

# MCs & Precision QCD at Future $e^+e^-$ Machines

Peter Skands (Monash U)

## Perturbative QCD: **High Accuracy**

Expect a new generation of precision showers merged through (N)NLO

## Nonperturbative QCD: **High Resolution**

Next generation of  $e^+e^-$  machines  $\rightarrow$  **trial by fire** not just for any post-LHC advanced hadronisation models, but also for *any future solution* (or *systematically improvable approximation*) **to the problem of confinement.**

$\rightarrow$  Need **Good PID & Good Momentum Resolution**  $\ll O(\Lambda_{\text{QCD}}) \sim 100 \text{ MeV}$

**+ Synergies with EW & Higgs Physics Goals** (MC uncertainties)



CEPC Workshop  
October 2020, Shanghai

# MC Generators — Perturbative Processes

Slide borrowed from A. Hoang (yesterday's EW session)

- Fast machinery from LHC, just change initial state
- Less modeling for color neutralization processes needed
- NLO-matched MC generators standard.

**Validation of NLO QCD for  $e^+e^-$  Collisions**

Process	$\sigma^{\text{LO}}$ [fb]	MG5_AMC $\sigma^{\text{NLO}}$ [fb]	$K$	$\sigma^{\text{LO}}$ [fb]	WHIZARD $\sigma^{\text{NLO}}$ [fb]	$K$
$e^+e^- \rightarrow jj$	622.3(5)	639.3(1)	1.02733	622.73(4)	639.41(9)	1.02678
$e^+e^- \rightarrow jjj$	340.1(2)	317.3(8)	0.93297	342.4(5)	318.6(7)	0.9305
$e^+e^- \rightarrow jjjj$	104.7(1)	103.7(3)	0.99045	105.1(4)	103.0(6)	0.98003
$e^+e^- \rightarrow jjjjj$	22.11(6)	24.65(4)	1.11488	22.80(2)	24.35(15)	1.06798
$e^+e^- \rightarrow jjjjjj$	N/A	N/A	N/A	3.62(2)	0.0(0)	0.0
$e^+e^- \rightarrow b\bar{b}$	92.37(6)	94.89(1)	1.02728	92.32(1)	94.78(7)	1.02664
$e^+e^- \rightarrow b\bar{b}b\bar{b}$	$1.644(3) \cdot 10^{-1}$	$3.60(1) \cdot 10^{-1}$	2.1897	$1.64(2) \cdot 10^{-1}$	$3.67(4) \cdot 10^{-1}$	2.2378
$e^+e^- \rightarrow t\bar{t}$	166.2(2)	174.5(3)	1.04994	166.4(1)	174.53(6)	1.04886
$e^+e^- \rightarrow t\bar{t}j$	48.13(5)	53.36(1)	1.10867	48.3(2)	53.25(6)	1.10248
$e^+e^- \rightarrow t\bar{t}jj$	8.614(9)	10.49(3)	1.21777	8.612(8)	10.46(6)	1.21458
$e^+e^- \rightarrow t\bar{t}jjj$	1.044(2)	1.420(4)	1.3601	1.040(1)	1.414(10)	1.3595
$e^+e^- \rightarrow t\bar{t}t\bar{t}$	$6.45(1) \cdot 10^{-4}$	$11.94(2) \cdot 10^{-4}$	1.85117	$6.463(2) \cdot 10^{-4}$	$11.91(2) \cdot 10^{-4}$	1.8428
$e^+e^- \rightarrow t\bar{t}t\bar{t}j$	$2.719(5) \cdot 10^{-5}$	$5.264(8) \cdot 10^{-5}$	1.93602	$2.722(1) \cdot 10^{-5}$	$5.250(14) \cdot 10^{-5}$	1.92873
$e^+e^- \rightarrow t\bar{t}b\bar{b}$	0.1819(3)	0.292(1)	1.60533	0.186(1)	0.293(2)	1.57527
$e^+e^- \rightarrow t\bar{t}H$	2.018(3)	1.909(3)	0.94601	2.022(3)	1.912(3)	0.9456
$e^+e^- \rightarrow t\bar{t}Hj$	$0.2533(3) \cdot 10^{-0}$	$0.2665(6) \cdot 10^{-0}$	1.05212	0.2540(9)	0.2664(5)	1.04889
$e^+e^- \rightarrow t\bar{t}Hjj$	$2.663(4) \cdot 10^{-2}$	$3.141(9) \cdot 10^{-2}$	1.1795	$2.666(4) \cdot 10^{-2}$	$3.144(9) \cdot 10^{-2}$	1.17928
$e^+e^- \rightarrow t\bar{t}\gamma$	12.7(2)	13.3(4)	1.04726	12.71(4)	13.78(4)	1.08418
$e^+e^- \rightarrow t\bar{t}Z$	4.642(6)	4.95(1)	1.06636	4.64(1)	4.94(1)	1.06467
$e^+e^- \rightarrow t\bar{t}Zj$	0.6059(6)	0.6917(24)	1.14168	0.610(4)	0.6927(14)	1.13565
$e^+e^- \rightarrow t\bar{t}Zjj$	$6.251(28) \cdot 10^{-2}$	$8.181(21) \cdot 10^{-2}$	1.30875	$6.233(8) \cdot 10^{-2}$	$8.201(14) \cdot 10^{-2}$	1.31573
$e^+e^- \rightarrow t\bar{t}W^\pm jj$	$2.400(4) \cdot 10^{-4}$	$3.714(8) \cdot 10^{-4}$	1.54747	$2.41(1) \cdot 10^{-4}$	$3.695(9) \cdot 10^{-4}$	1.5332
$e^+e^- \rightarrow t\bar{t}\gamma\gamma$	0.383(5)	0.416(2)	1.08618	0.382(3)	0.420(3)	1.09952
$e^+e^- \rightarrow t\bar{t}\gamma Z$	0.2212(3)	0.2364(6)	1.06873	0.220(1)	0.240(2)	1.09094
$e^+e^- \rightarrow t\bar{t}\gamma H$	$9.75(1) \cdot 10^{-2}$	$9.42(3) \cdot 10^{-2}$	0.96614	$9.748(6) \cdot 10^{-2}$	$9.58(7) \cdot 10^{-2}$	0.98277
$e^+e^- \rightarrow t\bar{t}ZZ$	$3.788(4) \cdot 10^{-2}$	$4.00(1) \cdot 10^{-2}$	1.05597	$3.756(4) \cdot 10^{-2}$	$4.005(2) \cdot 10^{-2}$	1.0663
$e^+e^- \rightarrow t\bar{t}W^+W^-$	0.1372(3)	0.1540(6)	1.1225	0.1370(4)	0.1538(4)	1.12257
$e^+e^- \rightarrow t\bar{t}HH$	$1.358(1) \cdot 10^{-2}$	$1.206(3) \cdot 10^{-2}$	0.888	$1.367(1) \cdot 10^{-2}$	$1.218(1) \cdot 10^{-2}$	0.8909
$e^+e^- \rightarrow t\bar{t}HZ$	$3.600(6) \cdot 10^{-2}$	$3.58(1) \cdot 10^{-2}$	0.99445	$3.596(1) \cdot 10^{-2}$	$3.581(2) \cdot 10^{-2}$	0.9958

Just pick what you need!

Not so fast..

# MC Generators — How precise are they?

Slide borrowed from A. Hoang (yesterday's EW session)

- Multipurpose MC generators (Pythia, Herwig, Whizard, Sherpa) can simulate all aspects of particle production and decay at the observable level

## How precise are they?

- The theoretical precision is tied to the precision of the parton showers, for a few very simple observable NLL, mostly LL or less. (Though showers do include some further all-orders aspects, such as exact conservation of energy and momentum, not accounted for in this counting.)
- Tuned hadronization models compensate (partly) for the deficiency but scale differently with  $\sqrt{s} \implies$  scaling studies
- In general we have 

observable precision	>	theoretical precision
----------------------	---	-----------------------

  
 CEPC  $\triangleright$  high statistics from 10 - 250 GeV (via ISR from Z pole)
- MCs are <sup>currently</sup> not very precise tools to extract QCD parameters or provide estimate of hadronization corrections to high-order perturbative analytical calculations
- NLO-matching does only improve the first hard gluon radiation. Does not improve observables governed by parton shower dynamics.

My additions

# MC Generators ► Next Generation

Slide borrowed from A. Hoang (yesterday's EW session)

- NLL precise parton showers with full coherence and improved models are an important step that needs to be taken (many different aspects, work already ongoing).

e.g. second order kernel

double emission

amplitude evolution (full coherence,  
non-global logs, color reconnection)

Li, Skands '16

Höche Prestel' 14, '15

Forshaw, Holguin, Plätzer '19

Gieseke, Kirchgaesser, Plätzer, Siodmok '19

Martinez, Forshaw, De Angelis, Plätzer,  
Seymour '18

New generation of MCs needed!

→ Definitely possible, community should support it more enthusiastically.

First shower models (Leading Log, Leading Colour) ~ 1980.  
40 years later, now at the threshold of the next **major** breakthrough!

# Second-Order Shower Kernels?



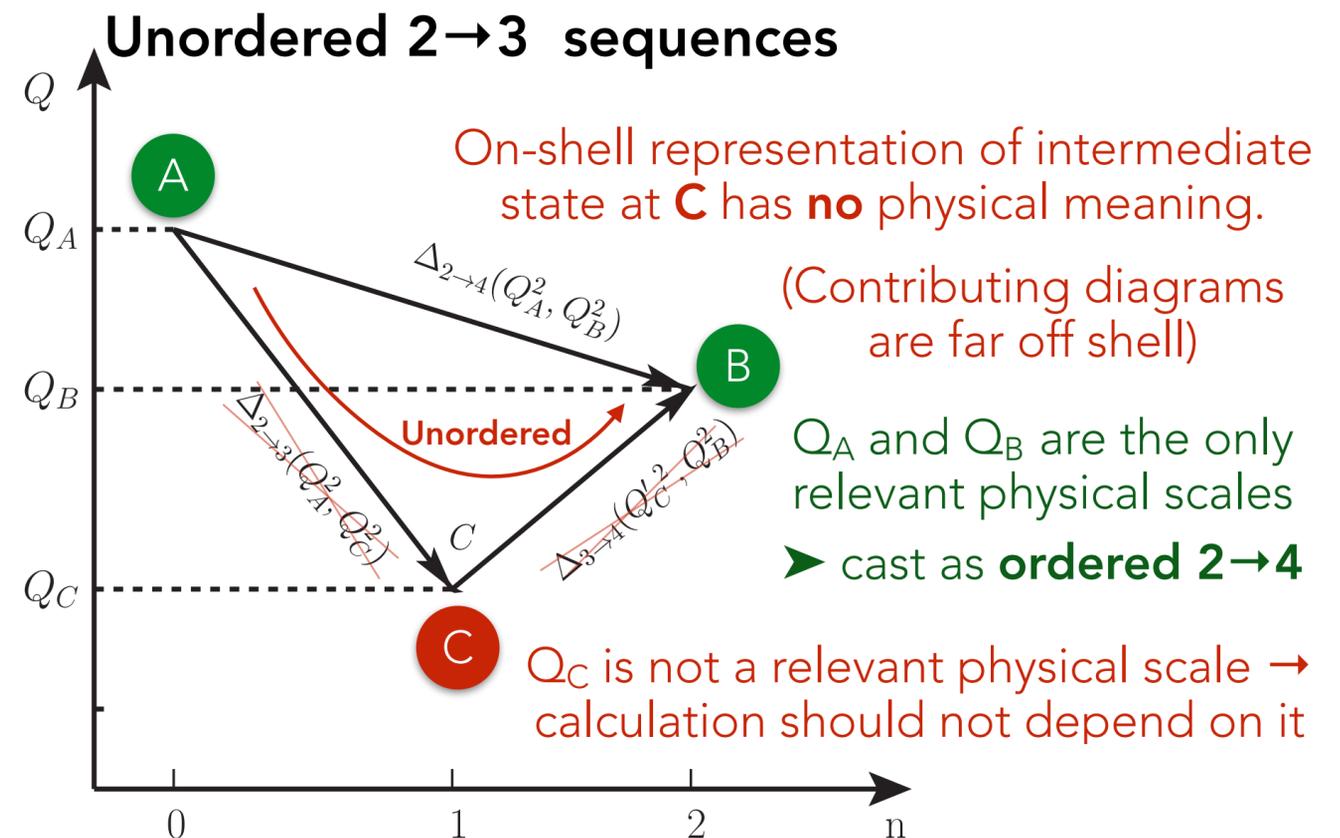
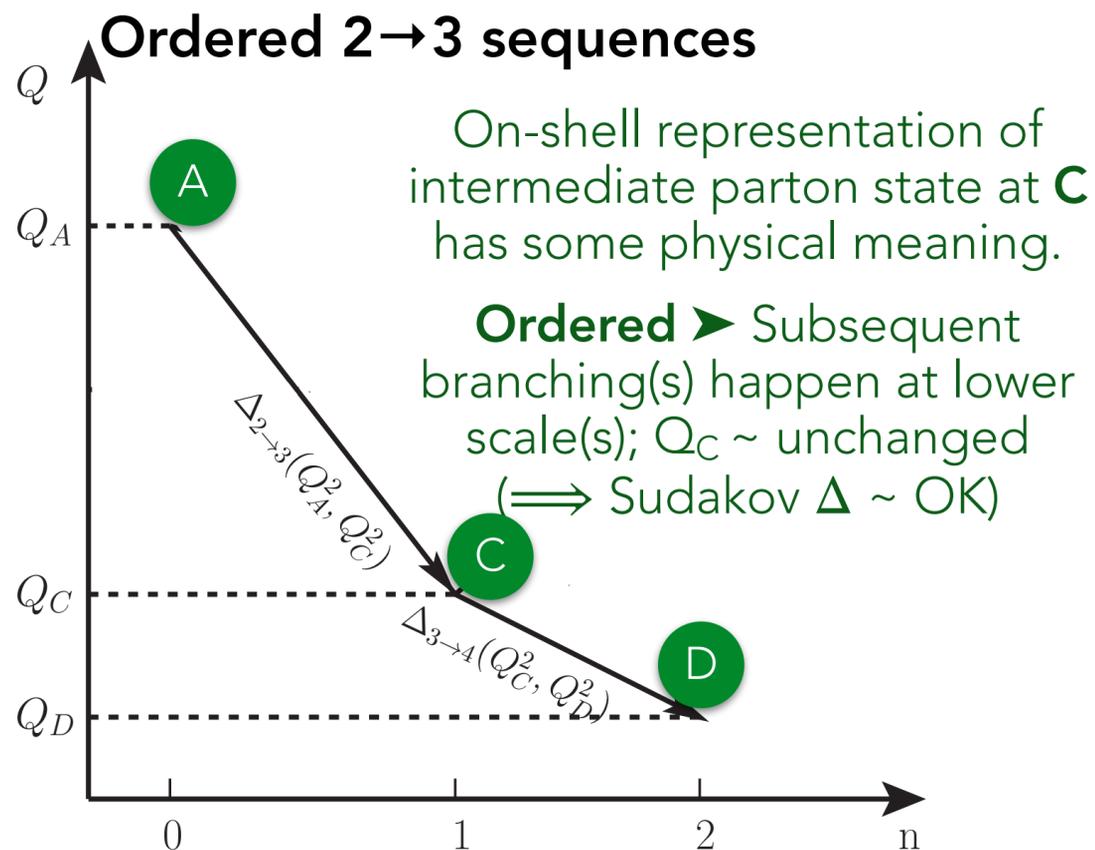
Li & PS, *PLB* 771 (2017) 59 (arXiv:1611.00013) + ongoing work

## Elements

Iterated dipole-style  $2 \rightarrow 3$  and new "direct  $2 \rightarrow 4$ " branchings populate complementary phase-space regions.

Ordered clustering sequences  $\rightarrow$  iterated  $2 \rightarrow 3$  (+ virtual corrections  $\sim$  differential K-factors)

Unordered clustering sequences  $\rightarrow$  direct  $2 \rightarrow 4$  (+ in principle higher  $2 \rightarrow n$ , ignored for now)



**Our approach:** continue to exploit iterated on-shell  $2 \rightarrow 3$  factorisations ...

... but in **unordered region** let  $Q_B$  define evolution scale for double-branching (integrate over  $Q_C$ )

# Second-Order Shower Evolution Equation

Li & PS, *PLB* 771 (2017) 59 (arXiv:1611.00013) + ongoing work

Putting 2→3 and 2→4 together ⇔ evolution equation for dipole-antenna with  $\mathcal{O}(\alpha_s^2)$  kernels:

~ POWHEG inside exponent  
(Hoeche, Krauss, Prestel ~ MC@NLO inside exponent)

Iterated 2→3  
with (finite) one-loop correction

Direct 2→4  
(as sum over "a" and "b" subpaths)

$$\frac{d\Delta(Q_0^2, Q^2)}{dQ^2} = \int d\Phi_{\text{ant}} \left[ \delta(Q^2 - Q^2(\Phi_3)) a_3^0 \right. \\ \left. \times \left( 1 + \frac{a_3^1}{a_3^0} + \sum_{s \in a, b} \int_{\text{ord}} d\Phi_{\text{ant}}^s R_{2 \rightarrow 4} s_3' \right) \Delta(Q_0^2, Q^2) \right. \\ \left. + \sum_{s \in a, b} \int_{\text{unord}} d\Phi_{\text{ant}}^s \delta(Q^2 - Q^2(\Phi_4)) R_{2 \rightarrow 4} s_3 s_3' \Delta(Q_0^2, Q^2) \right]$$

(2→)3→4 antenna function  
(2→)3→4 MEC  
2→4 as explicit product x MEC

Only generates double-unresolved singularities, not single-unresolved

Note: the equation is formally identical to:

$$\frac{d}{dQ^2} \Delta(Q_0^2, Q^2) = \int \frac{d\Phi_3}{d\Phi_2} \delta(Q^2 - Q^2(\Phi_3)) (a_3^0 + a_3^1) \Delta(Q_0^2, Q^2) \\ + \int \frac{d\Phi_4}{d\Phi_2} \delta(Q^2 - Q^2(\Phi_4)) a_4^0 \Delta(Q_0^2, Q^2), \quad (3)$$

poles → poles

But on this form, the pole  
cancellation happens  
between the two integrals

Limited manpower but expect this in PYTHIA within the next ~ 2 years.

# Opportunities & Requirements

**Expect current developments (if sustained) to produce new generation of highly precise perturbative MC models by 2030.**

Standalone fixed-order calculations probably very limited applicability, e.g. for accuracy beyond NNLO.

For all other cases, expect (N)NLO matched and merged with next-generation showers or inclusive resummations (not covered here).

## Tests and Validations

Require observables sensitive to subtle sub-LL differences.

E.g., sensitive to "direct"  $n \rightarrow n + 2$  branchings, multi-parton correlations (e.g., triple-energy correlations, cf Komiske's talk) and multi-parton coherence, subleading  $N_C$ , ...

Scaling studies with  $\sqrt{s}$  ➤ can disentangle power corrections, beta function, ...

CEPC/FCC-ee ➤ statistics to focus on small but "clean" corners of phase space

Important to develop a battery of such tests; relevant also for LHC

## Requirements (?)

Excellent **resolution** of **jet substructure**, and excellent **jet flavour tagging** (+  $Z \rightarrow 4b, 4c, 2b2c$ )

Forward coverage, to access **low  $\sqrt{s} \sim 10\text{-}20 \text{ GeV}$**  via ISR from Z pole?



# $e^+e^- \rightarrow WW$ : Resonance Decays

## Current MC Treatment ~ Double-Pole Approximation

~ First term in double-pole expansion (cf. Schwinn's talk in yesterday's EW session)

+ Some corrections, e.g., in PYTHIA:

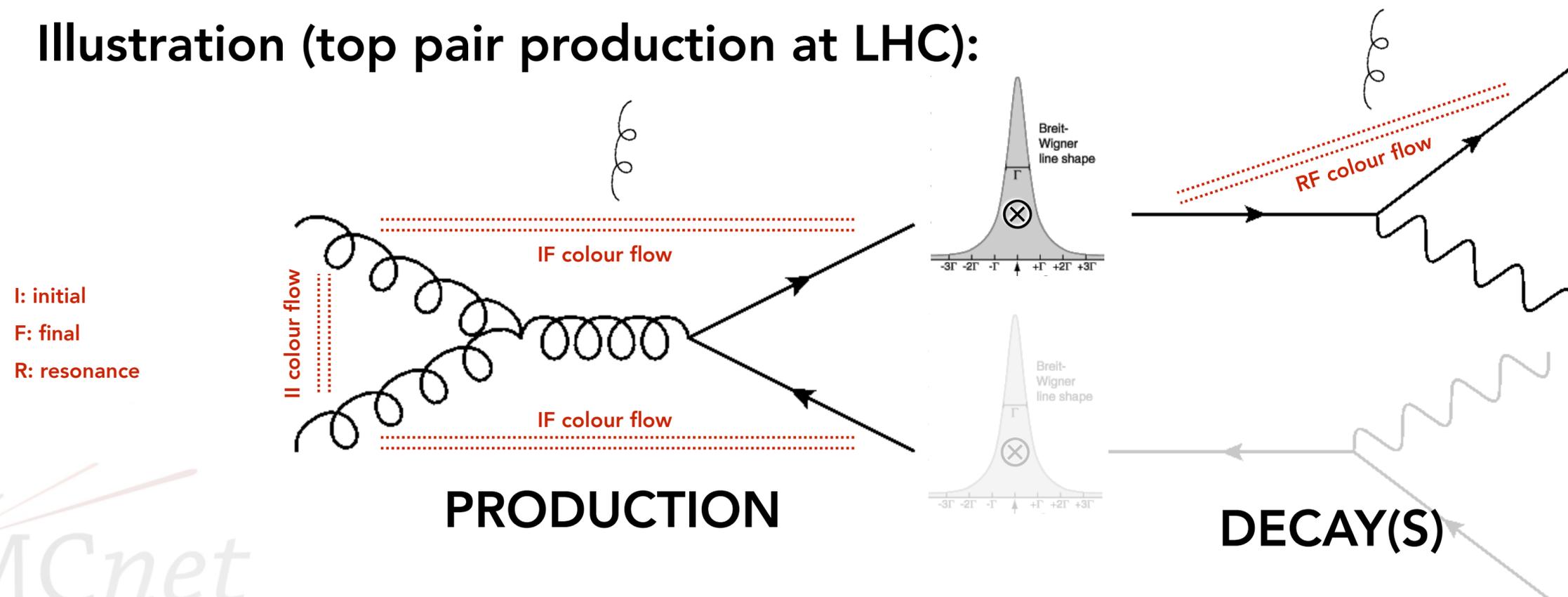
Independent Breit-Wigners for each of the  $W$  bosons, with running widths.

4-fermion ME used to generate correlated kinematics for the  $W$  decays.

Each  $W$  decay treated at NLO + shower accuracy.

No interference / coherence between ISR, and each of the  $W$  decay showers

## Illustration (top pair production at LHC):



# Interleaved Resonance Decays

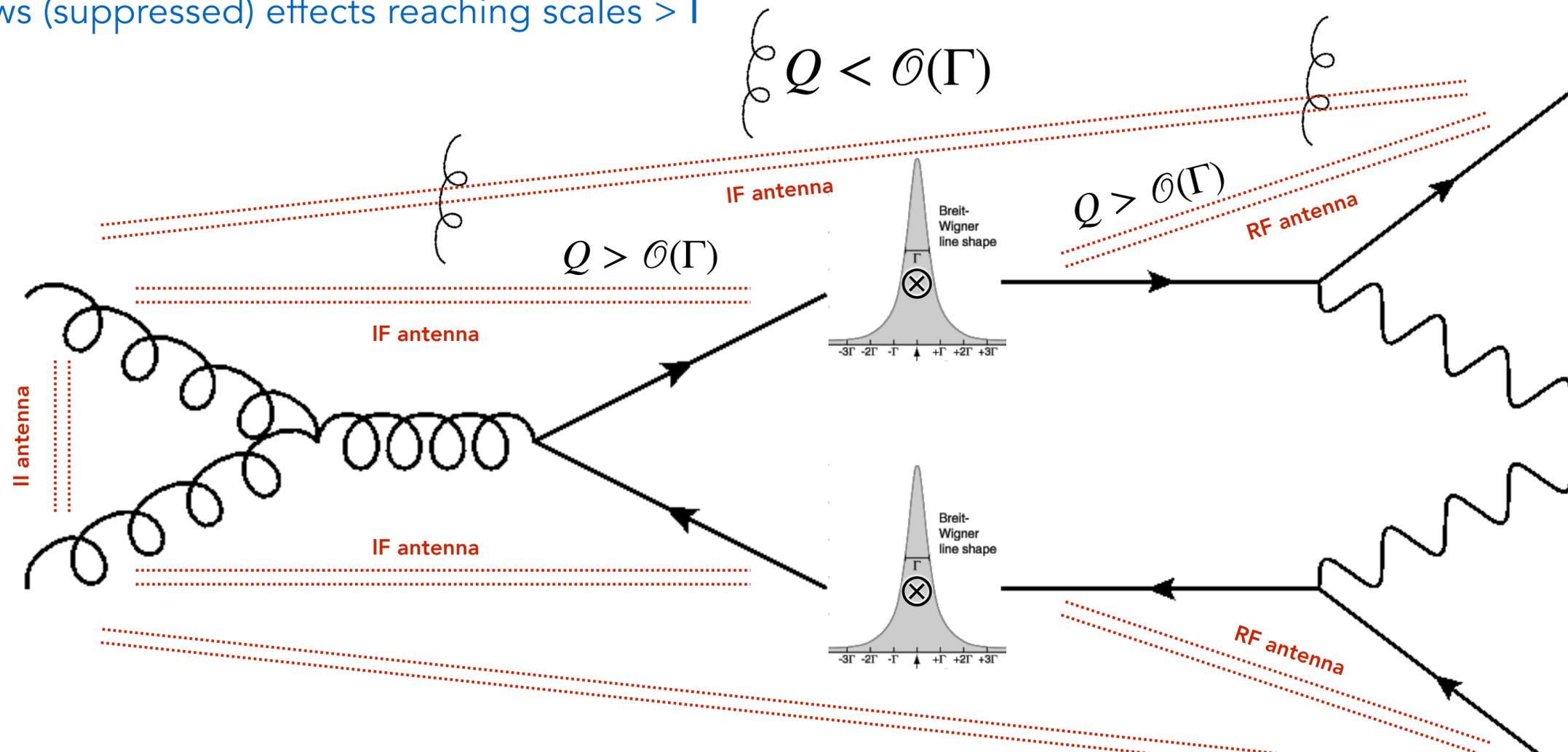
**Decays of unstable resonances introduced in shower evolution at an average scale  $Q \sim \Gamma$**

Cannot act as emitters or recoilors below that scale; only their decay products can do that.

The more off-shell a resonance is, the higher the scale at which it disappears.

Roughly corresponds to strong ordering (as measured by propagator virtualities) in rest of shower.

Allows (suppressed) effects reaching scales  $> \Gamma$



**Automatically provides a natural treatment of finite- $\Gamma$  effects.**

Expect in next Pythia release (8.304)

# Hadronisation (and low $z$ )

**Confinement** wasn't solved last century

Models **inspired by QCD** (hadronisation models) explore the non-perturbative quagmire (until it is solved and **uninspired** models can move in) FFs and IR safety (power corrs) observe from a safe distance

**Can do track reconstruction (3 hits) down to 30-40 MeV  $\ll \Lambda_{\text{QCD}}$  ?**

Below  $\Lambda_{\text{QCD}}$   $\rightarrow$  can study genuine non-perturbative dynamics

**Handles:** mass, strangeness, and spin. Need at least one of each meson & baryon isospin multiplet. Flavour separation crucial. (LEP  $|p_K| > 250$  MeV)

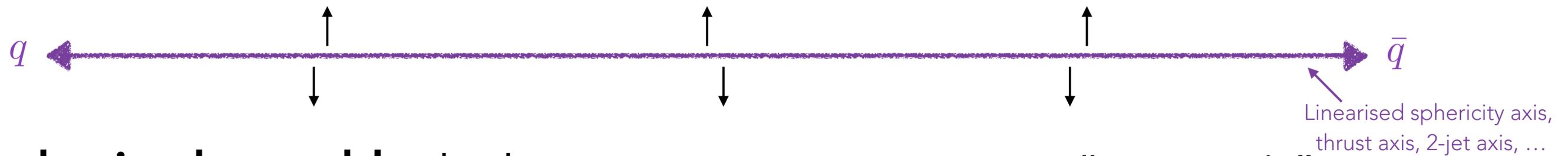
**QUESTIONS:** detailed mechanisms of hadron production. Is strangeness fraction constant or dynamic? Thermal vs Gaussian spectra. Debates rekindled by LHC observations of strangeness enhancement.

**Bonus: high(er)-precision jet calibration (particle flow) ?**

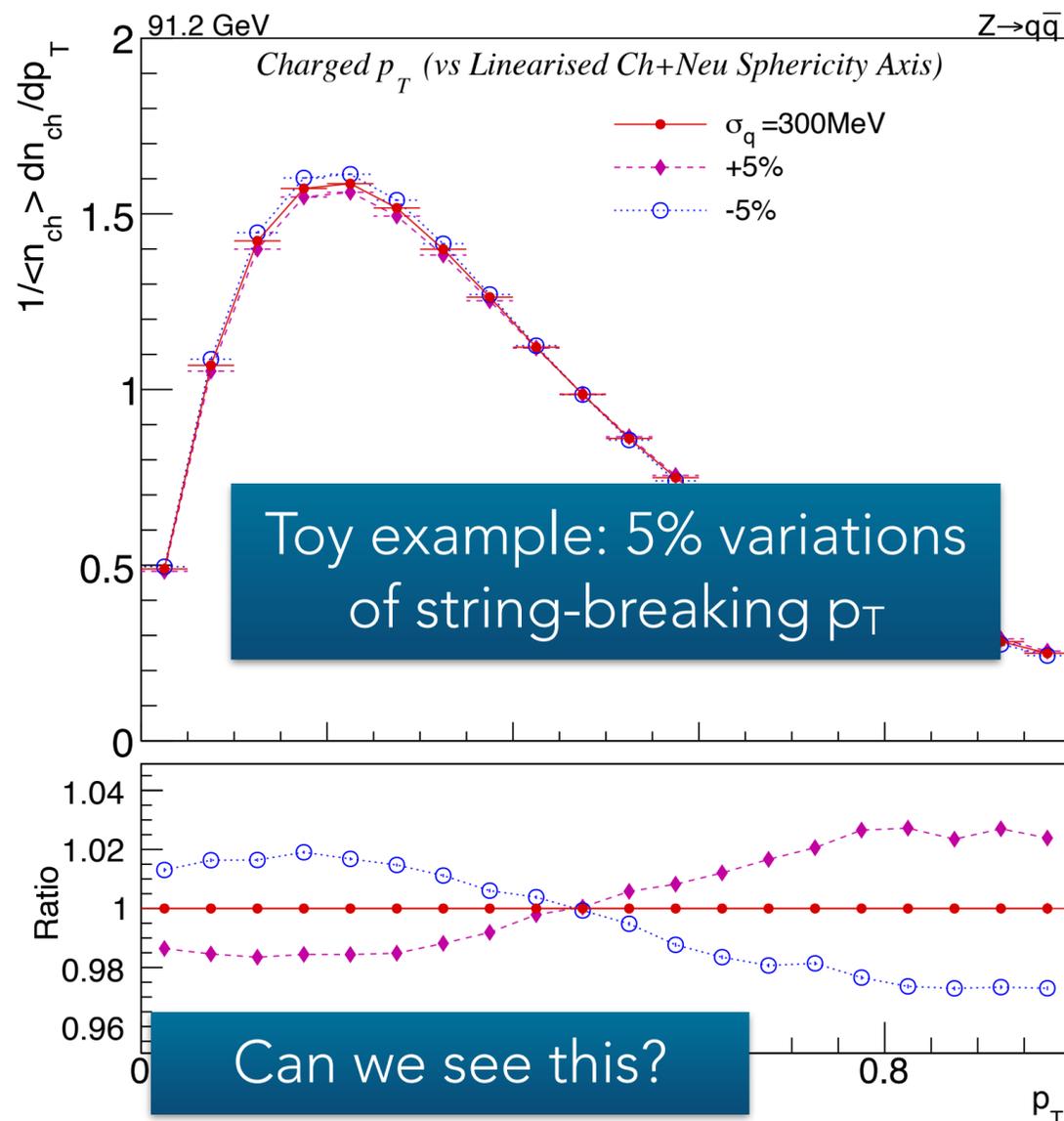
Accurate knowledge (+ modeling) of particle composition & spectra



# Transverse Fragmentation $\Leftrightarrow$ Momentum Resolution

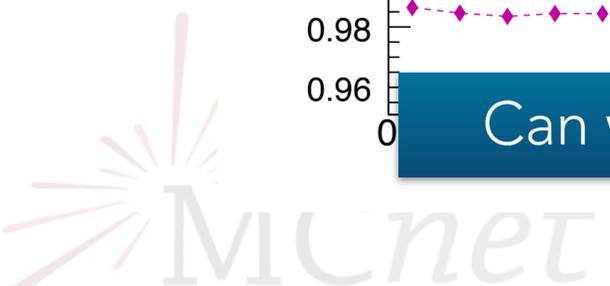
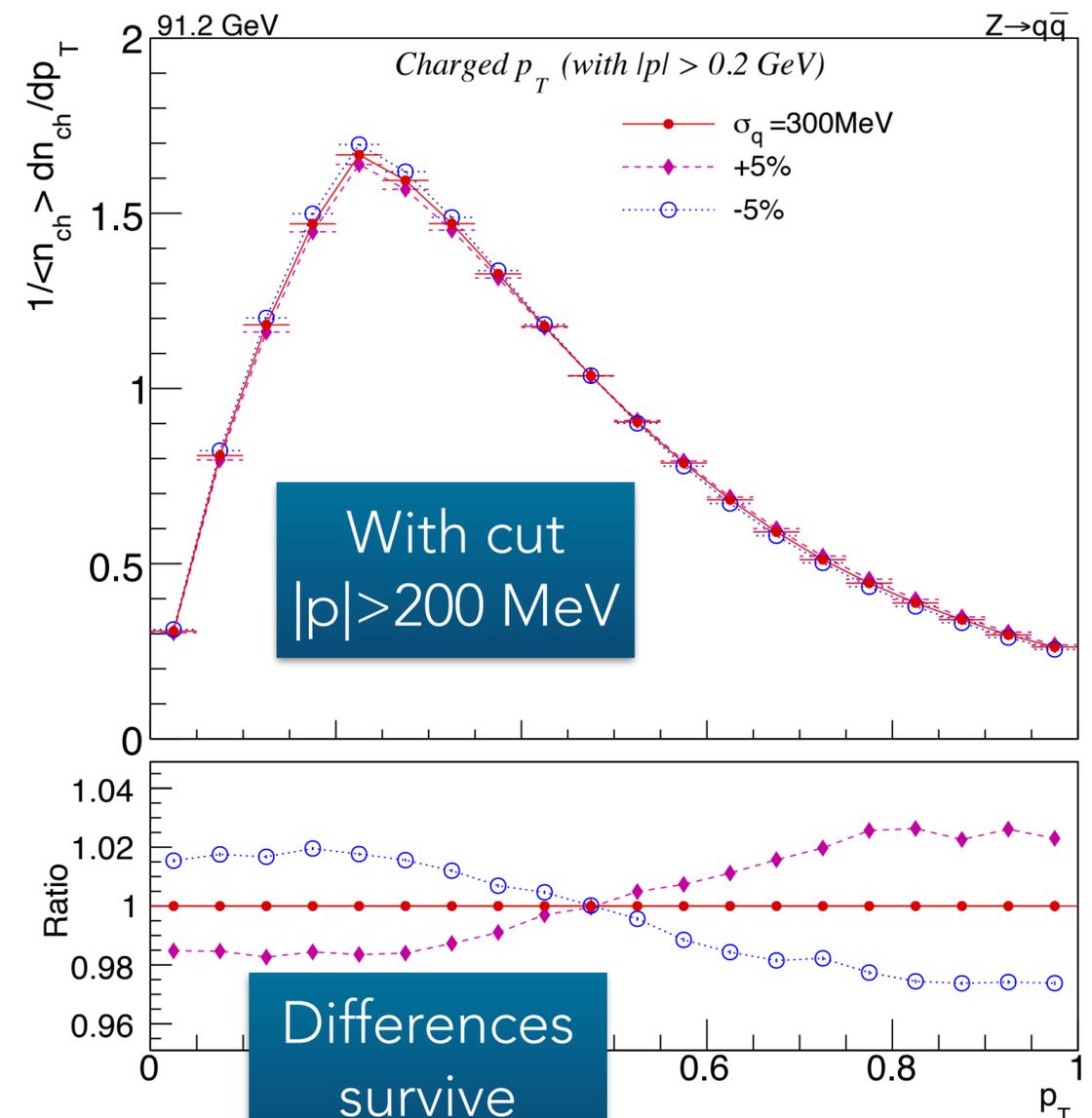


**Most basic observable:** hadron  $p_T$  spectra, transverse to "event axis"



Toy Example

Perturbatively dominated power-law tail



# Effects of order $\Lambda_{QCD} \sim 100 \text{ MeV} \leftrightarrow$ Coverage for $|p| < \Lambda_{QCD}$ ?

## $p_T$ kicks from hadronisation

Pythia  $\sim$  Gaussian  $\sim 300 \text{ MeV}$  (+  $\rho$  decays)

Acts as a sort of lower bound on hadron  $p_T$ .

Difficult for any hadron to have  $|p| < 300 \text{ MeV}$ .

To check this, look for pions with  $|p| < 300 \text{ MeV}$

► Probe of confinement mechanism for non-relativistic pions

## Data from both LEP and LHC indicate more soft pions; why?

Thermal vs Gaussian spectra?

Unresolved perturbative effects vs genuine string-breaking effects?

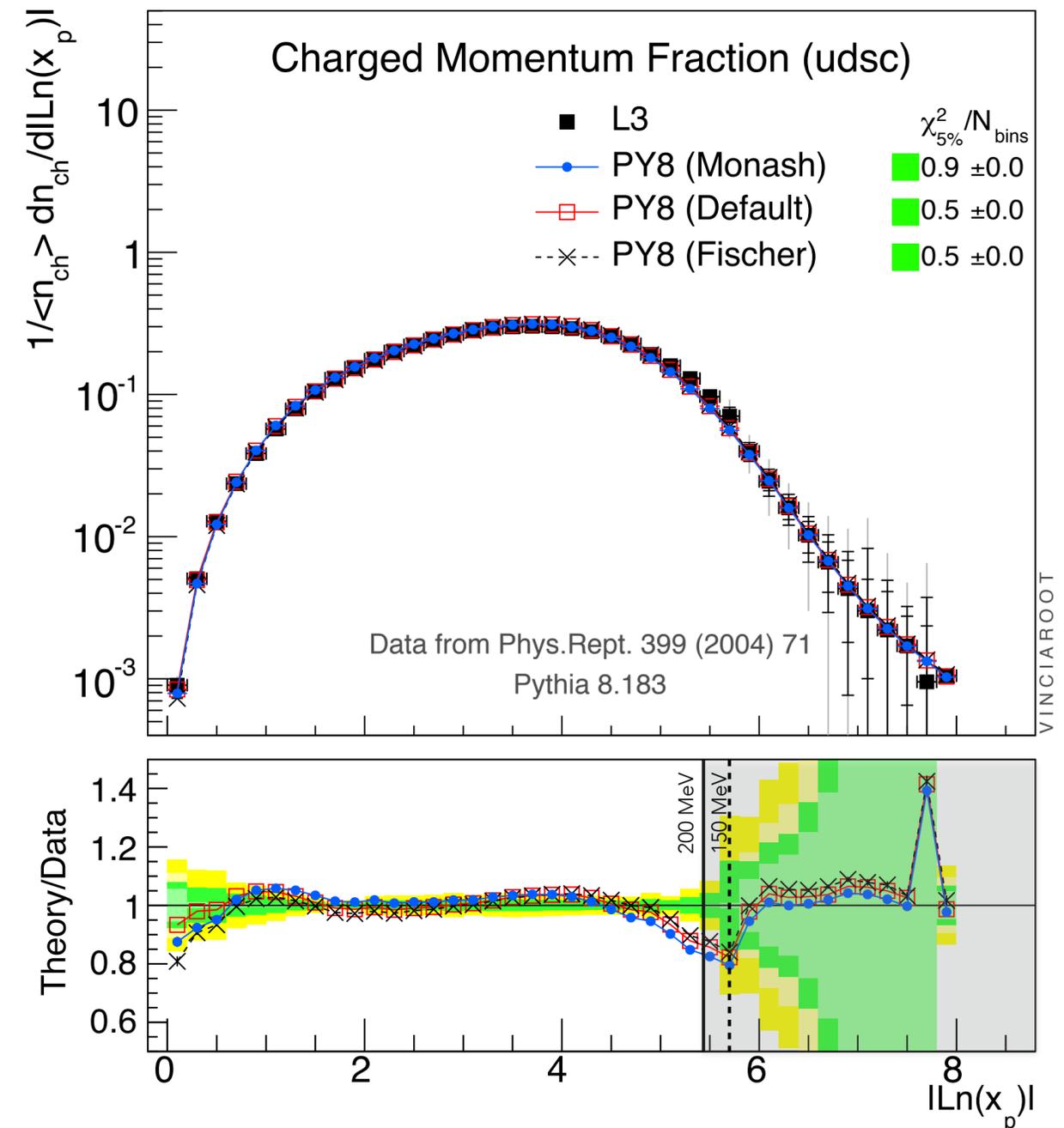
Mismodelled resonance decays?

## Cut at $|p| = 200 \text{ MeV}$ makes this tough to examine clearly

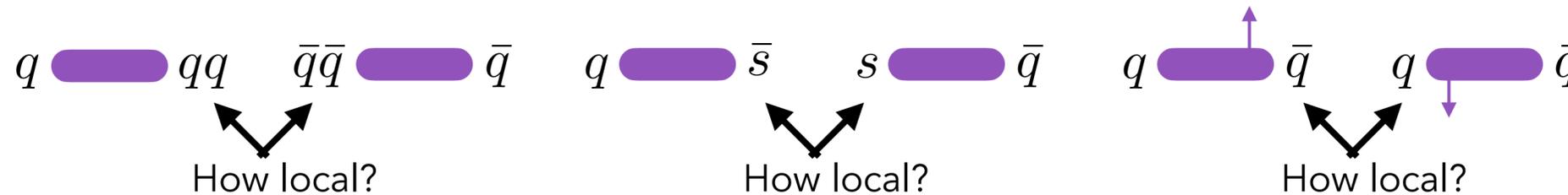
3 hits down to  $\sim 50 \text{ MeV}$  ?

Special runs / setups with lower thresholds?

## Example from LEP



The **point** of MC generators: address more than one hadron at a time!



Further precision non-perturbative aspects: **How local is hadronisation?**

Baryon-Antibaryon correlations — both OPAL measurements were statistics-limited (Kluth); would reach OPAL systematics at  $10^8$  Z decays ( $\rightarrow 10^9$  with improved systematics?)

+ Strangeness correlations,  $p_T$ , spin/helicity correlations (“screwiness”?)

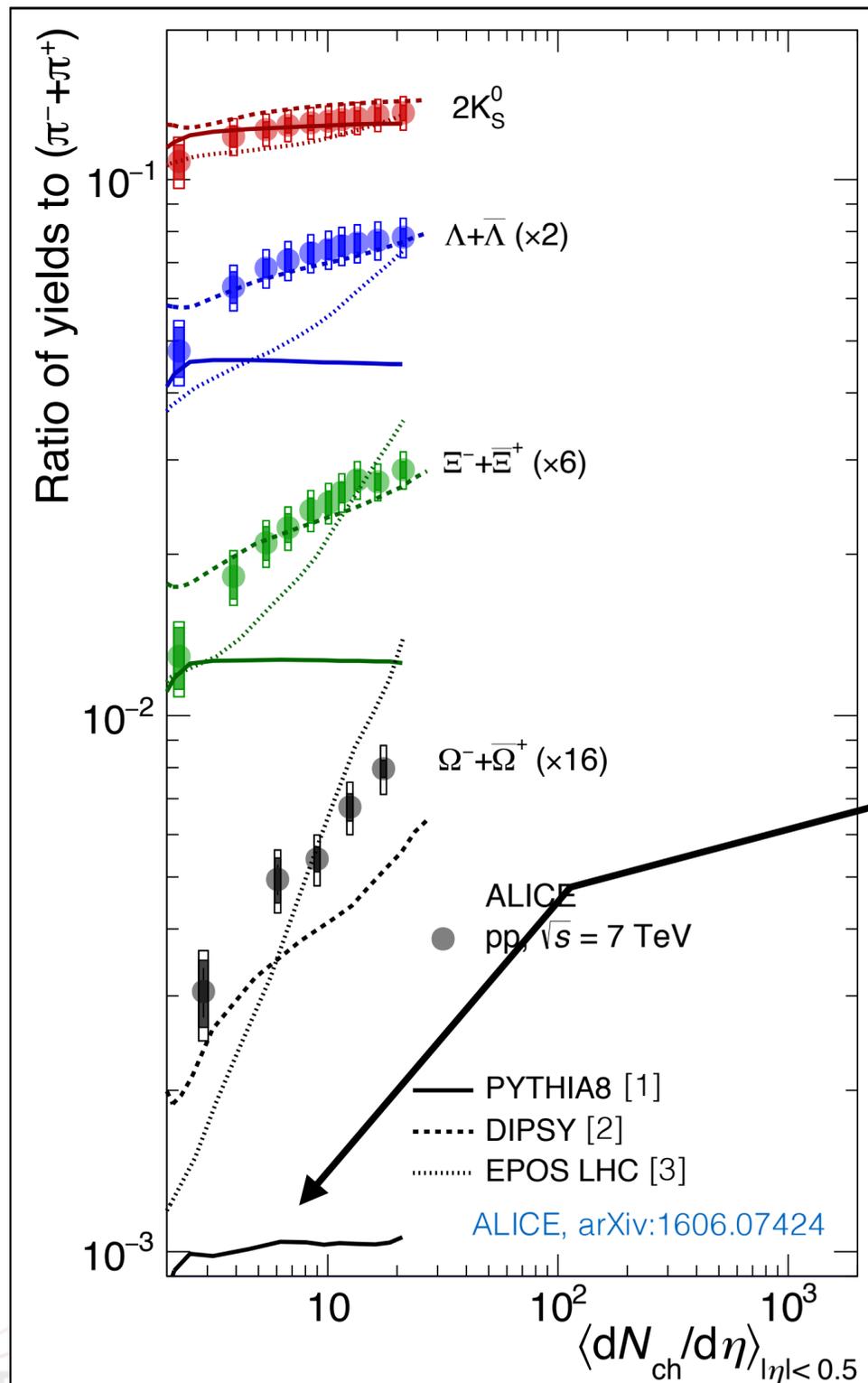
+ Bose-Einstein Correlations & Fermi-Dirac Correlations

Identical baryons (pp,  $\Lambda$ ) **highly** non-local in string picture — puzzle from LEP; correlations across multiple exs & for both pp and  $\Lambda\Lambda \rightarrow$  Fermi-Dirac radius  $\sim 0.1$  fm  $\ll r_p$  (Metzger)

Octet neutralisation? (zero-charge gluon jet with rapidity gaps)  $\rightarrow$  **neutrals**  
Colour reconnections, glueballs, ...

Leading baryons in g jets?  
(discriminates between string/cluster models)  
**High-x baryons**

# Strangeness (in PP)



D.D. Chinellato – 38th International Conference on High Energy Physics

**ALICE:** clear enhancement of strangeness with (pp) event multiplicity

No corresponding enhancement for protons (not shown here but is in ALICE paper) → must really be a strangeness effect

**Jet universality:** jets at LHC modelled the same as jets at LEP

→ Flat line ! (cf PYTHIA)

Some models **anticipated** the effect!

DIPSY (high-tension overlapping strings)

EPOS (thermal hydrodynamic "core")

**Is it thermal? Or stringy? (or both?)**

**Basic check in  $ee \rightarrow WW$ : two strings**

Requires **good PID** + high statistics

(LEP: total  $\Omega$  rate only known to  $\pm 20\%$ )

## At LEP 2: hot topic (by QCD standards): 'string drag' effect on W mass

**Non-zero effect** convincingly demonstrated at LEP-2

No-CR **excluded** at 99.5% CL [Phys.Rept. 532 (2013) 119]

**But not much detailed (differential) information**

Thousand times more WW at CEPC / FCC-ee

Turn the W mass problem around; use threshold scan + huge sample of semi-leptonic events to measure  $m_W$

→ **input** as constraint to **measure CR in hadronic WW**

## Has become even hotter topic at LHC

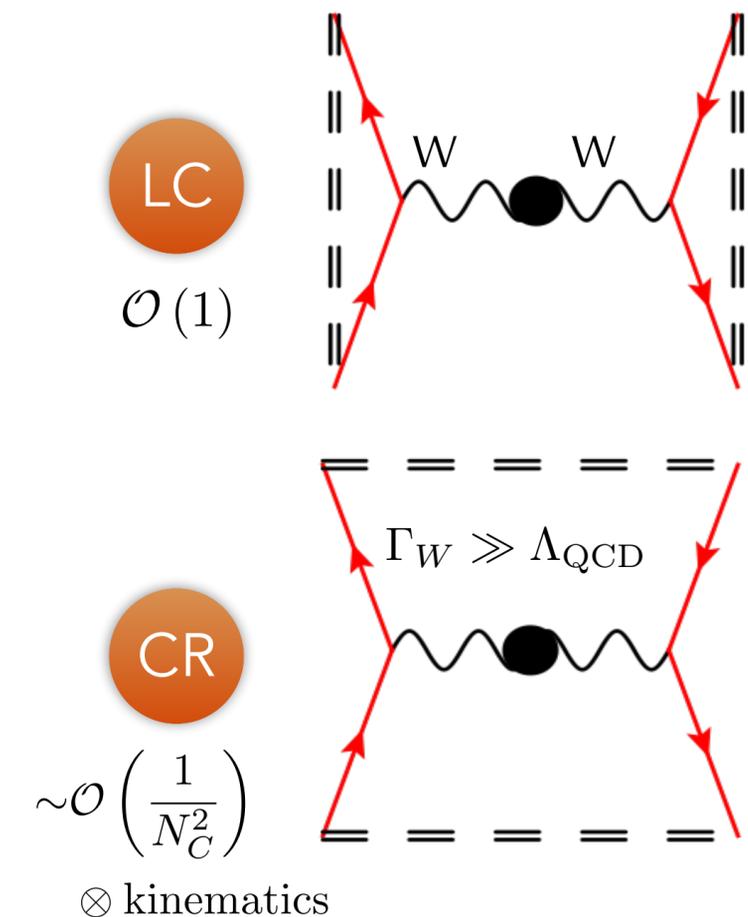
It appears jet universality is under heavy attack.

Fundamental to understanding & modeling hadronisation

Follow-up studies now underway at LHC.

## High-stats ee → other side of story

Also relevant in (hadronic)  $ee \rightarrow tt$ , and  $Z \rightarrow 4$  jets



+ Overlaps → interactions? increased tensions (strangeness)? breakdown of string picture?

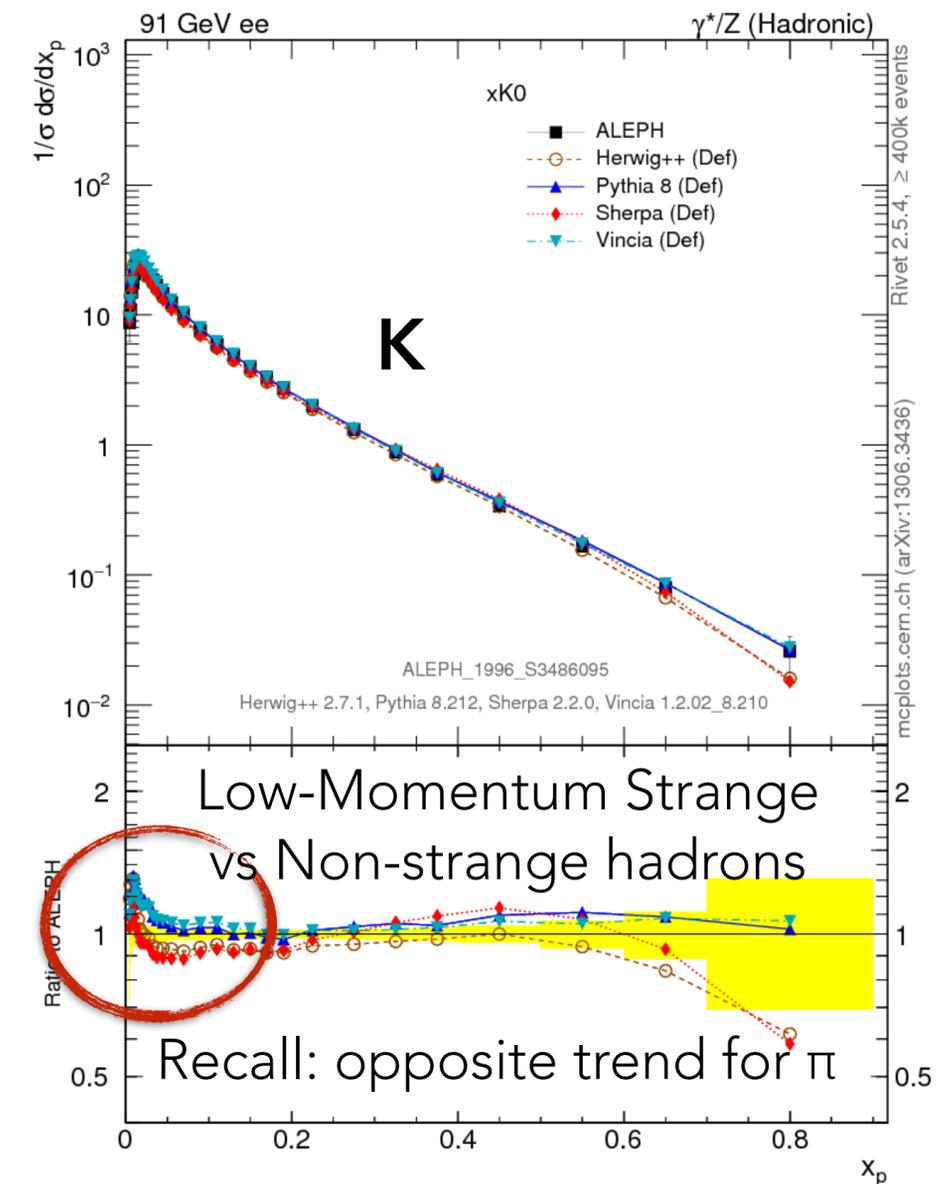
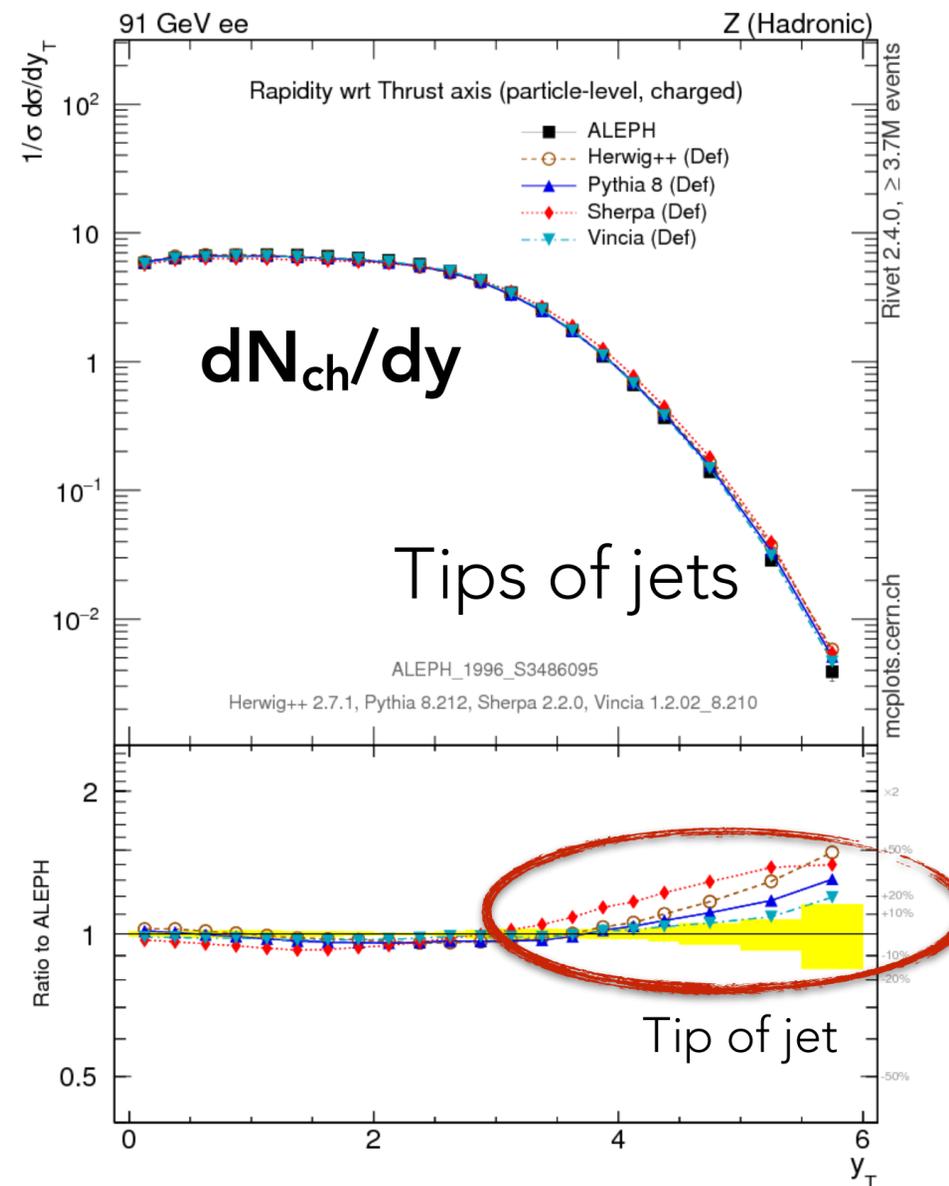
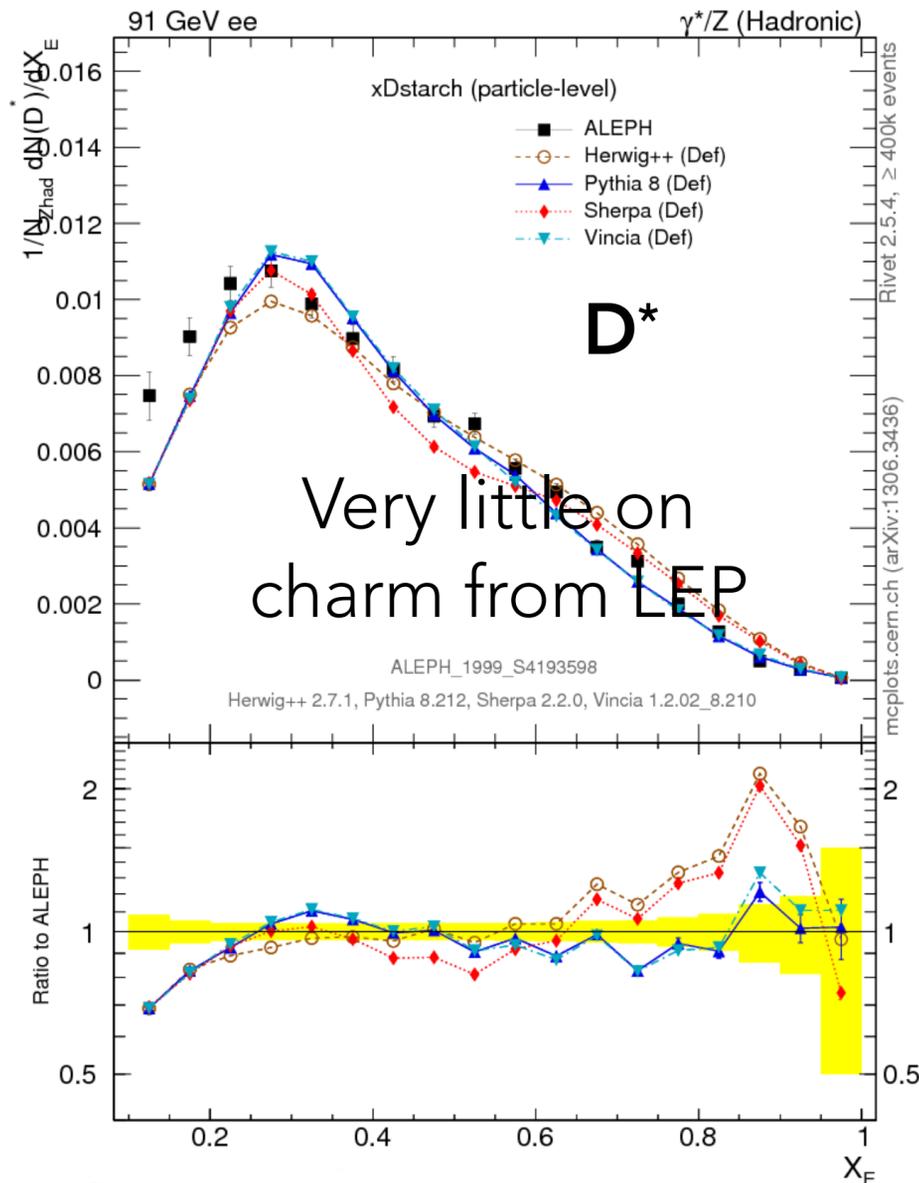
Little done for CEPC/FCC-ee so far ... (to my knowledge)  
Plenty of room to play with models, observables, ...

Overviews of recent models:  
[arXiv:1507.02091](https://arxiv.org/abs/1507.02091), [arXiv:1603.05298](https://arxiv.org/abs/1603.05298)

# Plenty of other interesting detailed features

(plots from [mcplots.cern.ch](http://mcplots.cern.ch))

Just a few examples



Capabilities for hadrons from decays ( $\pi^0$ ,  $\eta$ ,  $\eta'$ ,  $\rho$ ,  $\omega$ ,  $K^*$ ,  $\varphi$ ,  $\Delta$ ,  $\Lambda$ ,  $\Sigma$ ,  $\Sigma^*$ ,  $\Xi$ ,  $\Xi^*$ ,  $\Omega$ , ...)

+ heavy-flavour hadrons

Very challenging; conflicting measurements from LEP

# Example of recent reexamination of String Basics

## Cornell potential

Potential  $V(r)$  between **static** (lattice) and/or **steady-state** (hadron spectroscopy) colour-anticolour charges:

$$V(r) = -\frac{a}{r} + \kappa r$$

Coulomb part

String part

Dominates for  $r \gtrsim 0.2 \text{ fm}$

Lund string model built on the asymptotic large- $r$  linear behaviour

**But intrinsically only a statement about the late-time / long-distance / steady-state situation. Deviations at early times?**

Coulomb effects in the grey area between shower and hadronization?

**Low- $r$  slope  $> \kappa$**  favours "early" production of quark-antiquark pairs?

+ Pre-steady-state thermal effects from a (rapidly) **expanding string?**

Berges, Floerchinger, and Venugopalan JHEP 04(2018)145)

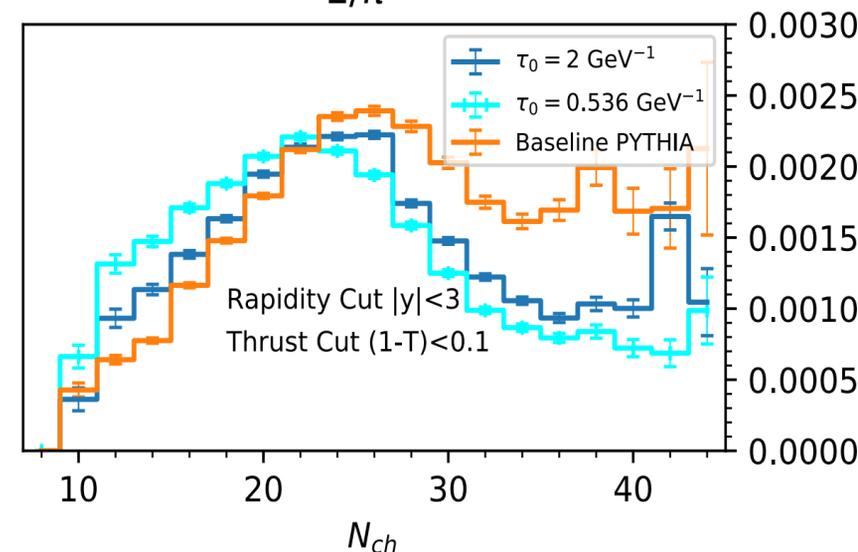
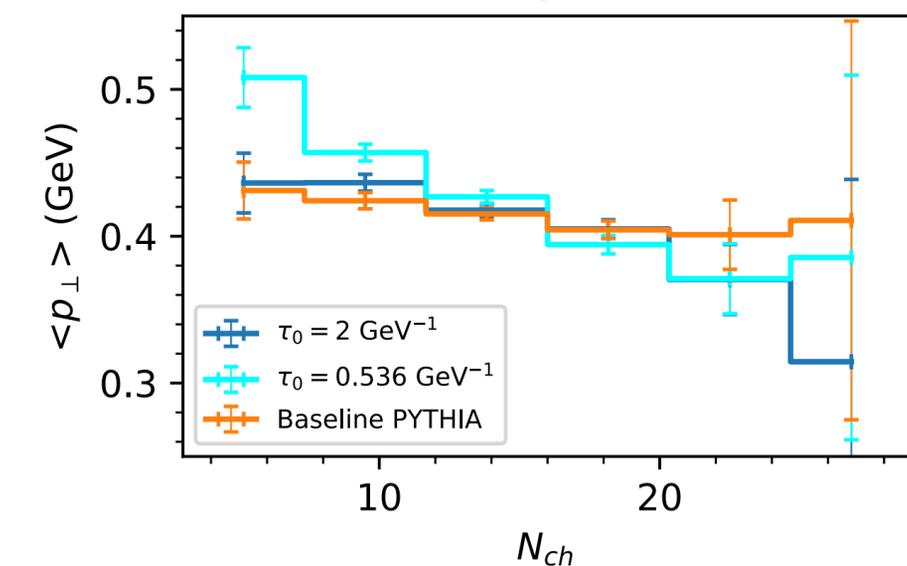
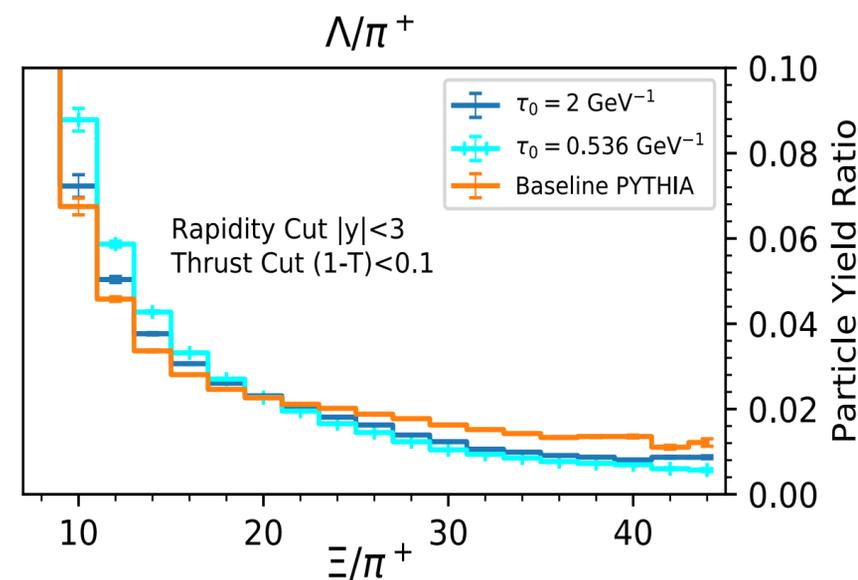
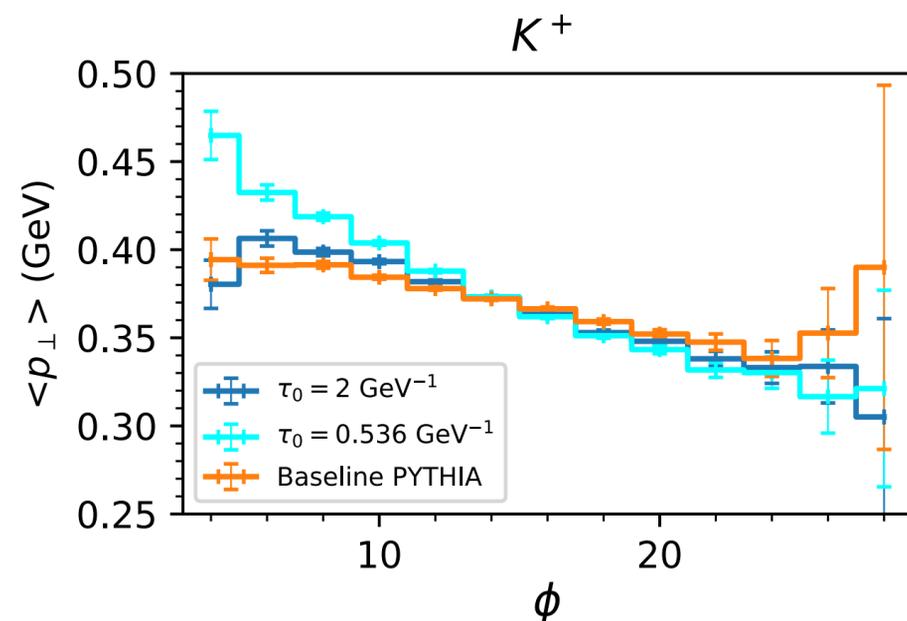
# Toy Model with Time-Dependent String Tension

N. Hunt-Smith & PS arxiv:[2005.06219](https://arxiv.org/abs/2005.06219)

**Model constrained to have same average tension as Pythia's default "Monash Tune"**

➤ same average  $N_{ch}$  etc ➤ main LEP constraints basically unchanged.

But expect different fluctuations / correlations, e.g. with multiplicity  $N_{ch}$ .



- Want to study (suppressed) tails with very low and very high  $N_{ch}$ .
- These plots are for LEP-like statistics.
- Would be crystal clear at CEPC/ FCC-ee

# MCs & Precision QCD at Future $e^+e^-$ Machines

## Perturbative QCD: **High Precision**

Measurements of  $\alpha_s$  with unprecedented accuracy (not covered here)

**Good jet substructure & flavour tagging** crucial to vet **N<sup>n</sup>LO QCD** + **Next Generation** of **Showers**

➔ Accurate starting point for non-perturbative modelling of **Hadronisation**

## Interplays with EW & Higgs Physics Goals

Impact of (in)accurate MC predictions?  $\Leftrightarrow$  Identify & Communicate crucial areas for improvements?

## Nonperturbative QCD: **High Resolution**

Confinement / Non-perturbative QFT remains fundamentally unsolved

Next generation of  $e^+e^-$  machines ➔ **trial by fire** not just for any post-LHC advanced models, but also for any future *solution* or systematically improvable approximation.

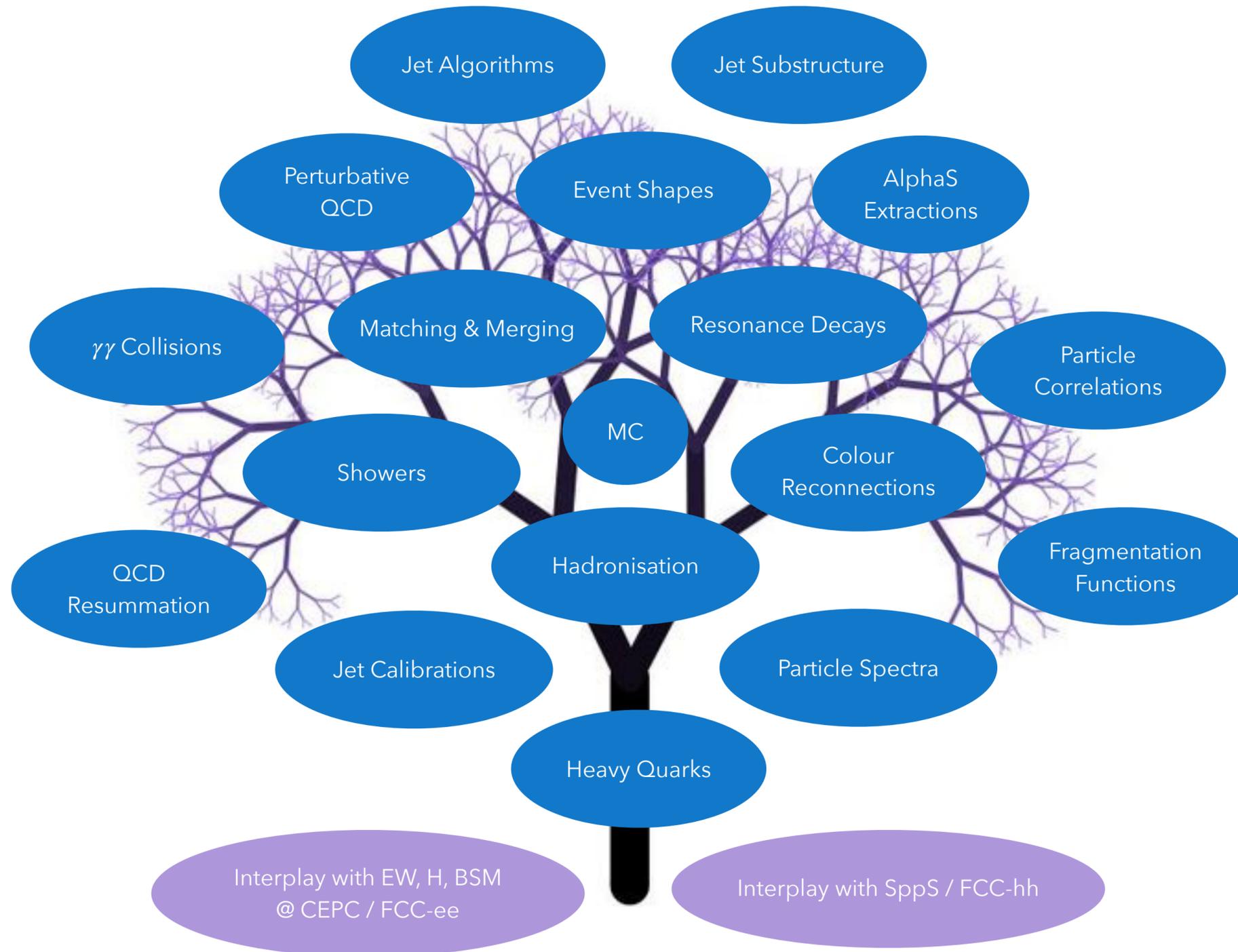
➔ **Good PID** crucial to reveal **details of final states**  $\Leftrightarrow$  **disentangle** strangeness, baryons, mass, spin

➔ **Good Momentum Resolution** crucial to measure  $\mathcal{O}(\Lambda_{\text{QCD}}) \sim 100$  MeV effects with **high precision**

Theory keeps evolving long after beams are switched off ➤ **Aim high!**



# Summary — QCD at EE Colliders



Extra Slides

# Themes

## Measure $\alpha_S$

High-Precision Z (and W) widths

High-Precision Event Shapes, Jet Rates, ... (IR safe observables sensitive to  $\alpha_S$ )

## Single-Inclusive Hadron Production and Decays

Fragmentation Functions; Hadron Spectra; (+ polarisation)

Exotic /rare hadrons, quarkonium, rare decays, ...

+ Interplay with flavour studies (+ Interplay with DM annihilation)

## Understanding Confinement (Multi-hadronic / Exclusive)

In high-energy processes  $\rightarrow$  hadronisation

Hadron correlations, properties with respect to global ("string") axes

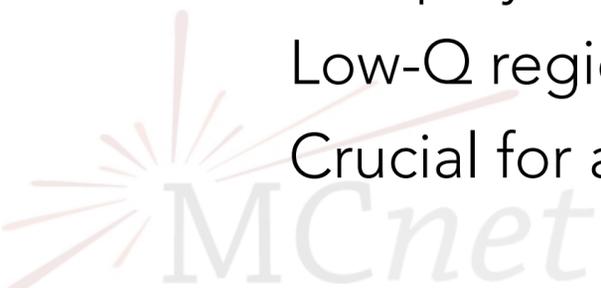
Dependence on (global and local) environment (distance to jets, hadronic density, flavours)

## Power Corrections / Hadronisation Corrections

Interplay with high- $p_T$  physics program

Low- $Q$  region of event shapes, jet rates, jet substructure; jet flavour tagging, ...

Crucial for  $\alpha_S$  measurements; also for jet calibration?



# Precision $\alpha_s$ Measurements

(see FCC-ee QCD workshops & writeups)

CURRENT STATE OF THE ART: O(1%)

**LEP:** Theory keeps evolving long after the beams are switched off

Recently, NNLO programs for 3-jet calculations

[Weinzierl, PRL 101, 162001 (2008)]; EERAD [Gehrmann-de-Ridder, Gehrmann, Glover, Heinrich, CPC185(2014)3331]

+ New resummations  $\rightarrow$  new  $\alpha_s(m_Z)$  extractions

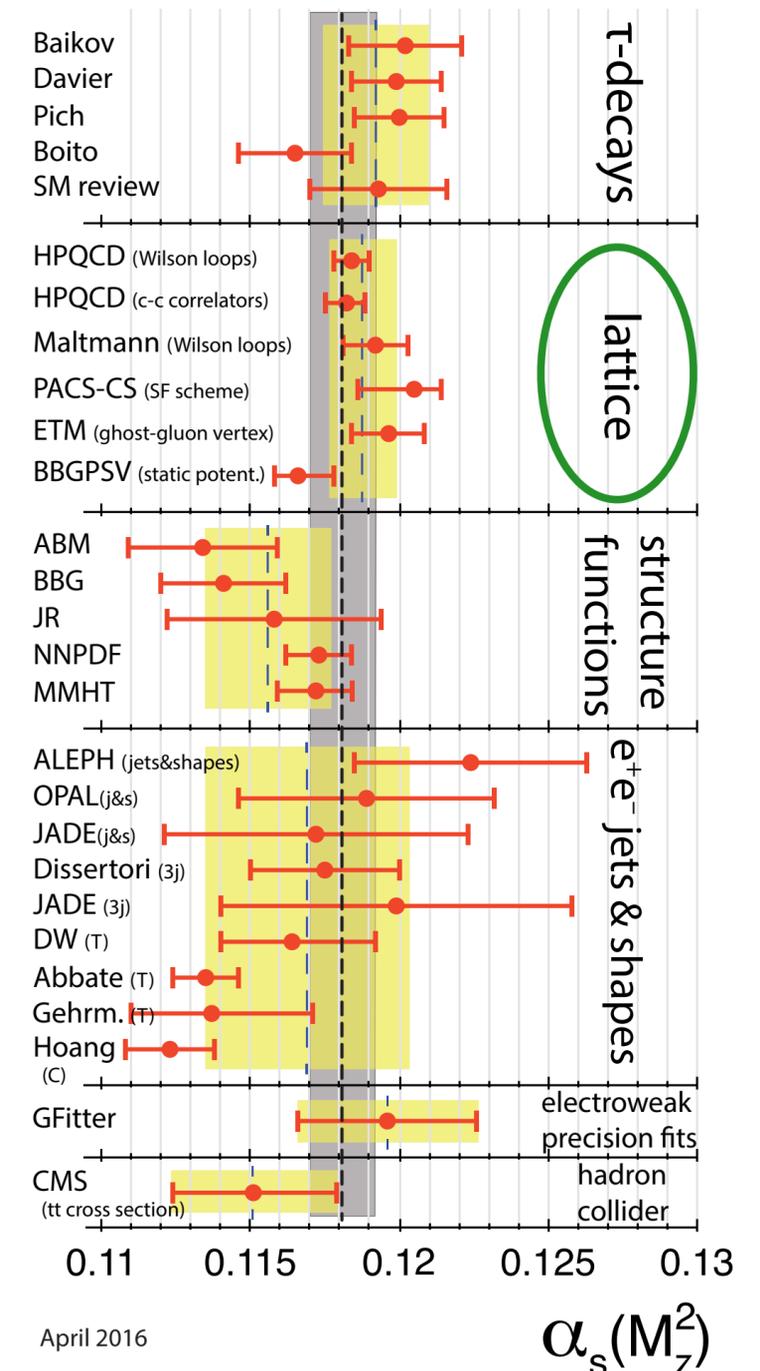
E.g., 2015 SCET-based C-parameter reanalysis

$N^3LL' + O(\alpha_s^3) + NPPC: \alpha_s(m_Z) = 0.1123 \pm 0.0015$

[Hoang, Kolodubretz, Mateu, Stewart, PRD91(2015)094018]

ee currently the least precise subclass (due to large spread between individual extractions)

Subclass	PDG 2016	$\alpha_s(M_Z^2)$
$\tau$ -decays		$0.1192 \pm 0.0023$
lattice QCD		$0.1188 \pm 0.0011$
structure functions		$0.1156 \pm 0.0021$
$e^+e^-$ jets & shapes		$0.1169 \pm 0.0034$
hadron collider		$0.1151 \pm 0.0028$
ewk precision fits		$0.1196 \pm 0.0030$



See also PDG QCD review and references therein

+ 2016 Moriond  $\alpha_s$  review [d'Enterria]: arXiv:1606.04772

+ 2015 FCC-ee  $\alpha_s$  workshop proceedings: arXiv:1512.05194

Maximum a factor 3 further reduction possible (without FCC-ee). [Some participants believed less.]

STATISTICS ALLOW TO AIM FOR  $\delta\alpha_s/\alpha_s < 0.1\%$

## Main Observable:

$$R_\ell^0 = \frac{\Gamma_{\text{had}}}{\Gamma_\ell} \quad \text{LO} \quad \Gamma_f \propto (g_{V,f}^2 + g_{A,f}^2) \quad g_{V,f} = g_{A,f}(1 - 4|q_f| \sin^2 \theta_W)$$

QCD corrections to  $\Gamma_{\text{had}}$  known to 4<sup>th</sup> order

Kuhn: Conservative QCD scale variations  $\rightarrow O(100 \text{ keV}) \rightarrow \delta\alpha_s \sim 3 \times 10^{-4}$

Comparable with the target for CEPC / FCC-ee

Electroweak beyond LO  $g_{A,f} \rightarrow \sqrt{1 + \Delta\rho_f} g_{A,f}$   $\sin^2 \theta_W \rightarrow \sqrt{1 + \Delta\kappa_f} \sin^2 \theta_W = \sin^2 \theta_{\text{eff}}^f$ ,

Can be calculated (after Higgs discovery) or use measured  $\sin^2 \theta_{\text{eff}}$

Mönig (Gfitter) assuming  $\Delta m_Z = 0.1 \text{ MeV}$ ,  $\Delta \Gamma_Z = 0.05 \text{ MeV}$ ,  $\Delta R_1 = 10^{-3}$

$\rightarrow \delta\alpha_s \sim 3 \times 10^{-4}$  ( $\delta\alpha_s \sim 1.6 \times 10^{-4}$  without theory uncertainties)

Better-than-LEP statistics also for  $W \rightarrow$  high-precision  $R_W$  ratio !

Srebre & d'Enterria: huge improvement in  $BR(W_{\text{had}})$  at FCC-ee (/CEPC?)

Combine with expected  $\Delta \Gamma_W = 12 \text{ MeV}$  from LHC (high- $m_T$   $W$ ) & factor-3 improvement in  $|V_{cs}| \rightarrow$  similar  $\alpha_s$  precision to extraction from  $Z$  decays?

# Fragmentation Functions

(see FCC-ee QCD workshops & writeups)

S. Moch (& others): field now moving towards NNLO accuracy: **1% errors** (or better)

## FFs from Belle to FCC-ee [A. Vossen]

**Precision** of TH and EXP big advantage

Complementary to pp and SIDIS

**Evolution:**

Belle has FCC-ee like stats at 10 GeV.

FCC-ee: very fine binning all the way to  $z=1$  with 1%  $l_{pl}$  resolution (expected)

**Flavour structure** for FFs of hyperons and other hadrons that are difficult to reconstruct in pp and SIDIS.

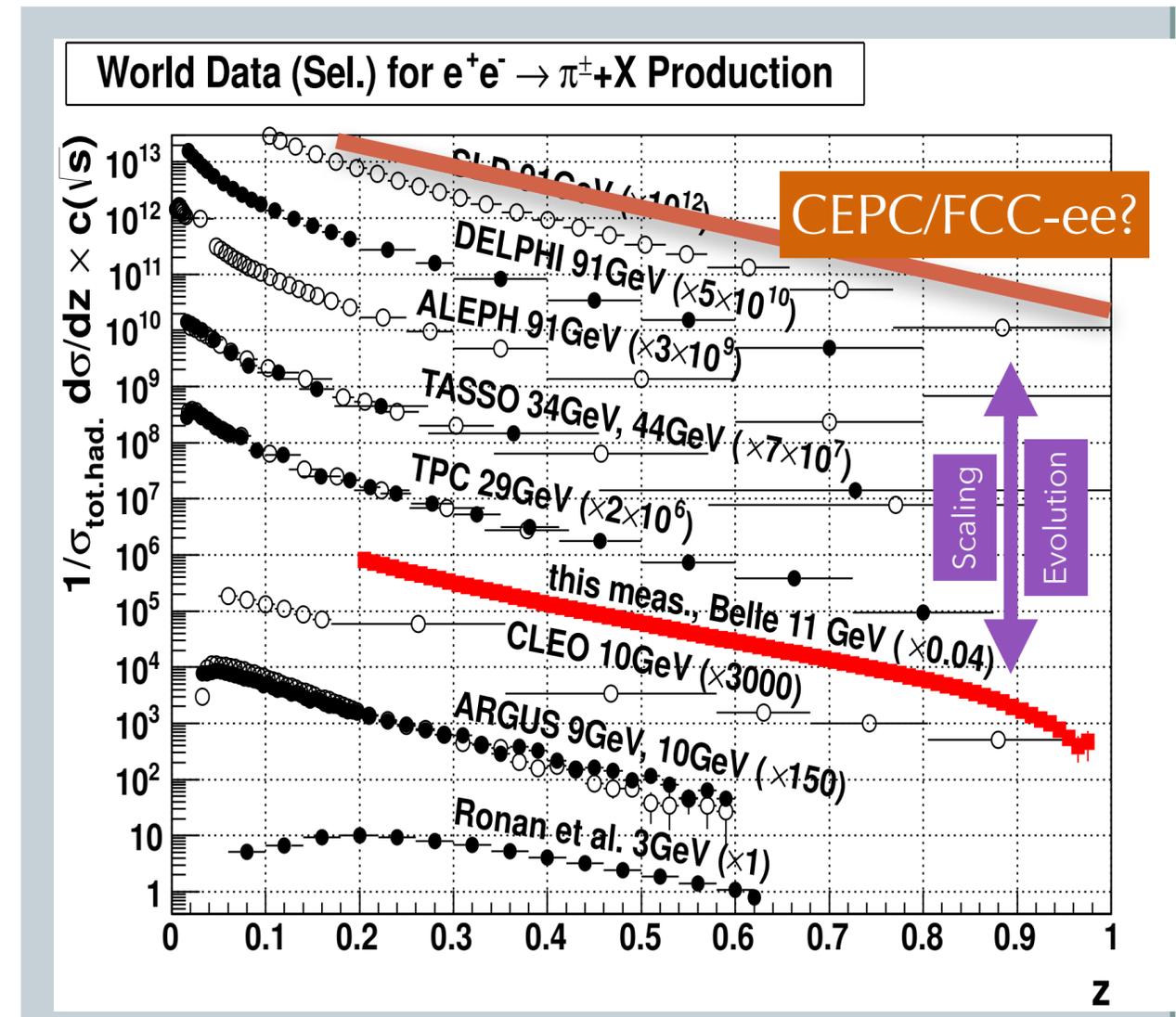
Will depend on Particle Identification capabilities.

**Low Z:** Higher ee energy (than Belle) → smaller mass effects at low  $z$ .

3 tracker hits down to 30-40 MeV allows to reach  $z = 10^{-3}$  ( $\ln(z) = -7$ )

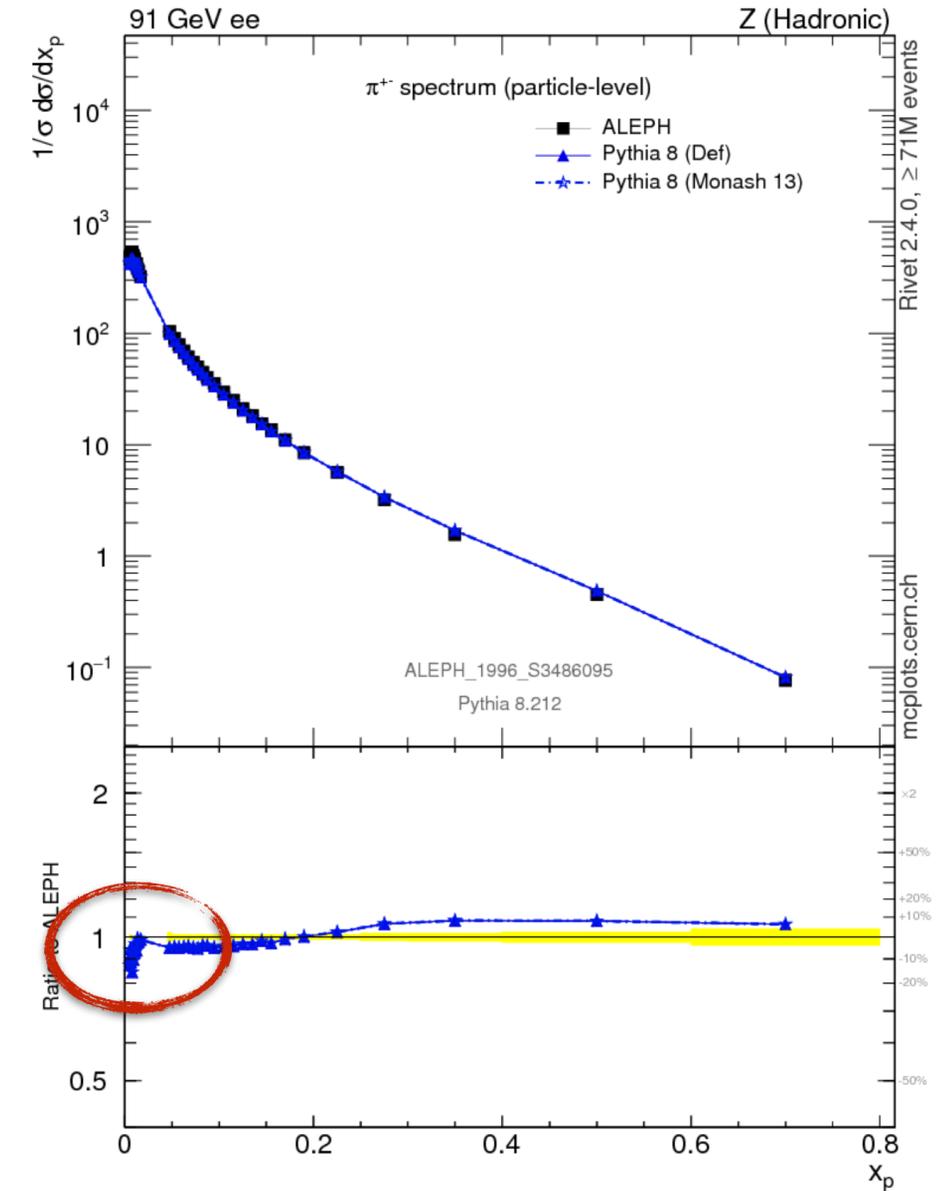
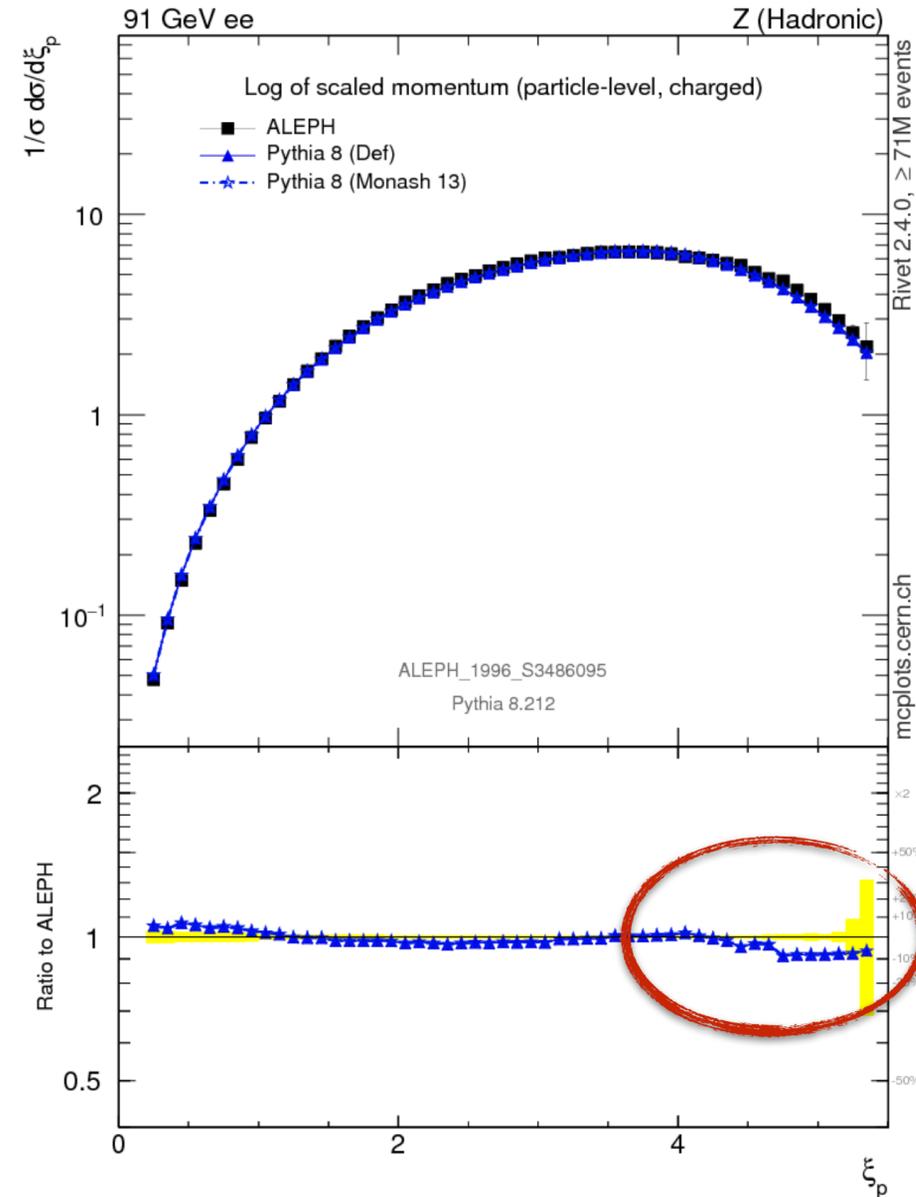
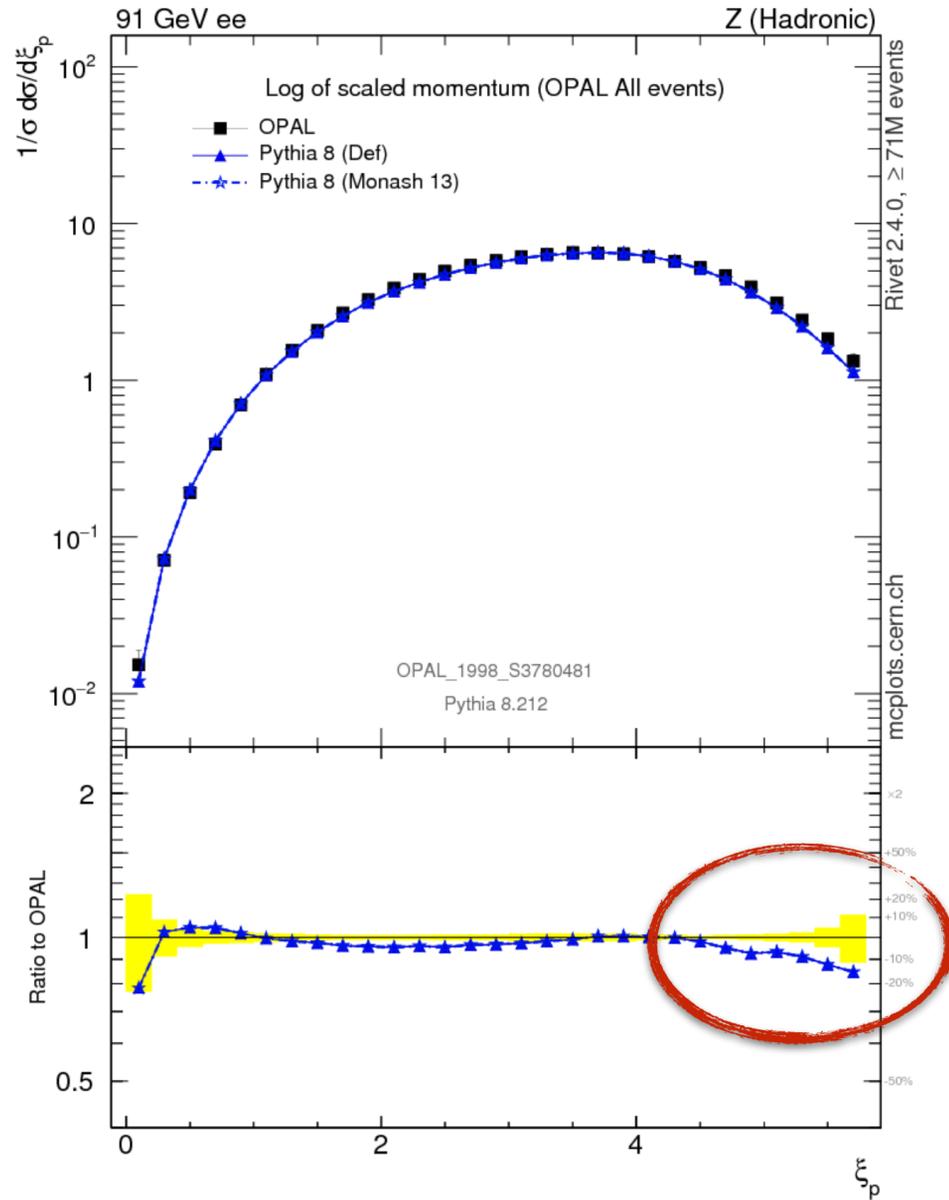
Kluth: if needed, could get O(LEP) sample in ~ 1 minute running with lower B-field

**gluon FFs, heavy-quark FFs,  $p_T$  dependence in hadron + jet, polarisation,...**



# L3 are you crazy?

(plots from [mcplots.cern.ch](http://mcplots.cern.ch))



Point of view A: small effects, and didn't you say toy model anyway?

Point of view B: this illustrates the kinds of things we can examine, with precise measurements

Flavour (in)dependence? (Controlling for feed-down?) Gauss vs Thermal?

# Jet (Sub)Structure

## LEP: mainly 45-GeV quark jet fragmentation

Inclusive: gluon FF only appears at NLO

3-jet events. Game of low sensitivity (3<sup>rd</sup> jet) vs low statistics ( $Z \rightarrow bbg$ )

(Initially only “symmetric” events; compare q vs g jets directly in data)

Naive  $C_A/C_F$  ratios between quarks and gluons verified

Many subtleties. Coherent radiation  $\rightarrow$  no ‘independent fragmentation’, especially at large angles. Parton-level “gluon” only meaningful at LO.

## ▣▣▣▣ Quark/gluon separation/tagging

Note: highly relevant interplay with Q/G sep @ LHC & FCC-hh: S/B

Language evolved: Just like “a jet” is inherently ambiguous, “quark-like” or “gluon-like” jets are ambiguous concepts [See Les Houches arXiv:1605.04692](#)

Define taggers (**adjective**: “q/g-LIKE”) using only final-state observables

Optimise tagger(s) using clean (theory) references, like  $X \rightarrow qq$  vs  $X \rightarrow gg$



G. SOYEZ, K. HAMACHER, G. RAUCO, S. TOKAR, Y. SAKAKI

## Handles to split degeneracies

$H \rightarrow gg$  vs  $Z \rightarrow qq$

Can we get a sample of  $H \rightarrow gg$  pure enough for QCD studies?

Requires good  $H \rightarrow gg$  vs  $H \rightarrow bb$ ;

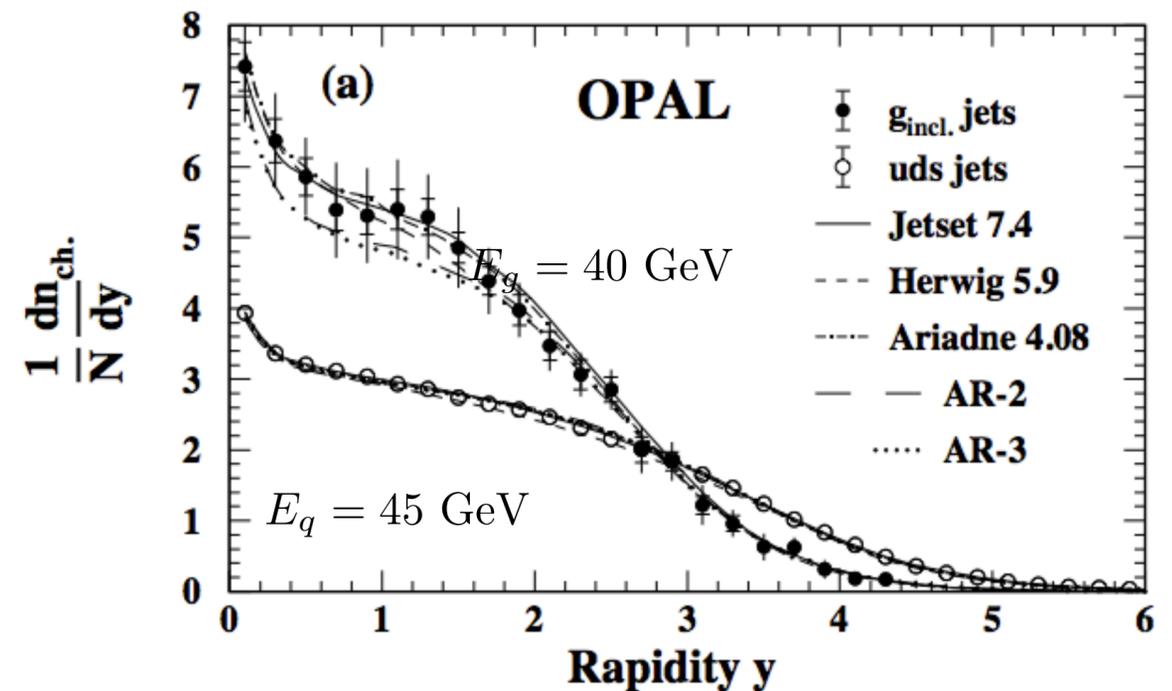
Driven by Higgs studies requirements?

$Z \rightarrow bbg$  vs  $Z \rightarrow qq(g)$

$g$  in one hemisphere recoils against  $b$ -jets in other hemisphere: **b tagging**

Study differential shape(s):  $N_{ch}$  (+low-R calo)

( $R \sim 0.1$  also useful for jet substructure)



## Scaling: radiative events $\rightarrow$ Forward Boosted

Scaling is **slow**, logarithmic  $\rightarrow$  prefer large lever arm

$E_{CM} > E_{Belle} \sim 10$  GeV [ **$\sim 10$  events / GeV at LEP**];

Useful benchmarks could be  $E_{CM} \sim 10$  (cross checks with Belle), 20, **30** (geom. mean between Belle and  $m_Z$ ), 45 GeV ( $=m_Z/2$ ) and 80 GeV =  $m_W$

(Also useful for FFs & general scaling studies)

# Unordered Clusterings of 4-Jet Events ( $ee$ $k_T$ , E scheme)

$$\frac{y_{34}}{y_{34} + y_{23}}$$

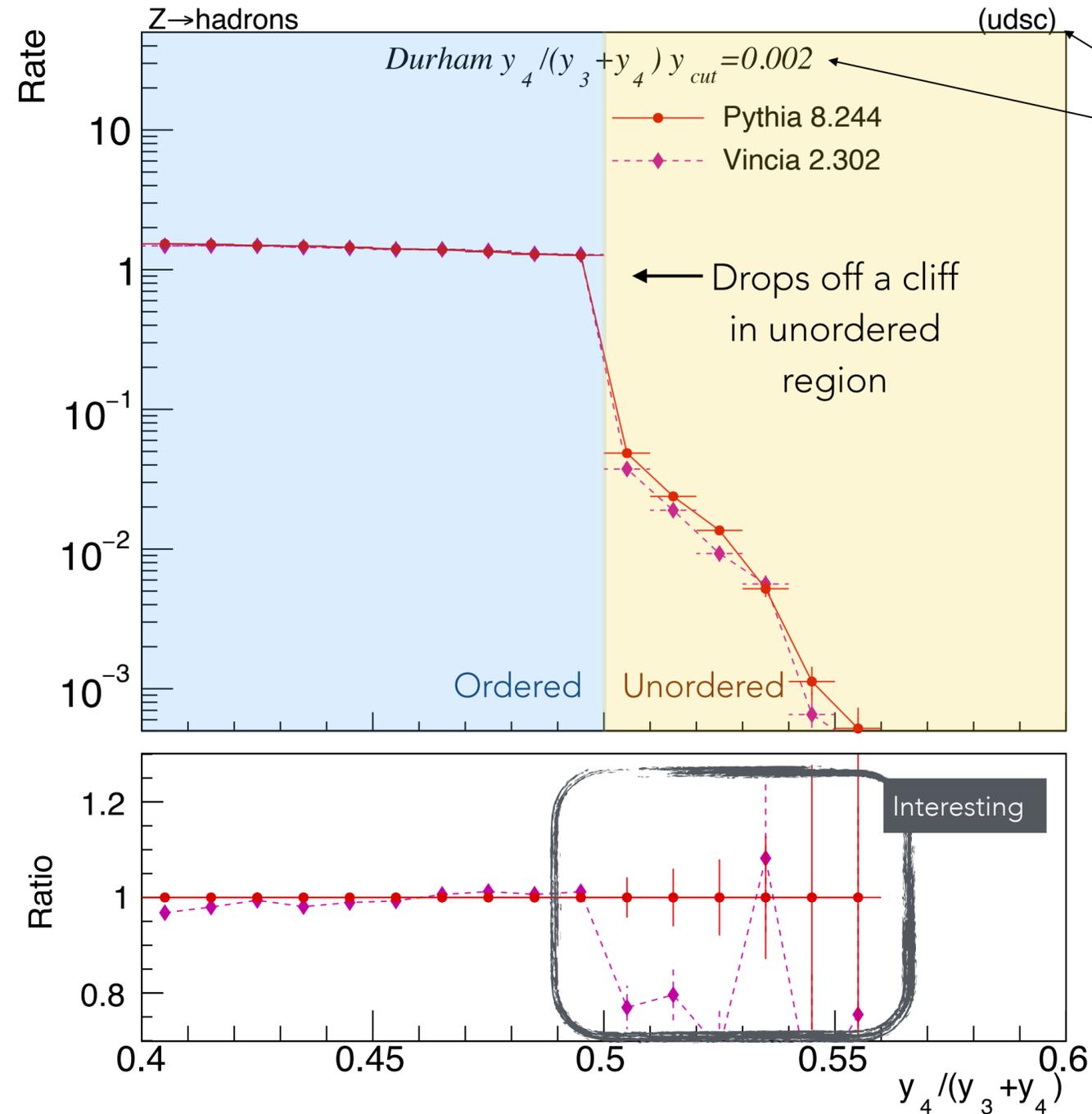
Rate normalised to total 4-jet rate

Off-the-shelf versions of Pythia and Vincia

Very similar results on individual jet rates.

Neither includes direct  $2 \rightarrow 4$ .

$4 \rightarrow 3 \rightarrow 2$



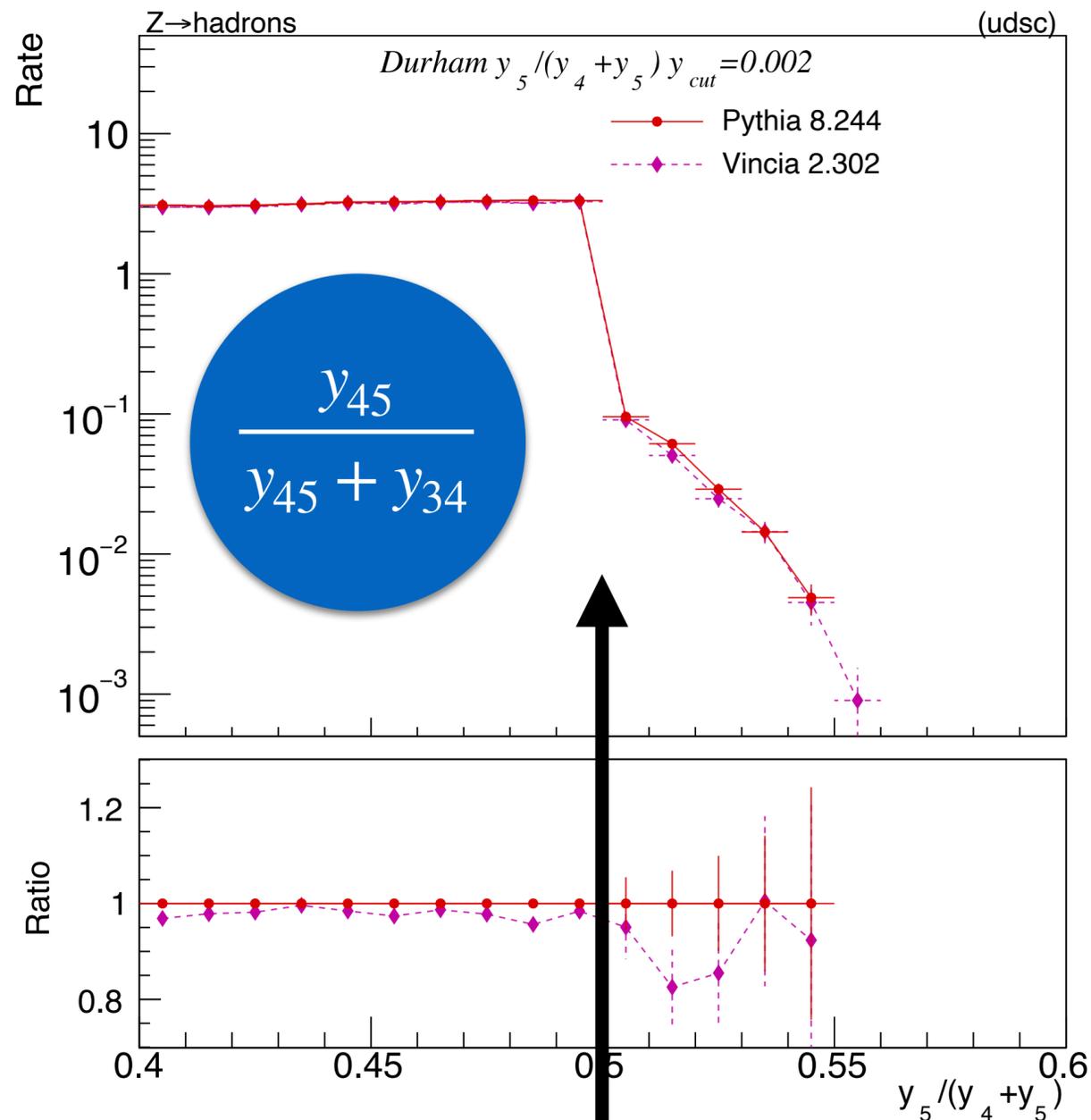
Small  $y_{cut} = 0.002$   
 ( $\leftrightarrow k_{\perp} \sim 4$  GeV) to maximise statistics  
 Excluded  $Z \rightarrow b\bar{b}$  to avoid contamination from B decays  
 4M events ( $\sim$  LEP 1)

(did not check the "interference" version of this observable here)

Q: could also be done for jet (sub)structure at the **LHC**?

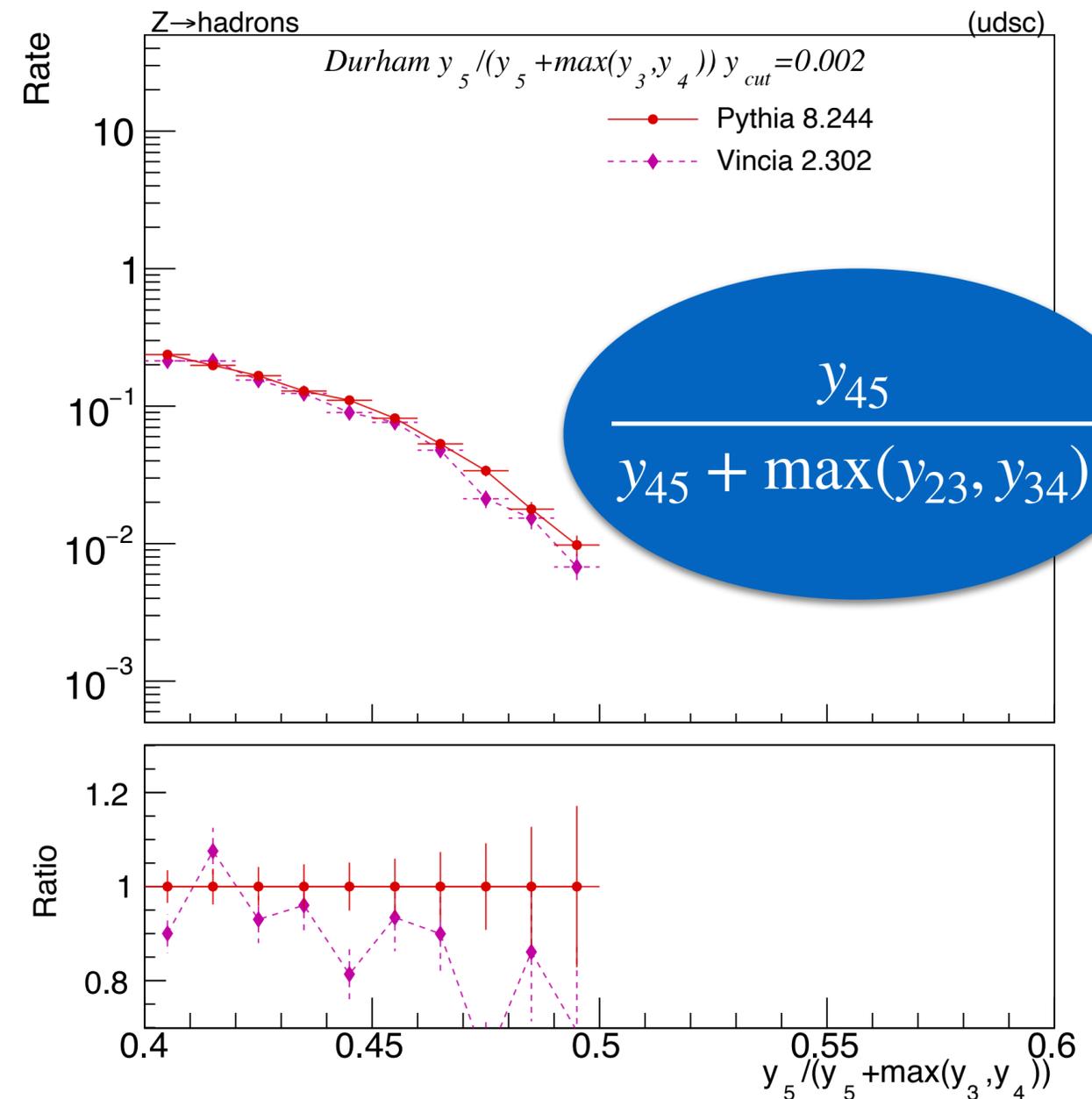
# 5-Jet Events

5 → 4 → 3



Same structure for 3 → 5 as for 2 → 4.  
(→ Combine to increase statistics?)

5 → 4 → 3 → 2



Limited power to probe 2 → 5  
(in this way) but worth an attempt?

# Triple-Energy Correlations

Suggested by Pier Monni, cf also 1912.11050

Generalisation of usual EEC, with relatively simple log structure.

Sensitive to triple-collinear?

I so far took a look at two triple-energy correlators:

“Equilateral”: all angles equal

“Planar”: two angles equal, the last one twice as large.

