# XYZ states：window to subatomic structure 

Sun Zhentian（孙振田）

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## Frontier of human knowledge



Micro-scale: What are we composed of? Whether matter is infinitely dividable?


Cosmic scale: Where are we from? Where to go?

## Well accommodated Atoms



## Quark model




Murray Gell-Mann 1929-2019
'quark'

1964
u,d, s quarks Proposed, SU(3)


George Zweig 1937-
'Ace'

Meson:
Integer Spin two quarks

Baryon:
Half Integer Spin Three quarks


## Discovery of $c$ quark



SU(3) Octet of Vector $\rightarrow$ SU(4) 16-plet

$\mathrm{SU}(3)$ Octet of hadron $\rightarrow \mathrm{SU}(4)$ 20-plet



Burton Richter


## exotic states - XYZ states

- The exotic states: Molecule, Tetraquark, Hadro-quarkonium, Glueball, Hybrid, Penta-Quark
- Here are two pics from two papers

N. Brambilla, S. Eidelman, C. Hanhart et al. / Physics Reports 873 (2020) 1-154


> | Part I: Y states |
| :--- |
| $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \psi\left(1^{--}\right)($well established charmonium $) \rightarrow$... |
| or |
| $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \mathrm{Y}\left(1^{-}\right)($Charmonium like, maybe exotic $) \rightarrow$... |



## $Y(4260), Y(4360)$ : some history




Y(4360) PDG value without BES result:
mass $=4346 \pm 6 \mathrm{MeV}$
Width $=102 \pm 10 \mathrm{MeV}$

$$
Y(4260) \sim Y(4230) \sim Y(4220)
$$

## Why $Y(4260)$ exotic?

$\square Y(4260), Y(4360)$ don't have corresponding level with ( $c \bar{c}$ ) potential model. $\mathrm{J} / \psi(1 \mathrm{~S}), \boldsymbol{\psi}(3686)(2 \mathrm{~S}), \boldsymbol{\psi}(4040)(3 \mathrm{~S}), \psi(4415)(4 \mathrm{~S})$. $\psi(3770)(1 \mathrm{D}), \psi(4160)(2 \mathrm{D})$

Cornell potential: $V^{(0)}(r)=-\frac{\kappa}{r}+\sigma r+C$.



9

## Why $Y(4260)$ exotic?

$\square \mathrm{Y}(4260), \mathrm{Y}(4360)$ doesn't correspond to a peak in $R$ scan spectrum.

$\square \mathrm{Y}(4260)$ has much smaller coupling to
open charm compare with observed $\psi$ states, which is not an expected behavior of charmonium in open charm range

## Cross section of $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \pi^{+} \pi^{-J} / \psi$

Belle: ee $\rightarrow Y_{\text {ISR }} \pi^{+} \pi^{-} \mathrm{J} / \Psi$
Phys. Rev. Lett. 110, 252002(2013)


$\mathrm{Y}(4260):$ Mass=4251 $\pm 9 \mathrm{MeV}$, width $=120 \pm 12 \mathrm{MeV} \rightarrow \mathrm{Y}(4220): \mathrm{M}=(4222.0 \pm 3.1 \pm 1.4) \mathrm{MeV}, \quad \Gamma=(44.1 \pm 4.3 \pm 2.0) \mathrm{MeV}$ $\mathrm{Y}(4360):$ Mass $=4346 \pm 6 \mathrm{MeV}$, width $102 \pm 10 \mathrm{MeV} \quad \rightarrow \quad \mathrm{M}=(4320.0 \pm 10.4 \pm 7) \mathrm{MeV}, \quad \Gamma=(101.4 \pm 25 \pm 10) \mathrm{MeV}$
$3.86 \mathrm{fb}-1$ at 8 energy points taken at 2017 , and $3.9 \mathrm{fb}-1$ at 8 energy points taken at 2019 can give more precise result.

## Cross section of e $\mathrm{e}^{+} \rightarrow \pi^{0} \pi^{0} \mathrm{~J} / \psi$



The measured cross section compared with $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \pi^{+} \pi^{-J} / \psi$ Satisfy the isospin symmetry.

In the fitting, the parameter of $Y(4320)$ are fixed. And the measured parameter of $Y(4220)$ also agree with that in $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \pi^{+} \pi^{-\mathrm{J}} / \psi$.
$\mathrm{Y}(4220): \mathrm{M}=(4220.4 \pm 2.4 \pm 2.3) \mathrm{MeV}, \quad \Gamma=(46.2 \pm 4.7 \pm 2.1) \mathrm{MeV}$

## $Y(4230)$ in other channels




The mass of $Y(4230)$ now is much smaller than the threshold of DD1(2420)~4285MeV, this reduce the probability (DD1(2420)) molecular description but can't rule it out.

Tetraquark mode
$(q Q)(\bar{q} \bar{Q})=4.244 \mathrm{GeV}$
$(s Q)(\bar{s} \bar{Q})=4.466 \mathrm{GeV}$

## Coupled channels fit

arXiv:1805.03565








## Extraction of $\Gamma_{e e}$ of $\mathrm{Y}(4230)$

- $\Gamma_{e e}$ is the decay width of $\mathrm{Y}(4230)$ to $\mathrm{e}^{+} \mathrm{e}^{-}$.
- Different model predicted $\Gamma_{e e}$ : Hybrid: $\sim 40 \mathrm{eV}($ (CPC40, 081002(2006))

DD1(2420) Molecule: ~500eV (PRD94, 054035(2016))
$\mathbf{4}^{\mathbf{3}} \boldsymbol{S}_{\mathbf{1}} \boldsymbol{c} \overline{\boldsymbol{c}}$ states: $\sim \mathbf{1 K e V}$ (PRD79,094004 (2009))
$3{ }^{3} \mathrm{D}_{1} \boldsymbol{c} \bar{c}$ states: $\sim 44 \mathrm{eV}$

- In Arxiv:2002.05641, they performed a fit to all the $\mathrm{Y}(4230)$ decay channels

$$
\begin{aligned}
\Gamma_{t o t}(s)= & \Gamma_{J / \psi \pi \pi}(s)+\Gamma_{h_{c} \pi \pi}(s)+\Gamma_{D \bar{D}^{*} \pi}(s)+\Gamma_{\psi(2 S) \pi \pi}(s)+\Gamma_{\omega \chi_{c 0}}+\Gamma_{J / \psi \eta} \\
& +\Gamma_{D_{z}^{*} \bar{D}_{s}^{*}}+\Gamma_{D \bar{D}}+\Gamma_{D \bar{D}^{*}}+\Gamma_{D^{*} \bar{D}^{*}}+\Gamma_{0} . \\
\Gamma_{e^{+} e^{-}}= & \frac{4 \alpha}{3} \frac{g_{0}^{2}}{M_{X}} .
\end{aligned}
$$

$$
\sigma_{e^{+} e^{-} \rightarrow X(4260) \rightarrow f}=\frac{3 \pi}{k^{2}}\left|\frac{\sqrt{s \Gamma_{e e} \Gamma_{f}}}{s-M_{X}^{2}+i \sqrt{s \Gamma_{t o t}(s)}}+\sum_{i} \frac{c_{i} e^{i \phi_{i}}}{s-M_{i}^{2}+i \sqrt{s \Gamma_{i}}}+\tilde{c}\right|^{2}
$$

Result : $\Gamma_{e e}=1.302 \mathrm{KeV}$ with Ds*Ds*,$\Gamma_{e e}=0.466 \mathrm{KeV}$ without Ds*Ds*
Exclude the Hybrid or pure ${ }^{3} \mathrm{D}_{1}$ model

## $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow$ open charm mesons



$\square \mathrm{Y}(4230)$ has much smaller coupling to open charm compare with observed $\psi$ states, which is not an expected behavior of charmonium states.

## About the fitting-Unitarity

- Sum of Breit-Wigners doesn't satisfy the S-matrix's unitarity, while K-matrix does.



Timofey Uglov
arXiv:1611.07582v2
Fitting Belle's open charm result with K -matrix method

Fitted well with only $\psi$ states

## Part II: Zc states

Isospin non-zero charmonium like states


## $Z_{c}(4430)^{-}$


$\square$ The first Zc state is observed by Belle in $\mathrm{B} \rightarrow \boldsymbol{\operatorname { C o }} \boldsymbol{\pi}^{-} \boldsymbol{\psi}^{\prime}$
$\mathrm{M}=4433 \pm 4 \pm 2 \mathrm{MeV}$
$\Gamma=45_{-13}^{+18+30} \mathrm{MeV}$
$\square \mathrm{Jp}^{\mathrm{p}}$ prefer to be $\mathbf{1}^{+}$

Other Zc states observed in B decay
$\square \mathrm{Zc}(4050), \mathrm{Zc}(4250)$ in $\mathrm{B} \rightarrow \mathrm{K} \boldsymbol{\pi}^{-} \chi_{c 1}$, Belle, PRD 78 (2008) 072004
$\square \mathrm{Zc}(4200)$ in $\mathrm{B} \rightarrow \mathrm{K} \boldsymbol{\pi}^{-} \mathrm{J} / \boldsymbol{\psi}$, Belle, PRD 90 (2014) 112009

## $\mathbf{Z c}(\mathbf{3 9 0 0})^{ \pm, 0}$ in $\pi^{+} \pi^{-} \mathbf{J} / \Psi, \pi^{0} \pi^{0} \mathrm{~J} / \Psi$



-The mass of $\mathrm{Zc}(3900)$ is in opencharm range and strongly coupled to charm $\rightarrow$ it should contain a (ccbar) pair.
$\cdot Z c(3900)^{ \pm}$is charged $\rightarrow$ need at least two more quarks to form a charge unit.
$Z_{c}(3900)$ is a four quark states?
TTetraquark states?
Phys. Rev. D89,054019(2014);
Phys. Rev. D90,054009(2014);
$\square \mathrm{Zc}(3900)$ is near the threshold of (DD*) $\rightarrow$ A molecular states?
Arxiv:1303.6608, 1304.2882
OR other explanation?


## Zc(3900), Zc(4020)



## $\Gamma\left(\mathrm{Z}_{\mathrm{c}}(3900) \rightarrow \mathrm{DD}^{*}\right) / \Gamma\left(\mathrm{Z}_{\mathrm{c}}(3900) \rightarrow \pi \mathrm{J} / \psi\right)$

- This ratio is important for discriminating the Zc model.
- Experiment result without interference considered

$$
\Gamma\left(\mathrm{Zc} \rightarrow \mathrm{DD}^{*}\right) / \Gamma(\mathrm{Zc} \rightarrow \pi J / \psi)=6.2 \pm 2.7
$$

- Theoretical work, PRD94, 094017 (2016)
$>$ TetraQuark

$$
\begin{aligned}
\Gamma\left(Z_{c}^{+} \rightarrow J / \psi+\pi^{+}\right) & =\left(4.3_{-0.6}^{+0.7}\right) \mathrm{MeV}, \\
\Gamma\left(Z_{c}^{+} \rightarrow \eta_{c}+\rho^{+}\right) & =\left(8.0_{-1.0}^{+1.2}\right) \mathrm{MeV}, \\
\Gamma\left(Z_{c}^{+} \rightarrow \bar{D}^{0}+D^{*+}\right) & \propto 10^{-9} \mathrm{MeV}, \\
\Gamma\left(Z_{c}^{+} \rightarrow \bar{D}^{* 0}+D^{+}\right) & \propto 10^{-9} \mathrm{MeV} .
\end{aligned}
$$

>Molecule

$$
\begin{aligned}
\Gamma\left(Z_{c}^{+} \rightarrow J / \psi+\pi^{+}\right) & =(1.8 \pm 0.3) \mathrm{MeV}, \\
\Gamma\left(Z_{c}^{+} \rightarrow \eta_{c}+\rho^{+}\right) & =\left(3.2_{-0.4}^{+0.5}\right) \mathrm{MeV}, \quad \text { prefered } \\
\Gamma\left(Z_{c}^{+} \rightarrow \bar{D}^{0}+D^{*+}\right) & =\left(10.0_{-1.4}^{+1.7}\right) \mathrm{MeV}, \\
\Gamma\left(Z_{c}^{+} \rightarrow \bar{D}^{* 0}+D^{+}\right) & =\left(9.0_{-1.3}^{+1.6}\right) \mathrm{MeV} .
\end{aligned}
$$

## Determination of Jp of $\mathrm{Zc}(3900)$

PRL 119, 072001 (2017)


$\mathrm{Vs}=4.23 \mathrm{GeV}$


-PWA with helicity formalism taking $\pi^{+} \pi \pi^{-J} / \psi$ as final states
$\square$ Simultaneous fit to data samples at 4.23 GeV and 4.26 GeV
$\square \pi^{+} \pi$ spectrum is parameterized with $\sigma, f_{0}(980), f_{2}(1270)$ and $f_{0}(1370)$

## Determination of Jp of $\mathrm{Zc}(3900)$

- Zc is parameterized with Flatte formula

$$
B W\left(s, M, g_{1}^{\prime}, g_{2}^{\prime}\right)=\frac{1}{s-M^{2}+i\left[g_{1}^{\prime} \rho_{1}(s)+g_{2}^{\prime} \rho_{2}(s)\right]}
$$

- $\mathrm{M}=(3901.5 \pm 2.7 \pm 38.0) \mathrm{MeV}, \mathrm{g}_{1}{ }^{\prime}=(0.075 \pm 0.006 \pm 0.025) \mathrm{GeV}^{2}$,

$$
\mathrm{g}_{2}^{\prime} / \mathrm{g}_{1}^{\prime}=27.1 \pm 2.0 \pm 1.9
$$

Which corresponding to pole Mass= $(3881.2 \pm 4.2 \pm 52.7) \mathrm{MeV}$, pole width $=(51.8 \pm 4.6 \pm 36.0) \mathrm{MeV}$

- $\mathrm{J}^{\mathrm{p}}$ of Zc favor to be $1^{+}$with statistical significance larger than7 $\mathbf{\sigma}$ over other quantum numbers
- The significance of $\mathbf{Z c}(\mathbf{4 0 2 0})$ process is found to be $\mathbf{3 \sigma}$


## PWA of $\mathbf{e}^{+} \mathbf{e}^{-} \rightarrow \boldsymbol{\pi}^{0} \boldsymbol{\pi}^{0} \mathbf{J} / \boldsymbol{\psi}$

PRD 102, 012009 (2020)


$M_{\mathrm{r}^{2} / \operatorname{lig}_{\varphi}}\left(\mathrm{GeV} / \mathrm{C}^{2}\right)$

$M_{\mathrm{r}^{2} \mathrm{~K}^{0}}\left(\mathrm{GeV} / \mathrm{c}^{2}\right)$

Simultaneous fit of 4 energy points near Y(4230).
$\mathbf{J}^{\mathrm{p}}$ of $\mathbf{Z c}^{\mathbf{0}}$ favor to be $\mathbf{1}^{+}$confirmed

## $\sigma\left(e^{+} e^{-} \rightarrow \pi^{0} Z_{c}^{0} \rightarrow \pi^{0} \pi^{0} \mathrm{~J} / \psi\right)$



- A hint of correlation between $\mathrm{Y}(4220)$ and $\mathrm{Zc}(3900)$

| Parameters | Solution I | Solution II |
| :--- | :---: | :---: |
| $p_{0}\left(c^{2} / \mathrm{MeV}\right)$ | $0.0 \pm 11.3$ |  |
| $p_{1}$ | $(1.8 \pm 1.9) \times 10^{-2}$ |  |
| $M(R)\left(\mathrm{MeV} / c^{2}\right)$ | $4231.9 \pm 5.3$ |  |
| $\left.\Gamma_{\text {tot }}(R)\right)(\mathrm{MeV})$ | $41.2 \pm$ | 16.0 |
| $\Gamma_{\mathrm{ee}} \mathcal{B}_{R \rightarrow \pi^{0} Z_{c}(3900)^{\circ}}(\mathrm{eV})$ | $0.53 \pm 0.15$ | $0.22 \pm 0.25$ |
| $\phi(R)$ | $(-103.9 \pm 33.9)^{\circ}$ | $(112.7 \pm 43.0)^{\circ}$ |

## Evidence of $\mathbf{e}^{+} \mathbf{e}^{-} \rightarrow \boldsymbol{\pi} \boldsymbol{Z}_{\boldsymbol{c}}^{(1)}, \boldsymbol{Z}_{\boldsymbol{c}}^{(\prime)} \rightarrow \boldsymbol{\rho}^{ \pm} \boldsymbol{\eta}_{\boldsymbol{c}}$


$\square$ Nine $\eta_{c}$ channels are used to reconstruct $\eta_{c}$.
$\square$ After the $\eta_{c}$ and $\rho$ mass window, a hint of $Z_{c}(3900)$ peak can be seen on the recoiled mass of the bachelor $\pi$.
$\square$ The blue histogram is $\eta_{c}$ sideband. $Z_{c}$ parameter are fixed to latest measurement.
$\square$ Strong evidence of $\mathrm{Zc}(3900) \rightarrow \rho \eta_{c}$ is observed at $\mathrm{Vs}=4.23 \mathrm{GeV}$, with statistical
significance $4.3 \sigma$ ( $3.9 \sigma$ including systematic uncertainty)
$\square$ No significant $Z^{\prime}(4020) \rightarrow \rho \eta_{c}$ observed. (statistical significance 1.0 $\sigma$ )

## $\mathrm{R}_{\mathrm{Zc}}=\operatorname{Br}\left(\mathrm{Z}_{\mathrm{c}} \rightarrow \boldsymbol{\rho} \boldsymbol{\eta}_{\boldsymbol{c}}\right) / \operatorname{Br}\left(\mathrm{Z}_{\mathrm{c}} \rightarrow \boldsymbol{\pi} \mathrm{J} / \boldsymbol{\psi}\right)$

TABLE III. Comparison of the measured $R_{Z_{c}(3900)}$ and $R_{Z_{c}(4020)}$ with the theoretical predictions.

| Ratio | Measurement | Tetraquark | Molecule |
| :--- | :---: | :---: | :---: |
| $R_{Z_{c}(3900)}$ | $2.3 \pm 0.8[29]$ | $230_{-140}^{+330}[12]$ | $0.046_{-0.017}^{+0.025}[12]$ |
|  |  | $0.27_{-0.17}^{+0.40}[12]$ | $1.78 \pm 0.41[17]$ |
|  |  | $0.66[13]$ | $6.84 \times 10^{-3}[18]$ |
|  |  | $0.56 \pm 0.24[14]$ | $0.12[19]$ |
|  |  | $0.95 \pm 0.40[15]$ |  |
|  |  | $1.08 \pm 0.88[16]$ |  |
|  |  | $1.28 \pm 0.37[17]$ |  |
|  |  | $1.86 \pm 0.41[17]$ |  |
| $R_{Z_{c}(4020)}$ | $<1.2[4]$ | $6.6_{-5.8}^{+56.8}[12]$ | $0.010_{-0.004}^{+0.006}[12]$ |

[^0]A.Esposito et.al., PLB 746(2015), 194-201


Tetraquak Type-1
Ietraquak rype-1

## Lineshape at different energy points



The line shape of Zc agree Better with Zc as molecule.

## Challenge in the PWA of XYZ data

- Ambiguity in $\pi \pi$ parametrization:

Sum of Breit-wigners, N/D method, K-matrix method. Different method causing big difference to Zc result.


- Parameterization of Zc: BW, Flatte, Theoretical model dependent line-shape.
- Coherently understanding of all energy points and all data channels.


## Part III: X(3872)



## About X(3872)



- X(3872) was first observed in 2003 by Belle, PRL 91.262001 (2003)
$\ln \mathrm{B}^{+} \rightarrow \mathrm{K}^{+} \boldsymbol{\pi}^{+} \boldsymbol{\pi}^{-} \mathbf{J} / \boldsymbol{\psi}$
- JPC $\mathbf{1}^{++}$(CDF, LHCb)
$\square$ Most recent measurement of mass/width by LHCb: arXiv:2005.13419 Breit-wigner

$$
\begin{aligned}
& m_{\chi_{c 1}(3872)}=3871.695 \pm 0.067 \pm 0.068 \pm 0.010 \mathrm{MeV} \\
& \Gamma_{\mathrm{BW}}=1.39 \pm 0.24 \pm 0.10 \mathrm{MeV}
\end{aligned}
$$

$\bullet$ Very close to threshold : $\mathrm{M}\left(D^{0} \bar{D}^{* 0}\right)=3871.70 \pm 0.11 \mathrm{MeV}$


## Determination of the absolute branching fractions of $X(3872)$ decays

- PRD.100.094003(2019)

Fitting to the measured branching fraction give by different Collaborations.
The average branching fraction of $X(3872)$ decay.
$\chi_{c 1}(3872)$ DECAY MODES

| Mode |  | Fraction $\left(\Gamma_{i} / \Gamma\right)$ |
| :--- | :--- | :--- |
| $\Gamma_{1}$ | $e^{+} e^{-}$ |  |
| $\Gamma_{2}$ | $\pi^{+} \pi^{-} J / \psi(1 S)$ | $>3.2 \%$ |
| $\Gamma_{3}$ | $\rho^{0} J / \psi(1 S)$ |  |
| $\Gamma_{4}$ | $\omega J / \psi(1 S)$ | $>2.3 \%$ |
| $\Gamma_{5}$ | $D^{0} \overline{D^{0}} \pi^{0}$ | $>40 \%$ |
| $\Gamma_{6}$ | $\bar{D}^{* 0} D^{0}$ | $>30 \%$ |
| $\Gamma_{7}$ | $\gamma \gamma$ |  |
| $\Gamma_{8}$ | $D^{0} \bar{D}^{0}$ |  |
| $\Gamma_{9}$ | $D^{+} D^{-}$ |  |
| $\Gamma_{10}$ | $\gamma \chi_{c 1}$ |  |
| $\Gamma_{11}$ | $\gamma \chi_{c 2}$ |  |
| $\Gamma_{12}$ | $\pi^{0} \chi_{c 2}$ |  |
| $\Gamma_{13}$ | $\pi^{0} \chi_{c 1}$ |  |
| $\Gamma_{14}$ | $\pi^{0} \chi_{c 0}$ |  |
| $\Gamma_{15}$ | $\gamma J / \psi$ |  |
| $\Gamma_{16}$ | $\gamma \psi(2 S)$ |  |
| $\Gamma_{11}$ | $\pi^{+} \pi^{-} \eta_{c}(1 S)$ |  |
| $\Gamma_{18}$ | $\pi^{+} \pi^{-} \chi_{c 1}$ |  |
| $\Gamma_{19}$ | $p \bar{p}$ |  |


| Parameter index | Decay mode | Branching fraction |
| :--- | :--- | :--- |
| 1 | $X(3872) \rightarrow \pi^{+} \pi^{-} J / \psi$ | $\left(4.1_{-1.1}^{+1.9}\right) \%$ |
| 2 | $X(3872) \rightarrow D^{* 0} \bar{D}^{0}+$ c.c. | $\left(52.4_{-14.3}^{+25.3}\right) \%$ |
| 3 | $X(3872) \rightarrow \gamma J / \psi$ | $\left(1.1_{-0.6}^{+0.6}\right) \%$ |
| 4 | $X(3872) \rightarrow \gamma \psi(3686)$ | $\left(2.4_{-0.8}^{+1.3}\right) \%$ |
| 5 | $X(3872) \rightarrow \pi^{0} \chi_{c 1}$ | $\left(3.6_{-1.6}^{+2.6} \%\right.$ |
| 6 | $X(3872) \rightarrow \omega J / \psi$ | $\left(4.4_{-1.3}^{+2.3}\right) \%$ |
| 7 | $B^{+} \rightarrow X(3872) K^{+}$ | $(1.9 \pm 0.6) \times 10^{-4}$ |
| 8 | $B^{0} \rightarrow X(3872) K^{0}$ | $\left(1.1_{-0.4}^{+0.5}\right) \times 10^{-4}$ |
|  | $X(3872) \rightarrow$ unknown | $\left(31.9_{-31.5}^{+18.1}\right) \%$ |

## X(3872) exclusive decay modes evidence of $\mathrm{X}(3872) \rightarrow \gamma J / \psi$



PRL124, 242001 (2020)
TABLE I. Relative branching ratios and UL on branching ratios compared with $X(3872) \rightarrow \pi^{+} \pi^{-} J / \psi[18,27]$, where systematic uncertainties have been taken into account.

| Mode | Ratio | UL |
| :--- | :---: | :---: |
| $\gamma J / \psi$ | $0.79 \pm 0.28$ | $\ldots$ |
| $\gamma \psi^{\prime}$ | $-0.03 \pm 0.22$ | $<0.42$ |
| $\gamma D^{0} \overline{D^{0}}$ | $0.54 \pm 0.48$ | $<1.58$ |
| $\pi^{0} D^{0} \overline{D^{0}}$ | $-0.13 \pm 0.47$ | $<1.16$ |
| $D^{* 0} \overline{D^{0}}+$ c.c. | $11.77 \pm 3.09$ | $\cdots$ |
| $\gamma D^{+} D^{-}$ | $0.00_{-0.00}^{+0.48}$ | $<0.99$ |
| $\omega J / \psi$ | $1.6_{-0.3}^{+0.4} \pm 0.2[18]$ | $\cdots$ |
| $\pi^{0} \chi_{c 1}$ | $0.88_{-0.27}^{+0.33} \pm 0.10[27]$ | $\cdots$ |

## Evidence of $\mathrm{X}(3872) \rightarrow \gamma \mathrm{J} / \psi$



- Improved uplimit of
$\boldsymbol{R}_{\psi}=\frac{B\left[X(3872) \rightarrow \gamma \psi^{\prime}\right]}{B[X(3872) \rightarrow \gamma J \psi]}<0.59$
- Previous measurement of $\boldsymbol{R}_{\psi}$

BaBar: $3.4 \pm 1.4$
LHCb: $2.46 \pm 0.64 \pm 0.29$
Belle: <2.1 (CL. 90\%)

- If BaBar or LHCb's result are correct, X(3872) prefer to be a charmonium
- If the result of BESIII is correct, $X(3872)$ prefere to be molecule or a combination of molecule and charmonium.


## Observation of $\mathrm{X}(3872) \rightarrow \boldsymbol{\pi}^{0} \boldsymbol{\chi}_{\boldsymbol{c 1}}(\mathbf{1 P})$



PRL122,202001 (2019)
$\square$ Data sets used:
$9.0 \mathrm{fb}^{-1}$ for $4.15<\mathrm{E}_{\mathrm{cm}}<4.30 \mathrm{GeV}$
$0.7 \mathrm{fb}^{-1}$ for $4.00<\mathrm{E}_{\mathrm{cm}}<4.15 \mathrm{GeV}$
$2.8 \mathrm{fb}^{-1}$ for $4.30<\mathrm{E}_{\mathrm{cm}}<4.60 \mathrm{GeV}$
$\square$ With in range of $4.15<\mathrm{E}_{\mathrm{cm}}<4.30 \mathrm{GeV}$ For the sum of events in all the three $\chi_{c J}$ range, a clear $\mathrm{X}(3872)$ signal is seen with events number $=16.9_{-4.5}^{+5.2}$, and Significance= $4.8 \sigma$
$\square$ No evidence of $X(3872)$ in other $E_{c m}$

## Observation of $\mathrm{X}(3872) \rightarrow \boldsymbol{\pi}^{0} \boldsymbol{\chi}_{\boldsymbol{c 1}}(\mathbf{1 P})$


$\square 10.8_{-3.1}^{+3.8} \mathrm{X}(3872)$ signal observed in $\chi_{c 1}$ range with statistical significance $\mathbf{5 . 2 \sigma}$
$\square$ The branching ratio

$$
\begin{aligned}
& R_{J}=B\left(X \rightarrow \pi^{0} \chi_{c 1}\right) / B\left(X \rightarrow \pi^{+} \pi^{-} J / \psi\right) \\
& \mathbf{R}_{0}<19(90 \% \text { U.L. }) \\
& \mathbf{R}_{1}=\mathbf{0 . 8 8}+-0.23 \pm 0.10 \\
& \mathbf{R}_{2}<\mathbf{1 . 1}(90 \% \text { U.L. })
\end{aligned}
$$



## Comparison between experiment and theory

$\square$ If we use the previous fitted

$$
\operatorname{Br}\left(X(3872) \rightarrow \pi^{0} \chi_{c 1}\right)
$$

| Parameter index | Decay mode | Branching fraction |
| :--- | :--- | :--- |
| 1 | $X(3872) \rightarrow \pi^{+} \pi^{-} J / \psi$ | $\left(4.1_{-1.1}^{+1.9}\right) \%$ |
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| 3 | $X(3872) \rightarrow \gamma J / \psi$ | $\left(1.1_{-0.3}^{+0.6}\right) \%$ |
| 4 | $X(3872) \rightarrow \gamma \psi(3686)$ | $\left(2.4_{-0.8}^{+1.3}\right) \%$ |
| 5 | $X(3872) \rightarrow \pi^{0} \chi_{c 1}$ | $\left(3.6_{-1.6}^{+2.2}\right) \%$ |

$\square$ If $X(3872)$ were the $\chi_{c 1}(2 p)$ state of charmonium, then
From the estimation of [Dubynskiy, Voloshin, PRD 77, 014013 (2008)],

$$
\Gamma\left(X(3872) \rightarrow \pi^{0} \chi_{c J}\right) \sim 0.06 \mathrm{keV}
$$

Which would imply an unrealistically small

$$
\Gamma_{\text {тот }}(X(3872)) \sim 1.7 \mathrm{keV}
$$

$\square$ So this measurement disfavor the $\chi_{c 1}(2 p)$ interpretation of the $X(3872)$.

## $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \gamma \mathrm{X}(3872), \mathrm{X}(3872) \rightarrow \pi^{+} \pi \mathrm{J} / \psi$


$\square$ BESIII observed $e^{+} e^{-} \rightarrow \gamma X(3872), X(3872) \rightarrow \pi^{+} \pi J / \psi$.
$\square e^{+} e^{-} \rightarrow \gamma X(3872) \rightarrow$ Charge parity of $X(3872)=+1$.
$\square$ It seems that $X(3872)$ is from the radiative transition of $Y(4260)$

## $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \gamma \gamma(3872) \rightarrow \gamma \omega \mathrm{J} / \psi$


(1). Cross section measurement of $\mathrm{e}^{+} \mathrm{e}^{-} \rightarrow \gamma \mathrm{X}$ (3872) for (mid) $\omega \mathrm{J} / \psi$ and (right) $\pi^{+} \pi^{-} \mathrm{J} / \psi$ channel
(2). Simultaneous fit to the cross section with a single Breit-Wigner resonance

$$
\begin{aligned}
& M[Y(4200)]=4200.6_{-13.3}^{+7.9} \pm 3.0 \mathrm{MeV} / c^{2} \\
& \Gamma[Y(4200)]=115_{-26}^{+38} \pm 12 \mathrm{MeV} \\
& \mathcal{R}=\frac{\mathcal{B}[X(3872) \rightarrow \omega J / 4]}{\mathcal{B}\left[X(3872) \rightarrow \pi^{+} \pi^{-J / \psi]}\right.}=1.6_{-0.3}^{+0.4} \pm 0.2
\end{aligned}
$$

## prospect

- At the summer of 2019, BEPCII has made a small upgrade to increase the beam energy up to $2.45 \mathrm{GeV}(\mathrm{Ecm}=4.9 \mathrm{GeV})$. And the upgrade of luminosity by a factor of 2 is under discussion

- BESIII has taken more data points around $\mathrm{E}_{\mathrm{cm}}=4.23 \mathrm{GeV}$ and above 4.6 GeV .

More precise cross section shape can be obtained.

- And more states can be searched in higher energy region.

Search for $\mathrm{Zcs}, \mathrm{Zc}(4430), \mathrm{Y}(4660)$ at BESIII

## Summary

- A lot of observation of structures from experiment
- More precise measurement of states parameters
- More decay channels observed
- Now we have more constrains for the theoretical models.
- 1777(Oxygen named) $\rightarrow 1869$ (the periodic table of
 elements) $\rightarrow 1897$ (electron discovered) $\rightarrow 1918$ (proton discovered) $\rightarrow 1932$ (neutron discovered)
- Which stage are we at?


[^0]:    BESIII result
    Theoretical prediction

