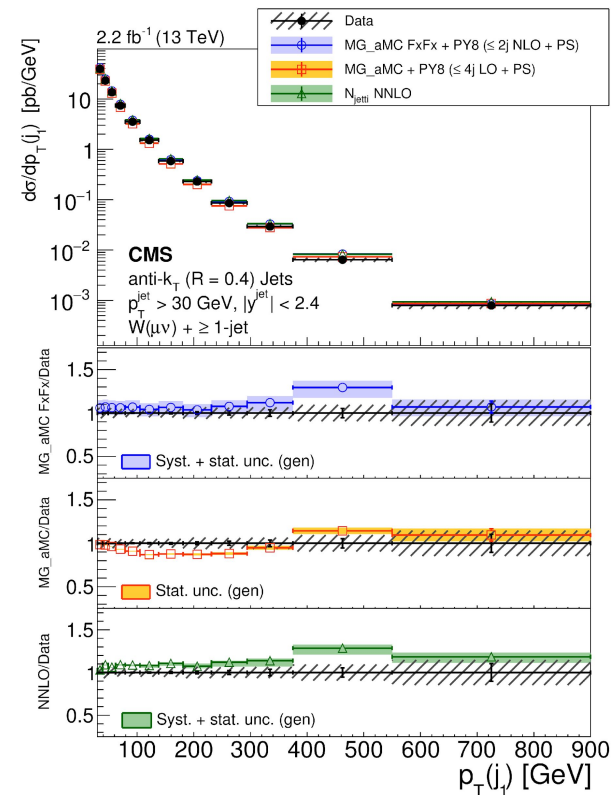
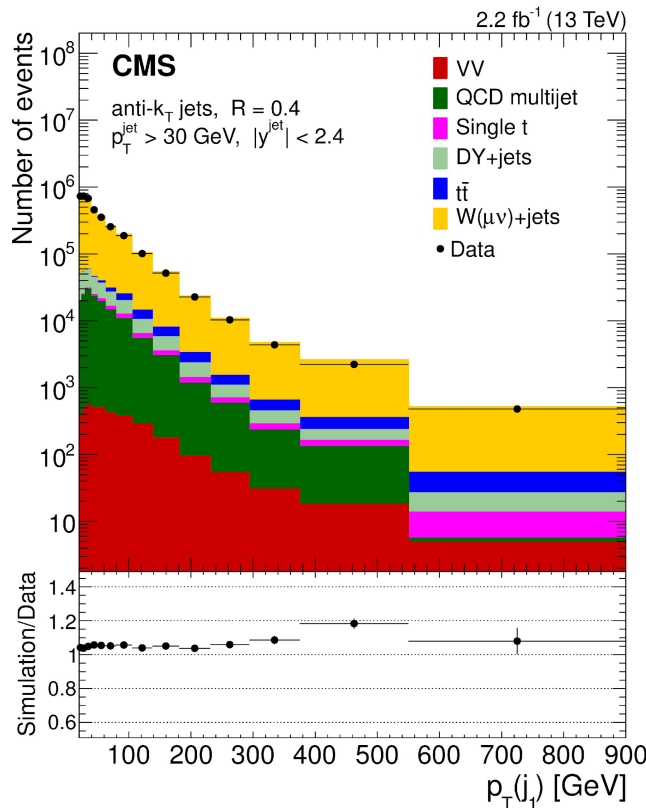
# Examples: CMS $W$ +jets

Phys. Rev. D 96 (2017) 072005  
Rivet: CMS\_2017\_11610623

## Fiducial phase-space:

(Follow data-analysis cuts closely)

- One dressed ( $\Delta R < 0.1$ ) prompt muon with  $p_{T\mu} > 25$  GeV,  $|\eta| < 2.4$
- $m_T > 50$  GeV (using muon and prompt truth neutrino)
- Jets (exclude neutrinos and above muon): anti- $k_T$  ( $R=0.4$ ) with  $p_{Tj} > 30$  GeV,  $|\eta| < 2.4$ , and  $\Delta R > 0.4$  from the muon



# Examples: CMS EWK W+dijet

[arXiv:1903.04040](https://arxiv.org/abs/1903.04040) (Submitted to EPJC)

No Rivet routine

## Data analysis cuts:

jet 1  $p_T > 50$  GeV, jet 2  $p_T > 30$  GeV

$m_{jj} > 200$  GeV

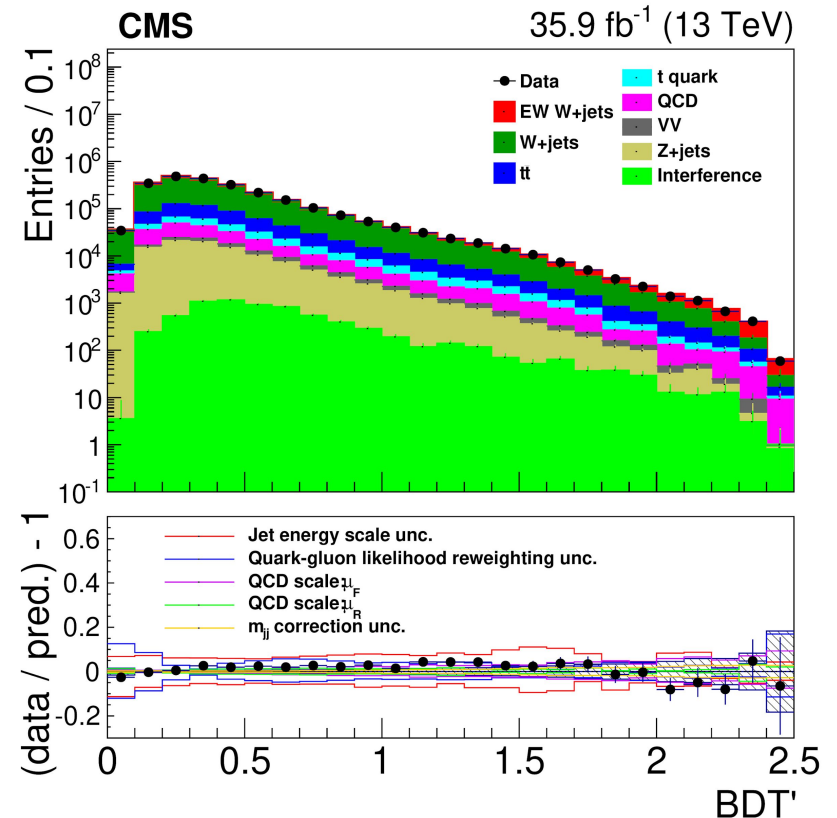
$p_T$  balance cut

$p_T$  lepton  $> 20$  GeV,  $|\eta| < 2.4$

## Fiducial phase-space:

$lv + dijets$  final state

$m_{jj} > 120$  GeV and jet  $p_T > 25$  GeV



$$\sigma_{EW}(Wjj) = 6.23 \pm 0.12 \text{ (stat)} \pm 0.61 \text{ (syst) pb}$$



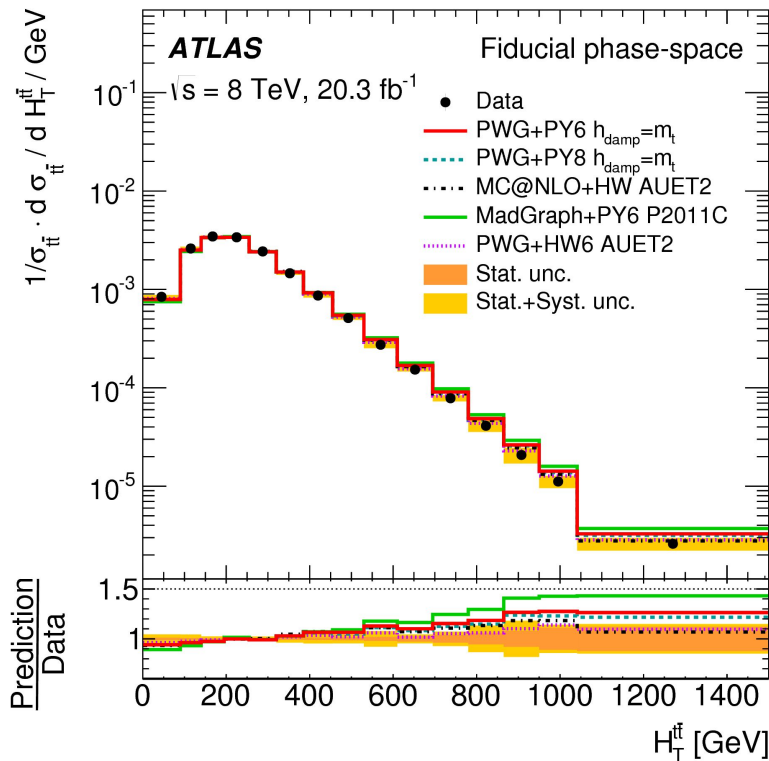
Why such a large extrapolation?



# Examples: ATLAS top-quark analyses

[Eur. Phys. J. C76 \(2016\) 538](#)

Rivet: [ATLAS\\_2015\\_I1404878](#)



Eur. Phys. J. C 78 (2018) 487

DOI: [10.1140/epjc/s10052-018-5904-z](#)



CERN-EP-2017-276

27th June 2018

## Measurement of the inclusive and fiducial $t\bar{t}$ production cross-sections in the lepton+jets channel in $pp$ collisions at $\sqrt{s} = 8 \text{ TeV}$ with the ATLAS detector

The ATLAS Collaboration

The inclusive and fiducial  $t\bar{t}$  production cross-sections are measured in the lepton+jets channel using  $20.2 \text{ fb}^{-1}$  of proton–proton collision data at a centre-of-mass energy of 8 TeV recorded with the ATLAS detector at the LHC. Major systematic uncertainties due to the modelling of the jet energy scale and  $b$ -tagging efficiency are constrained by separating selected events into three disjoint regions. In order to reduce systematic uncertainties in the most important background, the  $W$ +jets process is modelled using  $Z$ +jets events in a data-driven approach. The inclusive  $t\bar{t}$  cross-section is measured with a precision of 5.7% to be  $\sigma_{\text{inc}}(t\bar{t}) = 248.3 \pm 0.7 \text{ (stat.)} \pm 13.4 \text{ (syst.)} \pm 4.7 \text{ (lumi.) pb}$ , assuming a top-quark mass of 172.5 GeV. The result is in agreement with the Standard Model prediction. The cross-section is also measured in a phase space close to that of the selected data. The fiducial cross-section is  $\sigma_{\text{fid}}(t\bar{t}) = 48.8 \pm 0.1 \text{ (stat.)} \pm 2.0 \text{ (syst.)} \pm 0.9 \text{ (lumi.) pb}$  with a precision of 4.5%.

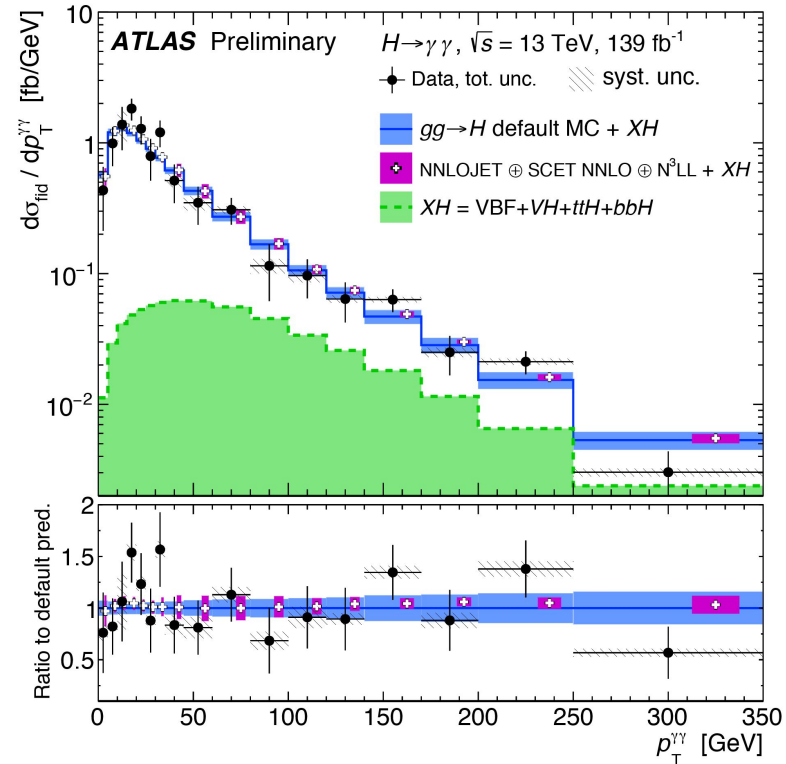
# Examples: ATLAS Higgs $\rightarrow \gamma\gamma$

Table 2: Particle-level selections for the fiducial measurements. The photon isolation,  $\sum p_T^i/p_T^\gamma$ , is defined as the sum of the  $p_T$  of charged particles within  $\Delta r < 0.2$  of the photon.

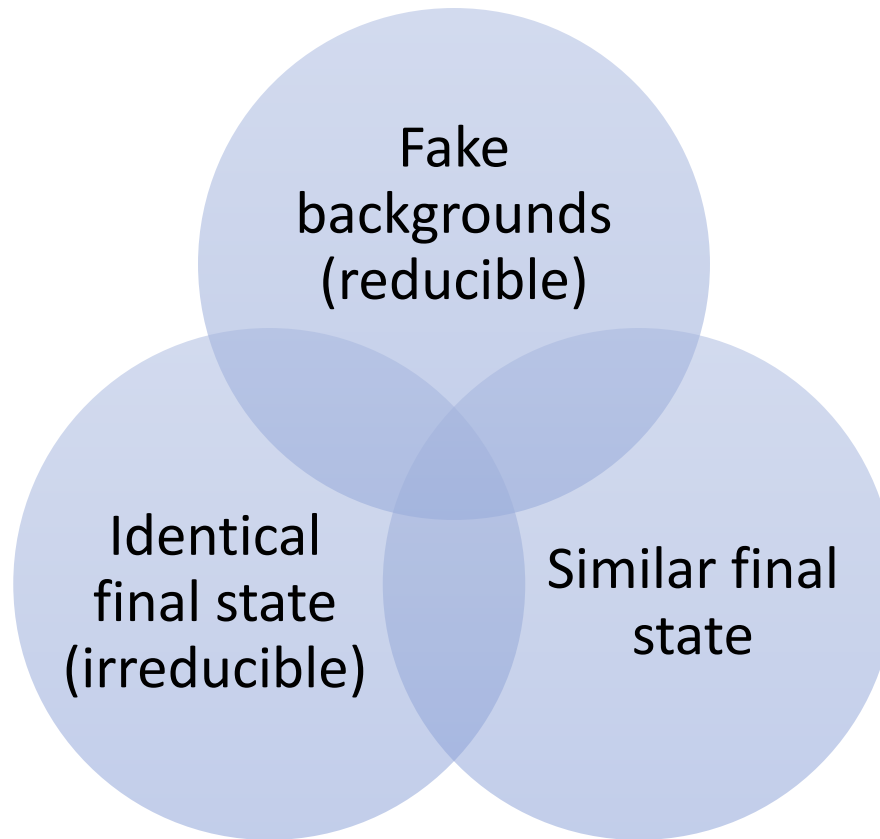
Objects	Fiducial definition
Photons	$ \eta  < 2.37$ (excluding $1.37 <  \eta  < 1.52$ ), $\sum p_T^i/p_T^\gamma < 0.05$
Jets	anti- $k_t$ , $R = 0.4$ , $p_T > 30$ GeV, $ y  < 4.4$
Diphoton	$N_\gamma \geq 2$ , $105 \text{ GeV} < m_{\gamma\gamma} < 160 \text{ GeV}$ , $p_T^{\gamma_1}/m_{\gamma\gamma} > 0.35$ , $p_T^{\gamma_2}/m_{\gamma\gamma} > 0.25$

## Fiducial Higgs!

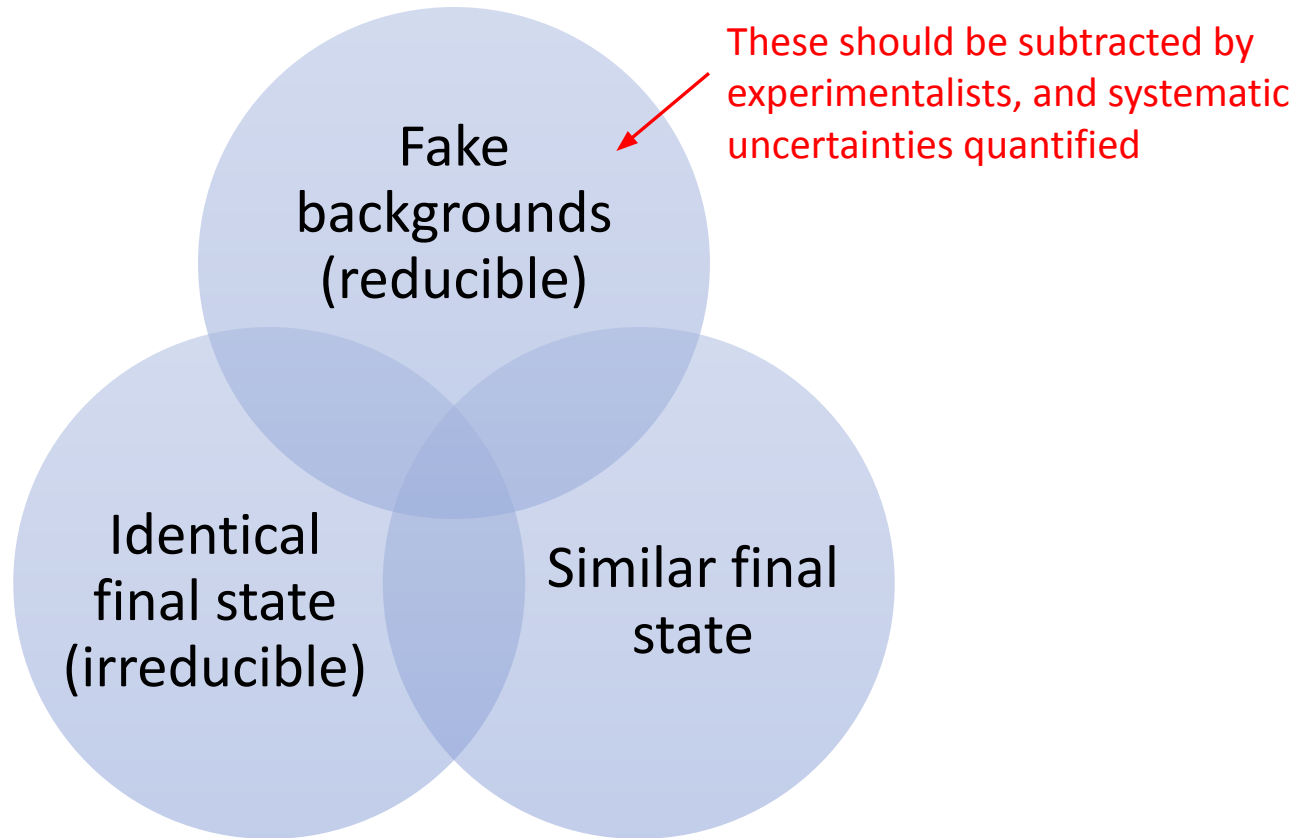
An important  
model-independent  
complement to other, more  
BSM-oriented approaches like  
STXS bins (per-production  
mode, summed over decays,  
often derived by unfolding from  
BDTs, etc.)



# Background-subtraction... or not?



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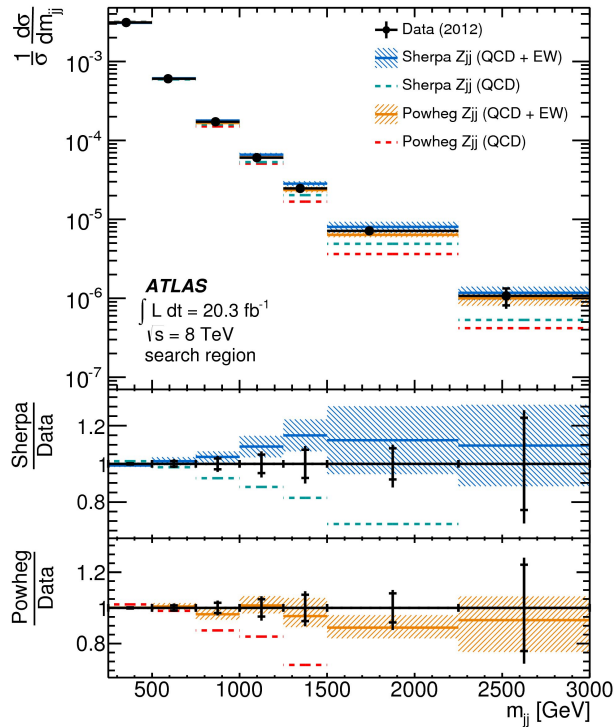
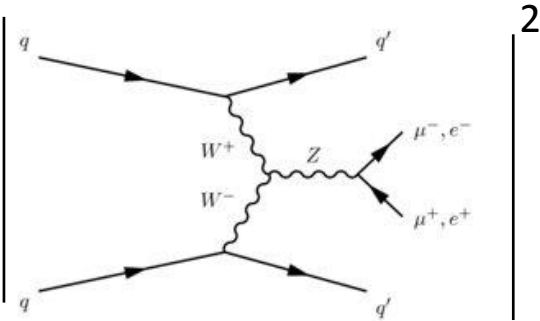


# Identical final-states

- Quantum mechanics tells us that processes with identical final states will interfere and cannot be calculated separately
  - Usually. Defining “identical” can be subtle, cf. colour indices
- Sometimes this is a huge effect and separating out diagrams breaks gauge-invariance
- Other times the effect is quite small and attempts are made to isolate certain processes (e.g. diagram removal/subtraction)



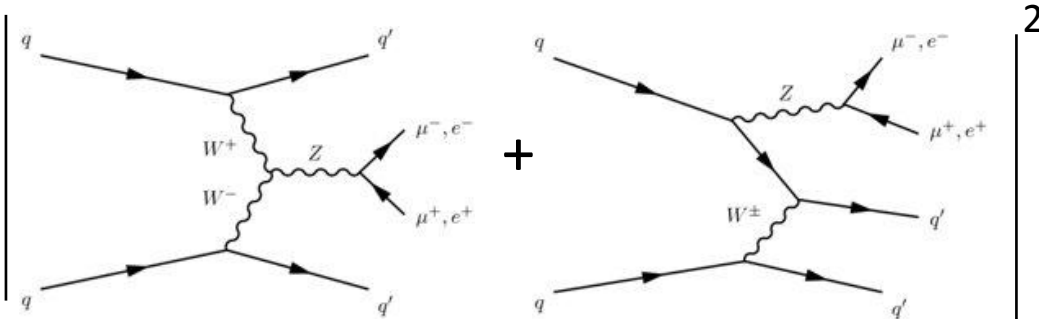
# Identical final-states example: $\ell^+ \ell^-$ VBF



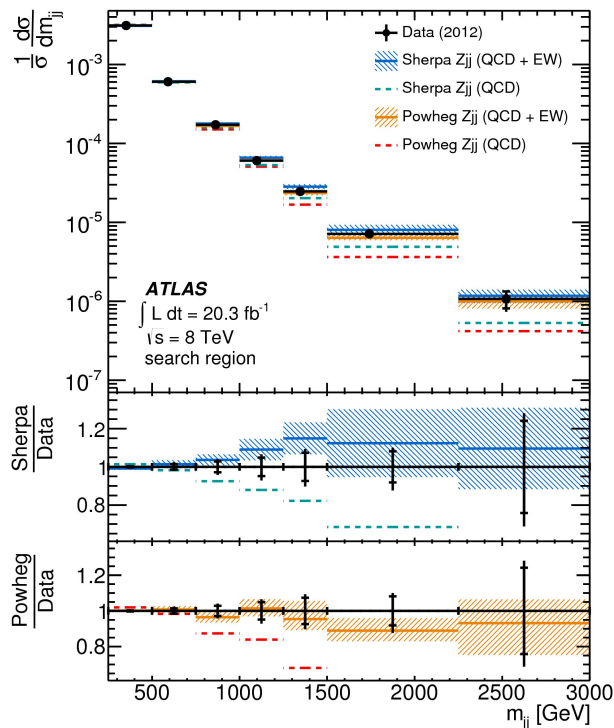
[JHEP04\(2014\)031](#)

[Rivet: ATLAS\\_2014\\_I1279489](#)

# Identical final-states example: ... is $\ell^+ \ell^- + \text{dijet}$



2

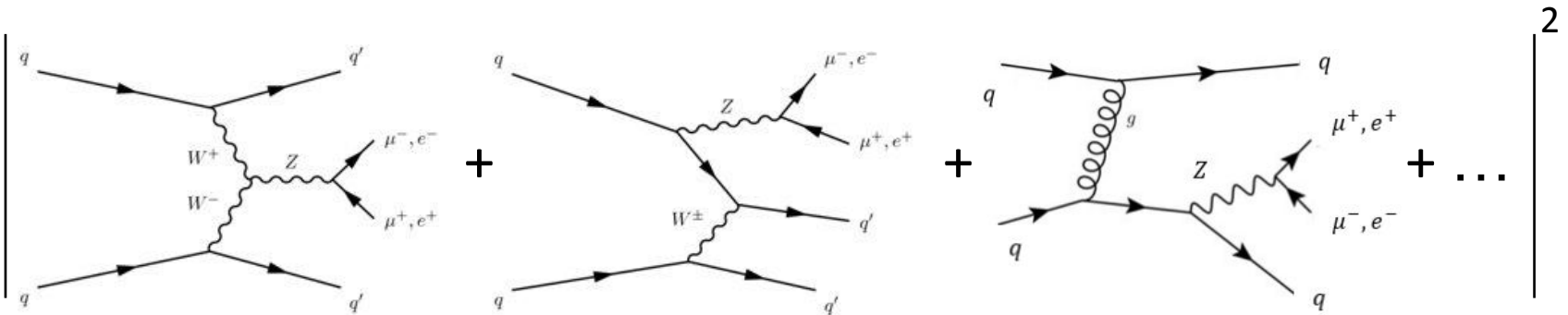


[JHEP04\(2014\)031](#)

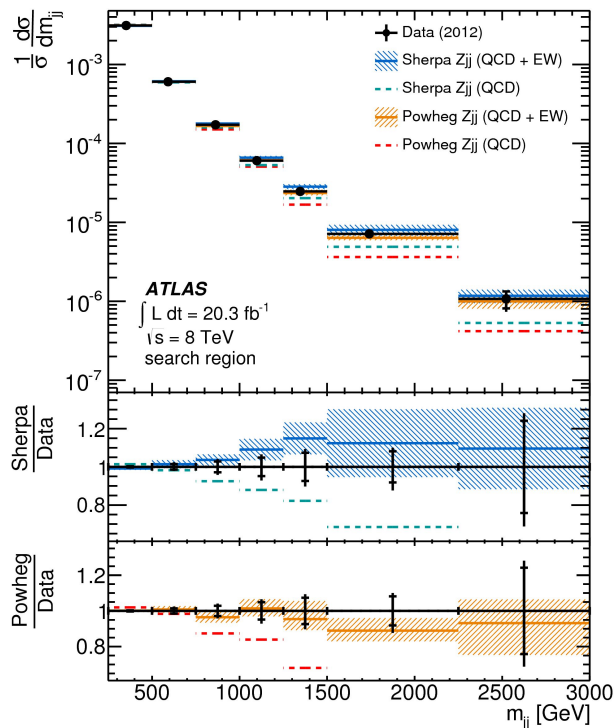
Rivet: ATLAS\_2014\_I1279489

“The VBF process cannot be isolated due to a large destructive interference with the electroweak Z-boson bremsstrahlung process.”

# Identical final-states example: ... etc.!



2

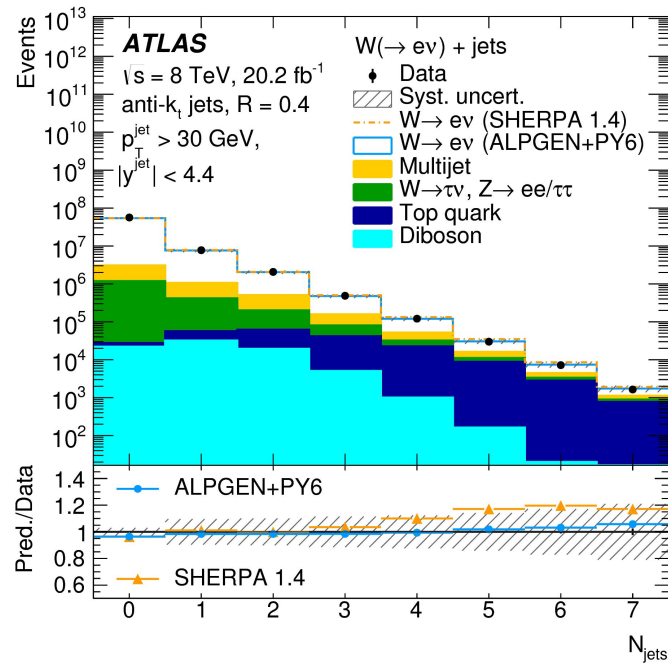
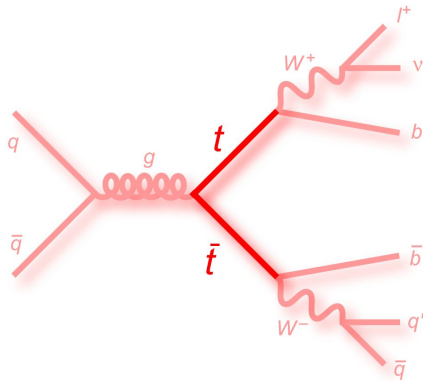


[JHEP04\(2014\)031](#)

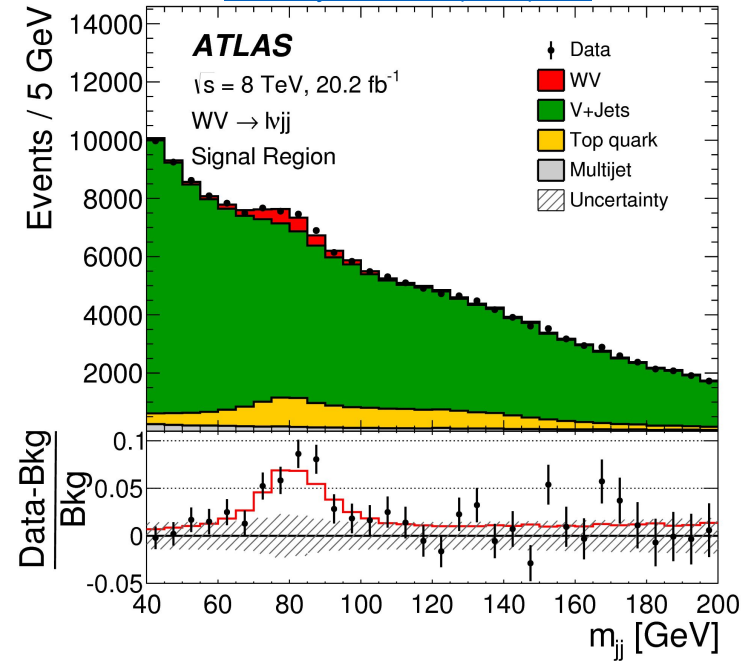
Rivet: ATLAS\_2014\_I1279489

It's usually good to also include an inclusive measurement with no assumptions or subtractions made

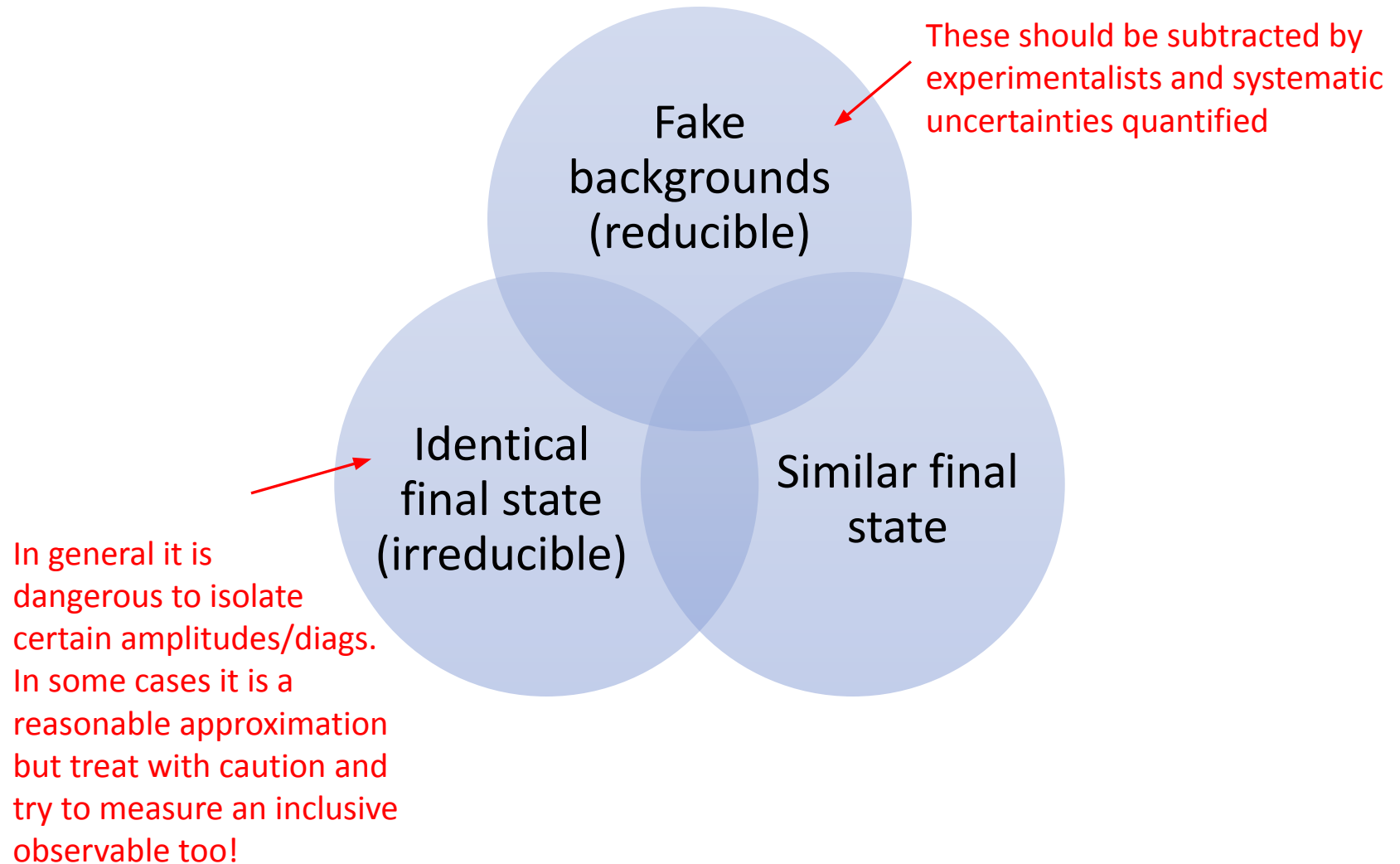
# And $t\bar{t}$ is $WW$ +jets is $W$ +jets...



[Eur. Phys. J. C 77 \(2017\) 563](#)

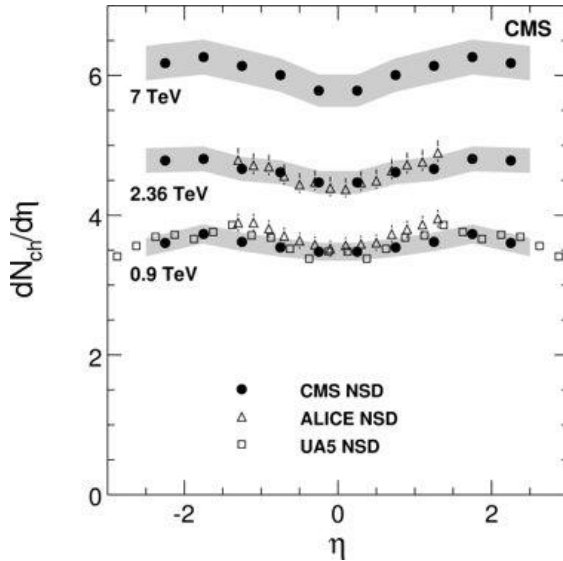


# So, background-subtraction... or not??





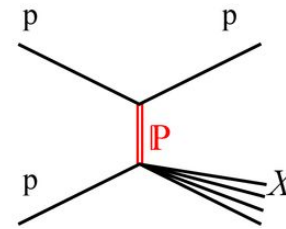
# Similar final-states: soft-QCD



[Phys. Rev. Lett. 105 \(2010\) 022002](#)

Rivet: CMS\_2010\_S8656010

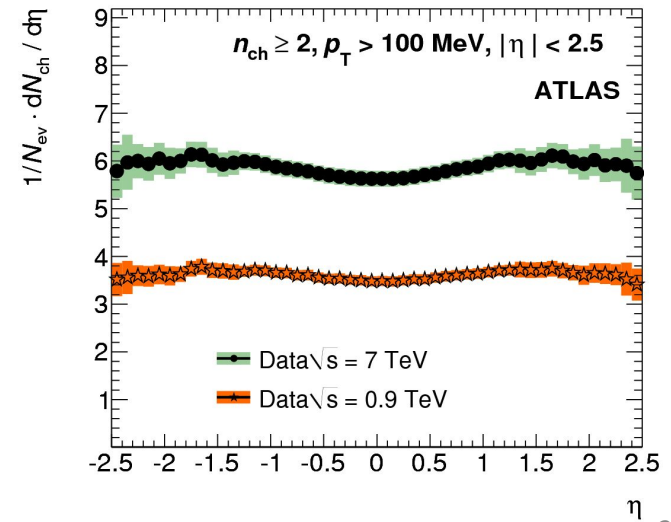
Single-diffractive subtracted!!!



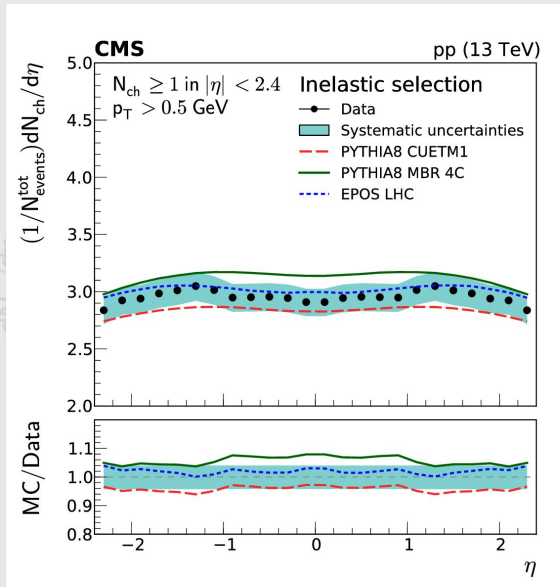
[New J. Phys. 13 \(2011\) 053033](#)

Rivet: ATLAS\_2010\_S8918562

Final-state particle definition



# Similar final-states: soft-QCD



[Eur. Phys. J. C 78 \(2018\) 697](#)

Rivet: CMS\_2018\_I1680318

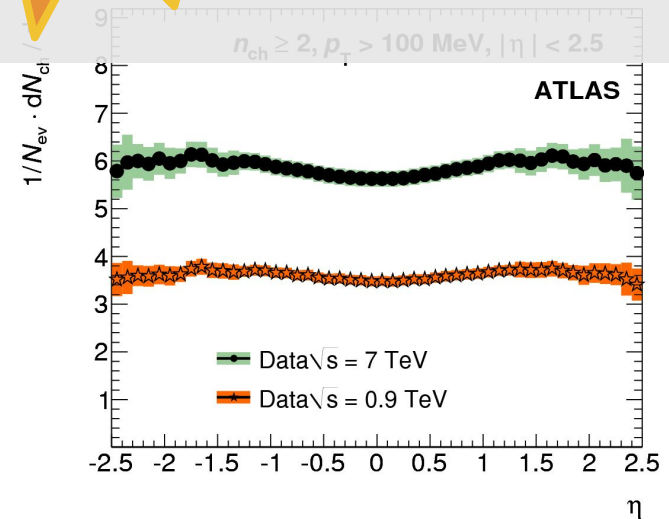
Final-state particle definition !!



[New J. Phys. 13 \(2011\) 053033](#)

Rivet: ATLAS\_2010\_S8918562

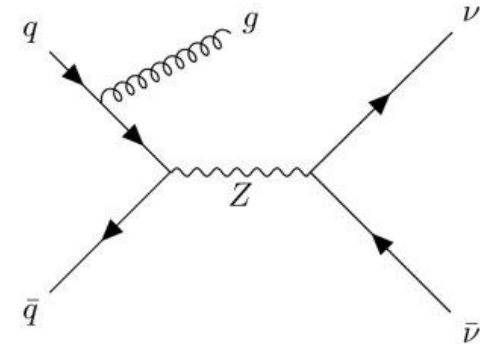
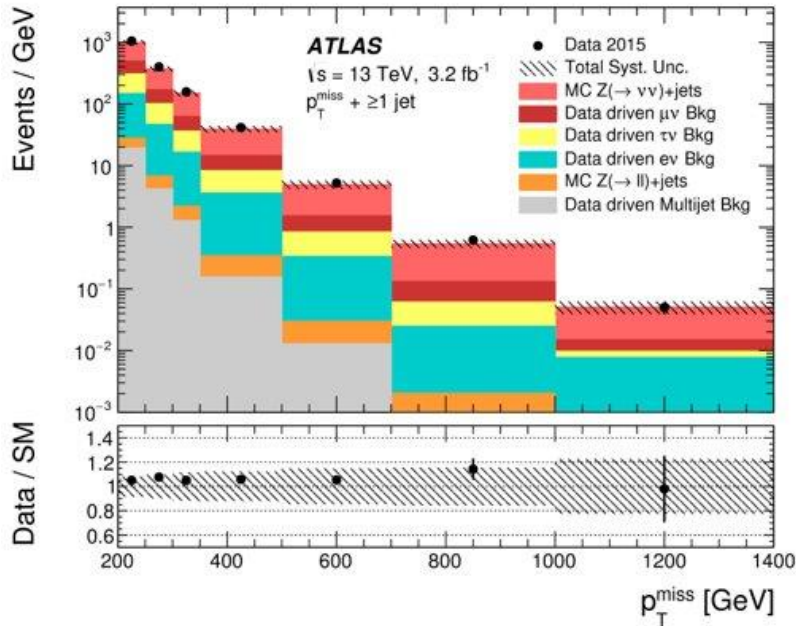
Final-state particle definition



# Similar final-states: $p_{T,\text{miss}} + \text{jets}$

[Eur. Phys. J. C 77 \(2017\) 765](#)

Rivet: ATLAS\_2017\_I1609448



Fiducial phase-space:

$p_{T,\text{miss}} + \text{jet(s)}$

No charged leptons with  $|\eta| < 2.5$ ,  $p_T > 7 \text{ GeV}$   
 + others

**$W \rightarrow l\nu$  with “out of acceptance” leptons contribute  $\sim$  the same as  $Z \rightarrow \nu\nu$ !**

- In this paper: background determined using control regions+MC, and subtracted
- Perhaps these  $W$ 's should be included as part of the “signal” definition?
  - This leaves the data uncontaminated and as close to “what we see” as possible.
  - Removes dependence on control regions and MC extrapolation between regions
- But be careful of fiducial phase-space definitions: e.g. out-of-acceptance muons should be included as *invisible* in a particle-level  $p_{T,\text{miss}}$  definition!

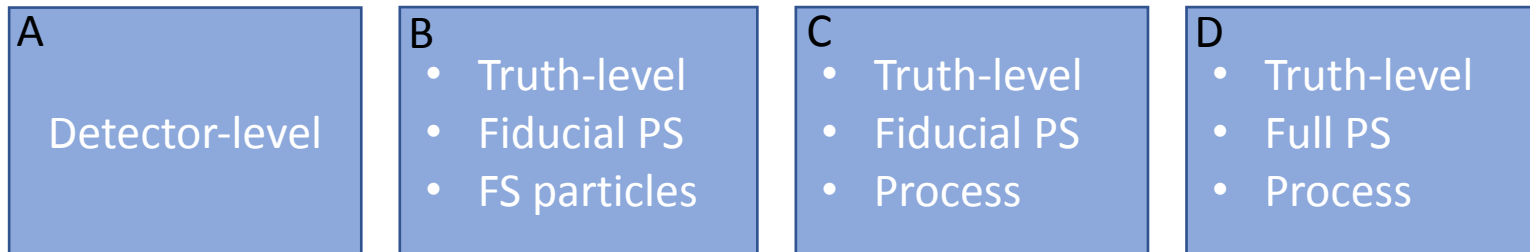
# What is the perfect measurement?

optimal

What is the ~~perfect~~ measurement?

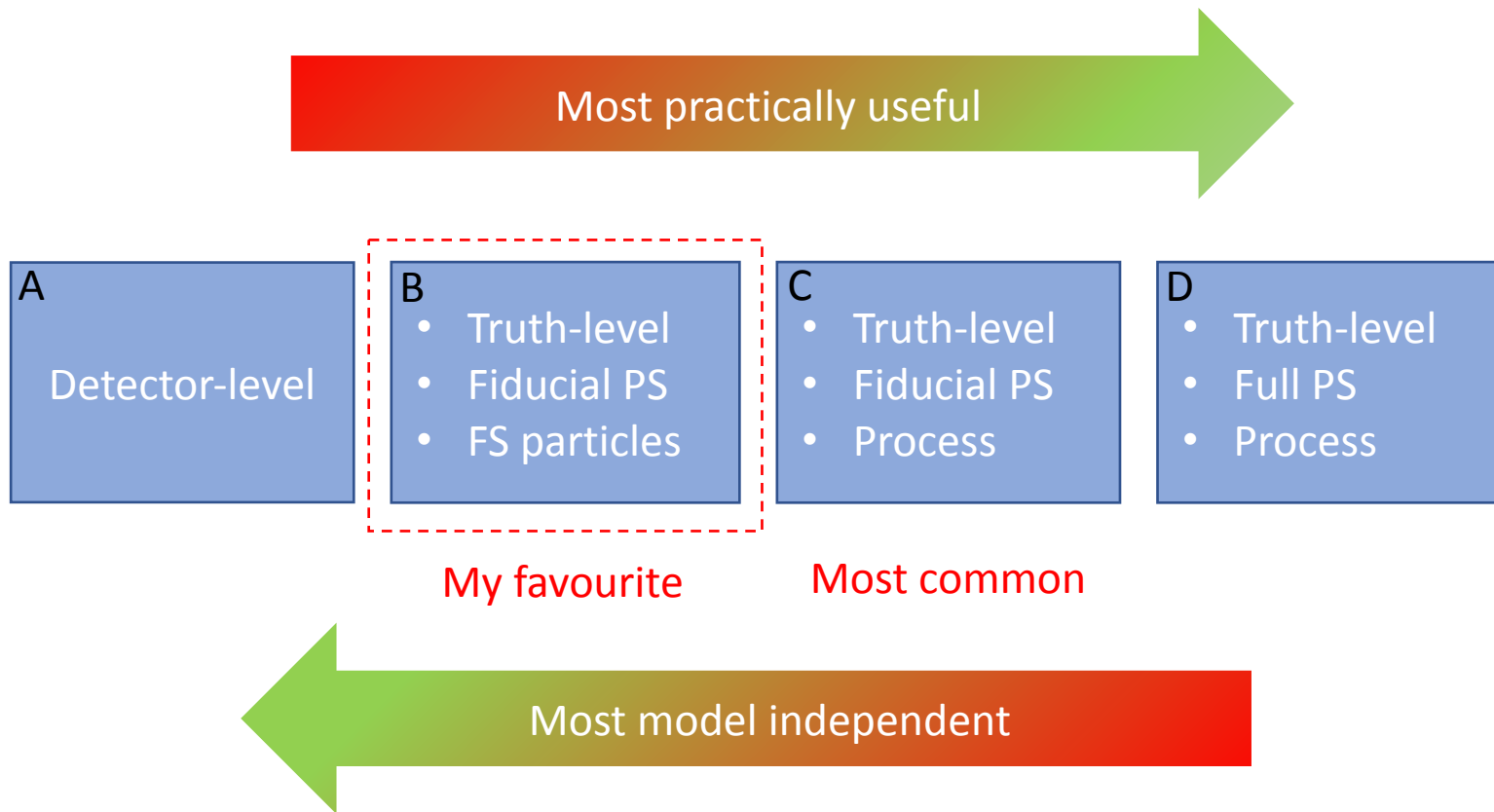
optimal

# What is the ~~perfect~~ measurement?



optimal

# What is the ~~perfect~~ measurement?



# BSM searches & detector-correction

- Typically done at reco-level, searches set limits on parameters in a given model by comparing to *reco-level* MC predictions

- But data in a given analysis can be sensitive to *many* BSM theories.

*How to re-interpret the measurements?*

- Many people working on how to reinterpret reco-level results, e.g. by using fast detector simulation (can be interfaced with Rivet)
- Another option is to correct for detector effects and allow comparisons with “truth-level” predictions. Some sensitivity may be lost due to binning but much easier to reinterpret

Distinction between BSM search and SM measurement becomes blurred:

- We measure the data in certain final-states and compare to the best SM predictions.
- We should do this more in regions particularly sensitive to new physics
- Important to stick to the “measuring a final-state” philosophy

See Contur <https://contur.hepforge.org> and its tutorial:

[arXiv:2102.04377](https://arxiv.org/abs/2102.04377)

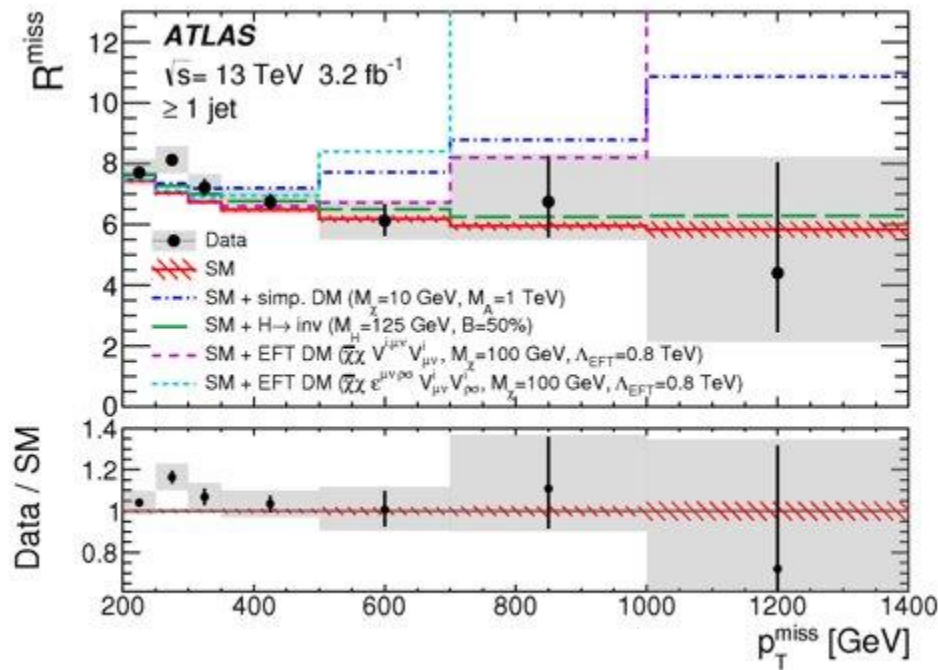
⇒ idea is to use all analyses in Rivet to constrain BSM parameter spaces

Sometimes surprises occur and a certain model pops up in multiple final states we haven't thought of yet!

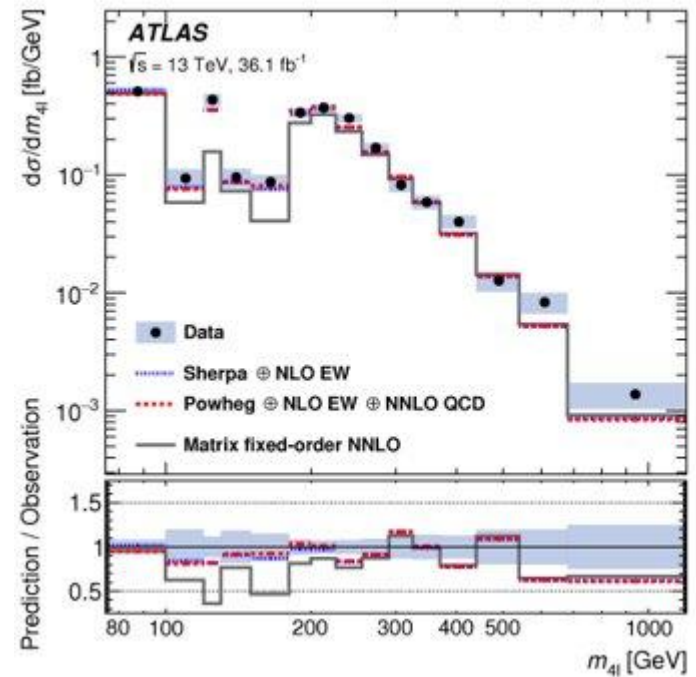


# BSM searches & detector-correction

Unfolding model-independent observables in BSM-search phase-space is a pretty new, and very exciting thing to do at the LHC!



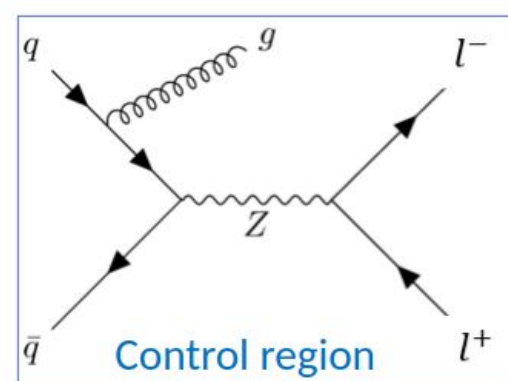
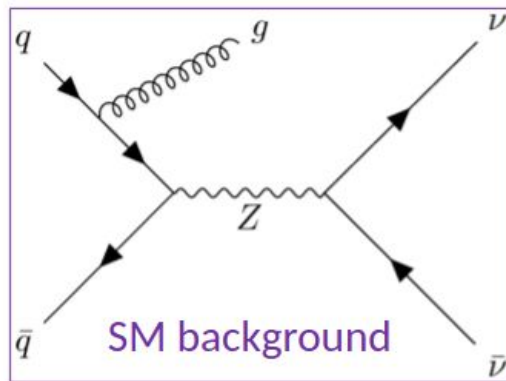
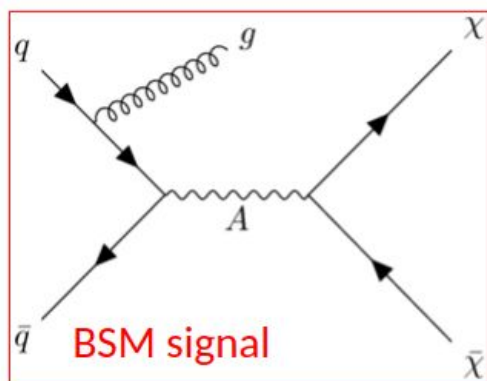
Eur. Phys. J. C 77 (2017) 765



JHEP 04 (2019) 048

# BSM searches & backgrounds

- Often backgrounds to BSM searches are predicted using constraints from “control regions”. These can be
  - similar final states, or
  - the same final-state with different kinematic cuts.
- This can be very useful, especially when modelling is bad, and can reduce systematic uncertainties a lot.



But it can limit re-interpretation: what if another BSM theory leads to final-state particles in the control region too?

“Everyone’s signal is someone else’s background”!

# BSM searches: backgrounds & unfolding

A possible solution:

- Unfold and publish the signal region and the control region with correlation information
- Control region constraints can then be made for models that allow it but not for others

Sometimes, maybe we should just say what we see...

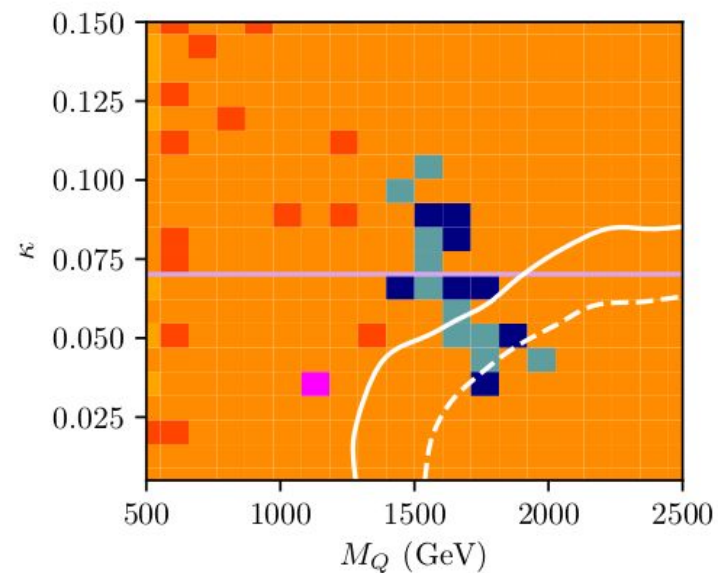
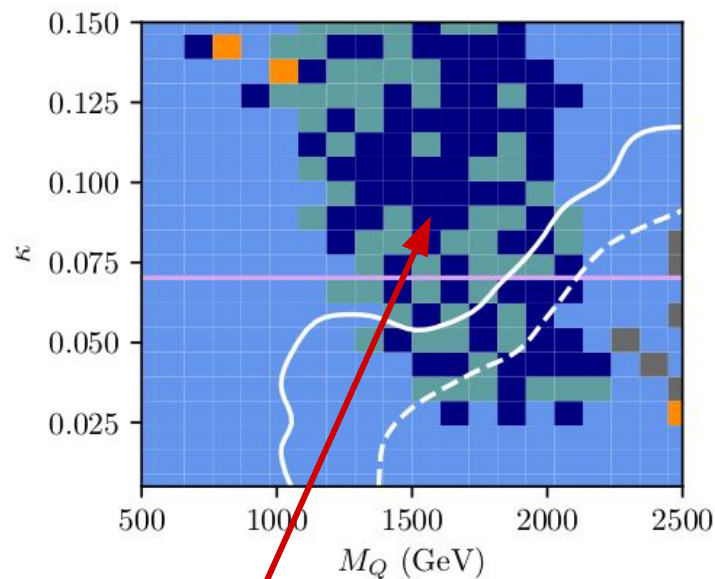
**These are all (weirdly) new ideas.**

**There's lots of room for studies and analyses — get involved!**

# BSM searches & control-region unfolding

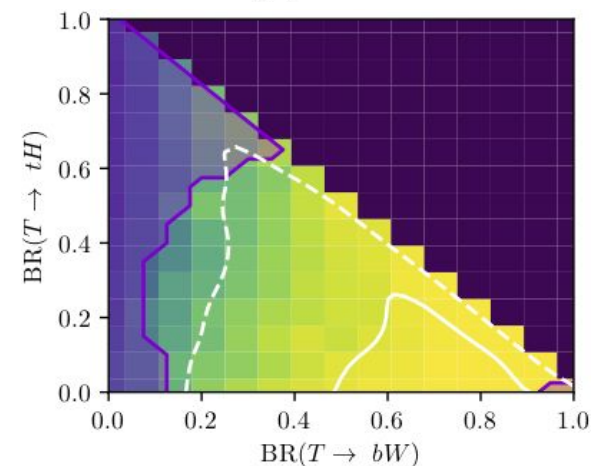
SciPost Phys. 9 (2020) 5, 069

[arxiv:2006.07172](https://arxiv.org/abs/2006.07172)



VLQ reinterpretation: unexpected  
“resonant” injection from a  $Wjj$  unfolded  
control-region in Rivet!

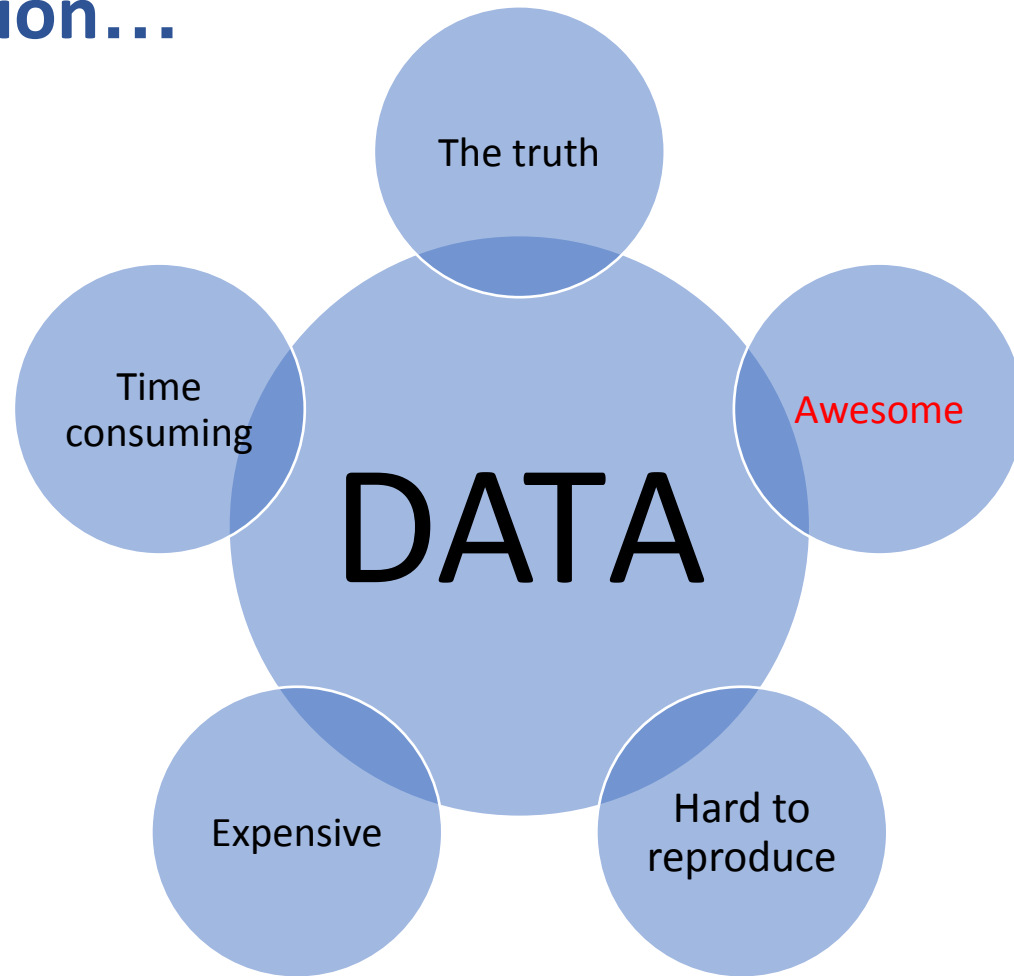
Another *search control-region* was also  
important in adding sensitivity!



# Better experimental analysis: a summary

- ✓ Correct (carefully) for detector effects  
(maybe even for BSM searches)
- ✓ Measure your fiducial phase-space
- ✓ Think carefully about subtracting “backgrounds”
- ✓ Keep the data as clean and model-independent as possible

**In conclusion...**



**... data is still awesome: look after it!**