



上海交通大学
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CP Violating Dark Photon Kinetic Mixing and Type-III Seesaw

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Model building

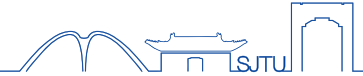
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Introduction



Dark photon



Abelian/Non-Abelian kinetic mixing

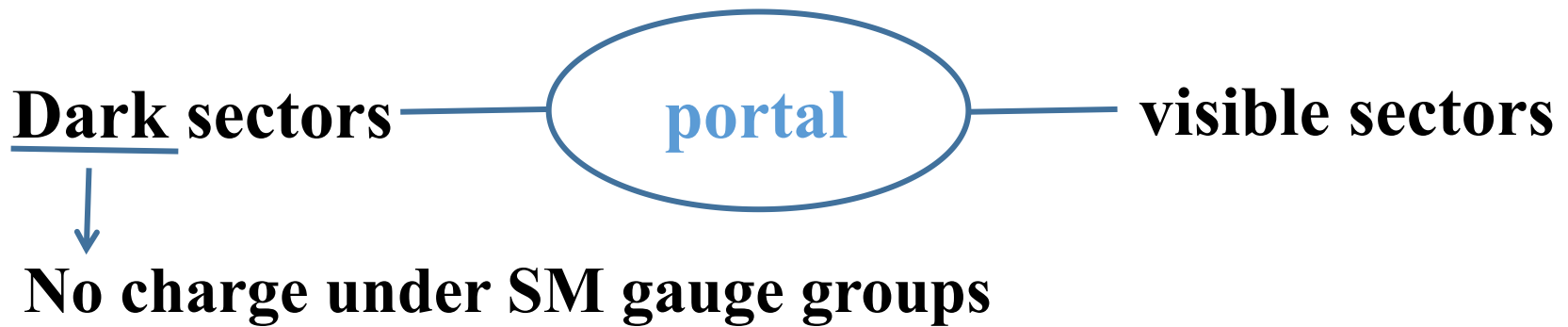
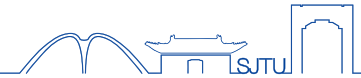


CP violating kinetic mixing



Renormalizable CP violating kinetic mixing

Dark photon

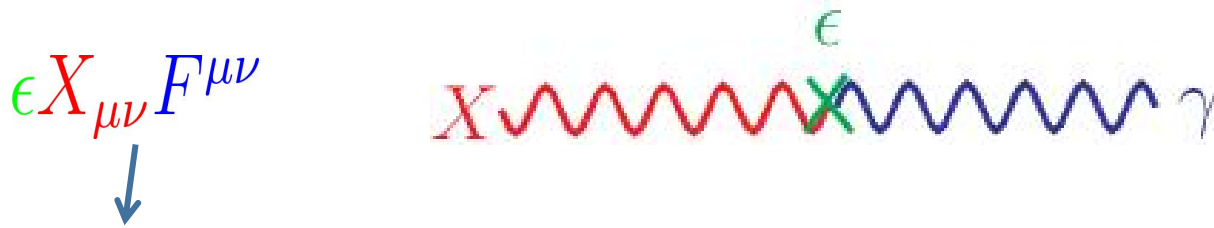
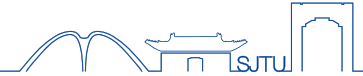


Portal	Spinor: Neutrino	$\bar{L}NH$
	Scalar: Higgs	$SH^\dagger H, S^2 H^\dagger H$
	Pseudo-scalar: Axion	$(a/f_a)F\tilde{F}$
	Vector: dark photon	$U(1)_Y \times U(1)_X$

$$L = -\frac{1}{4}Y_{\mu\nu}Y^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} + j_Y^\mu Y_\mu + \color{red}{j_X^\mu X_\mu}$$

No direct interaction with SM particle

Abelian kinetic mixing



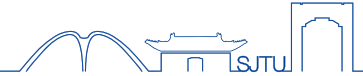
1. Renormalizable dimension 4 operator
2. Enlighten : dark photon interacts with SM particles

Gauge group $SU(3) \times SU(2)_L \times U(1)_Y \times U(1)_X$

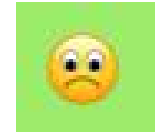
$$L = -\frac{1}{4}Y_{\mu\nu}Y^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} + j_Y^\mu Y_\mu + j_X^\mu X_\mu - \frac{\epsilon}{2}X_{\mu\nu}Y^{\mu\nu}$$

Rewrite in the canonical form to identify physical gauge boson (remove mixing term)

Non-Abelian kinetic mixing



$W_{\mu\nu}^a X^{\mu\nu} \longrightarrow$ **Not gauge invariant**



How to contract index a

1. triplet scalar Σ^a $\langle \Sigma^a \rangle = v_\Sigma / \sqrt{2}$

gauge singlet $W_{\mu\nu}^a X^{\mu\nu} \Sigma^a$

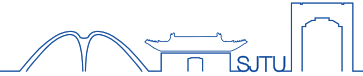
non-renormalizable dimension 5

2. Higher order $W_{\mu\nu}^a X^{\mu\nu} (H^\dagger \tau^a H)$

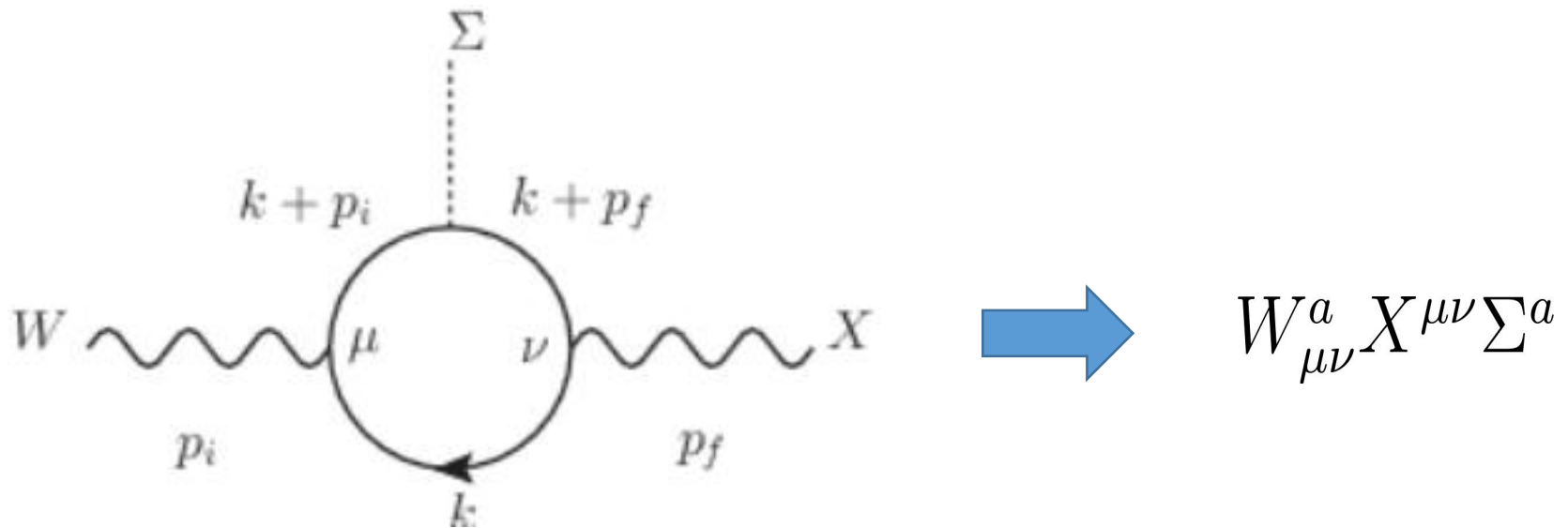
non-renormalizable dimension 6

Construct kinetic mixing between abelian and non-Abelian

Non-Abelian kinetic mixing



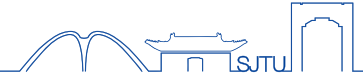
Generate the kinetic mixing from renormalizable theory?



Generate the kinetic mixing at loop level

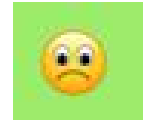
Non-abelian kinetic mixing $W_{\mu\nu}^a Y^{b\mu\nu} \Sigma^{ab}$

CP violation



1. Abelian

$$\left\{ \begin{array}{ll} \text{CP conserving} & X_{\mu\nu} Y^{\mu\nu} \\ \text{CP violating} & Y^{\mu\nu} \tilde{X}_{\mu\nu} = -\frac{\epsilon_{\mu\nu\alpha\beta}}{2} \partial^\alpha (Y^{\mu\nu} X^\beta) \end{array} \right.$$



2. Non-Abelian

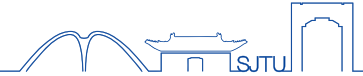
$$\mathcal{L}_X = -\frac{\beta_X}{\Lambda} \text{Tr}(W_{\mu\nu} \Sigma) X^{\mu\nu} - \frac{\tilde{\beta}_X}{\Lambda} \text{Tr}(W_{\mu\nu} \Sigma) \tilde{X}^{\mu\nu}$$



$$\begin{aligned} & -\frac{\beta_X}{2\Lambda} X^{\mu\nu} (s_W F_{\mu\nu} + c_W Z_{\mu\nu} + ig(W_\mu^- W_\nu^+ - W_\mu^+ W_\nu^-)) (v_\Sigma + \Sigma^0) \\ & -\frac{\tilde{\beta}_X}{2\Lambda} \tilde{X}^{\mu\nu} [(s_W F_{\mu\nu} + c_W Z_{\mu\nu}) \Sigma^0 + ig(W_\mu^- W_\nu^+ - W_\mu^+ W_\nu^-) (v_\Sigma + \Sigma^0)] \end{aligned}$$

CP violating non-Abelian kinetic mixing

Renormalizable CP violation

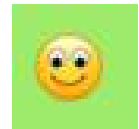
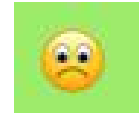


Renormalizable CPV non-Abelian kinetic mixing?



New fields generate the operator from loop order

$\left\{ \begin{array}{l} \text{Scalar: No generate CPV due to tensor} \\ \text{Fermion: SU(2)}_L \text{ multiplet} \end{array} \right.$



Type-III seesaw model

$$f = \begin{pmatrix} \frac{\Sigma^0}{\sqrt{2}} & \Sigma^+ \\ \Sigma^- & -\frac{\Sigma^0}{\sqrt{2}} \end{pmatrix} \longrightarrow (\bar{\nu}_L, \bar{f}^c) \begin{pmatrix} 0 & M_D \\ M_D^T & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ f \end{pmatrix}$$

Lepton triplet

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Renormalizable CPV kinetic mixing model

Gauge group $SU(3)_C \times SU(2)_L \times U(1)_Y \times U(1)_X$

1. SM particles has no $U(1)_X$ charge

2. Triplet scalar $\Sigma^a = (1, 3, 0)(0) \longrightarrow \epsilon^{\alpha\beta\mu\nu} X_{\alpha\beta} W_{\mu\nu}^a \Sigma^a$

3. Triplet fermion (No scalar) \longrightarrow **Type-III seesaw model**

$$f_1 = (1, 3, 0)(x_f), \quad f_2 = (1, 3, 0)(-x_f), \quad f_3 = (1, 3, 0)(0)$$

Gauge anomaly free

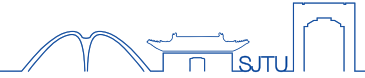
4. Singlet scalar \longrightarrow dark photon and heavy neutrino mass

$$S_X = (1, 1, 0)(-2x_f),$$

5. Higgs scalar \longrightarrow Neutrino mass matrix

$$H'_{1,2} : (1, 2, -1/2)(\mp x_f)$$

Component field



$$f_R = \frac{1}{2} \begin{pmatrix} f_R^0 & \sqrt{2}f_R^+ \\ \sqrt{2}f_R^- & -f_R^0 \end{pmatrix}, \quad f_L = f_R^c = \frac{1}{2} \begin{pmatrix} (f_R^0)^c & \sqrt{2}(f_R^-)^c \\ \sqrt{2}(f_R^+)^c & -(f_R^0)^c \end{pmatrix}$$

$$\Sigma = \frac{1}{2} \begin{pmatrix} \Sigma^0 & \sqrt{2}\Sigma^+ \\ \sqrt{2}\Sigma^- & -\Sigma^0 \end{pmatrix}, \quad W_\mu = \frac{1}{2} \begin{pmatrix} W_\mu^0 & \sqrt{2}W_\mu^+ \\ \sqrt{2}W_\mu^- & -W_\mu^0 \end{pmatrix}$$

$$-(\beta_X/\Lambda) \text{Tr}(W_{\mu\nu}\Sigma) X^{\mu\nu}$$

$$-(\tilde{\beta}_X/\Lambda) \text{Tr}(W_{\mu\nu}\Sigma) \tilde{X}^{\mu\nu}$$

$$-\frac{1}{2} \left[(s_W F_{\mu\nu} + c_W Z_{\mu\nu}) + ig(W_\mu^- W_\nu^+ - W_\mu^+ W_\nu^-) \right] (v_\Sigma + \Sigma^0)$$

Coupling Lagrangian



Quark Yukawa coupling

$$\mathcal{L}_Y(q) = -\bar{Q}_L Y_u H U_R - \bar{Q}_L Y_d \tilde{H} D_R + \text{H.c.}$$

Lepton Yukawa coupling

$$\begin{aligned} \mathcal{L}_Y(l) = & -\bar{L}_L Y_e \tilde{H} E_R - \bar{L}_L Y_{fL3} \tilde{H} f_{R3} - \bar{f}_{R1}^c Y_{fs1} S_X f_{R1} \\ & - \bar{f}_{R2}^c Y_{fs2} S_X^\dagger f_{R2} - \bar{f}_{R1}^c m_{12} f_{R2} - \bar{f}_{R3}^c m_{33} f_{R3} + \text{H.c.} \end{aligned}$$

New added term

$$-\bar{L}_L Y_{fL1} H'_1 f_{R1} - \bar{L}_L Y_{fL2} H'_2 f_{R2}$$

Interaction Lagrangian



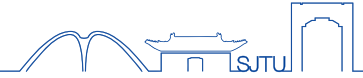
$W^0/X - f_R$ coupling

$$\begin{aligned}\mathcal{L}_{int} = & X_\mu \left[\bar{f}_{R1}^+ \gamma^\mu f_{R1}^+ + \bar{f}_{R1}^0 \gamma^\mu f_{R1}^0 + \bar{f}_{R1}^- \gamma^\mu f_{R1}^- \right. \\ & \left. - \bar{f}_{R2}^+ \gamma^\mu f_{R2}^+ - \bar{f}_{R2}^0 \gamma^\mu f_{R2}^0 - \bar{f}_{R2}^- \gamma^\mu f_{R2}^- \right] g_X x_f \\ & + W_\mu^0 \left[\bar{f}_{R1}^+ \gamma^\mu f_{R1}^+ - \bar{f}_{R1}^- \gamma^\mu f_{R1}^- \right. \\ & \left. + \bar{f}_{R2}^+ \gamma^\mu f_{R2}^+ - \bar{f}_{R2}^- \gamma^\mu f_{R2}^- \right] g\end{aligned}$$

$\Sigma - f$ coupling

$$\mathcal{L}_{Y_{f\sigma}} = Y_{f\sigma 12} \left(\overline{(f_{R1}^+)^c} f_{R2}^+ - \overline{(f_{R1}^-)^c} f_{R2}^- \right) (v_\Sigma + \Sigma^0) + \text{H.c.}$$

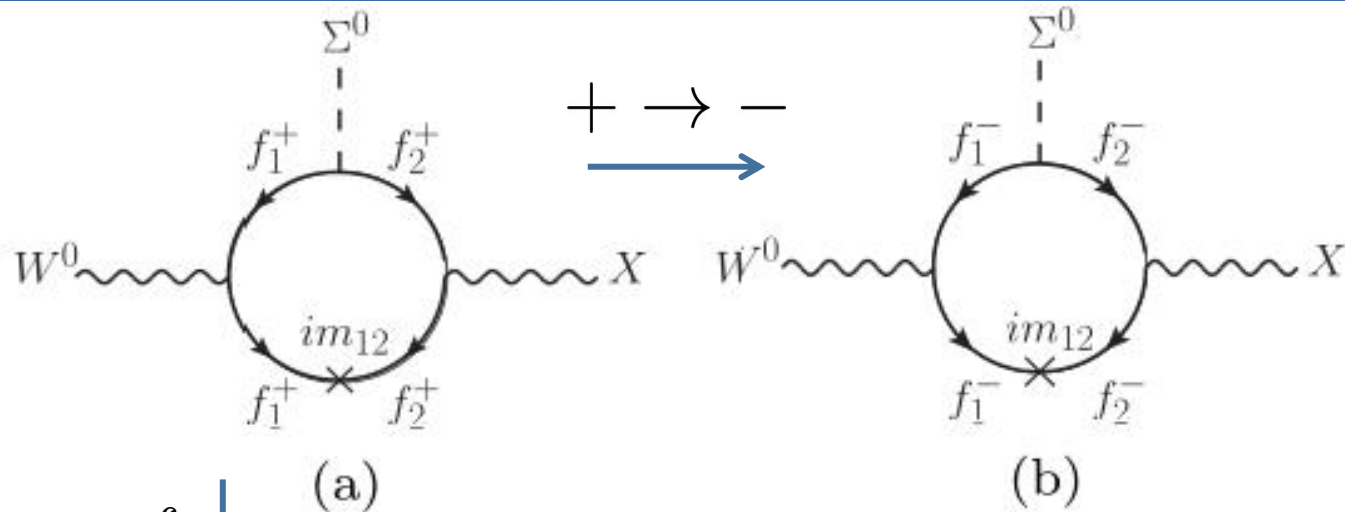
Lepton masses



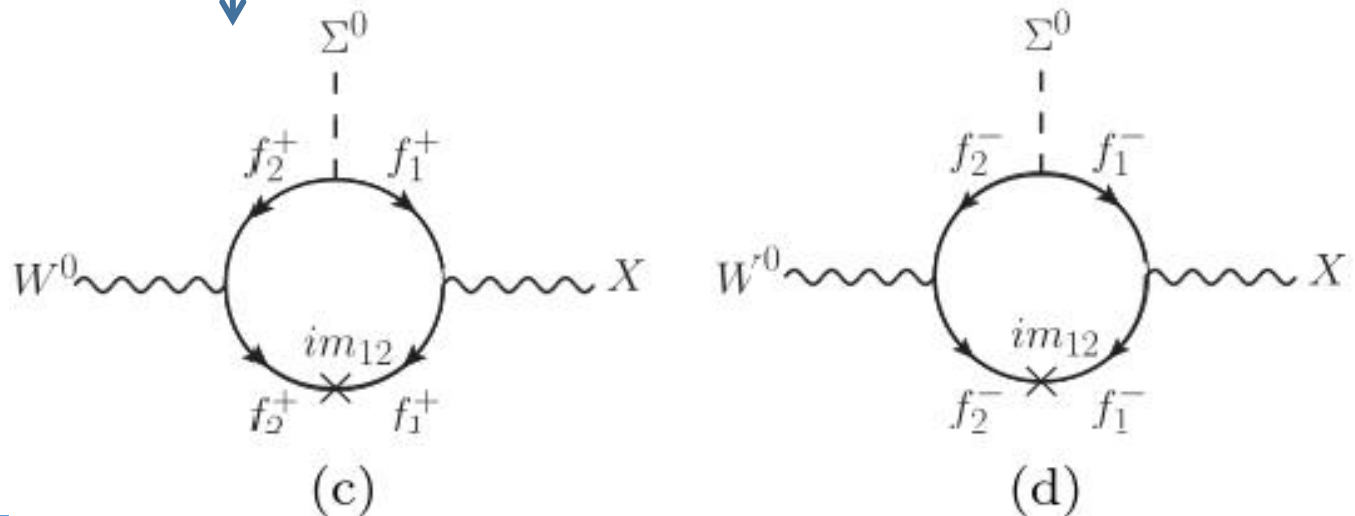
$$\mathcal{L}_m = -\frac{1}{2}(\bar{\nu}_L, \bar{\nu}_R^c) \begin{pmatrix} 0 & M_D \\ M_D^T & M_R \end{pmatrix} \begin{pmatrix} \nu_L^c \\ \nu_R \end{pmatrix} \\ -(\bar{E}_L, \bar{f}_L) \begin{pmatrix} m_e & \sqrt{2}M_D \\ 0 & M_R \end{pmatrix} \begin{pmatrix} E_R \\ f_R \end{pmatrix}$$

$$M_D = \begin{pmatrix} \frac{Y_{fL11}v'_1}{\sqrt{2}} & \frac{Y_{fL12}v'_2}{\sqrt{2}} & \frac{Y_{fL13}v}{\sqrt{2}} \\ \frac{Y_{fL21}v'_1}{\sqrt{2}} & \frac{Y_{fL22}v'_2}{\sqrt{2}} & \frac{Y_{fL23}v}{\sqrt{2}} \\ \frac{Y_{fL31}v'_1}{\sqrt{2}} & \frac{Y_{fL32}v'_2}{\sqrt{2}} & \frac{Y_{fL33}v}{\sqrt{2}} \end{pmatrix} \quad M_R = \begin{pmatrix} \frac{Y_{fs1}v_s}{\sqrt{2}} & m_{12} & 0 \\ m_{12} & \frac{Y_{fs2}v_s}{\sqrt{2}} & 0 \\ 0 & 0 & m_{33} \end{pmatrix}$$

One loop kinetic mixing



$f_1 \leftrightarrow f_2 \downarrow$



$$\longrightarrow \text{Tr}(W_{\mu\nu}\Sigma)\tilde{X}^{\mu\nu}$$

One loop kinetic mixing



$$M^{\mu\nu} = \frac{i}{4\pi^2} g g_X x_f m_{12} Y_{f\sigma}^* \epsilon^{\mu\nu\alpha\beta} p_{X\alpha} p_{W\beta}$$

No scalar

$$\times (f(m_1, m_2, p_W, p_X) + f(m_2, m_1, p_W, p_X))$$

$$f(m_1, m_2, p_W, p_X) = \int_0^1 dx \int_0^{1-x} dy (1 - x - y)$$

$$\times [D(m_1, m_2, p_W, p_X) + D(m_2, m_1, p_X, p_W)]$$

$$D(m_1, m_2, p_W, p_X) = \frac{m_1^2}{(m_1^2 - m_2^2)}$$

$$\times \frac{1}{(m_1^2 - y(m_1^2 - m_2^2) - x p_W^2 - y p_X^2 + (x p_W - y p_X)^2)}$$

CPC and CPV kinetic mixing



$$\mathcal{L}_X \rightarrow -\frac{\tilde{\beta}_X}{2\Lambda} \tilde{X}^{\mu\nu} [(s_W F_{\mu\nu} + c_W Z_{\mu\nu})] (v_\Sigma + \Sigma^0)$$

$$\downarrow 2\epsilon^{\mu\nu\alpha\beta} p_{X\alpha} p_{W\beta} \epsilon_{W\mu} \epsilon_{X\nu}^* \rightarrow -\tilde{X}^{\mu\nu} W_{\mu\nu}^0$$

CP conserving

$$\frac{\beta_X}{\Lambda} = \frac{1}{2\pi^2} g g_X x_f \text{Re}(m_{12} Y_{f\sigma}^*)$$

$$\times [f(m_1, m_2, p_W, p_X) + f(m_2, m_1, p_W, p_X)]$$

CP violating

$$\frac{\tilde{\beta}_X}{\Lambda} = \frac{1}{2\pi^2} g g_X x_f \text{Im}(m_{12} Y_{f\sigma}^*)$$

$$\times [f(m_1, m_2, p_W, p_X) + f(m_2, m_1, p_W, p_X)] \quad 19/30$$

Gauge field Lagrangian



$$\mathcal{L}_{KM} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} - \frac{1}{4}Z_{\mu\nu}Z^{\mu\nu} - \frac{1}{4}X_{\mu\nu}X^{\mu\nu} \quad \textbf{kinetic term}$$

$$+ \frac{1}{2}m_Z^2 Z_\mu Z^\mu + \frac{1}{2}m_X^2 X_\mu X^\mu \quad \textbf{mass term}$$

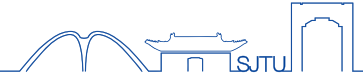
$$- \frac{1}{2}\epsilon_{AX} F_{\mu\nu} X^{\mu\nu} - \frac{1}{2}\epsilon_{ZX} Z_{\mu\nu} X^{\mu\nu} \quad \textbf{kinetic mixing}$$

$$\epsilon_{AX} = \alpha_{XY} c_W + \beta_X s_W v_\Sigma / \Lambda$$

$$\epsilon_{ZX} = -\alpha_{XY} s_W + \beta_X c_W v_\Sigma / \Lambda$$

$$- (1/2) \alpha_{XY} X^{\mu\nu} B_{\mu\nu}$$

Mass eigen-state



leading order

$$\begin{pmatrix} A \\ Z \\ X \end{pmatrix} = \begin{pmatrix} 1 & 0 & -\epsilon_{AX} \\ 0 & 1 & -\xi - \epsilon_{ZX} \\ 0 & \xi & 1 \end{pmatrix} \begin{pmatrix} A^m \\ Z^m \\ X^m \end{pmatrix}$$

$$\xi \approx -m_Z^2 \epsilon_{ZX} / (m_Z^2 - m_X^2)$$

$$\mathcal{L} = J_{em}^\mu A_\mu + J_Z^\mu Z_\mu + J_X^\mu X_\mu$$

$$\downarrow$$

$$\mathcal{L}^m$$

Numerical analysis



$$\frac{\tilde{\beta}_X}{\Lambda} = \frac{1}{2\pi^2} g g_X x_f \text{Im}(m_{12} Y_{f\sigma}^*)$$

$$\times [f(m_1, m_2, p_W, p_X) + f(m_2, m_1, p_W, p_X)]$$

↓ **degenerate** $m_1 \approx m_2 \approx |m_{12}| \approx m$

$$\tilde{\beta}_X v_\Sigma / \Lambda \approx \underbrace{g g_X x_f}_{\text{cos } \delta} \underbrace{|Y_{f\sigma}^*|}_{\text{sin } \delta} \underbrace{v_\Sigma}_{\text{sin } \delta} / 6\pi^2 m < 5 \sin \delta \times 10^{-4}$$

$$\rho - 1 = \frac{4v_\Sigma^2}{v^2} = 0.00038 \pm 0.00020 \longrightarrow v_\Sigma < 3\text{GeV}$$

$$m > 790\text{GeV}(\text{ATLAS}), 880\text{GeV}(\text{CMS})$$

$$g_X x_f \approx Y_{f\sigma} \approx \sqrt{4\pi} \quad \text{EPJC81(2021)3,218, JHEP03(2020), 051} \quad 22/30$$

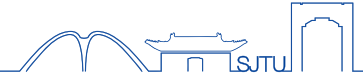
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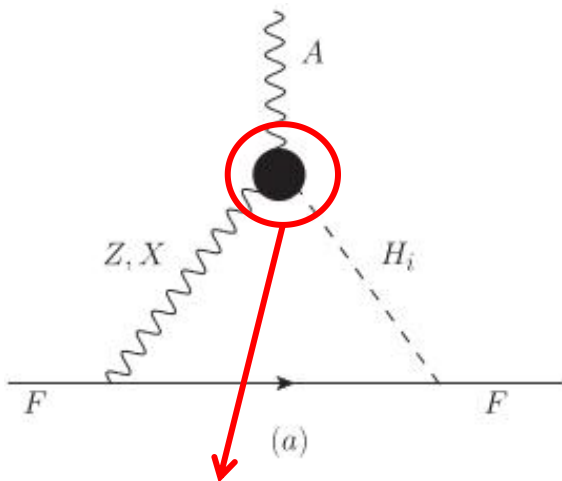
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EDM



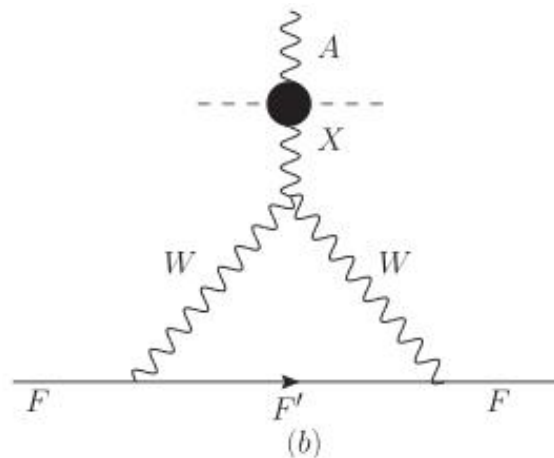
The EDM d_F for a SM fermion F



$\tilde{\beta}_X$ at one loop level



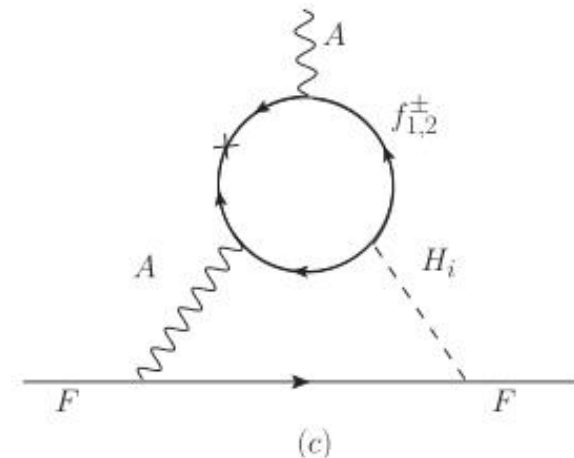
Two loop(**dominant**)



No A-X mixing



$d_q \approx 0$



Barr-Zee



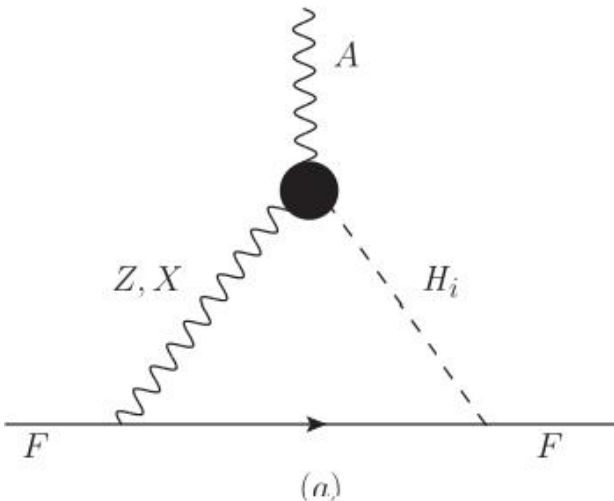
Cancel

EDM

$$L^{EDM} = -\frac{i}{2}d_F\bar{F}\sigma^{\mu\nu}\gamma_5FF_{\mu\nu}$$



$$d_F = \frac{e}{8\pi^2} \frac{m_F}{v} \sum_{i=[1, N-1]} \underbrace{V_{\Sigma i} V_{hi}}_{\text{mixing}} [C_Z V_Z^F f(m_Z^2/m_{H_i}^2, m_Z^2/m_{H_\Sigma}^2) + C_X V_X^F f(m_X^2/m_{H_i}^2, m_X^2/m_{H_\Sigma}^2)]$$



$$V_Z^F = (c_\xi - \epsilon_{ZX} s_\xi) \frac{g_Z^F}{c_W s_W} - Q_F \epsilon_{AX} s_\xi$$

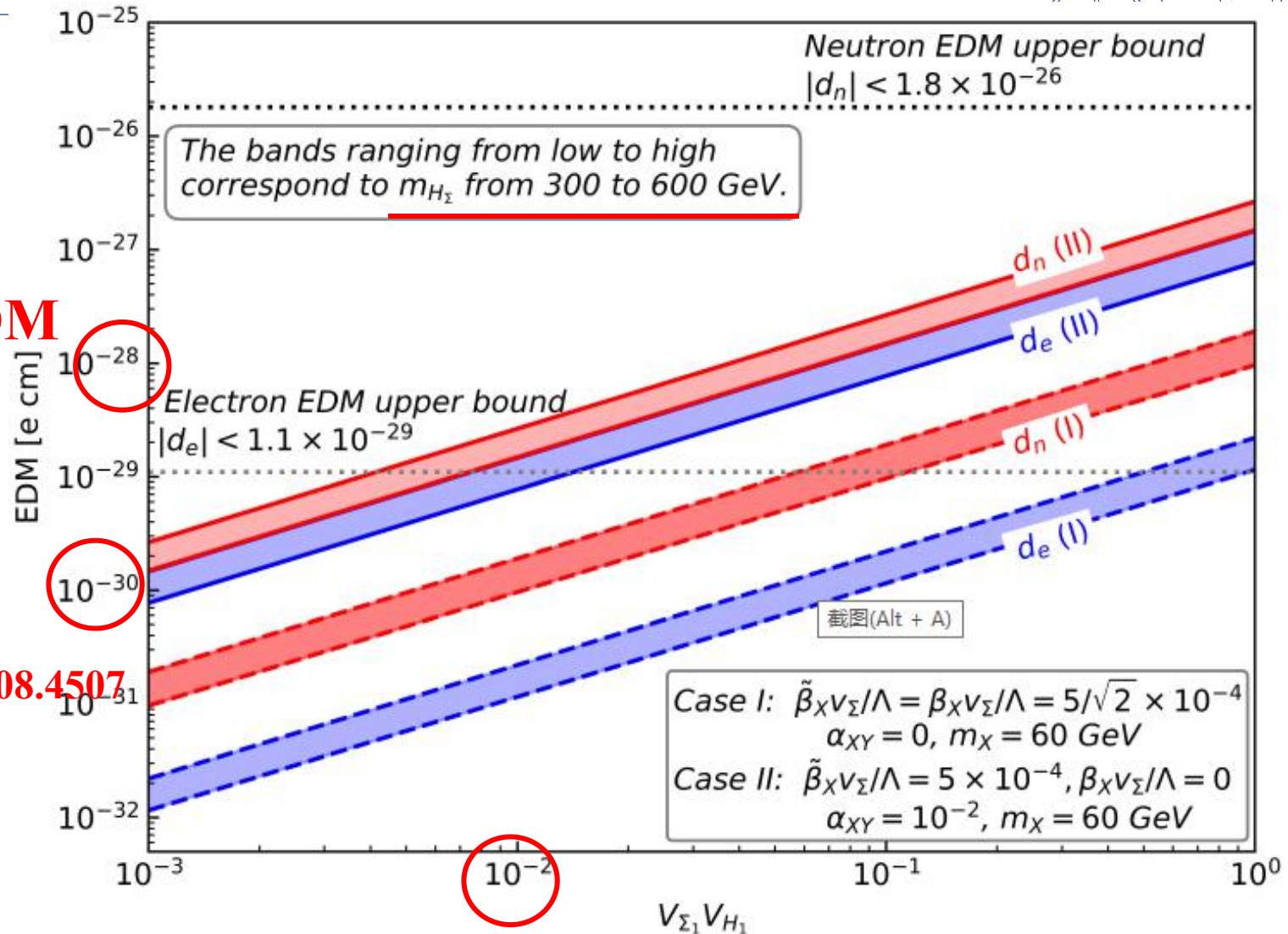
$$V_X^F = -(\underbrace{s_\xi}_{\text{red circle}} + \epsilon_{ZX} c_\xi) \frac{g_Z^F}{c_W s_W} - Q_F \epsilon_{AX} c_\xi$$

$$s_\xi \approx -\epsilon_{ZX} m_Z^2 / (m_Z^2 - m_X^2)$$

$$C_Z = \frac{\tilde{\beta}_X}{\Lambda} s_W s_\xi, \quad C_X = \frac{\tilde{\beta}_X}{\Lambda} s_W c_\xi, \quad g_Z^F = \frac{I_3^F}{2} - Q_F s_W^2$$

$$f(x, y) = (1/2)(\ln(y/x) - (x \ln x / (1-x) - y \ln y / (1-y)))$$

EDM



n2EDM

arXiv:1208.4507

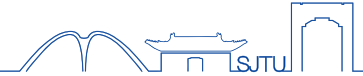
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Conclusions



- 1. We construct for the first time a renormalizable dark photon model with CP violating kinetic mixing in combination with the type-III seesaw model.**
- 2. We find that CP violating kinetic mixing induced interaction dominates the contribution to electron EDM which can be as large as experimental bound.**
- 3. The model provides a bridge connecting dark photon, neutrino physics and also CP violation, and can be directly tested by near future EDM measurements.**



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李政道先生
九十五华诞
The 95th Birthday of
Tsung-Dao Lee

<https://indico-tdli.sjtu.edu.cn/event/769/>



论道

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上海交通大学物理与天文学科
博士生学术创新论坛

SHANGHAI JIAO TONG UNIVERSITY PHYSICS AND ASTRONOMY
ACADEMIC INNOVATION FORUM FOR DOCTORAL STUDENTS



2021年12月10-13日

上海交通大学李政道研究所张江实验楼

李政道研究所
Tsung-Dao Lee Institute

SJTUPA
上海交通大学物理与天文学院



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Thanks!

