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Gamma-rays from young SNRs in dense circumstellar environments

Robert Brose, Jonathan Mackey, Iurii Sushch TeVPa, 25-29 October 2021



Galactic cosmic-rays Possible sources



Figure: The cosmic ray spectrum (Blümer et al. 2019)

(Funk et al. 2013)



Figure: IC443 – multi wavelength image (credit: Dieter Willasch)



SNRs as cosmic-ray sources Experimental evidence

- Some CC supernovae are bright in radio, optical and X-rays (Type-IIn)
- Bright Radio and X-ray SNe often show signs of circumstellar interaction (Chevalier 1982)
- Radio emission is synchrotron → also expect gamma-ray emission
- Theoretical models predict high-energy radiation if particle acceleration is efficient (e.g., Murase+2011, Marcowith+2018)

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Radio Lightcurves for All SNe with Known Types

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SNRs as cosmic-ray sources Experimental evidence

- H.E.S.S. Collaboration (2019) obtained upper limits on TeV emission from nearby CCSNe
- Ongoing HESS ToO programme so far has no detections
- Xi+(2020) detected γ-rays from the location of SN 2004dj, a bright and nearby SN IIP, with FERMI-LAT
- A recent variability analysis of FERMI-LAT data found evidence in support of 2 further detections (Prokhorov+2021).
- Also suggestion of increasing flux from SN1987A with FERMI-LAT (Malyshev+2019)



Fermi acceleration Coupled equations

Cosmic-ray transport equation

Hydro equations

Magnetic Turbulence

Magnetic field

Standard DSA

Powered by

Non-linear DSA NDSA + high MF

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- 10+ papers
- 140+ citations

Other contributions using RATPaC:

[18] SNRs: CR escape and gamma-ray halo formation

[88] Spectral softening in CC-SNRs



Fermi acceleration The equations

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$$\frac{\partial E_W}{\partial t} = - \left(v \nabla_r E_W + c \nabla_r v E_W \right) + k^3 \nabla_k D_k \nabla_k \frac{E_W}{k^3} + 2 \left(\Gamma_g - \Gamma_d \right) E_W$$

Advection + Compression Cascading Growth + Damping $\frac{\partial}{\partial t} \begin{pmatrix} \varrho \\ m \\ E \end{pmatrix} + \nabla \begin{pmatrix} \varrho v \\ mv + (P)I \\ (E+P)v \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ L \end{pmatrix} \qquad \frac{\rho v^2}{2} + \frac{P}{\gamma - 1} = E$

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The equations are solved:

- One dimensional
- Assuming spherical symmetry
- Including Synchrotron and inverse-Compton cooling for electrons
- On a comoving, expanding grid for turbulence and CRs → no free escape boundary

Hydrodynamic Setup SNe expanding into a steady wind medium

"Luminous Blue Variable (LBV)" model

- Dense and powerful wind from a massive star
- $\dot{M} = 10^{-2} M_{\odot} / yr$
- $v_{\infty} = 100 \ km/s$

Red Supergiant (RSG) Model

- Slow dense wind from RSG progenitor, at upper end of mass range
- $\dot{M} = 8 \times 10^{-5} M_{\odot}/yr$
- $v_{\infty} = 15 \ km/s$
- Model 1: fixed 5µG magnetic field in the wind at all radii
- Model 2: **1** G at 1000 Rsun stellar field with $B \propto r^{-1}$ at large r
- Both cases: Ejecta-distribution following Chevalier (1982) with: **LBV:** $E = 10^{51} erg$, $M_{ej} = 10 M_{\odot}$, density power-law index n = 10**RSG:** $E = 10^{51} erg$, $M_{ej} = 3 M_{\odot}$, density power-law index n = 9
- Initial ejecta-radius: $10^{14} cm$.

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Results

Thermal X-ray emission Model vs. measurement





- Continuum X-ray emissivity in post-shock medium
- Unabsorbed X-ray luminosity far above observations
- Considering local absorption in stellar wind and shocked shell, obtain better agreement with observations (Dwarkadas 2014)
- Peak of lightcurve later in denser wind because of absorption



Non-thermal particle distribution Time evolution of particle spectrum



- Fit momentum spectrum with power-law plus exponential cutoff → E_{max}
- Acceleration time is ~1 month to get to quasi-steady state
- $E_{max} \approx 250$ TeV constant (and weak) upstream B-field
- $E_{max} \approx 700 \text{ TeV}$ for $B \propto r^{-1}$ ambient-field



(stellar field is already quite strong)





Figure: Maximum energy over time.

Non-thermal particle distribution Time evolution of gamma-ray emission





- Gamma-ray luminosity in FERMI-LAT energy range rises fastest after explosion
- Type IIn approx. 100x more luminous then IIP.
- TeV emission peaks later than GeV emission, for a weak field
- Gamma-gamma absorption not yet included!

Non-thermal particle distribution Detectability of gamma-ray emission



- Optimal detection prospects ≈ 1.5 weeks post explosion
- Type IIn detectable out to 10-30 Mpc
- Type IIP about 100x fainter → must be 10x closer for detection.
- All cases have GeV emission fading strongly after a few months.



Non-thermal particle distribution Radio emission

- Radio emission risetimes and peakluminosities consistent with observations
- Injection-fraction based on historical remnants (e.g. SN1006) and magnetic fieldamplification automatically produce the right radio-flux

No additional downstream-field generation needed.



 $\log(\frac{t}{1 \text{ days}})$

Non-thermal particle distribution Interaction with inhomogeneous media: LBV

- Adding a 1M_☉ shell to the background wind of the LBV at 6 × 10¹⁵ cm / 0.002pc
- Gamma-ray luminosity increases by about 10x after the shell collision (for shell with 10x higher density ratio)
- TeV luminosity peaks shortly after GeV and then declines.
- Accompanying signatures in X-ray, Radio and other wavebands? Trigger criteria for IACTs are needed!



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Figure: Gamma-ray luminosity over time for a LBV interacting with a $1M_{sol}$ shell.

Conclusions and future work

1D time-dependent modelling of hydrodynamics, particle acceleration, magnetic-turbulence and high-energy radiation, from SN expanding into LBV and RSG winds

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- Verify that thermal X-ray emission consistent with observations of Type-IIn/IIP SN
- Maximum particle energy of $E_{max} \le 700$ TeV from 0.1-10 year timescales \rightarrow Unlikely to be PeVatrons
- GeV emission peaks on days to week timescales
- TeV on week to month timescale
- LBV progenitors about 100x more luminous in thermal X-rays and GeV/TeV gamma-rays.
- Circumstellar shells can produce late-time re-brightening at GeV and TeV.
- A few SN could be detectable with FERMI-LAT over its mission lifetime

Thank you for your attention!