



# Strong-phase measurement of D decays at BESIII

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# Outline

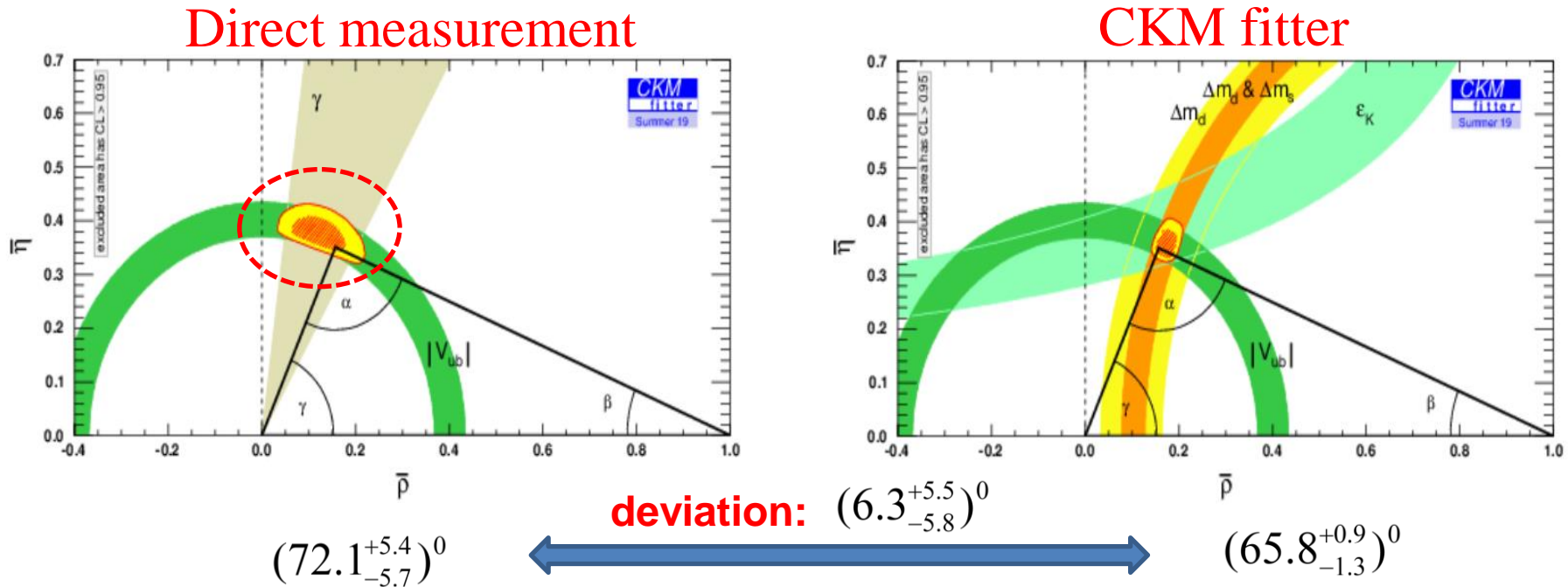
## ■ Strong-phase measurements at BESIII

- $D \rightarrow K_{S/L} h^+ h^-$  ( $h = \pi, K$ )
- $D \rightarrow K \pi \pi^0, K \pi \pi \pi$
- $D \rightarrow K \pi$

## ■ Summary

# Measurements of strong-phase parameters

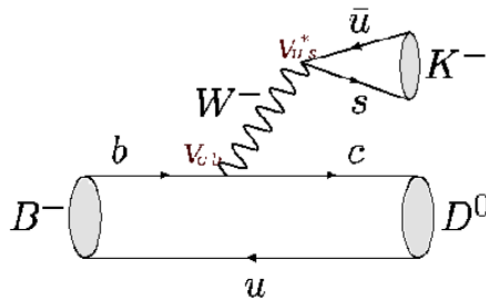
- Phase angle  $\gamma/\phi_3$  is the only CKM angle that can be measured in tree-level processes, in which the contribution of non-SM effects is expected to be small [JHEP 01(2014)051].
- Measurement of  $\gamma$  provides a benchmark of the SM with negligible theoretical uncertainty.



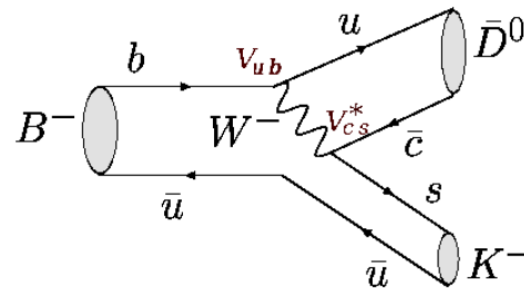
- An improved knowledge of the measurement of  $\gamma$  (**precision:  $\sim 1^\circ$ ,  $>5\sigma$** ) is important to further test the SM and probe for new physics.

# Measurements of strong-phase parameters

- Phase angle  $\gamma/\phi_3$  can be measured by studying the interference between  $B^- \rightarrow D^0 K^-$  and  $B^- \rightarrow \bar{D}^0 K^-$ .



$$A(B^- \rightarrow D^0 K^-) = A_B A_D$$



$$A(B^- \rightarrow \bar{D}^0 K^-) = A_B r_B e^{i(\delta_B - \gamma)} A_{\bar{D}}$$

where  $r_B$  is ratio of suppressed to favored amplitudes,  $\delta_B$  is the strong-phase difference between the favoured and suppressed amplitudes.

- Generally, three methods were proposed to measure  $\gamma/\phi_3$ :

- ✓ GLW <sup>[1]</sup>: via  $D^0 \rightarrow$  CP eigenstate,  $K^+ K^-$ ,  $\pi^+ \pi^-$ ,  $K_S^0 \pi^0$  etc.
- ✓ ADS <sup>[2]</sup>: via  $D^0 \rightarrow$  CF and DCS, such as  $K^+ \pi^-$ ,  $K^+ \pi^- \pi^0$ ,  $K^+ \pi^- \pi^- \pi^+$  etc.
- ✓ GGSZ <sup>[3]</sup>: via with  $D^0 \rightarrow$  Multi-body self-conjugate decays,  $K_S^0 \pi^+ \pi^-$  etc.

strong-phase parameters are key inputs.

[1] M. Gronau, D. London, Phys. Lett. B 253, 483 (1991); M. Gronau, D. Wyler, Phys. Lett. B 265, 172 (1991).

[2] D. Atwood, I. Dunietz and A. Soni, Phys. Rev. Lett. 78, 3257 (1997).

[3] A. Giri, Y. Grossman, A. Soffer and J. Zupan, Phys. Rev. D 68, 054018 (2003).

□ **ADS approach** [D. Atwood, I. Dunietz and A. Soni, Phys. Rev. Lett. 78, 3257 (1997) ] :

$B^- \rightarrow DK^-$  with  $D \rightarrow K n \pi$ , such as:  $D \rightarrow K \pi \pi \pi$ ,  $K \pi \pi^0$  etc.

$$\Gamma(B^- \rightarrow (K^+ 3\pi)_D K^-) \propto r_B^2 + (r_D^{K3\pi})^2 + 2R_{K3\pi} r_B r_D^{K3\pi} \cdot \cos(\delta_B + \delta_D^{K3\pi} - \gamma)$$

where  $R_{K3\pi}$  is coherence factor, and  $\delta_D^{K3\pi}$  is averaged strong-phase difference.

□ **GGSZ approach** [A. Giri, Y. Grossman, A. Soffer and J. Zupan, Phys. Rev. D 68, 054018] :

$B^- \rightarrow DK^-$  with  $D \rightarrow$  Multi-body self-conjugate decays,  $K_{s/L}^0 \pi^+ \pi^-$ ,  $K_{s/L}^0 K^+ K^-$  etc.

**Amplitude:**  $d\Gamma(B^\pm \rightarrow \tilde{D}^0 K^\pm) = |f_D(m_\pm^2, m_\mp^2)|^2 + r_B^2 |f_D(m_\mp^2, m_\pm^2)|^2 + 2r_B |f_D(m_\pm^2, m_\mp^2)| |f_D(m_\mp^2, m_\pm^2)|$   
 $\times [\cos \Delta \delta_D \cos(\delta_B \pm \phi_3) + [\sin \Delta \delta_D \sin(\delta_B \pm \phi_3)],$

$$c_i = \frac{1}{\sqrt{F_i F_{-i}}} \int_i |f_D(m_+^2, m_-^2)| |f_D(m_-^2, m_+^2)| \cos[\Delta \delta_D(m_+^2, m_-^2)] dm_+^2 dm_-^2$$

$$s_i = \frac{1}{\sqrt{F_i F_{-i}}} \int_i |f_D(m_+^2, m_-^2)| |f_D(m_-^2, m_+^2)| \sin[\Delta \delta_D(m_+^2, m_-^2)] dm_+^2 dm_-^2$$

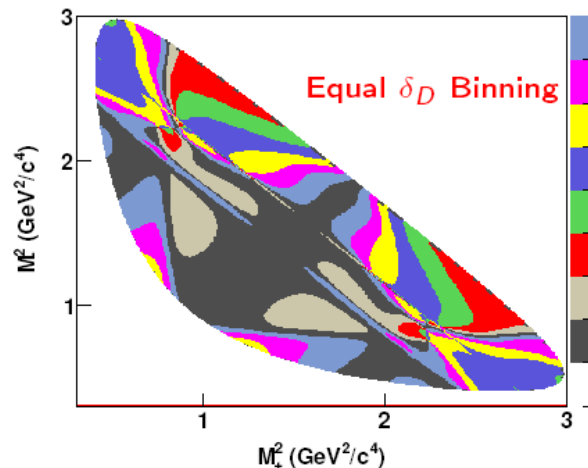
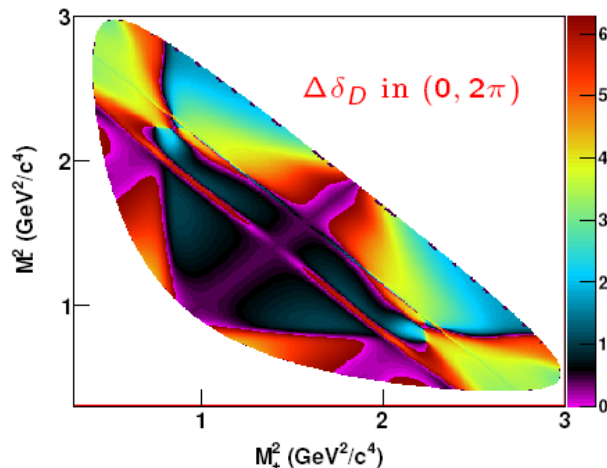
*strong-phase difference:*

$$\Delta \delta_D = \delta_D(m_+^2, m_-^2) - \delta_D(m_-^2, m_+^2)$$

where  $c_i$  and  $s_i$  are the amplitude-weighted averages of  $\cos \Delta \delta_D$  and  $\sin \Delta \delta_D$  over each Dalitz-plot bin.

• **Equal  $\Delta \delta_D$  binning scheme (N=8):**

$$2\pi(i - 3/2)/\mathcal{N} < \Delta \delta_D(m_+^2, m_-^2) < 2\pi(i - 1/2)/\mathcal{N}$$

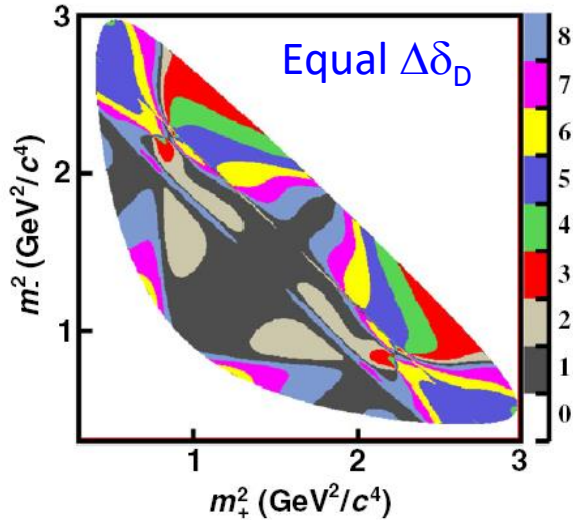


*The strong-phase difference ( $\Delta \delta_D$ ) calculated based on "Babar 2008" Model.*  
 Phys. Rev. D 78, 034023 (2008)

**Model-independent:**  
 Bondar and Poluektov  
 [EPJC47, 347(2006), EPJC55, 51(2008)].

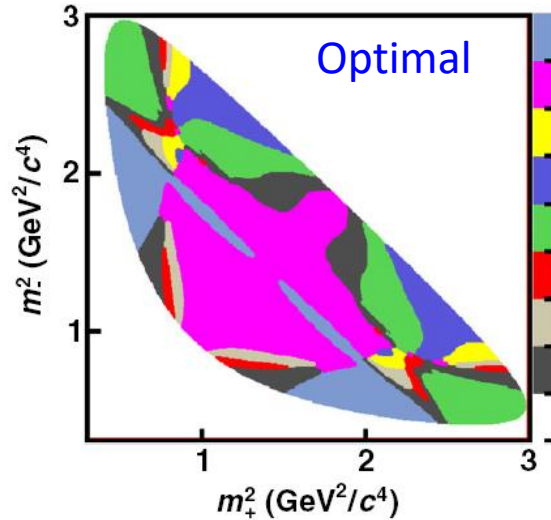
# Strong-phase parameters in $D \rightarrow K_S^0 \pi^+ \pi^-$

## □ Three typical binning schemes [Phys. Rev. D 82, 112006 (2010)]



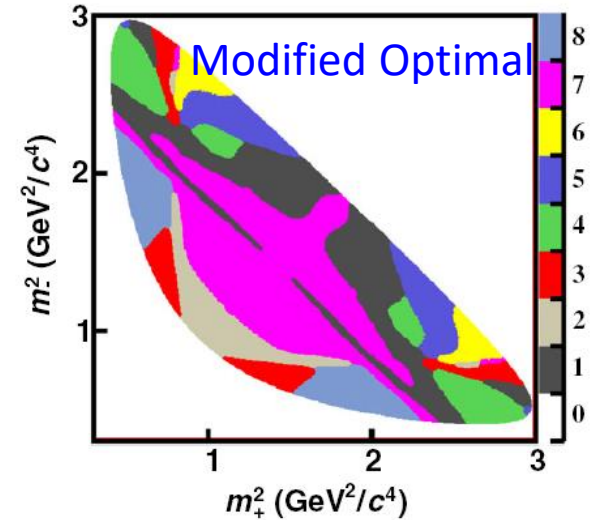
DD-mixing<sup>1</sup>,  $\beta$  measurements<sup>2</sup>

[minimum variation in  $\Delta\delta_D$ ]



$\gamma$  measurements<sup>3,4</sup>

[Optimized sensitivity,  
no background included]



$\gamma$  in Low yields

[Optimization including  
backgrounds]

- ✓ “BaBar K-matrix”  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$  model as in Ref. [Phys. Rev. D 78, 034023 (2008)].
- ✓ It should be noted that although the choice of binning is model-dependent, however, a poor choice of model results only in a loss of precision, instead of bias in measuring  $\gamma/\phi_3$ .

[1] R. Aaij et al. (LHCb Collaboration), Phys. Rev. Lett. 122, 231802 (2019); JHEP 04(2016) 033.

[2] V. Vorobyev et al. (Belle Collaboration), Phys. Rev. D 94, 052004 (2016).

[3] R. Aaij et al. (LHCb Collaboration), Phys. Lett. B 718, 43 (2012); JHEP 10 (2014) 097; JHEP 06 (2016) 131; JHEP 08 (2018) 176.

[4] H. Aihara et al. (Belle Collaboration), Phys. Rev. D 85, 112014 (2012).

# The Quantum Correlated $D\bar{D}$ meson pairs

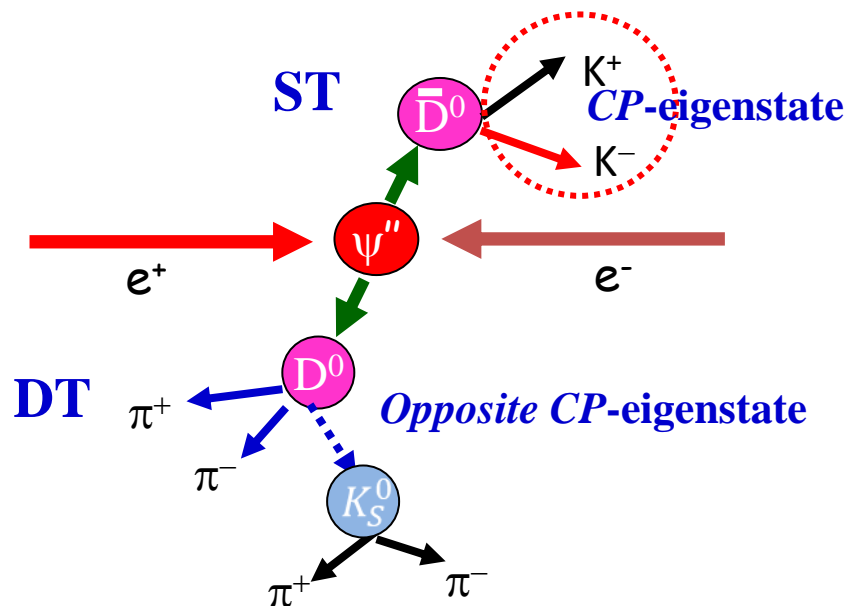
□  $\psi(3770)$  is a spin  $-1$  state and therefore the amplitude of  $\psi(3770) \rightarrow D^0 \bar{D}^0$ :

$$(|D^0\rangle|\bar{D}^0\rangle - |\bar{D}^0\rangle|D^0\rangle)/\sqrt{2} \quad [\text{anti-symmetric wave function}]$$

The amplitude for two D mesons to decay to states  $F$  and  $G$  is [PRD68, 033003 (2003)]:

$$\Gamma(F|G) = \Gamma_0 [A_F^2 \bar{A}_G^2 + \bar{A}_F^2 A_G^2 - 2R_F R_G A_F \bar{A}_F A_G \bar{A}_G \cos[\delta_D^F - \delta_D^G]]$$

□ Hence, the coherence factors  $R_F$ , the strong-phase difference  $\delta_D^F$ , can be extracted based on the study of the quantum correlated DD meson pairs.



The DT mode  $K^+K^-$  vs.  $K_S^0\pi^+\pi^-$  is selected as an example.

✓ **Single tag (ST) samples:**

decay products of only one D meson are reconstructed

✓ **Double tag (DT) samples:**

decay products of both D mesons are reconstructed

✓ **Some typical reconstructed D decay modes**

Tag group	
Flavor	$K^+\pi^-, K^+\pi^-\pi^0, K^+\pi^-\pi^-\pi^+, K^+e^-\bar{\nu}_e$
CP-even	$K^+K^-, \pi^+\pi^-, K_S^0\pi^0\pi^0, K_L^0\pi^0, \pi^+\pi^-\pi^0$
CP-odd	$K_S^0\pi^0, K_S^0\eta, K_S^0\omega, K_S^0\eta', K_L^0\pi^0\pi^0$
Mixed-CP	$K_S^0\pi^+\pi^-$

✓ **Expected events for  $K_S\pi^+\pi^-$  vs. CP-eigenstate:**

$$f_{CP\pm} = \frac{1}{\sqrt{2}} [f_D(m_+^2, m_-^2) \pm f_D(m_-^2, m_+^2)]$$

$$M_i^\pm = h_{CP\pm}(K_i \pm 2c_i \sqrt{K_i K_{-i}} + K_{-i})$$



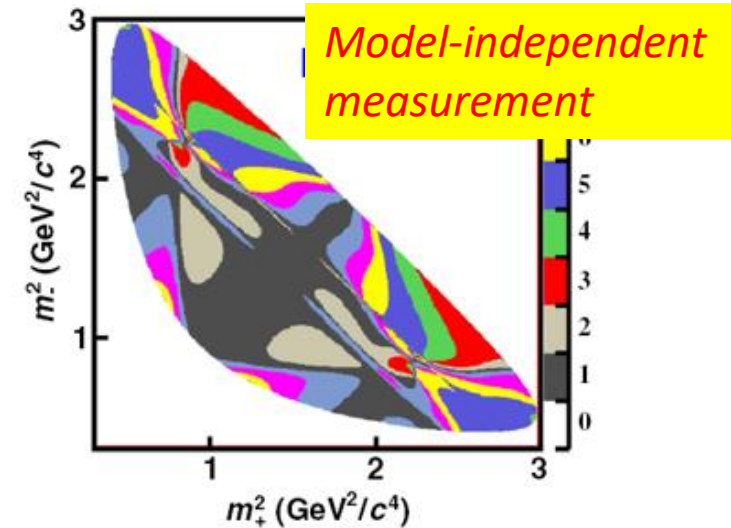
# Strong-phase parameters ( $c_i, s_i$ ) in $D \rightarrow K_S^0 \pi^+ \pi^-$

□ For CP-tagged  $K_S^0 \pi^+ \pi^-$ , the amplitude:

$$f_{CP\pm} = \frac{1}{\sqrt{2}} [f_D(m_+^2, m_-^2) \pm f_D(m_-^2, m_+^2)]$$

*The expected number of events in DP bins:*

➔  $M_i^\pm = \frac{s^\pm}{2s_f} (K_i \pm 2c_i \sqrt{K_i K_{-i}} + K_{-i})$  for  $D^0 \rightarrow K_S^0 \pi^+ \pi^-$



□ For  $K_S^0 \pi^+ \pi^-$  .vs.  $K_S^0 \pi^+ \pi^-$ , its amplitude is expressed by:

$$f(m_+^2, m_-^2, m'^2_+, m'^2_-) = \frac{f_D(m_+^2, m_-^2) f_D(m'^2_+, m'^2_-) - f_D(m'^2_+, m'^2_-) f_D(m_-^2, m_+^2)}{\sqrt{2}}$$

*The expected number of events in DP bins:*

➔  $M_{ij}^\pm = \frac{N_{D\bar{D}}}{2s_f^2} (K_i K_{-j} + K_{-i} K_j - 2\sqrt{K_i K_{-j} K_{-i} K_j} (c_i c_j + s_i s_j))$



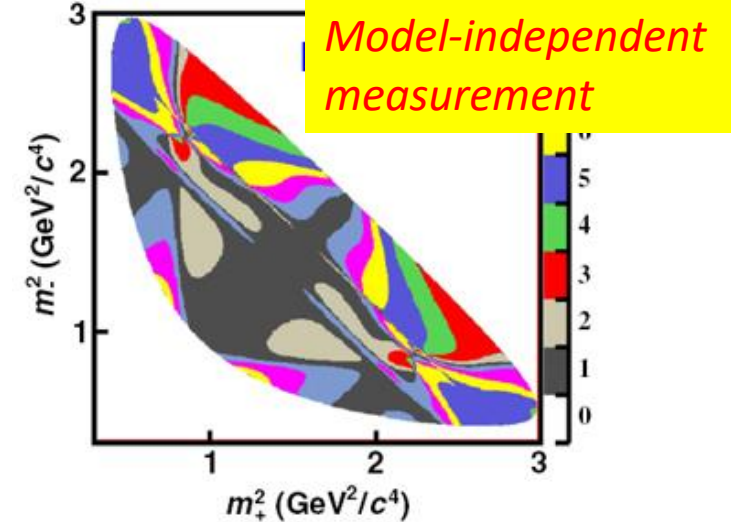
# Strong-phase parameters ( $c'_i, s'_i$ ) in $D \rightarrow K_L^0 \pi^+ \pi^-$

□ For CP-tagged  $K_L^0 \pi^+ \pi^-$ , the amplitude:

$$f_{CP\pm} = \frac{1}{\sqrt{2}} [f_D(m_+^2, m_-^2) \mp f_D(m_-^2, m_+^2)]$$

The expected number of events in DP bins:

➔  $M_i'^{\pm} = \frac{s_{\pm}^{\pm}}{2S_f} (K_i' \mp 2c_i' \sqrt{K_i' K_{-i}'} + K_{-i}') \text{ for } D^0 \rightarrow K_L^0 \pi^+ \pi^-$



□ For  $K_S^0 \pi^+ \pi^-$  .vs.  $K_L^0 \pi^+ \pi^-$ , its amplitude is expressed by:

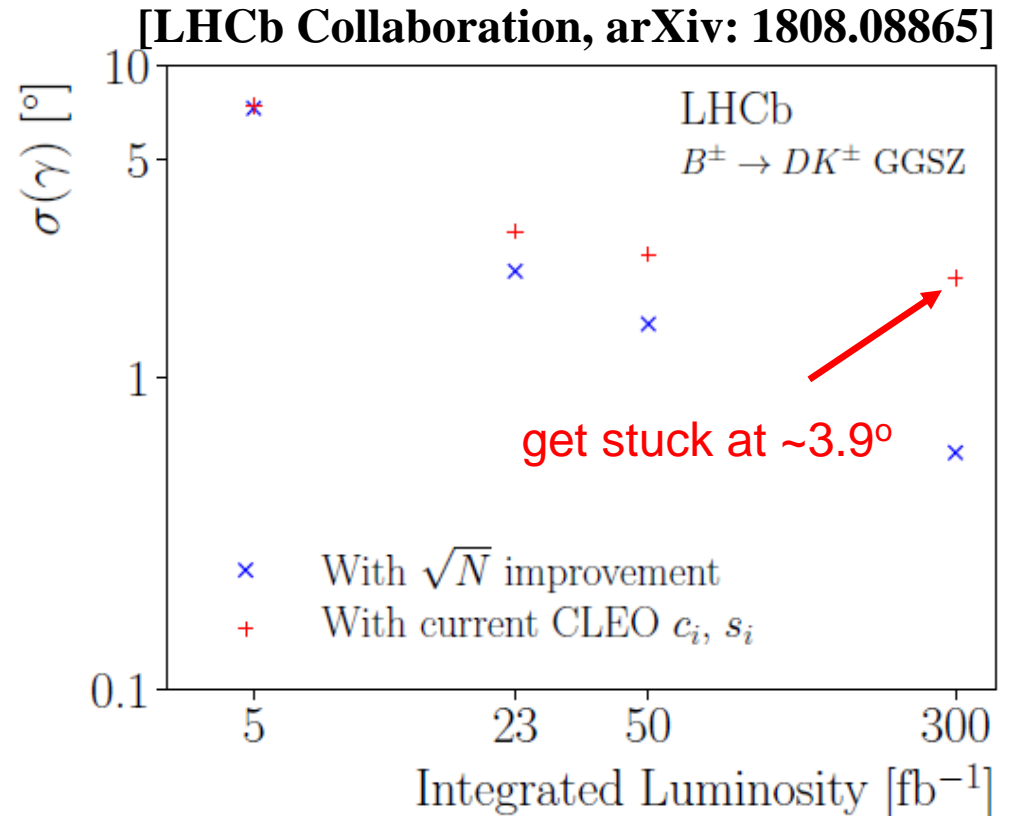
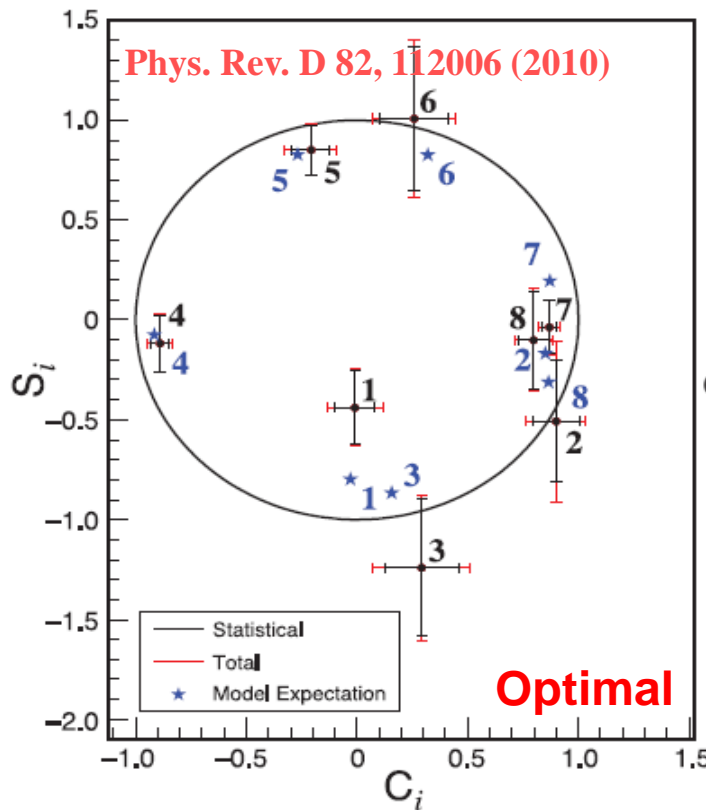
$$f(m_+^2, m_-^2, m_+^{\prime 2}, m_-^{\prime 2}) = \frac{f_D(m_+^2, m_-^2) f_D(m_-^{\prime 2}, m_+^{\prime 2}) + f_D(m_+^{\prime 2}, m_-^{\prime 2}) f_D(m_-^2, m_+^2)}{\sqrt{2}}$$

The expected number of events in DP bins:

➔  $M_{ij}'^{\pm} = \frac{N_{D\bar{D}}}{2S_f^2} (K_i K_{-j}' + K_{-i} K_j' + 2\sqrt{K_i K_{-j}' K_{-i} K_j'} (c_i c_j' + s_i s_j'))$

# Strong-phase parameters in $D \rightarrow K_S^0 \pi^+ \pi^-$

✓ Results of  $c_i$  and  $s_i$  in optimal binning from CLEO experiments.



✓ The systematic uncertainty in measurement of  $\gamma$  due to the input of strong-phase parameters is  $3.9^\circ$  for optimal binning. The overall sensitivity of  $\gamma$  is systematically limited to  $\sim 3.9^\circ$  for model-independent GGSZ approach.

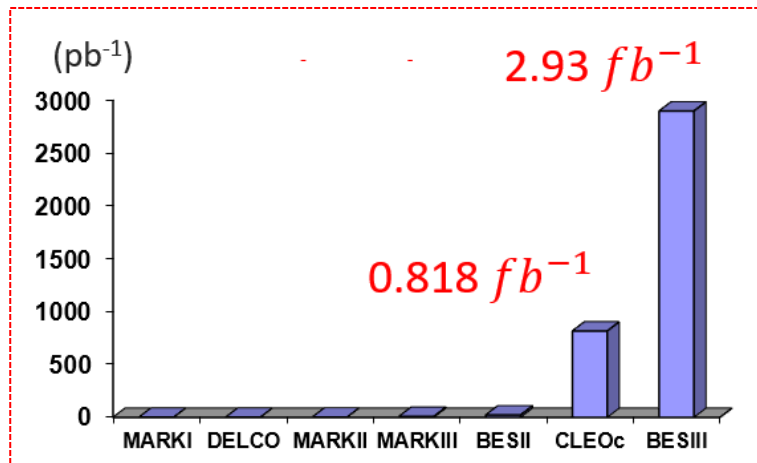
✓ Therefore, improved measurements in  $c_i$  &  $s_i$  from BESIII are essential for degree-level precision of measuring  $\gamma$  via model-independent GGSZ approach.

# $\psi(3770) \rightarrow D^0 \bar{D}^0$ samples at BESIII

- BESIII is the only machine running at  $\tau$ -charm energy region. The quantum-correlated studies are key to constrain the  $\gamma/\phi_3$  measurement at LHCb upgrades 1(2) and Belle II experiments.

**2.93/fb @ 3.773 GeV**

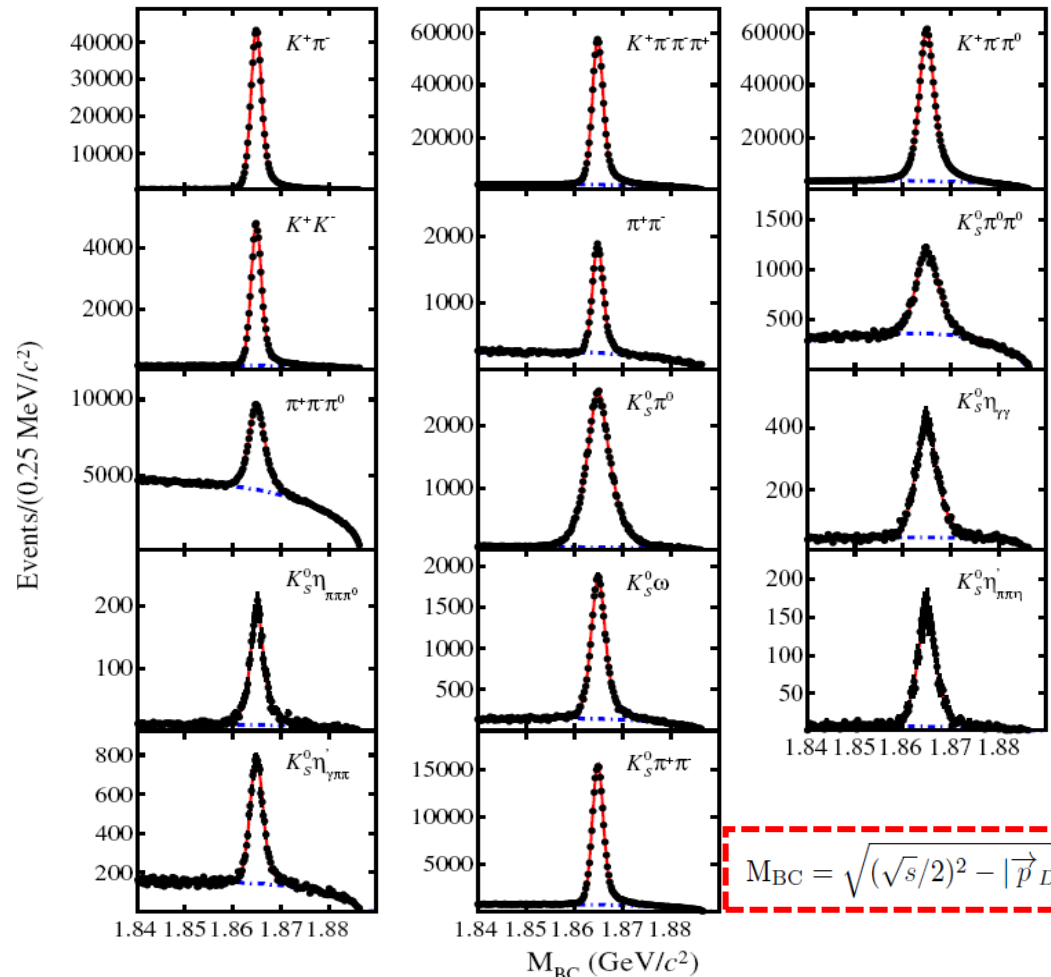
The largest  $\psi(3770)$  data sample



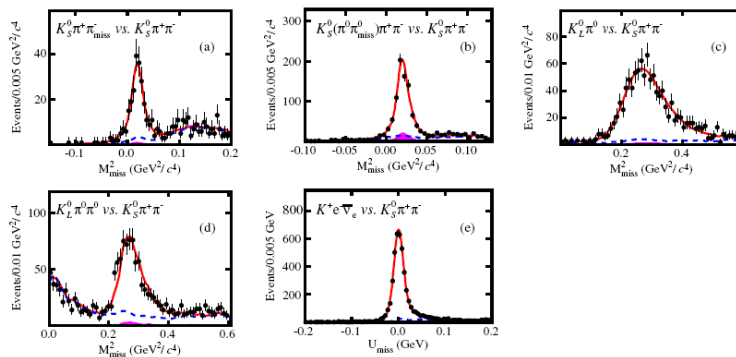
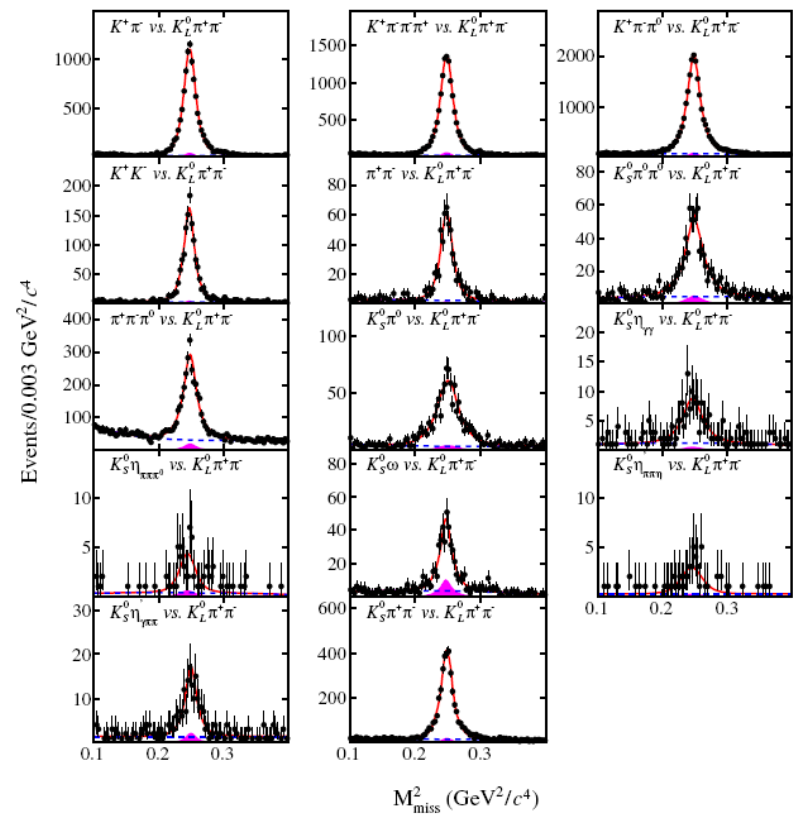
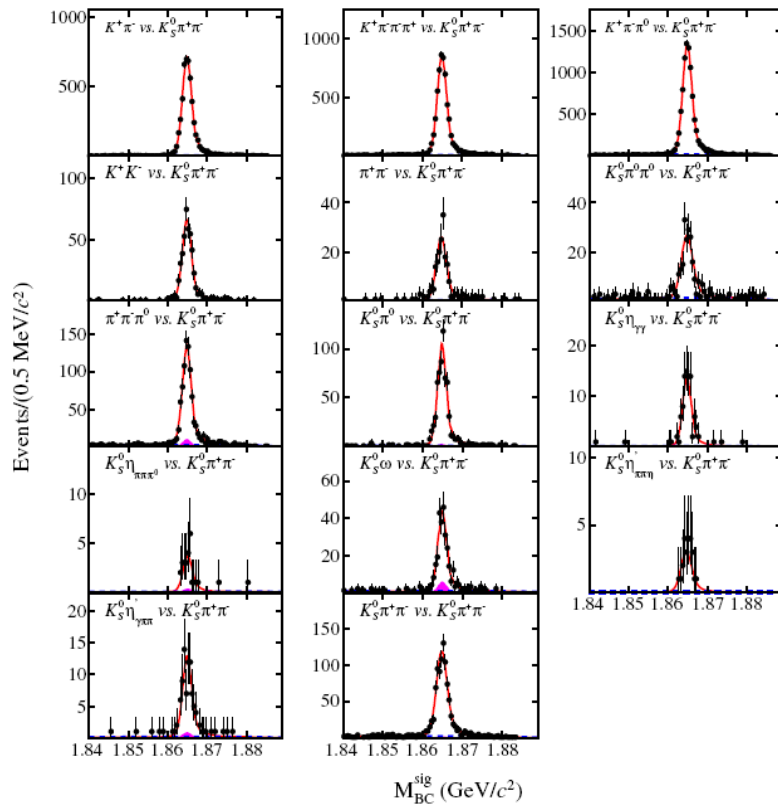
- ✓ Good performance of BESIII detector: high tracking & PID efficiencies; high purity samples.

Phys. Rev. Lett. 124, 241802(2020)  
Phys. Rev. D 101, 112002(2020)

**beam-constrained mass distributions in data**



# DT events for $D \rightarrow K_{S/L}^0 \pi^+ \pi^-$ in data



**DT events at BESIII:**

*CP – eigenstate vs.  $K_S^0 \pi^+ \pi^-$ :  $5.3 \times \text{CLEO}$*

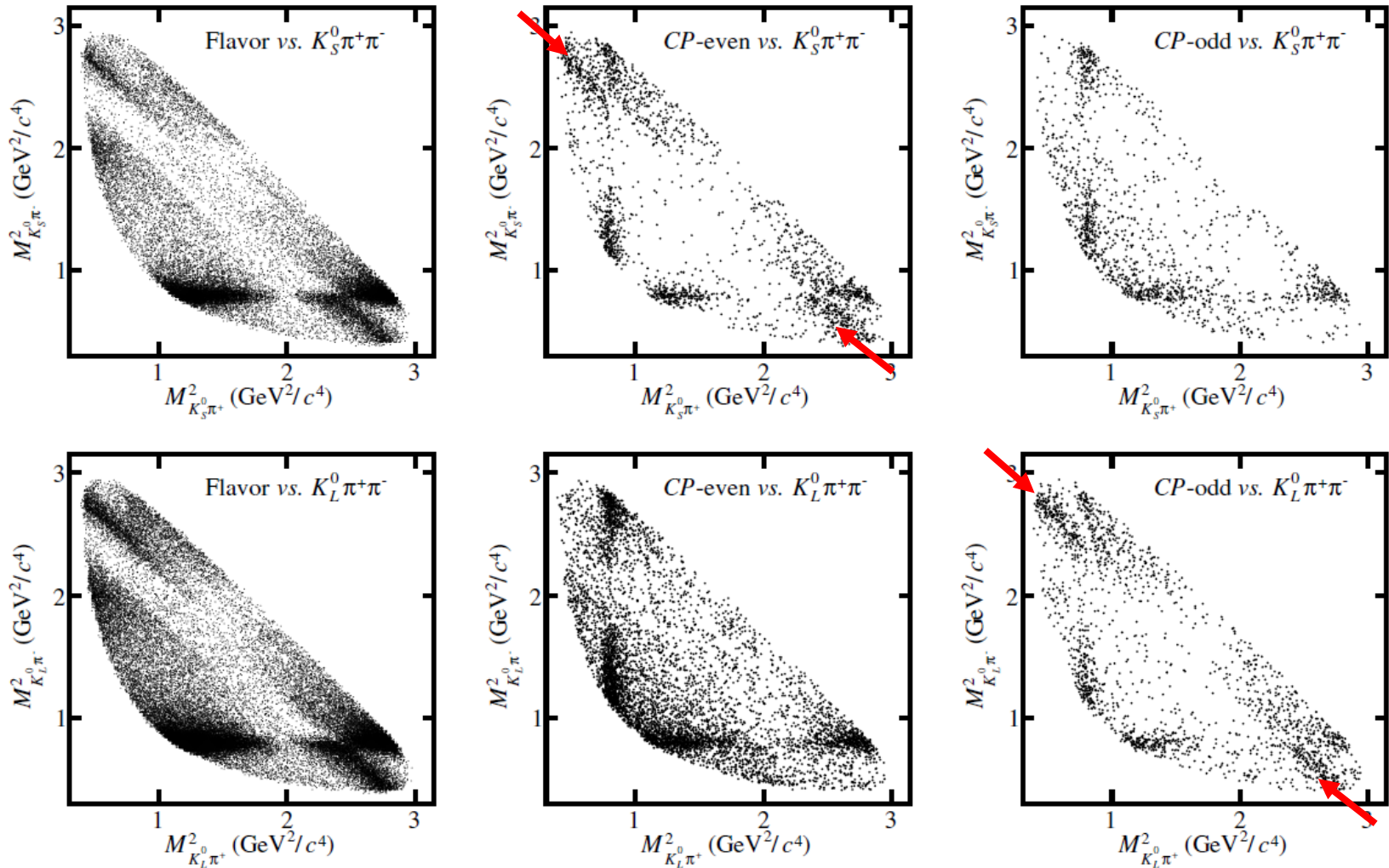
*CP – eigenstate vs.  $K_L^0 \pi^+ \pi^-$ :  $9.2 \times \text{CLEO}$*

*$K_S^0 \pi^+ \pi^-$  vs.  $K_S^0 \pi^+ \pi^-$ :  $3.9 \times \text{CLEO}$*

*$K_L^0 \pi^+ \pi^-$  vs.  $K_S^0 \pi^+ \pi^-$ :  $2.9 \times \text{CLEO}$*

# Dalitz plots for $D \rightarrow K_{S/L} \pi^+ \pi^-$ observed in data

❑ Effect of quantum correlation is immediately seen in Dalitz plots.

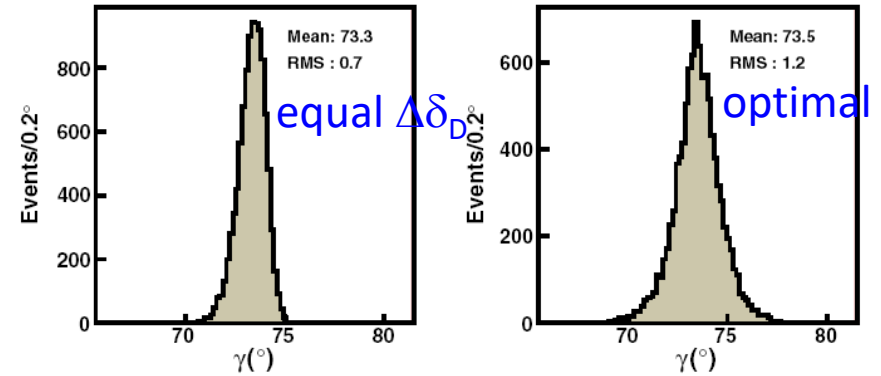
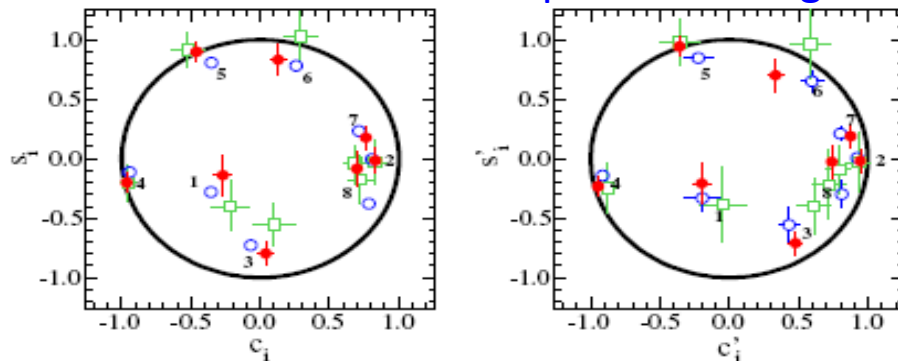
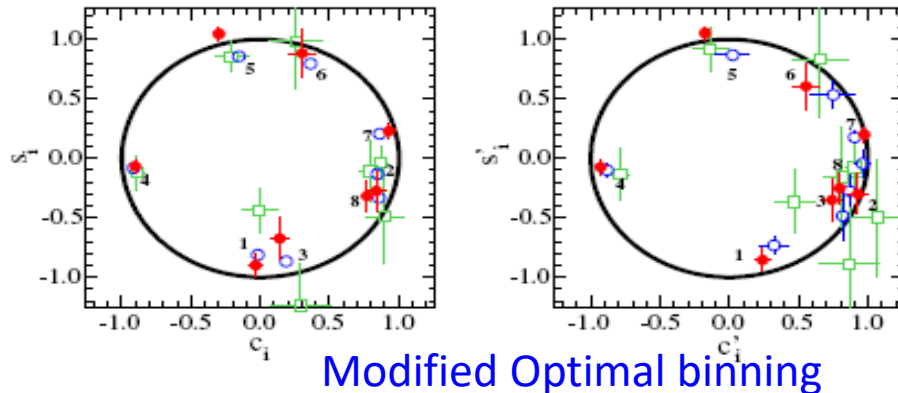
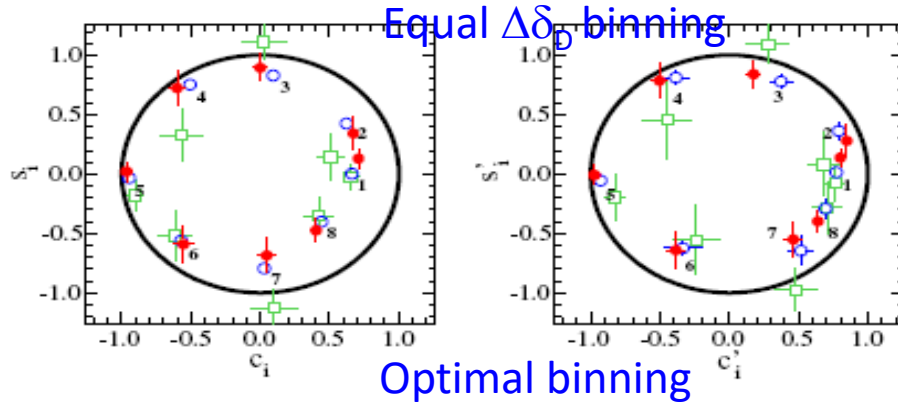


The CP-odd component  $K_S^0 \rho(770)^0$  is visible in CP-even tagged  $K_S^0 \pi^+ \pi^-$  decays, but is absent in CP-odd tagged  $K_S^0 \pi^+ \pi^-$  decays.

# The strong-phase parameters ( $c_i^{(\prime)}$ , $s_i^{(\prime)}$ ) in $D \rightarrow K_{S/L}^0 \pi^+ \pi^-$

✓ The  $c_i^{(\prime)}$  and  $s_i^{(\prime)}$  measured in this work, the expected results and the CLEO results.

✓ The expected uncertainty for  $\gamma$  with  $B^- \rightarrow DK^-$ ,  $D \rightarrow K_S^0 \pi^+ \pi^-$ .



Phys. Rev. Lett. 124, 241802(2020)  
Phys. Rev. D 101, 112002(2020)

✓ On average a factor of  $\sim 2.5$  ( $2.0$ ) more precise for  $c_i$  ( $s_i$ ) and  $\sim 2.8$  ( $2.2$ ) more precise for  $c'_i$  ( $s'_i$ ) than CLEO.

✓ The associated uncertainty on  $\gamma/\phi_3$  is expected to be roughly a factor of three smaller than that from CLEO analysis.

✓ The improved precision on  $c'_i$  and  $s'_i$  are important for Belle-II in  $\gamma/\phi_3$  measurement with  $B^- \rightarrow DK^-$ ,  $D \rightarrow K_L^0 \pi^+ \pi^-$ .



# The strong-phase parameters in $D \rightarrow K_{S/L}^0 K^+ K^-$

[Phys. Rev. D 102, 152008(2020)]

Equal  $\Delta\delta_D$  binning

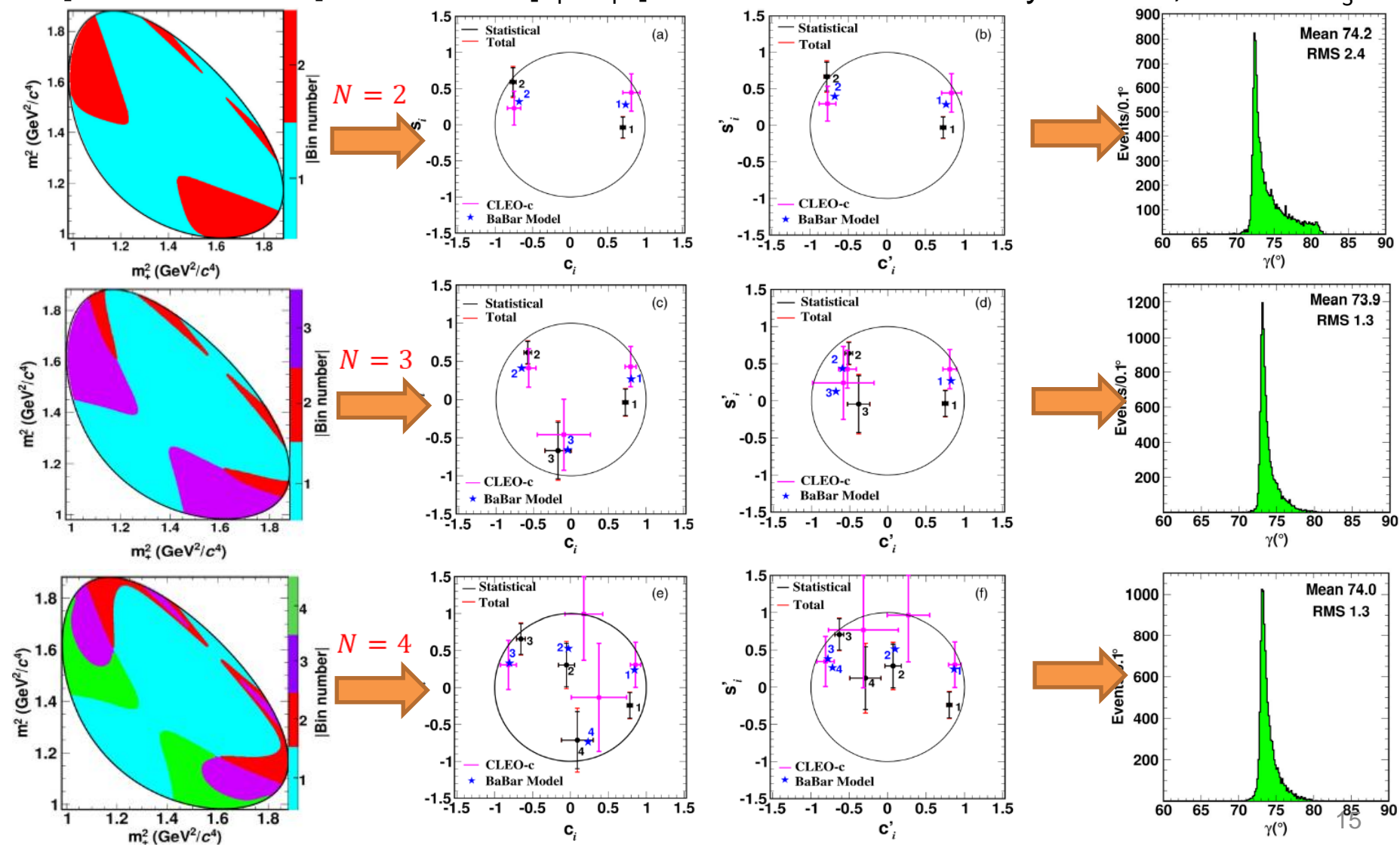
[PRD78,034023(2008)]

Strong-phase parameters

$[c_i^{(i)}, s_i^{(i)}]$  measured in data

expected uncertainties on  $\gamma$

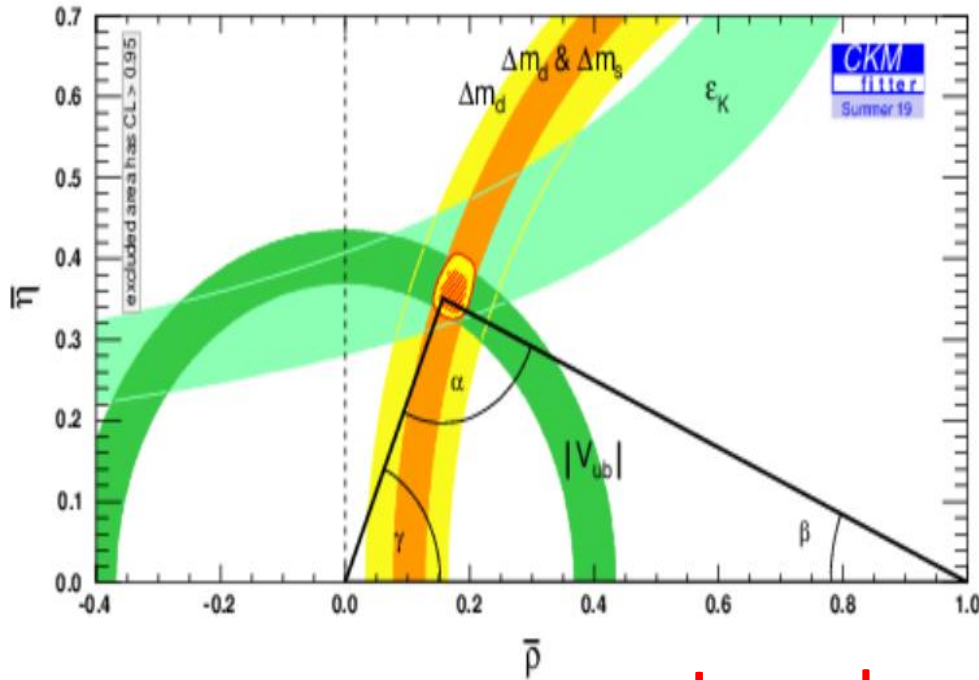
by  $B^- \rightarrow DK^-$ , with  $D \rightarrow K_S K^+ K^-$



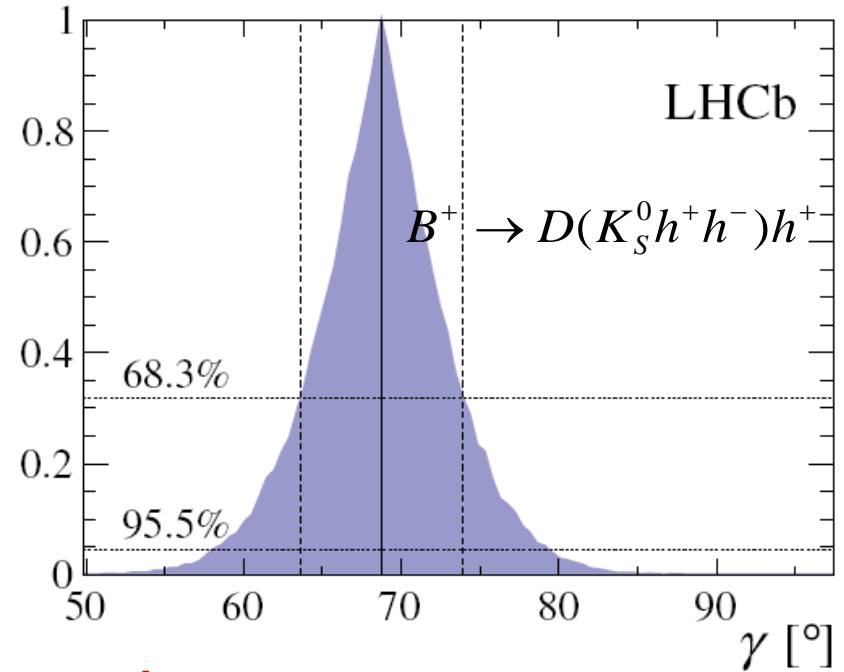


# $\gamma/\phi_3$ measurements with inputs of strong-phase parameters from $D \rightarrow K_S^0 h^+ h^-$

CKM fitter



JHEP02(2021)169



Indirect:  $(65.8^{+0.9}_{-1.3})^0$  In good agreement  $\longleftrightarrow \gamma/\phi_3 = (68.7^{+5.2}_{-5.1})^0$

The uncertainty from strong phases is roughly  $1^\circ$ , which is significantly improved.

□ Previous LHCb result,  $\gamma = (80 \pm 10)^\circ$  [JHEP08(2018)176] with CLEO inputs.

# Measurements of coherence factors in $D \rightarrow K^- \pi^+ \pi^+ \pi^-$ and $D \rightarrow K^- \pi^+ \pi^0$

[JHEP05(2021)164]

- For  $D \rightarrow K^- n \pi$ , the coherence factor  $R_S$ , the amplitude ratio  $r_D$  and strong-phase difference  $\delta_D$  between the CF and DCS amplitude averaged over phase space:

$$R_S e^{-i\delta_D^S} = \frac{\int \mathcal{A}_S^*(x) \mathcal{A}_{\bar{S}}(x) dx}{A_S A_{\bar{S}}} \quad \text{and} \quad r_D^S = A_{\bar{S}}/A_S \quad \text{A(s) is the decay amplitude of } D \rightarrow K^- n \pi.$$

- The amplitude for two D mesons decay to states S(single tag) and T(double tag):

$$\Gamma(S|T) = A_S^2 A_T^2 \left[ (r_D^S)^2 + (r_D^T)^2 - 2R_S R_T r_D^S r_D^T \cos(\delta_D^T - \delta_D^S) \right]$$

Flavour	Like sign	$K^- \pi^+ \pi^+ \pi^-, K^- \pi^+ \pi^0, K^- \pi^+$
	Opposite sign	$K^+ \pi^- \pi^- \pi^+, K^+ \pi^- \pi^0, K^+ \pi^-$
CP	Even	$K^+ K^-, \pi^+ \pi^-, K_S^0 \pi^0 \pi^0, K_L^0 \pi^0, K_L^0 \omega, \pi^+ \pi^- \pi^0$
	Odd	$K_S^0 \pi^0, K_S^0 \eta, K_S^0 \omega, K_S^0 \eta', K_S^0 \phi, K_L^0 \pi^0 \pi^0$
Self-conjugate		$K_S^0 \pi^+ \pi^-$

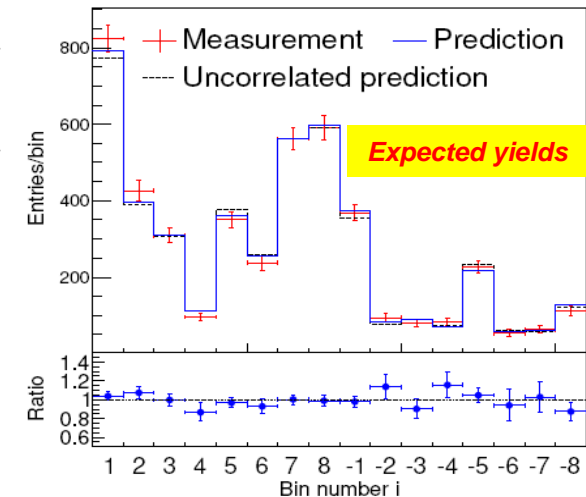
- ✓ For CP tags:  $\Gamma(S|CP) = A_S^2 A_{CP}^2 \left( 1 + (r_D^S)^2 - 2\lambda R_S r_D^S \cos \delta_D^S \right)$
- ✓ For Like-sign tags:  $\Gamma(S|S) = A_S^2 A_{\bar{S}}^2 [1 - R_S^2]$
- ✓ For Like-sign tags:  $Y_i^S = H \left( K_i + (r_D^S)^2 K_{-i} - 2r_D^S R_S \sqrt{K_i K_{-i}} \left[ c_i \cos \delta_D^S - s_i \sin \delta_D^S \right] \right)$

# Measurements of coherence factors in $D \rightarrow K^- \pi^+ \pi^+ \pi^-$ and $D \rightarrow K^- \pi^+ \pi^0$

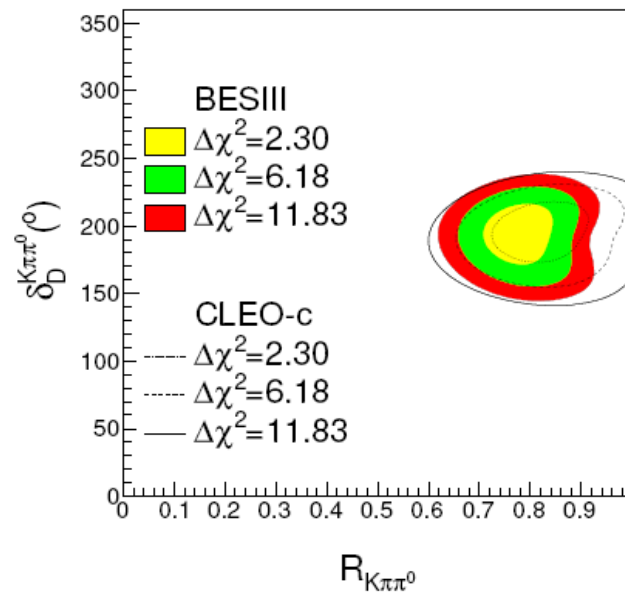
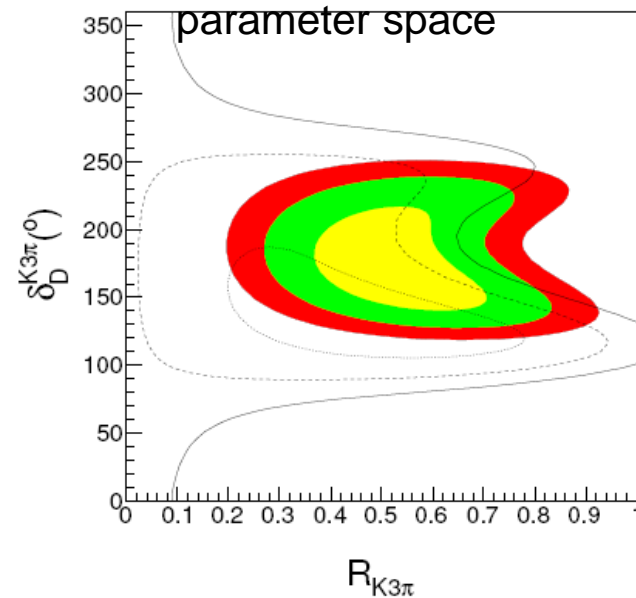
[JHEP05(2021)164]

▣ Fitted central values for strong-phase parameters.

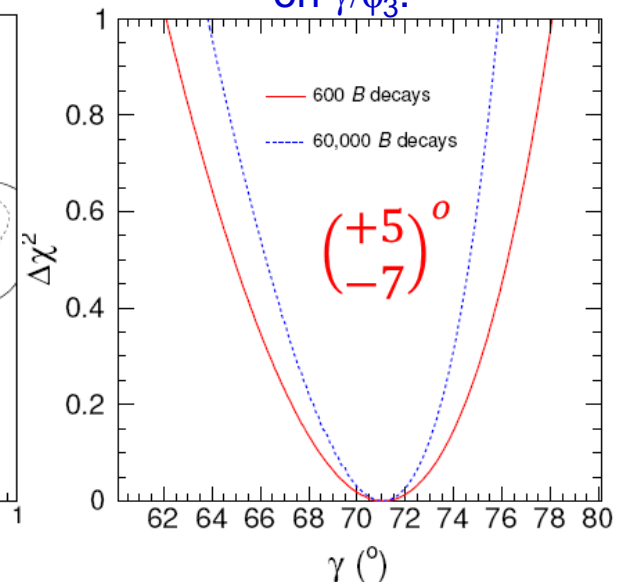
Parameter	Global fit	Binned fit			
		Bin 1	Bin 2	Bin 3	Bin 4
$R_{K3\pi}$	$0.52^{+0.12}_{-0.10}$	$0.58^{+0.25}_{-0.33}$	$0.78^{+0.50}_{-0.21}$	$0.85^{+0.15}_{-0.12}$	$0.45^{+0.33}_{-0.37}$
$\delta_D^{K3\pi}$	$(167^{+31}_{-19})^\circ$	$(131^{+124}_{-16})^\circ$	$(150^{+37}_{-39})^\circ$	$(176^{+57}_{-21})^\circ$	$(274^{+19}_{-30})^\circ$
$r_D^{K3\pi} (\times 10^{-2})$	$5.46 \pm 0.09$	$5.44^{+0.45}_{-0.14}$	$5.80^{+0.14}_{-0.13}$	$5.75^{+0.41}_{-0.14}$	$5.09^{+0.14}_{-0.14}$
$R_{K\pi\pi^0}$	$0.78 \pm 0.04$	$0.80 \pm 0.04$			
$\delta_D^{K\pi\pi^0}$	$(196^{+14}_{-15})^\circ$	$(200 \pm 11)^\circ$			
$r_D^{K\pi\pi^0} (\times 10^{-2})$	$4.40 \pm 0.11$	$4.41 \pm 0.11$			



✓ Scans of  $\Delta\chi^2$  in the global  $(R^{K3\pi}, \delta_D^{K3\pi})$  and  $(R_{K\pi\pi^0}, \delta_D^{K\pi\pi^0})$

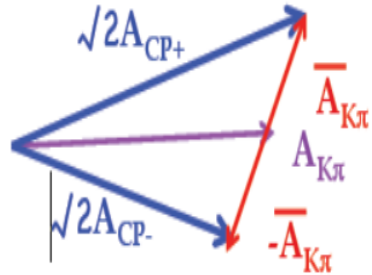


✓ Expected uncertainty on  $\gamma/\phi_3$ .



# Strong phase difference ( $\delta$ ) of $D \rightarrow K^- \pi^+$

Simplified Picture: (simple = no mixing)



Amplitude triangle:

$$CP_{\pm} = CF \pm DCSD$$

[DCSD enhanced for visibility!]

Complex ratio  
DCSD/CF amplitude

$$\frac{\langle K^- \pi^+ | \bar{D}^0 \rangle}{\langle K^- \pi^+ | D^0 \rangle} = -r e^{-i\delta_{K\pi}}$$

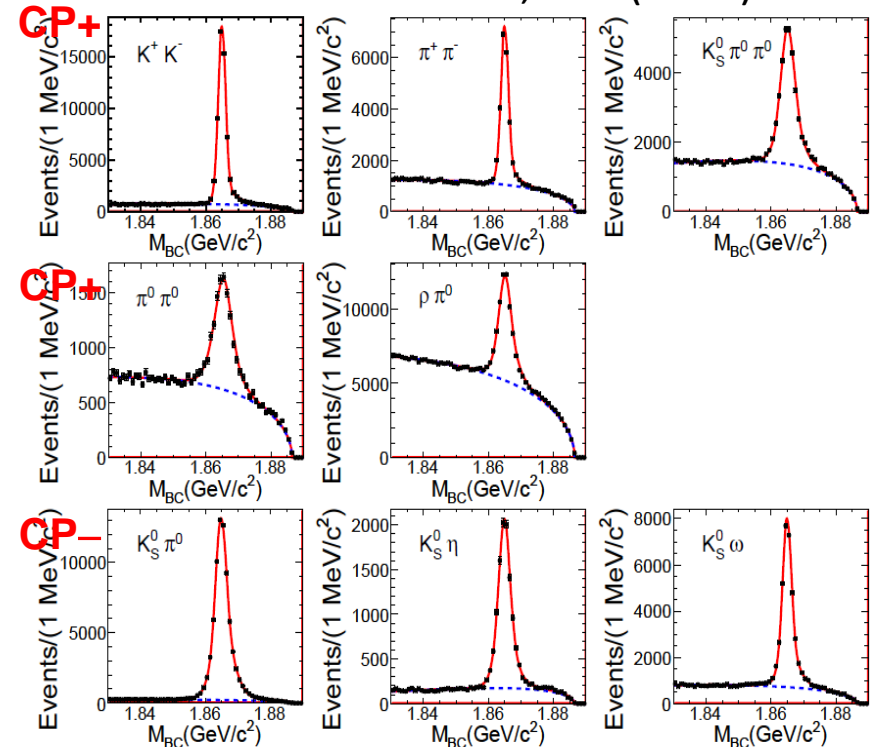
$$A_{K\pi}^{CP} = \frac{B_{D^{CP-} \rightarrow K^- \pi^+} - B_{D^{CP+} \rightarrow K^- \pi^+}}{B_{D^{CP-} \rightarrow K^- \pi^+} + B_{D^{CP+} \rightarrow K^- \pi^+}} = (12.7 \pm 1.3 \pm 0.7)\%$$

$$2r \cos \delta_{K\pi} + y = (1 + R_{WS}) A_{K\pi}^{CP}$$

With external inputs ( $r$ ,  $y$ ,  $R_{WS}$ ) and  $A_{K\pi}$ :

$$\cos \delta_{K\pi} = 1.02 \pm 0.11 \pm 0.06 \pm 0.01$$

BESIII: PLB734, 227 (2014)



CLEO-c: PRD86, 112001 (2012)

$$\cos \delta_{K\pi}^{\text{without input}} = 0.81_{-0.18-0.05}^{+0.22+0.07}$$

$$\cos \delta_{K\pi}^{\text{with input}} = 1.15_{-0.17-0.08}^{+0.19+0.00}$$

Benefit ADS method of extracting  $\gamma_9$

# Other topics: $y_{CP}$

## CP eigenstates of D:

$$|D_{CP+}\rangle \equiv \frac{|D^0\rangle + |\bar{D}^0\rangle}{\sqrt{2}}, \quad |D_{CP-}\rangle \equiv \frac{|D^0\rangle - |\bar{D}^0\rangle}{\sqrt{2}}$$

If no direct CPV, the difference of effective lifetimes between CP and flavor eigenstates:

$$y_{CP} = \frac{1}{2} \left[ y \cos \phi \left( \left| \frac{q}{p} \right| + \left| \frac{p}{q} \right| \right) - x \sin \phi \left( \left| \frac{q}{p} \right| - \left| \frac{p}{q} \right| \right) \right]$$

$$y = \Delta\Gamma/2\Gamma, \quad |D_{1,2}\rangle = p|D^0\rangle \pm q|\bar{D}^0\rangle \quad \phi = \arg(q/p)$$

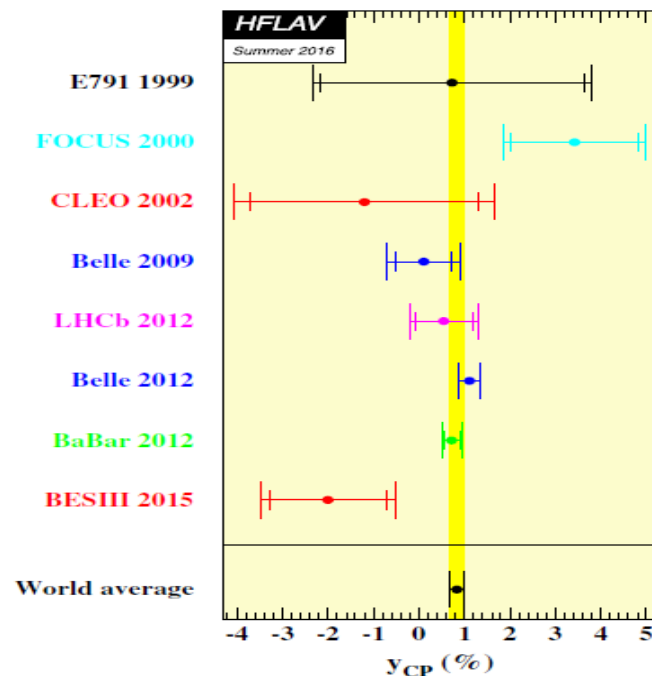
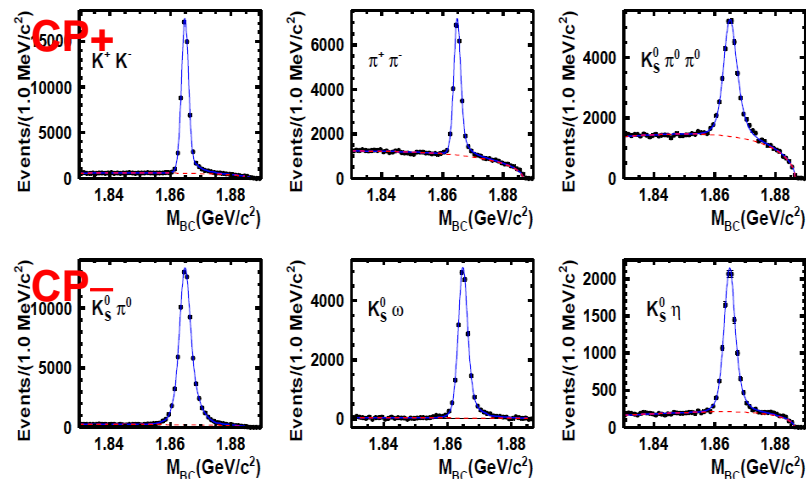
If no CPV,  $\phi=0$  and  $|p/q|=1$ :  $y_{CP}=y$

Measurements with  $D \rightarrow K\ell\nu$  &  $K\mu\nu$

$$y_{CP} \approx \frac{1}{4} \left( \frac{B_{D^{CP-} \rightarrow K\ell\nu}}{B_{D^{CP+} \rightarrow K\ell\nu}} - \frac{B_{D^{CP+} \rightarrow K\ell\nu}}{B_{D^{CP-} \rightarrow K\ell\nu}} \right) = (-2.0 \pm 1.3 \pm 0.7)\%$$

Need more data or global fit

BESIII, PLB744,339(2015)



# Desirable Quantum Correlated measurements

LHCb-PUB-2016

- Uncertainty of strong-phase inputs from CLEO-c contribute  $\sim 2^\circ$  to  $\gamma$ , and will be comparable with the experimental systematic uncertainty at LHCb RUN2

- BESIII is only machine running at  $\tau$ -charm energy region. Related QC studies are key to constrain the  $\gamma$  measurement at LHCb upgrades 1(2)

Decay mode	Quantity of interest	Comments
$D \rightarrow K_s^0 \pi^+ \pi^-$	$c_i$ and $s_i$ ✓	Binning schemes as those used in the CLEO-c analysis. With future, very large $\psi(3770)$ data sets, it might be worthwhile to explore alternative binning.
$D \rightarrow K_s^0 K^+ K^-$	$c_i$ and $s_i$ ✓	Binning schemes as those used in the CLEO-c analysis. With future, very large $\psi(3770)$ data sets, it might be worthwhile to explore alternative binning.
$D \rightarrow K^\pm \pi^\mp \pi^+ \pi^-$	$R, \delta$ ✓	In bins guided by amplitude models, currently under development by LHCb.
$D \rightarrow K^+ K^- \pi^+ \pi^-$	$c_i$ and $s_i$	Binning scheme can be guided by the CLEO model [18] or potentially an improved model from LHCb in the future.
$D \rightarrow \pi^+ \pi^- \pi^+ \pi^-$	$F_+$ or $c_i$ and $s_i$	Unbinned measurement of $F_+$ . Measurements of $F_+$ in bins or $c_i$ and $s_i$ in bins could be explored.
$D \rightarrow K^\pm \pi^\mp \pi^0$	$R, \delta$ ✓	Simple 2-3 bin scheme could be considered.
$D \rightarrow K_s^0 K^\pm \pi^\mp$	$R, \delta$	Simple 2 bin scheme where one bin encloses the $K^*$ resonance.
$D \rightarrow \pi^+ \pi^- \pi^0$	$F_+$	No binning required as $F_+ \sim 1$ .
$D \rightarrow K_s^0 \pi^+ \pi^- \pi^0$	$F_+$ and $c_i$ and $s_i$	Unbinned measurement of $F_+$ required. Additional measurements of $F_+$ or $c_i$ and $s_i$ in bins could be explored.
$D \rightarrow K^+ K^- \pi^0$	$F_+$	Unbinned measurement required. Extensions to binned measurements of either $F_+$ or $c_i$ and $s_i$ possible.
$D \rightarrow K^\pm \pi^\mp$	$\delta$	Of low priority due to good precision available through charm-mixing analyses.

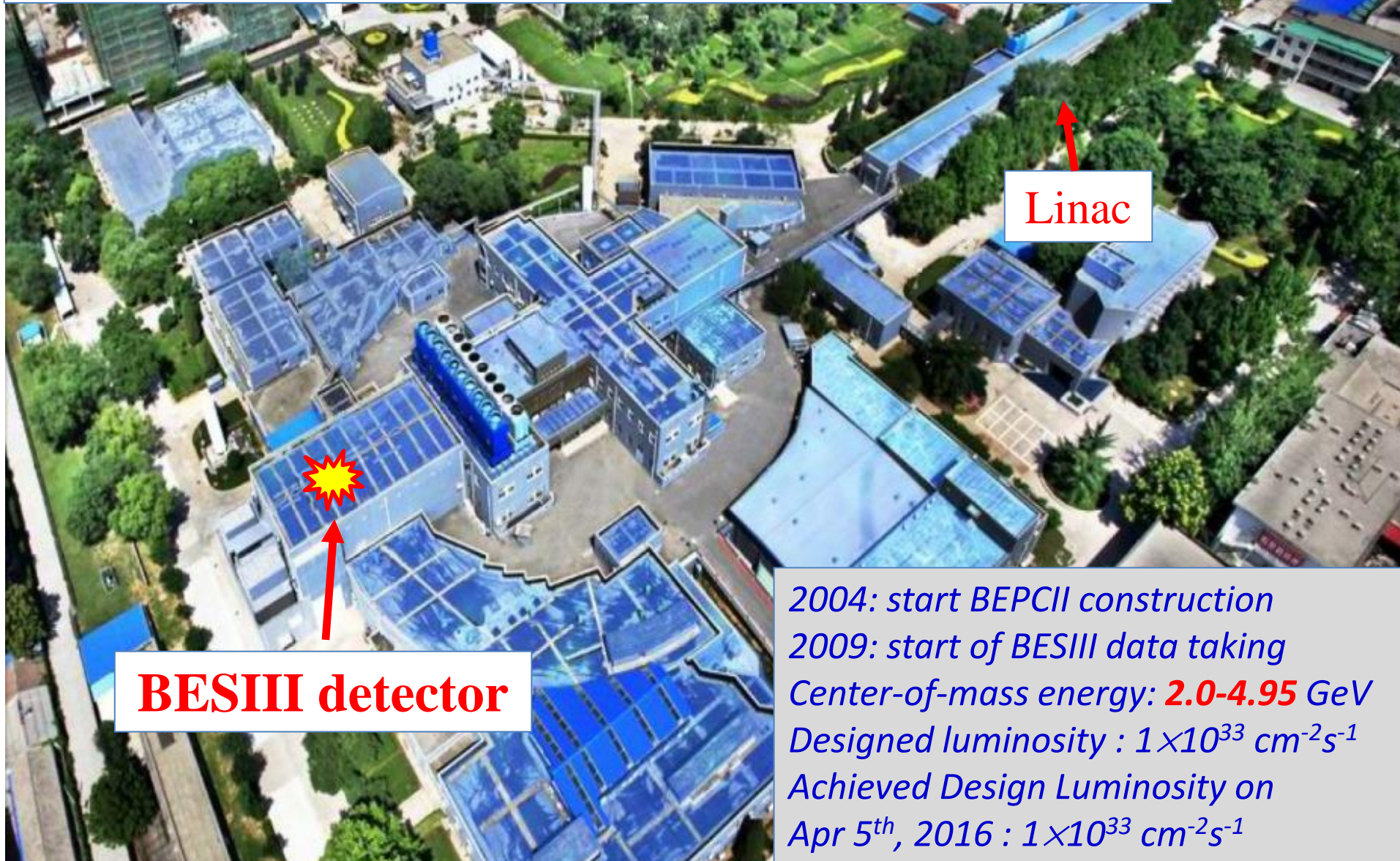
# Summary

- BESIII had collected the largest data samples for study of charmed hadrons, including  $D^0$ ,  $D^+$ ,  $D_s^+$  and  $\Lambda_c^+$  etc.
- A range of important and unique results had been published in recent years.
- BESIII will collect more 17/fb data at  $\psi(3770)$  resonance peak. The constraint on  $\gamma$  from strong phases is expected to be  $<0.5\%$ .
- More important results are expected in charmed hadron decays at BESIII.

# Thanks!



# Beijing Electron Positron Collider (BEPC)

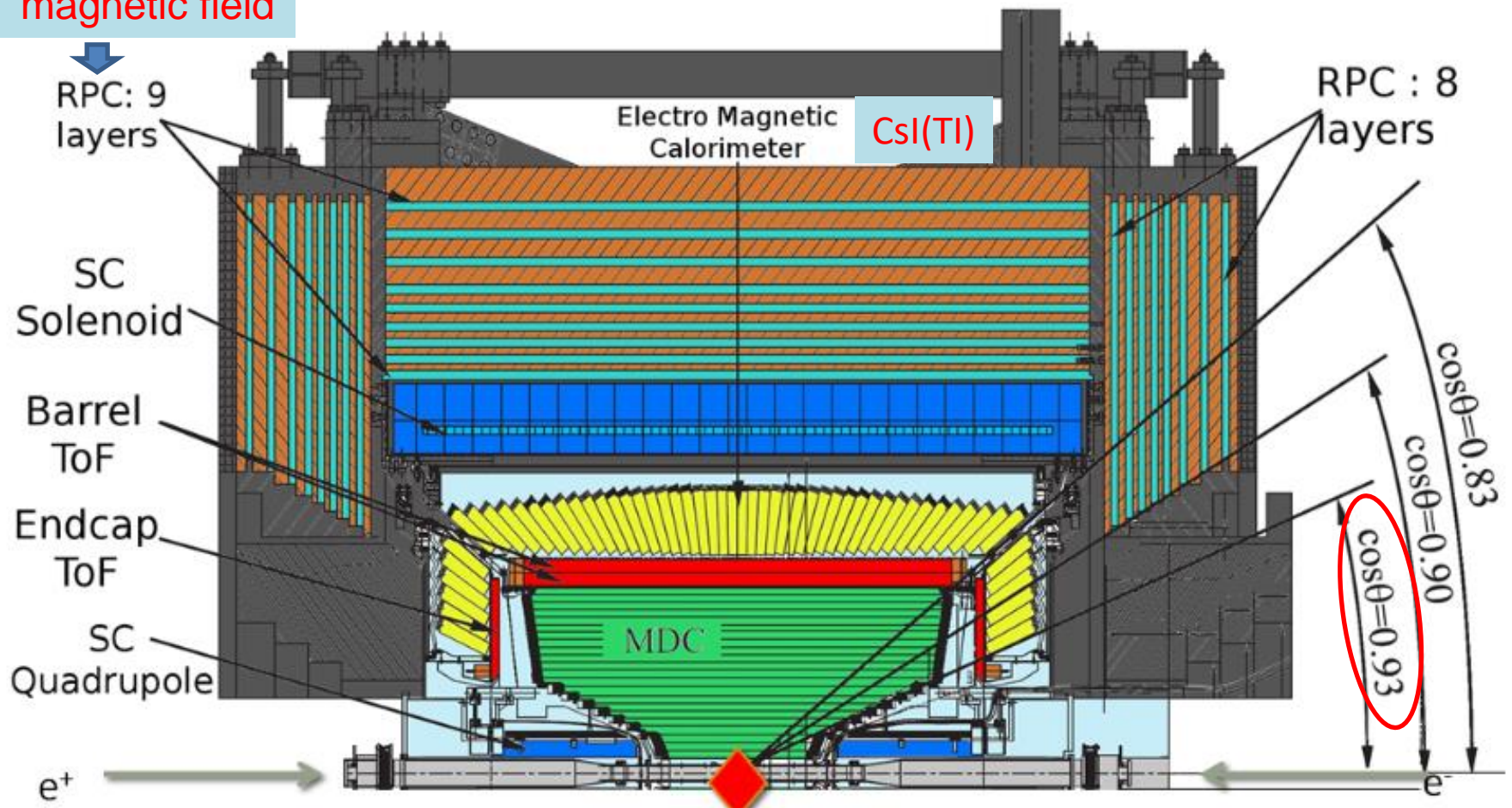


2004: start BEPCII construction  
2009: start of BESIII data taking  
Center-of-mass energy: **2.0-4.95 GeV**  
Designed luminosity :  $1 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$   
Achieved Design Luminosity on  
Apr 5<sup>th</sup>, 2016 :  $1 \times 10^{33} \text{ cm}^{-2} \text{ s}^{-1}$



# BESIII Detector

1 T magnetic field



□ MDC:

$\sigma_{p/p}$ : 0.5% @ 1 GeV/c  
 $\sigma_{dE/dx}$ : 6% for electrons

□ TOF:

$\sigma_T$ : 68 ps (Barrel)  
 110 ps (Endcap)

□ EMC:

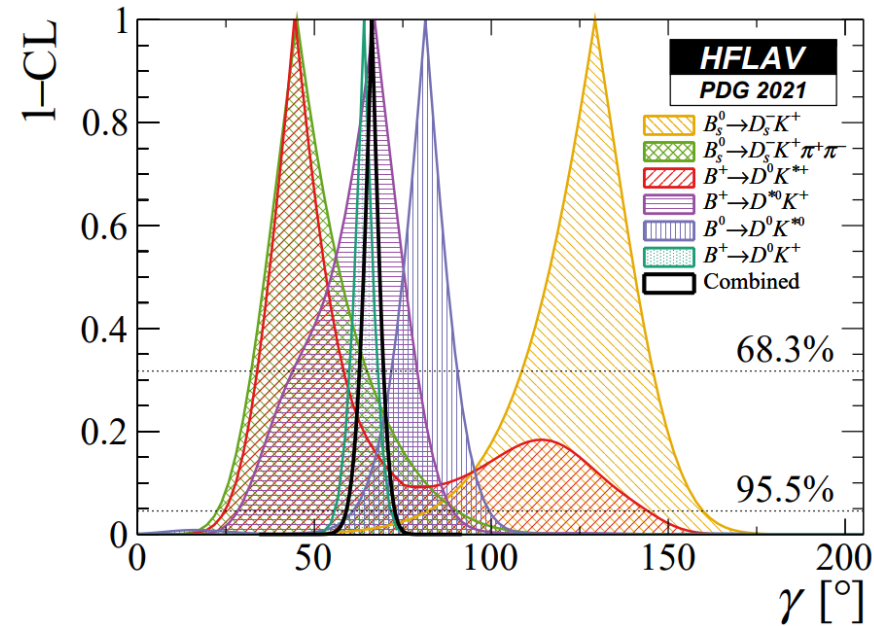
$\sigma_E/\sqrt{E}$ : 2.5(5.0)% @ 1 GeV  
 for barrel(endcap)

# BESIII Collaboration



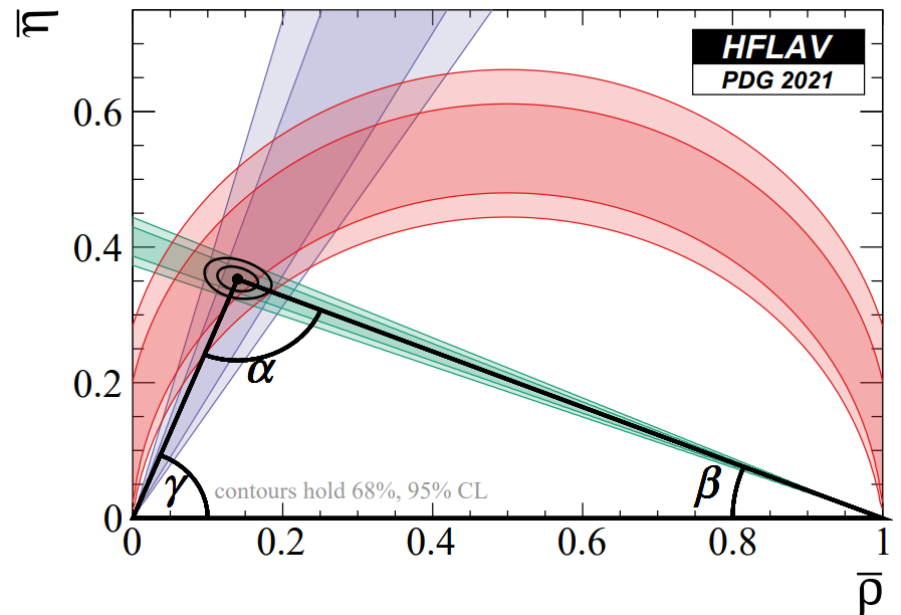
**BESIII: ~500 members from 80 institutes in 17 countries.**

# The latest world averages of the CKM Unitarity Triangle Angles



$$\gamma/\phi_3 = (66.2^{+3.4}_{-3.6})^\circ$$

$$\alpha + \beta + \gamma = (173.6^{+5.9}_{-5.6})^\circ$$



$$\beta/\phi_1 = (22.2 \pm 0.7)^\circ$$

$$\alpha/\phi_2 = (85.2^{+4.8}_{-4.3})^\circ$$



■ For  $B \rightarrow D(K_S^0 \pi^+ \pi^-) K$  decays, the relations are followed:

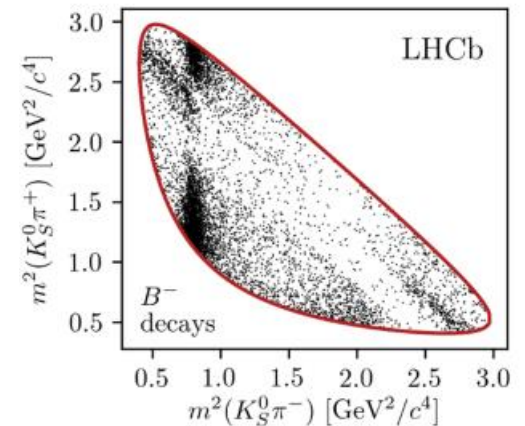
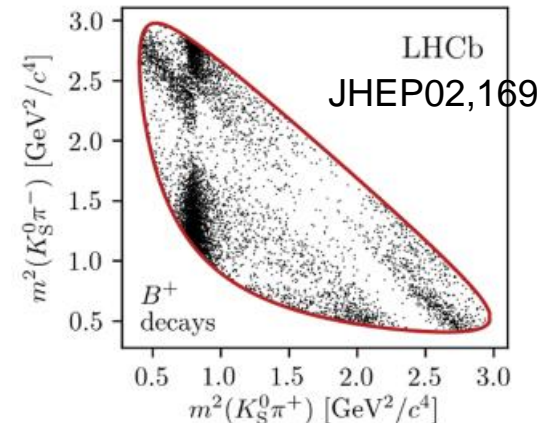
✓ bin +i:

$$\begin{aligned}\hat{\Gamma}_i^- &\equiv \int_i d\hat{\Gamma}(B^- \rightarrow (K_S \pi^- \pi^+)_D K^-) \\ &= T_i + r_B^2 T_i^- + 2r_B [\cos(\delta_B - \gamma) c_i + \sin(\delta_B - \gamma) s_i], \\ \hat{\Gamma}_i^+ &\equiv \int_i d\hat{\Gamma}(B^+ \rightarrow (K_S \pi^- \pi^+)_D K^+) \\ &= T_i^- + r_B^2 T_i + 2r_B [\cos(\delta_B + \gamma) c_i - \sin(\delta_B + \gamma) s_i],\end{aligned}$$

✓ bin -i:

$$\begin{aligned}\hat{\Gamma}_i^- &\equiv \int_i d\hat{\Gamma}(B^- \rightarrow (K_S \pi^- \pi^+)_D K^-) \\ &= T_i^- + r_B^2 T_i + 2r_B [\cos(\delta_B - \gamma) c_i - \sin(\delta_B - \gamma) s_i], \\ \hat{\Gamma}_i^+ &\equiv \int_i d\hat{\Gamma}(B^+ \rightarrow (K_S \pi^- \pi^+)_D K^+) \\ &= T_i + r_B^2 T_i^- + 2r_B [\cos(\delta_B + \gamma) c_i + \sin(\delta_B + \gamma) s_i].\end{aligned}$$

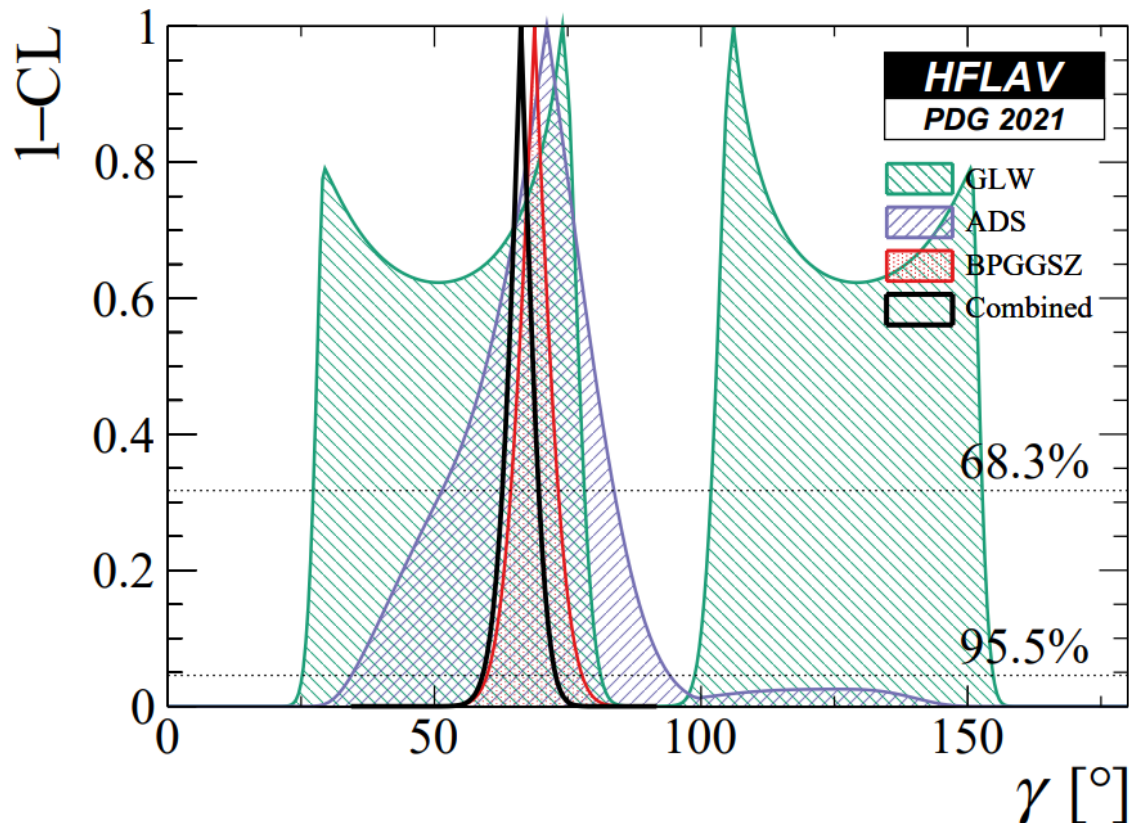
Assuming  $k$  bins in DP,  
4k relations and 2k+3 unknowns;  
The bin number:  $k \geq 2$ ;



Charm inputs ( $c_i$  and  $s_i$ ) are essential to constrain the  $\gamma$  measurement!

- $D \rightarrow K_S^0 \pi^+ \pi^-$  is the most important channel to measure strong-phases
- The  $(c_i, s_i)$  measured in  $D \rightarrow K_S^0 \pi^+ \pi^-$  also provide critical inputs of measuring strong-phase in other D decays.

- ✓ GLW <sup>[1]</sup>: via  $D^0 \rightarrow$  CP eigenstate,  $K^+ K^-$ ,  $\pi^+ \pi^-$ ,  $K_S^0 \pi^0$  etc. ↖ No strong-phase input for  $\gamma$  measurements for GLW.
- ✓ ADS <sup>[2]</sup>: via  $D^0 \rightarrow$  CF and DCS, such as  $K^+ \pi^-$ ,  $K^+ \pi^- \pi^0$ ,  $K^+ \pi^- \pi^- \pi^+$ ,  $K_S^0 K \pi$ , etc.
- ✓ GGSZ <sup>[3]</sup>: via with  $D^0 \rightarrow$  Multi-body self-conjugate decays,  $K_S^0 \pi^+ \pi^-$ ,  $K_S^0 K^+ K^-$  etc.



The precision of  $\gamma$  measurement from three methods.