

Based on Phys.Rev. D101, 111901(R) (2020)

# Contributions of $K^+K^-$ from resonances $\rho$ and $\omega$ for the $B^\pm \rightarrow \pi^\pm K^+K^-$ decays

Wang Wen-Fei

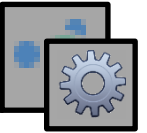
PQCD Group Meeting

2020-08-14



# Outline

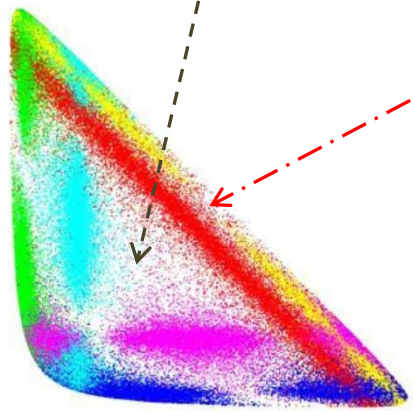
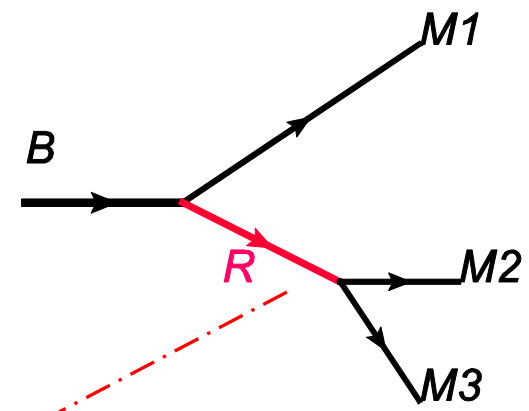
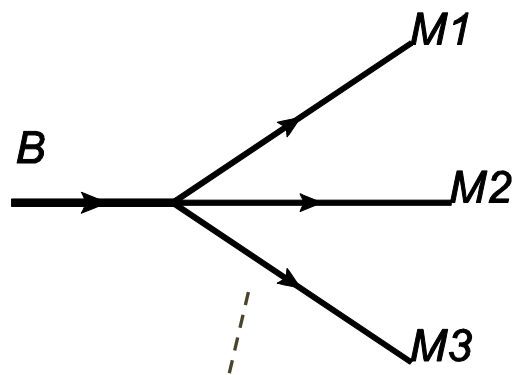
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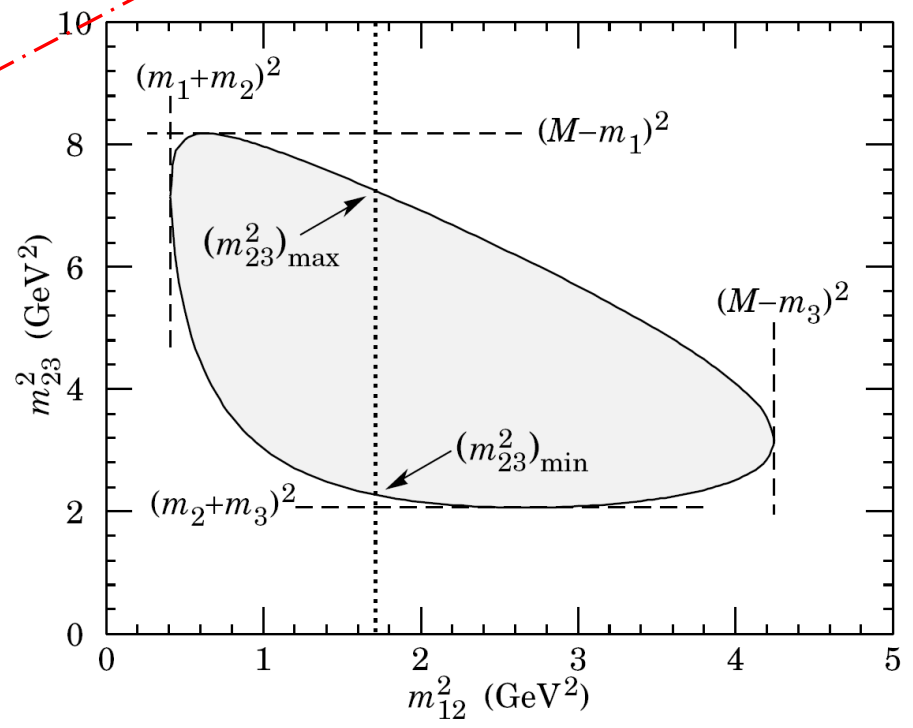
- Introduction for 3-body B decays
- Motivation for the  $\rho/\omega \rightarrow K^+ K^-$
- Results and discussions
- Summary



# Introduction for 3-body B decays



Dalitz Plot





## The weak effective Hamiltonian:

$$\mathcal{H}_{\text{eff}} = \frac{G_F}{\sqrt{2}} \left\{ V_{ub} V_{uq}^* [C_1(\mu) Q_1^u(\mu) + C_2(\mu) Q_2^u(\mu)] - V_{tb} V_{tq}^* \left[ \sum_{i=3}^{10} C_i(\mu) Q_i(\mu) \right] \right\} + \text{H.c.}, \quad (2)$$

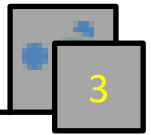
1.

where  $q = d, s$ . The functions  $Q_i$  ( $i = 1, \dots, 10$ ) are the local four-quark operators:

2.

Then what about the **other** strong interactions?

.....



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Physics Letters B 561 (2003) 258–265

PHYSICS LETTERS B

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## Three-body nonleptonic $B$ decays in perturbative QCD

Chuan-Hung Chen, Hsiang-Nan Li

### Abstract

We develop perturbative QCD formalism for three-body nonleptonic  $B$  meson decays. Leading contributions are identified by defining power counting rules for various topologies of amplitudes. The analysis is simplified into the one for two-body decays by introducing **two-meson distribution amplitudes**. This formalism predicts both nonresonant and resonant contributions, and can be generalized to baryonic decays.

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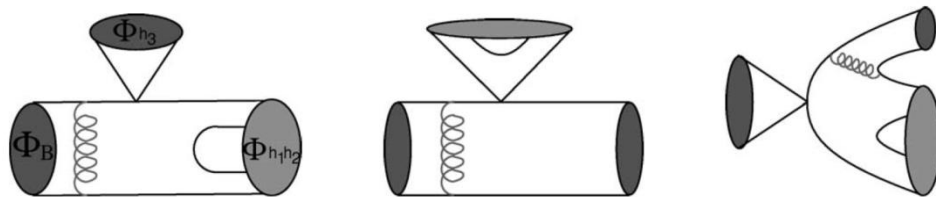
PHYSICS LETTERS B

Physics Letters B 561 (2003) 258–265

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## Three-body nonleptonic $B$ decays in perturbative QCD

Chuan-Hung Chen, Hsiang-Nan Li



$$B \rightarrow h_1 h_2 h_3$$

$$\Rightarrow \mathcal{M} = \Phi_B \otimes H \otimes \Phi_{h_1 h_2} \otimes \Phi_{h_3}$$

Physics Letters B 763 (2016) 29–39



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## Physics Letters B

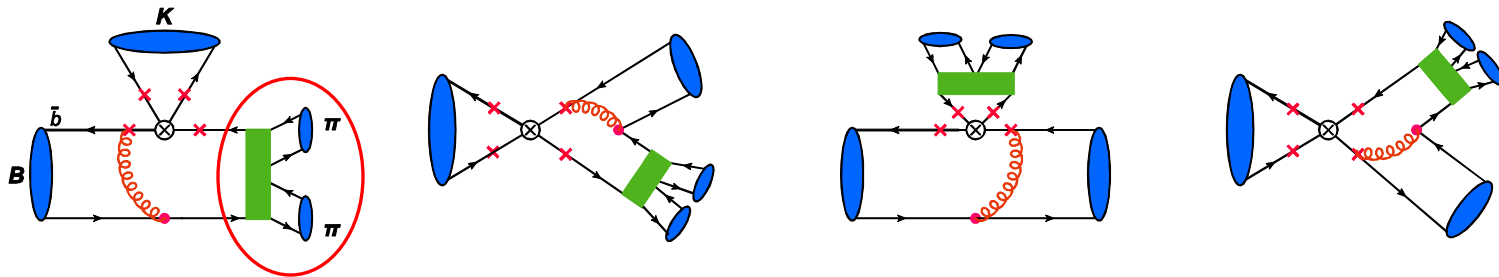
[www.elsevier.com/locate/physletb](http://www.elsevier.com/locate/physletb)



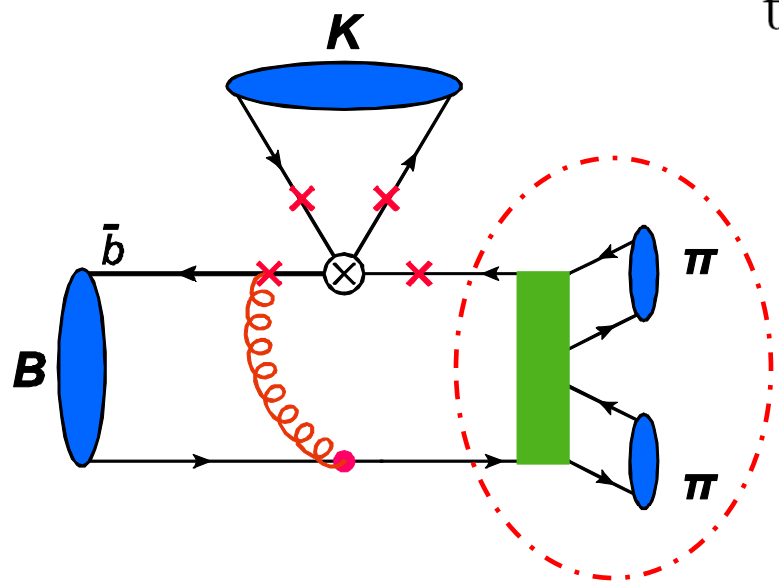
## Quasi-two-body decays $B \rightarrow K\rho \rightarrow K\pi\pi$ in perturbative QCD approach



Wen-Fei Wang<sup>a,b</sup>, Hsiang-nan Li<sup>a,\*</sup>



From the definition of the vector current time-like form factor  $F_\pi$  we have

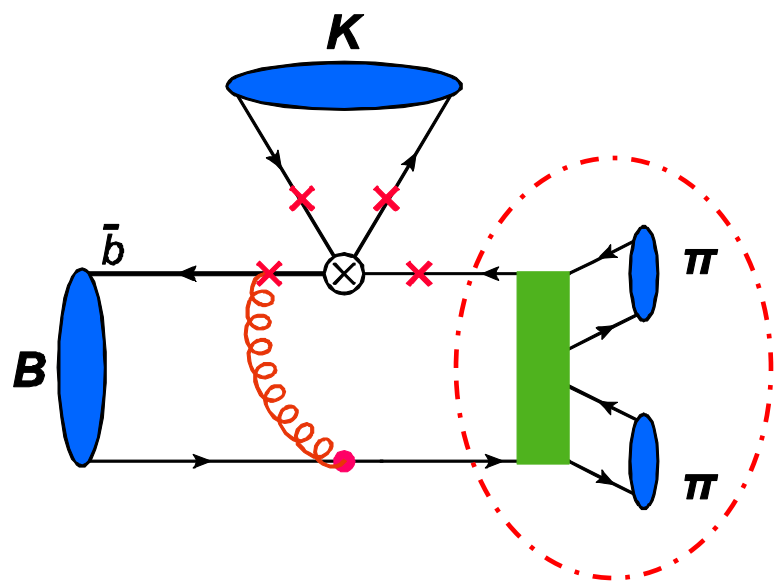


$$F_\pi^\rho(w^2) \approx \frac{g_{\rho\pi\pi} w f_\rho}{\sqrt{2} D_\rho(w^2)}$$



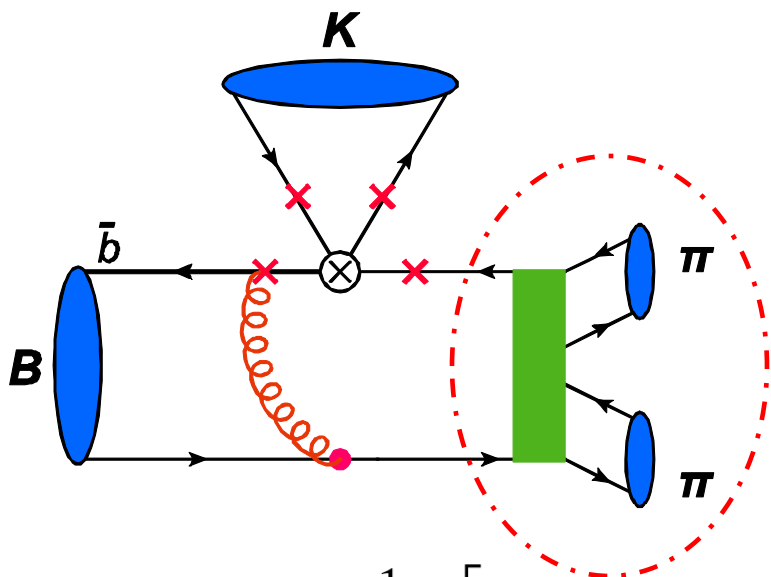
$$\mathcal{M} = \Phi_B \otimes H \otimes \Phi_K \otimes \left[ \Phi_\rho \frac{F_\pi^\rho(s)}{f_\rho} (2\zeta - 1) \right]$$





$$\mathcal{M} = \Phi_B \otimes H \otimes \Phi_K \otimes \left[ \Phi_\rho \frac{F_\pi^\rho(s)}{f_\rho} (2\zeta - 1) \right]$$

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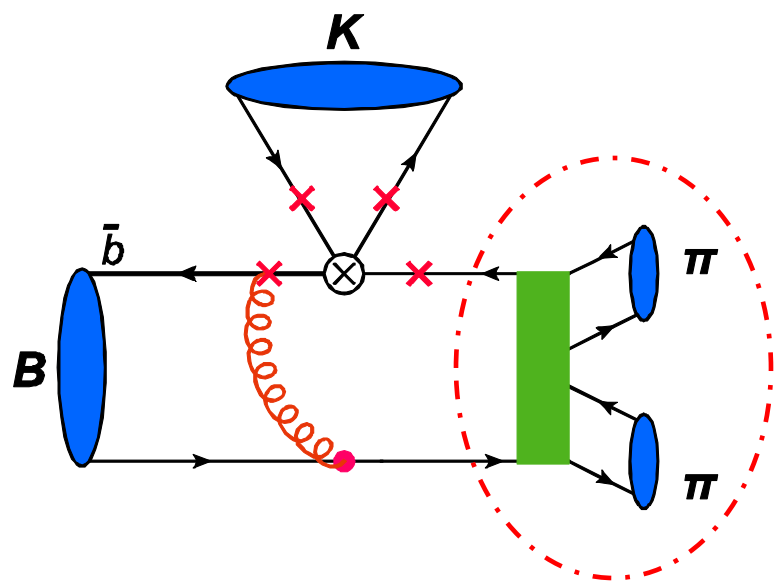
Legendre polynomial  $P_1(2\zeta - 1) = 2\zeta - 1$

$$\phi_{\pi\pi}^{l=1} = \frac{1}{\sqrt{2N_c}} \left[ \not{p} \phi_{vv=-}^{l=1}(z, \zeta, w^2) + w \phi_s^{l=1}(z, \zeta, w^2) + \frac{\not{p}_1 \not{p}_2 - \not{p}_2 \not{p}_1}{w(2\zeta - 1)} \phi_{tv=+}^{l=1}(z, \zeta, w^2) \right],$$

$$\phi_{vv=-}^{l=1}(z, \zeta, w^2) \equiv \phi^0(z, \zeta, w^2) = \frac{3F_\pi(w^2)}{\sqrt{2N_c}} z(1-z) \left[ 1 + a_2^0 C_2^{3/2}(1-2z) \right] P_1(2\zeta - 1),$$

$$\phi_s^{l=1}(z, \zeta, w^2) \equiv \phi^s(z, \zeta, w^2) = \frac{3F_s(w^2)}{2\sqrt{2N_c}} (1-2z) \left[ 1 + a_2^s (1 - 10z + 10z^2) \right] P_1(2\zeta - 1),$$

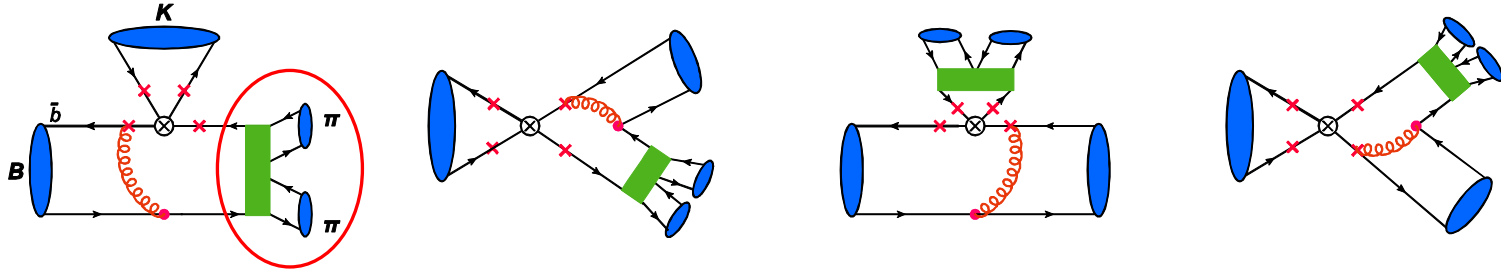
$$\phi_{tv=+}^{l=1}(z, \zeta, w^2) \equiv \phi^t(z, \zeta, w^2) = \frac{3F_t(w^2)}{2\sqrt{2N_c}} (1-2z)^2 \left[ 1 + a_2^t C_2^{3/2}(1-2z) \right] P_1(2\zeta - 1),$$



$$F_{\pi}(s) = \frac{\text{GS}_{\rho}(s, m_{\rho}, \Gamma_{\rho}) \frac{1+c_{\omega} \text{BW}_{\omega}(s, m_{\omega}, \Gamma_{\omega})}{1+c_{\omega}} + \sum c_i \text{GS}_i(s, m_i, \Gamma_i)}{1 + \sum c_i}$$

G. Gounaris, J.J. Sakurai  
PRL21-244(1968)

$$\frac{m_{\rho}^2 + dm_{\rho} \Gamma_{\rho}}{(m_{\rho}^2 - s) + \Gamma_{\rho} (m_{\rho}^2 / k_{\rho}^3) \{ k^2 [h(s) - h(m_{\rho}^2)] + k_{\rho}^2 h'(m_{\rho}^2) (m_{\rho}^2 - s) \} - im_{\rho} \Gamma_{\rho} (k/k_{\rho})^3 (m_{\rho} / \sqrt{s})}$$



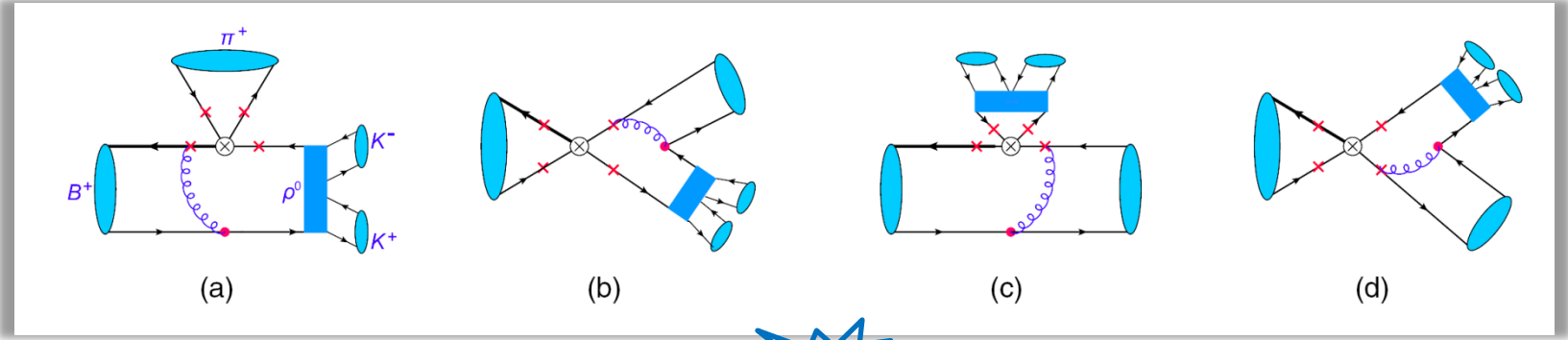
$$\phi_{\pi\pi}^{I=1} = \frac{1}{\sqrt{2N_c}} \left[ \not{p} \phi_{vv=-}^{I=1}(z, \zeta, w^2) + w \phi_s^{I=1}(z, \zeta, w^2) + \frac{\not{p}_1 \not{p}_2 - \not{p}_2 \not{p}_1}{w(2\zeta - 1)} \phi_{tv=+}^{I=1}(z, \zeta, w^2) \right],$$

$$\phi_{vv=-}^{I=1}(z, \zeta, w^2) \equiv \phi^0(z, \zeta, w^2) = \frac{3F_\pi(w^2)}{\sqrt{2N_c}} z(1-z) \left[ 1 + a_2^0 C_2^{3/2}(1-2z) \right] P_1(2\zeta - 1),$$

$$\phi_s^{I=1}(z, \zeta, w^2) \equiv \phi^s(z, \zeta, w^2) = \frac{3F_s(w^2)}{2\sqrt{2N_c}} (1-2z) \left[ 1 + a_2^s (1 - 10z + 10z^2) \right] P_1(2\zeta - 1),$$

$$\phi_{tv=+}^{I=1}(z, \zeta, w^2) \equiv \phi^t(z, \zeta, w^2) = \frac{3F_t(w^2)}{2\sqrt{2N_c}} (1-2z)^2 \left[ 1 + a_2^t C_2^{3/2}(1-2z) \right] P_1(2\zeta - 1),$$

→  $a_\rho^0 = 0.25, a_\rho^t = -0.60, a_\rho^s = 0.75$



**2.**

**1.**

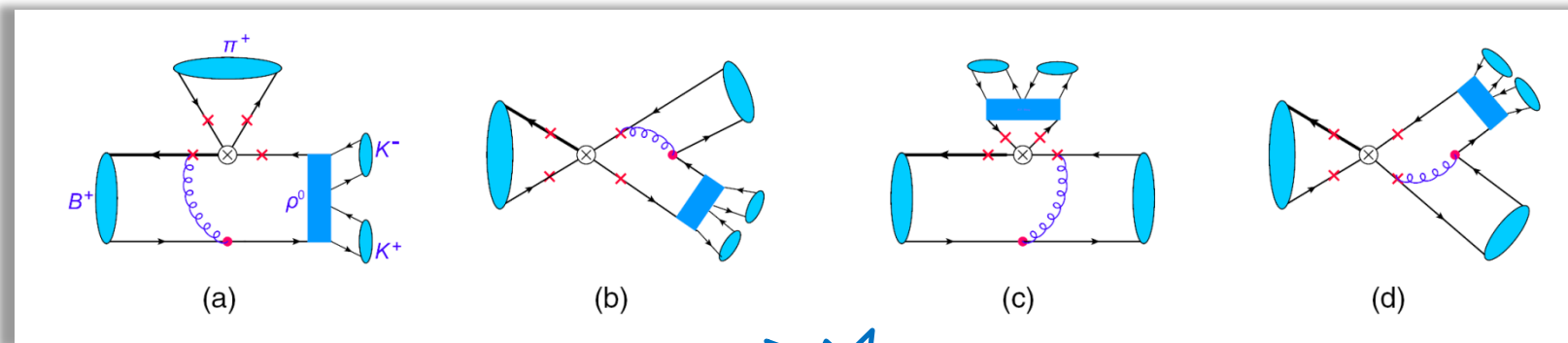
$$\frac{d\mathcal{B}}{d\eta} = \tau_B \frac{|\vec{q}_\pi|^3 |\vec{q}_K|^3}{12\pi^3 m_B^5} |\mathcal{A}|^2,$$

*In the rest frame of KK pair!*



$$|\vec{q}_K| = \frac{1}{2} \sqrt{s - 4m_K^2},$$

$$|\vec{q}_\pi| = \frac{1}{2\sqrt{s}} \sqrt{(m_B^2 - m_\pi^2)^2 - 2(m_B^2 + m_\pi^2)s + s^2},$$



**2.**

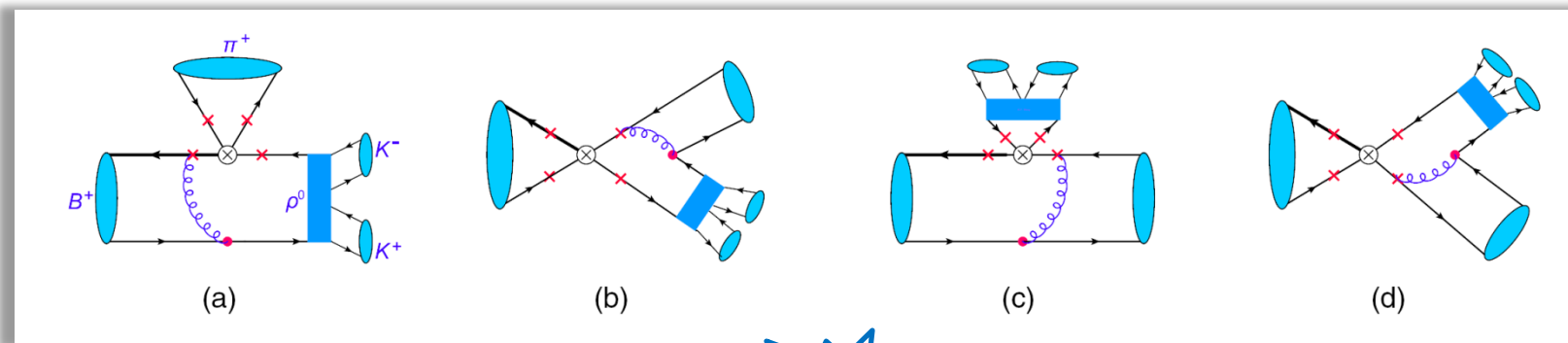
$$\frac{d\mathcal{B}}{d\eta} = \tau_B \frac{|\vec{q}_\pi|^3 |\vec{q}_K|^3}{12\pi^3 m_B^5} |\mathcal{A}|^2,$$

$$\Delta_{K\pi} = (m_K^2 - m_\pi^2)$$

$$\langle K^+(p_1)\pi^-(p_2) | \bar{d}\gamma_\mu(1 - \gamma_5)s | 0 \rangle = \left[ (p_1 - p_2)_\mu - \frac{\Delta_{K\pi}}{p^2} p_\mu \right] F_+^{K\pi}(s) + \frac{\Delta_{K\pi}}{p^2} p_\mu F_0^{K\pi}(s),$$

$$\langle \pi(q_2) | \bar{b}(0)\gamma_\mu(1 - \gamma_5)q(0) | B(q_1) \rangle = \left[ (q_1 + q_2)_\mu - \frac{m_B^2 - m_\pi^2}{q^2} q_\mu \right] F_+^{B\pi}(q^2) + \frac{m_B^2 - m_\pi^2}{q^2} q_\mu F_0^{B\pi}(q^2)$$

$$q = q_1 - q_2$$



2.

$$\frac{d\mathcal{B}}{d\eta} = \tau_B \frac{|\vec{q}_\pi|^3 |\vec{q}_K|^3}{12\pi^3 m_B^5} |\mathcal{A}|^2,$$

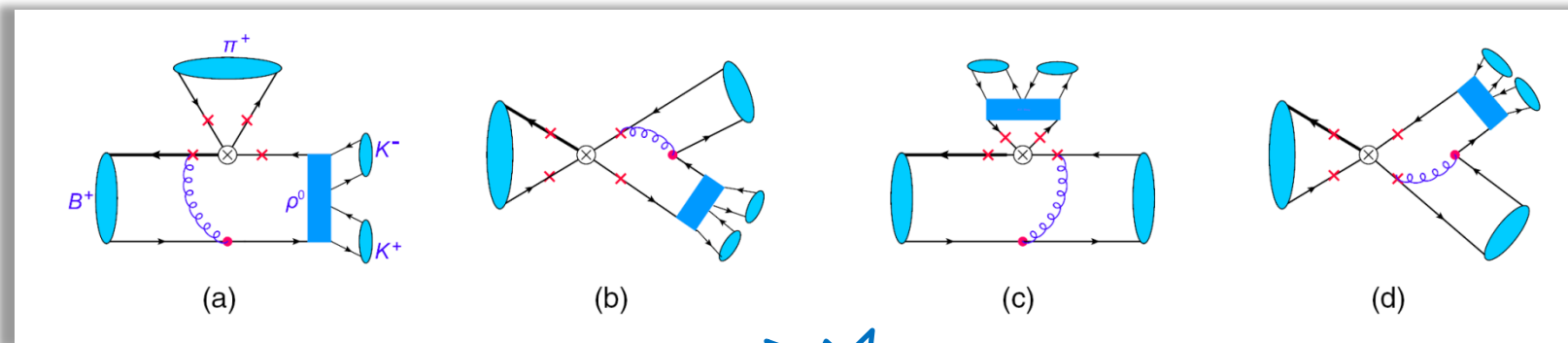
$$\Delta_{K\pi} = (m_K^2 - m_\pi^2)$$

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$$q = q_1 - q_2$$



**2.**

$$\frac{d\mathcal{B}}{d\eta} = \tau_B \frac{|\vec{q}_\pi|^3 |\vec{q}_K|^3}{12\pi^3 m_B^5} |\mathcal{A}|^2,$$

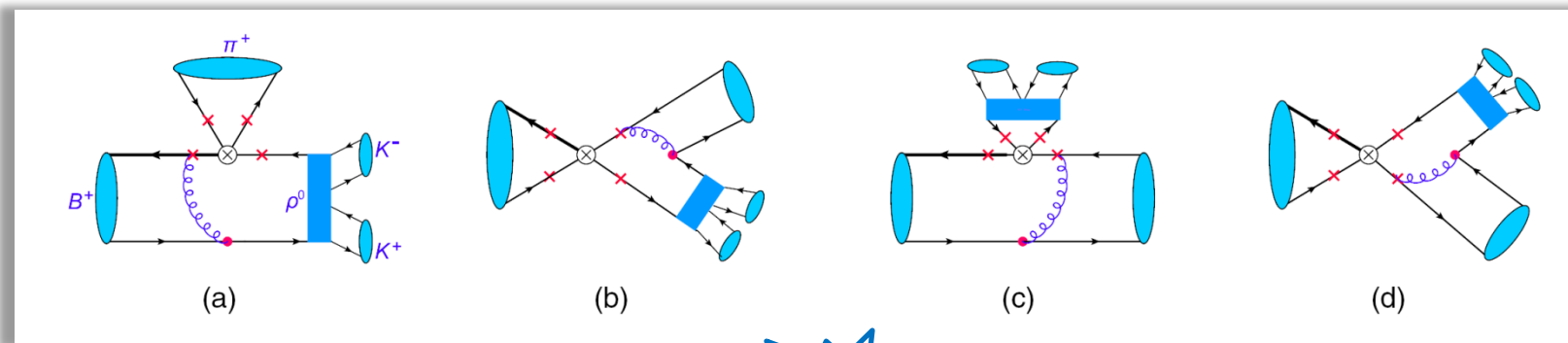
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$$\langle \pi(q_2) | \bar{b}(0)\gamma_\mu(1 - \gamma_5)q(0) | B(q_1) \rangle = \left[ (q_1 + q_2)_\mu - \frac{m_B^2 - m_\pi^2}{q^2} q_\mu \right] F_+^{B\pi}(q^2) - \frac{m_B^2 - m_\pi^2}{q^2} q_\mu F_0^{B\pi}(q^2)$$

$$q = q_1 - q_2$$





**2.**

$$\frac{d\mathcal{B}}{d\eta} = \tau_B \frac{|\vec{q}_\pi|^3 |\vec{q}_K|^3}{12\pi^3 m_B^5} |\mathcal{A}|^2,$$

$$m_{K^-\pi^-}^2 - m_{\pi^+\pi^-}^2 - (M_B^2 - m_\pi^2) \frac{m_K^2 - m_\pi^2}{q^2} = 4\mathbf{p}_{\pi^+} \cdot \mathbf{p}_{\pi^-} = 4|\mathbf{p}_{\pi^+}| |\mathbf{p}_{\pi^-}| \cos\theta,$$

$$\langle K^+(p_1)\pi^-(p_2) | \bar{d}\gamma_\mu(1 - \gamma_5)s | 0 \rangle = \left[ (p_1 - p_2)_\mu - \frac{\Delta_{K\pi}}{p^2} p_\mu \right] F_+^{K\pi}(s) + \frac{\Delta_{K\pi}}{p^2} p_\mu F_0^{K\pi}(s),$$

$$\langle \pi(q_2) | \bar{b}(0)\gamma_\mu(1 - \gamma_5)q(0) | B(q_1) \rangle = \left[ (q_1 + q_2)_\mu - \frac{m_B^2 - m_\pi^2}{q^2} q_\mu \right] F_+^{B\pi}(q^2) - \frac{m_B^2 - m_\pi^2}{q^2} q_\mu F_0^{B\pi}(q^2)$$

PHYSICAL REVIEW D **79**, 094005 (2009)



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## Physics Letters B

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## Virtual contributions from $D^*(2007)^0$ and $D^*(2010)^\pm$ in the $B \rightarrow D\pi h$ decays



Wen-Fei Wang<sup>a,b,\*</sup>, Jian Chai<sup>a,b</sup>

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<sup>b</sup> State Key Laboratory of Quantum Optics and Quantum Optics Devices, Shanxi University, Taiyuan, Shanxi 030006, China

**Table 2**

The PQCD predictions of the virtual contributions from  $D^*$  state in the  $D\pi$  invariant mass region  $\sqrt{s} > 2.1$  GeV for the  $B \rightarrow D^*h \rightarrow D\pi h$  decays.

Mode	Unit	$\mathcal{B}_v$
$\bar{B}^0 \rightarrow D^{*+}\pi^- \rightarrow D^0\pi^+\pi^-$	$(10^{-4})$	$0.87_{-0.27}^{+0.43}(\omega_B)_{-0.07}^{+0.08}(f_{D^*})_{-0.06}^{+0.08}(a_{D\pi}) \pm 0.03(A)_{-0.01}^{+0.02}(\omega_{D\pi})$
$\bar{B}^0 \rightarrow D^{*+}K^- \rightarrow D^0\pi^+K^-$	$(10^{-5})$	$0.72_{-0.22}^{+0.35}(\omega_B)_{-0.06}^{+0.07}(f_{D^*}) \pm 0.06(a_{D\pi}) \pm 0.03(A) \pm 0.02(\omega_{D\pi})$
$B^- \rightarrow D^{*0}\pi^- \rightarrow D^+\pi^-\pi^-$	$(10^{-4})$	$1.91_{-0.59}^{+0.86}(\omega_B)_{-0.16}^{+0.17}(f_{D^*})_{-0.10}^{+0.12}(a_{D\pi}) \pm 0.07(A)_{-0.05}^{+0.04}(\omega_{D\pi})$
$B^- \rightarrow D^{*0}K^- \rightarrow D^+\pi^-K^-$	$(10^{-5})$	$1.48_{-0.46}^{+0.65}(\omega_B) \pm 0.13(f_{D^*})_{-0.08}^{+0.09}(a_{D\pi}) \pm 0.05(A)_{-0.03}^{+0.02}(\omega_{D\pi})$



# Motivation for the $\rho/\omega \rightarrow K^+ K^-$

Physics Letters B 791 (2019) 342–350

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## Physics Letters B

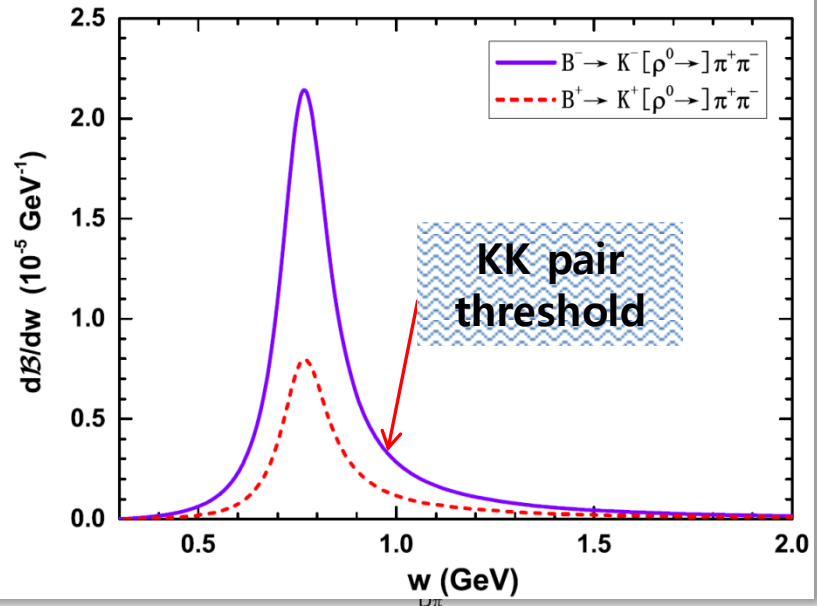
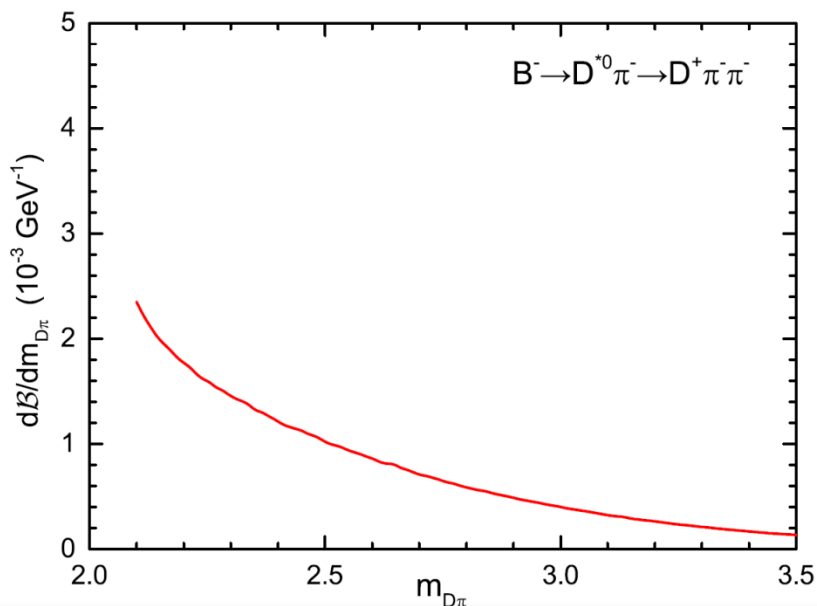
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### Virtual contributions from $D^*(2007)^0$ and $D^*(2010)^\pm$ in the $B \rightarrow D\pi h$



W.-F. Wang, J. Chai / Physics Letters B 791 (2019) 342–350





# Motivation for the $\rho/\omega \rightarrow K^+ K^-$

$$\rho(770,1450) \rightarrow \pi^+ \pi^- \quad \checkmark$$

$$\rho/\omega \rightarrow K^+ K^- \quad ?!$$

$\pi^- p \rightarrow K^- K^+ n$   
 $\pi^+ n \rightarrow K^- K^+ p$   
 $e^+ e^- \rightarrow K^+ K^-$



PHYSICAL REVIEW LETTERS **123**, 231802 (2019)

## Amplitude Analysis of $B^\pm \rightarrow \pi^\pm K^+ K^-$ Decays

R. Aaij *et al.*\*  
(LHCb Collaboration)

 (Received 12 June 2019; revised manuscript received 15 October 2019; published 6 December 2019)

TABLE I. Results of the Dalitz plot fit, where the first uncertainty is statistical and the second systematic. The fitted values of  $c_i$  ( $\bar{c}_i$ ) are expressed in terms of magnitudes  $|c_i|$  ( $|\bar{c}_i|$ ) and phases  $\arg(c_i)$  [ $\arg(\bar{c}_i)$ ] for each  $B^+$  ( $B^-$ ) contribution. The top row corresponds to  $B^+$  and the bottom to  $B^-$  mesons.

Contribution	Fit fraction (%)	$A_{CP}$ (%)	Magnitude ( $B^+/B^-$ )	Phase [ $^\circ$ ] ( $B^+/B^-$ )
$K^*(892)^0$	$7.5 \pm 0.6 \pm 0.5$	$+12.3 \pm 8.7 \pm 4.5$	$0.94 \pm 0.04 \pm 0.02$	0 (fixed)
$K_0^*(1430)^0$	$4.5 \pm 0.7 \pm 1.2$	$+10.4 \pm 14.9 \pm 8.8$	$0.74 \pm 0.09 \pm 0.09$	$-176 \pm 10 \pm 16$
Single pole	$32.3 \pm 1.5 \pm 4.1$	$-10.7 \pm 5.3 \pm 3.5$	$0.82 \pm 0.09 \pm 0.10$	$136 \pm 11 \pm 21$
$\rho(1450)^0$	$30.7 \pm 1.2 \pm 0.9$	$-10.9 \pm 4.4 \pm 2.4$	$2.19 \pm 0.13 \pm 0.17$	$-138 \pm 7 \pm 5$
$f_2(1270)$	$7.5 \pm 0.8 \pm 0.7$	$+26.7 \pm 10.2 \pm 4.8$	$1.97 \pm 0.12 \pm 0.20$	$166 \pm 6 \pm 5$
Rescattering	$16.4 \pm 0.8 \pm 1.0$	$-66.4 \pm 3.8 \pm 1.9$	$2.14 \pm 0.11 \pm 0.07$	$-175 \pm 10 \pm 15$
$\phi(1020)$	$0.3 \pm 0.1 \pm 0.1$	$+9.8 \pm 43.6 \pm 26.6$	$1.92 \pm 0.10 \pm 0.07$	$140 \pm 13 \pm 20$
			$0.86 \pm 0.09 \pm 0.07$	$-106 \pm 11 \pm 10$
			$1.13 \pm 0.08 \pm 0.05$	$-128 \pm 11 \pm 14$
			$1.91 \pm 0.09 \pm 0.06$	$-56 \pm 12 \pm 18$
			$0.86 \pm 0.07 \pm 0.04$	$-81 \pm 14 \pm 15$
			$0.20 \pm 0.07 \pm 0.02$	$-52 \pm 23 \pm 32$
			$0.22 \pm 0.06 \pm 0.04$	$107 \pm 33 \pm 41$



# Motivation for the $\rho/\omega \rightarrow K^+ K^-$

PHYSICAL REVIEW LETTERS **123**, 231802 (2019)

## Amplitude Analysis of $B^\pm \rightarrow \pi^\pm K^+ K^-$ Decays

R. Aaij *et al.*\*  
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TABLE I. Results of the Dalitz plot fit, where the first uncertainty is statistical and the second systematic. The fitted values of  $c_i$  ( $\bar{c}_i$ ) are expressed in terms of magnitudes  $|c_i|$  ( $|\bar{c}_i|$ ) and phases  $\arg(c_i)$  [ $\arg(\bar{c}_i)$ ] for each  $B^+$  ( $B^-$ ) contribution. The top row corresponds to  $B^+$  and the bottom to  $B^-$  mesons.

Contribution	$ c_i $	$ \bar{c}_i $	$\arg(c_i)$	$\arg(\bar{c}_i)$	$\int  c_i ^2$	$\int  \bar{c}_i ^2$
$K^*(892)^0$	$\rho(1450)^0 \pi^+, \rho^0 \rightarrow \pi^+ \pi^-$	$(1.4 \pm 0.6 \pm 0.9) \times 10^{-6}$			2434	—
$K_0^*(1430)^0$	$\rho(1450)^0 \pi^+, \rho^0 \rightarrow K^+ K^-$	$(1.60 \pm 0.14) \times 10^{-6}$			—	—
Single pole						
$\rho(1450)^0$	$30.7 \pm 1.2 \pm 0.9$	$4.5 \pm 0.7 \pm 1.2$	$-10.9 \pm 4.4 \pm 2.4$	$+10.4 \pm 14.9 \pm 8.8$	$2.14 \pm 0.11 \pm 0.07$	$1.97 \pm 0.12 \pm 0.20$
$f_2(1270)$	$7.5 \pm 0.8 \pm 0.7$		$+26.7 \pm 10.2 \pm 4.8$		$0.86 \pm 0.09 \pm 0.07$	$1.13 \pm 0.08 \pm 0.05$
Rescattering	$16.4 \pm 0.8 \pm 1.0$		$-66.4 \pm 3.8 \pm 1.9$		$1.91 \pm 0.09 \pm 0.06$	$0.86 \pm 0.07 \pm 0.04$
$\phi(1020)$	$0.3 \pm 0.1 \pm 0.1$		$+9.8 \pm 43.6 \pm 26.6$		$0.20 \pm 0.07 \pm 0.02$	$0.22 \pm 0.06 \pm 0.04$

PDG2020

Eur. Phys. J. C 39, 41–54 (2005)  
 Digital Object Identifier (DOI) 10.1140/epjc/s2004-02064-3

**THE EUROPEAN  
 PHYSICAL JOURNAL C**

## Modeling the pion and kaon form factors in the timelike region

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Received: 7 October 2004 /

$$\begin{aligned}
 & F_{K^+}(s) \\
 &= \frac{1}{2} (c_\rho^K \text{BW}_\rho(s) + c_{\rho'}^K \text{BW}_{\rho'}(s) + c_{\rho''}^K \text{BW}_{\rho''}(s)) \\
 &+ \frac{1}{6} (c_\omega^K \text{BW}_\omega(s) + c_{\omega'}^K \text{BW}_{\omega'}(s) + c_{\omega''}^K \text{BW}_{\omega''}(s)) \\
 &+ \frac{1}{3} (c_\phi \text{BW}_\phi(s) + c_{\phi'} \text{BW}_{\phi'}(s)), \tag{64}
 \end{aligned}$$

$$\begin{aligned}
 g_{\rho^0 \pi^+ \pi^-} &\equiv g_{\rho \pi \pi} = \sqrt{2} A_{qq}, \\
 g_{\rho^0 K^+ K^-} &= g_{\omega K^+ K^-} = \frac{1}{\sqrt{2}} A_{qs}, \\
 g_{\rho^0 K^0 \bar{K}^0} &= -g_{\omega K^0 \bar{K}^0} = -\frac{1}{\sqrt{2}} A_{qs}, \\
 g_{\phi K^+ K^-} &= g_{\phi K^0 \bar{K}^0} = -A_{sq}. \tag{59}
 \end{aligned}$$

## Will the subprocesses $\rho(770, 1450)^0 \rightarrow K^+ K^-$ contribute large branching fractions for $B^\pm \rightarrow \pi^\pm K^+ K^-$ decays?

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$$\begin{aligned} \mathcal{B}(B^\pm \rightarrow \pi^\pm \rho(1450)^0 \rightarrow \pi^\pm K^+ K^-) \\ = (8.96 \pm 1.58(a_2^s + a_2^0 + a_2^t) \pm 0.83(\omega_B) \\ \pm 0.79(m_0^\pi + a_2^\pi) \pm 1.73(c_{\rho(1450)}^K)) \times 10^{-8}, \end{aligned}$$

$$\begin{aligned} \mathcal{B}_v(B^\pm \rightarrow \pi^\pm \rho(770)^0 \rightarrow \pi^\pm K^+ K^-) \\ = (1.31 \pm 0.22(a_2^s + a_2^0 + a_2^t) \pm 0.12(\omega_B) \\ \pm 0.11(m_0^\pi + a_2^\pi) \pm 0.02(c_{\rho(770)}^K)) \times 10^{-7}. \end{aligned}$$



## Will the subprocesses $\rho(770, 1450)^0 \rightarrow K^+ K^-$ contribute large branching fractions for $B^\pm \rightarrow \pi^\pm K^+ K^-$ decays?

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$\rho(1450)^0 \pi^+, \rho^0 \rightarrow \pi^+ \pi^-$	$(1.4 \pm_{-0.9}^{0.6}) \times 10^{-6}$	PDG2020 <sup>2434</sup>
$\rho(1450)^0 \pi^+, \rho^0 \rightarrow K^+ K^-$	$(1.60 \pm 0.14) \times 10^{-6}$	

$$\begin{aligned} & \mathcal{B}(B^\pm \rightarrow \pi^\pm \rho(1450)^0 \rightarrow \pi^\pm \pi^+ \pi^-) \\ &= (9.97 \pm 1.81(a_2^s + a_2^0 + a_2^t) \pm 0.98(\omega_B) \\ & \quad \pm 0.91(m_0^\pi + a_2^\pi)) \times 10^{-7}, \end{aligned}$$



$$g_{\rho(1450)^0 K^+ K^-} \approx \frac{1}{2} g_{\rho(1450)^0 \pi^+ \pi^-}$$

$$|g_{VPP}| = \left[ \frac{6 \pi m_V^2 \Gamma_{VB}(V \rightarrow PP)}{q_{PP}^3(m_V)} \right]^{1/2}$$

$$\begin{aligned} R_{\rho(1450)} &= \frac{\mathcal{B}(\rho(1450)^0 \rightarrow K^+ K^-)}{\mathcal{B}(\rho(1450)^0 \rightarrow \pi^+ \pi^-)} \\ &\approx \frac{g_{\rho(1450)^0 K^+ K^-}^2 (m_{\rho(1450)}^2 - 4m_K^2)^{3/2}}{g_{\rho(1450)^0 \pi^+ \pi^-}^2 (m_{\rho(1450)}^2 - 4m_\pi^2)^{3/2}} \\ &= 0.107. \end{aligned}$$

PHYS. REV. D **101**, 111901 (2020)

$$R_{\rho(1450)} = \frac{\mathcal{B}(\rho(1450)^0 \rightarrow K^+ K^-)}{\mathcal{B}(\rho(1450)^0 \rightarrow \pi^+ \pi^-)} = 0.090 \pm 0.017,$$

$$\Gamma(B^\pm \rightarrow \rho^0 \pi^\pm \rightarrow h^+ h^- \pi^\pm) \approx \Gamma(B^+ \rightarrow \rho^0 \pi^\pm) \mathcal{B}(\rho^0 \rightarrow h^+ h^-)$$



PHYSICAL REVIEW D **95**, 072007 (2017)

## Dalitz plot analyses of $J/\psi \rightarrow \pi^+ \pi^- \pi^0$ , $J/\psi \rightarrow K^+ K^- \pi^0$ , and $J/\psi \rightarrow K_s^0 K^\pm \pi^\mp$ produced via $e^+ e^-$ annihilation with initial-state radiation

J. P. Lees,<sup>1</sup> V. Poireau,<sup>1</sup> V. Tisserand,<sup>1</sup> E. Grauges,<sup>2</sup> A. Palano,<sup>3,†</sup> G. Eigen,<sup>4</sup> D. N. Brown,<sup>5</sup> Yu. G. Kolomensky,<sup>5</sup> M. Fritsch,<sup>6</sup>

(*BABAR* Collaboration)

$$R_{\rho(1450)} = \frac{\mathcal{B}(\rho(1450)^0 \rightarrow K^+ K^-)}{\mathcal{B}(\rho(1450)^0 \rightarrow \pi^+ \pi^-)} = 0.307 \pm 0.084(\text{stat}) \pm 0.082(\text{sys}),$$

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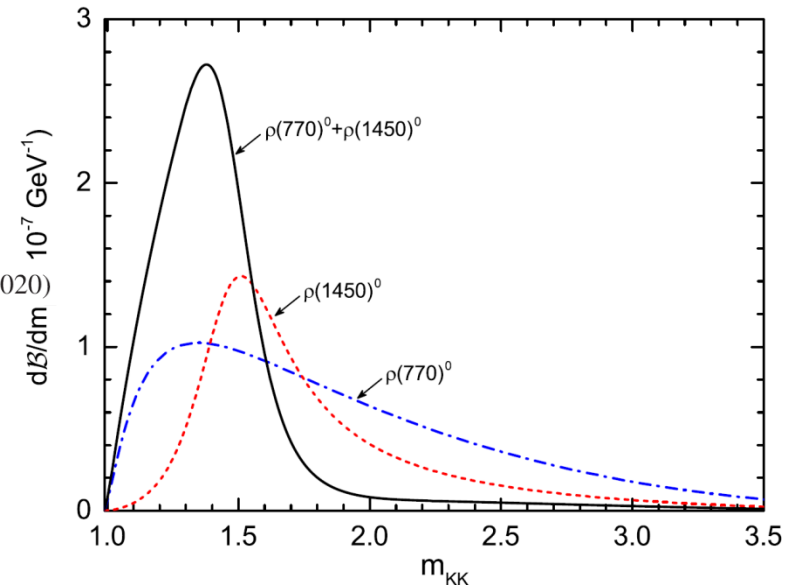
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$$\begin{aligned}
 R_{\rho(1450)} &= \frac{\mathcal{B}(\rho(1450)^0 \rightarrow K^+ K^-)}{\mathcal{B}(\rho(1450)^0 \rightarrow \pi^+ \pi^-)} \\
 &= 0.307 \pm 0.084(\text{stat}) \pm 0.082(\text{sys}),
 \end{aligned}$$

$$\begin{aligned}
 B^+ &\rightarrow \pi^+ \omega(782) \rightarrow \pi^+ K^+ K^- \\
 \mathcal{B}_v &\approx 2.32 \times 10^{-8}
 \end{aligned}$$

PHYS. REV. D **101**, 111901 (2020)

$$\begin{aligned}
 R_{\rho(1700)} &= \frac{\mathcal{B}(\rho(1700)^0 \rightarrow K^+ K^-)}{\mathcal{B}(\rho(1700)^0 \rightarrow \pi^+ \pi^-)} \\
 &\approx \frac{g_{\rho(1700)^0 K^+ K^-}^2 (m_{\rho(1700)}^2 - 4m_K^2)^{3/2}}{g_{\rho(1700)^0 \pi^+ \pi^-}^2 (m_{\rho(1700)}^2 - 4m_\pi^2)^{3/2}} \\
 &= 0.143,
 \end{aligned}$$





arXiv:2007.02558

**Branching Fractions and  $CP$  Violation  
in  $B^- \rightarrow K^+ K^- \pi^-$  and  $B^- \rightarrow \pi^+ \pi^- \pi^-$  Decays**

Hai-Yang Cheng<sup>1</sup>, Chun-Khiang Chua<sup>2</sup>

arXiv:2007.13141

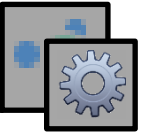
**Is  $f_X(1500)$  observed in the  $B \rightarrow \pi(K)KK$  decays  $\rho^0(1450)$ ?**

Zhi-Tian Zou<sup>1,\*</sup>, Ying Li<sup>1,2,†</sup> and Hsiang-nan Li<sup>3,‡</sup>



# Summary

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- ✓ Introduction for 3-body B decays
- ✓ Motivation for the  $\rho/\omega \rightarrow K^+ K^-$
- ✓ Results and discussions

*Thank You*