

Lattice calculation powered by factorization theory

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Outline

I. Hadron structure from experimental data

II. Hadron structure from lattice: ideas

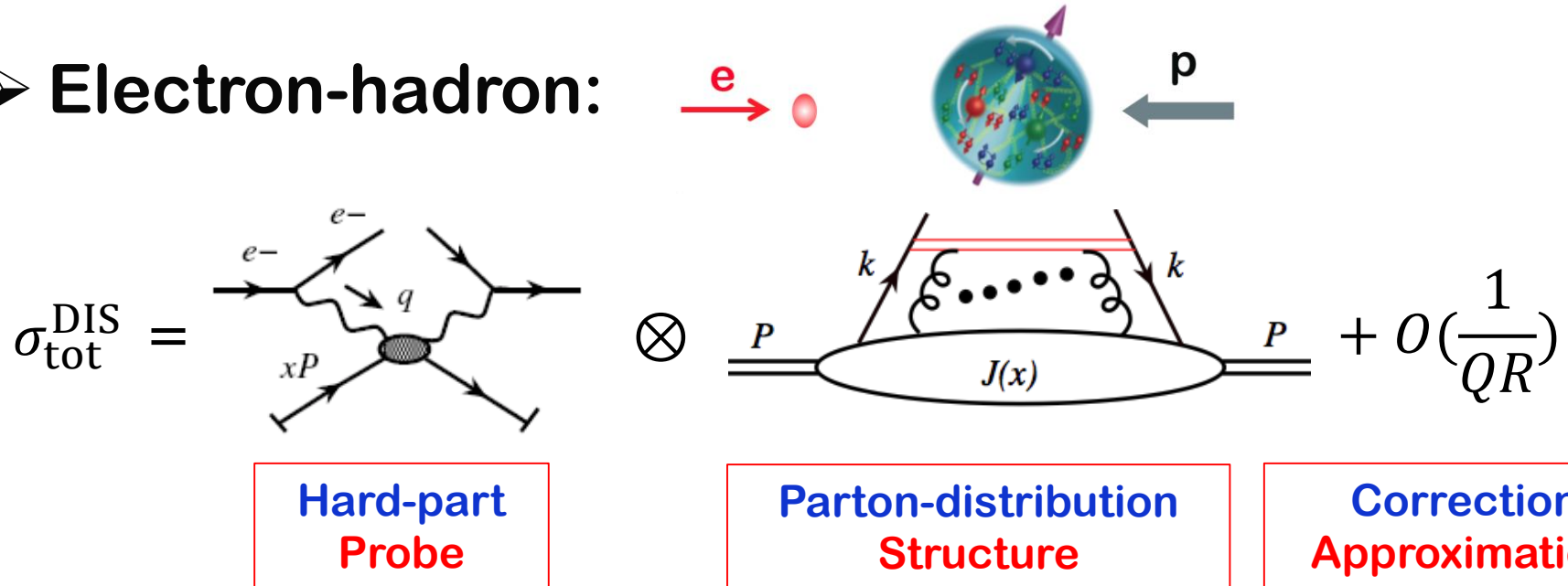
III. Hadron structure from lattice: details

IV. Summary and outlook

QCD factorization

➤ The key and a first principle method to relate experimental data to QCD theory

➤ Electron-hadron:



➤ PDFs/TMDs/...: encoding nonperturbative information in hadrons



Properties of PDFs

- **Well defined in QCD; process independent**

Spin-averaged quark distribution

$$f_{q/p}(x, \mu^2) = \int \frac{d\xi_-}{4\pi} e^{-ix\xi_- P_+} \langle P | \bar{\psi}(\xi_-) \gamma_+ \exp \left\{ -ig \int_0^{\xi_-} d\eta_- A_+(\eta_-) \right\} \psi(0) | P \rangle$$

- **Defined in infinite frame: not direct physical observable**

Determined by using factorization

- **Logarithmic UV divergent, renormalizable**

DGLAP evolution equations

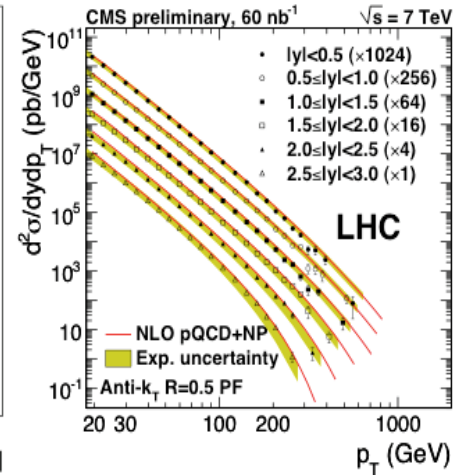
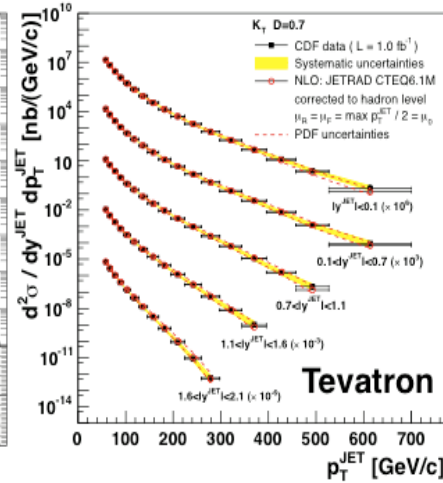
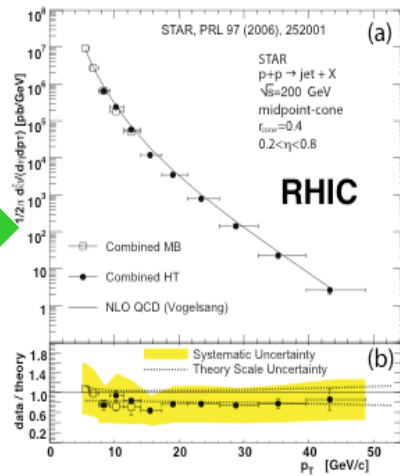
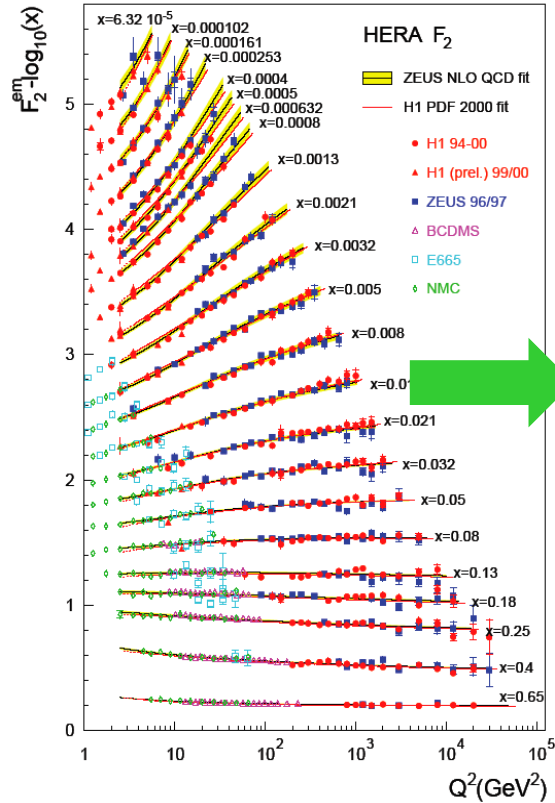
- **Operator defining PDFs: time dependent!**



Extract PDFs by fitting data

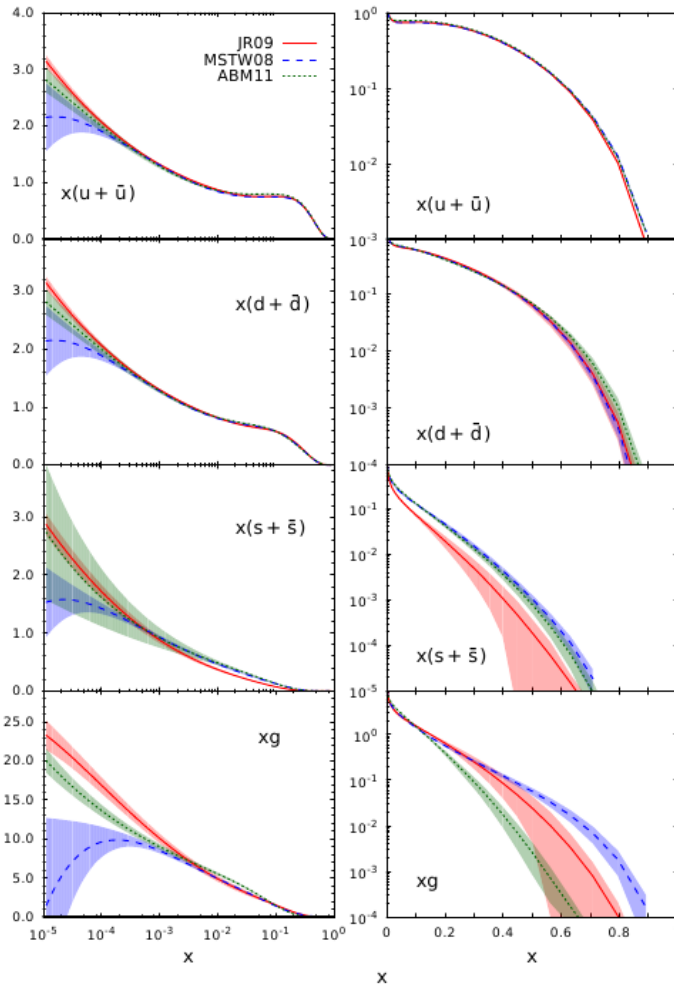
➤ **Successful**

Measure e-p at 0.3 TeV (HERA)
 Predict p-p at 0.2, 1.96, and 7 TeV





Uncertainty of PDFs



Large uncertainty in both small-x and large-x region



Questions

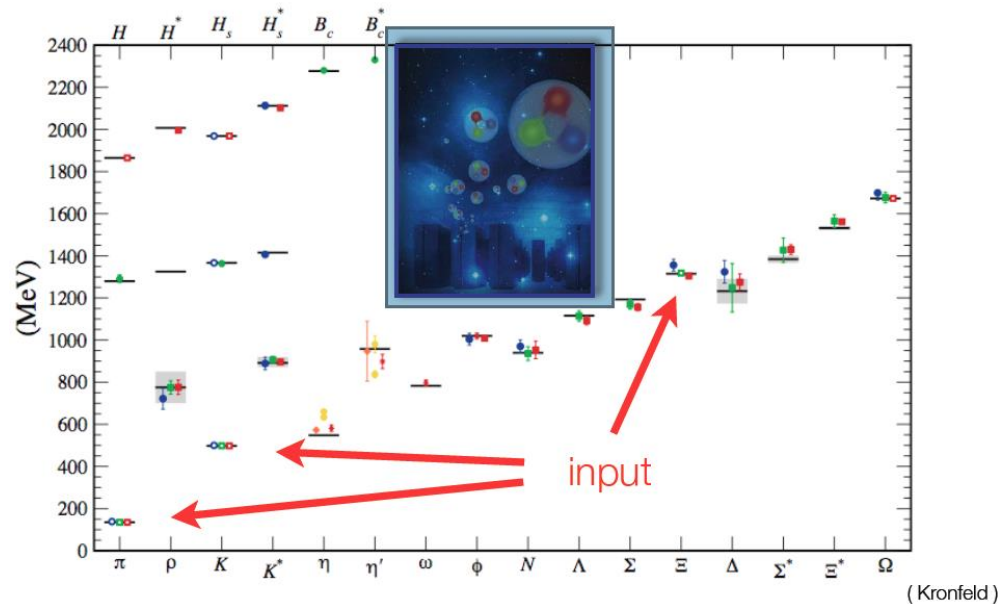
- **Can theoretical calculation verify the extracted values?**
- **Can theoretical calculation provide more information for hadron structure?**

**Need to calculate PDFs/TMDs/...
nonperturbatively from first principle!**



Lattice QCD

- The main nonperturbative approach to solve QCD
- Predict the hadron mass

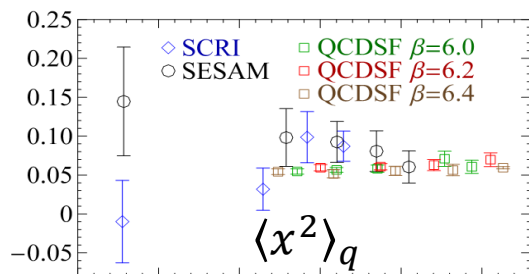


- Intrinsically Euclidean time: $\tau = i t$

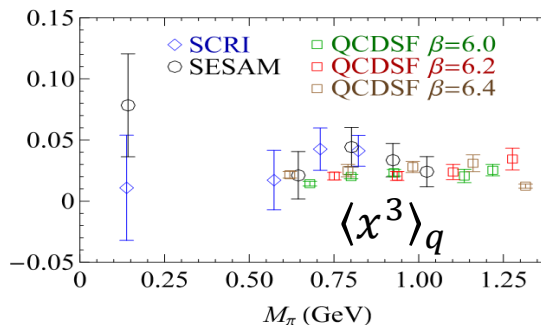


Difficulty of PDFs calculation

- PDFs have (Minkowski) time dependence, lattice QCD cannot calculate PDFs directly
- Very limited moments of PDFs



Dolgov et al., hep-lat/0201021



Gockeler et al., hep-ph/0410187

$$\langle x^n(\mu^2) \rangle_q \equiv \int_0^1 dx x^n f_{q/p}(x, \mu^2)$$



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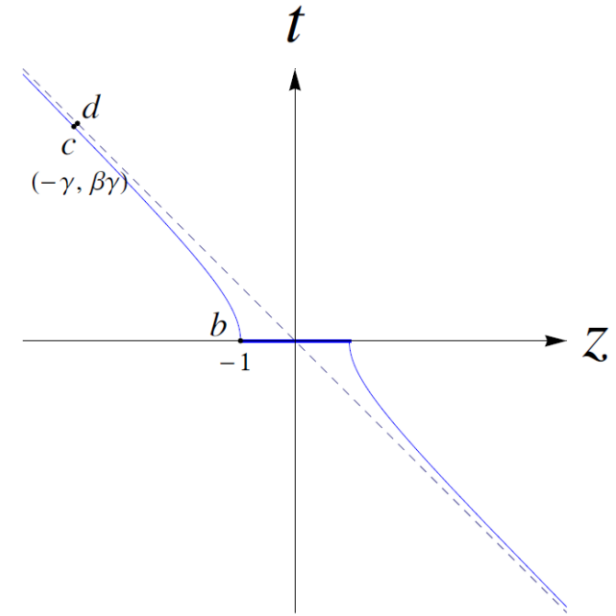
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Quasi-PDFs

➤ What if quark bilinear is slightly off light cone?

- Exist a frame where quark bilinear is equal time, but proton is moving fast



➤ Quasi-PDFs Ji, 1305.1539

$$\tilde{f}_{q/p}(x, \mu^2, P_z) = \int \frac{d\xi_z}{4\pi} e^{-ix\xi_z P_z} \langle P | \bar{\psi}(\xi_z) \gamma_z \exp \left\{ -ig \int_0^{\xi_z} d\eta_z A_z(\eta_z) \right\} \psi(0) | P \rangle$$

- **Classical picture:** When proton is moving fast enough, i.e. $P_z \rightarrow \infty$, there exists a frame so that quark bilinear in quasi-PDFs are only slightly off light cone, approach PDFs



➤ Advantage

- Fields separated along the z-direction, no time dependence, calculable using standard lattice method

➤ Difficulty

- Quantum fluctuation, UV divergences, does the simple classical picture still hold in QFT?
- P_z on lattice is very limited
- In fact, relies on the existence of factorization

$$q(x, \mu^2, P^z) = \int_x^1 \frac{dy}{y} Z\left(\frac{x}{y}, \frac{\mu}{P^z}\right) q(y, \mu^2) + \mathcal{O}\left(\Lambda^2/(P^z)^2, M^2/(P^z)^2\right)$$

Ji, 1305.1539

$$\sigma(P \cdot \xi, \xi^2, \tilde{\mu}^2) \approx \sum_{i=q, \bar{q}, g} \int_0^1 \frac{dx}{x} f_i(x, \mu^2) C_i(xP \cdot \xi, \xi^2, \tilde{\mu}^2, \mu^2) + \mathcal{O}(\xi^2)$$

$$\tilde{f}_{q/h}^{\text{ren}} = f_{i/h} \otimes C_{q/i} + \mathcal{O}((\tilde{x}P_z)^{-2})$$

YQM, Qiu, 1404.6860

$$\text{quasi } \mathcal{Q}(x, p) = q(x) \left\{ 1 + \mathcal{O}\left(\frac{\Lambda^2}{p^2 x^2 (1-x)}\right) \right\},$$

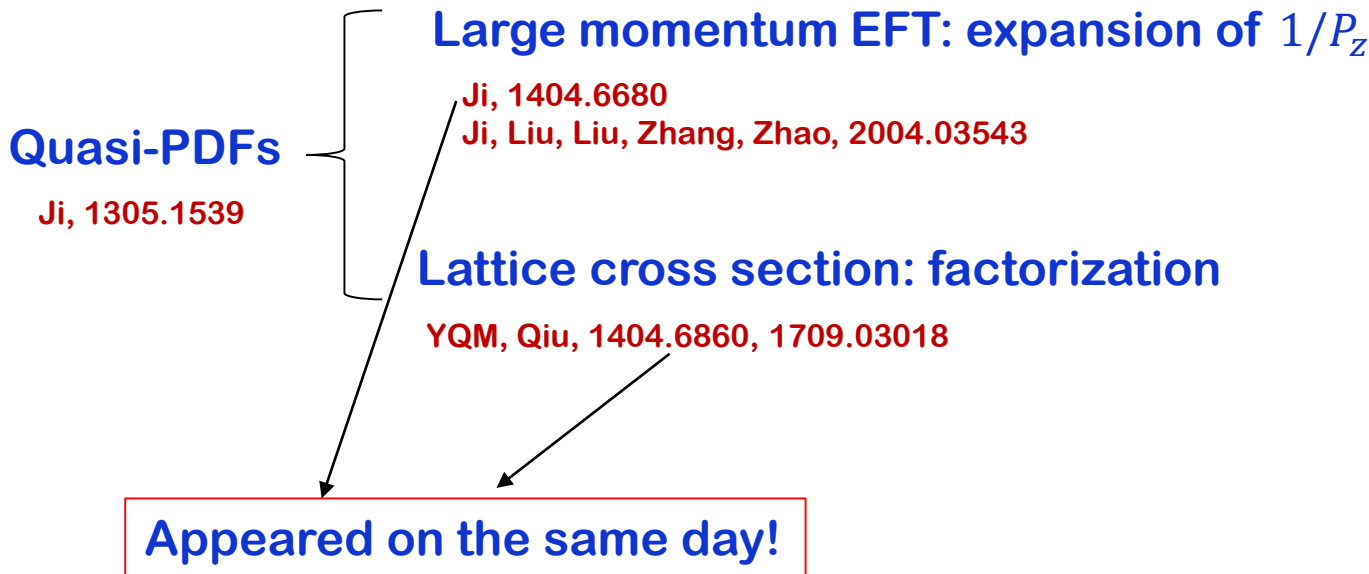
$$\text{pseudo } \mathcal{P}(x, z) = q(x) \left\{ 1 + \mathcal{O}(z^2 \Lambda^2 (1-x)) \right\},$$

Braun, Vladimirov, Zhang, 1810.00048



Generalization

- Why quasi-PDFs work? Is there a more general way to calculate PDFs?
- From quasi-PDFs to LaMET and LCS





The idea of lattice cross sections

YQM, Qiu, 1404.6860, 1709.03018

➤ Direct calculation is hard

- Lattice: time-dependence of operators defining PDFs
- LaMET: P_z cannot be too large on lattice

➤ The most general way of indirect calculation

- Relating PDFs, **not calculable** on lattice QCD, to some quantities **calculable** on lattice QCD
- The relation is nothing but a factorization!

$$\sigma_n(\xi^2, \omega, P^2) = \sum_a \int_{-1}^1 \frac{dx}{x} f_a(x, \mu^2) K_n^a(\xi^2, x\omega, x^2 P^2, \mu^2) + O(\xi^2 \Lambda_{\text{QCD}}^2)$$

Parton-distribution
Structure

Hard-part
Probe

Correction
Approximation

- Analogy: measured experimental cross section!
- We call these quantities: “lattice cross sections”



What kinds of LCSs are useful?

➤ Conditions for a good LCS

- ① Calculable on Euclidean lattice QCD
- ② Renormalizable for UV divergences
- ③ Factorizable for CO divergence with IR safe coefficients

$$\sigma_n(\xi^2, \omega, P^2) = \sum_a \int_{-1}^1 \frac{dx}{x} f_a(x, \mu^2) K_n^a(\xi^2, x\omega, x^2 P^2, \mu^2) + O(\xi^2 \Lambda_{\text{QCD}}^2)$$

- The last condition relates LCSs to, e.g., PDFs

➤ First condition: at least two possibilities

$$\sigma_n(\xi^2, \omega, P^2) = \langle P | T \{ \mathcal{O}_n(\xi) \} | P \rangle \quad \text{with } \xi_0 = 0, \text{ e.g., quasi-PDFs}$$

Lin, et. al. , 1402.1462

$$\tilde{\sigma}_n(q^2, \tilde{\omega}, P^2) = \int \frac{d^4 \xi}{\xi^4} e^{iq \cdot \xi} \sigma_n(\xi^2, P \cdot \xi, P^2) \quad \text{with } q_0 = 0, \text{ e.g., OPE w/o OPE}$$

Chambers, et. al. , 1703.01153

➤ Second and third conditions: need to prove



LCSs to determine PDFs

➤ LCSs in coordinate space or in momentum space

$$\sigma_n(\xi^2, \omega, P^2) = \langle P | T \{ \mathcal{O}_n(\xi) \} | P \rangle \quad \tilde{\sigma}_n(q^2, \tilde{\omega}, P^2) = \int \frac{d^4 \xi}{\xi^4} e^{iq \cdot \xi} \sigma_n(\xi^2, P \cdot \xi, P^2)$$

- $1/\xi^2$ and q^2 : hard scales to enable factorization

➤ Possible choices of nonlocal operators

Parton correlators

$$\mathcal{O}_q(\xi) = Z_q(\xi^2) \bar{\psi}_q(\xi) \not{x} \Phi(\xi, 0) \psi_q(0)$$

$$\Phi(\xi, 0) = \mathcal{P} e^{-ig \int_0^1 \xi \cdot A(\lambda \xi) d\lambda}$$

Ji, 1305.1539

- Path ordered gauge link needed
- Renormalization is complicated

Current-current operators

$$\mathcal{O}_S(\xi) = \xi^4 Z_S^2 [\bar{\psi}_q \psi_q](\xi) [\bar{\psi}_q \psi_q](0),$$

$$\mathcal{O}_V(\xi) = \xi^2 Z_V^2 [\bar{\psi}_q \not{x} \psi_q](\xi) [\bar{\psi}_q \not{x} \psi_q](0),$$

$$\mathcal{O}_{\tilde{V}}(\xi) = -\frac{\xi^4}{2} Z_V^2 [\bar{\psi}_q \gamma_\nu \psi_q](\xi) [\bar{\psi}_q \gamma^\nu \psi_q](0),$$

$$\mathcal{O}_{V'}(\xi) = \xi^2 Z_{V'}^2 [\bar{\psi}_q \not{x} \psi_{q'}](\xi) [\bar{\psi}_{q'} \not{x} \psi_q](0), \dots,$$

YQM, Qiu, 1709.03018

- Renormalization is very simple

Straight forward to construct much more operators with both quark fields and gluon fields



Different methods

➤ Inverse Laplace transform for structure functions

Liu, Dong, 9306299, Liu, 9910306, 1603.07352

➤ Coordinate-space method for LDA

Aglietti, et. al., 9806277; Abada, et.al., 0105221; Braun, Mueller, 0709.1348

➤ Quasi-PDFs

Ji, 1305.1539, 1404.6680

➤ Pseudo-DFs

Orginos, Radyushkin, Karpie, Zafeiropoulos, 1706.05373

➤ OPE without OPE

Chambers, et. al. , 1703.01153

➤ Current-current correlators

YQM, Qiu, 1709.03018

All existed methods rely on factorization, can be interpreted as constructing LCSs

Analogy: meat, vegetables, fruits are all food



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Renormalization

➤ Why proof is important?

- Need to take continuum limit for lattice calculation
- All-order proof of factorization needs multiplicative renormalization
- Find out all operators mixing under renormalization YQM, Qiu, 1404.6860

Lessons from gluon-gluon correlator: proof is crucial!

➤ Current-current operators: known

➤ Parton correlators: difficult

- Because of z -direction dependence, Lorentz symmetry is broken, hard to exhaust all possible UV divergences
- Renormalization of nonlocal composite operator



All-order proofs for renormalization

- **Diagrammatic method: key is to show that UV divergences are local in space-time**
 - Nontrivial conclusion! E.g. UV divergences for normal PDFs are non-local in “–” direction
 - Ishikawa YQM, Qiu, Yoshida, 1707.03107
 - Li, YQM, Qiu, 1809.01836
 - The most difficult part in this proof
- **Auxiliary field method: gauge link obtained by integration out an auxiliary field in HQET**
 - Taking advantage of the renormalizability of heavy-light current in HQET
 - Ji, Zhang, Zhao, 1706.08962
 - Green, Jansen, Steffens, 1707.07152
 - Zhang, et. al., 1808.10824



All-order proofs for factorization

➤ Diagrammatic method:

YQM, Qiu, 1404.6860

- Rigorous, but hard to resum target mass corrections

➤ OPE method:

YQM, Qiu, 1709.03018
Izubuchi, et. al., 1801.03917

- For quasi-PDFs, no proof of OPE: only a derivation of factorization

$$\sigma_n(\xi^2, \omega, P^2) = \sum_a \int_{-1}^1 \frac{dx}{x} f_a(x, \mu^2) K_n^a(\xi^2, x\omega, x^2 P^2, \mu^2) + O(\xi^2 \Lambda_{\text{QCD}}^2)$$

$$\tilde{\sigma}_n = \sum_a f_a \otimes \tilde{K}_n^a + O(\Lambda_{\text{QCD}}^2/q^2)$$

$$f_{\bar{a}/h}(x, \mu^2) = -f_{a/h}(-x, \mu^2)$$

where

$$K_n^a = \sum_J 2W_n^{(J,a)}(\xi^2, \mu^2) \Sigma_J(x\omega, x^2 P^2 \xi^2)$$

$$\tilde{K}_n^a = \int \frac{d^4\xi}{\xi^4} e^{iq \cdot \xi} K_n^a(\xi^2, xP \cdot \xi, x^2 P^2, \mu)$$



Theoretical status

➤ Conditions satisfied up to now

- With $\xi_0 = 0$ or $q_0 = 0$, calculable on Euclidean lattice
- Renormalizable
- Factorizable to PDFs

With these conditions, σ_n and $\tilde{\sigma}_n$ are good LCSs to extract PDFs

Questions still need to answer:

- 1) OPE for parton-parton correlators?
- 2) Renormalization using lattice regulator?

➤ What next?

- Calculate matching coefficients perturbatively
- Calculate LCSs nonperturbatively (including renormalization methods)



Matching coefficients

➤ Obtained by calculating Feynman diagrams

- Most quantities: calculated to one-loop order
- Flavor nonsinglet quark-quark correlator (qPDF, pPDF): two-loop order

Li, YQM, Qiu, 2006.12370
Chen, Wang, Zhu, 2006.14825

High-order calculations for more quantities are needed,
make lattice data more useful

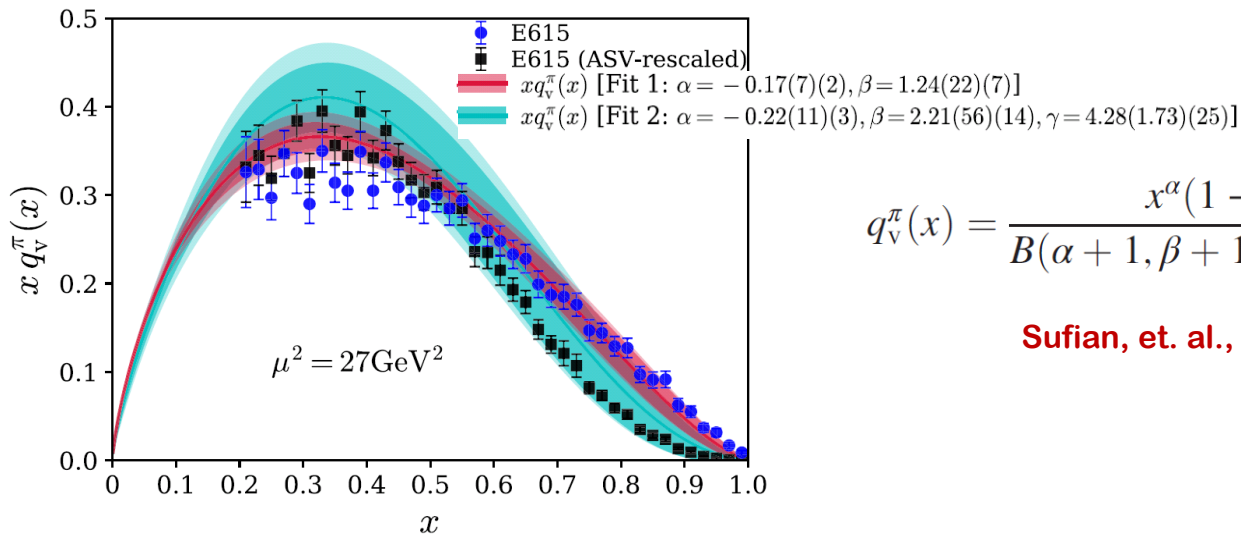
➤ Renormalization schemes

See Jian-Hui Zhang's talk



Lattice results

- LaMET: see many talks in this meeting
- Many interesting results from pseudo-PDFs
- An example using current-current correlator



$$q_V^\pi(x) = \frac{x^\alpha(1-x)^\beta(1+\gamma x)}{B(\alpha+1, \beta+1) + \gamma B(\alpha+2, \beta+1)}$$

Sufian, et. al., 2001.04960



Summary

- Hadronic matrix elements that are calculable + renormalizable + factorizable
 - Like experimental cross section, can constrain hadron structure
- Quasi-PDFs, pseudo-PDFs, OPE w/o OPE, current-current correlators: good LCSs
- Construct better operators?
- Many (non-)perturbative works to do

With world (lattice and experimental) data, a global fit to determine PDFs/TMDs/...

Thank you!