

HEFT vs SMEFT

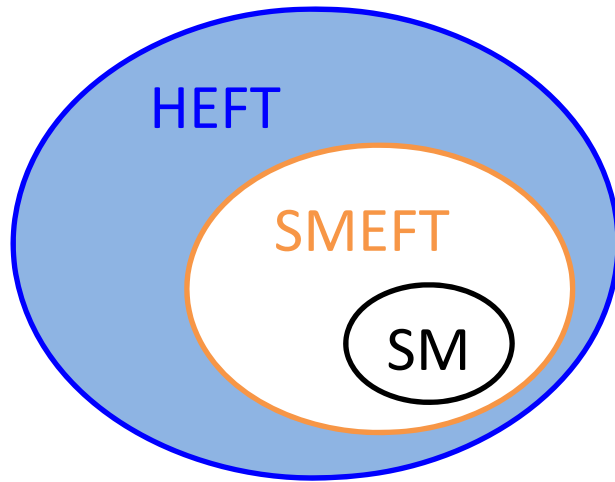
CEPC2020, Oct 26-28, 2020

Xiaochuan Lu

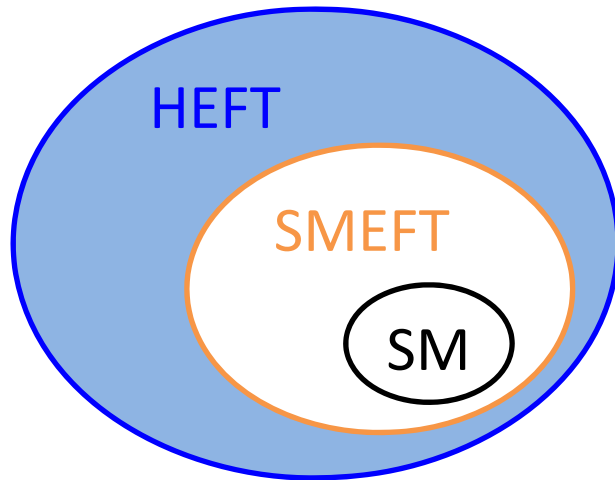
University of Oregon

arXiv: 2008.08597,

with Timothy Cohen, Nathaniel Craig, and Dave Sutherland

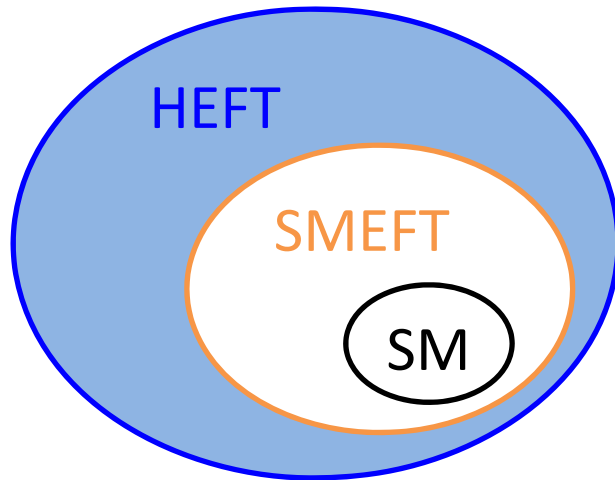


$$\mathcal{L}_{\text{SMEFT}} = \mathcal{L}_{\text{SM}} + \frac{C_H}{\Lambda^2} |H|^6 + \frac{C_{H\Box}}{\Lambda^2} |H|^2 \partial^2 |H|^2 + \frac{C_R}{\Lambda^2} |H|^2 |DH|^2 + \dots$$



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$$H \rightarrow \left\{ h, U \equiv \exp(i\pi^a t^a / v) \right\}$$

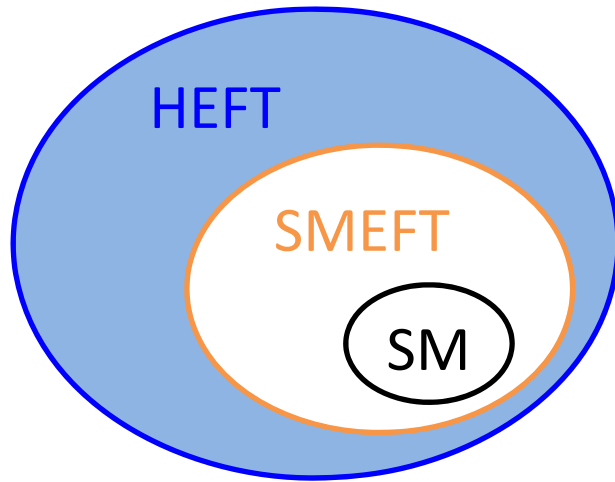


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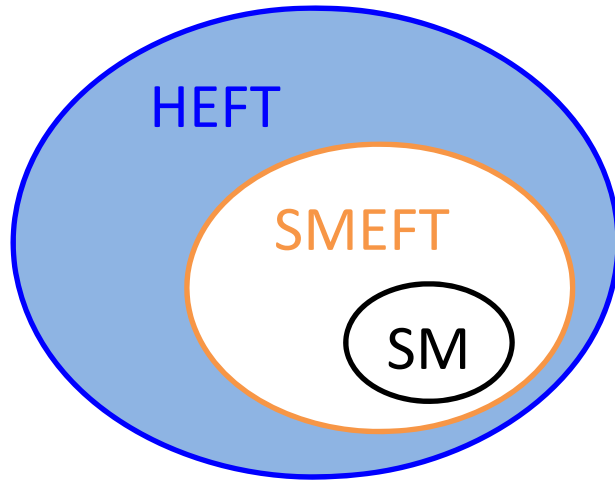
$$\begin{cases} \Sigma^\dagger \Sigma = |H|^2 \mathbf{1}_{2 \times 2} \\ \det(\Sigma) = |H|^2 = \frac{1}{2}(v+h)^2 \end{cases}$$



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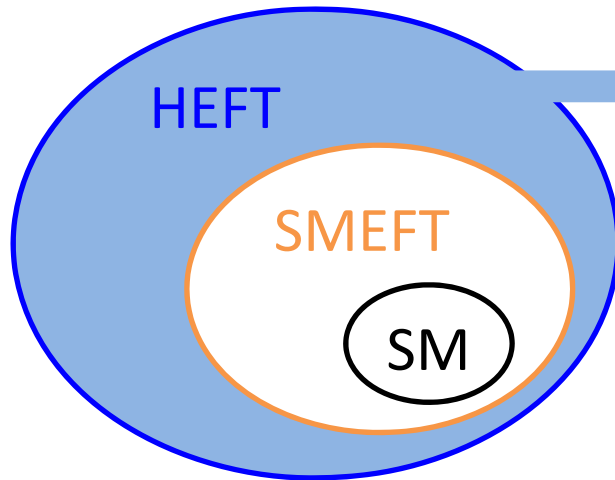


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$$\mathcal{L}_{\text{HEFT}} = \frac{1}{2} (\partial h)^2 - V(h) - \frac{v^2}{4} F(h) \text{tr} \left[(U^\dagger D_\mu U) (U^\dagger D^\mu U) \right] + \dots$$



Outline

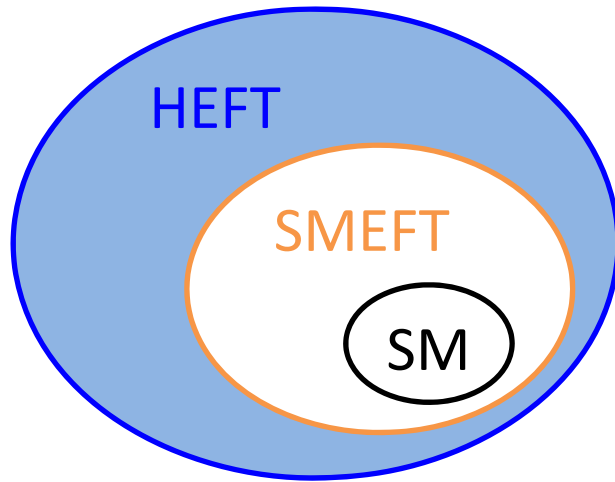
- What is HEFT?
- What UV physics generate HEFT?

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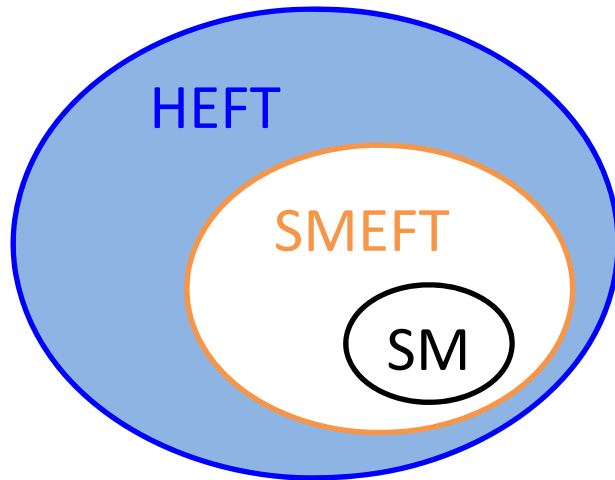


A confusion

$$\Sigma = \frac{1}{\sqrt{2}}(v+h)U \quad \Rightarrow \quad \begin{cases} v+h = \sqrt{2|H|^2} \\ U = \frac{\sqrt{2}}{v+h}\Sigma \end{cases}$$

$$\mathcal{L}_{\text{HEFT}}(h, U)$$

$$\mathcal{L}_{\text{SMEFT}}(H)$$



A confusion

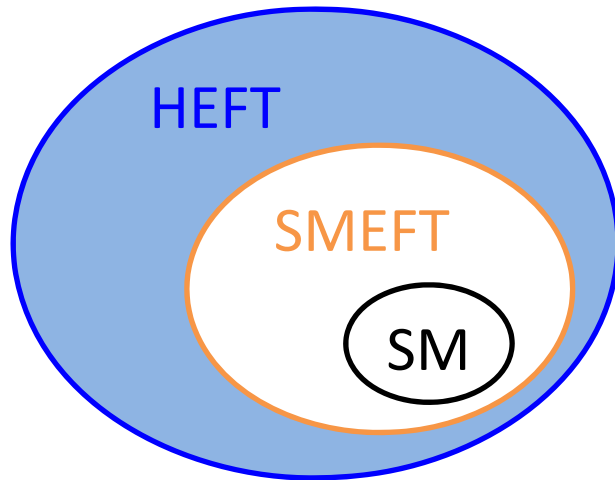
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$$\begin{cases} \text{SMEFT: } (x, y) \\ \text{HEFT: } (r, \theta) \end{cases} \Rightarrow r = \sqrt{x^2 + y^2}$$

Non-analytic



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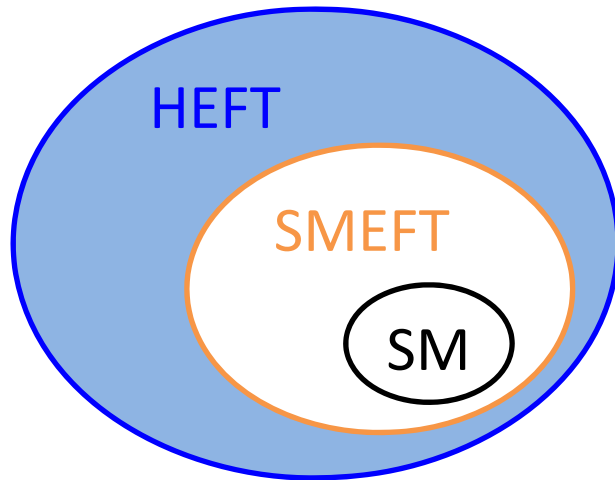
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$$vh + \frac{1}{2}h^2 = |H|^2 - \frac{1}{2}v^2$$

$$\begin{aligned} \sqrt{x} &= \sqrt{x_0 + (x - x_0)} \\ &= \sqrt{x_0} + \frac{1}{2\sqrt{x_0}}(x - x_0) + \dots \end{aligned}$$

All order

Adam Falkowski and Riccardo Rattazzi: (arXiv: 1902.05936)

UV cut-off. Our distinction between analytic and non-analytic lagrangians coincides with the distinction, in use in the Higgs EFT community, between linear (so-called SMEFT) and non-linear (so-called HEFT) effective theory, or equivalently between h being or not being part of a $SU(2)_W$ doublet. We however believe our classification is more adequate and enlightening from a physical point of view.

$$V(H) \supset \sqrt{2H^\dagger H} = v + h$$

$$(v + h)^{2k+1} = \left(\sqrt{2H^\dagger H} \right)^{2k+1}$$

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Field redefinition:
$$\begin{cases} v_1 \equiv \frac{3}{4}v \\ h_1 \equiv h + \frac{1}{4v}h^2 \end{cases}$$

$$V(H) \propto \frac{1}{4}(v+h)^2 + \frac{1}{4v}(v+h)^3 + \frac{1}{16v^2}(v+h)^4$$

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$$\uparrow = \left[\frac{1}{2}(v+h) + \frac{1}{4v}(v+h)^2 \right]^2$$

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$$\begin{aligned} &= \left[\frac{1}{2}(v+h) + \frac{1}{4v}(v+h)^2 \right]^2 \\ &= (v_1 + h_1)^2 = 2H_1^\dagger H_1 \end{aligned}$$

One can convert the SMEFT Lagrangian to HEFT form using Eq. (2.11) to switch from Cartesian and polar coordinates. One can attempt to convert from HEFT to SMEFT form using

$$\frac{\phi}{(\phi \cdot \phi)^{1/2}} = n \quad (2.30)$$

with $(\phi \cdot \phi)^{1/2}$ some function of h . This substitution gives a Lagrangian $L(\phi)$ that need not be analytic in ϕ . However, if there is an $O(4)$ fixed point, then there is a suitable change of variables such that the resulting Lagrangian is analytic in ϕ .

$$H = \frac{1}{\sqrt{2}} \begin{pmatrix} \phi^1 + i\phi^2 \\ \phi^3 + i\phi^4 \end{pmatrix}$$

$$\mathcal{L}_{\text{Kinetic}} = \frac{1}{2} g_{ij}(\phi) (\partial_\mu \phi^i) (\partial^\mu \phi^j)$$



$$ds^2 \equiv g_{ij}(\phi) d\phi^i d\phi^j$$

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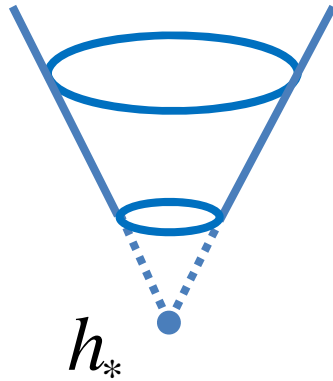
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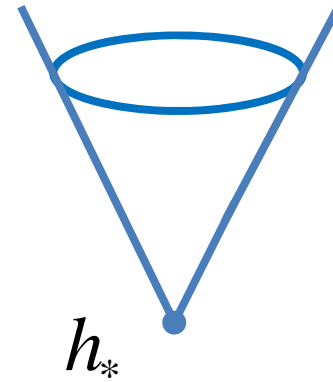
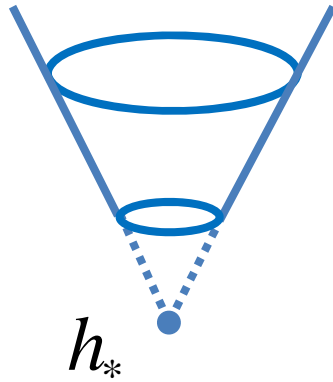
$$ds^2 \equiv g_{ij}(\phi) d\phi^i d\phi^j$$

$$\exists h_* \text{ such that } F(h_*) = 0$$



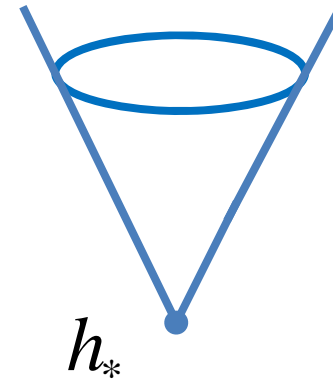
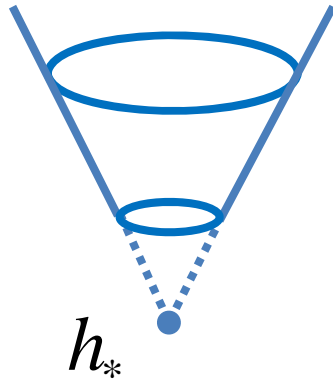
AJM (arXiv: 1605.03602)

$\exists h_*$ such that $F(h_*) = 0$



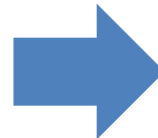
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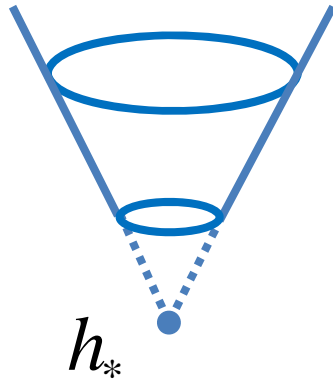
us (arXiv: 2008.08597)

$$F = 0$$

$$R, \nabla^2 R, \nabla^4 R, \dots < \infty$$

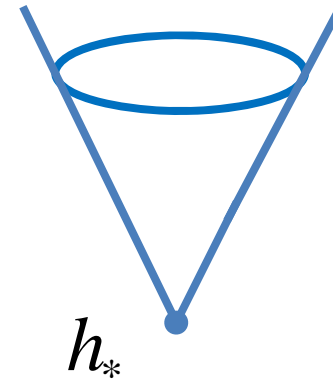
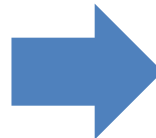
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UV theories that will generate HEFT?



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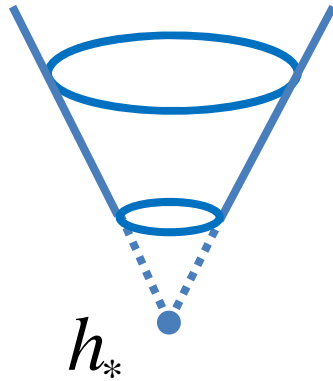
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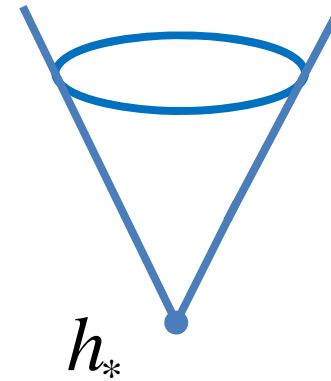
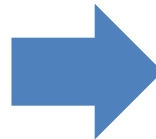
UV theories that will generate HEFT?

- Extra electroweak symmetry breaking



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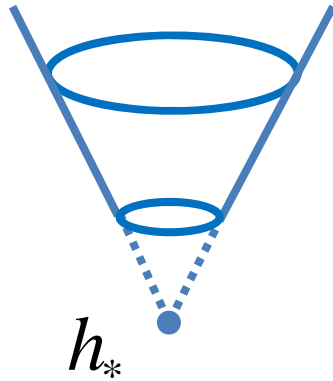
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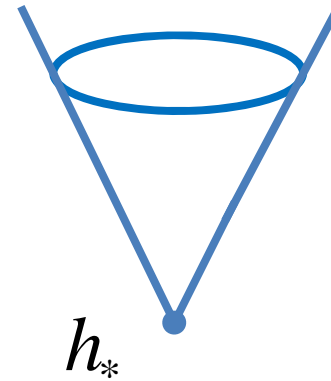
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AJM (arXiv: 1605.03602)

$\exists h_*$ such that $F(h_*) = 0$

- Mass fully from electroweak symmetry breaking



us (arXiv: 2008.08597)

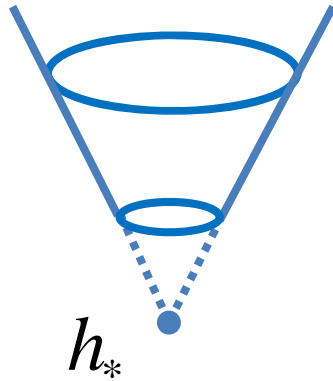
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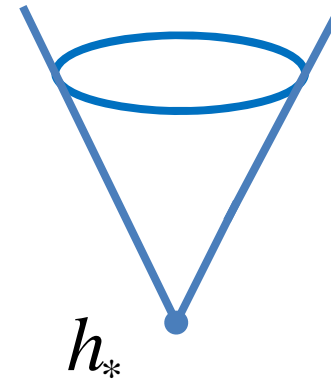
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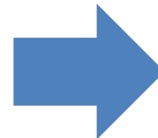
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Leading Order
Criterion

AJM (arXiv: 1605.03602)

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us (arXiv: 2008.08597)

$F = 0$

$R, \nabla^2 R, \nabla^4 R, \dots < \infty$

$V, \nabla^2 V, \nabla^4 V, \dots < \infty$

Matching to **all-order in fields**

Coleman-Weinberg: $\mathcal{L}_{\text{UV}}[\phi, \Phi] \supset -\frac{1}{2}\Phi[\partial^2 + M^2 + U(\phi)]\Phi$

$$\begin{aligned}\mathcal{L}_{\text{EFT}}[\phi] &= \frac{i}{2} \ln \det(\partial^2 + M^2 + U) = \frac{i}{2} \text{Tr} \ln(\partial^2 + M^2 + U) \\ &= \frac{i}{2} \int d^4x \int \frac{d^4p}{(2\pi)^4} \text{tr} \left[\ln(-p^2 + M^2 + U) + \frac{1}{6(p^2 - M^2 - U)^3} (\partial U)^2 \right] \\ &= \frac{1}{2} \int d^4x \frac{1}{16\pi^2} \text{tr} \left[\frac{1}{2} (M^2 + U)^2 \left(\ln \frac{\mu^2}{M^2 + U} + \frac{3}{2} \right) + \frac{1}{12} \frac{1}{M^2 + U} (\partial U)^2 + \dots \right]\end{aligned}$$

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If $[U, \partial_\mu U] \neq 0$: see App. D in 2008.08597

Example: A heavy Singlet

$$\mathcal{L}_{\text{UV}} = \mathcal{L}_{\text{SM}} - \frac{1}{2} S \left(\partial^2 + M^2 + \kappa |H|^2 \right) S \quad , \quad m_S^2 = M^2 + \frac{1}{2} \kappa v^2$$

$M^2 > 0$, SMEFT $M^2 = 0$, HEFT

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$$vF(h) = v + h \quad , \quad K^2 \equiv 1 + \frac{1}{96\pi^2} \frac{\kappa^2 |H|^2}{M^2 + \kappa |H|^2}$$

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$$R = \frac{1}{16\pi^2} \left[\frac{1}{2K^4} \frac{\kappa^2 M^2}{\left(M^2 + \kappa |H|^2 \right)^2} + \frac{1}{2K^2} \frac{\kappa^2}{M^2 + \kappa |H|^2} \right]$$

Summary

- A geometric criterion to tell a HEFT Lagrangian from SMEFT
- An understanding of what UV theories would generate HEFT
- A functional method to match to all order in *fields* (not *derivatives* yet)
- A few UV examples show that the leading order criterion works