

# Lattice calculation of the hadronic light-by-light contribution to the muon magnetic moment

Thomas Blum (UConn / RBRC)

Norman Christ (Columbia)

**靳路昶 Luchang Jin** (UConn / RBRC)

Masashi Hayakawa (Nagoya)

Taku Izubuchi (BNL / RBRC)

Chulwoo Jung (BNL)

Christoph Lehner (Regensburg / BNL)

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<https://meeting.tencent.com/p/6733913824>

# The RBC & UKQCD collaborations

## [BNL and BNL/RBRC](#)

Yasumichi Aoki (KEK)

**Taku Izubuchi**

Yong-Chull Jang

**Chulwoo Jung**

Meifeng Lin

Aaron Meyer

Hiroshi Ohki

Shigemi Ohta (KEK)

Amarjit Soni

## [UC Boulder](#)

Oliver Witzel

## [CERN](#)

Mattia Bruno

## [Columbia University](#)

Ryan Abbot

**Norman Christ**

Duo Guo

Christopher Kelly

Bob Mawhinney

Masaaki Tomii

Jiqun Tu

Bigeng Wang

Tianle Wang

Yidi Zhao

## [University of Connecticut](#)

**Tom Blum**

Dan Hoying (BNL)

**Luchang Jin (RBRC)**

**Cheng Tu**

## [Edinburgh University](#)

Peter Boyle

Luigi Del Debbio

Felix Erben

Vera Gülpers

Tadeusz Janowski

Julia Kettle

Michael Marshall

Fionn Ó hÓgáin

Antonin Portelli

Tobias Tsang

Andrew Yong

Azusa Yamaguchi

[Masashi Hayakawa \(Nagoya\)](#)

## [KEK](#)

Julien Frison

## [University of Liverpool](#)

Nicolas Garron

## [MIT](#)

David Murphy

## [Peking University](#)

Xu Feng

## [University of Regensburg](#)

**Christoph Lehner (BNL)**

## [University of Southampton](#)

Nils Asmussen

Jonathan Flynn

Ryan Hill

Andreas Jüttner

James Richings

Chris Sachrajda

## [Stony Brook University](#)

Jun-Sik Yoo

Sergey Syritsyn (RBRC)

- **Introduction**
- Connected diagrams: exact photon propagator, the moment method
- Disconnected diagrams
- Results @  $L = 5.5\text{fm}$ ,  $1/a = 1.73\text{GeV}$
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# Muon $g - 2$ : experiments

$$\vec{\mu} = -g \frac{e}{2m} \vec{s}$$

$$a = \frac{g - 2}{2}$$

Authors	Lab	Muon Anomaly
Garwin et al. '60	CERN	0.001 13(14)
Charpak et al. '61	CERN	0.001 145(22)
Charpak et al. '62	CERN	0.001 162(5)
Farley et al. '66	CERN	0.001 165(3)
Bailey et al. '68	CERN	0.001 166 16(31)
Bailey et al. '79	CERN	0.001 165 923 0(84)
Brown et al. '00	BNL	0.001 165 919 1(59) ( $\mu^+$ )
Brown et al. '01	BNL	0.001 165 920 2(14)(6) ( $\mu^+$ )
Bennett et al. '02	BNL	0.001 165 920 4(7)(5) ( $\mu^+$ )
Bennett et al. '04	BNL	0.001 165 921 4(8)(3) ( $\mu^-$ )

World Average dominated by BNL

$$a_\mu = (11659208.9 \pm 6.3) \times 10^{-10}$$

In comparison, for electron

$$a_e = (11596521.8073 \pm 0.0028) \times 10^{-10}$$

# Muon $g - 2$ : Fermilab E989, J-PARC E34 3 / 43

SM (Model HLbL)  $11659182.2 \pm 3.8$

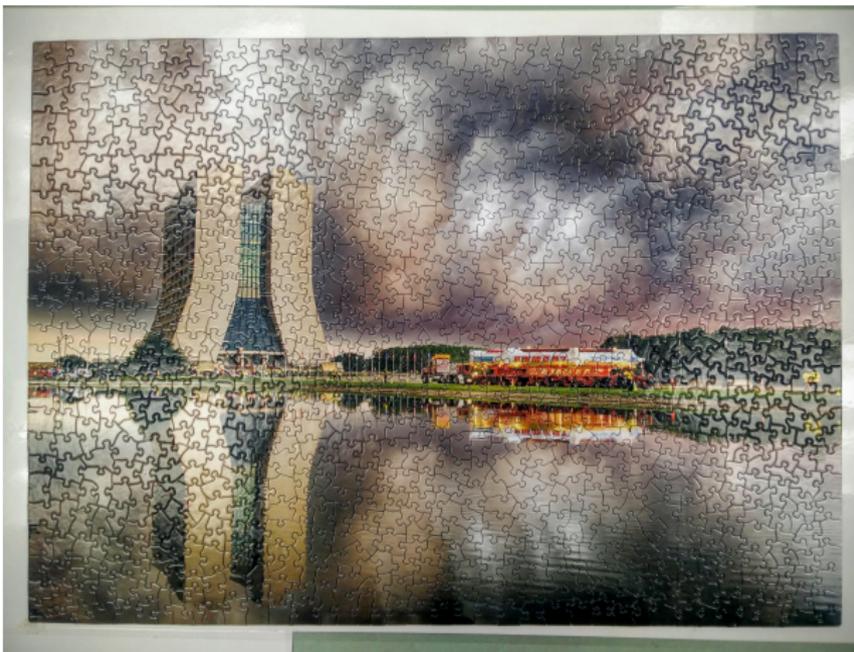
BNL E821 Exp  $11659208.9 \pm 6.3$

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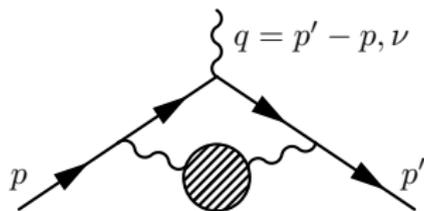
Diff (Exp - SM)  $26.7 \pm 7.4$

3.6 $\sigma$  deviations

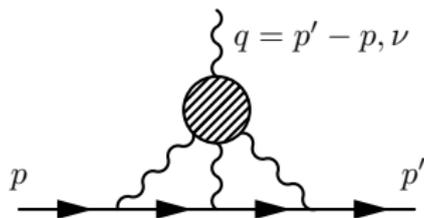
New Physics?



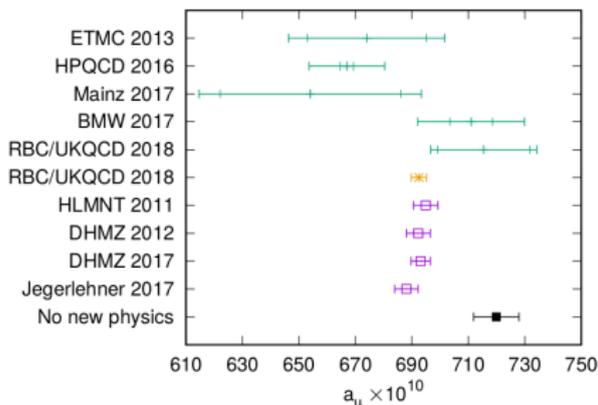
	$a_\mu$	$\times$	$10^{10}$	
QED incl. 5-loops	11658471.9	$\pm$	0.0	Aoyama, et al, 2012
Weak incl. 2-loops	15.4	$\pm$	0.1	Gnendiger et al, 2013
HVP	693.1	$\pm$	4.0	WP2020
HVP NLO&NNLO	-8.6	$\pm$	0.1	KNT2020
HLbL	9.0	$\pm$	1.7	WP2020
HLbL NLO	0.2	$\pm$	0.1	Colangelo, et al 2014
Standard Model	11659181.0	$\pm$	4.3	WP2020
Experiment	11659208.9	$\pm$	6.3	E821, The $g - 2$ Collab. 2006
Difference (Exp-SM)	27.9	$\pm$	7.6	



HVP: Hadronic Vacuum  
Polarization



HLbL: Hadronic Light by Light

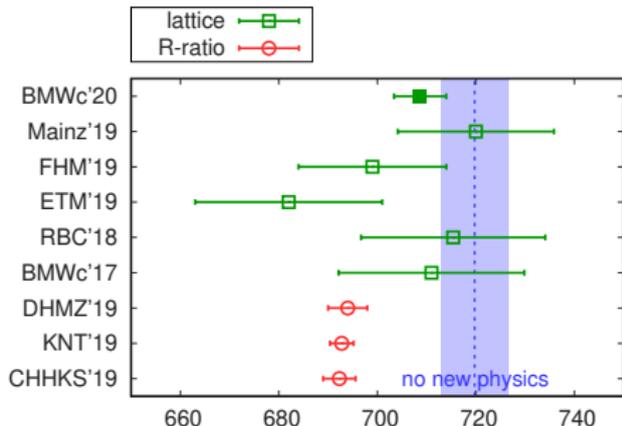


C. Lehner et al. 2018

RBC-UKQCD

(PRL 121, 022003)

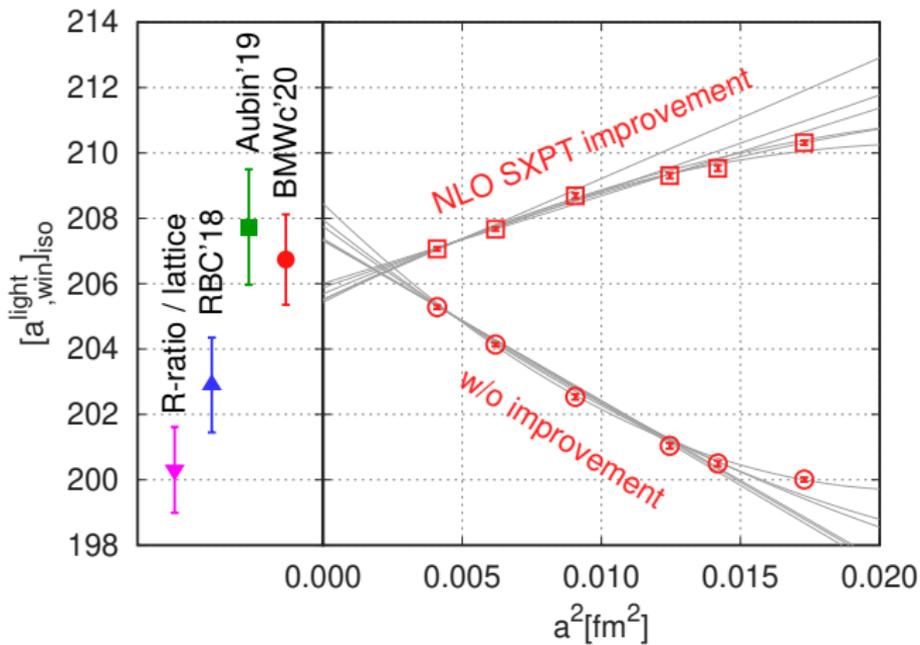
- Accuracy of lattice has caught up.
- BMW  $2.4\sigma$  tension with R-ratio.
- More results from different collaborations will appear.



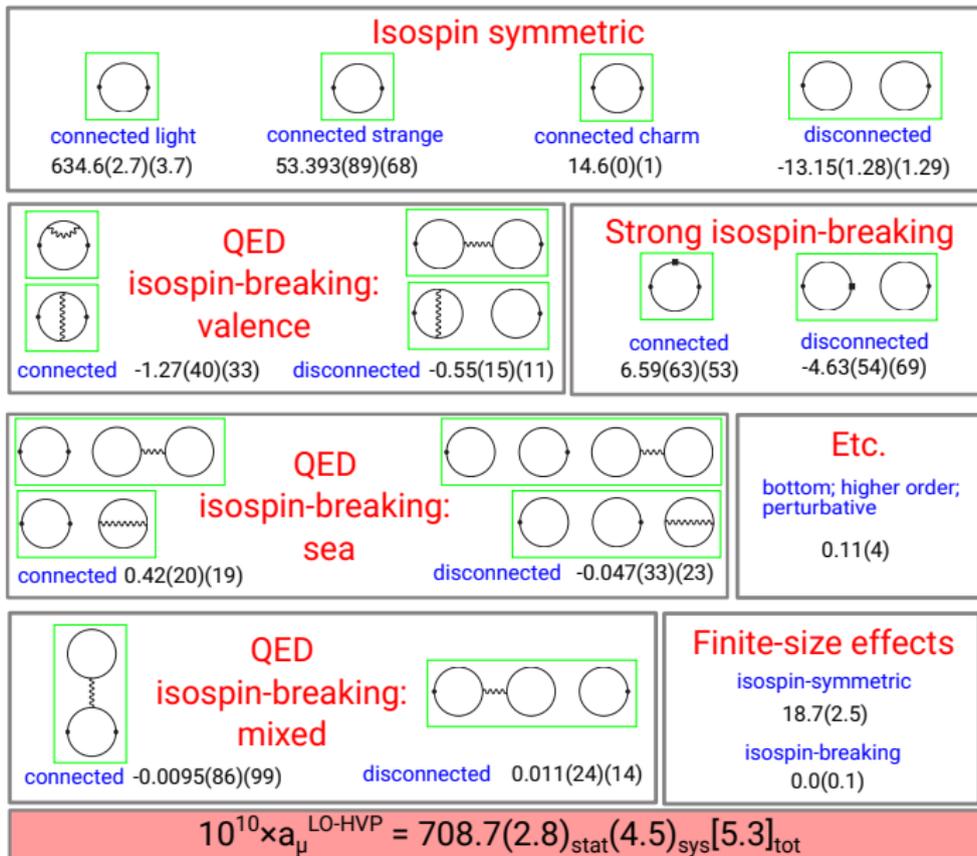
Sz. Borsanyi et al. 2020

BMW

(2002.12347)



- Light quark connected diagram contribution in a window (from 0.4 fm to 1 fm).

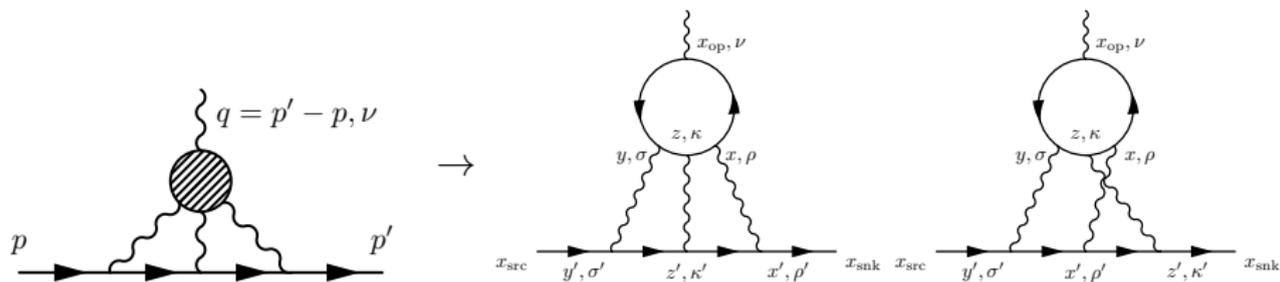


Contribution	PdRV(09) [471]	N/JN(09) [472, 573]	J(17) [27]	Our estimate
$\pi^0, \eta, \eta'$ -poles	114(13)	99(16)	95.45(12.40)	93.8(4.0)
$\pi, K$ -loops/boxes	-19(19)	-19(13)	-20(5)	-16.4(2)
$S$ -wave $\pi\pi$ rescattering	-7(7)	-7(2)	-5.98(1.20)	-8(1)
subtotal	88(24)	73(21)	69.5(13.4)	69.4(4.1)
scalars	-	-	-	} - 1(3)
tensors	-	-	1.1(1)	
axial vectors	15(10)	22(5)	7.55(2.71)	6(6)
$u, d, s$ -loops / short-distance	-	21(3)	20(4)	15(10)
$c$ -loop	2.3	-	2.3(2)	3(1)
total	105(26)	116(39)	100.4(28.2)	92(19)

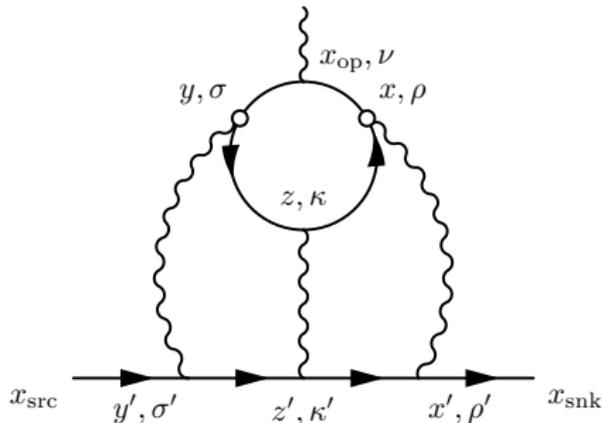
Table 15: Comparison of two frequently used compilations for HLbL in units of  $10^{-11}$  from 2009 and a recent update with our estimate. Legend: PdRV = Prades, de Rafael, Vainshtein (“Glasgow consensus”); N/JN = Nyffeler / Jegerlehner, Nyffeler; J = Jegerlehner.

- Values in the table is in unit of  $10^{-11}$ .
- We will use the unit in  $10^{-10}$  in the rest of the talk.
- The total HLbL contribution is on the order of  $10 \times 10^{-10}$ .

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- Gluons and sea quark loops (not directly connected to photons) are included automatically to all orders!
- There are additional four different permutations of photons not shown.
- The photons can be connected to different quark loops. These are referred to as the disconnected diagrams. They will be discussed later.
- First results are obtained by [T. Blum et al. 2015 \(PRL 114, 012001\)](#).



$$\frac{a_\mu}{m_\mu} \bar{u}_{s'}(\vec{0}) \frac{\Sigma}{2} u_s(\vec{0}) = \sum_{r=x-y} \sum_z \sum_{x_{\text{op}}} \frac{1}{2} (\vec{x}_{\text{op}} - \vec{x}_{\text{ref}}) \times \bar{u}_{s'}(\vec{0}) i \vec{\mathcal{F}}^C(\vec{0}; x, y, z, x_{\text{op}}) u_s(\vec{0})$$

$$\vec{\mu} = \sum_{\vec{x}_{\text{op}}} \frac{1}{2} (\vec{x}_{\text{op}} - \vec{x}_{\text{ref}}) \times \vec{J}(\vec{x}_{\text{op}})$$

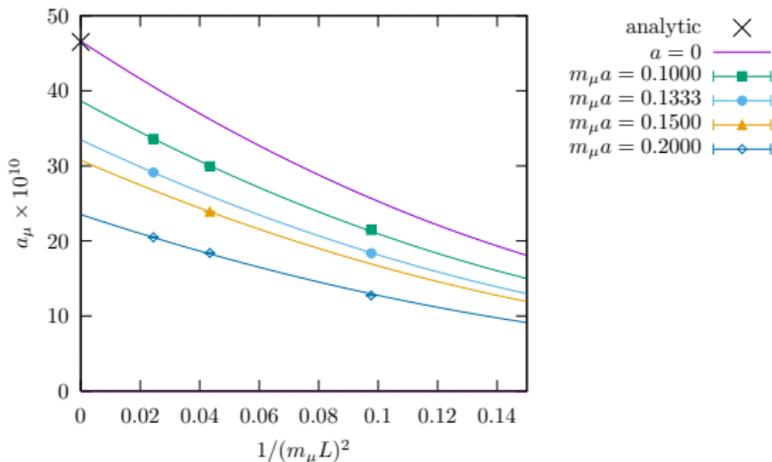
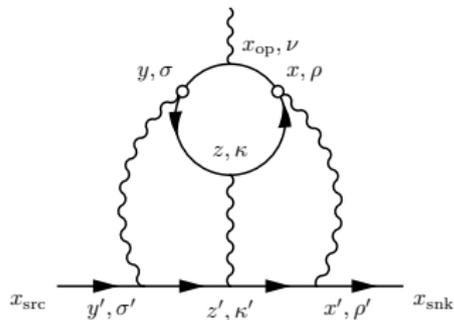
Reorder summation

(will discuss later).

- Two point sources at  $x, y$ : randomly sample  $x$  and  $y$ .
- Importance sampling: focus on small  $|x - y|$ .
- Complete sampling for  $|x - y| \leq 5a$  upto discrete symmetry.

- Muon is plane wave,  $x_{\text{ref}} = (x + y)/2$ .
- Sum over time component for  $x_{\text{op}}$ .
- Only sum over  $r = x - y$ .

- We test our setup by computing **muon leptonic light by light** contribution to muon  $g - 2$ .

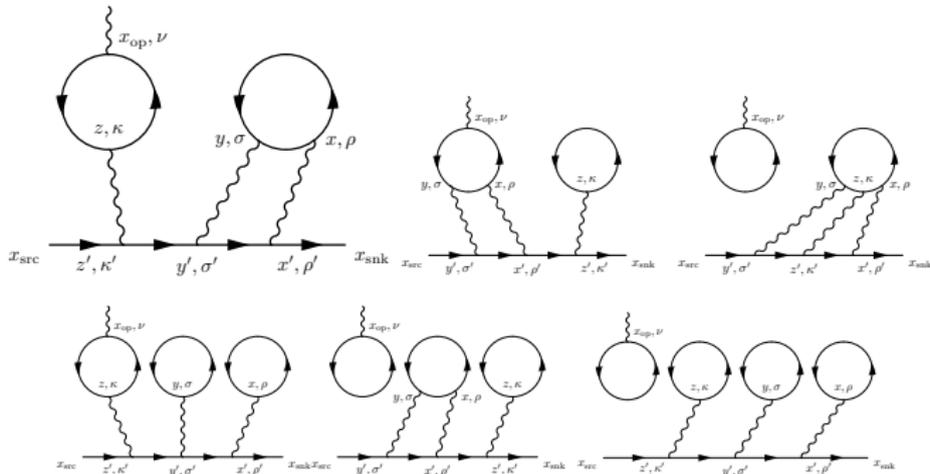


$$F_2(a, L) = F_2 \left( 1 - \frac{c_1}{(m_\mu L)^2} + \frac{c'_1}{(m_\mu L)^4} \right) (1 - c_2 a^2 + c'_2 a^4) \rightarrow F_2 = 46.6(2) \times 10^{-10} \quad (19)$$

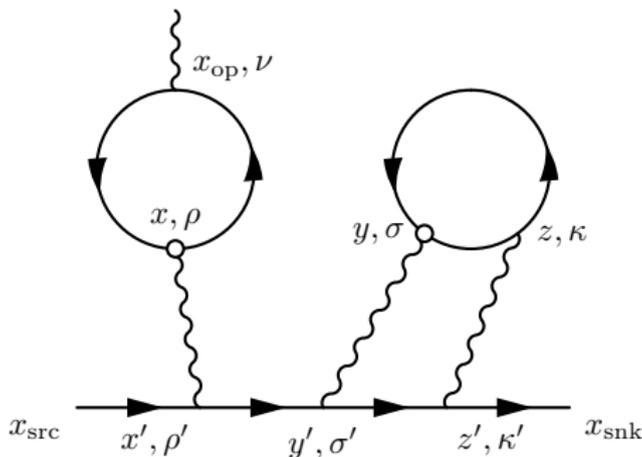
- Pure QED computation.** Muon leptonic light by light contribution to muon  $g - 2$ . Phys.Rev. D93 (2016) 1, 014503. arXiv:1510.07100.
- Analytic results:  $0.371 \times (\alpha/\pi)^3 = 46.5 \times 10^{-10}$ .
- $\mathcal{O}(1/L^2)$  finite volume effect, because the photons are emitted from a conserved loop.

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- One diagram (the biggest diagram below) do not vanish even in the  $SU(3)$  limit.
- We extend the method and computed this leading disconnected diagram as well.



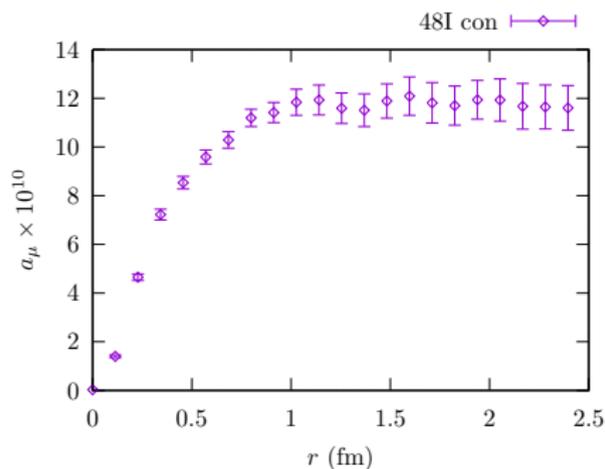
- Permutations of the three internal photons are not shown.
- **Gluons exchange between and within the quark loops are not drawn.**
- We need to make sure that the loops are connected by gluons by “vacuum” subtraction. So the diagrams are 1-particle irreducible.



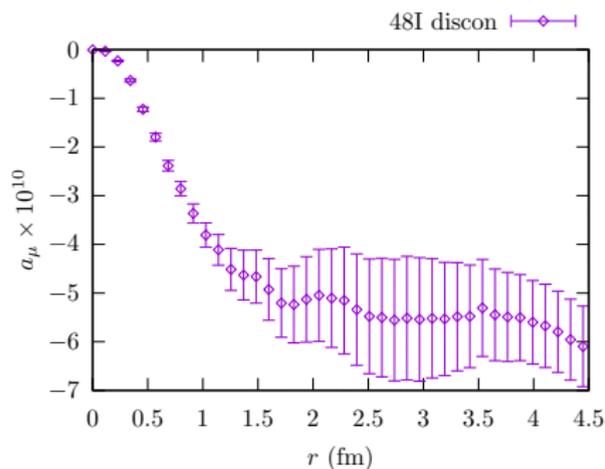
- Point  $x$  is used as the reference point for the moment method.
- We can use two point source photons at  $x$  and  $y$ , which are chosen randomly. The points  $x_{\text{op}}$  and  $z$  are summed over exactly on lattice.
- Only point source quark propagators are needed. We compute  $M$  point source propagators and all  $M^2$  combinations of them are used to perform the stochastic sum over  $r = x - y$ .

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Connected diagrams

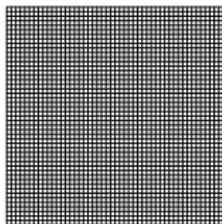
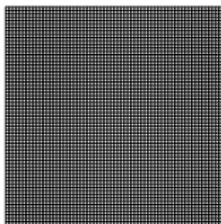


Disconnected diagrams

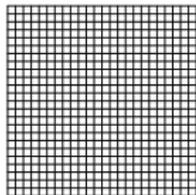
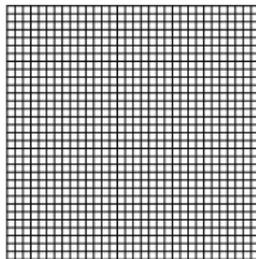
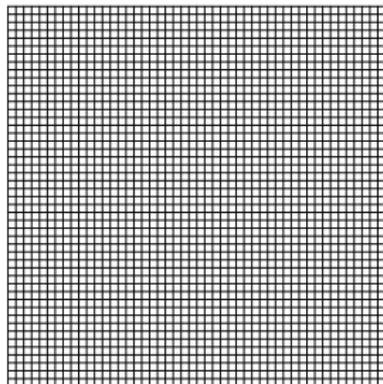
Partial sum is plotted above. Full sum is the right most data point.

$$a_\mu = 5.35(1.35)_{\text{stat}} \times 10^{-10} \text{ @ } L = 5.5\text{fm}, 1/a = 1.73\text{GeV}, m_\pi = 139\text{MeV}.$$

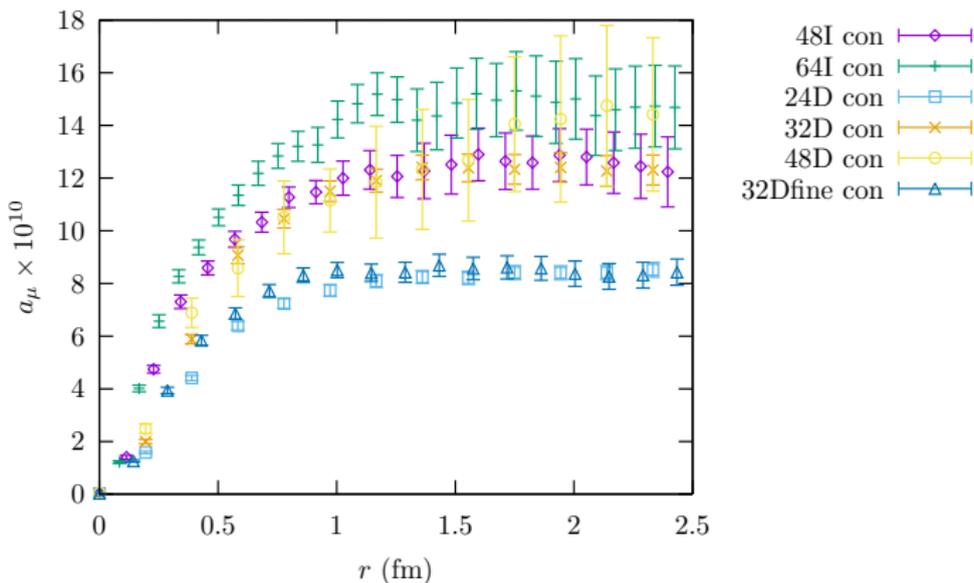
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48l:  $48^3 \times 96$ , 5.5fm box64l:  $64^3 \times 128$ , 5.5fm box

Phys. Rev. D 93, 074505  
(2016)

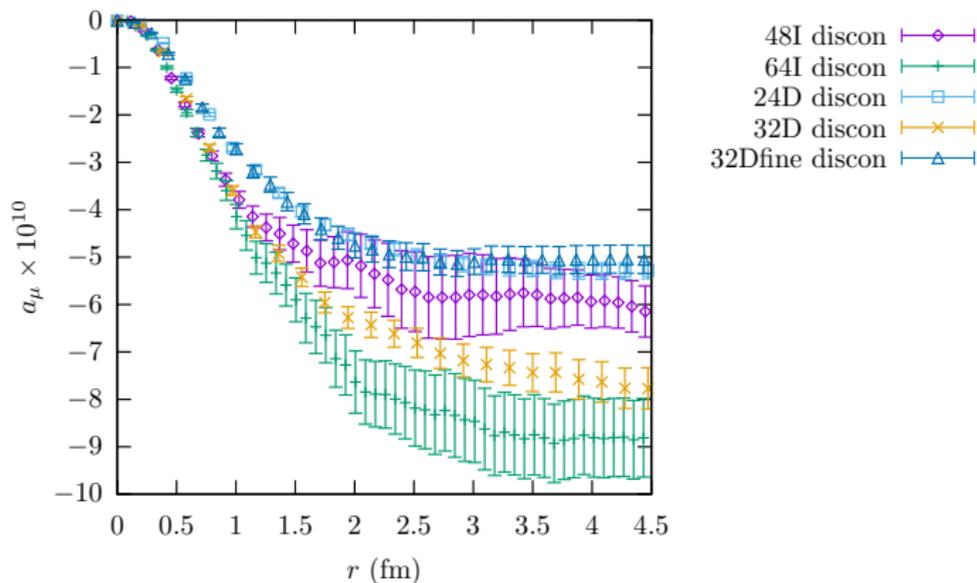
24D:  $24^3 \times 64$ , 4.8fm box32D:  $32^3 \times 64$ , 6.4fm box48D:  $48^3 \times 64$ , 9.6fm box32Dfine:  $32^3 \times 64$ , 4.8fm box

$$\frac{a_\mu}{m_\mu} \bar{u}_{s'}(\vec{0}) \frac{\Sigma}{2} u_s(\vec{0}) = \sum_{r=x-y} \sum_z \sum_{x_{\text{op}}} \frac{1}{2} (\vec{x}_{\text{op}} - \vec{x}_{\text{ref}}) \times \bar{u}_{s'}(\vec{0}) i \vec{\mathcal{F}}^C(\vec{0}; x, y, z, x_{\text{op}}) u_s(\vec{0})$$



Partial sum is plotted above. Full sum is the right most data point.

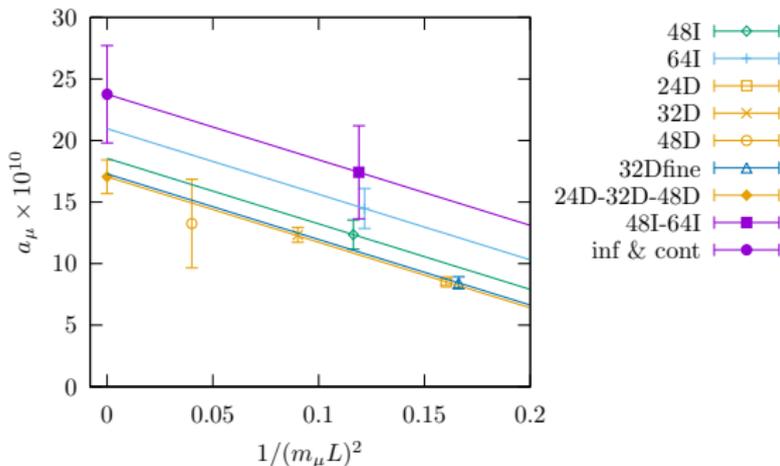
$$\frac{a_\mu}{m_\mu} \bar{u}_{s'}(\vec{0}) \frac{\Sigma}{2} u_s(\vec{0}) = \sum_{r=x-y} \sum_z \sum_{x_{\text{op}}} \frac{1}{2} (\vec{x}_{\text{op}} - \vec{x}_{\text{ref}}) \times \bar{u}_{s'}(\vec{0}) i \vec{\mathcal{F}}^C(\vec{0}; x, y, z, x_{\text{op}}) u_s(\vec{0})$$



Partial sum is plotted above. Full sum is the right most data point.

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} \right. \\ \left. - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$

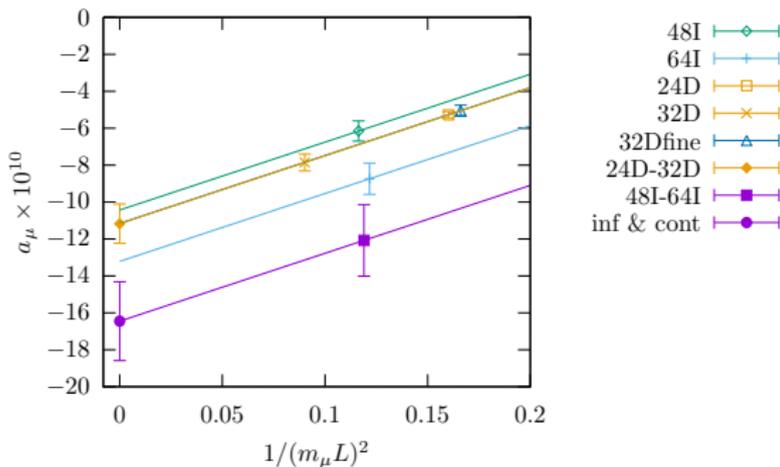
I-DSDR and Iwasaki ensembles have different  $\mathcal{O}(a^2)$  coefficients.



$$a_\mu = 23.76(3.96)_{\text{stat}} \times 10^{-10}$$

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$

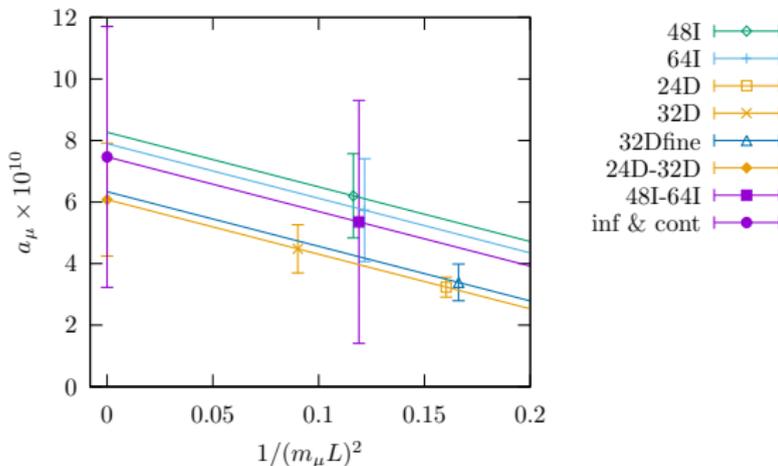
I-DSDR and Iwasaki ensembles have different  $\mathcal{O}(a^2)$  coefficients.



$$a_\mu = -16.45(2.13)_{\text{stat}} \times 10^{-10}$$

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} \right. \\ \left. - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$

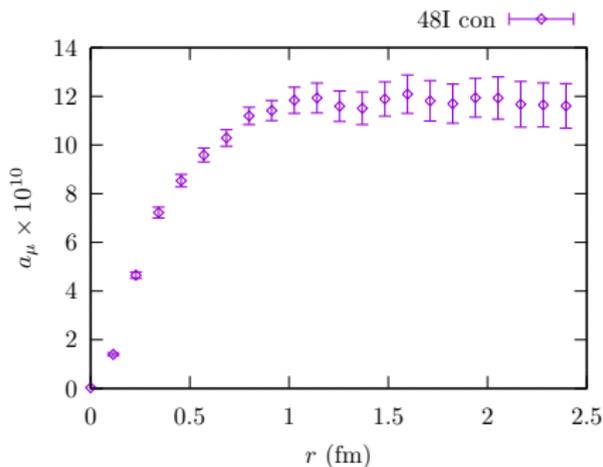
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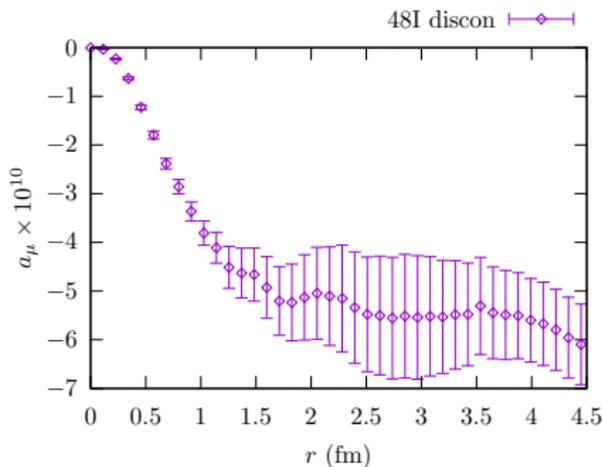
$$a_\mu = 7.47(4.24)_{\text{stat}} \times 10^{-10}$$

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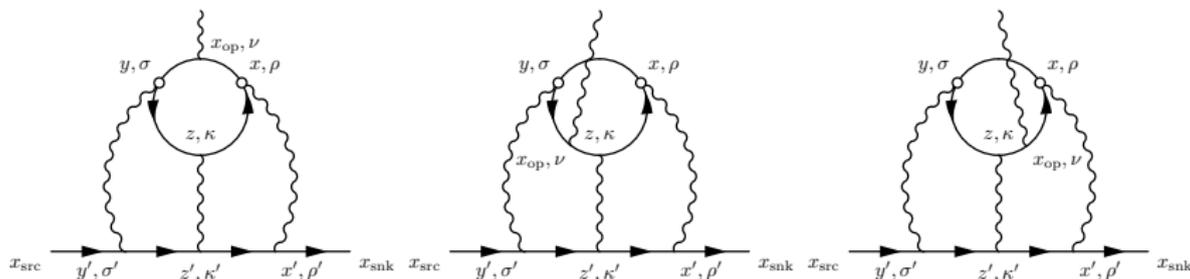
Connected diagrams



Disconnected diagrams

Partial sum is plotted above. Full sum is the right most data point.

Contribution to the connected diagrams mostly from small  $r$  ( $r < 1$  fm).



- The three internal vertex attached to the quark loop are equivalent (all permutations are included).
- We can pick the closer two points as the point sources  $x, y$ .

$$\sum_{x, y, z} \rightarrow \sum_{x, y, z} \begin{cases} 3 & \text{if } |x - y| < |x - z| \text{ and } |x - y| < |y - z| \\ 3/2 & \text{if } |x - y| = |x - z| < |y - z| \\ 3/2 & \text{if } |x - y| = |y - z| < |x - z| \\ 1 & \text{if } |x - y| = |y - z| = |x - z| \\ 0 & \text{others} \end{cases}$$

Split the  $a_\mu^{\text{con}}$  into two parts:

$$a_\mu^{\text{con}} = a_\mu^{\text{con,short}} + a_\mu^{\text{con,long}}$$

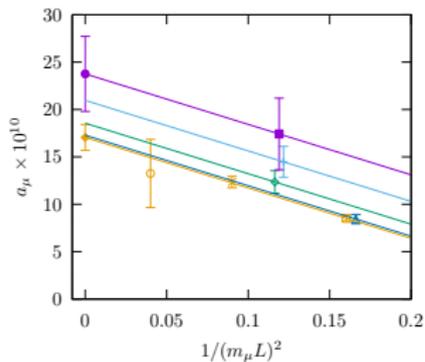
- $a_\mu^{\text{con,short}} = a_\mu^{\text{con}}(r \leq 1\text{fm})$ :  
most of the contribution, small statistical error.
- $a_\mu^{\text{con,long}} = a_\mu^{\text{con}}(r > 1\text{fm})$ :  
small contribution, large statistical error.

Perform continuum extrapolation for short and long parts separately.

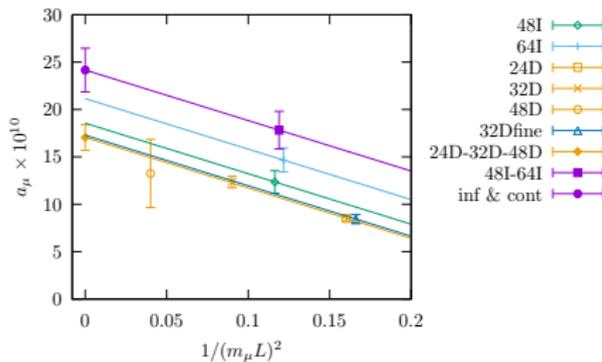
- $a_\mu^{\text{con,short}}$ : conventional  $a^2$  fitting.
- $a_\mu^{\text{con,long}}$ : simply use 48l value.  
Conservatively estimate the relative  $\mathcal{O}(a^2)$  error: it may be as large as for  $a_\mu^{\text{con,short}}$  from 48l.

$$a_\mu(L, a^l, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} - c_1^l (a^l \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$

I-DSDR and Iwasaki ensembles have different  $\mathcal{O}(a^2)$  coefficients.



Conventional continuum limit

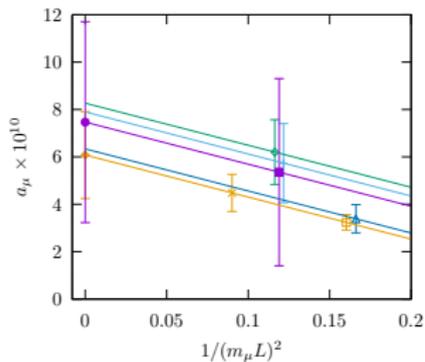


Hybrid continuum limit

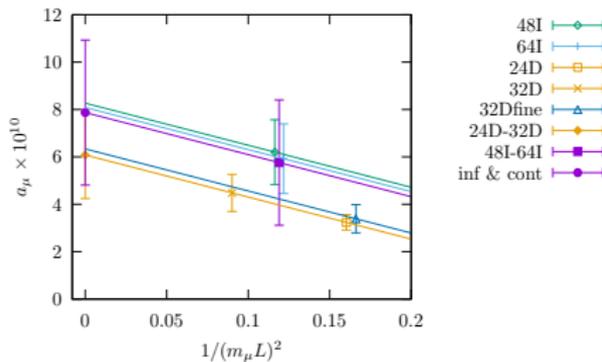
$$a_\mu = 23.76(3.96)_{\text{stat}} \times 10^{-10} \rightarrow 24.16(2.30)_{\text{stat}}(0.20)_{\text{sys}, a^2} \times 10^{-10}$$

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} \right. \\ \left. - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$

I-DSDR and Iwasaki ensembles have different  $\mathcal{O}(a^2)$  coefficients.



Conventional continuum limit



Hybrid continuum limit

$$a_\mu = 7.47(4.24)_{\text{stat}} \times 10^{-10} \rightarrow 7.87(3.06)_{\text{stat}}(0.20)_{\text{sys}, a^2} \times 10^{-10}$$

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	con	discon	tot
$a_\mu$	24.16(2.30)	-16.45(2.13)	7.87(3.06)
sys hybrid $\mathcal{O}(a^2)$	0.20(0.45)	0	0.20(0.45)
sys $\mathcal{O}(1/L^3)$	2.34(0.41)	1.72(0.32)	0.83(0.56)
sys $\mathcal{O}(a^4)$	0.88(0.31)	0.71(0.28)	0.95(0.92)
sys $\mathcal{O}(a^2 \log(a^2))$	0.23(0.08)	0.25(0.09)	0.02(0.11)
sys $\mathcal{O}(a^2/L)$	4.43(1.38)	3.49(1.37)	1.08(1.57)
sys strange con	0.30	0	0.30
sys sub-discon	0	0.50	0.50
sys all	5.11(1.32)	3.99(1.29)	1.77(1.13)

- Systematic error has some cancellation between the connected and disconnected diagrams.

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$

$\mathcal{O}(1/L^3)$

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} + \frac{b_2}{(m_\mu L)^3} - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$

$\mathcal{O}(a^2 \log(a^2))$

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} - \left( c_1^I (a^I \text{ GeV})^2 + c_1^D (a^D \text{ GeV})^2 - c_2^D (a^D \text{ GeV})^4 \right) \times \left( 1 - \frac{\alpha_S}{\pi} \log((a \text{ GeV})^2) \right) \right)$$

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} \right. \\ \left. - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$

$\mathcal{O}(a^4)$  (maximum of the following two)

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} \right. \\ \left. - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2 (a \text{ GeV})^4 \right)$$

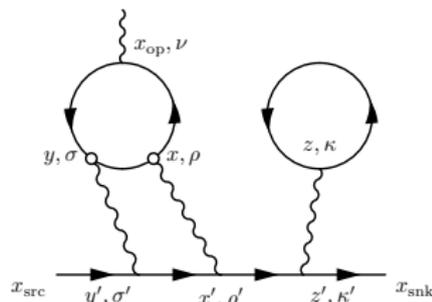
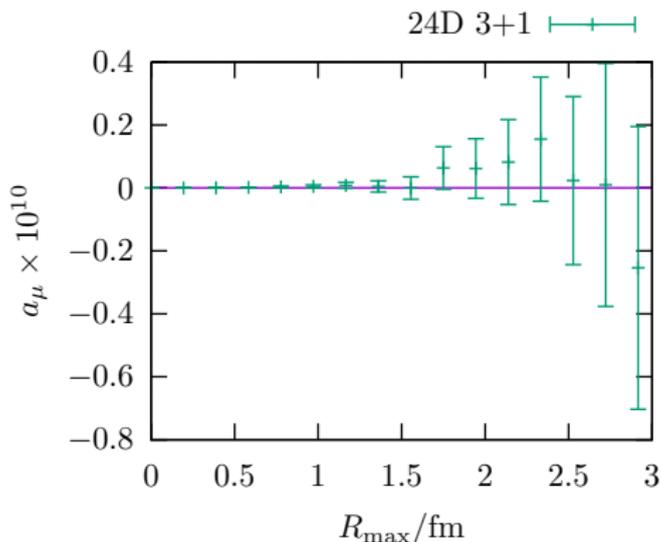
$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} \right. \\ \left. - c_1 (a \text{ GeV})^2 + c_2^I (a^I \text{ GeV})^4 + c_2^D (a^D \text{ GeV})^4 \right)$$

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} \right. \\ \left. - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$

$\mathcal{O}(a^2/L)$  (maximum of the following two)

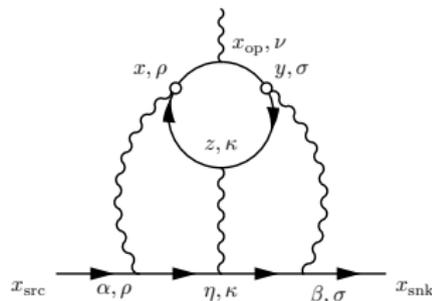
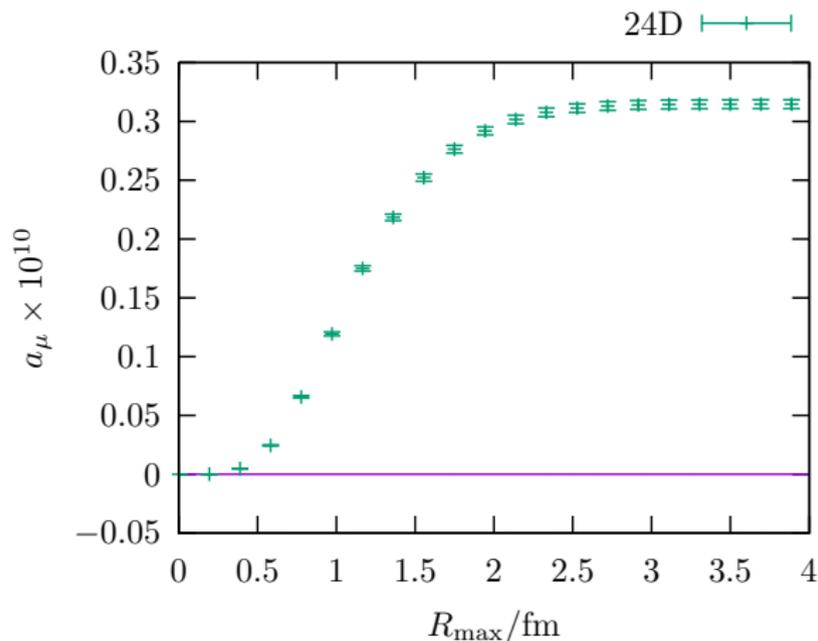
$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} \right. \\ \left. - \left( c_1^I (a^I \text{ GeV})^2 + c_1^D (a^D \text{ GeV})^2 - c_2^D (a^D \text{ GeV})^4 \right) \left( 1 - \frac{1}{m_\mu L} \right) \right)$$

$$a_\mu(L, a^I, a^D) = a_\mu \left( 1 - \frac{b_2}{(m_\mu L)^2} \right) \\ \times \left( 1 - c_1^I (a^I \text{ GeV})^2 - c_1^D (a^D \text{ GeV})^2 + c_2^D (a^D \text{ GeV})^4 \right)$$



- Partial sum upto  $R_{\max}$   
 $R_{\max} = \max(|x - y|, |x - z|, |y - z|)$
- The tadpole part comes from [C. Lehner et al. 2016 \(PRL 116, 232002\)](#)
- Systematic error (subdiscon):  $0.5 \times 10^{-10}$

- 24D:  $24^3 \times 64$   
 $L = 4.8$  fm
- $a^{-1} = 1.015$  GeV  
 $M_\pi = 142$  MeV  
 $M_K = 512$  MeV



- Partial sum upto  $R_{\max}$   
 $R_{\max} = \max(|x - y|, |x - z|, |y - z|)$
- Systematic error (strange con):  $0.3 \times 10^{-10}$

- 24D:  $24^3 \times 64$   
 $L = 4.8 \text{ fm}$
- $a^{-1} = 1.015 \text{ GeV}$   
 $M_{\pi} = 142 \text{ MeV}$   
 $M_K = 512 \text{ MeV}$

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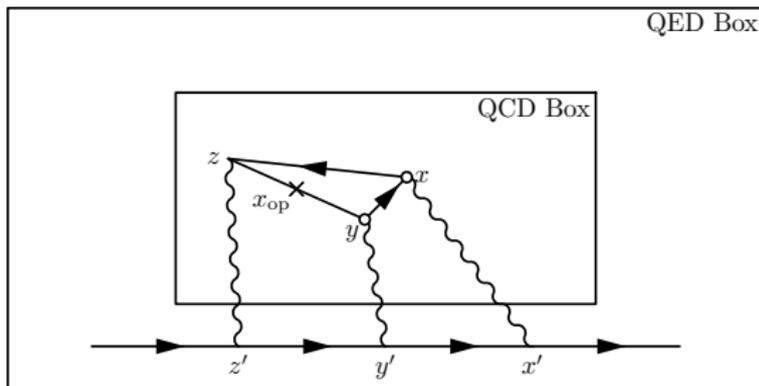
- $a_\mu = 7.87(3.06)_{\text{stat}}(1.77)_{\text{sys}} \times 10^{-10}$ .
- Consistent with hadronic model estimate:  
 $10.3(2.9) \times 10^{-10}$  (compiled by Fred Jegerlehner 2017).
- Leaves little room for the HLbL contribution to explain the difference between the Standard Model and the BNL experiment.
- Better accuracy is desired to compare with the on-going Fermilab muon  $g - 2$  experiments. Initial experimental result (using portion of the statistics) is expected to release later this year.
- Plan to invest in the infinite volume QED approach.

- Mainz group initially proposed the idea of calculating QED part of the process in infinite volume.

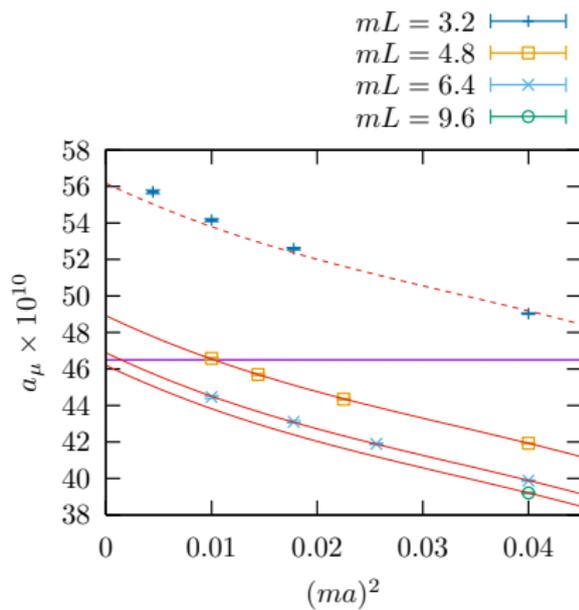
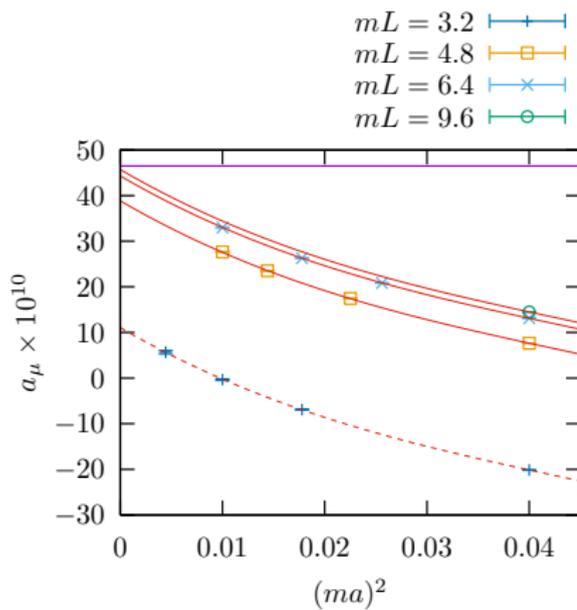
[N. Asmussen, J. Green, H. Meyer, A. Nyffeler 2016](#)

- Motivated by Mainz group, we have also started to work on this approach.

[T. Blum et al, PRD 96, 034515](#)

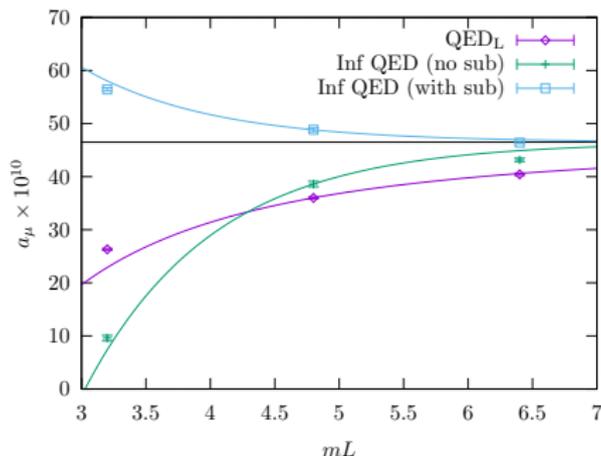


- Compare the two  $\mathfrak{G}_{\rho,\sigma,\kappa}(x, y, z)$  in **pure QED computation**.



- Notice the vertical scales in the two plots are different.

- Compare the finite volume effects in different approaches in **pure QED computation**,



- QED<sub>L</sub>:  $\mathcal{O}(1/L^2)$  finite volume effect, because the photons are emitted from a conserved loop. Phys.Rev. D93 (2016) 1, 014503.
- Inf QED (no sub):  $\mathcal{O}(e^{-mL})$  finite volume effect. Everything except the four-point-correlation function is evaluated in infinite volume. arXiv:1705.01067.
- Inf QED (with sub): smaller  $\mathcal{O}(e^{-mL})$  finite volume effect. arXiv:1705.01067.

- En-Hung Chao, Antoine Gerardin, Jeremy R. Green, Renwick J. Hudspith, and Harvey B. Meyer. [arXiv:2006.16224](https://arxiv.org/abs/2006.16224)
- Connected diagram:  $a_\mu = 9.89(25) \times 10^{-10}$ .
- Disconnected diagram:  $a_\mu = -3.35(42) \times 10^{-10}$ .
- Total:  $a_\mu = 6.54(49)(66)_{\text{sys-cont}} \times 10^{-10}$ .
- Adjust to physical pion/kaon mass:  $a_\mu = 10.41(91) \times 10^{-10}$ .  
Subtracting the  $\pi^0$ -pole contribution in this unphysical setup and add back the physical  $\pi^0$ -pole contribution.



Thank You!