



# BESIII上自旋极化相关课题的研究

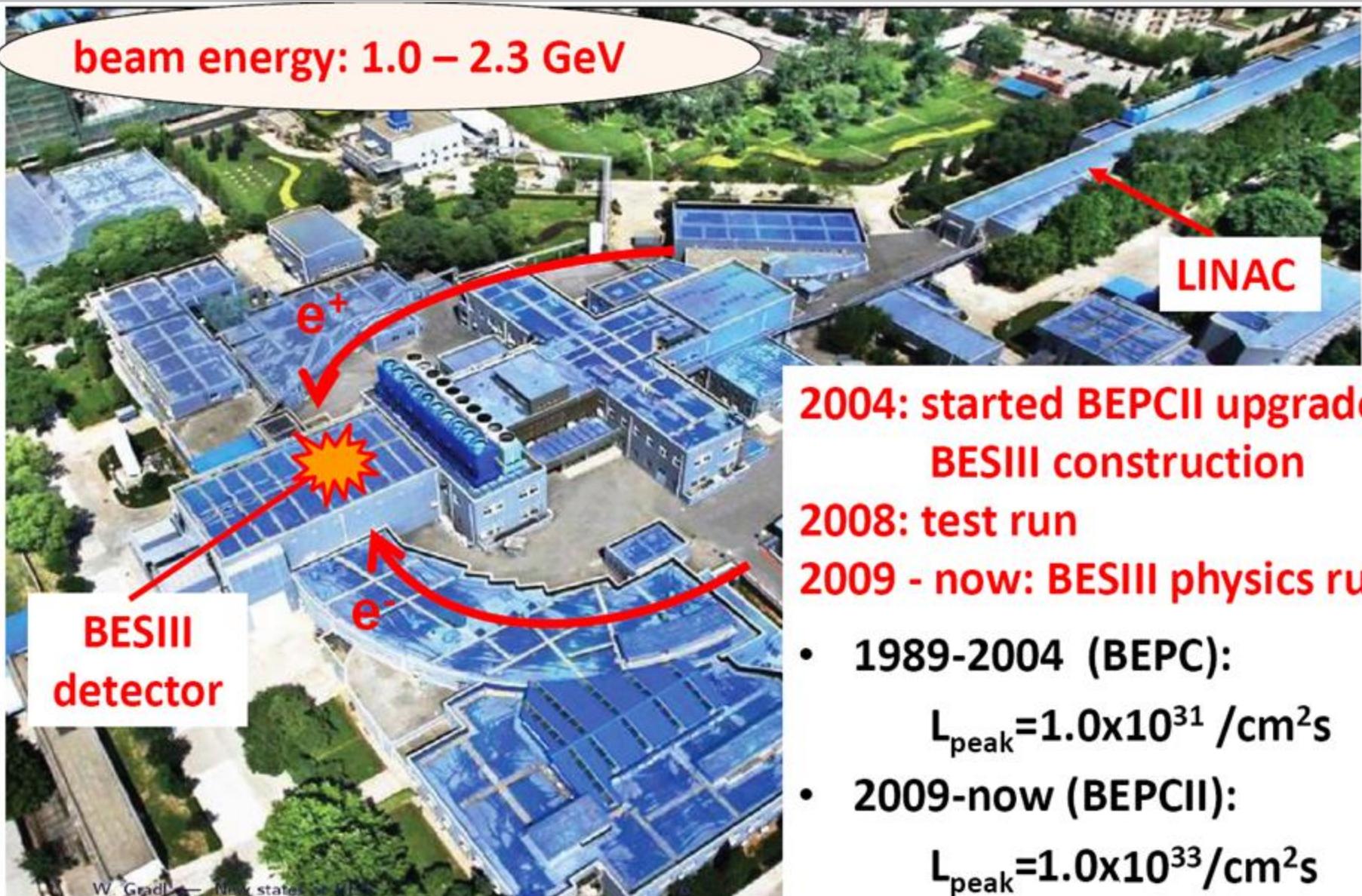
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强子物理新发展研讨会 (Online), 24-26 April 2020

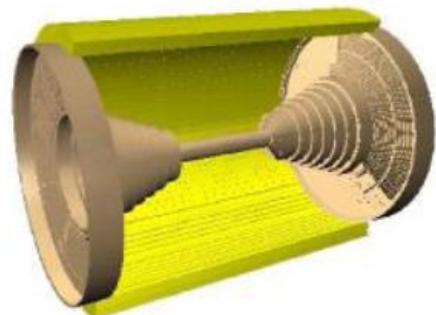
# Beijing Electron Positron Collider (BEPC)

beam energy: 1.0 – 2.3 GeV

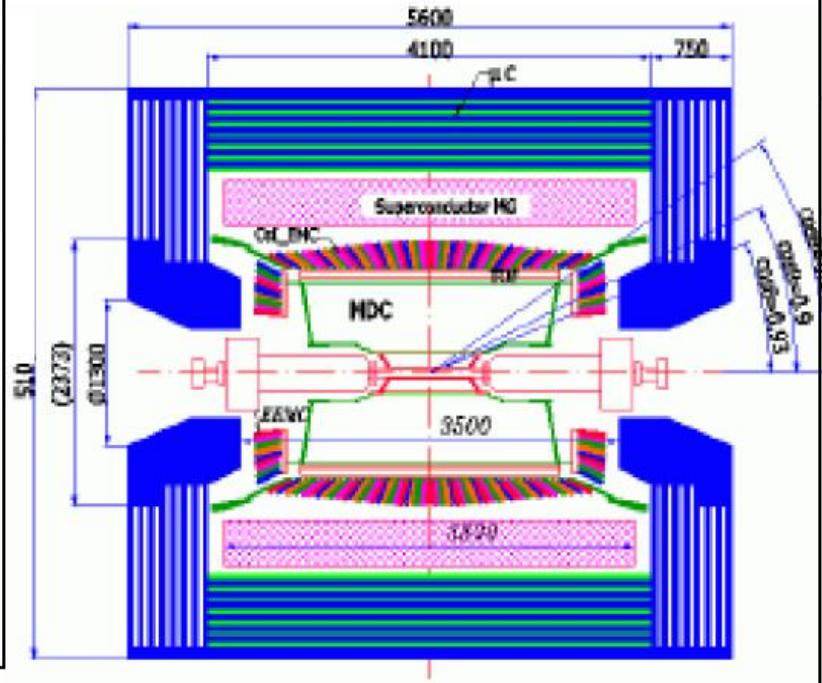


# BESIII Detector

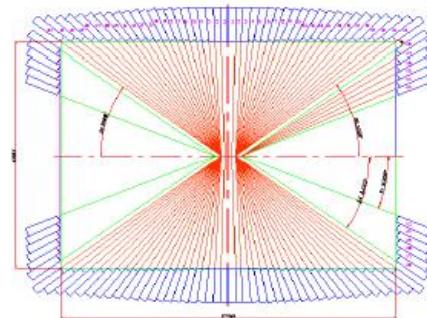
MDC



R inner: 63mm ;  
R outer: 810mm  
Length: 2582 mm  
Layers: 43

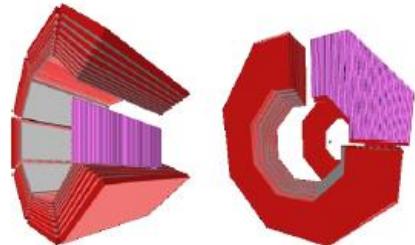


CsI(Tl) EMC



Crystals: 28 cm( $15 X_0$ )  
Barrel:  $|cos\theta| < 0.83$   
Endcap:  
 $0.85 < |cos\theta| < 0.93$

RPC MUC



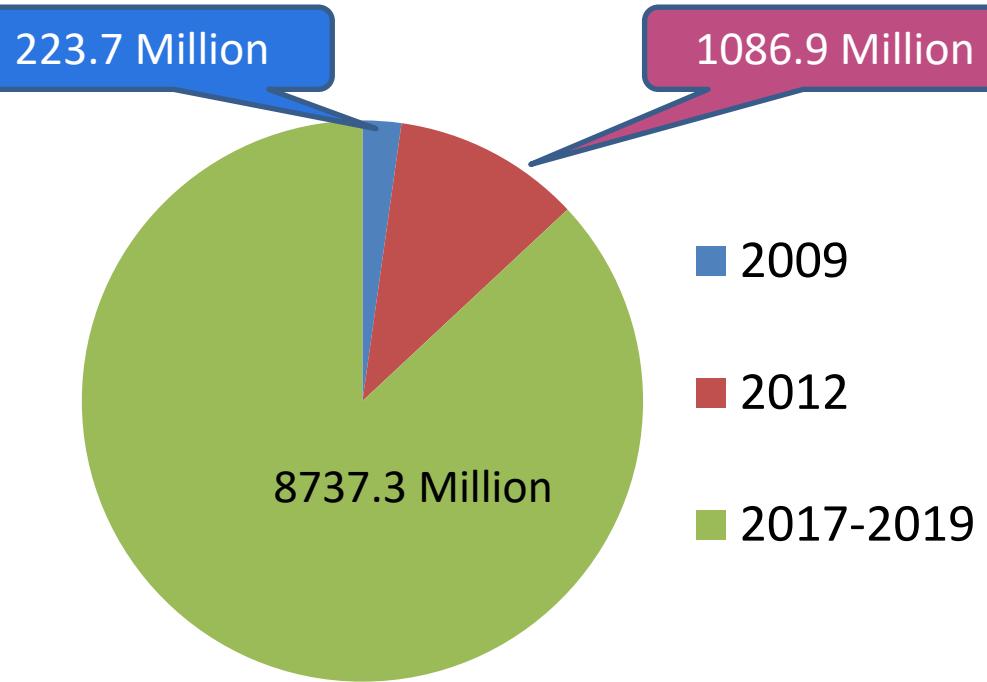
BMUC: 9 layers – 72 modules  
EMUC: 8 layers – 64 modules

TOF

BTOF: two layers  
ETOFT: 48 scintillators for each  
MRPC --- new ETOF



# BESIII $J/\psi, \psi'$ data sets



The analysis based on  
2009+2012  $J/\psi$  data:

~1.31 billion decays

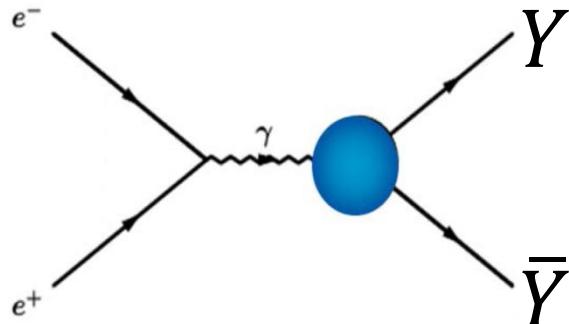
$J/\psi$  : Total 10.047 billion  $J/\psi$  decays

$\psi'$ : 448 million decays

Continuum: 12/fb ( $\sqrt{s} > 3.8$  GeV)

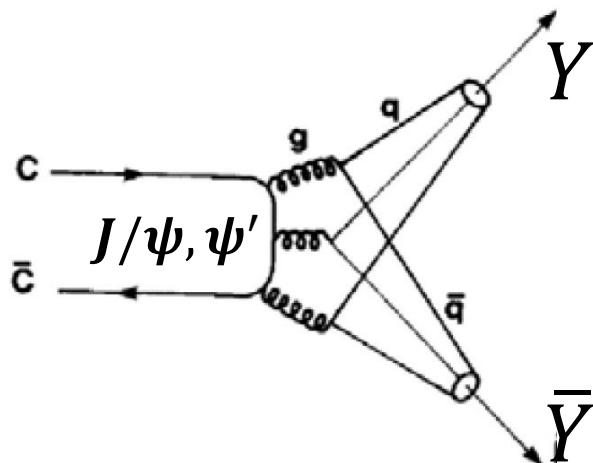
# Hyperon pair production at BESIII

- $e^+e^- \rightarrow \gamma^* \rightarrow \Lambda\bar{\Lambda}, \Sigma\bar{\Sigma}, \Xi\bar{\Xi}, \Omega\bar{\Omega}, \Lambda_c^+\bar{\Lambda}_c^-, @ \sqrt{s} = 2.0 \sim 4.6 \text{ GeV}$ , or update for  $\Lambda_c^+\bar{\Sigma}_c^-, \Sigma_c\bar{\Sigma}_c$



- Threshold enhancement
- Form factor
- Excited hyperon
- .....

- $J/\psi, \psi' \rightarrow \Lambda\bar{\Lambda}, \Sigma\bar{\Sigma}, \Xi\bar{\Xi}, \Omega\bar{\Omega}$



10 billion  $J/\psi$ :

$16.1 \times 10^6$	$\Lambda\bar{\Lambda}$
$12.0 \times 10^6$	$\Xi^0\bar{\Xi}^0$
$12.9 \times 10^6$	$\Sigma^0\bar{\Sigma}^0$
$8.6 \times 10^6$	$\Xi^-\bar{\Xi}^+$

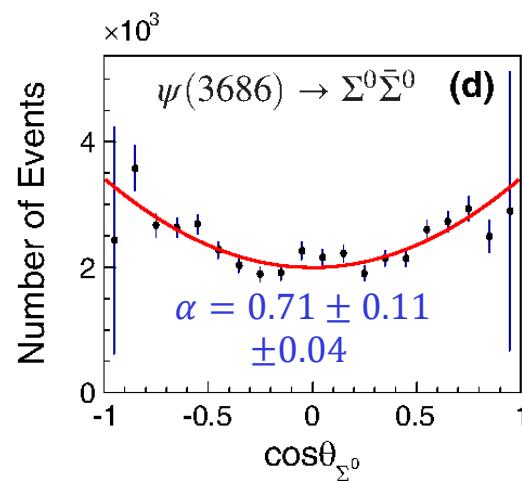
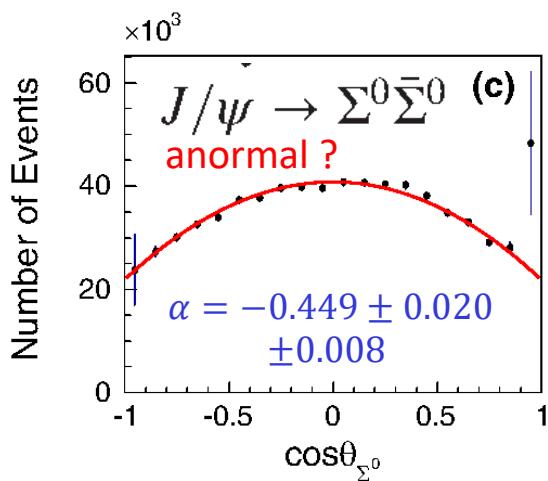
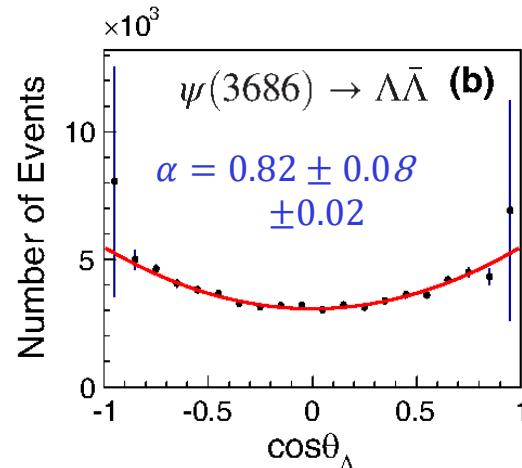
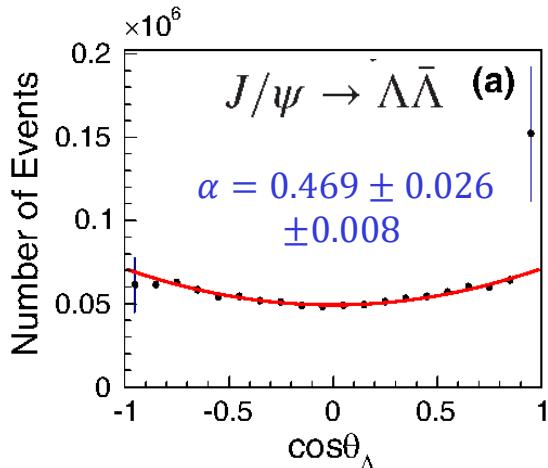
- **Hyperon polarization**
- CP test
- Semileptonic decay
- Radiative decay
- Rare decay

# Role of polarization physics

- Probing of spin degree freedom
  - Test the dynamic of SM and low energy hadron interaction
  - Existant exp. : RHIC, Jlab, GRAAL, CERN and DESY
  - Spin observable, spin-dependent structure function and parton distribution
  - Spin crisis at eighties
- BEPCII/BESIII, unpolarized beam, inaccessible polarization of final state by BESIII
  - Polarized beam for post-BEPCII options, CPV in tau decay, Hyperon weak decay,.....
  - Useful tool: transverse polarization of hyperon, spontaneous production at  $e^+e^-$  collision
  - $\Lambda \rightarrow p\pi^-$  decay plays important role in particle physics

# Polarization puzzle in charmonium decays

Angular distribution:  $\frac{dN}{dcos\theta} \propto (1 + \alpha \cos^2 \theta)$



## □ Branching ratios

$$\frac{\mathcal{B}(\psi(3686) \rightarrow \Lambda\bar{\Lambda})}{\mathcal{B}(J/\psi \rightarrow \Lambda\bar{\Lambda})} = (20.43 \pm 0.11 \pm 0.58)\%$$

$$\frac{\mathcal{B}(\psi(3686) \rightarrow \Sigma^0\bar{\Sigma}^0)}{\mathcal{B}(J/\psi \rightarrow \Sigma^0\bar{\Sigma}^0)} = (20.96 \pm 0.27 \pm 0.92)\%$$

## □ Angular distribution

$$\alpha = \frac{|H_{\uparrow\downarrow}|^2 - 2|H_{\uparrow\uparrow}|^2}{|H_{\uparrow\downarrow}|^2 + 2|H_{\uparrow\uparrow}|^2}$$

Helicity selection rule

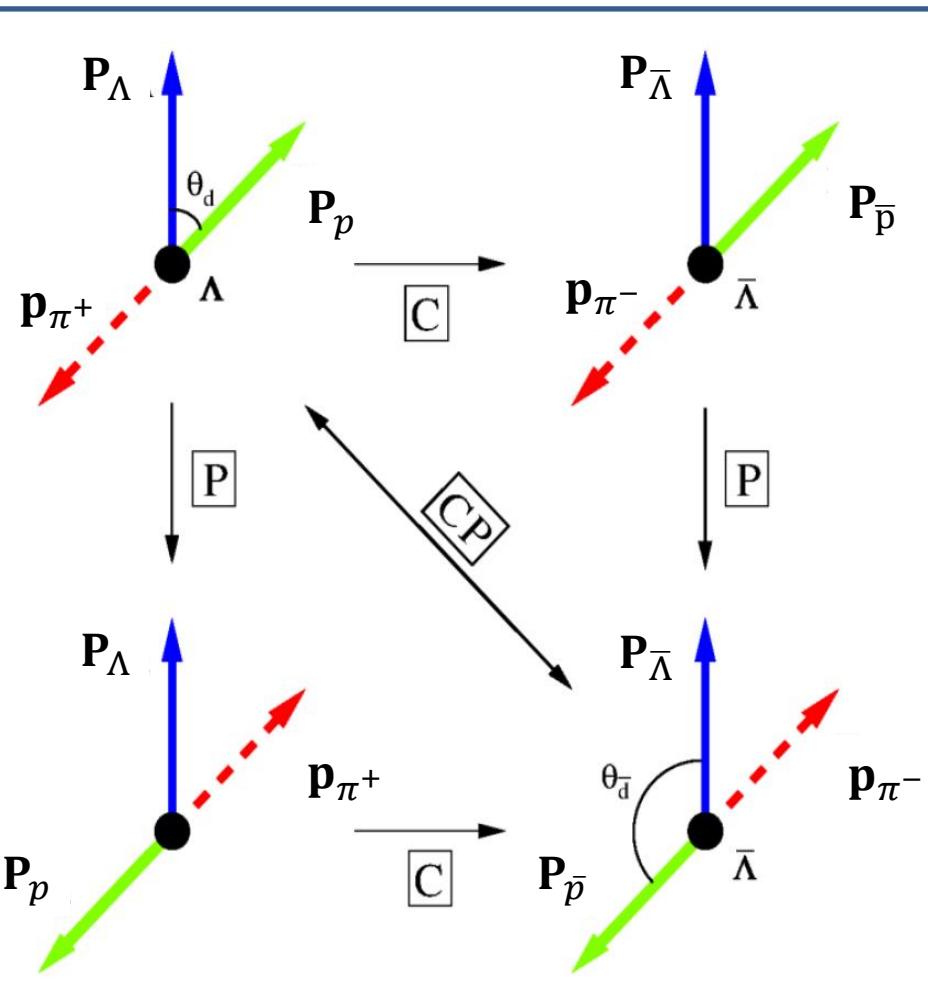
$$|H_{\uparrow\uparrow}|/|H_{\uparrow\downarrow}| \sim 0$$

Measurements:

$$|H_{\uparrow\uparrow}|/|H_{\uparrow\downarrow}| = \begin{cases} \sim 1.3 & (J/\psi \rightarrow \Sigma^0\bar{\Sigma}^0) \\ \sim 0.08 & (\psi' \rightarrow \Sigma^0\bar{\Sigma}^0) \end{cases}$$

Why enhanced for  $J/\psi$ ?

# C,P- transformation in $\Lambda \rightarrow p\pi^-$



$$\alpha = \frac{|B_+|^2 - |B_-|^2}{|B_+|^2 + |B_-|^2},$$

$$\bar{\alpha} = \frac{|\bar{B}_+|^2 - |\bar{B}_-|^2}{|\bar{B}_+|^2 + |\bar{B}_-|^2}$$

CP invariance :

$$\bar{B}_{-\lambda_p} = \eta_\Lambda \eta_p \eta_\pi (-1)^{s_\Lambda - s_p - s_\pi} B_{\lambda_p}$$

$$= -B_{\lambda_p}$$

$$\alpha = -\bar{\alpha}$$

CP odd-variable:

$$\Delta = \frac{\Gamma - \bar{\Gamma}}{\Gamma + \bar{\Gamma}}, A = \frac{\alpha + \bar{\alpha}}{\alpha - \bar{\alpha}}, B = \frac{\beta + \bar{\beta}}{\beta - \bar{\beta}}, B' = \frac{\beta + \bar{\beta}}{\alpha - \bar{\alpha}}$$

# Ancient history

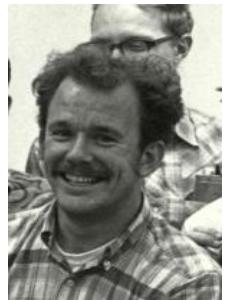
Phys. Rev. Lett. 24, 843 (1970)

## ASYMMETRY PARAMETER FOR $\Lambda^0 \rightarrow n\pi^0 \dagger^*$

S. Olsen, L. Pondrom, R. Handler, and P. Limon†  
University of Wisconsin, Madison, Wisconsin 53706

and

J. A. Smith and O. E. Overseth  
University of Michigan, Ann Arbor, Michigan 48104  
(Received 18 February 1970)



The asymmetry parameter  $\alpha_0$  for  $\Lambda^0 \rightarrow n\pi^0$  has been measured relative to  $\alpha_-$  by comparing the neutron distribution with the proton distribution from the decay  $\Lambda^0 \rightarrow p\pi^-$  for polarized  $\Lambda^0$  hyperons. A sample of 4760 neutron decay events and 8500 proton decay events gave  $\alpha_0/\alpha_- = 1.000 \pm 0.068$  in good agreement with the  $|\Delta I| = \frac{1}{2}$  rule.



Phys. Rev. Lett. 24, 165 (1970)

## PROTON POLARIZATION IN $\Sigma^+ \rightarrow p\pi^0 \dagger^*$

F. Harris and O. E. Overseth  
University of Michigan, Ann Arbor, Michigan 48104

and

L. Pondrom and E. Dettmann  
University of Wisconsin, Madison, Wisconsin 53706  
(Received 10 November 1969)

The polarization of protons from the decay of polarized  $\Sigma^+$  hyperons has been measured by scattering the protons in a carbon-plate spark chamber. A sample of 1335 useful scatters gave  $\alpha_0 = -0.98 \pm 0.05$  and  $\varphi_0 = 22^\circ \pm 90^\circ$ , where  $\tan\varphi_0 = \beta_0/\gamma_0$ . Using the data on  $\Sigma^+ \rightarrow n\pi^+$  and  $\Sigma^- \rightarrow n\pi^-$  and fitting to the  $|\Delta I| = \frac{1}{2}$  rule gave  $\chi^2 = 0.3$  for 2 degrees of freedom.

Phys. Rev. 184, 1663 (1969)

## Final-State Interactions in Nonleptonic Hyperon Decay

O. E. OVERSETH\*

The University of Michigan, Ann Arbor, Michigan 48104

AND

S. PAKVASA†

University of Hawaii, Honolulu, Hawaii 96822

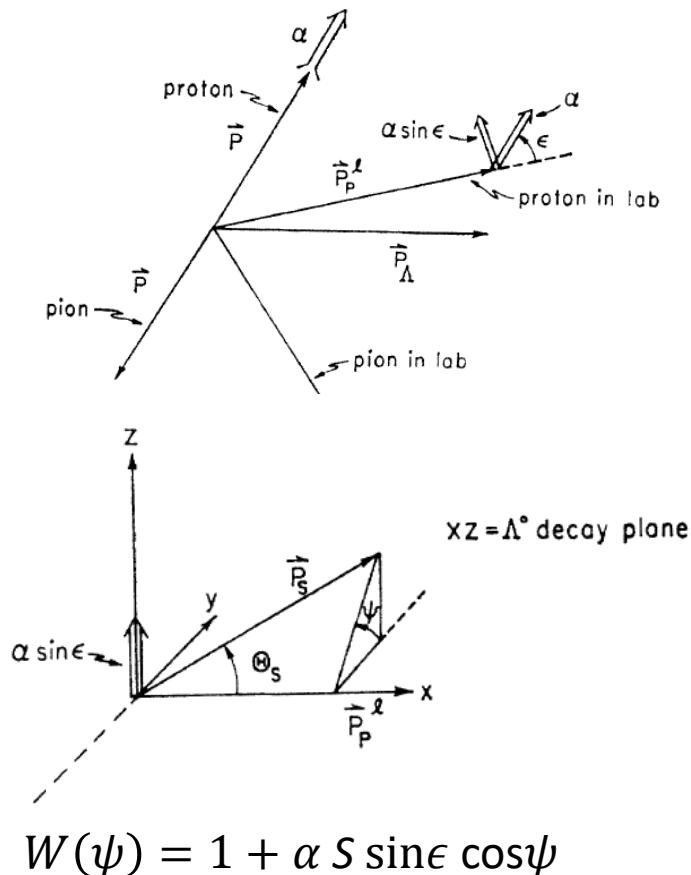
(Received 1 April 1969)



We discuss the consequences of including the final-state interactions in the analysis of  $\Lambda^0$ ,  $\Sigma$ , and  $\Xi$  non-leptonic decays. Emphasis is on the role that the final-state interactions play in tests for  $T$  invariance, in tests of the  $\Delta I = \frac{1}{2}$  rule (including the resolution of sign ambiguities), and in the determination of the decay amplitudes for these processes.

# Most earlier measurement on $\alpha_-$

- CNTR exp.,  $\pi^- + p \rightarrow \Lambda + K^0$
- Unpolarized  $\Lambda \rightarrow p\pi^-$
- Proton polarization measured with carbon-plate spark chamber



Phys.Rev. 129 (1963) 1795-1807

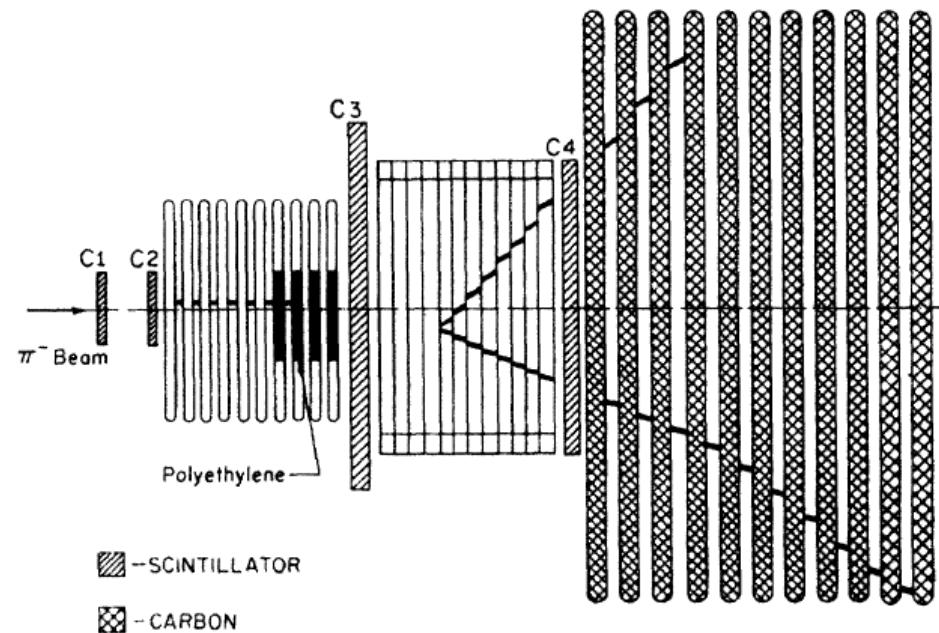


FIG. 1. Schematic diagram showing arrangement of apparatus.  
An example of an event has been sketched in.

$$\alpha = -\frac{2}{\pi} \frac{1}{\langle S \rangle \langle \sin \epsilon \rangle} \frac{N_+ - N_-}{N_+ + N_-},$$

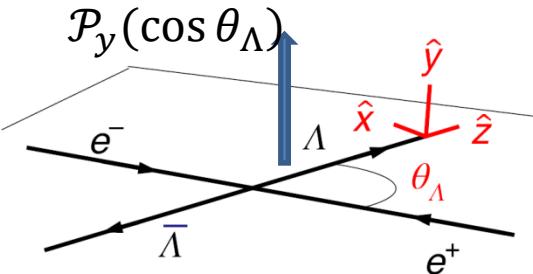
1156 events  
 $\langle S \rangle = 0.565$   
 $\langle \sin \epsilon \rangle = 0.84,$   
 $\alpha = 0.62.$

# Observation of $\Lambda$ polarization and entanglement in $J/\psi \rightarrow \Lambda\bar{\Lambda}$

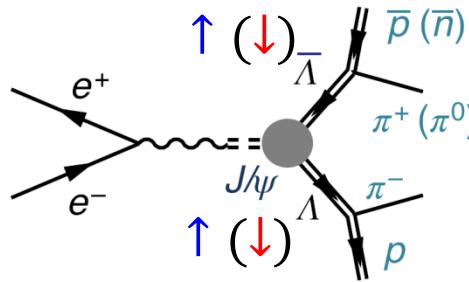


## Highlights!

Published in Nature Physics, 15, 631 (2019)



Helicity system

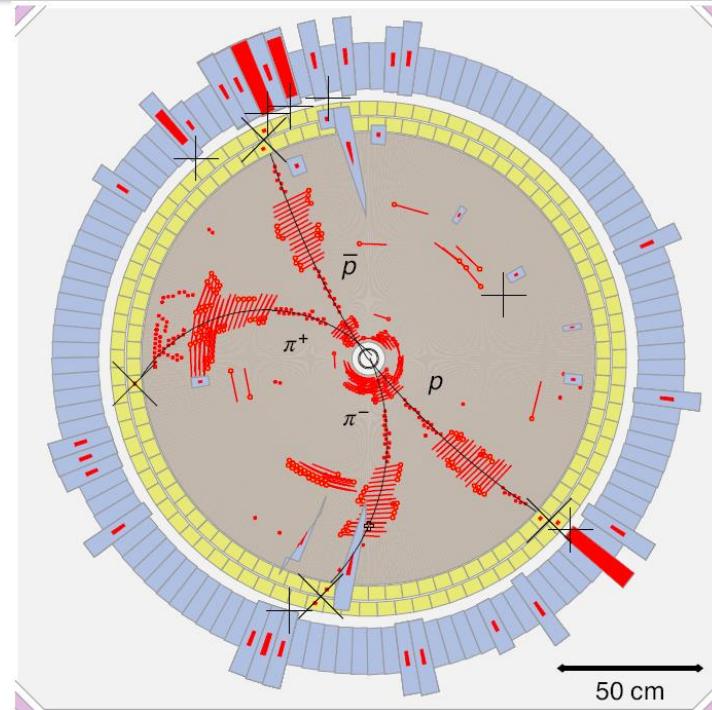


Spin entanglement

### ➤ Joint angular distribution

$$\begin{aligned} & \mathcal{W}(\xi; \alpha_\psi, \Delta\Phi, \alpha_-, \alpha_+) \\ &= 1 + \alpha_\psi \cos^2 \theta_\Lambda + \textcircled{a_- \alpha_+} [\sin^2 \theta_\Lambda (n_{1,x} n_{2,x} - \alpha_\psi n_{1,y} n_{2,y}) \\ &+ (\cos^2 \theta_\Lambda + \alpha_\psi) n_{1,z} n_{2,z}] \quad \text{Entanglement term} \\ &+ \textcircled{a_- \alpha_+} \sqrt{1 - \alpha_\psi^2} \cos(\Delta\Phi) \sin \theta_\Lambda \cos \theta_\Lambda (n_{1,x} n_{2,z} + n_{1,z} n_{2,x}) \\ &+ \sqrt{1 - \alpha_\psi^2} \sin(\Delta\Phi) \sin \theta_\Lambda \cos \theta_\Lambda (\alpha_- n_{1,y} + \alpha_+ n_{2,y}) \end{aligned}$$

Transverse polarization



Event display in BESIII detector

### ➤ Transverse polarization

$$P_y(\cos \theta_\Lambda) = \frac{\sqrt{1 - \alpha_\psi^2} \sin(\Delta\Phi) \cos \theta_\Lambda \sin \theta_\Lambda}{1 + \alpha_\psi \cos^2 \theta_\Lambda}$$

# Observation of $\Lambda$ polarization and entanglement in



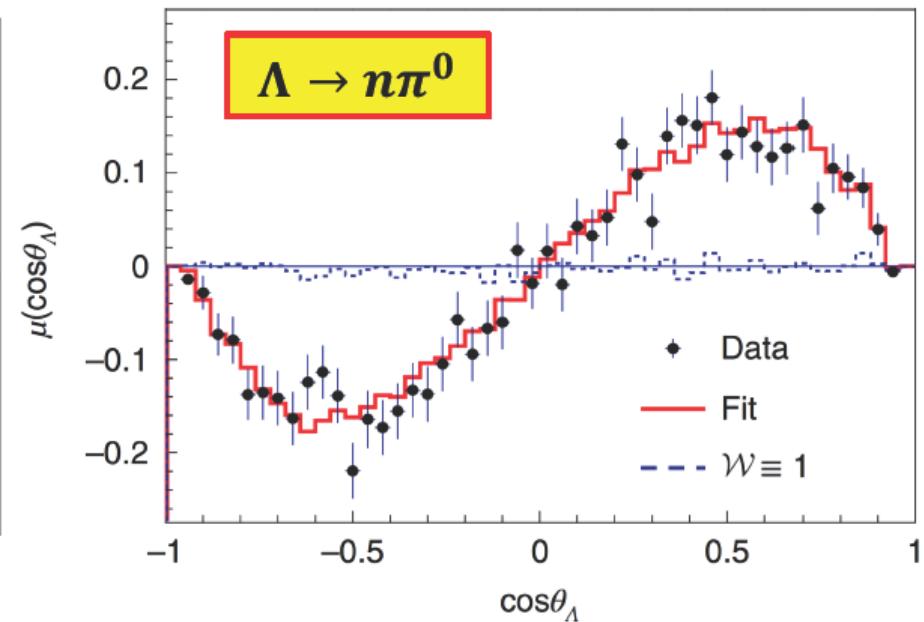
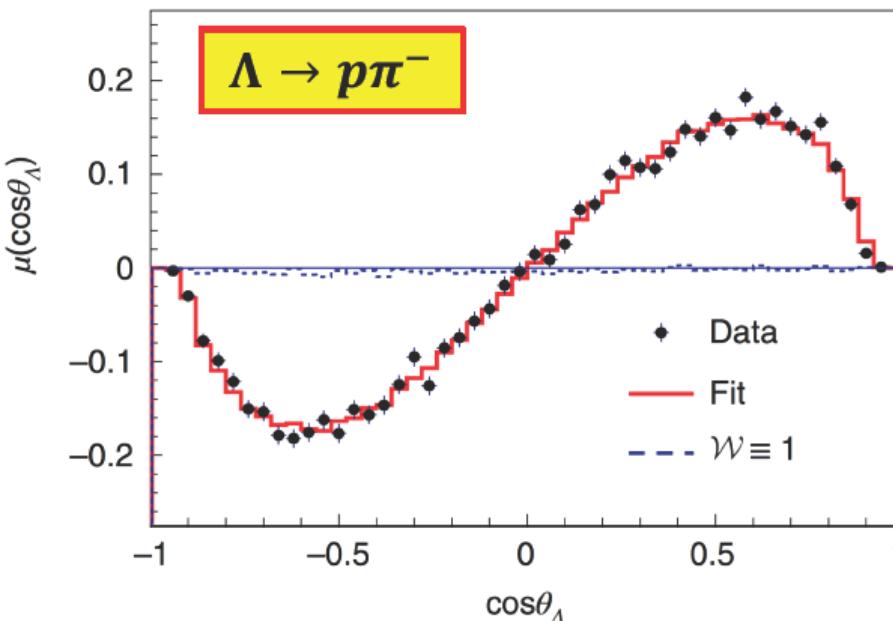
## Highlights!

$J/\psi \rightarrow \Lambda\bar{\Lambda}$

Published in Nature Physics, 15, 631 (2019)

$J/\psi$  data sample:  $1.31 \times 10^9$

- Observation of  $\Lambda / \bar{\Lambda}$  transverse polarization



- spin observable

$$\mu(\cos \theta_\Lambda) = \frac{m}{N} \sum_{i=1}^{N_k} (n_{1,y}^{(i)} - n_{2,y}^{(i)})$$

- Maximum  $\Lambda / \bar{\Lambda}$  polarization: ~25%
- Net polarization vanishing

# Observation of $\Lambda$ polarization and entanglement in $J/\psi \rightarrow \Lambda\bar{\Lambda}$



## Highlights!

Published in Nature Physics, 15, 631 (2019)

$J/\psi$  data sample:  $1.31 \times 10^9$

First observation of the  
 $\Lambda / \Lambda^-$  transverse  
polarization

Higher than PDG value 17%,  
deviation with sig.>  $5\sigma$

CP odd observable:  
 $A_{CP} = \frac{\alpha_\Lambda - \alpha_{\bar{\Lambda}}}{\alpha_\Lambda + \alpha_{\bar{\Lambda}}}$

Test selection rule  $\Delta I = \frac{1}{2}$

**Table 1 | Summary of the results**

Parameters	This work	Previous results
$\alpha_\psi$	$0.461 \pm 0.006 \pm 0.007$	$0.469 \pm 0.027$ (ref. <sup>14</sup> )
$\Delta\Phi$	$42.4 \pm 0.6 \pm 0.5^\circ$	-
$\alpha_-$	$0.750 \pm 0.009 \pm 0.004$	$0.642 \pm 0.013$ (ref. <sup>6</sup> )
$\alpha_+$	$-0.758 \pm 0.010 \pm 0.007$	$-0.71 \pm 0.08$ (ref. <sup>6</sup> )
$\bar{\alpha}_0$	$-0.692 \pm 0.016 \pm 0.006$	-
$A_{CP}$	$-0.006 \pm 0.012 \pm 0.007$	$0.006 \pm 0.021$ (ref. <sup>6</sup> )
$\bar{\alpha}_0/\alpha_+$	$0.913 \pm 0.028 \pm 0.012$	-

- Verification  $\alpha_\Lambda / \alpha_{\bar{\Lambda}}$ :  $J/\psi \rightarrow \Xi^- \bar{\Xi}^+ \rightarrow (\Lambda \pi^-)(\bar{\Lambda} \pi^+)$
- Transverse polarization allows to improve CP test precision over previous measurements.
- Standard model precision  $A_{CP} \sim 10^{-4}$ .

# Observation of $\Lambda$ polarization and entanglement in



## Highlights!

$J/\psi \rightarrow \Lambda\bar{\Lambda}$

$\alpha_-$  FOR  $\Lambda \rightarrow p\pi^-$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$0.750 \pm 0.009 \pm 0.004$	420k	ABLIKIM	2018AG	BES3 $J/\psi$ to $\Lambda\bar{\Lambda}$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
0.584 $\pm 0.046$	8500	ASTBURY	1975	SPEC
0.649 $\pm 0.023$	10325	CLELAND	1972	OSPK
0.67 $\pm 0.06$	3520	DAUBER	1969	HBC From $\Xi$ decay
0.645 $\pm 0.017$	10130	OVERSETH	1967	OSPK $\Lambda$ from $\pi^- p$
0.62 $\pm 0.07$	1156	CRONIN	1963	CNTR $\Lambda$ from $\pi^- p$

$\alpha_+$  FOR  $\bar{\Lambda} \rightarrow \bar{p}\pi^+$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
$-0.758 \pm 0.010 \pm 0.007$	420k	ABLIKIM	2018AG	BES3 $J/\psi$ to $\Lambda\bar{\Lambda}$
• • • We do not use the following data for averages, fits, limits, etc. • • •				
$-0.755 \pm 0.083 \pm 0.063$	$\approx 8.7k$	ABLIKIM	2010	BES $J/\psi \rightarrow \Lambda\bar{\Lambda}$
$-0.63 \pm 0.13$	770	TIXIER	1988	DM2 $J/\psi \rightarrow \Lambda\bar{\Lambda}$

# Where does the TP come from?

## ❑ From the $e^+/e^-$ beam ?

✗ No, BEPC beams unpolarized

## ❑ From the $e^+/e^-$ natural polarization when circulating in the BEPCII storage ring ?

✗ Sokolov-Ternov effects: 4.3 hrs @ $\psi'$  peak, but beam lifetime  $\sim 2.0$  hrs

## ❑ From the $J/\psi$ spin transfer ?

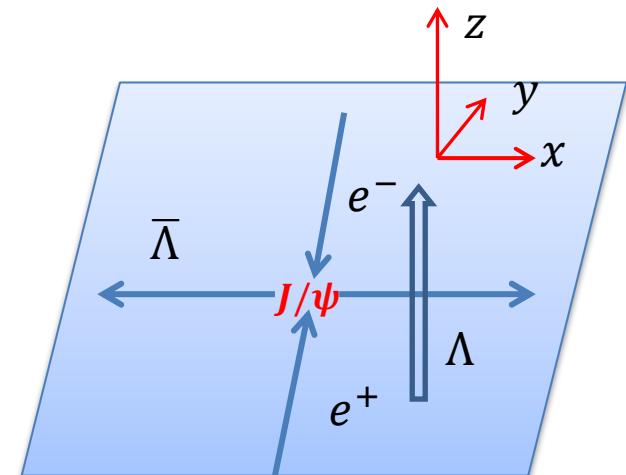
✓ Yes, it does from the  $J/\psi$  tensor polarization

$$J/\psi \text{ polarization: } \mathcal{P}_z = 0, \quad T_{zz} = \frac{1}{\sqrt{6}}$$

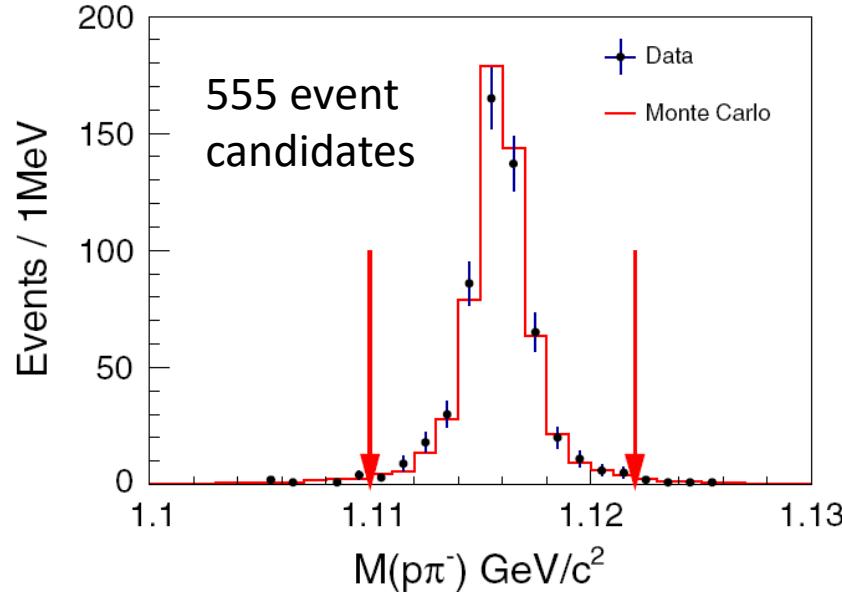
$\Lambda$  transverse polarization:

$$\mathcal{P}_y = \sqrt{6} \frac{T_{zz} \sin \theta \cos \theta \sin \Delta \sqrt{1 - \alpha_\psi^2}}{1 + \alpha_\psi \left[ \frac{1}{3} + \frac{1}{\sqrt{6}} T_{zz} (1 + 3 \cos 2\theta) \right]}$$

$\mathcal{P}_y$  manifest if  $\sin \Delta \neq 0$

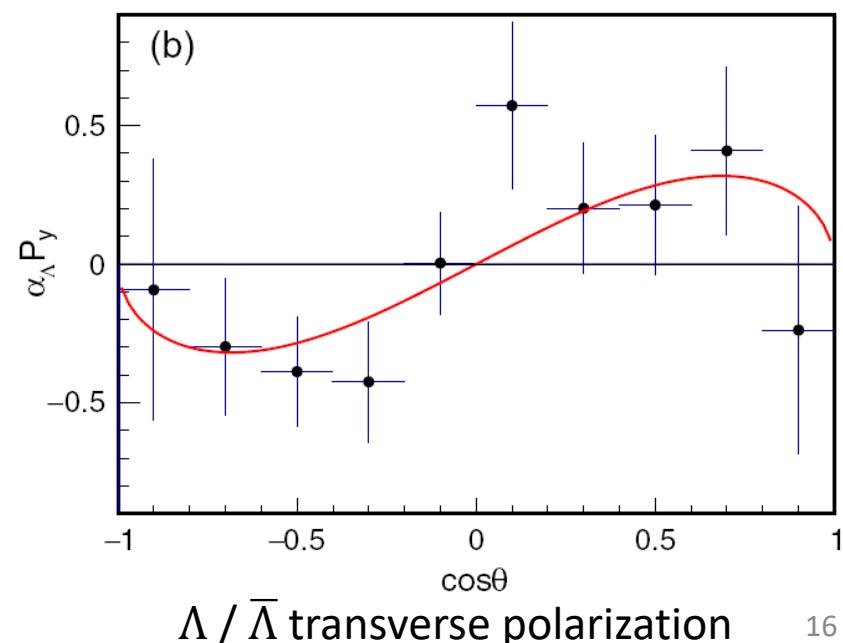
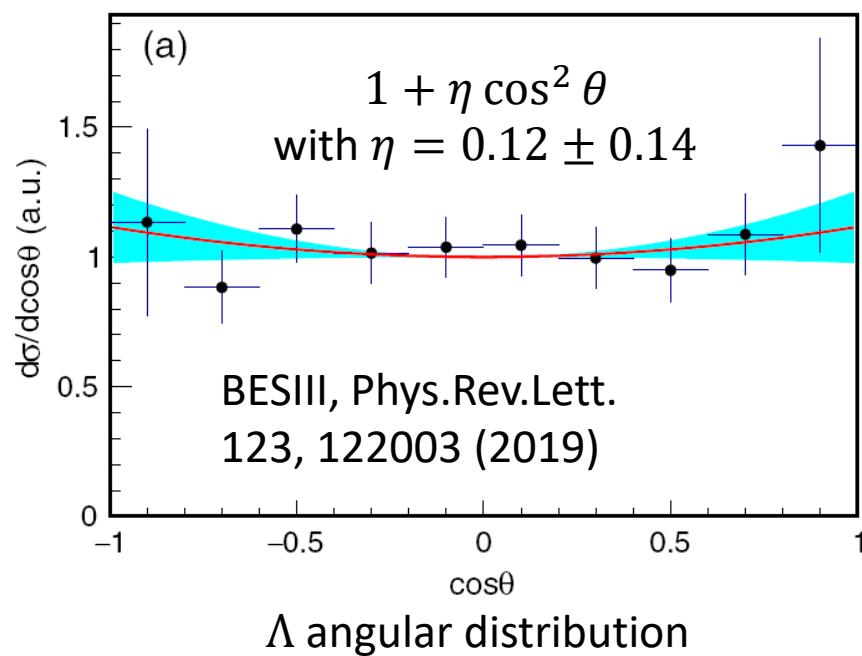


# $\Lambda / \bar{\Lambda}$ polarization in continuum production

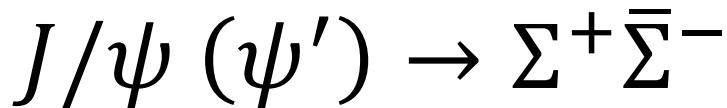


- ❑ data set: 2.396 GeV,  $L=66.9 \text{ pb}^{-1}$ 
$$e^+ e^- \rightarrow \gamma^* \rightarrow \Lambda \bar{\Lambda}$$
- ❑ 555 candidate events
- ❑  $\Delta\Phi = 37^\circ \pm 12^\circ (\text{stat}) \pm 6^\circ (\text{syst})$
- ❑ Maximum polarization degree is  $\sim 30\%$

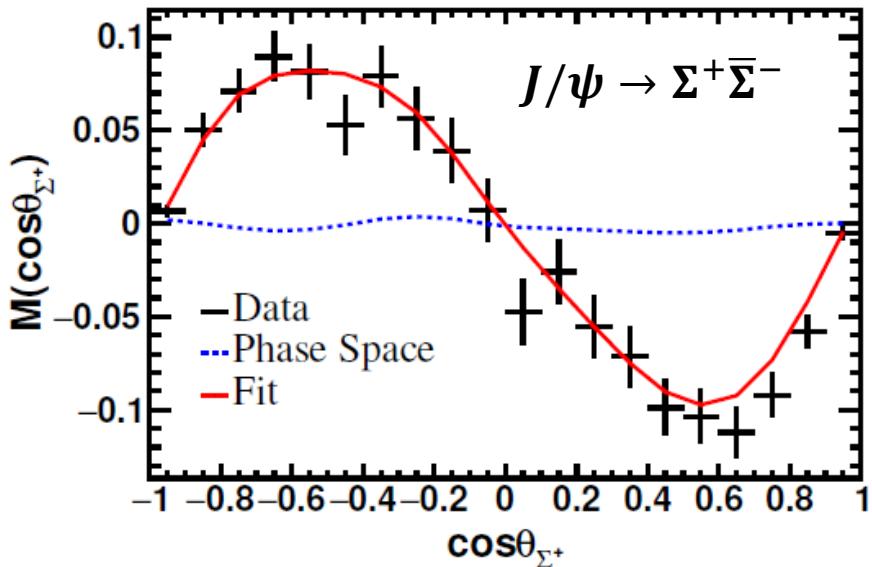
$$P_y = \frac{\sqrt{1-\eta^2} \sin \theta \cos \theta}{1+\eta \cos^2 \theta} \sin(\Delta\Phi).$$



# $\Sigma^+ / \bar{\Sigma}^-$ transverse polarization

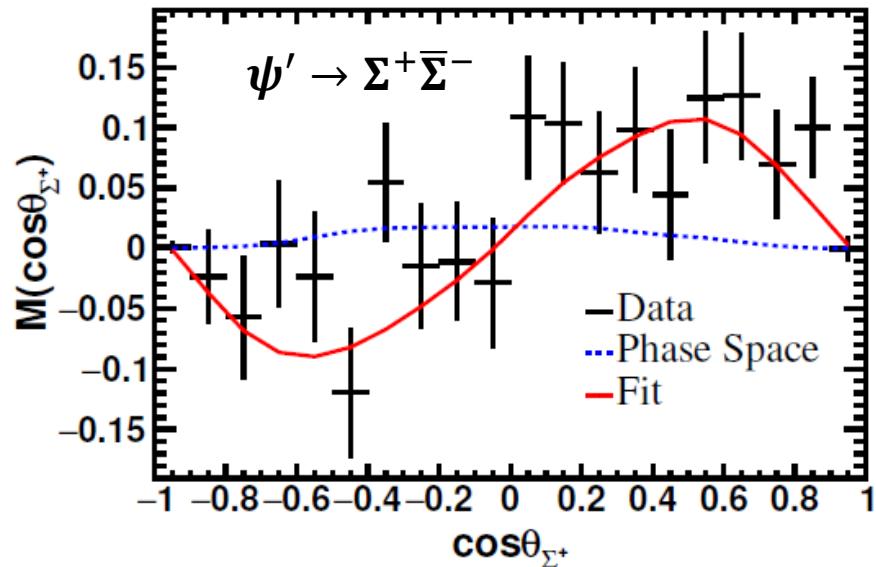


BESIII: arXiv: 2004.07701



$$P_y(\cos\theta_\Sigma) = \frac{\sqrt{1 - \alpha_\psi^2} \sin(\Delta\Phi) \cos\theta_\Sigma \sin\theta_\Sigma}{1 + \alpha_\psi \cos^2\theta_\Sigma}$$

$$\begin{aligned} \Sigma^+ &\rightarrow p\pi^0 \quad (\alpha_0 = -0.998 \pm 0.037 \pm 0.009) \\ \bar{\Sigma}^- &\rightarrow \bar{p}\pi^0 \quad (\bar{\alpha}_0 = 0.990 \pm 0.037 \pm 0.011) \end{aligned}$$



$$\Delta\Phi = (-15.5 \pm 0.7 \pm 0.5)^\circ \text{ for } J/\psi$$

$$\Delta\Phi = (+21.7 \pm 4.0 \pm 0.7)^\circ \text{ for } \psi'$$

$$A_{CP} = \frac{\alpha_0 + \bar{\alpha}_0}{\alpha_0 - \bar{\alpha}_0} = -0.004 \pm 0.037 \pm 0.010$$

**PDG:**

$$\alpha_0 \text{ FOR } \Sigma^+ \rightarrow p\pi^0 \quad -0.980^{+0.017}_{-0.015}$$

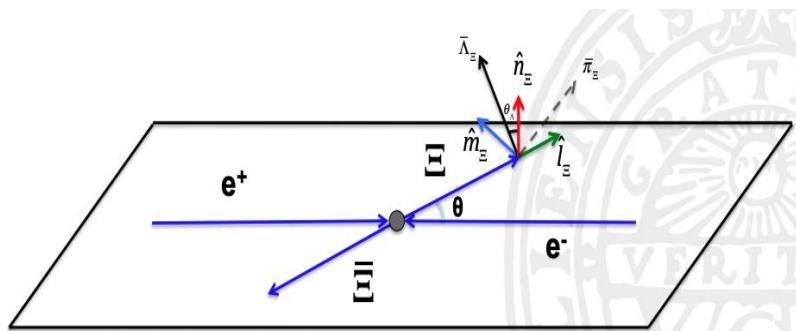
# CPV with $J/\psi \rightarrow \Xi^-(\Lambda\pi^-)\bar{\Xi}^+(\bar{\Lambda}\pi^+)$ +c.c.

缺点:

complicated topology: 9-dimensions

$\theta_\Xi, \theta_\Lambda, \phi_\Lambda, \theta_{\bar{\Lambda}}, \phi_{\bar{\Lambda}}, \theta_p, \phi_p, \theta_{\bar{p}}, \phi_{\bar{p}}$

72 terms, 8 parameters to determine



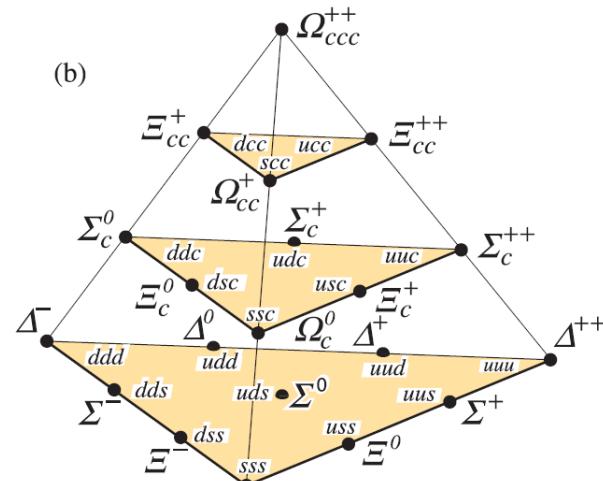
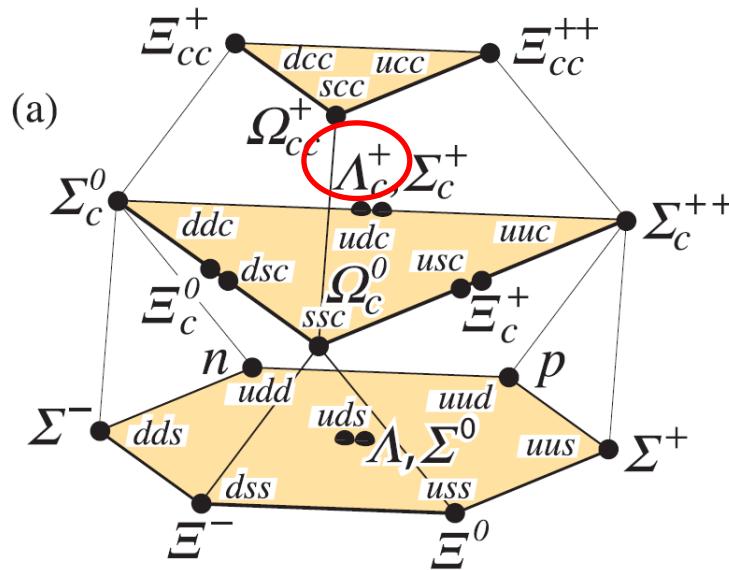
优点:

- $\Lambda\bar{\Lambda}$  polarizations are measureable via their parity-violating  $p\pi^-$  ( $p\pi^+$ ) decays;
- $\beta_-$  and  $\beta_0$  parameters can be determined.
- Preliminary results indicate that the  $\Xi$ s are even more polarized than the  $\Lambda$ s.

Low rate compared to  $\Lambda\bar{\Lambda}$

1.3B  $J/\psi$ : 420K  $\Lambda(p\pi^-)\bar{\Lambda}(\bar{p}\pi^+)$  evts  
61K  $\Xi(\Lambda\pi^-) \Xi(\Lambda\pi^+)$

# $\Lambda_c$ spin and decay asymmetry parameter



$\Lambda_c^+$

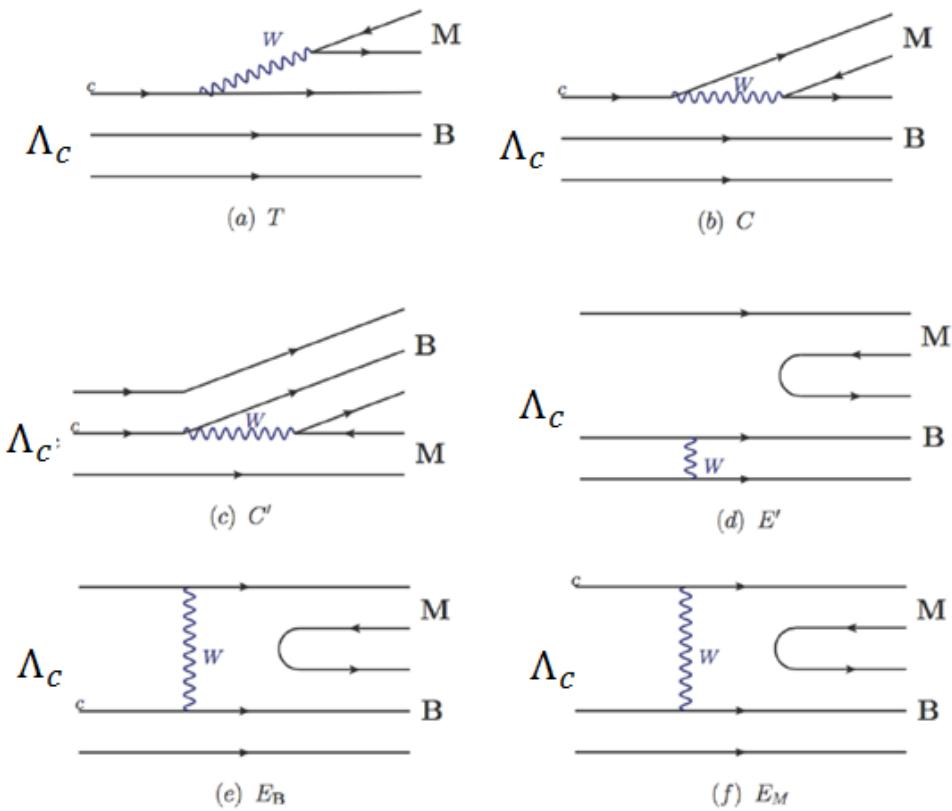
$I(J^P) = 0(\frac{1}{2}^+)$  Status: \* \* \* \*

The parity of the  $\Lambda_c^+$  is defined to be positive (as are the parities of the proton, neutron, and  $\Lambda$ ). The quark content is  $u d c$ . Results of an analysis of  $p K^- \pi^+$  decays (JEZABEK 92) are consistent with  $J = 1/2$ . Nobody doubts that the spin is indeed  $1/2$ .

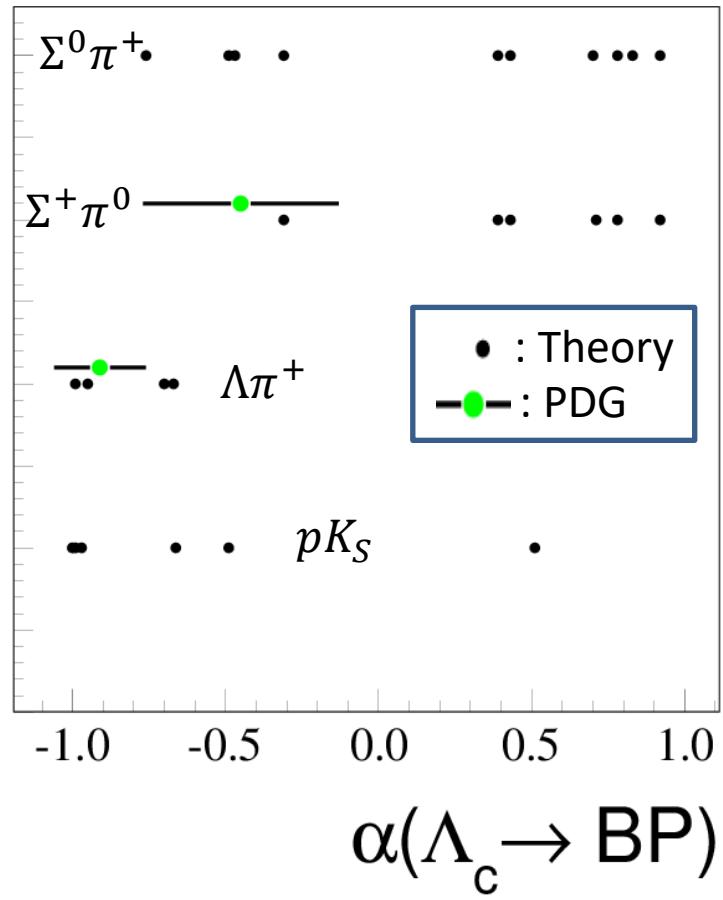
We have omitted some results that have been superseded by later experiments. The omitted results may be found in earlier editions.

# Predictions on $\Lambda_c$ asymmetry parameters

- $W$ -exchange complicate  $\Lambda_c \rightarrow B M$

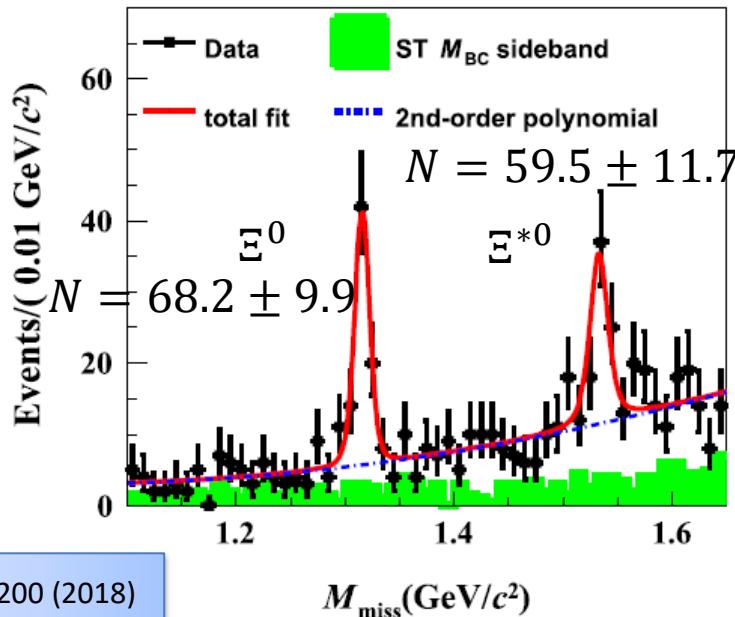
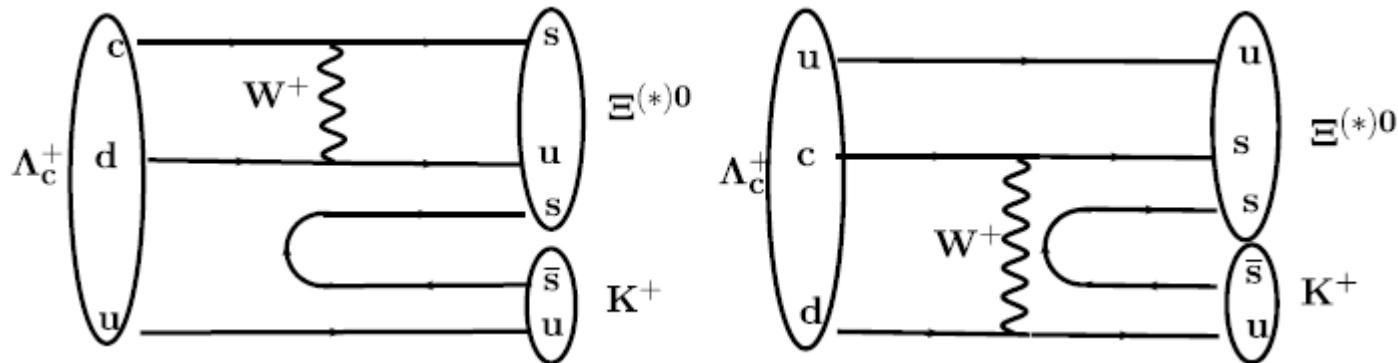


$\Lambda_c^+ \rightarrow BP$  decay asymmetry



# Predictions on $\Lambda_c^+$ asymmetry parameters

- $W$ -exchange decay:  $\Lambda_c^+ \rightarrow \Xi^{(*)0} K^+$

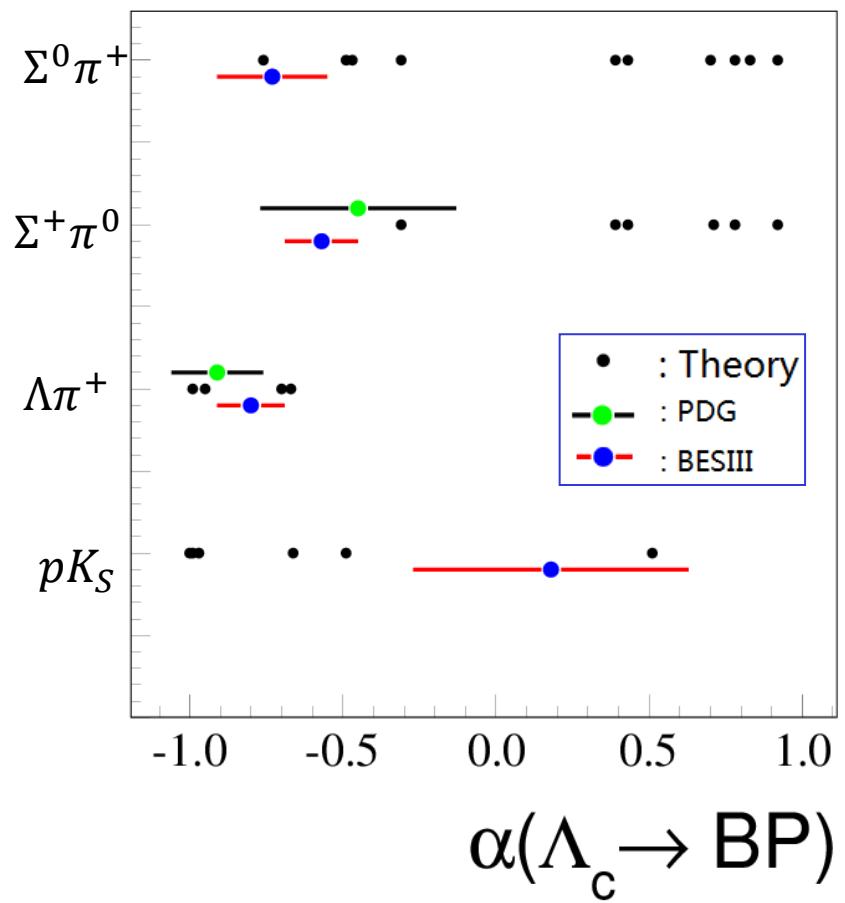
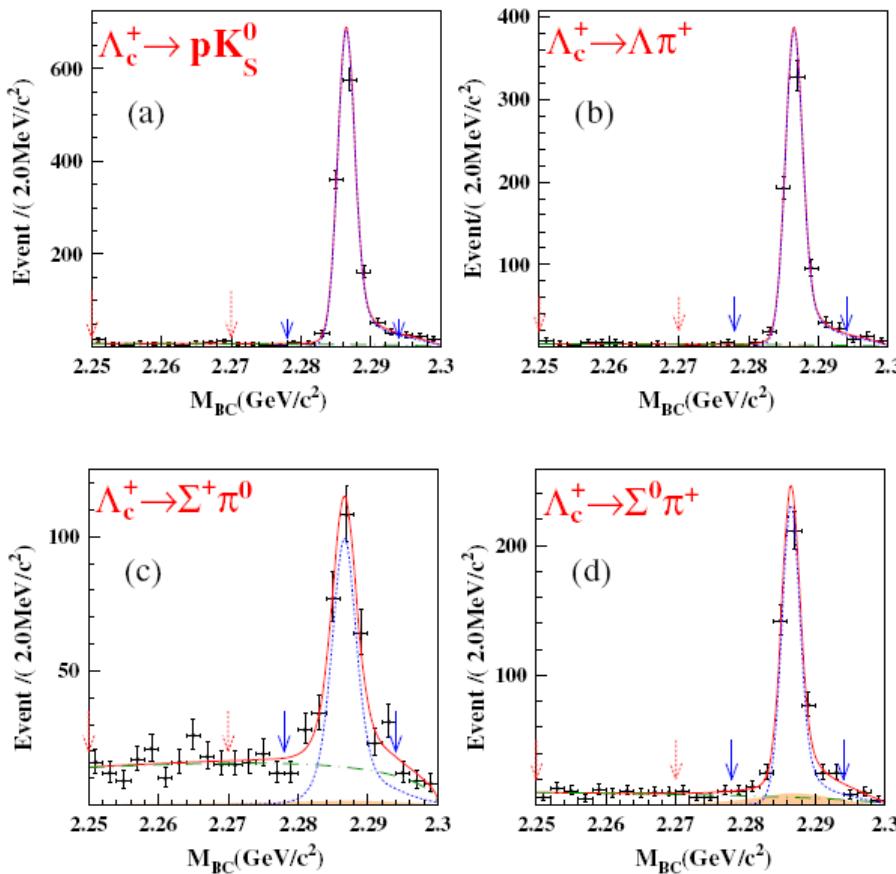


$$\mathcal{B}(\Lambda_c^+ \rightarrow \Xi^0 K^+) = (5.90 \pm 0.86 \pm 0.39) \times 10^{-3}$$
$$\mathcal{B}(\Lambda_c^+ \rightarrow \Xi^{*0} K^+) = (5.02 \pm 0.99 \pm 0.31) \times 10^{-3}$$

Compared to  
 $\mathcal{B}(\Lambda_c^+ \rightarrow \Lambda \pi^+) = (1.29 \pm 0.07)\%$

# $\Lambda_c$ decay asymmetry parameter

- $L = 567\text{pb}^{-1}$  @  $\sqrt{s} = 4.6 \text{ GeV}$  with single tag



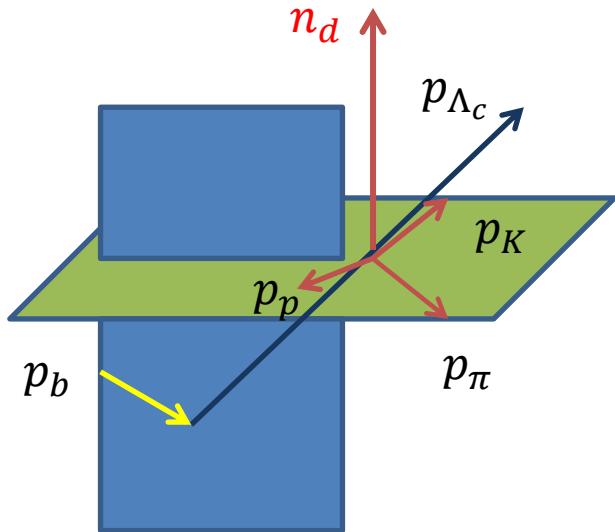
Phys. Rev. D 100, 072004 (2019)

- Significance for  $\Lambda_c$  transverse polarization:  $2.1\sigma$

# $\Lambda_c$ spin

- NA2 experiment at CERN-SPS
- $\pi \text{ Cu} \rightarrow \Lambda_c^+ \bar{D} X$
- 160 events for  $\Lambda_c \rightarrow p K^- \pi^+$

Phys.Lett., B286, 175 (1992)



- Angular distribution of three – body decay

$$I(\theta, \phi) = \frac{2J+1}{4\pi}$$

$$\times \sum_{M,M'} \varrho_{MM'}^J \sum_{\mu} f_{\mu}^J D_{M\mu}^{J*}(\phi, \theta, 0) D_{M'\mu}^J(\phi, \theta, 0) ,$$

$$I(\theta) = \frac{1}{2} \left( 1 + \sum_{l=1}^{2J} d_l P_l(\cos \theta) \right),$$

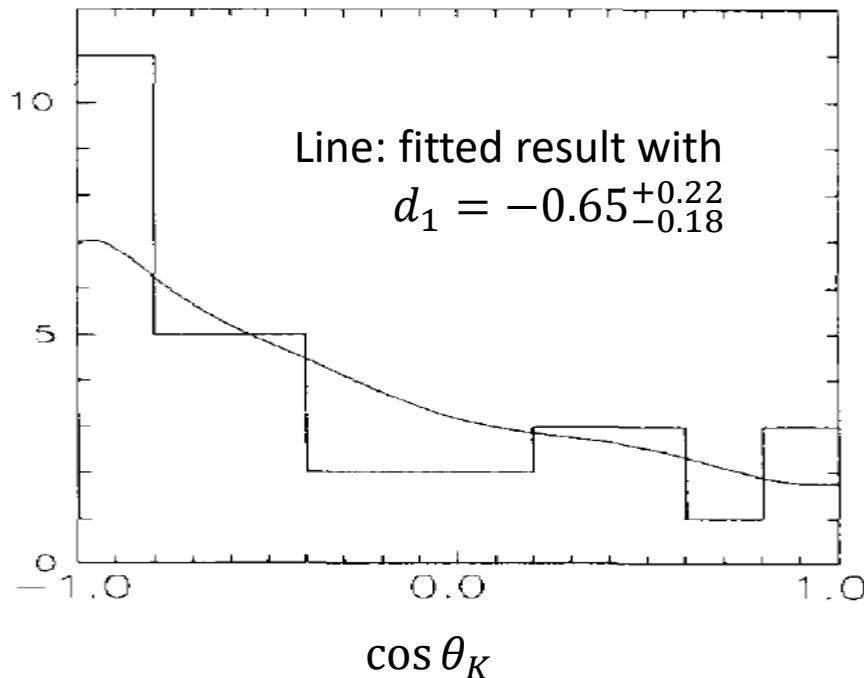
Simplified as:

$$d_l \equiv p_l a_l ,$$

$$p_l = \sqrt{2l+1} \sum_m \varrho_{MM}^J \langle J, M; l, 0 | J, M \rangle ,$$

$$a_l = \sqrt{2l+1} \sum_{\mu} f_{\mu}^J \langle J, \mu; l, 0 | J, \mu \rangle .$$

# $\Lambda_c$ spin



Assume  $J = 1/2$ , then

$$p_1 = \rho_{1/2,1/2}^{1/2} - \rho_{-1/2,-1/2}^{-1/2}$$

$$a_1 = f_{1/2}^{1/2} - f_{-1/2}^{1/2}$$

They concluded:

- Results consistent with assumption  $J = 1/2$
- Unable to establish  $J = 1/2$  due to low statistics

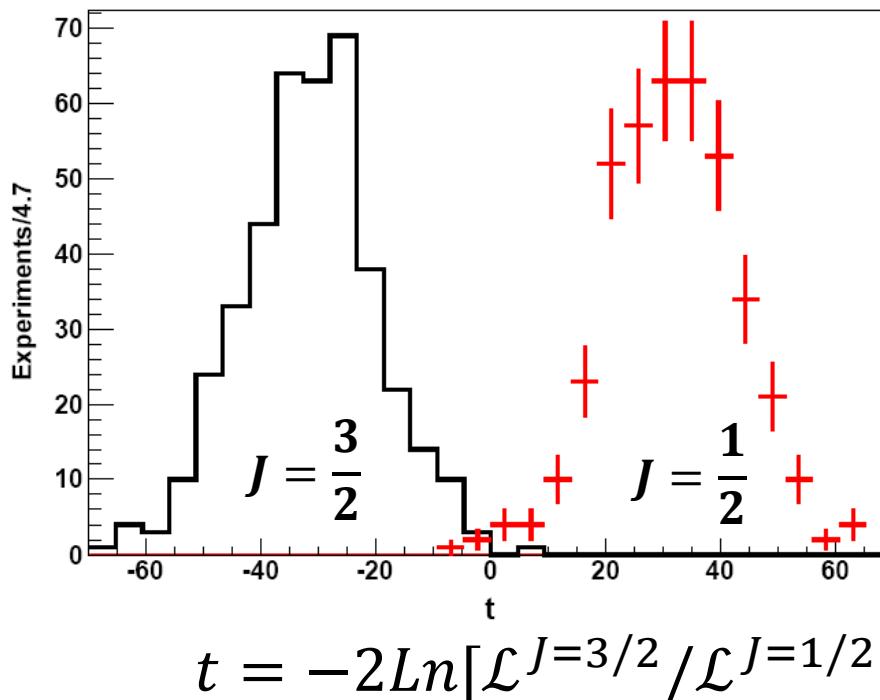
Remarks:

- decay parameter should be model momentum dependent
- resonance structure in 3-body decay should be considered.
- only  $\Lambda_c$  longitudinal polarization is used (see  $\rho_{M,M}$ ), and transverse part is missed

# Determination of $\Lambda_c$ spin

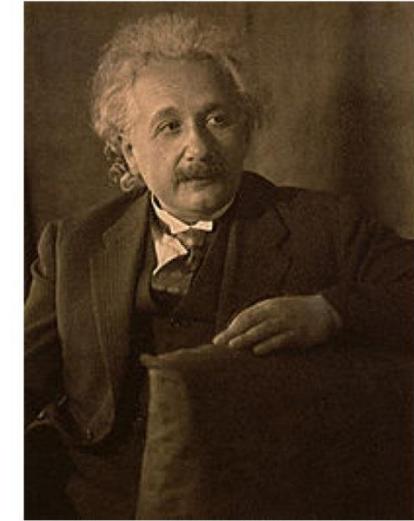
- Toy MC study: ( $L = 567 pb^{-1}$  @  $\sqrt{s} = 4.6$  GeV)

Decays	Br [9]	$\alpha_{[BM]}^{[\Lambda_c^+]}$	$N^{\text{obs}}$ [1]
$\Lambda_c^+ \rightarrow p K_S^0$	$(1.58 \pm 0.08)\%$	-0.75	1243
$\Lambda_c^+ \rightarrow \Lambda \pi^+$	$(1.29 \pm 0.07)\%$	$-0.91 \pm 0.15$ [9]	706
$\Lambda_c^+ \rightarrow \Sigma^0 \pi^+$	$(1.28 \pm 0.07)\%$	$-0.45 \pm 0.32$ [9]	522



# Einstein-Podolsky-Rosen (EPR) paradox

- EPR paradox<sup>[1]</sup> (1935) initiated an intense discussion on the philosophical foundations of quantum mechanics.
- Physics quantity corresponding to “element of reality” can be predicted without any uncertainty.
- A given example: two particle system with sum of momentum ( $p_1 + p_2$ ) and position difference ( $x_1 - x_2$ )
- They concluded: the premises of “element of reality” is wrong or Q.M. is incomplete.

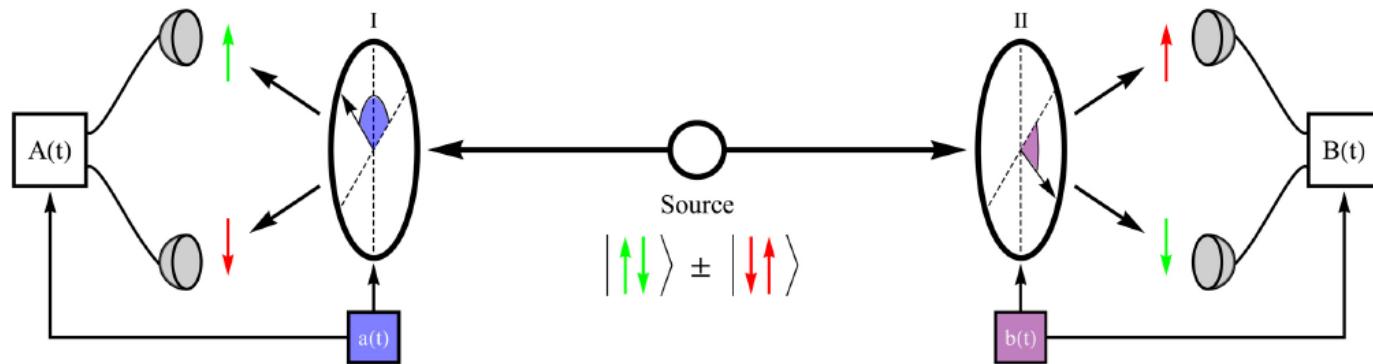


Albert Einstein

Ref.[1] . Einstein, A; B Podolsky and N. Rosen, Phys. Rev. 47, 10 (1935), “Can Quantum-Mechanical description of Physical Reality be considered Complete?”

# Local hidden variable theory (LHVT)

Bohm  
Spin  
setup



- Local hidden variable theory:  
A spin singlet formed by two spin-1/2 particles  
 $|\psi_0\rangle = \frac{1}{\sqrt{2}}(|\uparrow\rangle|\downarrow\rangle - |\downarrow\rangle|\uparrow\rangle)$ . In Q.M.  
 $A: \vec{\sigma}_1 \cdot \vec{a} \quad B: \vec{\sigma}_2 \cdot \vec{a}$
- In LHVT,  $A(\vec{a}, \vec{\lambda}) = \text{sign}(\vec{a} \cdot \vec{\lambda})$ ,  $B(\vec{b}, \vec{\lambda}) = \text{sign}(\vec{a} \cdot \vec{\lambda})$ , probability describing correlation between them  
 $P(\vec{a}, \vec{b}) = \int d\lambda \rho(\lambda) A(\vec{a}, \vec{\lambda}) B(\vec{b}, \vec{\lambda})$

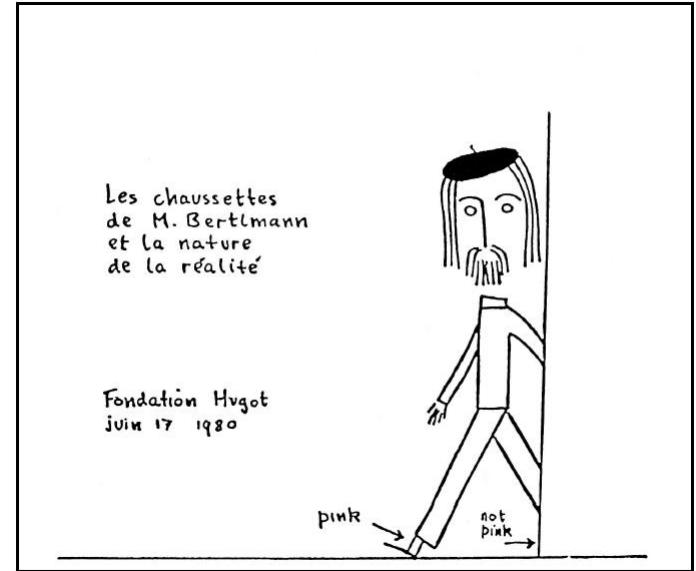
# BELL inequality(BI)

John S. Bell



John Stewart Bell, CERN, 1973

Cartoon of  
Berlmann's socks  
*by Bell (1980)*



- BELL inequality  
(J.S. Bell, Physics, 1, 195(1964) ):
$$1 + P(\vec{b}, \vec{c}) \geq |P(\vec{a}, \vec{b}) - P(\vec{a}, \vec{c})|$$
- BI developed many variants, and tested in optics and high energy physics.

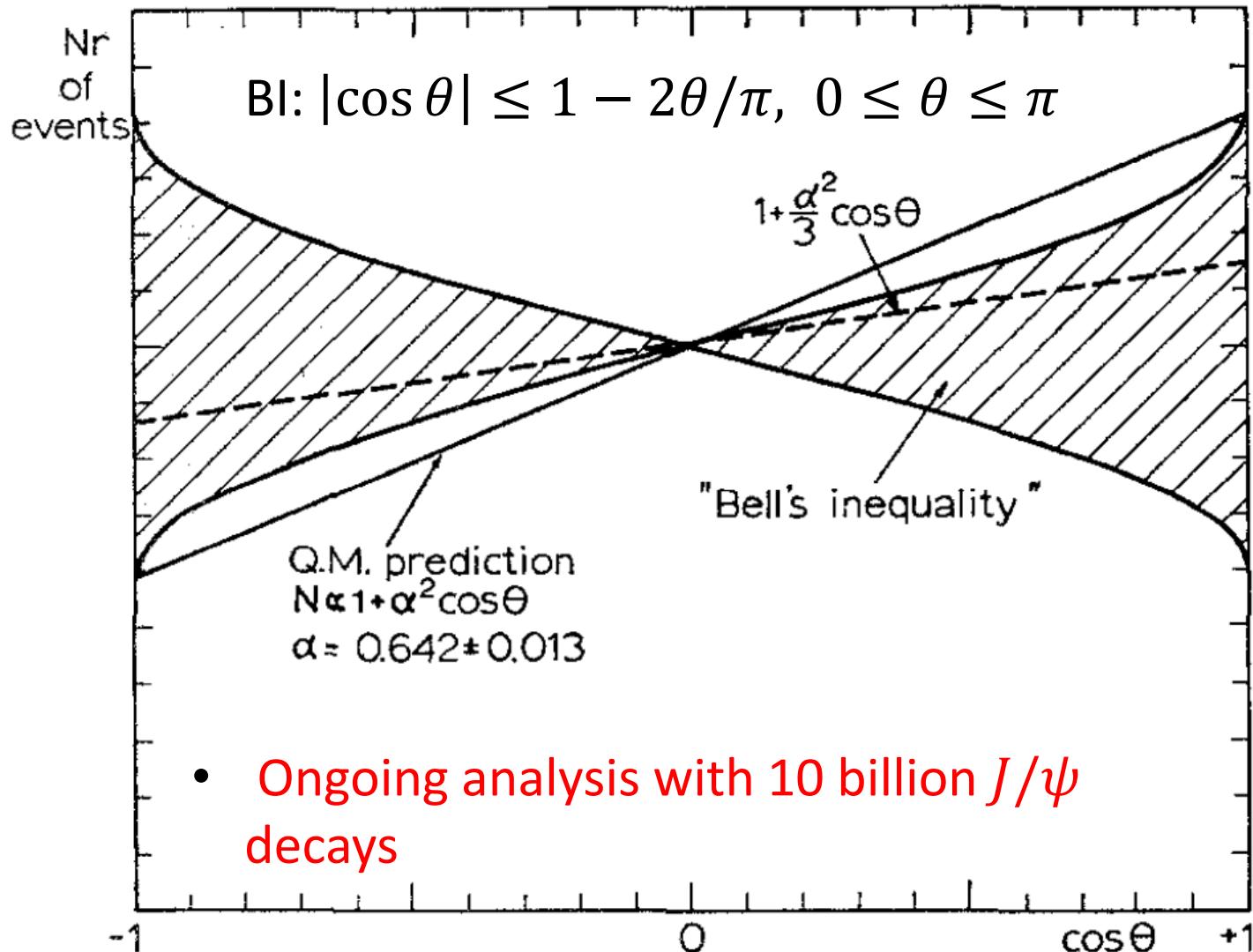
# Test BI in charmonium decays

- $\eta_c \rightarrow \Lambda\bar{\Lambda} \rightarrow p\bar{p}\pi^+\pi^-$ 
  - Q.M.:  $\Lambda: M_A = (4\pi)^{-1/2}(S + P\vec{\sigma}_A \cdot \vec{A}),$   
 $\bar{\Lambda}: M_B = (4\pi)^{-1/2}(S + P\vec{\sigma}_B \cdot \vec{B});$   
 $I(\vec{A}, \vec{B}) = \left( \frac{|S|^2 + |P|^2}{4\pi} \right)^2 (1 + \alpha_\Lambda^2 \vec{A} \cdot \vec{B})$
  - In LHVT

$$I(\vec{A}, \vec{B}) = \left( \frac{|S|^2 + |P|^2}{4\pi} \right)^2 \frac{1}{4\pi} \int (1 + \alpha_\Lambda X \cdot A)(1 + \alpha_\Lambda X \cdot B) d\Omega$$
$$= \left( \frac{|S|^2 + |P|^2}{4\pi} \right)^2 (1 + \frac{1}{3} \alpha_\Lambda^2 \vec{A} \cdot \vec{B})$$

# Test BI in $\eta_c \rightarrow \Lambda\bar{\Lambda}$ decays

N.A. Tornqvist, Foundation of Physics, 11, 171(1981)



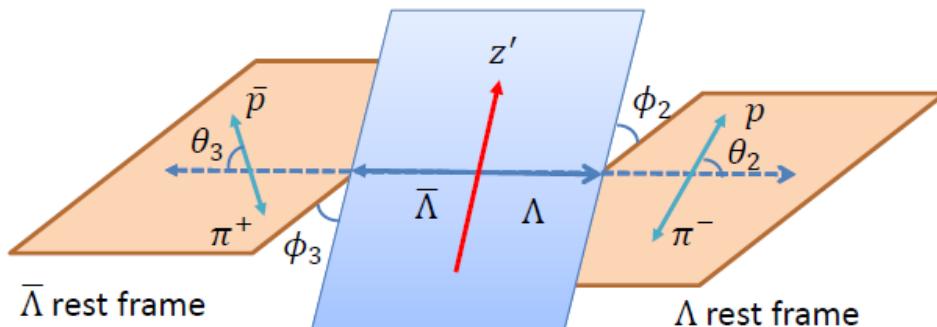
# Test CH inequality in $\eta_c \rightarrow \Lambda\bar{\Lambda}$ decays

Generalized CH inequality: C. Qian, et al. arXiv:2002.04283

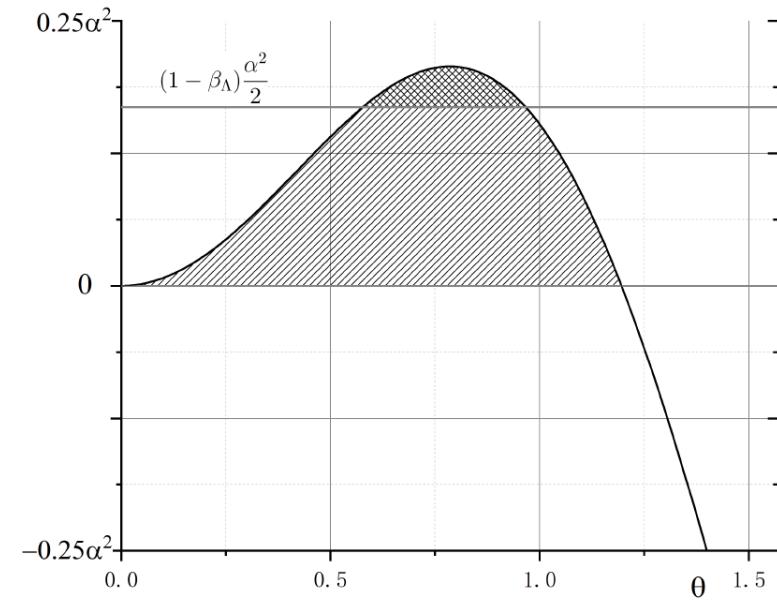
$$P(\vec{n}_1, \vec{n}_2) - P(\vec{n}_1, \vec{n}'_2) + P(\vec{n}'_1, \vec{n}_2) + P(\vec{n}'_1, \vec{n}'_2)$$

$$-(a_2 + b_2)P(\vec{n}'_1) - (a_1 + b_1)P(\vec{n}_2) + a_1 b_2 + b_1 a_2 \leq 0 .$$

$$P(\vec{n}_1, \vec{n}_2) = \int_{\Gamma} p_m(\lambda, \vec{n}_1) p_m(\lambda, \vec{n}_2) \rho(\lambda) d\lambda$$



$$\alpha^2 \left[ \frac{3 \cos \theta - \cos(3\theta)}{4} - \frac{1}{2} \right] \leq 0 .$$



# Ongoing polarization analyses at BESIII

- $J/\psi, \psi' \rightarrow \Sigma^-\bar{\Sigma}^+$
- $J/\psi \rightarrow \Xi^-\bar{\Xi}^+$  ( $\Xi, \Lambda$  asymmetry par., CP test)
- $\psi' \rightarrow \Omega^-\bar{\Omega}^+$  (polarization analysis )
- $\eta_c \rightarrow \Lambda\bar{\Lambda}$  (EPR test )
- $e^+e^- \rightarrow \Lambda_c^+\bar{\Lambda}_c^-$ 
  - $\Lambda_c^+ \rightarrow \Lambda\pi^+, \Sigma^+\pi^0, \Sigma^0\pi^+, K_S p$  (update)
  - $\Lambda_c$  spin
- future charmed baryon program
  - Update Ecms up to 4.9 GeV
  - Taken data around  $\Lambda_c\bar{\Lambda}_c$  threshold this year

# Summary

- $\Lambda/\bar{\Lambda}$  transverse polarization significantly observed at BESIII in  $J/\psi$  or continuum processes
- BESIII 10 billion  $J/\psi$  data provides us chances to access hyperon physics.
- Extension study to charmed hyperon are ongoing.
- Polarized beam in the future super-tau charm facility (STCF) help to improve the precision.

敬请指教！

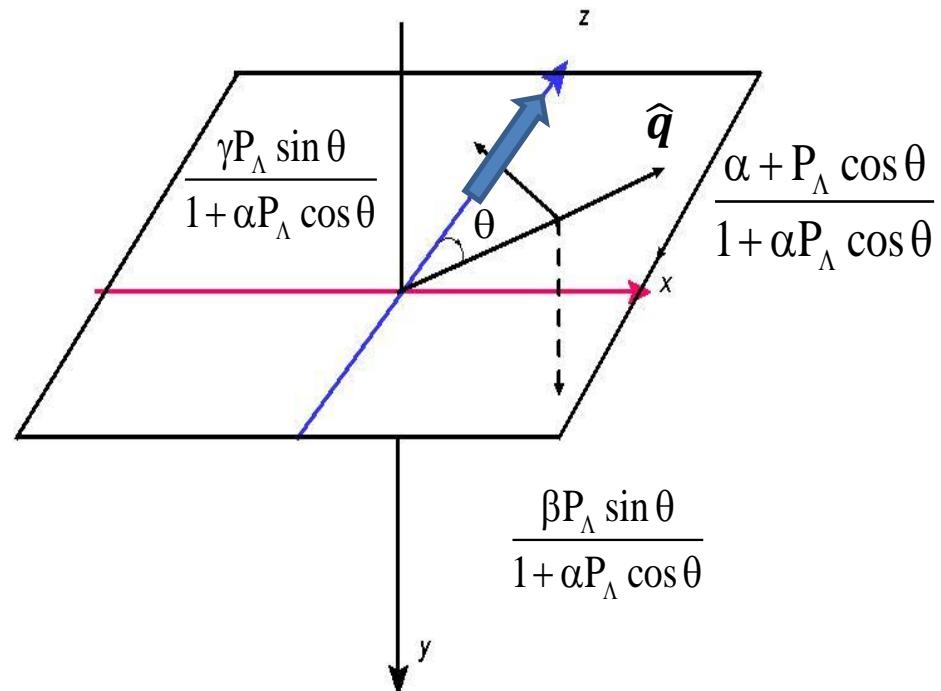
# backup

# Proton polarization from $\Lambda \rightarrow p\pi^-$

$$\vec{P}_p = \frac{(\alpha + \vec{P}_\Lambda \cdot \hat{q})\hat{q} + \beta (\vec{P}_\Lambda \times \hat{q}) + \gamma \hat{q} \times (\vec{P}_\Lambda \times \hat{q})}{(1 + \alpha \vec{P}_\Lambda \cdot \hat{q})}$$

- If  $P_\Lambda = 0$  then  $P_p = \alpha P_\Lambda \cdot q$
- T-odd transverse polarization  
 $\beta \neq 0$
- If CP is conserved :
 
$$\alpha = -\bar{\alpha}, \beta = -\bar{\beta},$$

$$\gamma = \bar{\gamma} \text{ and } \Gamma = \bar{\Gamma}$$



# Previous Measurements

## 2018 PDG list

$\alpha_-$  FOR  $\Lambda \rightarrow p\pi^-$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>0.642 \pm 0.013</math></b>	<b>OUR AVERAGE</b>			
$0.584 \pm 0.046$	8500	ASTBURY	1975	SPEC
$0.649 \pm 0.023$	10325	CLELAND	1972	OSPK
$0.67 \pm 0.06$	3520	DAUBER	1969	HBC
$0.645 \pm 0.017$	10130	OVERSETH	1967	OSPK
$0.62 \pm 0.07$	1156	CRONIN	1963	$\Lambda$ from $\pi^- p$
				$\Lambda$ from $\pi^- p$

$\alpha_+$  FOR  $\bar{\Lambda} \rightarrow \bar{p}\pi^+$

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b><math>-0.71 \pm 0.08</math></b>	<b>OUR AVERAGE</b>			
$-0.755 \pm 0.083 \pm 0.063$	$\approx 8.7k$	ABLIKIM	2010	BES
$-0.63 \pm 0.13$	770	TIXIER	1988	DM2
				$J/\psi \rightarrow \Lambda\bar{\Lambda}$
				$J/\psi \rightarrow \Lambda\bar{\Lambda}$

# Previous Measurements

2018 PDG list

## $(\alpha + \bar{\alpha})/(\alpha - \bar{\alpha})$ in $\Lambda \rightarrow p\pi^-$ , $\bar{\Lambda} \rightarrow \bar{p}\pi^+$

Zero if  $CP$  is conserved;  $\alpha_-$  and  $\alpha_+$  are the asymmetry parameters for  $\Lambda \rightarrow p\pi^-$  and  $\bar{\Lambda} \rightarrow \bar{p}\pi^+$  decay. See also the  $\Xi^-$  for a similar test involving the decay chain  $\Xi^- \rightarrow \Lambda\pi^-$ ,  $\Lambda \rightarrow p\pi^-$  and the corresponding antiparticle chain.

VALUE	EVTS	DOCUMENT ID	TECN	COMMENT
<b>0.006 ± 0.021 OUR AVERAGE</b>				
-0.081 ± 0.055 ± 0.059	≈ 8.7k	ABLIKIM	10	BES $J/\psi \rightarrow \Lambda\bar{\Lambda}$
+0.013 ± 0.022	96k	BARNES	96	CNTR LEAR $\bar{p}p \rightarrow \bar{\Lambda}\Lambda$
+0.01 ± 0.10	770	TIXIER	88	DM2 $J/\psi \rightarrow \Lambda\bar{\Lambda}$
-0.02 ± 0.14	10k	<sup>1</sup> CHAUVAT	85	CNTR $p\bar{p}, \bar{p}p$ ISR
• • • We do not use the following data for averages, fits, limits, etc. • • •				
-0.07 ± 0.09	4063	BARNES	87	CNTR See BARNES 96

<sup>1</sup> CHAUVAT 85 actually gives  $\alpha_+(\bar{\Lambda})/\alpha_-(\Lambda) = -1.04 \pm 0.29$ . Assumes polarization is same in  $\bar{p}p \rightarrow \bar{\Lambda}X$  and  $p\bar{p} \rightarrow \Lambda X$ . Tests of this assumption, based on  $C$ -invariance and fragmentation, are satisfied by the data.