

# Precision QCD at future $e^+e^-$ colliders (mostly FCC-ee)

**Intl Workshop on the high-energy  
Circular  $e^-e^+$  Collider (CEPC)  
Shanghai/Virtual, Oct. 2020**

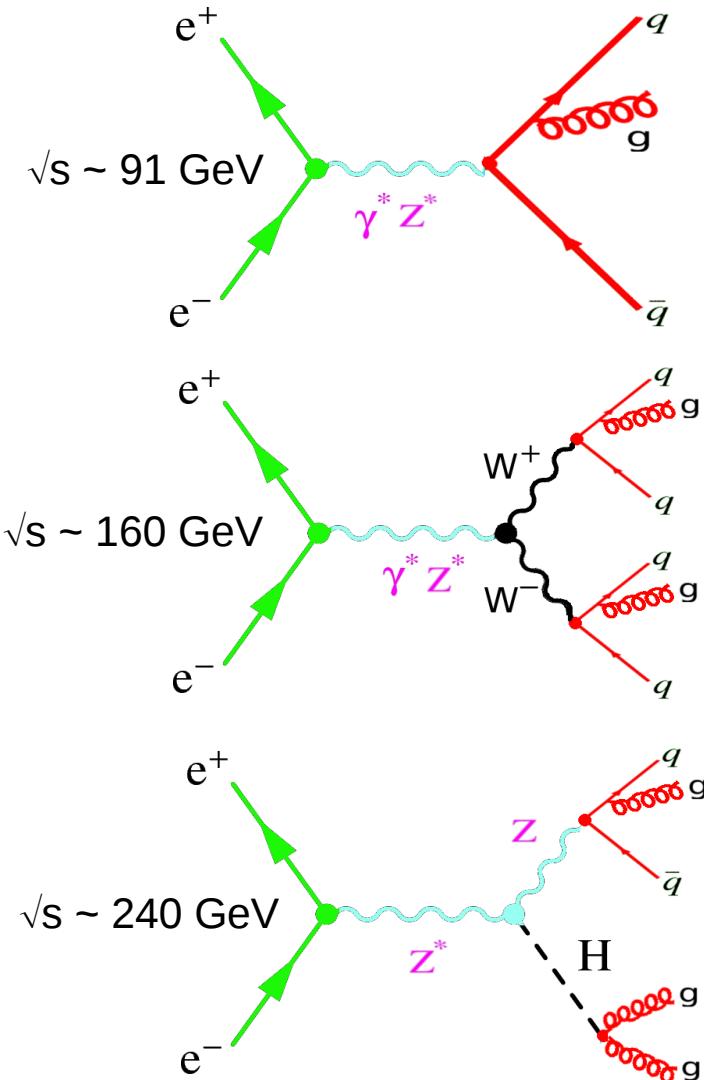
**David d'Enterria  
(CERN)**

# QCD = Key piece at future ee, pp colliders

- ▶ Though QCD is *not per se* the main driving force behind future colliders, QCD is crucial for many pp, ee measurements (signals & backgrounds):
  - **High-precision  $\alpha_s$** : Affects all x-sections & decays (esp. Higgs, top, EWPOs).
  - **N^nLO corrs., N^nLL resummations**: Affects all pQCD x-sections & decays.
  - **High-precision PDFs**: Affects all precision W,Z,H (**mid-x**) measurements & all searches (**high-x**) in pp collisions.
- **Heavy-Quark/Quark/Gluon separation** (jet substructure, boosted topologies..):  
Needed for all **precision SM** measurements & **BSM** searches with final jets.
- **Semihard QCD** (**low-x** gluon saturation, **multiple hard parton interactions,...**):  
Leading x-sections at FCC-pp (Note:  $Q_0 \sim 10$  GeV at 100 TeV).
- **Non-perturbative QCD**: Affects final-states with jets: **Colour reconnection**,  $e^+e^- \rightarrow Z, WW$ ,  $t\bar{t}$  → 4j, 6j... ( $m_W, m_{top}$  extractions). **Parton hadronization,...**

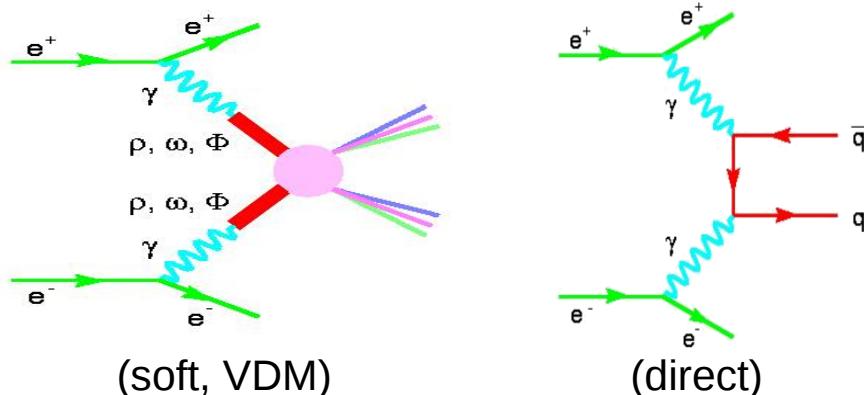
# Precision QCD in $e^+e^-$ collisions

- $e^+e^-$  collisions provide an **extremely clean** environment with fully-controlled initial-state to very precisely probe q,g dynamics:



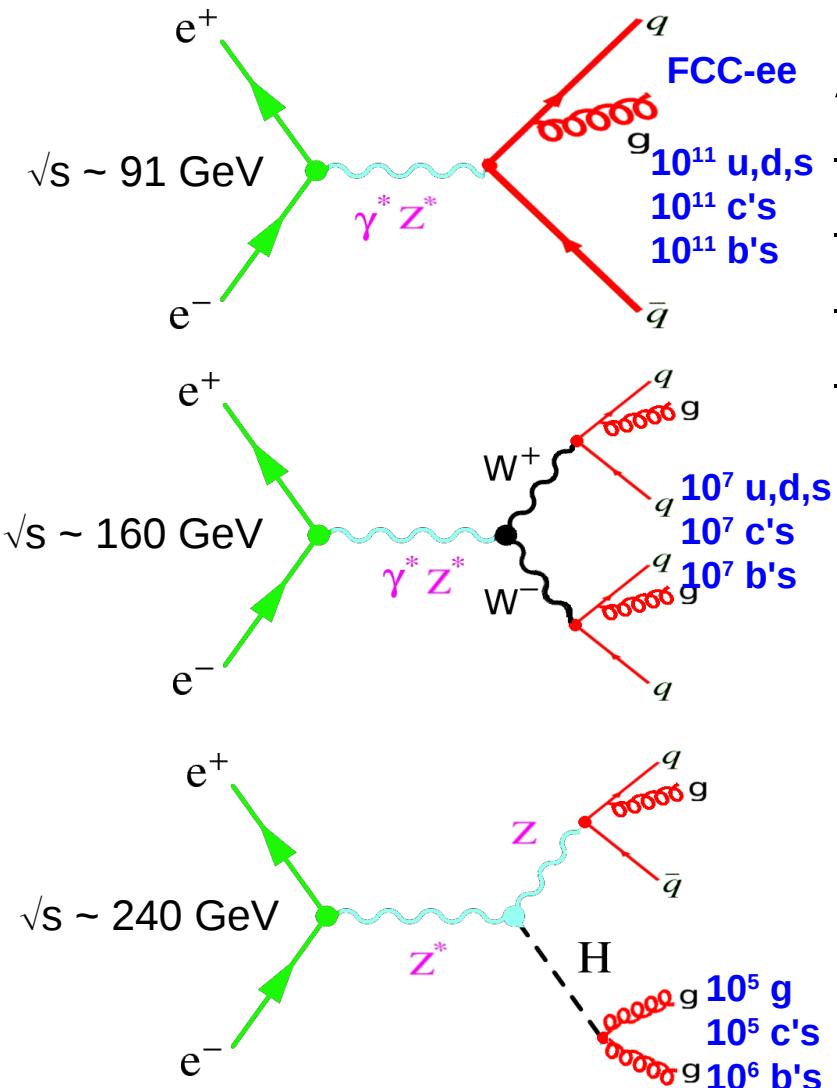
Advantages compared to p-p collisions:

- QED initial-state with **known kinematics**
  - **Controlled QCD radiation** (only in final-state)
  - Well-defined **heavy-Q, quark, gluon jets**
  - **Smaller non-pQCD uncertainties:**  
no PDFs, no QCD “underlying event”,...
- Direct clean parton fragmentation & hadroniz.
- Plus **QCD physics** in  $\gamma\gamma$  (EPA) collisions:

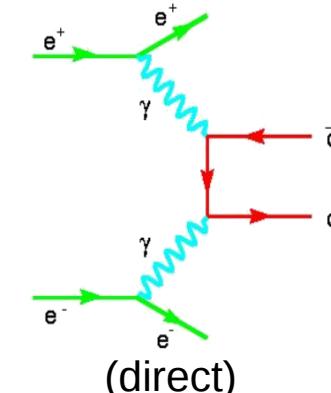
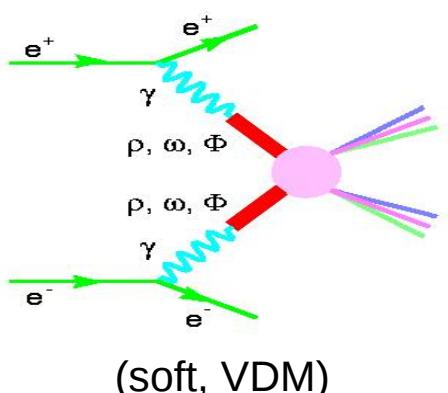


# Precision QCD in $e^+e^-$ collisions (FCC-ee)

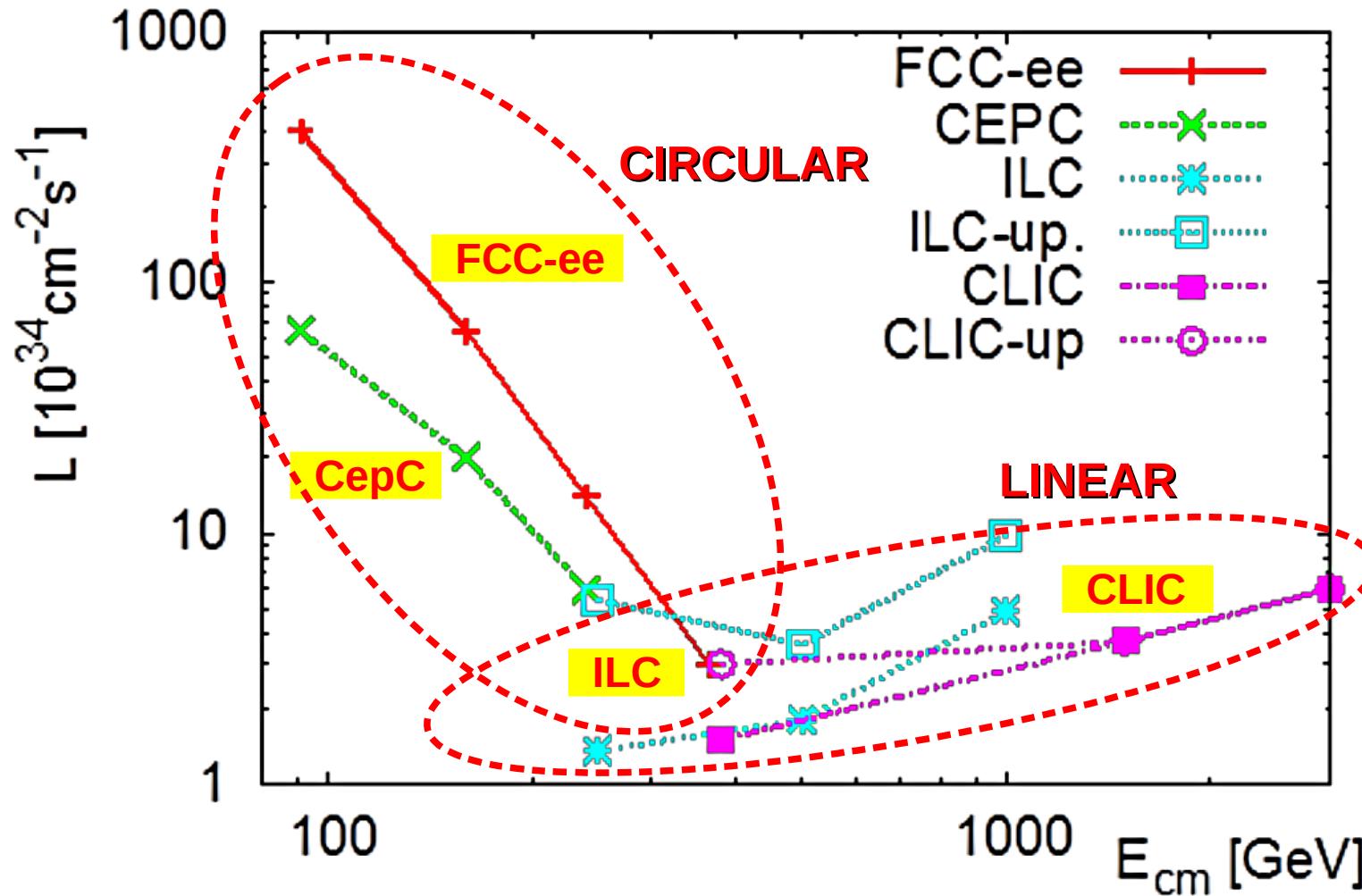
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  - Smaller non-pQCD uncertainties:
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- Direct clean parton fragmentation & hadroniz.
- Plus **QCD physics in  $\gamma\gamma$  (EPA) collisions:**



# Future $e^+e^-$ colliders under discussion



- FCC-ee features lumis a few times larger than other machines over 90–300 GeV
- Unparalleled Z, W, jets,  $\tau$ ,... data sets: Negligible stat. uncertainties

# **QCD physics at future $e^+e^-$ machines**

**(1) QCD coupling**

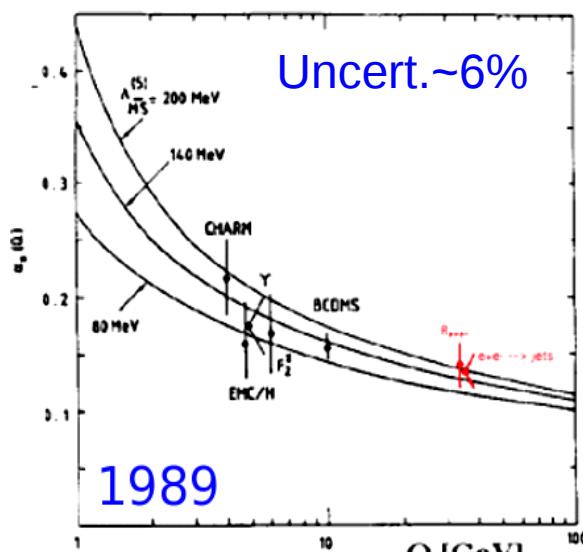
**(2) Jet substructure & flavour tagging**

**(3) Parton shower & Non-perturbative QCD**

NOTE: Only UNIQUE QCD measurements, inaccessible at any current machine, are covered.

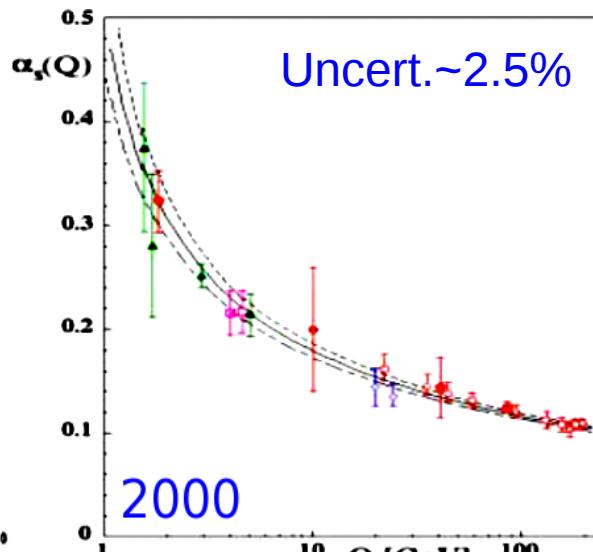
# QCD coupling $\alpha_s$

- Determines **strength of the strong interaction** between quarks & gluons.
- Single free parameter of QCD in the  $m_q \rightarrow 0$  limit.
- Determined at a ref. scale ( $Q=m_Z$ ), decreases as  $\alpha_s \sim \ln(Q^2/\Lambda^2)^{-1}$ ,  $\Lambda \sim 0.2$  GeV



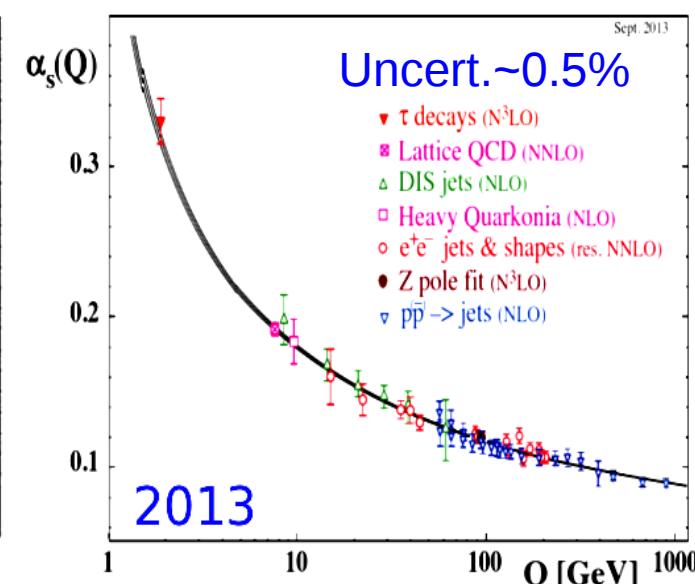
$$\alpha_s(M_Z) = 0.110^{+0.006}_{-0.008} \text{ (NLO)}$$

G. Altarelli, Ann. Rev. Nucl. Part. Sci. 39, 1989



$$\alpha_s(M_Z) = 0.1184 \pm 0.0031 \text{ (NNLO)}$$

S. B., J. Phys. G 26, 2000

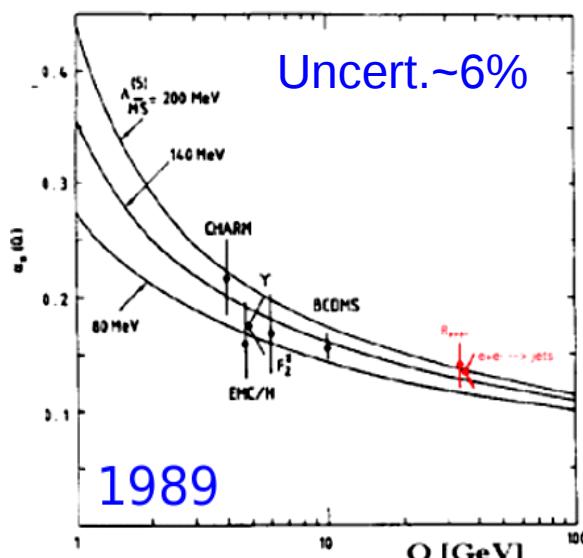


$$\alpha_s(M_Z) = 0.1185 \pm 0.0006 \text{ (NNLO)}$$

David d'Enterria (CERN)

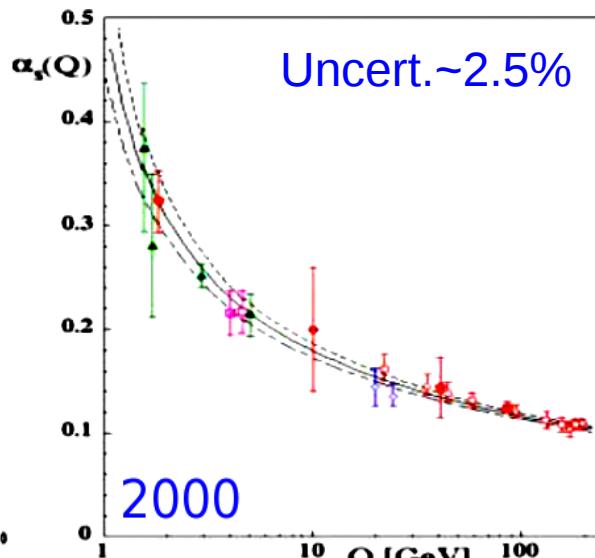
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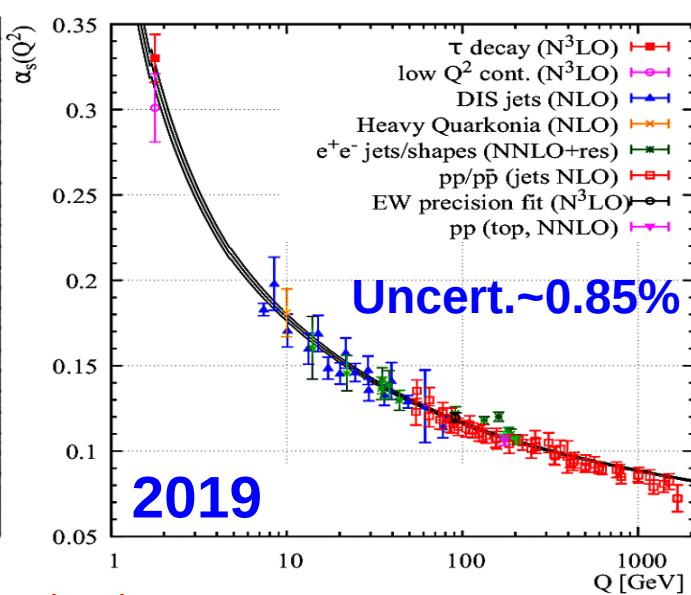
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$$\equiv \alpha_s(M_Z^2) = 0.1179 \pm 0.0010$$

- Least precisely known of all interaction **couplings** !

$$\delta\alpha \sim 10^{-10} \ll \delta G_F \ll 10^{-7} \ll \delta G \sim 10^{-5} \ll \delta\alpha_s \sim 10^{-3}$$

# Importance of the QCD coupling $\alpha_s$

→ Impacts all QCD x-sections & decays (H), precision top & parametric EWPO:

Process	$\sigma$ (pb)	$\delta\alpha_s$ (%)	PDF + $\alpha_s$ (%)	Scale(%)
ggH	49.87	$\pm 3.7$	-6.2 +7.4	-2.61 + 0.32
tH	0.611	$\pm 3.0$	$\pm 8.9$	-9.3 + 5.9

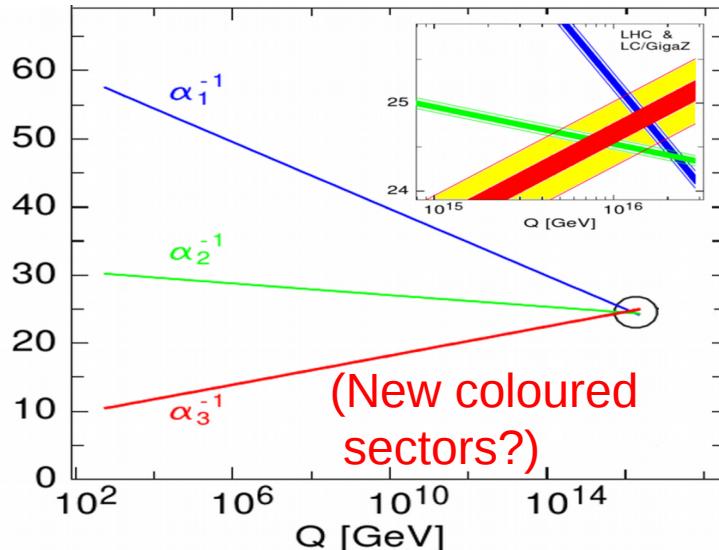
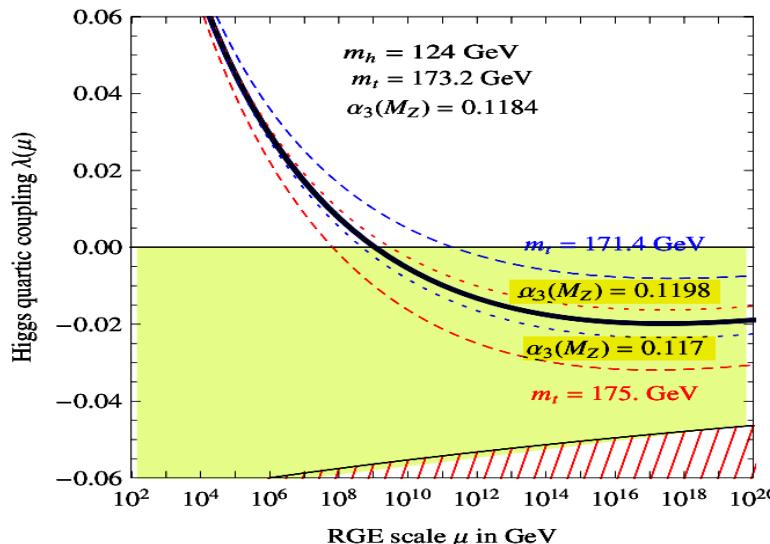
Channel	$M_H$ [GeV]	$\delta\alpha_s$ (%)	$\Delta m_b$	$\Delta m_c$
$H \rightarrow c\bar{c}$	126	$\pm 7.1$	$\pm 0.1\%$	$\pm 2.3\%$
$H \rightarrow gg$	126	$\pm 4.1$	$\pm 0.1\%$	$\pm 0\%$

Msbar mass error budget (from threshold scan)			
$(\delta M_t^{\text{SD-low}})^{\text{exp}}$	$(\delta M_t^{\text{SD-low}})^{\text{theo}}$	$(\delta \bar{m}_t(\bar{m}_t))^{\text{conversion}}$	$(\delta \bar{m}_t(\bar{m}_t))^{\alpha_s}$
40 MeV	50 MeV	7 – 23 MeV	70 MeV
⇒ improvement in $\alpha_s$ crucial			$\delta\alpha_s(M_z) = 0.001$

Quantity	FCC-ee	future param. unc.	Main source
$\Gamma_Z$ [MeV]	0.1	0.1	$\delta\alpha_s$
$R_b$ [ $10^{-5}$ ]	6	< 1	$\delta\alpha_s$
$R_\ell$ [ $10^{-3}$ ]	1	1.3	$\delta\alpha_s$

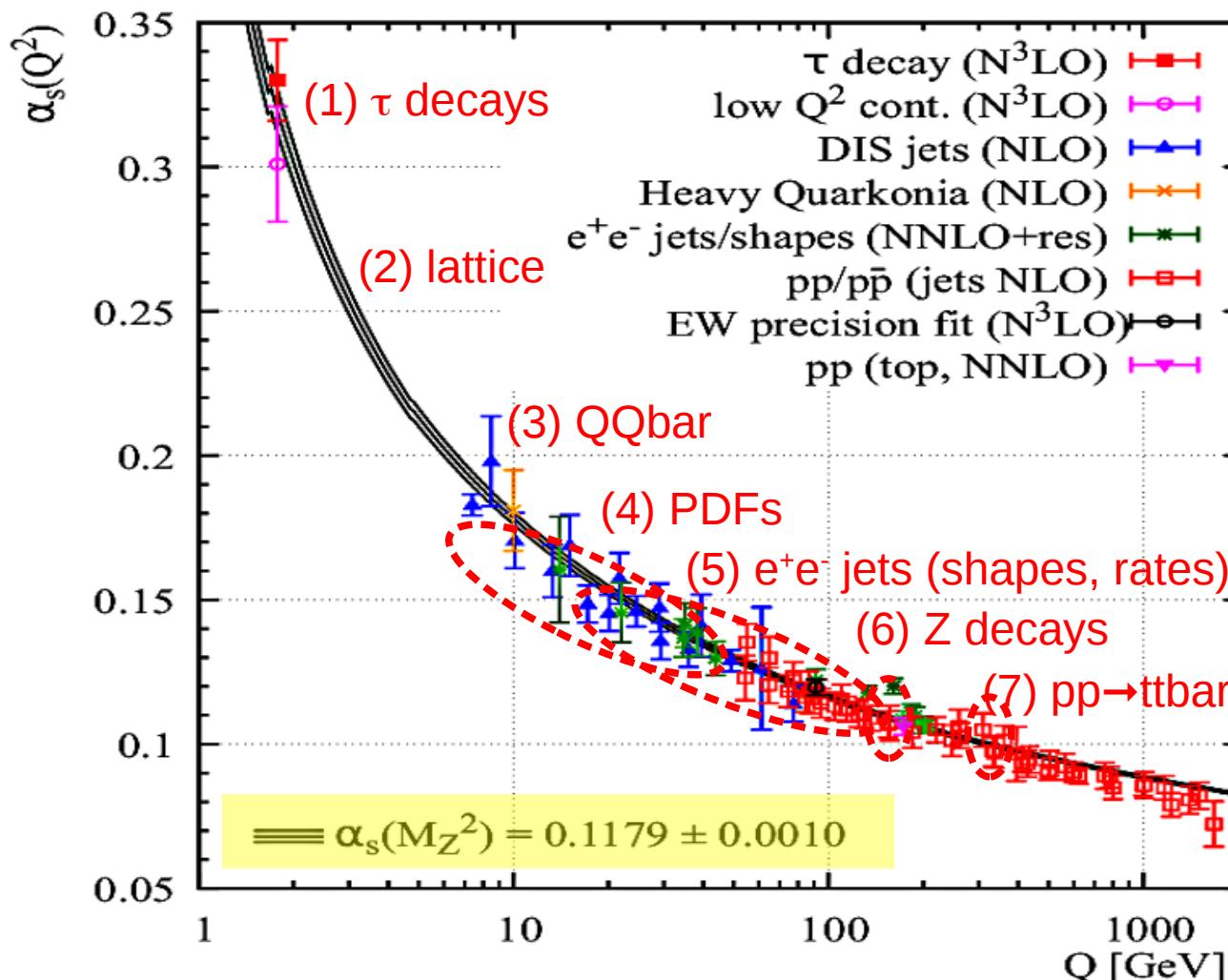
Sven Heinemeyer – 1st FCC physics workshop, CERN, 17.01.2017

→ Impacts physics approaching Planck scale: EW vacuum stability, GUT



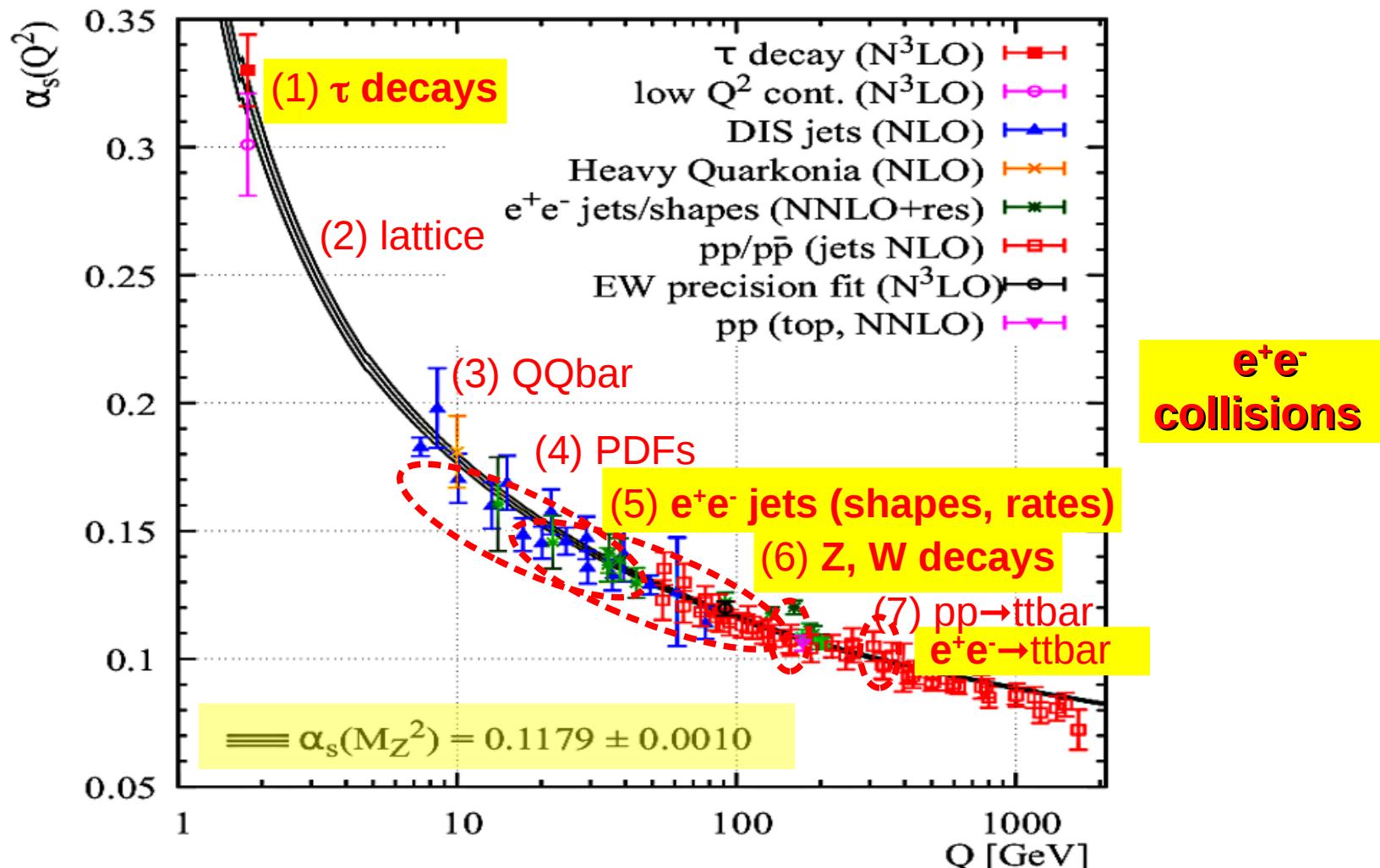
# World $\alpha_s$ determination (PDG 2019)

- Determined today by comparing 7 experimental observables to pQCD NNLO,N<sup>3</sup>LO predictions, plus global average at the Z pole scale:



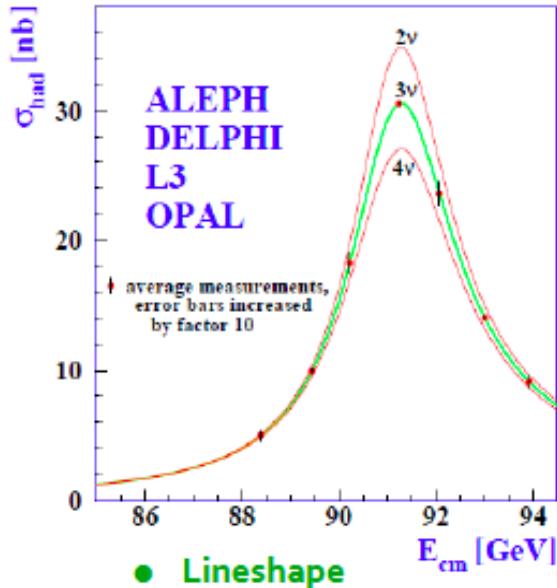
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# Ultra-precise W, Z, top physics at FCC-ee

$\sqrt{s}=91 \text{ GeV}, 10^{12} Z's$



- Exquisite  $E_{\text{beam}}$  (unique!)

- $m_Z, \Gamma_Z$  to 10 keV (stat.)

- Asymmetries 100 keV (syst.)

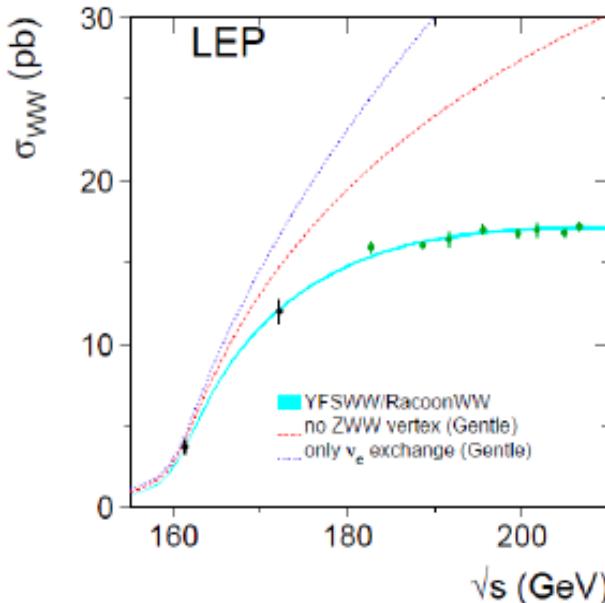
- $\sin^2 \theta_W$  to  $5 \times 10^{-6}$

- Branching ratios,  $R_b, R_{\text{had}}$

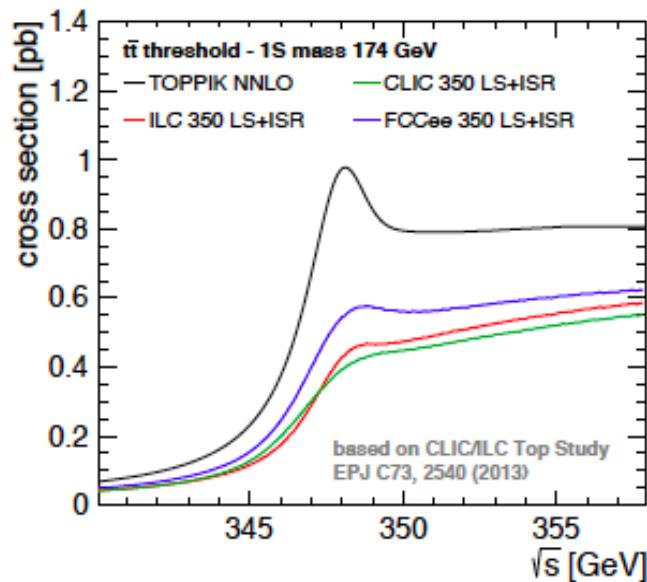
- $\alpha_s(m_Z)$  to 0.0002

- Predict  $m_{\text{top}}, m_W$  in SM

$\sqrt{s}=161 \text{ GeV}, 10^8 W's$



$\sqrt{s}=350 \text{ GeV}, 10^6 \text{ tops}$



- Threshold scan

- $m_W$  to 500 keV

- Branching ratios  $R_b, R_{\text{had}}$

- $\alpha_s(m_W)$  to 0.0002

- Radiative returns  $e^+e^- \rightarrow \gamma Z$  ( $Z \rightarrow \nu\nu, \mu^+\mu^-$ )

- $N_\nu$  to 0.001

- Threshold scan + 4D fit

- $m_{\text{top}}$  to 10 MeV (stat.)  
40 MeV (th.)

- $\lambda_{\text{top}}$  to 13%

- EWK couplings to 1–10%

■ Mostly thanks to:

(i) Huge statistics

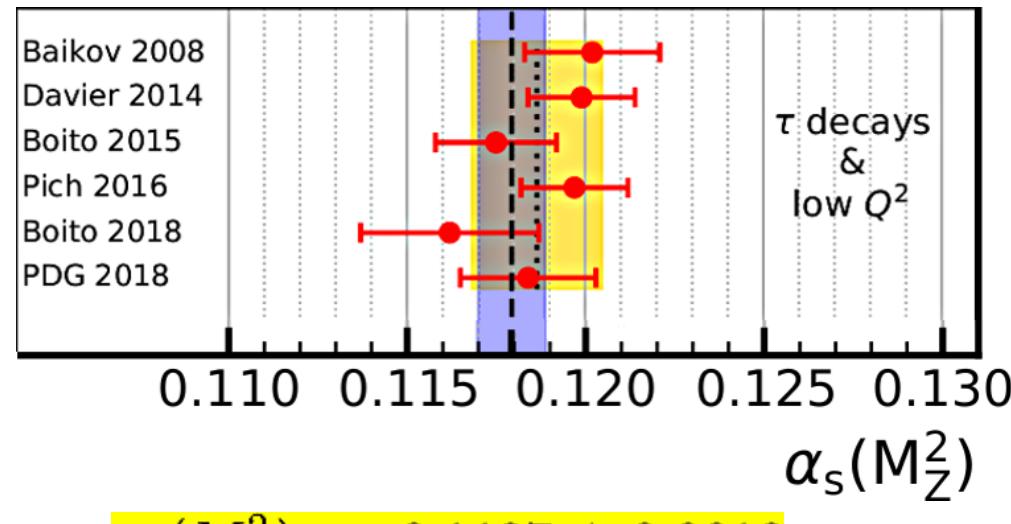
(ii) Threshold scans with  $\delta E_{\text{cm}} \sim 0.1, 0.3, 2., 4.$  MeV ( $Z, W, H, t$ )

# $\alpha_s$ from hadronic $\tau$ -lepton decays

→ Computed at  $N^3\text{LO}$ :  $R_\tau \equiv \frac{\Gamma(\tau^- \rightarrow \nu_\tau + \text{hadrons})}{\Gamma(\tau^- \rightarrow \nu_\tau e^- \bar{\nu}_e)} = S_{\text{EW}} N_C (1 + \sum_{n=1}^4 c_n \left(\frac{\alpha_s}{\pi}\right)^n + \mathcal{O}(\alpha_s^5) + \delta_{\text{np}})$

→ Experimentally:  $R_{\tau, \text{exp}} = 3.4697 \pm 0.0080 (\pm 0.23\%)$

→ Various pQCD approaches (FOPT vs CIPT) & treatment of non-pQCD corrections ( $(\Lambda/m_\tau)^2 \sim 2\%$ , yield different results).



Uncertainty slightly increased:  
2013 ( $\pm 1.3\%$ ) → 2019 ( $\pm 1.5\%$ )

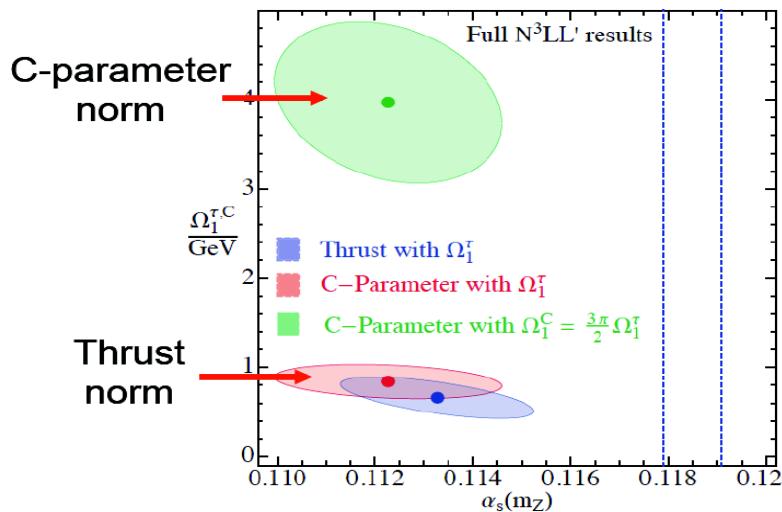
$$\alpha_s(M_Z^2) = 0.1187 \pm 0.0018 (\pm 1.5\%)$$

→ Future:

- TH: Better understanding of FOPT vs CIPT differences.
- Better spectral functions needed (high stats & better precision): B-factories (BELLE-II)?
- High-stats:  $\mathcal{O}(10^{11})$  from  $Z \rightarrow \tau\tau$  at FCC-ee(90) :  $\delta\alpha_s/\alpha_s \ll 1\%$

# $\alpha_s$ from $e^+e^-$ event shapes & jet rates (today)

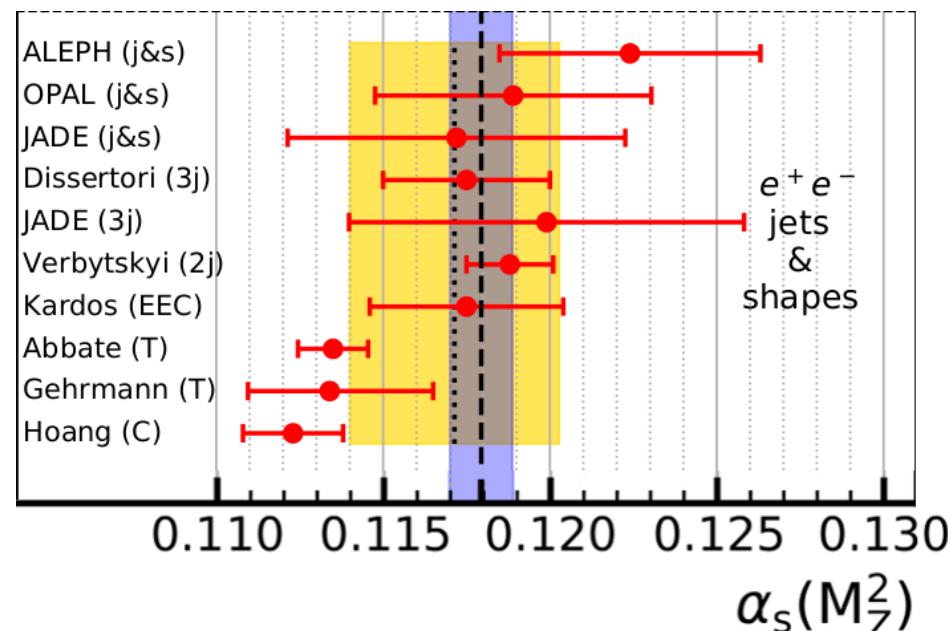
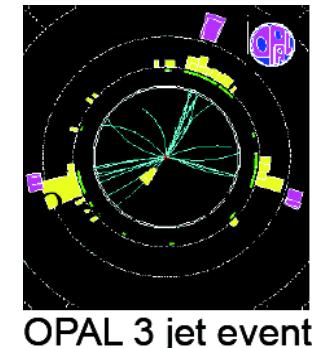
- Computed at  $N^{2,3}\text{LO} + N^{(2)}\text{LL}$  accuracy.
- Experimentally (LEP):
  - Thrust, C-parameter, jet shapes
  - n-jet x-sections
- Results sensitive to non-pQCD (hadronization) accounted for via MCs or analytically:



- Wide span of TH extractions...

$$\tau = 1 - \max_{\hat{n}} \frac{\sum |\vec{p}_i \cdot \hat{n}|}{\sum |\vec{p}_i|}$$

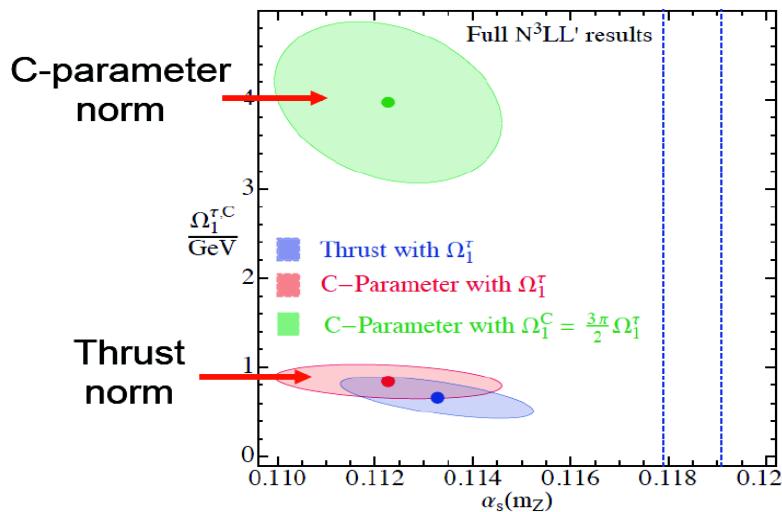
$$C = \frac{3}{2} \frac{\sum_{i,j} |\vec{p}_i||\vec{p}_j| \sin^2 \theta_{ij}}{\left(\sum_i |\vec{p}_i|\right)^2}$$



$$\alpha_s(M_Z^2) = 0.1171 \pm 0.0031 \quad (\pm 2.6\%)$$

# $\alpha_s$ from $e^+e^-$ event shapes & jet rates (future)

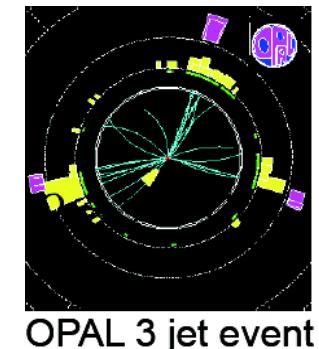
- Computed at  $N^{2,3}\text{LO} + N^{(2)}\text{LL}$  accuracy.
- Experimentally (LEP):
  - Thrust, C-parameter, jet shapes
  - 3-jet x-sections
- Results sensitive to non-pQCD (hadronization) accounted for via MCs or analytically:



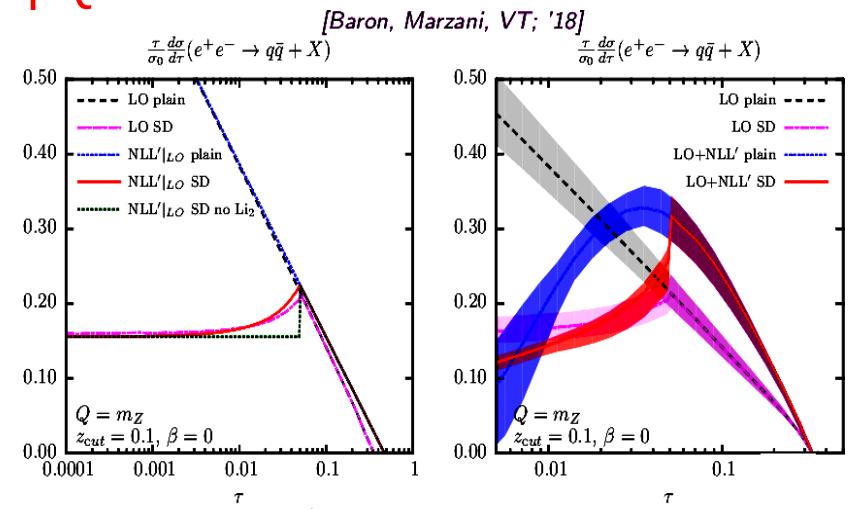
- Future:  $\delta\alpha_s/\alpha_s < 1\%$ 
  - FCC-e<sup>+</sup>e<sup>-</sup>: Lower- $\sqrt{s}$  (ISR) for shapes, higher- $\sqrt{s}$  for jet rates
  - TH: Improved ( $N^{2,3}\text{LL}$ ) resummation for rates, hadronization for shapes

$$\tau = 1 - \max_{\hat{n}} \frac{\sum |\vec{p}_i \cdot \hat{n}|}{\sum |\vec{p}_i|}$$

$$C = \frac{3}{2} \frac{\sum_{i,j} |\vec{p}_i||\vec{p}_j| \sin^2 \theta_{ij}}{\left(\sum_i |\vec{p}_i|\right)^2}$$



- Modern jet substructure techniques:  
“Soft drop” can help reduce non-pQCD corrections for thrust:



# $\alpha_s$ from hadronic Z, W decays

→ Z & W observables theoretically known at N<sup>3</sup>LO accuracy:

- The W and Z hadronic widths :

$$\Gamma_{W,Z}^{\text{had}}(Q) = \Gamma_{W,Z}^{\text{Born}} \left( 1 + \sum_{i=1}^4 a_i(Q) \left( \frac{\alpha_S(Q)}{\pi} \right)^i + \mathcal{O}(\alpha_S^5) + \delta_{\text{EW}} + \delta_{\text{mix}} + \delta_{\text{np}} \right)$$

- The ratio of W, Z hadronic-to-leptonic widths :

$$R_{W,Z}(Q) = \frac{\Gamma_{W,Z}^{\text{had}}(Q)}{\Gamma_{W,Z}^{\text{lep}}(Q)} = R_{W,Z}^{\text{EW}} \left( 1 + \sum_{i=1}^4 a_i(Q) \left( \frac{\alpha_S(Q)}{\pi} \right)^i + \mathcal{O}(\alpha_S^5) + \delta_{\text{mix}} + \delta_{\text{np}} \right)$$

- In the Z boson case, the hadronic cross section at the resonance peak in  $e^+e^-$ :

$$\sigma_Z^{\text{had}} = \frac{12\pi}{m_Z} \cdot \frac{\Gamma_Z^e \Gamma_Z^{\text{had}}}{(\Gamma_Z^{\text{tot}})^2}$$

DdE, Jacobsen:  
arXiv:2005.04545 [hep-ph]

TH uncertainties:

( $\alpha^2, \alpha^3$  included for Z):

±0.015–0.03% (Z)

±0.015–0.04% (W)

Param. uncerts.:

( $m_{Z,W}, \alpha, V_{cs,ud}$ ):

±0.01–0.03% (Z)

±1.1–1.7% (W)

±0.03% (W, CKM unit)

→ Measured at LEP with ±0.1–0.3% (Z), ±0.9–2% (W) exp. uncertainties:

	theory			experiment			Recent update of LEP luminosity bias(*) change the Z values by few permil
	previous	new (this work)	change	previous [6]	new [20, 21]	change	
$\Gamma_Z^{\text{tot}}$ (MeV)	$2494.2 \pm 0.8_{\text{th}}$	$2495.2 \pm 0.6_{\text{par}} \pm 0.4_{\text{th}}$	+0.04%	$2495.2 \pm 2.3$	$2495.5 \pm 2.3$	+0.012%	
$R_Z$	$20.733 \pm 0.007_{\text{th}}$	$20.750 \pm 0.006_{\text{par}} \pm 0.006_{\text{th}}$	+0.08%	$20.767 \pm 0.025$	$20.7666 \pm 0.0247$	-0.040%	
$\sigma_Z^{\text{had}}$ (pb)	$41490 \pm 6_{\text{th}}$	$41494 \pm 5_{\text{par}} \pm 6_{\text{th}}$	+0.01%	$41540 \pm 37$	$41480.2 \pm 32.5$	-0.144%	

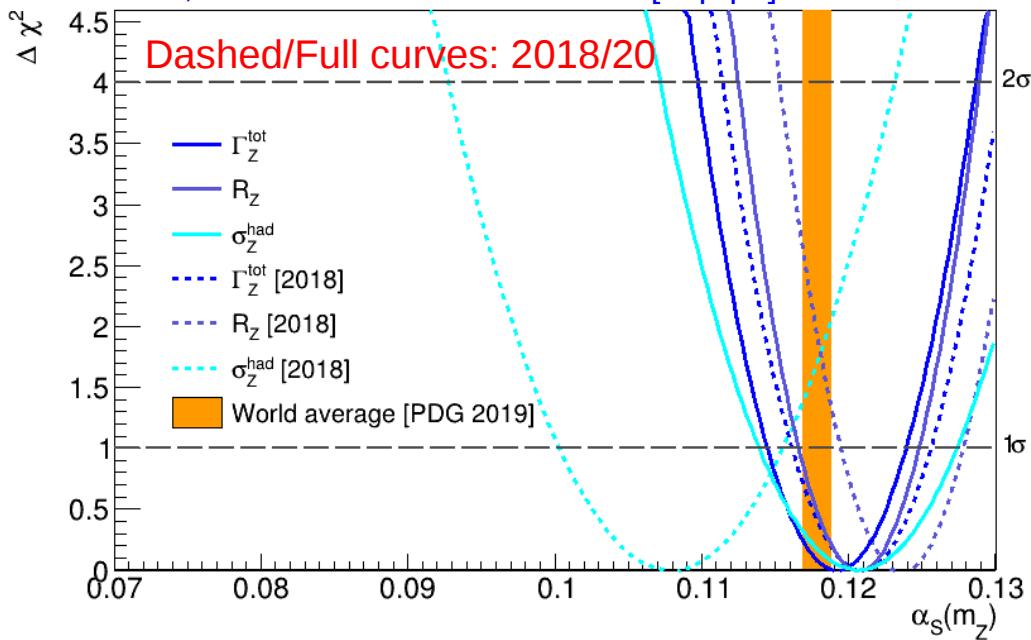
W boson observables	GFITTER 2.2 (NNLO)			this work (N <sup>3</sup> LO)		experiment
	(exp. CKM)			(CKM unit.)		
$\Gamma_W^{\text{had}}$ (MeV)	—			$1440.3 \pm 23.9_{\text{par}} \pm 0.2_{\text{th}}$	$1410.2 \pm 0.8_{\text{par}} \pm 0.2_{\text{th}}$	$1405 \pm 29$
$\Gamma_W^{\text{tot}}$ (MeV)	$2091.8 \pm 1.0_{\text{par}}$			$2117.9 \pm 23.9_{\text{par}} \pm 0.7_{\text{th}}$	$2087.9 \pm 1.0_{\text{par}} \pm 0.7_{\text{th}}$	$2085 \pm 42$
$R_W$	—			$2.1256 \pm 0.0353_{\text{par}} \pm 0.0008_{\text{th}}$	$2.0812 \pm 0.0007_{\text{par}} \pm 0.0008_{\text{th}}$	$2.069 \pm 0.019$

(\*) Voutsinas et al.  
arXiv:1908.01704,  
Janot et al.  
arXiv:1912.02067

# $\alpha_s$ from hadronic Z decays (today)

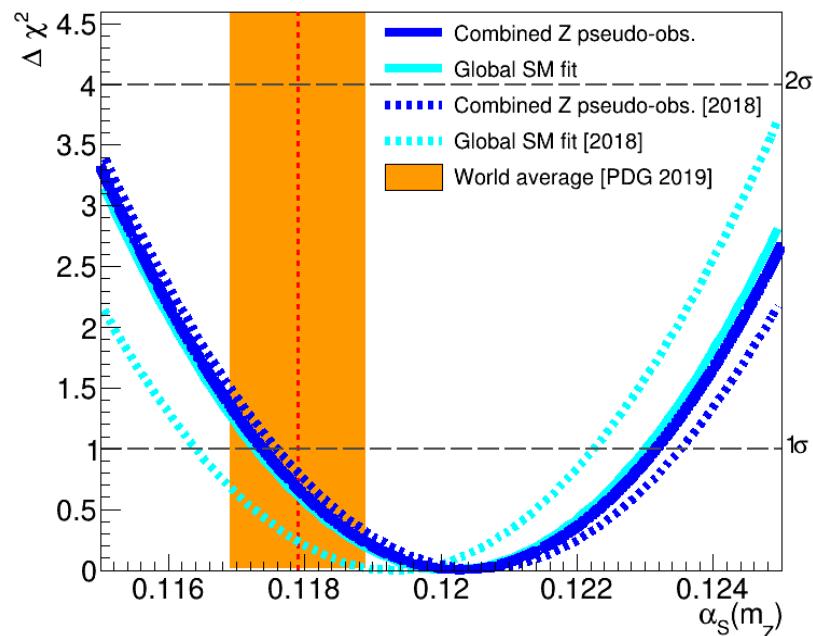
- QCD coupling extracted from:
  - (i) Combined fit of 3 Z pseudo-observ.
  - (ii) Full SM fit (with  $\alpha_s$  free parameter)

DdE, Jacobsen: arXiv:2005.04545 [hep-ph]



- LEP lumi-bias updates lead to much better agreement among  $\Gamma_z$ ,  $R_z$ ,  $\sigma_0$  extractions:
- Improved  $\alpha_s(m_z) = 0.1203 \pm 0.0028$  ( $\pm 2.3\%$ )  
 $PDG'19: \alpha_s(m_z) = 0.1205 \pm 0.0030$  ( $\pm 2.5\%$ )

Z boson observable	$\alpha_s(m_Z)$ extraction	uncertainties		
	exp.	param.	theor.	
$\Gamma_Z^{\text{tot}}$	$0.1192 \pm 0.0047$	$\pm 0.0046$	$\pm 0.0005$	$\pm 0.0008$
$R_Z$	$0.1207 \pm 0.0041$	$\pm 0.0041$	$\pm 0.0001$	$\pm 0.0009$
$\sigma_Z^{\text{had}}$	$0.1206 \pm 0.0068$	$\pm 0.0067$	$\pm 0.0004$	$\pm 0.0012$
All combined	$0.1203 \pm 0.0029$	$\pm 0.0029$	$\pm 0.0002$	$\pm 0.0008$
Global SM fit	$0.1202 \pm 0.0028$	$\pm 0.0028$	$\pm 0.0002$	$\pm 0.0008$



- EXP/TH updates lead to better agreement with full SM fit:
- $\alpha_s(m_z) = 0.1202 \pm 0.0028$   
 $PDG'19: \alpha_s(m_z) = 0.1194 \pm 0.0029$

# $\alpha_s$ from hadronic Z decays (FCC-ee)

→ QCD coupling extracted from:

- (i) Combined fit of 3 Z pseudo-observ:
- (ii) Full SM fit (with  $\alpha_s$  free parameter)

Z boson observable	$\alpha_s(m_Z)$ extraction	uncertainties		
		exp.	param.	theor.
All combined	$0.1203 \pm 0.0029$	$\pm 0.0029$	$\pm 0.0002$	$\pm 0.0008$
Global SM fit	$0.1202 \pm 0.0028$	$\pm 0.0028$	$\pm 0.0002$	$\pm 0.0008$
All combined (FCC-ee)	$0.12030 \pm 0.00026$	$\pm 0.00013$	$\pm 0.00005$	$\pm 0.00022$
Global SM fit (FCC-ee)	$0.12020 \pm 0.00026$	$\pm 0.00013$	$\pm 0.00005$	$\pm 0.00022$

→ FCC-ee:

- Huge Z pole stats. ( $\times 10^5$  LEP)
- Exquisite systematic/parametric precision (stat. uncert. much smaller):

$$\Delta R_Z = 10^{-3}, \quad R_Z = 20.7500 \pm 0.0010$$

$$\Delta \Gamma_Z^{\text{tot}} = 0.1 \text{ MeV}, \quad \Gamma_Z^{\text{tot}} = 2495.2 \pm 0.1 \text{ MeV}$$

$$\Delta \sigma_Z^{\text{had}} = 4.0 \text{ pb}, \quad \sigma_Z^{\text{had}} = 41494 \pm 4 \text{ pb}$$

$$\Delta m_Z = 0.1 \text{ MeV}, \quad m_Z = 91.18760 \pm 0.00001 \text{ GeV}$$

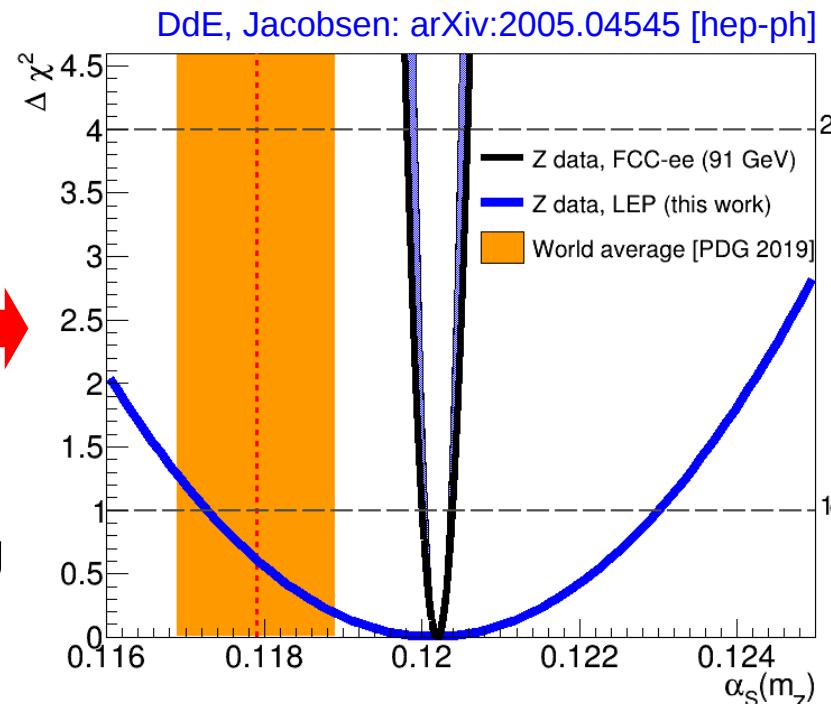
$$\Delta \alpha = 3 \cdot 10^{-5}, \quad \Delta \alpha_{\text{had}}^{(5)}(m_Z) = 0.0275300 \pm 0.0000009$$

- TH uncertainty reduced by  $\times 4$  computing missing  $\alpha_s^5, \alpha^3, \alpha \alpha_s^2, \alpha \alpha_s^2, \alpha^2 \alpha_s$  terms

→ 10 times better precision than today:

$$\delta \alpha_s / \alpha_s \sim \pm 0.2\% \text{ (tot), } \pm 0.1\% \text{ (exp)}$$

Strong (B)SM consistency test.



$$\alpha_s(m_Z) = 0.12030 \pm 0.00028 \quad (\pm 0.2\%)$$

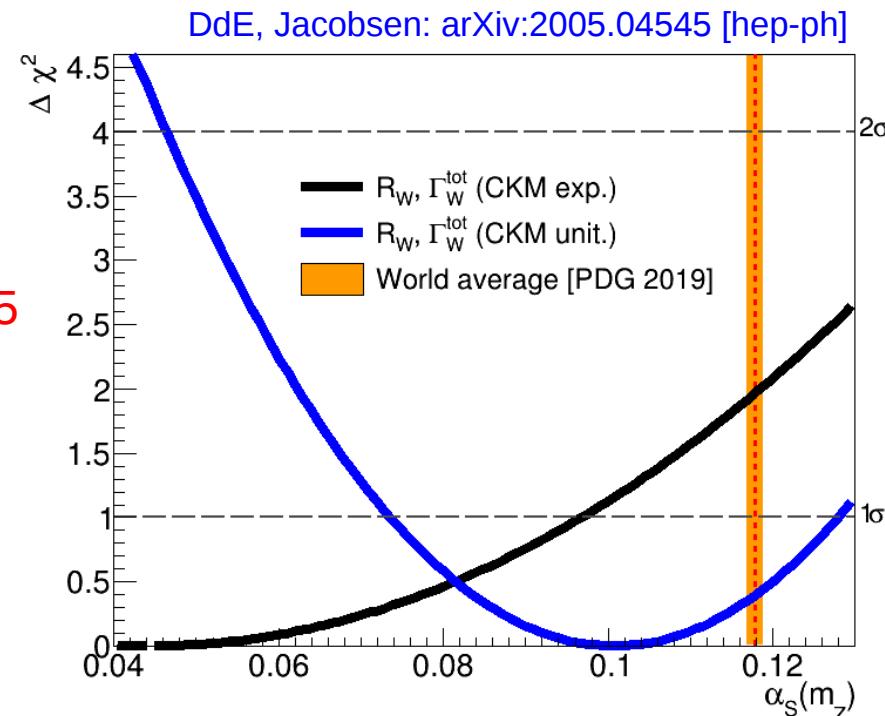
# $\alpha_s$ from hadronic W decays (today)

- QCD coupling extracted from new N<sup>3</sup>LO fit of combined  $\Gamma_w$ ,  $R_w$  pseudo-observ.:

W boson observables	$\alpha_S(m_Z)$ extraction	uncertainties		
		exp.	param.	theor.
$\Gamma_W^{\text{tot}}$ , $R_w$ (exp. CKM)	$0.044 \pm 0.052$	$\pm 0.024$	$\pm 0.047$	( $\pm 0.0014$ )
$\Gamma_W^{\text{tot}}$ , $R_w$ (CKM unit.)	$0.101 \pm 0.027$	$\pm 0.027$	( $\pm 0.0002$ )	( $\pm 0.0016$ )
$\Gamma_W^{\text{tot}}$ , $R_w$ (FCC-ee, CKM unit.)	$0.11790 \pm 0.00023$	$\pm 0.00012$	$\pm 0.00004$	$\pm 0.00019$

► Very imprecise extraction:

- Large propagated parametric uncert. from poor  $V_{cs}$  exp. precision ( $\pm 2\%$ ): QCD coupling unconstrained:  $0.04 \pm 0.05$
- Imposing CKM unitarity: large exp. uncertainties from  $\Gamma_w$ ,  $R_w$  (0.9–2%): QCD extracted with ~27% precision
- Propagated TH uncertainty much smaller today: ~1.5%



$$\alpha_s(m_Z) = 0.101 \pm 0.027 \quad (\pm 27\%)$$

# $\alpha_s$ from hadronic W decays (FCC-ee)

► QCD coupling extracted from new N<sup>3</sup>LO fit of combined  $\Gamma_W$ ,  $R_W$  pseudo-observ.:

W boson observables	$\alpha_S(m_Z)$ extraction	uncertainties		
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$\Gamma_W^{\text{tot}}$ , $R_W$ (FCC-ee, CKM unit.)	$0.11790 \pm 0.00023$	$\pm 0.00012$	$\pm 0.00004$	$\pm 0.00019$

► FCC-ee extraction:

- Huge W pole stats. ( $\times 10^4$  LEP-2).
- Exquisite syst./parametric precision:

$$\Gamma_W^{\text{tot}} = 2088.0 \pm 1.2 \text{ MeV}$$

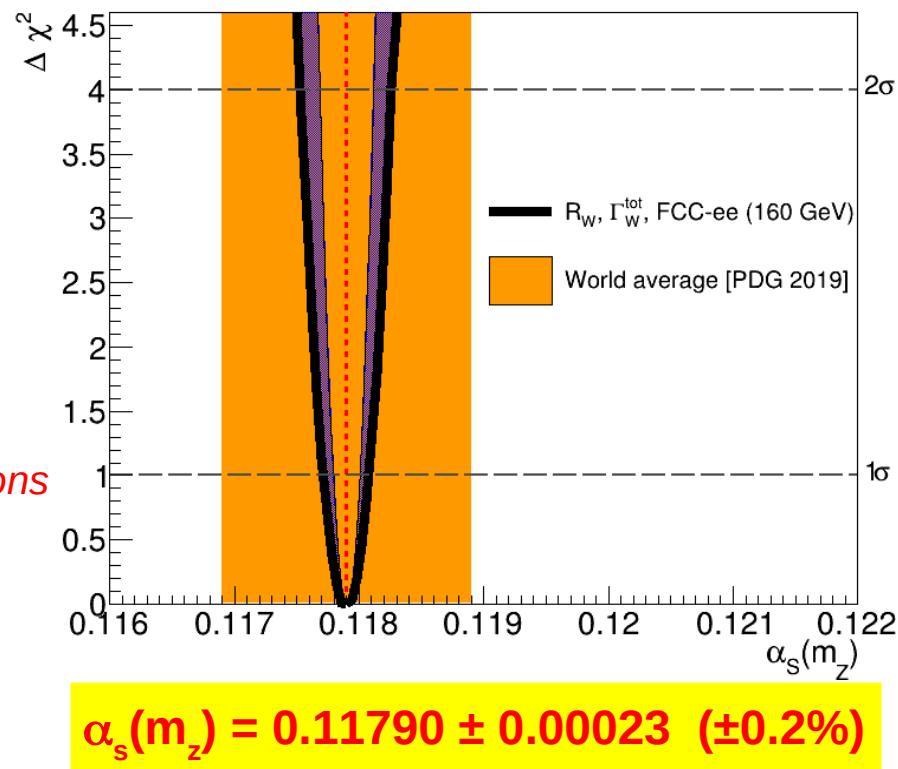
$$R_W = 2.08000 \pm 0.00008$$

$$m_W = 80.3800 \pm 0.0005 \text{ GeV}$$

$$|V_{cs}| = 0.97359 \pm 0.00010 \leftarrow O(10^{12}) D \text{ mesons}$$

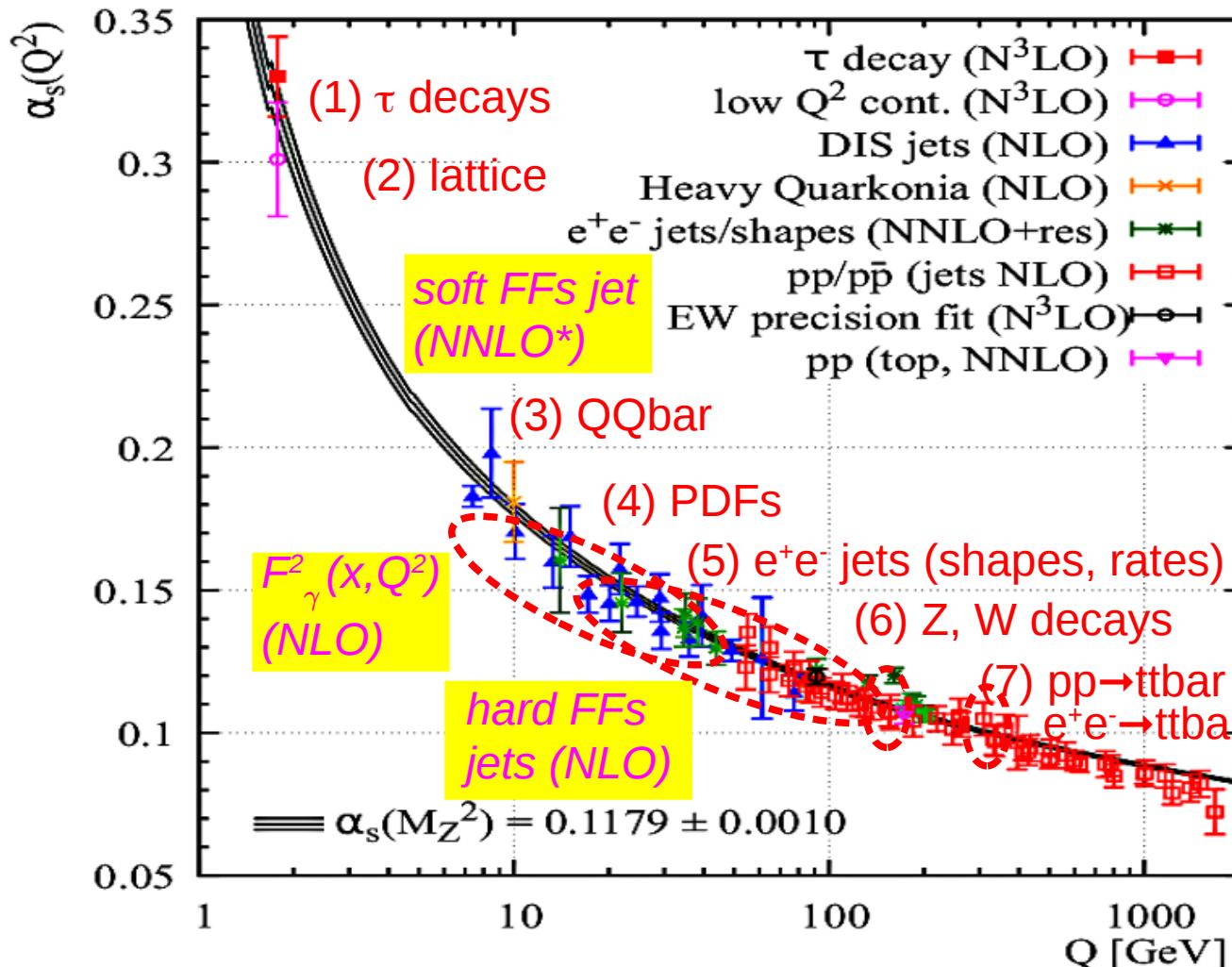
- TH uncertainty reduced by  $\times 10$  after computing missing  $\alpha_s^5$ ,  $\alpha^2$ ,  $\alpha^3$ ,  $\alpha\alpha_s^2$ ,  $\alpha\alpha_s^2$ ,  $\alpha^2\alpha_s$  terms

DdE, Jacobsen: arXiv:2005.04545 [hep-ph]



# Other $\alpha_s$ extractions (not yet in world average)

- There are few other classes of  $e^+e^-$  observables, computed today at lower accuracy (NLO, NNLO\*), that can be used to extract the QCD coupling:



# $\alpha_s$ from photon QCD structure function (NLO)

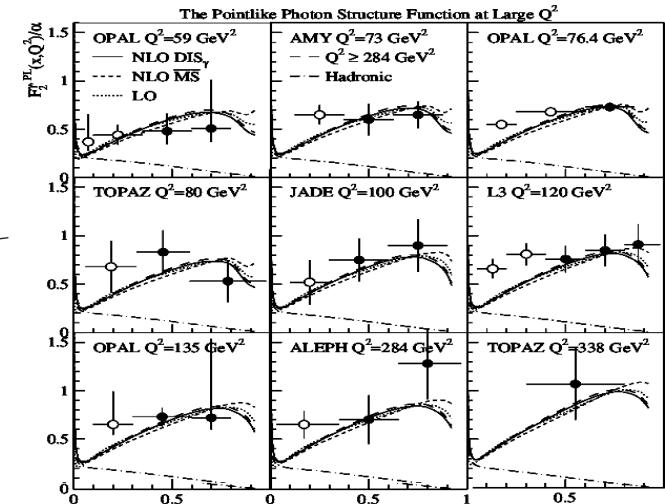
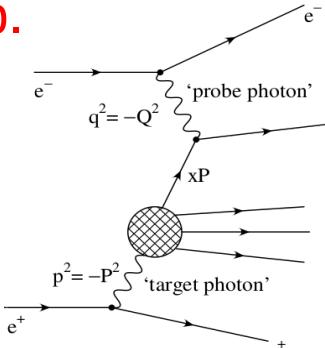
→ Computed at NNLO:  $\int_0^1 dx F_2^\gamma(x, Q^2, P^2) = \frac{\alpha}{4\pi} \frac{1}{2\beta_0} \left\{ \frac{4\pi}{\alpha_s(Q^2)} c_{LO} + c_{NLO} + \frac{\alpha_s(Q^2)}{4\pi} c_{NNLO} + \mathcal{O}(\alpha_s^2) \right\}$

→ Poor  $F_\gamma^2(x, Q^2)$  experimental measurements:

→ Extraction (NLO) with large exp. uncertainties today:

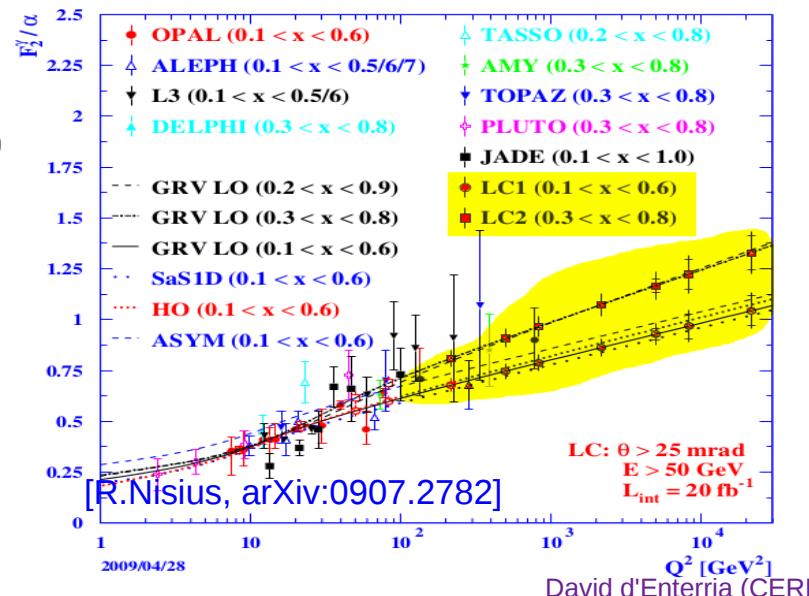
$$\alpha_s(m_z) = 0.1198 \pm 0.0054 \\ (\pm 4.5\%)$$

[M.Klasen et al. PRL89 (2002)122004]



→ Future prospects:

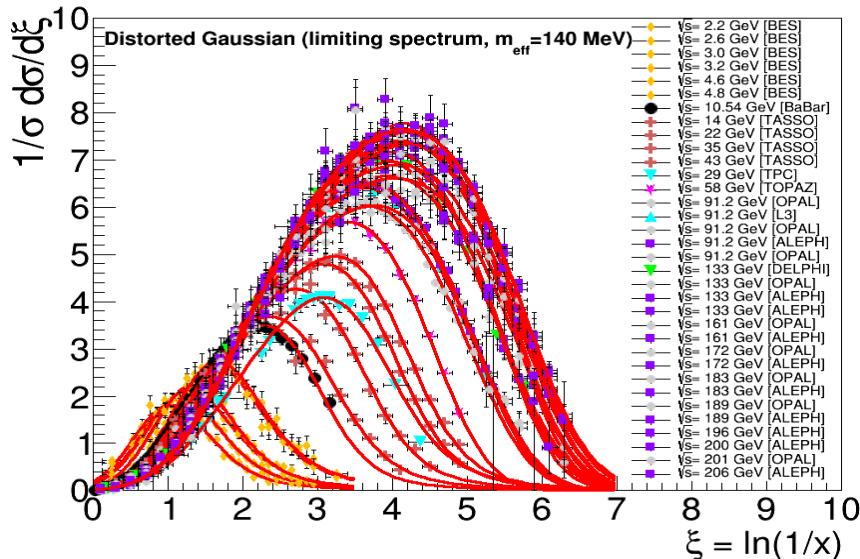
- Fit with NNLO  $F_\gamma^2$  evolution (ongoing)
- Better data badly needed: Belle-II ?
- Dedicated simul. studies at ILC exist:
- Huge  $\gamma\gamma$  (EPA) stats at FCC-ee will lead to:  $\delta\alpha_s/\alpha_s < 1\%$



# $\alpha_s$ extractions from jet fragmentation (NLO,NNLO\*)

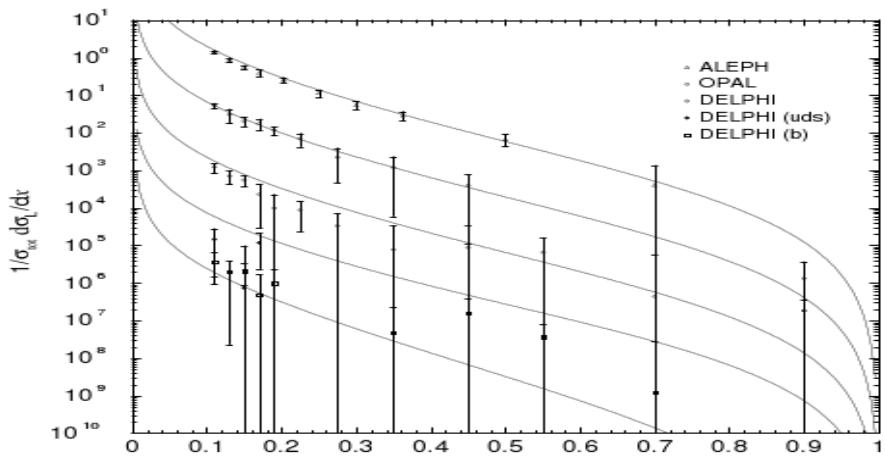
→ Soft parton-to-hadron FFs (NNLO\*+NNLL):

[D.d'E.,R.Perez-Ramos, arXiv:1505.02624 ]



→ Hard parton-to-hadron FFs (NLO):

$$\alpha_s(m_z) = 0.1176 \pm 0.0055 (\pm 4.7\%)$$



[AKK, B. Kniehl et al., NPB 803(2008)42]

Combined fit of the jet-energy evolution of the FF moments (multiplicity, peak, width,...)

with  $\alpha_s$  as single free parameter:

$$\alpha_s(m_z) = 0.1205 \pm 0.0022 (\pm 2\%)$$

(full-NNLO corrections missing)

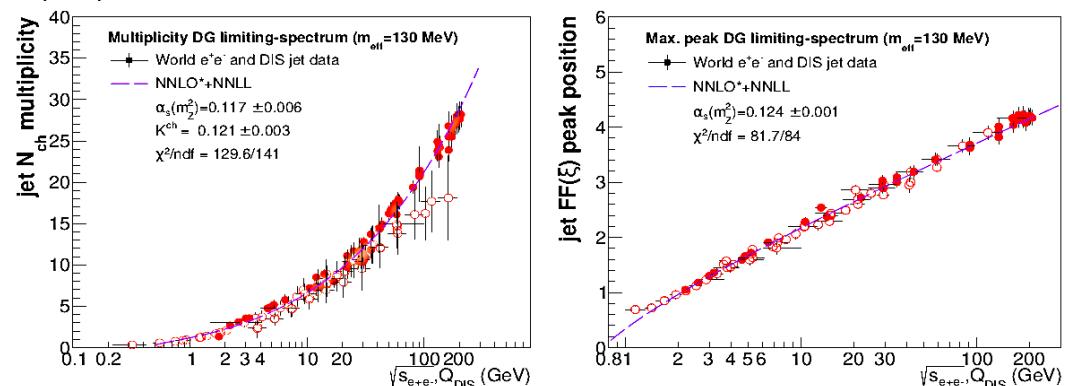


Figure 3: Energy evolution of the charged-hadron multiplicity (left) and of the FF peak position (right) measured in  $e^+e^-$  and DIS data fitted to the NNLO\*+NNLL predictions. The obtained  $\mathcal{K}_{\text{ch}}$  normalization constant, individual NNLO\*  $\alpha_s(m_z)$  values, and the goodness-of-fit per degree-of-freedom  $\chi^2/\text{ndf}$ .

# $\alpha_s$ at future $e^+e^-$ colliders (summary)

- World-average QCD coupling at  $N^{2,3}\text{LO}$  today:
  - Determined from *7 observables* with combined *0.85% uncertainty* (least well-known gauge coupling).
  - Impacts *all LHC QCD x-sections & decays*.
  - Role *beyond SM*: GUT, EWK vacuum stability, New colored sectors?
- $e^+e^-$  extractions:
  - Hadronic tau decays:  $\pm 1\%$  TH
  - Event shapes, jet rates:  $\pm 1\%$  TH
  - Z&W pseudo-observ.:  $\pm 0.1\%$  TH
- State-of-the-art Z, W extractions:
  - Z boson: New fit with high-order EW corrections + updated LEP data:  $\sim 2.3\%$  (exp.) uncertainty today.
  - W boson: New  $N^3\text{LO}$  fit to  $\Gamma_W$ ,  $R_W$   $\sim 27\%$  (exp.) uncertainty today.
- Permil uncertainty only possible with a machine like FCC-e $^+e^-$ 

$\alpha_s(m_z) = 0.12030 \pm 0.00028$   
 $\pm 0.2\% \text{ (tot)}, \pm 0.1\% \text{ (exp)}$

$\alpha_s(m_z) = 0.11790 \pm 0.00023$   
 $\pm 0.2\% \text{ (tot)}, \pm 0.1\% \text{ (exp)}$

What are the detector design improvements needed to bring propagated syst. uncert. on W,Z pseudo-observ. below 0.1% ?

# **QCD physics at future $e^+e^-$ machines**

**(1) QCD coupling**

**(2) Jet substructure & flavour tagging**

**(3) Parton shower & Non-perturbative QCD**

NOTE: Only UNIQUE QCD measurements, inaccessible at any current machine, are covered.

# Jet substructure & flavour tagging

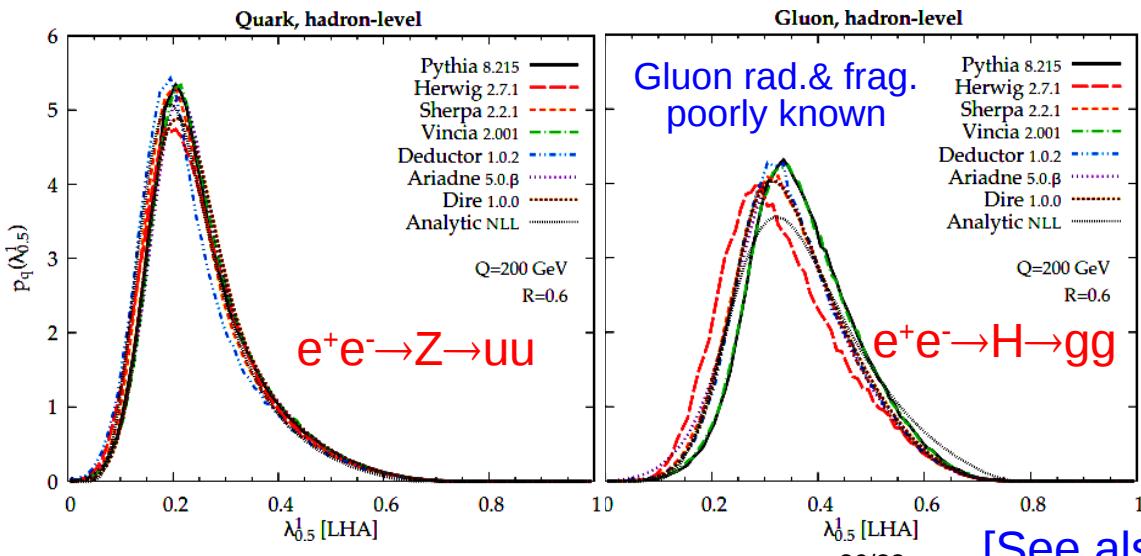
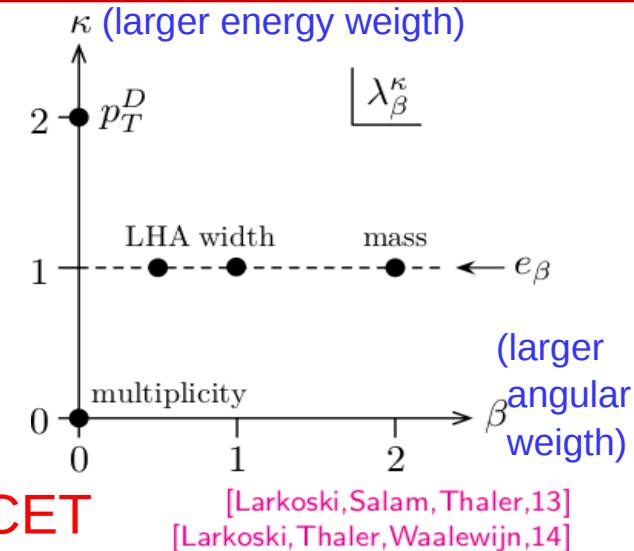
- State-of-the-art jet substructure studies based on **angularities**

$$\lambda_{\beta}^{\kappa} = \sum_{i \in \text{jet}} z_i^{\kappa} \theta_i^{\beta}, \quad (\text{normalized } E^n \times \theta^n \text{ products})$$

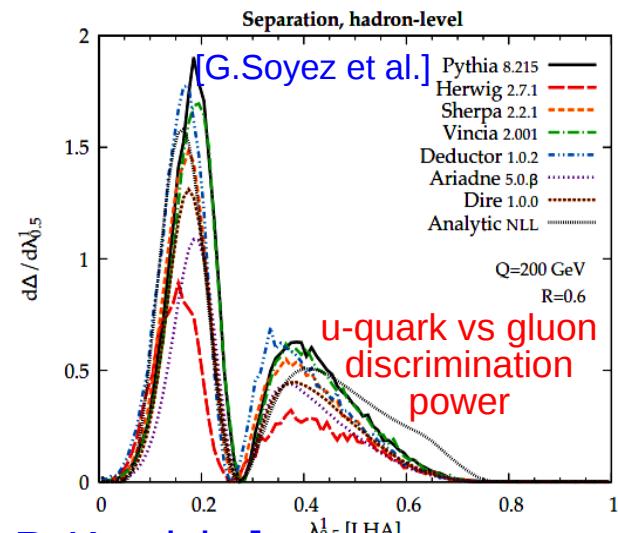
- "Sudakov"-safe variables of **jet constituents**: multiplicity, LHA, width/broadening, mass/thrust, C-parameter,...

- $k=1$ : IRC-safe computable ( $N^n\text{LO}+N^n\text{LL}$ ) via SCET  
(but uncertainties from non-pQCD effects)

- MC parton showers differ on gluon (less so quark) radiation patterns:

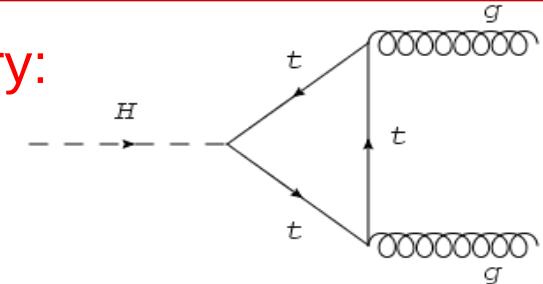


[See also P. Komiske]



# High-precision gluon & quark jet studies

- Exploit FCC-ee  $H(gg)$  as a "pure gluon" factory:  
 $H \rightarrow gg$  (BR~8% accurately known) provides  
120.000 extra-clean digluon events.

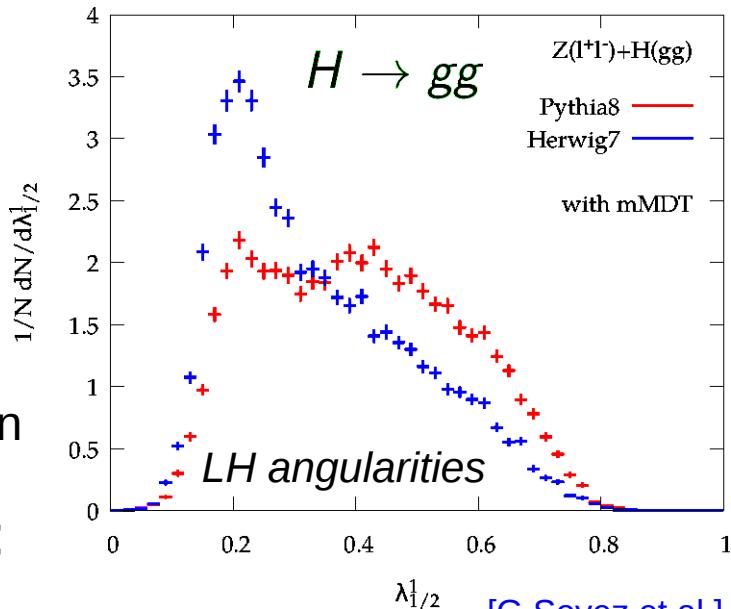


- Multiple handles to study gluon radiation & g-jet properties:

- ▶ Gluon vs. quark via  $H \rightarrow gg$  vs.  $Z \rightarrow qq$   
(Profit from excellent g,b separation)
- ▶ Gluon vs. quark via  $Z \rightarrow bbg$  vs.  $Z \rightarrow qq(g)$   
(g in one hemisphere recoiling  
against 2-b-jets in the other).
- ▶ Vary  $E_{jet}$  range via ISR:  $e^+e^- \rightarrow Z^*, \gamma^* \rightarrow jj(\gamma)$
- ▶ Vary jet radius: small-R down to calo resolution

- Multiple high-precision analyses at hand:

- BSM: Improve  $q/g/Q$  discrimination tools
- pQCD: Check  $N^nLO$  antenna functions. High-precision QCD coupling.
- non-pQCD: Gluon fragmentation: Octet neutralization? (zero-charge gluon jet with rap gaps). Colour reconnection? Glueballs ? Leading  $\eta$ 's,baryons?



[See also P. Komiske]

David d'Enterria (CERN)

# **QCD physics at future $e^+e^-$ machines**

**(1) QCD coupling**

**(2) Jet substructure & flavour tagging**

**(3) Parton shower & Non-perturbative QCD**

NOTE: Only UNIQUE QCD measurements, inaccessible at any current machine, are covered.

# QCD uncertainties on EWK observables

- With  $\times 10^5$  more Z's than LEP, EWK uncertainties at FCC-ee will be dominated by syst. (QCD).

Example:  $e^+e^- \rightarrow b\bar{b}$  forward–backward asymmetry

- 8 measurements at LEP:
  - 4 lepton-based, 4 jet-charge-based
  - Exp. observable with **largest discrepancy today wrt. the SM:  $2.8\sigma$**

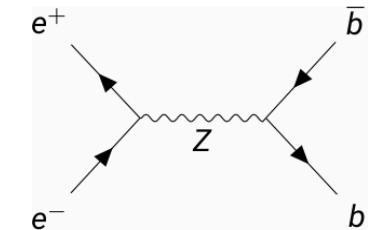
- Exp. Uncertainties:  $\sim 1.6\%$

- Statistical:  $\pm 1.5\%$  ( $\sim 0.05\%$  at FCC-ee)
- Systematics:  $\pm 0.6\%$  (QCD-related:  $\pm 0.4\%$ )

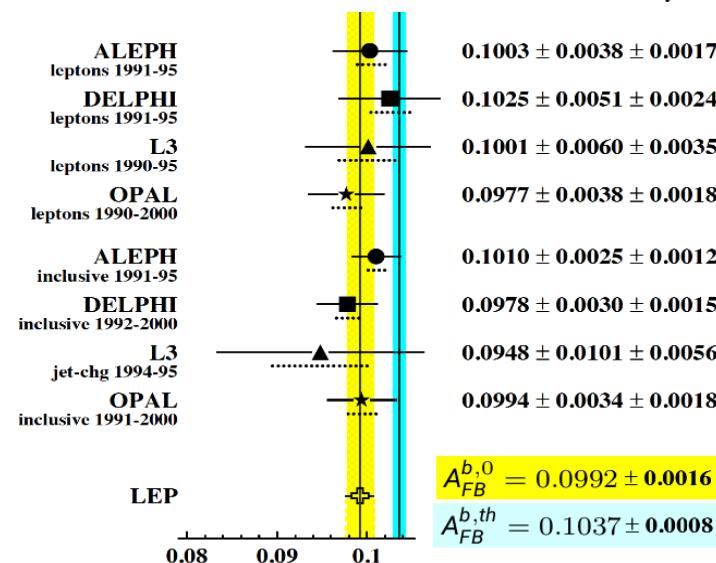
- QCD effects on  $A_{FB}^{0,b}$  (depending strongly on exp. selection procedure):
  - Gluon splitting (TH control:  $\alpha_s^2$  corrections)
  - Smearing of b-jet/thrust axis
  - b and c radiation & fragmentation. B and D decay models.

[Uncertainties estimated by Abbaneo et al., EPJC 4 (1998)]

- We have **revisited** the impact of QCD effects on  $A_{FB}^{0,b}$  implementing original analyses in up-to-date retuned parton-shower+hadronization MCs

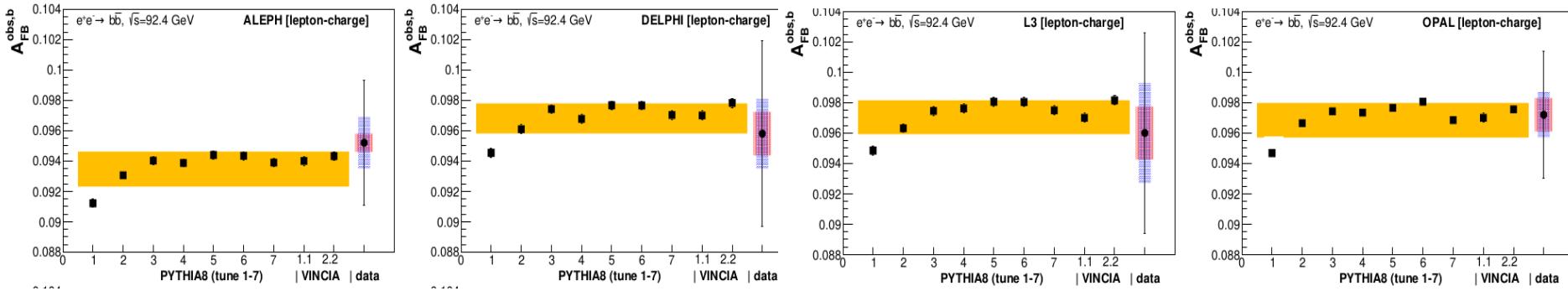


$$A_{FB}^b = \frac{N_F - N_B}{N_F + N_B} \quad A_{FB} = \frac{\sigma_A}{\sigma_S} \propto \frac{-g_{\mu\nu}T^{\mu\nu}}{ie_{\mu\nu\rho}n^\lambda Q^\rho n\cdot Q}$$

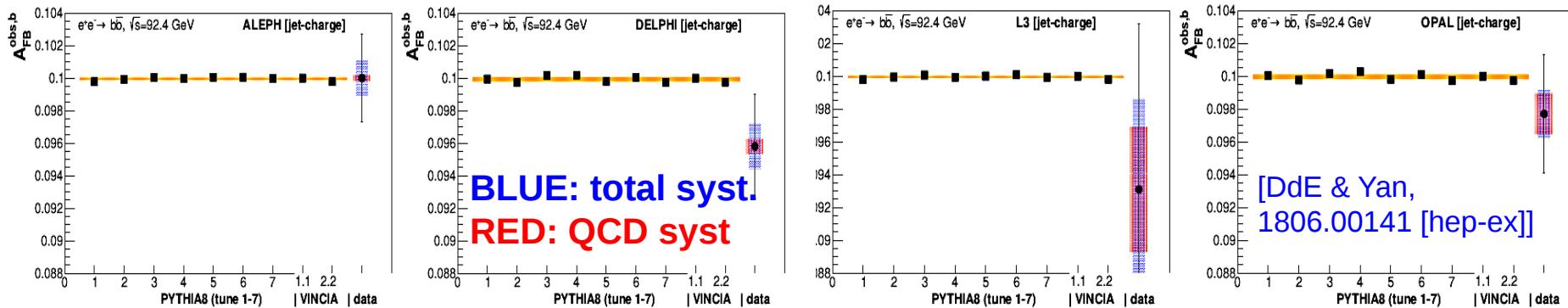


# Reduced QCD uncertainties on $A_{FB}$ at Z pole

- QCD uncertainties recomputed from PYTHIA8.226 (7 tunes) & VINCIA2.2
- $e^+e^- \rightarrow b\bar{b}$  forward–backward asymmetry for lepton-based analyses:



- $e^+e^- \rightarrow b\bar{b}$  forward–backward asymmetry for jet-charge-based analyses:

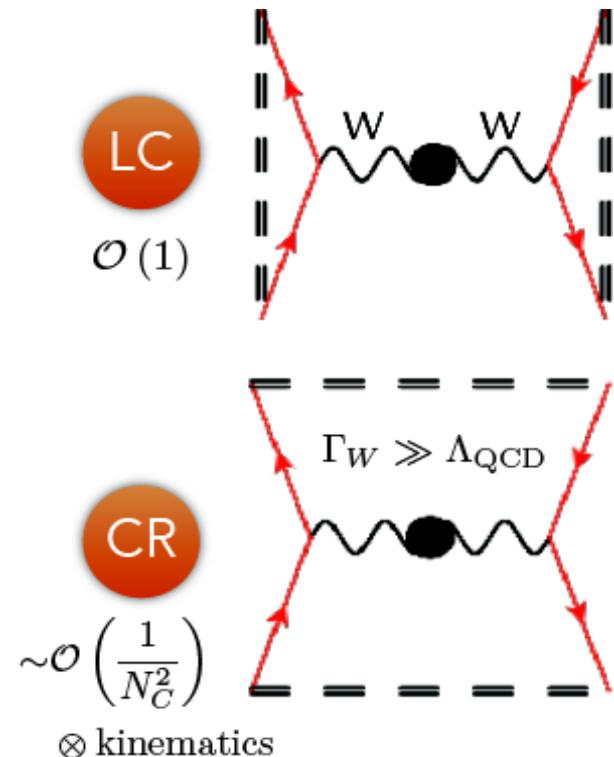


- 2020 vs. 1998 parton shower+hadronization uncertainties:
  - Lepton-based: Consistent for ALEPH, slightly smaller for DELPHI, L3, OPAL.
  - Jet-charge-based: Much smaller for all experiments.
- Improved PS & non-pQCD tunes w/  $e^+e^-$  data needed to reduce syst. uncert.

# Colour reconnection

- Colour reconnection among partons is source of uncertainty in  $m_W$ ,  $m_{top}$ , aGC extractions in multijet final-states. Especially in pp (MPI cross-talk).

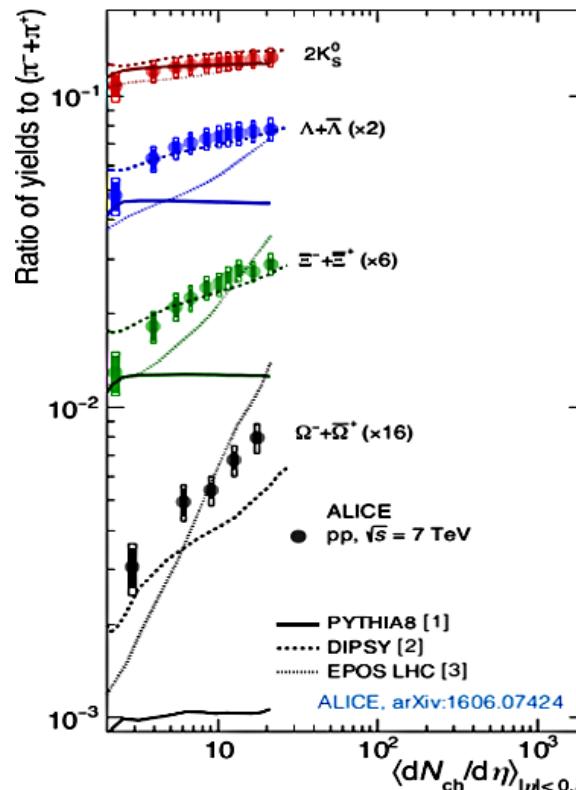
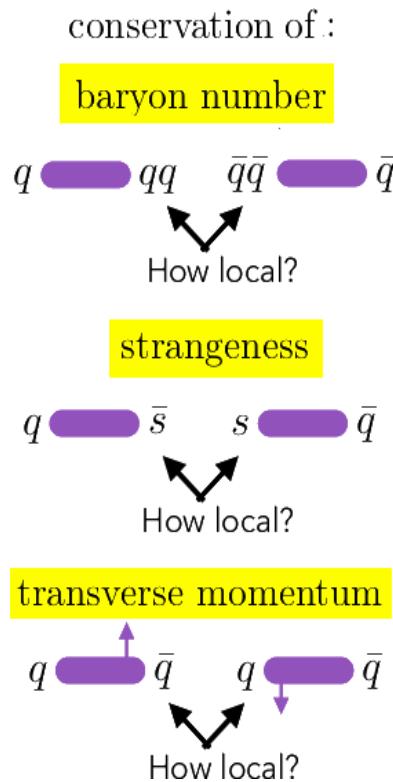
- ▶ CR impacts all FCC-ee multi-jet final-states (potentially shifted angular correlations):  
 $e^+e^- \rightarrow WW(4j)$ ,  $Z(4j)$ ,  $t\bar{t}\text{bar}$ ,  
 $H(2j,4j)$  CP studies,...  
String-drag effect on W mass (hinted at LEP).
- ▶ Exploit huge W stats ( $\times 10^4$  LEP) to measure  $m_W$  leptonically & hadronically and constrain CR in hadronic WW.



[See also P. Skands slides]

# Detailed hadronization studies

- High-precision low- $p_T$  PID hadrons in  $e^+e^-$  required for detailed studies:
  - Baryon & strangeness production. Colour string dynamics.
  - Final-state correlations (spin: BoseEinstein, FermiDirac; momenta; space)
  - Bound state formation: Onia, multi-quark states, glueballs, ...

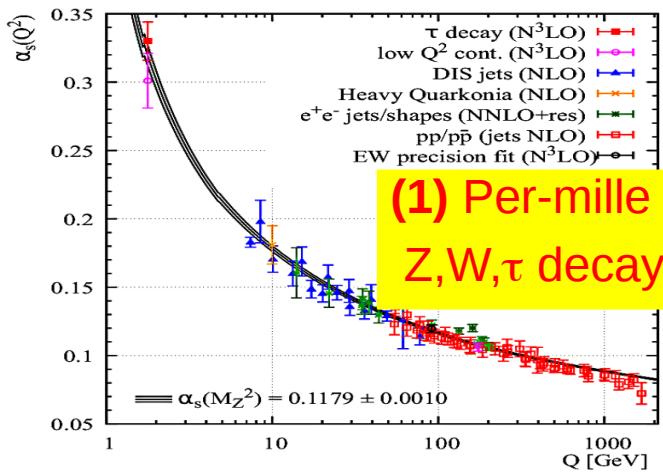


- ▶ Understand breakdown of universality of parton hadronization observed at LHC.
- Baseline vacuum  $e^+e^-$  studies for high-density QCD in small & large systems.

[See also P. Skands slides]

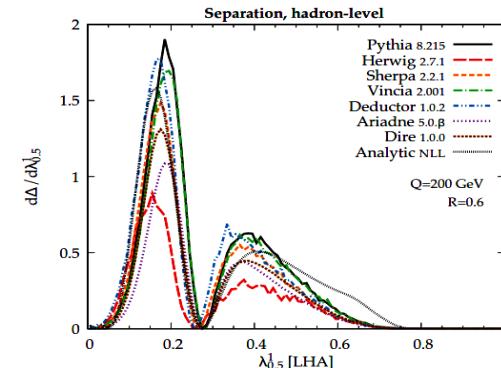
# Summary: QCD at future e<sup>+</sup>e<sup>-</sup> colliders

- The precision needed to fully exploit all future ee/pp/ep/eA/AA SM & BSM programs requires exquisite control of pQCD & non-pQCD physics.
- Unique QCD precision studies accessible at FCC-ee (CEPC, ILC):



(1) Per-mille  $\alpha_s$  via hadronic Z,W, $\tau$  decays, evt shapes...

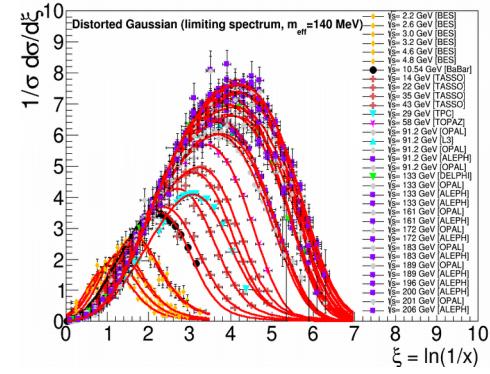
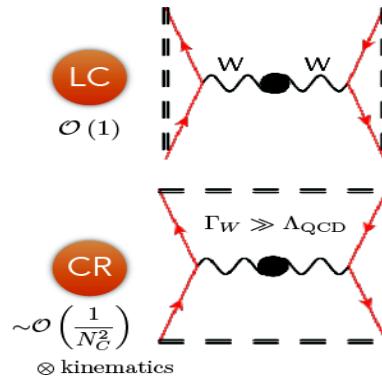
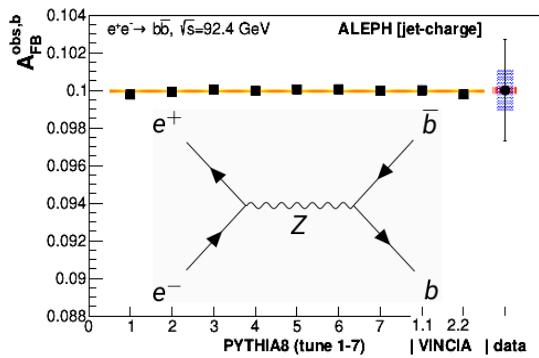
(2)  $N^n\text{LO}+N^n\text{LL}$  jet structure  
High g/q/Q discrimination



(3) Reduced PS+hadroniz. uncert. of EWK observ.

(4) <1% control of colour reconnection

(5) High-precision hadronization:



conservation of :  
baryon number

$q \text{ } \textcolor{purple}{qq}$     $\bar{q} \bar{q}$     $\textcolor{purple}{q} \bar{q}$

How local?

strangeness

$q \text{ } \textcolor{purple}{\bar{s}s}$     $s \text{ } \textcolor{purple}{\bar{q}q}$

How local?

transverse momentum

How local?

# Backup slides

# CERN FCC-ee project

- $e^+e^-$  operation before pp at  $\sqrt{s} = 90, (125), 160, 240, 350$  GeV

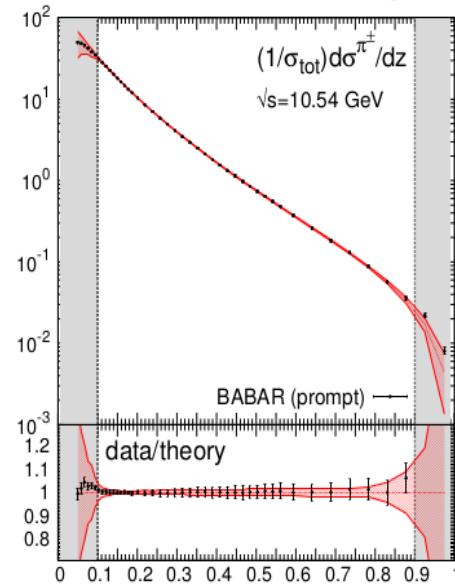
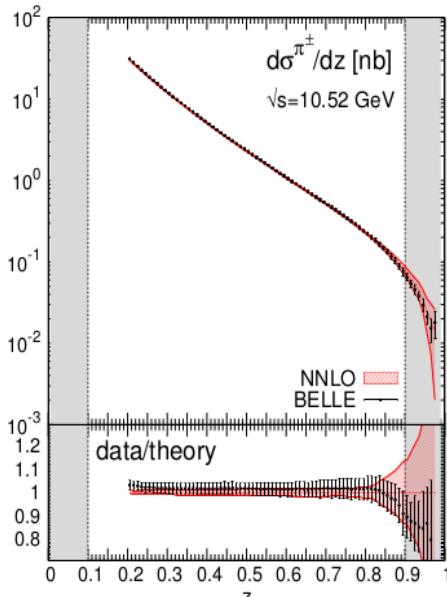


$\sqrt{s}$ (GeV):	90 (Z)	125 (eeH)	160 (WW)	240 (HZ)	350 ( $t\bar{t}$ )	350 (WW → H)
$\sigma$	43 nb	290 ab	4 pb	200 fb	0.5 pb	25 fb
$\mathcal{L}/IP$ ( $\text{cm}^{-2} \text{s}^{-1}$ )	$10^{36}$	$5 \cdot 10^{35}$	$10^{35}$	$7 \cdot 10^{34}$	$1.5 \cdot 10^{34}$	$1.5 \cdot 10^{34}$
$\mathcal{L}_{\text{int}}$ ( $\text{ab}^{-1}/\text{yr}$ , 2 IPs)	50	10	8	1.8	0.5	0.35
Events/year (2 IPs)	$10^{12}$	$3 \cdot 10^3$	$3 \cdot 10^7$	$3 \cdot 10^5$	$2.5 \cdot 10^5$	$10^4$
Years needed (2 IPs)	4	1.5	1	3	0.5	4
# of light-q jets/year:	$\mathcal{O}(10^{12})$	–	$\mathcal{O}(10^7)$	$\mathcal{O}(10^5)$	–	–
# of gluon-jets/year:	$\mathcal{O}(10^{11})$	$\mathcal{O}(10^2)$	$\mathcal{O}(10^6)$	$\mathcal{O}(10^4)$	–	$\mathcal{O}(10^3)$
# of heavy-Q jets/yr:	$\mathcal{O}(10^{12})$	$\mathcal{O}(10^3)$	$\mathcal{O}(10^7)$	$\mathcal{O}(10^5)$	$\mathcal{O}(10^5)$	$\mathcal{O}(10^4)$

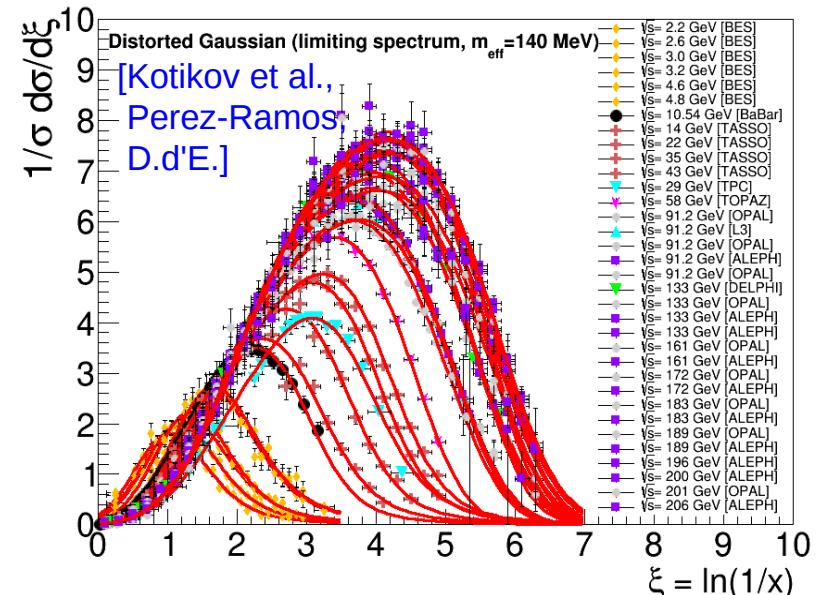
# High-precision parton FFs

■ Parton-to-hadron fragment. functions evolution known at NNLO at high-z &

[D.Anderle et al., A.Vossen et al., B.Kniehl et al.,  
V.Bertone et al., N.Sato et al., D.deFlorian et al.,...]



at NNLO\*+NNLL at low-z:



provide additional QCD coupling extractions:

Method	Current $\delta\alpha_s(m_z^2)/\alpha_s(m_z^2)$ uncertainty (theory & experiment state-of-the-art)	Future $\delta\alpha_s(m_z^2)/\alpha_s(m_z^2)$ uncertainty (theory & experiment progress)
soft FFs	$1.8\%_{\text{th}} \oplus 0.7\%_{\text{exp}} \approx 2\%$ (NNLO* only (+NNLL), npQCD small)	$0.7\%_{\text{th}} \oplus 0.7\%_{\text{exp}} \approx 1\%$ ( $\sim 2$ yrs), $< 1\%$ (FCC-ee) (NNLO+NNLL. More precise $e^+e^-$ data: 90–350 GeV)
hard FFs	$1\%_{\text{th}} \oplus 5\%_{\text{exp}} \approx 5\%$ (NLO only. LEP data only)	$0.7\%_{\text{th}} \oplus 2\%_{\text{exp}} \approx 2\%$ (+B-factories), $< 1\%$ (FCC-ee) (NNLO. More precise $e^+e^-$ data)

■ FCC-ee (much broader z range) allows for  $\alpha_s$  extraction with  $\delta\alpha_s < 1\%$