

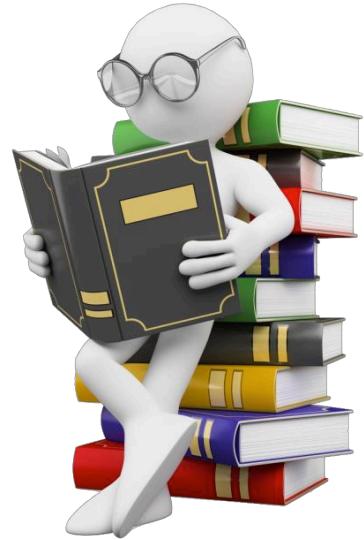


Study of the charmed-antistrange and anticharmed-strange meson pairs at Belle and status of Belle II

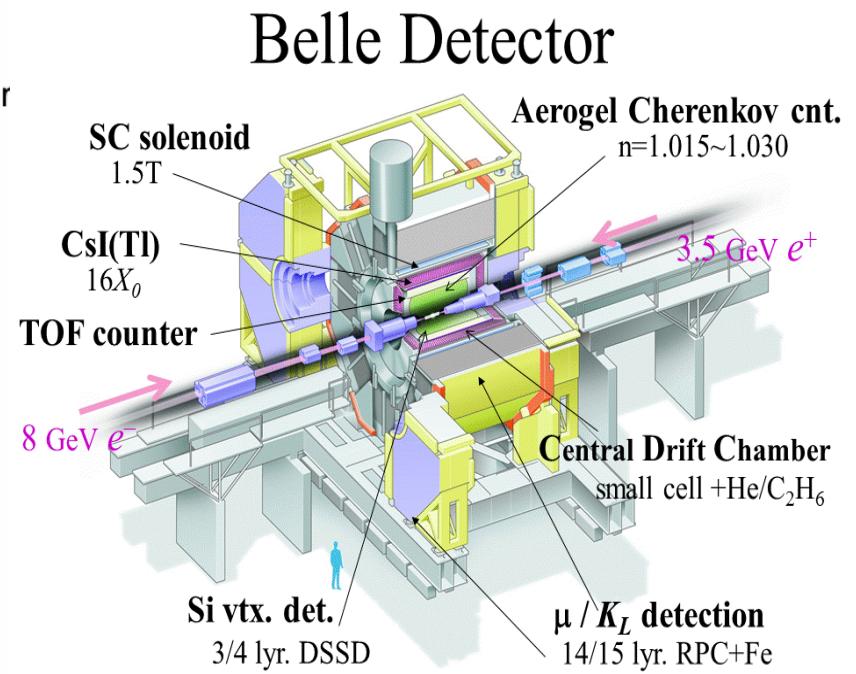
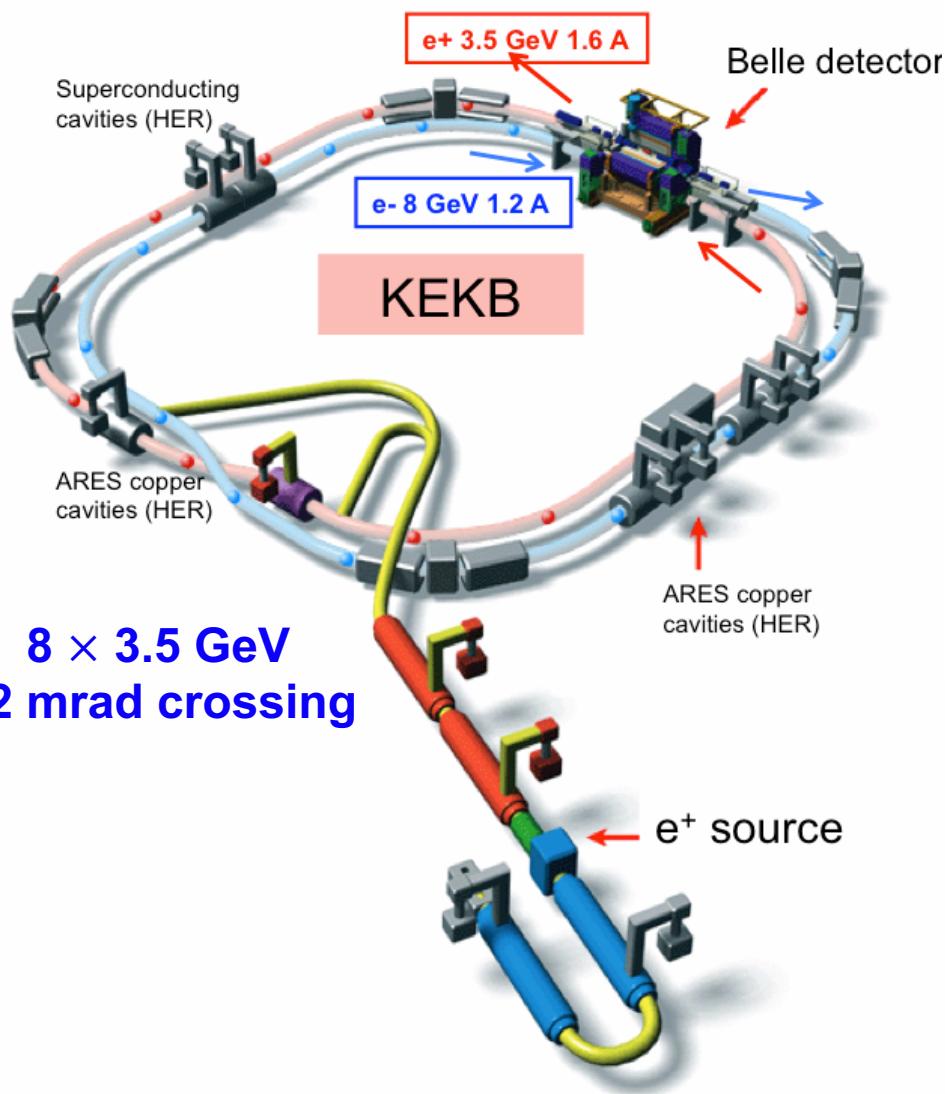
贾森
复旦大学
2020年4月24日
强子物理新发展研讨会

Outline

- *Belle experiments*
- $e^+ e^- \rightarrow D_s^+ D_{s1}(2536)^-$
- $e^+ e^- \rightarrow D_s^+ D_{s2}^*(2573)^-$
- *Summary and discussion*
- *Status and Prospects of Belle II*

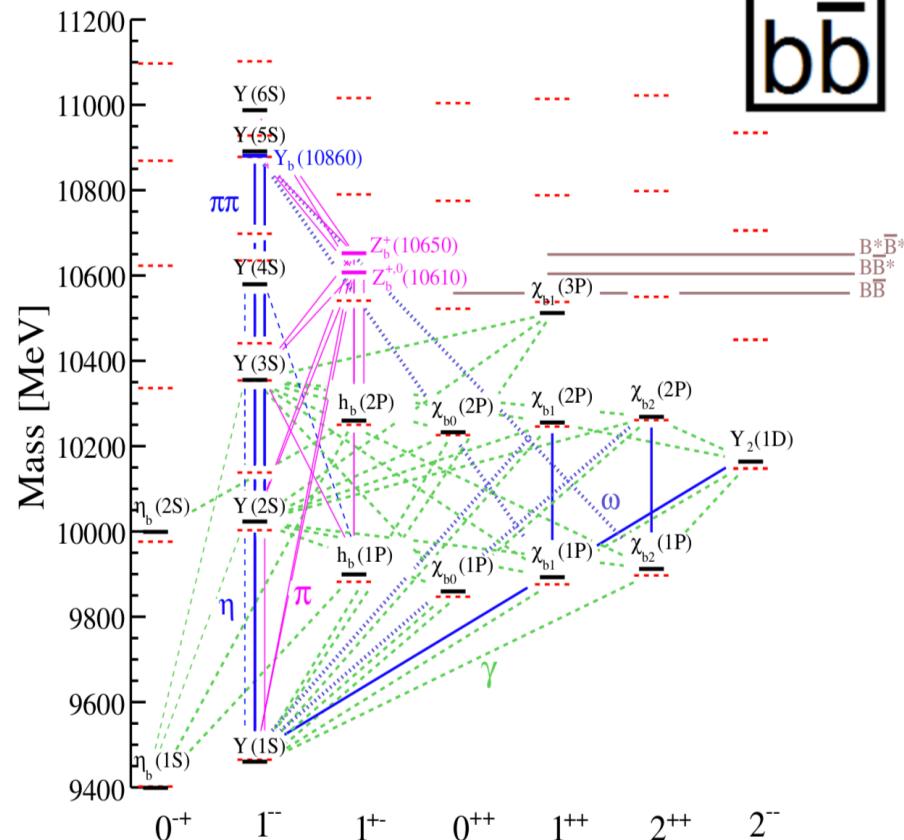
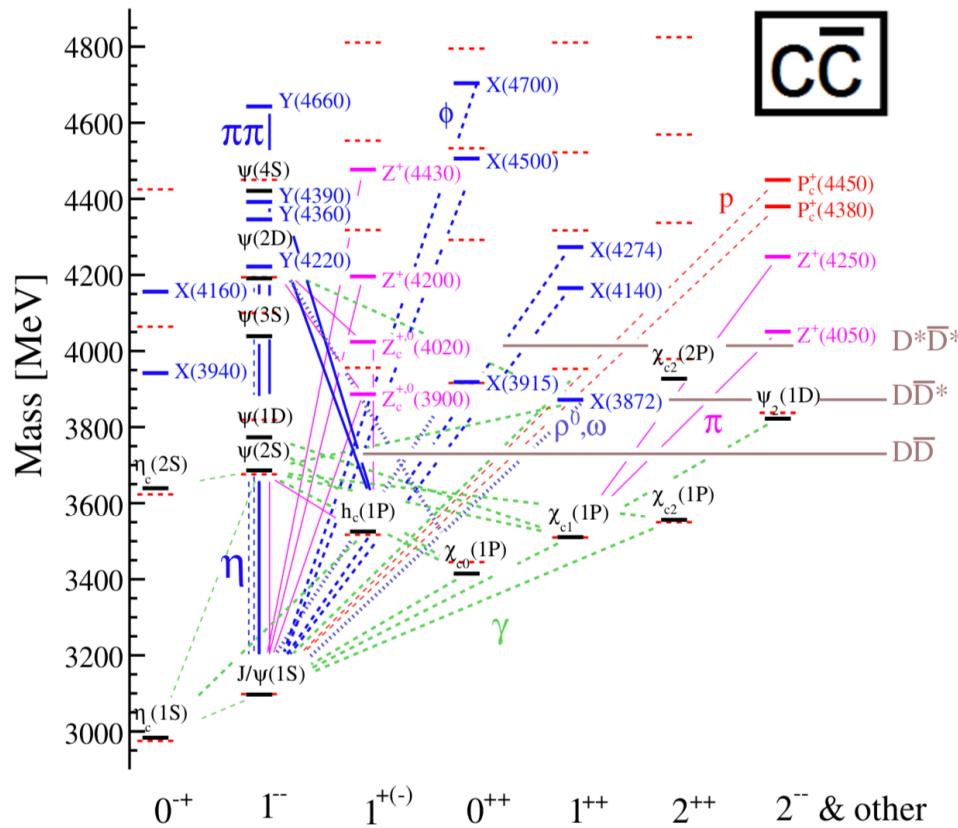


Belle experiment and data samples



Data taking: 1999 – 2010
On/off/Scan $\Upsilon(nS)$ peaks
Total luminosity: 980 fb^{-1}
772M $B\bar{B}$ events @ $\Upsilon(4S)$

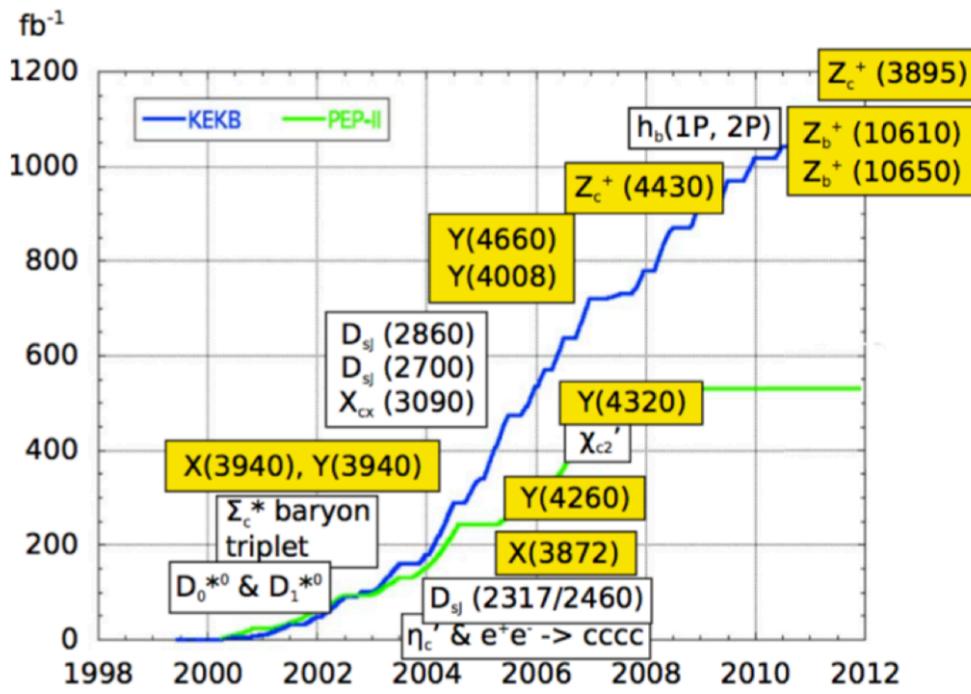
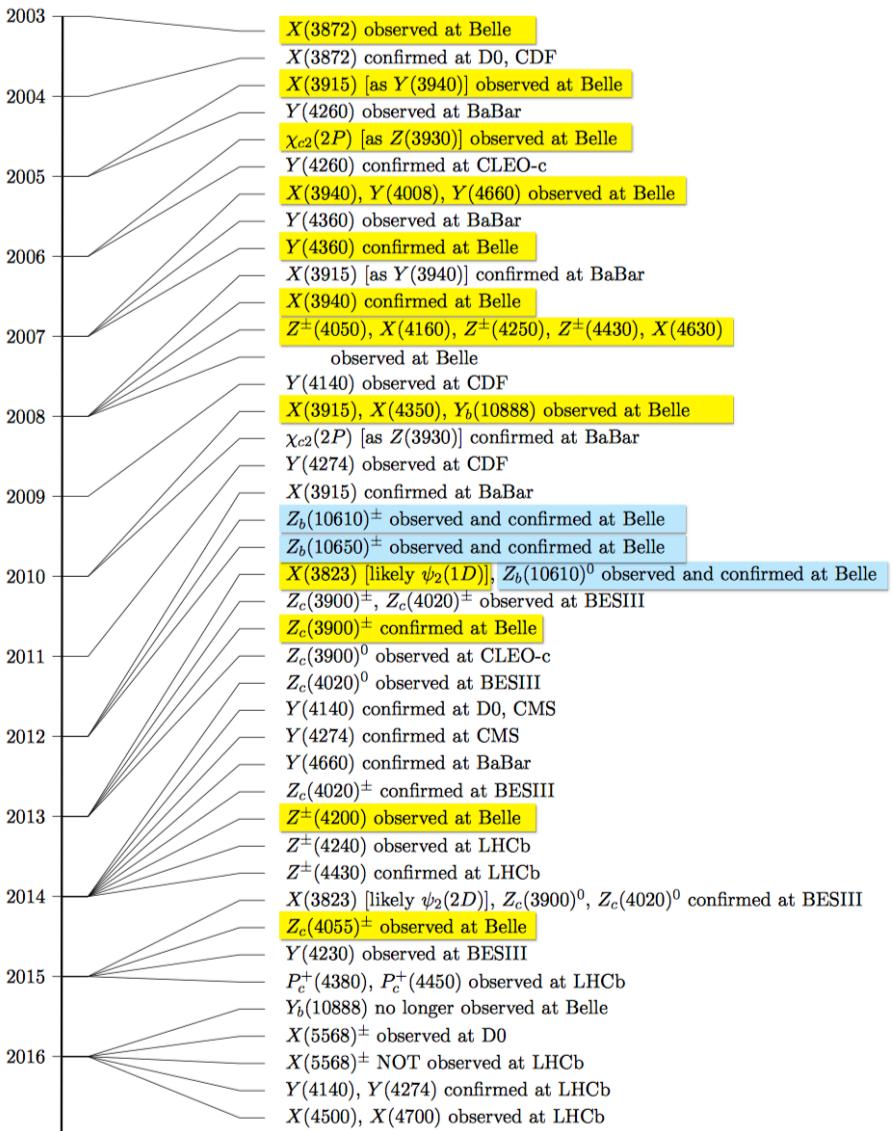
Quarkonium



Rev. Mod. Phys. 90, 015003 (2018)

- Quarkonium: $q\bar{q}$, the simplest system of a hadron.
- Below $D\bar{D}/B\bar{B}$ thresholds – both charmonium and bottomonium are successful stories of QCD.
- But there are many exotic states observed in the past decade, and they are hard to fit in the two families.

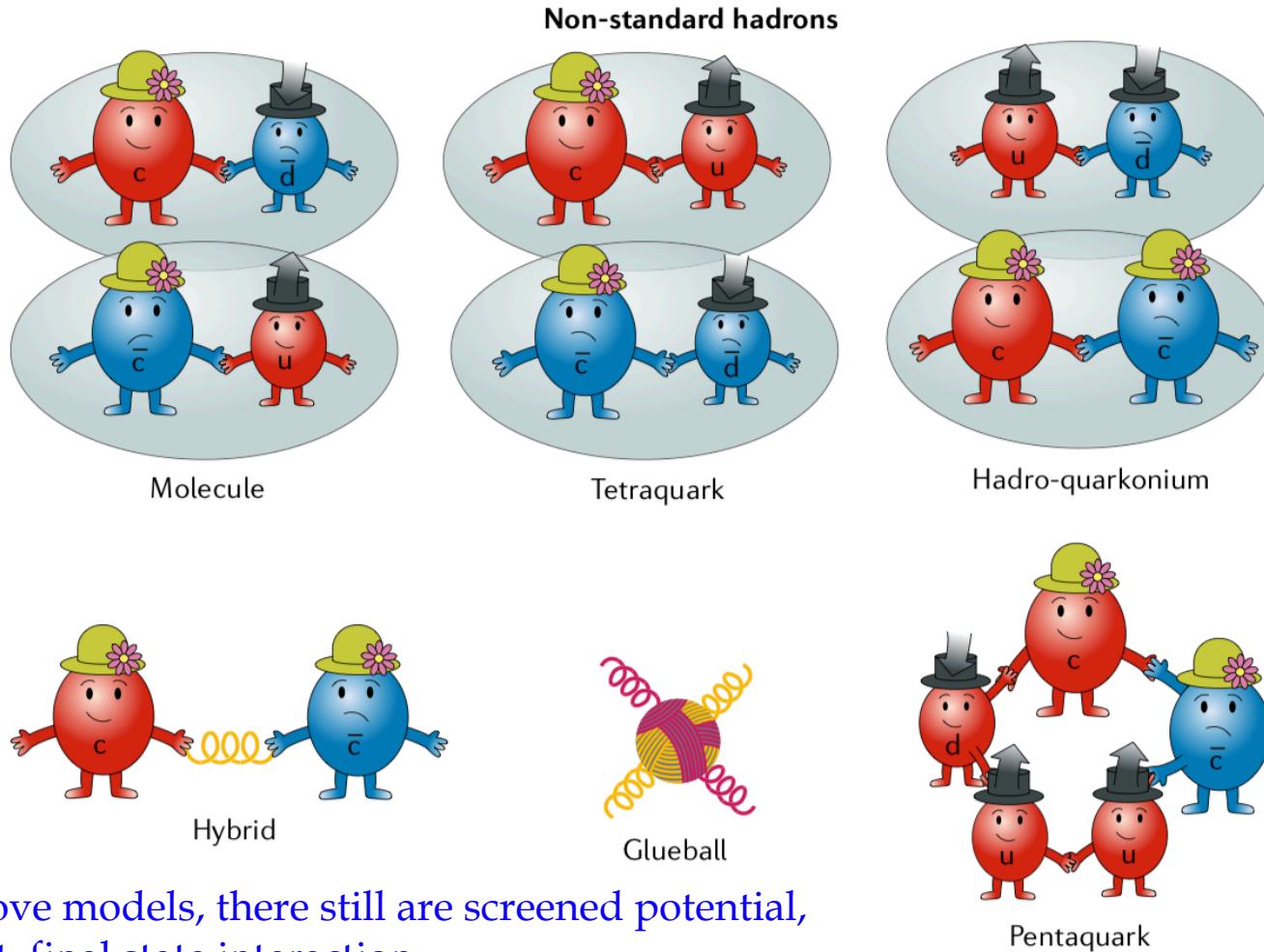
Discoveries of XYZ states at Belle



Coloured boxes: exotic candidates

- History of the quarkonium-like exotic states
- Belle accounts for ~1/2 of the discoveries, including the very first one, X(3872)

Various interpretations of the exotic states



Besides above models, there still are screened potential, cusps effect, final state interaction ...

High Priority:

- Identify most prominent component in wave function
- Seek unique picture describing all XYZ states, not state-by-state

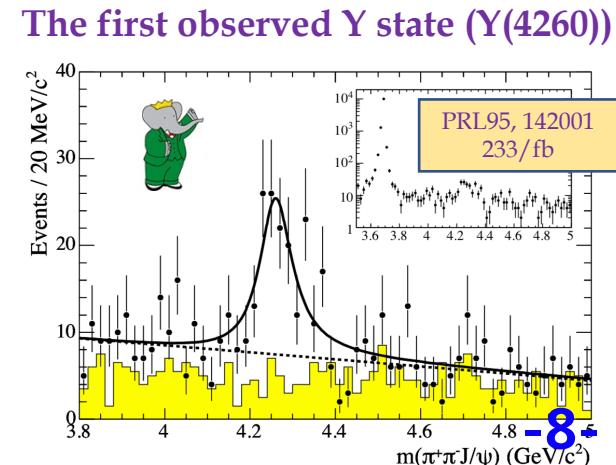
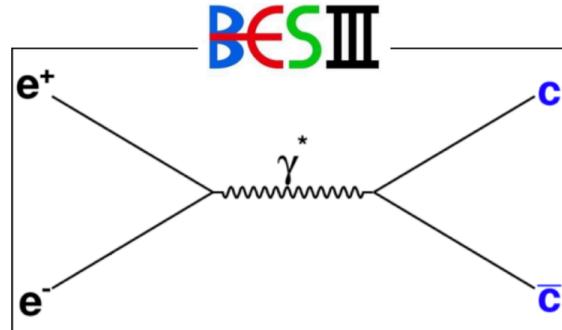
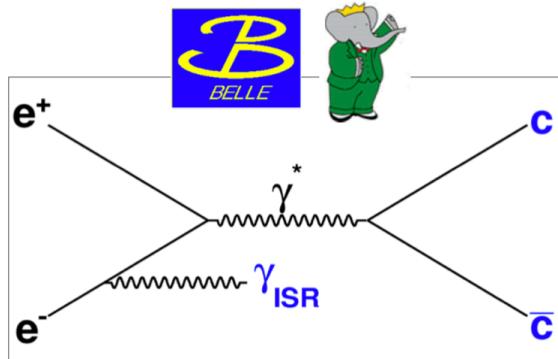
Nature Reviews Physics 1, 480 (2019)

$$e^+ e^- \rightarrow D_s^+ D_{s1}(2536)^-$$

[PHYSICAL REVIEW D 100, 111103(R) (2019)]

Y states

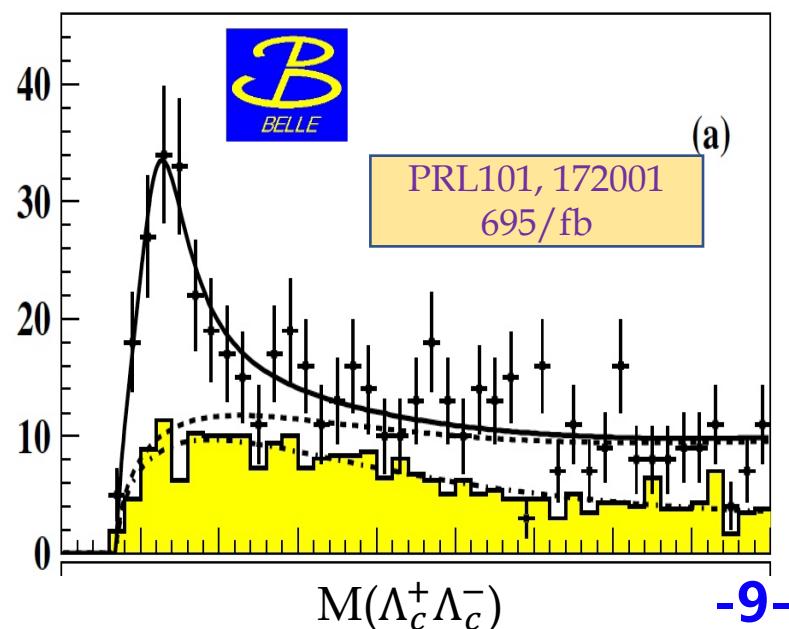
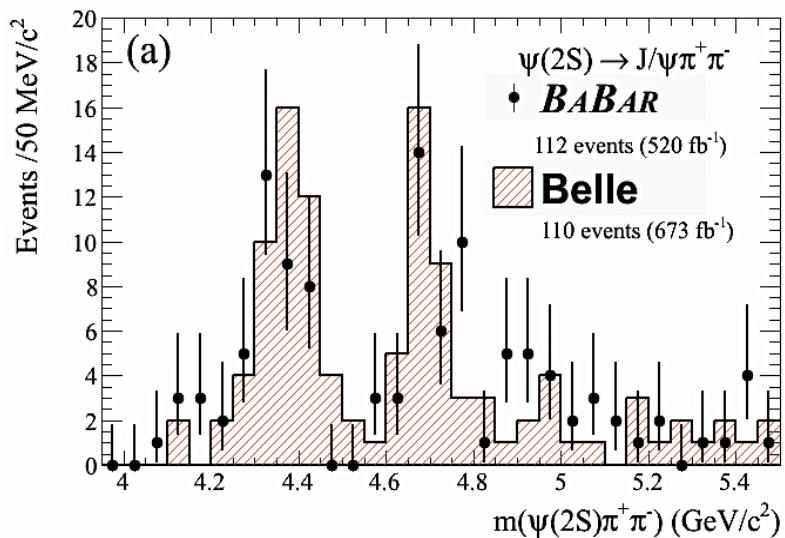
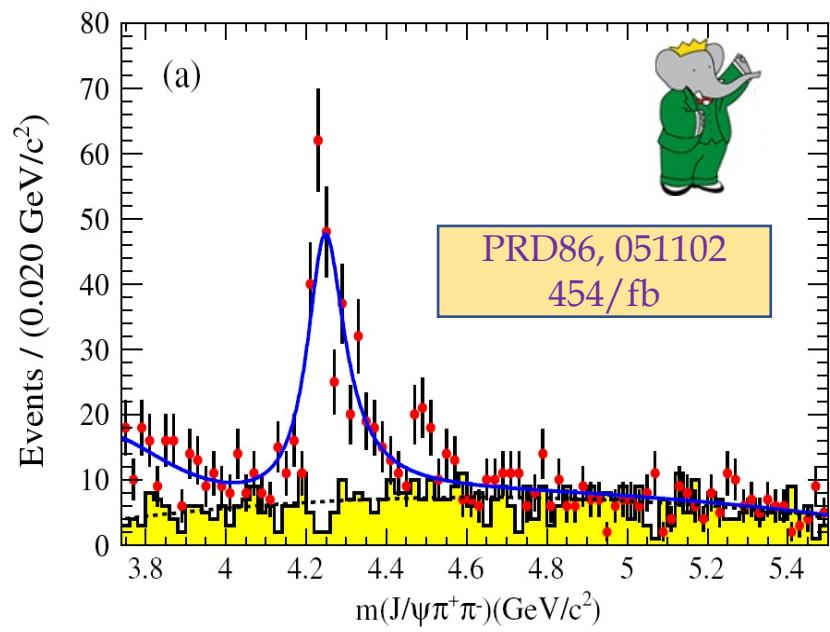
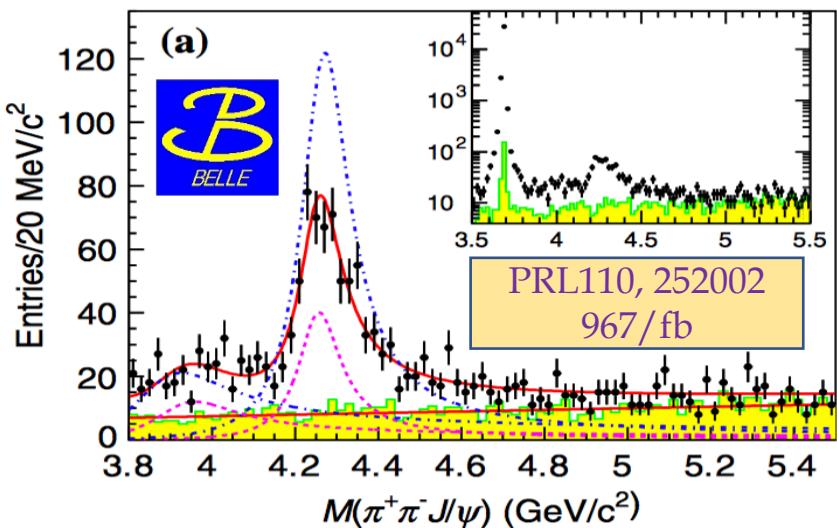
- $\Upsilon(4008), \Upsilon(4260), \Upsilon(4360), \Upsilon(4630), \Upsilon(4660)$: $J^{PC} = 1^{--}$
- Strong coupling to hidden-charm final states in contrast to the vector charmonium states in the same energy region [$\psi(4040), \psi(4160), \psi(4415)$], which couple dominantly to open-charm meson pairs.
- Many theoretical interpretations: tetraquark, molecule, hybrids, or hadrocharmonia?
- Observed in Initial state radiation processes (Belle and Babar) and e^+e^- annihilations in the charmonium energy region (BESIII).



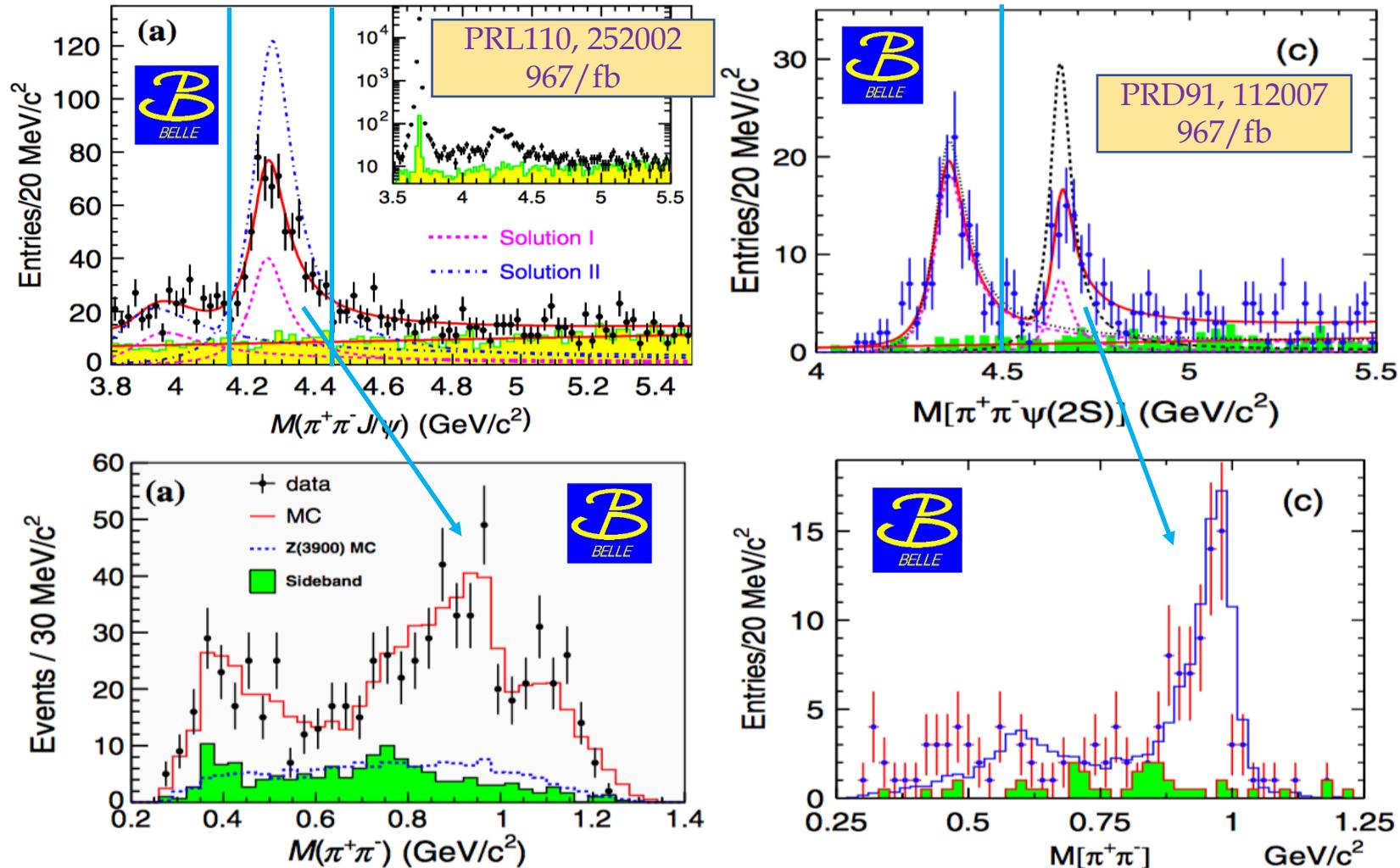
Y states via ISR

Belle: PRL99, 142002, 670/fb
 BaBar: PRD89, 111103, 670/fb

Y(4008)
 Y(4260)
 Y(4360)
 Y(4630)
 Y(4660)



$M(\pi^+\pi^-)$ in $\Upsilon(4260)$ and $\Upsilon(4660)$ signal region



- $\Upsilon(4260) \rightarrow f_0(980)(\rightarrow \pi^+\pi^-)J/\psi, \Upsilon(4660) \rightarrow f_0(980)(\rightarrow \pi^+\pi^-)\psi(2S)$
 $f_0(980)$ has a $s\bar{s}$ component, and J/ψ has a $c\bar{c}$ component.
- It is natural to search for such Υ states with a quark component of $(c\bar{s})(\bar{c}s)$, e.g., $D_s D_{s1}(2536)$.

Analysis method

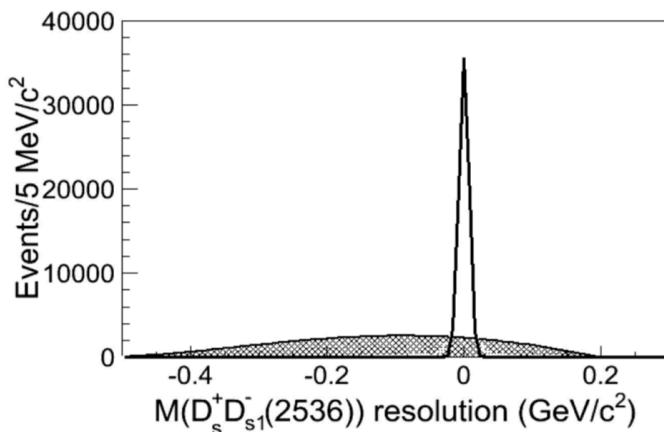
$$e^+ e^- \rightarrow \gamma_{\text{ISR}} D_s^+ D_{s1}(2536)^- (\rightarrow \bar{D}^{*0} K^- / D^{*-} K_S^0)$$

We require full reconstruction of the γ_{ISR} , D_s^+ , and K^- / K_S^0 .

- $D_s^+ \rightarrow \phi \pi^+, \bar{K}^{*0} K^+, K_S^0 K^+, K^+ K^- \pi^+ \pi^0, K_S^0 \pi^0 K^+, K^{*+} K_S^0, \eta \pi^+$, and $\eta' \pi^+$
- For the signals, the spectrum of the mass recoiling against the $D_s^+ K^- \gamma_{\text{ISR}}$ system should be accumulated at the \bar{D}^{*0} / D^{*-} nominal mass.

$$M_{\text{rec}}(\gamma_{\text{ISR}} D_s^+ K^- / K_S^0) = \sqrt{(E_{\text{c.m.}}^* - E_{\gamma_{\text{ISR}} D_s^+ K^- / K_S^0}^*)^2 - (p_{\gamma_{\text{ISR}} D_s^+ K^- / K_S^0}^*)^2}$$

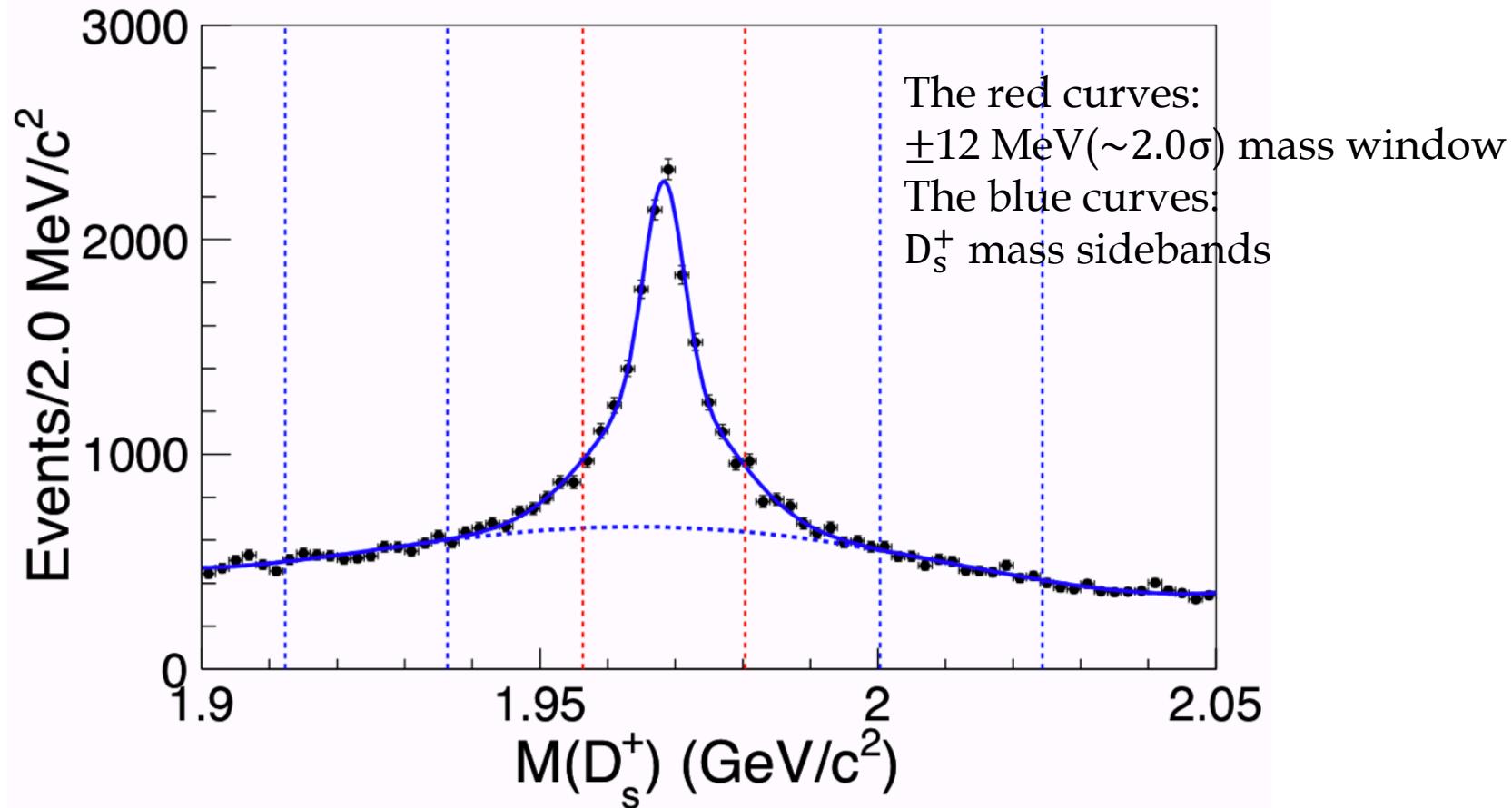
- To improve the $M_{\text{rec}}(\gamma_{\text{ISR}})$ resolution, $M_{\text{rec}}(\gamma_{\text{ISR}} D_s^+ K^- / K_S^0)$ is constrained to be the nominal mass of the \bar{D}^{*0} / D^{*-} . As a result, the resolution of $M_{\text{rec}}(\gamma_{\text{ISR}}) \equiv M(D_s^+ D_{s1}(2536)^-)$ is drastically improved ($\sim 180 \text{ MeV} \rightarrow \sim 5 \text{ MeV}$).



Data samples:

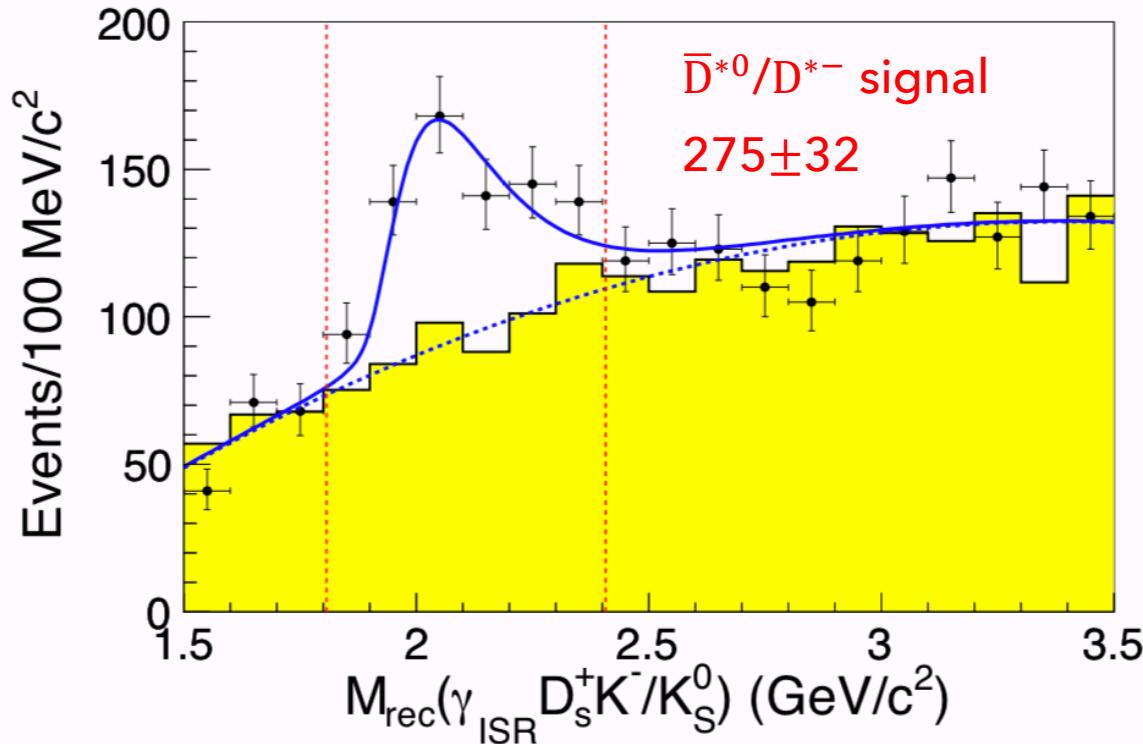
\sqrt{s} (GeV)	Luminosity (fb $^{-1}$)
10.52	89.5 ± 1.3
10.58	711 ± 10
10.867	121.4 ± 1.7
Total	921.9 ± 12.9

The invariant mass distribution for D_s^+ candidates



- Since the intrinsic width of the D_s^+ could be neglected, a double Gaussian function is used to fit the D_s^+ mass spectrum.
- The purity is $N_{\text{sig}}/(N_{\text{sig}} + N_{\text{bkg}}) = 64\%$.

The recoil mass spectrum against $\gamma_{\text{ISR}} D_s^+ K^- / K_S^0$

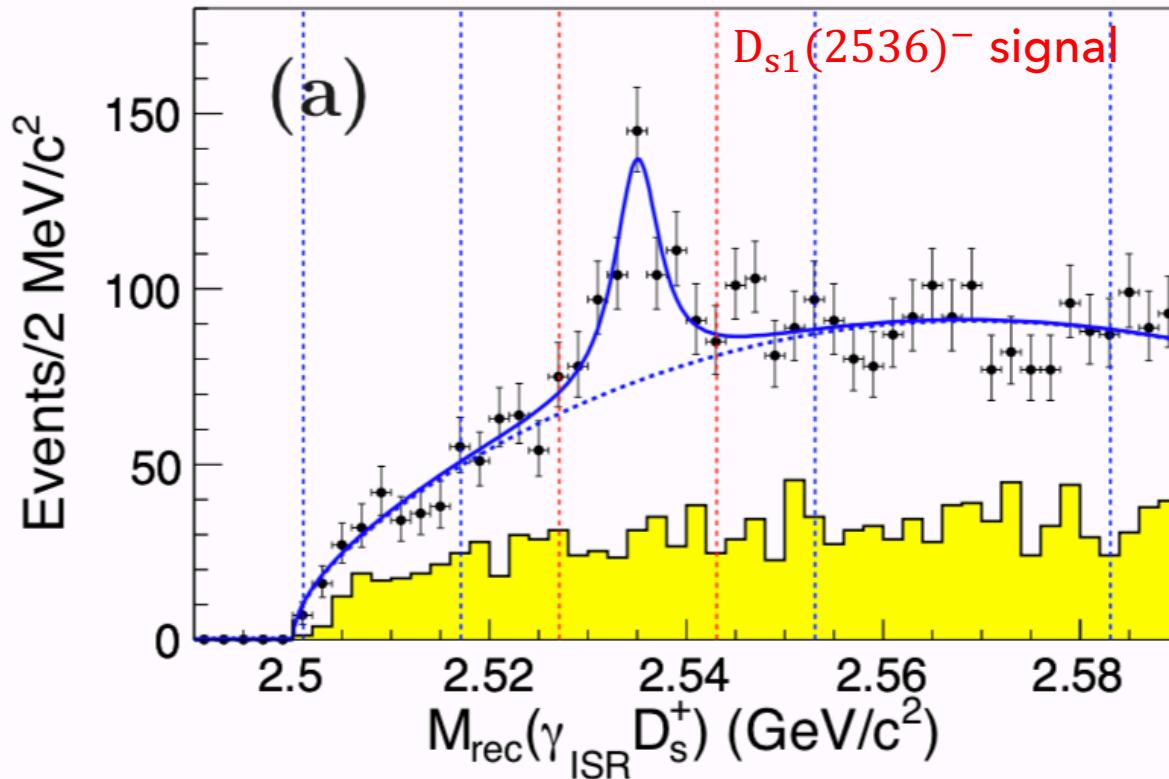


\bar{D}^{*0}/D^{*-} signal:
Gaussian \otimes Novosibirsk

The combinatorial
backgrounds:
a second-order polynomial

- $M_{\text{rec}}(\gamma_{\text{ISR}} D_s^+ K^- / K_S^0)$ distribution is making **before** applying the \bar{D}^{*0}/D^{*-} mass constraint.
- The yellow histogram shows the normalized $D_{s1}(2536)^-$ mass sidebands (see below).
- Due to the poor mass resolution, the \bar{D}^{*0}/D^{*-} signal is very wide.

The recoil mass spectrum against $\gamma_{\text{ISR}} D_s^+$



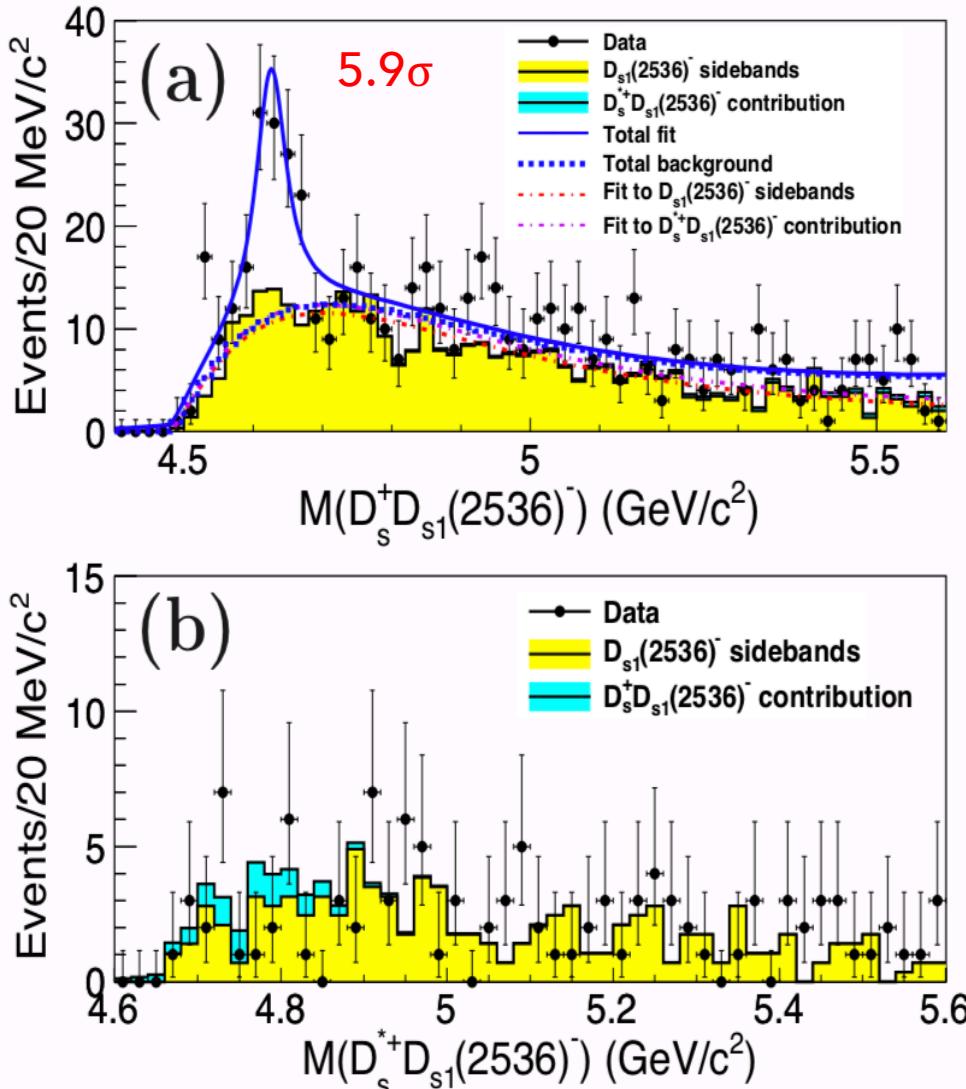
$D_{s1}(2536)^-$ signal:
Double Gaussian

The combinatorial
backgrounds:
threshold function

- $M_{\text{rec}}(\gamma_{\text{ISR}} D_s^+)$ distribution is making after applying the \bar{D}^{*0}/D^{*-} mass constraint.
- The yellow histogram shows the normalized D_s^+ mass sidebands.
- The fit yields 254 ± 36 $D_{s1}(2536)^-$ signal events with the statistical significance of 8.0σ .

Final mass spectrum $M(D_s^+ D_{s1}(2536)^-)$

After applying the \bar{D}^{*0}/D^{*-} mass constraint



An unbinned simultaneous likelihood fit:

- Signal: a BW convolved with a Gaussian function, then multiplied by an efficiency function
- $D_{s1}(2536)^-$ mass sidebands: a threshold function
- $e^+ e^- \rightarrow D_s^{*+} D_{s1}(2536)^-$ background contribution: a threshold function
- A non-resonant contribution: a two-body phase space form

$$M = (4625.9^{+6.2}_{-6.0}(\text{stat.}) \pm 0.4(\text{syst.}) \text{ MeV}/c^2)$$

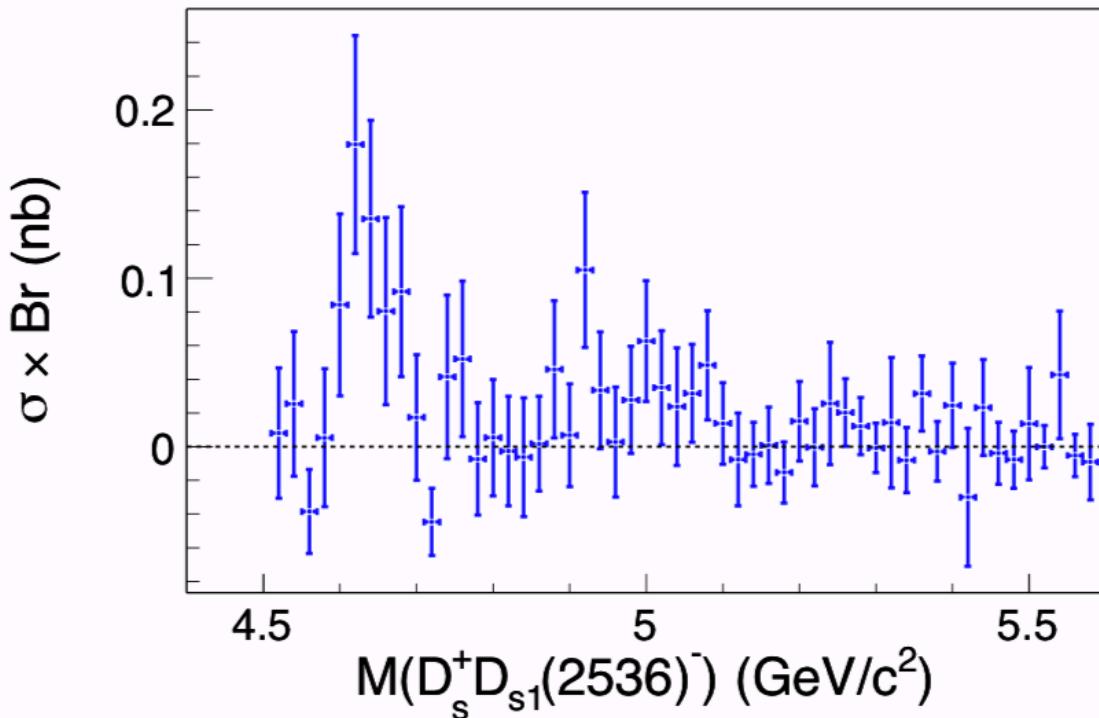
$$\Gamma = (49.8^{+13.9}_{-11.5}(\text{stat.}) \pm 4.0(\text{syst.}) \text{ MeV})$$

$$\Gamma_{ee} \times \mathcal{B}(Y \rightarrow D_s^+ D_{s1}(2536)^-) \times \mathcal{B}(D_{s1}(2536)^- \rightarrow \bar{D}^{*0} K^-) = (14.3^{+2.8}_{-2.6}(\text{stat.}) \pm 1.5(\text{syst.}) \text{ eV})$$

One possible background is from $e^+ e^- \rightarrow D_s^{*+} (\rightarrow D_s^+ \gamma) D_{s1}(2536)^-$, where the photon from the D_s^{*+} remains undetected. No obvious structure is observed in the $e^+ e^- \rightarrow D_s^{*+} (\rightarrow D_s^+ \gamma) D_{s1}(2536)^-$. 15

Cross section

$$\sigma(e^+e^- \rightarrow D_s^+ D_{s1}(2536)^-) \mathcal{B}(D_{s1}(2536)^- \rightarrow \bar{D}^{*0} K^-) =$$
$$\frac{N_{\text{fit}}^{D_{s1}(2536)^-}}{d\mathcal{L} \times [\sum_i (\varepsilon_i^{\bar{D}^{*0} K^-} \times \mathcal{B}_i) + R_{\bar{D}^{*0} K^-}^{D^* - K_S^0} \times \sum_i (\varepsilon_i^{D^* - K_S^0} \times \mathcal{B}_i)]},$$



The yield of fitted $D_{s1}(2536)^-$ signal events after subtracting the $e^+e^- \rightarrow D_s^{*+} D_{s1}(2536)^-$ background contribution.

The peak value of the $\sigma \times \text{Br}$ at $M(D_s^+ D_{s1}(2536)^-) \sim 4.63 \text{ GeV}/c^2$ is about $(0.18 \pm 0.06) \text{ nb}$.

$$e^+ e^- \rightarrow D_s^+ D_{s2}^*(2573)^-$$

arXiv:2004.02404 (2020)

Analysis method

$$e^+ e^- \rightarrow \gamma_{\text{ISR}} D_s^+ D_{s2}^*(2573)^-(\rightarrow \bar{D}^0 K^-)$$

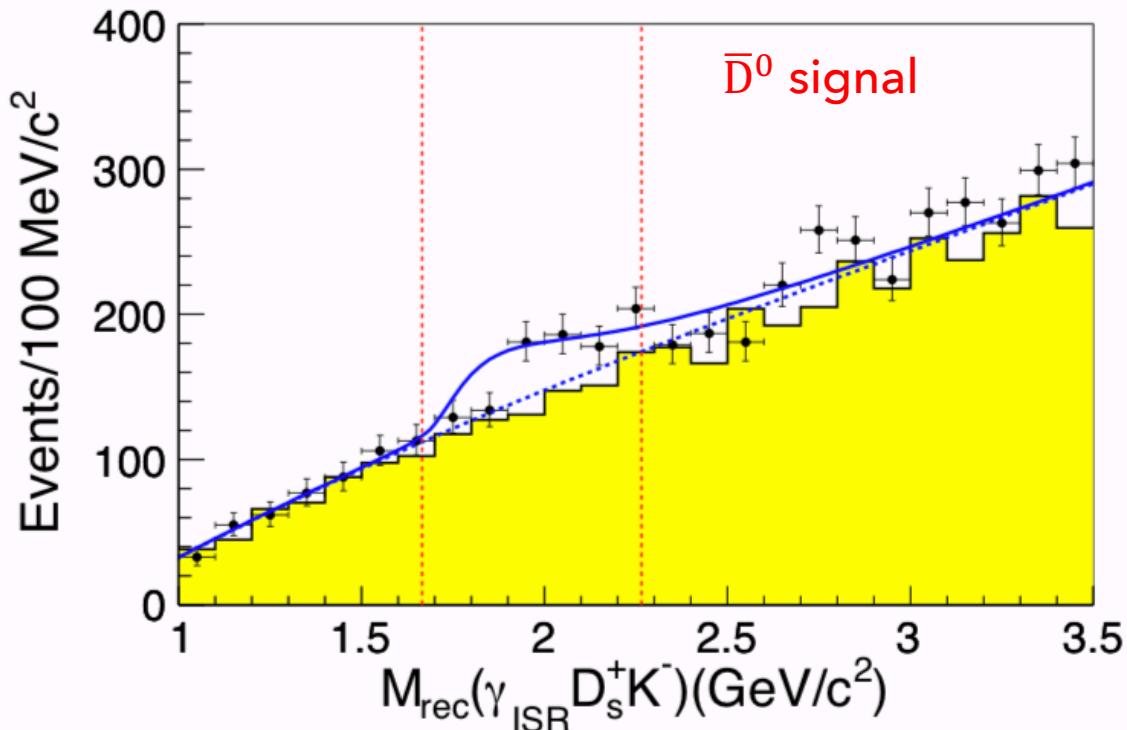
We require full reconstruction of the γ_{ISR} , D_s^+ , and K^- .

- $D_s^+ \rightarrow \phi\pi^+, K_s^0 K^+, \bar{K}^{*0} K^+, \phi\rho^+, K^{*+} \bar{K}^{*0}, K^{*+} K_s^0, K_s^0 K^+ \pi^+ \pi^-, \eta\pi^+, \text{ and } \eta'\pi^+$.
- For the signals, the spectrum of the mass recoiling against the $D_s^+ K^- \gamma_{\text{ISR}}$ system should be accumulated at the \bar{D}^0 nominal mass.

$$M_{\text{rec}}(\gamma_{\text{ISR}} D_s^+ K^-) = \sqrt{(E_{\text{c.m.}}^* - E_{\gamma_{\text{ISR}} D_s^+ K^-}^*)^2 - (p_{\gamma_{\text{ISR}} D_s^+ K^-}^*)^2}$$

- To improve the γ_{ISR} resolution we refit recoil mass against $D_s^+ K^- \gamma_{\text{ISR}}$ into the \bar{D}^0 mass. As a result, the resolution of $M_{\text{rec}}(\gamma_{\text{ISR}}) \equiv M(D_s^+ D_{s2}^{*-}(2573))$ is drastically improved ($\sim 180 \text{ MeV} \rightarrow \sim 5 \text{ MeV}$).

The recoil mass spectrum against $\gamma_{\text{ISR}} D_s^+ K^-$

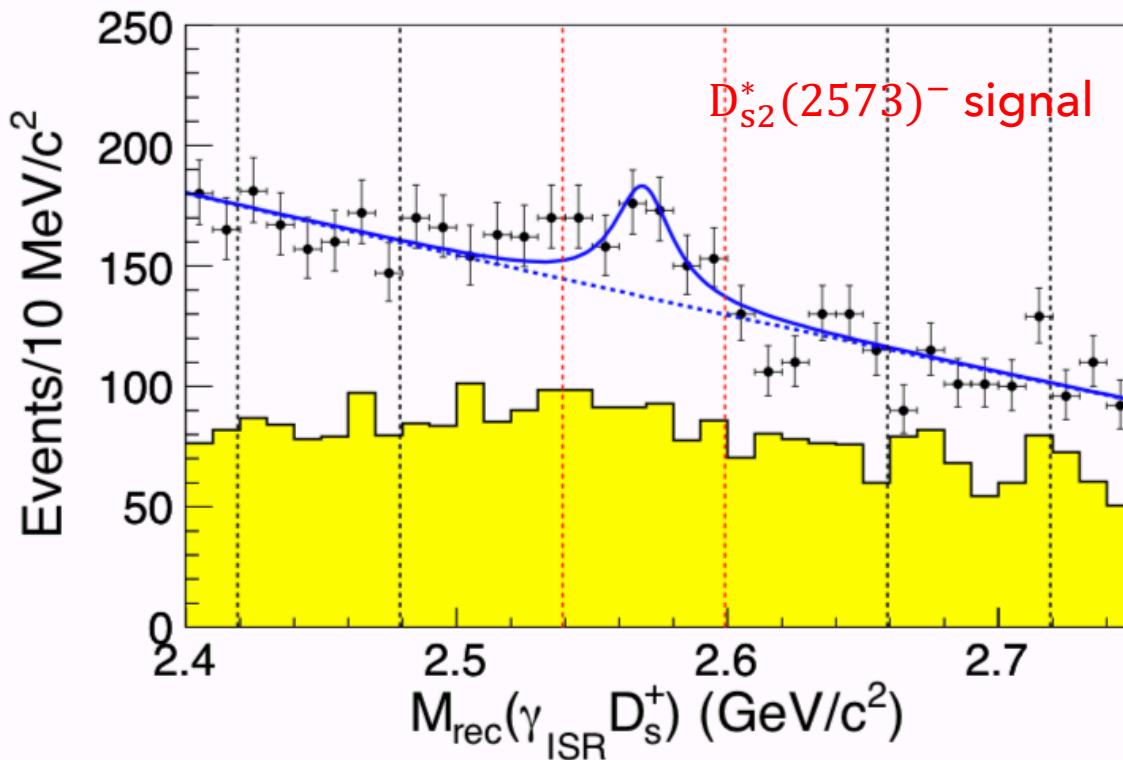


A simultaneous likelihood fit to the $M_{\text{rec}}(\gamma_{\text{ISR}} D_s^+ K^-)$ of $D_{s2}^*(2573)^-$ signal candidates and the normalized $D_{s2}^*(2573)^-$ mass sidebands is applied.

The best fit yields 224 ± 42 $D_{s2}^*(2573)^-$ signal events.

- $M_{\text{rec}}(\gamma_{\text{ISR}} D_s^+ K^-)$ distribution is making before applying the \bar{D}^0 mass constraint.
- The yellow histogram shows the normalized $D_{s2}^*(2573)^-$ mass sidebands (see below).
- The \bar{D}^0 signal is wide and asymmetric due to the asymmetric resolution function of the ISR photon energy and higher order ISR corrections. -19-

The recoil mass spectrum against $\gamma_{\text{ISR}} D_s^+$



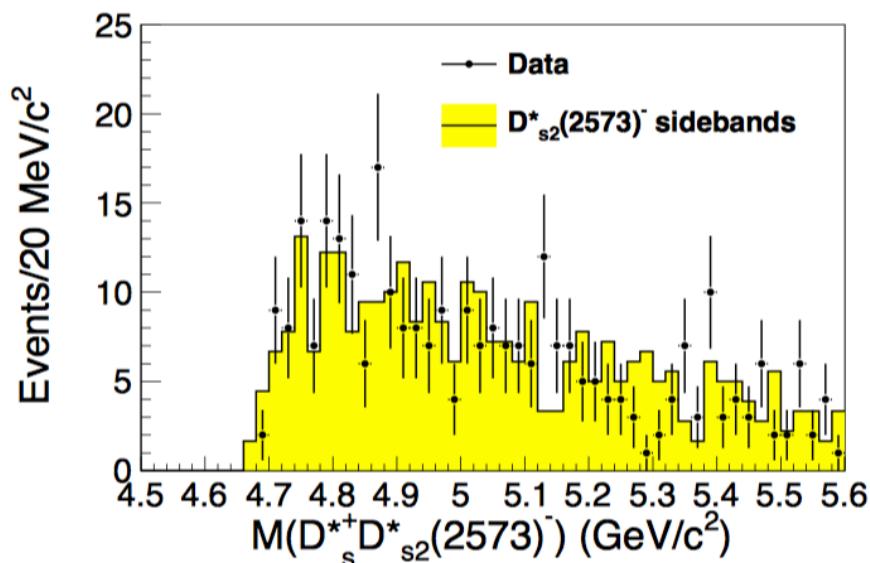
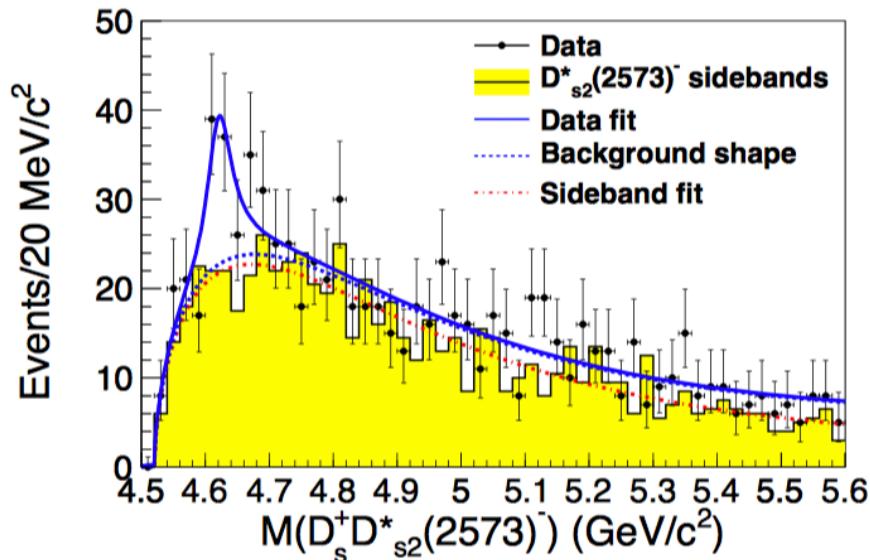
D_{s2}^{*}(2573)⁻ signal:
BW \otimes Gaussian

The combinatorial
backgrounds:
A second-order polynomial

- M_{rec}(γ_{ISR} D_s⁺) distribution is making after applying the \bar{D}^0 mass constraint.
- The yellow histogram shows the normalized D_s⁺ mass sidebands.
- The fit yields 182 ± 47 D_{s2}^{*}(2573)⁻ signal events with the statistical significance of 4.1σ .

Final mass spectrum $M(D_s^+ D_{s2}^*(2573)^-)$

After applying the \bar{D}^0 mass constraint



An unbinned simultaneous likelihood fit:

- Signal: a BW convolved with a Gaussian function, then multiplied by an efficiency function
- $D_{s2}^*(2573)^-$ mass sidebands: a threshold function
- A non-resonant contribution: a two-body phase space form

$$M = (4619.8^{+8.9}_{-8.0}(\text{stat.}) \pm 2.3(\text{syst.}) \text{ MeV}/c^2)$$

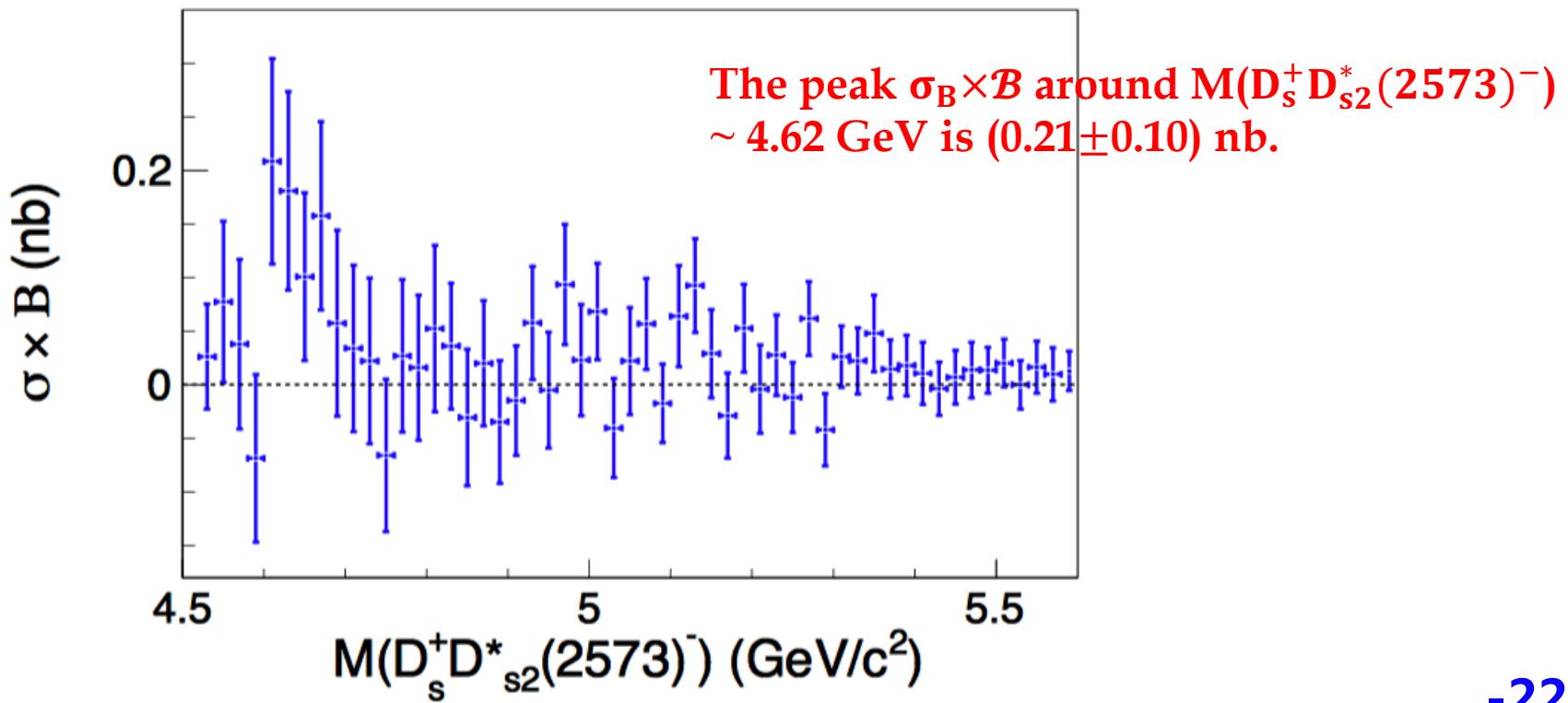
$$\Gamma = (47.0^{+31.3}_{-14.8}(\text{stat.}) \pm 4.6(\text{syst.}) \text{ MeV}$$

$$G_{ee} \times \mathcal{B}(Y \rightarrow D_s^+ D_{s2}^*(2573)^-) \times \mathcal{B}(D_{s2}^*(2573)^- \rightarrow \bar{D}^0 K^-) = (14.7^{+5.9}_{-4.5}(\text{stat.}) \pm 3.6(\text{syst.}) \text{ eV}$$

Cross section

$$\sigma_B(e^+e^- \rightarrow D_s^+ D_{s2}^*(2573)^-) \mathcal{B}(D_{s2}^*(2573)^- \rightarrow \bar{D}^0 K^-) =$$
$$\frac{N^{obs}}{\sum_i (\varepsilon_i \times \mathcal{B}_i) \times \Delta \mathcal{L}}$$

N^{obs} : the number of observed $e^+e^- \rightarrow D_s^+ D_{s2}^*(2573)^-$ signal events after subtracting the normalized $D_{s2}^*(2573)^-$ mass sidebands

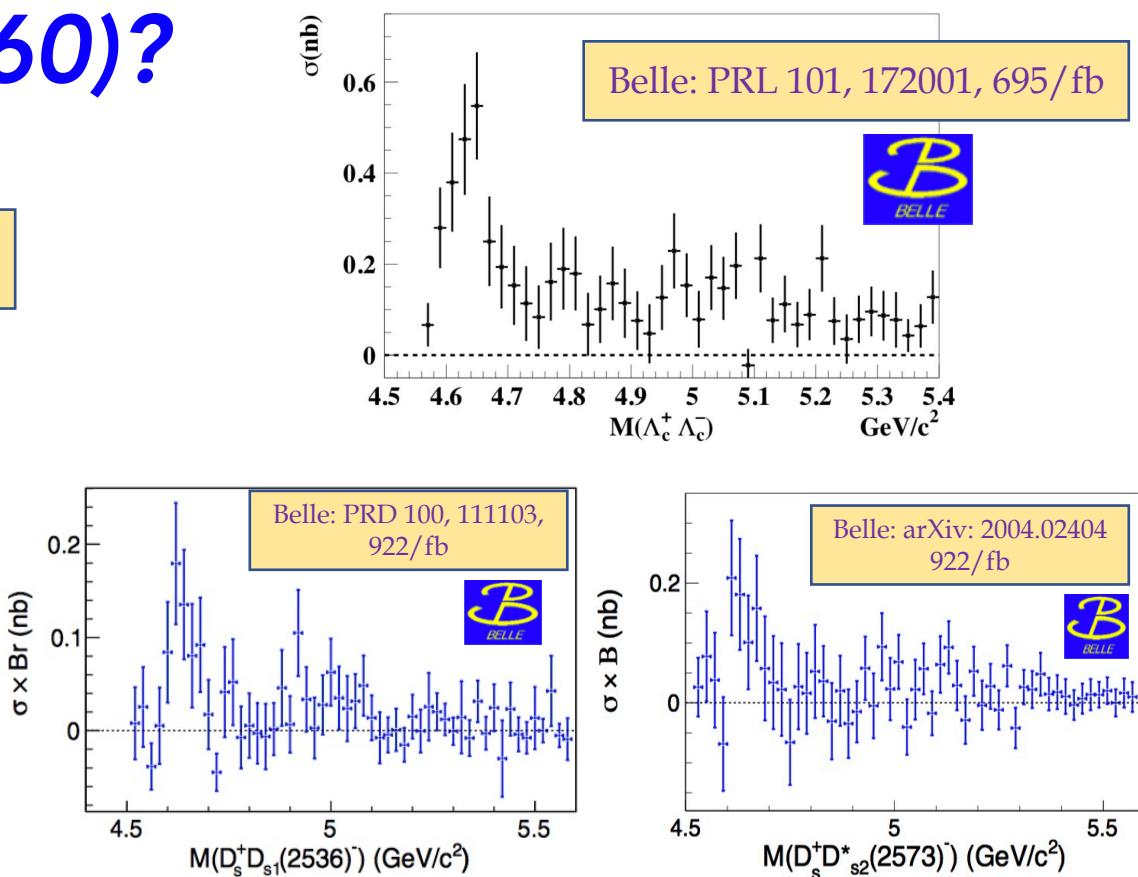
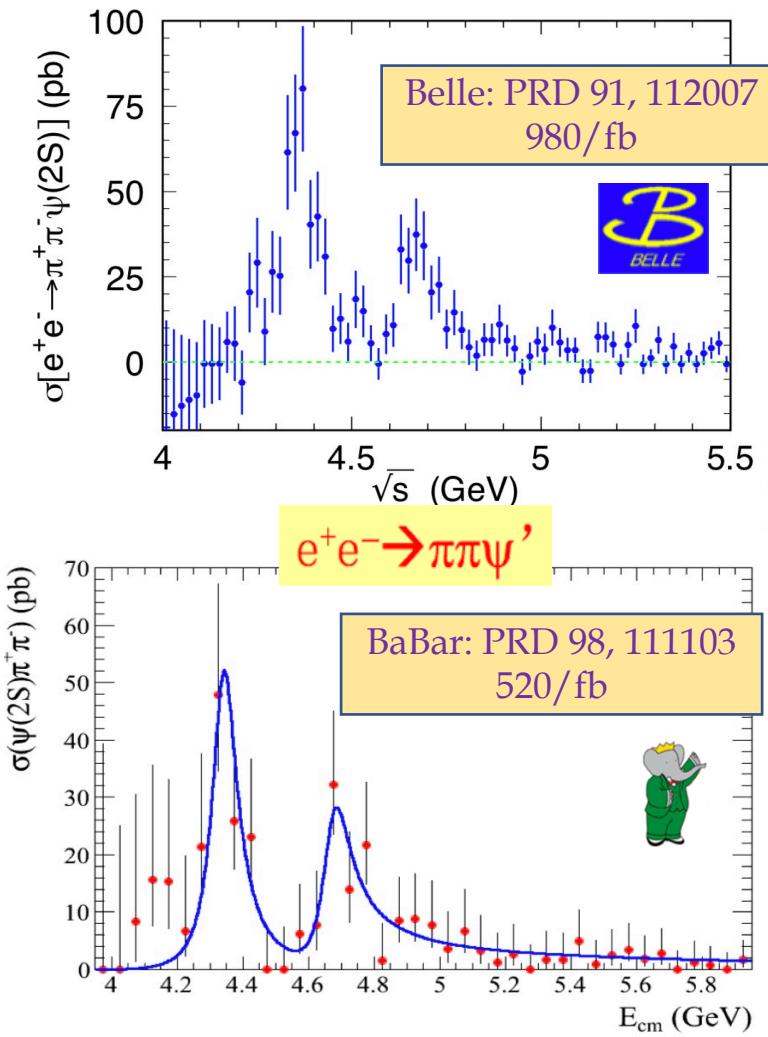


Summary and discussion

Interpretations of Y(4626)

- A **tetraquark state** in a chiral constituent quark model with a scaling method [PRD101, 054010 (2020)].
- A **P-wave tetraquark state** $[cs][\bar{c}\bar{s}]$ with 1^{--} in the multiquark color flux-tube model [PRD 101, 054039 (2020)].
- A hidden-strange **molecular state** from $\Lambda_c^+ \Lambda_c^-$ interaction [arXiv:1911.03706 (2020)].
- A **molecular state** from interaction $D_s^* \bar{D}_{s1}(2536) - D_s \bar{D}_{s1}(2536)$ [EPJC 80, 246 (2020)].
- A tetraquark and etc instead of $D_s^* \bar{D}_{s1}(2536)$ molecular within the Bethe-Salpeter framework [arXiv:2004.03167 (2020)].
- A higher charmonium [PRD 101, 034001 (2020)].
- A P-Wave hidden-charm exotic mesons in the diquark model [arXiv:2003.02802 (2020)].

$\text{Y}(4630) = \text{Y}(4660)?$



Experiment	Mass (MeV)	Width (MeV)
Belle, $\Lambda_c^+\Lambda_c^-$	4634^{+8+5}_{-7-8}	92^{+40+10}_{-24-21}
Belle, $\pi^+\pi^-\psi(2S)$	$4652 \pm 10 \pm 8$	$68 \pm 11 \pm 1$
BaBar, $\pi^+\pi^-\psi(2S)$	$4669 \pm 21 \pm 3$	$104 \pm 48 \pm 10$
Belle, $D_s^+D_{s1}(2536)^-$	$4626^{+7}_{-7} \pm 1$	$49.8^{+14}_{-12} \pm 4$
Belle, $D_s^+D_{s2}^*(2573)^-$	$4620^{+9}_{-8} \pm 3$	$47.0^{+32}_{-15} \pm 5$

What is Y(4660)?

- Charmonium?
- Molecule [$f_0(980)\psi'$, $\Lambda_c^+\Lambda_c^-$]?
- Hadron-charmonium?
- Tetraquark state?
- Hybrid?
- ...

Experimental measurements:

$Y(4660) \rightarrow$

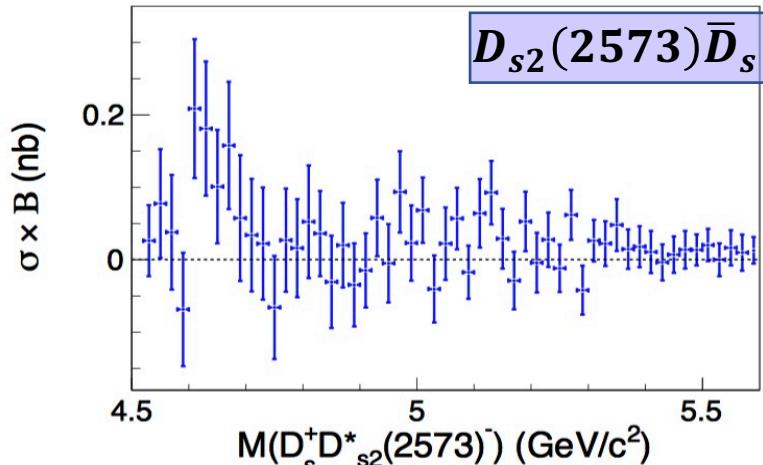
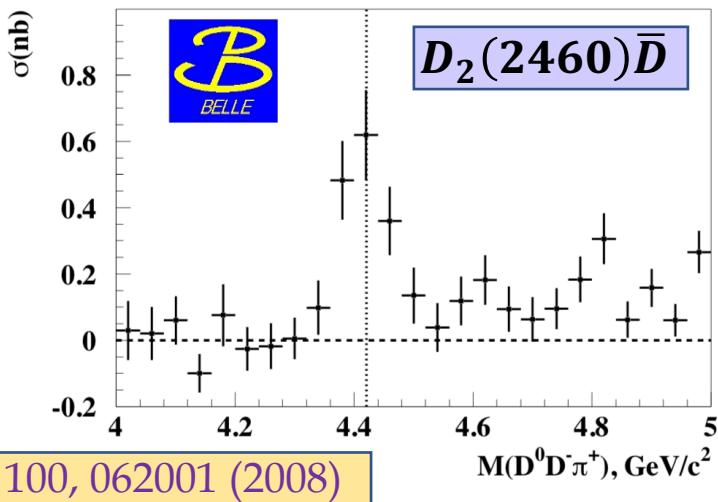
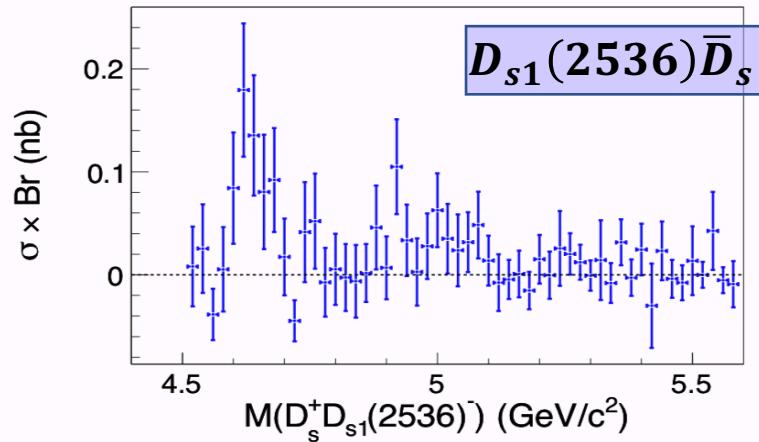
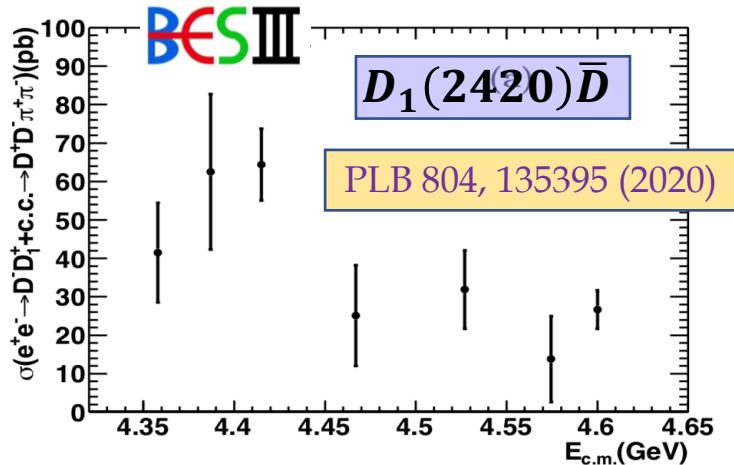
- $D_s^* D_{s0}^*(2317)$
- $D_s D_{s1}(2460)$
- $D_s^* D_{s1}(2460)$
- $D_s^* D_{s1}(2536)$
- $D_s^* D_{s2}^*(2573)$

May these rates be estimated according to $D_s^+ D_{s1}(2536)^-$ and $D_s^+ D_{s2}^*(2573)^-$?

Why does Y(4660) couple to charmed baryon strongly?
Why does Y(4660) couple to $s\bar{s}$ strongly?

Extend to the charm meson pairs

- $D_1(2420)$, $D_2(2460)$, $D_{s1}(2536)$, $D_{s2}(2573)$ are $j_l^P = \frac{3}{2}^+$ heavy-light states.
- The widths of the $D_1(2420)$, $D_2(2460)$, $D_{s1}(2536)$, $D_{s2}(2573)$ are narrow.



Do $Y(4360)/Y(4390)/\Psi(4415)$ & $Y(4630)/Y(4660)$ have similar structures?-27-

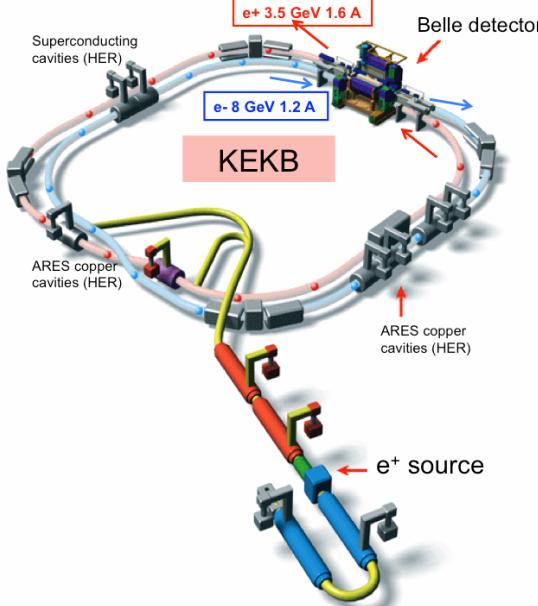
Status and Prospects of Belle II

SuperKEKB

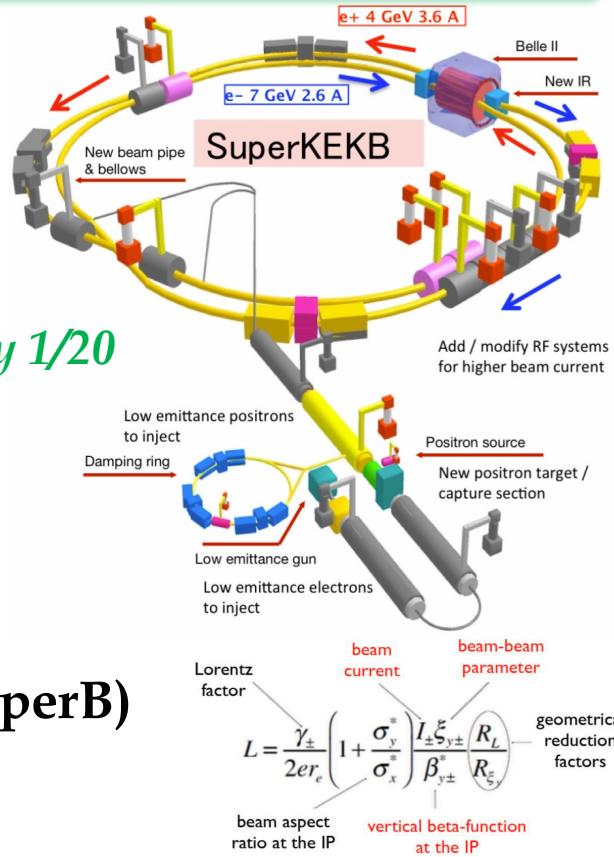
- 1st Vs. 2nd generation B-factory

goal

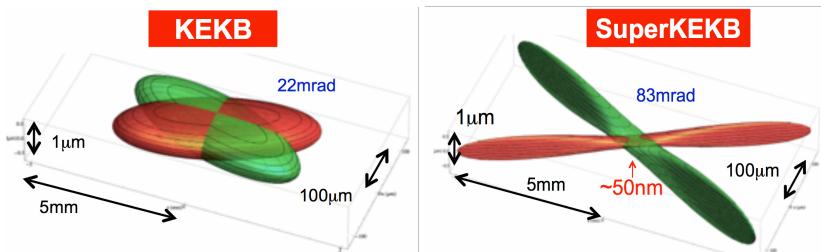
$$\int \mathcal{L} dt = 50 \text{ ab}^{-1} = 50 \times \mathcal{L}_{\text{Belle}}^{\text{int}}$$



- *Double beam currents*
- *Squeeze beams @IP by 1/20*
- *Reduced CM boost*



- Nano-beam design (by P. Raimondi for SuperB)



	E_{\pm} (GeV) LER/HER	Cross Angle (mrad)	I_{\pm} (A) LER/HER	β_y^* (mm) LER/HER	\mathcal{L} ($\text{cm}^{-2}\text{s}^{-1}$)
KEKB	3.5/8.0	22	1.64/1.19	5.9/5.9	2.1×10^{34}
SuperKEKB	4.0/7.0	83	3.60/2.60	0.27/0.31	80×10^{34}

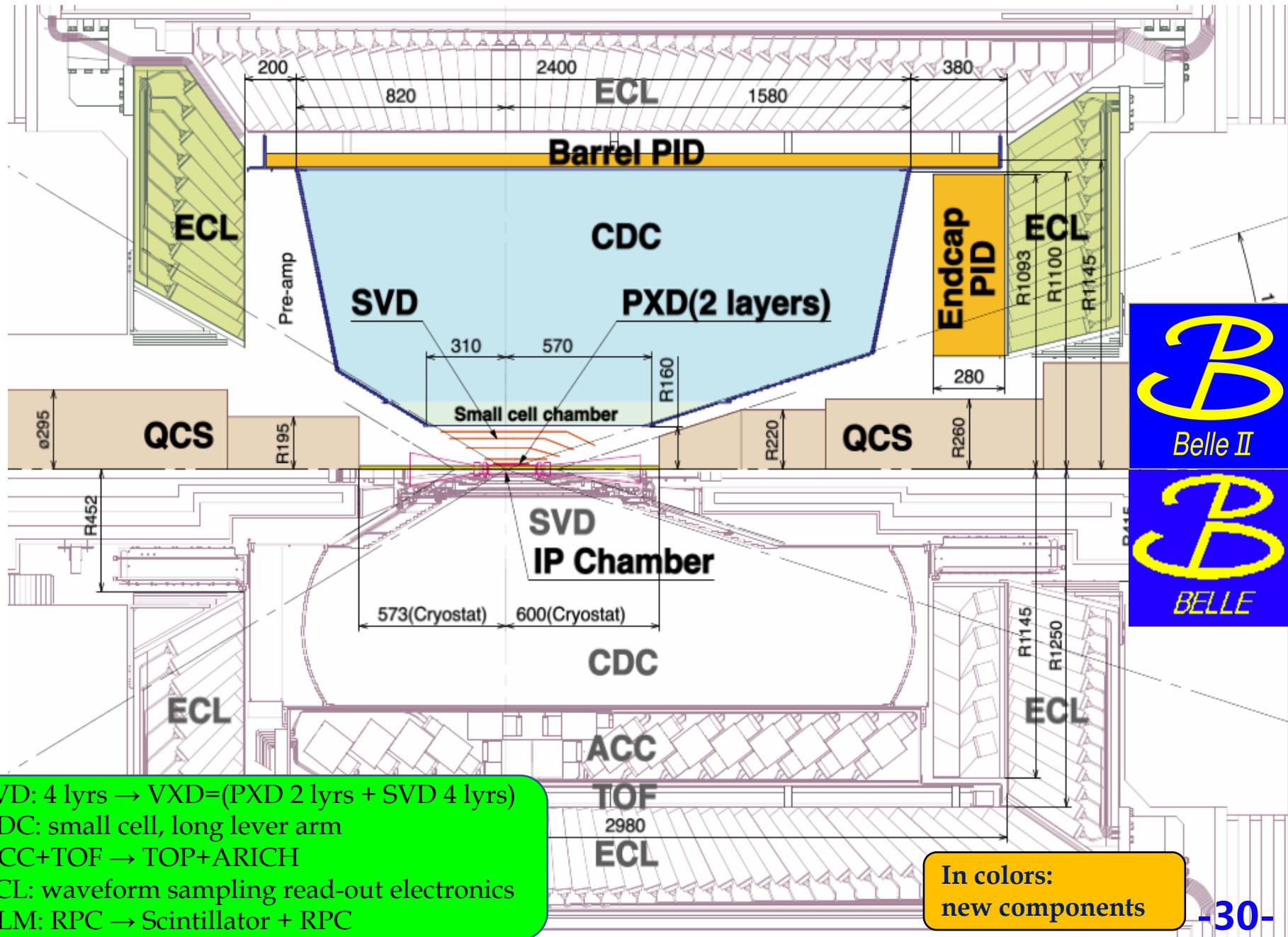
$\beta\gamma \sim 2/3$

$\times 2$

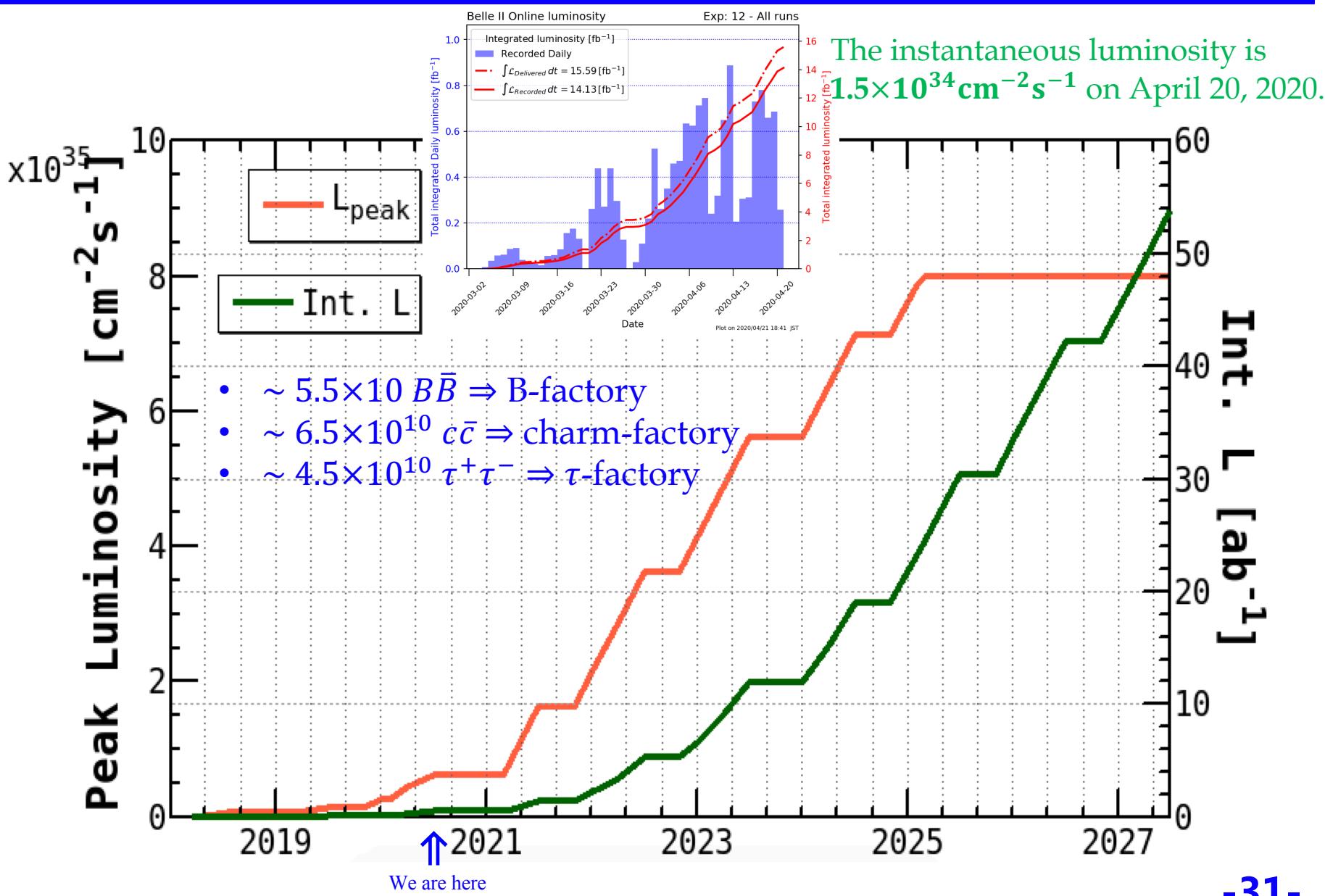
$\times 20$

$\times 40$

Detector: Belle Vs. Belle II



Dataset at Belle II



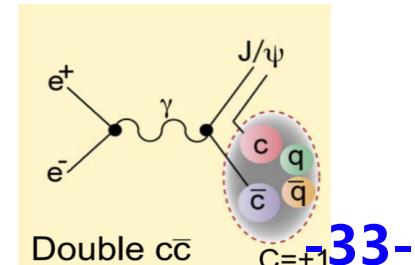
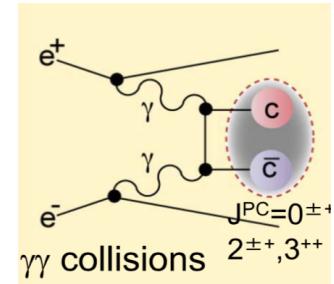
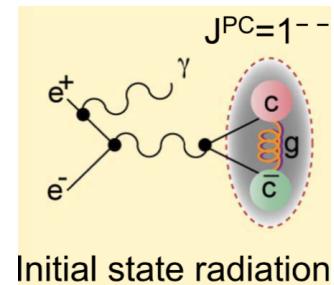
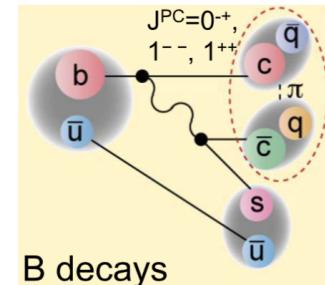
Exotic and Quarkonium Prospects at Belle II

Production of charmonium(-like) states at B-factory

- **B decay** ($B \rightarrow KX_{c\bar{c}}$)
 - CKM favored process, large branching fractions $10^{-3} \sim 10^{-4}$
 - $J^{PC} = 0^{-+}, 1^{--}, 1^{++}, \dots$
- **Initial-state radiation (ISR)**
 - $J^{PC} = 1^{--}$
- **Two-photon process**
 - $J^{PC} = 0^{-+}, 0^{++}, 2^{++}, 2^{-+}, \dots$
- **Double charmonium**
 - e.g. $e^+e^- \rightarrow J/\psi X(3940)$ [PRL 98, 082001 (2007)]

Expected statistics @50 ab⁻¹ of XYZ

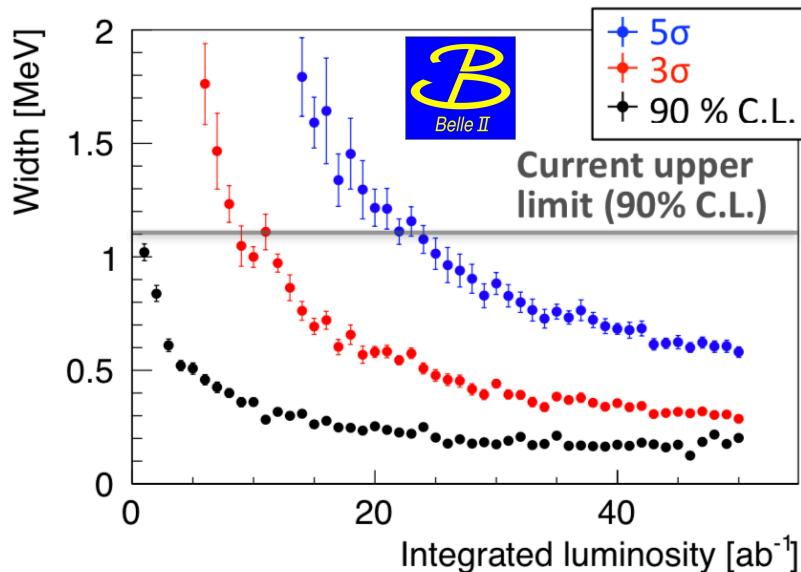
State	Production and Decay	N
X(3872)	$B \rightarrow KX(3872)$, $X(3872) \rightarrow J/\psi \pi^+ \pi^-$	$\simeq 14400$
Y(4260)	ISR, $Y(4260) \rightarrow J/\psi \pi^+ \pi^-$	$\simeq 29600$
Z(4430)	$B \rightarrow K^\mp Z(4430)$, $Z(4430) \rightarrow J/\psi \pi^\pm$	$\simeq 10200$



Further investigations at Belle II

B decays:

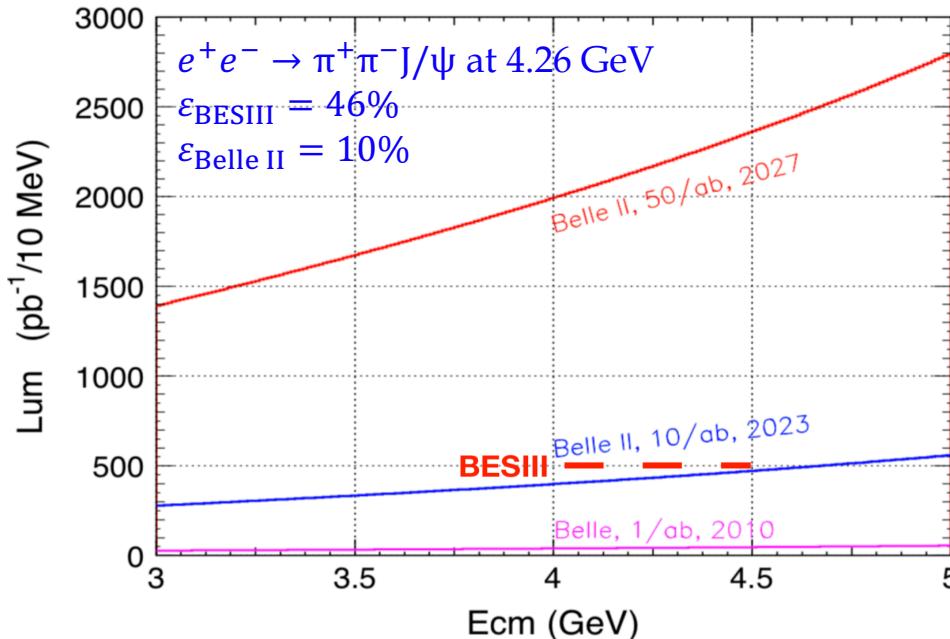
- Determination of the X(3872) total width
- The branching fraction of $B \rightarrow KX(3872)$
- $X(3872) \rightarrow D^0 \bar{D}^{*0}$ has been seen. Other open flavor decays? $B \rightarrow K D \bar{D}$, $K D^* \bar{D}^*$, $K D \bar{D}^{**}$, $K D^* \bar{D}^{**}$.
- Full amplitude analysis to $B \rightarrow K\omega J/\psi$ and $B \rightarrow K\pi\chi_{c1}$ to determine the spin-parities of X(3915), Z(4050), Z(4250) ...
- $B \rightarrow K + \text{anything}$, to discover new particles?



With the full data sample of Belle II (50 ab^{-1}), total width with values up to
[90% C.L.] $\sim 180 \text{ keV}$
[3σ significance] $\sim 280 \text{ keV}$
[5σ significant] $\sim 570 \text{ keV}$
can be measured.

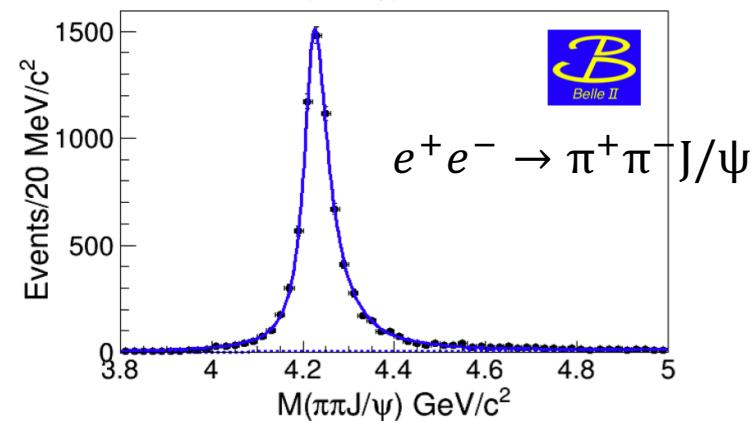
Further investigations at Belle II

ISR process



- Belle II $50 \text{ ab}^{-1} \sim \text{BESIII } 500 \text{ pb}^{-1}$, 10 MeV scan
- Line shape of the resonance and fine structures can be investigated.

PHOKHARA simulation:



Golden Channels	$E_{c.m.}$ (GeV)	Statistical error (%)	Related $X Y Z$ states
$\pi^+\pi^-J/\psi$	4.23	7.5 (3.0)	$Y(4008)$, $Y(4260)$, $Z_c(3900)$
$\pi^+\pi^-\psi(2S)$	4.36	12 (5.0)	$Y(4260)$, $Y(4360)$, $Y(4660)$, $Z_c(4050)$
K^+K^-J/ψ	4.53	15 (6.5)	Z_{cs}
$\pi^+\pi^-h_c$	4.23	15 (6.5)	$Y(4220)$, $Y(4390)$, $Z_c(4020)$, $Z_c(4025)$
$\omega\chi_{c0}$	4.23	35 (15)	$Y(4220)$

10 ab^{-1} 35 (15) 50 ab^{-1}

Bottomonium(-like) at B -factories

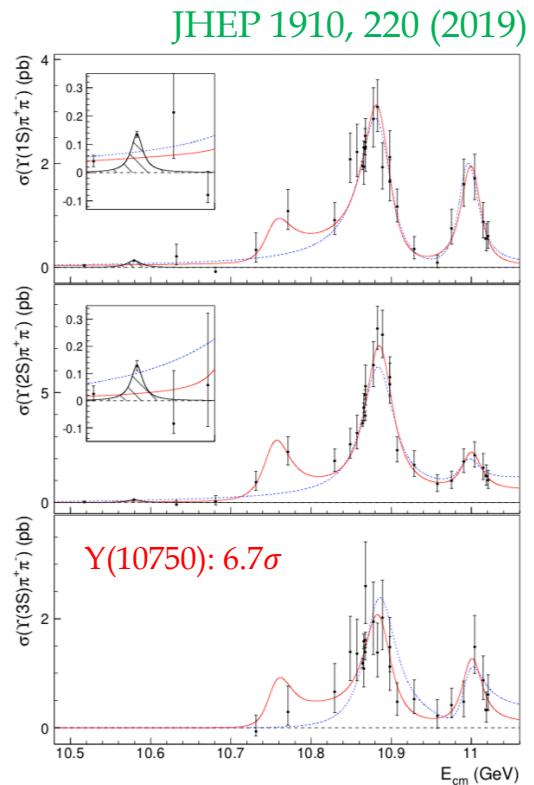
Three ways to access bottomonia below $B\bar{B}$ threshold:

- Decays of higher mass states (e.g. $\Upsilon(4S,5S,6S)$)
- Production of 1^{--} states via initial-state radiation
- Direct production via operation at a lower C.M. energy.

Predicted Missing bottomonium levels below $B\bar{B}$ threshold

Name	L	S	J^{PC}	Mass, MeV/ c^2	Emitted hadrons [Threshold, GeV/ c^2]
$\eta_b(3S)$	0	0	0^{-+}	10336	ω [11.12], ϕ [11.36]
$h_b(3P)$	1	0	1^{+-}	10541	$\pi^+\pi^-$ [10.82], η [11.09], η' [11.50]
$\eta_{b2}(1D)$	2	0	2^{-+}	10148	ω [10.93], ϕ [11.17]
$\eta_{b2}(2D)$	2	0	2^{-+}	10450	ω [11.23], ϕ [11.47]
$\Upsilon_J(2D)$	2	1	$(1, 2, 3)^{--}$	10441 – 10455	$\pi^+\pi^-$ [10.73], η [11.00], η' [11.41]
$h_{b3}(1F)$	3	0	3^{+-}	10355	$\pi^+\pi^-$ [10.63], η [10.90], η' [11.31]
$\chi_{bJ}(1F)$	3	1	$(2, 3, 4)^{++}$	10350 – 10358	ω [11.14], ϕ [11.38]
$\eta_{b4}(1G)$	4	0	4^{-+}	10530	ω [11.31], ϕ [11.55]
$\Upsilon_J(1G)$	4	1	$(3, 4, 5)^{--}$	10529 – 10532	$\pi^+\pi^-$ [10.81], η [11.08], η' [11.49]

Operation energies (in fb-1 (M events))

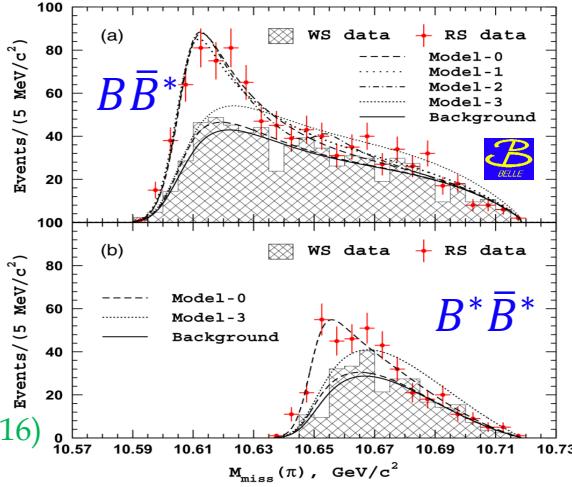


Experiment	$\Upsilon(1S)$	$\Upsilon(2S)$	$\Upsilon(3S)$	$\Upsilon(4S)$	$\Upsilon(5S)$	$\Upsilon(6S)$
CLEO	1.2 (21)	1.2 (10)	1.2 (5)	16 (17.1)	0.1 (0.4)	-
BaBar	-	14 (99)	30 (122)	433 (471)	R _b scan	R _b scan
Belle	6 (102)	25 (158)	3 (12)	711 (772)	121 (36)	5.5
Belle II	200	200	200	5×10^4 (5.4×10^4)	1000 (300)	100 + 400 (scan)

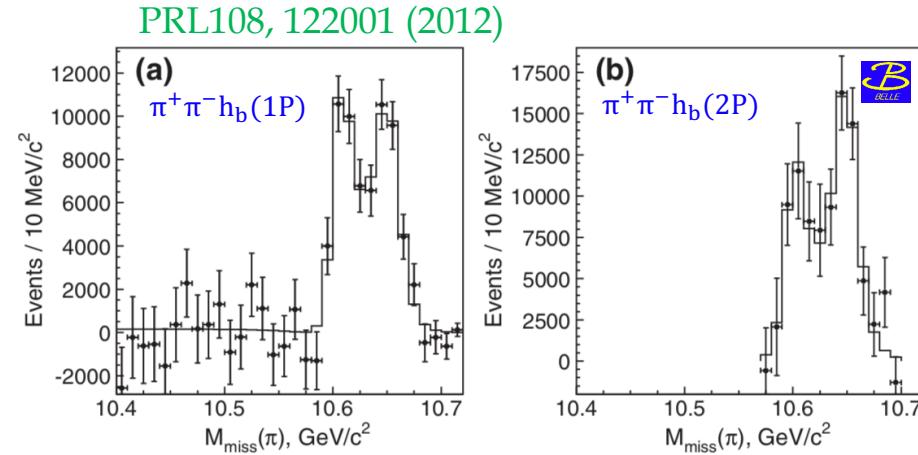
More data samples provide opportunities to explore bottomonium spectrum further. 36-

Search for new exotics at Belle II

- Observed $Z_b(10610)$ and $Z_b(10650)$ in $\Upsilon(5S,6S) \pi\pi$ transitions.
- The decays $Z_b(10610) \rightarrow B\bar{B}^*$ and $Z_b(10650) \rightarrow B^*\bar{B}^*$ are dominant.



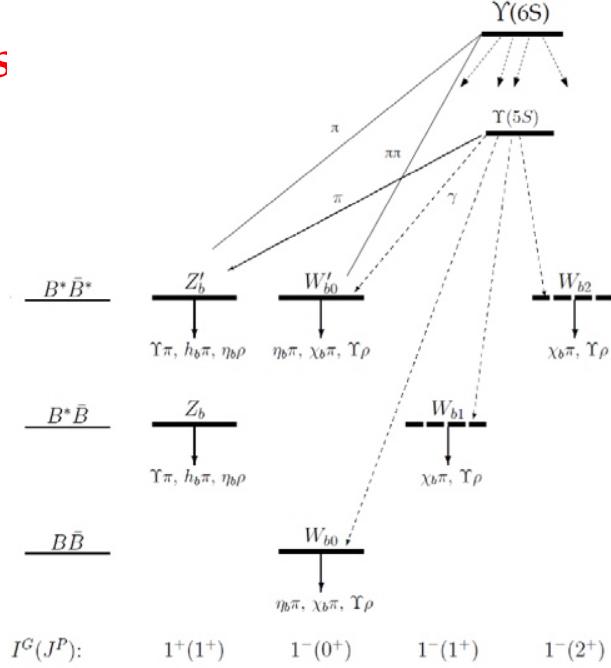
PRL116,
212001 (2016)



PRL108, 122001 (2012)

The expected molecular states with the structures $B\bar{B}$, $B\bar{B}^$, and $B^*\bar{B}^*$.*

$I^G(J^P)$	Name	Content	Co-produced particles [Threshold, GeV/c ²]	Decay channels
$1^+(1^+)$	Z_b	$B\bar{B}^*$	π [10.75]	$\Upsilon(nS)\pi, h_b(nP)\pi, \eta_b(nS)\rho$
$1^+(1^+)$	Z'_b	$B^*\bar{B}^*$	π [10.79]	$\Upsilon(nS)\pi, h_b(nP)\pi, \eta_b(nS)\rho$
$1^-(0^+)$	W_{b0}	$B\bar{B}$	ρ [11.34], γ [10.56]	$\Upsilon(nS)\rho, \eta_b(nS)\pi$
$1^-(0^+)$	W'_{b0}	$B^*\bar{B}^*$	ρ [11.43], γ [10.65]	$\Upsilon(nS)\rho, \eta_b(nS)\pi$
$1^-(1^+)$	W_{b1}	$B\bar{B}^*$	ρ [11.38], γ [10.61]	$\Upsilon(nS)\rho$
$1^-(2^+)$	W_{b2}	$B^*\bar{B}^*$	ρ [11.43], γ [10.65]	$\Upsilon(nS)\rho$
				<i>Mod. Phys. Lett. A</i> 32, 1750025 (2017)
$0^-(1^+)$	X_{b1}	$B\bar{B}^*$	η [11.15]	$\Upsilon(nS)\eta, \eta_b(nS)\omega$
$0^-(1^+)$	X'_{b1}	$B^*\bar{B}^*$	η [11.20]	$\Upsilon(nS)\eta, \eta_b(nS)\omega$
$0^+(0^+)$	X_{b0}	$B\bar{B}$	ω [11.34], γ [10.56]	$\Upsilon(nS)\omega, \chi_{bJ}(nP)\pi^+\pi^-, \eta_b(nS)\eta$
$0^+(0^+)$	X'_{b0}	$B^*\bar{B}^*$	ω [11.43], γ [10.65]	$\Upsilon(nS)\omega, \chi_{bJ}(nP)\pi^+\pi^-, \eta_b(nS)\eta$
$0^+(1^+)$	X_b	$B\bar{B}^*$	ω [11.39], γ [10.61]	$\Upsilon(nS)\omega, \chi_{bJ}(nP)\pi^+\pi^-$
$0^+(2^+)$	X_{b2}	$B^*\bar{B}^*$	ω [11.43], γ [10.65]	$\Upsilon(nS)\omega, \chi_{bJ}(nP)\pi^+\pi^-$



Summary

- Although Belle has stopped data taking for ~10 years ago, we are still producing exciting results.
- We reported the first vector charmonium-like state decaying to the charmed-antistrange and anticharmed-strange meson pairs $D_s^+ D_{s1}(2536)^-$ and $D_s^+ D_{s2}^*(2573)^-$ with a signal significance of 5.9σ and 3.4σ . The masses and widths are close to those of the Y(4660).
- Belle II will reach 50 ab^{-1} by 2027, which will provide greater sensitivity and precise measurements in hadron physics.

Thanks for your attentions!

Backup

Table 1

Possible S-wave resonances with two D_s mesons below 4.8 GeV. Only those states with $C = +$, permitted to couple to $J/\psi \phi$, are shown. Here $\Delta M \equiv M - M(J/\psi) - M(\phi)$. Exchanges are labeled “d” for direct ($A\bar{B} \rightarrow A\bar{B}$) and “x” for exchange ($A\bar{B} \rightarrow B\bar{A}$). “Rank” denotes anticipated ease of detection (see text). Allowed J^P values and rank are only shown for those states which can bind via η exchange.

States (J^P)	M (MeV)	ΔM (MeV)	Binding		Allowed J^P	Rank
			by η ?	by $f_0(980)$?		
$D_s^+(0^-) D_s^-(0^-)$	3936.6	-179.8	No	d	-	-
$D_s^+(0^-) D_s^{*-}(1^-)$	4080.4	-36.0	x	d	1^+	1
$D_s^{*+}(1^-) D_s^{*-}(1^-)$	4224.2	107.8	d	d	$0^+, 2^+$ ^a	2
$D_s^+(0^-) D_{s0}^{*-}(2317)(0^+)$	4286.0	169.6	x	d	0^-	2
$D_s^+(0^-) D_{s1}^-(2460)(1^+)$	4427.8	311.4	No ^b	d, x	$[1^-]$ ^b	2
$D_s^{*+}(1^-) D_{s0}^{*-}(2317)(0^+)$	4429.8	313.4	No ^b	d, x	$[1^-]$ ^b	3
$D_s^+(0^-) D_{s1}^-(2536)(1^+)$	4503.4	387.0	No	d, x	-	-
$D_s^+(0^-) D_{s2}^{*-}(2573)(2^+)$	4540.2	423.8	x	d	2^-	2
$D_s^{*+}(1^-) D_{s1}^-(2460)(1^+)$	4571.6	455.2	d, x	d, x	$0^-, 1^-, 2^-$	3
$D_{s0}^{*-}(2317)(0^+) D_{s0}^{*-}(2317)(0^+)$	4635.4	519.0	No	d	-	-
$D_s^{*+}(1^-) D_{s1}^-(2536)(1^+)$	4647.2	530.8	d, x	d, x	$0^-, 1^-, 2^-$	3
$D_s^{*+}(1^-) D_{s2}^{*-}(2573)(2^+)$	4684.0	567.6	d, x	d, x	$1^-, 2^-, 3^-$	3
$D_{s0}^{*-}(2317)(0^+) D_{s1}^-(2460)(1^+)$	4777.2	660.8	x	d	1^+	4

^a $J^P = 1^+$ forbidden by C symmetry.

^b Proximity of these two channels may lead to binding. See text.