





## OVERVIEW AND RESULTS THE 8-YEAR HIGH ENERGY NON-STANDARD INTERACTIONS ANALYSIS

TeVPA 2021 Parallel 10/28/21

GRANT K. PARKER grant.parker@mavs.uta.edu

## The IceCube Neutrino Observatory



**Figure**: The IceCube neutrino observatory. DeepCore is an additional collection of strings that allow for signals as low as 5 GeV.

- World's largest-volume neutrino detector and telescope, located under the ice at the South Pole
- Detection mechanism:
  - Neutrino interaction with ice produces charged products.
  - Interaction products have sufficient energy to generate Cherenkov radiation along their trajectory.
  - Collected light allows for neutrino direction and energy reconstruction.
  - Two event types: tracks (linear trajectories) and cascades (blob-like signal).
  - **DeepCore** is collection of specialized center strings that lower the event energy threshold to ~5 GeV



## Non-Standard Interactions

#### NSI arise in accounting for mass effects on neutrino oscillations

- Standard Model (SM) cannot account for neutrino masses
- If neutrino masses are Majorana, they can be added as a dim(5) Weinberg operator
- In many theories, a dim(6) operator accompanies the dim(5) operator, from which NSI arise
- In a basic model, coupling strength  $\varepsilon \sim g_X^2 m_W^2/m_X^2$ , so for natural values of *g*, the mediator > 1TeV



Figure: Diagramed neutrino-matter interactions for SM and NSI [Ohlsonn 2013 arXiv 1209.2710].

#### Neutrino oscillations are only sensitive to neutral-current NSI:

$$\mathcal{L}_{\mathrm{NSI}} = -2\sqrt{2}G_F \sum_{f,P,\alpha,\beta} \varepsilon_{\alpha\beta}^{f,P} (\bar{\nu}_{\alpha}\gamma^{\mu}P_L\nu_{\beta}) (\bar{f}\gamma_{\mu}Pf) \longrightarrow \varepsilon_{\alpha\beta} = \sum_{f,P} \varepsilon_{\alpha\beta}^{fP} \frac{N_f}{N_e} \longrightarrow H_{\mathrm{mat}} = \sqrt{2}G_F N_e(x) \begin{pmatrix} 1 + \varepsilon_{ee}(x) \ \varepsilon_{e\mu}(x) \ \varepsilon_{e$$



## Non-Standard Interactions

#### If NSI exist:

- NSI <u>fit better</u> than the SM to some measurements of data, suggesting the alleviation of tensions
  - Example Tension (Top Right): Constraints on *o*<sub>CP</sub> from T2K and NOvA [A. Himmel for the NOvA Collaboration (plenary)]
- However, NSI also introduce degeneracies to parameters such as  $\pmb{\theta}_{23}$  and  $\pmb{\delta}_{\rm CP}$ 
  - **Example Degeneracy (Bottom Right)**: Loss of  $\theta_{23}$  octant sensitivity as a function of  $|\varepsilon_{e\mu}|$  [Agarwalla et. al. 2016 <u>arXiv:1607.01745</u>]

## <u>Therefore, NSI are very</u> <u>compelling to study!</u>





#### NSI are a modification to the matter potential:

- IceCube atmospheric neutrinos
  - Various energies + various matter baselines (right) = optimal sample for BSM oscillations searches through muon track appearance and disappearance
- The rate of detected atmospheric neutrinos far exceeds backgrounds and other signal types (bottom)



Figure: (Top) Model of neutrinos reaching IceCube from multiple atmospheric sites (Credit: E. Lohfink for the IceCube Collaboration). (Bottom) Rates of neutrinos from different sources (S. Axani for the IceCube Collaboration).

downgoing



## Other IceCube Analyses:

- The 8-year DeepCore analysis (ongoing):
  - Fit on *all* complex NSI parameters
  - Different parameterization that allows for constraining the differences between the diagonal Hamiltonian elements
- Sample: ~5-300 GeV, 8.2 years of data, ~300,000 events
  - Binning in energy, zenith, and topology (tracks and cascades)
- DeepCore analysis 3-year (<u>Phys. Rev. D</u> <u>104, 072006</u>) : 50k events



**Figure**: Sensitivities for the upcoming 8-year low-energy NSI result. Credit: E. Lohfink for the IceCube collaboration.



## The 8-Year High-Energy Analysis

- Our analysis fits at much higher energies (500 GeV - 10 TeV) than DeepCore:
  - Only muon tracks, 300k events,7.6 years of data
  - This is to constrain a single parameter,  $\epsilon_{\mu\tau}$ 
    - ϵ<sub>μτ</sub> has the predominant effect on expected fluxes at high energies (right figure)



Figure: Neutrino oscillation probabilities with NSI. Red represents positive NSI values, while blue represents negative values. Credit: E. Lohfink for the IceCube Collaboration.



## Analysis Predecessor

## IceCube's latest sterile neutrino search published last year:

- 305,891 CC  $\nu$  and  $\overline{\nu}$  muon track events (7.64 years)
- Muon energy proxy: 500 9976 GeV
- Baseline MC: 500 years equivalent livetime
- Muon anti/neutrino disappearance shape signal

The larger data set was met with an updated analysis framework:

- MC Treatment with compactification abilities
- New analysis software for weighting and fitting
- Improved systematic treatment for larger sample
- Optimized event selection with reduced background and boosted statistics



**Figure**: Results of Analysis I (left) and Analysis II (right), each presenting best-fits and CL contours for their respective parameter spaces. [arXiv:2005.12943]



## Parameter of Interest for This Analysis:

- At the energies and baselines of this analysis,  $\nu_{\rm e}$  decouples from oscillation and vacuum terms are suppressed.
  - Leads to only mu-tau oscillation, which means the mu-tau NSI parameter becomes predominant (Right)
- This is seen in simulated fluxes— all parameters except  $\varepsilon_{\mu\tau}$  only enhance/suppress fluxes at O(10%) for parameter strength 0.01 (Below)

$$P(\nu_{\mu} \to \nu_{\tau}) \simeq \left(\sin 2\theta_{23} \frac{\Delta m_{31}^2}{2 E_{\nu}} + 2 V_d \varepsilon_{\mu\tau}\right)^2 \left(\frac{L}{2}\right)^2$$





When we simulate neutrino and antineutrino fluxes independently, we see:



This confirms the theory prediction: that as the sign of the matter potential changes between anti/neutrinos, signal shapes switch between +/-  $\text{Re}(\epsilon_{\mu\tau})$ 

#### Grant K. Parker



Combining the  $\nu$  and  $\overline{\nu}$  effects, the expected signal at IceCube is predominantly disappearance, and is much weaker than a pure  $\nu$  or  $\overline{\nu}$  flux.



UTA

Let's rescale to look at the signal shape:



UTA

## Inverted Hierarchy Prediction



## Sensitivity and Impact of Systematic Uncertainties

#### The 90% CL sensitivity is given in black ("central").

#### Impact of Systematic Uncertainties:

- Individual systematic uncertainties have have little impact on the sensitivity.
- For testing, we group uncertainties by type, then turn them off individually to measure their impact (right).

#### **Breakdown of Categorizations**

- Bulk Ice: Uncertainties from the optical properties of South Pole glacial ice.
- Hole Ice: Uncertainties from the optical properties of refrozen ice in the drilled sensor column
- DOM Efficiency: How well the light sensors operate post-installation
- Atmospheric Neutrino Flux: Uncertainties in the production factors for atmospheric neutrinos
- Cosmic Ray and Astrophysical Neutrino Flux: Uncertainties regarding the fluxes of cosmic rays and astrophysical neutrinos





## Results (Real-Only)

- Figure (Top)
  - $\operatorname{Re}(\boldsymbol{\varepsilon}_{\mu\tau})$  results for this analysis, including the -2LLH profile for the data and the obtained CL intervals
- Figure (Bottom)
  - Comparison of the  $\operatorname{Re}(\varepsilon_{\mu\tau})$  result from this analysis to other leading analyses that constrain  $\operatorname{Re}(\varepsilon_{\mu\tau})$  only
- Best-Fit:

 $\circ \boldsymbol{\varepsilon}_{\mu\tau} = -0.0029$ 

- Significance:
  - p-value = 0.252
  - Result is 0.68 **o** from the mean NSI value recovered from 1000 null hypothesis (no NSI) trials
- 90% CL Limits:

• -0.0041 <  $\varepsilon_{\mu\tau}$  < 0.0031

IceCube obtains a new world-leading limit





## Results (Complex)

#### • Figure:

- CL regions for complex  $\boldsymbol{\varepsilon}_{\mu\tau}$
- The Im( $\boldsymbol{\varepsilon}_{\mu\tau}$ ) component is degenerate for a given Re( $\boldsymbol{\varepsilon}_{\mu\tau}$ ) (see slide 21)
- Due to the degeneracy, we fit only along  $\operatorname{Re}(\varepsilon_{\mu\tau})$  and infer the  $\operatorname{Im}(\varepsilon_{\mu\tau})$  limits. This has the advantages of faster computing time and one less degree of freedom to fit to.





## **Global Comparison**

#### • Figure:

- Few analyses fit to complex  $\boldsymbol{\varepsilon}_{\mu\tau}$ 
  - We compare our result and sensitivity to those of IceCube DeepCore analyses
- DeepCore analyses are in the low-energy regime (5-300 GeV), where *v*<sub>e</sub> does not effectively decouple from atmospheric neutrino oscillations
- $\circ~$  While the Im( $\pmb{\varepsilon}_{\mu\tau}$ ) symmetry is still observed, the overall radial symmetry for the contours is lost in the 3-neutrino mixing regime





## Conclusions

In the era of high-precision neutrino physics, NSI must be constrained

With 8 years of high-energy muon neutrino data, our result sets 90% CL limits of  $\varepsilon_{\mu\tau}$  parameter: -0.0041 <  $\varepsilon_{\mu\tau}$  < 0.0031

This analysis sets the **tightest constraint on any NSI parameter in any channel** globally, with improvement on the nearest limits of  $\varepsilon_{\mu\tau}$  by a factor of ~2

# Thank You

# Backup Slides

## **Sterile Results**

• Below is the comparison of 90% CL results for Analysis I (left) and Analysis II (right) to other published limits.



## Symmetry Between +Im(NSI) and -Im(NSI

- For a given Re, +Im and -Im are identical in the two-neutrino calculation (right)
- Therefore, we can:
  - 1: Scan only the Re axis to get all the information
  - 2: Calculate statistics object for 1 DoF

\_\_\_\_

$$P(+) = \left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} + 2(a+bi)V_d \right|^2 \left(\frac{L}{2}\right)^2$$

$$P(-) = \left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} + 2(a-bi)V_d \right|^2 \left(\frac{L}{2}\right)^2$$

$$P(+)/P(-) = \frac{\left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} + 2(a+bi)V_d \right|^2}{\left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} + 2(a-bi)V_d \right|^2}$$

$$\frac{\left( \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} \right)^2 + 4(a)V_d \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} + 4(a^2+b^2)V_d^2}{\left( \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} \right)^2 + 4(a)V_d \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} + 4(a^2+b^2)V_d^2}$$

= 1

## **NSI LLH Offset**

- At sample energies, the electron flavor state decouples from atmospheric oscillation.
- Below is the approximate calculation of the difference in probabilities for NSI values with equal imaginary components, opposite-sign real components.

$$P(-) - P(+) =$$

$$\left[ \left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} + 2(-a+bi)V_d \right|^2 - \left| \sin 2\theta_{23} \frac{\Delta m_{31}^2}{2E_{\nu}} + 2(a+bi)V_d \right|^2 \right] \left(\frac{L}{2}\right)^2$$
$$= -aV_d L^2 \sin 2\theta_{23} \frac{\Delta m_{31}^2}{E_{\nu}}$$

= negative value

This confirms what we see in the  $-2\Delta$ LLH distribution, as  $-2\Delta$ LLH $(-a + bi) < -2\Delta$ LLH(a + bi).

## **Expected Distributions**

#### Signal in reconstruction space:



