

$b \rightarrow sl\ell$ angular analyses on CMS

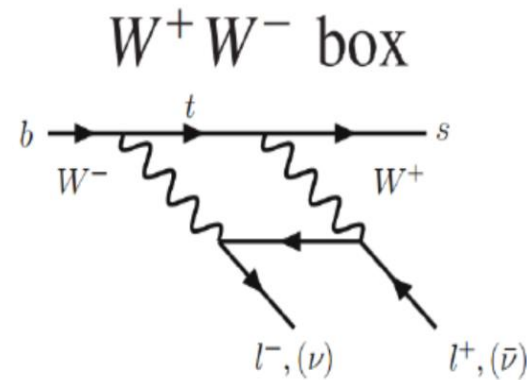
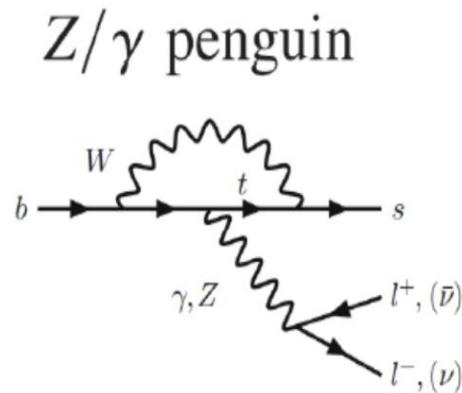
Chuang Jiang
on behalf of CMS collaboration

Outline

- Introduction
- $b \rightarrow sll$ angular analyses on CMS Run1
 - $B^+ \rightarrow K^+ \mu^+ \mu^-$
 - $B^0 \rightarrow K^{*0} \mu^+ \mu^-$
 - $B^+ \rightarrow K^{*+} \mu^+ \mu^-$
- $b \rightarrow sll$ angular analyses on CMS Run2
 - $B^+ \rightarrow K^+ \mu^+ \mu^-$
 - $B^0 \rightarrow K^{*0} \mu^+ \mu^-$
 - $B_s \rightarrow \Phi \mu^+ \mu^-$
- Prospect of HL-LHC
- Summary

Introduction: theoretical motivation

- New physics can be discovered by two ways:
 1. by producing new particles
 2. by searching discrepancies between measured observables and SM predictions.
- In SM, $b \rightarrow s ll$ is a flavor-changing neutral current (FCNC) process forbidden at tree level.



- New physics can contribute to the loop diagrams and make pronounced modifications.

Introduction: effective hamiltonian

- ❖ $b \rightarrow sll$ process can be described using effective theory[1], the effective Hamiltonian is:

$$H_{eff} = -\frac{4G_F}{\sqrt{2}} V_{tb} V_{ts}^* \sum_i C_i(\mu) O_i(\mu)$$

Where $O_i(\mu)$ indicate operator, $C_i(\mu)$ is Wilson coefficients, μ is the scale

- ❖ For SM, $i = 1, \dots, 10$. Terms 7, 9, 10 dominate in $b \rightarrow sl\bar{l}$

$$\mathcal{O}_7 = \frac{e}{(4\pi)^2} \bar{m}_b [\bar{s} \sigma^{\mu\nu} P_R b] F_{\mu\nu}, \mathcal{O}_9 = \frac{e^2}{(4\pi)^2} [\bar{s} \gamma_\mu P_L b] [\bar{l} \gamma^\mu l], \mathcal{O}_{10} = \frac{e^2}{(4\pi)^2} [\bar{s} \gamma_\mu P_L b] [\bar{l} \gamma^\mu \gamma_5 l],$$

- ❖ To describe new physics in $b \rightarrow sll$. Following terms are taken into consideration.

$$\mathcal{O}_S^l = \frac{e^2}{(4\pi)^2} [\bar{s} P_R b] [\bar{l} l],$$

$$\mathcal{O}_S^{ll} = \frac{e^2}{(4\pi)^2} [\bar{s} P_L b] [\bar{l} l],$$

$$\mathcal{O}_P^l = \frac{e^2}{(4\pi)^2} [\bar{s} P_R b] [\bar{l} \gamma_5 l],$$

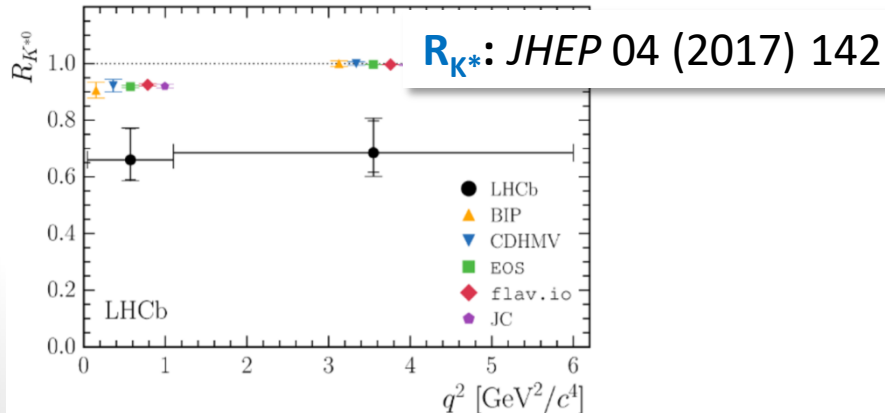
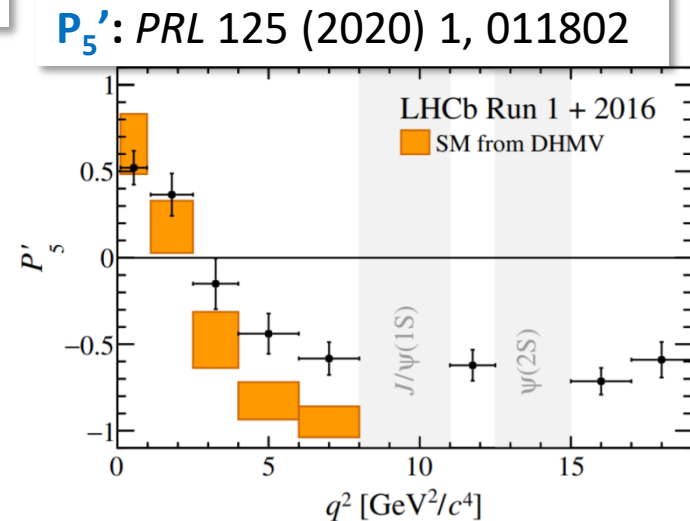
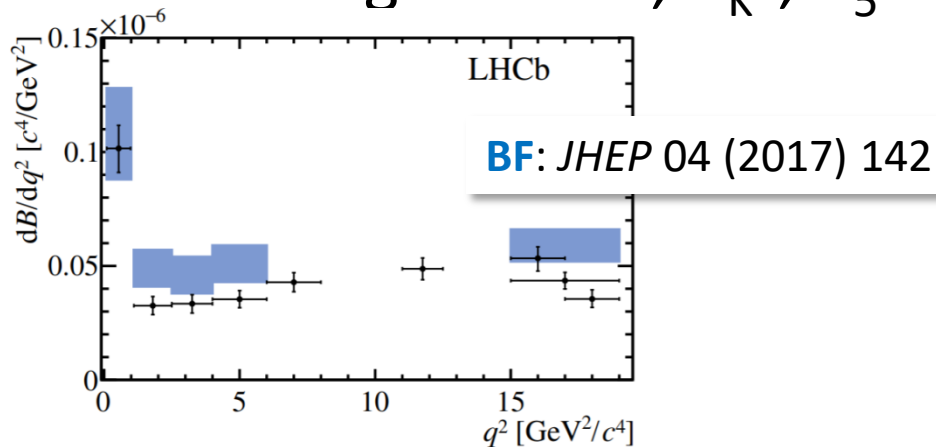
$$\mathcal{O}_P^{ll} = \frac{e^2}{(4\pi)^2} [\bar{s} P_L b] [\bar{l} \gamma_5 l],$$

$$\mathcal{O}_T^l = \frac{e^2}{(4\pi)^2} [\bar{s} \sigma_{\mu\nu} b] [\bar{l} \sigma^{\mu\nu} l],$$

$$\mathcal{O}_{T5}^l = \frac{e^2}{(4\pi)^2} [\bar{s} \sigma_{\mu\nu} b] [\bar{l} \sigma^{\mu\nu} \gamma_5 l],$$

Introduction: $b \rightarrow sll$ anomaly

- Some observables of $b \rightarrow sll$ process are found deviating from SM predictions by 2-3 σ , e.g. branching fraction, R_{K^*} , P_5' ...

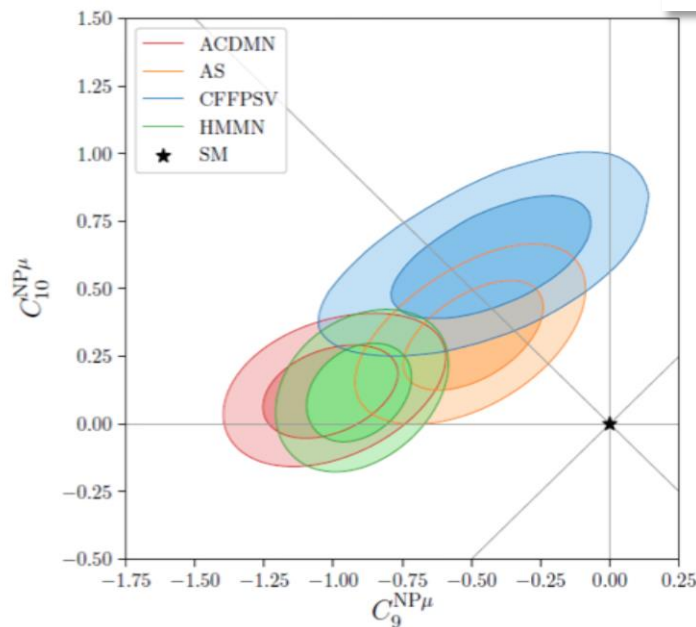


Introduction: $b \rightarrow sll$ anomaly

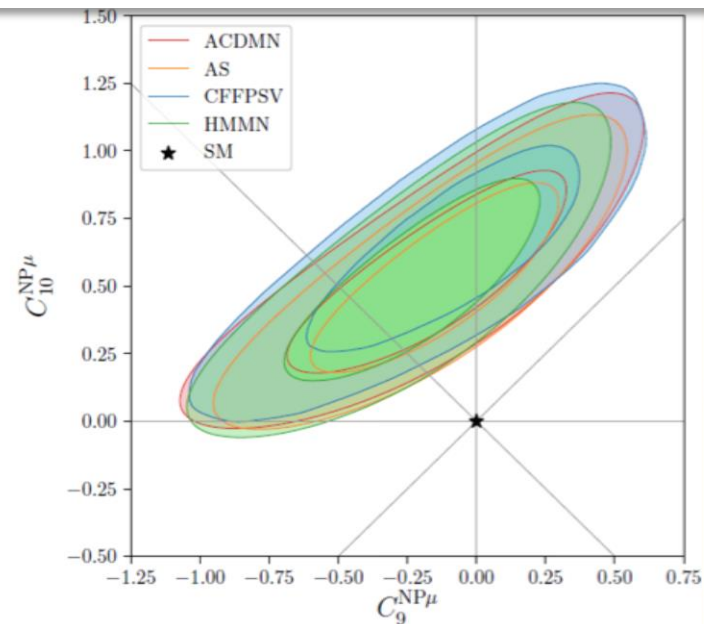
- Global fit of many observables from various measurements presents a result in tension with SM prediction by $>4\sigma$.

2-dimensional fits

*B. Capdevila, M. Fedele, S. Neshatpour, P. Stangl
@ Flavour Anomaly Workshop, 20 Oct. 2021*



global fit



fit to LFU observables + $B_s \rightarrow \mu\mu$

Introduction: angular analyses

- Investigate the angular distribution of the final state particles.
- **Pros:**
 - More utilizable observables than branching fraction measurement. More information.
 - Some observables can be constructed to reduce the hadronic uncertainty.
- **Cons:**
 - Low statistics. More dimensions to be investigate than branching fraction measurement
 - Difficult to determine the angular distribution of background

Introduction: CMS

CMS DETECTOR

Total weight : 14,000 tonnes
Overall diameter : 15.0 m
Overall length : 28.7 m
Magnetic field : 3.8 T

STEEL RETURN YOKE
12,500 tonnes

SILICON TRACKERS
Pixel ($100 \times 150 \mu\text{m}^2$) $\sim 1.9 \text{ m}^2 \sim 124\text{M}$ channels
Microstrips ($80\text{--}180 \mu\text{m}$) $\sim 200 \text{ m}^2 \sim 9.6\text{M}$ channels

SUPERCONDUCTING SOLENOID
Niobium titanium coil carrying $\sim 18,000 \text{ A}$

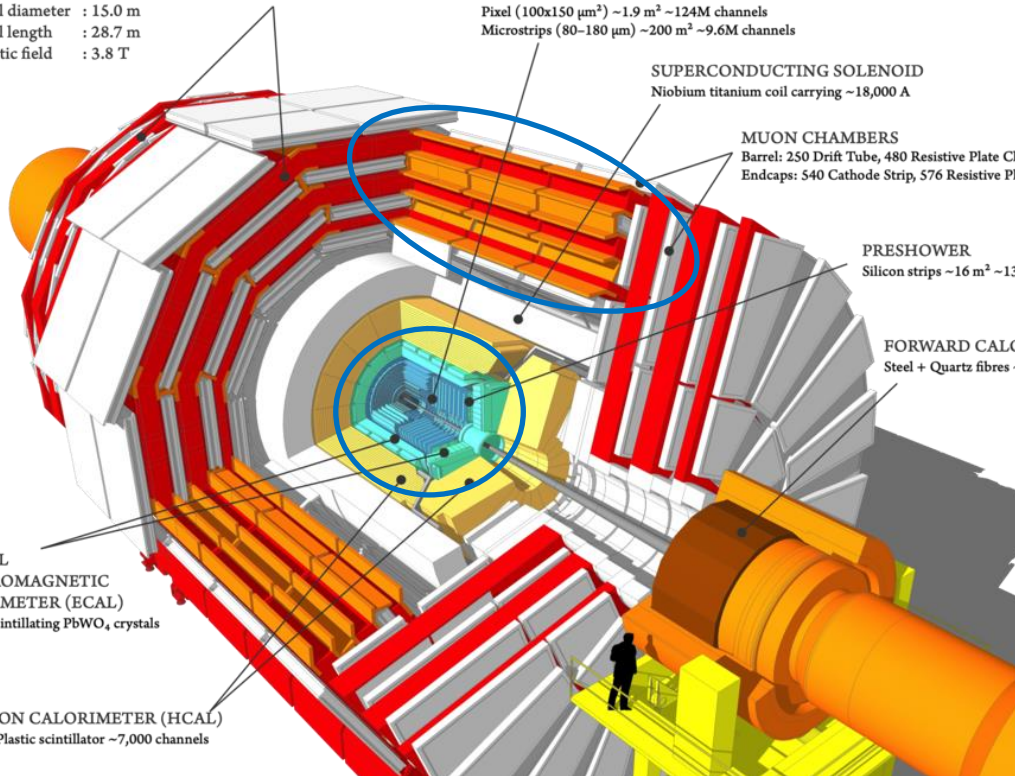
MUON CHAMBERS
Barrel: 250 Drift Tube, 480 Resistive Plate Chambers
Endcaps: 540 Cathode Strip, 576 Resistive Plate Chambers

PRESHOWER
Silicon strips $\sim 16 \text{ m}^2 \sim 137,000$ channels

FORWARD CALORIMETER
Steel + Quartz fibres $\sim 2,000$ Channels

CRYSTAL
ELECTROMAGNETIC
CALORIMETER (ECAL)
 $\sim 76,000$ scintillating PbWO_4 crystals

HADRON CALORIMETER (HCAL)
Brass + Plastic scintillator $\sim 7,000$ channels



- ❖ COM energy 13 TeV
- ❖ Lumi $\sim 10^{34} \text{ cm}^{-2}\text{s}^{-1}$
- ❖ $\Delta p_T/p_T \sim 1\%$ when $p_T < 100 \text{ GeV}$
- ❖ transverse impact parameter resolution $O(10 - 100 \mu\text{m})$
- ❖ Vertex resolution $O(20 - 100 \mu\text{m})$
- ❖ Covering most of the 4π solid angle

$b \rightarrow sll$ angular analyses on CMS Run1

- Several $b \rightarrow sll$ angular analyses have been studied on CMS, such as $B^0 \rightarrow K^{*0} \mu^+ \mu^-$, $B^+ \rightarrow K^+ \mu^+ \mu^-$, $B^+ \rightarrow K^{*+} \mu^+ \mu^-$
- Run1 analyses based on 20fb^{-1} 8TeV data have been completed.
- Publications:
 - Angular analysis of $B^+ \rightarrow K^+ \mu^+ \mu^-$: Physical Review D, 98(2018), 112011
 - Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (A_{FB}) : Physics Letters B, 753(2016), 424-448
 - Angular analysis of $B^0 \rightarrow K^{*0} \mu^+ \mu^-$ (P_5') : Physics Letters B, 781(2018), 517-541
 - Angular analysis of $B^+ \rightarrow K^{*+} \mu^+ \mu^-$: Journal of High Energy Physics, 04(2021), 124

Run1 angular analysis: $B^+ \rightarrow K^+ \mu^+ \mu^-$

- The decay for the process $B^+ \rightarrow K^+ \mu^+ \mu^-$ can be described by $\cos \vartheta_\ell$ and q^2
- Differential decay rate formula:

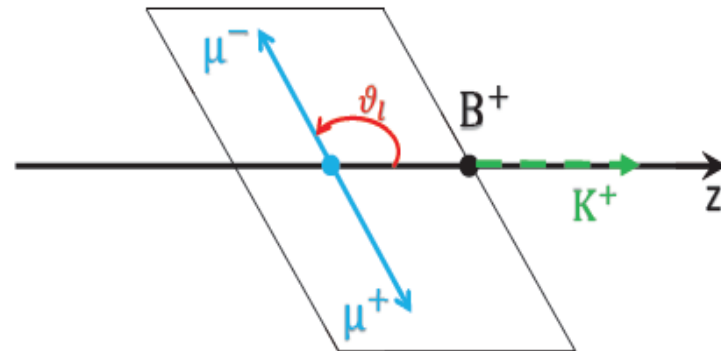
$$\frac{1}{\Gamma} \frac{d\Gamma[B^+ \rightarrow K^+ \mu^+ \mu^-]}{d \cos \theta_l} = \frac{3}{4} (1 - F_H) (1 - \cos^2 \theta_l) + \frac{1}{2} F_H + A_{FB} \cos \theta_l$$

$$0 \leq F_H \leq 3, A_{FB} \leq \min(1, F_H/2)$$

θ_l : the angle between the μ^+ (μ^-) and the K^+ (K^-) in the rest frame of the dimuon system.

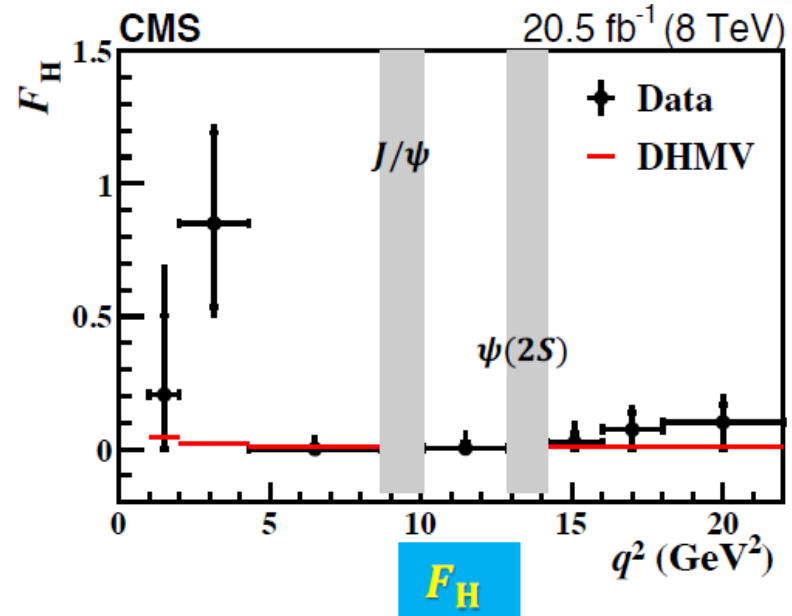
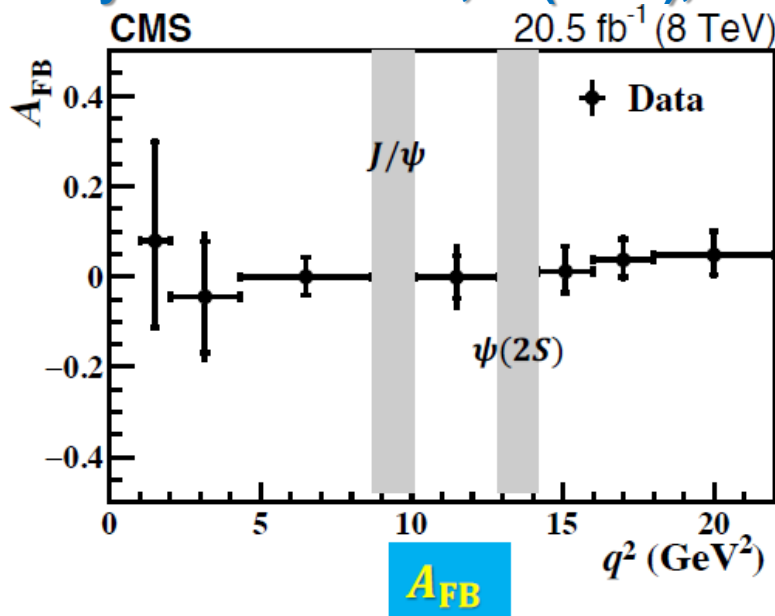
A_{FB} : $\mu^+ \mu^-$ forward-backward asymmetry.

F_H : a measure of the contribution from pseudoscalar, scalar and tensor amplitudes to the decay width Γ .



Run1 result: $B^+ \rightarrow K^+ \mu^+ \mu^-$

Physical Review D, 98(2018), 112011

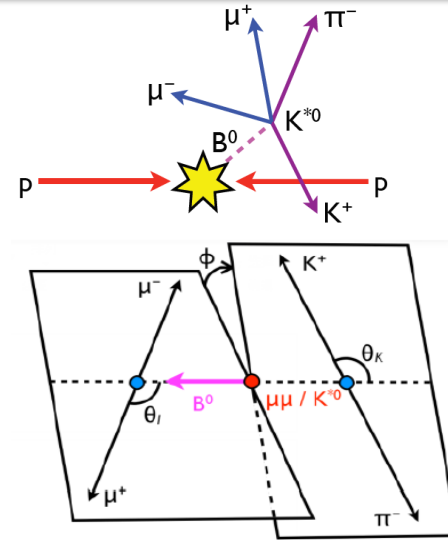


The events are measured in seven q^2 bins ranging from 1 to 22 GeV², 2286 signal events in total.

The measured A_{FB} and F_H are in agreement with the SM predictions within the uncertainties.

Run1 angular analysis: $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

- Final state: $K^+ \pi^- \mu^+ \mu^-$
- Fully described by the three angles (ϑ_ℓ , ϑ_K , ϕ) and q^2 .
- Robust SM predictions of several angular parameters, A_{FB} , F_L , P_1 and P_5' , are available.
- Angular observables are measured in q^2 bins from 1 to 19 GeV^2 .



❖ Differential decay rate :

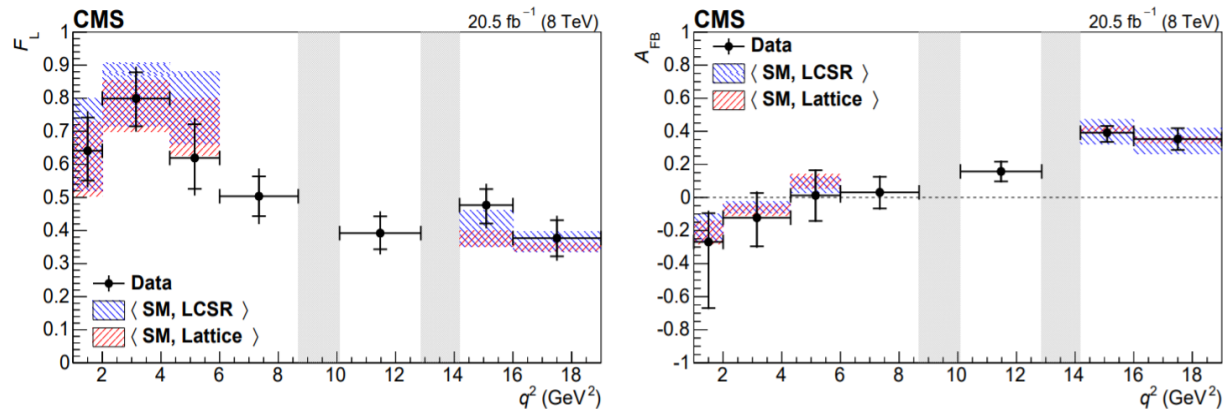
$$\frac{1}{\Gamma} \frac{d^3\Gamma}{d\cos\theta_K d\cos\theta_l dq^2} = \frac{9}{16} \left\{ \frac{2}{3} [F_S + A_S \cos\theta_K] (1 - \cos^2\theta_l) + (1 - F_S) [2F_L \cos^2\theta_K (1 - \cos^2\theta_l) + \frac{1}{2} (1 - F_L) (1 - \cos^2\theta_K) (1 + \cos^2\theta_l) + \frac{4}{3} A_{FB} (1 - \cos^2\theta_K) \cos\theta_l] \right\}.$$

$$\frac{1}{d\Gamma/dq^2} \frac{d^4\Gamma}{dq^2 d\cos\theta_l d\cos\theta_K d\phi} = \frac{9}{8\pi} \left\{ \frac{2}{3} [F_S + A_S \cos\theta_K] (1 - \cos^2\theta_l) + A_S^2 \sqrt{1 - \cos^2\theta_K} \sqrt{1 - \cos^2\theta_l} \cos\phi + (1 - F_S) [2F_L \cos^2\theta_K (1 - \cos^2\theta_l) + \frac{1}{2} (1 - F_L) (1 - \cos^2\theta_K) (1 + \cos^2\theta_l) + \frac{1}{2} P_1 (1 - F_L) (1 - \cos^2\theta_K) (1 - \cos^2\theta_l) \cos 2\phi + 2P_5' \cos\theta_K \sqrt{F_L} (1 - F_L) \sqrt{1 - \cos^2\theta_K} \sqrt{1 - \cos^2\theta_l} \cos\phi] \right\}$$

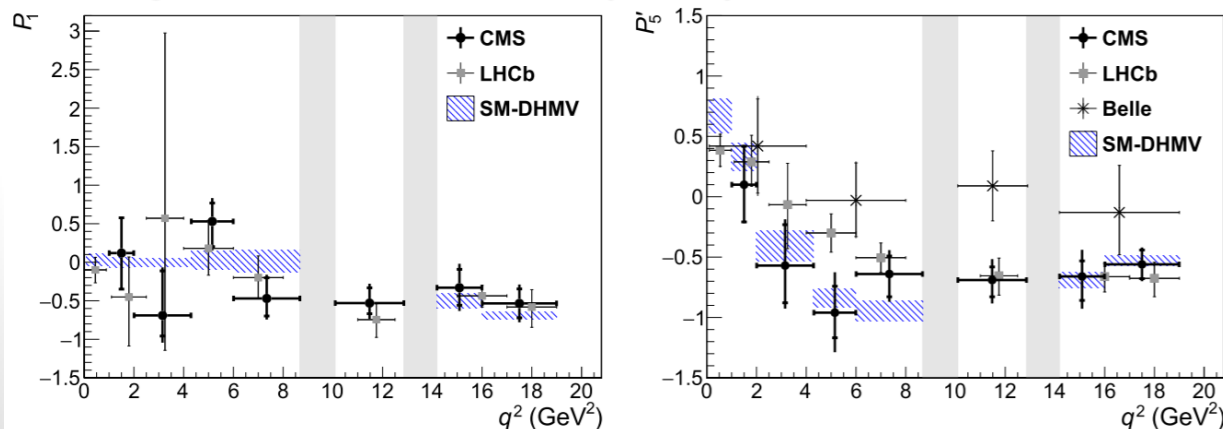
P-wave

Run1 result: $B^0 \rightarrow K^{*0} \mu^+ \mu^-$

Physics Letters B, 753(2016), 424-448



Physics Letters B, 781(2018), 517-541

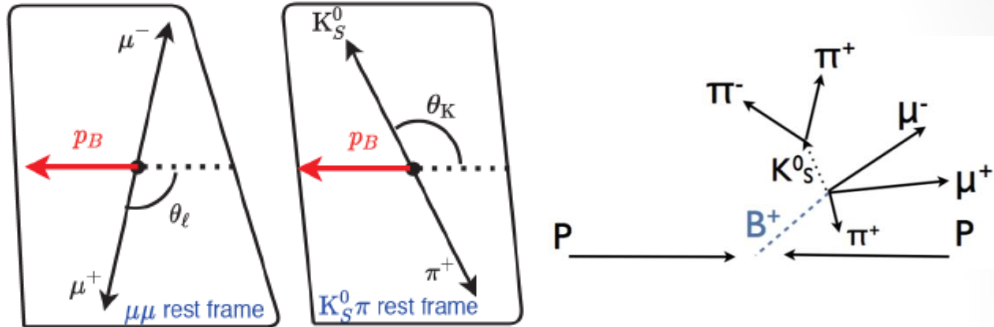


The events are fit in seven q^2 bins ranging from 1 to 19 GeV²

The measured (A_{FB} , F_H) and (P_1 , P_5') are in agreement with the SM predictions within the uncertainties.

Run1 angular analysis: $B^+ \rightarrow K^{*+} \mu^+ \mu^-$

- Final state: $\pi^+ \pi^- \pi^0 \mu^+ \mu^-$
- The decay can be fully described by the three angles (ϑ_ℓ , ϑ_K , ϕ) and q^2



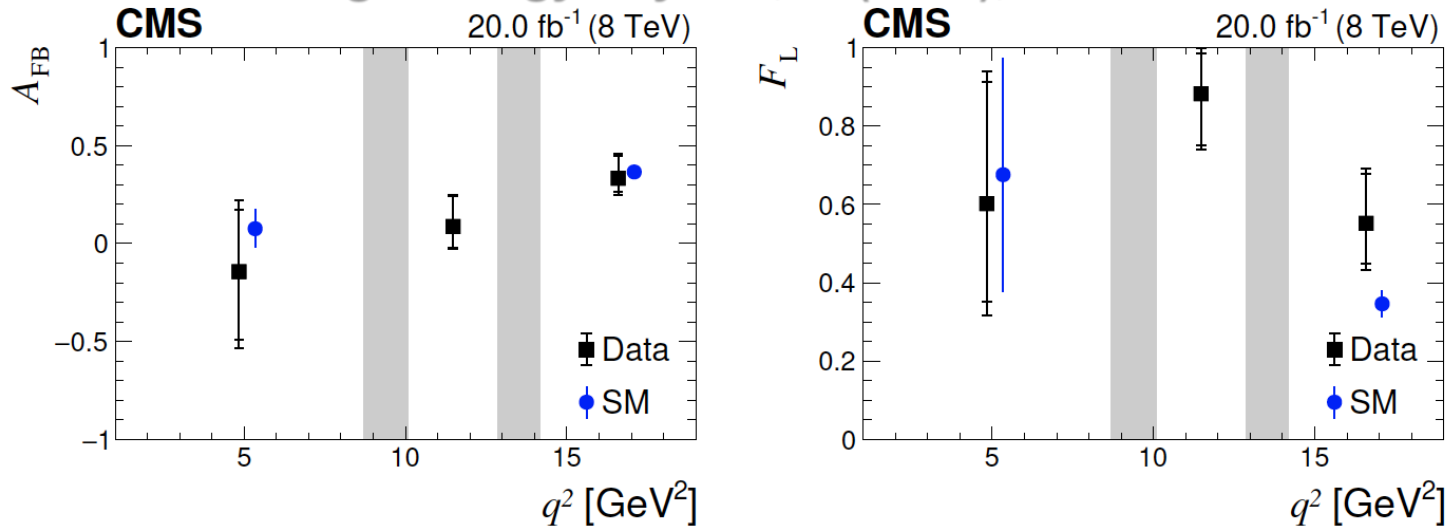
- Differential decay rate (integrate out ϕ):

$$\frac{1}{d\Gamma/dq^2} \frac{d^3\Gamma}{dq^2 d\cos\theta_l d\cos\theta_K} = \frac{9}{16} \left[\frac{1}{2} (1 - F_L) (1 - \cos^2\theta_K) (1 + \cos^2\theta_l) + 2F_L \cos^2\theta_K (1 - \cos^2\theta_l) + \frac{4}{3} A_{FB} (1 - \cos^2\theta_K) \cos\theta_l \right]$$

- Two observables, the forward-backward asymmetry of the muon (A_{FB}) and the longitudinal polarization fraction of the K^{*+} (F_L), are measured as a function of q^2

Run1 result: $B^+ \rightarrow K^{*+} \mu^+ \mu^-$

Journal of High Energy Physics, 04(2021), 124



q^2 (GeV ²)	Y_S	A_{FB}	F_L
1–8.68	22.1 ± 8.1	$-0.14^{+0.32}_{-0.35} \pm 0.17$	$0.60^{+0.31}_{-0.25} \pm 0.13$
10.09–12.86	25.9 ± 6.3	$0.09^{+0.16}_{-0.11} \pm 0.04$	$0.88^{+0.10}_{-0.13} \pm 0.05$
14.18–19	45.1 ± 8.0	$0.33^{+0.11}_{-0.07} \pm 0.05$	$0.55^{+0.13}_{-0.10} \pm 0.06$

A_{FB} and F_L in bins of q^2 are found to be consistent with a standard model prediction.

$b \rightarrow sll$ angular analyses on CMS Run2

- Based on 137fb^{-1} 13TeV data collected in CMS Run2, new $b \rightarrow sll$ angular analyses are being studied.
- The new angular analyses are: $B^0 \rightarrow K^{*0} \mu^+ \mu^-$, $B^+ \rightarrow K^+ \mu^+ \mu^-$, $B_s \rightarrow \Phi \mu^+ \mu^-$ based on Run2 data.
- The three analyses are ongoing.

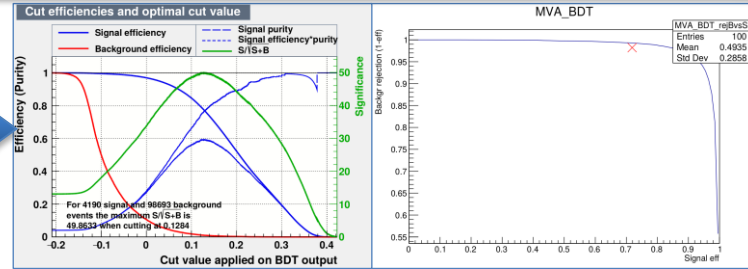
Updates of Run2 angular analyses

- **Luminosity**
 - From 20 fb^{-1} to 137 fb^{-1} , 7 times improvement
- **Selection**
 - Using new offline selection, improving the signal efficiency and the background rejection
- **Fitting**
 - Using new fitting PDF, stability improved
- **Statistic error**
 - Using new statistic error evaluation method

Example: $B^+ \rightarrow K^+ \mu^+ \mu^-$

Selection

Pre-selection;
Offline-selection (using AdaBoost BDT);

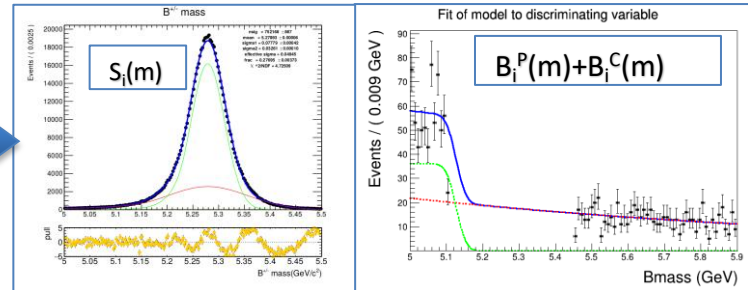


PDF components determination

Signal mass
shape
(using signal
MC);

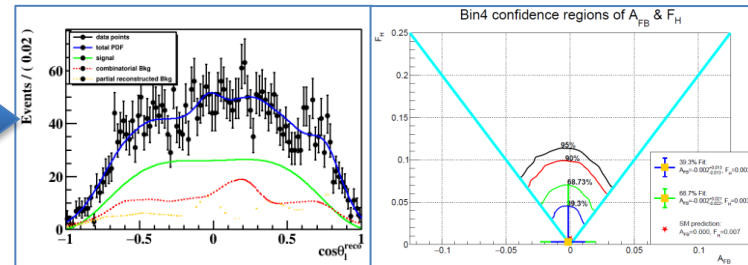
Bkg mass shape
(using data
sideband);

Bkg angular
shape
(using data
sideband);



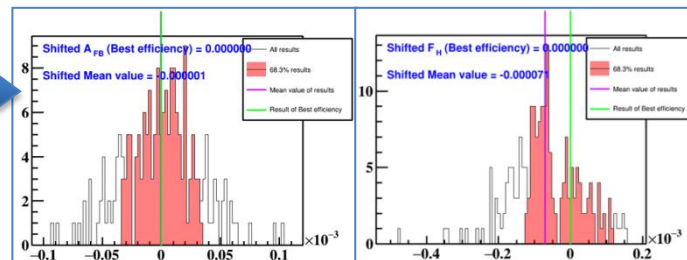
Fitting

2 dimensional fitting;
Statistic uncertainty (using 2D F-C method);



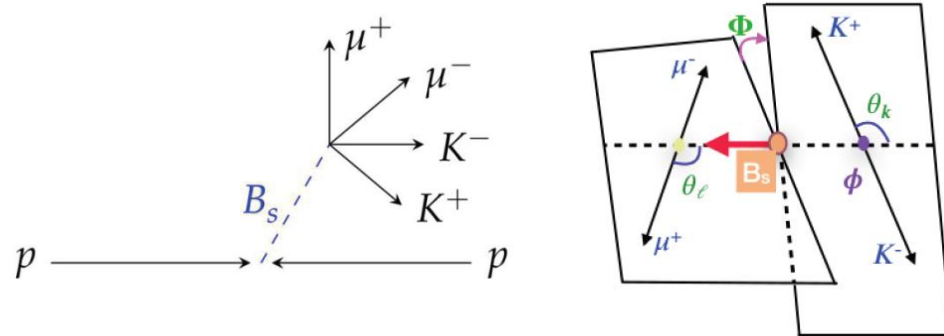
Systematic uncertainty

Fitting bias;
MC mismodelling;
PDF function forms ...



Run2 angular analysis: $B_s \rightarrow \Phi \mu^+ \mu^-$

- Final state: $K^+ K^- \mu^+ \mu^-$
- The decay can be described by the three angles (ϑ_ℓ , ϑ_K , ϕ) and q^2



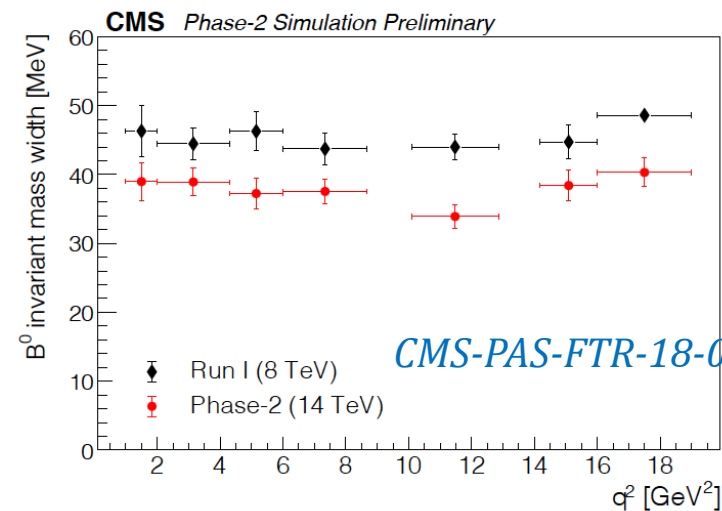
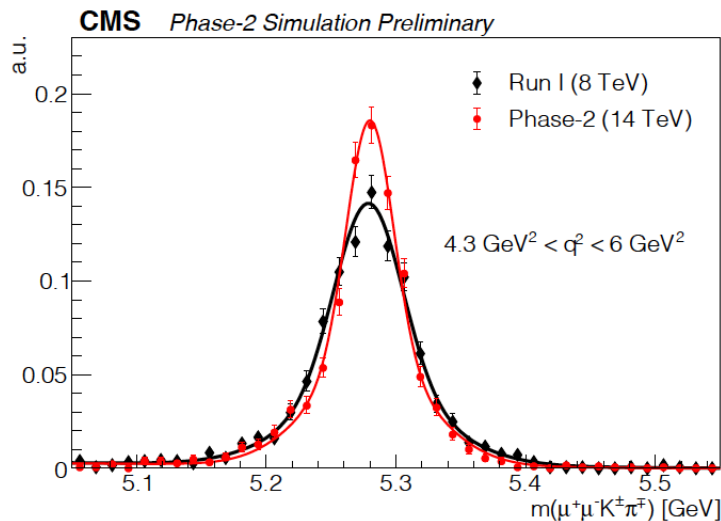
- Differential decay rate (integrate out ϕ):

$$\frac{1}{d\Gamma/dq^2} \frac{d^2\Gamma}{d\cos\theta_K d\cos\theta_l} = \frac{9}{16} \left(\frac{1}{2} (1 - F_L) \cdot (1 - \cos^2\theta_K) \cdot (1 + \cos^2\theta_l) \right. \\ \left. + 2F_L \cdot \cos^2\theta_K \cdot (1 - \cos^2\theta_l) + A_6 \cdot (1 - \cos^2\theta_K) \cdot \cos\theta_l \right)$$

- Two observables, the muon CP asymmetry (A_6) and the longitudinal polarization fraction (F_L), are measured as a function of q^2

Prospect of HL-LHC

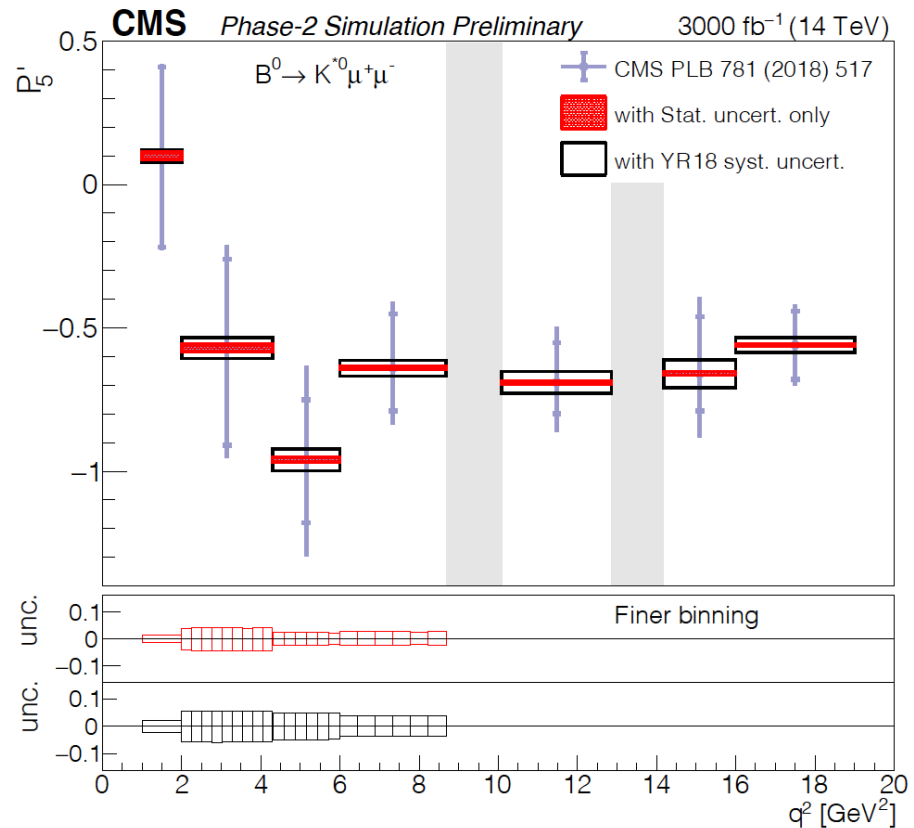
- Motivation: P'_5 discrepancy between experimental measurement and SM
- Run 1 results used as base line, the expected sensitivity of P'_5 with HL-LHC statistics at 3000 fb^{-1} was studied
 - Improved mass resolution with upgraded tracker detector
 - No changes in trigger performances and analysis strategy have been considered
 - Expected signal yield: $\sim 700\text{k}$ in full q^2 bin in 200 pileup scenario



CMS-PAS-FTR-18-033

Prospective uncertainties on P_5'

- Statistical uncertainty: scaled according to simulation yield
- Systematic uncertainty
 - Based on data control channel: scaled according to statistics
 - Others: scaled by factor of 2
- Total uncertainty: improve by up to a factor of 15 compared with Run 1 results
- Allow to split q^2 range in finer bins



Summary

- ❖ The angular analyses of FCNC B rare decays are important ways to search BSM phenomena.
- ❖ The three measurements done on CMS, $B^+ \rightarrow K^+ \mu^+ \mu^-$, $B^0 \rightarrow K^{*0} \mu^+ \mu^-$, and $B^+ \rightarrow K^{*+} \mu^+ \mu^-$, presented results consistent with standard model predictions.
- ❖ New angular analyses are ongoing, with larger amount of CMS datasets and updated procedure.
- ❖ More sensitive results can be expected on HL-LHC

END

Thank you!

Backup



Definition of observables $P_i^{(')}$

$$\frac{d^4\Gamma[\bar{B}^0 \rightarrow \bar{K}^{*0} \mu^+ \mu^-]}{dq^2 d\vec{\Omega}} = \frac{9}{32\pi} \sum_i I_i(q^2) f_i(\vec{\Omega}) \quad \text{and}$$

$$\frac{d^4\bar{\Gamma}[B^0 \rightarrow K^{*0} \mu^+ \mu^-]}{dq^2 d\vec{\Omega}} = \frac{9}{32\pi} \sum_i \bar{I}_i(q^2) f_i(\vec{\Omega}),$$

$$S_i = (I_i + \bar{I}_i) / \left(\frac{d\Gamma}{dq^2} + \frac{d\bar{\Gamma}}{dq^2} \right) \quad \text{and}$$

$$A_i = (I_i - \bar{I}_i) / \left(\frac{d\Gamma}{dq^2} + \frac{d\bar{\Gamma}}{dq^2} \right).$$

$$P_1 = \frac{2 S_3}{(1 - F_L)} = A_T^{(2)},$$

$$P_2 = \frac{2}{3} \frac{A_{FB}}{(1 - F_L)},$$

$$P_3 = \frac{-S_9}{(1 - F_L)},$$

$$P'_{4,5,8} = \frac{S_{4,5,8}}{\sqrt{F_L(1 - F_L)}},$$

$$P'_6 = \frac{S_7}{\sqrt{F_L(1 - F_L)}}.$$

S_i : CP averages; A_i : CP asymmetries;

$P_i^{(')}$: constructed observables for $B^0 \rightarrow K^{*0}$ leading form-factor uncertainty cancel