# Magnetic field amplification by turbulent dynamo in relativistic collisionless shocks

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TeVPA2021, Chengdu, China (hybrid), 25-29, Oct. 2021

# Magnetic field around Collisionless Shocks

**Collisionless shocks in high-energy astrophysical phenomena:** 

- Acceleration site for cosmic rays (CRs).
- High-energy emission regions.
- → b-fields around a shock have a crucial role, it's highly uncertain!!







#### It's the key to unveil

- the acceleration mechanism of CRs.
- the environment of astrophysical objects.

## Problems about b-fields around the shock

• When the interstellar magnetic fields ( $\sim \mu G$ ) are compressed by



- The structure of magnetic field ?
- Usually, there is no information about pre-shock region.

### Magnetic Field Amplification in Collisionless Shocks

Plasma Instabilities

Weibel1959; Lucek & Bell 2000; Bell 2004; Spitokovsky 2008; Drury & Downes 2012; Tomita et al. 2016, 2019

Both mechanisms are required.

Balsara & Kim 2001; Giacalone & Jokipii 2007; Inoue et al. 2009, 2011; Zhang et al. 2009; Mizuno et al. 2011

#### B-field amplification by turbulent dynamo in 3D MHD simulation (Inoue et al. 2011)

Kinetic energy of the<br/>downstream turbulence>Magnetic energy<br/>Sano, T. et al. 2013

#### However,

MHD simulation cannot solve b-fields amplification concerning high-energy particles.



In the downstream region of collisionless shocks, MHD approximation is always valid?

**Density structures are maintained in post-shock regions?** 

- Collisionless shocks generate non-thermal particles.
- Particle diffusion is negligible??

In order for the MHD approximation to be applicable in downstream regions of collisionless shocks,

"gyro radius < size of density fluctuations" ?

 $\ll$  How large ?

+ Amplitude of density fluctuations?

#### **Purpose of Our Study:**

We search the surrounding environment required for b-field amplification by the turbulent dynamo.

(+ particle acceleration via turbulence??)

#### Method:

Particle-in-Cell(PIC) simulation (L $\sim$  plasma skindepth) of a relativistic collisionless shock propagating into inhomogeneous media (L  $\gg$  plasma skindepth).

### Today's talk:

When the density fluctuation has a small amplitude ( $\delta n_1/n_1 \leq 0.5$ ), MHD turbulence does not work in the downstream region of the relativistic collsionless shocks.

### Simulation set up

\*Source code given by Matsumoto Y.(Chiba Univ.)

- Two-dimensional electromagnetic PIC code\* & Athena++ MHD code.
- $e^{\pm}$  plasmas ( $\equiv e^{-}$ -ion plasmas in the downstream regions of relativistic shocks.)
- Calculator: Cray XC50 (520 core) @NAOJ
- Box Size:  $L_x = 3120 \, {}^c/_{\omega_{
  m pe}}$ ,  $L_y = 1200 \, {}^c/_{\omega_{
  m pe}}$  ( $\Delta x = \Delta y = 0.1 \, {}^c/_{\omega_{
  m pe}}$ )



## Simulation set up

- Upstream magnetization  $\sigma_{\rm e}=B_0^2/4\pi\Gamma n_0m_{\rm e}c^2$ .
- + (Clump size  $2r_{
  m c}$ )/ (Gyroradius  $r_{
  m ge}$ ), Amplitude  $\delta$  :

	δ	$\sigma_{ m e}$	$2r_c/r_{ge}$
case1	0.5	<b>10</b> -3	9.5
case2	0.5	<b>10</b> -5	0.9
case3	10.0	<b>10</b> -3	9.5

(For the case of GRB afterglows)

The required ratio of the Clump size to the Gyroradius, for b-field amplification by the turbulent dynamo is 10-100 in the downstream region (downstream rest frame).

### The ratio of a Clump Size to a Gyroradius

The condition required for turbulent dynamo to work

"Eddy Turn Over Time  $\leq$  Decceleration time of the shock",

predicts the maximum size of the upstream density fluctuations.

For Gamma-ray bursts, 
$$t_{dec} \approx 10^4 \sec \left(\frac{E_{iso,53}}{n_{ISM,0}}\right)^{\frac{1}{3}} \Gamma_{sh,2}^{-\frac{5}{3}}$$
,  $(n_1 \approx const.)$ .  
Since  $t_{eddy} \approx \frac{\lambda}{c} \leq t_{dec}$ , (\*Downstream rest frame)  
 $\lambda \leq 10^{14} \left(\frac{E_{iso,53}}{n_{ISM,0}}\right)^{1/3} \Gamma_{sh,2}^{-5/3} cm \approx 10^7 \left(\frac{E_{iso,53}}{n_{ISM,0}}\right)^{1/3} \Gamma_{sh,2}^{-5/3} \frac{c}{\omega_{pp}}$ .

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the gyroradius of downstream thermal protons,  $r_{\rm gp} \approx 10^4 \left(\frac{\sigma}{10^{-9}}\right)^{-\frac{1}{2}} \frac{c}{\omega_{\rm ex}}$ . Thus, when the downstream clump size is  $\lambda = 10^5 - 10^6 \frac{c}{\omega_{nn}}$ , The ratio of the clump size to the gyroradius is  $\frac{\lambda}{r_{op}} \approx 10-100$ .

Results: MHD ( $\sigma_{
m e}=10^{-3}$ , Amplitude  $\delta=0.5$ )



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## Results: PIC ( $\sigma_{\rm e}=10^{-3}$ , Amplitude $\delta=0.5$ )



\* There is no injection particles in the white region in density distribution yet.

## PIC, Amplitude dependence: Density 2D plots



Particle diffusion is suppressed,
so that the b-field is amplified
by the turbulent dynamo!

Why does the particle diffusion is suppressed? Thermal velocity of particles in the clump becomes lower than that for  $n/n_0 = 0.5$ ?



 $\beta_{x,cd}$ : Velocity of the Shocked clump

 $\Rightarrow n/n_0 \gtrsim 3$ 

# Summary

PIC simulations of a relativistic collisionless shock propagating into  $e^{\pm}$  plasmas with the density clump:

Even if the clump size is larger than the gyroradius of the upstream plasma,

the density fluctuation with small amplitude  $(n/n_0 < 1)$  could not drive efficient magnetic field amplification by the turbulent dynamo.

#### Future work:

We are performing many PIC simulations to understand the parameter dependence on the downstream turbulence and e<sup>-</sup> energy distribution.

## $\sigma_e$ dependence: Density 2-dimensional(2D) plots



# Density 2D plots: Comparison between PIC & MHD



The ratio of a Clump Size to a Gyroradius The condition required for turbulent dynamo to work "Eddy Turn Over Time  $\leq$  Decceleration time of the shock", predicts the maximum size of the upstream density fluctuations.

\* For Gamma-ray bursts,  $t_{
m dec} \approx 10^3 \sec E_{
m iso,53} v_{
m WR,8.3} \dot{M}_{WR,-5}^{-1} \Gamma_{sh,2}^{-3}$ ,  $(n_1 \propto r^{-2})$ ,

$$\approx 10^{4} \sec \left(\frac{E_{\rm iso,53}}{n_{\rm ISM,0}}\right)^{\frac{1}{3}} \Gamma_{\rm sh,2}^{-\frac{5}{3}}, \qquad (n_{1} \approx const.).$$
(\*Downstream rest frame)
Since  $t_{\rm eddy} \approx \frac{\lambda}{c} \leq t_{\rm dec}, \ \lambda \leq 10^{14} \left(\frac{E_{\rm iso,53}}{n_{\rm ISM,0}}\right)^{1/3} \Gamma_{\rm sh,2}^{-5/3} cm \approx 10^{7} \frac{c}{\omega_{\rm pp}}.$ 
(\*Downstream rest frame)
If the upstream b-field is 3uG, the gyroradius of thermal protons.

If the upstream b-field is  $3\mu G$ , the gyroradius of thermal protons,

$$r_{\rm gp} \approx 10^4 \left(\frac{\sigma}{10^{-9}}\right)^{-1/2} c/\omega_{\rm pp}.$$

Thus, when the clump size is  $\lambda = 10^5 - 10^6 \frac{c}{\omega_{\rm pp}}$ ,

The ratio of the clump size to the gyroradius is  $\frac{\lambda}{r_{gp}} \approx$ 10-100.

# Particle-in-Cell(PIC) Simulation

### **Fundamental Equations:**

**Equation of motion of N** particles  $\frac{du_{\rm s}}{dt} = \frac{q_{\rm s}}{m_{\rm s}} \left( E + \frac{u_{\rm s}}{c\gamma_{\rm s}} \times B \right),$  $\frac{dx_{s}}{dx_{s}} = \frac{u_{s}}{dx_{s}}$ dt  $\gamma_{\rm s}$ 

### The Algorithm:

- **1.** Compute charge density in a grid point from velocities of particle in a cell.
- 2. Compute electromagnetic field in a grid point  $B_{A}$ from the charge density.
- 3. Update velocity and position.

**Maxwell equations** 

